

# Obtaining continuum physics from dynamical simulations of Hamiltonian lattice gauge theories

Speaker: Christopher Kane<sup>1,2</sup>

Collaborators: Siddarth Hariprakash<sup>3</sup>, Christian Bauer<sup>4</sup>

Based off work in: [\[arXiv:2506.16559\]](https://arxiv.org/abs/2506.16559)

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<sup>1</sup>Maryland Center for Fundamental Physics and Department of Physics

<sup>2</sup>Joint Center for Quantum Information and Computer Science

<sup>3</sup>IonQ

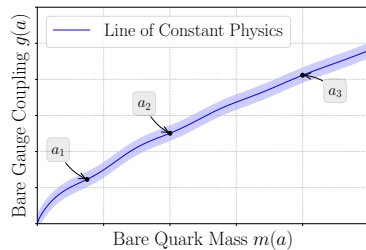
<sup>4</sup>Physics Division, Lawrence Berkeley National Laboratory

# Cost of a quantum simulation

- Calculating physical observables requires taking the continuum limit  $a \rightarrow 0$
- Fair comparisons between methods must include this cost
- Error from approximate time evolution can spoil continuum limit
- Procedure to remove time evolution errors when using Trotter methods developed  
[Marcela Carena, Henry Lamm, Ying-Ying Li, Wanqiang Liu, PRD, arXiv:2107.01166]  
→ not applicable to other algorithms
- **Our Goal:** Develop a general framework for controlling impact of time evolution errors on continuum limit applicable to any algorithm

# Renormalization: taking the continuum limit

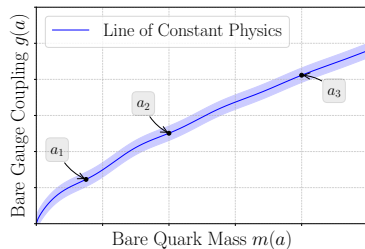
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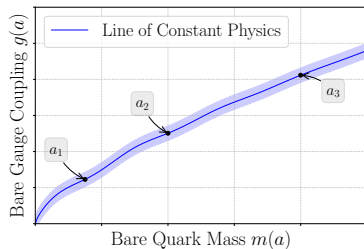
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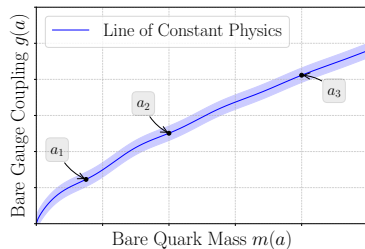
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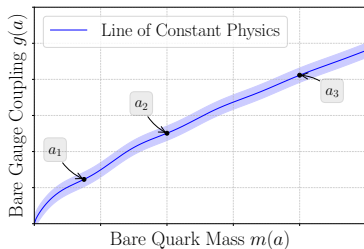
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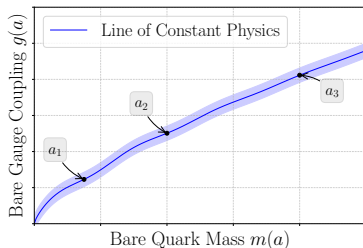


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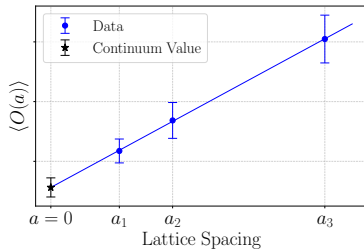
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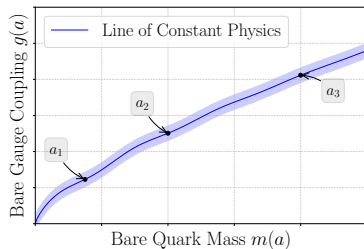
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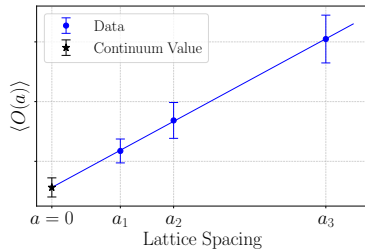
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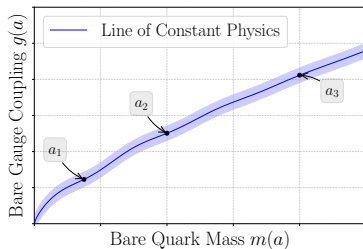
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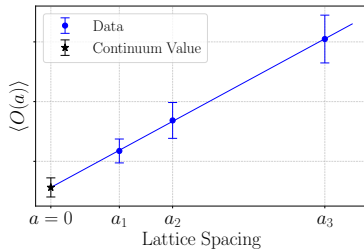
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- c Extrapolate to continuum



