

Nonperturbative Hadronic Structure of the Photon

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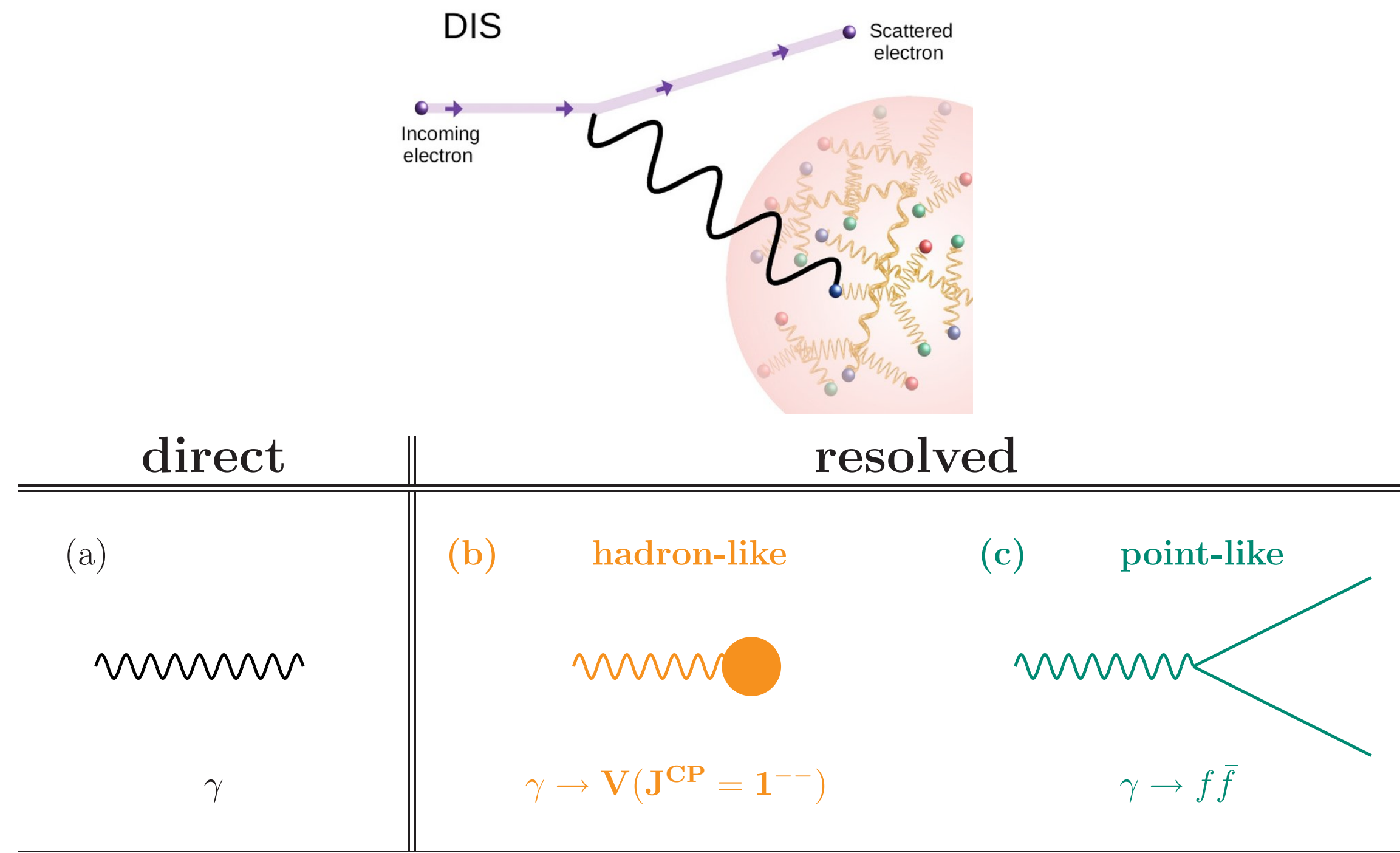


Funding acknowledgment:



Hadronic structure of the “resolved” photon from LQCD

- QCD allows photon to exhibit a hadronic structure due to quantum fluctuations



- $d\sigma = d\sigma_{\text{direct}} + d\sigma_{\text{resolved}}$
- Well controlled parametrization when photon appears as an active, partonic constituent of the proton.
 - Not so for “resolved” photon’s parton contents

QCD factorization

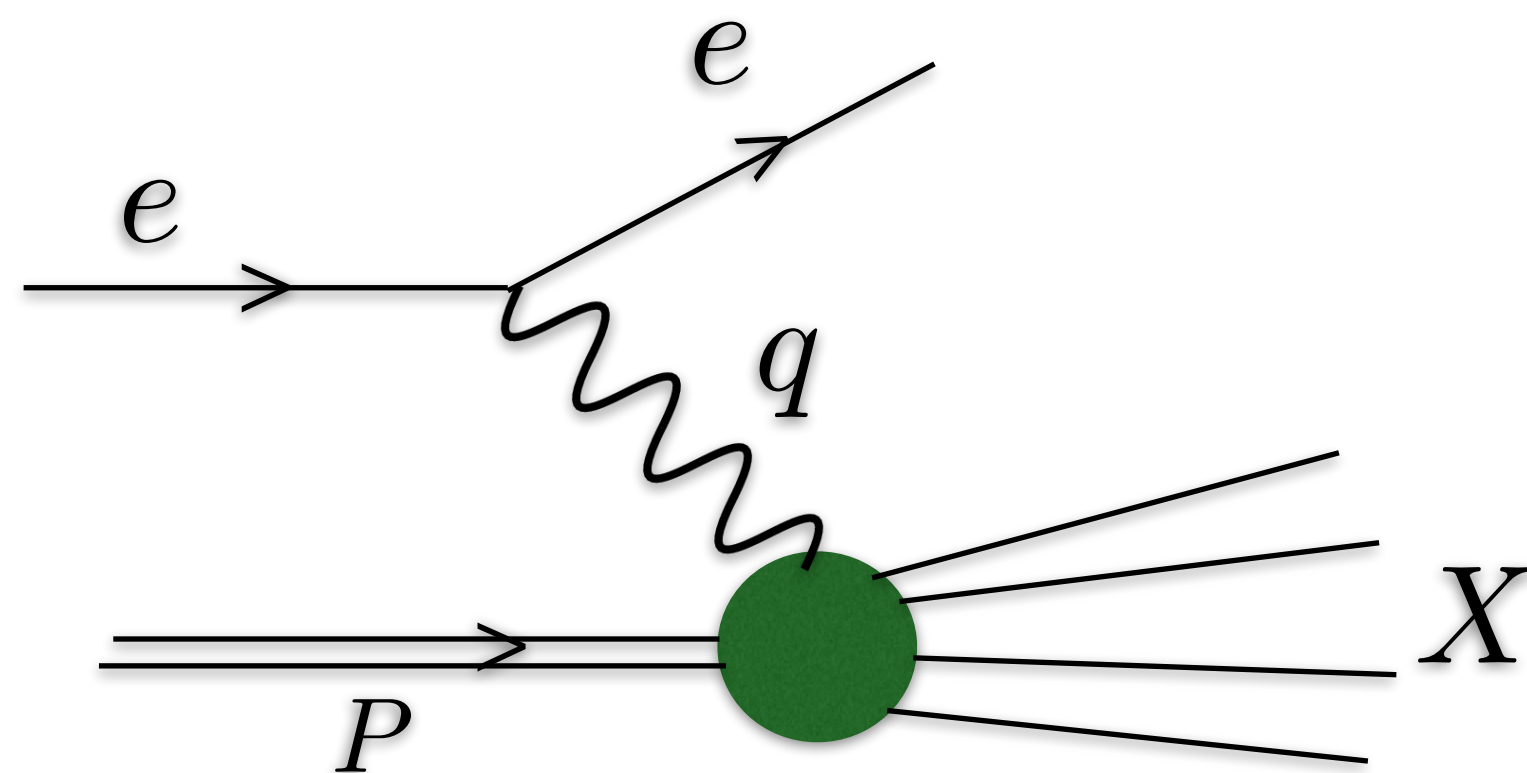
- QCD factorization is the cornerstone that connects hadronic scattering measurements into quantitative information on the quark and gluon structure of matter.

