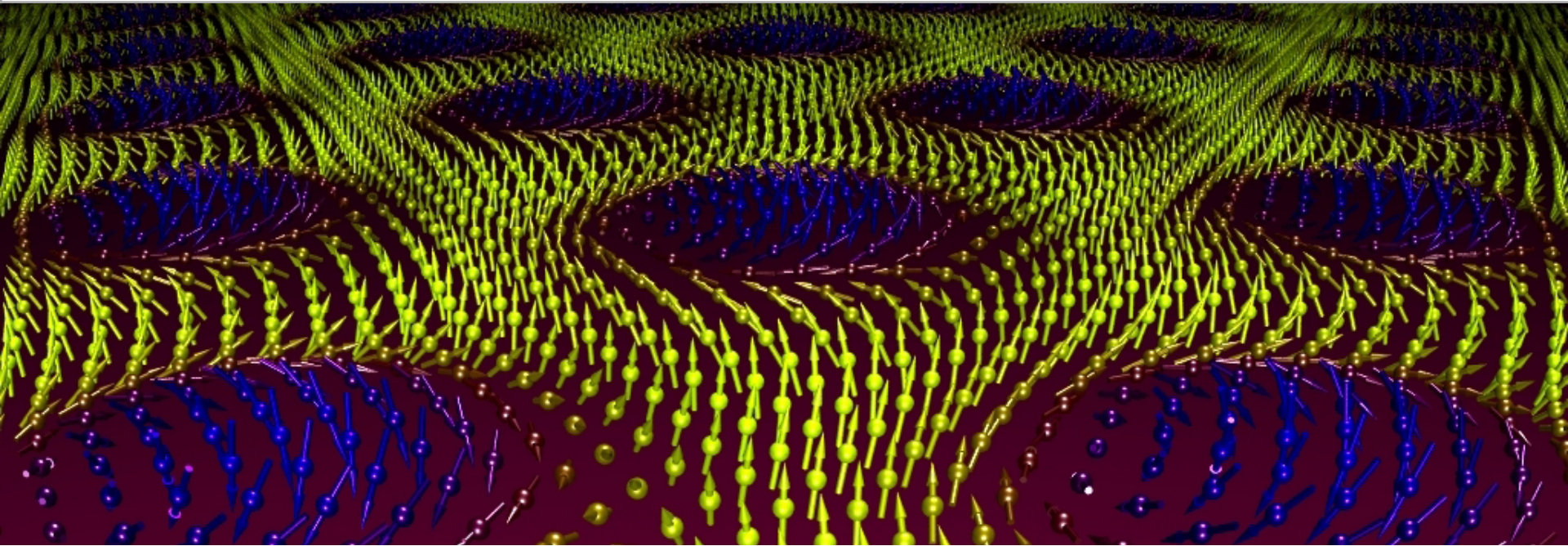


Cosmological Aspects of Axion-Monopole Interactions

© [Artificial magnetic monopoles discovered - phys.org]



Hengameh Bagherian

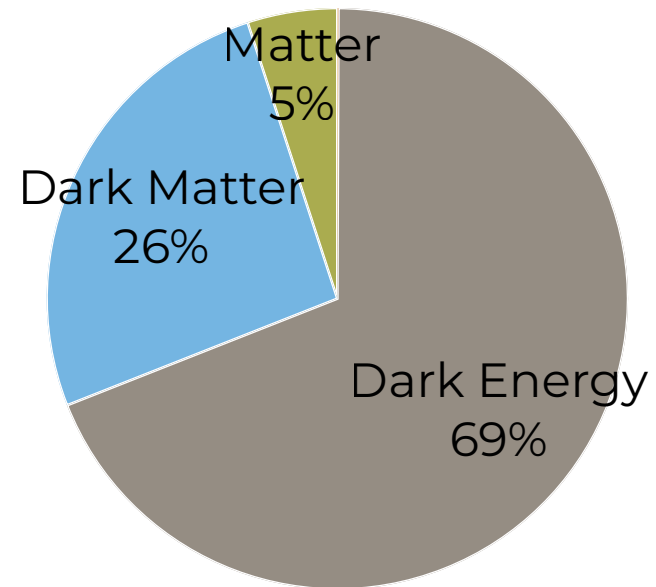
EFI and Leinweber Institute, University of Chicago

Pheno 2026 Symposium | May, 2026

[H. Bagherian, M. Reece, L.-T. Wang, and H. Xiao, [arXiv: 2026.XXXXX](#)]

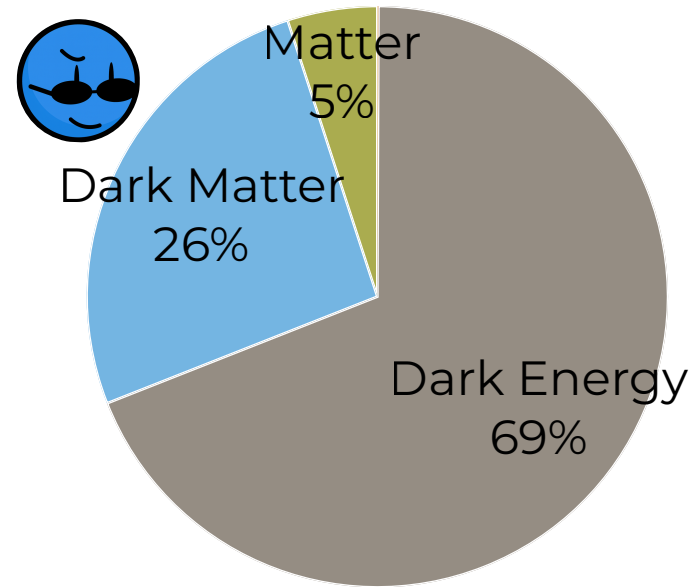
Dark sector today

- **Dark matter** candidates:
WIMPS, **Axions**, PBH, etc.
- **Dark energy** candidates:
Cosmological constant (Λ),
quintessence (**light axion**, scalar field), etc.
- **Question:**
Are these sectors **independent**?
What if they are **interacting**?
What if they have the same origin?
DESI DR2 data gives hints of
interaction!



