

Axiverse Baryogenesis[†]

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*Based on : 2512.20696, 2606.SOON
With : David Cyncynates, Stefania Gori*

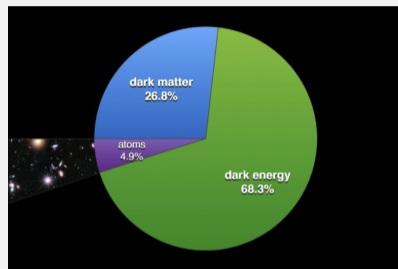
Talk Presented @ PHENO 2026

[†] With Color Commentary

May 11, 2026

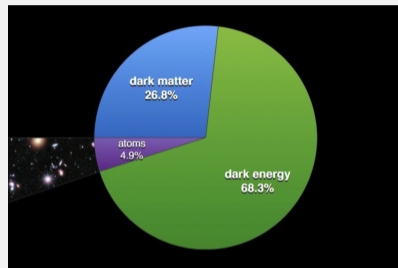
The Matter Coincidence Problem

- A curious pattern



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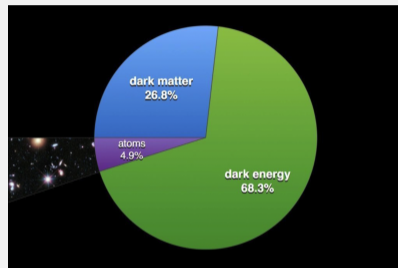
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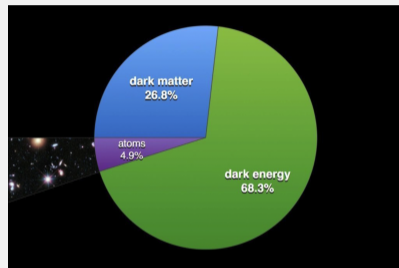
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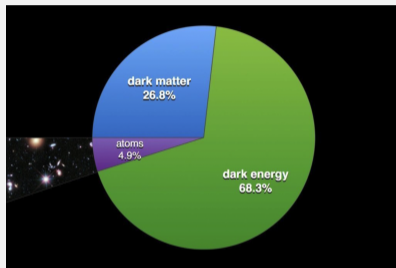
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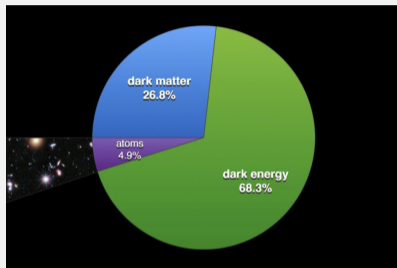
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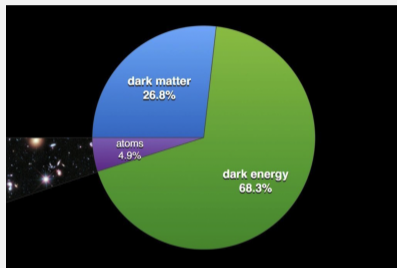
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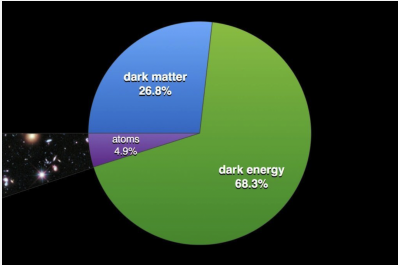
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- | Component | Percentage |
|-------------|------------|
| dark energy | 68.3% |
| dark matter | 26.8% |
| atoms | 4.9% |
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 - $n_\gamma/n_b \sim 10^{10}$: even interaction wont guarantee similar abundances.
 - Strongest hint for DM-SM interaction. Solutions can drive searches.
 - Axion is a motivated DM candidate. An axion cogenesis model?^{††}

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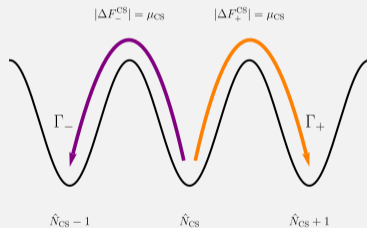
^{††} Proposal: An axion's main competitor should be called a *Nelson*.

