

# Improving the Wilson action

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in collaboration with A. Risch, M. Roost, S. Schaefer and R. Sommer



**Deutsches Elektronen  
-Synchrotron Zeuthen**

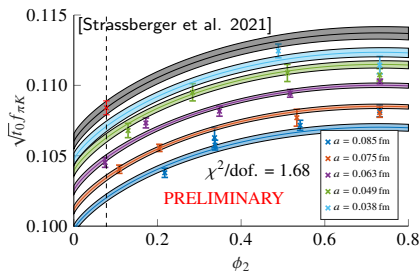
# Limitations of the Wilson action

**$\mathcal{O}(a)$ -improved Wilson action has been very successful** in many applications, but experience has shown us **it has some limitations**

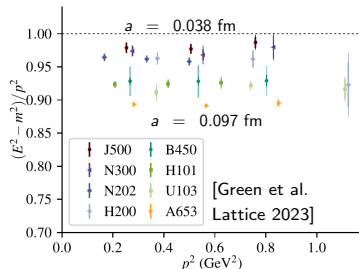
# Limitations of the Wilson action

$\mathcal{O}(a)$ -improved Wilson action has been very successful in many applications, but experience has shown us it has some limitations

➤ **Significant discretization effects** found for some observables



Meson decay constants

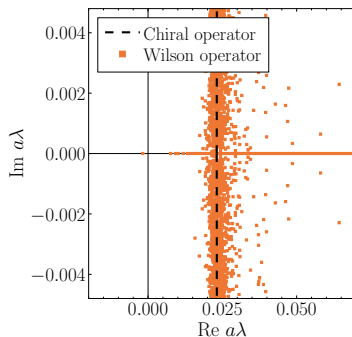


$D$ -meson dispersion relation

Main discretization error affecting two-particle spectroscopy [Hansen and Peterken, 2023]

# Limitations of the Wilson action

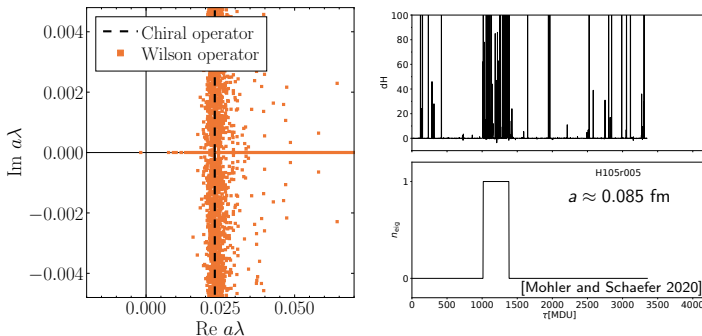
- **Small and negative real eigenvalues** in the spectrum of  $D$



- **Large condition numbers** and slow convergence

# Limitations of the Wilson action

- **Small and negative real eigenvalues** in the spectrum of  $D$



- **Large condition numbers** and slow convergence
- Potential barrier → **Ergodicity problems**  
(long autocorrelation times, reweighting factors...)
- Enhanced by **large values of improvement coeffs.**:  $c_{\text{SW}} \sim 2$  at  $a \approx 0.1$  fm

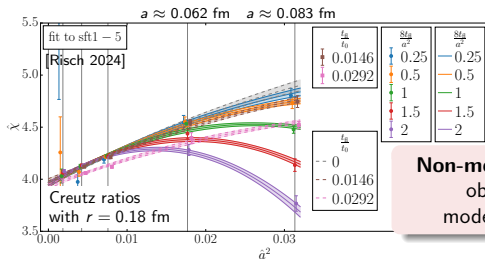
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We propose a **new action based on three modifications**

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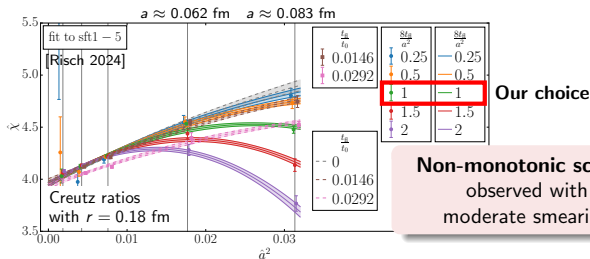
- **“Safe” smearing of links in  $D$**   
(single step of stout-smearing,  $\rho = 0.125$ )
- Reduces appearance of small real eigenvalues and magnitude of improvement coefficients
- Excessive smearing **may induce non-monotonic scaling**



