

The nucleon unpolarized generalized form factors and Mellin moments up to fourth order

Christian Kummer

with: C. Alexandrou, S. Bacchio, M. Constantinou, Y. Li, G. Spanouides

Second Moment:

See C. Alexandrou Thu 2 pm



C. Alexandrou *et al.*,
(2026), 2607.03458



AQTIVATE

Motivation

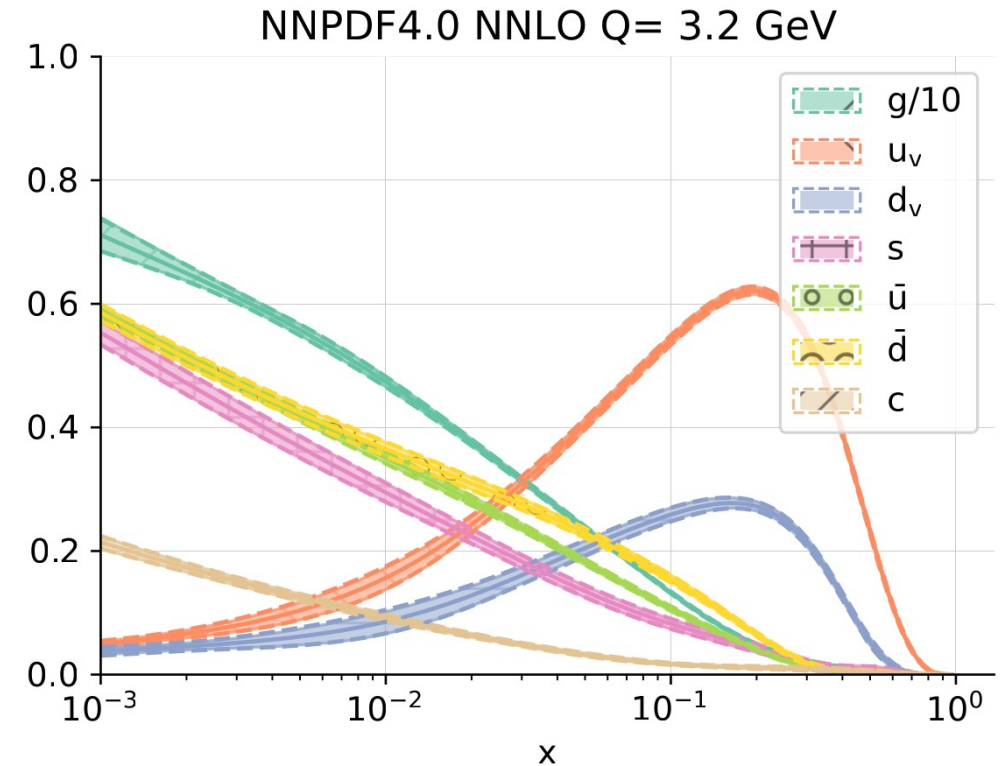
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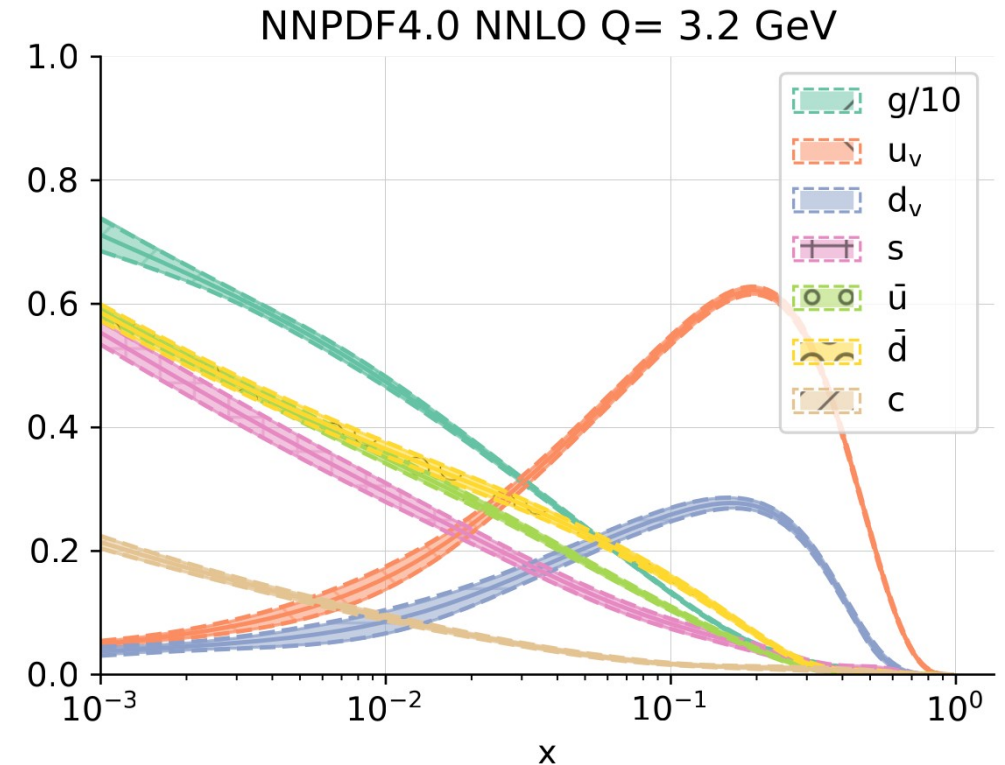
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- Internal structure highly non-perturbative → Lattice QCD
- Interest in generalized parton distributions (GPDs)



NNPDF, EPJC 82, 428 (2022)

Motivation

- Proton is an excellent probe for strong processes
- Internal structure highly non-perturbative → Lattice QCD
- Interest in generalized parton distributions (GPDs)
- Direct calculation of Mellin moments



NNPDF, EPJC 82, 428 (2022)

Definitions

- Leading twist operators

$$\mathcal{O}^{\{\mu_1\mu_2\cdots\mu_n\}} = \bar{\psi} \gamma^{\{\mu_1} i \overleftrightarrow{D}^{\mu_2} \cdots i \overleftrightarrow{D}^{\mu_n\}} \psi$$

- Matrix elements

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- Matrix elements

$$\begin{aligned} \langle N(p') | \mathcal{O}^{\{\mu_1 \mu_2 \dots \mu_n\}} | N(p) \rangle = & \\ \bar{u}(p') \left[\sum_{\substack{i=0 \\ \text{even}}}^{n-1} \left\{ q^{\{\mu_1} \dots q^{\mu_i} P^{\mu_{i+1}} \dots P^{\mu_{n-1}} \gamma^{\mu_n\}} A_{n,i}(q^2) \right. \right. & \\ + q^{\{\mu_1} \dots q^{\mu_i} P^{\mu_{i+1}} \dots P^{\mu_{n-1}} \frac{i \sigma^{\mu_n\}{}^\alpha q_\alpha}{2m_N} B_{n,i}(q^2) \Big\} & \\ + \left. \frac{q^{\{\mu_1} q^{\mu_2} \dots q^{\mu_n\}}}{m_N} C_{n,0}(q^2) \right]_{n \text{ even}} u(p) & \end{aligned}$$

X.-D. Ji, JPG 24, 1181 (1998)

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GFFs

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- Matrix elements

- Connection to GPDs

$$\int_{-1}^1 dx x^{n-1} H(x, \xi, q^2) = \sum_{\substack{i=0 \\ \text{even}}}^{n-1} A_{n,i}(q^2) (-2\xi)^i + C_{n0}(q^2) (-2\xi)^n \Big|_{n \text{ even}}$$

$$\int_{-1}^1 dx x^{n-1} E(x, \xi, q^2) = \sum_{\substack{i=0 \\ \text{even}}}^{n-1} B_{n,i}(q^2) (-2\xi)^i - C_{n0}(q^2) (-2\xi)^n \Big|_{n \text{ even}}$$

