

Testing RG Evolution of Neutrino Mixing with IceCube Double Bang Events

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arXiv:2511.07541, PLB 875 (2026) 140322

In Collaboration with

Carlos A. Argüelles, K.S. Babu, Vedran Brdar



Φαινó 2026

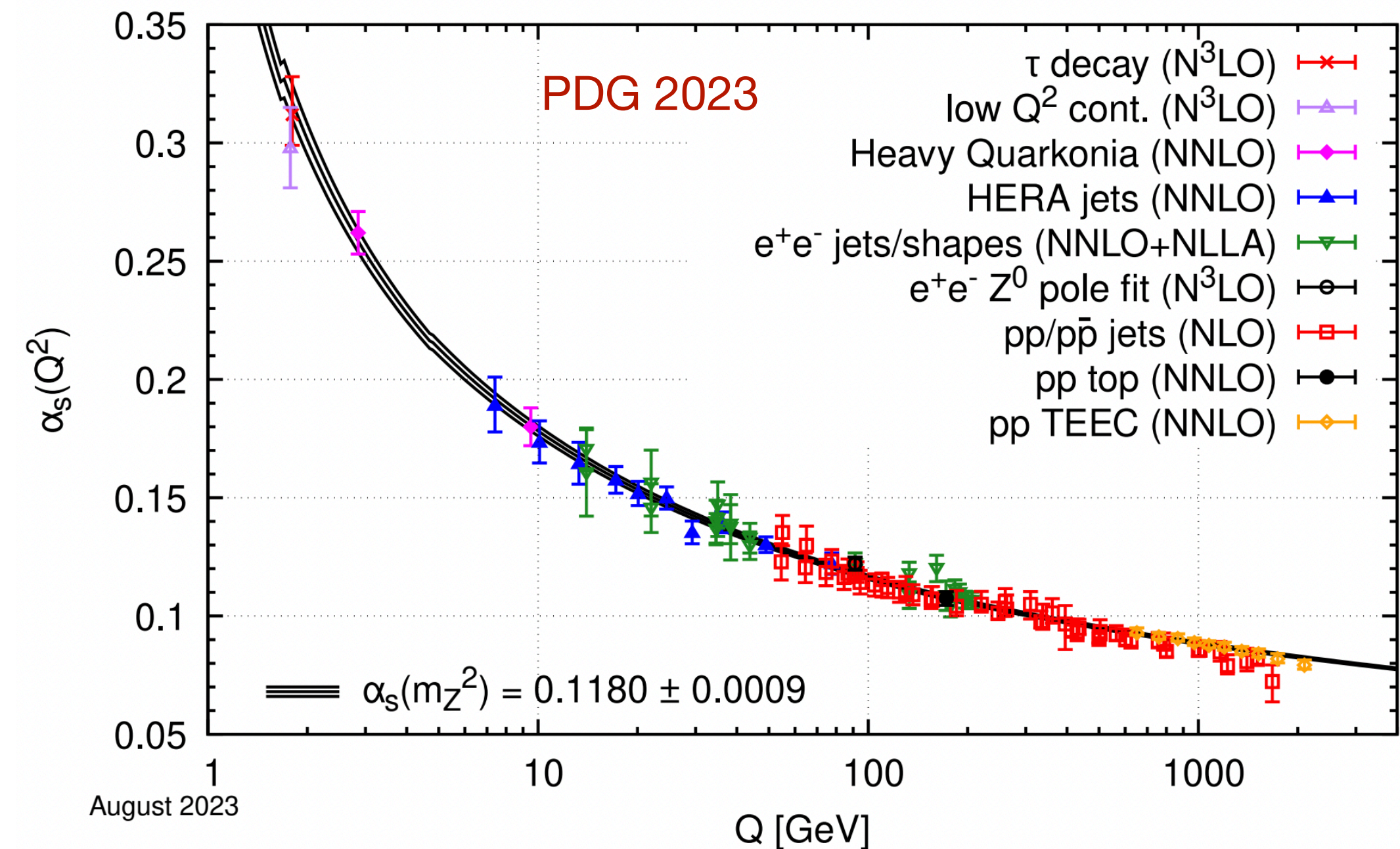
Testing Neutrino Mass Models

- There are many neutrino mass models to explain the origin of neutrino mass Mohapatra, Senjanovic, PRL1980; Zee PLB1980; Babu PLB1988
- The test of any model can be –
 - Unique signature of the new fields
e.g. single/double charged scalar
 - Direct detection of the new fields e.g. heavy sterile neutrino decaying to SM, detectable particles

The neutrino parameters can be subject to observable Renormalization Group (RG) running in a BSM scenario

Revisiting RG running:

Renormalization Group Evolution (RGE) effects are the changes of quantum field theory parameters with energy.



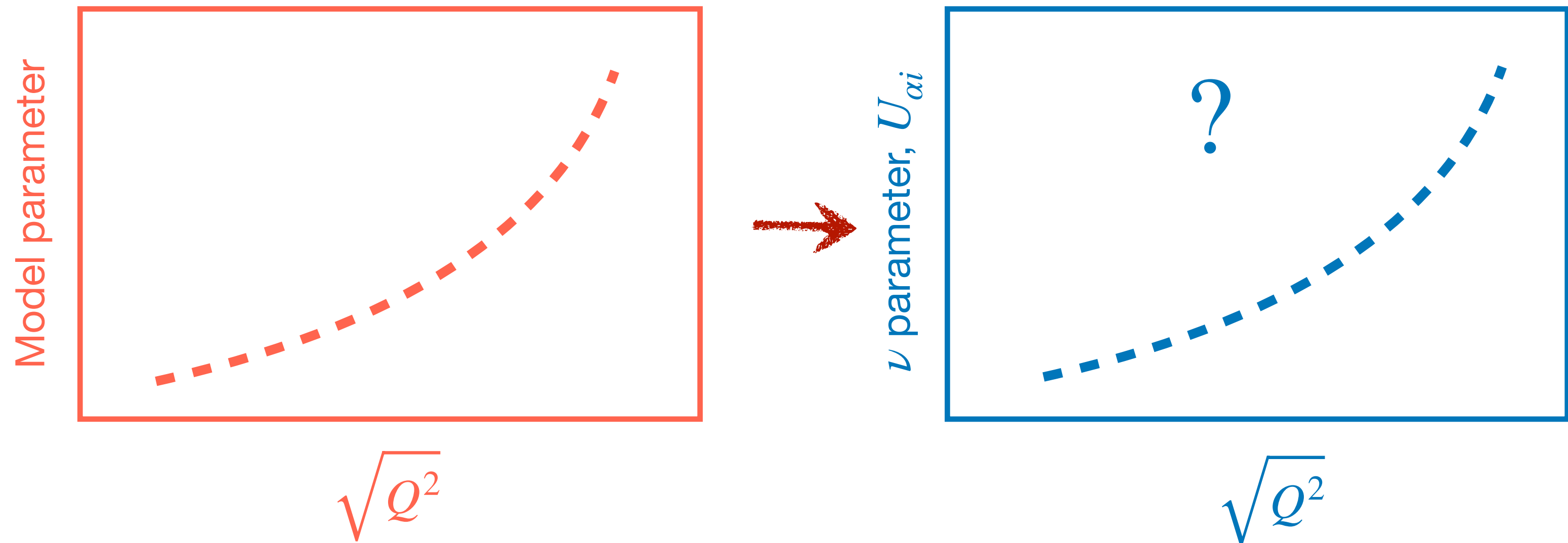
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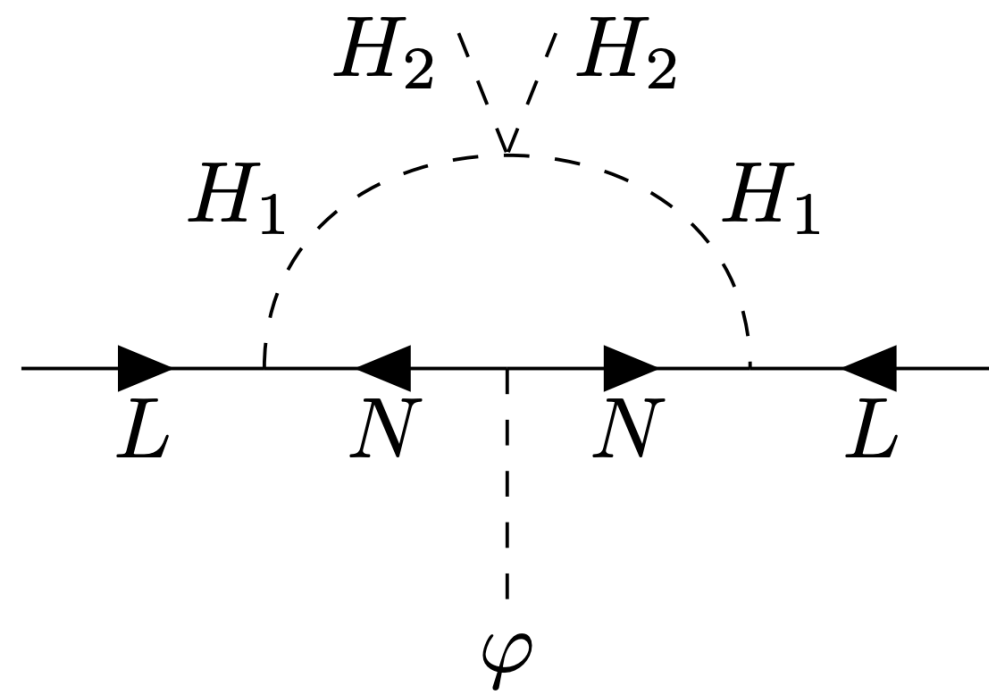


Model Specific Running: from $Y_N(Q^2)$ to $U(Q^2)$

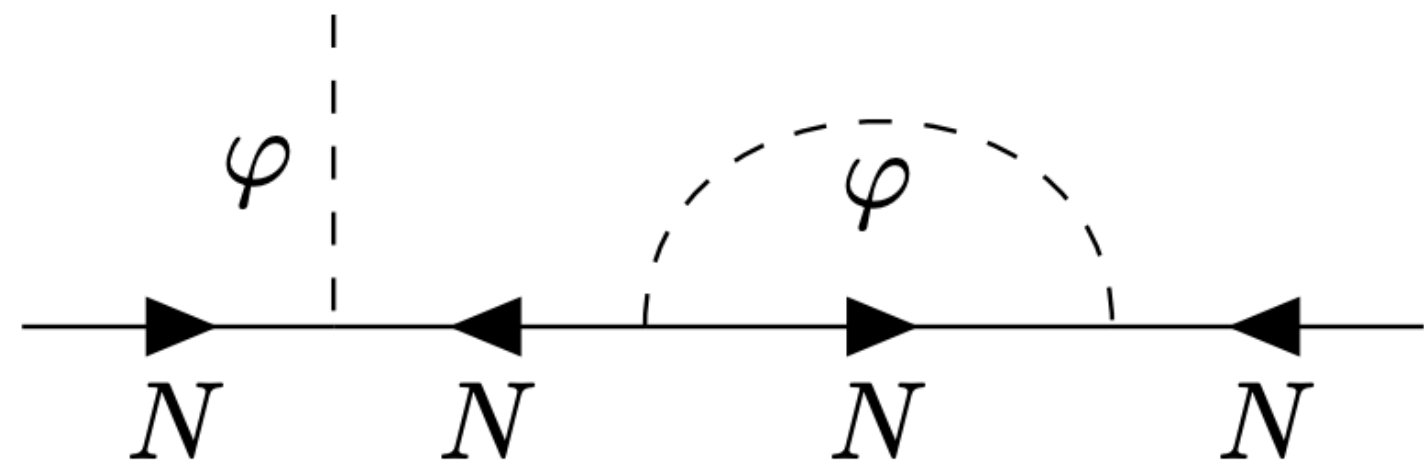
Variation of the scotogenic model

Ma, arXiv:hep-ph/0601225

Babu, Brdar, de Gouvêa, Machado, arXiv:2108.11961



$$-\mathcal{L}_\nu^{(1)} = \bar{L}Y_\nu\tilde{H}_1N_R + \varphi\bar{N}_R^cY_NN_R + \text{h.c.}$$

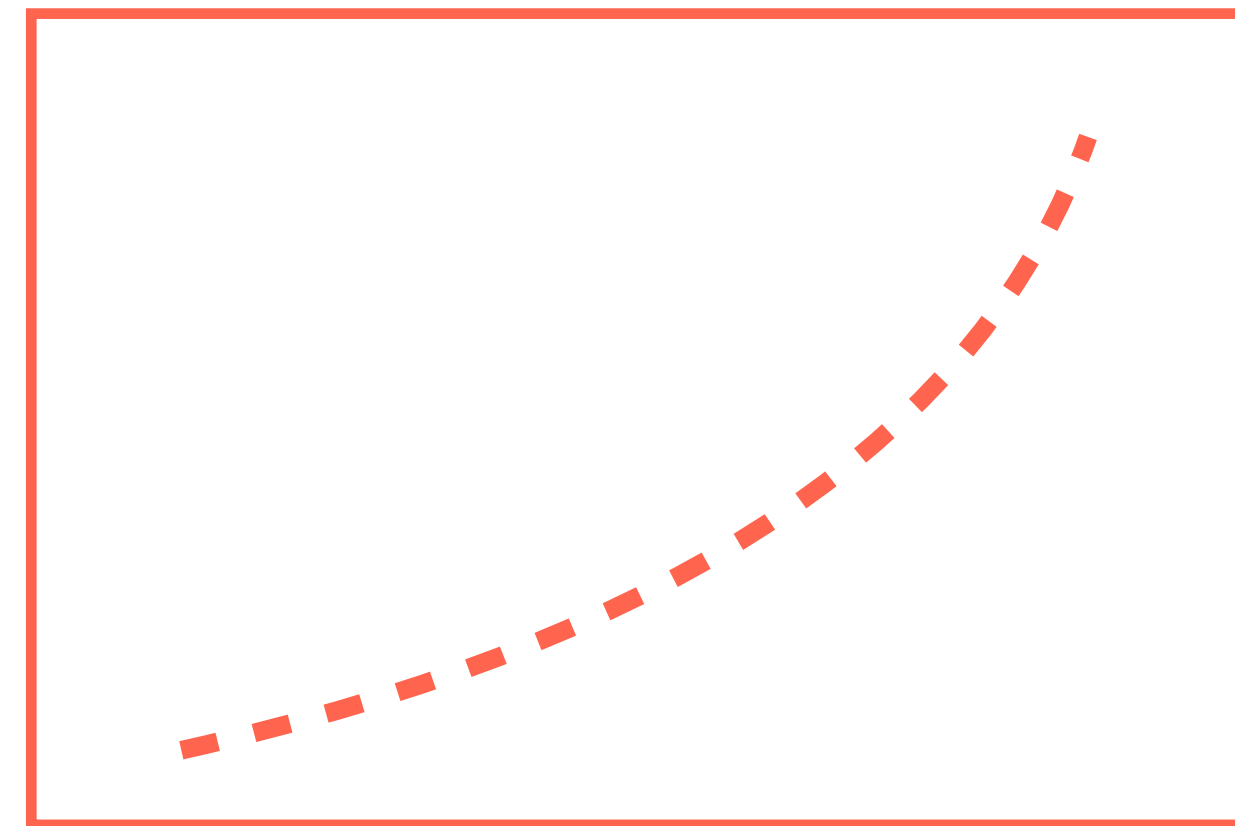


$$16\pi^2\beta(Y_N) \equiv 16\pi^2\frac{dY_N}{d\ln|Q|}$$

$$= 4Y_N\left[Y_N^2 + \frac{1}{2}\text{Tr}(Y_N^2)\right]$$

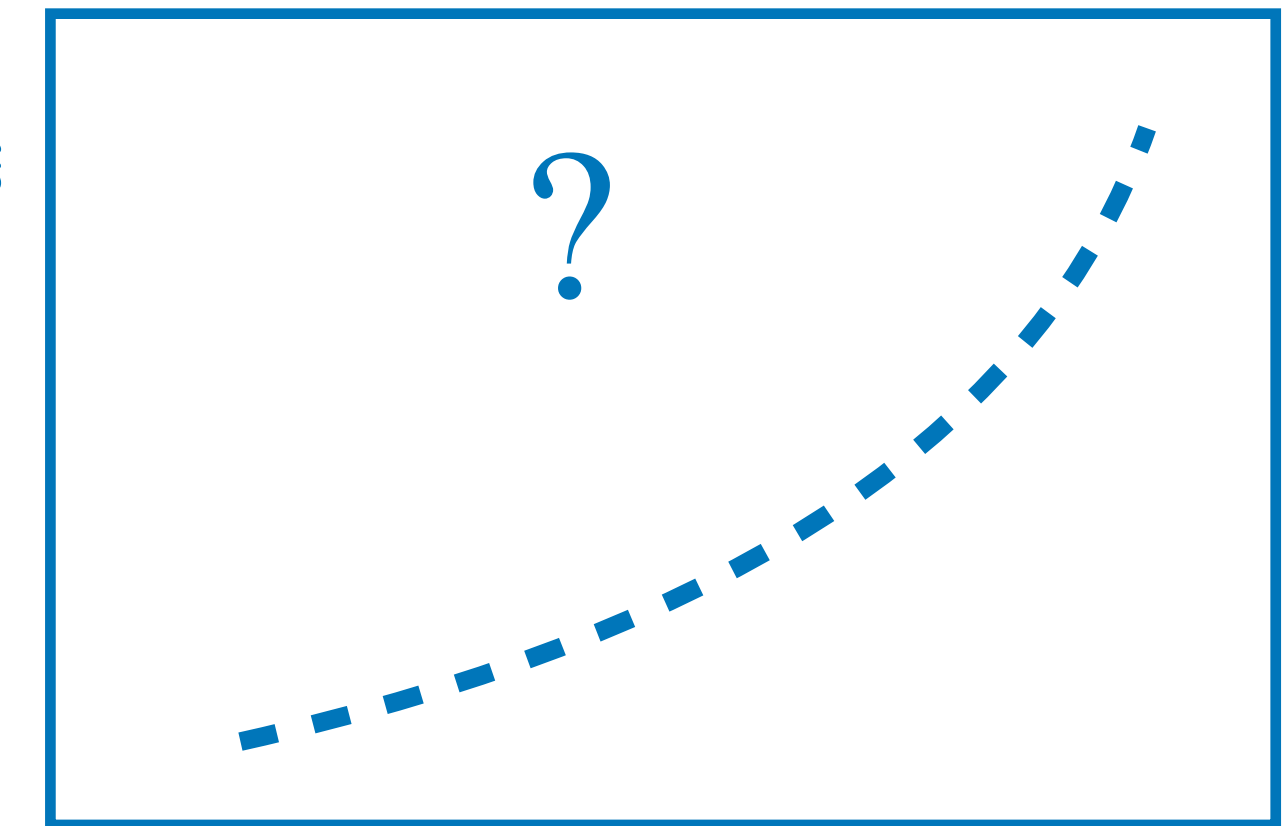
→ PMNS matrix $U(Q^2)$

Model parameter, Y_N



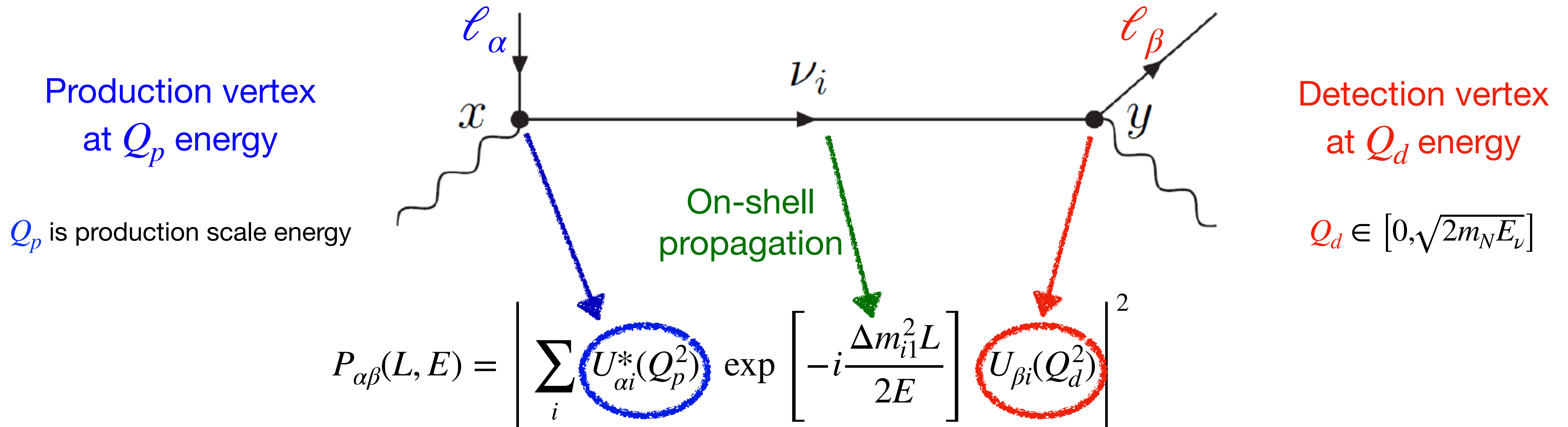
$\sqrt{Q^2}$

ν parameter, $U_{\alpha i}$



$\sqrt{Q^2}$

Oscillation with Mixing Mismatch



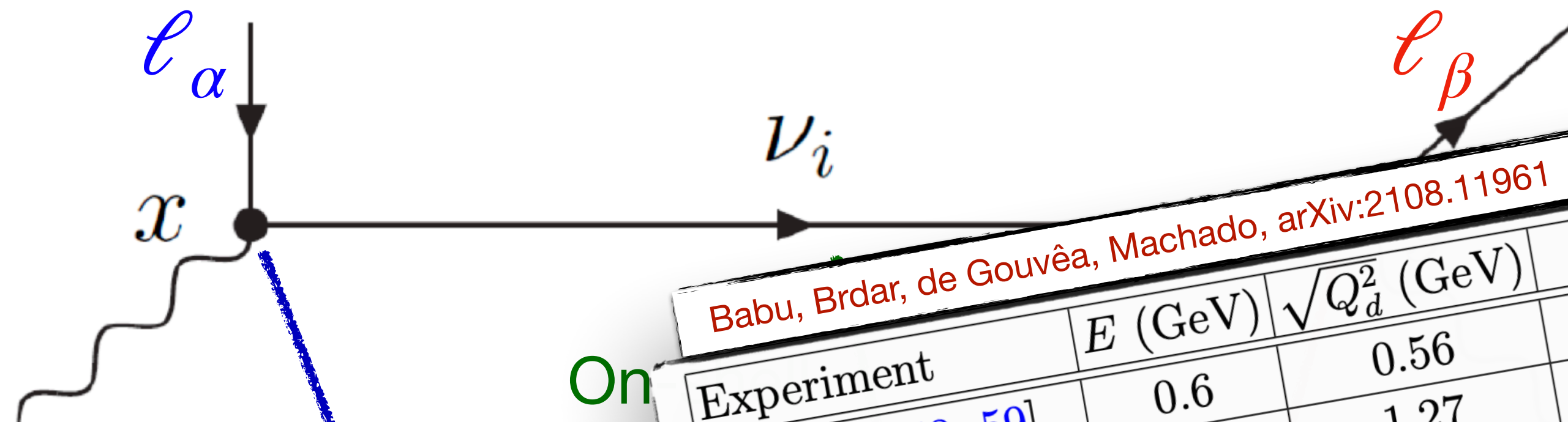
$U(Q_p^2) \neq U(Q_d^2)$ affects $P_{\alpha\beta}$.