$$\sigma^{DIS}(x, Q^2, \sqrt{s}) = \sum_{\alpha=q, \bar{q}, g} K_{\alpha}\left(x, \frac{Q^2}{\mu^2}, \sqrt{s}\right) \otimes f_{\alpha}(x, \mu^2) + \text{Power Corrections}$$

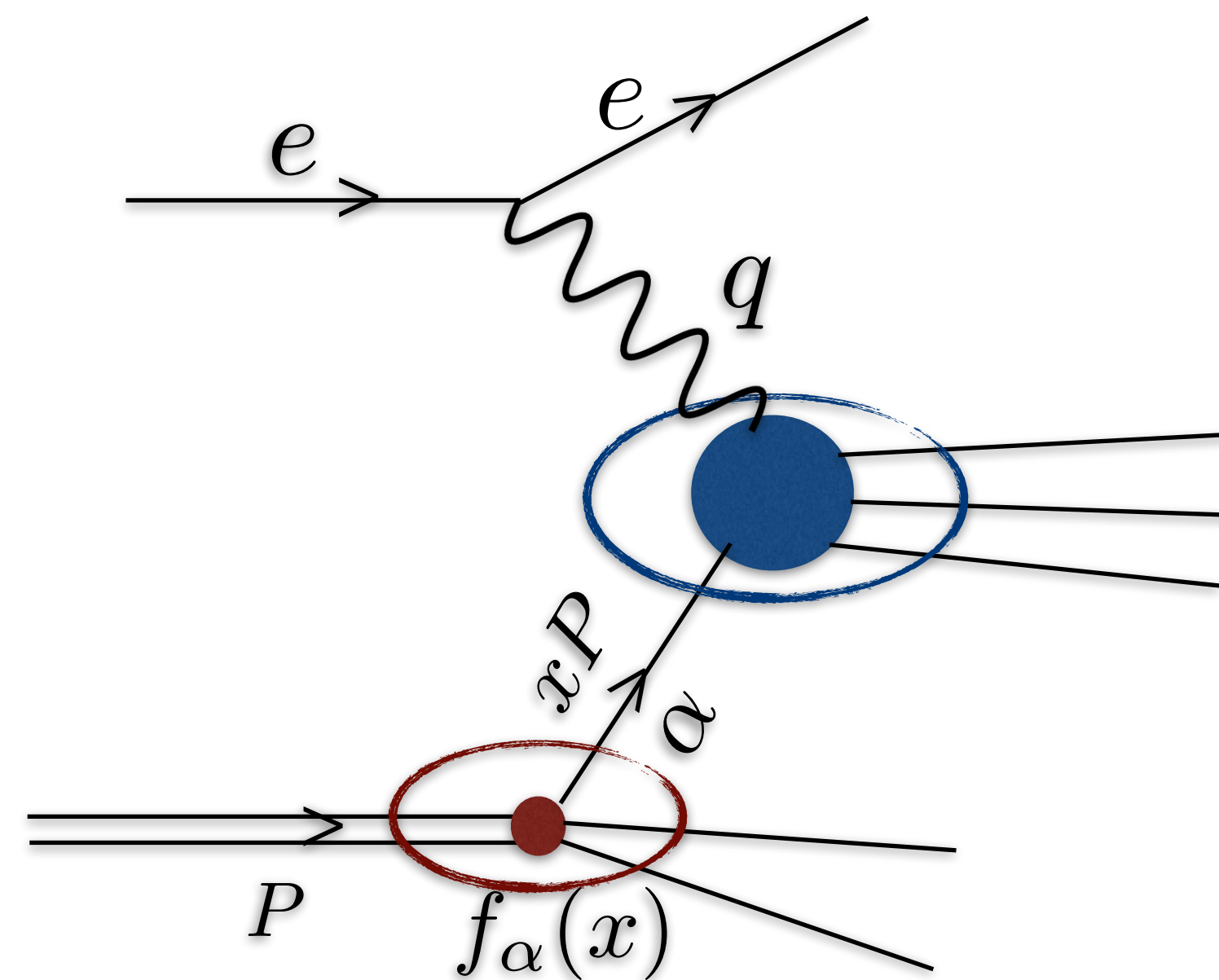
**Experimental
cross sections**
(Lattice QCD
matrix elements)