Operator Mixing

- Continuum: Lorentz group $SO(1,3)$
- Lattice: Hypercubic group $H(4)$

G. Beccarini *et al.*, NPB 456, 271 (1995)

M. Göckeler *et al.*, PRD 54, 5705 (1996)

Operator Mixing

- Continuum: Lorentz group $SO(1,3)$
- Lattice: Hypercubic group $H(4)$
- Mixing!
- In general: Avoid repeated indices
- Two derivatives $\mathcal{O}\{\mu\nu\rho\}$
with $\mu \neq \nu \neq \rho$
- Three derivatives
 $\mathcal{O}\{1234\}$ $\mathcal{O}\{44ij\}$ $\mathcal{O}\{44ij\} - \mathcal{O}\{kkij\}$

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- Special case: Forward limit
- Leads to boosted frame

G. Beccarini *et al.*, NPB 456, 271 (1995)

M. Göckeler *et al.*, PRD 54, 5705 (1996)

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Methods

- Ensemble at physical point

Ensemble	a [fm]	$(L/a)^3 \times T/a$	m_π [GeV]	L [fm]
cB211.072.64	0.0801(4)	$64^3 \times 128$	0.1393(7)	5.12(3)

C. Alexandrou *et al.*, PRD 98, 054518 (2018)



Methods

- Two- and three-point-functions
- Neglect disconnected

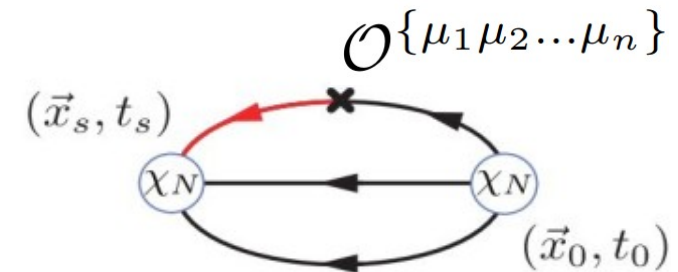
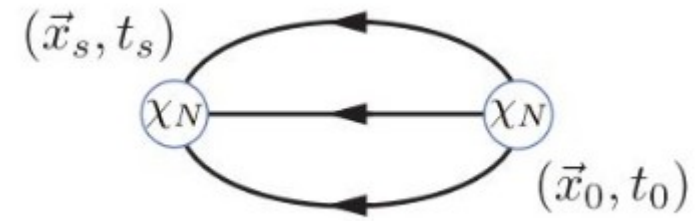
$$C(\Gamma_0, \vec{p}, t_s, t_0) = \sum_{\vec{x}_s} e^{-i(\vec{x}_s - \vec{x}_0) \cdot \vec{p}} \times$$

$$\text{tr}[\Gamma_0 \langle J_N(t_s, \vec{x}_s) \bar{J}_N(t_0, \vec{x}_0) \rangle]$$

$$C^{\{\mu_1 \mu_2 \dots \mu_n\}}(\Gamma, \vec{q}, \vec{p}', t_s, t_{\text{ins}}, t_0) =$$

$$\sum_{\vec{x}_{\text{ins}}, \vec{x}_s} e^{-i(\vec{x}_{\text{ins}} - \vec{x}_0) \cdot \vec{q}} e^{-i(\vec{x}_s - \vec{x}_0) \cdot \vec{p}'} \times$$

$$\text{tr}[\Gamma \langle J_N(t_s, \vec{x}_s) \mathcal{O}^{\{\mu_1 \mu_2 \dots \mu_n\}}(t_{\text{ins}}, \vec{x}_{\text{ins}}) \bar{J}_N(t_0, \vec{x}_0) \rangle]$$



Methods

- Statistics

$\vec{n}'$	N_{mom}	t_s/a	N_{conf}	N_{srcs}	N_{meas}
two derivatives					
(1, 1, 0)	12	8	735	1	8820
(1, 1, 0)	12	10	735	3	26460
(1, 1, 0)	12	12	735	9	79380
(1, 1, 0)	12	14	735	18	158760
three derivatives					
(1, 1, 1)	8	8	735	1	5880
(1, 1, 1)	8	10	735	3	17640
(1, 1, 0)	12	14	735	18	158760

Methods

- Construct Ratio to extract matrix element

$$R^{\mu\nu}(\Gamma; \vec{p}', \vec{p}; t_s, t_{\text{ins}}) = \frac{C^{\mu\nu}(\Gamma; \vec{p}', \vec{p}; t_s, t_{\text{ins}})}{C(\Gamma_0; \vec{p}'; t_s)} \times$$
$$\sqrt{\frac{C(\Gamma_0, \vec{p}; t_s - t_{\text{ins}})C(\Gamma_0, \vec{p}'; t_{\text{ins}})C(\Gamma_0, \vec{p}'; t_s)}{C(\Gamma_0, \vec{p}'; t_s - t_{\text{ins}})C(\Gamma_0, \vec{p}; t_{\text{ins}})C(\Gamma_0, \vec{p}; t_s)}}$$
$$R^{\mu\nu}(\Gamma; \vec{p}', \vec{p}; t_s, t_{\text{ins}}) \xrightarrow[t_{\text{ins}} \gg a]{t_s - t_{\text{ins}} \gg a} \Pi^{\mu\nu}(\Gamma; \vec{p}', \vec{p})$$

Methods

- Multiplicative renormalization

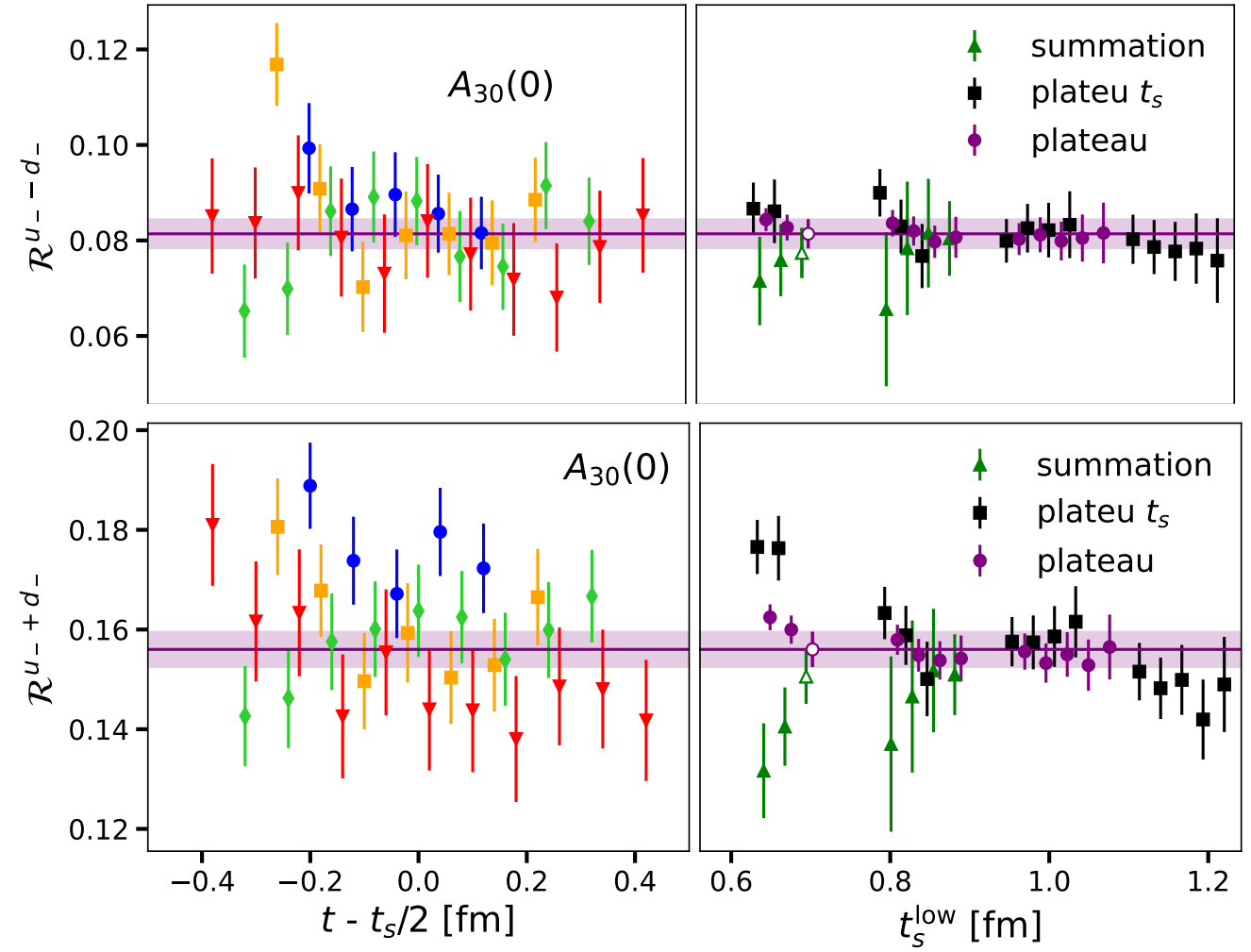
non-perturbative renormalization: RI'-MOM