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- **Tree-level improvement at  $\mathcal{O}(a^2)$**   
[Sheikholeslami and Wohlert 1985]

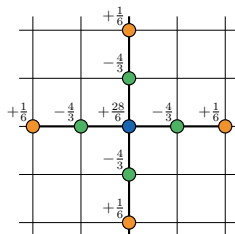
- ▶ Restores rotational invariance at  $\mathcal{O}(a^2)$  at tree level
- ▶ Reduces logarithmic lattice artifacts

$$\mathcal{O}(0) = \mathcal{O}(a) + \sum_i \hat{c}_i M_i a^2 [\alpha_s(a^{-1})]^{\Gamma_i}$$

with  $\Gamma_i \rightarrow \Gamma_i + 1$  (no dim-6 operator in SymEFT at tree level) [Husung 2024]

$D_{234}$  operator

[Alford et al. 1997]



Includes dim-6 and -7 operators, and  $r = 2/3$

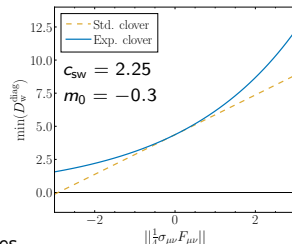


**Efficient implementation**

# Proposal for a new action

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- **Tree-level improvement at  $\mathcal{O}(a^2)$**   
[Sheikholeslami and Wohlert 1985]
- **Exponentiated clover term**  
[Francis et al. 2020]
- Expected to reduce small / negative eigenvalues,  
especially for large  $c_{\text{sw}}$



$$D_{234}^{\text{diag}} = \frac{28}{6} + m_0 + \frac{ic_{\text{sw}}}{6} \sigma_{\mu\nu} F_{\mu\nu} \longrightarrow \left( \frac{28}{6} + m_0 \right) \exp \left[ \frac{1}{\frac{28}{6} + m_0} \frac{ic_{\text{sw}}}{6} \sigma_{\mu\nu} F_{\mu\nu} \right]$$

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➔ **Discretization effects**

➔ **Spectrum of  $\gamma_5 D$**

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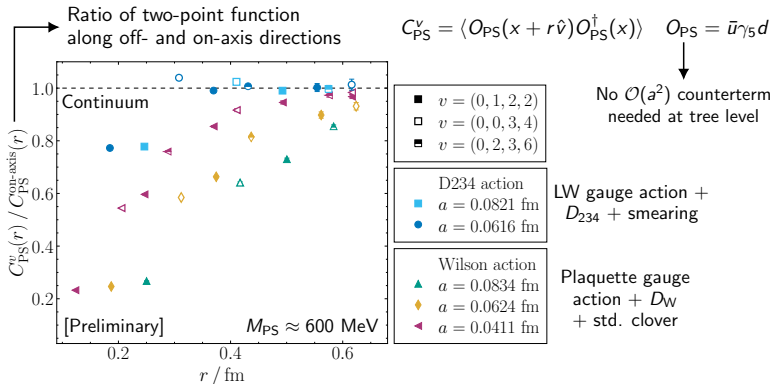
**Discretization effects**

**Spectrum of  $\gamma_5 D$**

- **Non-perturbative  $c_{\text{sw}}$**  to remove  $\mathcal{O}(a)$  effects in all cases
- Unless otherwise stated, exp. clover + Lüscher-Weisz

# Restoration of rotational invariance

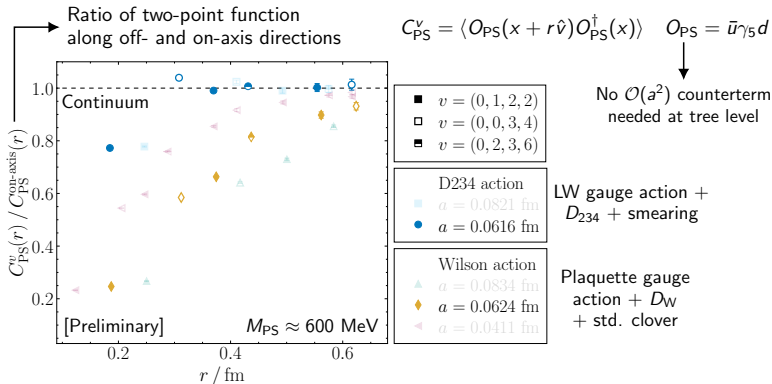
## First: Restoration of rotational symmetry at short distances