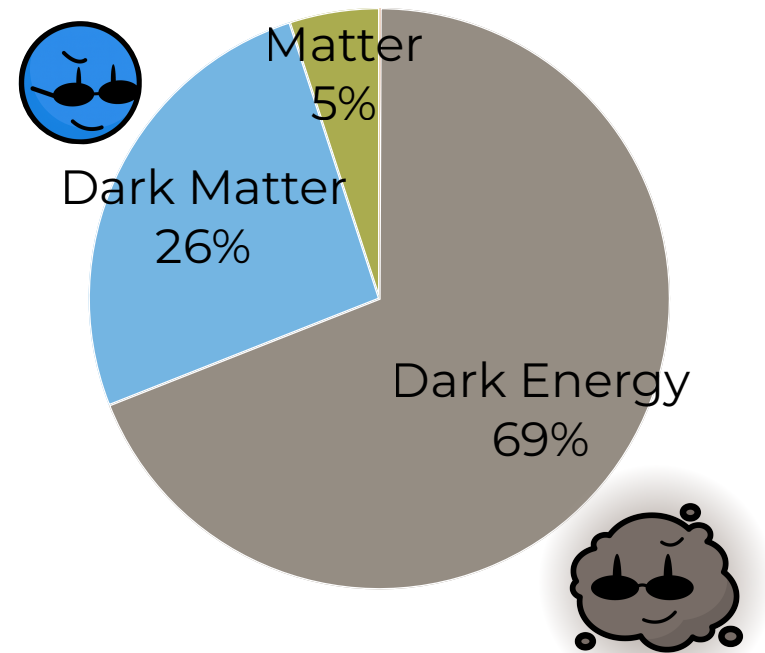
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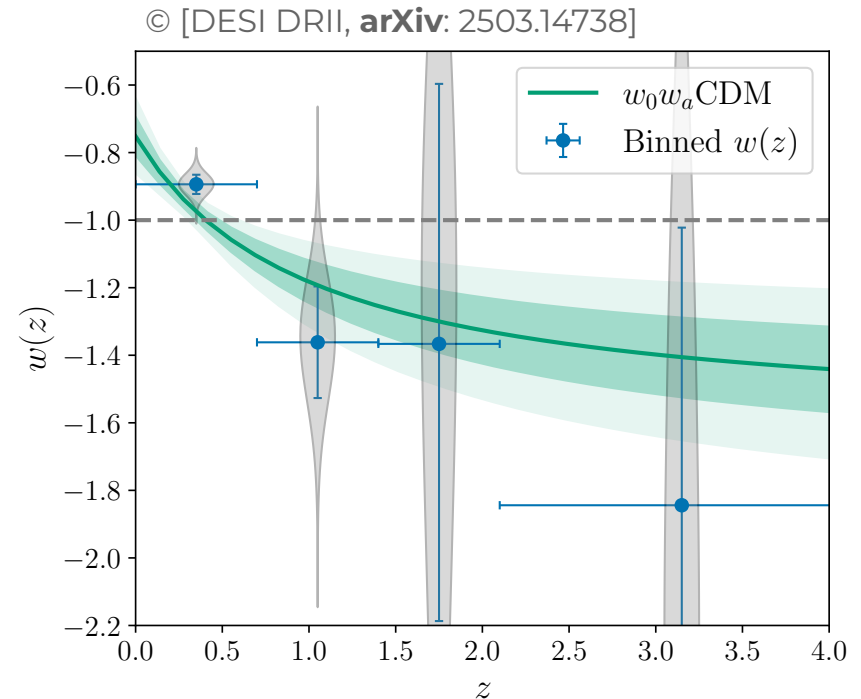
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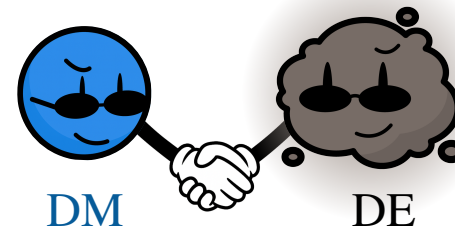
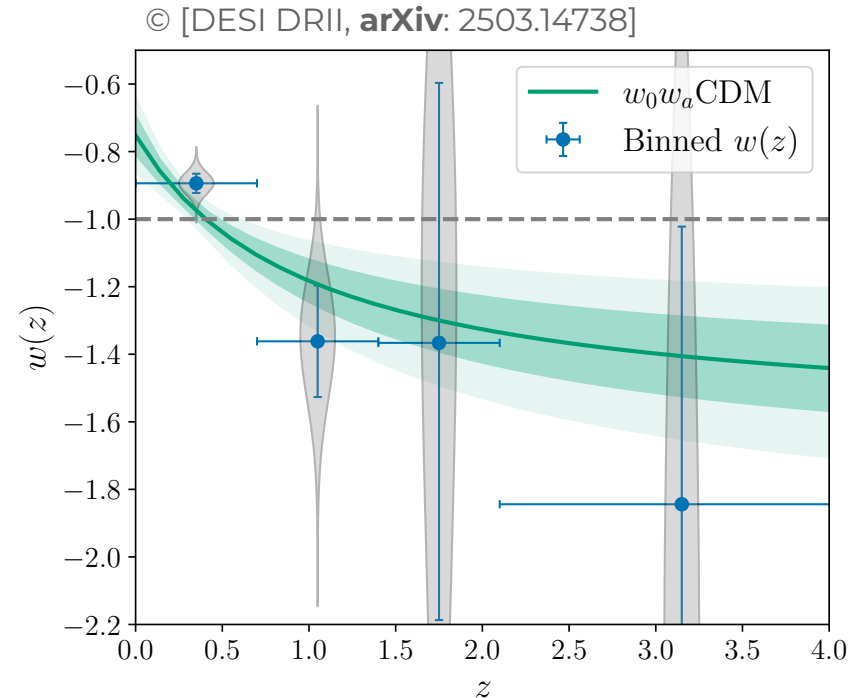
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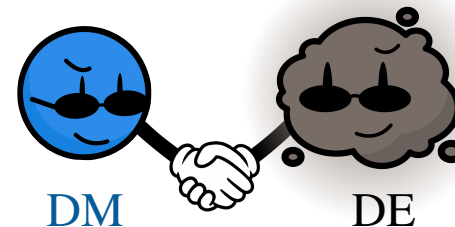
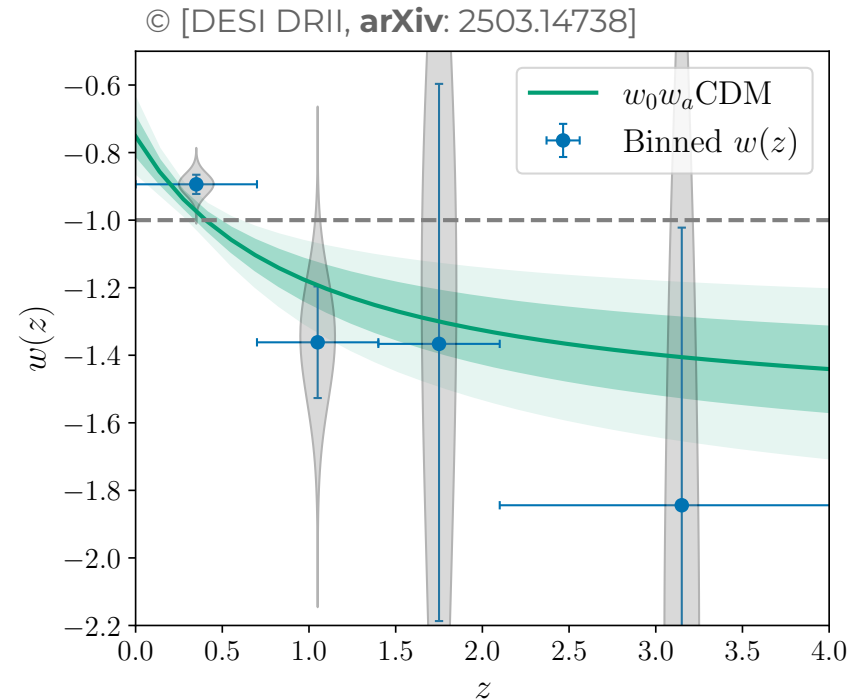
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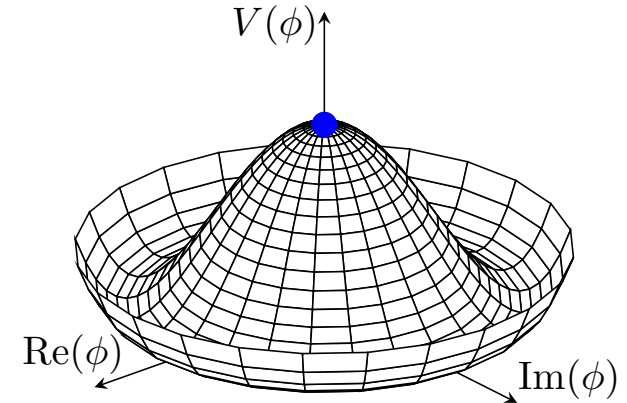
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I claim **axions naturally** provide an **interacting dark sector today!**

Axions can be both DM and DE

- Axions in this talk: a **scalar field** with **shift symmetry** + potential from **instanton effects**
- **Oscillating axions** with mass heavier than H_0 redshift as **dark matter**:
$$V(\theta) = m_a^2 f_a^2 (1 - \cos \theta)$$
$$\ddot{\theta} + 3H\dot{\theta} + m_a^2 \sin \theta \simeq 0$$
- **Ultra light axion** with mass smaller than H_0 behaves like **dark energy**
- Axions & Gauge Fields: **couple to monopoles** and detection signals (**birefringence**)