**Short-distance
physics**
(pen and paper)

Nonperturbative
long-distance
physics



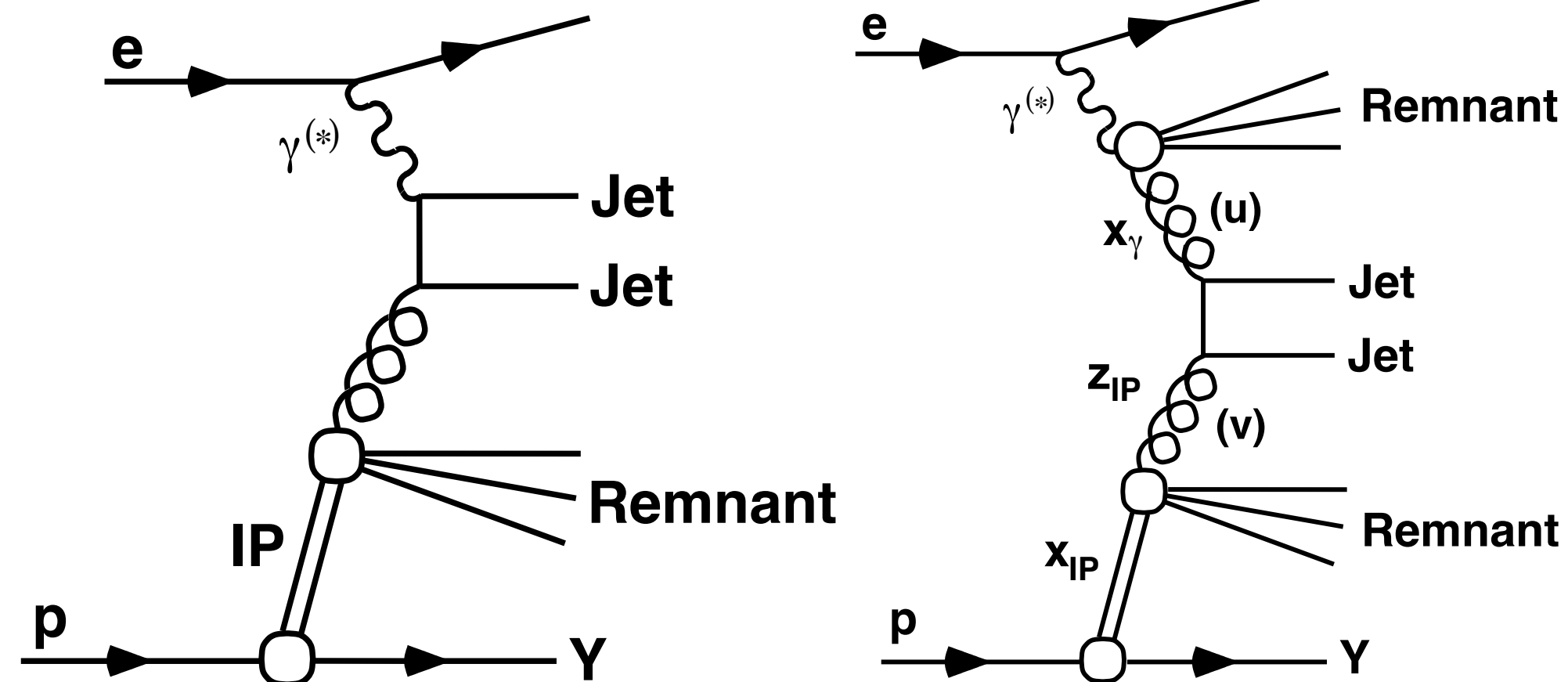
Scattering Process



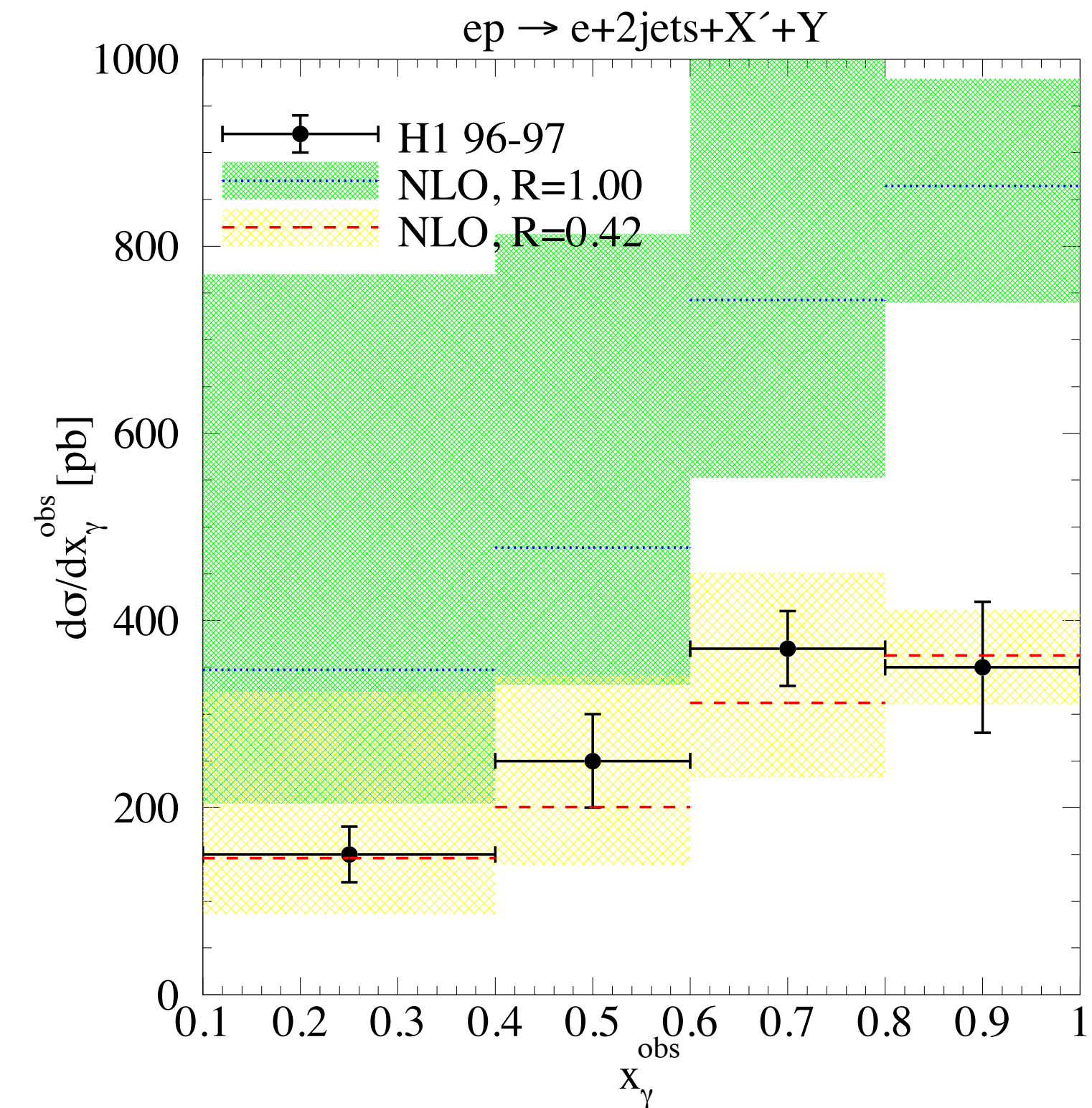
Parton Picture

Hadronic structure of the “resolved” photon from LQCD

- Test of QCD factorization “breaking” (or lack of understanding of resolved photon structure)



Differential cross sections for dijet photo production measured by H1 and ZEUS [Phys. Lett. 2008]



- Critically important for EIC and constraining gluon PDF [Chu, Aschenauer, Lee 2017]
- Experiments are only able to probe the structure of quasi-real or virtual photons, extrapolating to real-photon structure involves significant complications and systematic uncertainties

Hadronic structure of the “resolved” photon from LQCD

- LQCD calculation of moment

X. Ji, C. Jung [PRL 2001, PRD 2001]

- Using LSZ reduction formula

$$\blacktriangleright \langle \gamma(k\lambda') | \hat{O}(0) | \gamma(k\lambda) \rangle = - \lim_{k'^\mu \rightarrow k^\mu} \int d^4x d^4y e^{-ik'(y-x)} k'^2 \langle 0 | T \epsilon^*(\lambda') \cdot A(x) \hat{O}(0) \epsilon(\lambda) \cdot A(y) | 0 \rangle k'^2$$

- Evaluate electromagnetic part of the Green's function in the lowest-order perturbation theory and write the matrix element

$$\langle \gamma(k\lambda') | \hat{O}(0) | \gamma(k\lambda) \rangle = -e^2 \int d^4x d^4y e^{-ik(y-x)} \epsilon^{*\mu}(\lambda') \langle 0 | T J_\mu(x) \hat{O}(0) J_\nu(y) | 0 \rangle \epsilon^\nu(\lambda)$$

- ▶ Minkowski space correlation function

- Our goal is to calculate quasi-PDF matrix elements associated with the photon's hadronic structure

- ▶ Quasi-PDF formalism

X. Ji [PRL 2013]

Hadronic structure of the “resolved” photon from LQCD

- Quasi-PDF hadronic matrix element with a photon state

$$h_f(z) = \langle \gamma(\vec{p}) | O_f(z, 0) \exp \left(- i q_f e \int_0^z A_\mu(t) dt_\mu \right) | \gamma(\vec{p}) \rangle$$

$$O_f(z, 0) = Z(z) \bar{q}_f(z) \gamma_t U(z, 0) q_f(0)$$

- Including both QCD and QED

$$\blacktriangleright h_f(z, p_z) \rightarrow \sum_{\lambda=x,y} \sum_{\vec{w}, \vec{v}} e^{i\vec{p} \cdot (\vec{w} - \vec{v})} \langle T A_\lambda(T, \vec{v}) O(z, 0) \exp \left(- i q_f e \int_0^z A_\mu(t) dt_\mu \right) A_\lambda(-T, \vec{w}) \exp \left(- i e \int A_\mu(x) J_\mu(x) d^4x \right) \rangle$$

- After integrating out QED part:

$$h(z) = \frac{1}{2} (-ie)^2 \int d^4x \int d^4y \sum_{\lambda=x,y} (e^{|\vec{p}|x_t} e^{-i\vec{p} \cdot \vec{x}} - 1) (e^{-|\vec{p}|y_t} e^{i\vec{p} \cdot \vec{y}} - 1) \langle T J_\lambda(x) O_u(z, 0) J_\lambda(y) \rangle_{\text{QCD}}$$