Contributions to Fermion Asymmetries

We write down the Boltzmann equation for SM fermion asymmetries, in the presence of time-varying QCD axion θ .

Contributions to Fermion Asymmetries

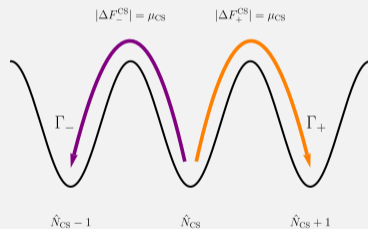
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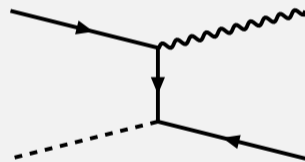
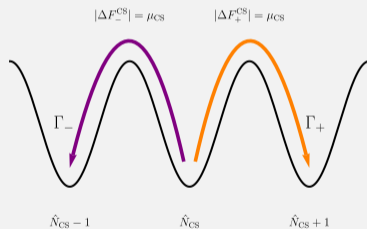
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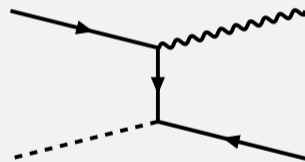
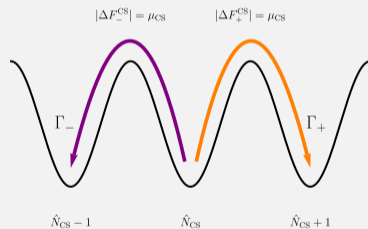


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- 2-to-2 scatterings with Higgs.

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- 2-to-2 scatterings with Higgs.
- Only act as wash-out terms.

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Generating this velocity requires PQ violation: quality problem.

Asymmetries Backreaction on Axion: Friction

$$\ddot{\theta} + \left(3H + \frac{\Upsilon_c}{f_a^2} \right) \dot{\theta} = 0$$

- Sphaleron and fermion asymmetries give rise to a friction on the axion.

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- $\dot{\theta}(T_{EW}) \sim 5$ keV: sustained to low T and **overproduces axion DM**.
- **A new confining sector can have the suitable friction!**^{††}

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^{††} Conjecture: There is no problem in life that can't be solved by a new confining force.

Axiogenesis in the Axiverse

$$\mathcal{L} \supset \frac{\alpha_c}{8\pi} \tilde{G}_c G_c \quad + \quad \frac{\alpha_d}{8\pi} \tilde{G}_d G_d \quad + V_{P/Q}$$

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$$\mathcal{L} \supset \frac{\alpha_c}{8\pi} \tilde{G}_c G_c (\theta_1 + \theta_2 + \theta_3) + \frac{\alpha_d}{8\pi} \tilde{G}_d G_d (\theta_2 + \theta_3) + V_{P/Q}(\theta_3)$$

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- Solves the strong CP problem + a cogenesis model.

Evolution below T_{EW}

We solve EOMs numerically. Evolution below T_{EW} : (θ_3 identified with θ_2).

$$0 = \ddot{\theta}_1 + 3H\dot{\theta}_1$$

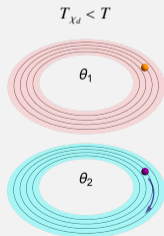
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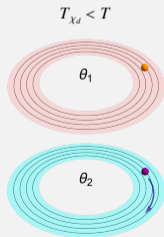


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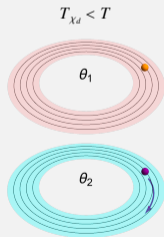


