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Heavy Vector Triplets at the Muon Collider

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Phenomenology 2026 Symposium



Outline

- A bridge theory
- How to discover HVT at a Muon Collider
- Interference
- Comparison with other colliders

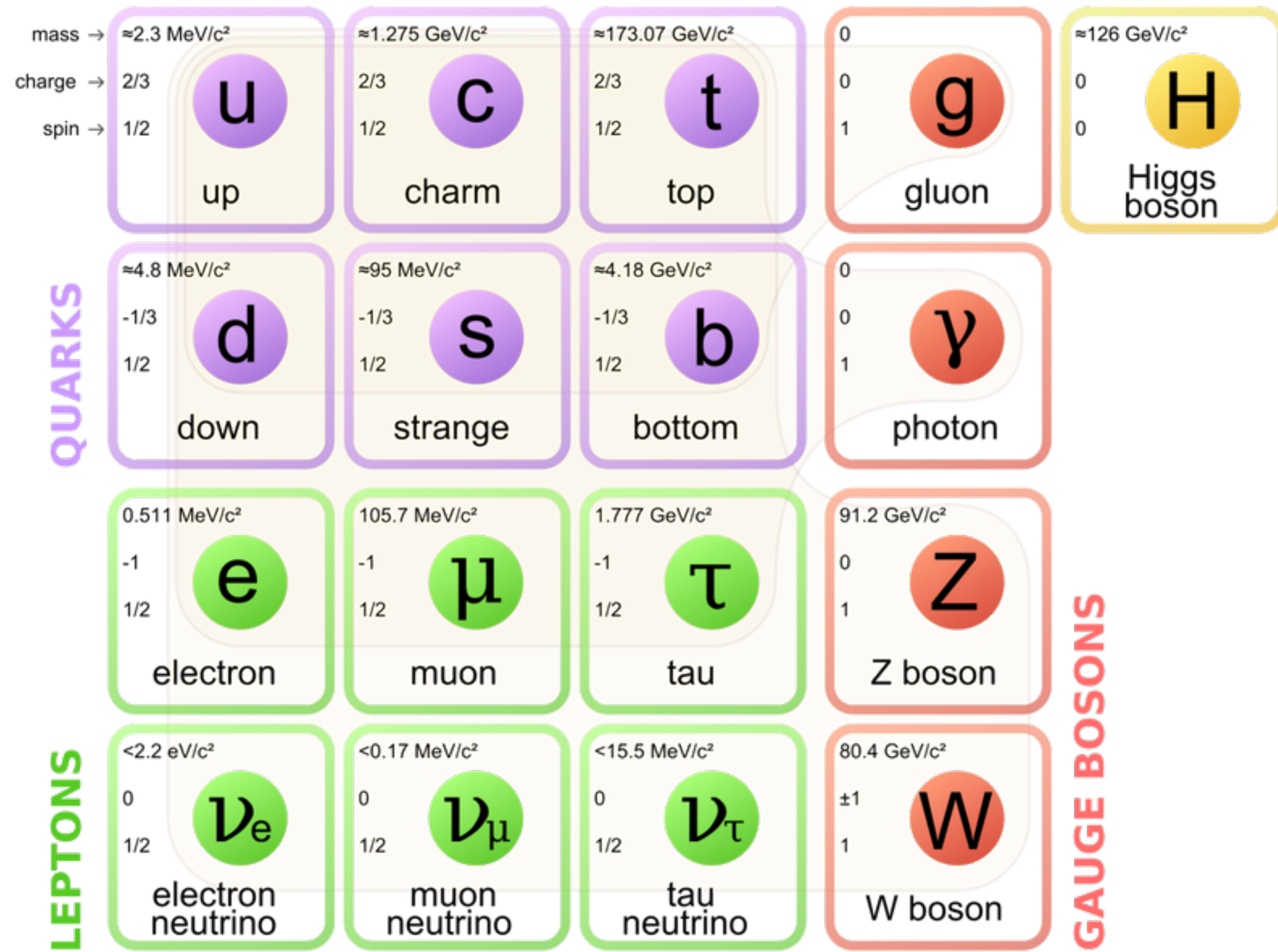
A bridge theory

mass →	$\approx 2.3 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 173.07 \text{ GeV}/c^2$	0	$\approx 126 \text{ GeV}/c^2$
charge →	$2/3$	$2/3$	$2/3$	0	0
spin →	$1/2$	$1/2$	$1/2$	1	0
	u up	c charm	t top	g gluon	H Higgs boson
QUARKS	$\approx 4.8 \text{ MeV}/c^2$	$\approx 95 \text{ MeV}/c^2$	$\approx 4.18 \text{ GeV}/c^2$	0	
	$-1/3$	$-1/3$	$-1/3$	0	
	$1/2$	$1/2$	$1/2$	1	
	d down	s strange	b bottom	γ photon	
	$0.511 \text{ MeV}/c^2$	$105.7 \text{ MeV}/c^2$	$1.777 \text{ GeV}/c^2$	$91.2 \text{ GeV}/c^2$	
	-1	-1	-1	0	
	$1/2$	$1/2$	$1/2$	1	
	e electron	μ muon	τ tau	Z Z boson	
LEPTONS	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 15.5 \text{ MeV}/c^2$	$80.4 \text{ GeV}/c^2$	
	0	0	0	± 1	
	$1/2$	$1/2$	$1/2$	1	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
				GAUGE BOSONS	

+

NP

- Extended gauge symmetries
- Composite Higgs models
- Extra dimensions
- ...



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	$0.511 \text{ MeV}/c^2$	$105.7 \text{ MeV}/c^2$	$1.777 \text{ GeV}/c^2$		
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	$1/2$	$1/2$	$1/2$		
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino		



Francesca Acanfora

HVT @ MuC

UMass Amherst

$$\begin{aligned}
\mathcal{L}_V \supset & -\frac{1}{4} D_{[\mu} \mathcal{V}_{\nu]}^a D^{[\mu} \mathcal{V}^{\nu]}_a + \frac{m_V^2}{2} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} \\
& + i g_V c_H \mathcal{V}_\mu^a H^\dagger \tau^a \overleftrightarrow{D}^\mu H + \frac{g^2}{g_V} c_q \mathcal{V}_\mu^a \sum_q \bar{q}_L \gamma^\mu \tau^a q_L + \frac{g^2}{g_V} c_\ell \mathcal{V}_\mu^a \sum_\ell \bar{\ell}_L \gamma^\mu \tau^a \ell_L \\
& + \frac{g_V}{2} c_{VVV} \epsilon_{abc} \mathcal{V}_\mu^a \mathcal{V}_\nu^b D^{[\mu} \mathcal{V}^{\nu]}_c + g_V^2 c_{VVHH} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} H^\dagger H - \frac{g}{2} c_{V VW} \epsilon_{abc} W^{\mu\nu a} \mathcal{V}_\mu^b \mathcal{V}_\nu^c
\end{aligned}$$

kinetic

$$\begin{aligned} \mathcal{L}_V \supset & \boxed{-\frac{1}{4} D_{[\mu} \mathcal{V}_{\nu]}^a D^{[\mu} \mathcal{V}^{\nu]}_a} + \frac{m_V^2}{2} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} \\ & + i g_V c_H \mathcal{V}_\mu^a H^\dagger \tau^a \overleftrightarrow{D}^\mu H + \frac{g^2}{g_V} c_q \mathcal{V}_\mu^a \sum_q \bar{q}_L \gamma^\mu \tau^a q_L + \frac{g^2}{g_V} c_\ell \mathcal{V}_\mu^a \sum_\ell \bar{\ell}_L \gamma^\mu \tau^a \ell_L \\ & + \frac{g_V}{2} c_{VVV} \epsilon_{abc} \mathcal{V}_\mu^a \mathcal{V}_\nu^b D^{[\mu} \mathcal{V}^{\nu]}_c + g_V^2 c_{VVHH} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} H^\dagger H - \frac{g}{2} c_{V VW} \epsilon_{abc} W^{\mu\nu a} \mathcal{V}_\mu^b \mathcal{V}_\nu^c \end{aligned}$$

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\end{aligned}$$

Higgs
current

$$\begin{aligned}
\mathcal{L}_V \supset & -\frac{1}{4} D_{[\mu} \mathcal{V}_{\nu]}^a D^{[\mu} \mathcal{V}^{\nu]}{}_a + \frac{m_V^2}{2} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} \\
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\end{aligned}$$

quarks

$$\begin{aligned}
\mathcal{L}_V \supset & -\frac{1}{4} D_{[\mu} \mathcal{V}_{\nu]}^a D^{[\mu} \mathcal{V}^{\nu]}{}_a + \frac{m_V^2}{2} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} \\
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leptons

$$\begin{aligned}
\mathcal{L}_V \supset & -\frac{1}{4} D_{[\mu} \mathcal{V}_{\nu]}^a D^{[\mu} \mathcal{V}^{\nu]}_a + \frac{m_V^2}{2} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} \\
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\end{aligned}$$

self-interaction

$$\begin{aligned}
\mathcal{L}_V \supset & -\frac{1}{4} D_{[\mu} \mathcal{V}_{\nu]}^a D^{[\mu} \mathcal{V}^{\nu]}_a + \frac{m_V^2}{2} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} \\
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& + \frac{g_V}{2} c_{VVV} \epsilon_{abc} \mathcal{V}_\mu^a \mathcal{V}_\nu^b D^{[\mu} \mathcal{V}^{\nu]}_c + \underbrace{g_V^2 c_{VVHH} \mathcal{V}_\mu^a \mathcal{V}^{\mu a} H^\dagger H}_{VVhh} - \frac{g}{2} c_{VW} \epsilon_{abc} W^{\mu\nu a} \mathcal{V}_\mu^b \mathcal{V}_\nu^c
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\end{aligned}$$

SM SU(2) current

- Dimension 4 operators
- Compatible with SM symmetry
- Assume CP
- Ignored 4V vertices because pheno irrelevant
- A bridge theory: only list pheno relevant operators
- Not an EFT: HVT is not integrated away

$$\begin{aligned}
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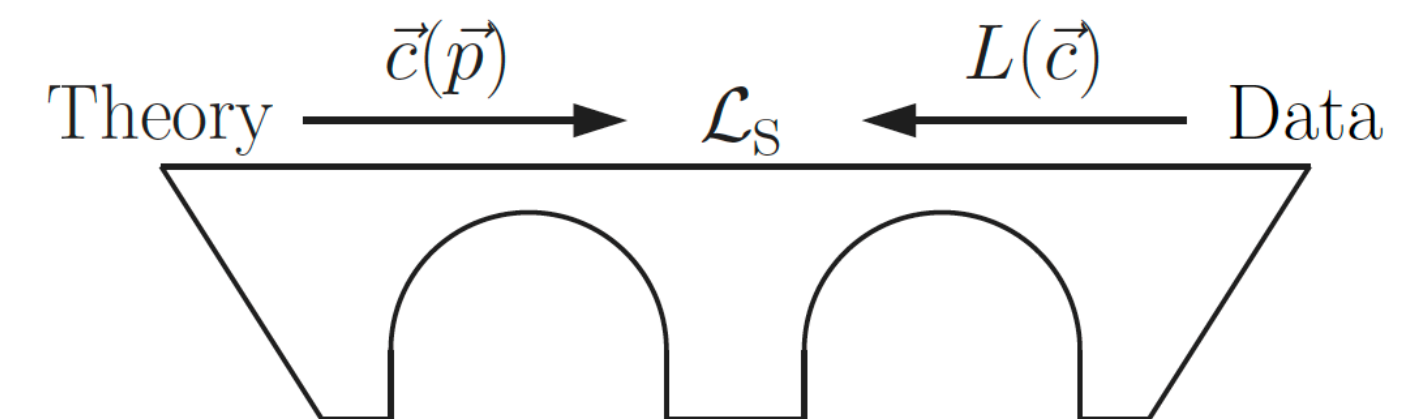


Figure 1.1: Pictorial view of the Bridge Method.

1402.4431

Bridge theory

Bridge theory



"Matching"

Express c_i in terms of g_v

Bridge theory



"Matching"

Express c_i in terms of g_ν



Specific BSM theory

Bridge theory

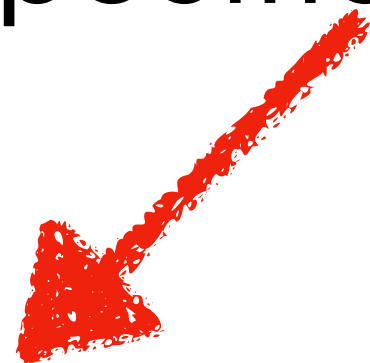


"Matching"

Express c_i in terms of g_v



Specific BSM theory



Model A

Bridge theory

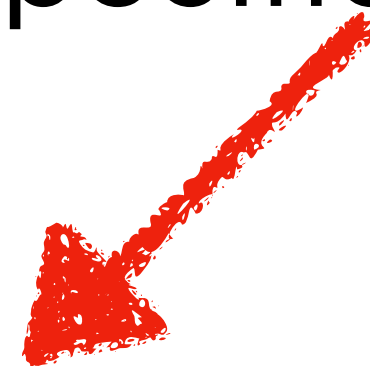


"Matching"

Express c_i in terms of g_v

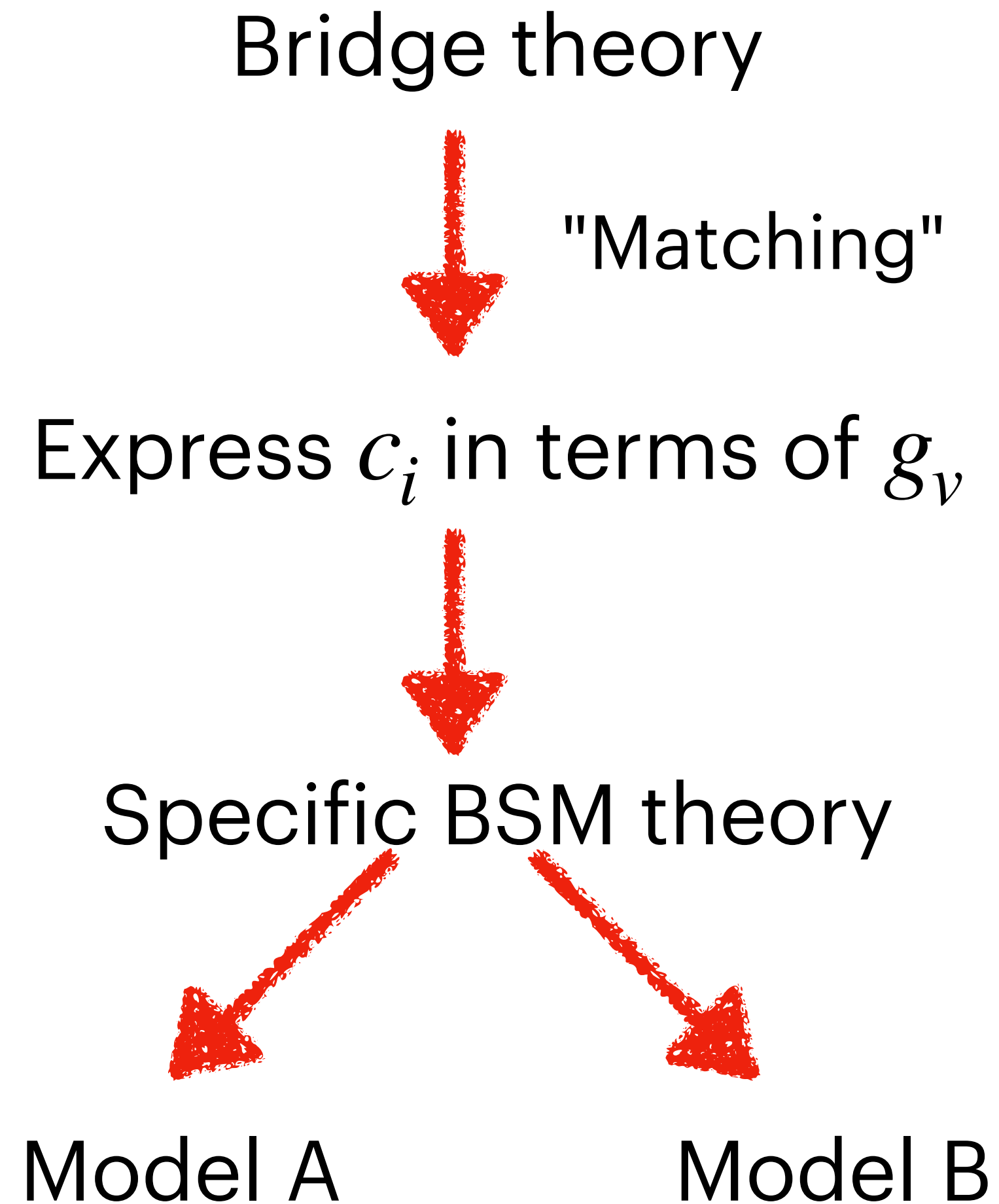


Specific BSM theory



Model A

Model B



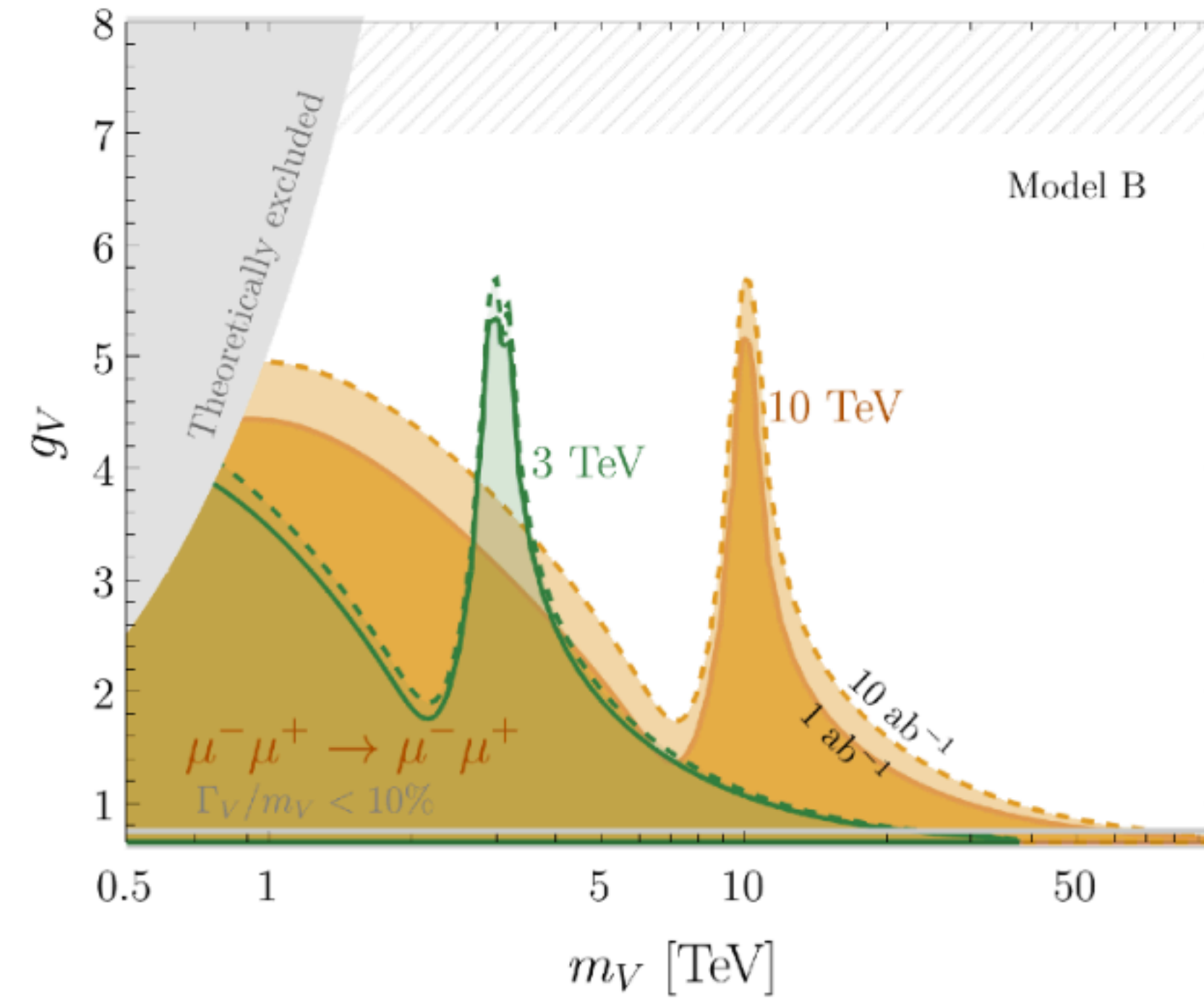
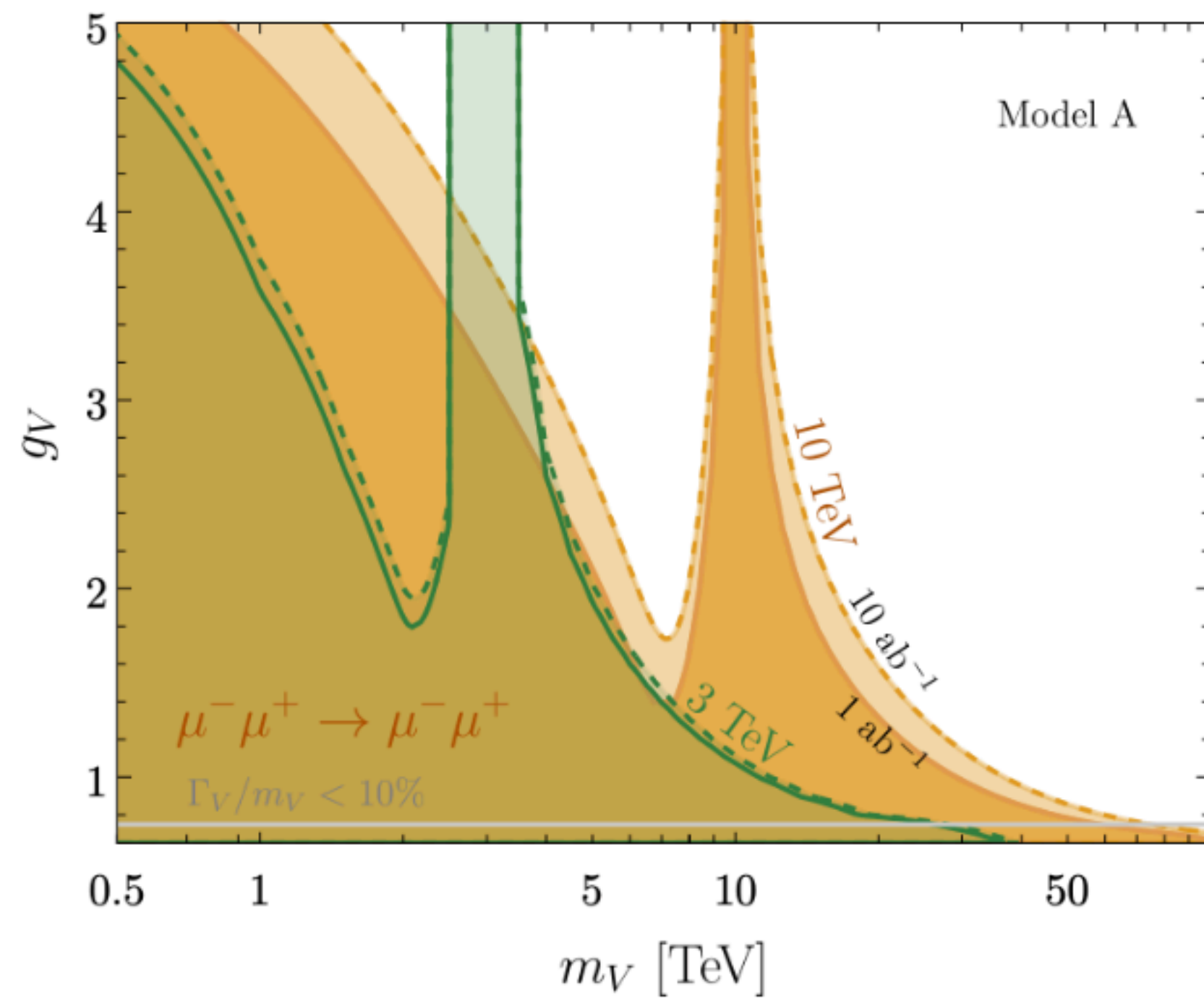
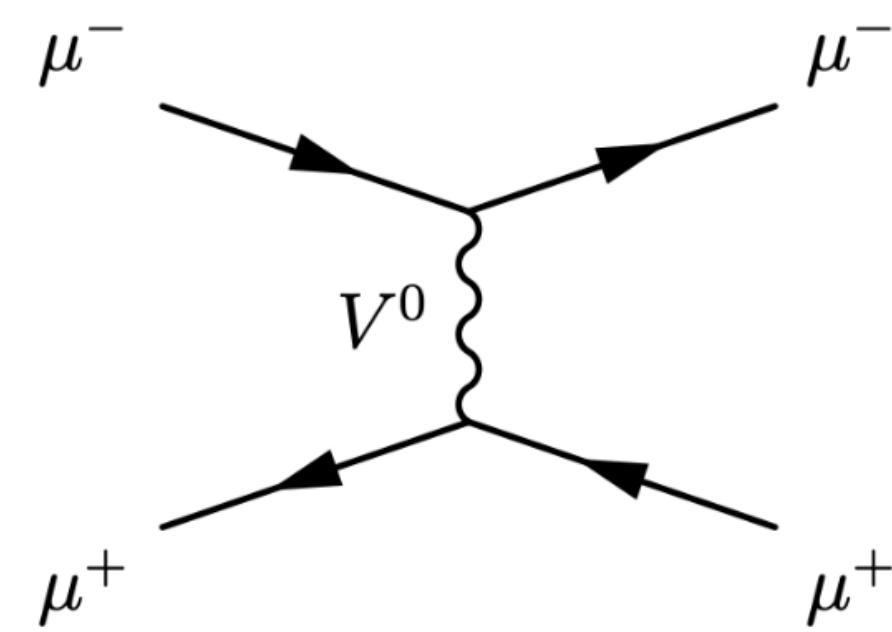
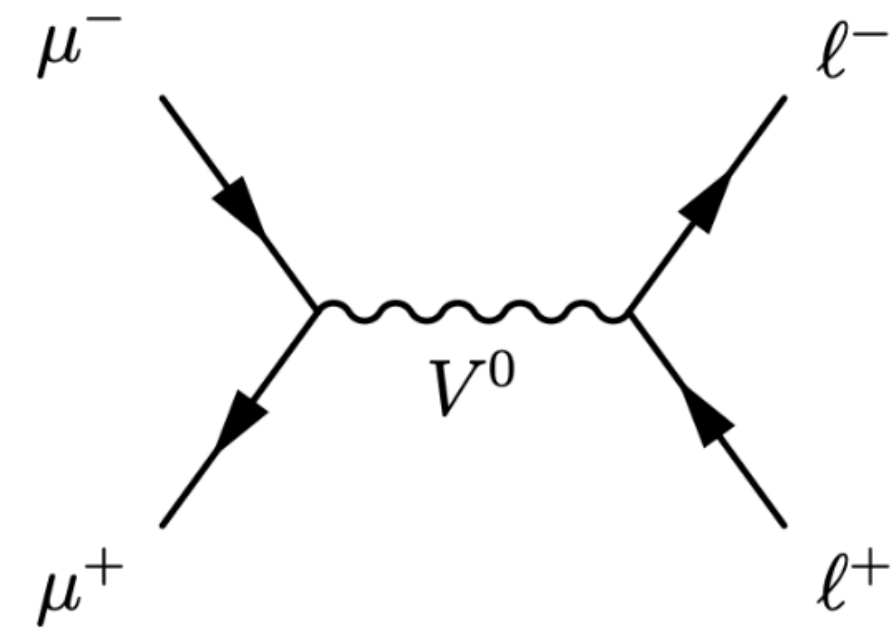
	Model A $su(2)_1 \times su(2)_2$	Model B $SO(5)/SO(4)$
g_V	g_2	$\pm g_\rho$
m_V	$\frac{ g_V u}{k_V}$	$\frac{m_\rho}{k_V}$
c_F	$-\frac{1}{k_V}$	$\frac{1}{k_V}$
c_H	$-\frac{g^2}{g_V^2 k_V}$	$-\frac{1}{k_V} \left(a_\rho^2 - \frac{g^2}{g_V^2} \right)$
c_{VVV}	$-\frac{1}{k_V} \left(1 - \frac{2g^2}{g_V^2} \right)$	$\frac{1}{k_V} \left(1 - \frac{2g^2}{g_V^2} \right)$
c_{VW}	1	1
c_{VVHH}	$\frac{g^4}{4g_V^4 k_V^2}$	$-\frac{g^2}{2g_V^2 k_V^2} \left(a_\rho^2 - \frac{g^2}{2g_V^2} \right)$

