



TENSOR RENORMALIZATION GROUP STUDY OF COLD AND DENSE TWO- AND THREE-COLOR QCD IN THE STRONG COUPLING LIMIT

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Introduction



Strong Coupling limit of SU(2) case



Strong Coupling limit of SU(3) case



Summary



1 Introduction



2 Strong Coupling limit of SU(2) case



3 Strong Coupling limit of SU(3) case

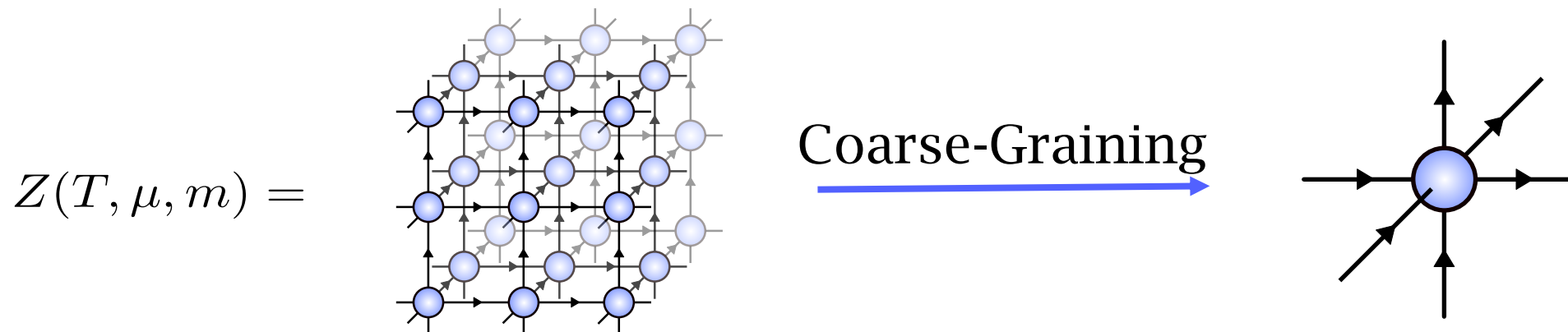
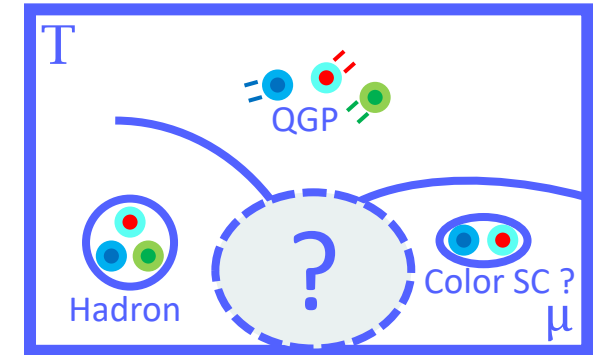


4 Summary

BACKGROUND

GOAL : First-principles calculations of finite density QCD

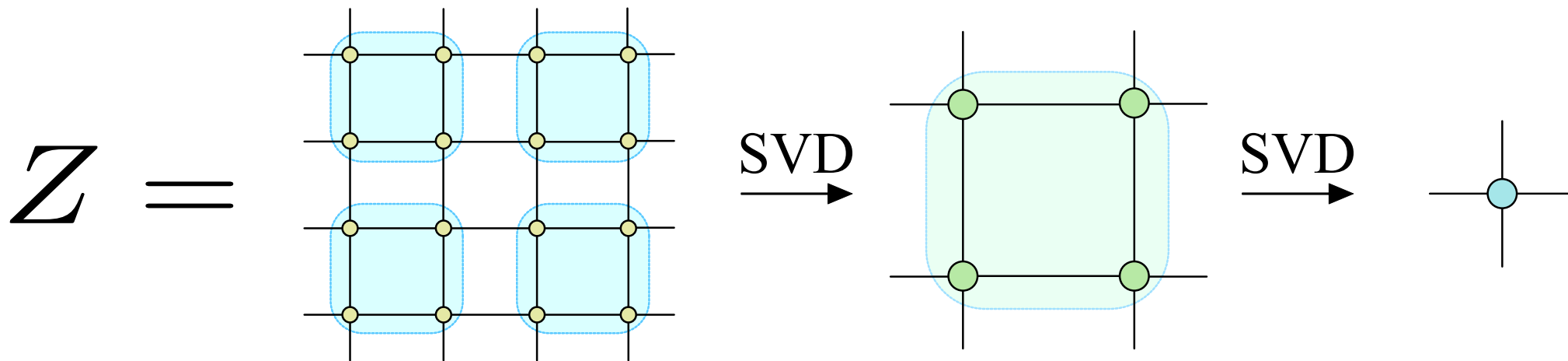
- Finite density is hot topic
- MC simulations are hindered by sign problem
- Tensor Network is next candidate
- We use Tensor Renormalization Group for finite density QCD



ABOUT TRG

[M. Levin and C. P. Nave, (2007).] [Xie et al, (2012).]

- 🔗 **Tensor Renormalization Group** : A computational method based on **SVD**
- 🔗 Rewrite partition function **Z** as a tensor network & approximate it
 - 🔗 Thermodynamic limit
 - 🔗 **NO sign problem !** → For **finite density QCD** ?



STRONG COUPLING LIMIT

S. Akiyama and D. Kadoh, JHEP. 2021, 188 (2021).

$$S[U, \chi, \bar{\chi}] = \underbrace{\beta \sum_n [\text{Gauge action}]}_{\text{Gauge action}} + \underbrace{\sum_n \left[\frac{1}{2} \sum_{\nu=1}^4 \eta_{\nu}(n) (e^{\mu \delta_{\nu,4}} \bar{\chi}(n) U_{\nu}(n) \chi(n + \hat{\nu}) - e^{-\mu \delta_{\nu,4}} \bar{\chi}(n + \hat{\nu}) U_{\nu}^{\dagger}(n) \chi(n)) + m \bar{\chi}(n) \chi(n) \right]}_{\text{Fermion action (staggered, finite } \mu)}$$

Strong coupling limit : $\beta=0$

All link variables at hopping term can be integrated out

$$\int_{SU(N)} dU U_{i_1 j_1} \cdots U_{i_k j_k} U_{j'_1 i'_1}^{\dagger} \cdots U_{j'_k i'_k}^{\dagger} = \sum_{\sigma, \tau \in S_k} \left(\prod_{r=1}^k \delta_{i_r, i'_{\sigma(r)}} \right) \left(\prod_{r=1}^k \delta_{j_r, j'_{\tau(r)}} \right) \text{Wg}_N(\sigma^{-1} \tau).$$

Fermion partition function can be simulated using [Grassmann TN](#)

