

Exotic Higgs Decays at a Muon Collider

Yanhan Wang

Brown University

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JiJi Fan, Lingfeng Li, Tao Liu, **Yanhan Wang**, Mingrui Zhou

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Motivation

Why exotic Higgs decays?

- The Higgs can act as a portal to light BSM particles.
- Exotic Higgs decays appear in scenarios related to: naturalness, dark matter, and electroweak phase transition.
- Comprehensive study and review: D. Curtin et al., Exotic decays of the 125 GeV Higgs boson, Phys. Rev. D 90 (2014) 075004 [1312.4992], M. Cepeda, S. Gori, V.M. Outskoorn and J. Shelton, Exotic Higgs Decays, 2111.12751.

Motivation

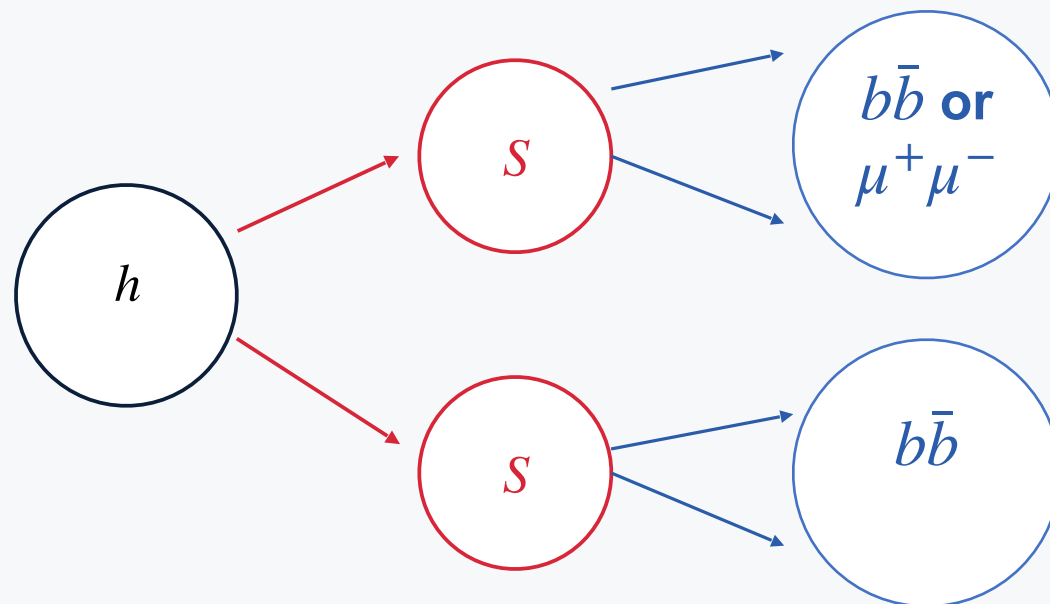
Can a future muon collider improve the sensitivity to exotic Higgs decays?

Muon collider advantages:

- cleaner than hadron colliders
- suppressed synchrotron radiation and reach higher centre of mass energy, compared to electron colliders
- large vector boson fusion (VBF) Higgs production at high energy
- potentially better sensitivities to hadronic final states

Model : SM + real singlet scalar S

$$V(H, S) = -\mu^2 |H|^2 + \lambda |H|^4 + \frac{1}{2} a_1 |H|^2 S + \frac{1}{2} a_2 |H|^2 S^2 + b_1 S + \frac{1}{2} b_2 S^2 + \frac{1}{3} b_3 S^3 + \frac{1}{4} b_4 S^4.$$



After electroweak symmetry breaking, h and S mix. In the small-mixing limit, the SM-like Higgs can decay exotically to the singlets S . The singlets S inherit Higgs-like couplings to SM particles. Therefore, the branching fractions of S follow those of a SM Higgs boson with mass m_S .