$$a_\rho = 1/\sqrt{2}, k_V = \sqrt{1 - g_w^2/g_V^2}$$

Pheno of HVT at a Muon Collider

The setup

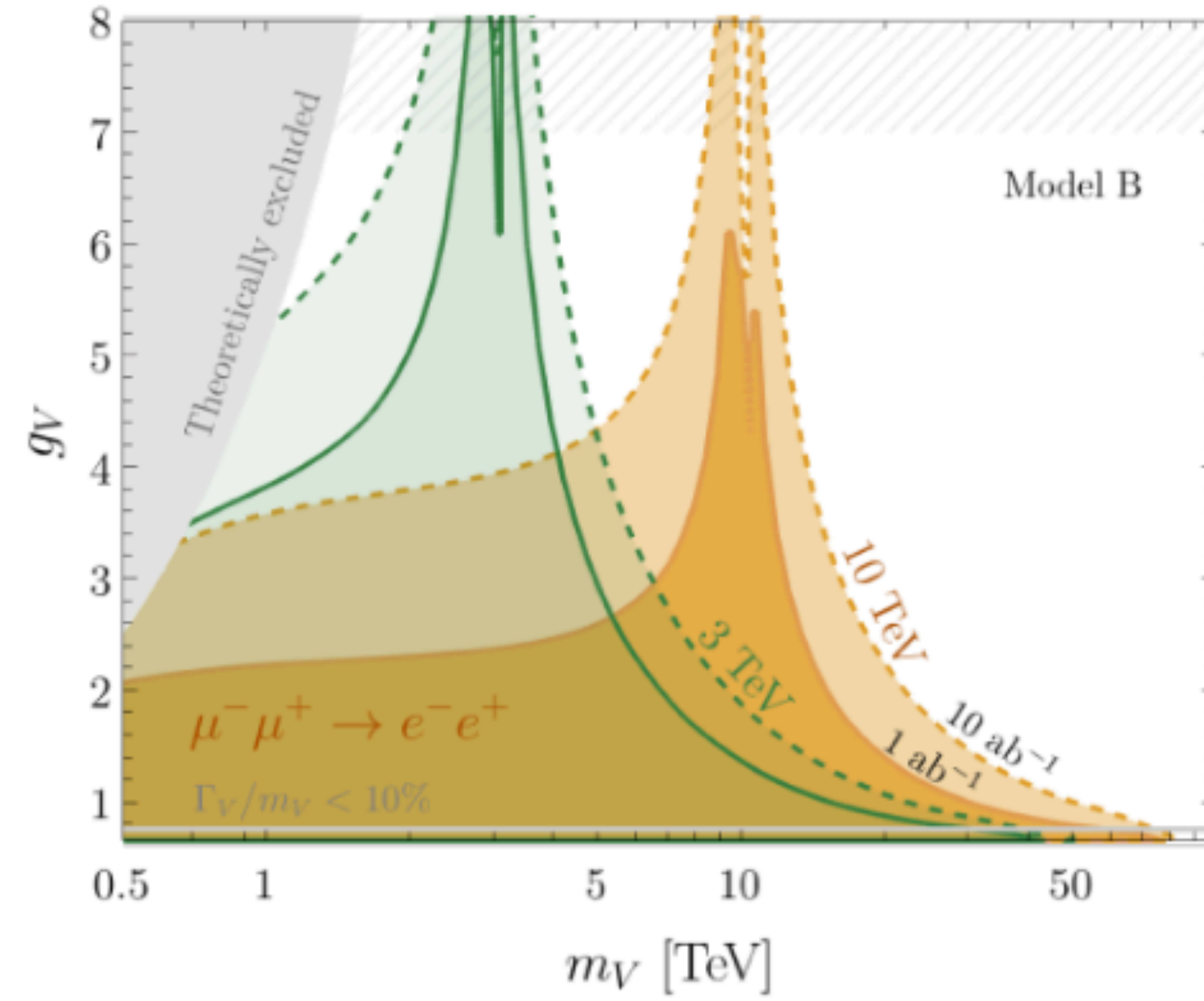
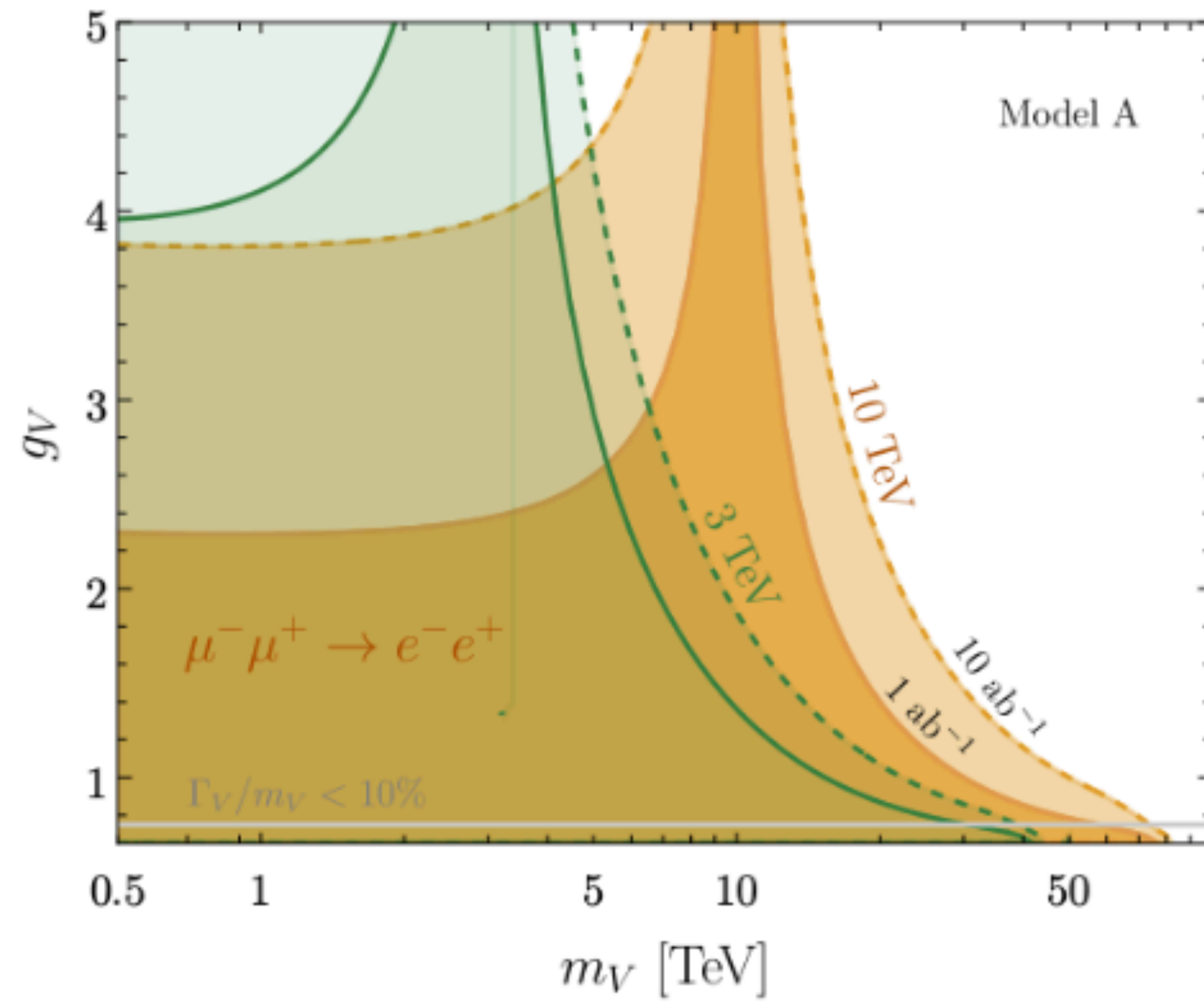
- Can not use **Narrow Width Approximation**:
 - Can have $\frac{\Gamma}{m_V} > 10\%$
 - Can have $m_V \sim \sqrt{s}$
 - Can have non resonant important channels
- Must simulate signal + background + **interference**
 - **MC errors** become important
- 3 or 10 TeV MuC, $E_X > 20$ GeV, $|\eta_X| < 2.5$, $\Delta R_{i,j} > 0.4$

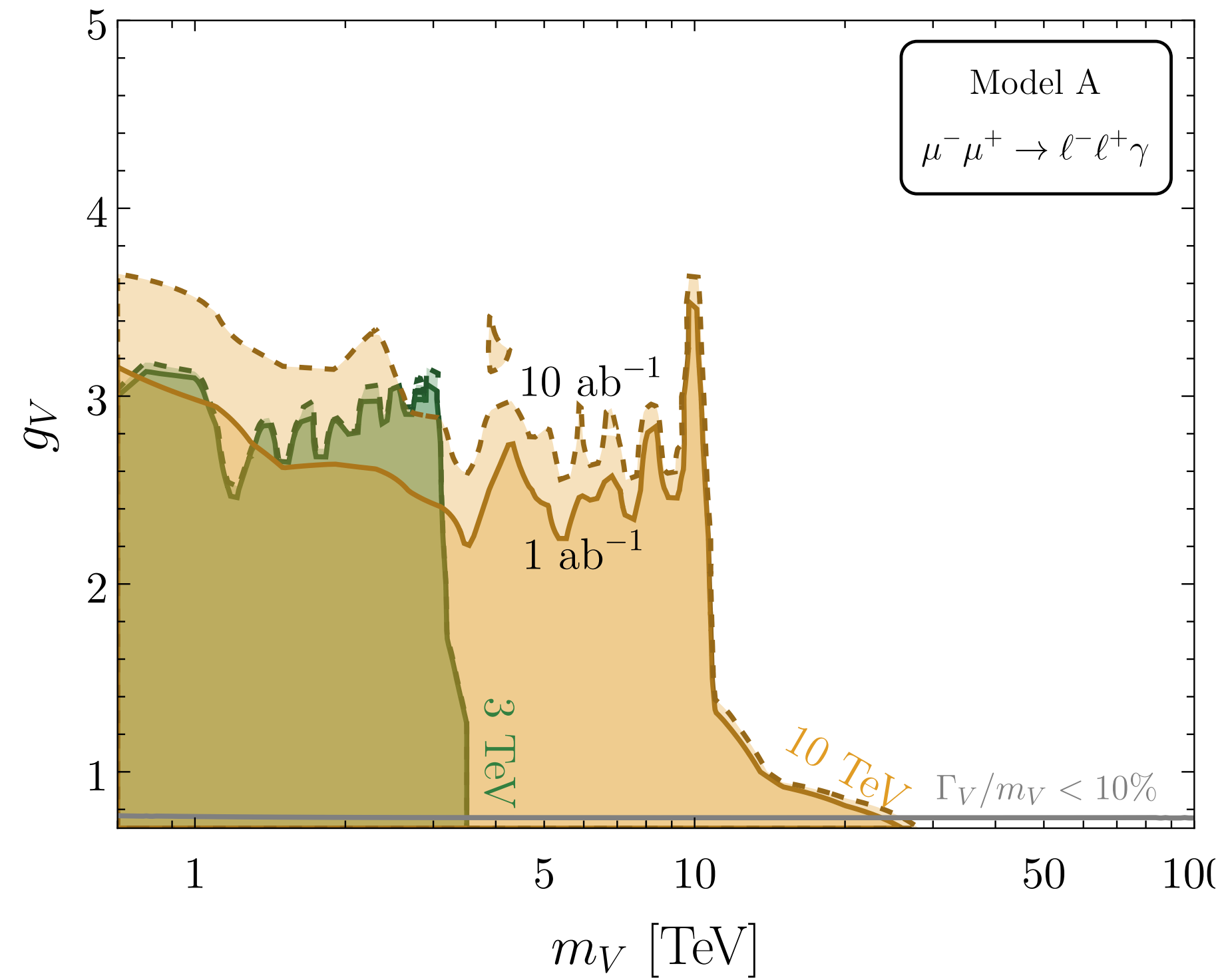
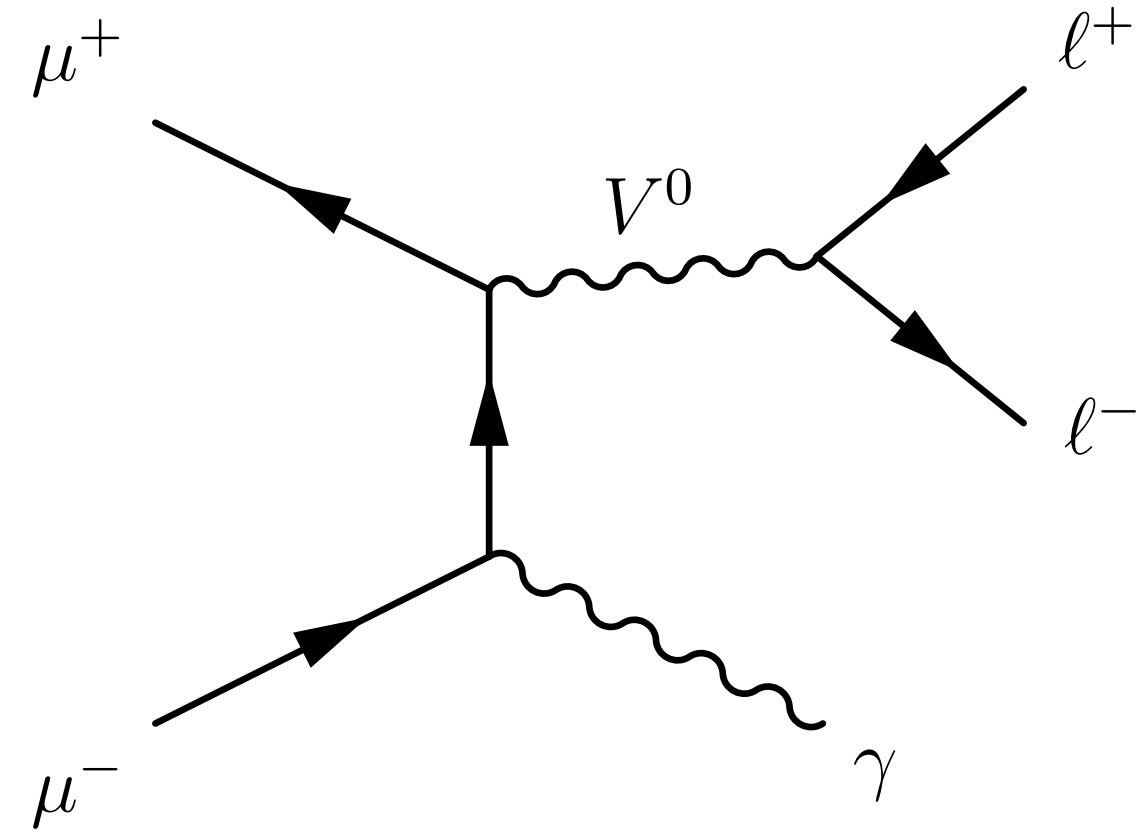


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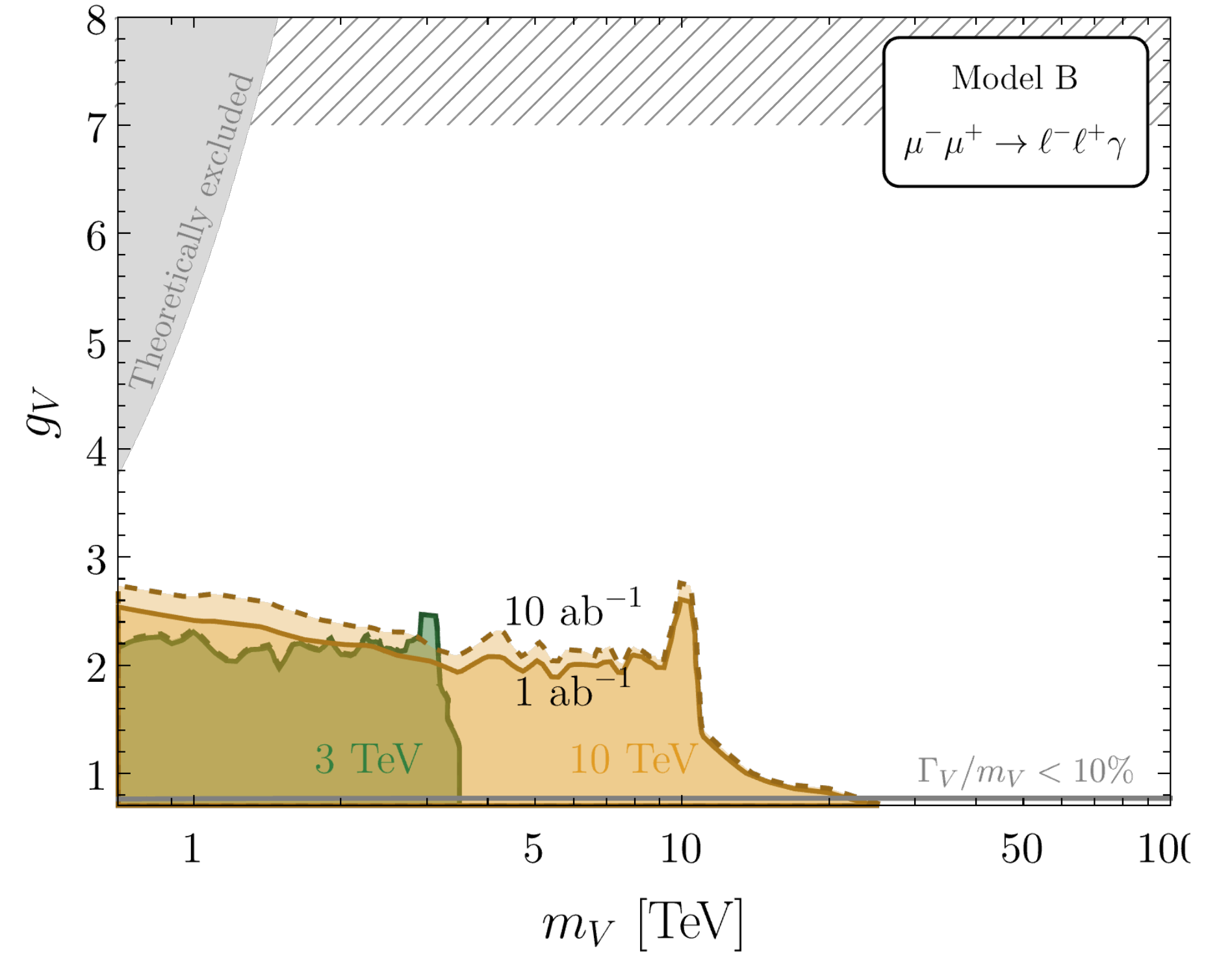
HVT @ MuC

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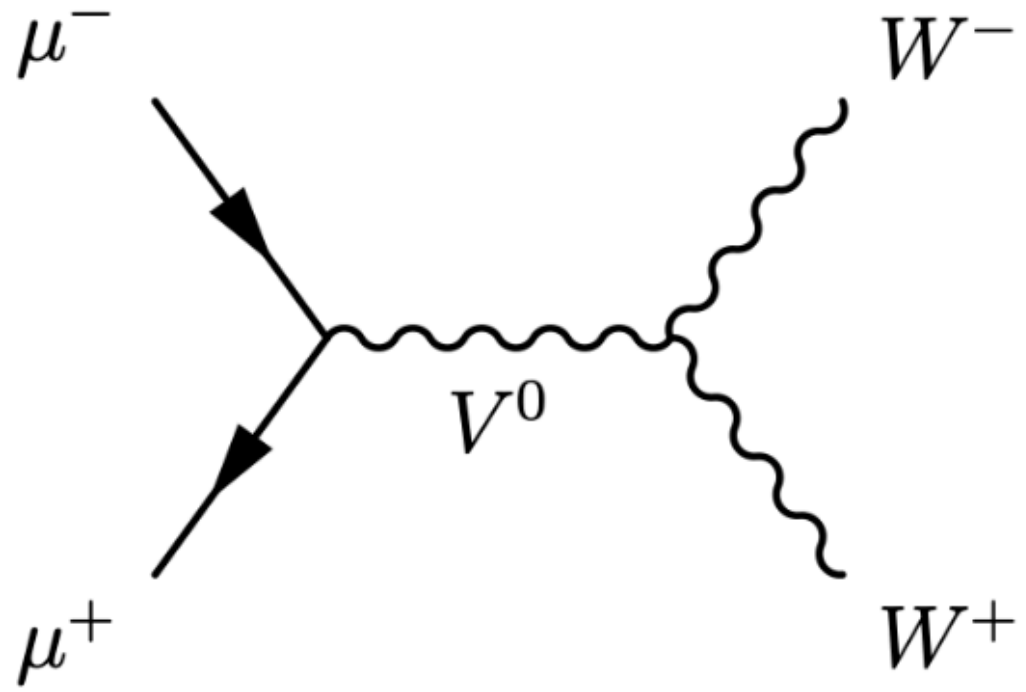


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HVT @ MuC

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Unitarity

$$|\mathcal{M}|_V^2 = \frac{e^6 c_W c_\ell^2}{s_W^6 g_V^2 m_W^4 (s - m_V^2)^2} \theta_C \theta_N \quad (3.10)$$

$$|\mathcal{M}|_\gamma^2 = \frac{e^4}{4s_W^4} \frac{t(s+t)}{m_W^4} [-8s_W^4] \quad (3.11)$$

$$|\mathcal{M}|_Z^2 = \frac{e^4}{4s_W^4} \frac{t(s+t)}{m_W^4} \left[-8s_W^4 + 4s_W^2 - 1 - \frac{2ec_\ell}{g_V s_W c_W} \theta_N - \frac{4es_W^3 c_\ell}{g_V c_W} \theta_N + \frac{6es_W c_\ell}{g_V c_W} \theta_N \right] \quad (3.12)$$

$$|\mathcal{M}|_\nu^2 = \frac{e^4}{4s_W^4} \frac{t(s+t)}{m_W^4} \left[-1 - \frac{4ec_\ell}{g_V s_W} \theta_C \right] \quad (3.13)$$

$$|\mathcal{M}|_{\gamma Z}^2 = \frac{e^4}{4s_W^4} \frac{t(s+t)}{m_W^4} \left[16s_W^4 - 4s_W^2 + \frac{4es_W^3 c_\ell}{g_V c_W} \theta_N - \frac{4es_W c_\ell}{g_V c_W} \theta_N \right] \quad (3.14)$$

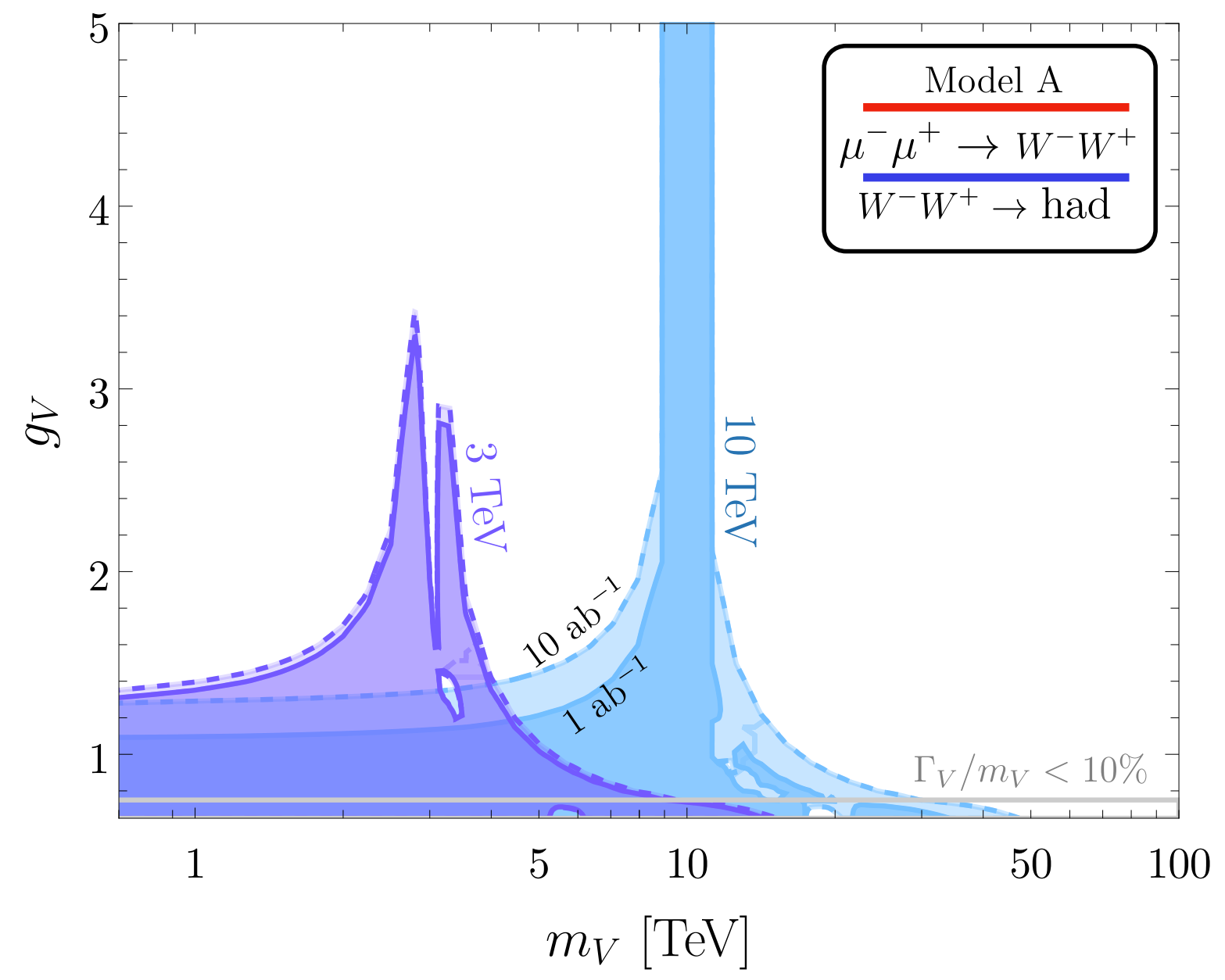
$$|\mathcal{M}|_{\gamma\nu}^2 = \frac{e^4}{4s_W^4} \frac{t(s+t)}{m_W^4} \left[4s_W^2 + \frac{8ec_\ell s_W}{g_V} \theta_C \right] \quad (3.15)$$