perturbative renormalization: $\overline{\text{MS}}$

matching at $\bar{\mu} = 2 \text{ GeV}$

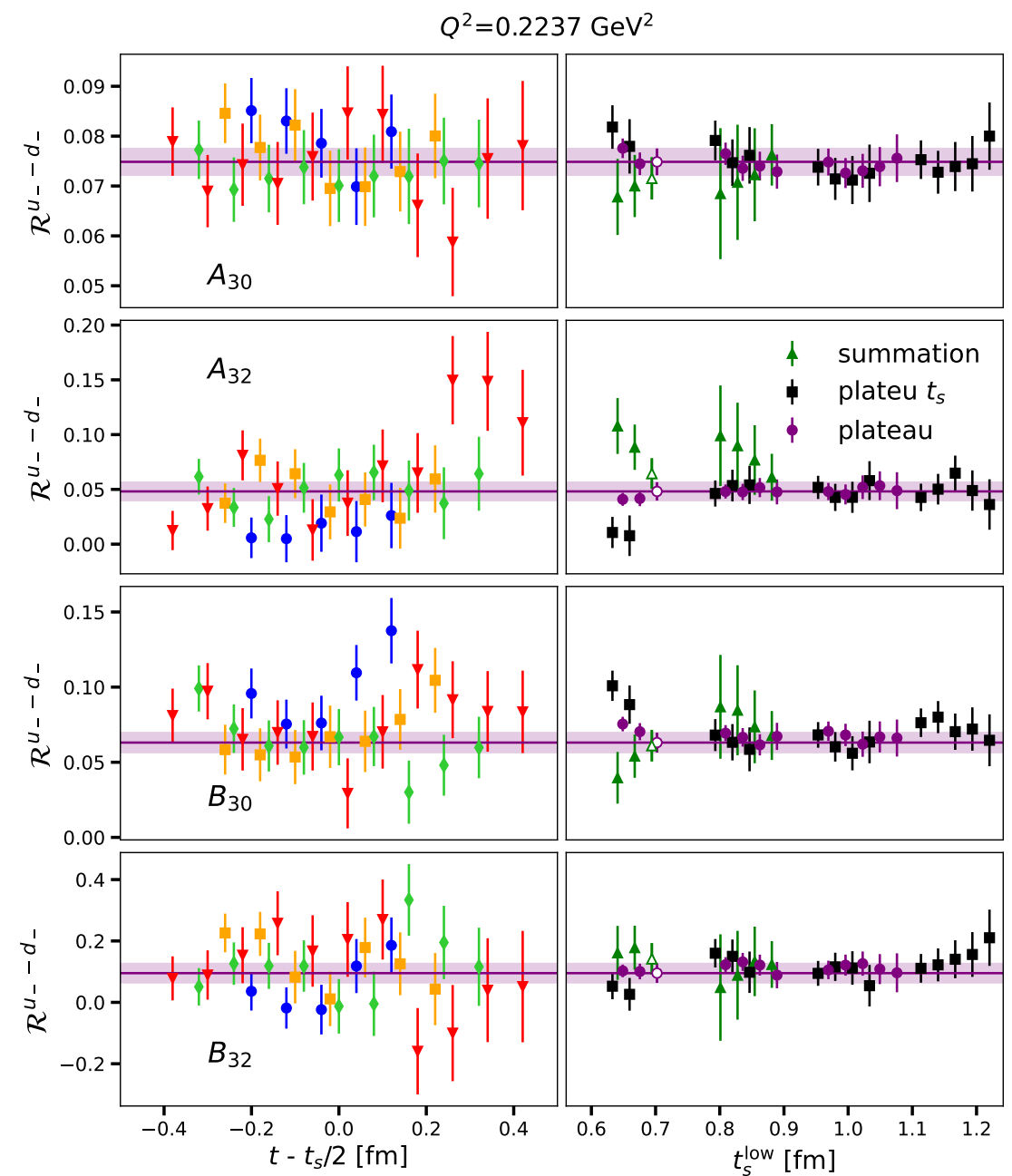
Results – Two-derivative operators

- Forward limit



Results – Two-derivative operators

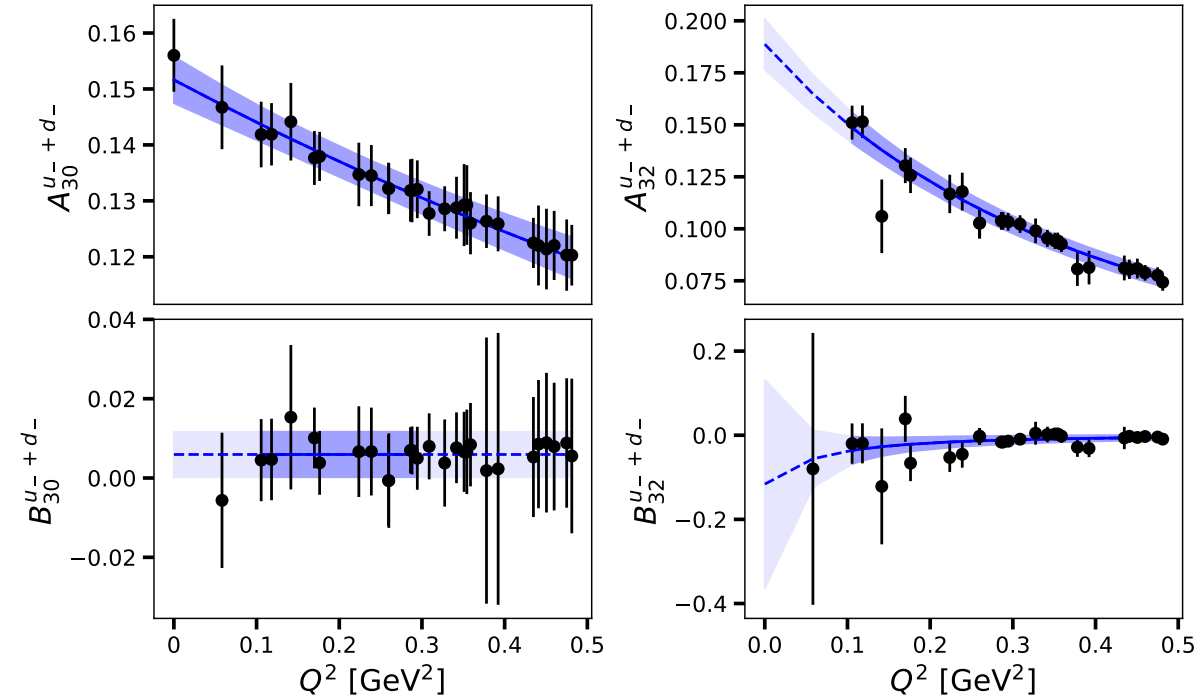
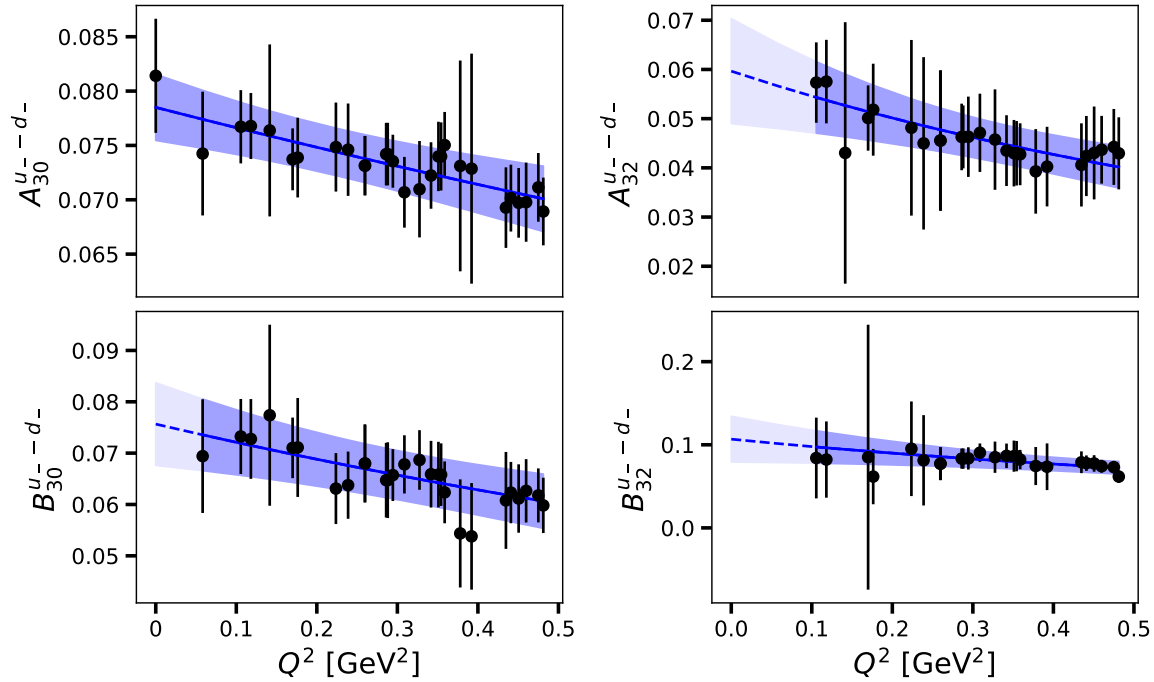
- Momentum transfer