Mismatch impacts flavor oscillation.

Oscillation with Mixing Mismatch

Production vertex
at Q_p energy

Q_p is production scale energy



On
propa

$$P_{\alpha\beta}(L, E) = \left| \sum_i U_{\alpha i}^*(Q_p^2) \exp\left[-i\sqrt{E^2 - Q_p^2} L\right] U_{\beta i}(Q_p^2) \right|^2$$

$$U(Q_p^2) \neq U$$

Mismatch impacts flavor oscillation.

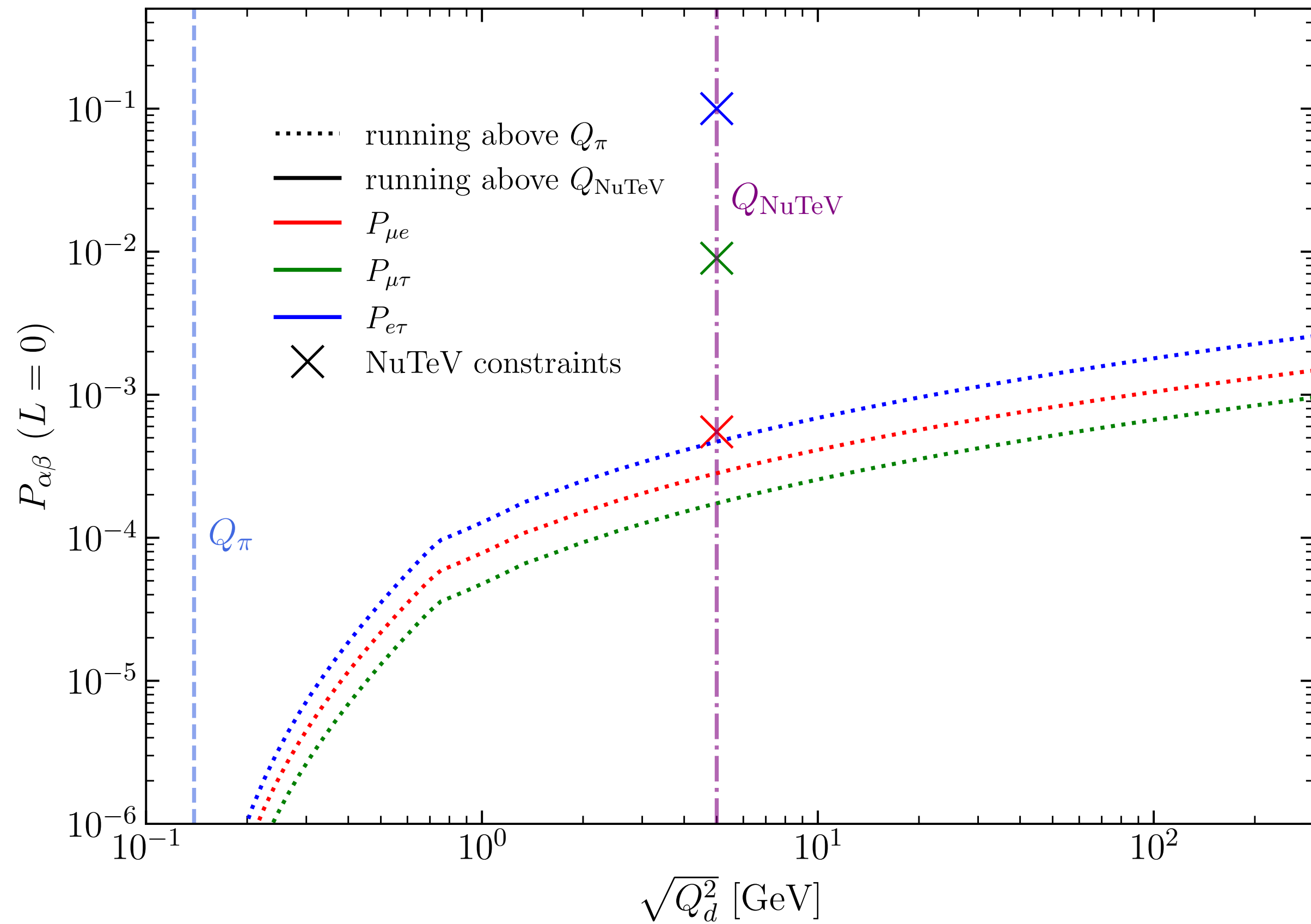
Babu, Brdar, de Gouvêa, Machado, arXiv:2108.11961

Experiment	E (GeV)	$\sqrt{Q_d^2}$ (GeV)	channel	constraint
T2K [47, 58, 59]	0.6	0.56	-	-
NOvA [48, 60, 61]	2.1	1.27	$\nu_\mu \rightarrow \nu_e$	3.4×10^{-3}
ICARUS [65]	17	3.94	$\nu_\mu \rightarrow \nu_e$	2.8×10^{-3}
CHARM-II [66]	24	4.70	$\nu_\mu \rightarrow \nu_e$	7.4×10^{-3}
NOMAD [62-64]	47.5	6.64	$\nu_\mu \rightarrow \nu_\tau$	1.63×10^{-4}
NuTeV [67, 68]	250	15.30	$\nu_\mu \rightarrow \nu_e$	5.5×10^{-4}
			$\nu_e \rightarrow \nu_\tau$	0.1
			$\nu_\mu \rightarrow \nu_\tau$	9×10^{-3}

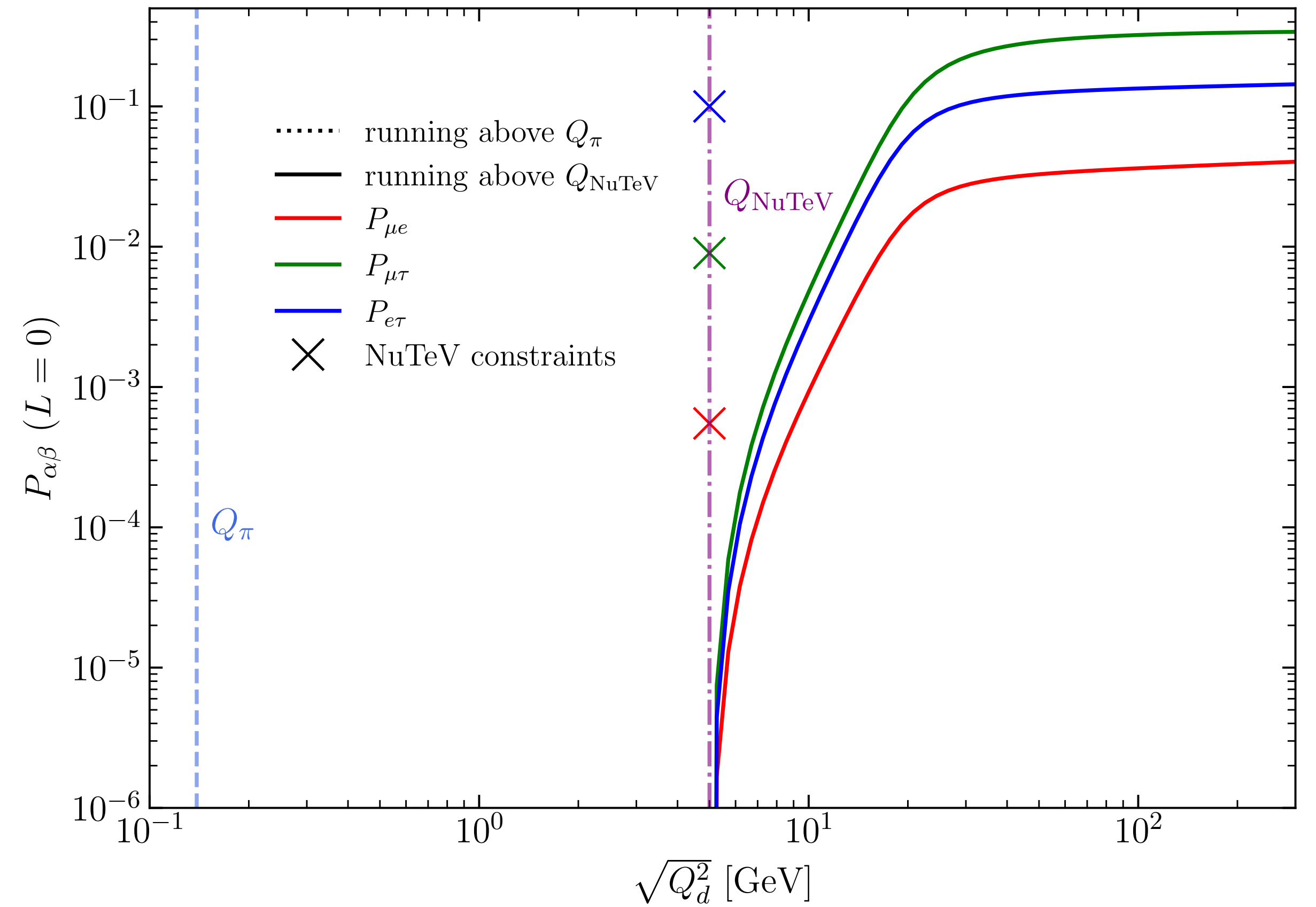
Stringent constraints come from short-baseline neutrino experiments, the highest energy constraints from NuTeV

BSM with Large Transition $P_{\alpha\beta}$

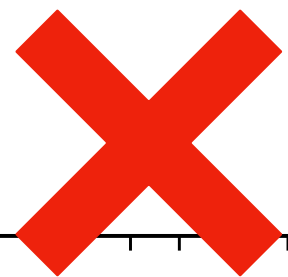
$$m_{\varphi,N} \gtrsim Q_p = m_\pi$$



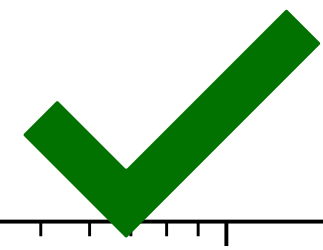
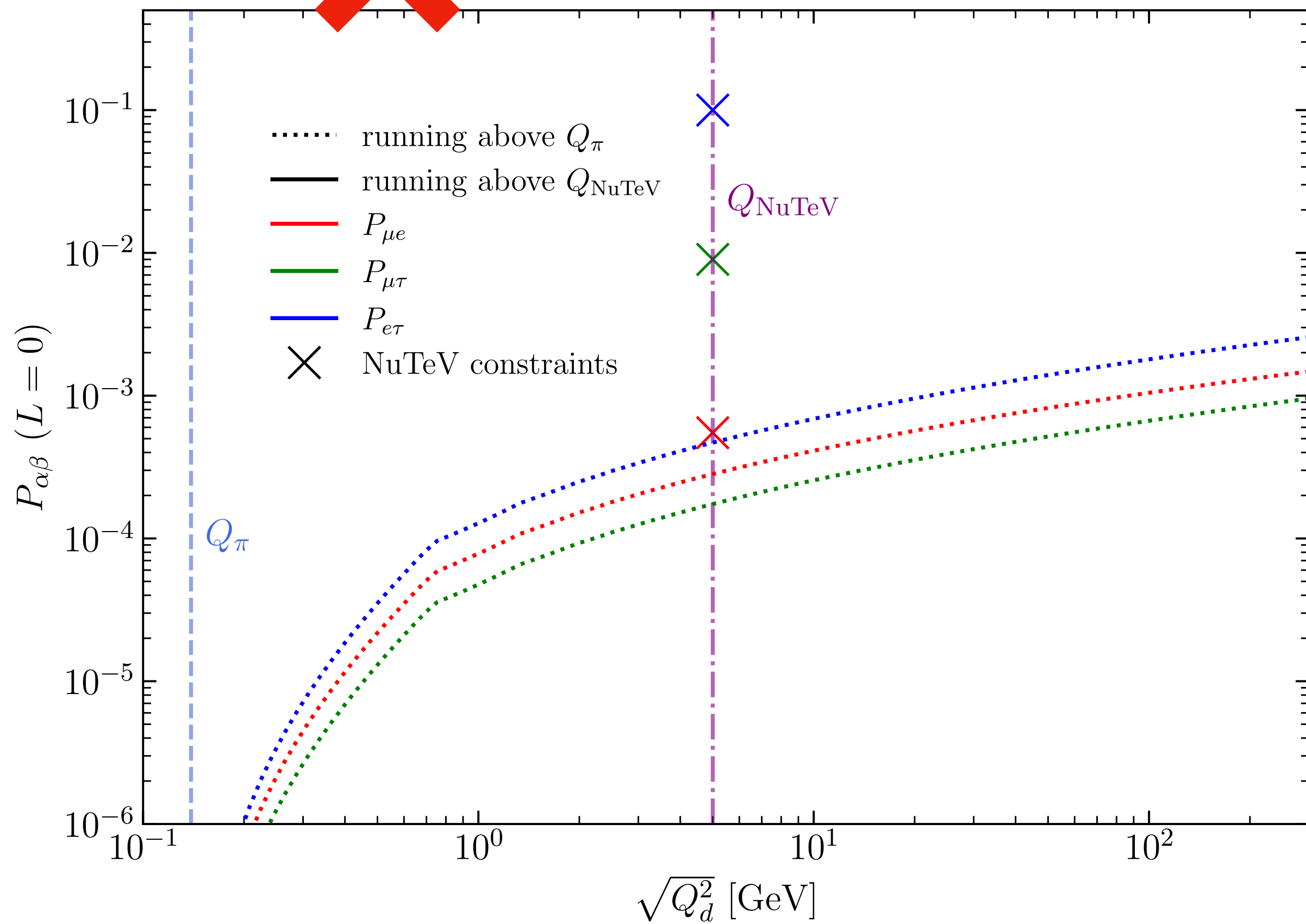
$$m_{\varphi,N} \gtrsim Q_{\text{NuTeV}}$$



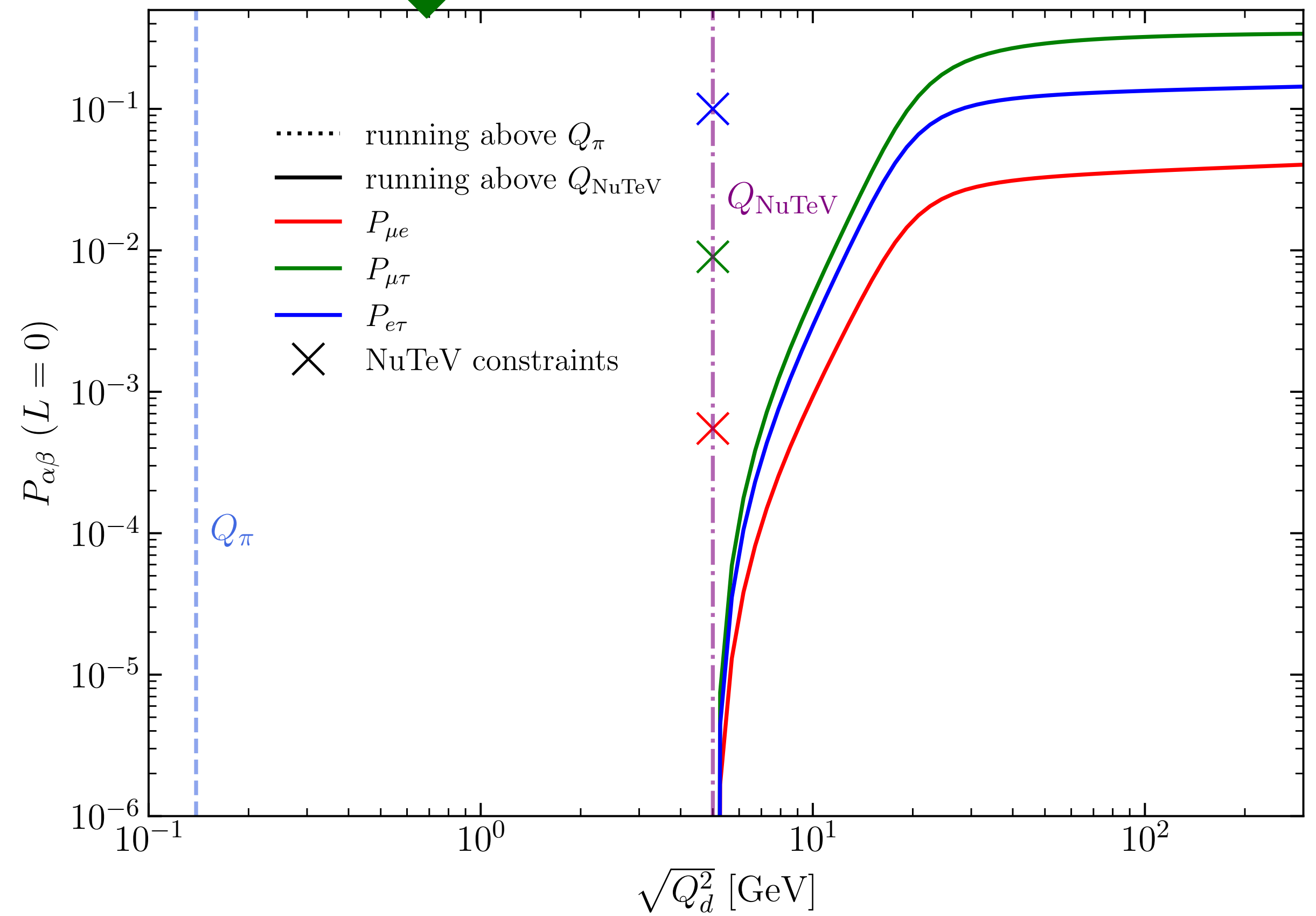
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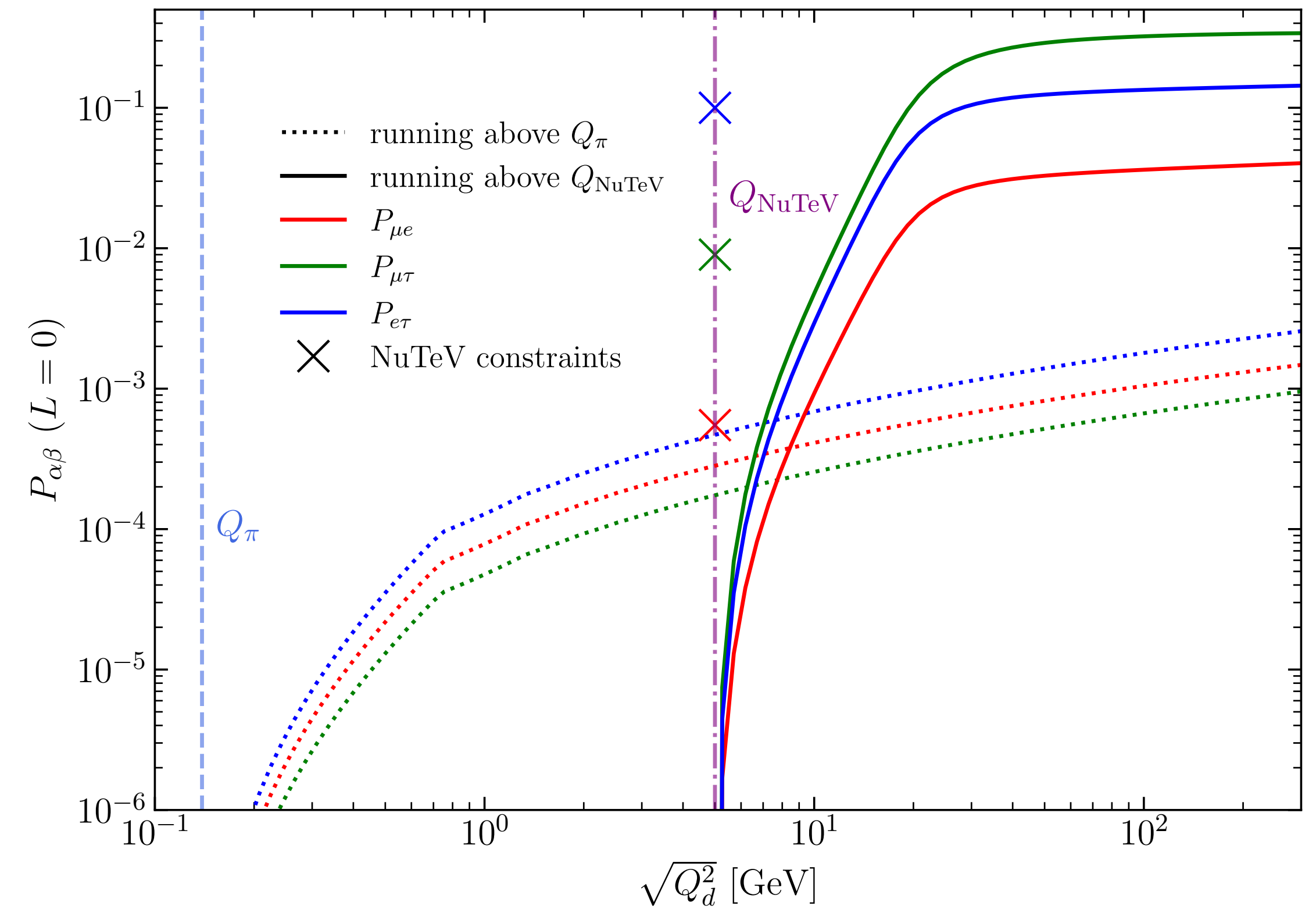
$$m_{\varphi,N} \gtrsim Q_{\text{NuTeV}}$$