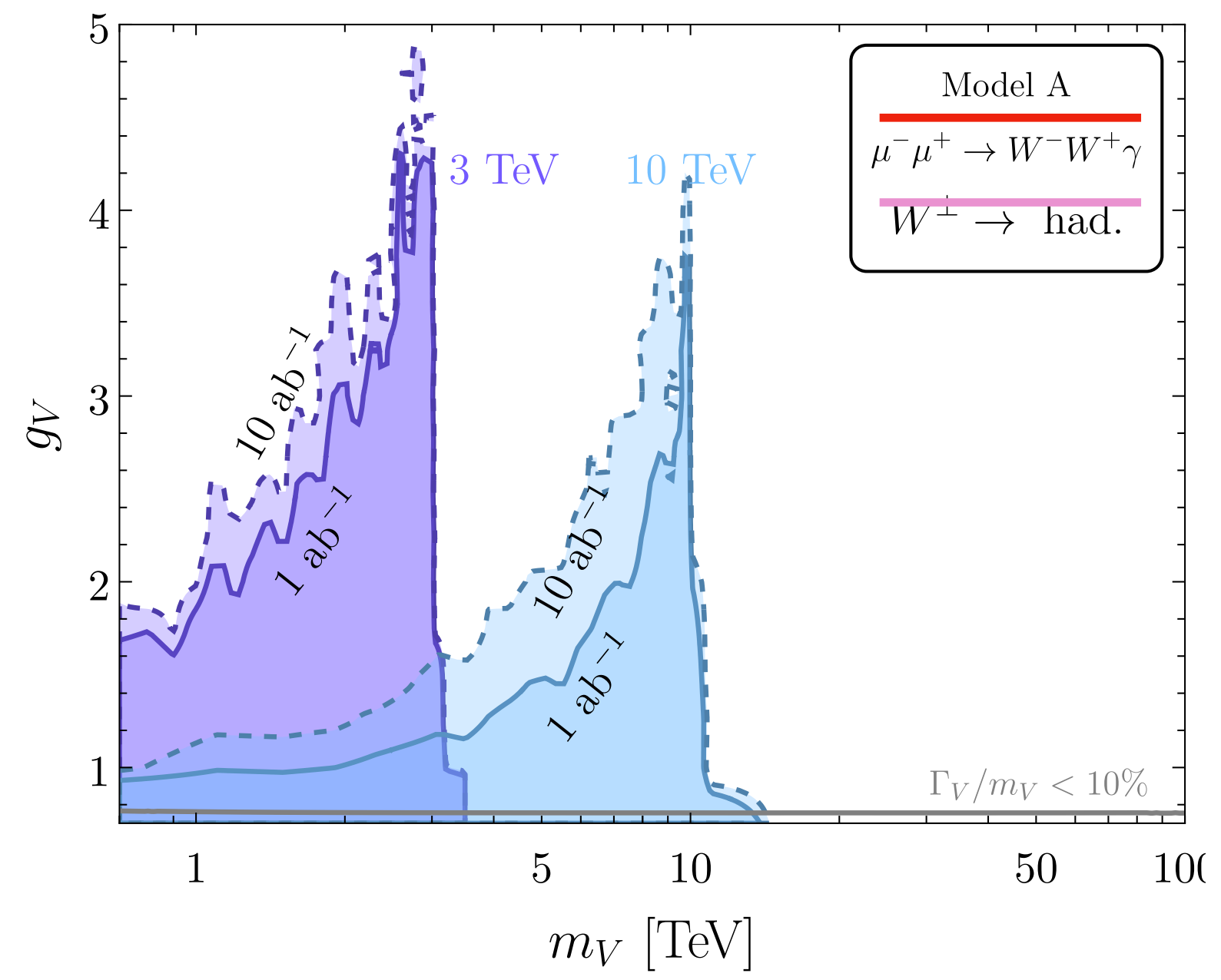
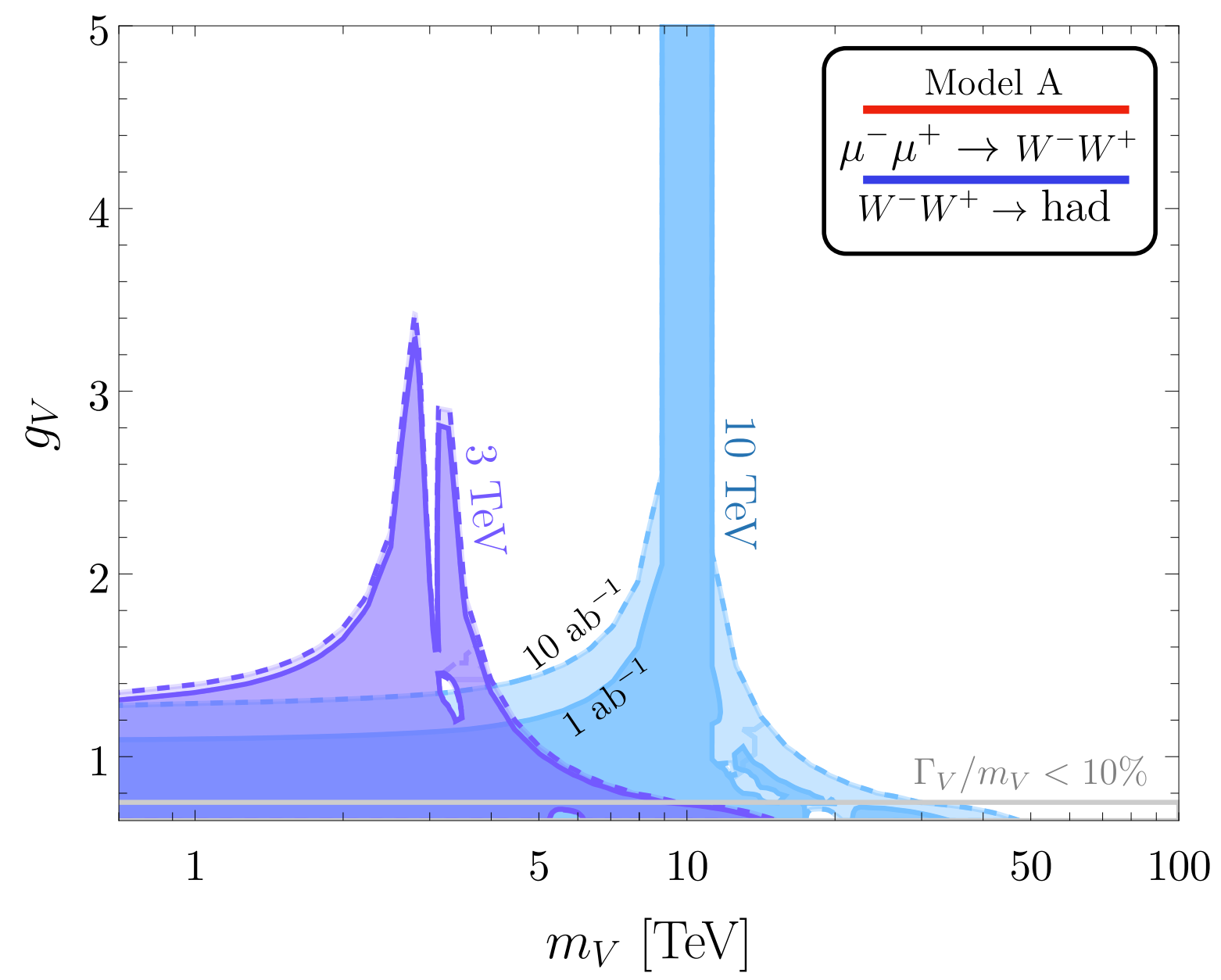
$$|\mathcal{M}|_{Z\nu}^2 = \frac{e^4}{4s_W^4} \frac{t(s+t)}{m_W^4} \left[2 - 4s_W^2 - \frac{2es_W c_\ell}{g_V c_W} \theta_N + \frac{2ec_\ell}{g_V c_W s_W} \theta_N - \frac{8ec_\ell s_W}{g_V} \theta_C + \frac{4ec_\ell}{g_V s_W} \theta_C \right] \quad (3.16)$$

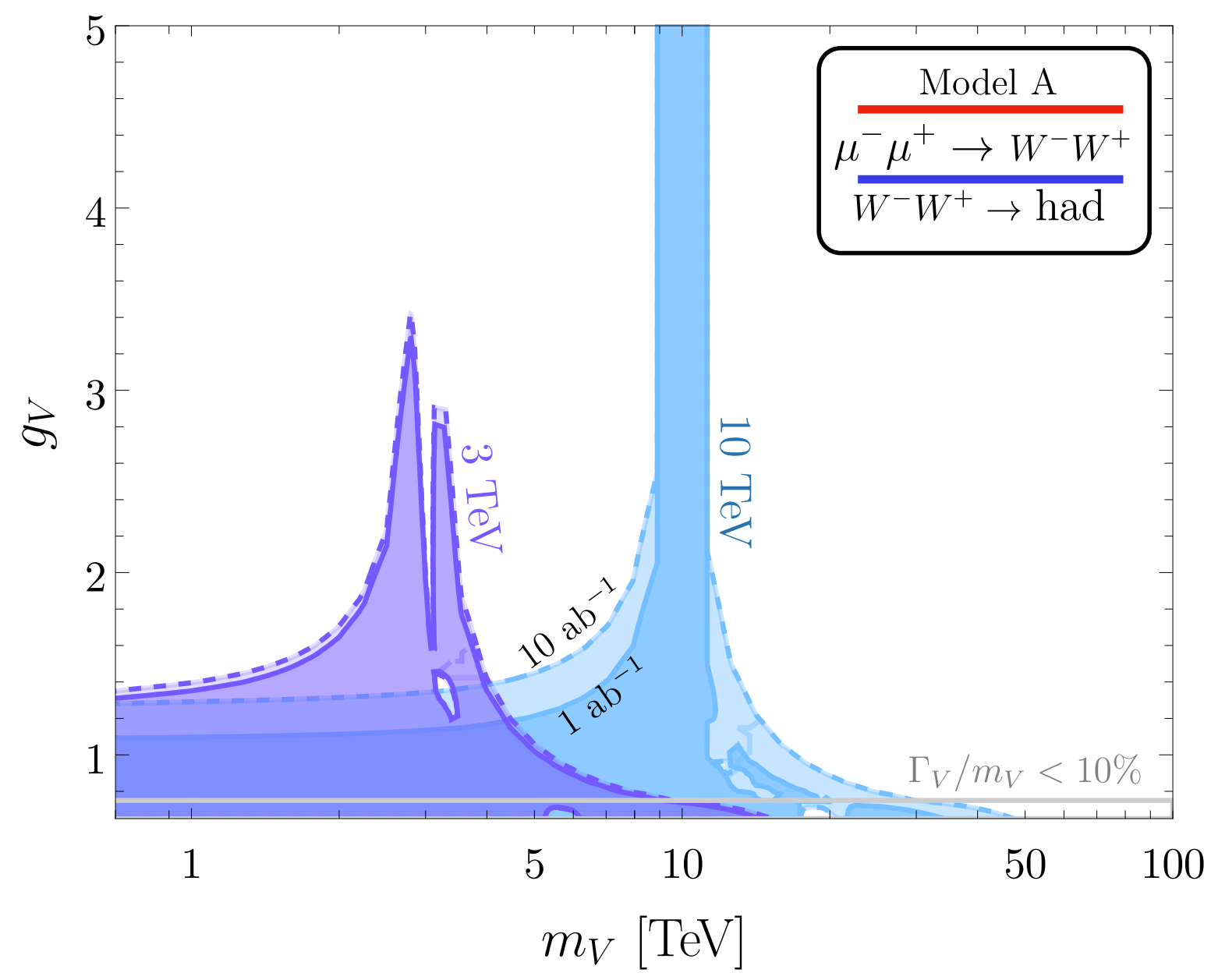
$$|\mathcal{M}|_{\gamma V}^2 = \frac{e^5 c_\ell}{s_W^5 2g_V m_W^4 (s - m_V^2)} \frac{st(s+t)}{m_W^4} \left[4c_W s_W^2 \theta_C - 2s_W^2 \theta_N + 2s_W^4 \theta_N \right] \quad (3.17)$$

$$|\mathcal{M}|_{ZV}^2 = \frac{e^5 c_\ell}{s_W^5 2g_V m_W^4 (s - m_V^2)} \frac{st(s+t)}{m_W^4} \left[2c_W \theta_C - 4c_W s_W^2 \theta_C - \theta_N - 2s_W^4 \theta_N + 3s_W^2 \theta_N \right] \quad (3.18)$$

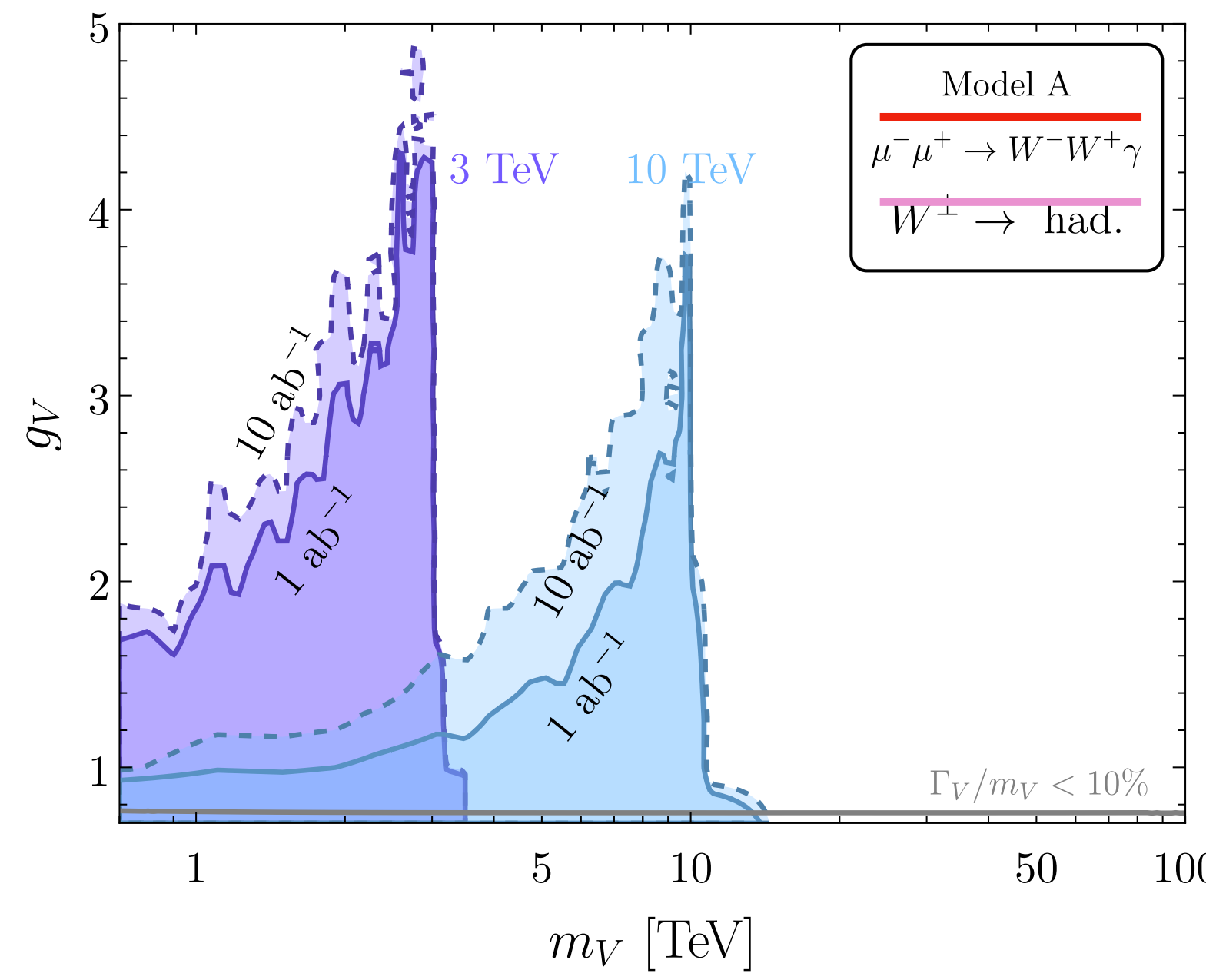
$$|\mathcal{M}|_{\nu V}^2 = c_C^3 c_N \frac{e^5 c_\ell}{s_W^5 2g_V m_W^4 (s - m_V^2)} \frac{st(s+t)}{m_W^4} \left[-2c_W \theta_C + \theta_N - s_W^2 \theta_N \right], \quad (3.19)$$