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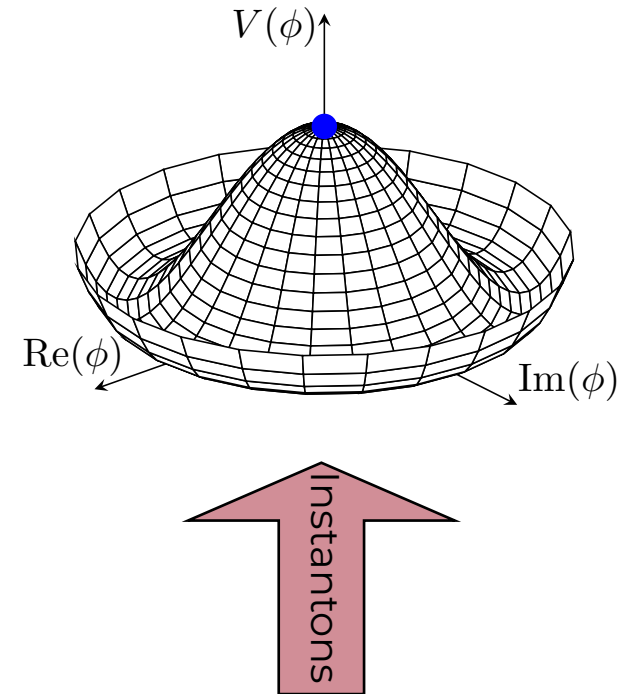
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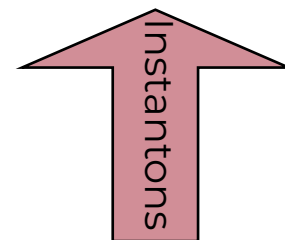
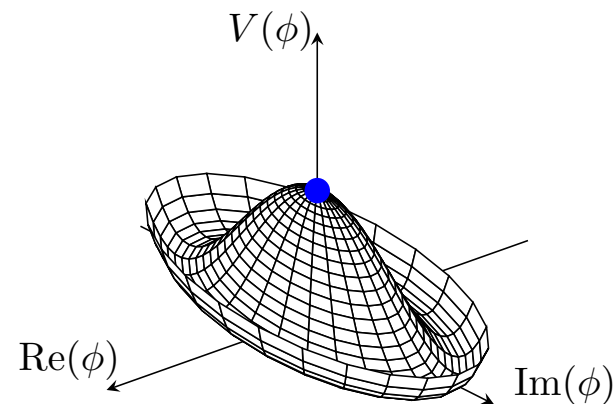
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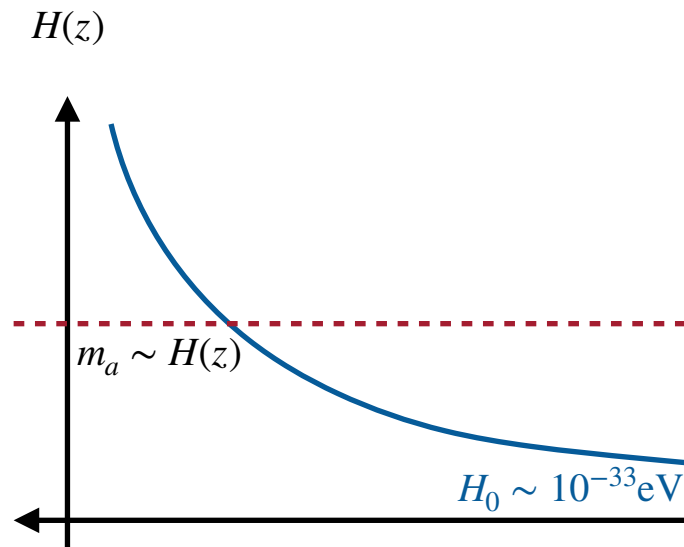
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[Marsh, Phys. Rept. 643 (2016), 1-79]

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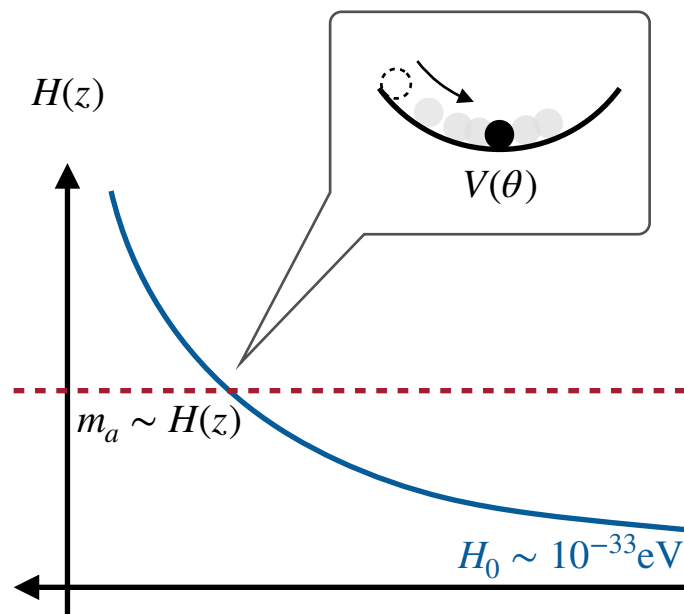
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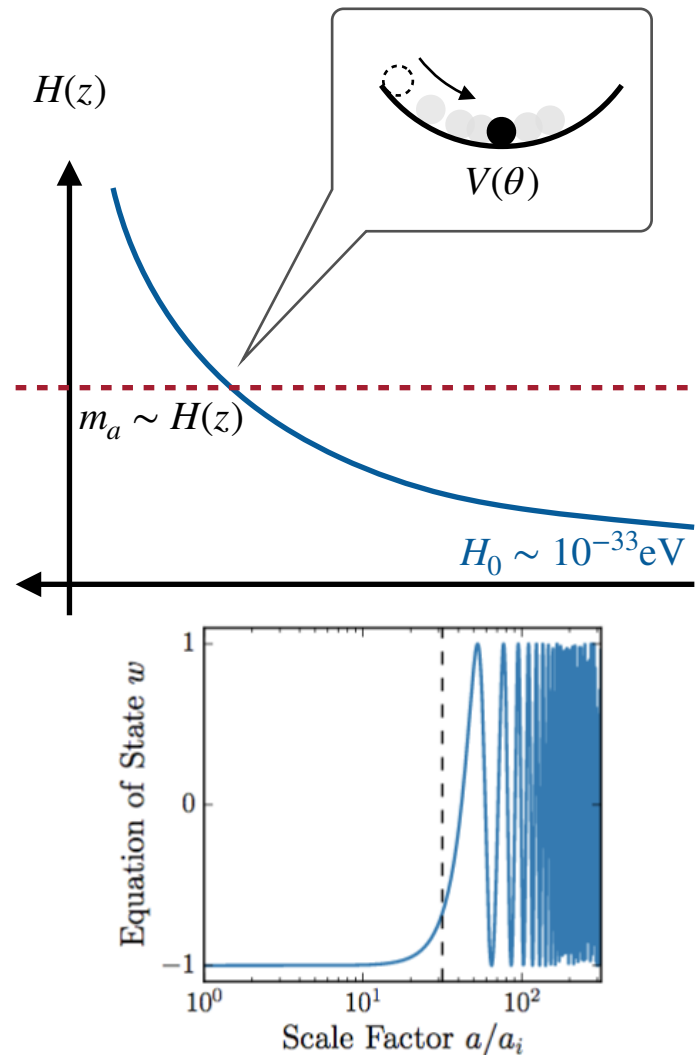
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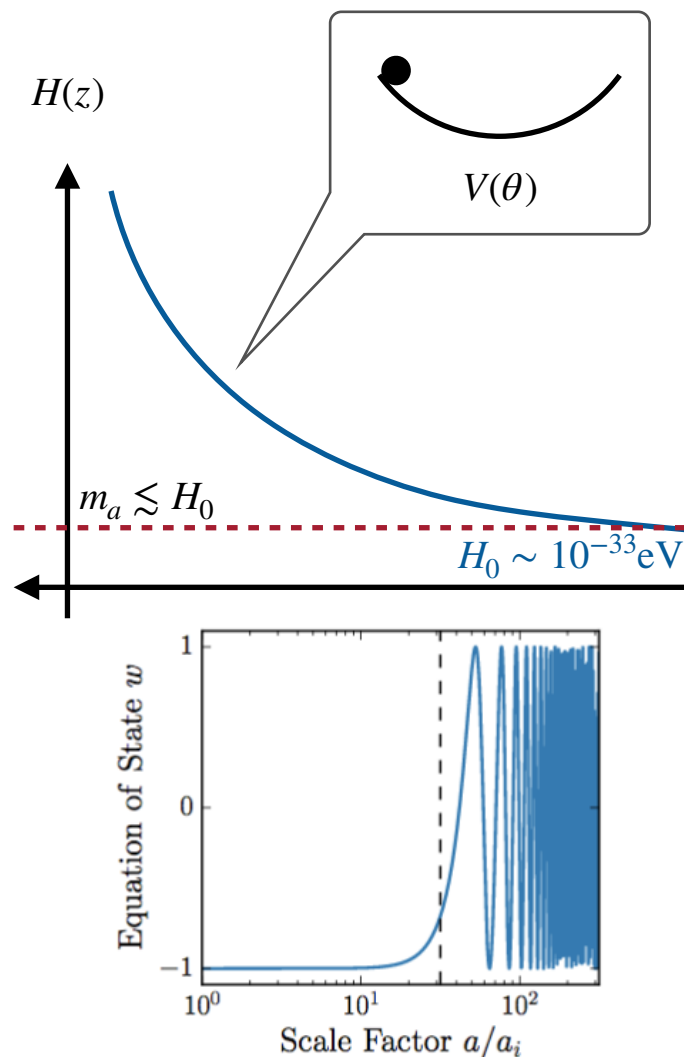
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Outline I