Results – Two-derivative operators

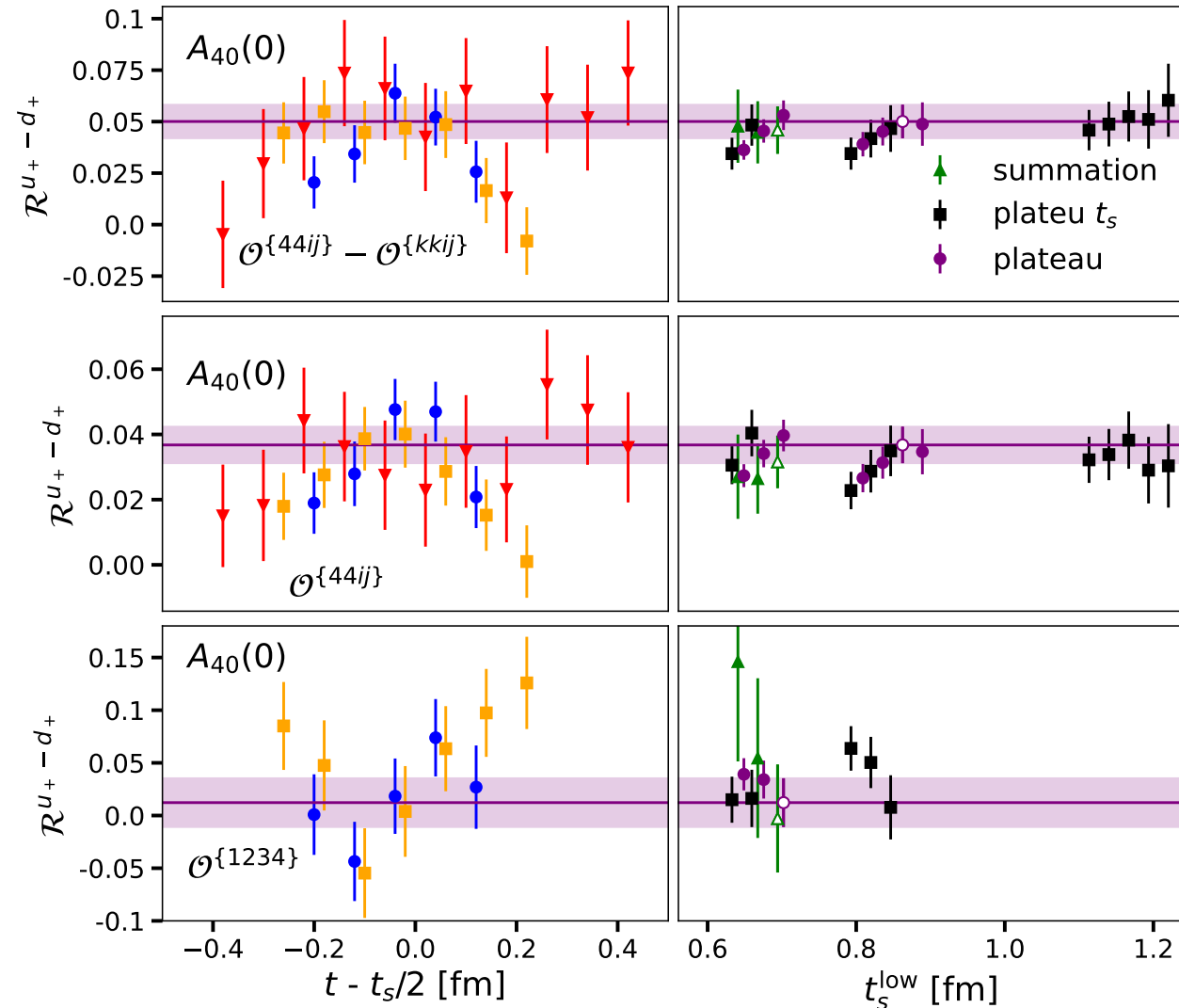
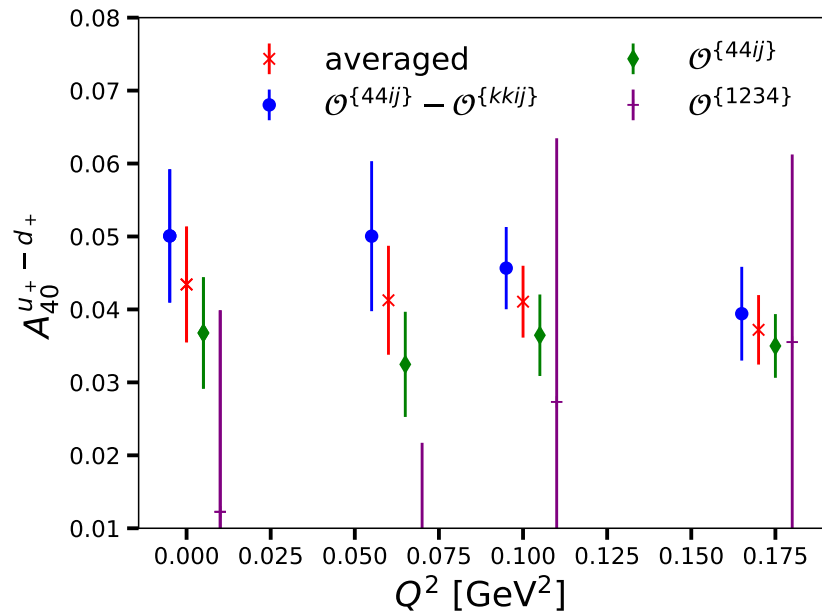
- Dipole fit