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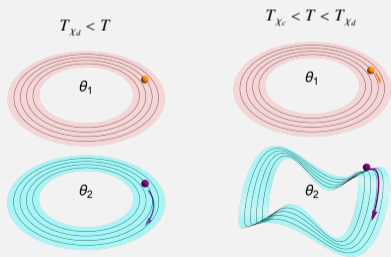


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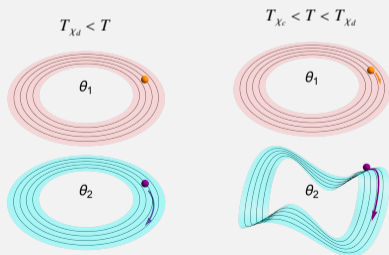
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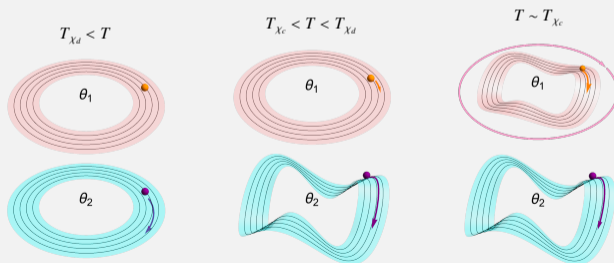
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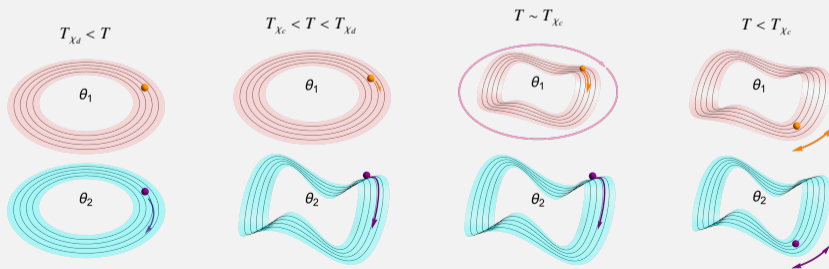
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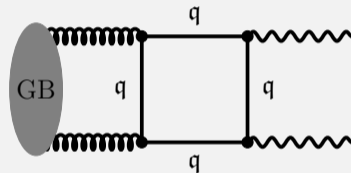
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We Further Need Dark Quarks

- Have to deplete the energy from the dark glueballs.

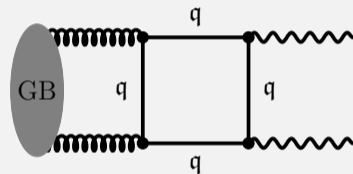
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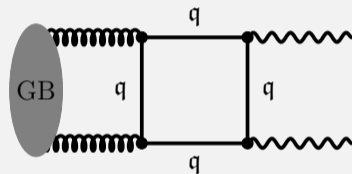
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 $m_\Psi \in (0.1, 10)$ TeV thanks to collider and BBN bounds.[†]