- **Axions & Gauge Fields**

- 't Hooft- Polyakov **monopoles** as topological defects (**DM**)
- **Witten effect** (Why axions couple to monopoles)
- 't Hooft-Polyakov **charge** + **mass** in axion background (**DM-DE interaction**)

- What is a **natural** DM-DE interaction?

- **Potential cosmological signatures**

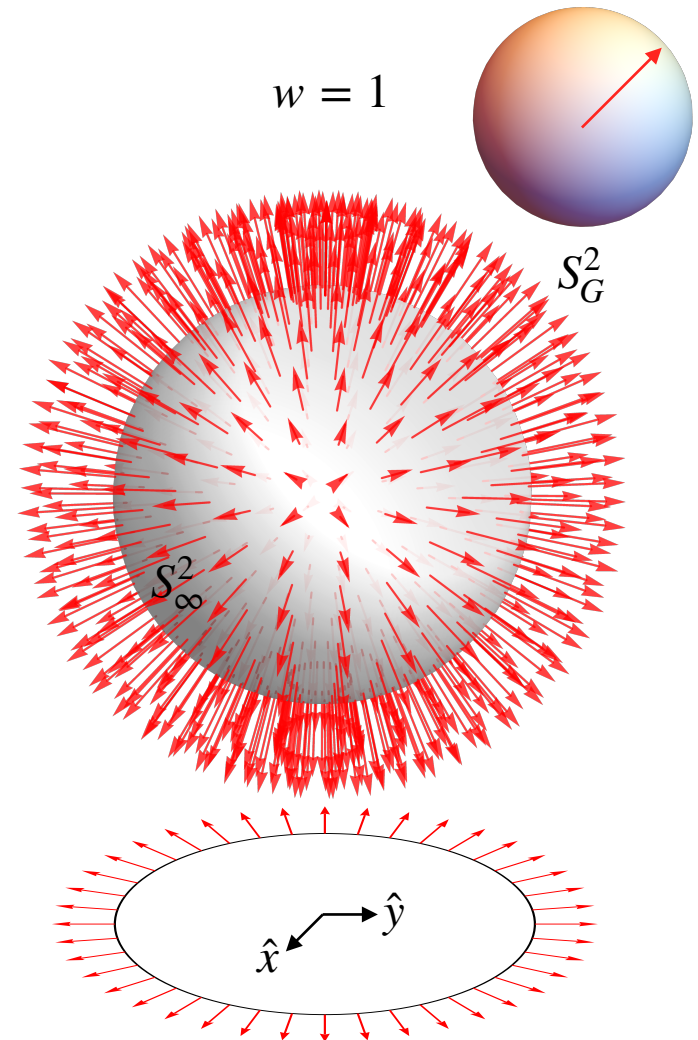
- Evolving DM-DE interaction (DESI)
- Cosmic Birefringence (Ask me after the talk)

Minimal example: 't Hooft-Polyakov

't Hooft-Polyakov monopole arises in the **Georgi-Glashow model**:

$$\mathcal{L} \supset -\frac{1}{4}G_{\mu\nu}^a G^{\mu\nu,a} + \frac{1}{2}(D_\mu \sigma^a)^2 + V(\sigma)$$

- Spontaneous breaking of $SU(2) \rightarrow U(1)$
- Higgs triplet σ^a gets a vev to minimize $V(\phi) = -\lambda(\sigma^a \sigma^a - v^2)^2$
- Any finite energy configuration obeys $\sigma^a \sigma^a \rightarrow v^2$ (this is S_G^2) at $|\vec{r}| \rightarrow \infty$ (which is S_∞^2)
- Monopole solutions exist bc $\pi_2(S_G^2) = \mathbb{Z}$
- Monopole **charge**: $q_M = 4\pi/e_D w$ ($\in \mathbb{Z}$)
- Monopole **mass**: $m_M = v |q_M|$

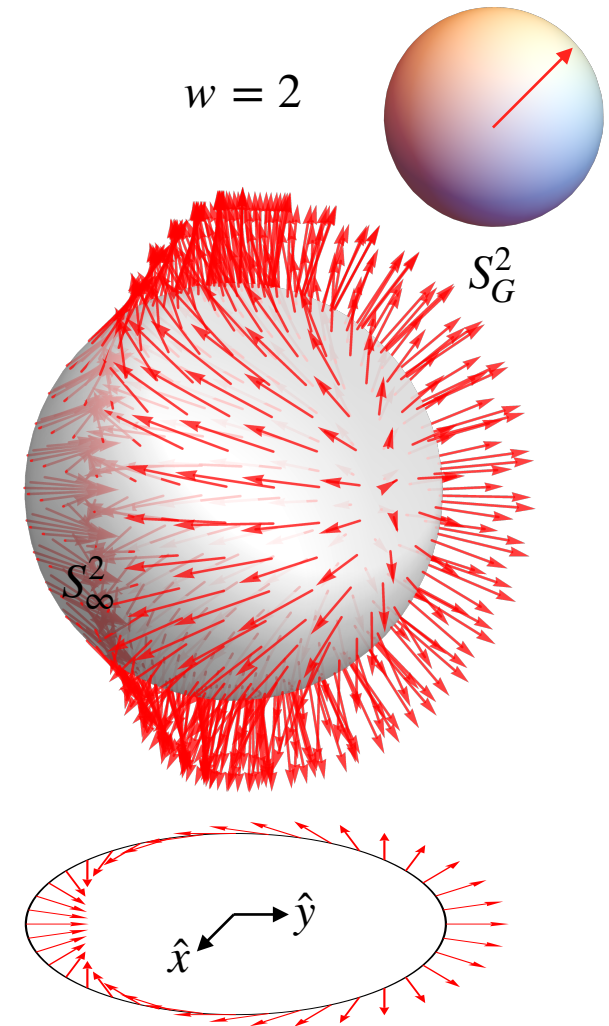


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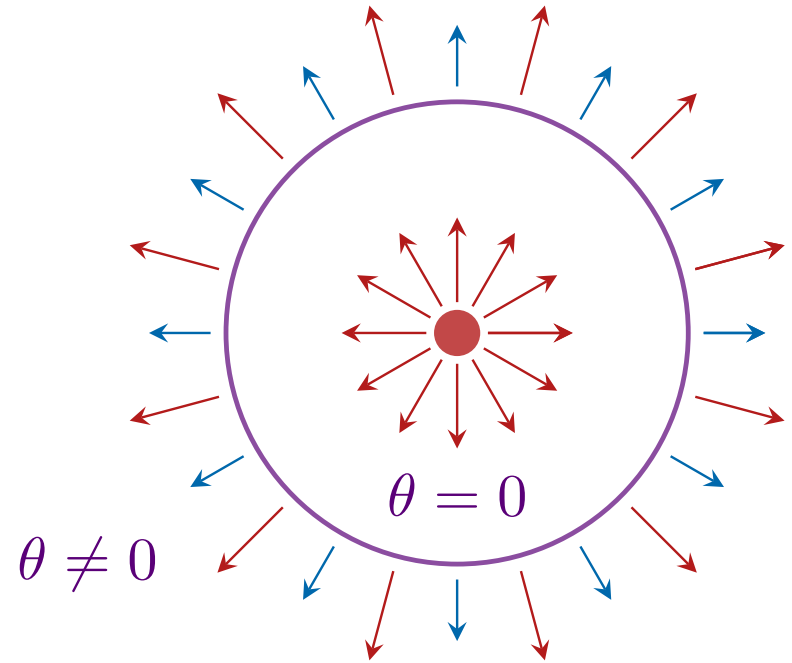