Production and final states

VBF production dominates at high-energy muon colliders

Muon collider benchmarks

3 TeV, 1 ab^{-1}
10 TeV, 10 ab^{-1}

Dominant Higgs production:
 $\mu^+ \mu^- \rightarrow \nu_\mu \bar{\nu}_\mu h$ via W-boson
fusion

$\sigma \approx 1 \text{ pb}$ at $\sqrt{s} = 10 \text{ TeV}$

Two analysis channels

$h \rightarrow SS \rightarrow 4b$

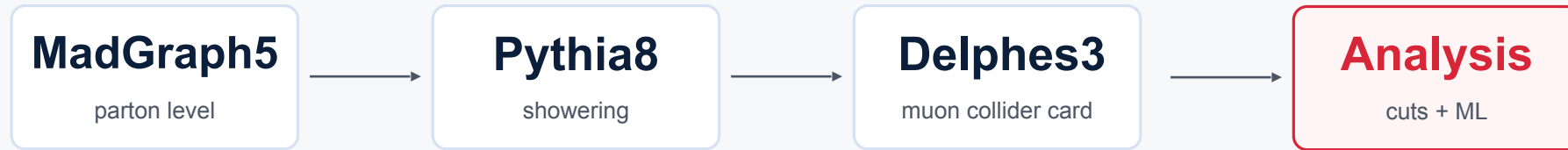
largest Higgs-portal rate
challenging jet combinatorics

$h \rightarrow SS \rightarrow 2b2\mu$

clean dimuon resonance
smaller Higgs-portal rate

Simulation and object reconstruction

From partons to detector-level analysis



Generator-level cuts:

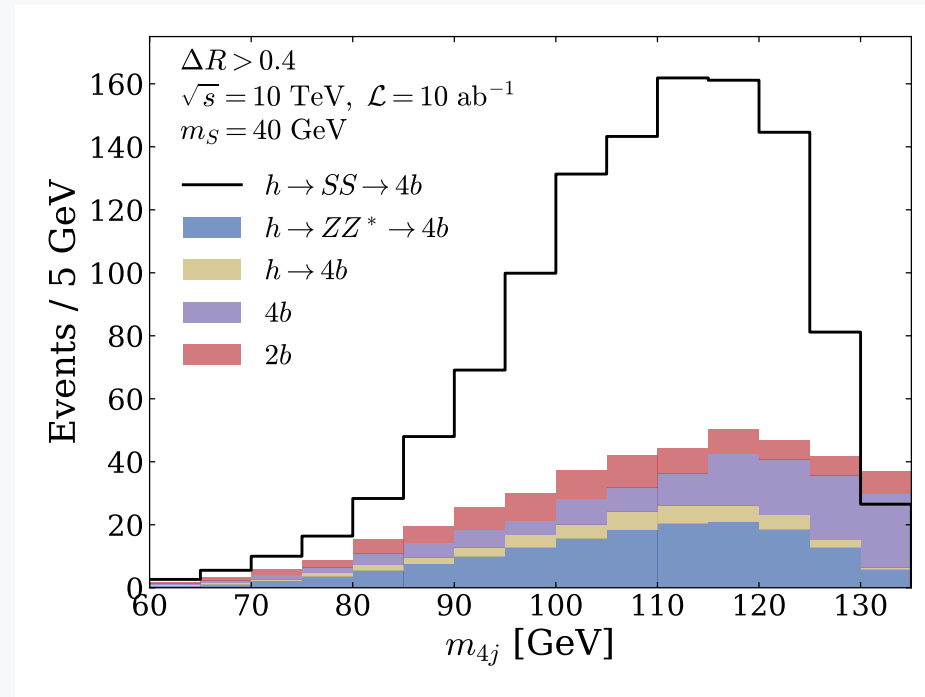
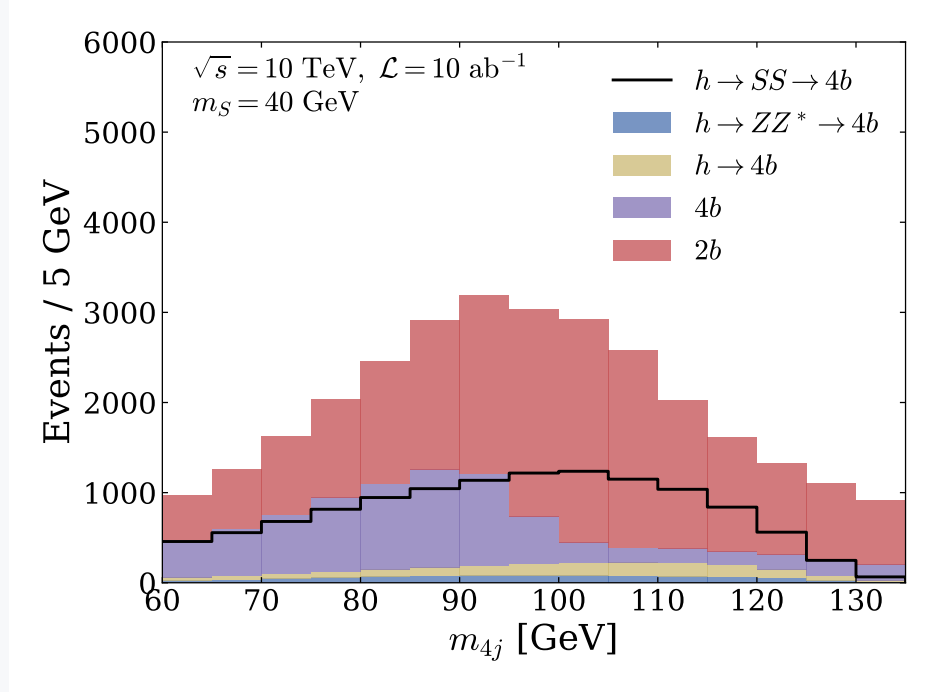
- b quarks: $p_T > 15$ GeV, $|\eta| < 2.5$, $\Delta R_{bb} > 0.25$
- muons: $p_T > 0.5$ GeV, $|\eta| < 8$, $\Delta R_{\mu\mu} > 0.1$