- Calculable on the lattice

Lattice ensembles

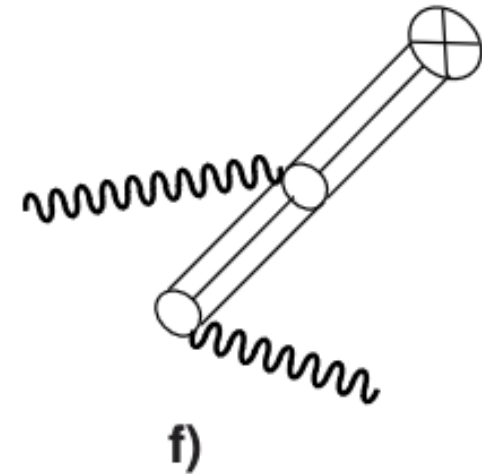
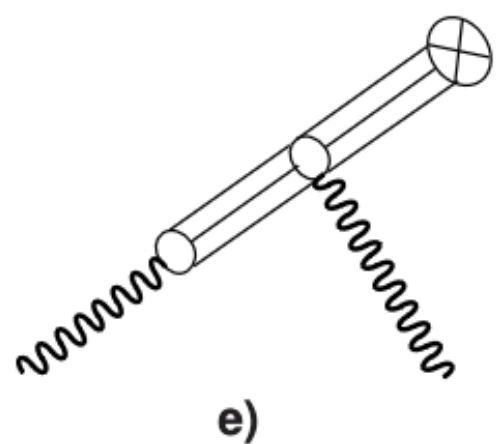
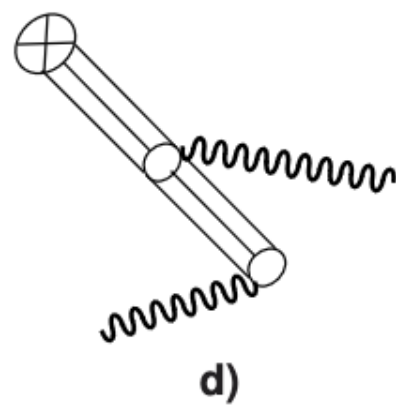
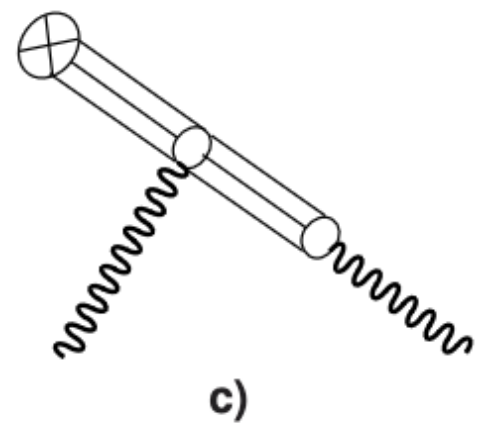
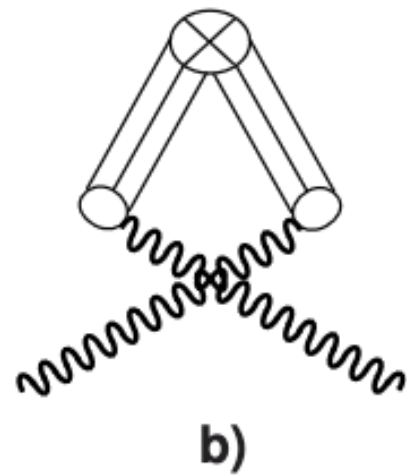
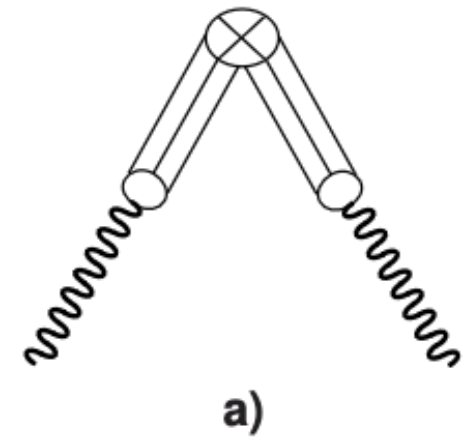
- RBC/UKQCD domain wall fermion

Ensemble-1:	$16^3 \times 32$ ($a \approx 0.12$ fm)
Ensemble-2:	$32^3 \times 64$ ($a \approx 0.063$ fm)
Ensemble-3:	$32^3 \times 64$ ($a \approx 0.088$ fm)

- Light- and strange-quark propagators already generated by the RBC and UKQCD collaborations with $O(1,000)$ randomly chosen point sources per configuration
- ~50 configurations per ensemble
- $m_\pi \approx 450 - 330 \text{ MeV}$