The Witten effect

Monopoles in the presence of a non-zero θ term ($\mathcal{L} \supset \theta F\tilde{F}$)

acquire electric charge:

- Put a magnetic monopole inside a region, surrounded by $\theta \neq 0$
- Monopole magnetic field **induces an electric field** when θ changes.
- Take size $\rightarrow 0$ limit, because electric charge is independent of the inner region size.



$$\vec{\nabla} \cdot \vec{E} = -\frac{\alpha_D}{2\pi} (\vec{\nabla} \theta \cdot \vec{B})$$

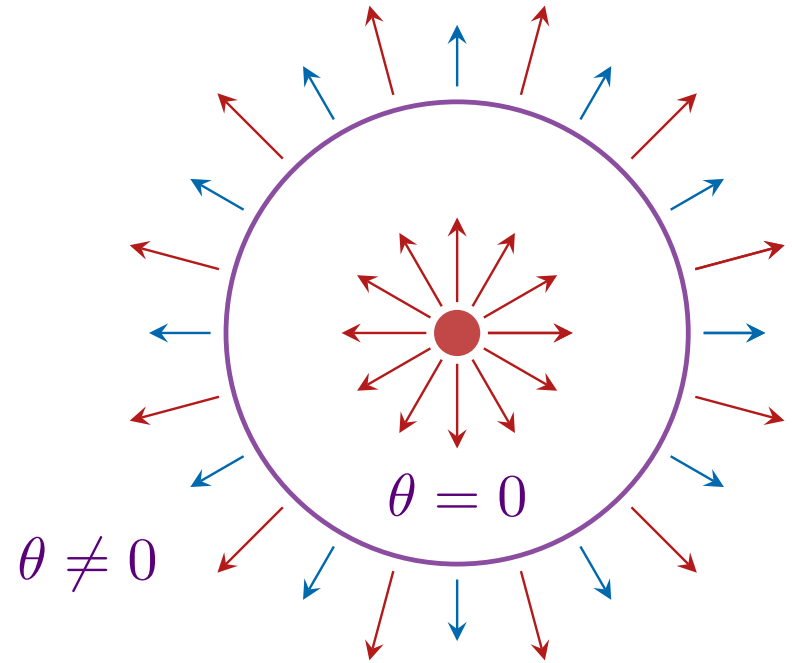
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How does this effect **generalize** to 't Hooft-Polyakov monopoles?

't Hooft-Polyakov charge + mass

- In the presence of θ , monopoles acquire electric charge and **become a tower of dyons** (Witten effect):

$$q_E = \alpha_D q_M \left(n - \frac{\theta}{2\pi} \right)$$

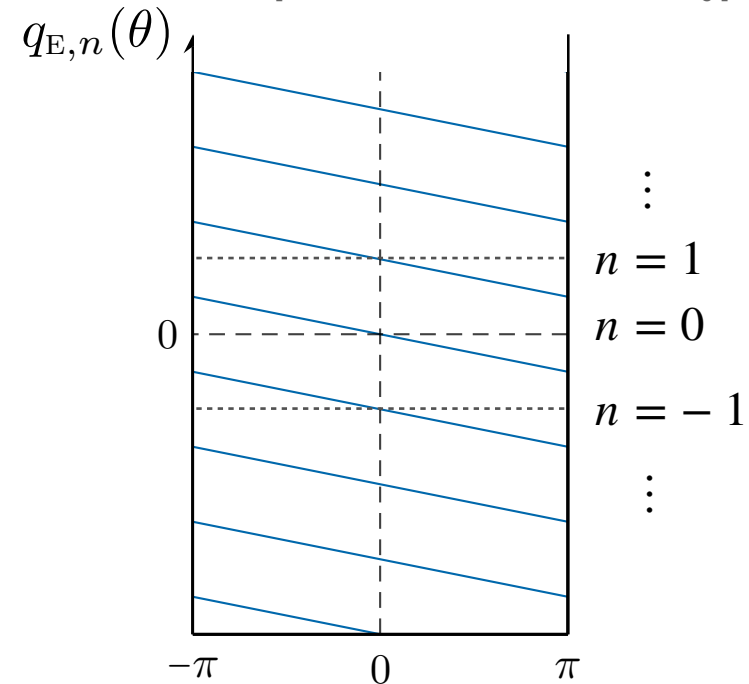
- The electric charge changes the dyon masses in a θ -dependent way:

$$m_D = v \sqrt{q_M^2 + q_E^2} \sim m_M \left(1 + \frac{\alpha_D^2}{2} \left(n - \frac{\theta}{2\pi} \right)^2 \right)$$

- This is equivalent to axion having a potential

$$V(\theta) = \Lambda^4 (1 - \cos \theta) + \rho_M \frac{\alpha_D^2}{2} \left(n - \frac{\theta}{2\pi} \right)^2$$

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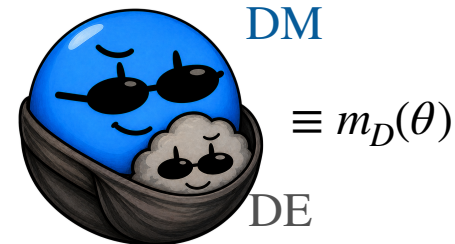
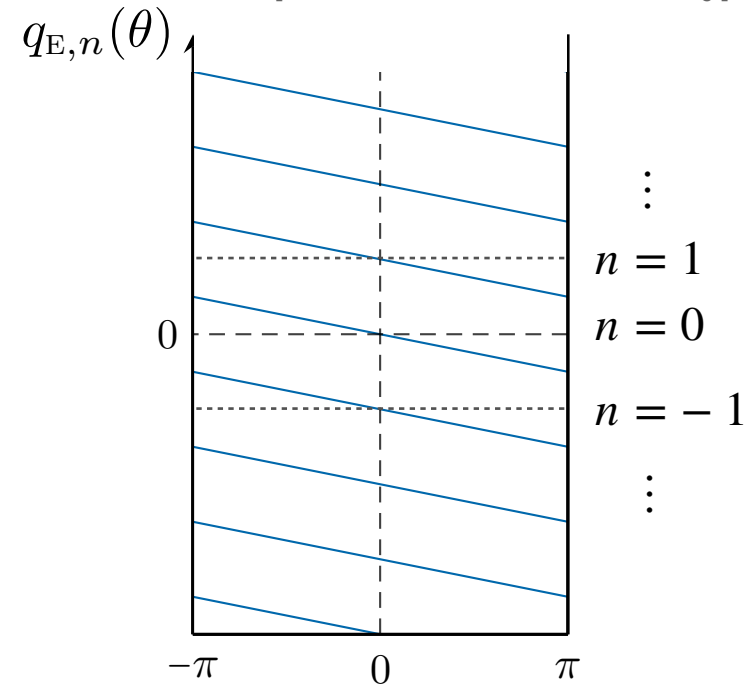
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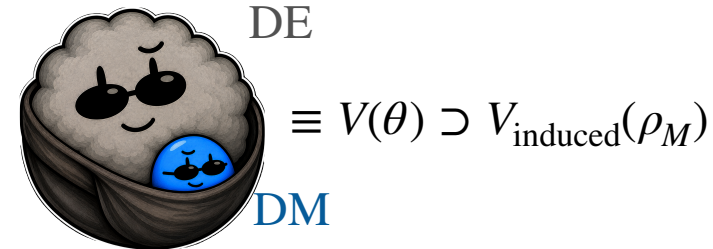
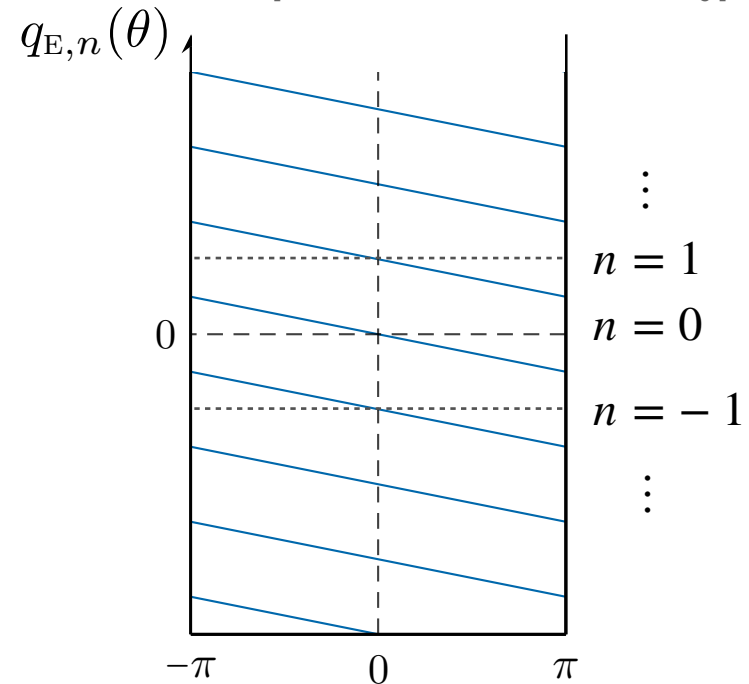
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Outline II