$$G(Q^2) = \frac{G(0)}{\left(1 + \frac{Q^2}{M^2}\right)^2}$$



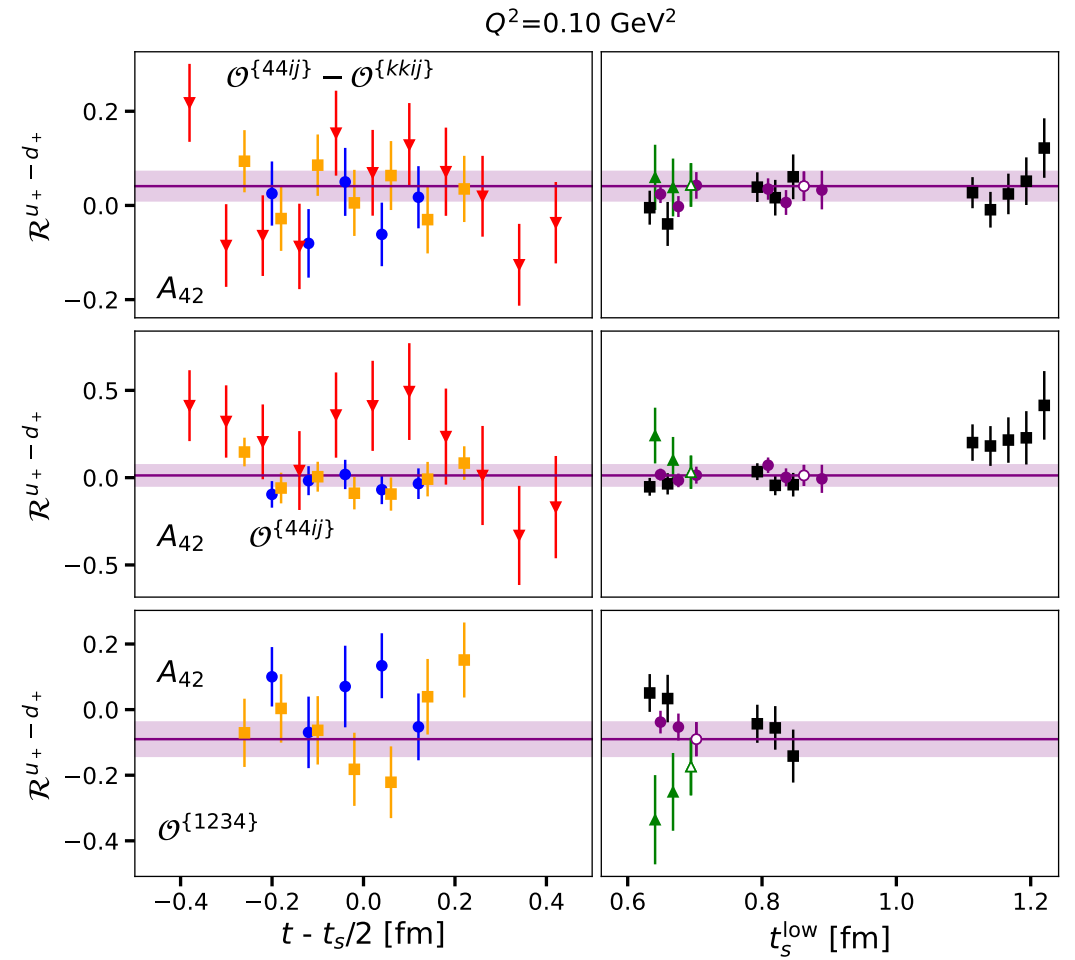
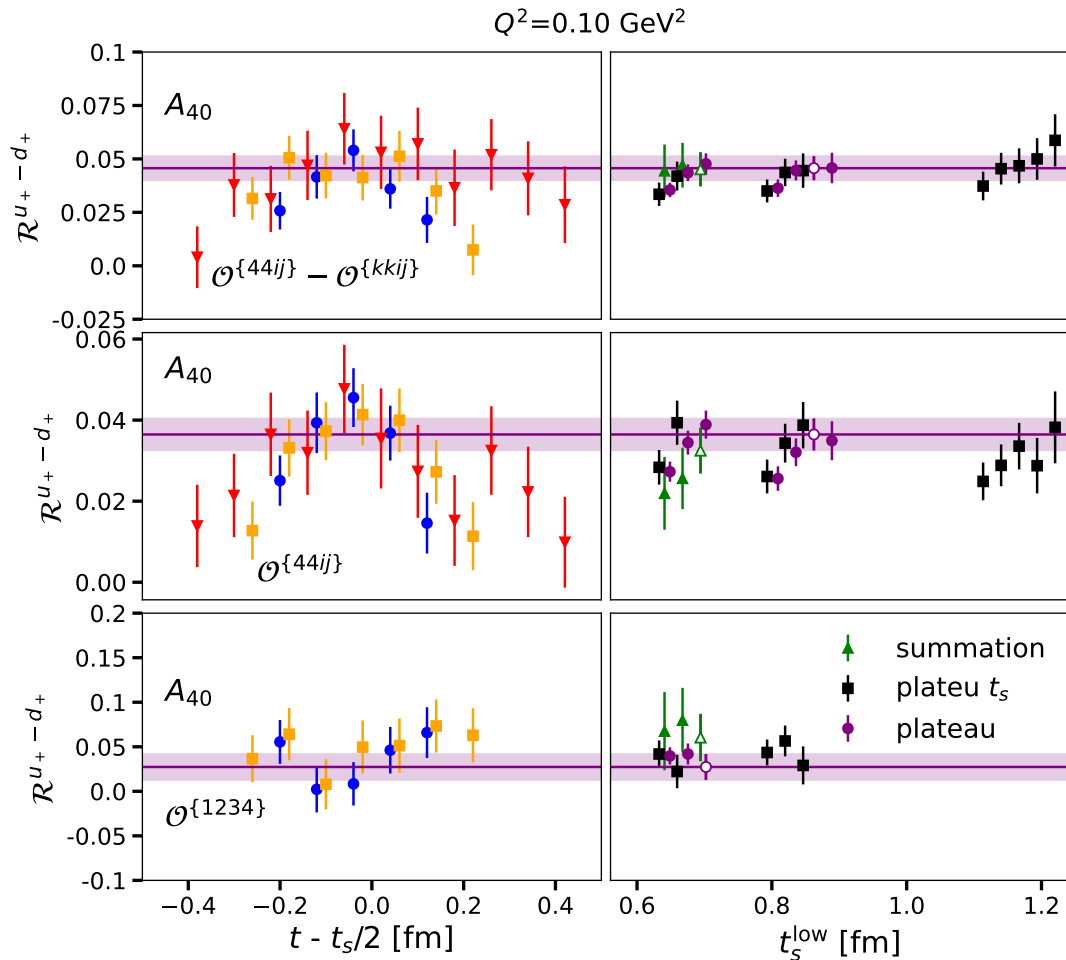
Results – Three-derivative operators

