

Quantum chaos with and without OTOCs

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Part 1

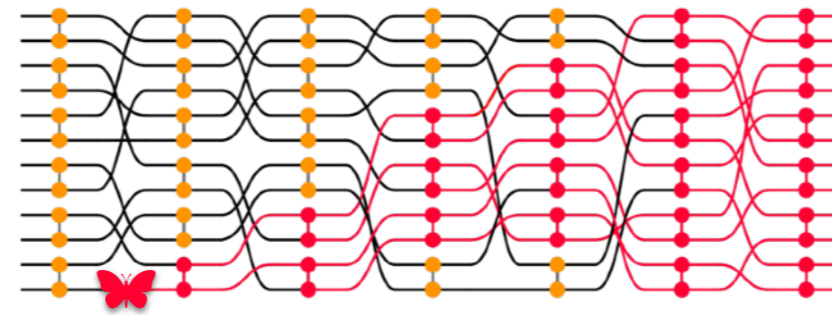
Motivation and
some recent developments

Thermalization & chaos

- ▶ Consider time-dependent, complex systems, esp. at finite temperature
- ▶ Different observables and timescales require different **contours**
- ▶ Via holography: relations to **black holes & quantum gravity**



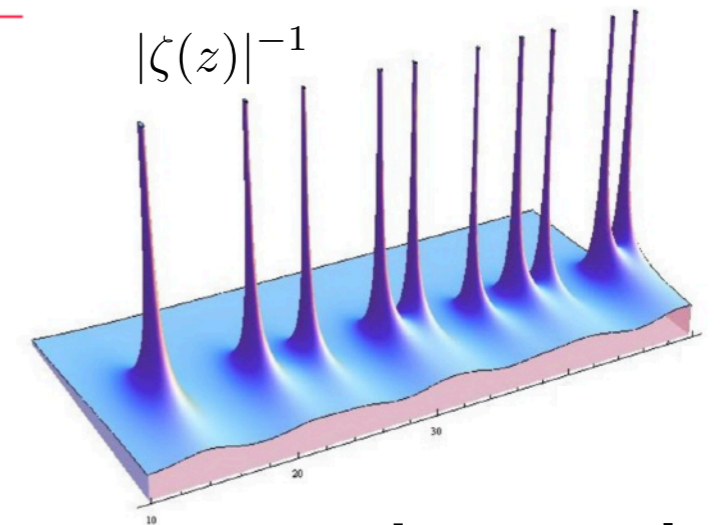
transport & dissipation



operator growth



quantum butterfly effect



random matrix universality

Effective field theory for quantum chaos



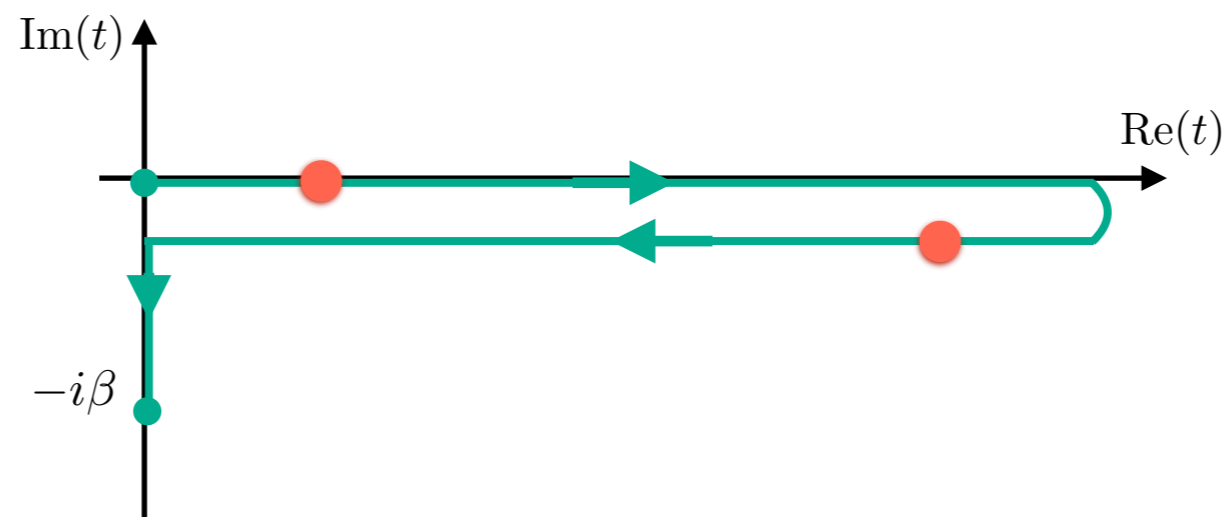
Transport & Dissipation

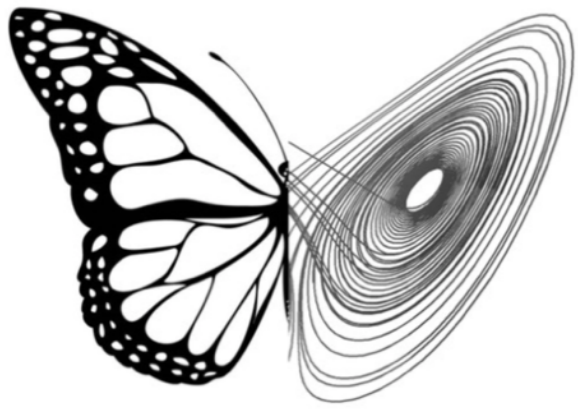
Hydrodynamics: paradigm for universal EFT, slow relaxation of conserved quantities, non-equilibrium path integrals

- ▶ Thermal Schwinger-Keldysh EFT recently understood systematically

[Nickel/Son] [Dubovsky/Hui/Nicolis/Son] [Kovtun/Moore/Romatschke] ...

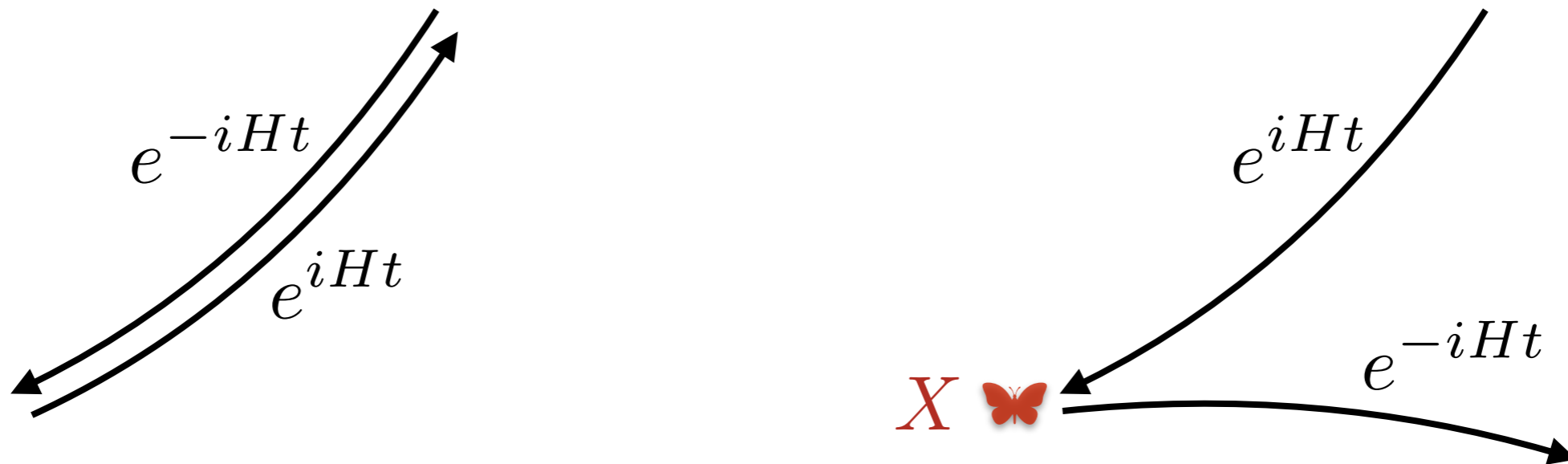
[FMH/Loganayagam/Rangamani] [Crossley/Glorioso/Liu] ... [Jensen/Pinzani-Fokeeva/Yarom] ...



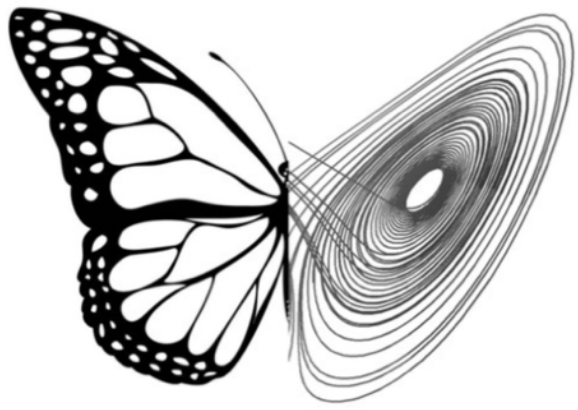


Quantum Butterfly Effect

► $X(t) = e^{iHt} X e^{-iHt}$ is 'complicated' even if X was 'simple':



$$X(t) = e^{iHt} X e^{-iHt} = X + it[H, X] - \frac{t^2}{2}[H, [H, X]] + \dots$$



Quantum Butterfly Effect

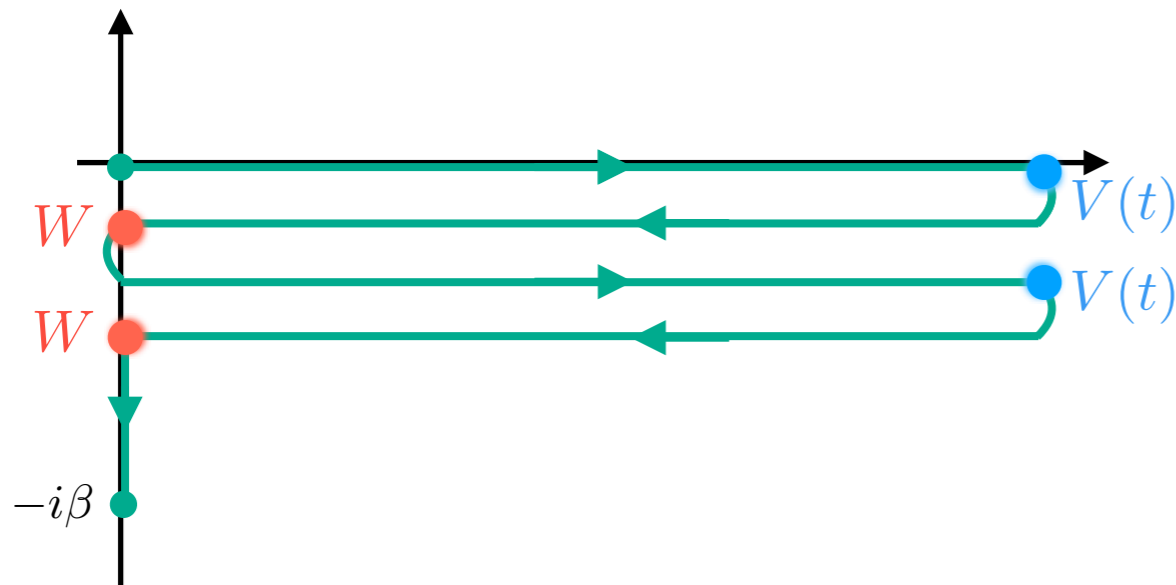
► OTOC = overlap of $|V(t)W(0)\rangle$ and $|W(0)V(t)\rangle$:

$$\langle W(0)V(t)|W(0)V(t)\rangle_\beta \sim \langle VV\rangle_\beta \langle WW\rangle_\beta \left(1 - \frac{c_0}{N} e^{\kappa t} + \dots\right)$$



quantum Lyapunov exponent κ

[Larkin/Ovchinnikov] [Kitaev] [Shenker/Stanford] ...