Detector-level choices

- muons selected with $|\eta_\mu| < 2.5$
- b-tagging working point: 70% flat efficiency
- b-flavor matching cone: $\Delta R < 0.3$

Preselection: resolve objects first

ΔR cuts suppress collinear radiation and wrong b -tagging



Preselection cuts

$$p_T^{\text{jets}} > 20 \text{ GeV}; p_T^\mu > 5 \text{ GeV}$$

$$\Delta R_{jj} > 0.4, \text{ and for } 2b2\mu \text{ also } \Delta R_{\mu\mu}, \Delta R_{\mu j} > 0.4$$

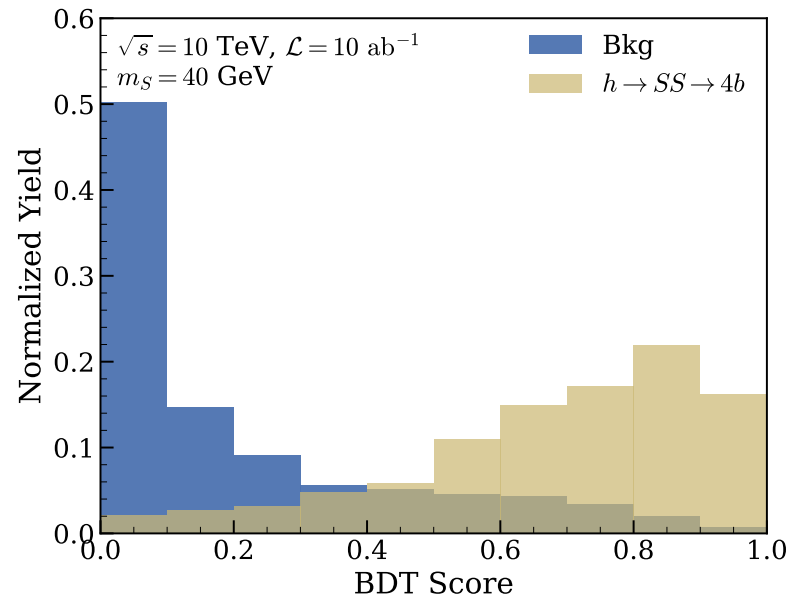
$$\text{Higgs window: } m_{4b} \text{ or } m_{2b2\mu} \in [100, 150] \text{ GeV}$$

Key effect in $4b$

After $\Delta R_{jj} > 0.4$, the m_{4j} peak moves closer to the Higgs mass, and reducible backgrounds drop significantly.