- Forward limit
- Comparison of different operators
- Note: y-axis scale



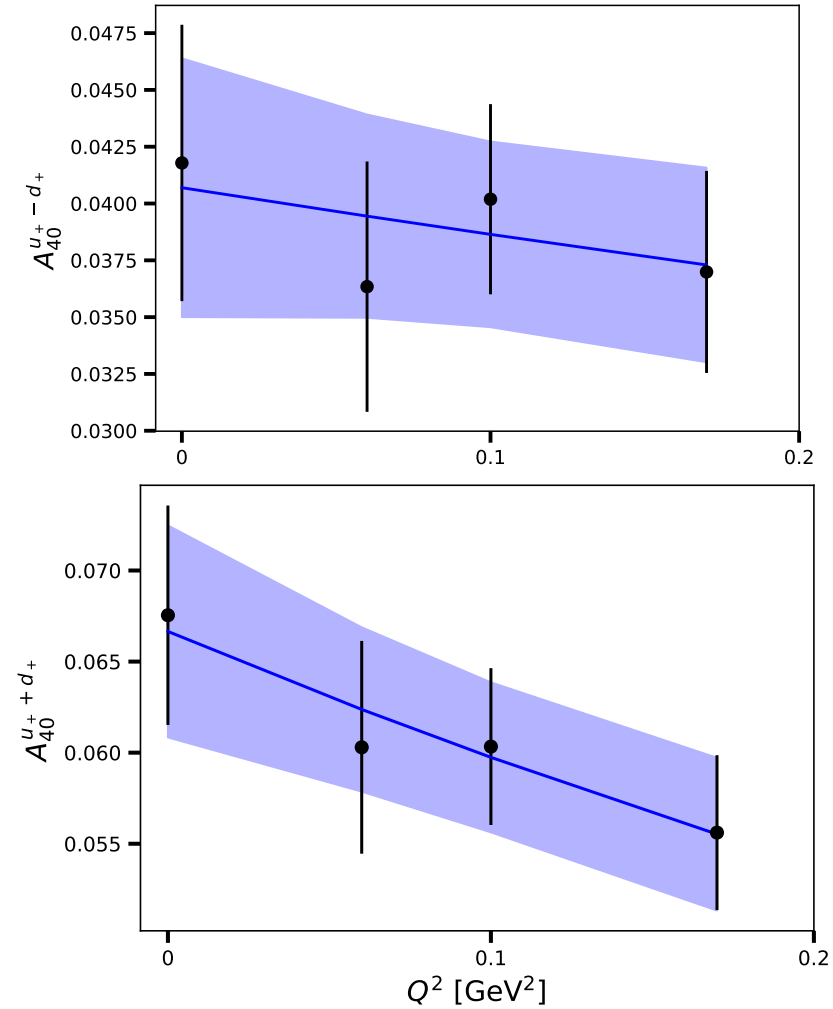
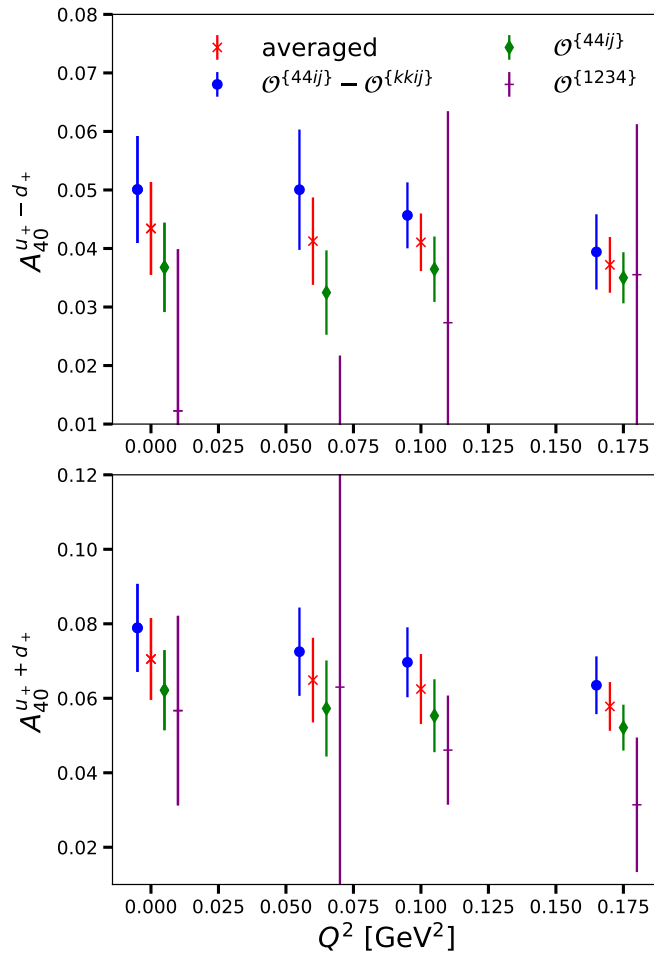
Results – Three-derivative operators

- Momentum transfer



Results – Three-derivative operators

- Only A_{40} nonzero



Comparison with literature

- JAM and NNPDF are phenomenological studies

JAM, PRD 106, L051502 (2022)

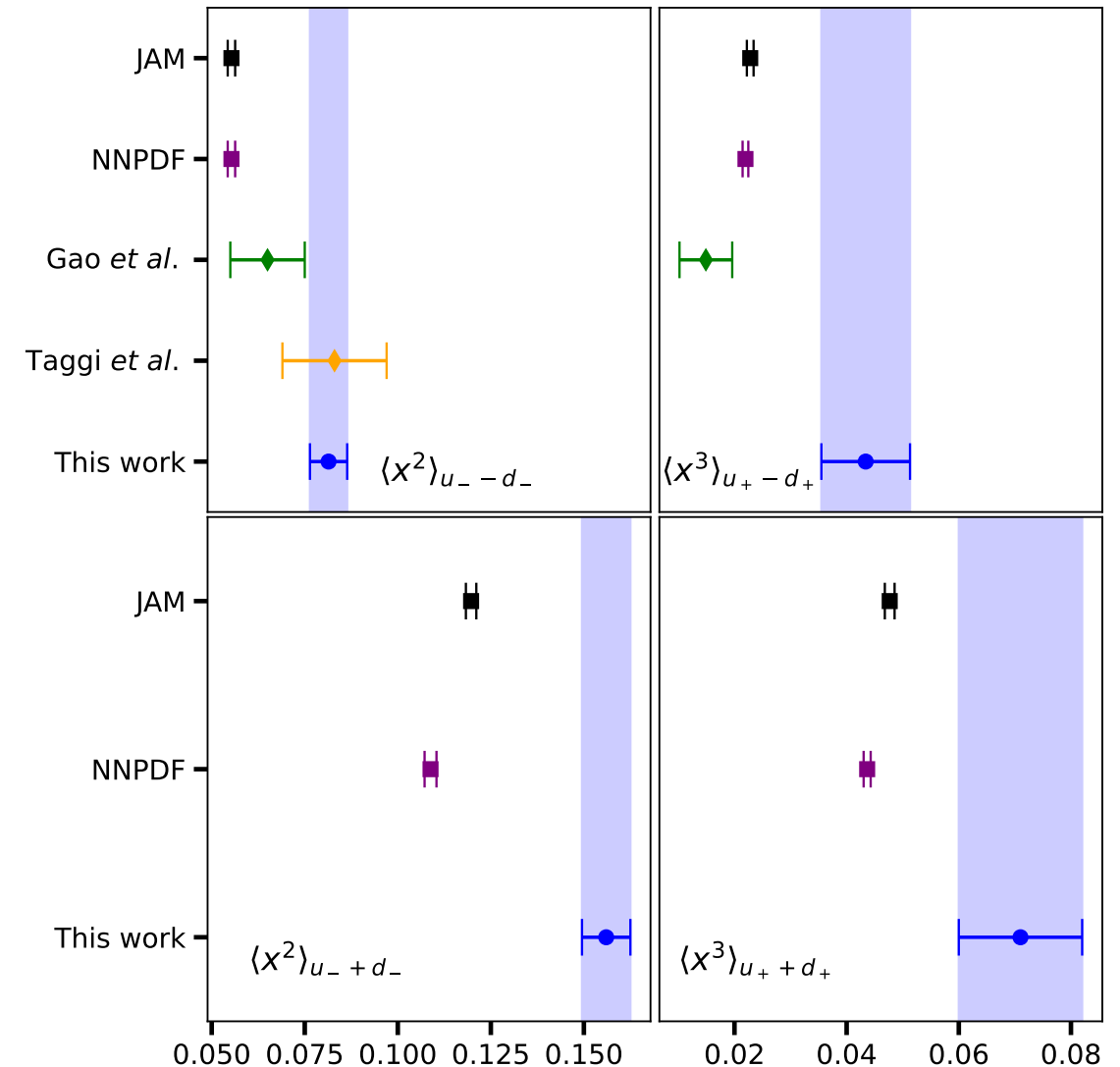
NNPDF, EPJC 82, 428 (2022)