- Smaller discretization effects at short distances
- Rotational invariance restored after few lattice spacings

# Restoration of rotational invariance

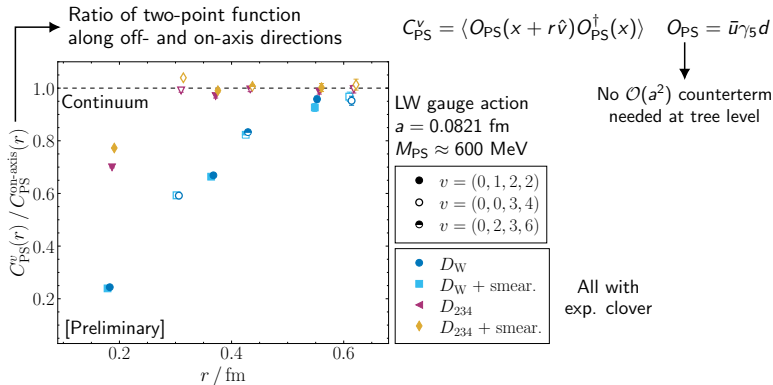
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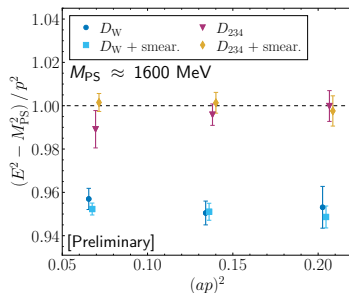
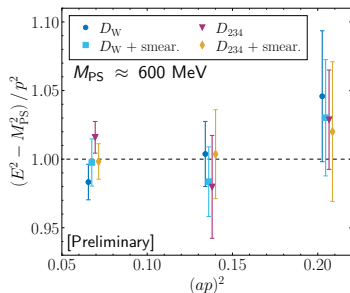
## First: Restoration of rotational symmetry at short distances



- Smaller discretization effects at short distances
- Rotational invariance restored after few lattice spacings
- $\mathcal{O}(a^2)$  improvement plays the crucial role

# Meson dispersion relation

**Second:** Dispersion relations of pseudoscalar mesons ( $a = 0.0821$  fm)

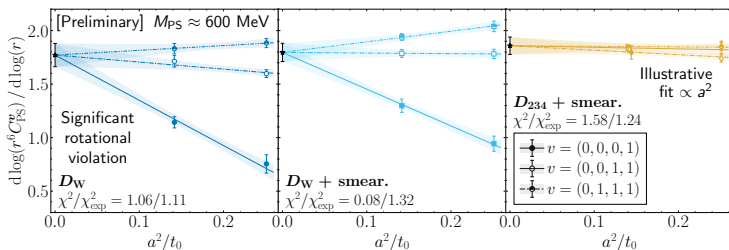


- Correct dispersion relations for *light* mesons
- **$D_{234}$  operator recovers continuum result for *heavier* mesons, even at  $M_{PS} \approx 1600$  MeV**

# Continuum limit of log-derivative

**Third: Continuum limit of log-derivative** at intermediate distances

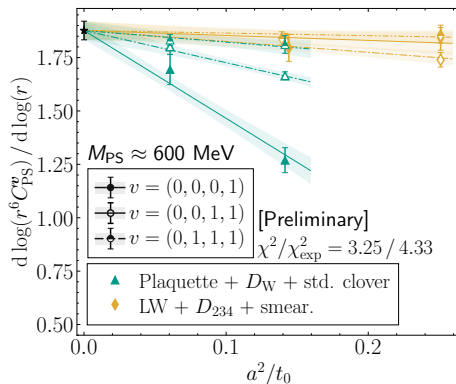
$$L_{PS}^V(r) = \left. \frac{d \log[r^6 C_{PS}^V(r)]}{d \log r} \right|_{r=r^*} \quad r^* \approx 0.5 \text{ fm}$$