- Chaos bound: $\kappa \leq 2\pi T$
[Maldacena/Shenker/Stanford]
- Holography: $\kappa^{(\text{black hole})} = 2\pi T$
[Sekino/Susskind] [Shenker/Stanford] ...

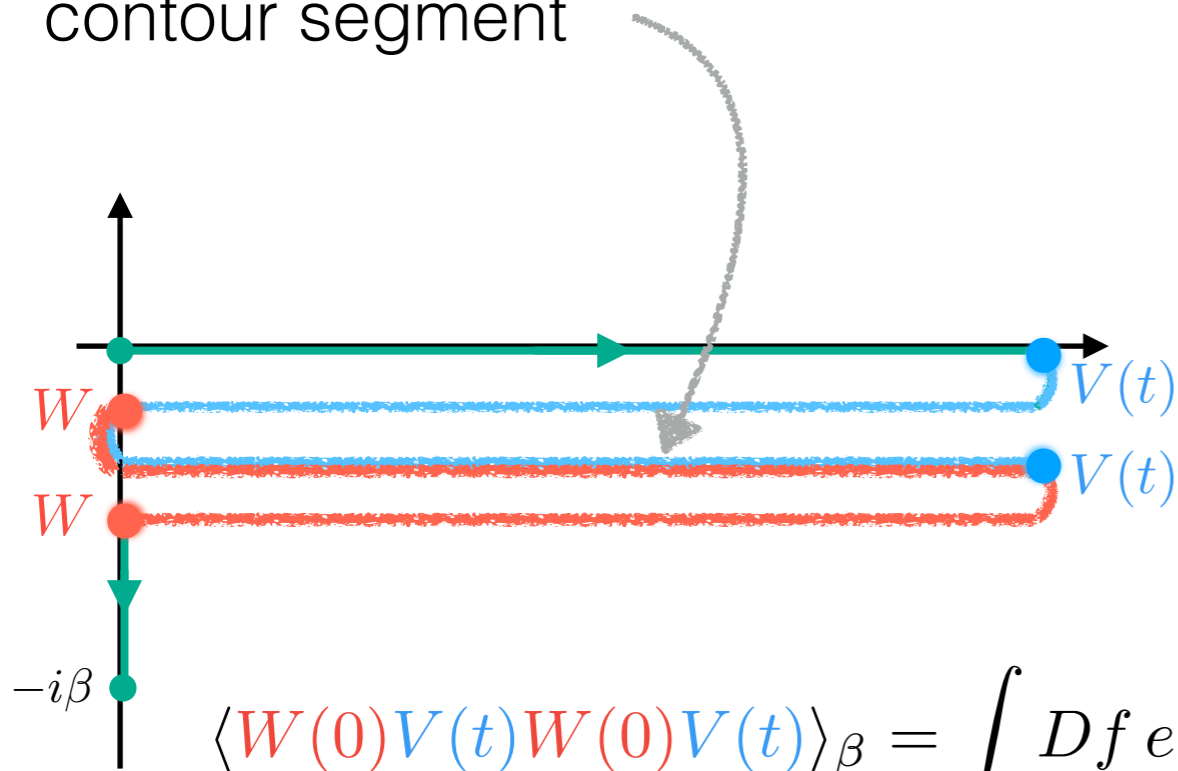
► **Hydro-like chaos EFT on four-fold contour: effective chaos mode (“scramblon”)**

[Maldacena/Stanford/Yang] [Jensen] ... [Blake/Lee/Liu] [FMH/Rozali] ...

► Can sometimes derive explicitly from microscopics, e.g. in SYK-type models

[Stanford/Yang/Yao] [Gu/Kitaev/Zhang] [Gao/Liu] [Choi/FMH/Mezei/Sarosi] ...

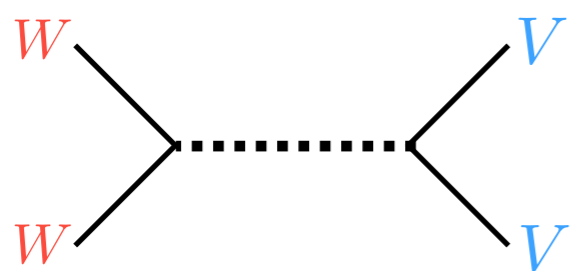
► **Effective d.o.f.:** soft modes of the mean-field action that interact on shared contour segment



► Large-N saddle point breaks down when $t \sim \log(N)$

► Exact path integral over soft modes leads to **eikonal effective action**

$$\langle W(0)V(t)W(0)V(t) \rangle_\beta = \int Df e^{-S_{\text{eff}}[f]} \langle VV \rangle_\beta^{(f)} \langle WW \rangle_\beta^{(f)}$$

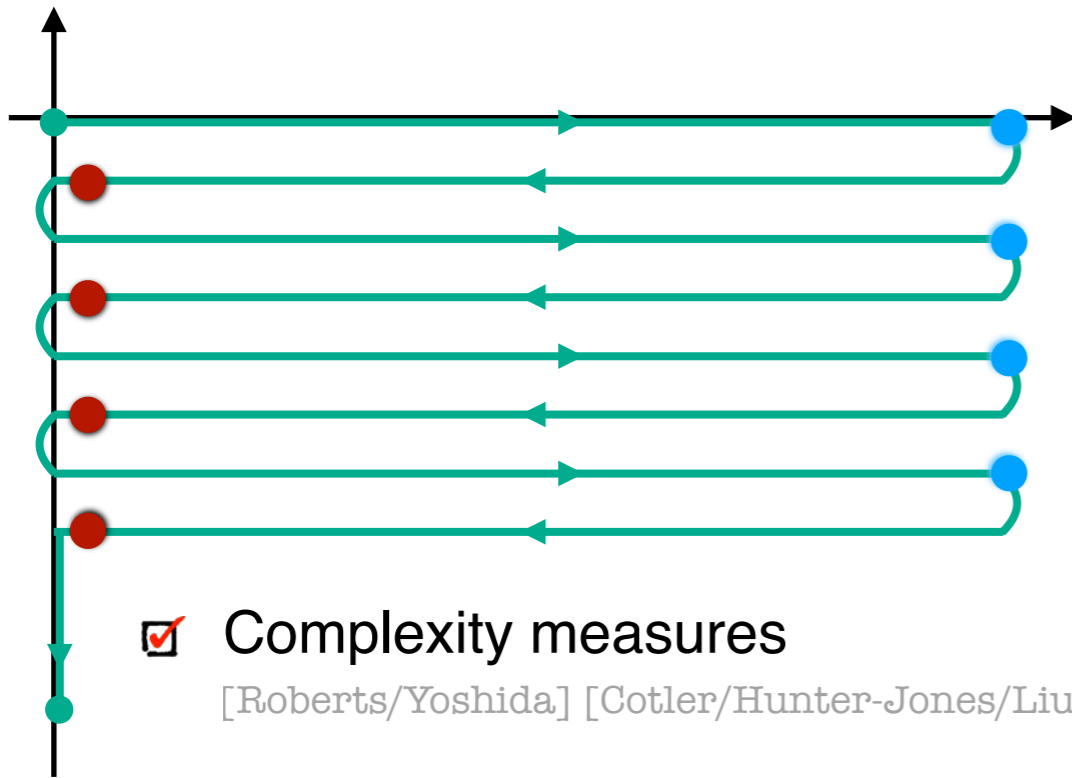


$$f(t, x) \sim f_0(t, x) + X^+(x) e^{-\kappa t} + X^-(x) e^{\kappa t} + \dots$$

$$S_{\text{eik}} \sim N \int dk e^{\kappa(\frac{i\pi}{2} - t)} (\dots) X^+(k) X^-(-k)$$