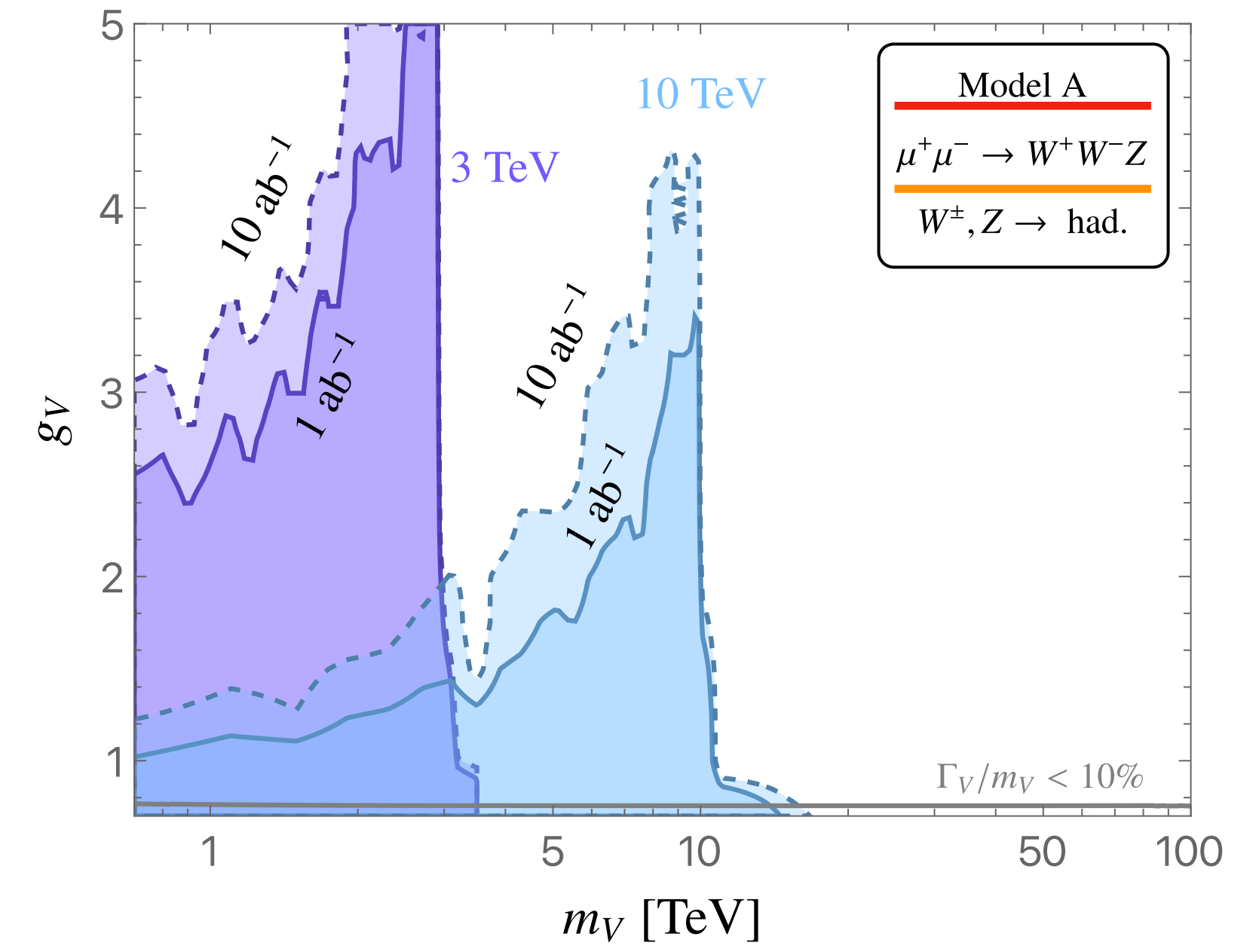




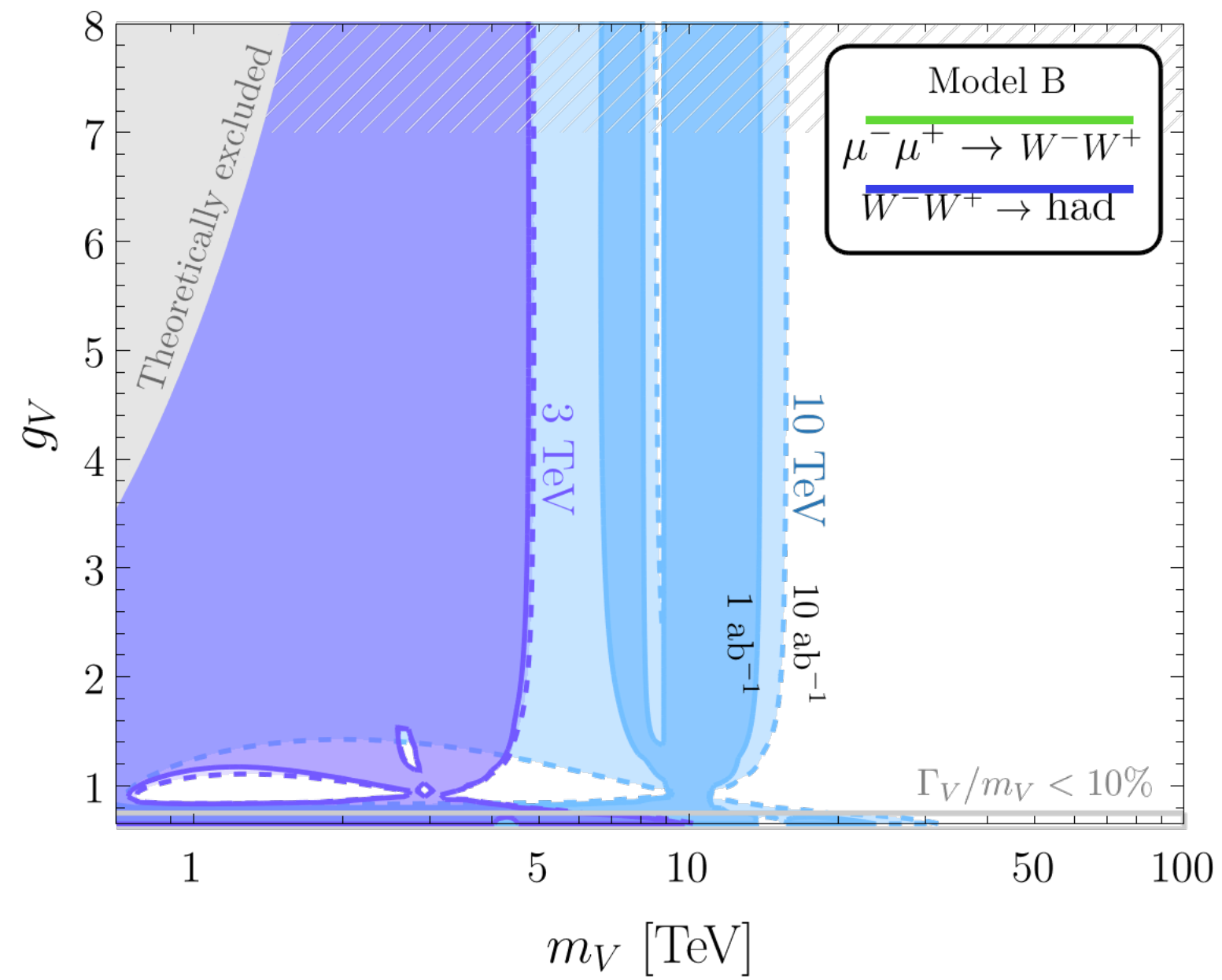
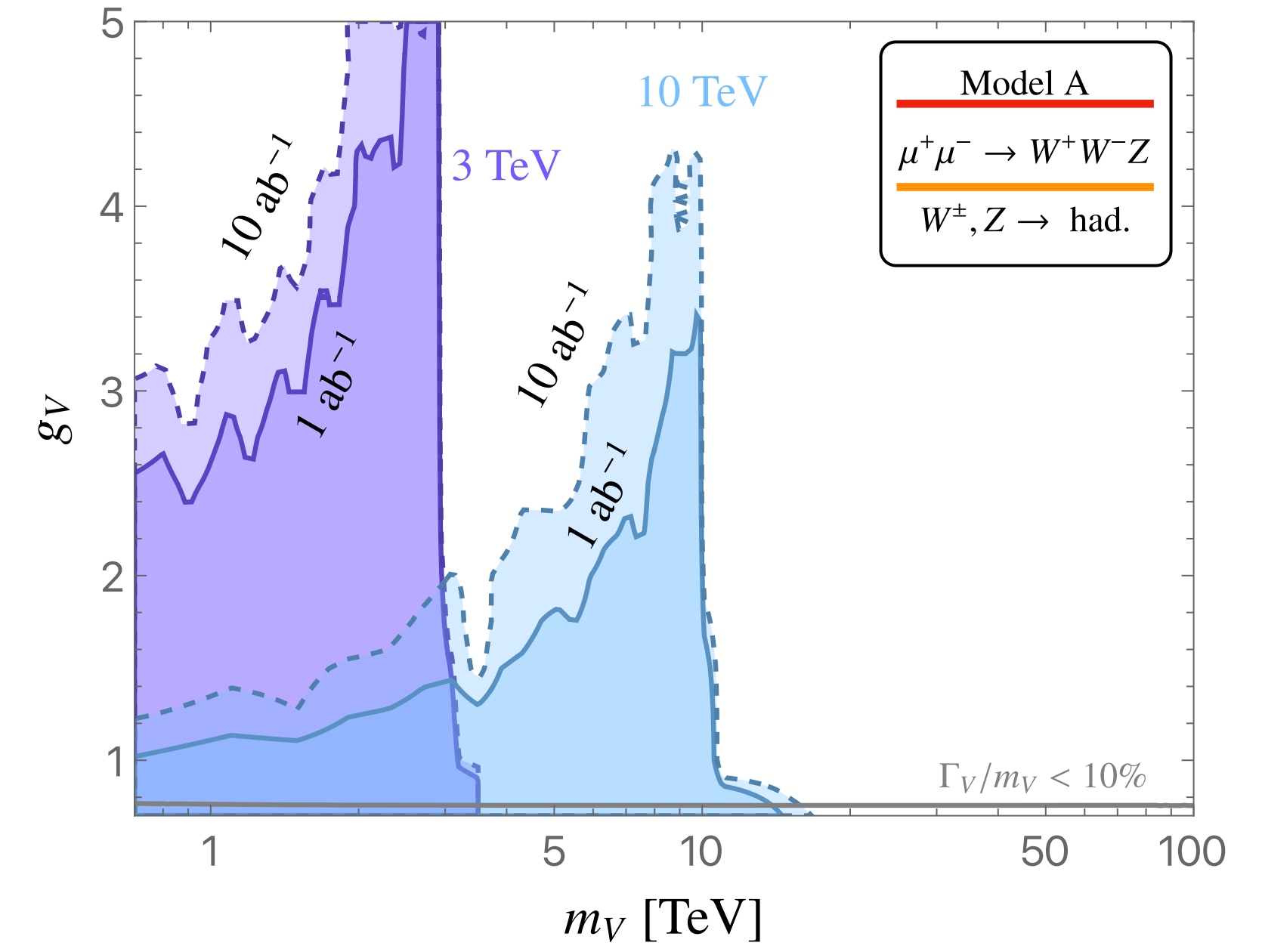
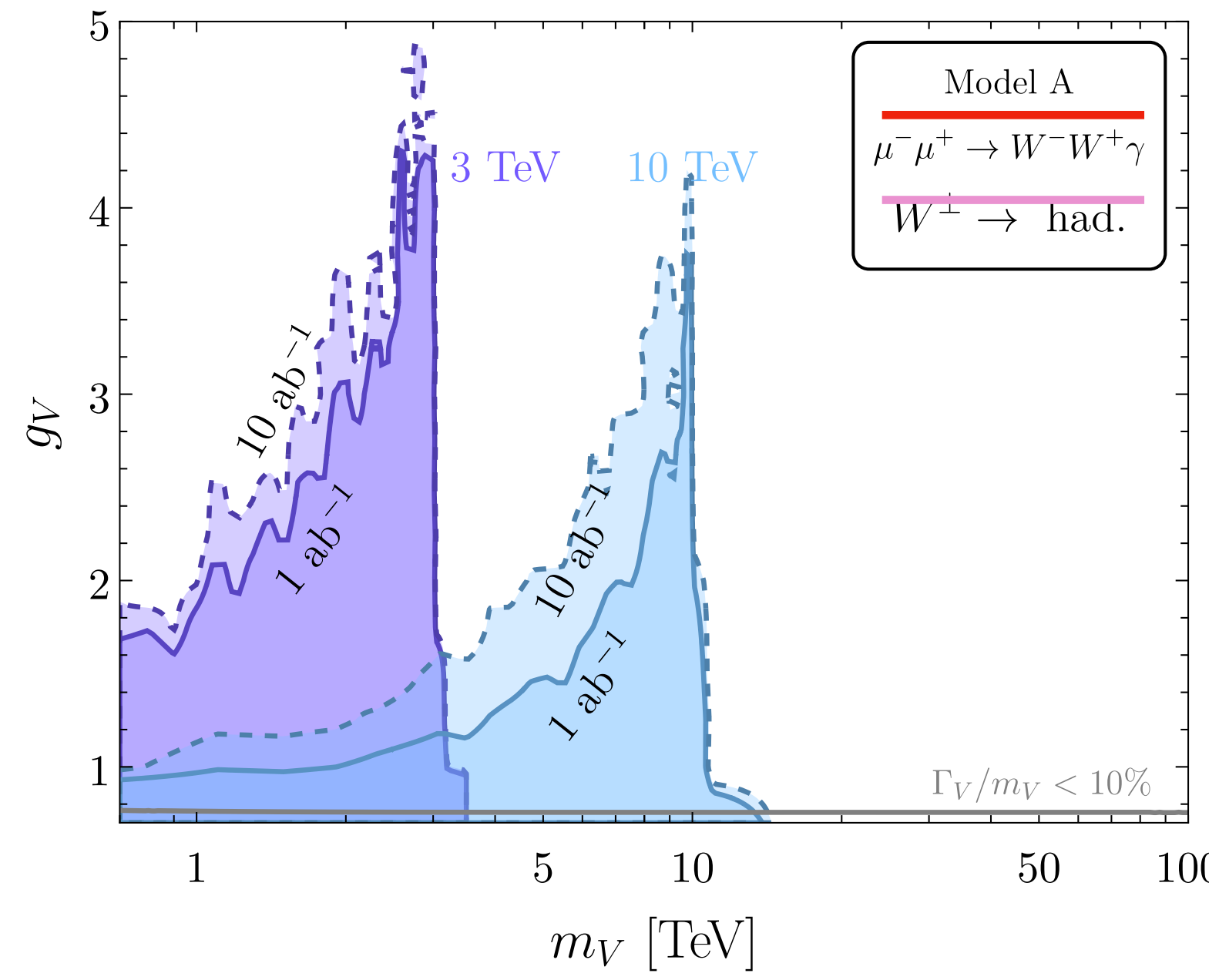
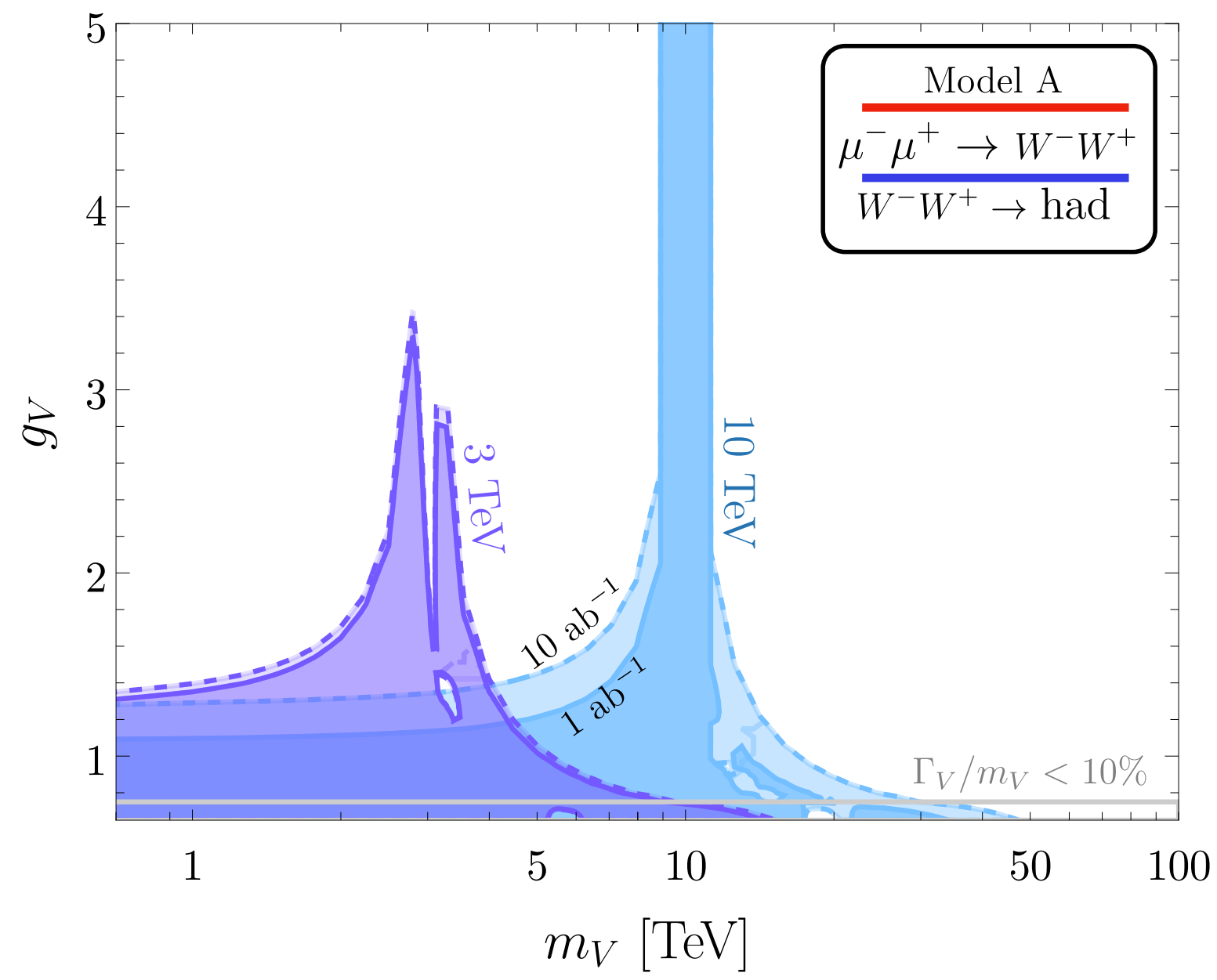
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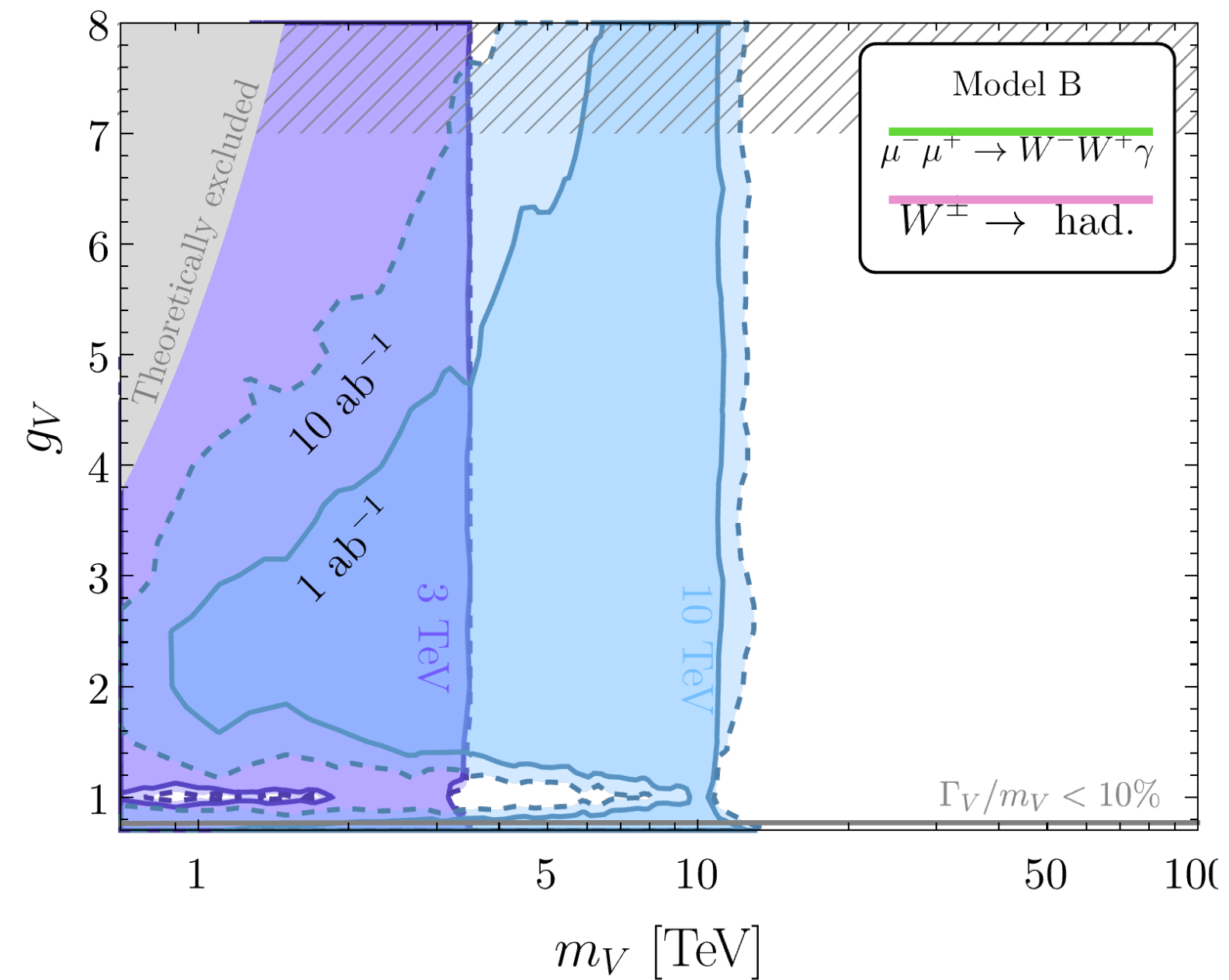
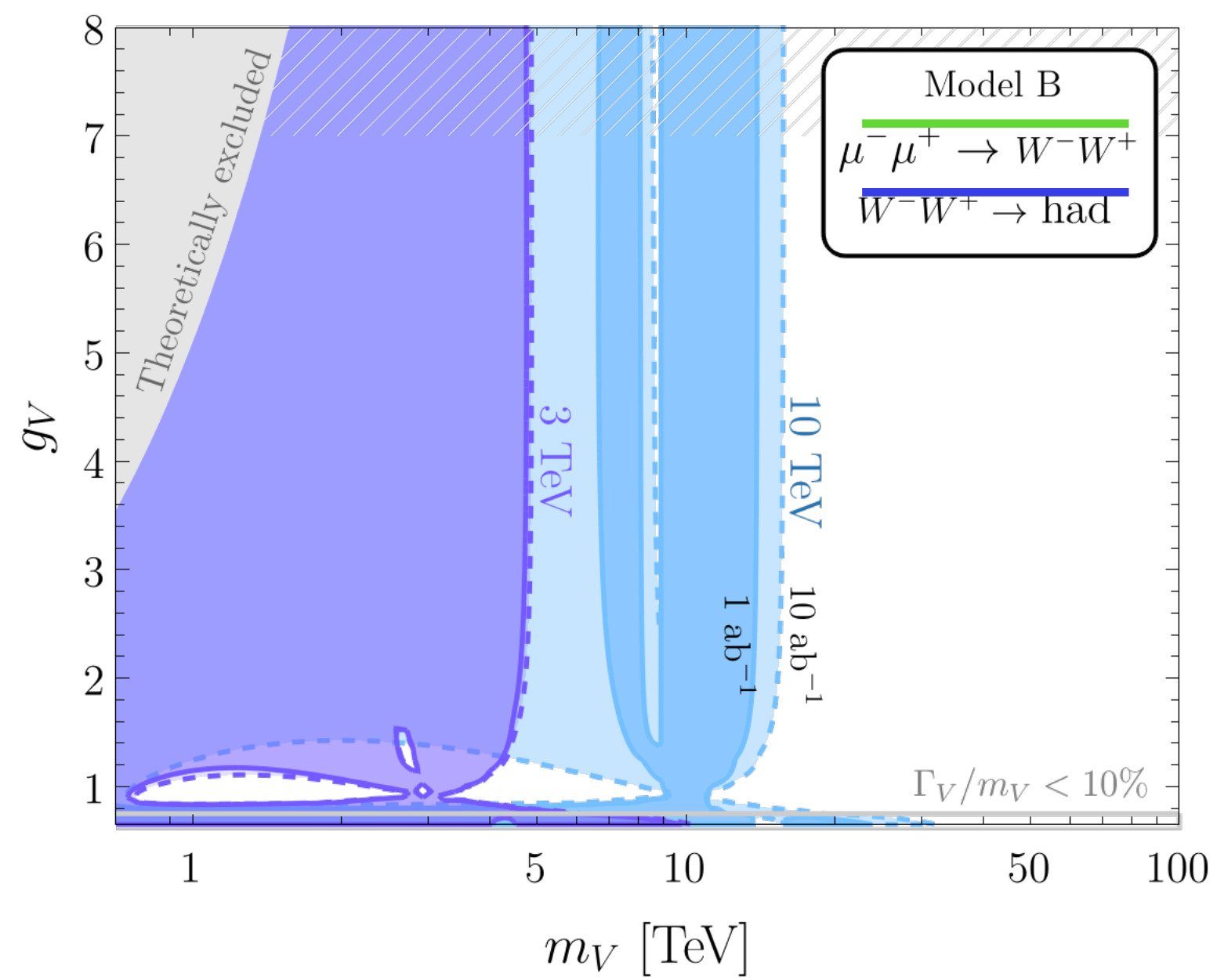
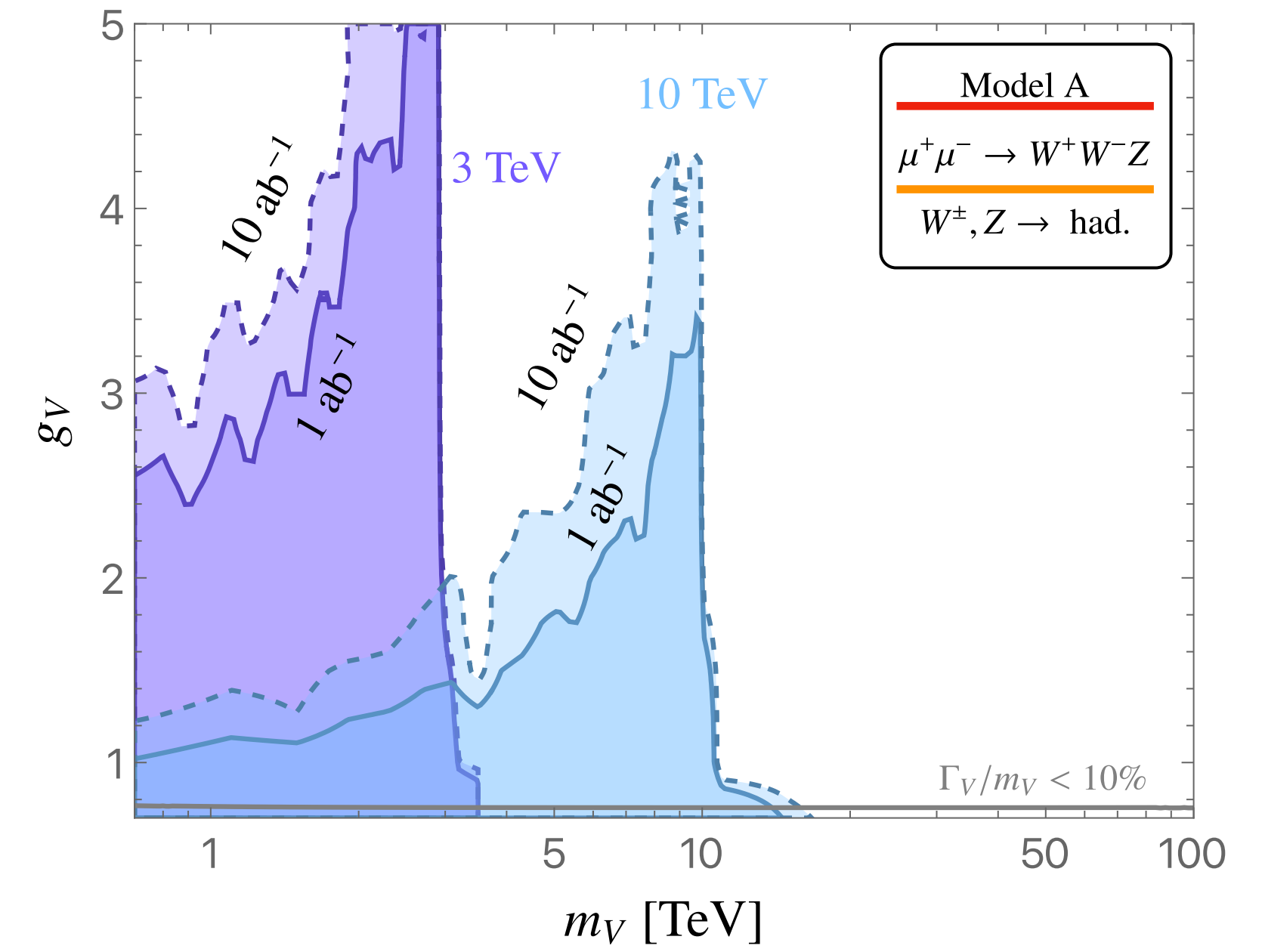
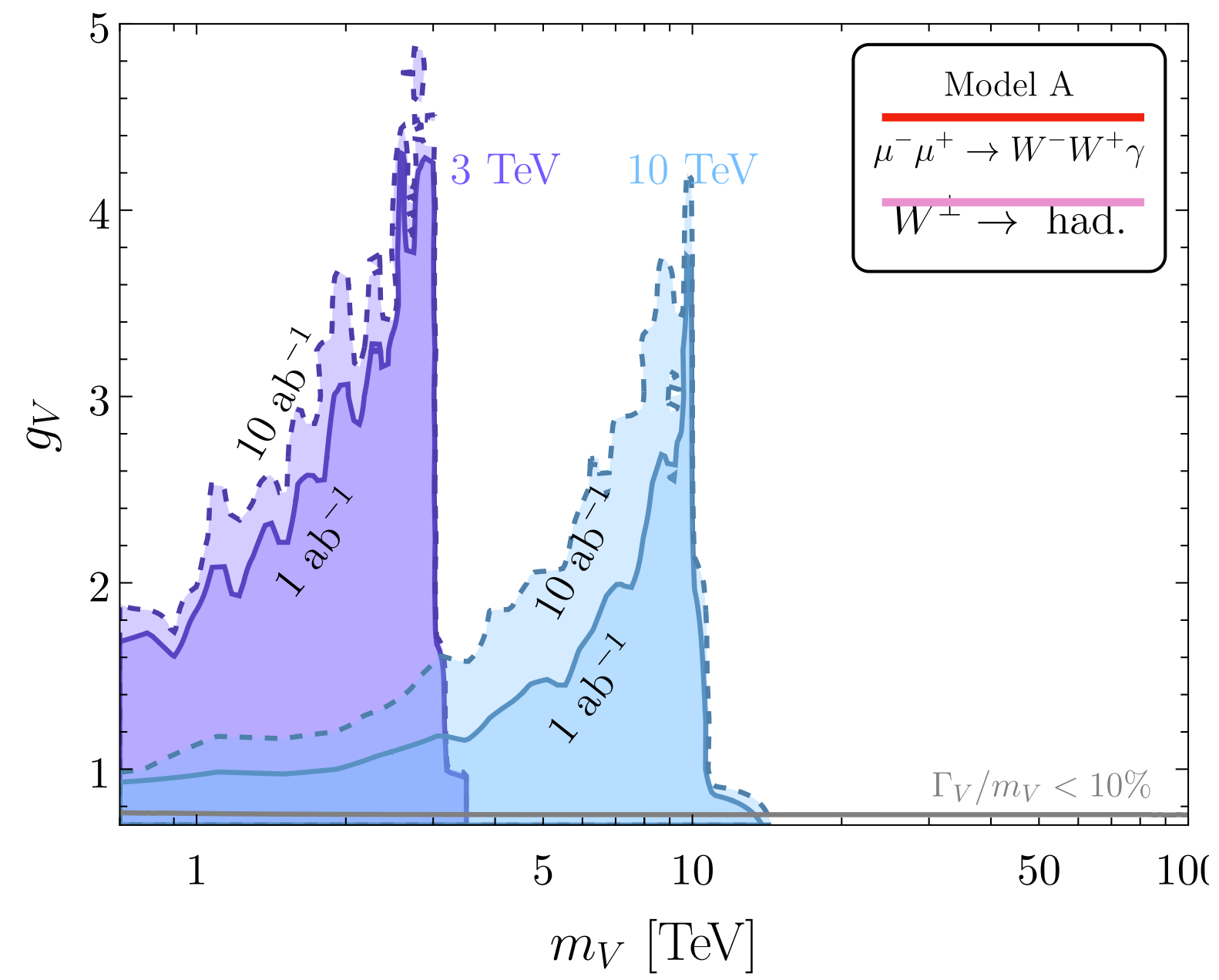
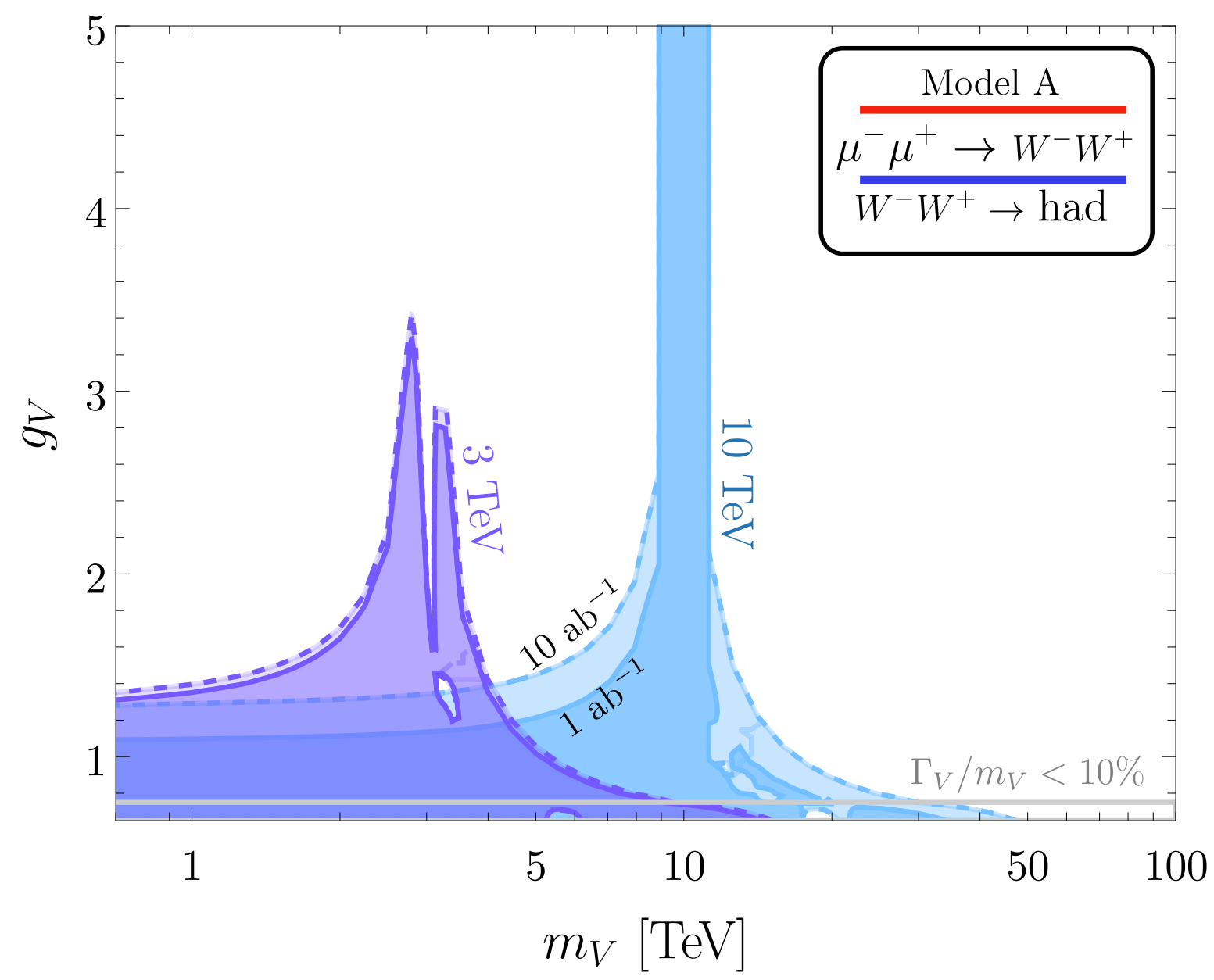
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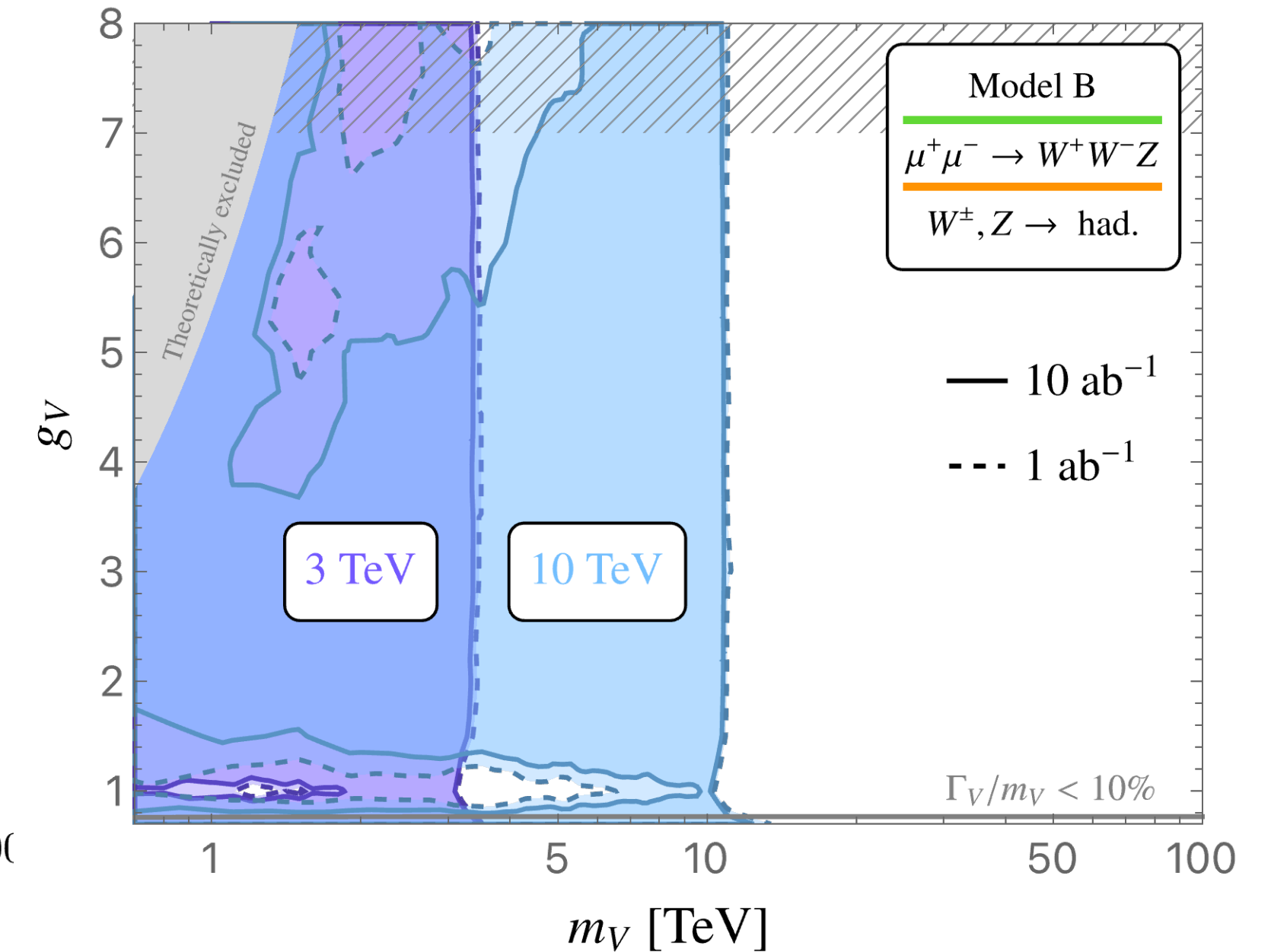