We consider both SU(2) and SU(3) case

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STRONG COUPLING QC₂D AT ZERO TEMPERATURE

$$S[U, \chi, \bar{\chi}] = \underbrace{\beta \sum_n \left[\frac{1}{2} \sum_{\nu=1}^4 \eta_\nu(n) U_\nu(n) \right]}_{\text{Gauge action}} + \underbrace{\sum_n \left[\frac{1}{2} \sum_{\nu=1}^4 \eta_\nu(n) \left(e^{\mu \delta_{\nu,4}} \bar{\chi}(n) U_\nu(n) \chi(n + \hat{\nu}) - e^{-\mu \delta_{\nu,4}} \bar{\chi}(n + \hat{\nu}) U_\nu^\dagger(n) \chi(n) \right) + m \bar{\chi}(n) \chi(n) \right]}_{\text{Fermion action (staggered, finite } \mu)}$$

Consider strong coupling limit ($\beta \rightarrow 0$)

MF prediction at $T = 0$ (Cold & dense)

low μ : $\langle \bar{\chi} \chi \rangle \neq 0$, $\langle \chi \chi \rangle = 0$

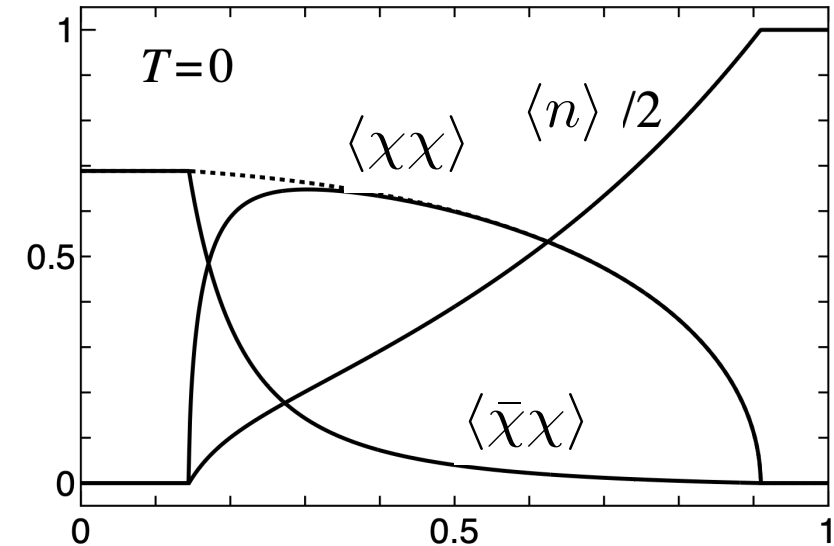
At $\mu_{\text{low}} \approx m_\pi/2$: $\langle \chi \chi \rangle \neq 0$

Higher μ : $\langle \bar{\chi} \chi \rangle = 0$, $\langle \chi \chi \rangle = 0$, $\langle n \rangle = 2$

$\langle n \rangle$: number density

$\langle \bar{\chi} \chi \rangle$: chiral cond.

$\langle \chi \chi \rangle$: diquark cond.



Y. Nishida, K. Fukushima, and T. Hatsuda (2004).

NUMERICAL SETUP I

[1] S. Akiyama, et al. (2021).

- Apply Grassmann-ATRG^[1] method to 4D Strong coupling QC₂D
- Employ multiple GPU parallelization [Y. Sugimoto and S. Sasaki (2025). 2507.21909]

-- Setup --

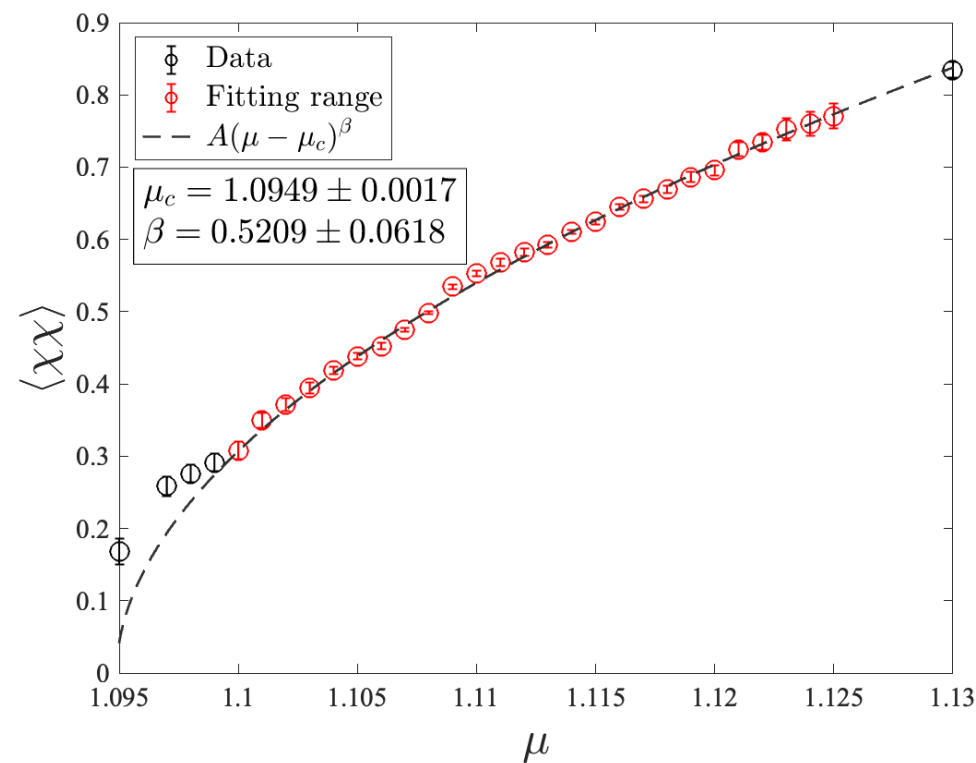
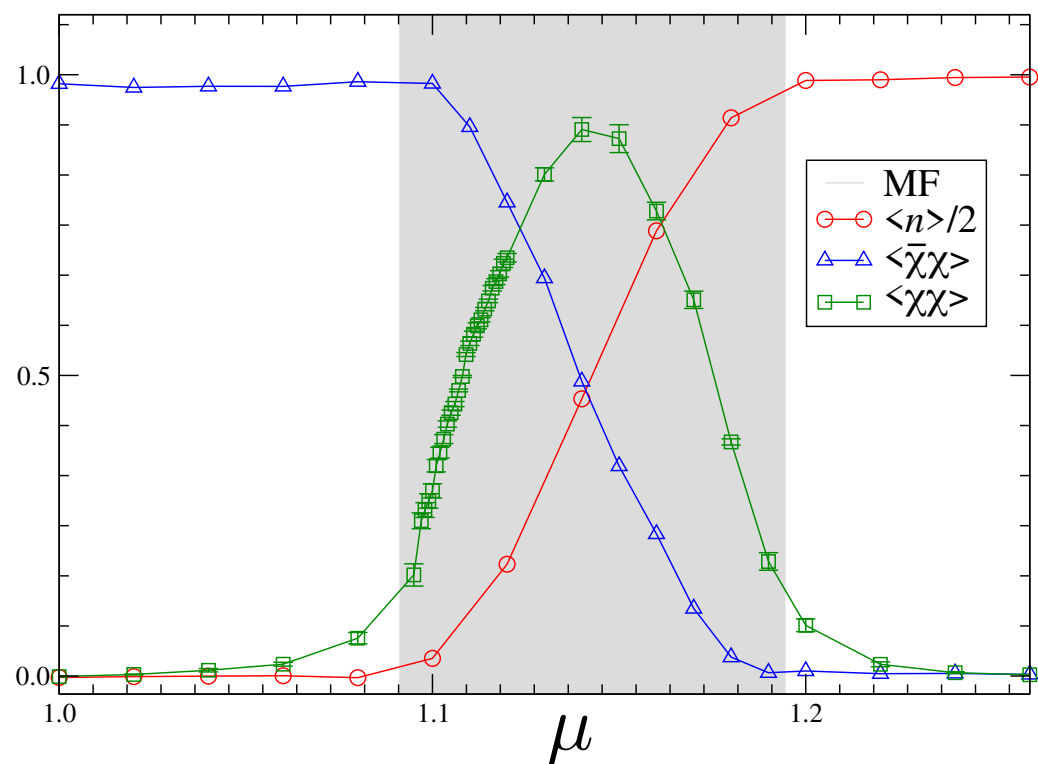
- parameter : $V = 1024^4$, $T = 0$, $m = 1.0$, $D = 55$
 D : bond dimensions (# of retained singular values)