- ▶ The **magnetic monopoles** are a theoretically well-motivated **DM candidate**. It is a generic consequence of spontaneous symmetry breaking of gauge theories when the vacuum manifold has nontrivial topology.
- ▶ They **interact naturally** with **dynamic axion field (DE)**
 - What is a *natural* DM-DE interaction?
 - **Potential cosmological signatures**
 - Evolving DM-DE interaction (DESI)
 - Cosmic Birefringence (Ask me after the talk)

DM-DE Interactions

Dynamic DE models with scalar potential $V(\theta)$ &

- DM particle χ with θ -dependent mass,

$$m_\chi(\theta) \gg m_\theta$$

→ **Suffer from a naturalness problem:**

Loops of χ give a **Coleman-Weinberg**

correction to θ -potential that need to stay

small: $\Delta V'_{\text{CW}}(\theta) \ll V'(\theta)$

- *Assume “Interesting physics”:*

χ -density induced effective potential

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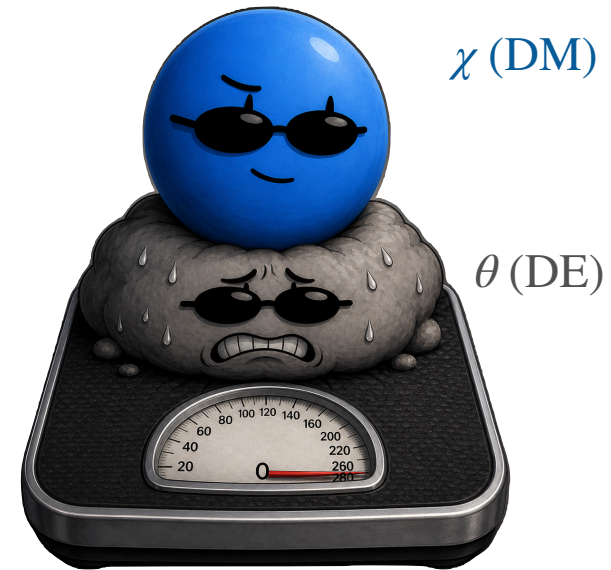
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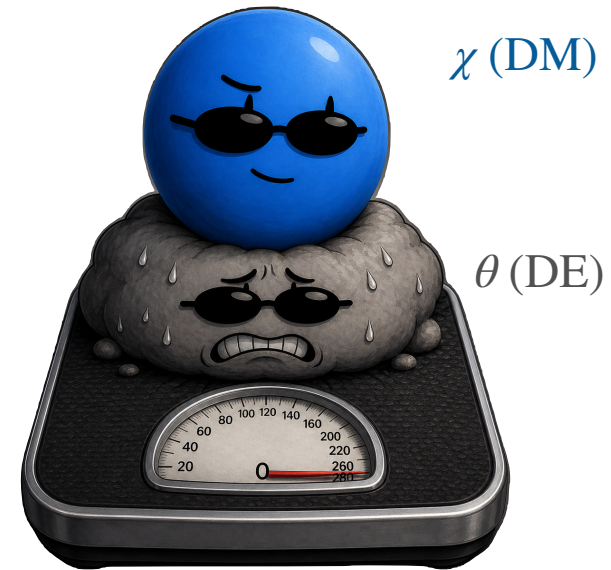
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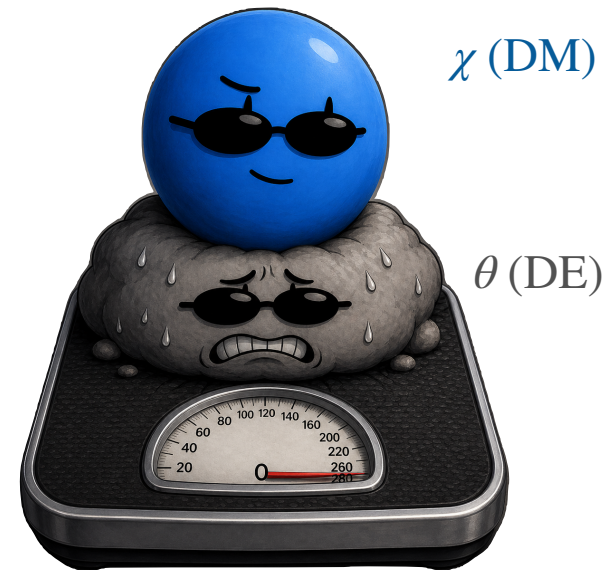
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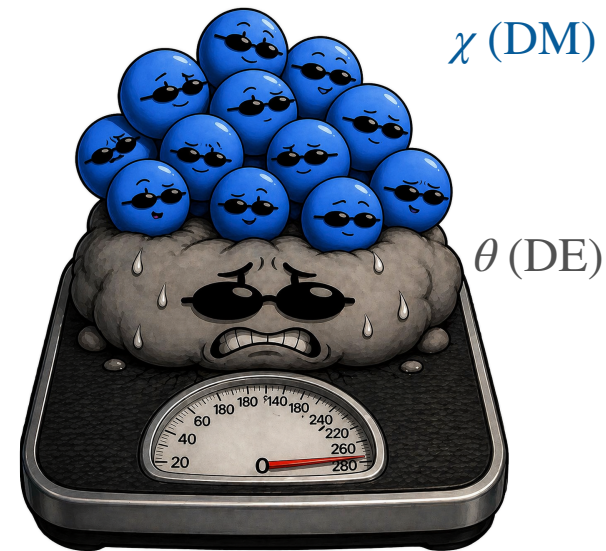
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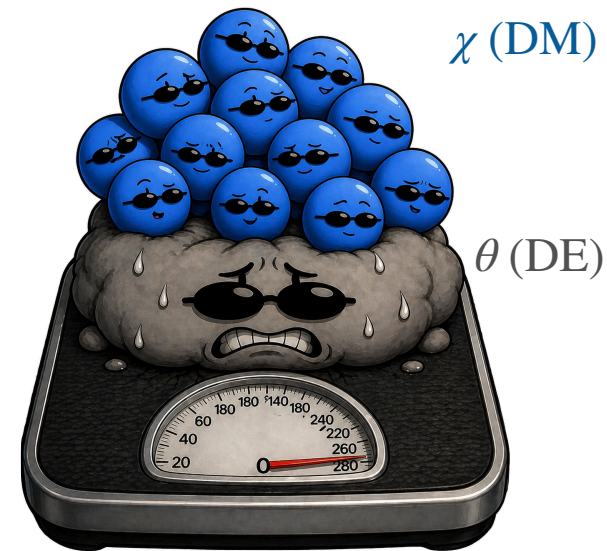
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What kinds of model can **solve this naturalness** problem?