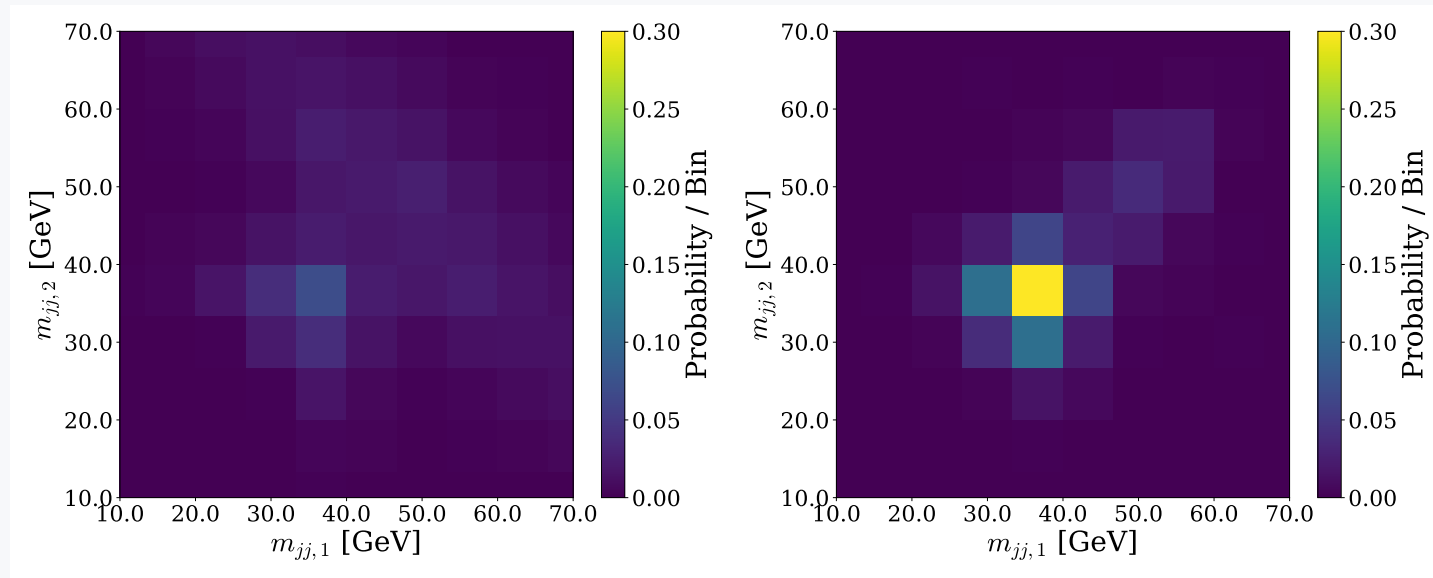
Machine learning selection

BDT learns resonance structure and jet pairing



Method

XGBoost boosted decision tree (BDT), 5 parallel models; select the threshold that maximizes S/\sqrt{B} for each m_S benchmark.



What the BDT uses

$4b$: the minimum mass difference to the m_S benchmark $|m_{jj} - m_S|$.