Testing RG Running



Which experiment to look at?

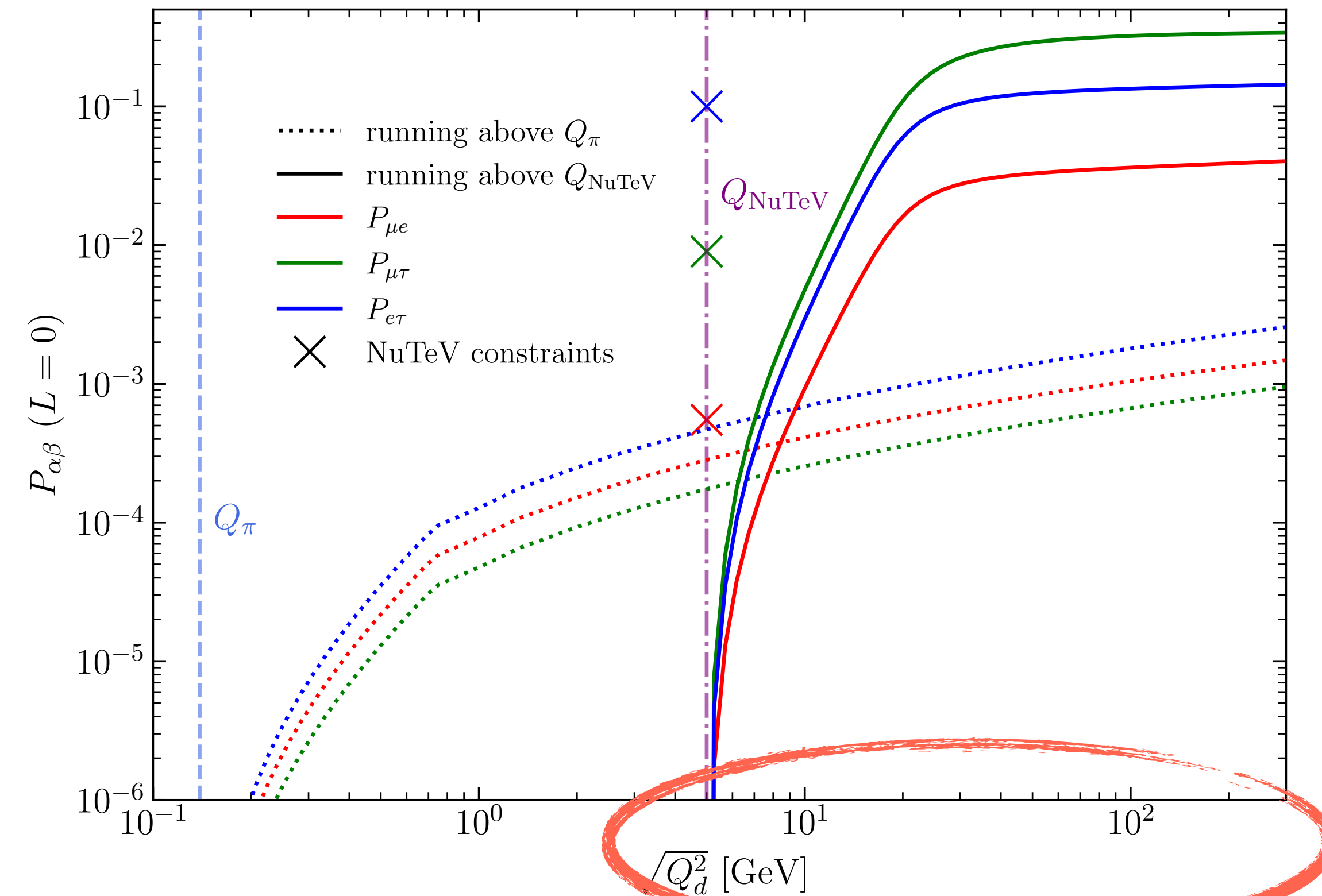


Testing RG Running



Which experiment to look at?

Access to high-energy neutrinos



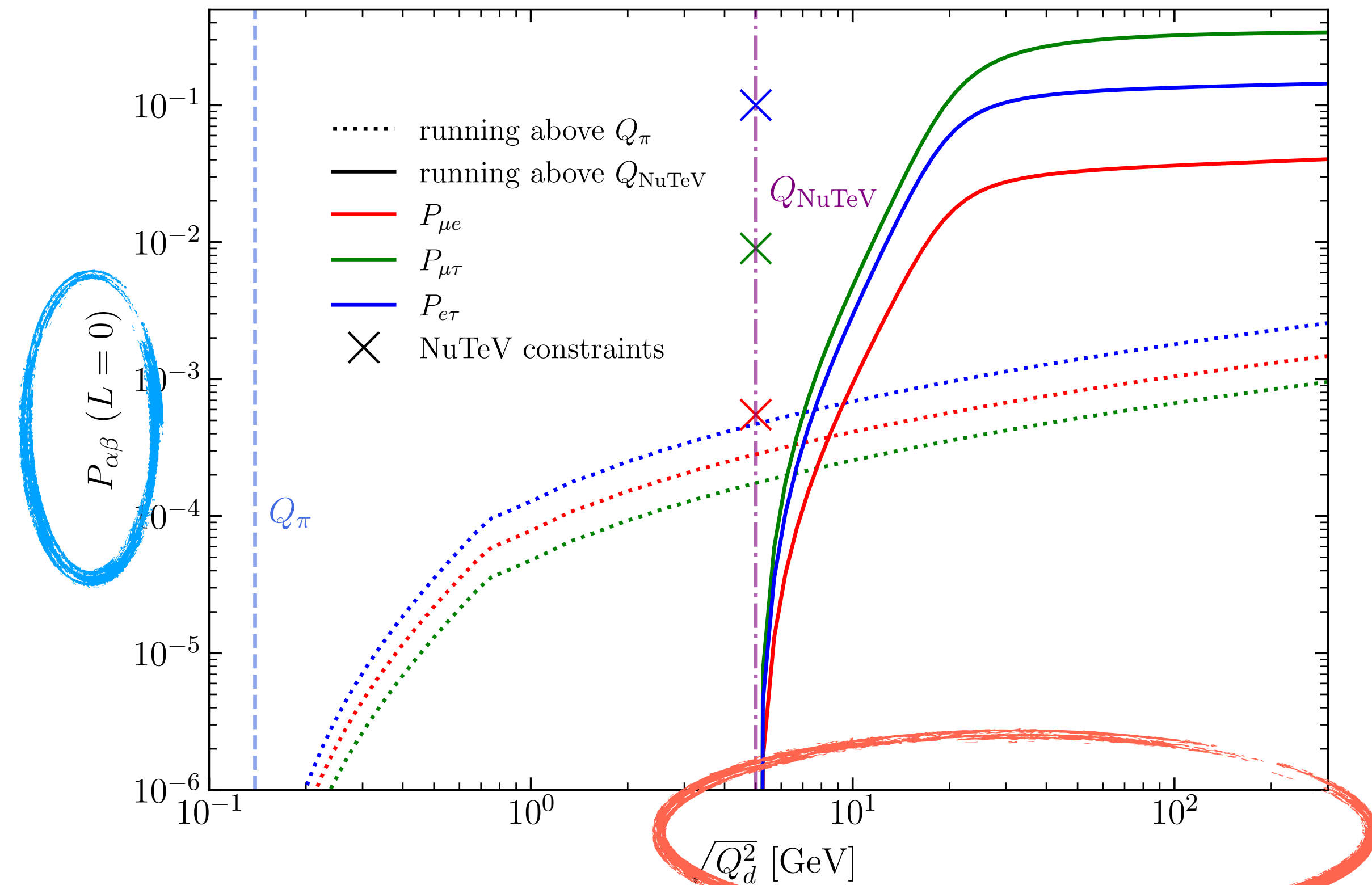
Testing RG Running



Which experiment to look at?

Access to high-energy neutrinos

Sensitive to flavor oscillation $P_{\alpha\beta}$



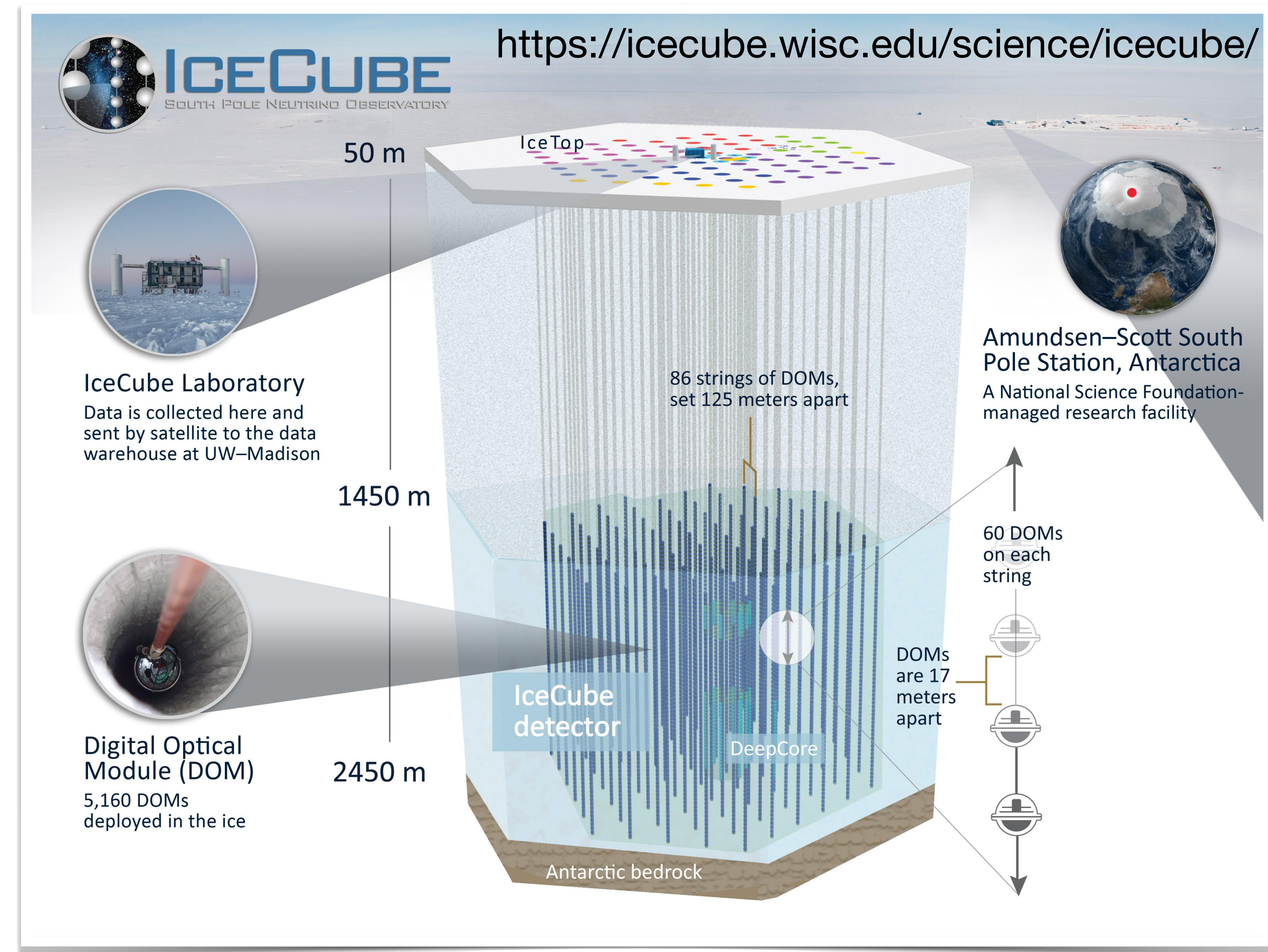
Testing RG Running



Which experiment to look at?

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Tau Neutrino Detection at IceCube

○ double bangs morphology at high-energy:

Learned and Pakvasa, arXiv:hep-ph/9405296

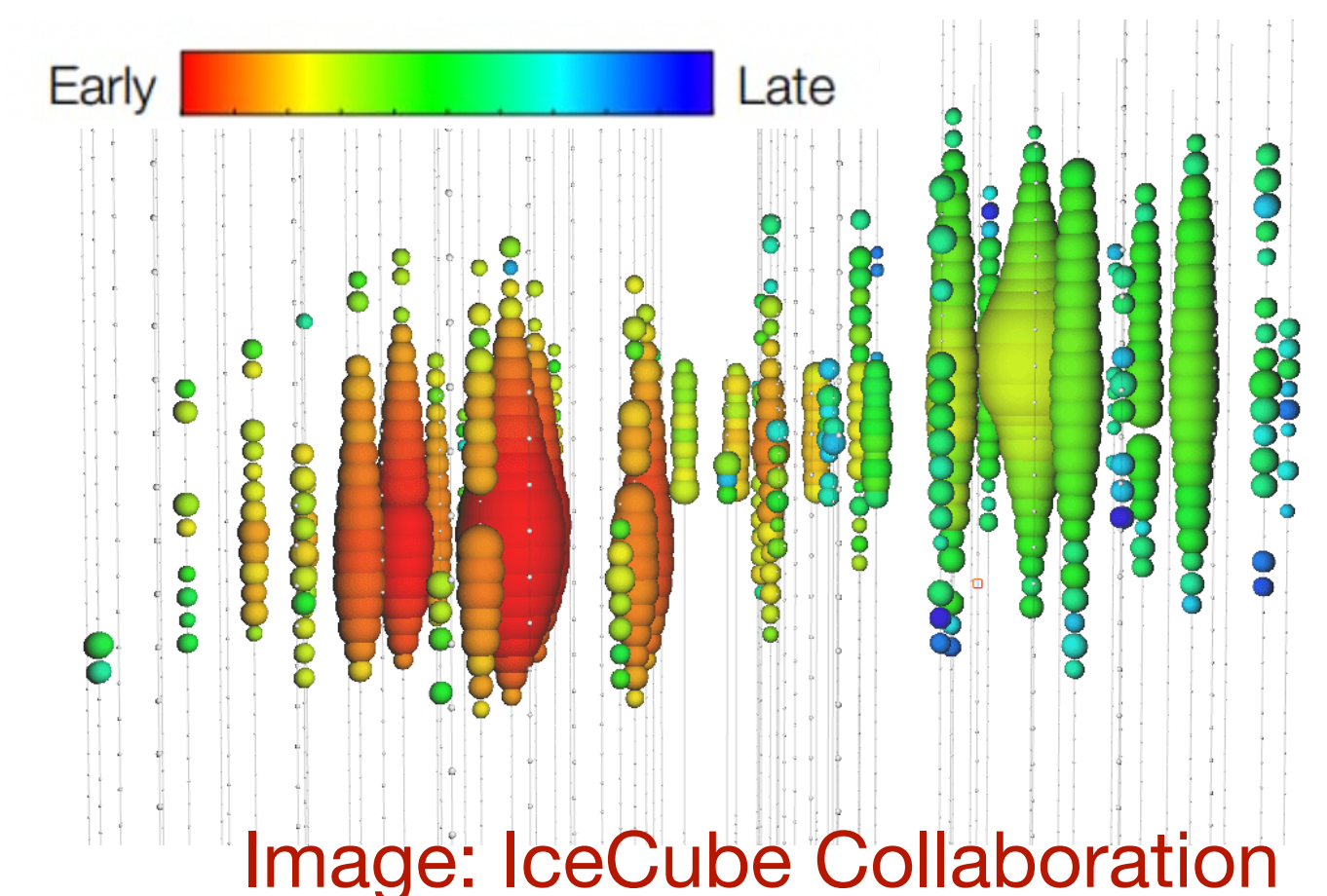
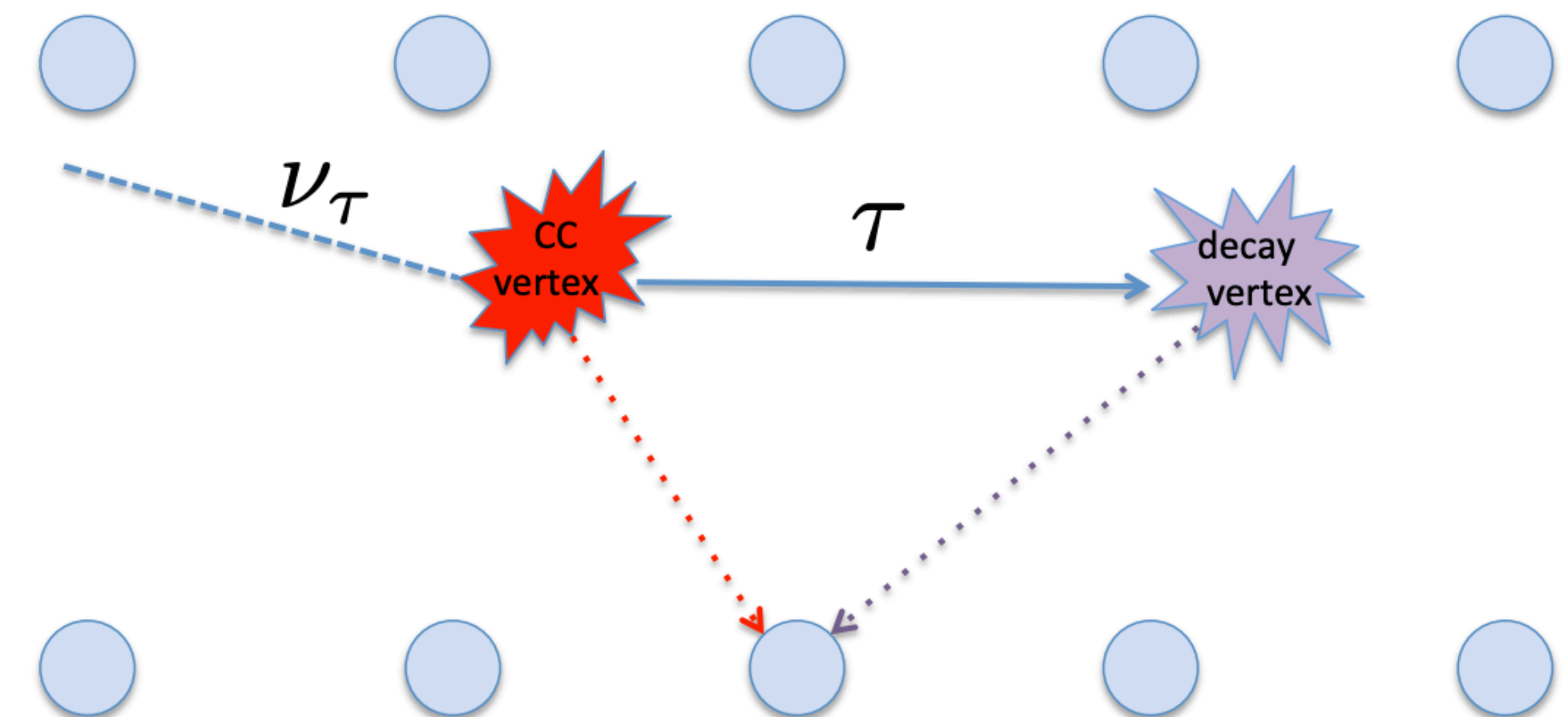
- ☑ Tau producing vertex — first shower
- ☑ Tau decaying vertex — second shower