Monodromic DM-DE Interactions

Dynamical DE models (at least one is true):

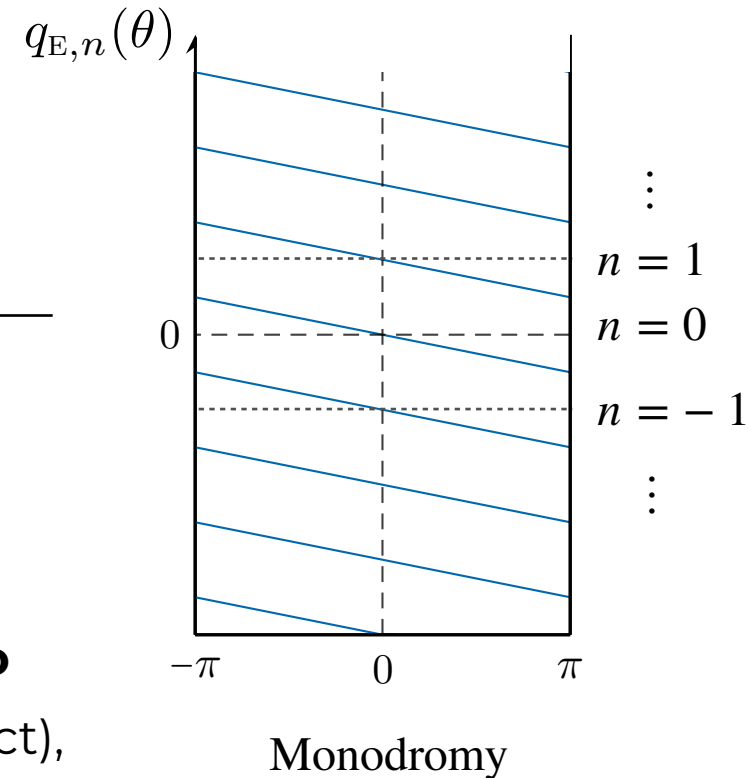
- Extremely fine tuned $V(\theta)$
- **Natural suppression** of CW-corrections
- $n_\chi \gg m_\chi^3$ and $m_\chi \ll \text{eV}$:

No fermions, χ is **wave dark-matter**

Few models with CW-suppression:

- Dark magnetic dyons (DM)
- Light axion (DE)
- An infinite tower of dyon states **shifted into each other when** $\theta \rightarrow \theta + 2\pi f$ (Witten effect), leads to a CW potential scaling $\exp(-2\pi/\alpha_D)$, [M. Reece, et al., **arXiv**: 2105.09950]

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Monodromic DM-DE Interactions

Dynamical DE models (at least one is true):

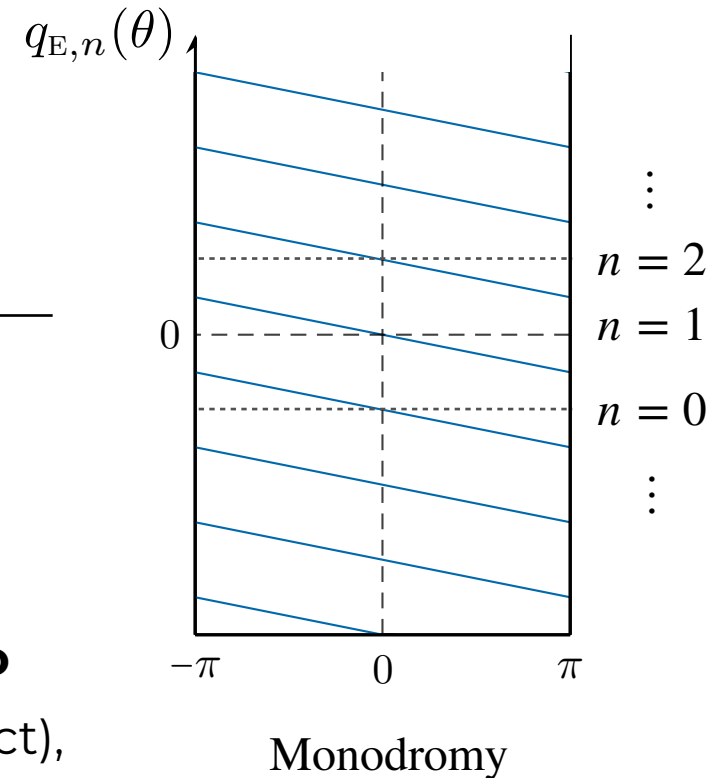
- Extremely fine tuned $V(\theta)$
- **Natural suppression** of CW-corrections
- $n_\chi \gg m_\chi^3$ and $m_\chi \ll \text{eV}$:

No fermions, χ is **wave dark-matter**

Few models with CW-suppression:

- Dark magnetic dyons (DM)
- Light axion (DE)
- An infinite tower of dyon states **shifted into each other when** $\theta \rightarrow \theta + 2\pi f$ (Witten effect), leads to a CW potential scaling $\exp(-2\pi/\alpha_D)$, [M. Reece, et al., **arXiv**: 2105.09950]

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Outline III

- ▶ The **magnetic monopoles** are a theoretically well-motivated **DM candidate**.
- ▶ They **interact naturally** with **dynamic axion field (DE)**
 - ▶ Summing over tower of dyon states **exponentially suppresses CW-corrections (Monodromy)**
- **Potential cosmological signatures**
 - Evolving DM-DE interaction (DESI)
 - Cosmic Birefringence (Ask me after the talk)

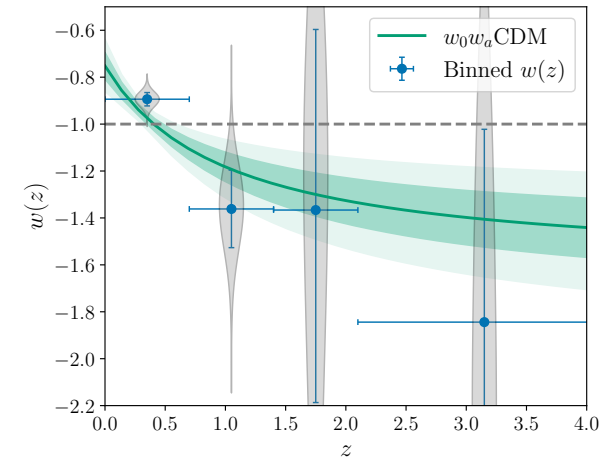
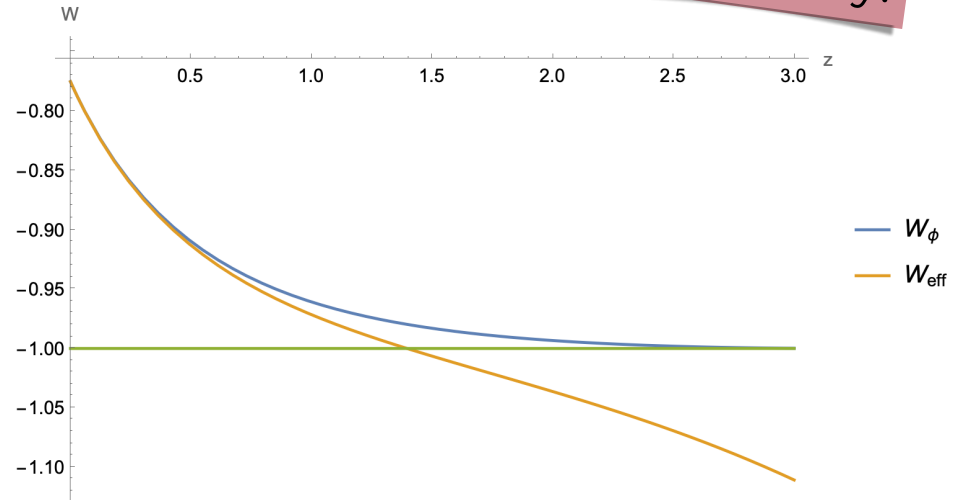
Effective equation of state

Preliminary!