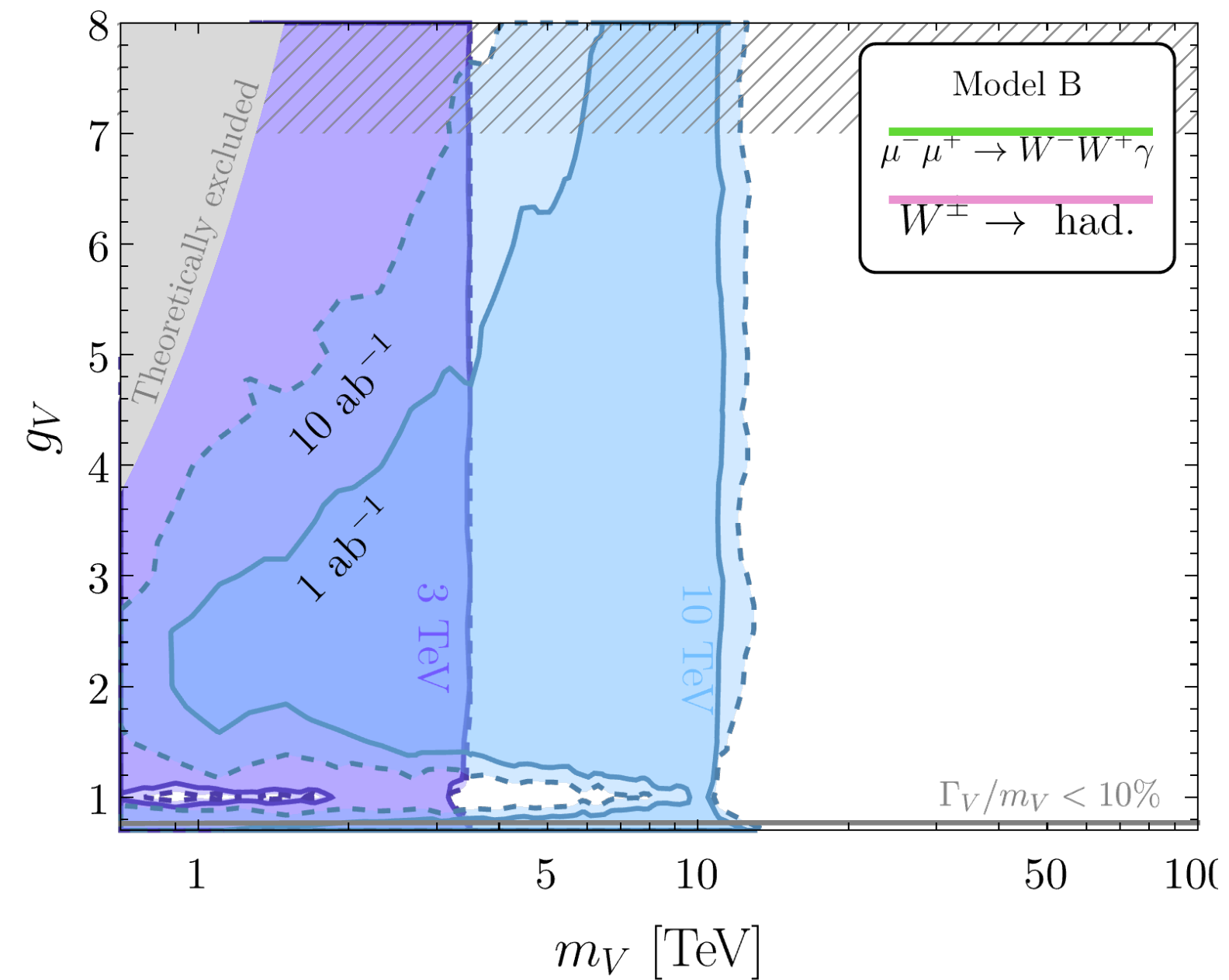
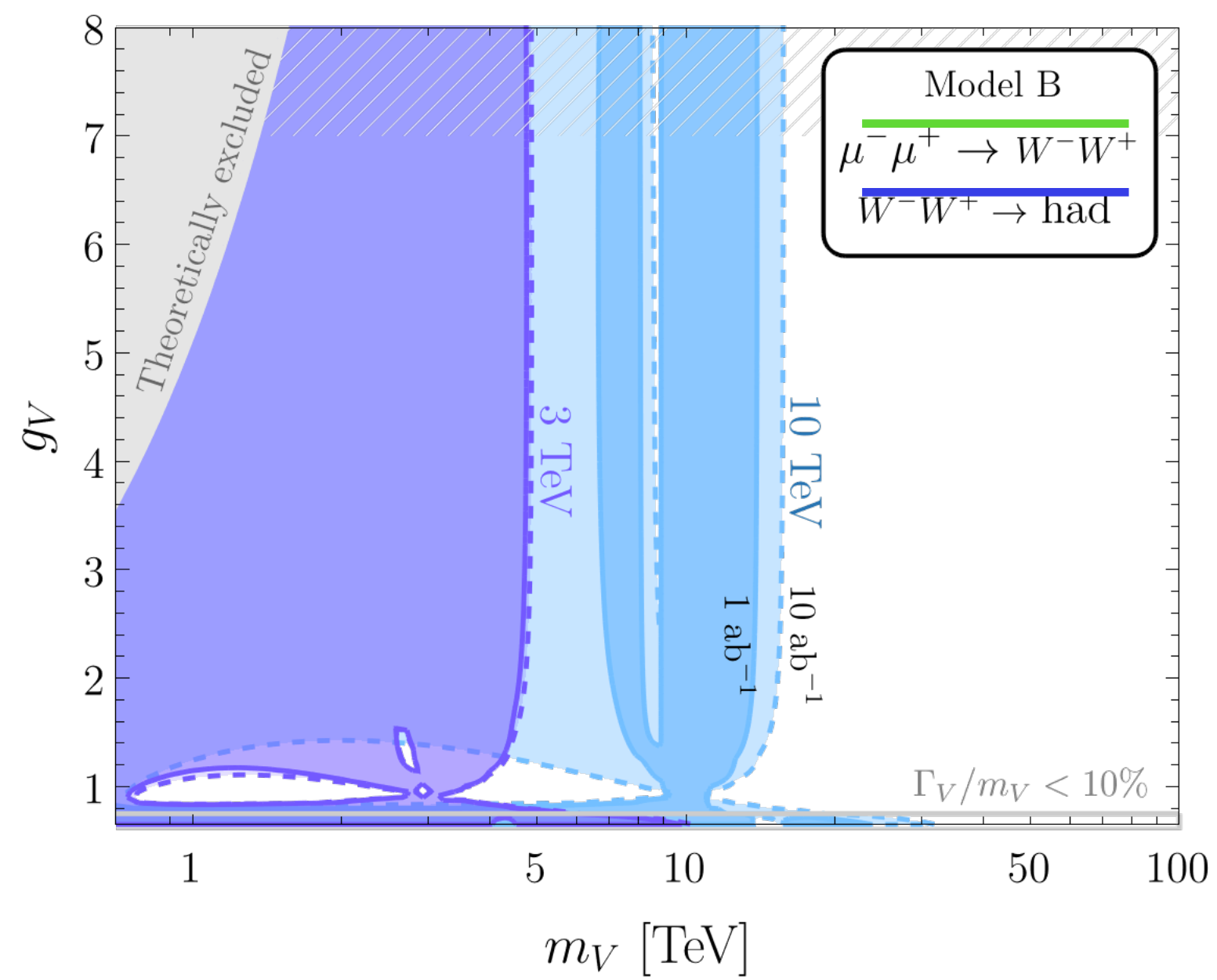
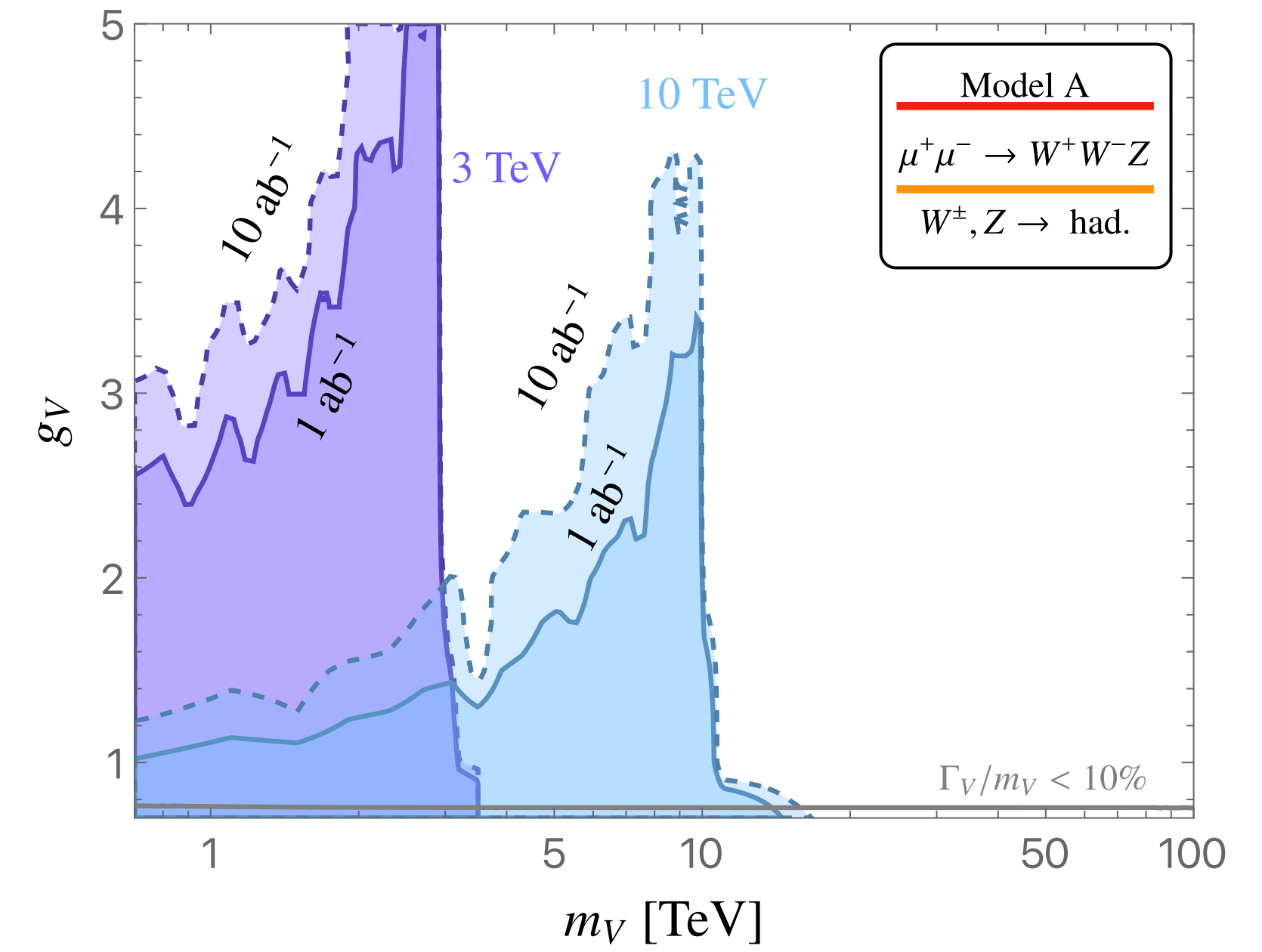
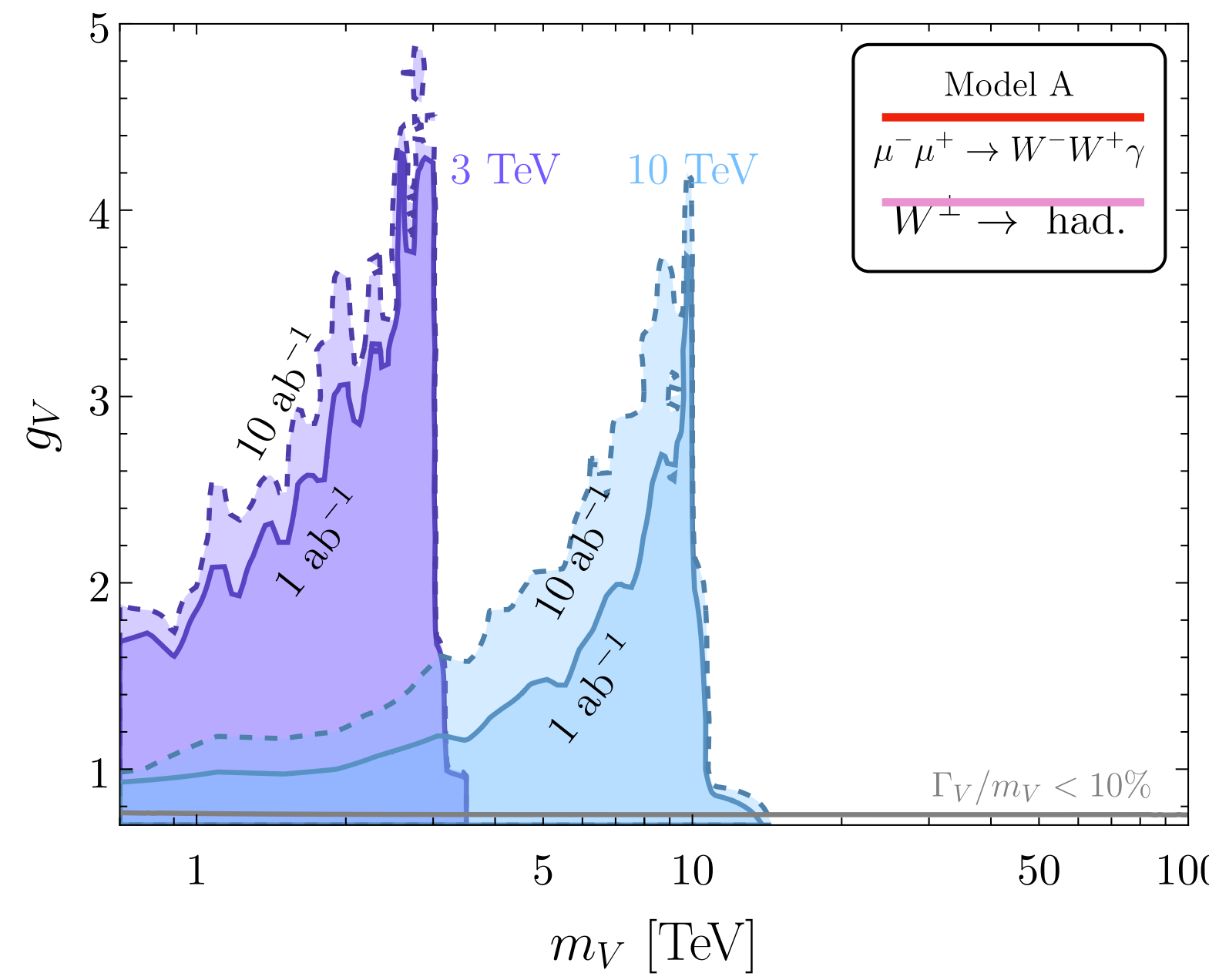
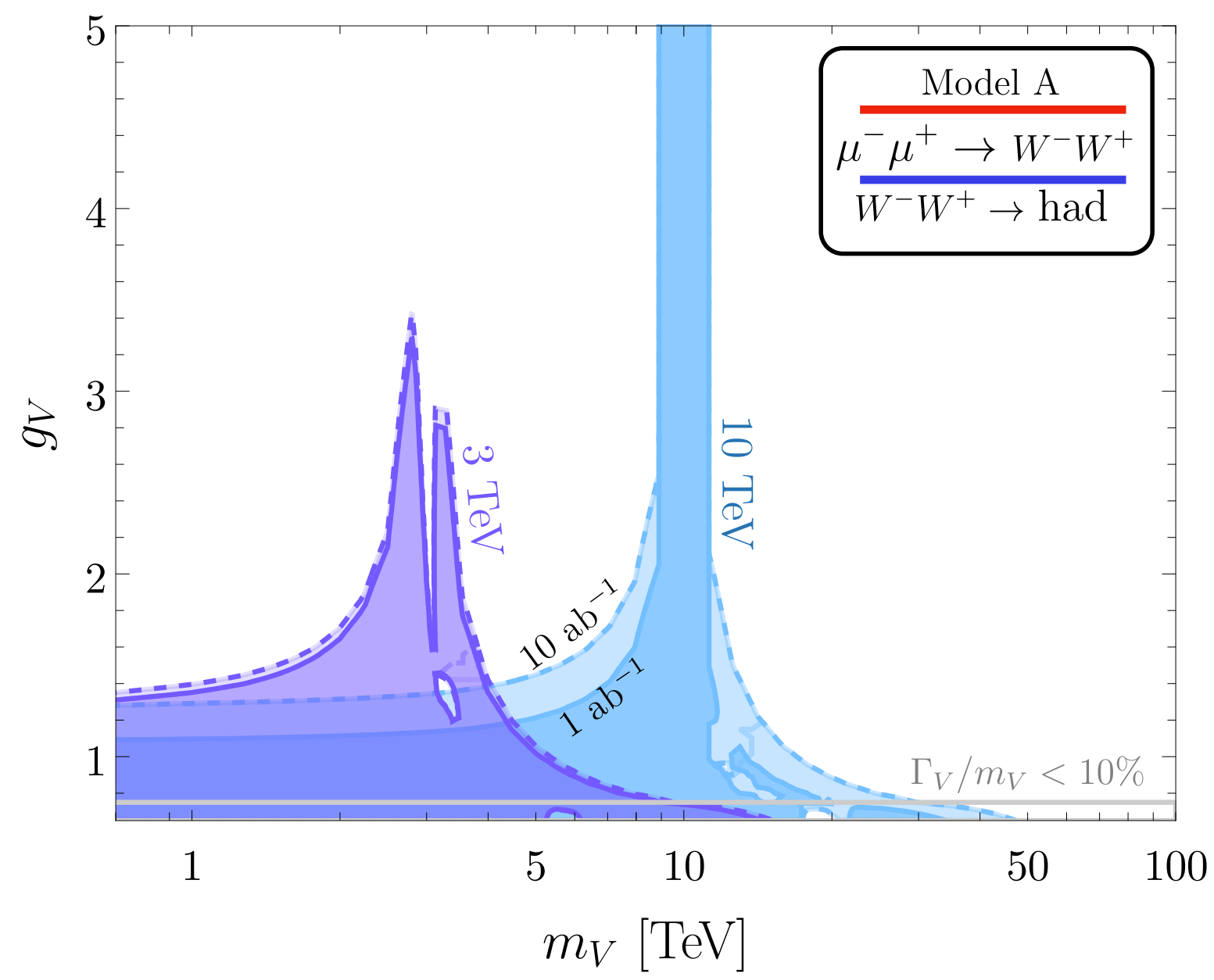
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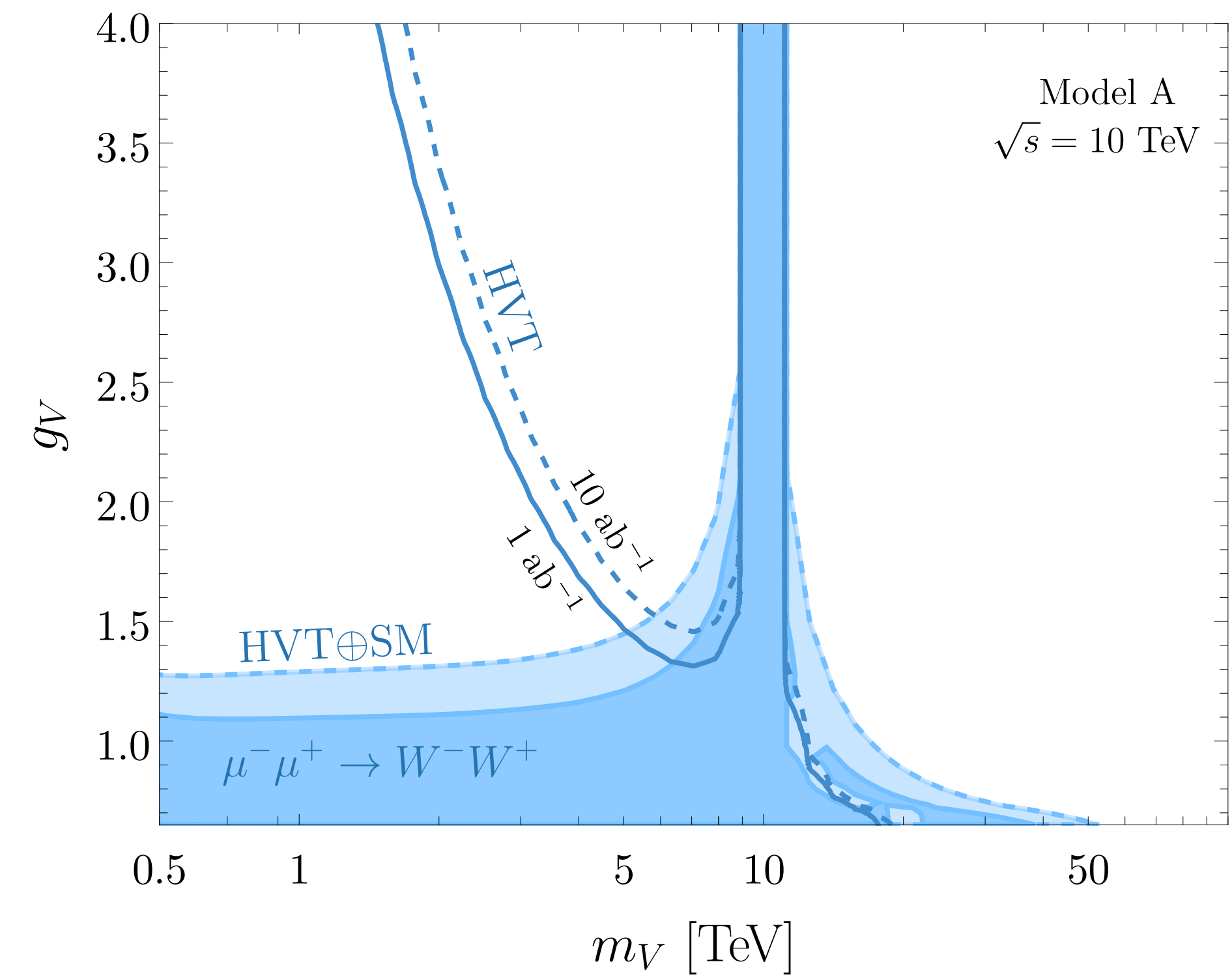
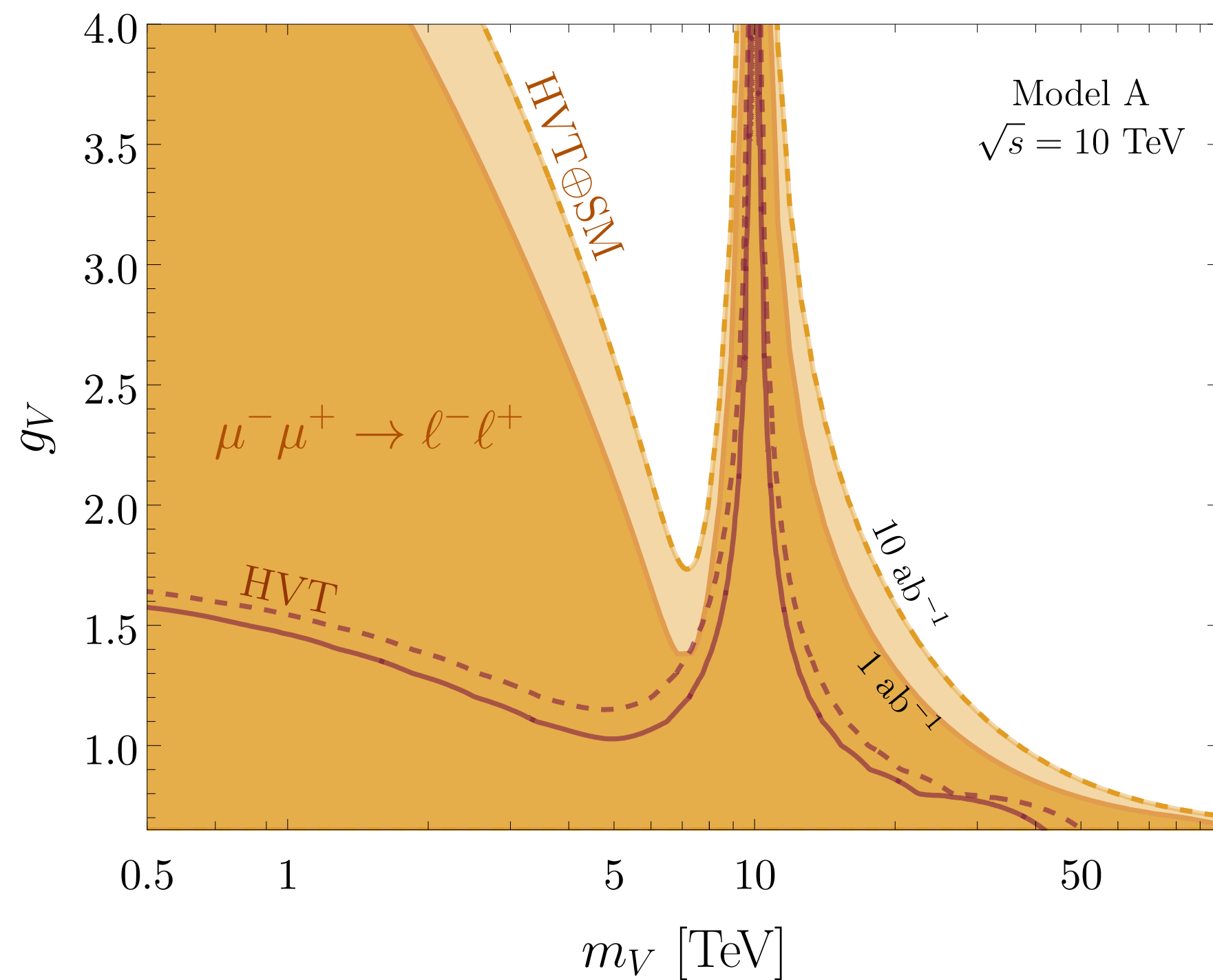
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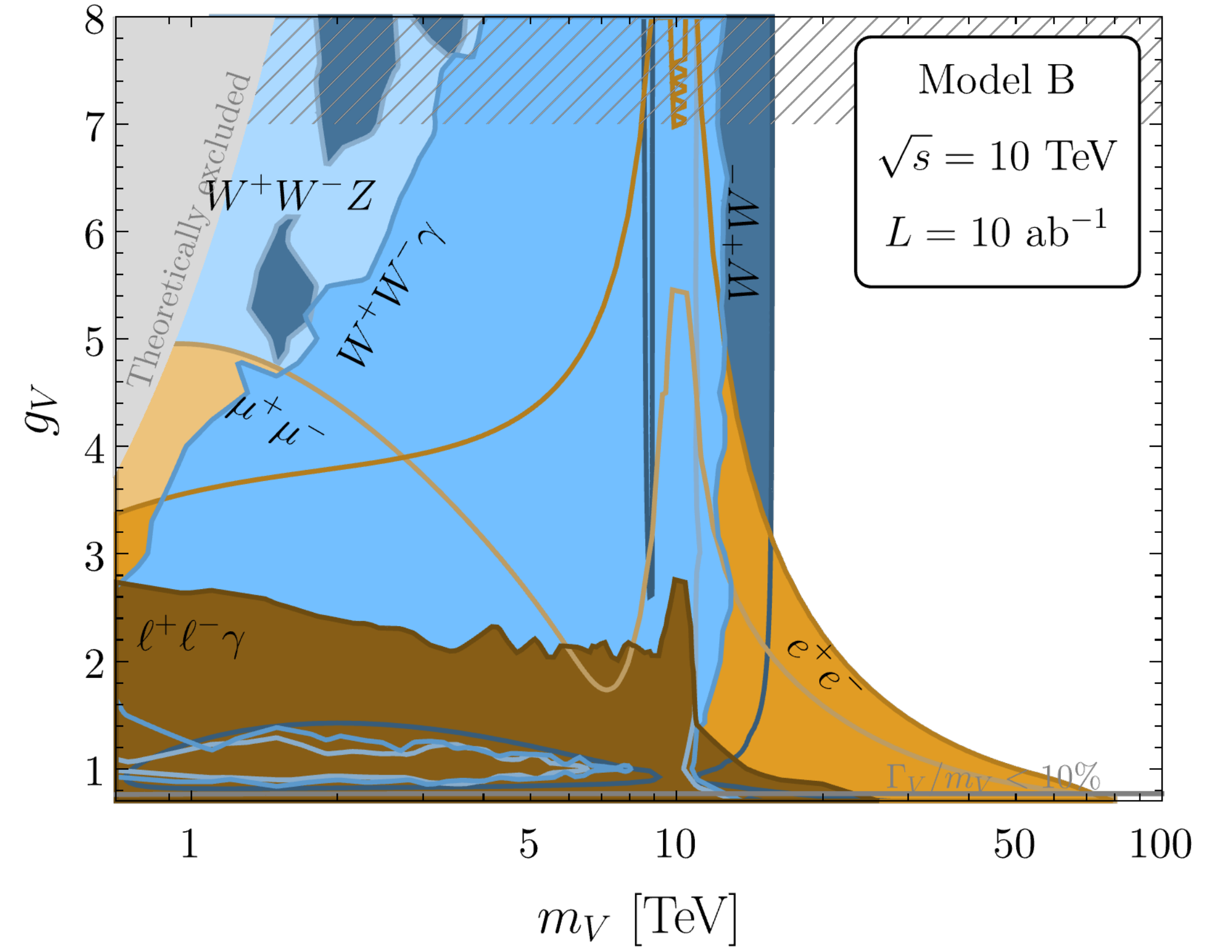
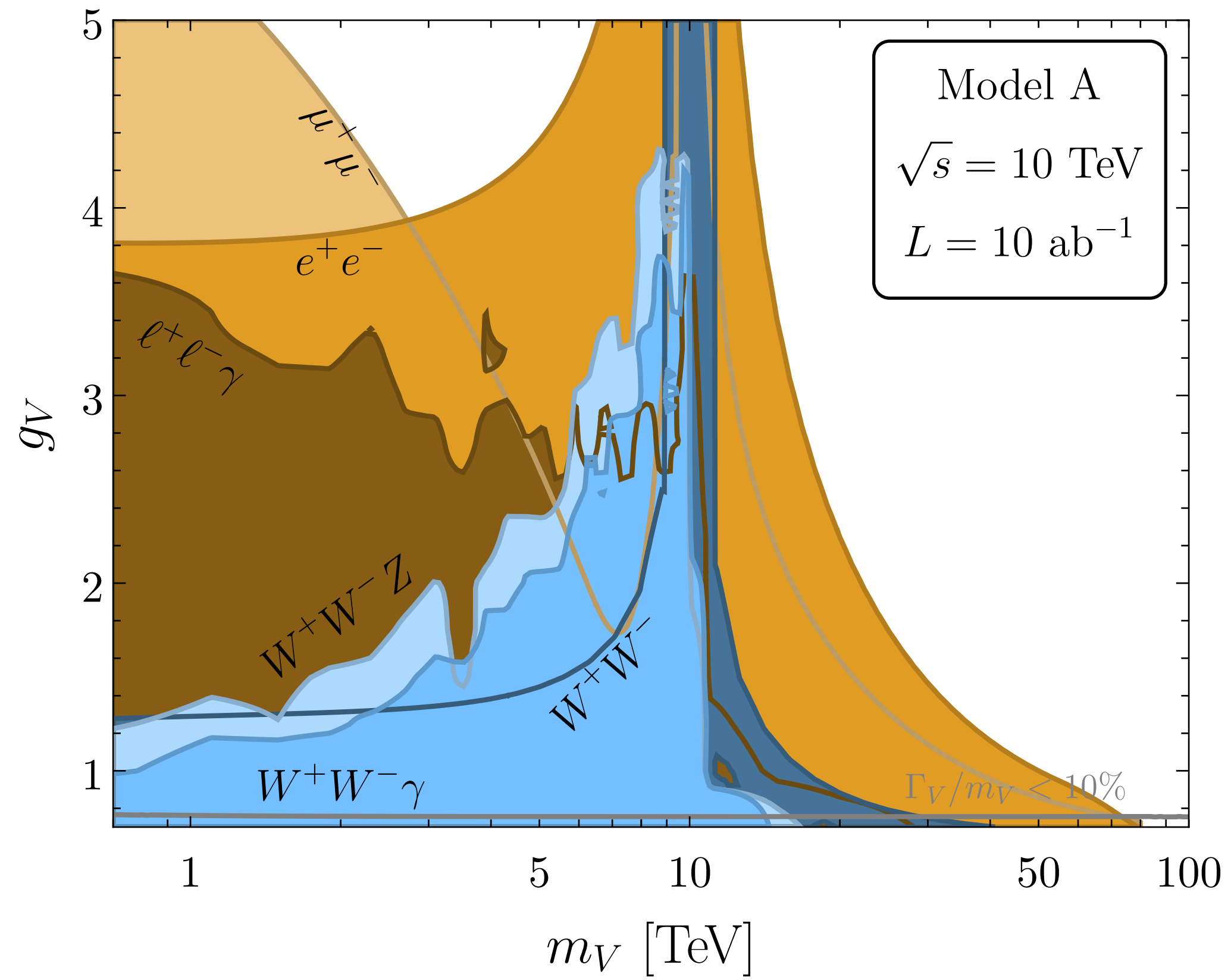
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Impact of interference