- Combine multiple directions for a constrained extrapolation
- Again, **significant improvement from  $D_{234}$  operator**
- Moderate error reduction from smearing

# Continuum limit of log-derivative

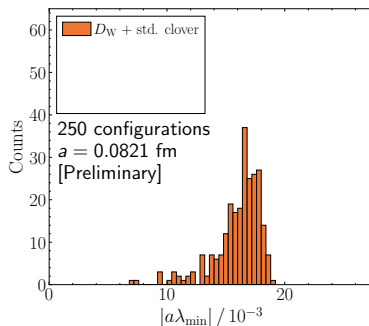
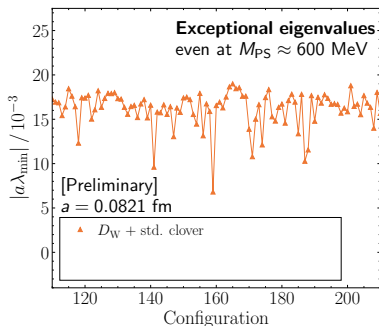
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Results for **old and new actions can be combined** for constrained continuum extrapolations

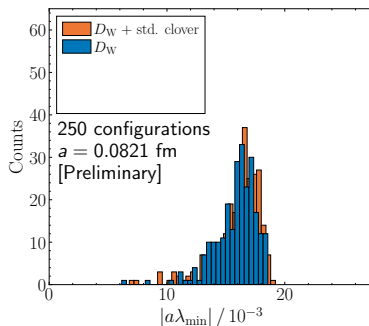
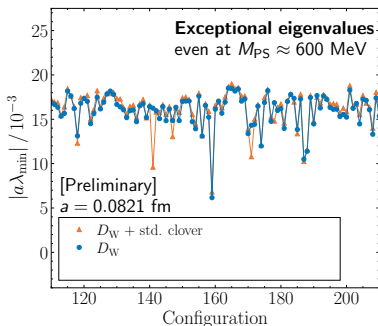
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We investigate lowest eigenvalue of  $\gamma_5 D$  at  $M_{PS} \approx 600$  MeV



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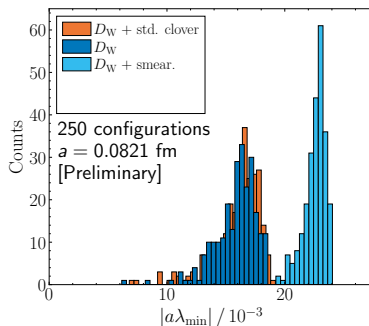
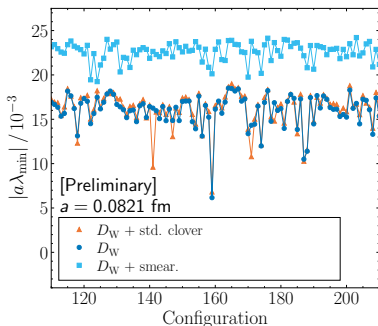
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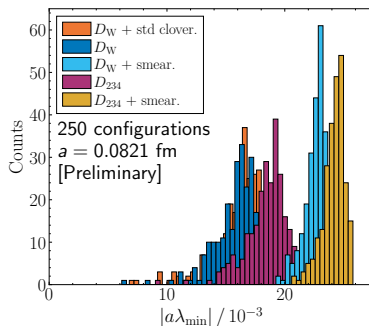
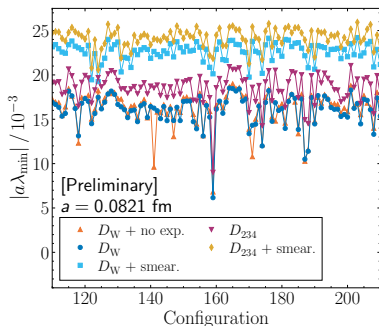
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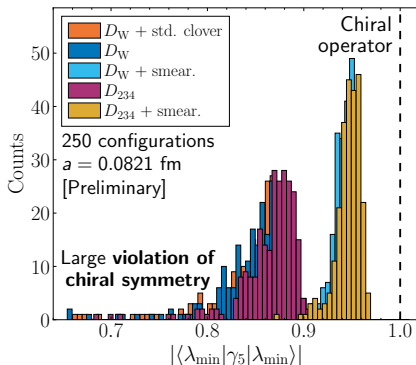