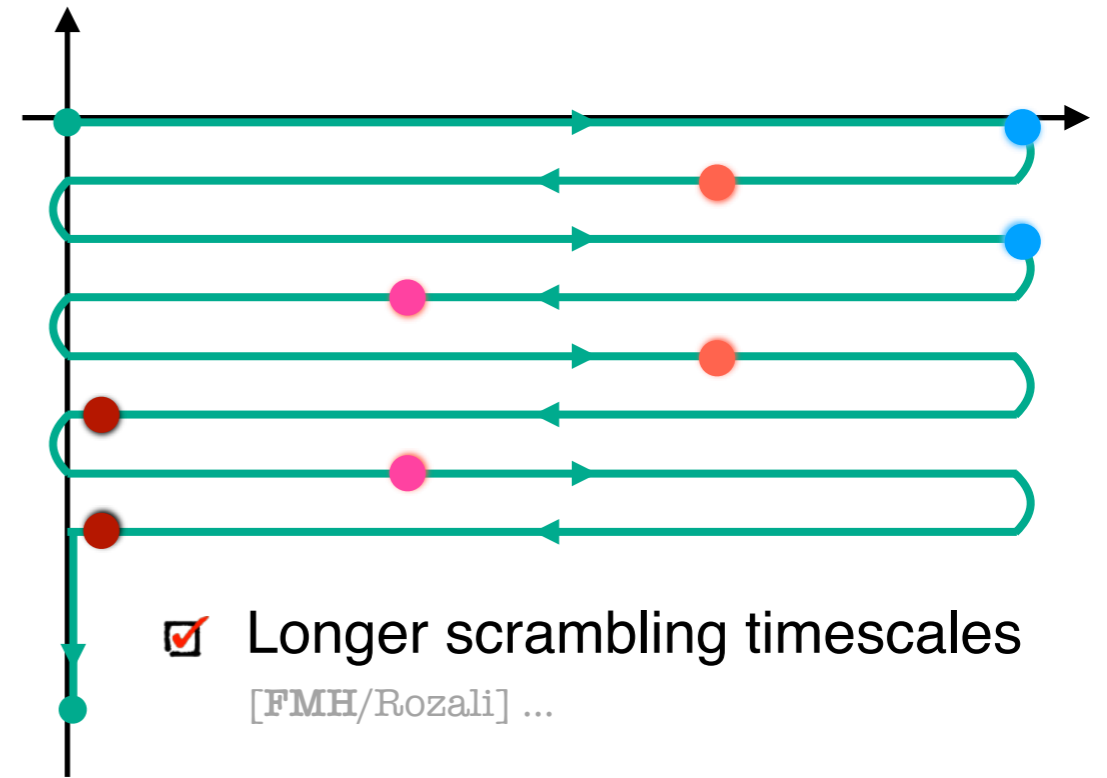
Chaos beyond the standard OTOC

► One way to generalise $\langle W V W V \rangle$: higher-order OTOCs



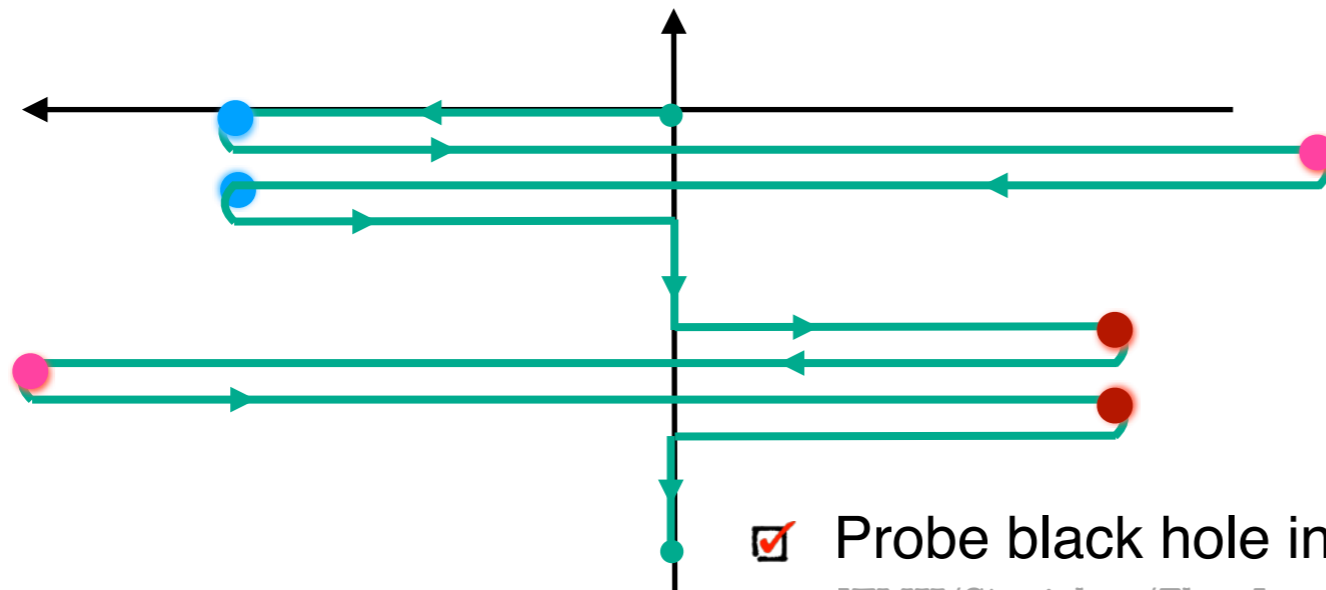
☑ Complexity measures

[Roberts/Yoshida] [Cotler/Hunter-Jones/Liu/Yoshida] ...



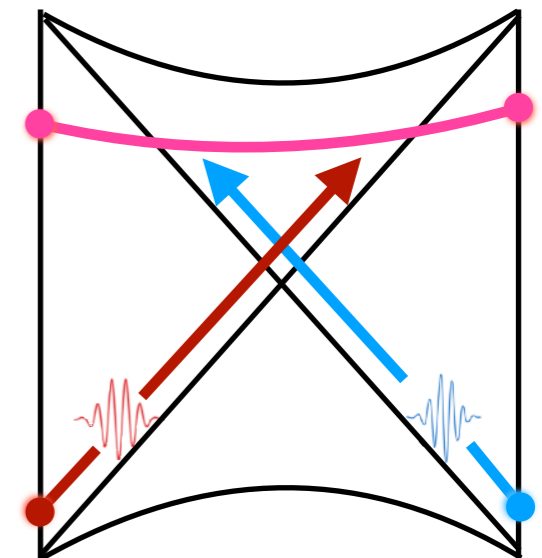
☑ Longer scrambling timescales

[FMH/Rozali] ...



☑ Probe black hole interior

[FMH/Streicher/Zhao] ...



☑ Higher order fluctuation-dissipation theorems

[Tsuji/Shitara/Ueda] [FMH/Loganayagam/Narayan/Nizami/Rangamani] ...

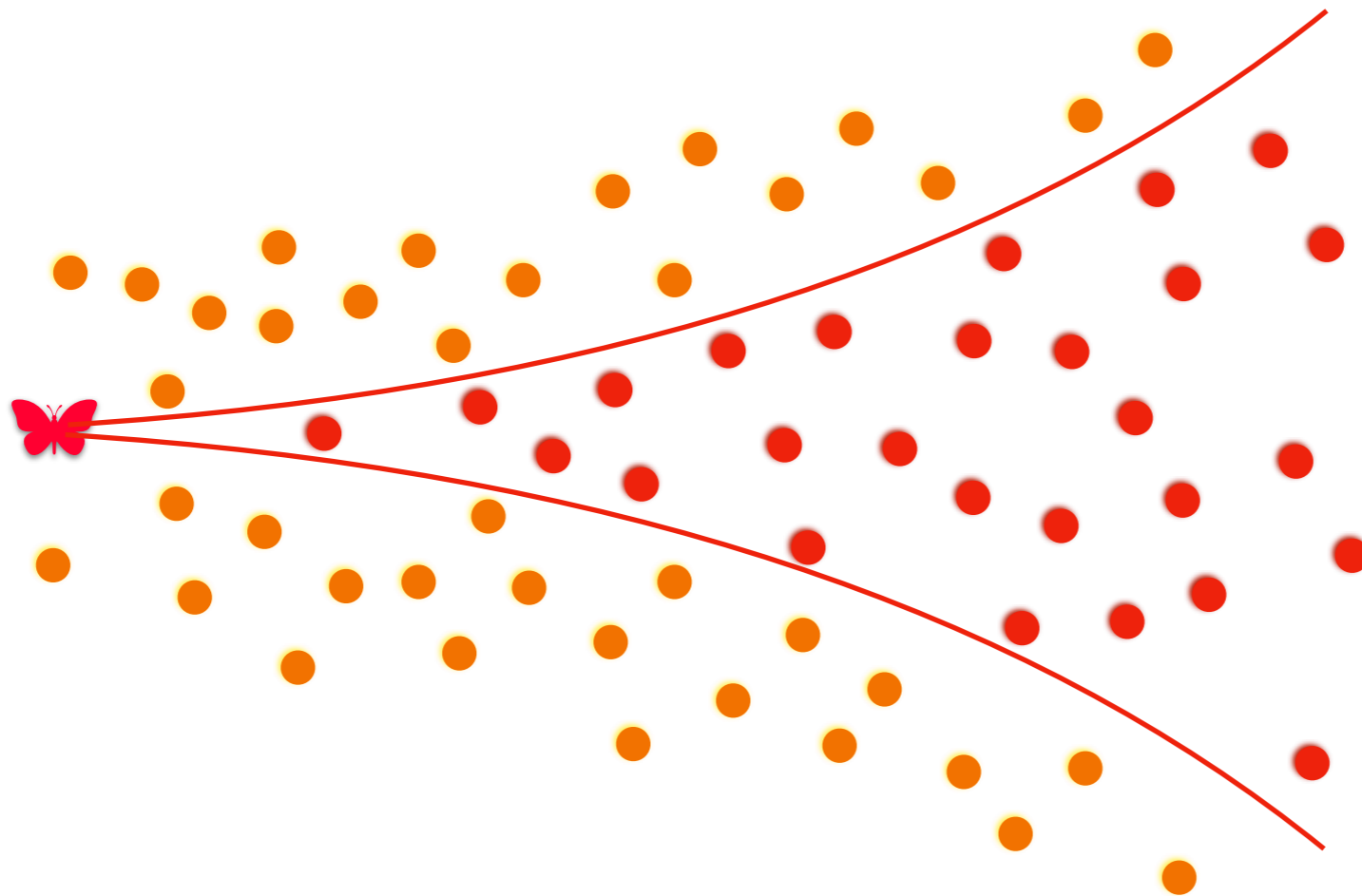
► Can also discover **quantum chaos signatures** in lower-order (Schwinger-Keldysh) correlators and related observables. E.g.:

► **Pole skipping**: signatures of quantum Lyapunov physics in $\langle TT \rangle_{\text{ret}}$

[Grozdanov/Schalm/Scopelliti] [Blake/Davison/Grozdanov/Liu] ... [FMH/Rozali] ...

► **Operator growth**: spread of support over a given operator basis

$$X(t) = X + it[H, X] - \frac{t^2}{2}[H, [H, X]] + \dots$$



- ▶ Can also discover **quantum chaos signatures** in **lower-order** (Schwinger-Keldysh) correlators and related observables. E.g.:

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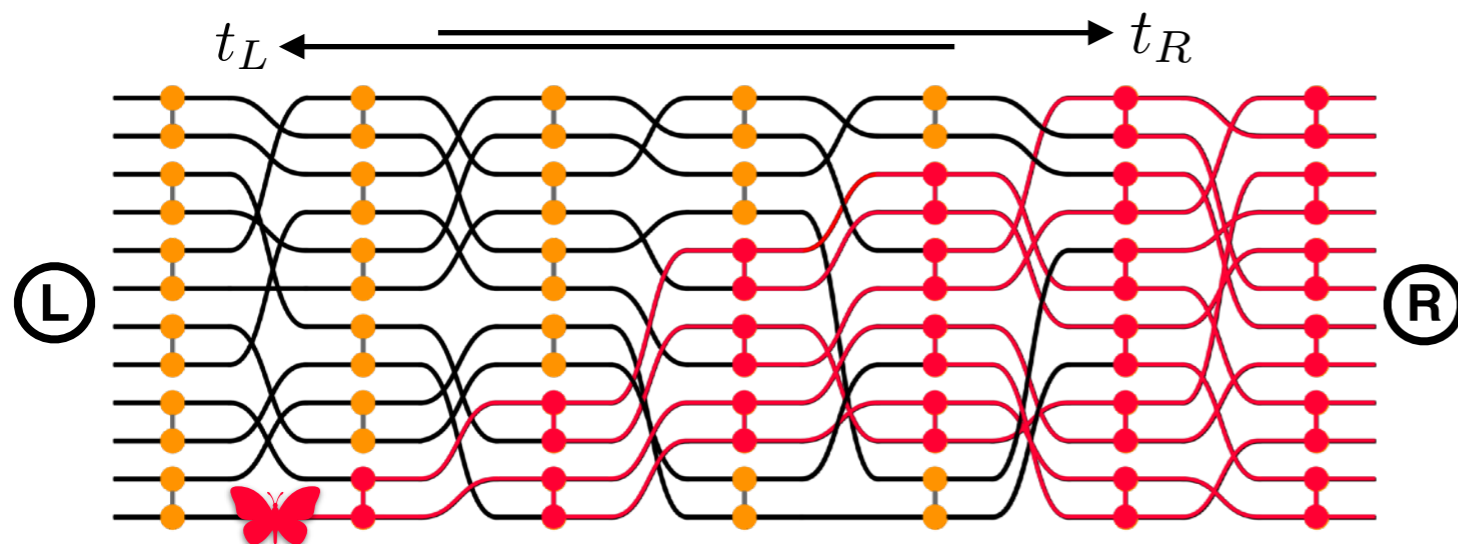
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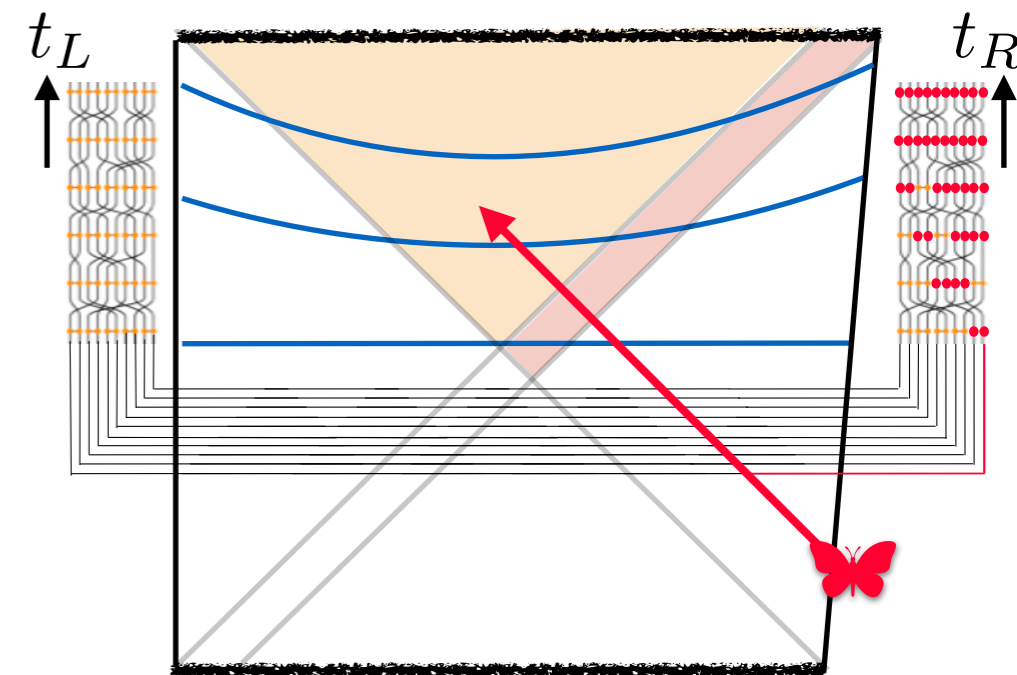
$$X(t) = X + it[H, X] - \frac{t^2}{2}[H, [H, X]] + \dots$$

- ▶ Toy model: random unitary **quantum circuit**

[Hayden/Preskill] [Susskind] ... [Zhao] [FMH/Zhao] ...



- ▶ Perturbation spreads exponentially via all-to-all interactions (“gates”) until saturation



- ▶ Toy model for holographic black hole evolution: perturbation = shockwave

Part 2

Operator growth from the OPE

[Chowdhury/**FMH**/Sanchez-Garrido/Zhao, PRL 136 (2026) 12]

Motivation

- ▶ Another aspect of operator growth:



“How complicated has a simple perturbation become?”



“How do two growing perturbations interact over time?”

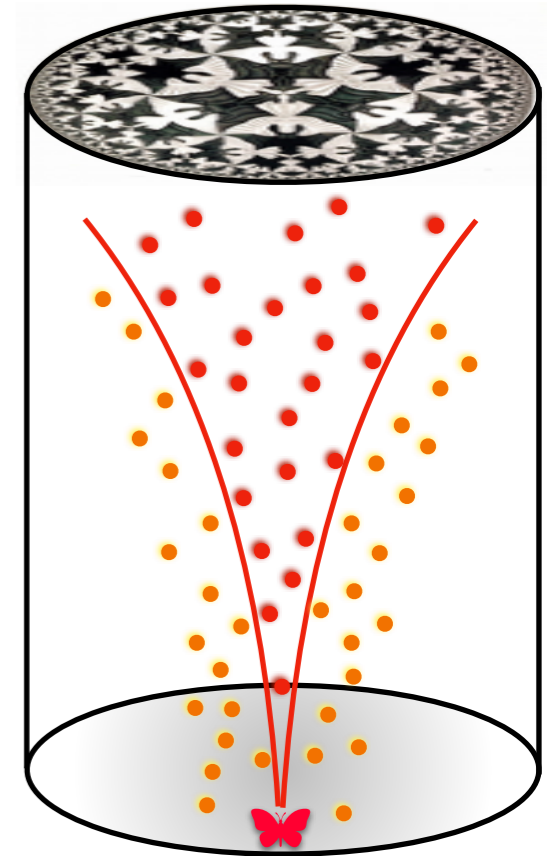
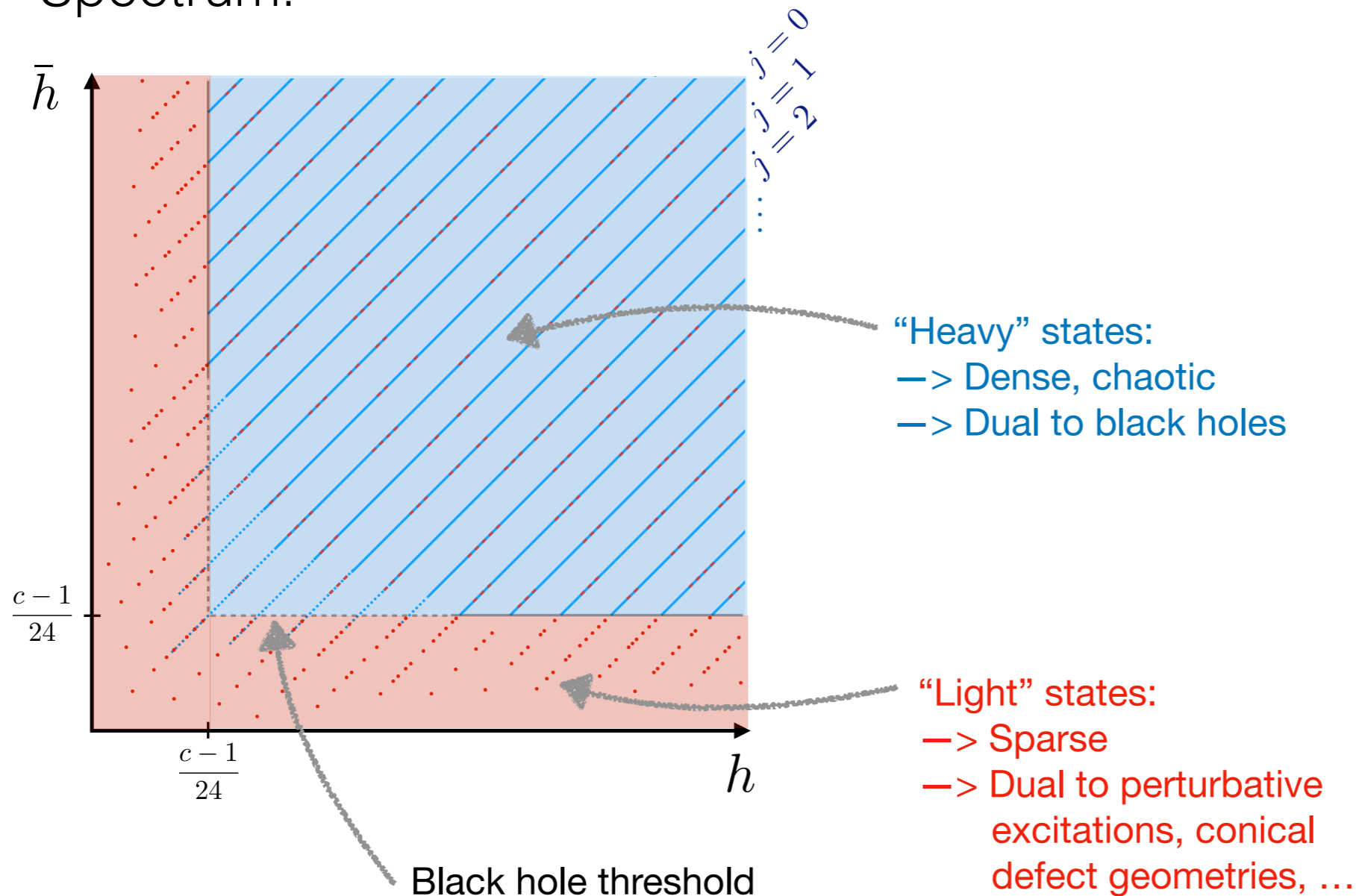
“What operator best describes their combined effect?”

Setup

- ▶ Consider chaotic, large c , holographic CFT₂
- ▶ Can make notions of operator growth very precise

[Dymarsky/Smolkin] ... [Caputa/Magan/Patramanis] [Caputa/Datta] ...

- ▶ Spectrum:

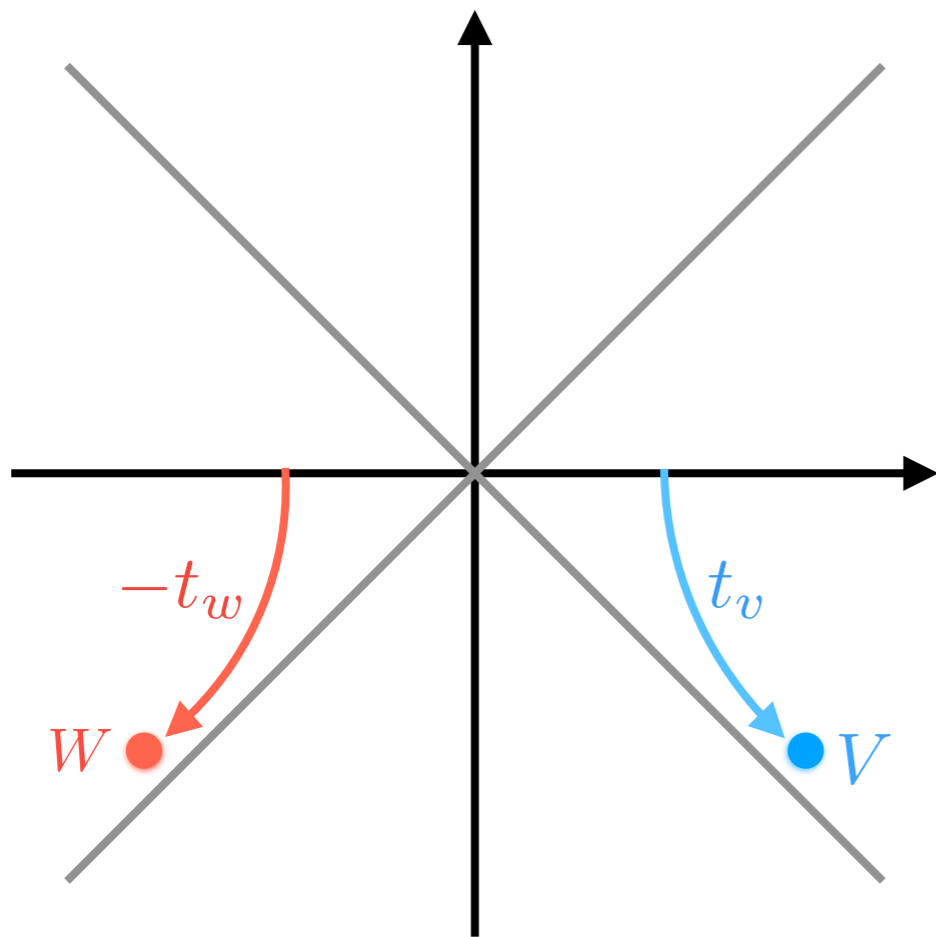


Setup

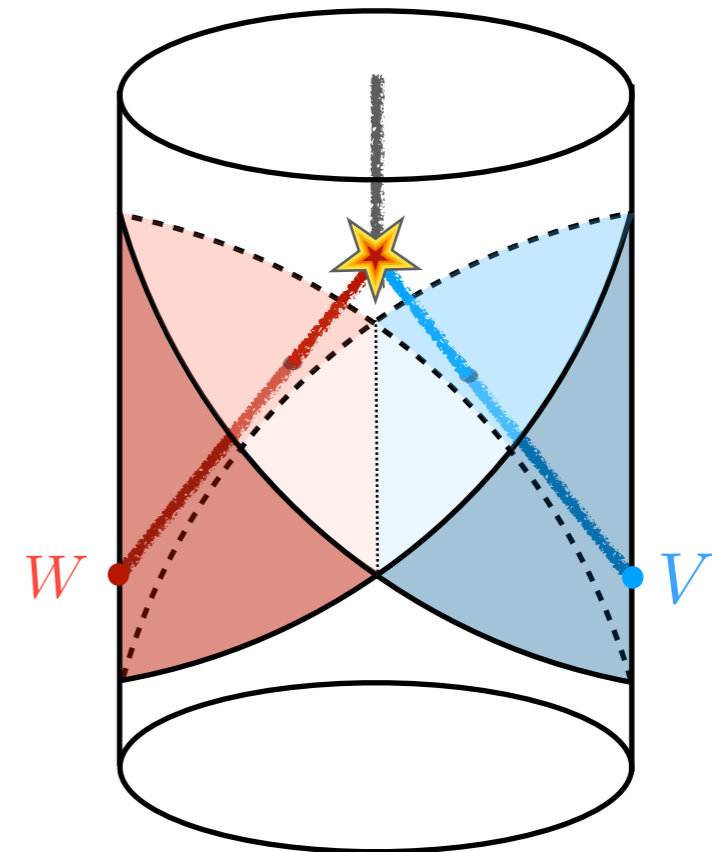
- Consider the two-operator state for light primary operators in CFT₂:

$$|\Psi_{W_L V_R}\rangle = W(-t_w - i\pi, x_w) V(t_v, x_v) |0\rangle \quad (t_v < 0, t_w < 0)$$

Rindler time evolution:



AdS₃ shockwaves:



Setup

- ▶ Consider the two-operator state for light primary operators in CFT₂:

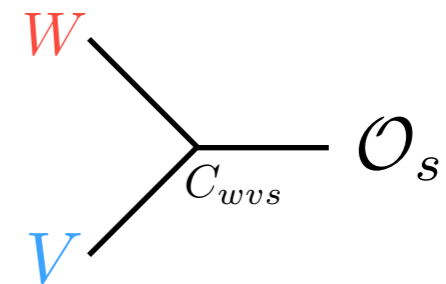
$$|\Psi_{W_L V_R}\rangle = W(-t_w - i\pi, x_w) V(t_v, x_v) |0\rangle \quad (t_v < 0, t_w < 0)$$

- ▶ Decompose into irreps of the Virasoro algebra (operator product expansion):

$$|\Psi_{W_L V_R}\rangle \propto \sum_{\mathcal{O}_s} C_{wvs} |\mathcal{B}_{WV\mathcal{O}_s}(t_w, x_w; t_v, x_v)\rangle$$

OPE coefficients

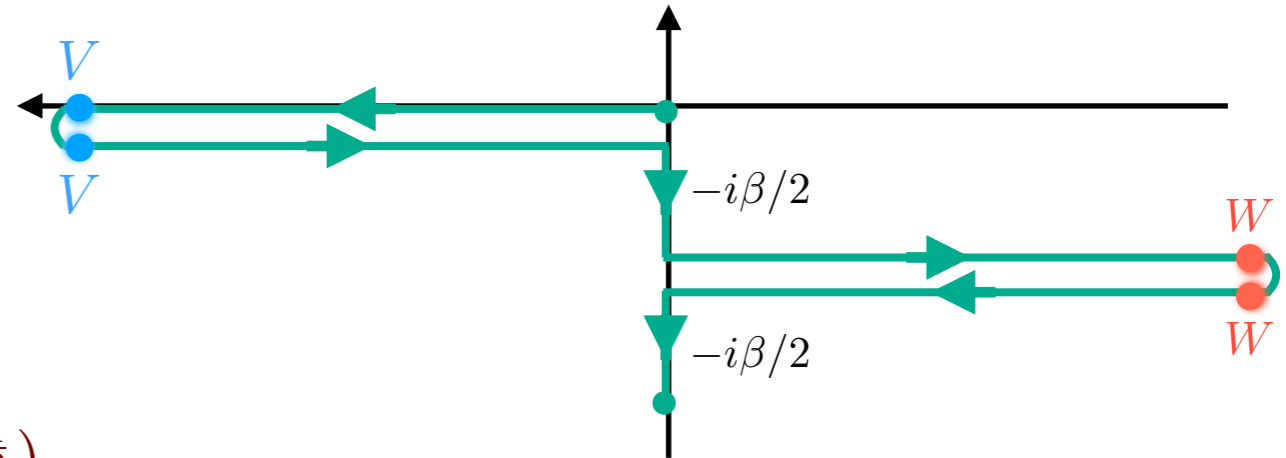
OPE blocks



- ▶ How does the “wavefunction” on the space of $\mathcal{B}_{WV\mathcal{O}_s}$ depend on $t \equiv -t_v - t_w$?

Self-overlap

► Compute self-overlap of $|\Psi_{W_L V_R}\rangle$:

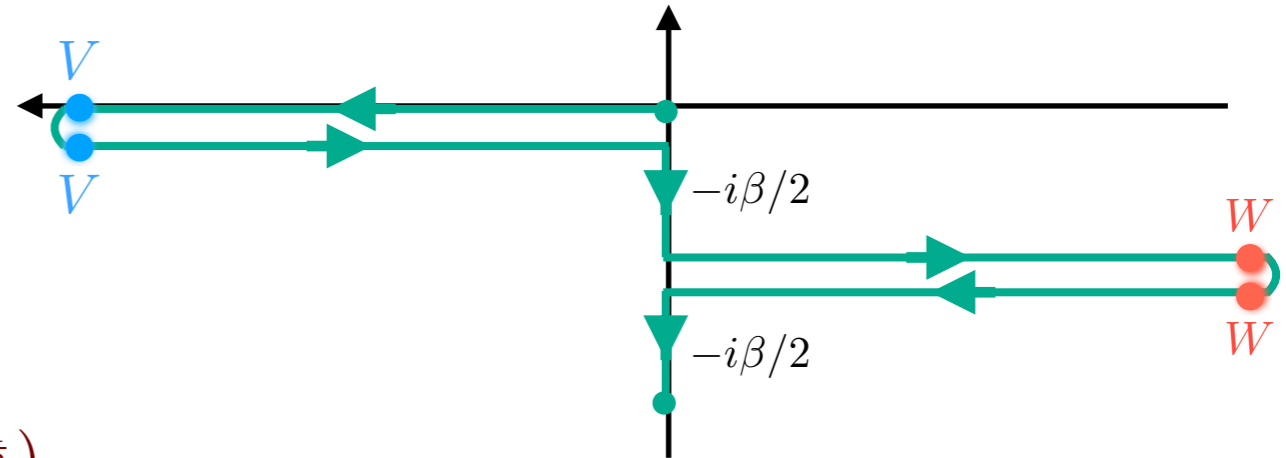


$$\mathcal{F} = \frac{\langle \Psi_{W_L V_R} | \Psi_{W_L V_R} \rangle}{\langle VV \rangle \langle WW \rangle} = \frac{\text{tr}(W^\dagger W \rho^{\frac{1}{2}} V^\dagger V \rho^{\frac{1}{2}})}{\text{tr}(V^\dagger V \rho) \text{tr}(W^\dagger W \rho)}$$