$$|\mathcal{M}|^2 = |\mathcal{M}_{\text{SM}} + \mathcal{M}_{\text{BSM}}|^2 = |\mathcal{M}_{\text{SM}}|^2 + \mathcal{M}_{\text{int}} + |\mathcal{M}_{\text{BSM}}|^2$$



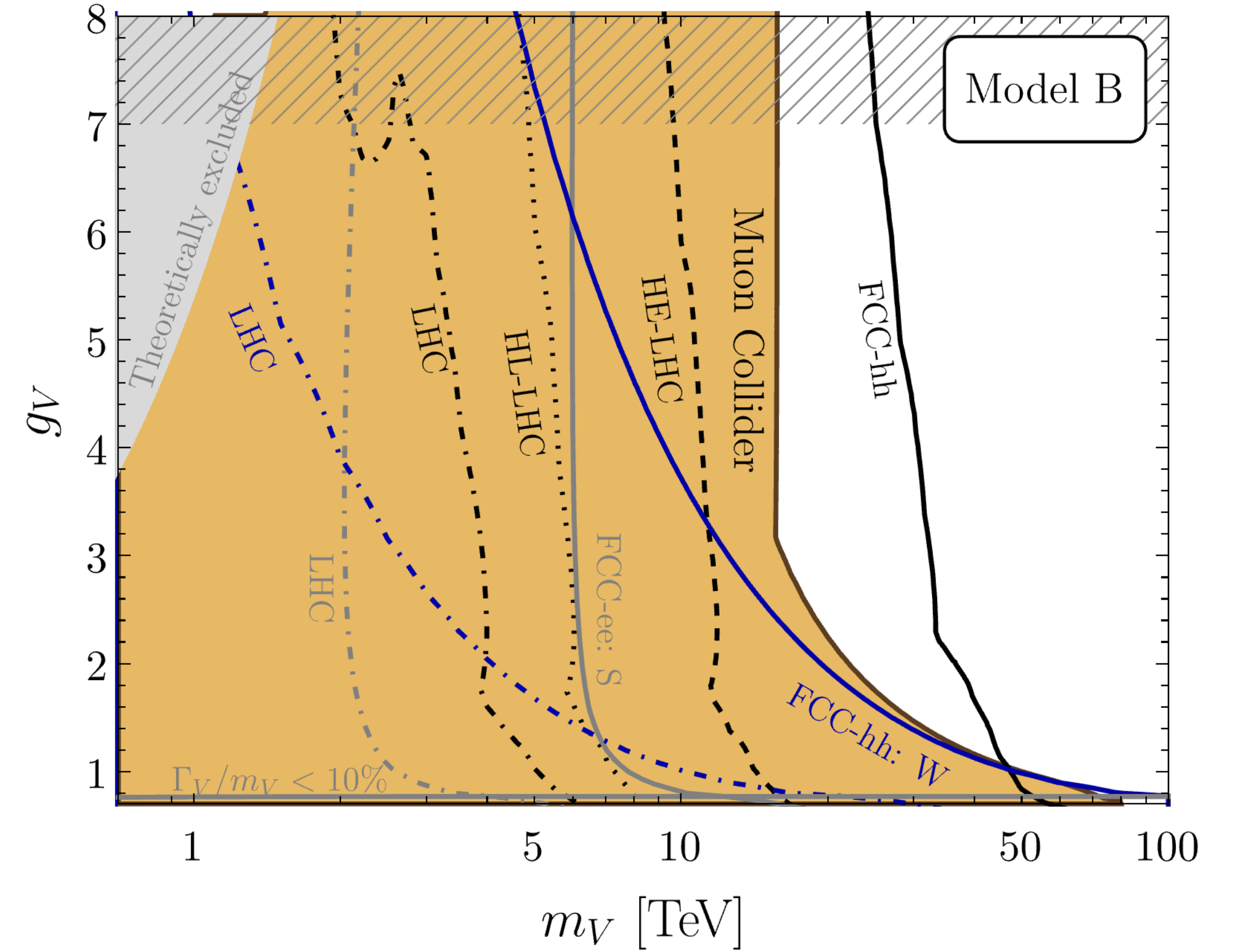
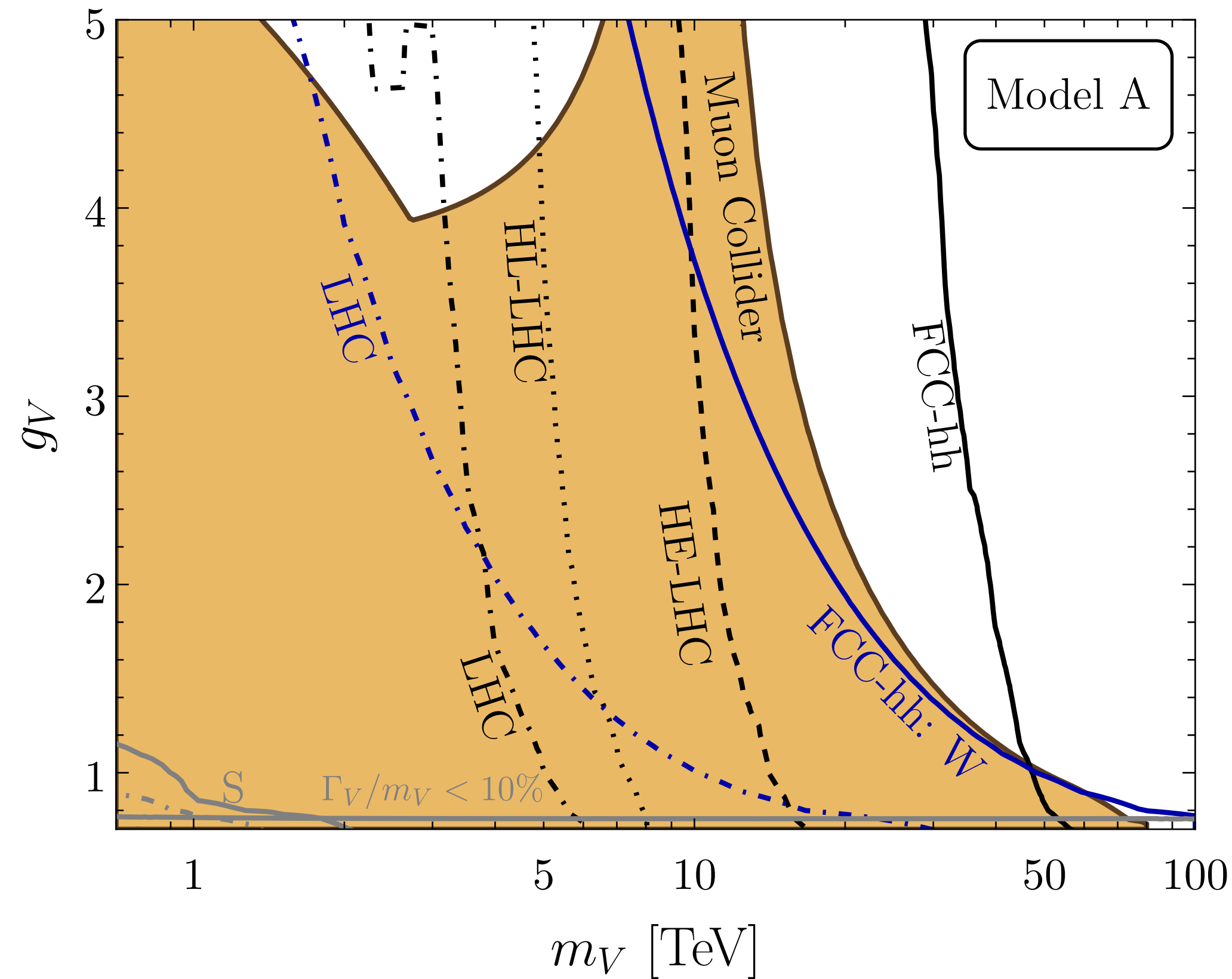


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Comparison with future colliders



Conclusions

- HVT can be found already with $2 \rightarrow 2$ scatterings
- Must include interference
- \sim all (m_V, g_V) plane excluded
- Overcomes all LHC stages, competitive with FCC

Conclusions

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- Must include interference
- \sim all (m_V, σ)
- Over all LHC stages, competitive with FCC

MONDAY ON YOUR ARXIV FEED!

Thank you!
Questions?

Backup slides

Muon Collider

MuC performance

Requirement	Baseline		Aspirational
	$\sqrt{s} = 3 \text{ TeV}$	$\sqrt{s} = 10 \text{ TeV}$	
Angular acceptance	$ \eta < 2.5$	$ \eta < 2.5$	$ \eta < 4$
Minimum tracking distance [cm]	~ 3	~ 3	< 3
Forward muons ($\eta > 5$)	–	tag	$\sigma_p/p \sim 10\%$
Track σ_{p_T}/p_T^2 [GeV^{-1}]	4×10^{-5}	4×10^{-5}	1×10^{-5}
Photon energy resolution	$0.2/\sqrt{E}$	$0.2/\sqrt{E}$	$0.1/\sqrt{E}$
Neutral hadron energy resolution	$0.5/\sqrt{E}$	$0.4/\sqrt{E}$	$0.2/\sqrt{E}$
Timing resolution (tracker) [ps]	$\sim 30 - 60$	$\sim 30 - 60$	$\sim 10 - 30$
Timing resolution (calorimeters) [ps]	100	100	10
Timing resolution (muon system) [ps]	~ 50 for $ \eta > 2.5$	~ 50 for $ \eta > 2.5$	< 50 for $ \eta > 2.5$
Flavour tagging	b vs c	b vs c	b vs c , s -tagging
Boosted hadronic resonance ID	h vs W/Z	h vs W/Z	W vs Z

From Interim report for the International Muon Collider Collaboration 2025

Higgs

- We need to study the Higgs:
 - couplings to itself:
 - couplings to **BSM**
 - is it **composite**?
- What would the best machine be? To have h **on shell** you need
 - An e^+e^- collider of $\sqrt{s} \sim 240$ GeV . Lepton precision
 - A pp collider of $\sqrt{s} \sim 100$ TeV can study 3h vertices.

$$\mathcal{L} = \frac{1}{2}(\partial_\mu h)^2 - \frac{m_h^2}{2}h^2 - \lambda_3 h^3 - \lambda_4 h^4$$

- Higgs was found at **pp** collider:
 - Dirty: richer but unpredictable initial state. Do not control soft gluon emission
- A **MuC** can:
 - have lepton **precision**
 - go to **high energy**
 - emit EW **radiation** controllably

Muon Collider

- $\sqrt{s} \sim 3$ or 10 TeV $\rightarrow v_{\mu} \sim (1 - 10^{-10})c$
- **Lumi**: 1 or 10 ab^{-1} because we are interested in 1 ab cross section
- A MuC of **10 Km** has the same effective CoM energy of a pp collider with 100 km ring.
- $\tau_{\mu} \sim 2.2 \cdot 10^{-6} \text{ s}$ (but goes up to 0.2 s at 10 TeV)

The strength of a Muon Collider

The strength of a Muon Collider

- **Synchrotron radiation:** energy loss per turn around the circle $\Delta E \sim \frac{\alpha}{R} \left(\frac{E}{m} \right)^4$

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- **Synergy** with neutrino physics because a μ beam will also make a ν beam
- A Higgs factory can probe BSM only in loops. A MuC can probe **both Higgs and BSM**

HVT

Mixing

- Covariant derivatives give VW terms:

- $D_{[\mu} \mathcal{V}_{\nu]}^a = D_{\mu} \mathcal{V}_{\nu}^a - D_{\nu} \mathcal{V}_{\mu}^a$

$$D_{\mu} \mathcal{V}_{\nu}^a = \partial_{\mu} \mathcal{V}_{\nu}^a + g \epsilon^{abc} W_{\mu}^b \mathcal{V}_{\nu}^c$$

- HVT mixes with SM. Let's disentangle

From gauge to mass base

0

B
 W^3
 W^2
 W^1
 γ^3
 γ^2
 γ^1

From gauge to mass base

0

B
 W^3
 W^2
 W^1
 \mathcal{V}^3
 \mathcal{V}^2
 \mathcal{V}^1

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1

B
 \hat{W}^0
 \hat{W}^+
 \hat{W}^-
 \mathcal{V}^0
 \mathcal{V}^+
 \mathcal{V}^-

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From gauge to mass base

0

B
 W^3
 W^2
 W^1
 \mathcal{V}^3
 \mathcal{V}^2
 \mathcal{V}^1

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1

B
 \hat{W}^0
 \hat{W}^+
 \hat{W}^-
 \mathcal{V}^0
 \mathcal{V}^+
 \mathcal{V}^-

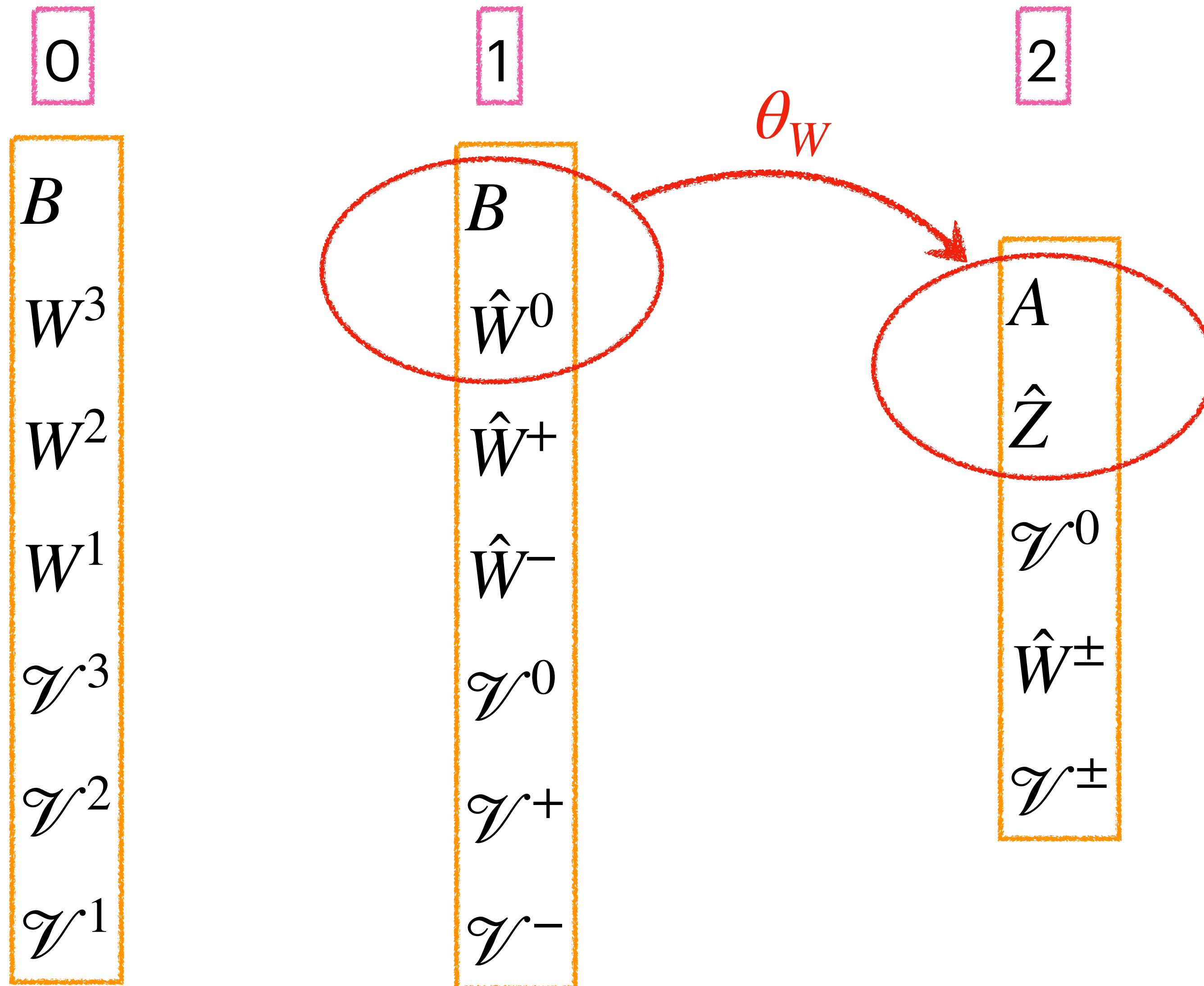
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2

A
 \hat{Z}
 \mathcal{V}^0
 \hat{W}^\pm
 \mathcal{V}^\pm

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From gauge to mass base



From gauge to mass base

0

B
 W^3
 W^2
 W^1
 \mathcal{V}^3
 \mathcal{V}^2
 \mathcal{V}^1

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1

B
 \hat{W}^0
 \hat{W}^+
 \hat{W}^-
 \mathcal{V}^0
 \mathcal{V}^+
 \mathcal{V}^-

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2

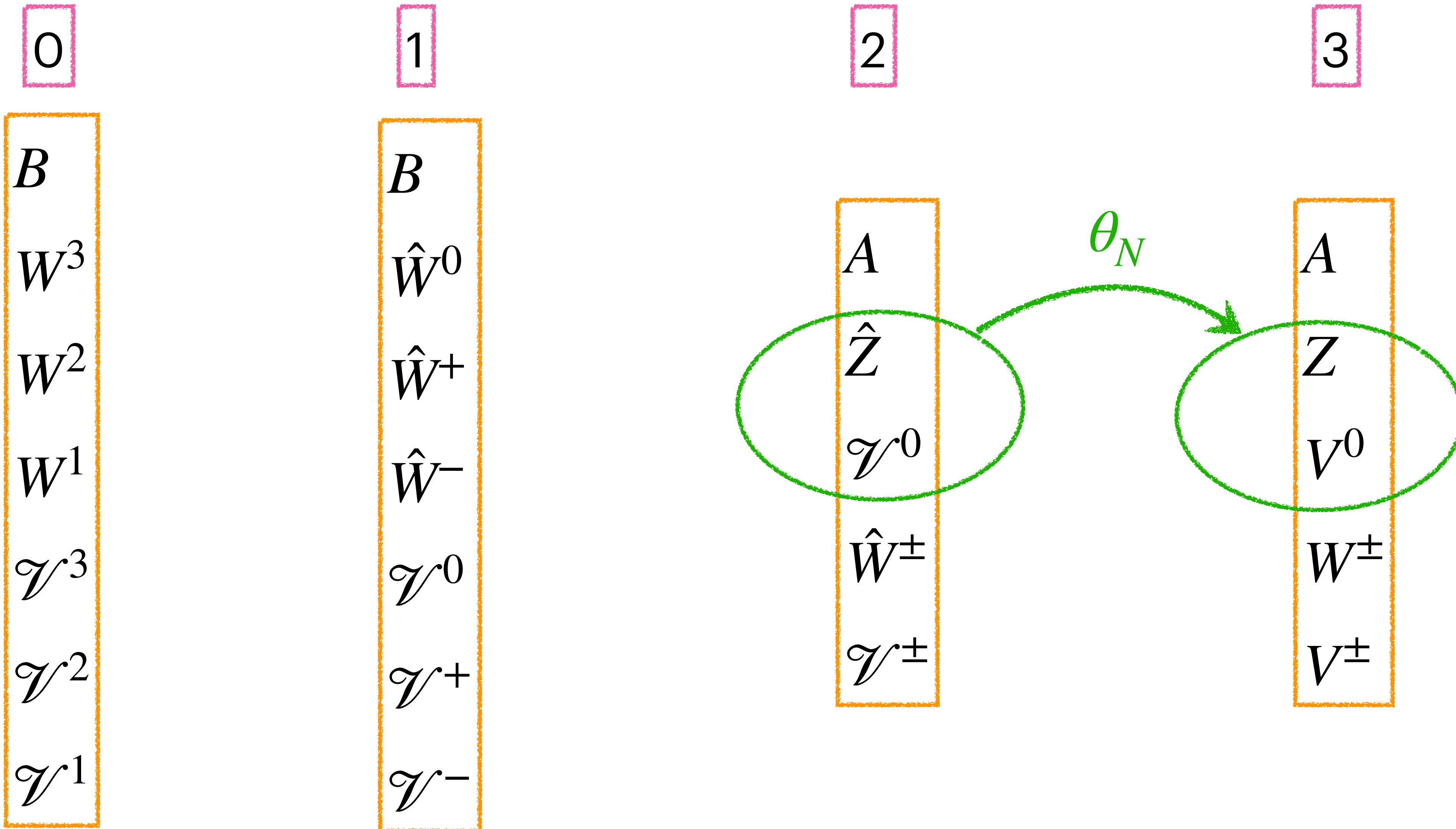
A
 \hat{Z}
 \mathcal{V}^0
 \hat{W}^\pm
 \mathcal{V}^\pm

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From gauge to mass base

0	1	2	3
B	B		
W^3	\hat{W}^0	A	A
W^2	\hat{W}^+	\hat{Z}	Z
W^1	\hat{W}^-	\mathcal{V}^0	V^0
\mathcal{V}^3	\mathcal{V}^0	\hat{W}^\pm	W^\pm
\mathcal{V}^2	\mathcal{V}^+	\mathcal{V}^\pm	V^\pm
\mathcal{V}^1	\mathcal{V}^-		

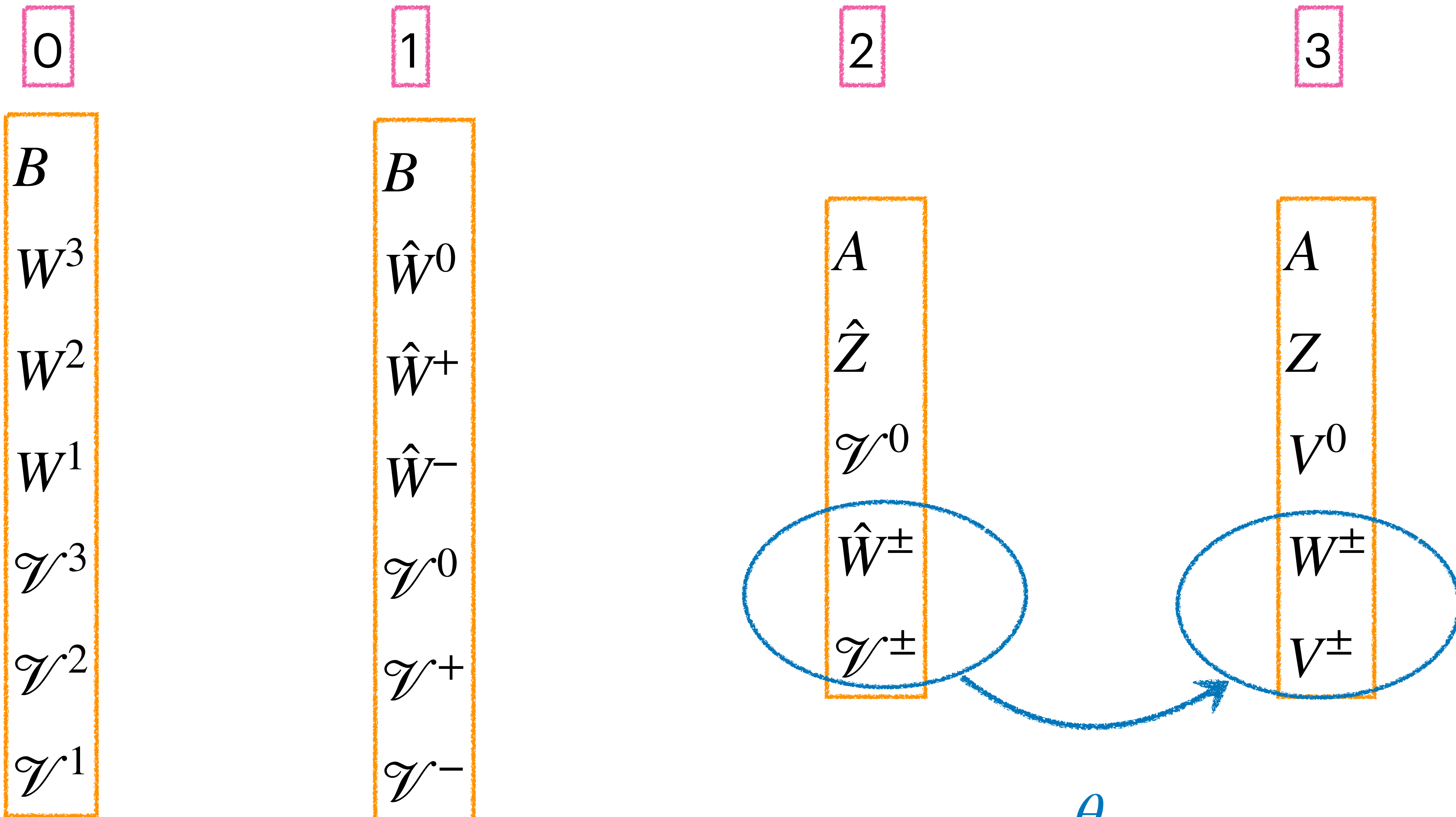
From gauge to mass base



From gauge to mass base

0	1	2	3
B	B		
W^3	\hat{W}^0	A	A
W^2	\hat{W}^+	\hat{Z}	Z
W^1	\hat{W}^-	\mathcal{V}^0	V^0
\mathcal{V}^3	\mathcal{V}^0	\hat{W}^\pm	W^\pm
\mathcal{V}^2	\mathcal{V}^+	\mathcal{V}^\pm	V^\pm
\mathcal{V}^1	\mathcal{V}^-		

From gauge to mass base



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From gauge to mass base

0	1	2	3
B	B		
W^3	\hat{W}^0	A	A
W^2	\hat{W}^+	\hat{Z}	Z
W^1	\hat{W}^-	\mathcal{V}^0	V^0
\mathcal{V}^3	\mathcal{V}^0	\hat{W}^\pm	W^\pm
\mathcal{V}^2	\mathcal{V}^+	\mathcal{V}^\pm	V^\pm
\mathcal{V}^1	\mathcal{V}^-		

- $\mathcal{V}_\mu^\pm = \frac{\mathcal{V}_\mu^1 \mp i\mathcal{V}_\mu^2}{\sqrt{2}}, \quad \mathcal{V}_\mu^0 = \mathcal{V}_\mu^3$

- Massless photon: $A_\mu = c_W B_\mu + s_W W_\mu^3$

- $\tan \theta_W = \frac{g'}{g}$

- $\tan 2\theta_N = g_V c_H \frac{\hat{v}}{\hat{m}_V} \frac{\hat{m}_Z \hat{m}_V}{\hat{m}_V^2 - \hat{m}_Z^2}$

- $\tan 2\theta_C = g_V c_H \frac{\hat{v}}{\hat{m}_V} \frac{\hat{m}_W \hat{m}_V}{\hat{m}_V^2 - \hat{m}_W^2}$

- To reproduce SM $\frac{\hat{m}_{W,Z}}{\hat{m}_V} \ll 1$

- Custodial symmetry: if $\hat{m}_W \rightarrow \hat{m}_Z$ then $\theta_C \rightarrow \theta_N$

- $m_W^2 M_+^2 = \cos^2 \theta_W m_Z^2 M_0^2$

i.e. $M_+^2 = M_0^2 (1 + \mathcal{O}(\%))$

A small correction to SM masse

- $\hat{m}_Z = \frac{e}{2 \sin \theta_W \cos \theta_W} \hat{v}$
- $\hat{m}_W = \cos \theta_W \hat{m}_Z$
- $\hat{m}_V = \sqrt{m_{\mathcal{V}}^2 + g_V^2 c_{VVHH} \hat{v}^2}$
- $\hat{v} = \frac{\mu}{\sqrt{\lambda}}$

$$m_Z^2 = \hat{m}_Z^2 (1 - c_H^2 \zeta^2) \left(1 + \mathcal{O}(\hat{m}_Z^2 / \hat{m}_V^2) \right)$$

$$m_W^2 = \hat{m}_W^2 (1 - c_H^2 \zeta^2) \left(1 + \mathcal{O}(\hat{m}_W^2 / \hat{m}_V^2) \right)$$

$$\zeta := \frac{g_V \hat{v}}{2 \hat{m}_V} \ll 1$$

Model A (weakly coupled)