- Exp. clover cures some exceptional eigenvalues, but not all
- **Significant improvement from smearing**
- Some **minor improvement from  $D_{234}$  operator**, possibly due to  $r = 2/3$

# Reduction of chiral symmetry violation

Lowest eigenvalue leads to largest **chiral-symmetry violation**

- For a **chiral operator**,  $\langle \lambda_{\min} | \gamma_5 | \lambda_{\min} \rangle = \pm 1$



- **Smearing reduces violation of chiral symmetry**
- Significant reduction of eigenvalues that violate chiral symmetry the most

# Summary and outlook

We propose a **new Wilson-like action** with “safe” smearing,  $\mathcal{O}(a^2)$  improvement at tree level and exp. clover

- We study the impact of all three features in the **quenched approximation**
- $\mathcal{O}(a^2)$  improvement at tree level **significantly reduces violations of rotational invariance and other discretization effects**
- Smearing has a **major impact on reducing exceptional eigenvalues and chiral-symmetry violation**

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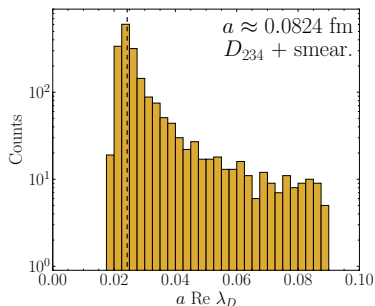
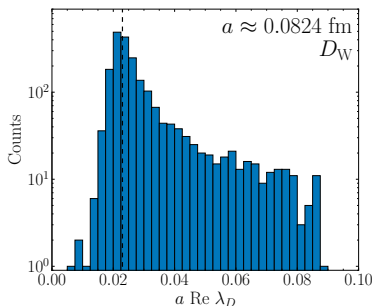
**Paper in preparation:** improvement coefficients, eigenvalues of  $D...$

**Next steps:** Dynamical simulations and improvement

# Thank you for your attention!

# Eigenvalues of $D$

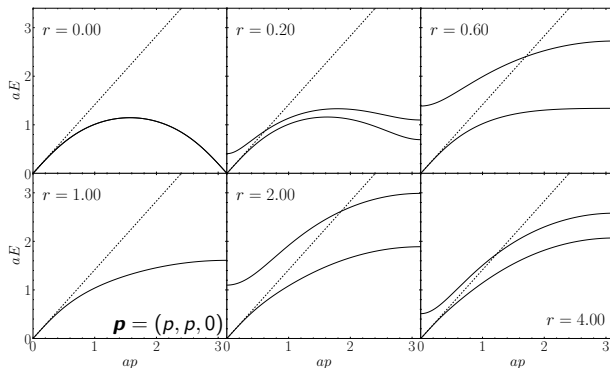
We have studied the **eigenvalues of the non-Hermitian Dirac operator** ( $M_{PS} \approx 600$  MeV)



**Significant reduction of small eigenvalues**

# The role of $r$ in the Wilson action

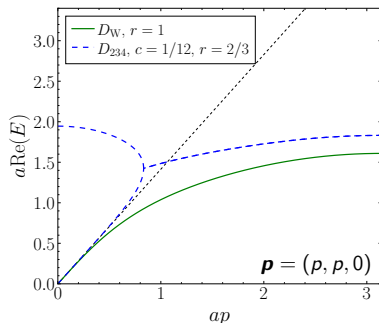
The choice  $r = 1$  **eliminates ghost branches in the free dispersion relation** and allows for the **projection trick**



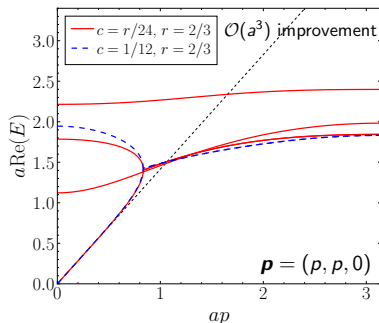
$$D_W^{\text{free}} \psi(x) = -P_\mu^- \psi(x + a\hat{\mu}) + \psi(x) - P_\mu^+ \psi(x - a\hat{\mu}) + m_0 \psi(x)$$