- Evaluate free energy : $f(\mu, m) = -\frac{1}{V} \ln Z(\mu, m)$

→ Calculate $\langle n \rangle$, $\langle \bar{\chi} \chi \rangle$, $\langle \chi \chi \rangle$

NUMERICAL RESULT

- Physical quantities behaves consistent with mean-field theory
- Fitting as power law for $\langle\chi\chi\rangle = A(\mu - \mu_c)^\beta$ gives $\beta = 0.521 \pm 0.062 \sim 0.5(\text{MF})$





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STRONG COUPLING QCD

$$S[U, \chi, \bar{\chi}] = \underbrace{\beta \sum_n \left[\frac{1}{2} \sum_{\mu, \nu} \text{Tr} [U_{\mu\nu}(n)] \right]}_{\text{Gauge action}} + \underbrace{\sum_n \left[\frac{1}{2} \sum_{\nu=1}^4 \eta_{\nu}(n) \left(e^{\mu \delta_{\nu,4}} \bar{\chi}(n) U_{\nu}(n) \chi(n + \hat{\nu}) - e^{-\mu \delta_{\nu,4}} \bar{\chi}(n + \hat{\nu}) U_{\nu}^{\dagger}(n) \chi(n) \right) + m \bar{\chi}(n) \chi(n) \right]}_{\text{Fermion action (staggered, finite } \mu)}$$

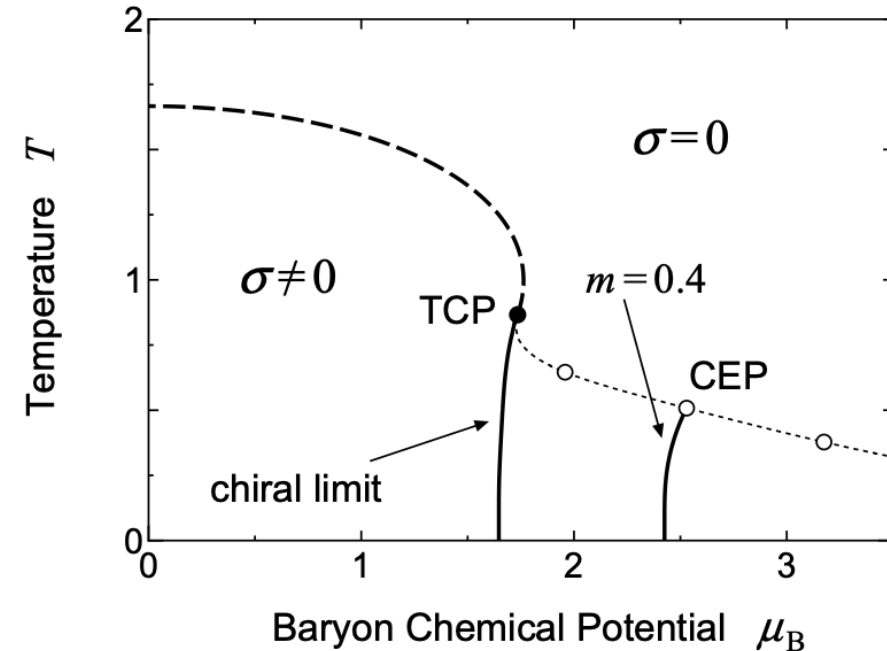
Consider strong coupling limit ($\beta \rightarrow 0$)

No dynamical gluons

Chiral & Nuclear transition

CEP at finite T & mass

1-st order transition at $T = 0$



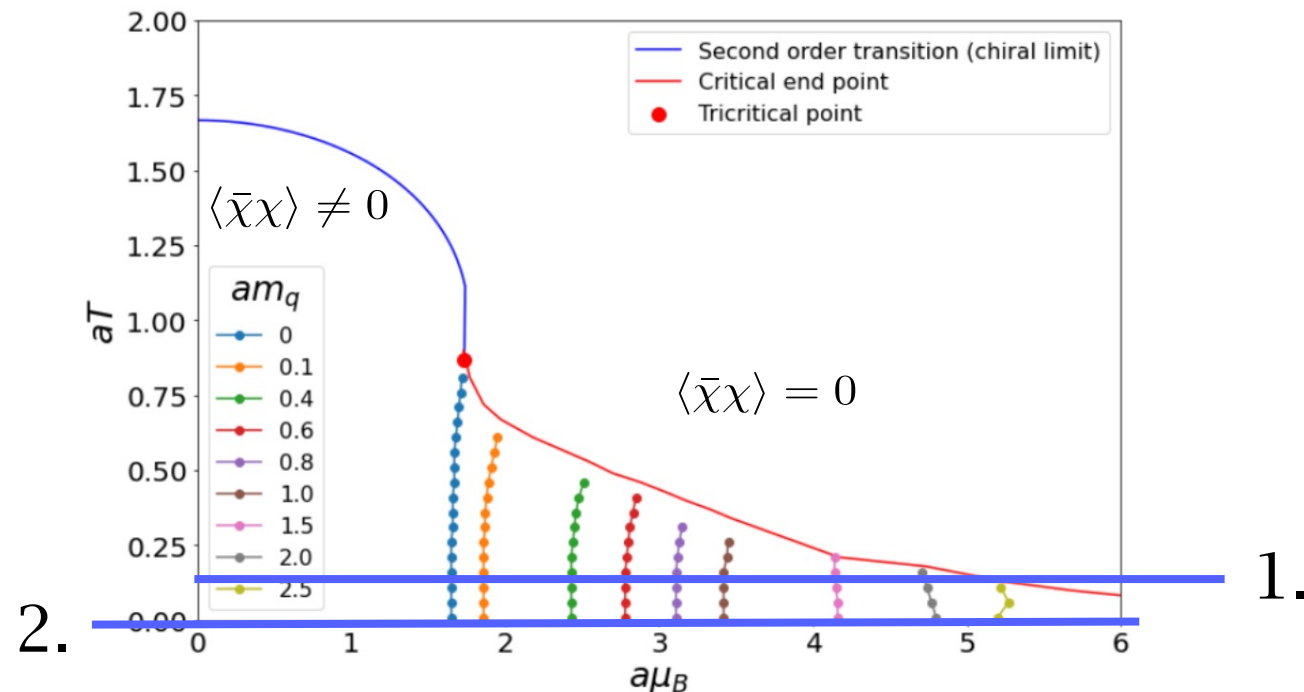
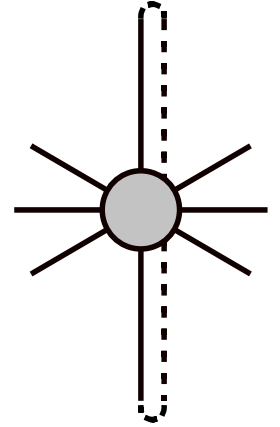
Phys. Rev. D **69**, 094501 (2004).