# Effective Hamiltonian with Trotter

\*(hatted quantity  $\leftrightarrow$  dimensionless)

Look at effective Hamiltonian simulated using Trotter [Marcela Carena, Henry Lamm, Ying-Ying Li, Wanqiang Liu, PRD, arXiv:2107.01166]

$$e^{-i\hat{\delta}_t(\hat{H}_E + \hat{H}_B)} \approx e^{-i\frac{\hat{\delta}_t}{2}\hat{H}_E} e^{-i\hat{\delta}_t\hat{H}_B} e^{-i\frac{\hat{\delta}_t}{2}\hat{H}_E}$$

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- View as “temporal lattice”, treat  $\hat{\delta}_t$  as parameter in the effective Hamiltonian
- $\hat{\delta}_t \neq 0$  changes physics which changes values used in tuning and scale setting

$$g(\textcolor{violet}{a}) \rightarrow g(\textcolor{violet}{a}, \delta_t), \quad m(\textcolor{violet}{a}) \rightarrow m(\textcolor{violet}{a}, \delta_t), \quad \dots$$

Continuum physics achieved taking simultaneous limit  $\lim_{a \rightarrow 0} \lim_{\delta_t \rightarrow 0}$

(Or, work at fixed anisotropy  $\xi = a/\delta_t$  and extrap  $a \rightarrow 0$ )

# Euclidean simulations to assist renormalization

Main idea: relate  $H_{\text{eff}}$  to Euclidean transfer matrix  $T$  on anisotropic lattice

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- Including fermions introduces  $\mathcal{O}(a_0)$  systematics

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- For  $p$ 'th order, Trotter errors go to zero as  $\mathcal{O}(a^p)$

## Extension to other algorithms?

Can we apply a similar procedure to simulations done using other simulation algorithms?

Consider Quantum Signal Processing as test case

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Provably optimal scaling with  $t$  and  $\epsilon$  [G. Low, I. Chuang, Quantum, arXiv:1606.02685]

$$\text{Calls to } U_H = \mathcal{O}\left(\lambda t + \log \frac{1}{\epsilon}\right)$$

## Breakdown of previous approach to QSP

Can we apply previous approaches to control time evolution errors as  $a \rightarrow 0$  for QSP?

Approximate time evolution operator

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Previous picture relies on effective Hamiltonian formalism

Unclear how to proceed:

→ approximate evolution is not unitary, not clear what  $H_{\text{eff}}$  is

We need an alternative more general approach

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**Key insight:** unlike taking  $a \rightarrow 0$ , taking  $\delta \rightarrow 0$  does not introduce divergences

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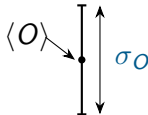
**Implication:** instead of treating  $\delta \neq 0$  with machinery of renormalization, simply treat as source of systematic uncertainty

→ driving time evolution error much below statistical precision effectively takes  $\delta \rightarrow 0$

→ renormalization procedure simplifies to exact time evolution case

# Statistically-Bounded Time Evolution (SBTE) Protocol

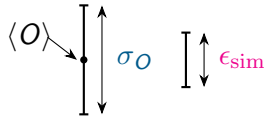
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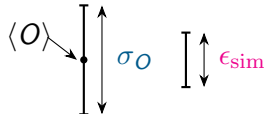
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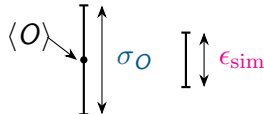
Let  $\alpha$  be ratio of errors:

$$\alpha = \frac{\epsilon_{\text{sim}}}{\sigma_O} = \frac{\text{time evolution error}}{\text{statistical error}}$$

- $\alpha = 1$ : equal error split
- $\alpha \ll 1$ : time evolution error negligible relative to statistical → SBTE

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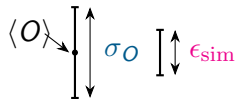
- $\alpha = 1$ : equal error split
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**Key question:** what is the additional computational cost to ensure  $\alpha \ll 1$ ?