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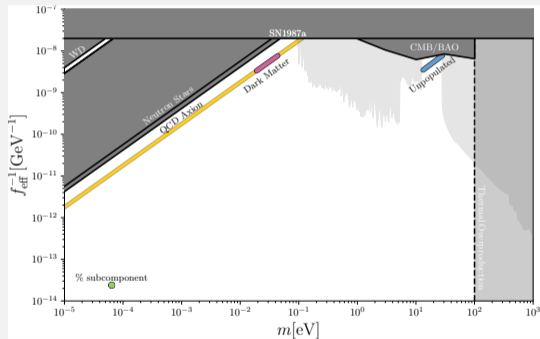
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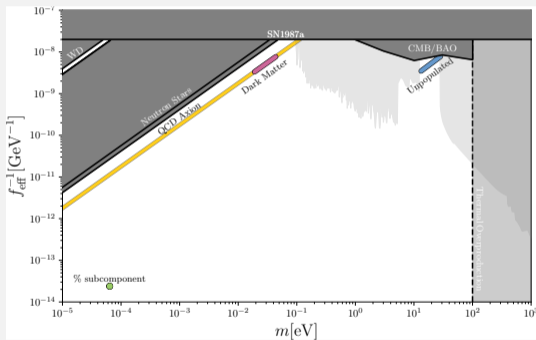
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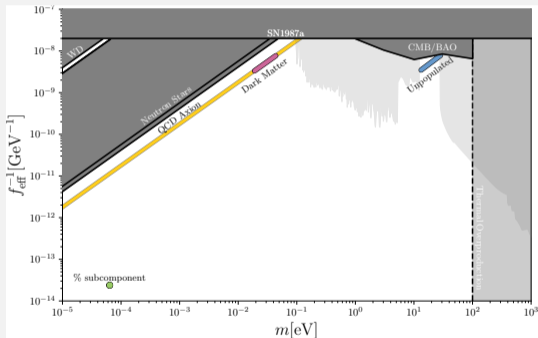
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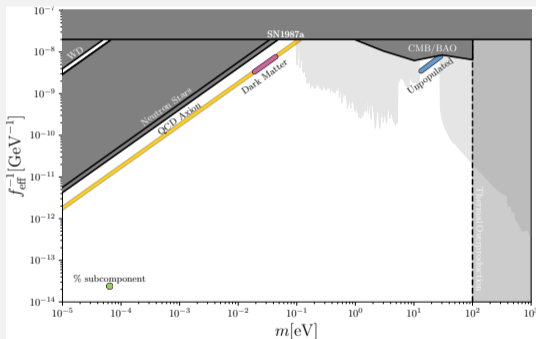
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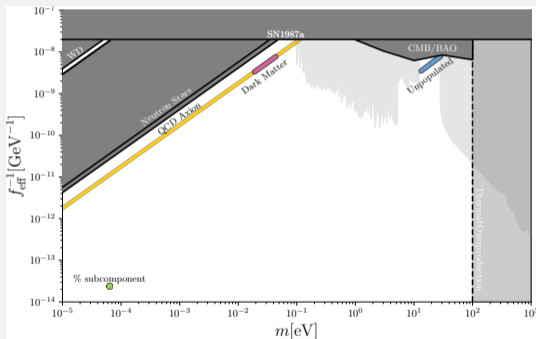
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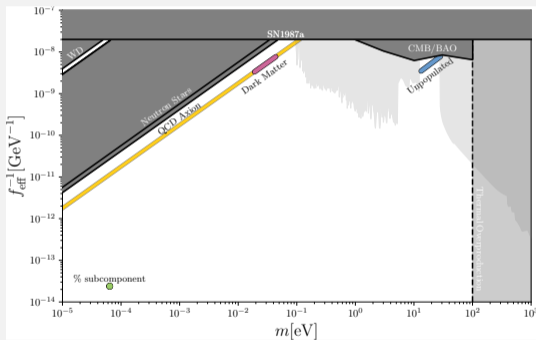
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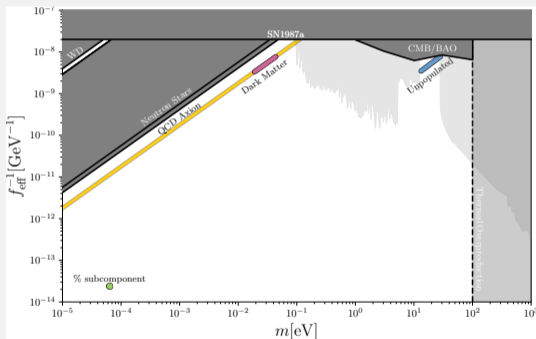
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- We have signals at colliders too.