- Gao uses non-local operators

X. Gao *et al.*, PRD 107, 074509 (2023)

- Taggi is the first direct computation third moment

E. Taggi *et al.*, (2026), 2605.02808



PDF reconstruction

- Ansatz

$$q(x) = Nx^{\alpha}(1-x)^{\beta}$$

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- Problem with antiquarks

$$\langle x^{n-1} \rangle = \int_0^1 dx x^{n-1} [q(x) - (-1)^{n-1} \bar{q}(x)]$$

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- Even powers give valence distribution

$$\langle 1 \rangle_{q-} = \langle 1 \rangle_{q-}$$

$$\langle x^2 \rangle_{q-} = \langle x^2 \rangle_{q-}$$

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- Even powers give valence distribution
- Odd powers do not

$$\langle 1 \rangle_{q-} = \langle 1 \rangle_{q-}$$

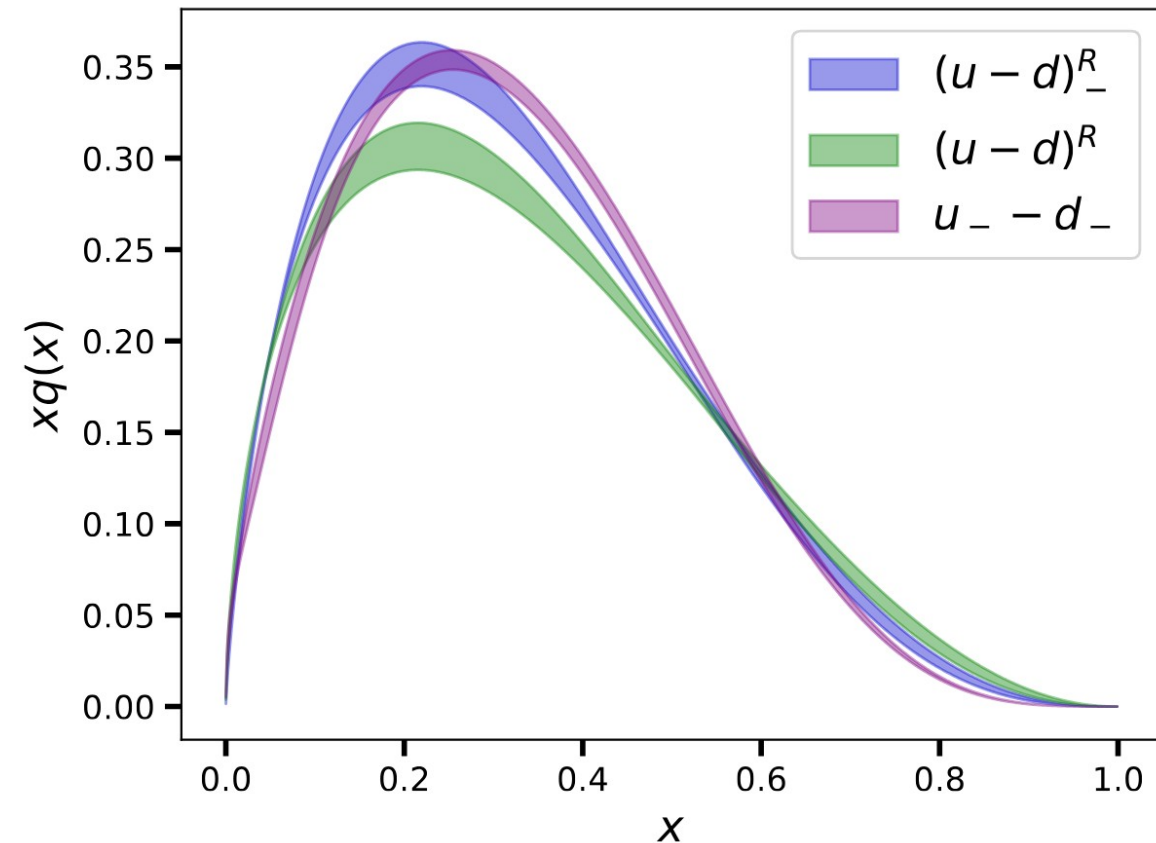
$$\langle x^2 \rangle_{q-} = \langle x^2 \rangle_{q-}$$

$$\langle x \rangle_{q-} = \langle x \rangle_{q+} - 2 \langle x \rangle_{\bar{q}}$$

$$\langle x^3 \rangle_{q-} = \langle x^3 \rangle_{q+} - 2 \langle x^3 \rangle_{\bar{q}}$$

PDF reconstruction

- Comparison between approaches with JAM data
- JAM data: $u_- - d_-$
- Calculate moments $\langle x \rangle_{q_-}$, $\langle x^2 \rangle_{q_-}$, $\langle x^3 \rangle_{q_-}$ and reconstruct PDF: $(u - d)_-^R$
- Calculate moments $\langle x \rangle_{q_+}$, $\langle x^2 \rangle_{q_-}$, $\langle x^3 \rangle_{q_+}$ and reconstruct PDF: $(u - d)^R$



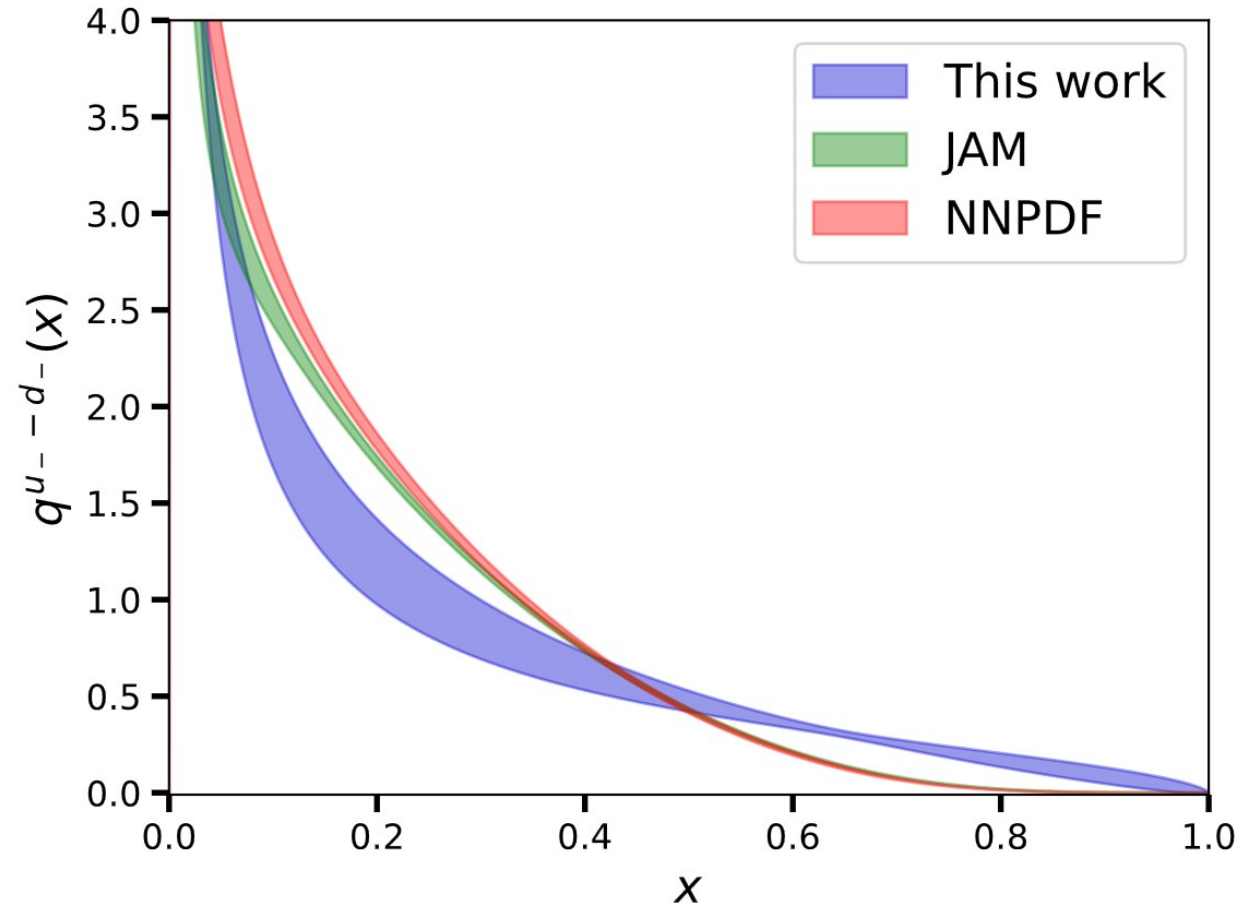
JAM, PRD 106, L051502 (2022)

PDF reconstruction

- Ansatz

$$q(x) = Nx^\alpha(1-x)^\beta$$

- Include JAM data for $\langle x^n \rangle_{\bar{q}}$ in reconstruction



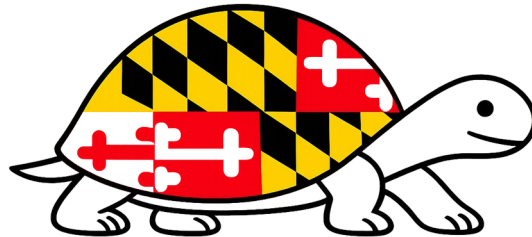
JAM, PRD 106, L051502 (2022)