TARGET

[1] S. Akiyama, et al. (2021).

Two setup with Grassmann-ATRG^[1]

1. $L_\tau = 8$: Coarse-graining along τ direction $\rightarrow$ 3D GATRG
2. $L_\tau = 1024$ (zero temperature) : 4D GATRG

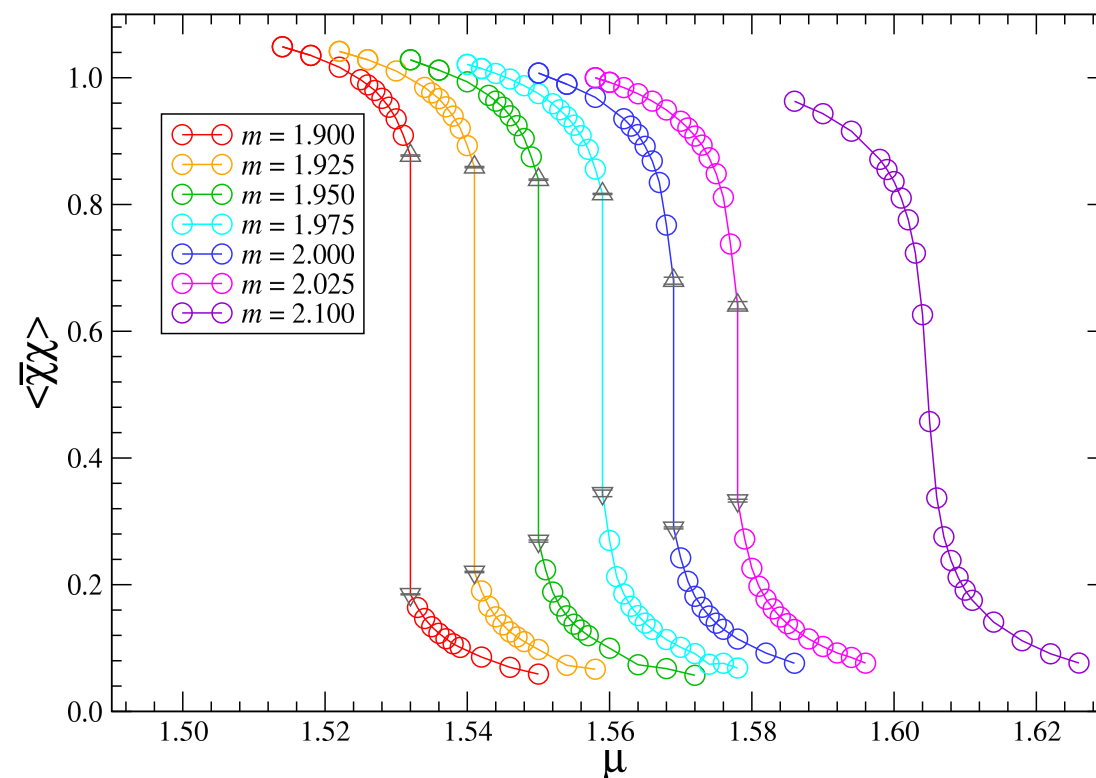
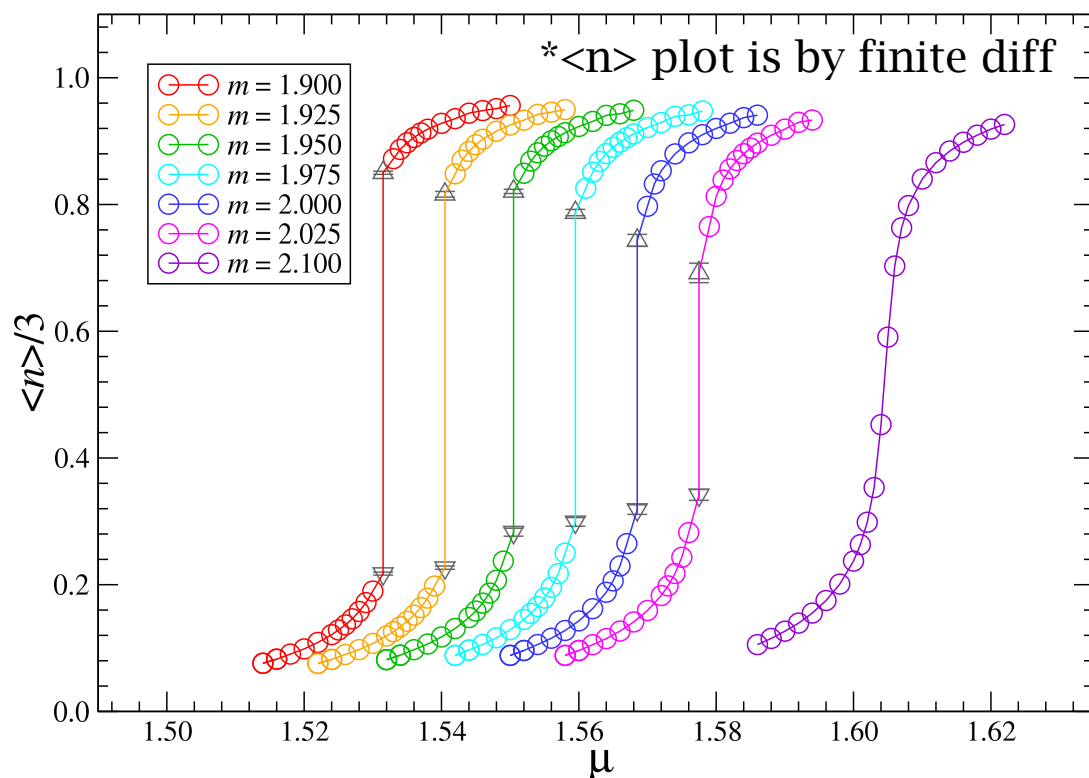


MF prediction [Phys. Rev. D 107, 094514 (2023)]

1. FINITE TEMPERATURE

Investigate $\langle n \rangle$ and $\langle \bar{\chi} \chi \rangle$ at $D_{4d} = 55, D_{3d} = 120$