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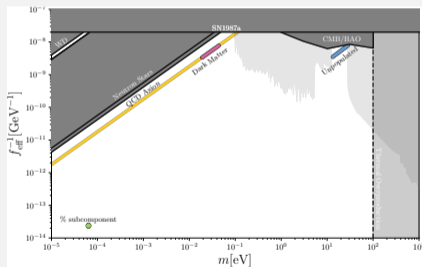
A vector-like confining sector and multiple axions from the axiverse solve the quality and the overclosure problem of the minimal axiogenesis model.[†]

[†] We have a UV completion that solves the flatness and fragmentation problems as well.

Summary

Coincidence problem is a strong guide for BSM model-building.

A vector-like confining sector and multiple axions from the axiverse solve the quality and the overclosure problem of the minimal axiogenesis model.†

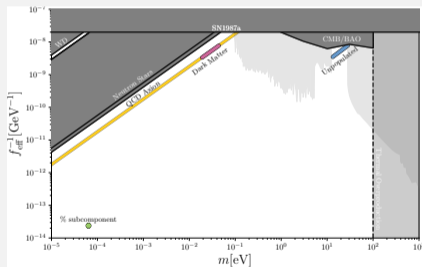


† We have a UV completion that solves the flatness and fragmentation problems as well.

Summary

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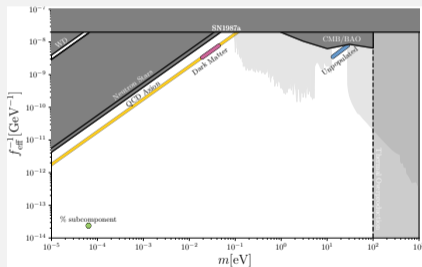
Defines new correlated targets for experiments (in collider and cosmology).

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Summary

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THANK YOU!

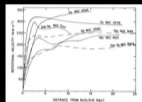
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Back up

- Evidence For DM
- Coincidence vs Cogenesis
- $\dot{\theta}$ as CS Chemical Potential
- Fermion Contribution to Free Energy
- Finite T Effects in Chirality flipping
- Friction and Kinetic Misalignment
- Susceptibility
- Quarks and Glueballs
- Other Scenarios
- UV Completion

Evidence for DM

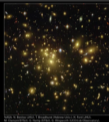
1. Rotation Curves



Bulge, Ford & Tremaine 1976

What we learn:
mass fraction
distribution

2. Cluster Dynamics

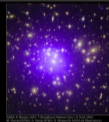


What we learn:
mass fraction
distribution



Clowe 1997

3. Cluster Gas



What we learn:
mass fraction
distribution

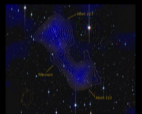
~90% of the luminous
matter in a cluster is
hot gas

4. Strong Gravitational Lensing



What we learn:
mass fraction
distribution

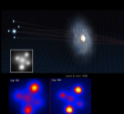
5. Weak Gravitational Lensing



Clowe et al. 2006

What we learn:
distribution
shape
structure

6. Cosmological Microlensing

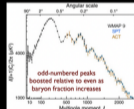


What we learn:
mass fraction
smoothness



Johnston & Braxton 2003

7. CMB Acoustic Peaks

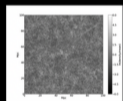


Hirshbuhl et al. 2012

What we learn:
ratio of DM/
collisional
matter
thermal history

odd-numbered peaks
boosted relative to even as
baryon fraction increases

9. Large Scale Structure



Peebles 1981, Tinker Simulation

What we learn:
ratio of DM/
collisional
matter
thermal history

Excellent agreement
between simulations
and galaxy distribution
on the largest scales

10. Galaxy/Cluster Collisions

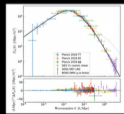


MADAU et al. 2008

What we learn:
distribution
separation from
collisional
matter
self-interaction

Difficult to explain
without
collisionless matter

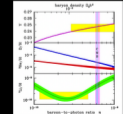
8. Matter Power Spectrum



Challinor et al. 2019

What we learn:
ratio of DM/
collisional
matter
thermal history

11. Big Bang Nucleosynthesis

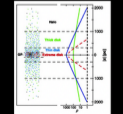


PDG 2018

What we learn:
amount of
baryonic matter

Remaining mystery:
lithium abundance
(that still needs low
baryon fraction)