$2b2\mu$: invariant mass of jet pairs m_{jj}

After ML, the two reconstructed S candidates concentrate near $m_S = 40 \text{ GeV}$

4b channel after ML

$$m_S = 40 \text{ GeV}, \sqrt{s} = 10 \text{ TeV}, \mathcal{L} = 10 \text{ ab}^{-1}$$

Process	σ [pb]	Preselection	ML selection	4b-tagging	Yield
Signal					
$h \rightarrow SS \rightarrow 4b$	$0.84 \times \text{BR}$	1.0×10^{-2}	7.0×10^{-3}	1.9×10^{-3}	$1.3 \times 10^4 \times \text{BR}$
Background					
$h \rightarrow ZZ^* \rightarrow 4b$	5.0×10^{-4}	2.3×10^{-2}	6.6×10^{-3}	1.4×10^{-3}	7.0
$h \rightarrow 4b$	1.0×10^{-3}	2.9×10^{-3}	1.3×10^{-3}	2.9×10^{-4}	2.9
2b	2.0×10^{-2}	8.4×10^{-4}	2.7×10^{-4}	$\leq 3.8 \times 10^{-6}$	≤ 0.76
4b	8.7×10^{-3}	1.1×10^{-3}	2.5×10^{-4}	1.7×10^{-5}	1.5

Why this works: small final background + sizable 4b branching fraction

Signal yield after all cuts:

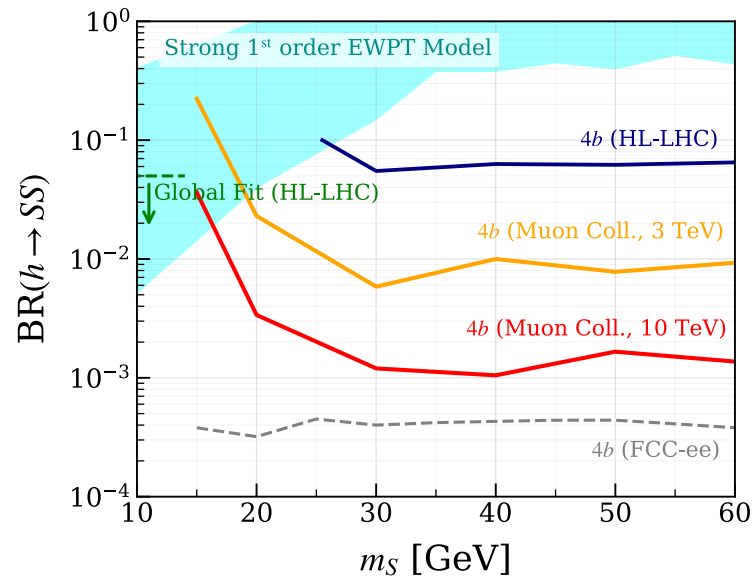
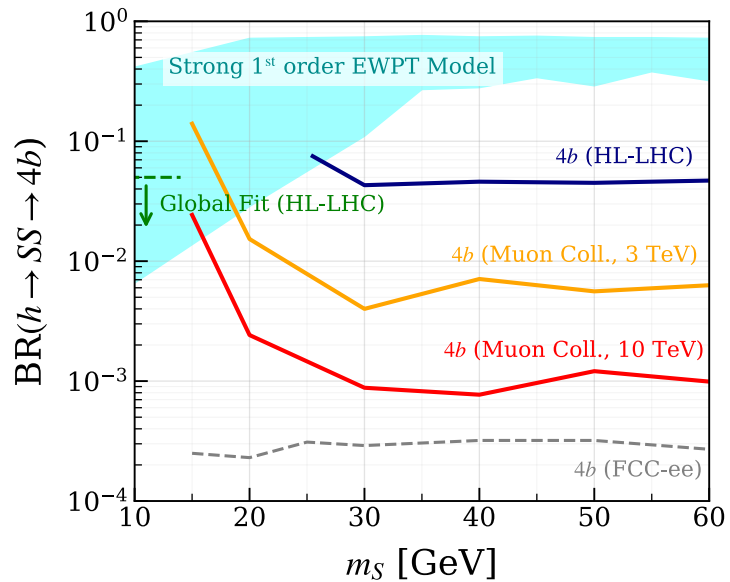
$$1.3 \times 10^4 \times \text{BR}(h \rightarrow SS \rightarrow 4b)$$

Total background: $\mathcal{O}(10)$ events

ML + 4b-tagging leaves the irreducible $h \rightarrow ZZ^* \rightarrow 4b$ background as the largest component.