○ 50 m/PeV separation

○ High-energy starting events (HESE) data in 7.5 years data [IceCube Collaboration, arXiv:2011.03545]

- Total 102 events (ν_e, ν_μ, ν_τ) → Only 4 ν_τ -candidates
- 4 ν_τ -candidates → 2 ν_τ -events with high likelihood of being ν_τ [IceCube Collaboration, arXiv:2011.03561]

IceCube Collaboration, arXiv:1509.06212



BSM weight and Number of Events N_β

Number of ν_β neutrino events,

$$N_\beta = T \int dE_\nu w_{\alpha \rightarrow \beta}^{\text{BSM}}(E_\nu) \epsilon_\beta(E_\nu) \Phi_{\nu_\alpha}(E_\nu) \sigma_{\nu N}(E_\nu)$$

T = Detector livetime

ϵ_α = Detector efficiency

Φ_{ν_α} = Incoming atmospheric and astrophysical flux

$\sigma_{\nu N}$ = Deep inelastic cross section

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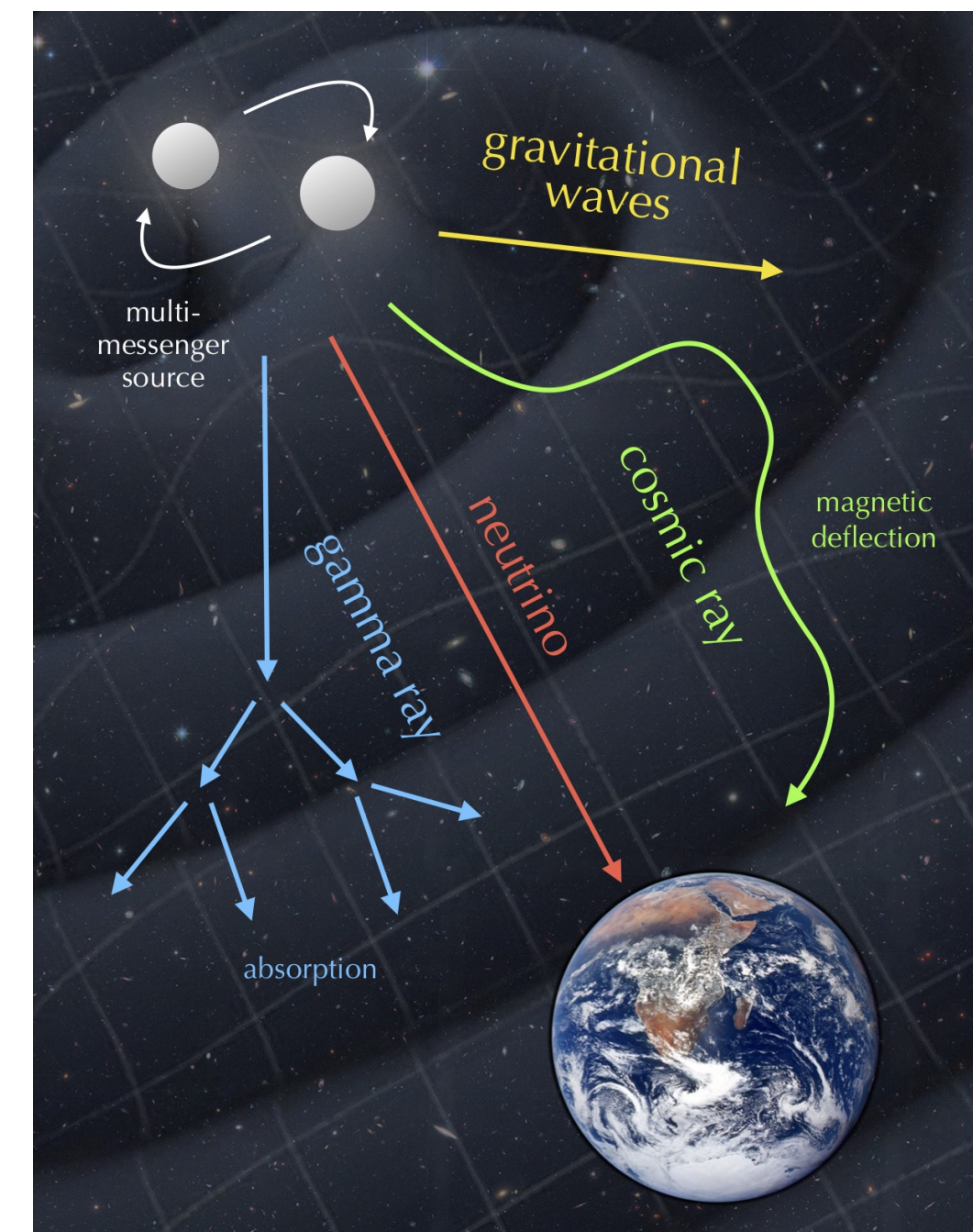
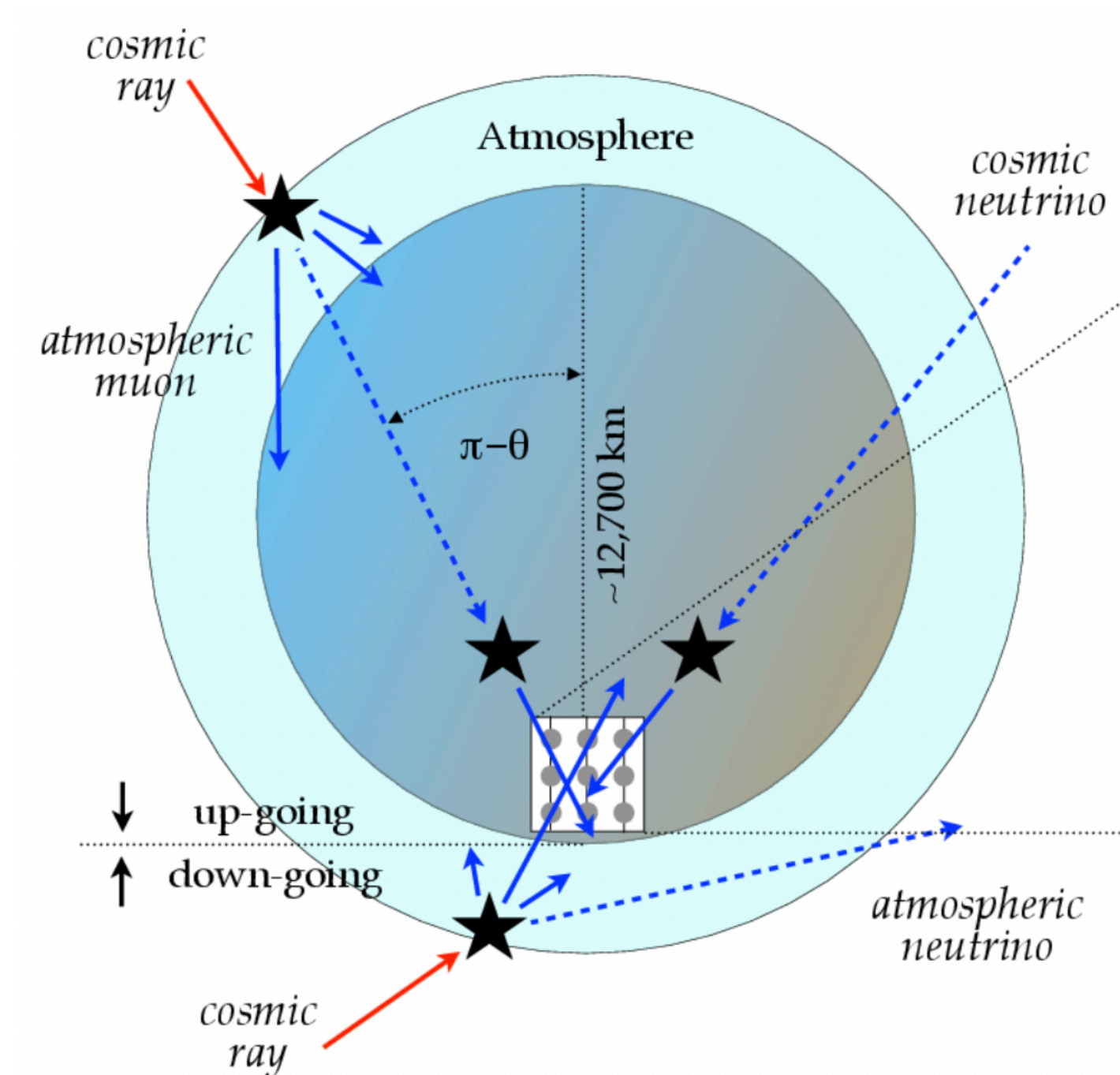
$\sigma_{\nu N}$ = Deep inelastic cross section

BSM weight from the deep inelastic scattering cross section weighted by $P_{\alpha\beta}(E_\nu, Q^2)$,

$$w_{\alpha \rightarrow \beta}^{\text{BSM}}(E_\nu) = \frac{\int_\nu P_{\alpha\beta}(E_\nu, Q^2) + \int_{\bar{\nu}} P_{\bar{\alpha}\bar{\beta}}(E_\nu, Q^2)}{\int_\nu + \int_{\bar{\nu}}},$$

$$\int_{\nu/\bar{\nu}} = \int dx dQ^2 \frac{d^2\sigma_{\nu/\bar{\nu}}}{dx dQ^2}$$

$P_{\alpha\beta}$ for Neutrino Fluxes



Atmospheric neutrinos: $L \ll L_{osc}$

$$P_{\alpha\beta} = P_{\alpha\beta}(L = 0) = \left| \sum_i U_{\alpha i}^*(Q_p^2) U_{\beta i}(Q_d^2) \right|^2$$

Astrophysical neutrinos: $L \gg L_{osc}$

$$P_{\alpha\beta}(L \gg L_{osc}) = \sum_i \left| U_{\alpha i}(Q_p^2) \right|^2 \left| U_{\beta i}(Q_d^2) \right|^2$$

Tau Neutrino Appearance at IceCube

Difference in the number of events,

$$\Delta N_{\alpha} = (N_{\text{running}} - N_{\text{no running}})_{\alpha}$$

$$\text{Standard oscillation: } \Delta N_{\mu} = \Delta N_{\tau} = 0$$

$$N_{\tau}^{\text{no running}} = 2\text{-}3 \text{ events}$$

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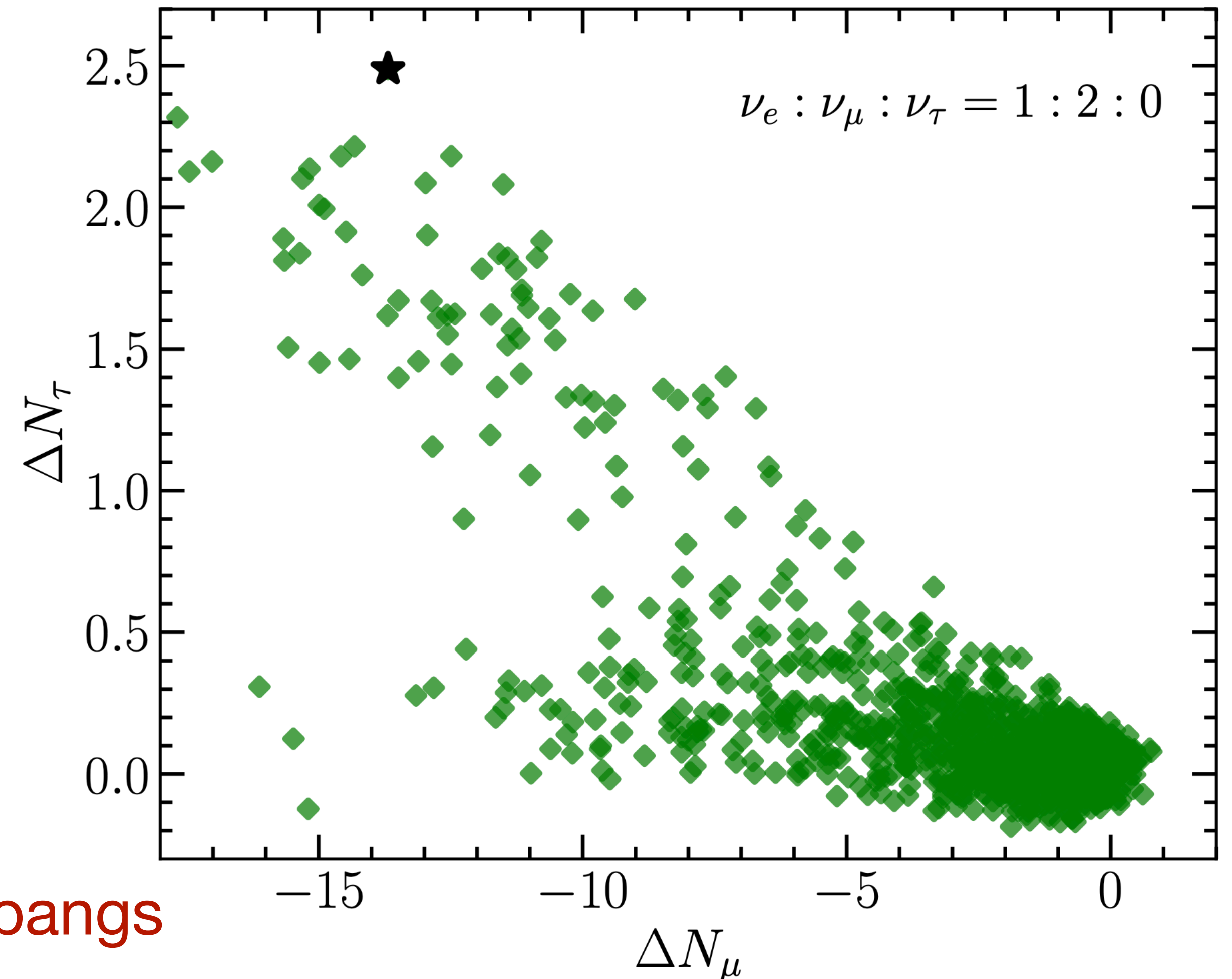
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Running \rightarrow **green points**

★ benchmark:

maximal number of BSM-induced double bangs



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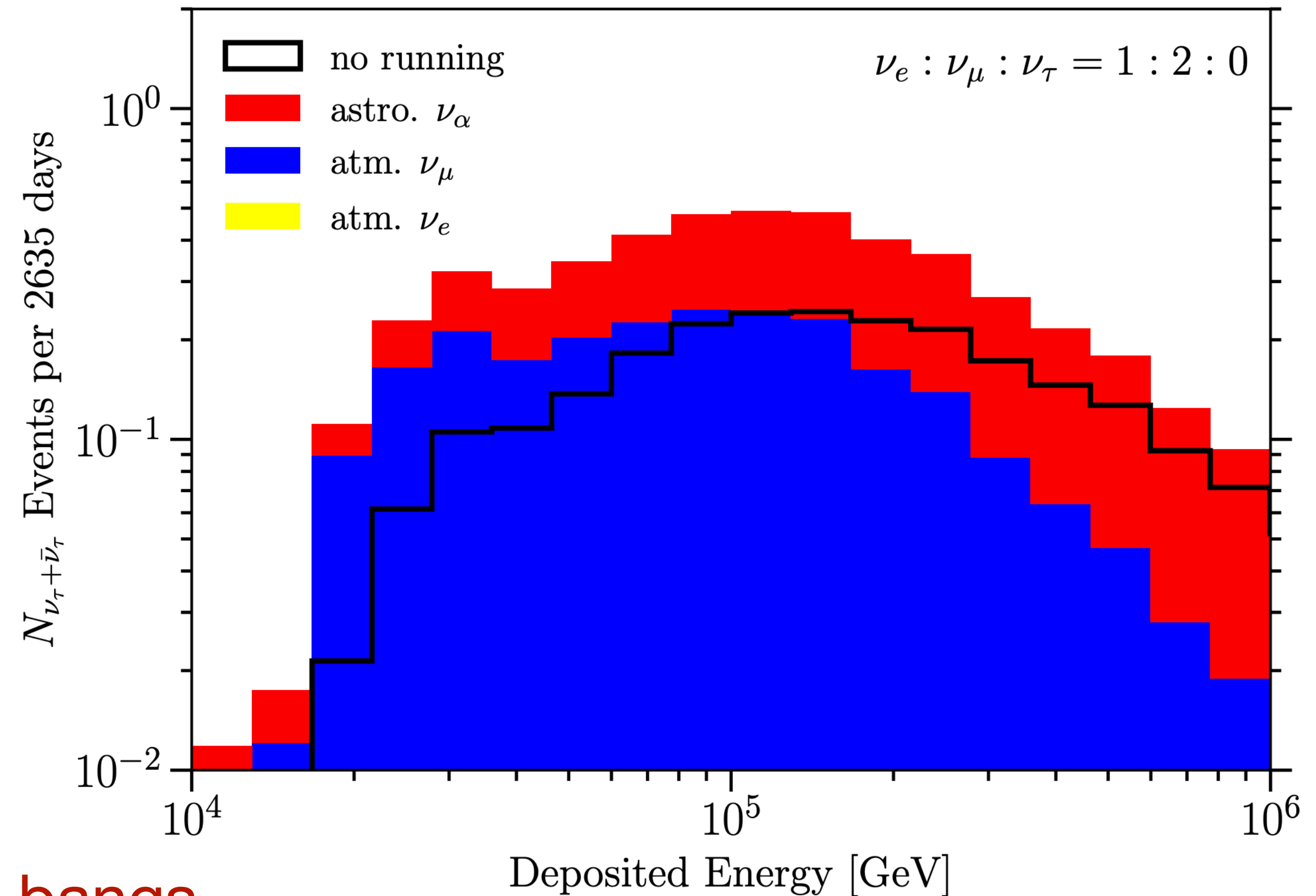
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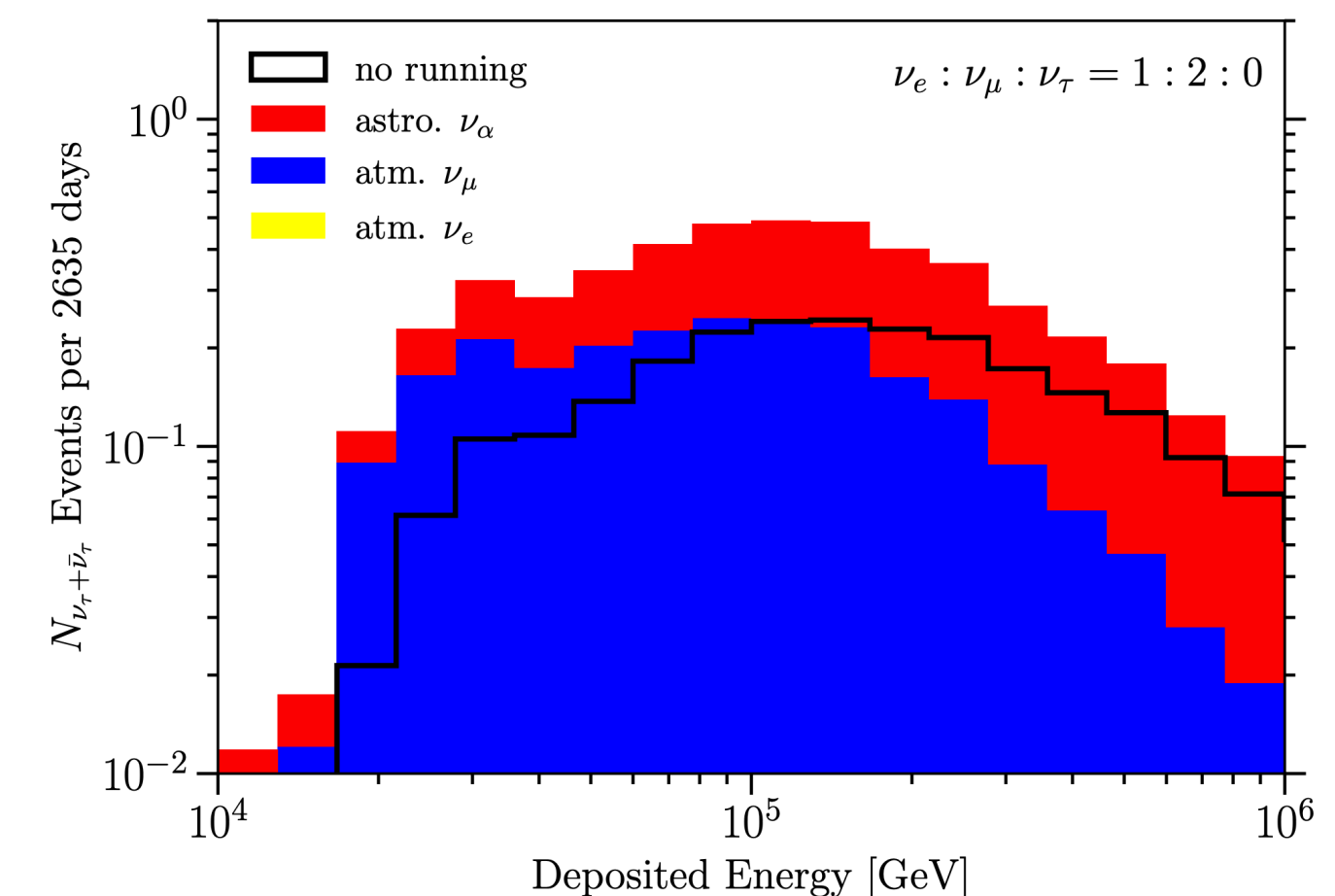
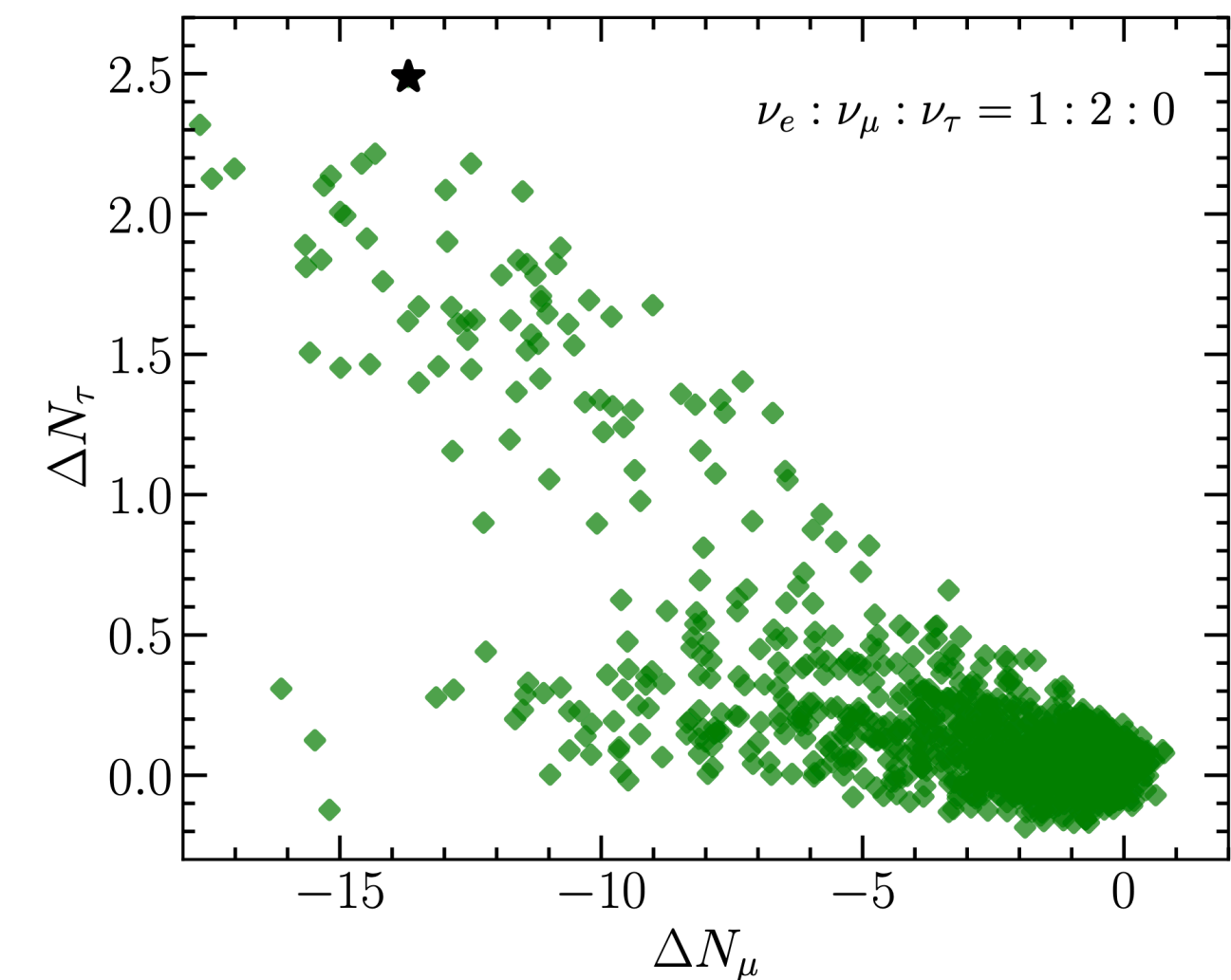
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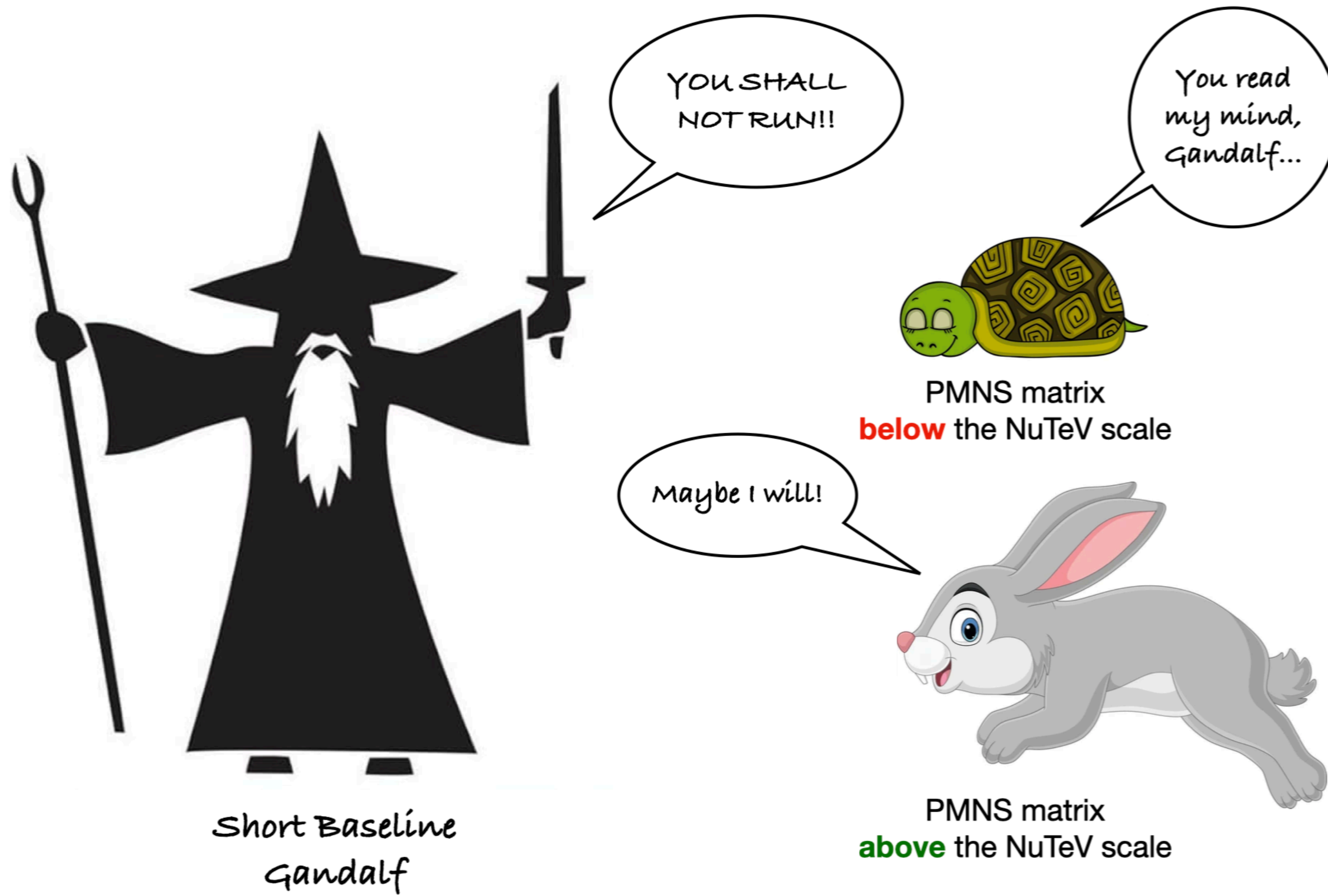
maximal number of BSM-induced double bangs



Summary

- BSM exhibiting RG running  of the neutrino mixing **below the electroweak scale**
- **Excess** in **double bang events** at IceCube
- Substantial contribution from **atmospheric ν_μ**
- **$\sim 6x$ larger statistics** at IceCube-Gen2 than IceCube, reducing fluctuations

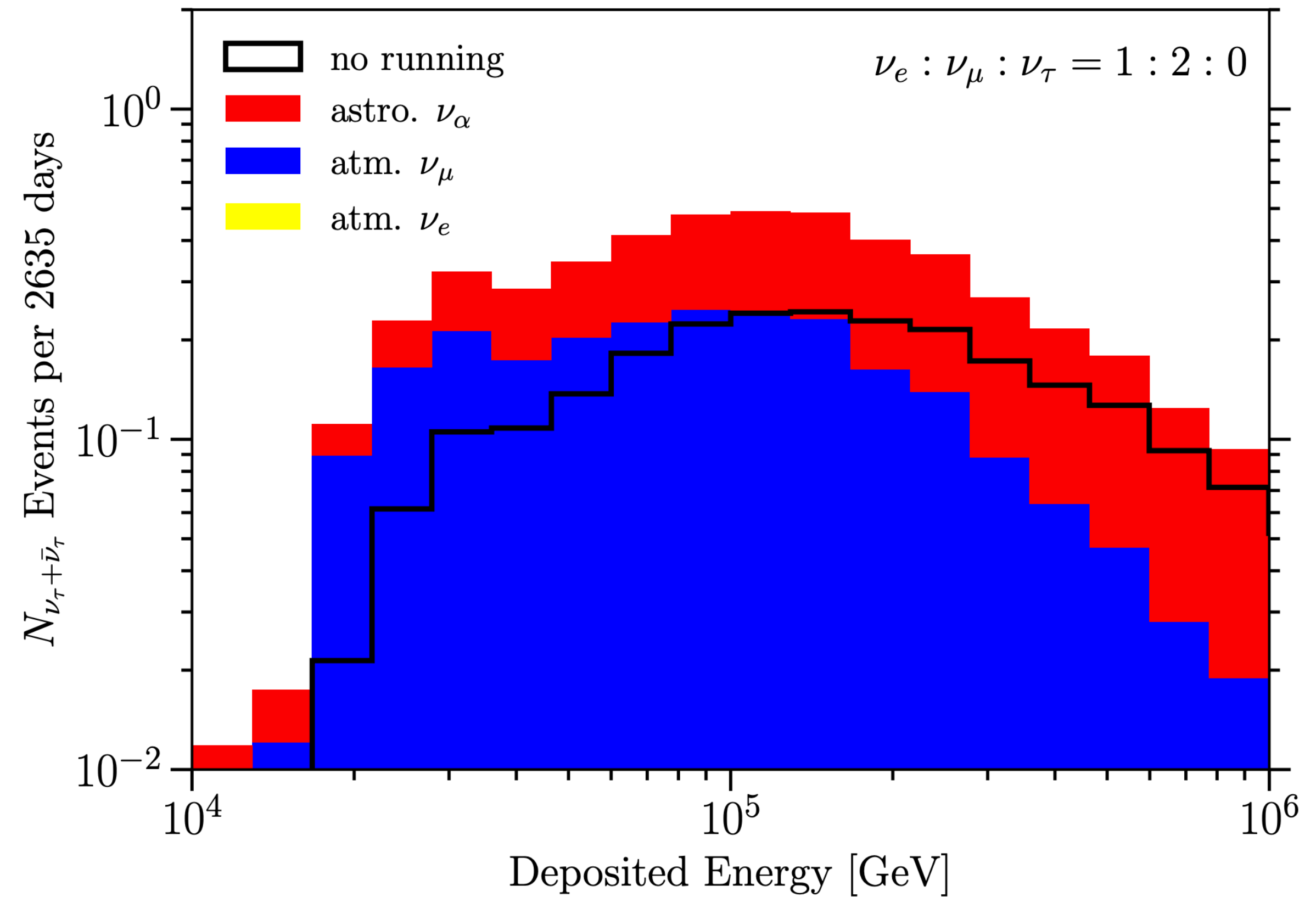
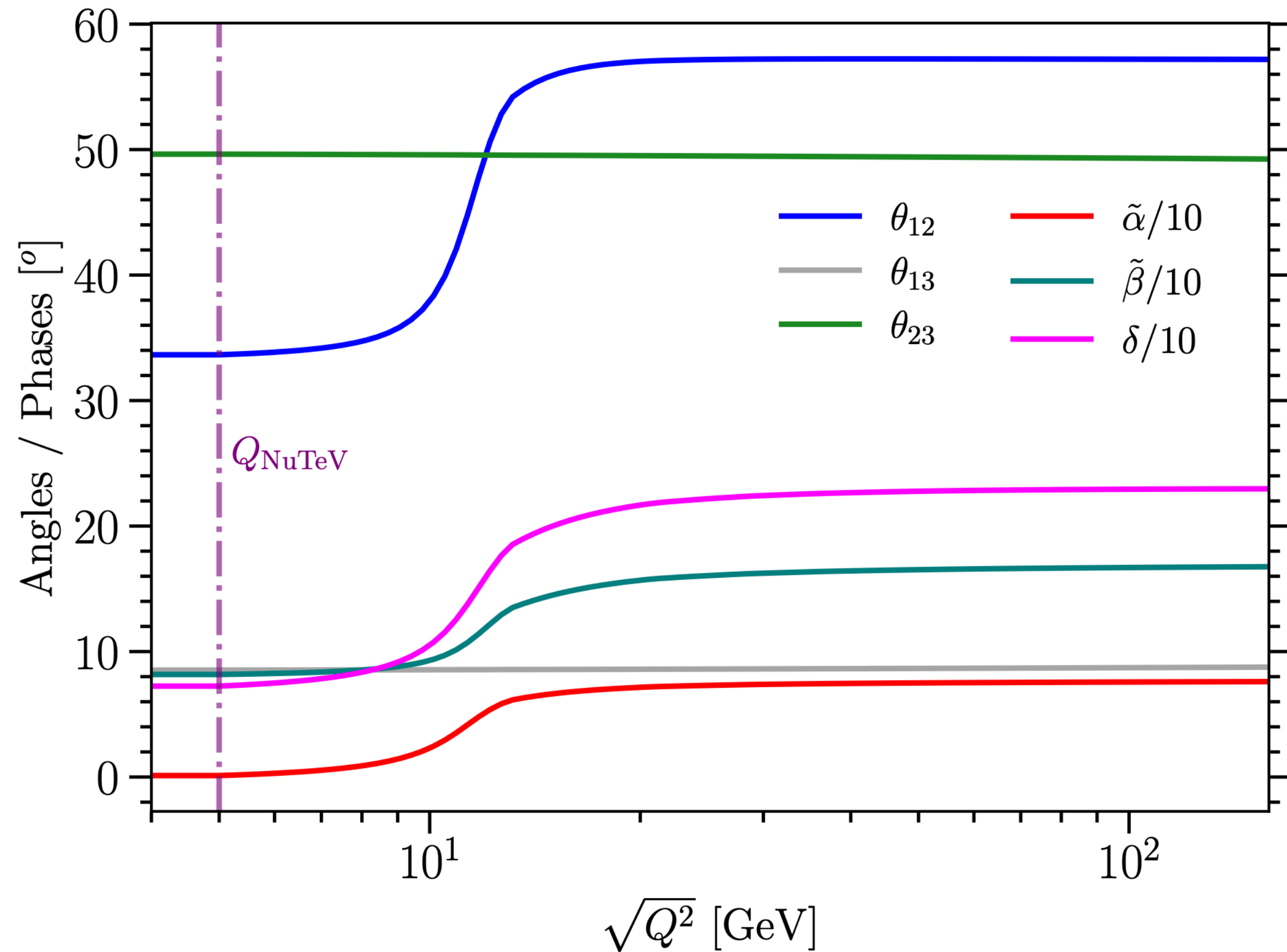




Backup

★ Benchmark at IceCube

★ 100% excess in double bangs ★ Dominant contribution from atm. $\nu_\mu \rightarrow \nu_\tau$