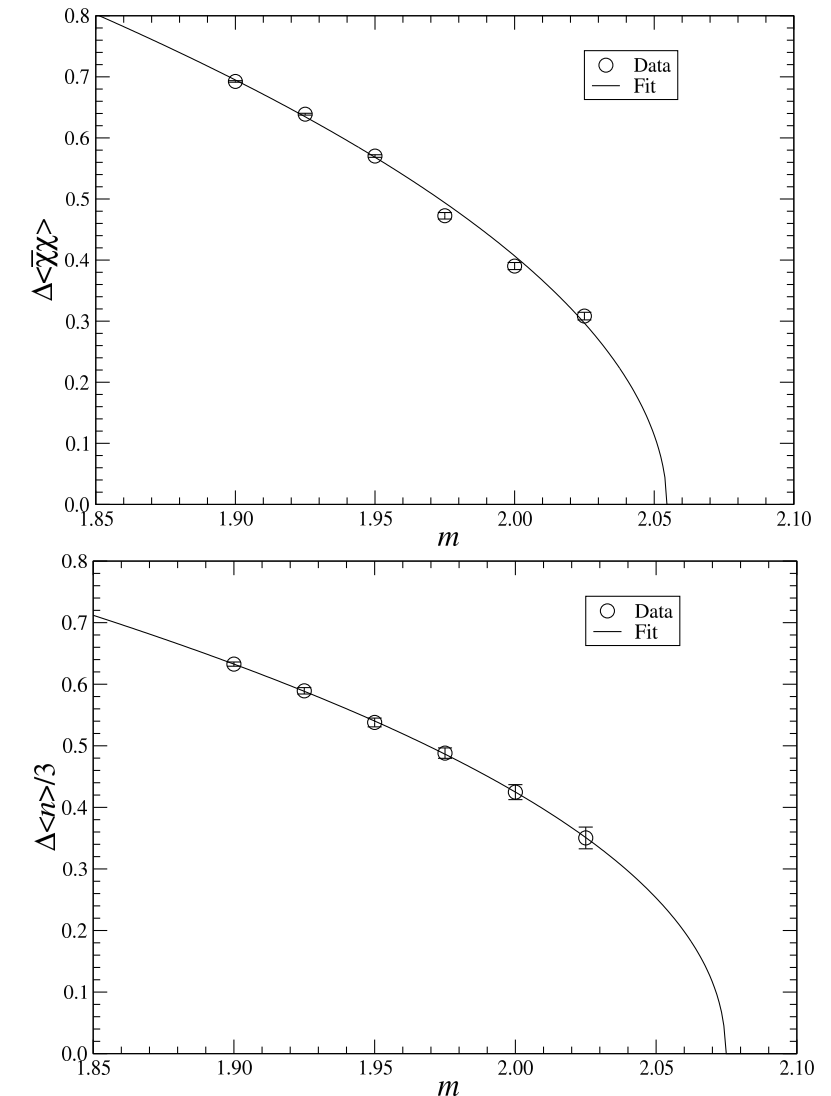
Increasing m , gaps become small



1. FINITE TEMPERATURE

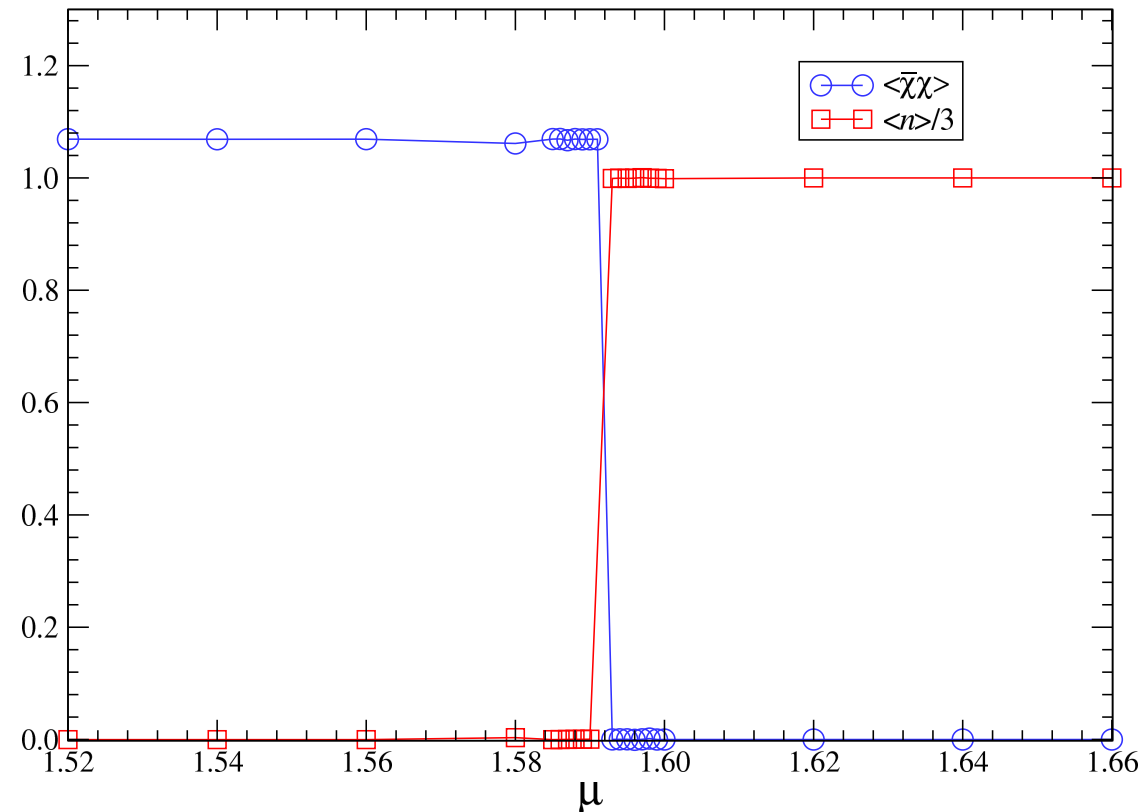
[1] Phys. Rev. D 107, 094514 (2023)
[2] Phys. Rev. D 69, 094501 (2004).

- Fit of critical exponent β and critical mass
 - $\Delta\langle n \rangle : m_c^n = 2.075(23), \beta_n = 0.47(10)$
 - $\Delta\langle \bar{\chi}\chi \rangle : m_c^\chi = 2.0545(34), \beta_\chi = 0.514(20)$
- m_c is consistent with each other
- deviation from previous research
 - MC in dual & reweighting = 1.7(1)
 - Mean field ~ 2.4



2. ZERO TEMPERATURE

- Investigate $\langle n \rangle$ and $\langle \bar{\chi}\chi \rangle$ at $D_{4d} = 55, m = 2.07$
- Clear 1-st order transition is observed $\rightarrow$ expected from MF