Main result: $4b$ sensitivity

Best channel for the Higgs-portal benchmark



Projected 95% CL reach

$$\sqrt{s} = 10 \text{ TeV}, \mathcal{L} = 10 \text{ ab}^{-1}:$$

$$\mathbf{BR}(h \rightarrow SS) \sim 10^{-3}$$

$$\mathbf{\text{for } m_S \gtrsim 20 \text{ GeV}}$$

$$\sqrt{s} = 3 \text{ TeV}, \mathcal{L} = 1 \text{ ab}^{-1}:$$

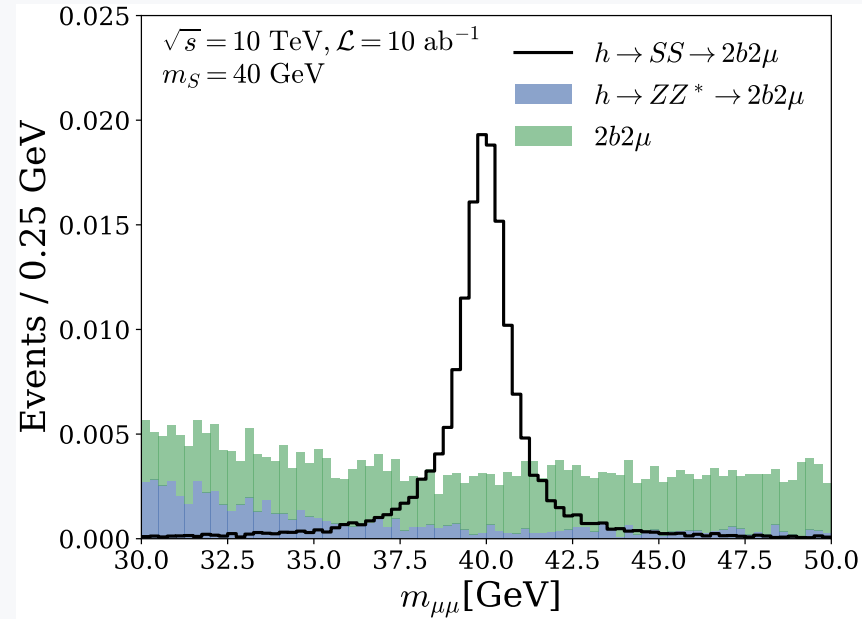
$$\mathbf{BR}(h \rightarrow SS) \sim 10^{-2} \text{ for}$$

$$\mathbf{m_S \gtrsim 20 \text{ GeV}}$$

Low- m_S sensitivity weakens because the jets become more collimated, making it difficult to get four resolved b -jets.

$2b2\mu$ channel: clean resonance

Use $m_{\mu\mu}$ only after ML to avoid classifier domination



Process	σ [pb]	Preselection	ML selection	2b-tagging	$m_{\mu\mu}$	Yield
Signal						
$h \rightarrow SS \rightarrow 2b2\mu$	$0.84 \times \text{Br}$	8.0×10^{-2}	7.8×10^{-2}	3.8×10^{-2}	3.7×10^{-2}	$3.1 \times 10^5 \times \text{Br}$
Background						
$h \rightarrow 2b2\mu$	4.4×10^{-4}	2.4×10^{-2}	4.2×10^{-3}	1.4×10^{-3}	2.2×10^{-4}	0.96
$2b2\mu$	1.2×10^{-3}	3.0×10^{-3}	9.8×10^{-4}	3.9×10^{-4}	1.0×10^{-4}	1.2