Preliminary matrix elements

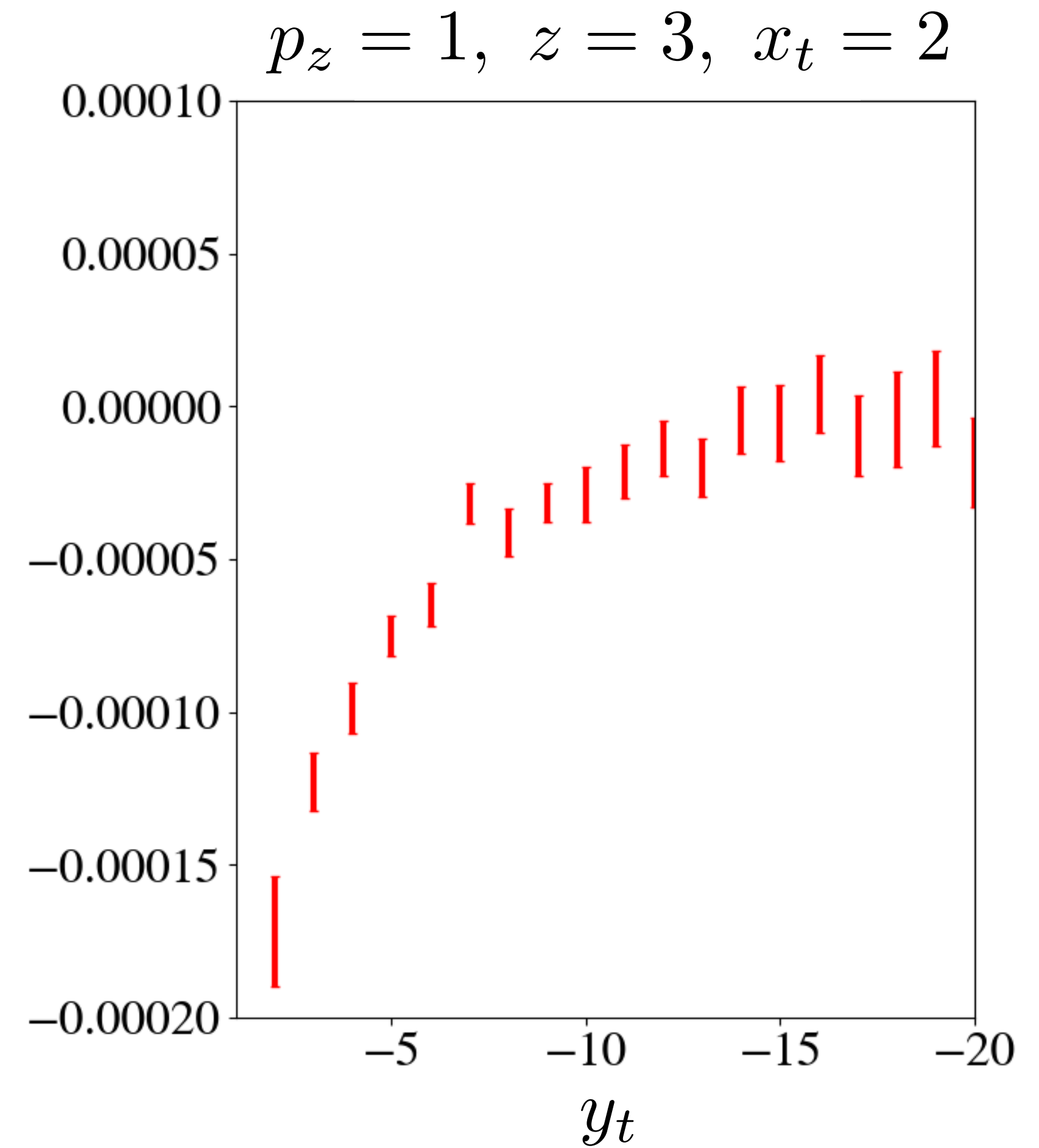
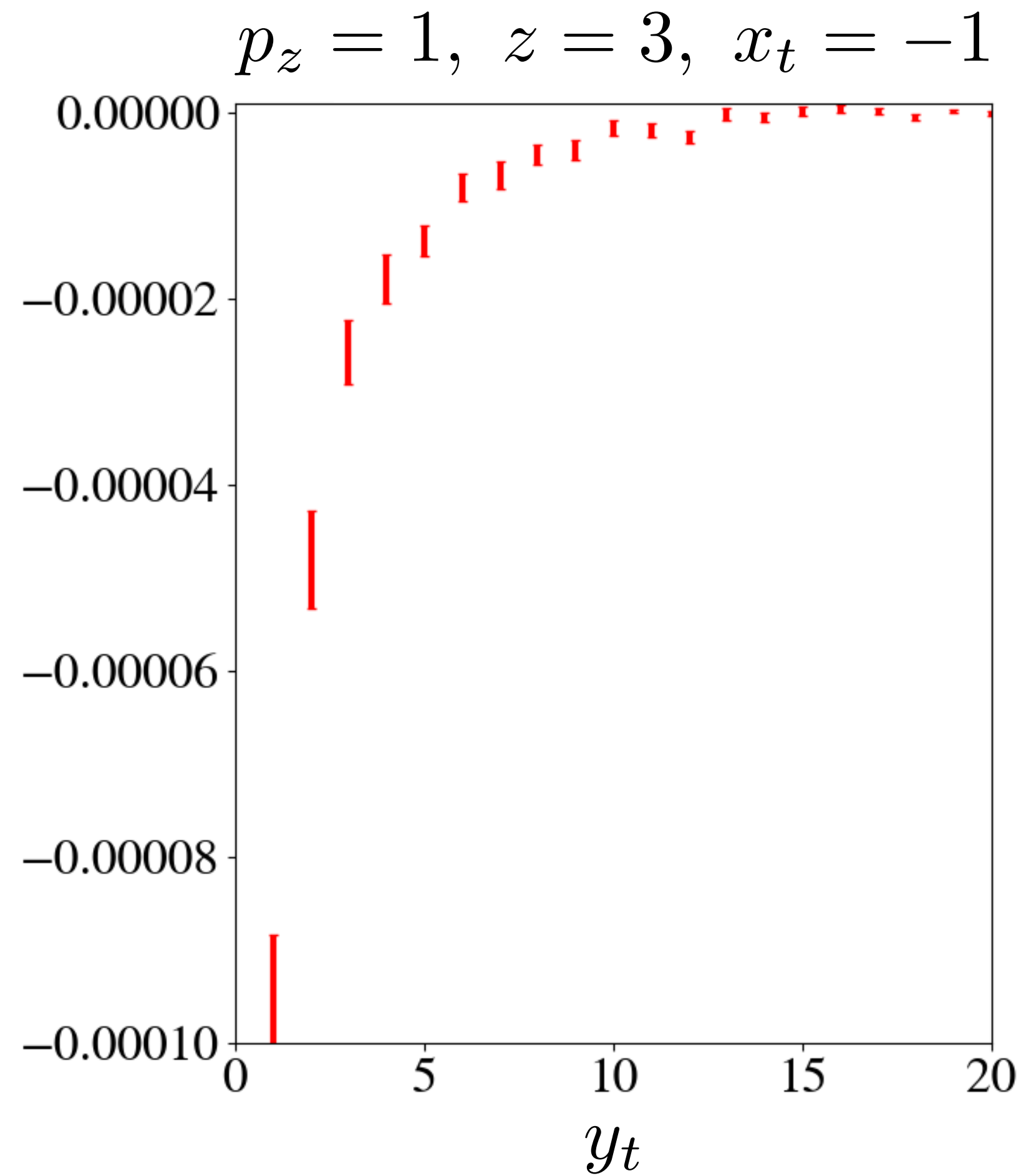
● Different diagrams