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SUMMARY & FUTURE WORK

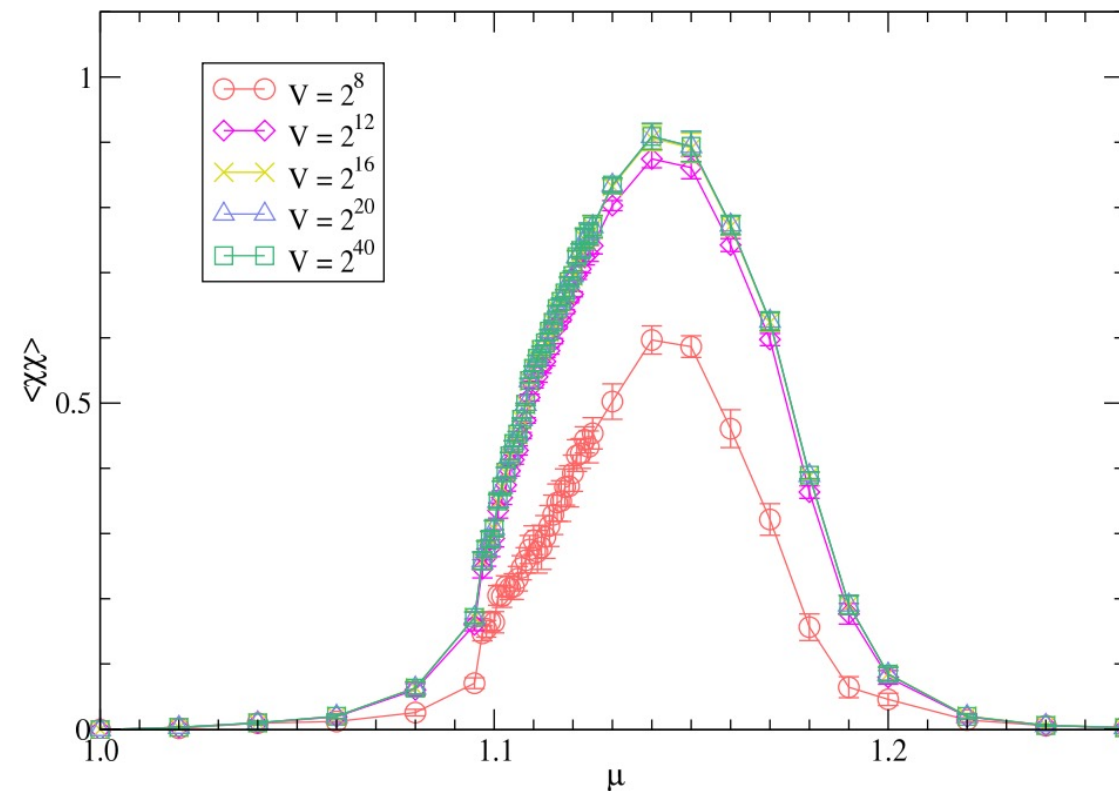
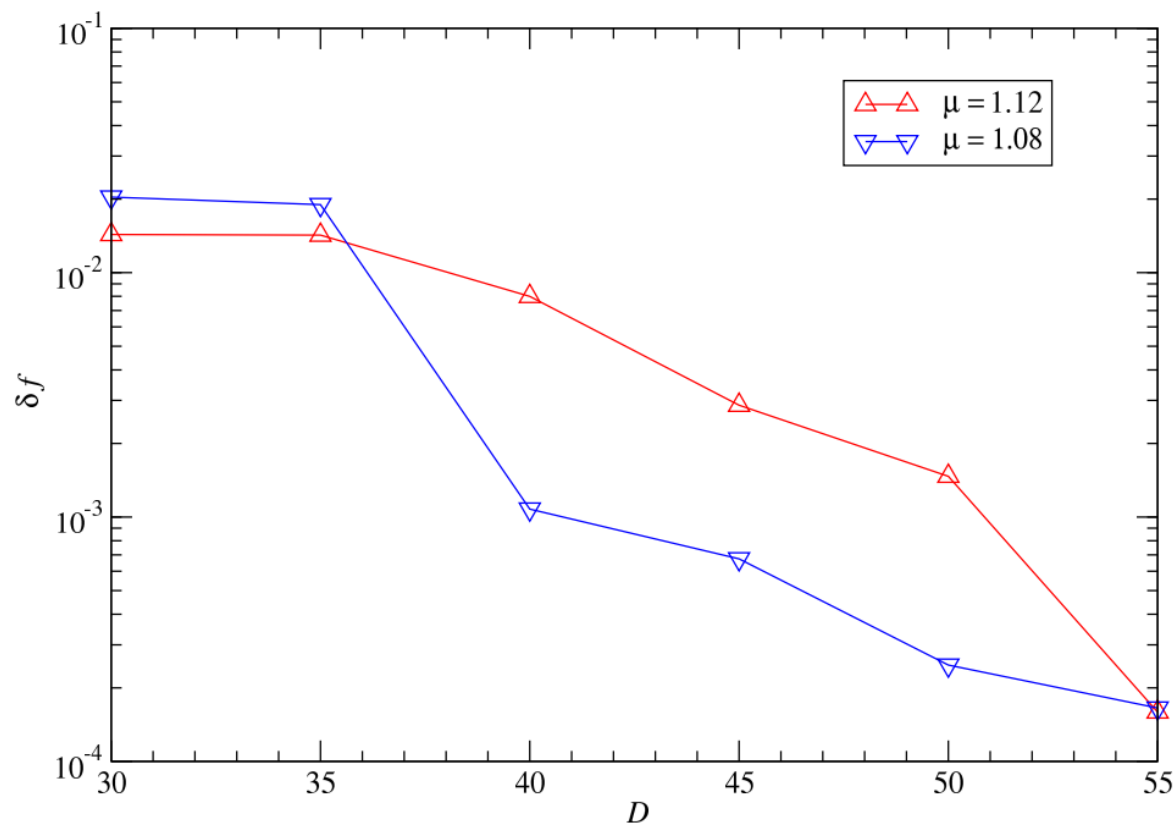
Summary

- 🔗 Performed (3+1)d **finite density QCD** calculations at strong coupling limit
- 🔗 SU(2) : T=0 calculation & estimated critical exponents
- 🔗 SU(3) : Observed 1-st order phase transition at zero temperature