# The $D_{234}$ Dirac operator

Reduce the number of ghosts in the **free-quark dispersion relation**



Small discretization effects  
for  $aE \lesssim 1$



Minimum number of ghosts,  
with  $aE_{\text{ghost}} \gg 1$

# The $D_{234}$ Dirac operator

Rewrite the operator using the **projection trick**

$$\begin{aligned} D_{234}\psi(x) = & \frac{1}{6} U_\mu(x) U_\mu(x + a\hat{\mu}) \mathbf{P}_\mu^- \psi(x + 2a\hat{\mu}) \\ & - \frac{4}{3} U_\mu(x) \mathbf{P}_\mu^- \psi(x + a\hat{\mu}) \\ & + \frac{1}{6} U_\mu^\dagger(x - a\hat{\mu}) U_\mu^\dagger(x - 2a\hat{\mu}) \mathbf{P}_\mu^+ \psi(x - 2a\hat{\mu}) \\ & - \frac{4}{3} U_\mu^\dagger(x - a\hat{\mu}) \mathbf{P}_\mu^+ \psi(x - a\hat{\mu}) \\ & + \frac{7}{6} \psi(x) + m_0 \psi(x) + \tilde{D}_{\text{clover}}^{\text{exp}} \psi(x). \end{aligned}$$

Only two spin components from neighbors needed

- Simple algorithmic implementation

$$\text{Cost}(D_{234}) \approx 2 \times \text{Cost}(D_W)$$

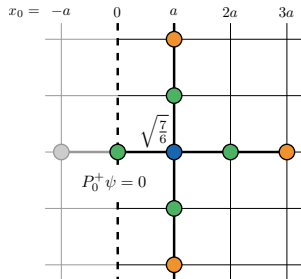
Expected **smaller factor for deflated solver**

- Convenient structure for SF boundaries (**ongoing work**)

# $D_{234}$ operator in the SF

**Complication:** cannot define two-hop derivative in the first time slice next to the boundaries

- ▶ Naive definition leads to  $\mathcal{O}(a)$  boundary errors



PT at tree level and  
ideas from [Lüscher 2006]

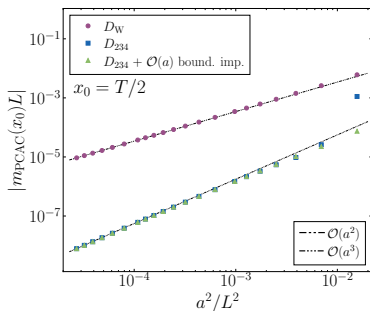
- ▶ Need to include **boundary counterterm**:

**Counterterm**  $\left( \tilde{c}_t^{234} - \frac{7}{6} \right) \bar{\psi}(a) P_0^+ D_0^- \psi(a) \rightarrow \tilde{c}_t^{234} = \sqrt{\frac{7}{6}}$

**Boundary field**  $\zeta = \sqrt{\frac{7}{6}} U_0^\dagger(0) P_0^- \psi(a) - \frac{1}{6} U_0^\dagger(0) U_0^\dagger(a) P_0^- \psi(2a)$

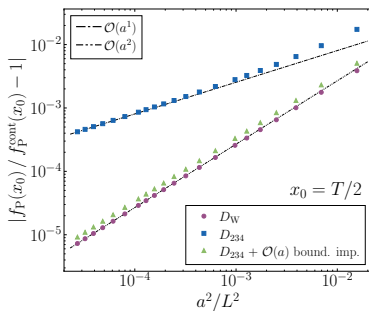
# SF in perturbation theory

We have studied the scaling of observables at **TL in perturbation theory**  
( $T = L$ ,  $am_0 = 0$ , background field from [Lüscher et al. 1993])