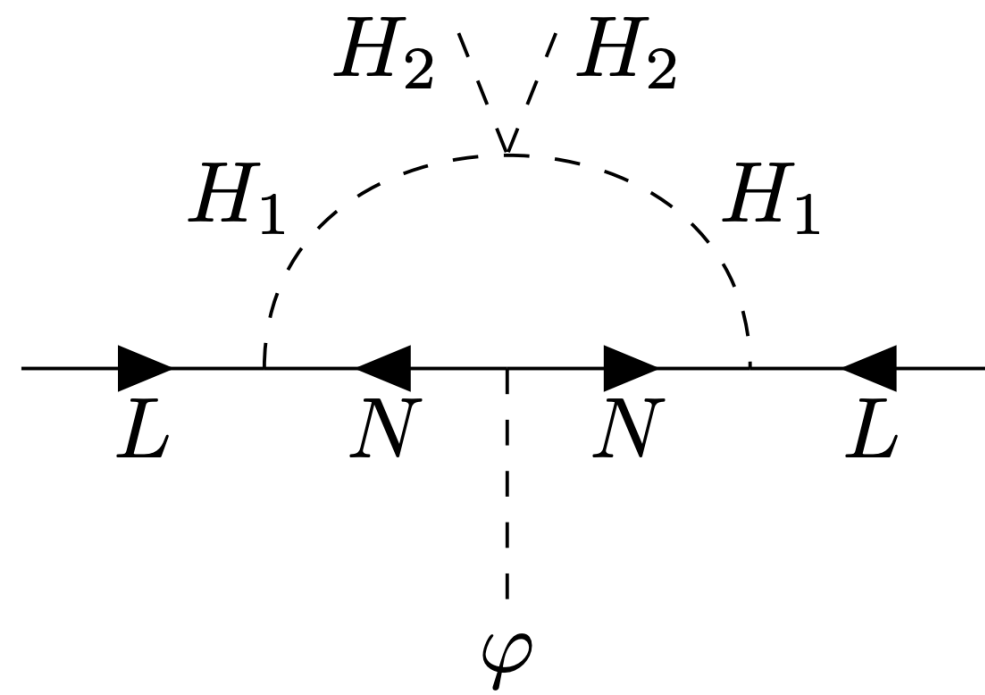


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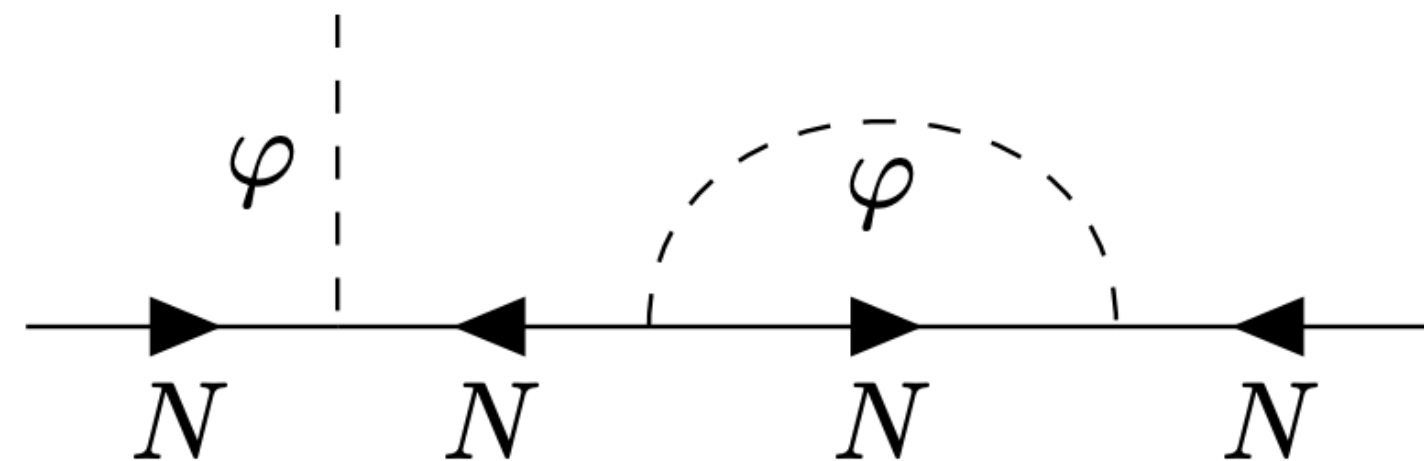
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$$16\pi^2\beta(Y_N) \equiv 16\pi^2\frac{dY_N}{d\ln|Q|} = 4Y_N\left[Y_N^2 + \frac{1}{2}\text{Tr}(Y_N^2)\right]$$



$$\text{Neutrino mass matrix, } M_\nu \simeq \frac{\lambda v_\varphi}{16\sqrt{2}\pi^2} \ln\left(\frac{M_H^2}{M_A^2}\right) Y_\nu Y_N(Q^2) Y_\nu^T$$



PMNS matrix $U(Q^2)$ diagonalizes
 $M_\nu \rightarrow U^T M_\nu U = \text{diag}(m_1, m_2, m_3)$

Oscillation with 2 Flavours

$$U(Q^2) = \begin{pmatrix} \cos \theta(Q^2) & \sin \theta(Q^2) \\ -\sin \theta(Q^2) & \cos \theta(Q^2) \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & e^{i\tilde{\beta}(Q^2)} \end{pmatrix} \quad \theta(Q_p^2) \equiv \theta_p, \quad \theta(Q_d^2) \equiv \theta_d, \quad \tilde{\beta}(Q_d^2) - \tilde{\beta}(Q_p^2) \equiv \beta$$

$\tilde{\beta}$ = CP-violating phase

$$P_{e\mu} = P_{\mu e} = \sin^2(\theta_p - \theta_d) + \sin 2\theta_p \sin 2\theta_d \sin^2\left(\frac{\Delta m^2 L}{4E} + \frac{\beta}{2}\right),$$

$$P_{ee} = P_{\mu\mu} = \cos^2(\theta_p - \theta_d) - \sin 2\theta_p \sin 2\theta_d \sin^2\left(\frac{\Delta m^2 L}{4E} + \frac{\beta}{2}\right).$$

At zero baseline $L = 0$ with $\theta_p \approx \theta_d$:

$$P_{\alpha\beta} = \sin^2 2\theta \sin^2\left(\frac{\Delta m^2 L}{4E} + \frac{\beta}{2}\right)$$

$$= \sin^2 2\theta \sin^2\left(\frac{\beta}{2}\right) \neq 0$$

$\beta = 0$ in the absence of running.

$\beta \neq 0$ gives oscillation contribution at zero baseline $L = 0$.

Tau Neutrino Detection at IceCube

- 4 tau-neutrino-candidate events reported out of 102 high-energy starting events (HESE) sample in 7.5 years data [[arXiv:2011.03545](#)]

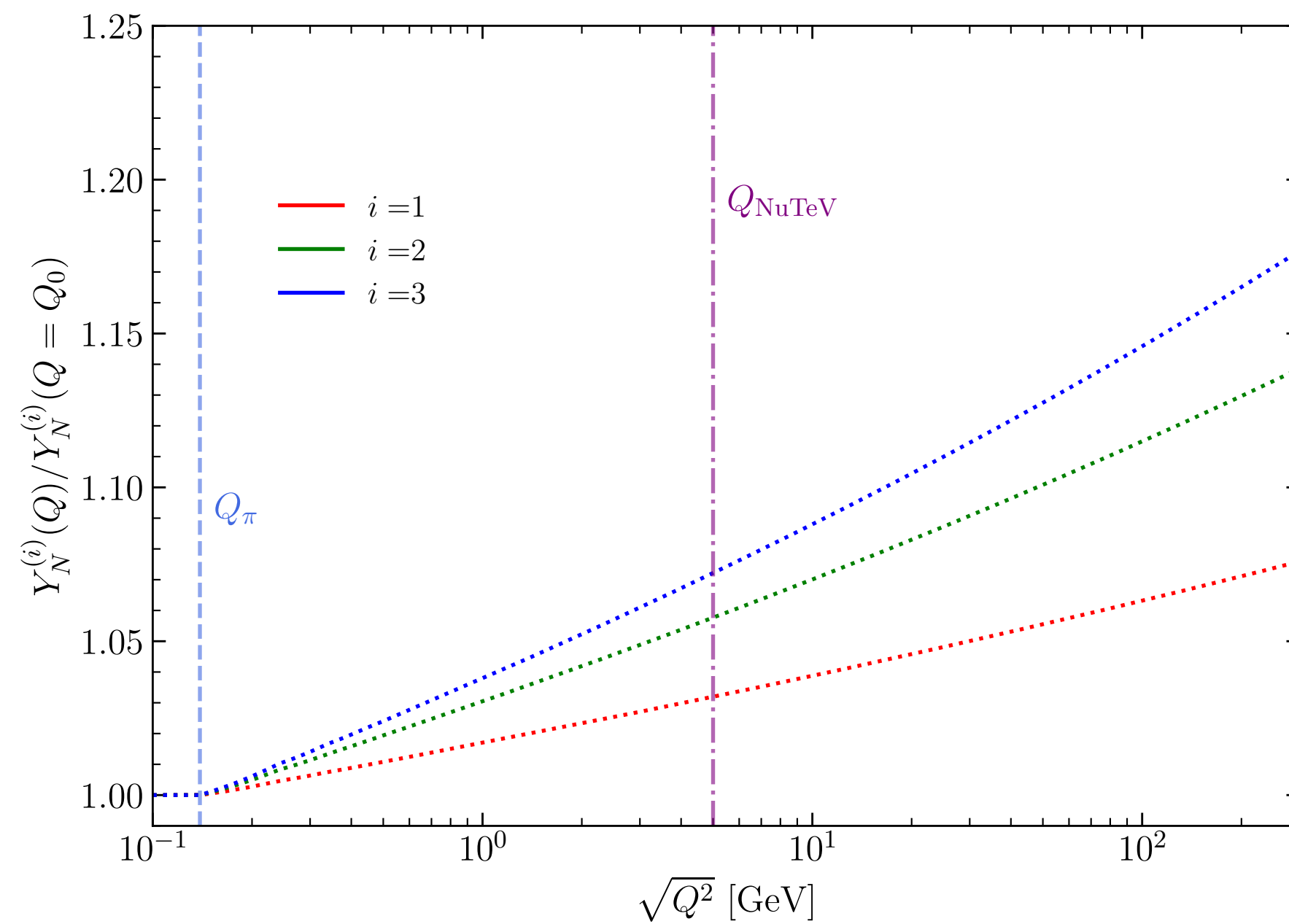
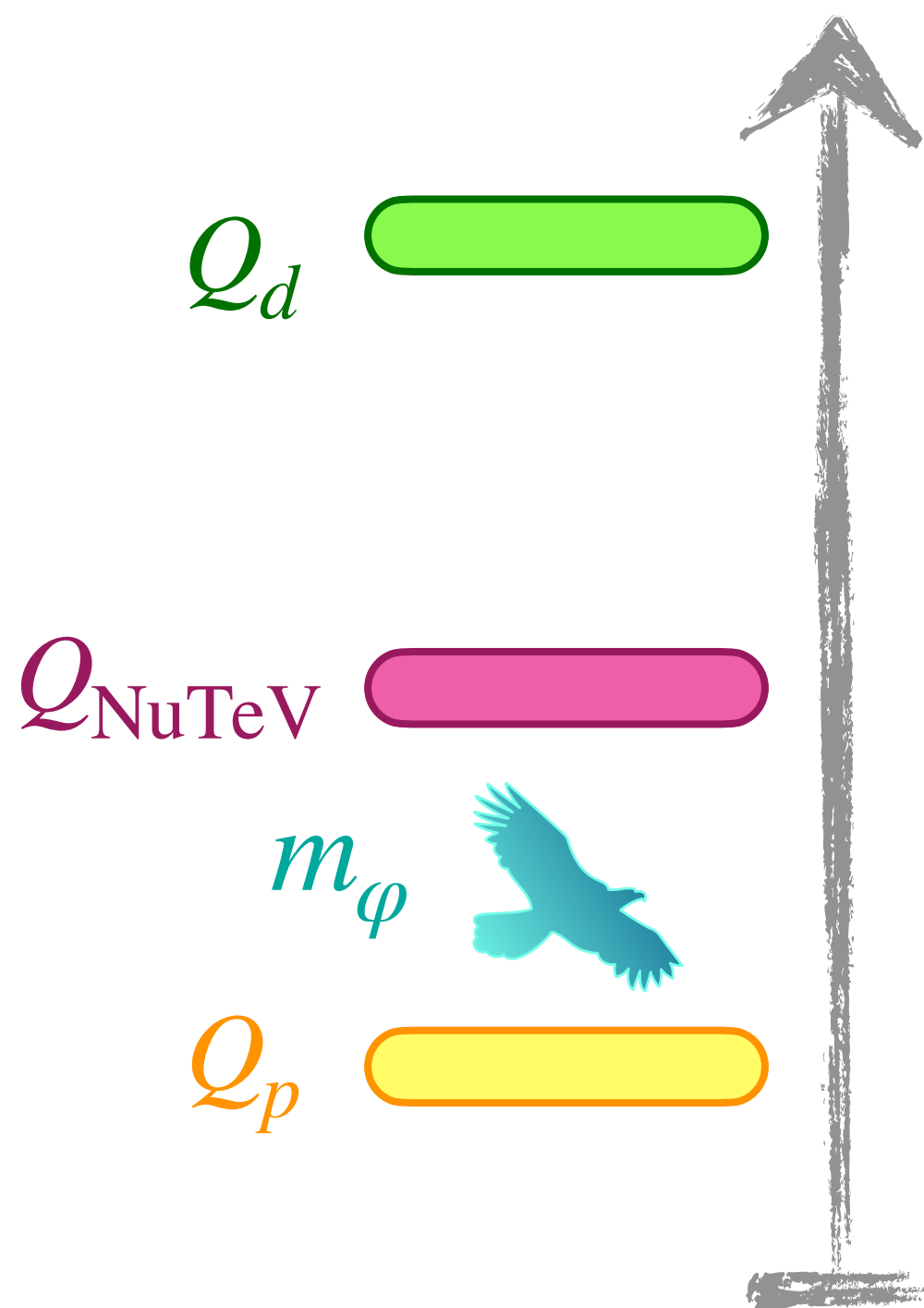
- Two event has the high likelihood of being tau neutrino [[arXiv:2011.03561](#)]

Category	$E < 60 \text{ TeV}$	$E > 60 \text{ TeV}$	Total
Total Events	42	60	102
Up	19	21	40
Down	23	39	62
Cascade	30	41	71
Track	10	17	27
Double Cascade	2	2	4

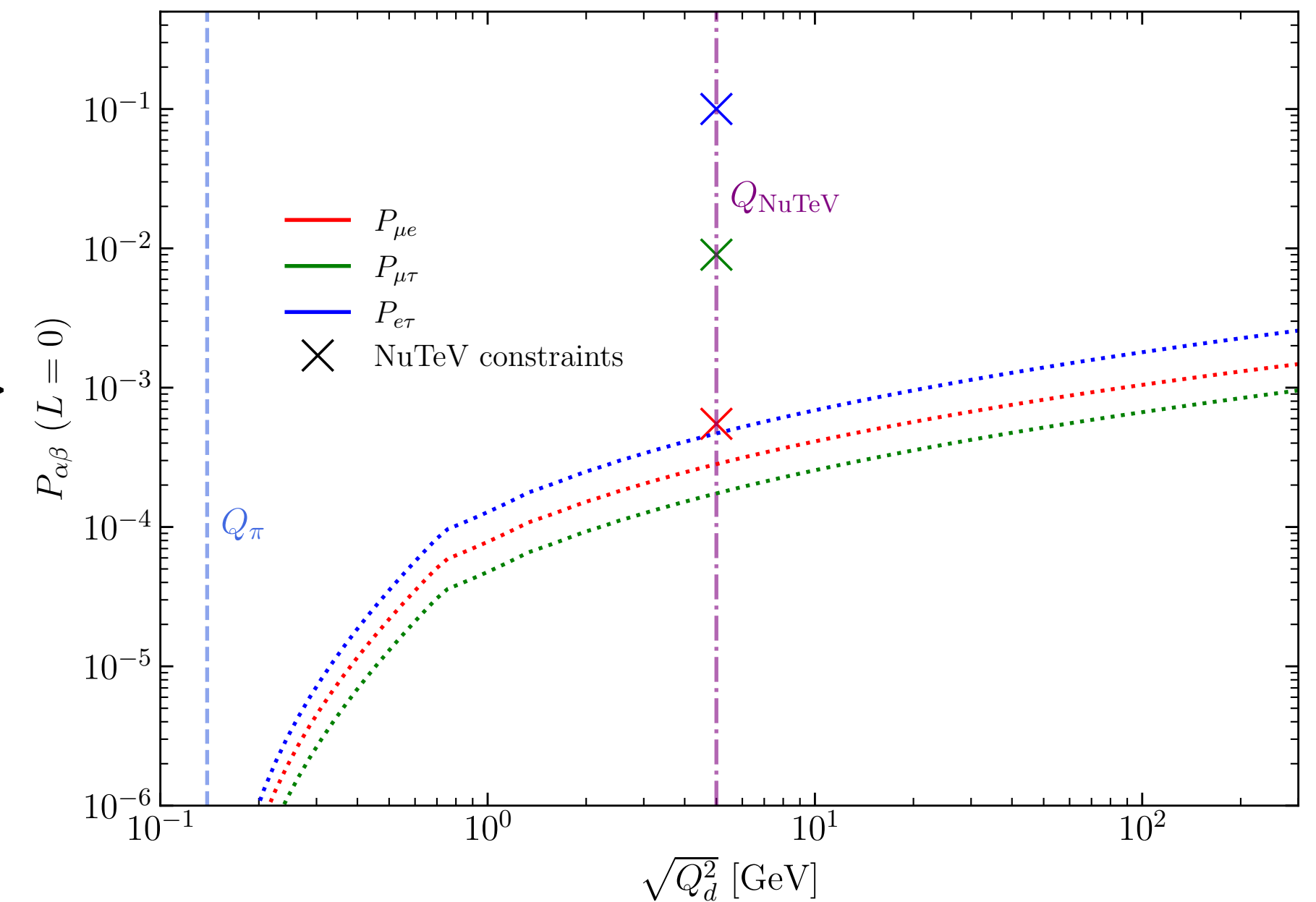
- Tauness of two events, $P_{\tau}^{DD} = 97.5 \%$, $P_{\tau}^{BB} = 76 \%$

RG Running

Case 1: $m_{\varphi, N} \sim Q_p = Q_\pi = m_\pi$ (pion decay production)



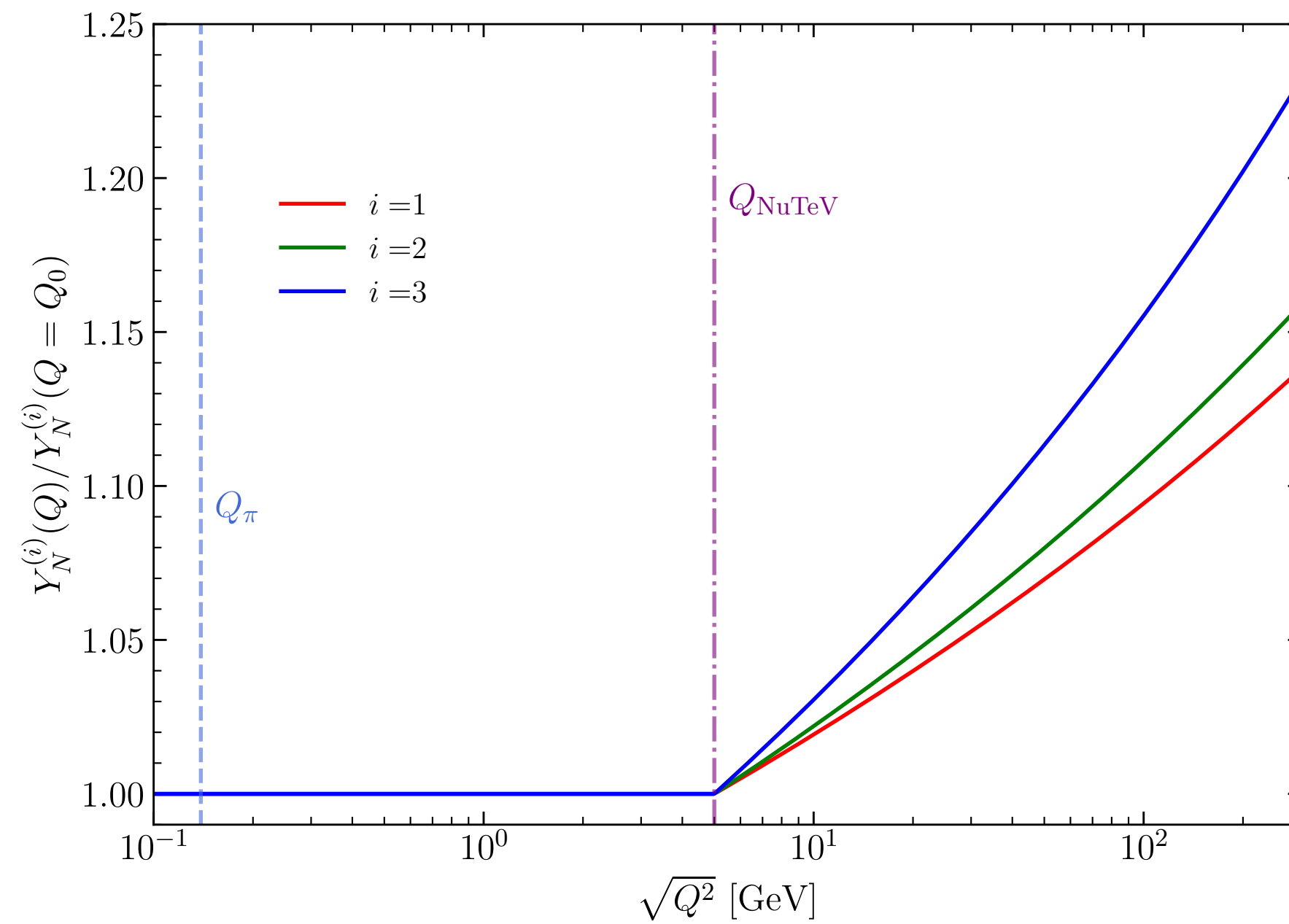
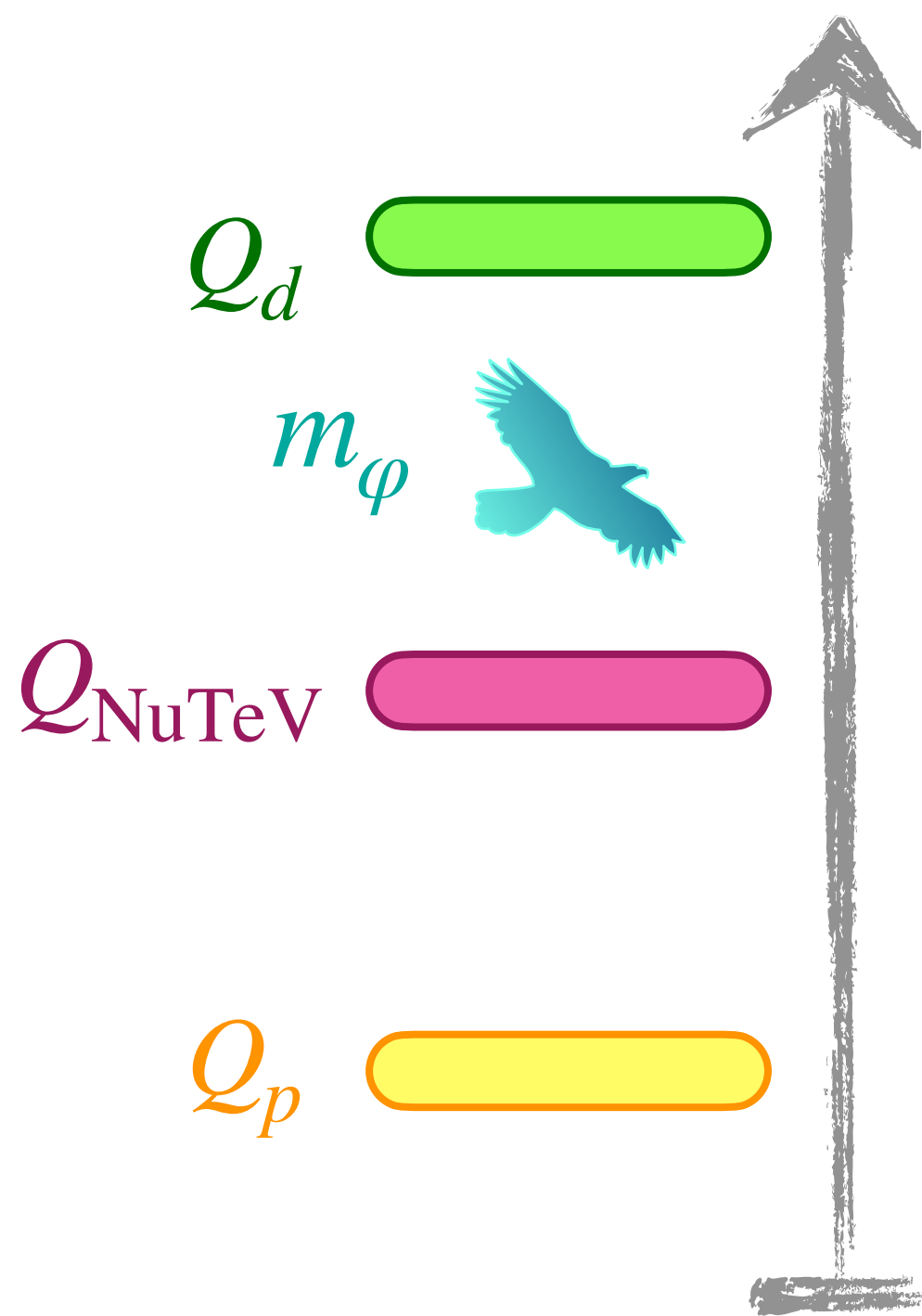
Y_N running