12. Local Stellar Motions



Bulge 2000

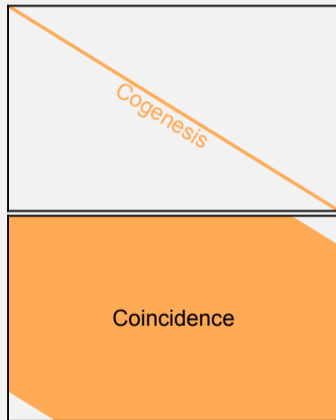
What we learn:
local dark
matter density

Estimate:
 $\rho_{DM} \sim 0.3 \text{ GeV/cm}^3$
 $\sim 0.008 \text{ M}_{\odot}/\text{pc}^3$

via Katie Mack (Aspen Center for Physics Colloquium 2019)

Coincidence vs. Cogenesis

- Coincidence models \neq Cogenesis models!
 - Is the ratio of abundances *generically* $\mathcal{O}(1)$?
- We put forward a cogenesis mechanism - future step: coincidence.
- Do any of motivated DM models fit in a cogenesis setup?
 - Let's explore this in QCD axion DM models.

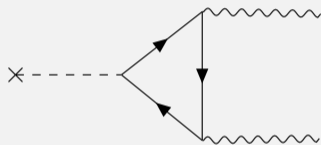


CS Chemical Potential

$$\begin{aligned} \mathcal{S} &= \int d^4x \mathcal{L} \supset \int d^4x \frac{\alpha}{8\pi} \zeta W_{\mu\nu}^a \tilde{W}^{\mu\nu a} \\ &= \int d^4x \zeta \partial_\mu j_{\text{CS}}^\mu \\ &= \int d^4x \zeta \partial_t j_{\text{CS}}^0 \\ &= - \int d^4x \partial_t \zeta j_{\text{CS}}^0 . \end{aligned}$$

Fermion Asymmetry and Free Energy

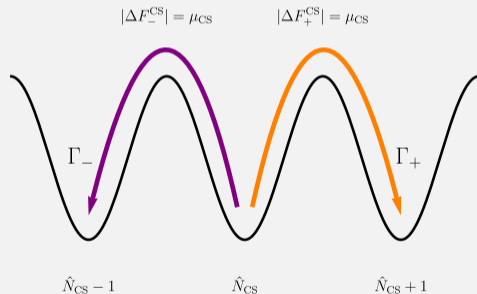
$$\partial_t N_{\text{CS}} = \Gamma_+ - \Gamma_-, \quad \frac{\Gamma_+}{\Gamma_-} = e^{-\Delta F/T}, \quad \Gamma_{\text{sph}} \equiv \frac{\Gamma_+ + \Gamma_-}{V},$$



$$\partial_\mu j_i^\mu = \mathcal{N}_i^G \frac{\alpha}{8\pi} G_{\mu\nu}^a \tilde{G}^{\mu\nu, a} \quad \delta n_i = \frac{\mathcal{N}_i^G}{V} \times \delta N_{\text{CS}}$$

$$F_i \sim \int \frac{d^3 k}{(2\pi)^3} \left[\log \left(1 + e^{-(E-\mu_i)/T} \right) + (\mu_i \rightarrow -\mu_i) \right]$$

$$F_i \supset \frac{3V}{g_i T^2} n_i^2 \implies \Delta F_i = 12 \mathcal{N}_i^G \frac{n_i}{g_i T^2}$$



Boltzmann Equations for Asymmetries

- $\dot{\theta}$ is the only source term.
- $n_i \propto \dot{\theta}$.