—> not an OTOC!
 —> time ordered

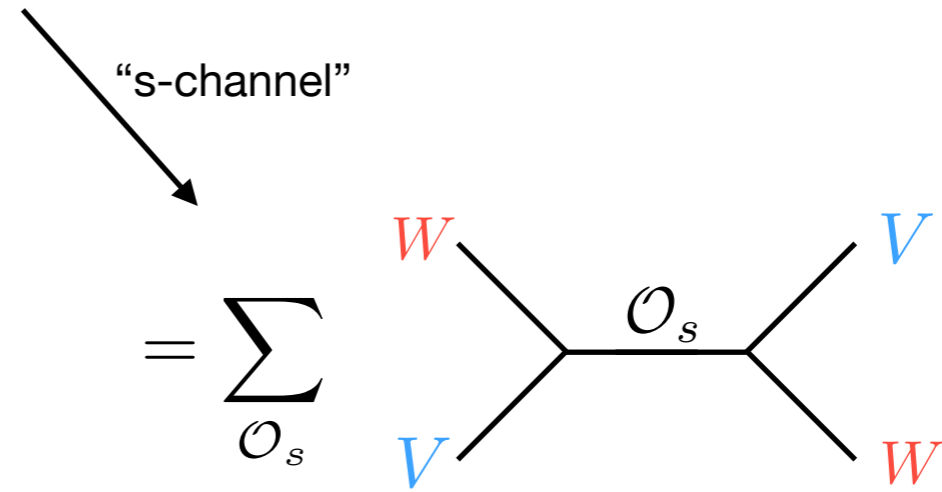
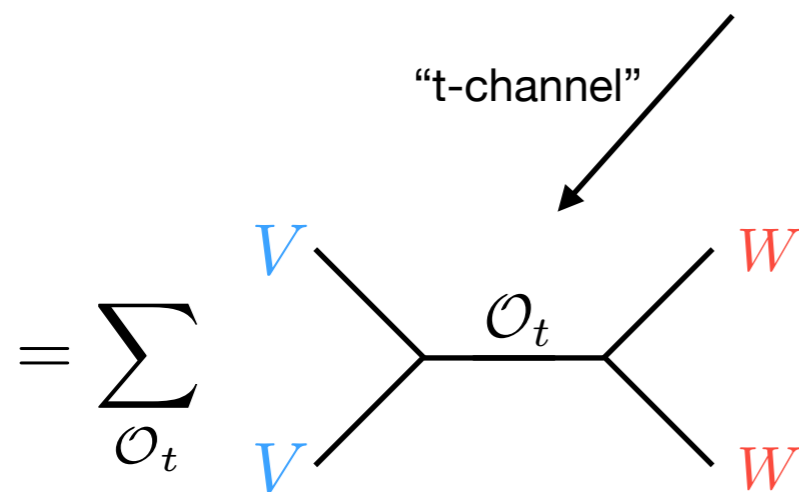
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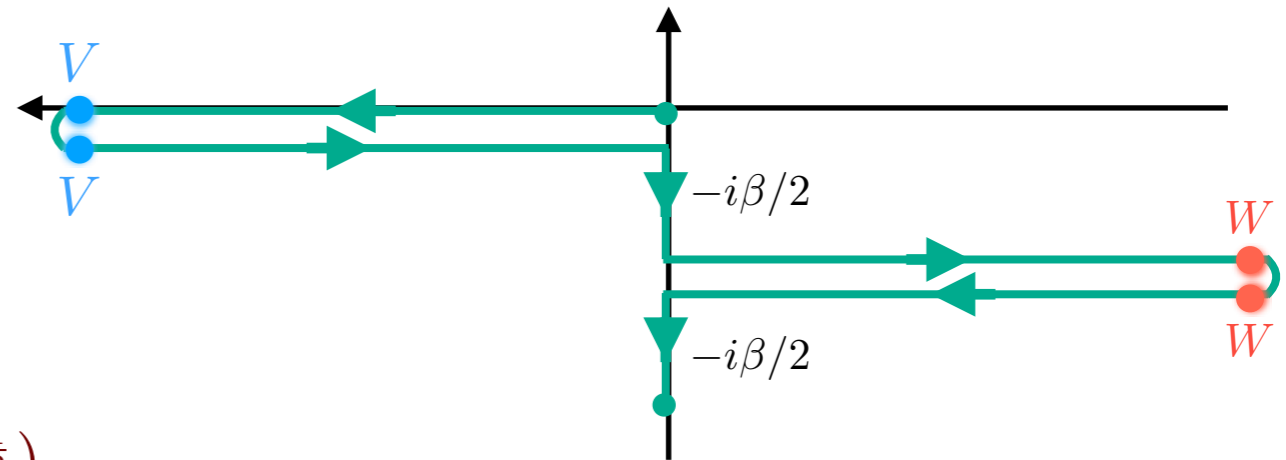
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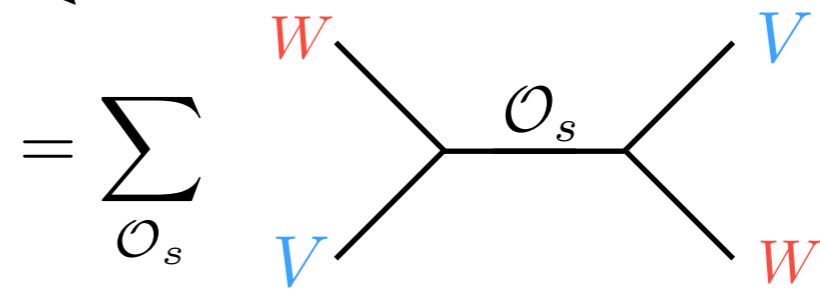
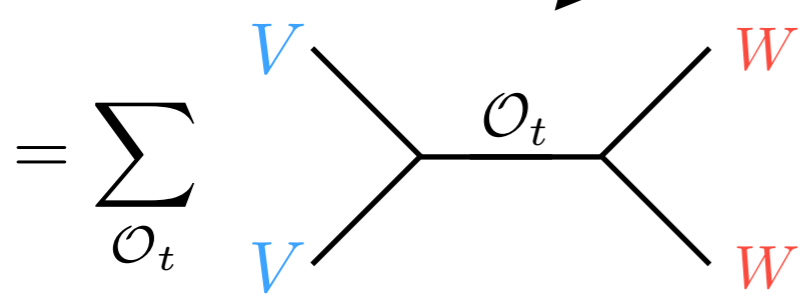


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→ not an OTOC!
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“t-channel”

“s-channel”

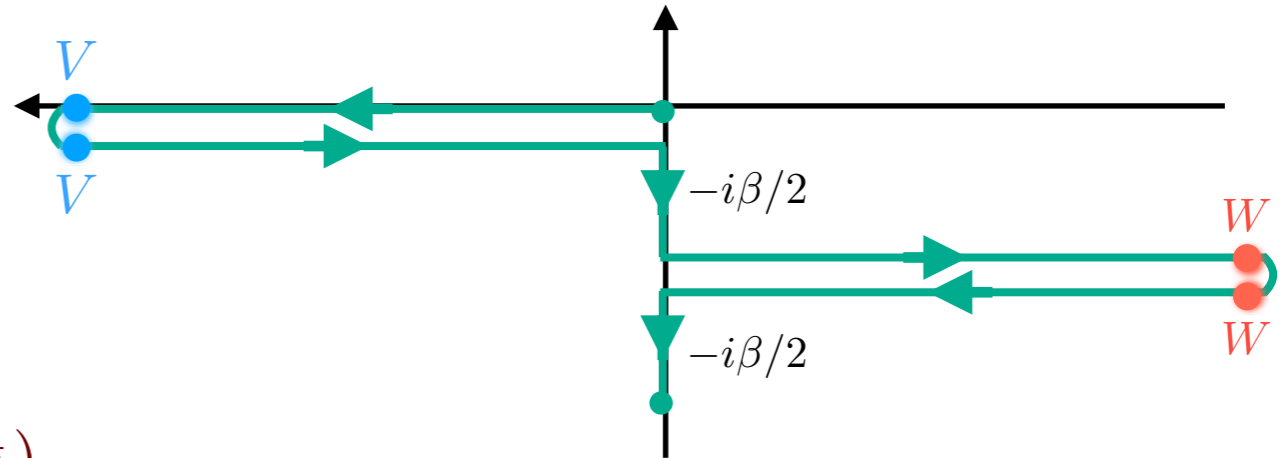


$$= \text{Diagram with } \mathbf{1} \text{ and } + \mathcal{O}(e^{-t\Delta_{\text{gap}}})$$

$$\approx 1$$

Self-overlap

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→ not an OTOC!
→ time ordered

“t-channel”

$$= \sum_{\mathcal{O}_t} \begin{array}{c} V \\ \diagdown \\ \mathcal{O}_t \\ \diagup \\ V \end{array} \begin{array}{c} W \\ \diagdown \\ \mathcal{O}_t \\ \diagup \\ W \end{array}$$

“s-channel”

$$= \sum_{\mathcal{O}_s} \begin{array}{c} W \\ \diagdown \\ \mathcal{O}_s \\ \diagup \\ V \end{array} \begin{array}{c} V \\ \diagdown \\ \mathcal{O}_s \\ \diagup \\ W \end{array}$$

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$$= \sum_{\mathcal{O}_s} C_{wvs}^2 \underbrace{\langle \mathcal{B}_{WV\mathcal{O}_s} | \mathcal{B}_{WV\mathcal{O}_s} \rangle}_{\text{conformal blocks } \mathcal{V}_s(z, \bar{z})}$$

$$\approx 1$$

$$1 = \sum_{\mathcal{O}_s} C_{wvs}^2 \underbrace{\langle \mathcal{B}_{WVO_s} | \mathcal{B}_{WVO_s} \rangle}_{\text{conformal blocks } \mathcal{V}_s(z, \bar{z})}$$

- How does the decomposition of 1 into s-channel conformal blocks change as a function of $t \equiv -t_v - t_w$?

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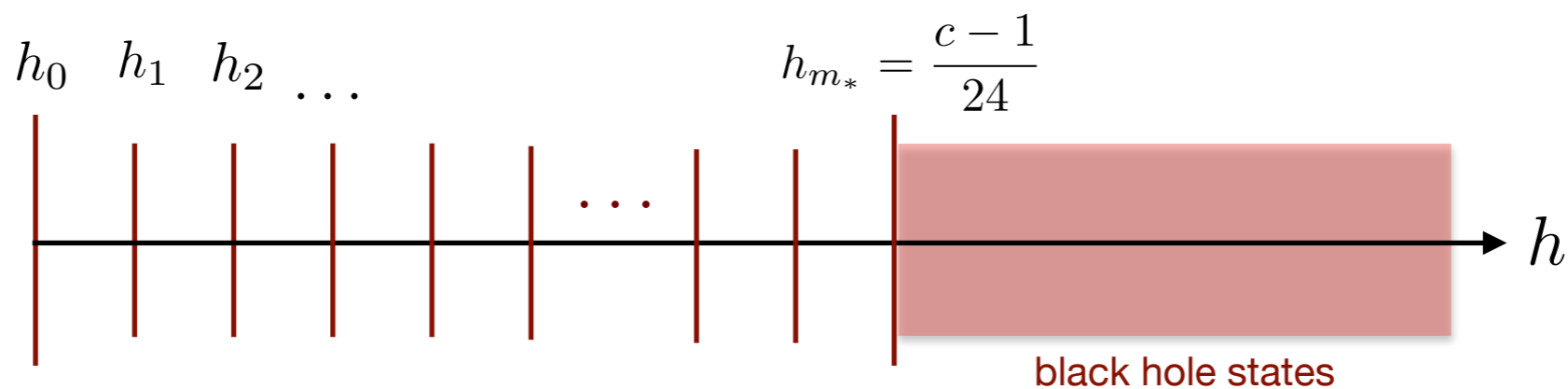
► Conformal bootstrap tells us the approximate spectrum of exchanges \mathcal{O}_s

[Ponsot/Teschner] ... [Caron-Huot] ... [Collier/Gobeil/Maxfield/Perlmutter] [Das/Datta/Pal] ...