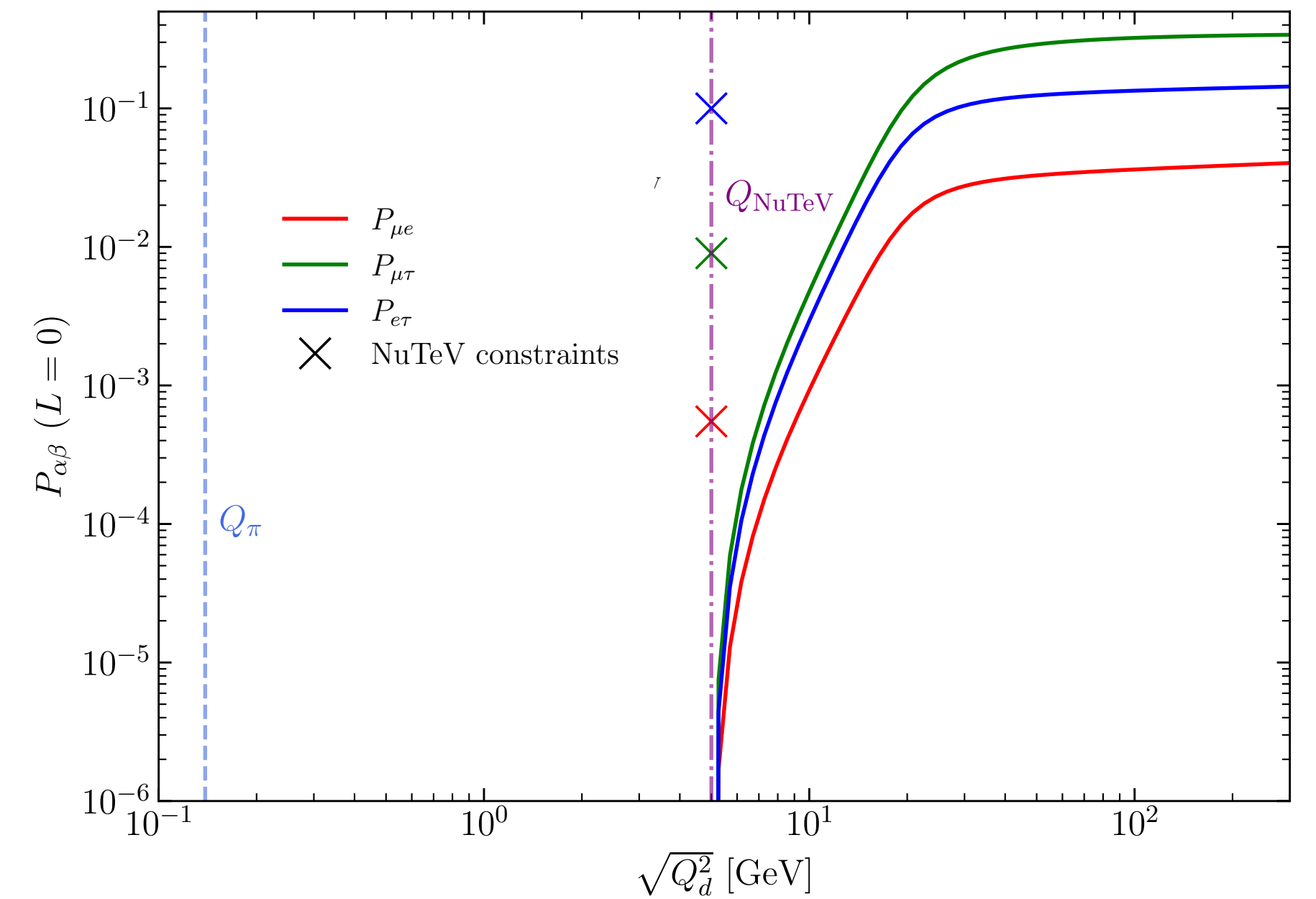
Oscillation at zero baseline

RG Running

Case 2: $m_{\varphi, N} \gtrsim Q_{\text{NuTeV}}$

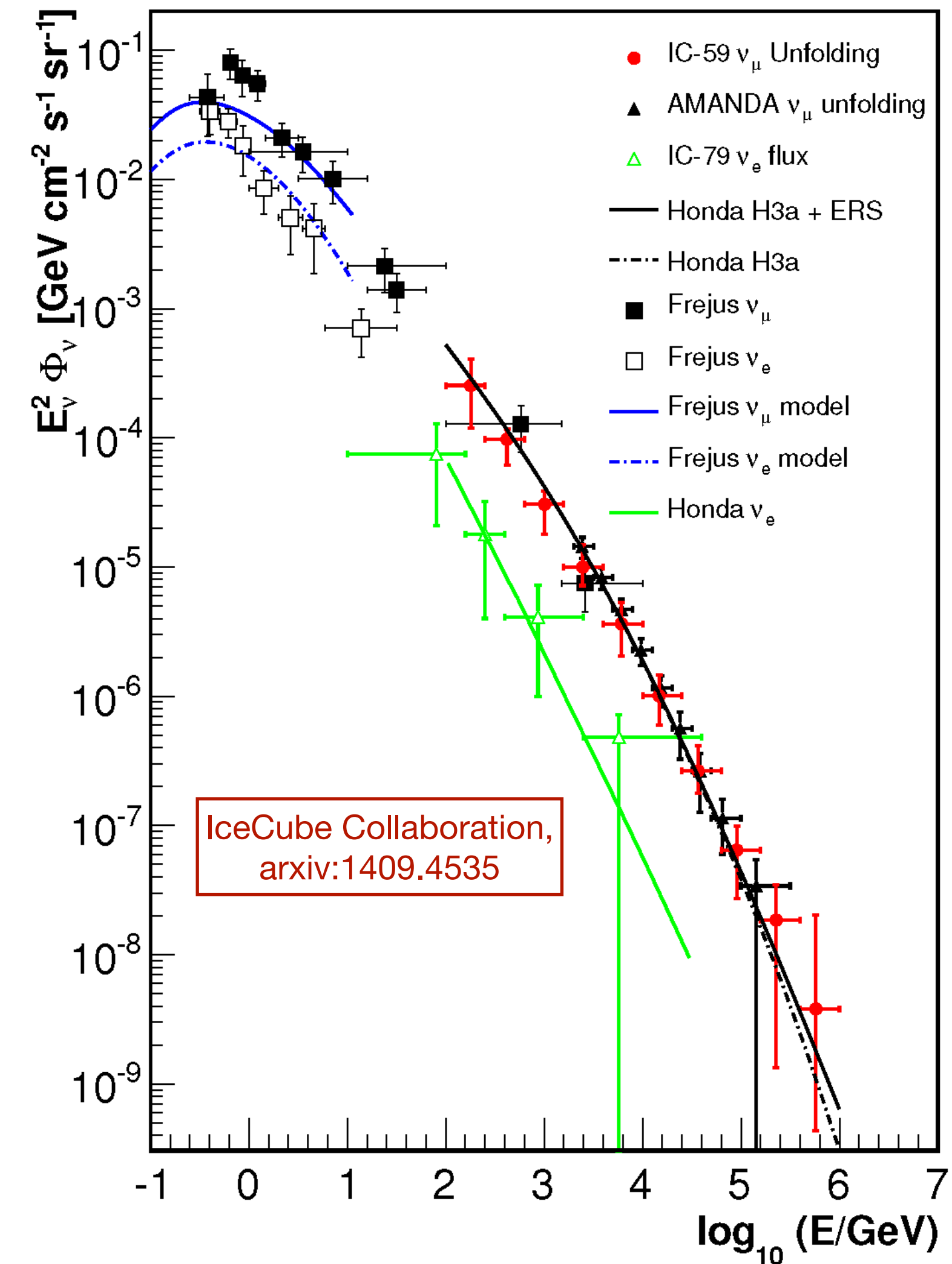
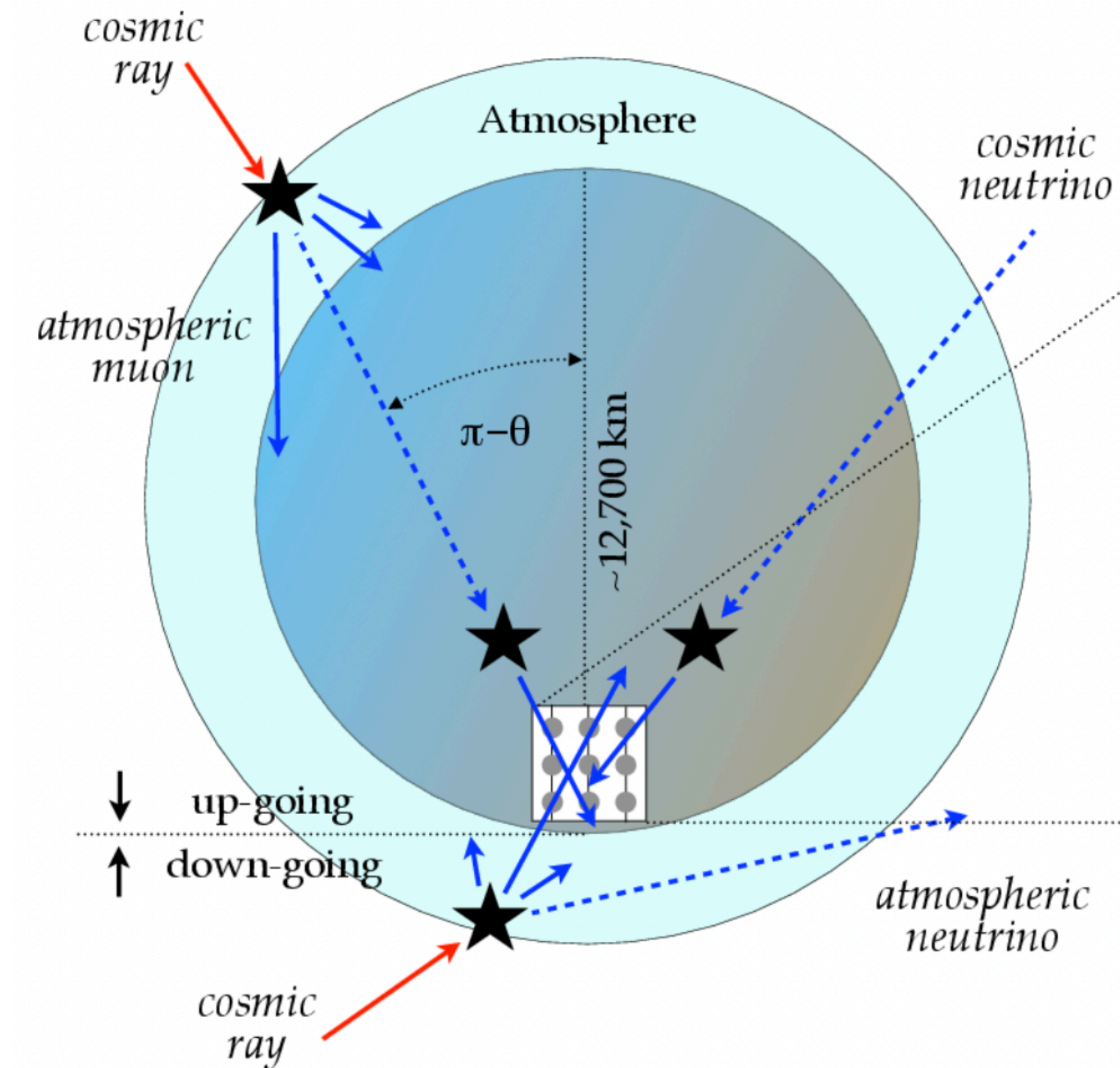


Y_N running



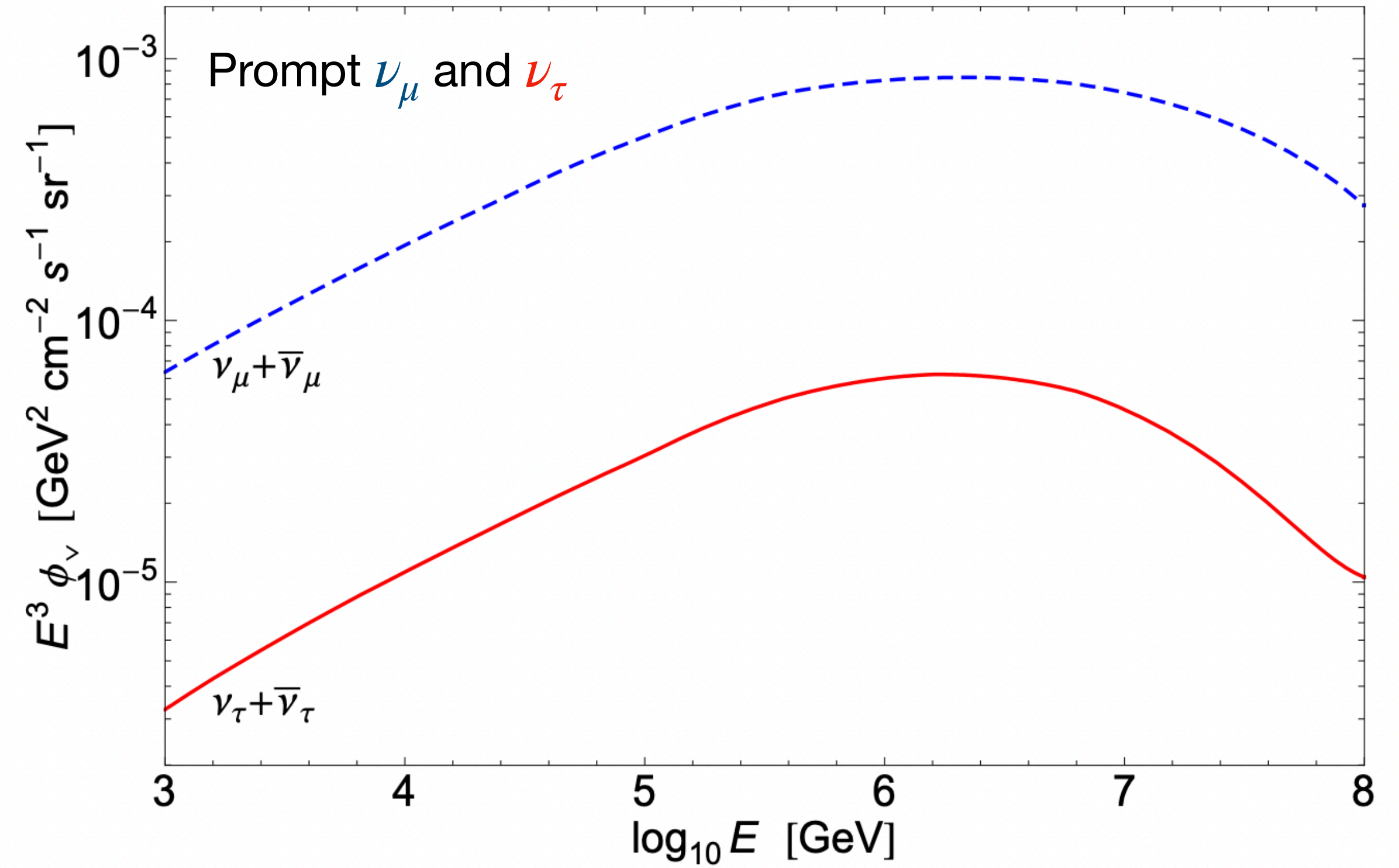
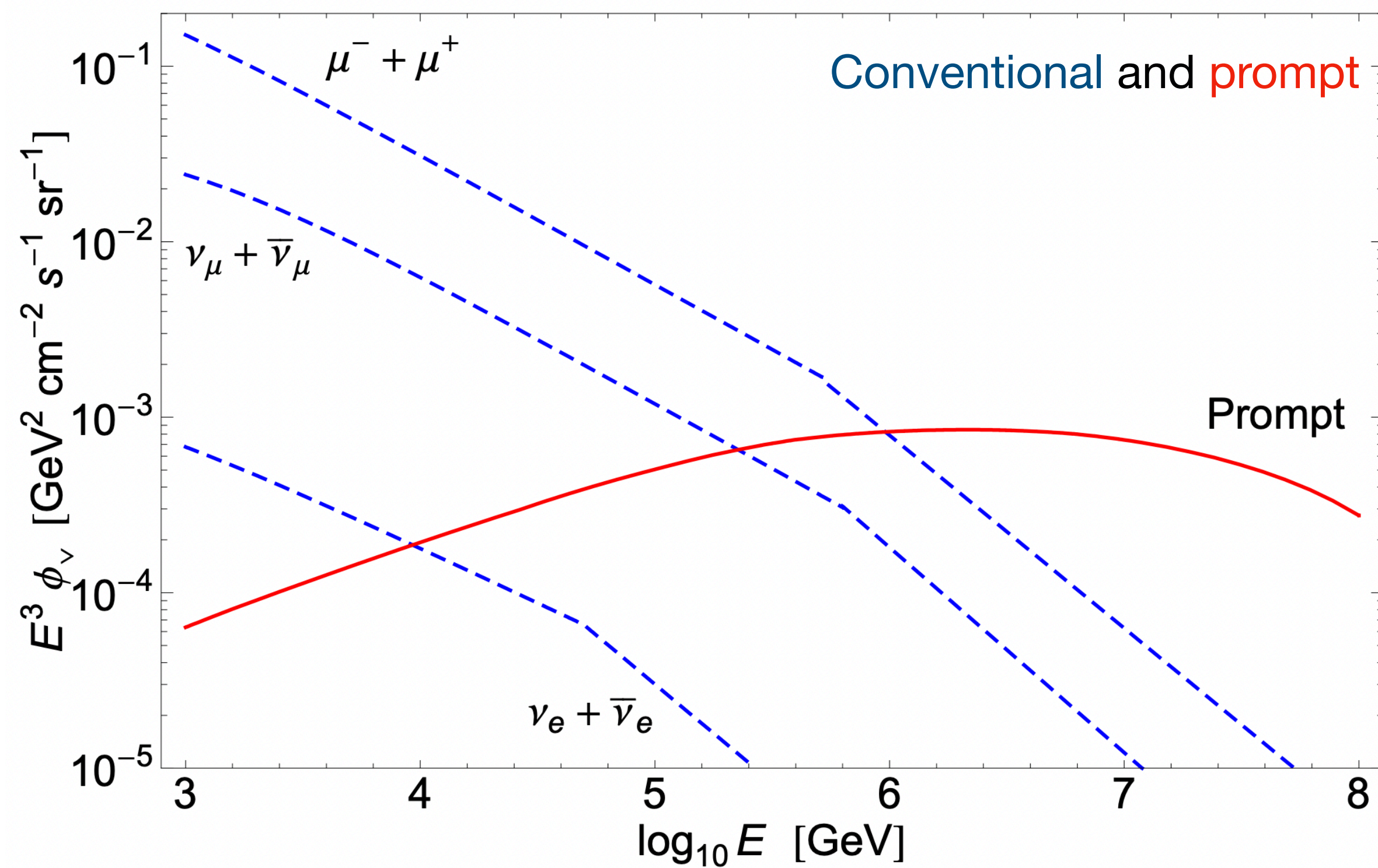
Oscillation at zero baseline