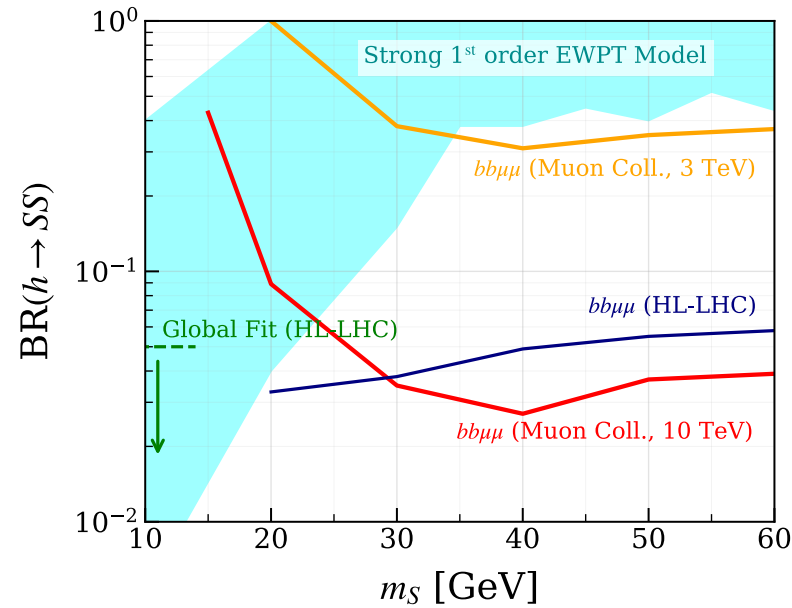
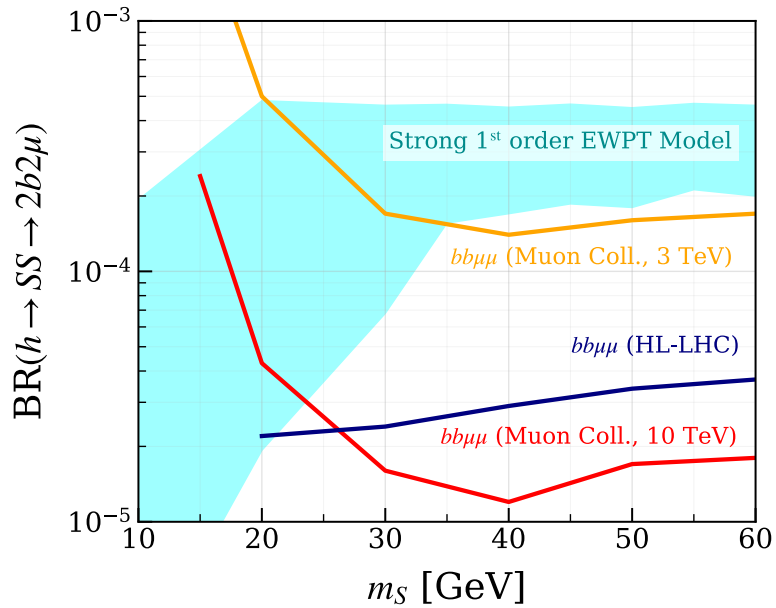
Why $2b2\mu$ helps

narrow dimuon resonance
 small final background
 strong model-independent reach

But in the Higgs-portal model,
 $S \rightarrow \mu\mu$ is rare, so this channel
 becomes rate-limited.

$2b2\mu$ sensitivity

Clean experimentally, but model-dependent rate is small



Model-independent final state

For $\sqrt{s} = 10$ TeV case can probe

$$BR(h \rightarrow SS \rightarrow 2b2\mu) \sim 10^{-5}$$

Higgs-portal model

$BR(S \rightarrow \mu\mu)$ is small, so the inferred $BR(h \rightarrow SS)$ reach is much weaker than $4b$.

Conclusions

- 1 A future muon collider can probe exotic Higgs decays with fully hadronic and semi-leptonic final states.
- 2 For $h \rightarrow SS$ in a Higgs-portal model, the $4b$ channel is the leading discovery mode.
- 3 At $\sqrt{s} = 10$ TeV with $\mathcal{L} = 10$ ab⁻¹, $4b$ reaches $\text{BR}(h \rightarrow SS) \sim 10^{-3}$ for $m_S \gtrsim 20$ GeV.
- 4 $2b2\mu$ gives a clean dimuon resonance and reaches $\text{BR}(h \rightarrow SS \rightarrow 2b2\mu) \sim 10^{-5}$ model-independently.

Thank you — questions?