PCAC relation is protected  
from boundary effects

$\mathcal{O}(a^2)$  impr. in the **bulk**



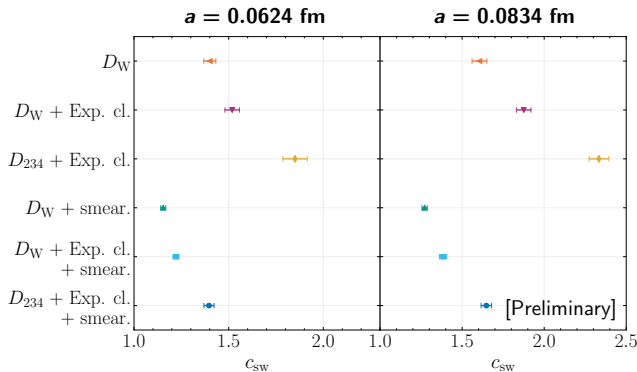
Boundary improvement is  
**required at TL**

$\mathcal{O}(a)$  impr. at the **boundary**

# Non-perturbative tuning of $c_{\text{SW}}$

Study the impact of modifications in the **non-perturbative** value of  $c_{\text{SW}}$

- Understood how to **define  $D_{234}$  operator in the Schrödinger functional**

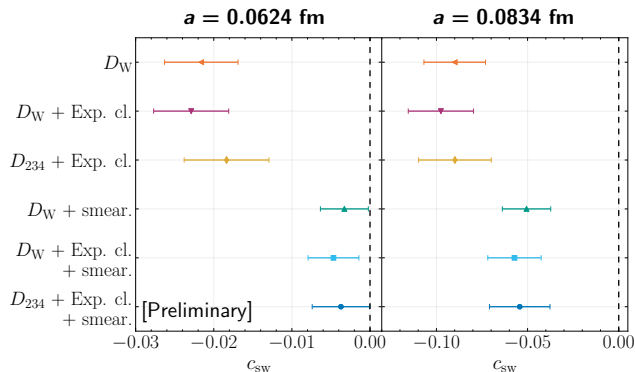


Effect of exp. clover and  $D_{234}$  operators **may change in full QCD**

- In full QCD,  $c_{\text{SW}}$  is smaller than with exp. clover [Francis et al. 2020]

# Non-perturbative tuning of $c_A$

Study the impact of modifications in the **non-perturbative** value of  $c_A$



Smearing significantly **reduces** the non-perturbative value of  $c_A$

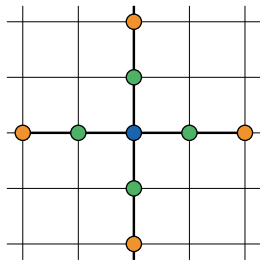
# Deflated solver for the $D_{234}$ operator

We have implemented a **deflated solver for the  $D_{234}$  operator**

$D_{234}$  only contains straight links



**Structure of little system remains the same**



SAP could be implemented, with a measured **cost**  $\sim 1.6\times$  that of  $D_W$

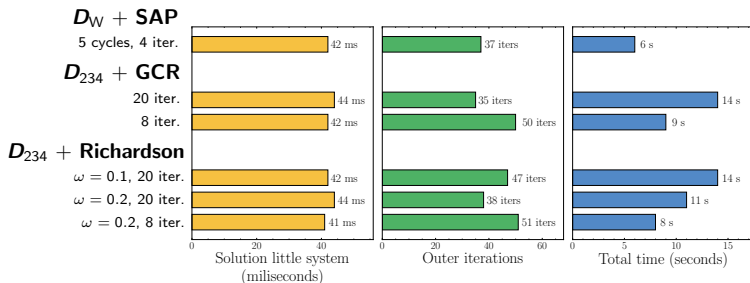
- ▶ We have tried **alternatives** suggested by A. Frommer
  - ▶ Fixed number of **GCR iterations**
  - ▶ **Richardson iteration** (no scalar products)

$$D\psi = \eta \quad \longrightarrow \quad \psi_{k+1} = \psi_k + \omega(\eta - D\psi_k)$$

# Tests of the deflated solver

We have **tested the deflated solver on the CLS N451r000 ensemble**

- ▶  $T \times L^3 = 128 \times 48^3$ ,  $a \approx 0.076$  fm,  $M_\pi \approx 286$  MeV
- ▶ Tuned quark mass for  $D_{234}$  operators with  $c_{\text{SW}} = 1.3$
- ▶ Results for 48 nodes of JUWELS, no SSE/AVX acceleration for  $D_W$



Solving for the  $D_{234}$  operator **only  $\sim 1.5 \times$  more expensive**

- ▶ Need further exploration (optimal parameters...)