For LH quarks:

$$Y_B \sim 10^{-3} \frac{\dot{\theta}}{T} \Big|_{T \rightarrow T_{EW}} \quad \dot{n}_{q_g} + 3Hn_{q_g} = 3\kappa_{\text{ch},u} \min(\alpha_3, y_{u_g}^2) T \left(-\frac{n_{q_g}}{6} - \frac{n_{\bar{u}_g}}{3} - \frac{n_H}{4} \right)$$

$$\text{Upshot : } \dot{\theta}(T_{EW}) \approx 5 \text{ keV} \rightarrow Y_B \approx 9 \times 10^{-11}$$

$$+ 3\kappa_{\text{ch},d} \min(\alpha_3, y_{d_g}^2) T \left(-\frac{n_{q_g}}{6} - \frac{n_{\bar{d}_g}}{3} + \frac{n_H}{4} \right)$$

$$+ \frac{9}{2} \frac{\Gamma_w}{T^3} \left(-\frac{c_W}{3} \dot{\theta} T^2 - \sum_{g'} (n_{q_{g'}} + n_{\ell_{g'}}) \right)$$

$$+ 2 \frac{\Gamma_c}{T^3} \left(-\frac{1}{2} \dot{\theta} T^2 - \sum_{g'} (n_{q_{g'}} + n_{\bar{u}_{g'}} + n_{\bar{d}_{g'}}) \right)$$

$$q_g \rightarrow (\bar{u}_g, \bar{d}_g, \bar{e}_g, \ell_g)$$

This velocity will be generated by PQ violation in the UV. In this talk we focus on IR dynamics.

Friction and Kinetic Misalignment

$$\frac{\Omega_a}{\Omega_{\text{DM}}} \approx \frac{\chi_c(0)}{\frac{5}{12} \frac{\pi^2}{30} g_*(T_{\text{eq}}) T_{\text{eq}}^4} \frac{\omega_{\text{ESFO}}}{m_a} \frac{\pi^2}{3} \left(\frac{g_*(T_{\text{eq}})}{g_*(T_{\text{ESFO}})} \right)^{3/4} \left(\frac{T_{\text{eq}}}{T_{\text{ESFO}}} \right)^3 \approx 345 \left(\frac{f_a}{5 \times 10^7 \text{ GeV}} \right),$$

$$\dot{\Theta}^2 = \omega_{\text{ESFO}}^2 \exp \left\{ - \int_{t_{\text{ESFO}}}^t d \log t' \gamma_{\text{roll}}(t') \right\}, \quad \frac{1}{2} f_a^2 \dot{\Theta}^2 > \chi_c,$$

$$\gamma_{\text{roll}} = 3 + \frac{\Upsilon_c}{H f_a^2}.$$

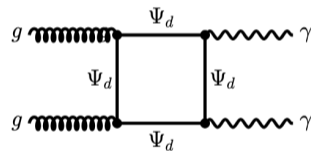
Susceptibility

$$\chi_{c,d}(T) \approx \chi_{c,d}(0) \left(\frac{\Lambda_{c,d}}{T} \right)^{\frac{11}{3}N_{c,d}-4} \prod_{m_{q_{c,d}} < T} \left(\frac{\max\{m_{q_{c,d}}, \Lambda_{c,d}\}}{T} \right)^{1/3}, \quad T > \Lambda_{c,d},$$

Quarks and Glueballs

Quark	SM Charges	Principal Role
ψ_d	(1,1,0)	Absorbing θ_2 overabundance
Ψ_d	(1,1, \mathbf{q}_d)	Depleting glueballs abundance