→ rigorous cost bounds in our paper [CFK, S. Hariprakash, C. Bauer, [arXiv:2506.16559]

## Schematic cost examples

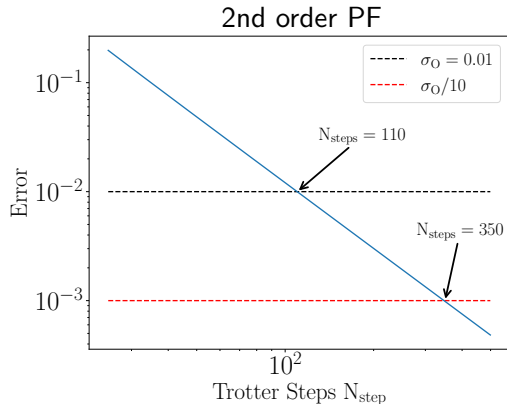
Set  $\sigma_O = 10^{-2}$  and precision factor  $\alpha = 0.1 \rightarrow \epsilon_{\text{sim}} = 10^{-3}$



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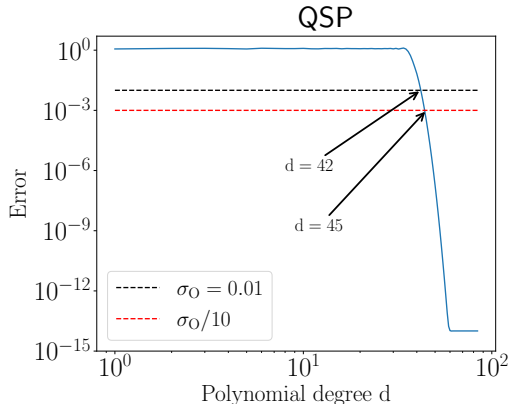
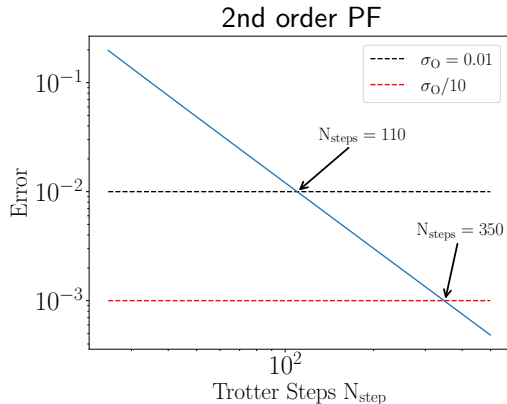
$$\langle O \rangle \quad \begin{array}{|c|} \hline \updownarrow \\ \hline \end{array} \sigma_O \quad \begin{array}{|c|} \hline \updownarrow \\ \hline \end{array} \epsilon_{\text{sim}}$$



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Set  $\sigma_O = 10^{-2}$  and precision factor  $\alpha = 0.1 \rightarrow \epsilon_{\text{sim}} = 10^{-3}$

$$\langle O \rangle \rightarrow \left| \begin{array}{c} \updownarrow \\ \bullet \\ \updownarrow \end{array} \right| \sigma_O \left| \begin{array}{c} \updownarrow \\ \bullet \\ \updownarrow \end{array} \right| \epsilon_{\text{sim}}$$



## Summary and conclusions

- Full cost of a quantum simulation requires cost of  $a \rightarrow 0$  limit
- Previously, controlling impact of approximate time evolution errors only understood for Trotter methods
- We introduced a simple, general approach applicable to any quantum algorithm  
→ Statistically-Bounded Time Evolution Protocol
- Opens the door to performing end-to-end cost comparisons between different algorithms

Backup slides

# Hamiltonian LGT

## Kogut-Susskind Hamiltonian

$$H_{\text{KS}} = \left[ \frac{g_t^2}{a} \tilde{H}_E - \frac{1}{a g_s^2} \tilde{H}_B + \frac{\kappa}{a} \tilde{H}_{\text{hop}} + m \tilde{H}_M \right]$$

Lorentz invariance broken  $\rightarrow$  gauge coupling for  $H_E$  and  $H_B$  different

- Speed of light:  $c = \frac{g_t}{g_s}$
- Gauge coupling  $g = \sqrt{g_s g_t}$
- Hopping coefficient  $\kappa$ 
  - $\rightarrow$  Euclidean simulations on anisotropic lattices have bare fermionic anisotropy factor  $\gamma_f$
  - $\rightarrow$  need  $\gamma_f$  so physical anisotropy  $\xi = a/a_0$  "seen" by gluons and fermions is the same
  - $\rightarrow$  keeping track of  $\gamma_f$  in Transfer matrix implies need for  $\kappa(a) \neq 1$