- Extend the SM with another SU(2):
 $SU(2)_1 \times SU(2)_2 \times U(1)_Y$
- You get two copies of Ws, one going to SM W and one making the HVT
- You break the symmetry down to SM thanks to the real bidoublet
 $\Phi = (2,2)_0$ with vev $\begin{pmatrix} f & 0 \\ 0 & f \end{pmatrix}$

Model B (strongly coupled)

- In general:
 - the Higgs boson is the pNGB of some broken symmetry of a strong new force
 - Feature heavy vector resonances with EW quantum numbers.
- Example: Minimal Composite Higgs Model (MCHM). $SO(5)$ spontaneously broken to $SO(4)$

ElectroWeak Precision Tests (EWPT)

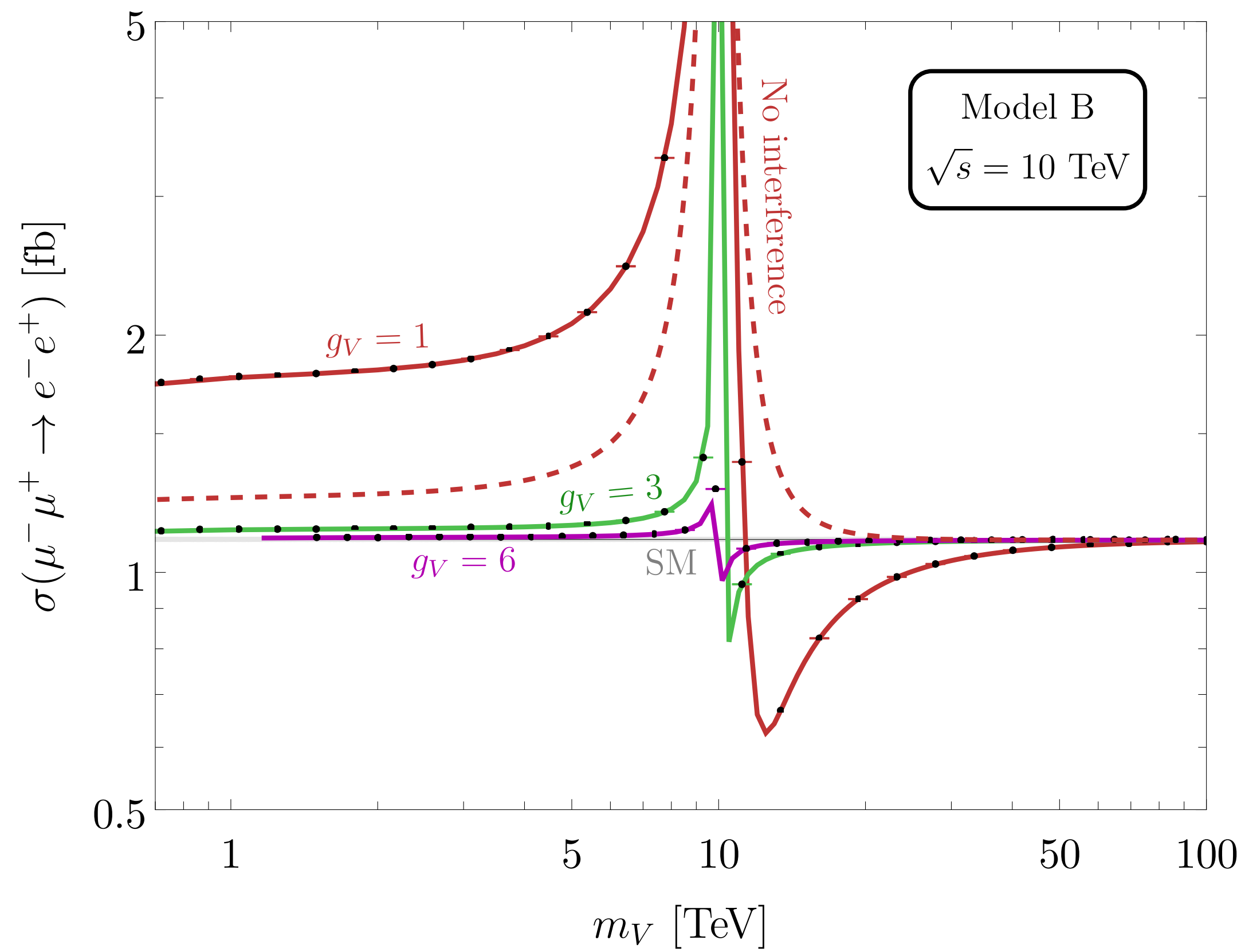
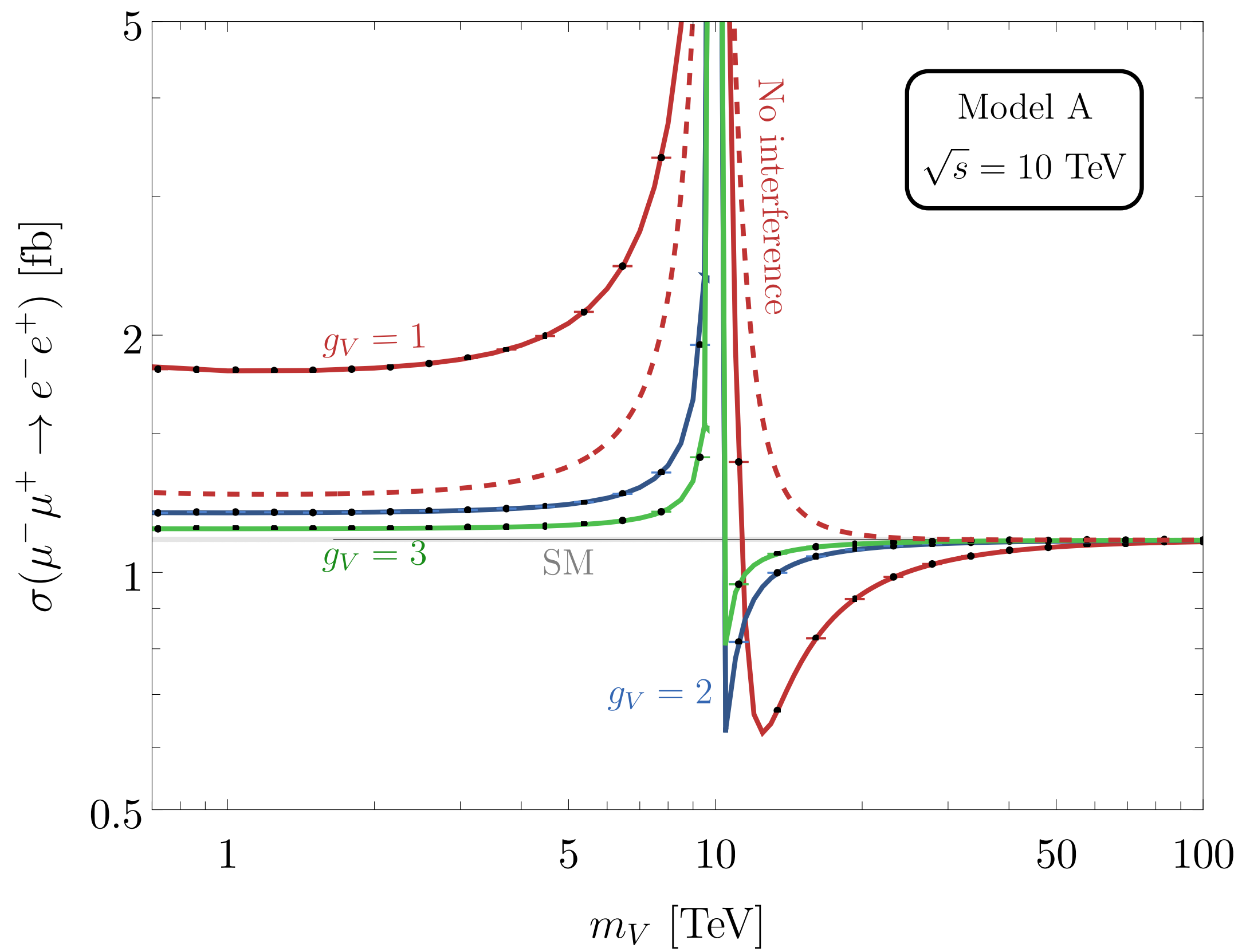
- We introduced a new heavy particle
- At low energies, can integrate away
- Effectively, W and B interactions are modified. Can express with form factors

$$\mathcal{L} = -\frac{1}{2}W_\mu^3\Pi_{33}(p^2)W^{\mu 3} - \frac{1}{2}B_\mu\Pi_{00}(p^2)B^\mu - W_\mu^3\Pi_{30}(p^2)B^\mu - W_\mu^+\Pi_\pm(p^2)W^{\mu-}$$

- Π_x are very well measured around $p^2 \sim 0$: need to reconstruct well known values

like $M_{W,Z}, \Gamma_{W,Z}, \alpha_{EM}$

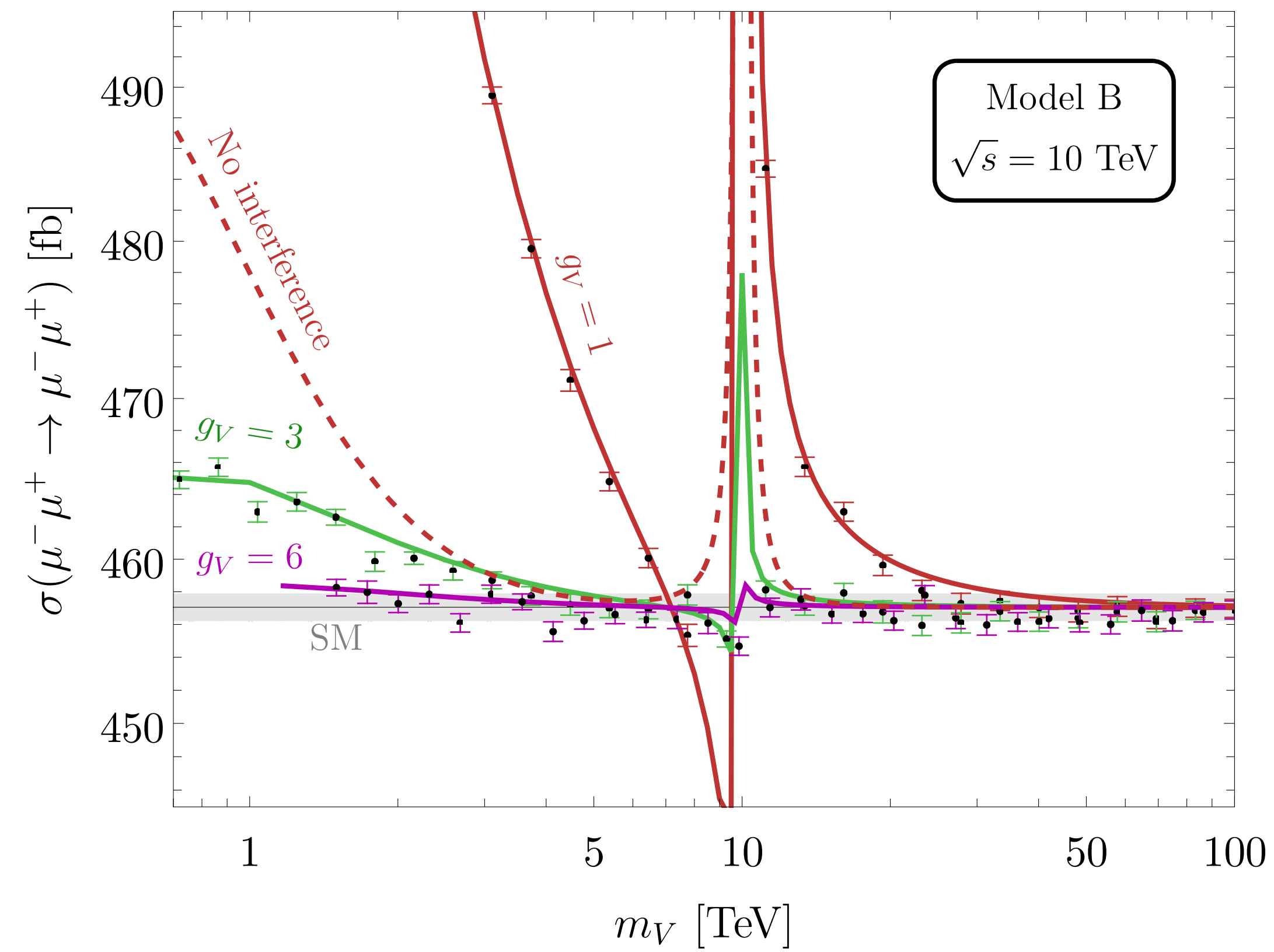
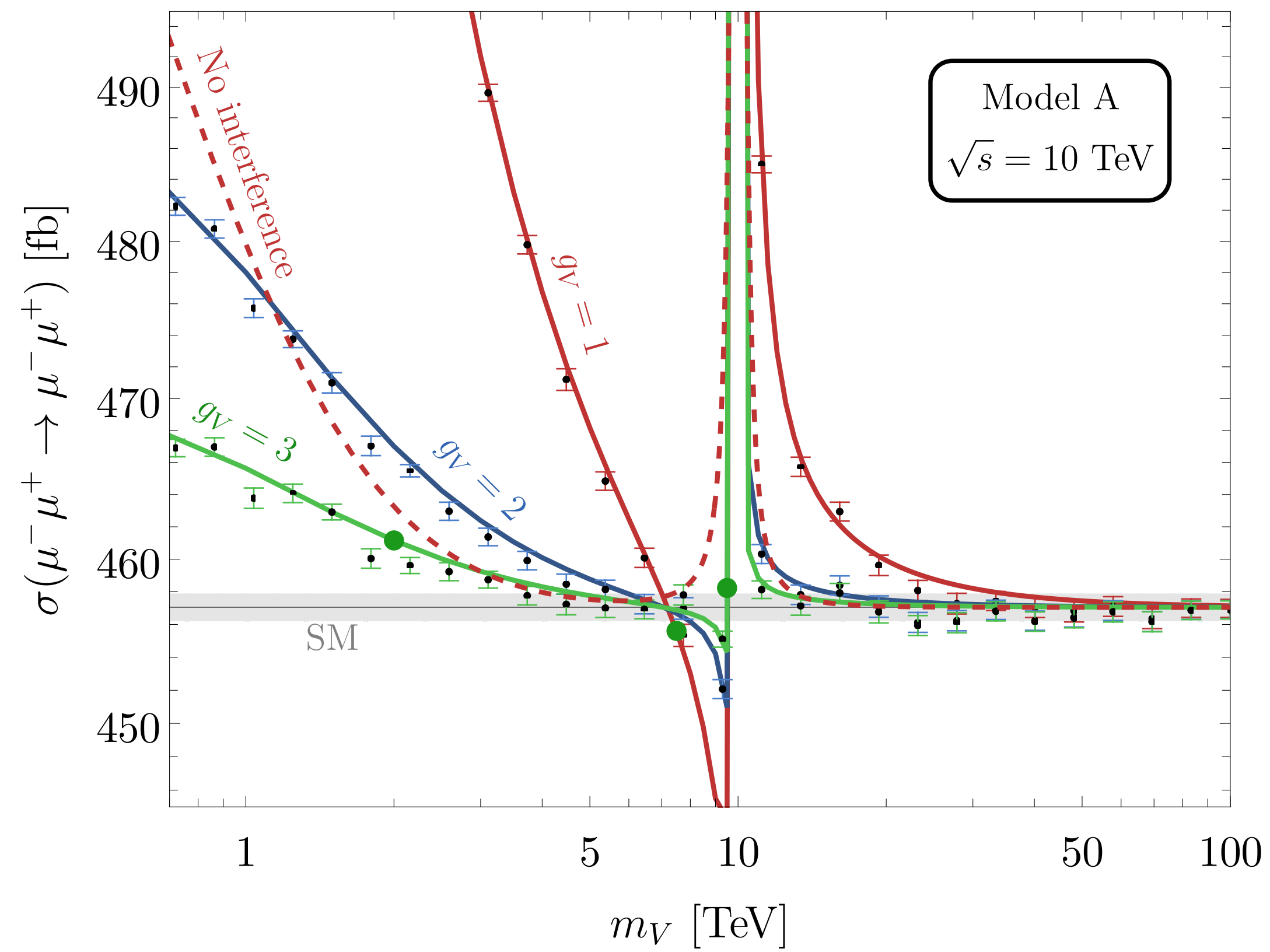
Cross sections