Future work

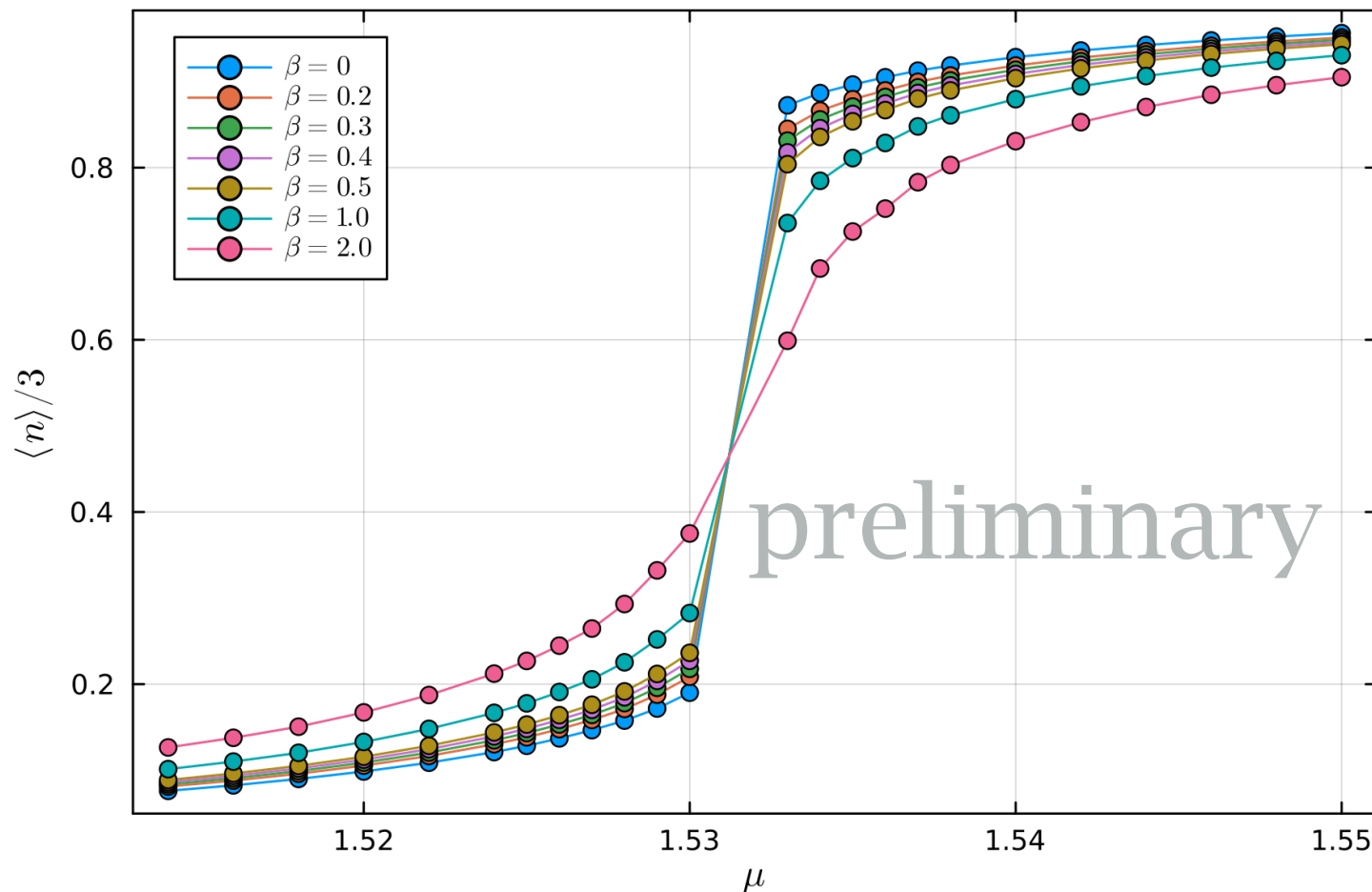
- 🔗 **finite β** in (3+1)d in strong coupling regime
- 🔗 Wilson fermion

- Convergence for SU2

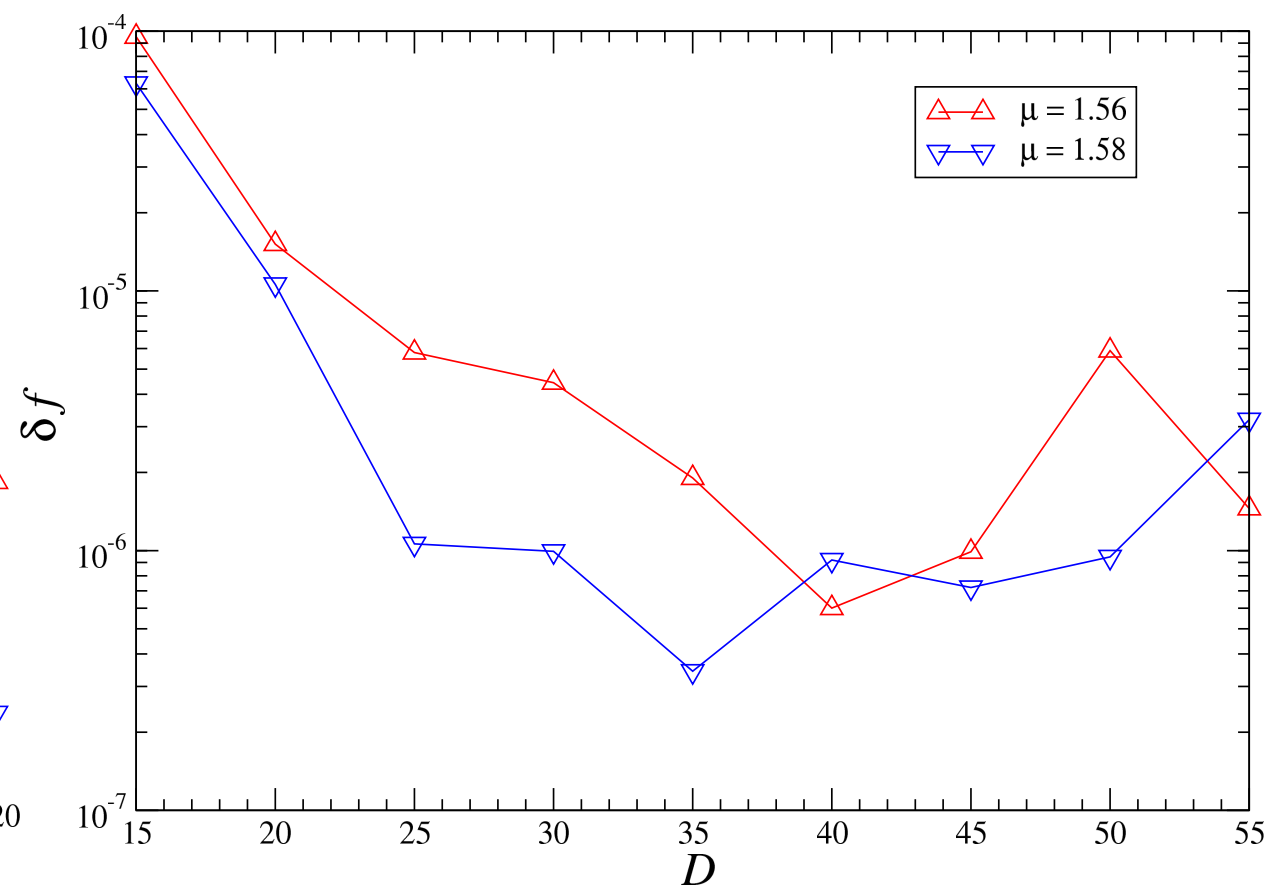
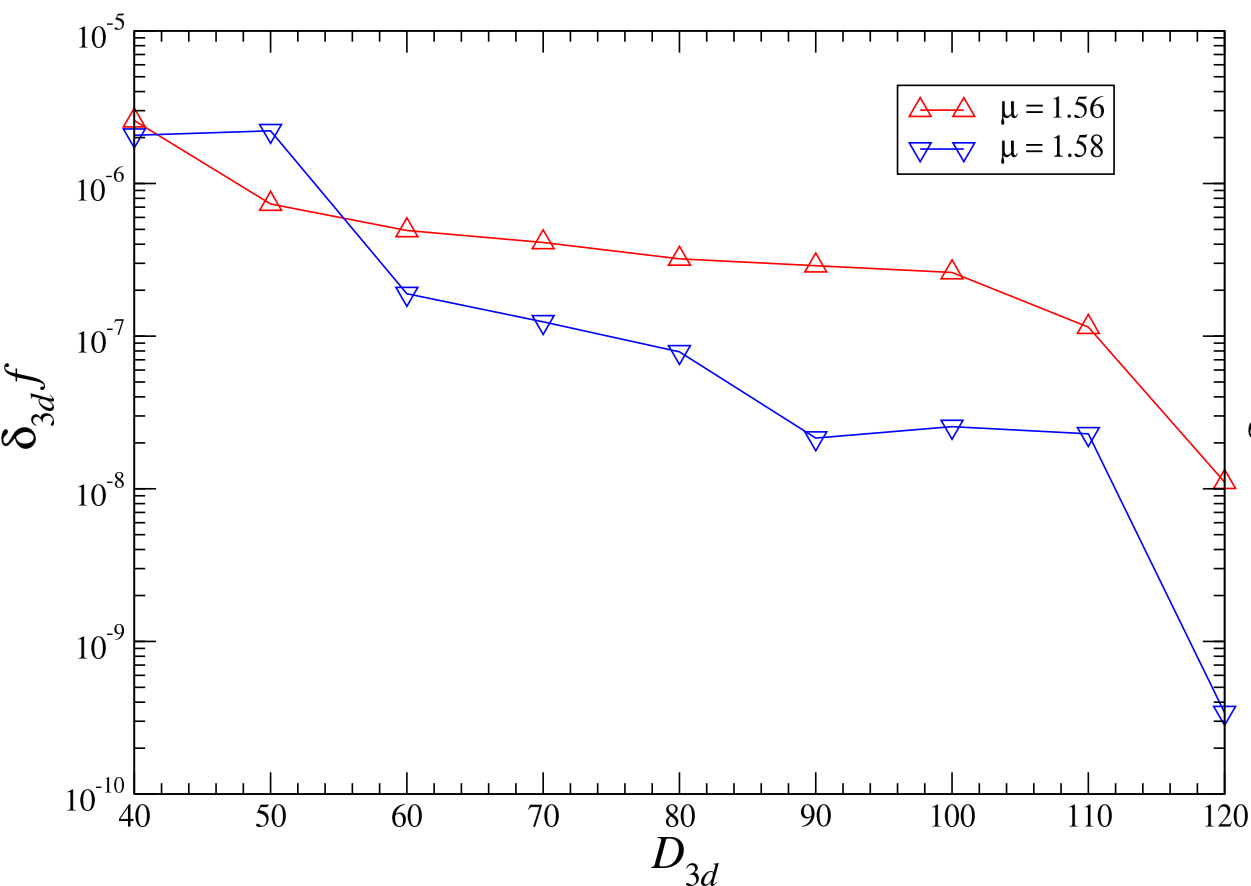