Speed of light  $c(a) \neq 1$  changes overall scale of  $H$ :

$$H_{\text{KS}} = \frac{c}{a} \left[ g^2 \tilde{H}_E - \frac{1}{g^2} \tilde{H}_B + \frac{\kappa}{c} \tilde{H}_{\text{hop}} + \left( \frac{a}{c} m \right) \tilde{H}_M \right] \equiv \frac{1}{a_t} \hat{H}_{\text{KS}}$$

# Hamiltonian LGT: Two Different Scales

Hamiltonian

$$\hat{H}_{\text{KS}} = a_t H_{\text{KS}}$$

**Temporal scale:**  $a_t = \frac{a}{c}$

Momentum Operator

$$\hat{P}_{\text{KS}} = a P_{\text{KS}}$$

**Spatial scale:**  $a$  (lattice spacing)

Different “rulers” for temporal and spatial quantities

- **Speed of light  $c(a)$ :** conversion factor between  $a$  and  $a_t$

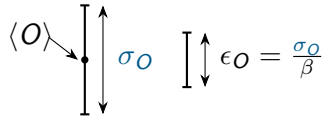
Dimensionless quantities measured on the lattice:

$$\hat{m} = a_t m, \quad \hat{t} = \frac{1}{a_t} t, \quad \hat{p} = a p, \quad \hat{x} = \frac{1}{a} x$$

(Notation: hatted quantities are dimensionless)

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Calculate time-dependent observable  $\langle \hat{O}(t, a) \rangle$  to uncertainty  $\sigma_O$   
→ sources of uncertainty from shot noise, device noise, etc.



If approximate time evolution error  $\epsilon_O \ll \sigma_O$ , we can neglect it

For some  $\beta > 1$ , we require

$$\epsilon_O \equiv \|e^{iHt} \hat{O}(0, a) e^{-iHt} - U_\delta^\dagger(t) \hat{O}(0, a) U_\delta(t)\| \leq \frac{\sigma_O}{\beta}$$

This is guaranteed by choosing time evolution operator error  $\epsilon_{\text{sim}}$  as:

$$\epsilon_{\text{sim}} = \|e^{-iHt} - U_\delta(t)\| \leq \left( \frac{\sigma_O}{2\beta \|\hat{O}(0, a)\|} \right)$$

**Key question:** what is the additional computational cost to ensure this?

## Examples: Product formula and Quantum Signal Processing

Set simulation error  $\epsilon_{\text{sim}} = \|e^{-iHt} - U_{\delta}(t)\| = \frac{\sigma_O}{2\beta\|\hat{O}(0,a)\|}$

$$\langle O \rangle \begin{array}{c} \updownarrow \\ \sigma_O \end{array} \begin{array}{c} \updownarrow \\ \epsilon_O = \frac{\sigma_O}{\beta} \end{array}$$

**SBTE applied to Product Formulas:** extra cost not obviously negligible

$$\begin{aligned} \text{Trotter number : } N_{\text{PF}} &\geq \left( \frac{1}{\epsilon_{\text{sim}}} \right)^{1/p} \tilde{\alpha}^{1/p} t^{1+1/p} && [\text{A. Childs, Y. Su, M. Tran, S. Zhu, PRX} \\ &&& \text{arXiv:1912.08854}] \\ &\geq (2\beta\|\hat{O}(0,a)\|)^{1/p} \times \frac{\tilde{\alpha}^{1/p} t^{1+1/p}}{\sigma_O^{1/p}} \end{aligned}$$

**SBTE applied to QSP:** extra cost negligible *a priori*

$$\begin{aligned} \text{Calls to Block Encoding : } N_{\text{QSP}} &\geq \frac{e\lambda t}{2} + \log \frac{1}{\epsilon_{\text{sim}}}, && \lambda > \|H\| \\ &\geq \frac{e\lambda t}{2} + \log\left(\frac{1}{\sigma_O}\right) + (\log 2\beta\|\hat{O}(0,a)\|) \end{aligned}$$

[G. Low, I. Chuang, Quantum, arXiv:1606.02685]