► Labelled by s-channel conformal weights (h_s, \bar{h}_s) or $(E_s, J_s) = (\frac{h_s + \bar{h}_s}{2}, \bar{h}_s - h_s)$:

Discrete light “double-twist” operators $\mathcal{O}_s^{(m, \bar{m})}$: $h_m = h_v + h_w + m$ ($m = 1, 2, \dots, m_*$)

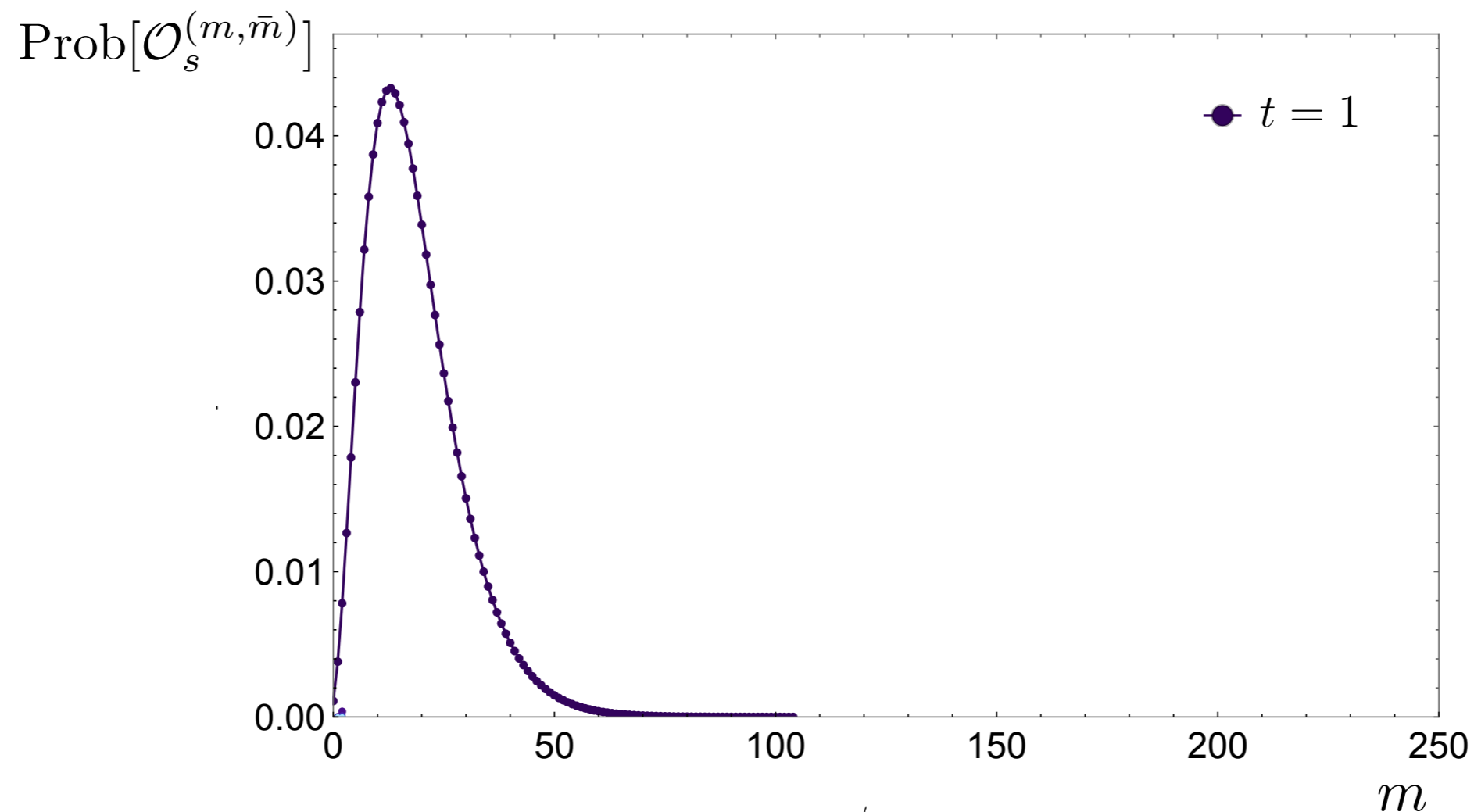
Continuum of heavy “mean-field” operators: $h_s > \frac{c-1}{24}$



► Higher boost t excites intermediate states \mathcal{O}_s with higher energy E_s

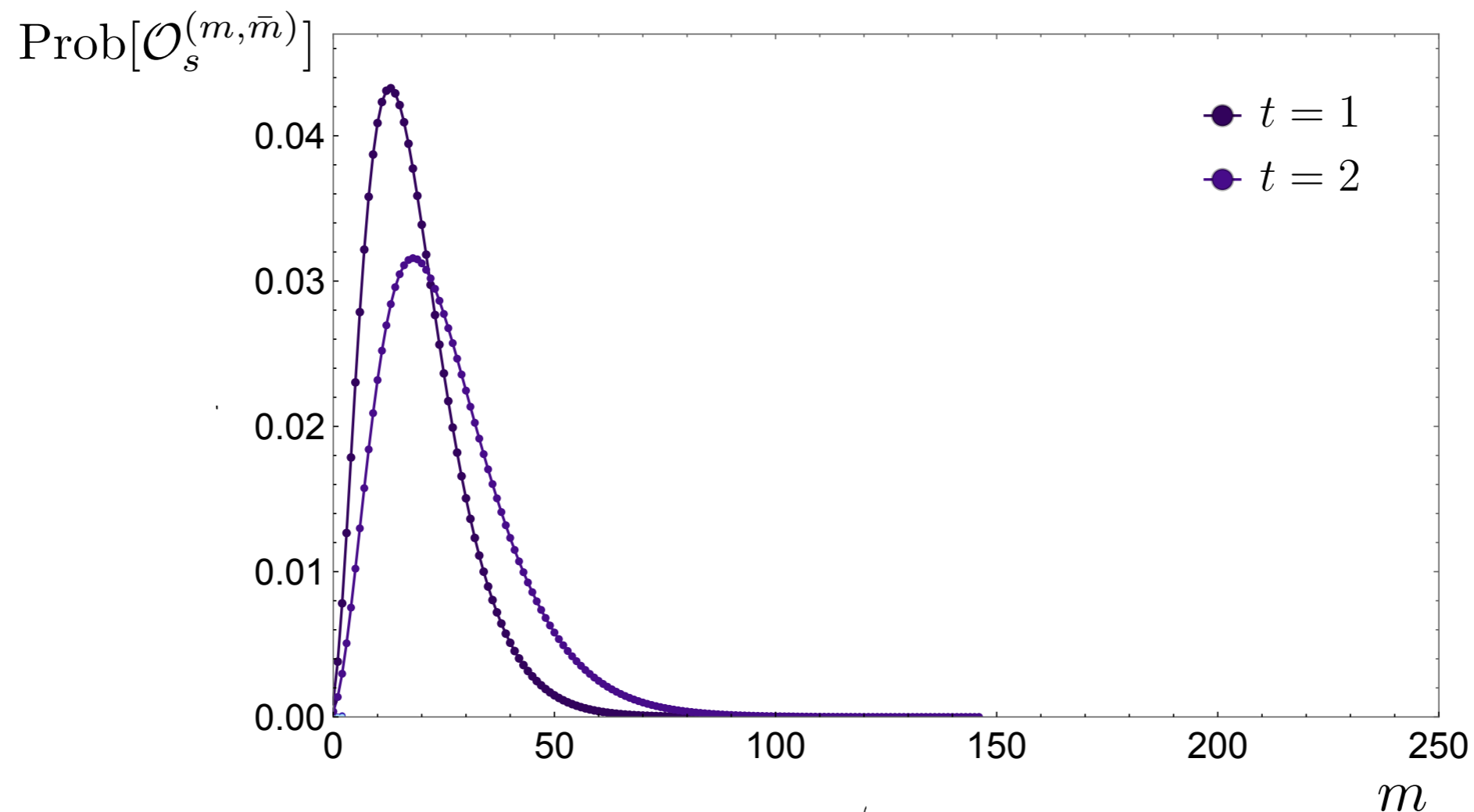
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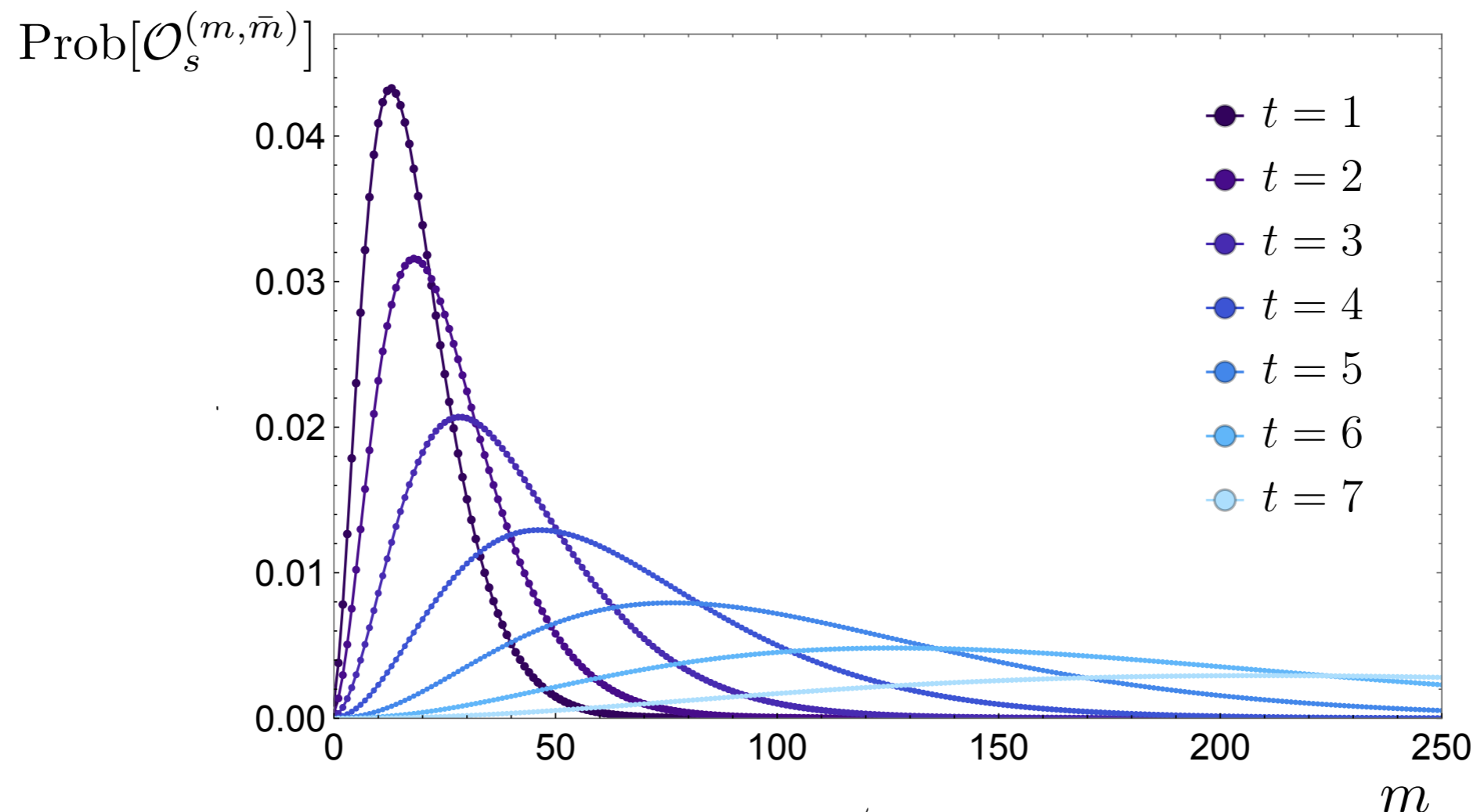
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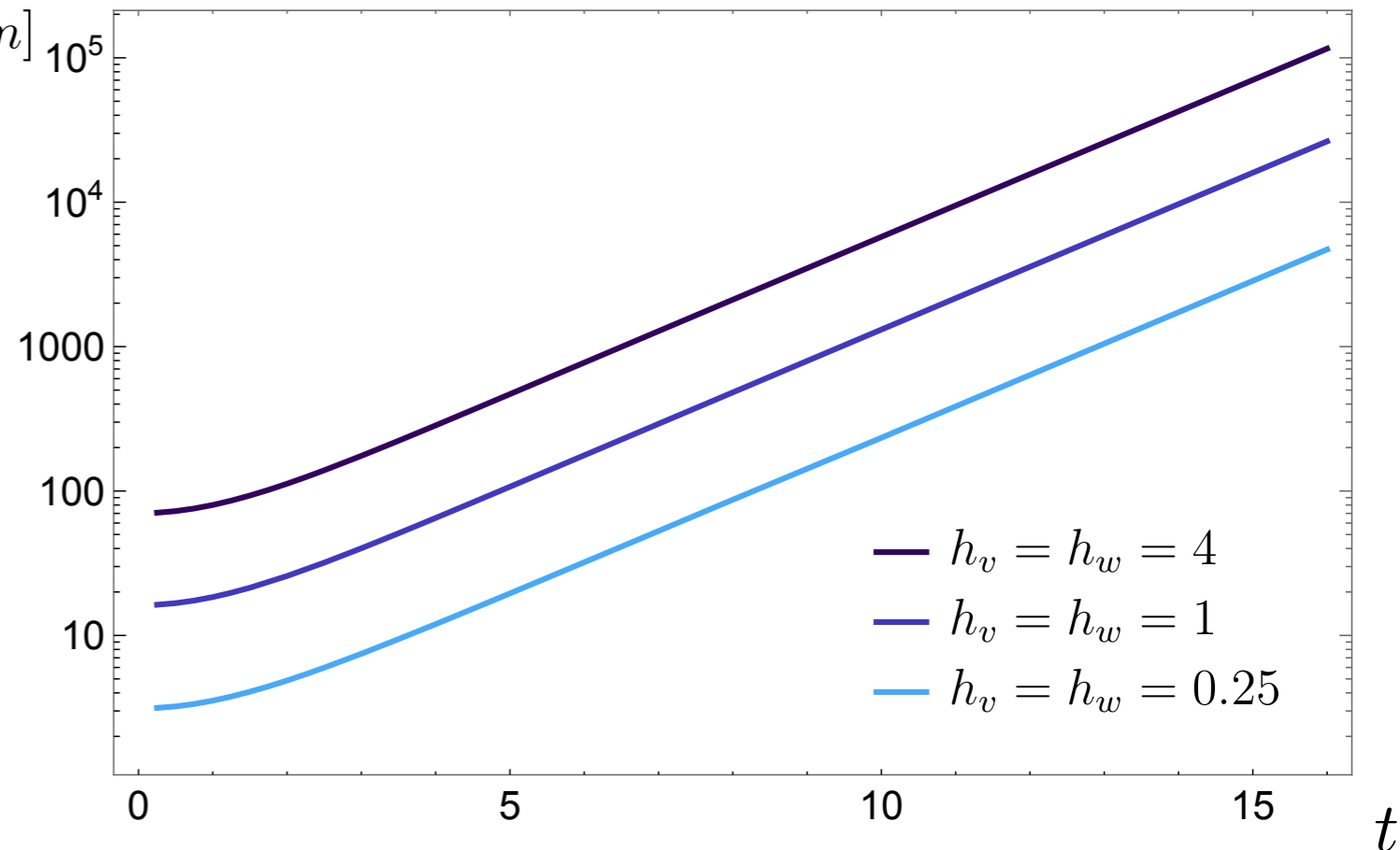
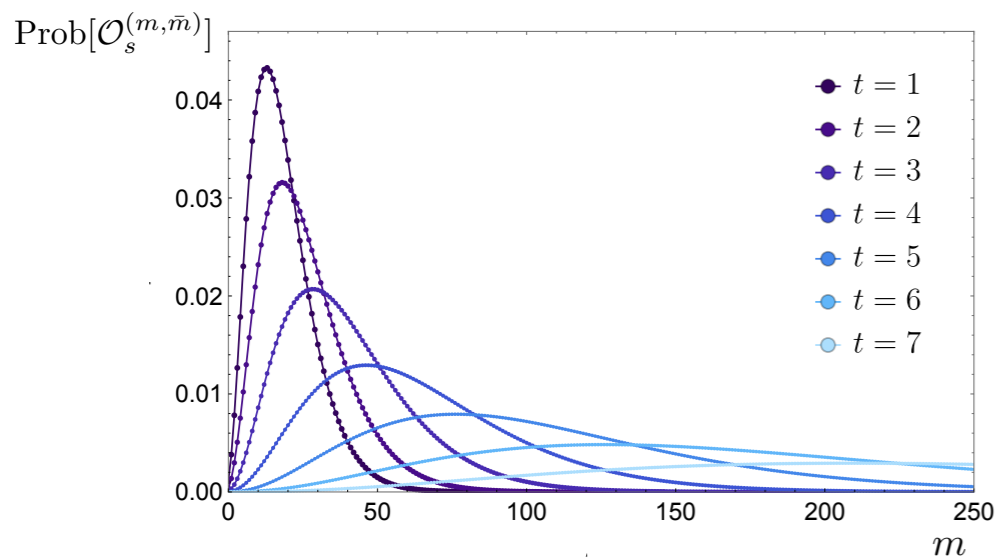
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$$\mathbb{E}[m] = \sum_m m \text{Prob}[m]$$