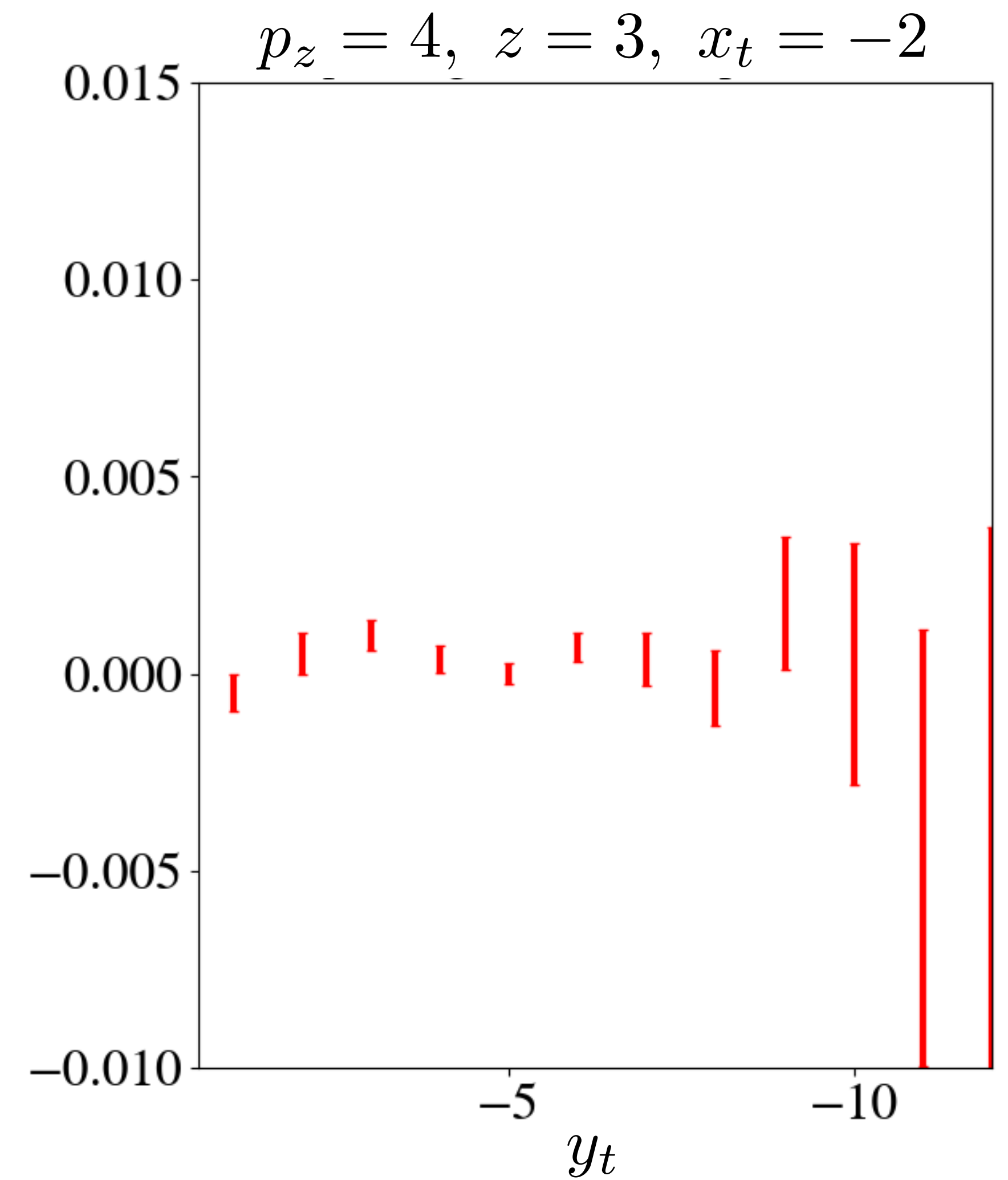
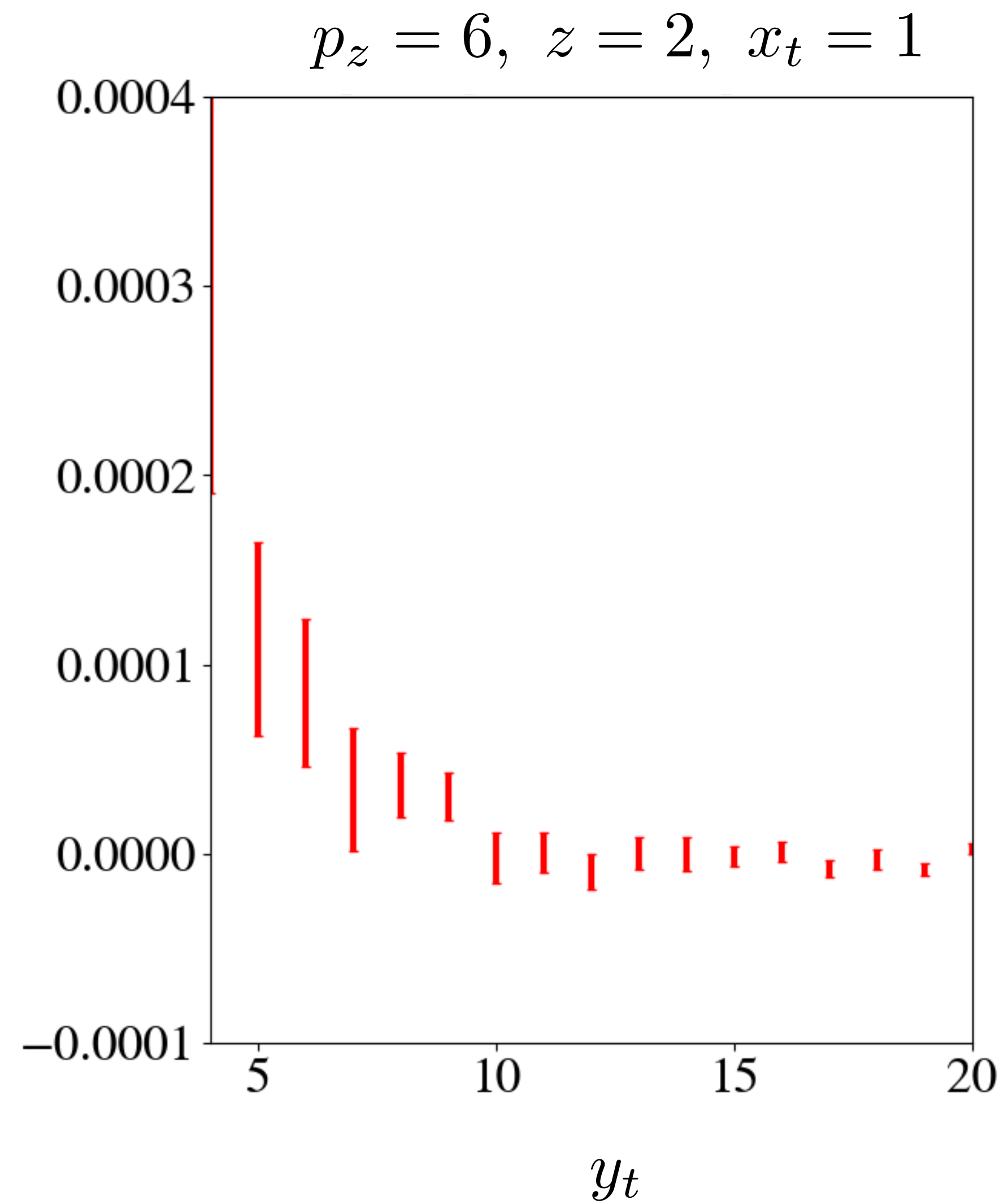
- ▶ The incoming photon converts through J_1 into a QCD eigenstate $|n\rangle$. The operator acts between hadronic states and the second current converts the hadronic state back into the outgoing photon.
- ▶ The operator first acts on the vacuum sector. Subsequent electromagnetic interactions generate and then remove the hadronic component associated with the photon.
- ▶ The incoming photon interacts twice before the operator insertion. At the time when the operator acts, the state can contain a QCD excitation and/or a photon-hadron mixture generated by the two electromagnetic interactions.

Preliminary matrix elements

- Ensemble-3: $32^3 \times 64$ ($a \approx 0.088$ fm)