Atmospheric Muon Neutrinos at IceCube

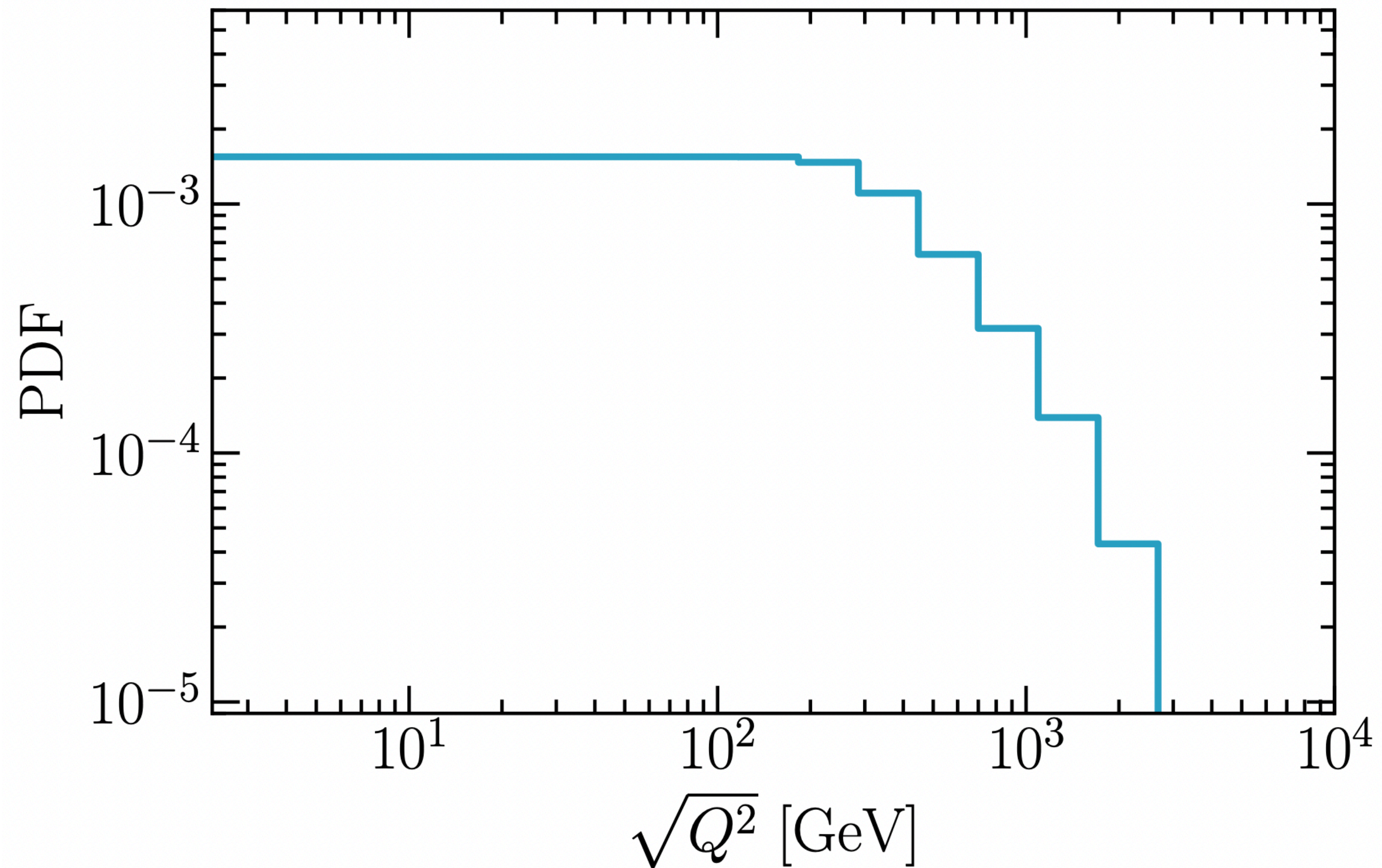


Atmospheric Neutrinos

Enberg, Reno, Sarcevic, arXiv:0806.0418



$\sqrt{Q^2}$ Distribution at HESE Sample



Atmospheric Neutrinos at IceCube

◆ Atmospheric neutrinos:
CR + atm. $\rightarrow \pi, K, \dots \rightarrow \nu + \ell^\pm$

◆ ~99% ν are muon neutrinos

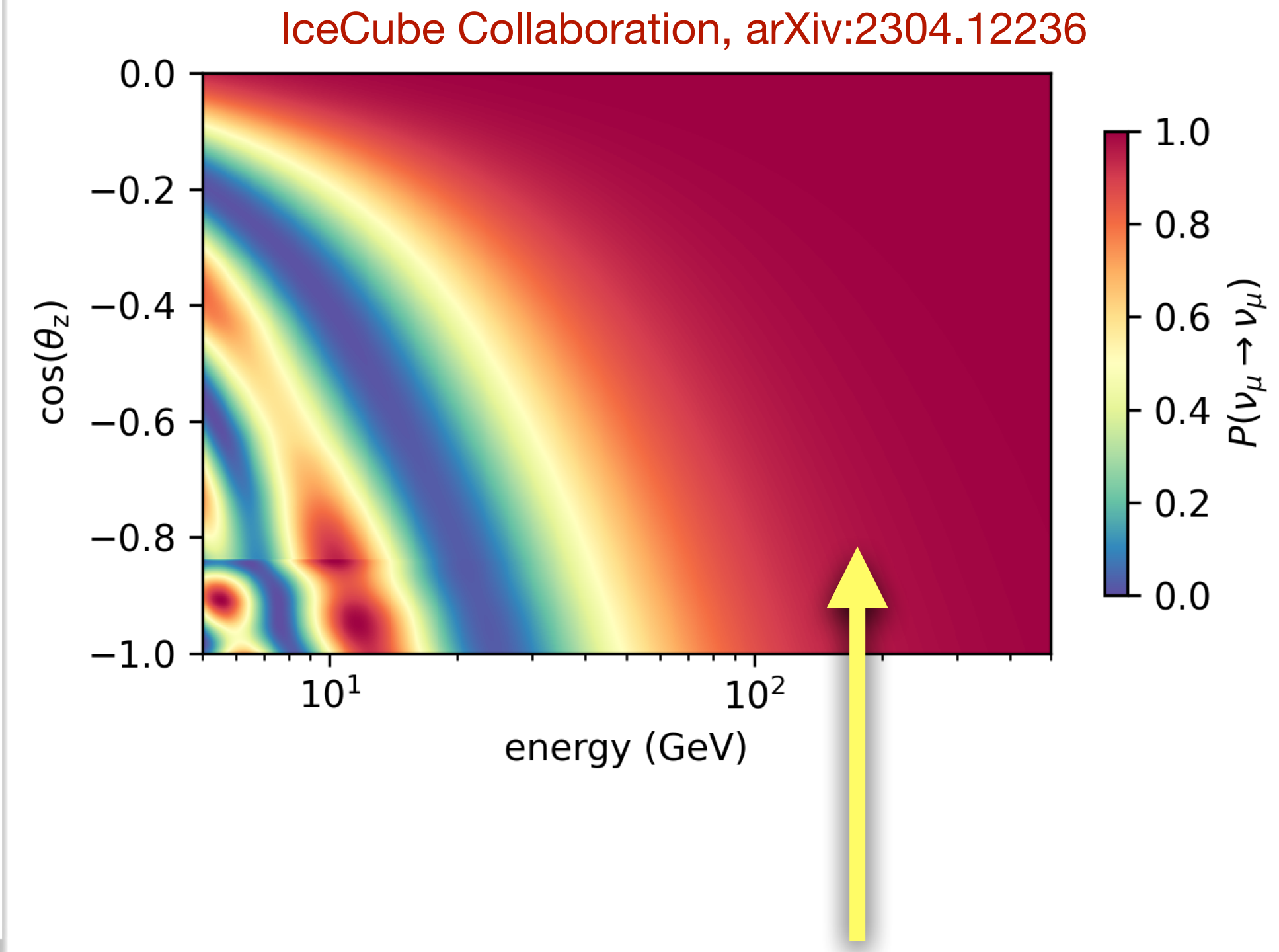
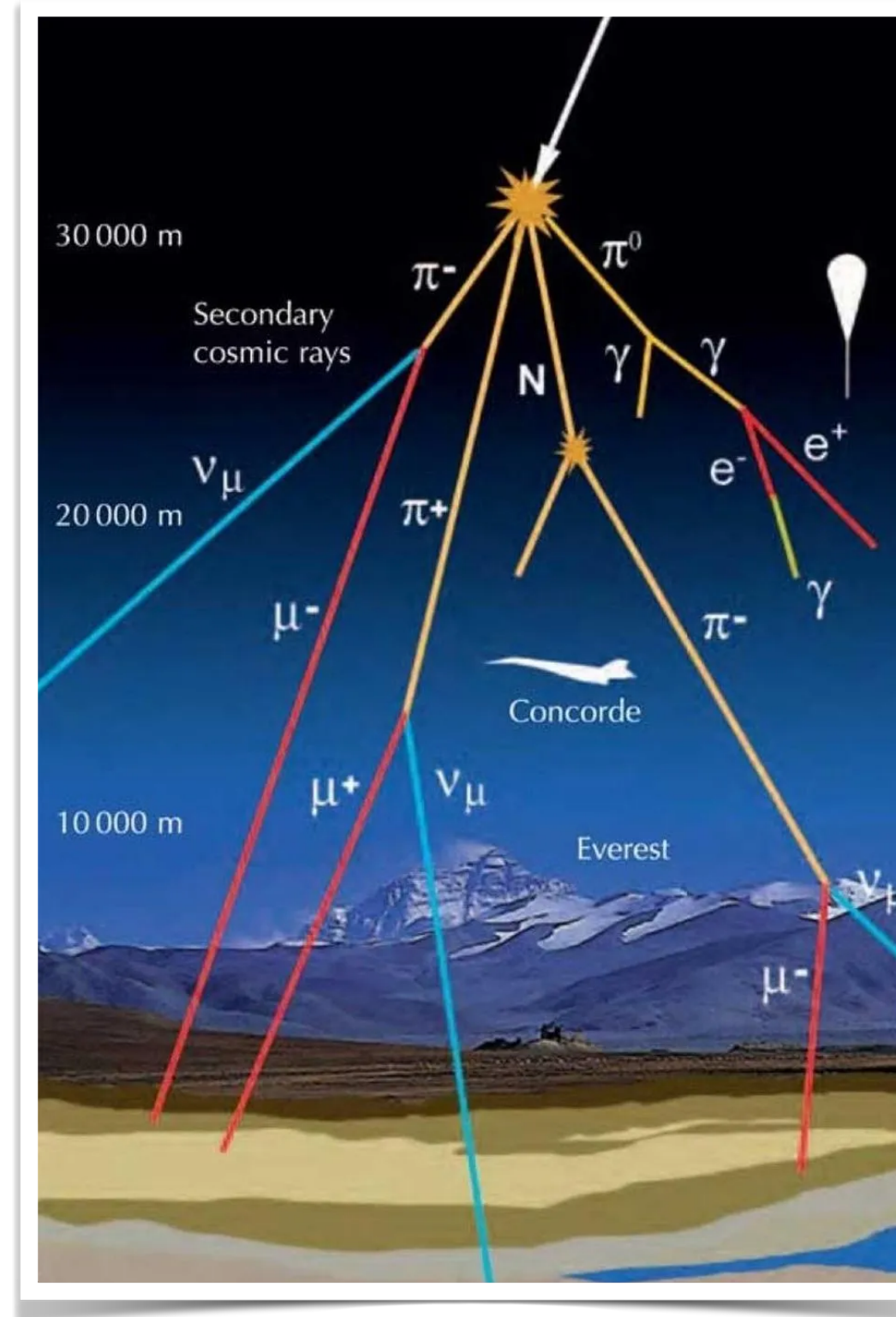
◆ ‘atmosphere to IceCube’
baseline $L \sim 10$ km to 10000 km

◆ $E_\nu = 100$ GeV, $L_{osc} = 10^5$ km

where transition probability $P \propto \sin^2 \left(\frac{L}{L_{osc}} \right)$

◆ Baseline is **not long enough to oscillate**
at high energy $E_\nu > 100$ GeV

◆ RG running may allow atmospheric $\nu_{e,\mu} \rightarrow \nu_\tau$



Standard oscillation \rightarrow no flavor transition
BSM running \rightarrow flavor transition leading to tau neutrino appearance

Tau Neutrino Appearance at IceCube: Astrophysical Neutrinos

● Flavor conversion, $P_{\alpha\beta}(L \gg L_{osc}) = \sum_i \left| U_{\alpha i}(Q_p^2) \right|^2 \left| U_{\beta i}(Q_d^2) \right|^2$

● This $P_{\alpha\beta}(L \gg L_{osc})$ enters the BSM weight,

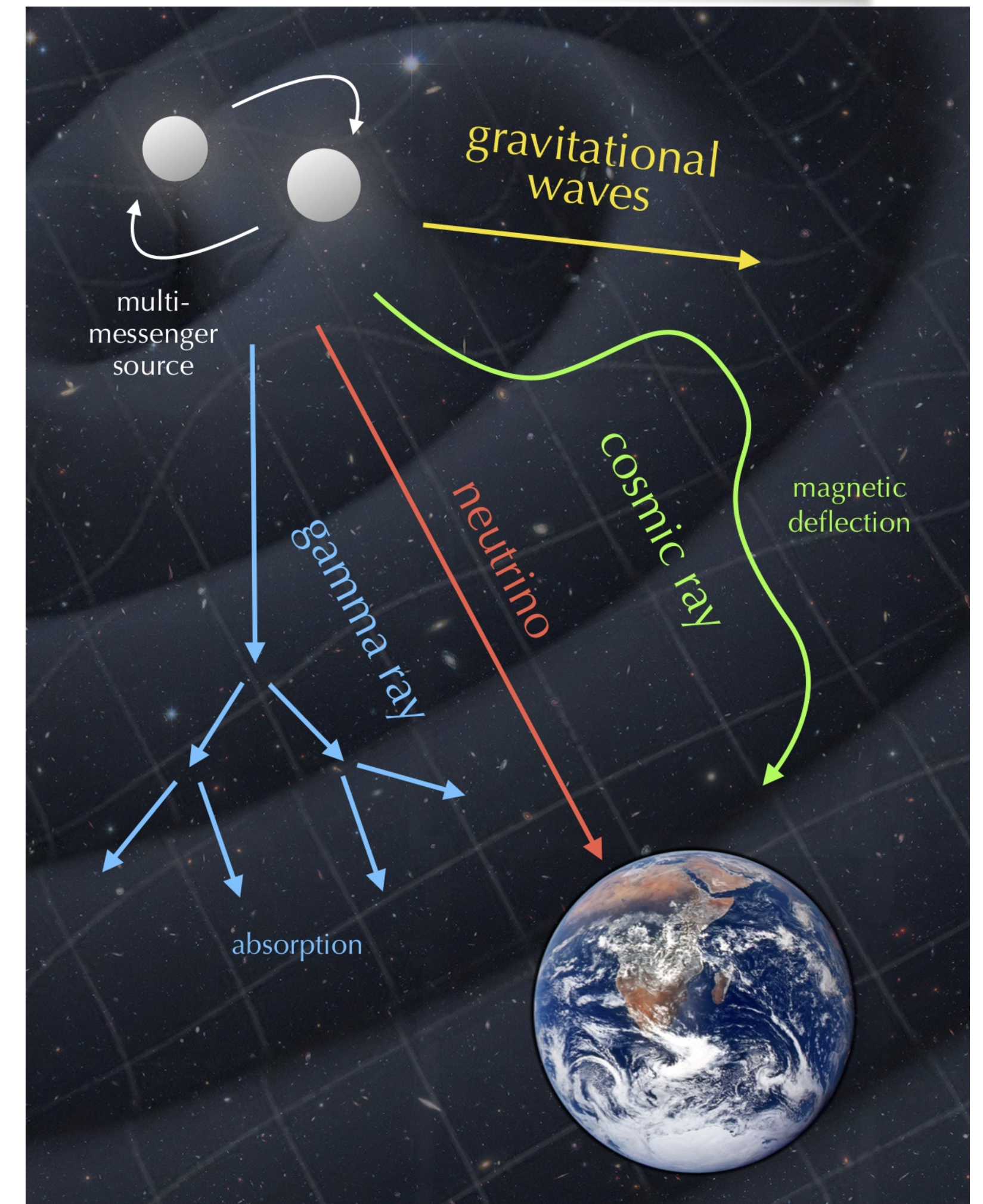
$$w_{\alpha \rightarrow \beta}^{\text{BSM}}(E_\nu) = \frac{\int_\nu P_{\alpha\beta}(E_\nu, Q^2) + \int_{\bar{\nu}} P_{\bar{\alpha} \rightarrow \bar{\beta}}(E_\nu, Q^2)}{\int_\nu + \int_{\bar{\nu}}}$$

● Flavor ratio on earth, $X_\beta^{\text{earth}} = P_{\alpha\beta} X_\alpha^{\text{prod}}$

where X_α^{prod} = astrophysical flavor ratio at the source

- * Pion decay:
 $(f_e : f_\mu : f_\tau)_S = 1 : 2 : 0$
- * Muon damped:
 $(f_e : f_\mu : f_\tau)_S = 0 : 1 : 0$
- * Neutron decay:
 $(f_e : f_\mu : f_\tau)_S = 1 : 0 : 0$

$$N_\tau^{\text{astro}} = T \int dE_\nu \epsilon_\tau(E_\nu) \Phi_{1\nu}^{\text{astro}}(E_\nu) \sigma_{\nu N}(E_\nu) \sum_{\alpha=e,\mu,\tau} w_{\alpha \rightarrow \tau}^{\text{BSM}}(E_\nu) X_\alpha^{\text{prod}}$$



IceCube Gen-2 and Fluctuation

<https://icecube-gen2.wisc.edu/science/publications/tdr/>

Table 5: Rate of upgoing tracks in the IceCube-Gen2 optical array [yr^{-1}]

Deposited energy	astrophysical- ν		atmospheric- ν	
	IceCube	IceCube-Gen2	IceCube	IceCube-Gen2
>1 TeV	1.2×10^2	4.0×10^2	2.9×10^5	3.4×10^5
>10 TeV	16	59	2.4×10^2	5.3×10^2
>100 TeV	1.4	3.8	0.67	0.99

Table 6: Rate of astrophysical neutrinos of various event categories (optical array only) [yr^{-1}]

Deposited energy	shower-type events		double cascades		starting tracks	
	IceCube	IC-Gen2	IceCube	IC-Gen2	IceCube	IC-Gen2
> 10 TeV	6.0	50	0.36	2.4	1.4	21
> 100 TeV	3.6	32	0.3	1.9	0.81	13

Expected number of events n	Poisson probability of observing 2n events
2	0.090
5	0.018
10	0.002