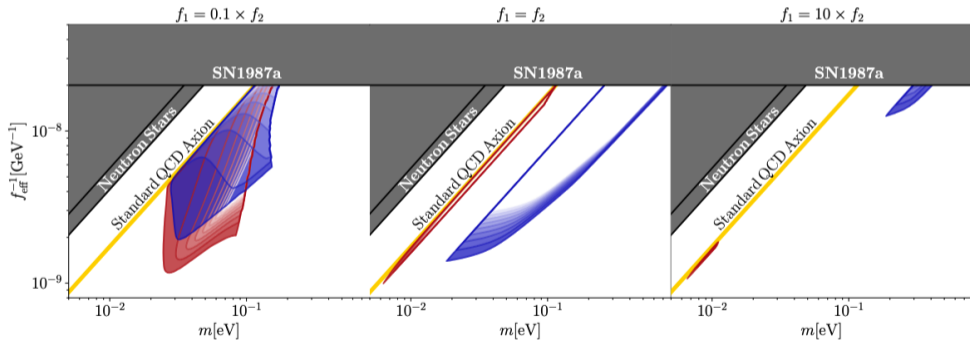
$$\tau_{0^{++}} \simeq 3600 \frac{32\pi}{(N_d^2 - 1) \mathbf{q}_d^4 \alpha_{\text{em}}^2 \alpha_d^2 (m_{0^{++}})} \left(\frac{m_\Psi^3}{F_{0^{++}}^S} \right)^2 \left(\frac{m_\Psi}{m_{0^{++}}} \right)^2 \frac{1}{m_{0^{++}}},$$



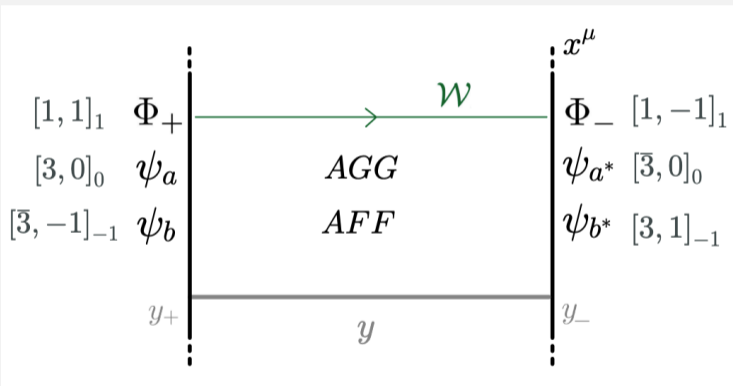
$$\tau_{0^{++}} \simeq 10^{-22} \text{ s} \times \left(\frac{m_\Psi}{\Lambda} \right)^8 \times \left(\frac{1 \text{ GeV}}{\Lambda} \right)$$

Other IR Scenarios

We can have a different boundary condition at T_{EW} .



Speculation on a High-Quality UV Completion

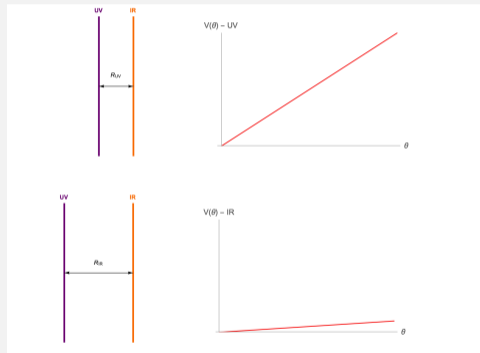


$$\mathcal{W} = \Phi_+ e^{ig_d q \int_{y_+}^{y_-} dy A_y} \Phi_-$$

$$V_{PQ} = \frac{(\mu_- \mu_+)^{\frac{3}{2}}}{2M} e^{-ML} \mathcal{W} + \text{h.c.} + \dots$$

$$V = \chi_c (1 - \cos(\theta_1 + \theta_2)) + \chi_d (1 - \cos(\theta_2)) + V_{PQ}(\theta_2)$$

Speculation on a High-Quality UV Completion



- $V_{PQ}^{UV} \propto e^{-MR_{UV}} \sim 1$: large PQ breaking.
- R_{UV} is away from the equilibrium point of the Goldberger-Wise mechanism.
- Branes are pushed to the GW equilibrium at R_{IR} .
- $V_{PQ}^{IR} \propto e^{-MR_{IR}} \ll 1$: negligible PQ breaking.