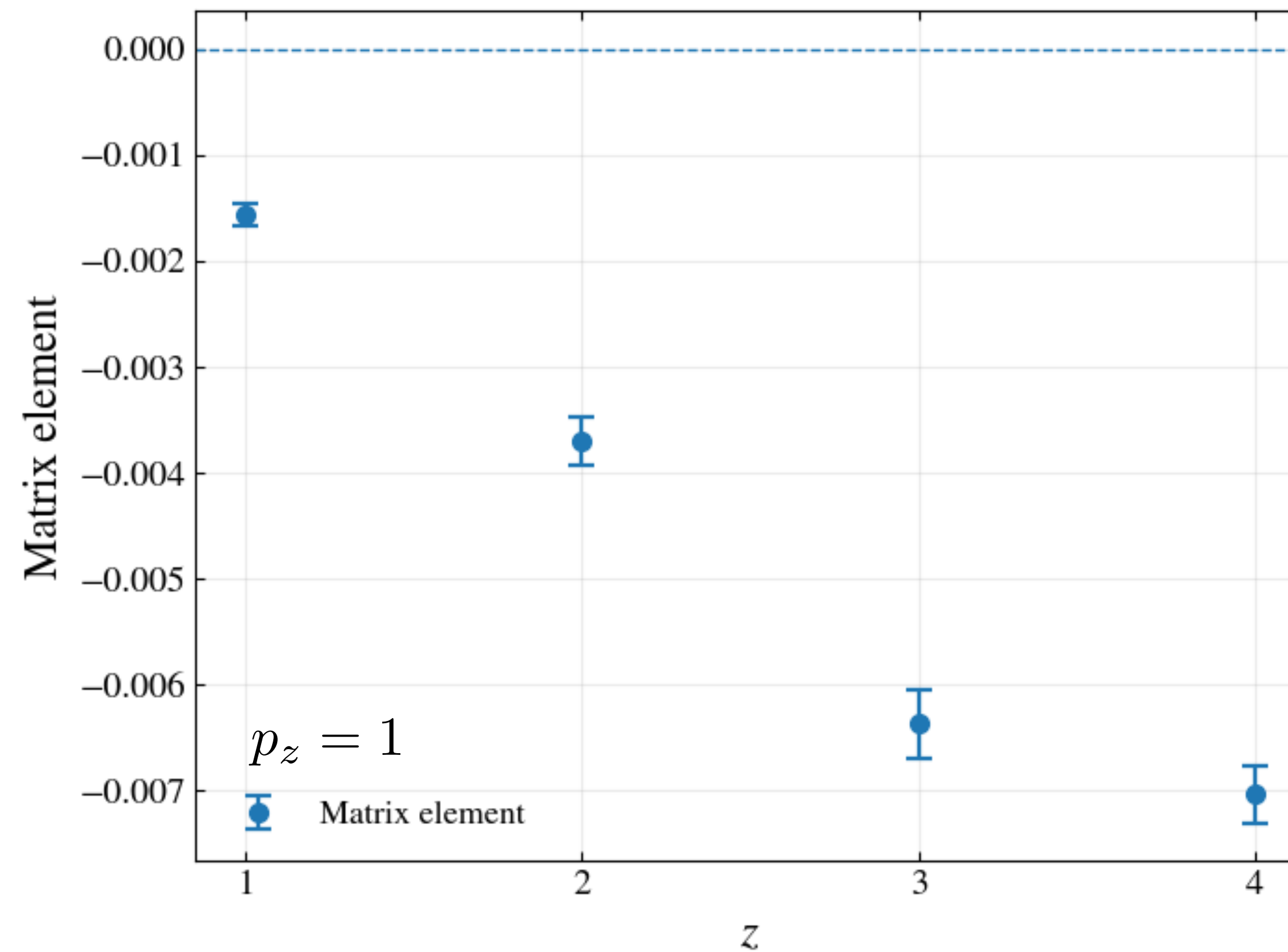
Preliminary matrix elements



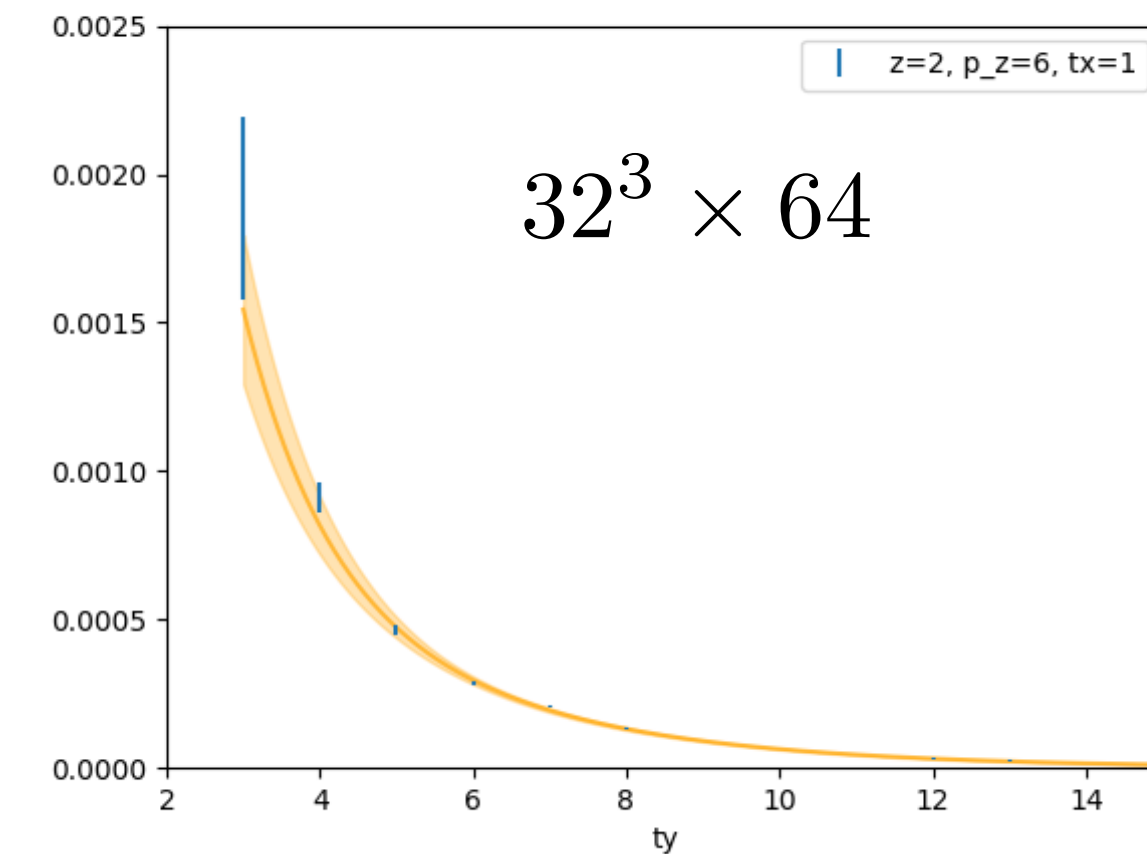
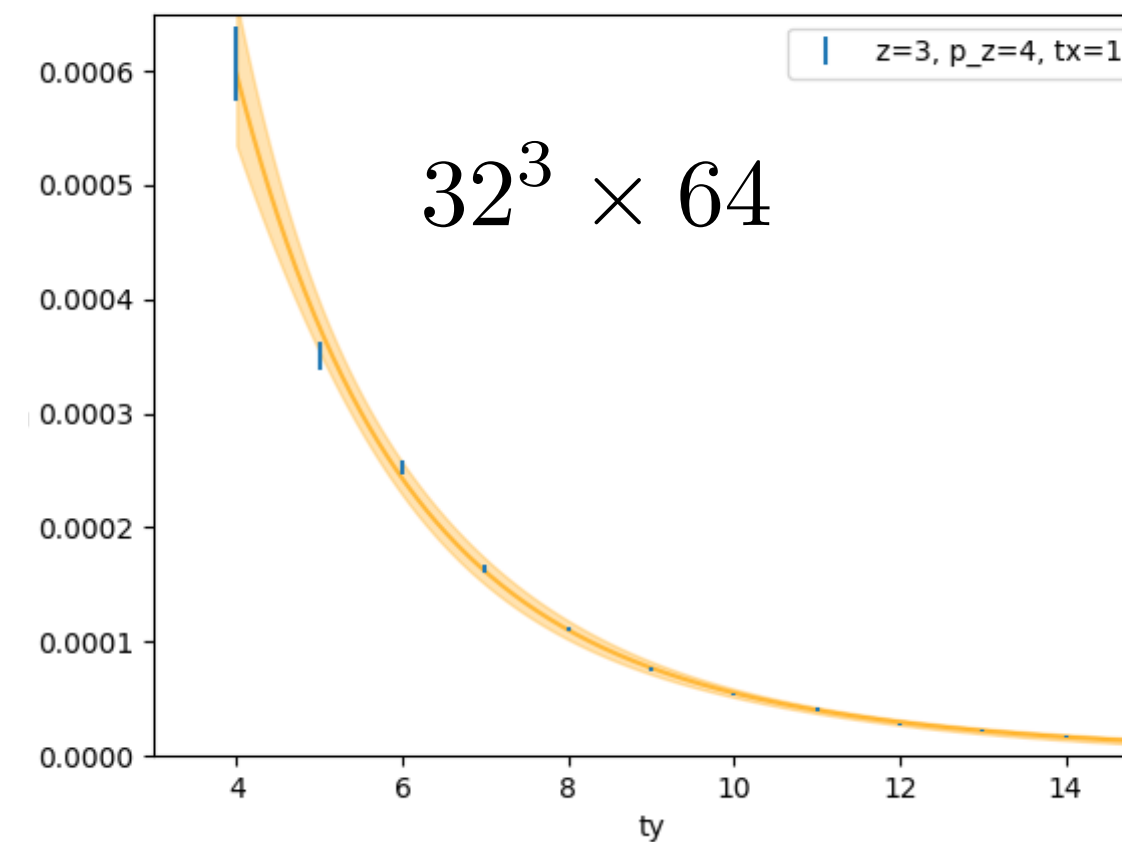
Toward hadronic structure of the “resolved” photon from LQCD