🔗 $O(\beta)$ correction makes gaps smaller

🔗 At some $\beta c(m)$, gap vanishes

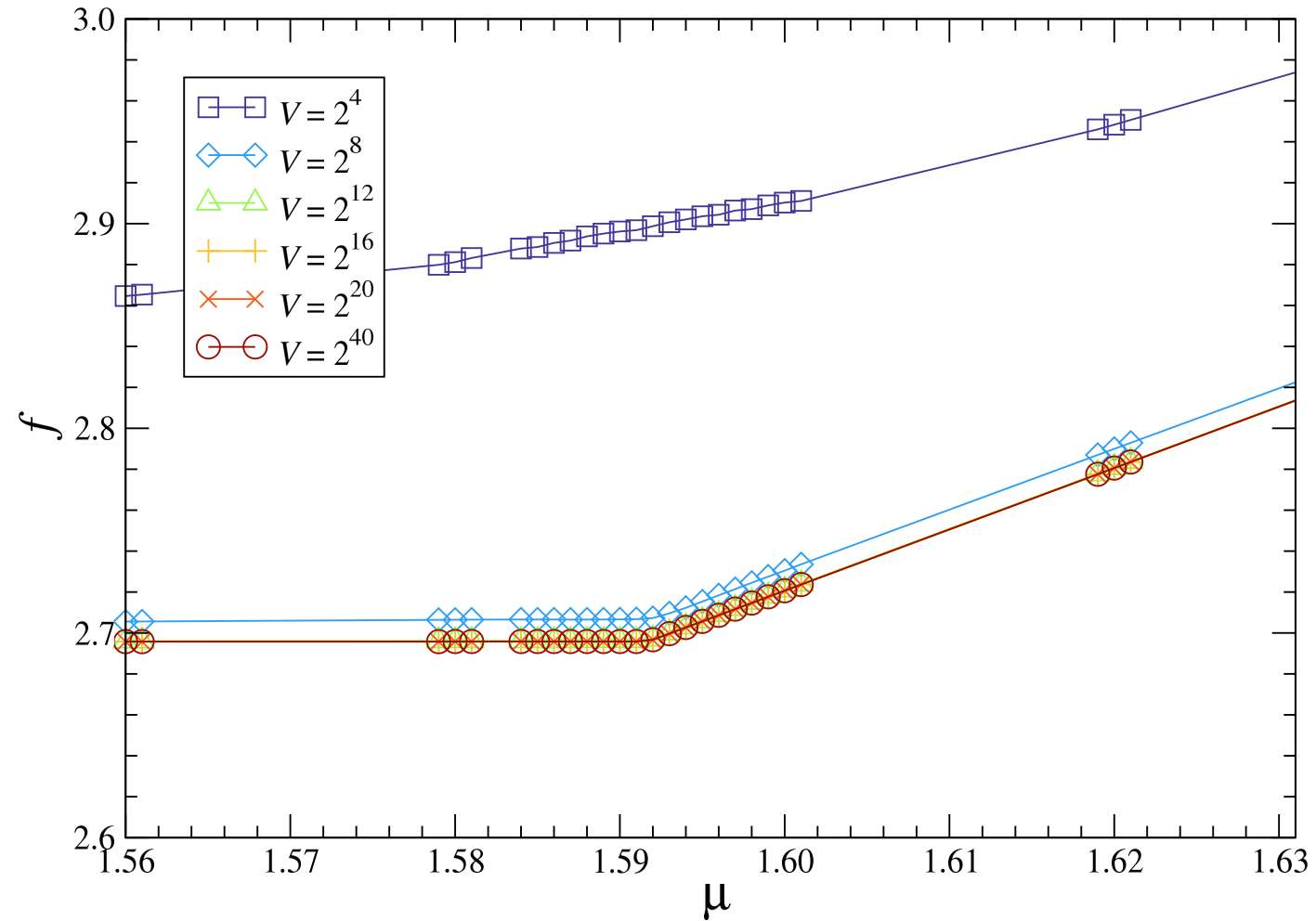


- Convergence for SU3



well converged

Free energy at zero temperature



Free energy at finite temperature