A model with bench mark values:
 $m_a = 1.5 H_0, \alpha_D \sim \mathcal{O}(0.1), f_a \sim \mathcal{O}(M_{\text{Pl}})$

- A local observer who does not assume DM-DE interaction will attribute the change on the DM density to an **effective component of DE**

$$W_{\text{eff}} = \frac{W_\phi}{1 + \left(\frac{A(\phi)}{A(\phi_0)} - 1\right) \frac{\rho_{DM}(0)}{a^3 \rho_\phi}}$$



© [DESI DR11, [arXiv: 2503.14738](https://arxiv.org/abs/2503.14738)]

Summary

- ▶ The **magnetic monopoles** are a theoretically well-motivated **DM candidate**.
- ▶ They interact naturally with **dynamic axion field (DE)**
 - ▶ Summing over infinite dyon states, **exponentially suppresses CW-corrections (Monodromy)**
- ▶ The model has potential **cosmological signatures** with **natural benchmark parameters**, including DESI and Birefringence.

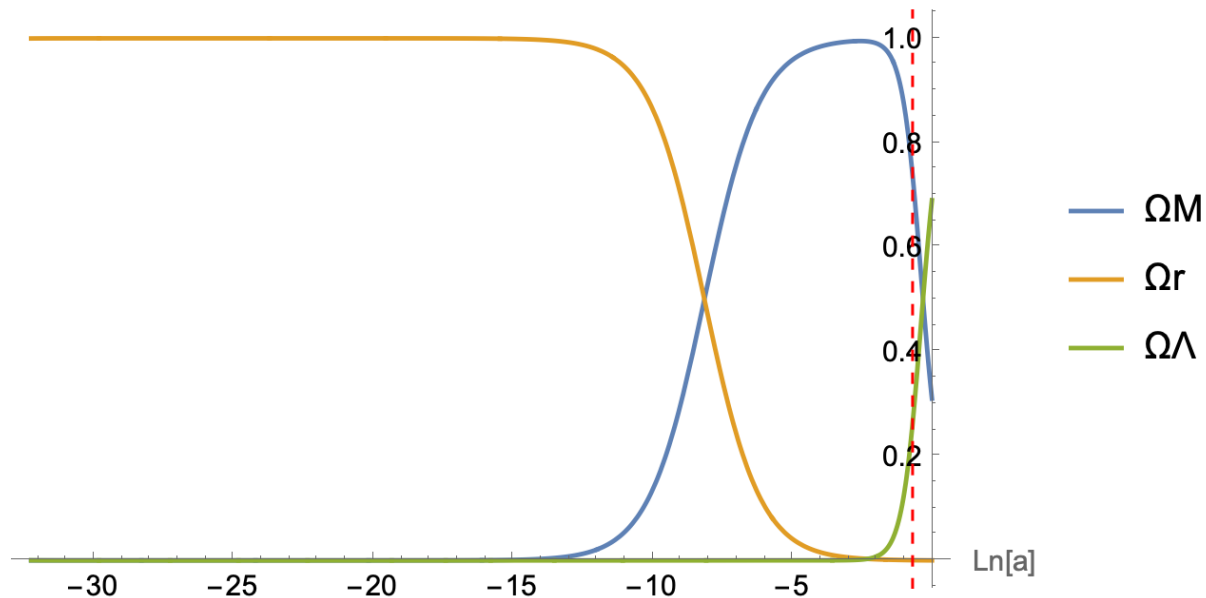
Thank you!

EOM and cosmic evolution

A model with bench mark values:

$$m_a = 1.5 H_0, \alpha_D \sim \mathcal{O}(0.1), f_a \sim \mathcal{O}(M_{\text{Pl}})$$

Dark Energy (Axion)	Dark Matter (Monopole)
$\rho_\theta = \Lambda^4(1 - \cos \theta) + \dot{\theta}^2 f_a^2 / 2$	$\rho_M = \frac{\rho_M^0}{a^3} \left(1 + \frac{\alpha_D^2}{8\pi^2} (\theta - 2\pi n)^2\right)$

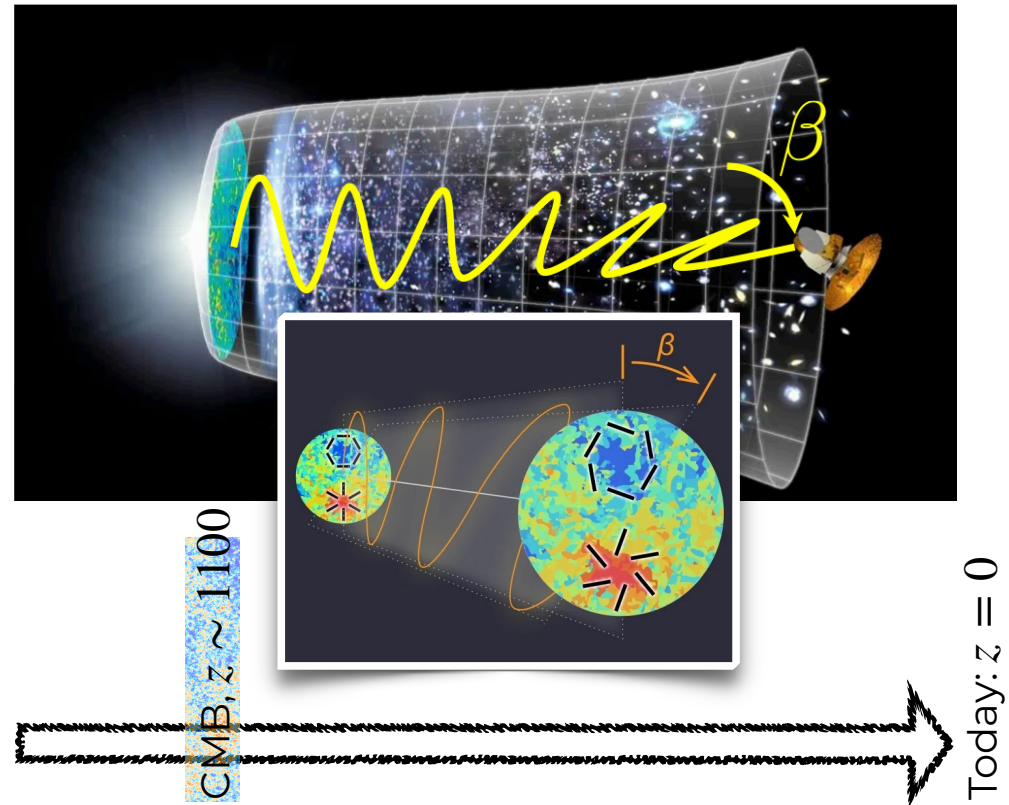


On-going: Axion DE & Birefringence

- Birefringence signal: Rotation of the photon polarization vector
- Turning on the coupling to photons: $g_{\theta\gamma} \theta F\tilde{F}$
- The **Birefringence angle** is computed via: $2\beta = g_{\theta\gamma} \Delta\theta$ (remember

$$c_{L/R} \equiv \frac{\omega_{L/R}}{k} = \sqrt{1 \pm \frac{g_{\theta\gamma} \dot{\theta}}{k}}$$

- To get $\beta_{\text{Planck}} \sim 0.35$ deg, we need $\alpha_{\text{vis}} \sim 0.1$



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Axions rotate polarization:

$$\ddot{A}_k^{L/R} + \omega_{L/R}^2 A_k^{L/R} = 0$$

$$c_{L/R} \equiv \frac{\omega_{L/R}}{k} = \sqrt{1 \pm \frac{g_{\theta\gamma} \dot{\theta}}{k}}$$