● Recall

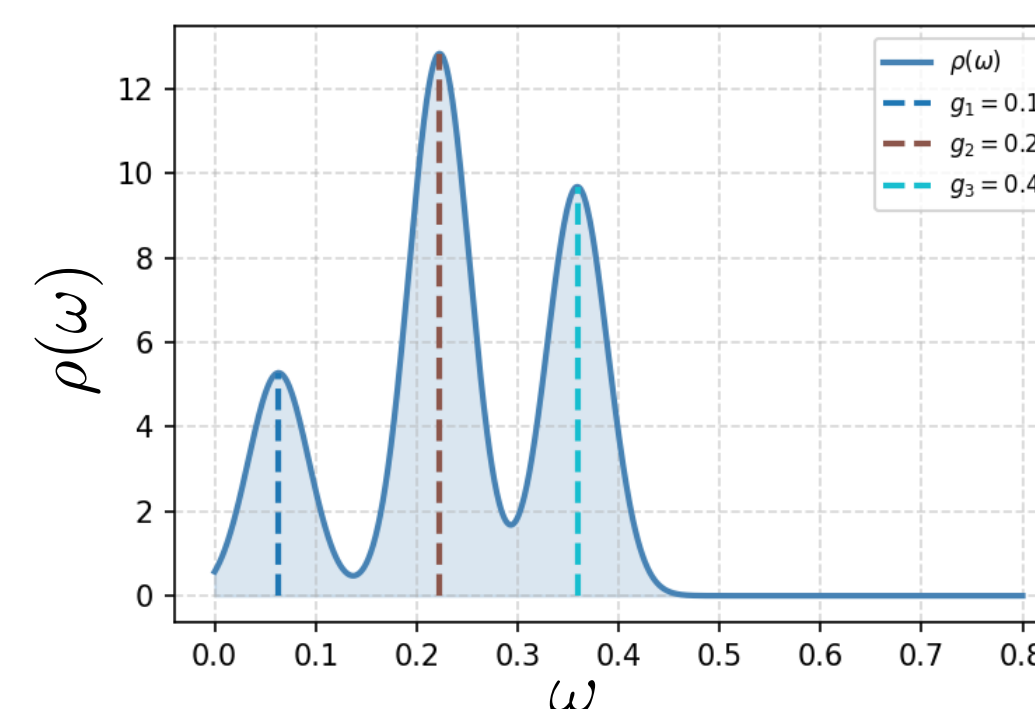
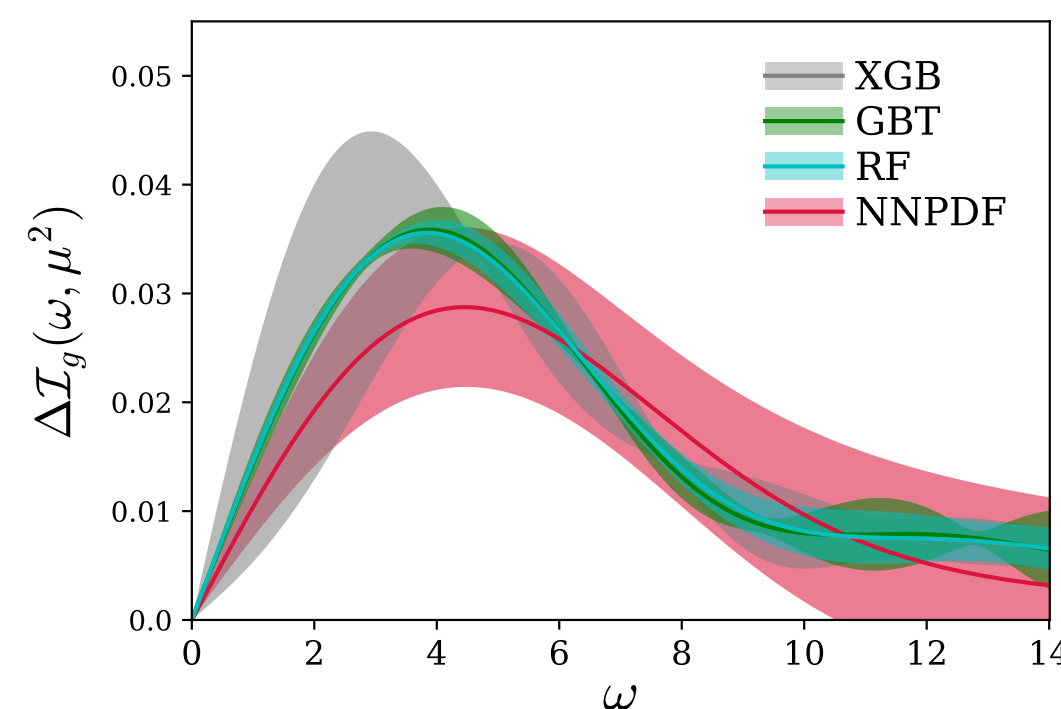
$$h(z) = \frac{1}{2}(-ie)^2 \int d^4x \int d^4y \sum_{\lambda=x,y} (e^{|\vec{p}|x_t} e^{-i\vec{p}\cdot\vec{x}} - 1)(e^{-|\vec{p}|y_t} e^{i\vec{p}\cdot\vec{y}} - 1) \langle T J_\lambda(x) O_u(z, 0) J_\lambda(y) \rangle_{\text{QCD}}$$



- One can do better
- Preliminary fit to matrix elements:
 $32^3 \times 64$ ($a \approx 0.063$ fm)



- Sum over fitted matrix elements
- Renormalization: Ratio-renormalization [Radyushkin 2017] & self renormalization [Ji, et. al. NPB 2021]
- Perturbative matching
- PDF extraction using generative ML (**earlier/ongoing** examples: gluon helicity PDF, reconstruction of spectral function from hadronic tensor)



[arXiv\[2409.17234\]](#)
[arXiv \[2311.04206\]](#)
 In preparation [2026]

- First LQCD determination of the resolved photon's hadronic structure is in progress
 - ▶ QCD allows photon to exhibit a hadronic structure due to quantum fluctuations
 - ▶ EIC will have the unique opportunity to explore (un)polarized photon's structure (essential but hard!)
 - ▶ Can we investigate it on the lattice? - Yes, clear advantage over phenomenology and experiment
 - ▶ Preliminary data indicates existence of nonzero PDF

Thank you!

$$h_f(z) = 2|\vec{p}|L^3 \frac{\sum_{\lambda=x,y} \sum_{\vec{w},\vec{v}} e^{i\vec{p}\cdot(\vec{w}-\vec{v})} \langle 0| T A_\lambda(T, \vec{v}) O_f(z, 0) \exp\left(-iq_f e \int_0^z A_\mu(t) dt_\mu\right) A_\lambda(-T, \vec{w}) |0\rangle_{\text{QCD+QED}}}{\sum_{\lambda=x,y} \sum_{\vec{w},\vec{v}} e^{i\vec{p}\cdot(\vec{w}-\vec{v})} \langle 0| T A_\lambda(T, \vec{v}) A_\lambda(-T, \vec{w}) |0\rangle}$$

$$d\sigma(e + A \rightarrow e' + 2 \text{ jets} + X) = \sum_i \int dy \int dx_\gamma \int dx_A f_{\gamma/e}(y) \times f_{a/\gamma}(x_\gamma, \mu^2) f_{b/A}(x_A, \mu^2) d\hat{\sigma}(ab \rightarrow \text{jets})$$

where $f_{\gamma/e}(y)$ is the photon flux of the electron with y being the momentum fraction carried by the photon; $f_{a/\gamma}(x_\gamma, \mu^2)$ are the photon PDFs in the resolved-photon case, which depend on the parton-in-photon momentum fraction x_γ and the scale μ ; $f_{b/A}(x_A, \mu^2)$ are nuclear PDFs depending on the parton momentum fraction x_A and the scale μ ; $d\hat{\sigma}(ab \rightarrow \text{jets})$ is the cross section of hard scattering of partons a and b into jets.