NNPDF, EPJC 82, 428 (2022)

Conclusions

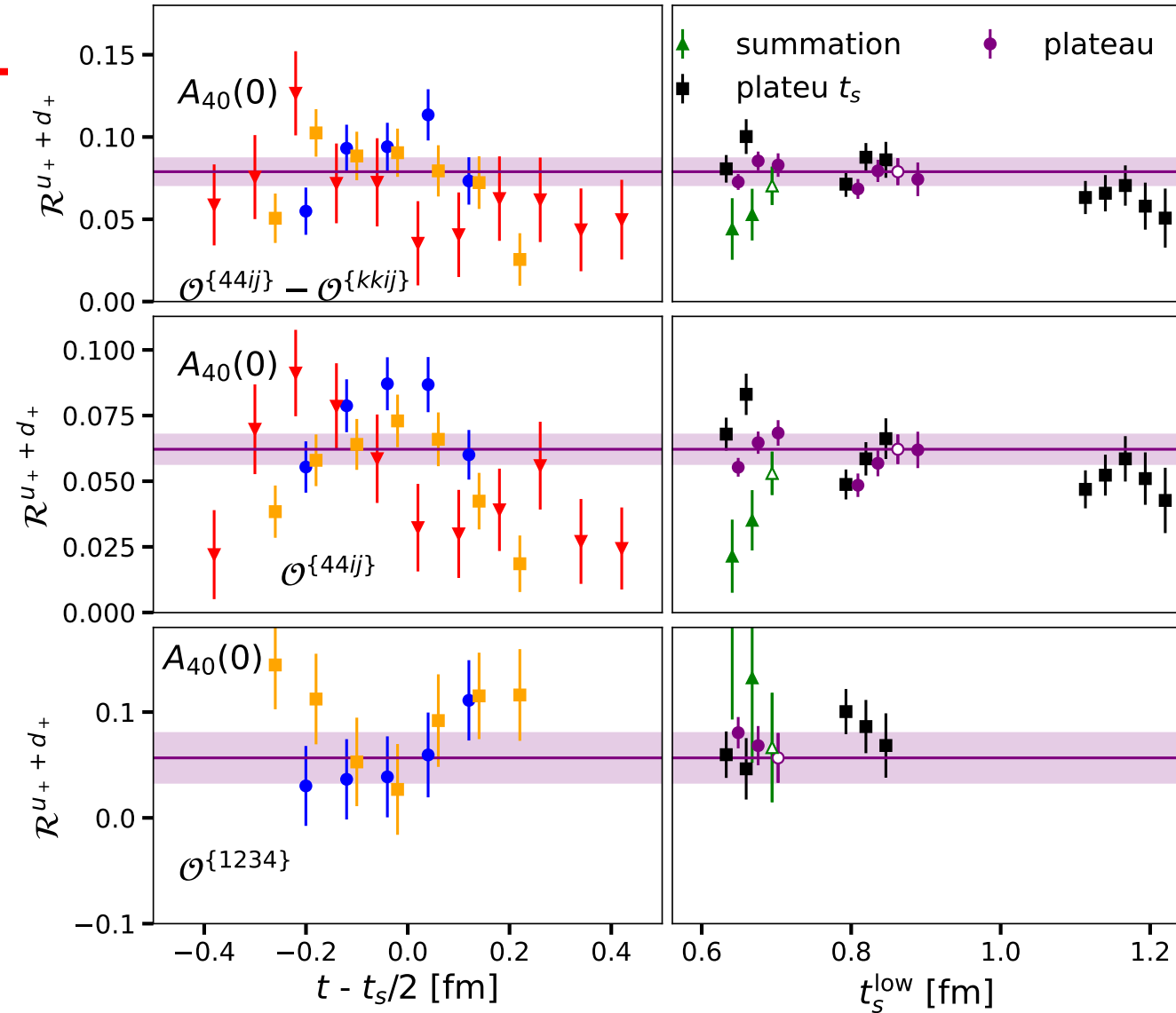
- Direct calculation of third nucleon moment
- Direct calculation of fourth moment
- Determination of GFFs
- Reconstruction of isovector PDF

Backup



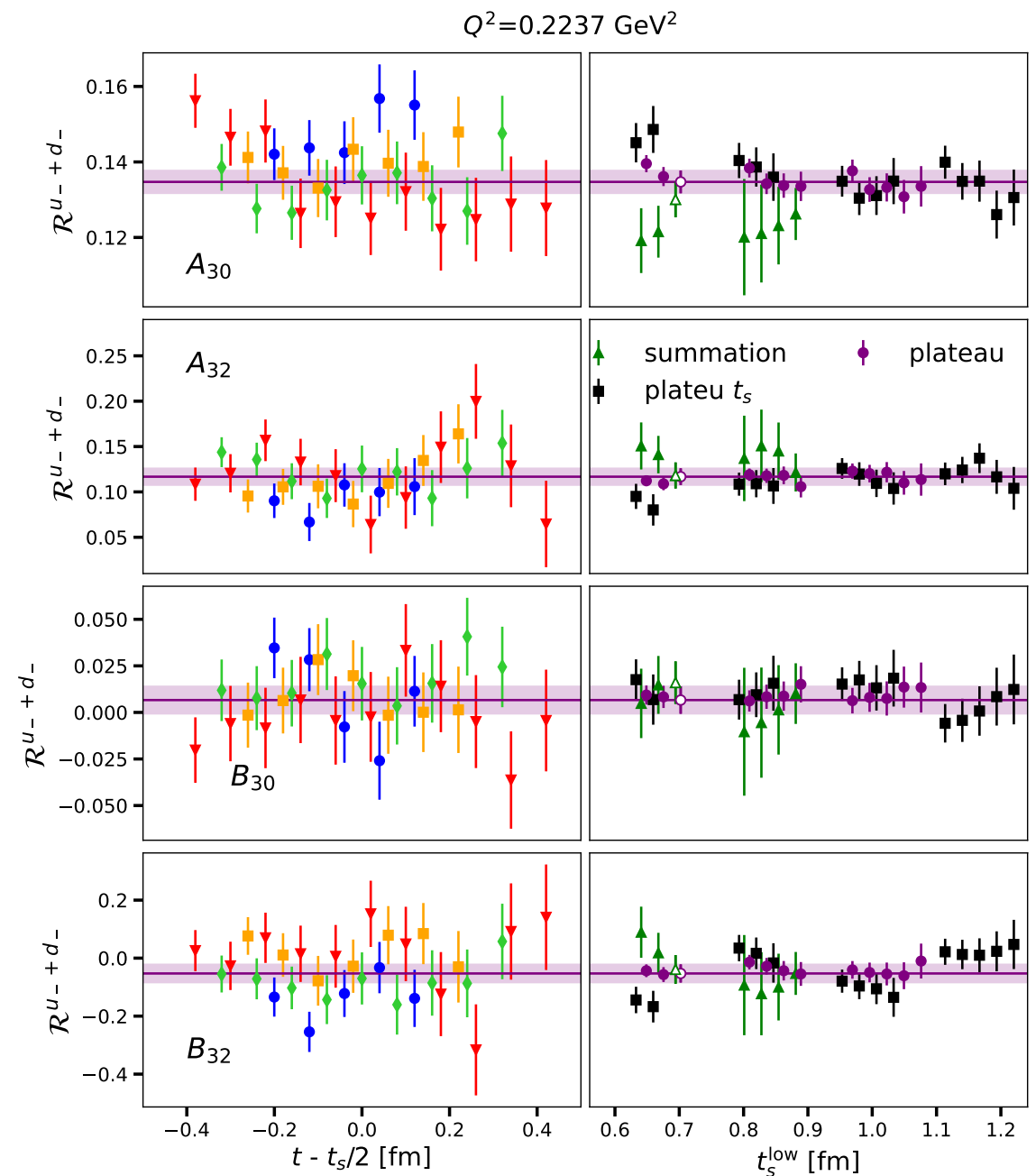
Results

- Forward limit
- Three derivative operators
- isoscalar



Results

- Momentum transfer
- Two-derivative operators
- isoscalar



Results

- Three-derivative operators, isoscalar