- ▶ Mean value of m grows exponentially: $\mathbb{E}[m] \sim e^{\frac{t-b}{2}}$
- $\mathbb{E}[\bar{m}] \sim e^{\frac{t+b}{2}}$
- $t \equiv -t_v - t_w$
- $b \equiv x_v - x_w$

$$1 = \sum_{\mathcal{O}_s} C_{wvs}^2 \underbrace{\langle \mathcal{B}_{WV\mathcal{O}_s} | \mathcal{B}_{WV\mathcal{O}_s} \rangle}_{\text{conformal blocks } \mathcal{V}_s(z, \bar{z})}$$

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 $\mathbb{E}[\bar{m}] \sim e^{\frac{t+b}{2}}$ $b \equiv x_v - x_w$

► Quantum butterfly effect / operator growth!

► $W \times V \rightarrow \mathcal{O}_s^{(m, \bar{m})}$: increasingly heavy & broad superpositions of $\mathcal{O}_s^{(m, \bar{m})}$

► Exponent = $\frac{1}{2} \times$ [maximal chaos exponent]

—> Expected for these kind of observables

—> Previously seen in 6-point OTOCs and quantum circuit models

$$1 = \sum_{\mathcal{O}_s} C_{wvs}^2 \underbrace{\langle \mathcal{B}_{WVO_s} | \mathcal{B}_{WVO_s} \rangle}_{\text{conformal blocks } \mathcal{V}_s(z, \bar{z})}$$

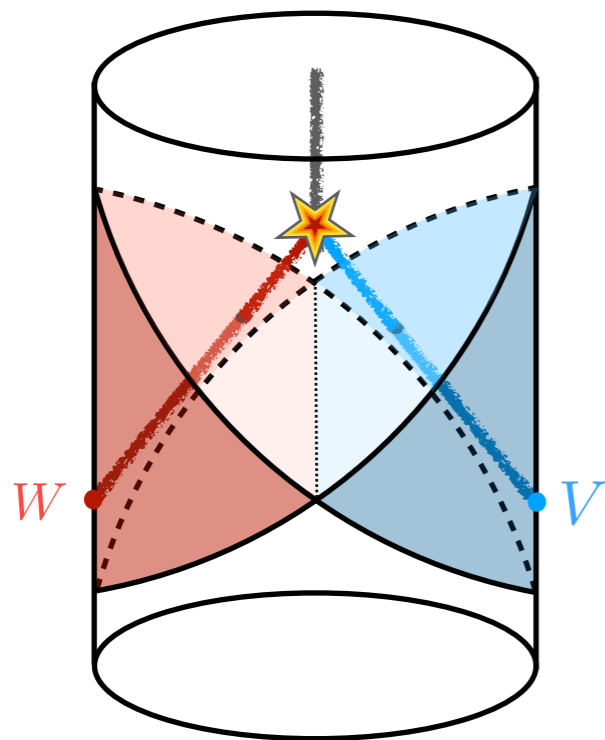
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$$\mathbb{E}[\bar{m}] \sim e^{\frac{t+b}{2}} \quad b \equiv x_v - x_w$$

► Fusion into discretum of light states **breaks down** when:



$$\mathbb{E}[m] \sim \mathbb{E}[\bar{m}] \sim \frac{c}{24}$$

$$\Leftrightarrow t - |b| \sim 2 \times \log c \quad (2 \times \text{“scrambling time”})$$

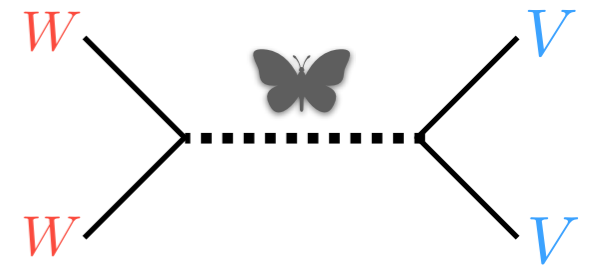
► In AdS₃/CFT₂, this is the threshold for **black hole formation**

► Shockwave collision creates black hole with mass $M = \mathbb{E}\left[\frac{m+\bar{m}}{2}\right] - \frac{c}{12}$ and spin $J = \mathbb{E}[m - \bar{m}]$

Summary

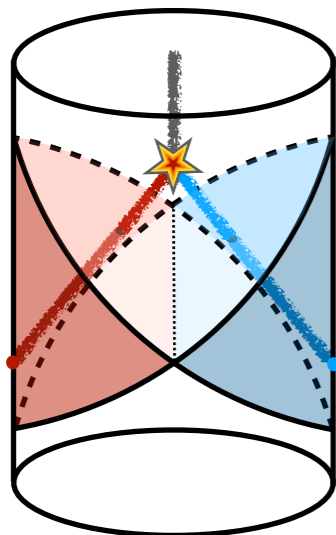
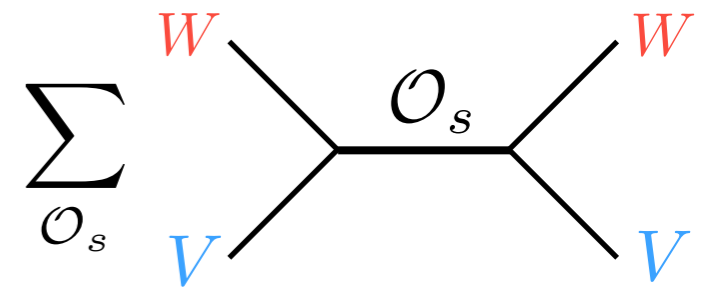
- ▶ The ‘EFT of quantum chaos’ is a powerful and useful paradigm

- ▶ Outlook: develop for more general theories (e.g. CFTs) and incorporate into wider framework of quantum chaos



- ▶ ‘Two-shockwave state’ $|\Psi_{W_L V_R}\rangle = W(-t_w - i\pi, x_w) V(t_v, x_v) |0\rangle$

- ▶ OPE decomposition depends on $e^{(t\pm b)/2}$
- ▶ Qualitative change after (2x) scrambling time
- ▶ New & precise notion of operator growth



- ▶ Holographic interpretation:
black hole formation via shockwave collision
- ▶ Outlook: mean-field theory of ‘heavy’ black hole states