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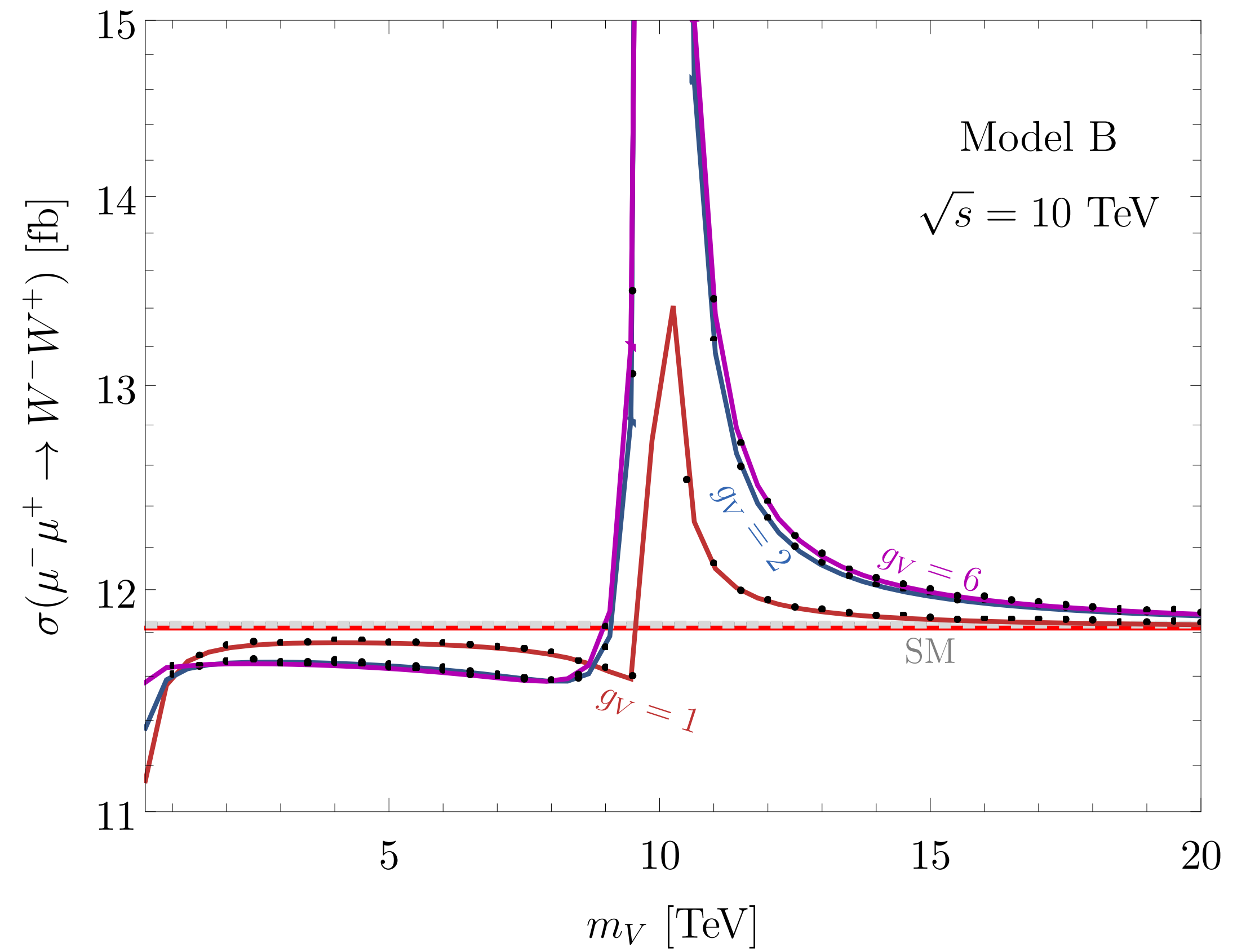
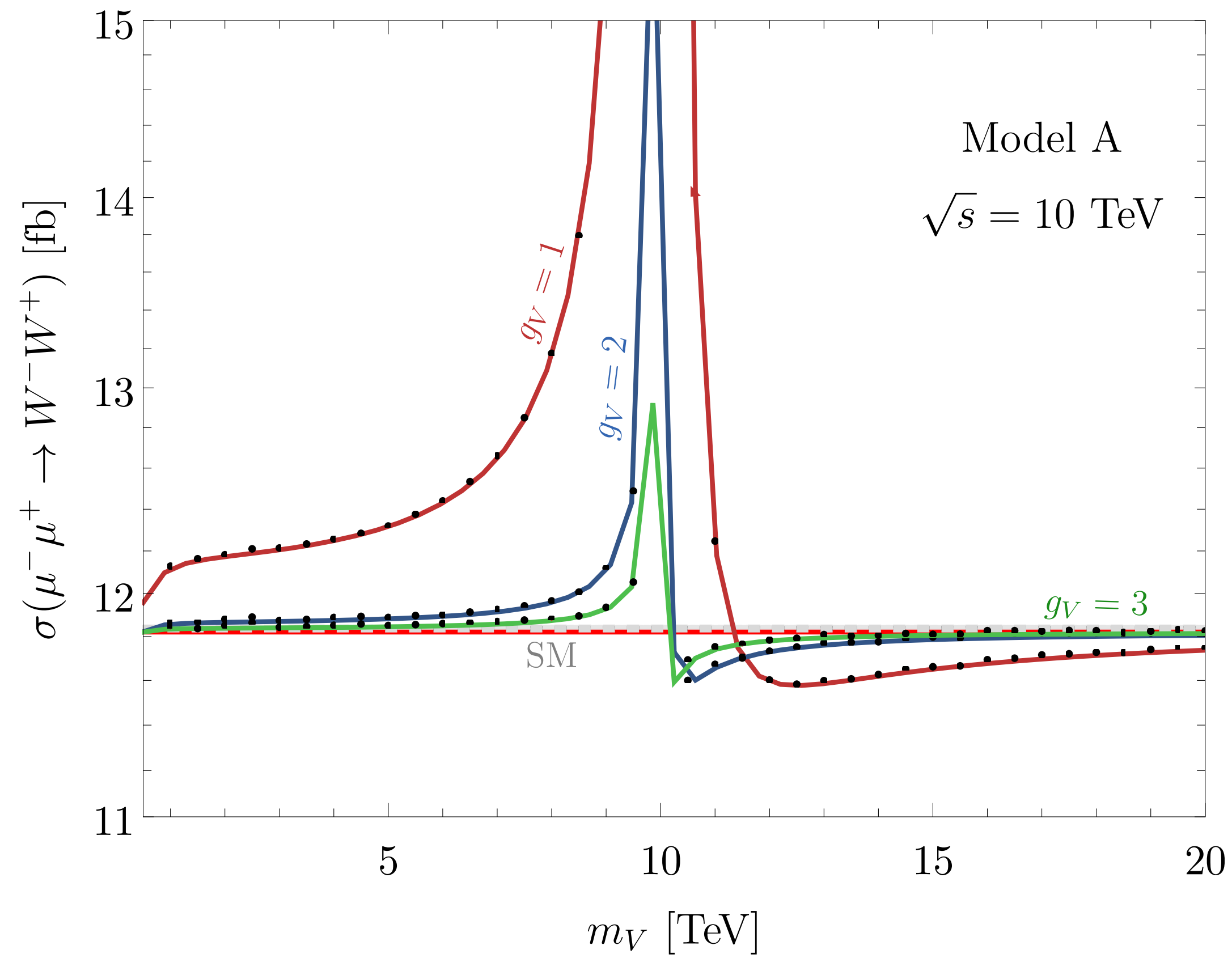
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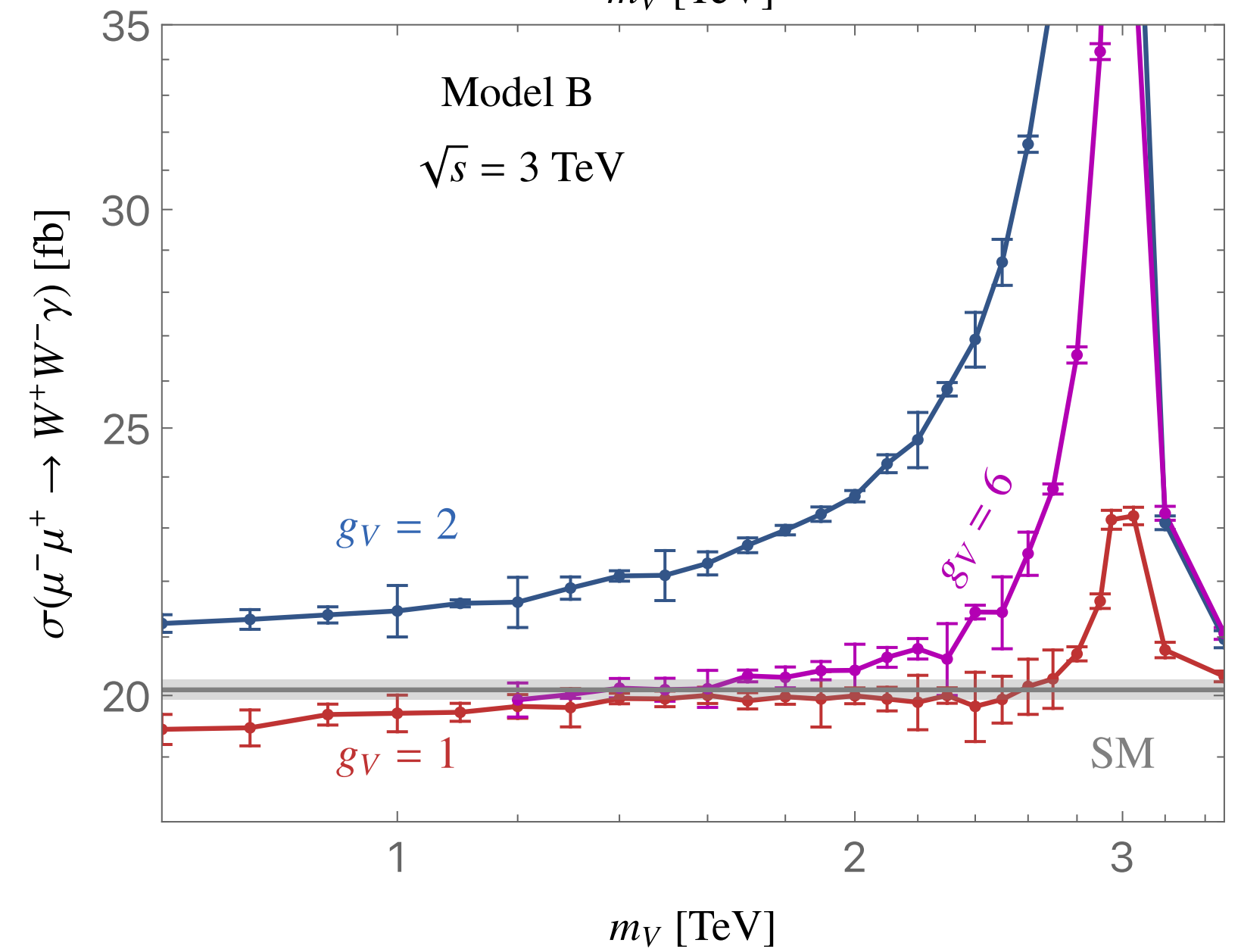
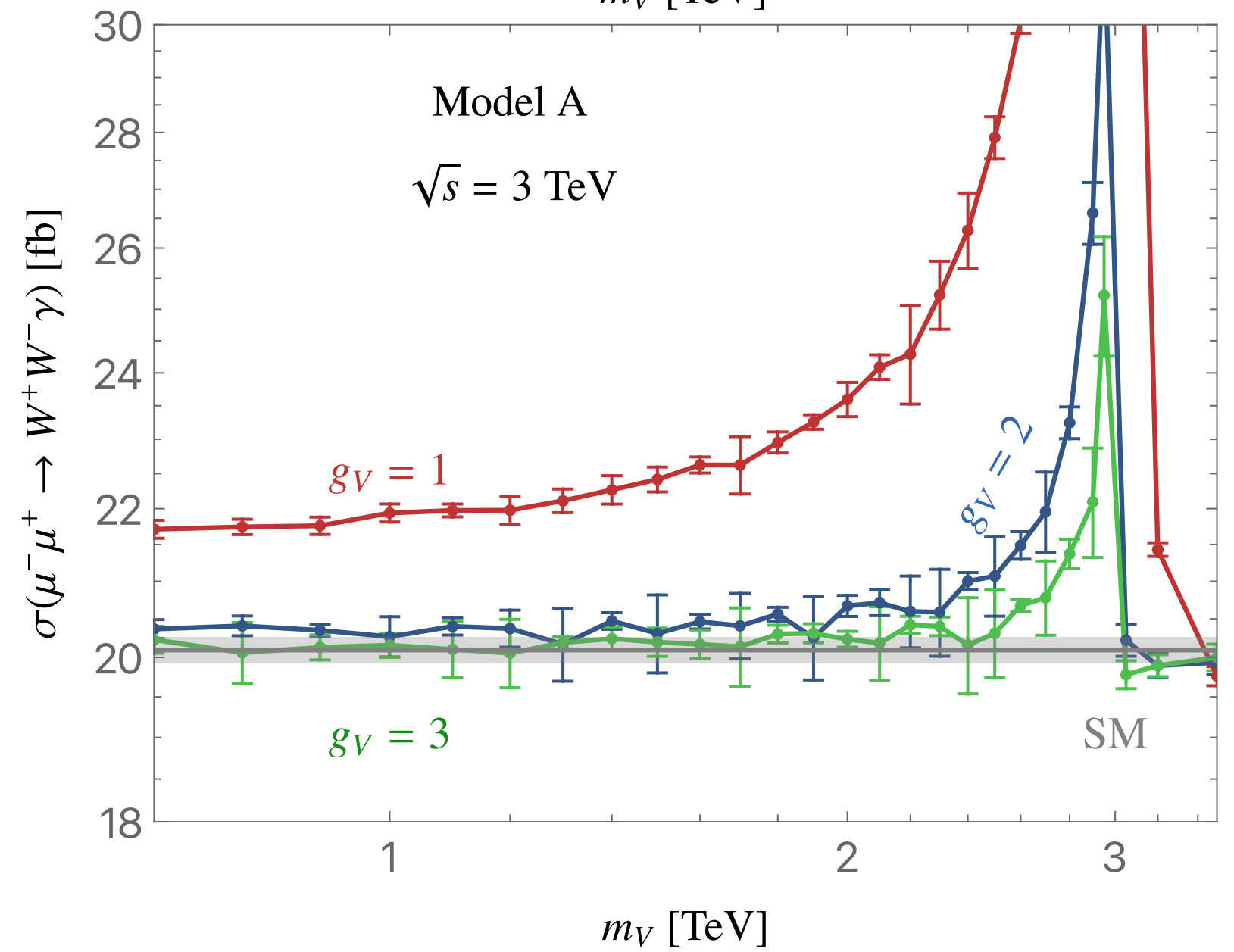
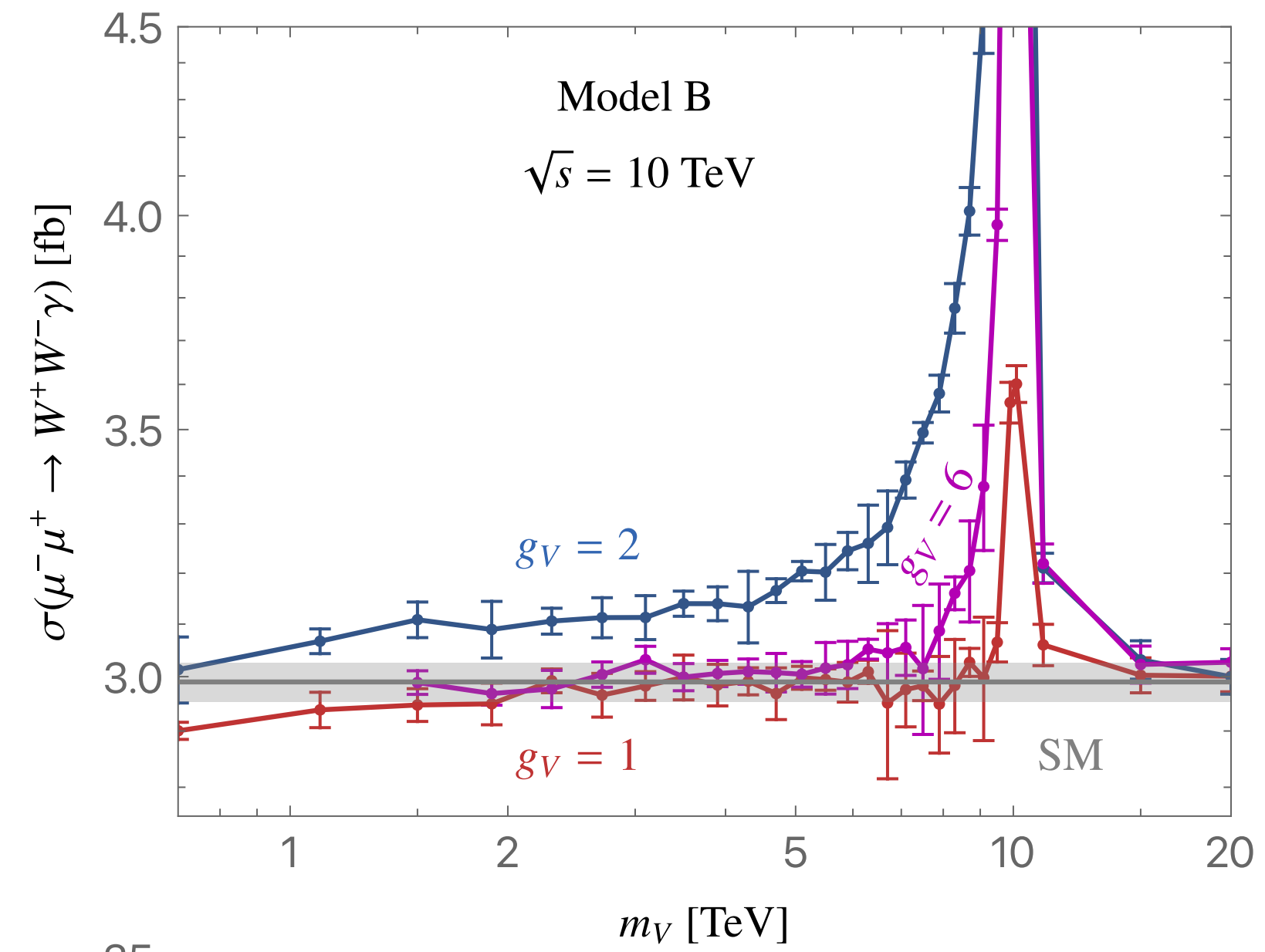
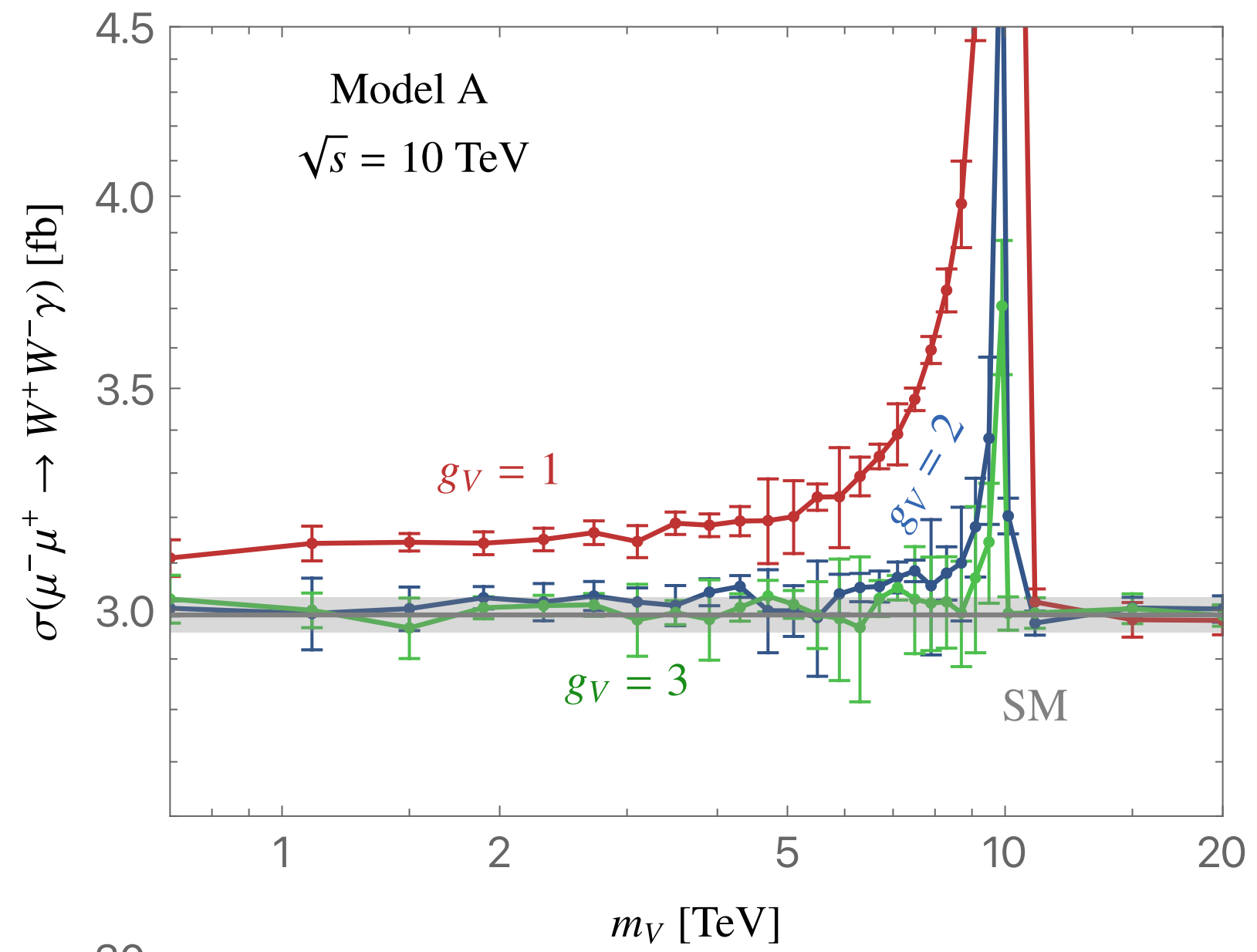
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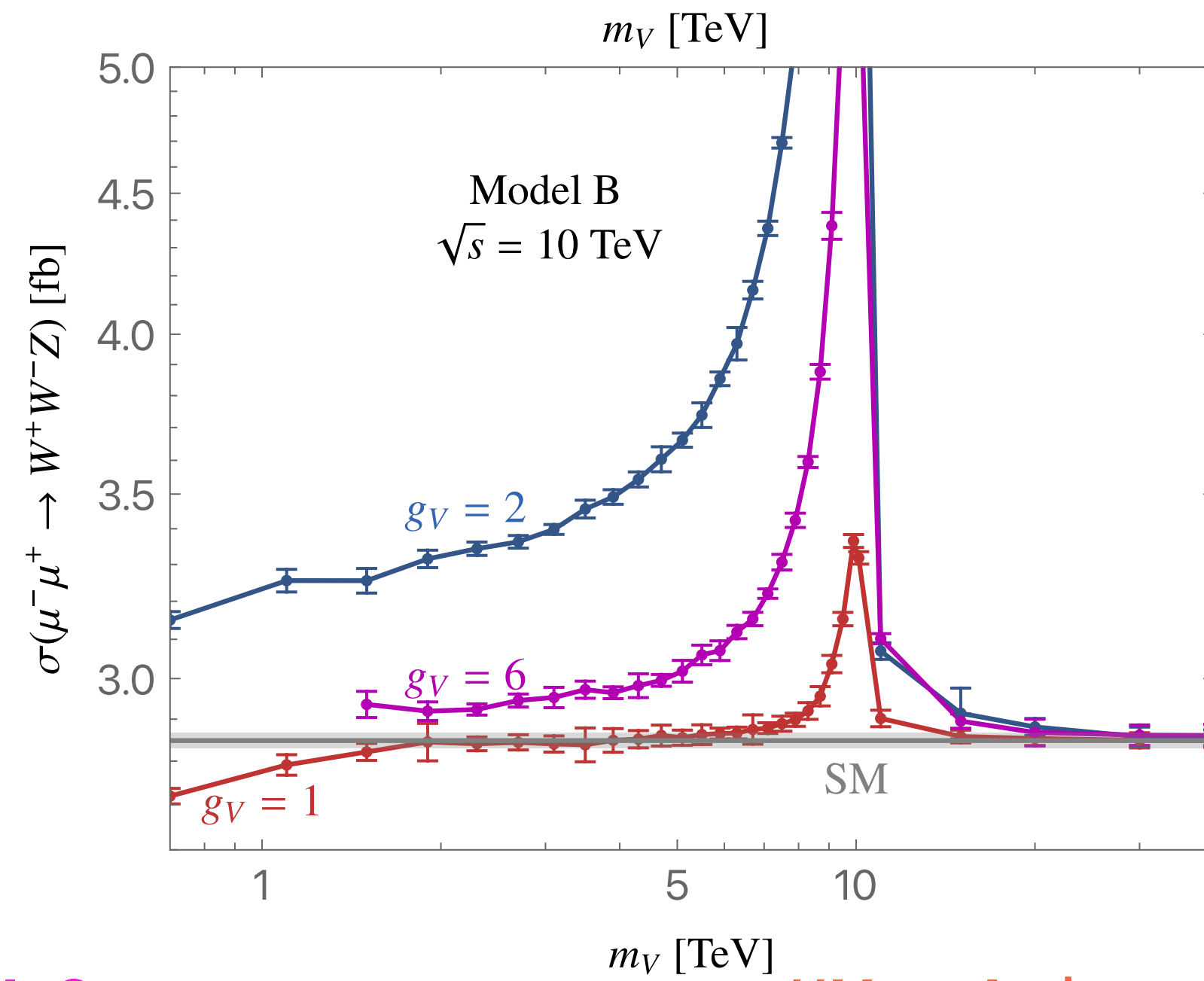
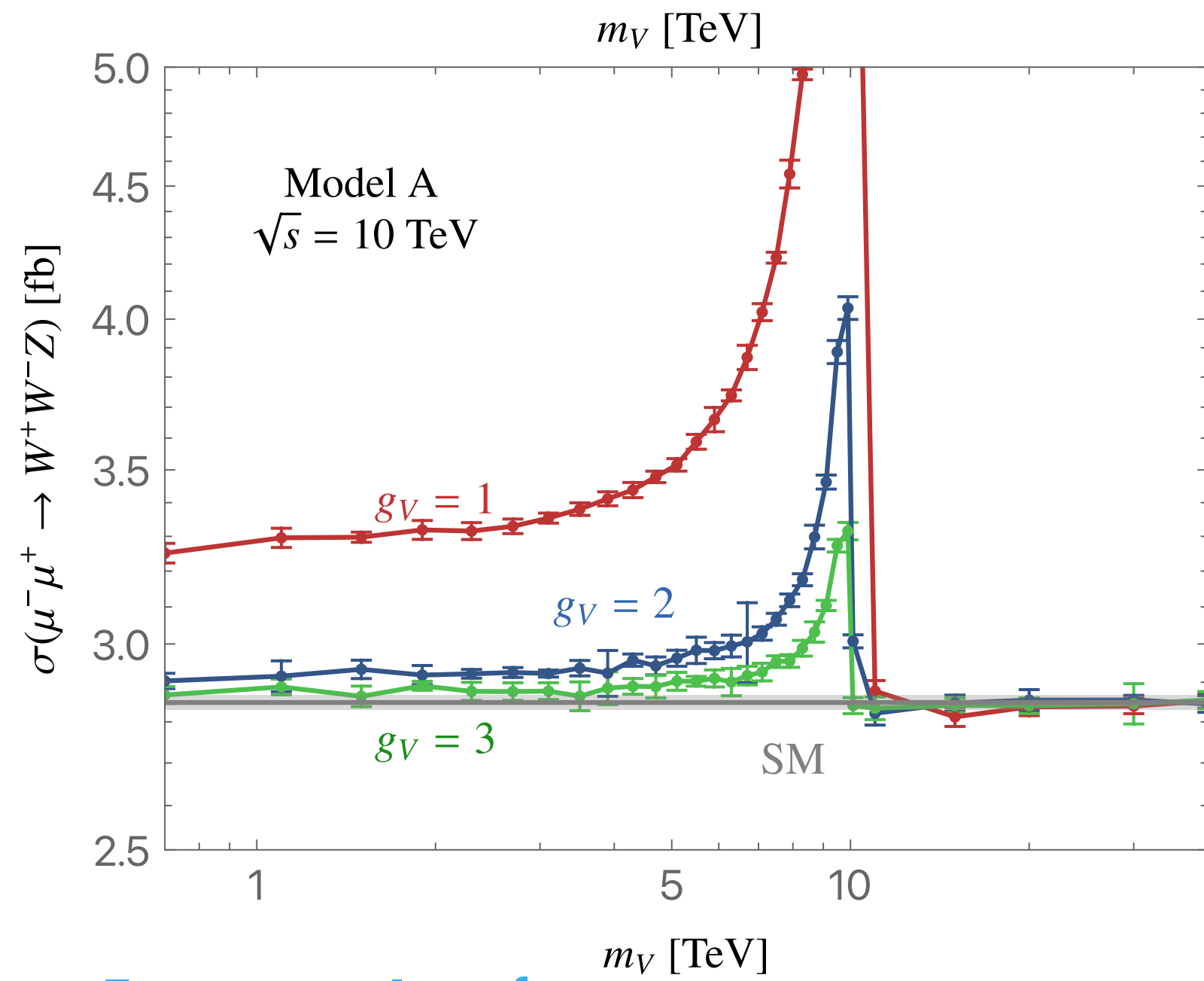
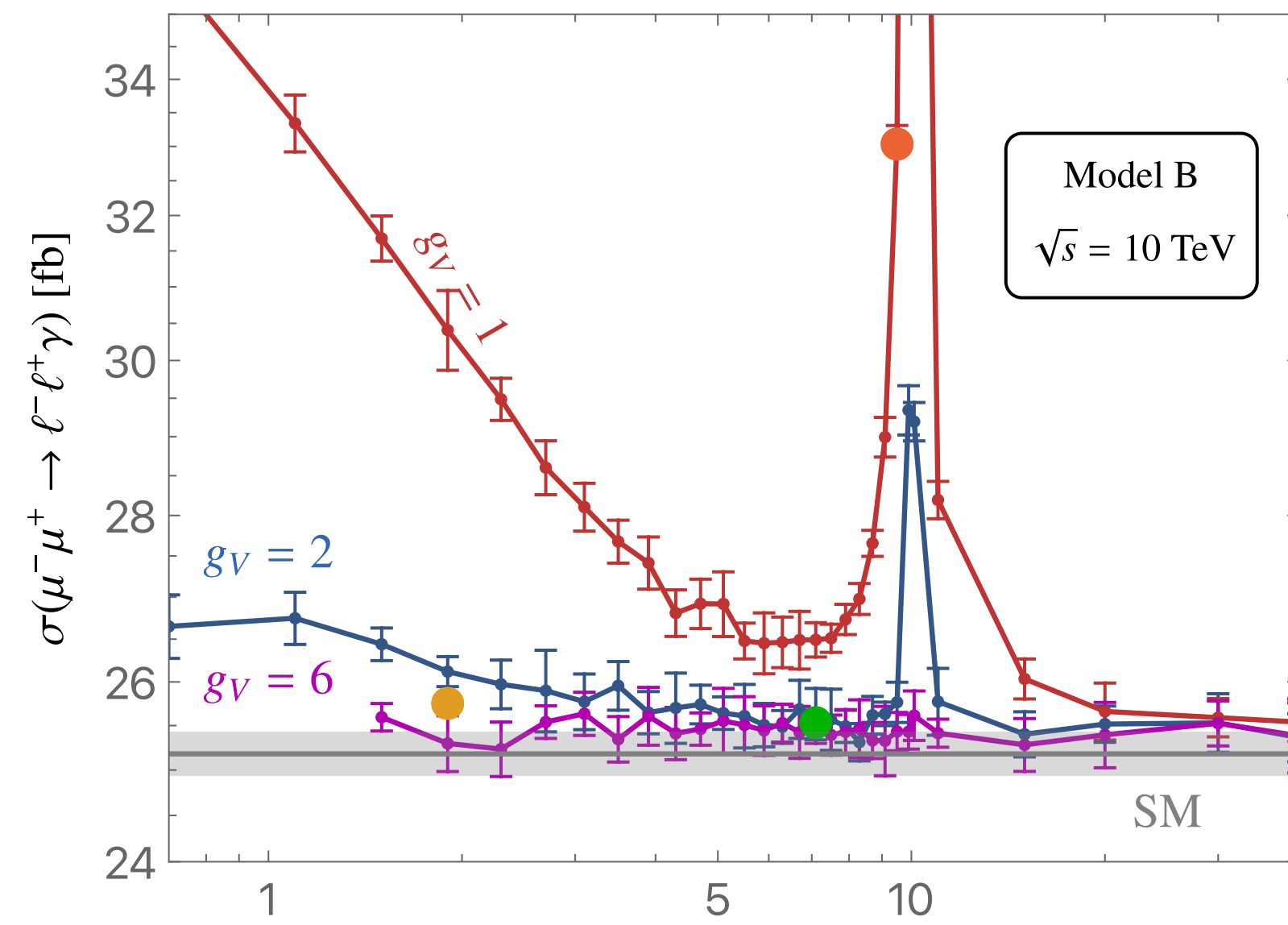
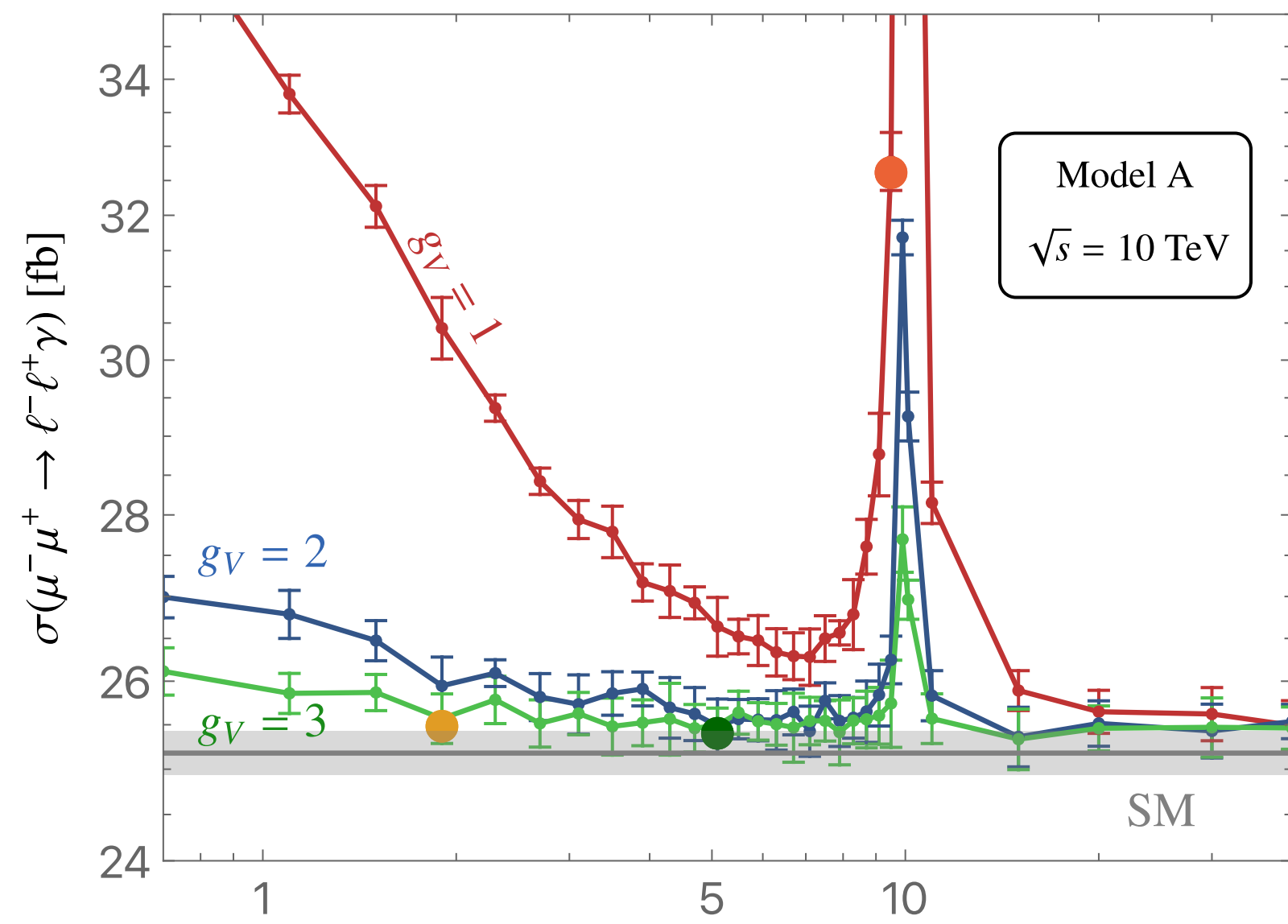
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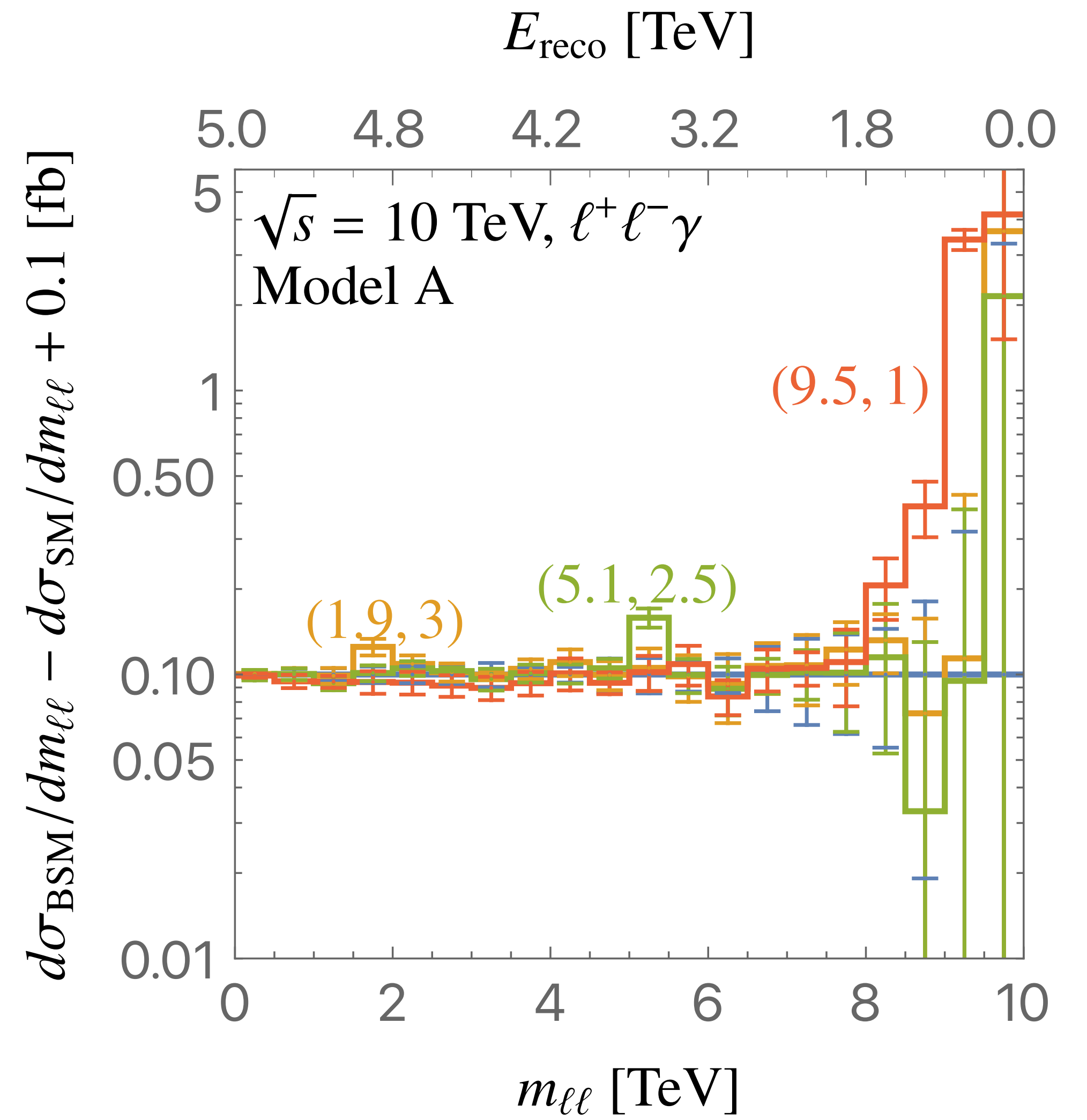
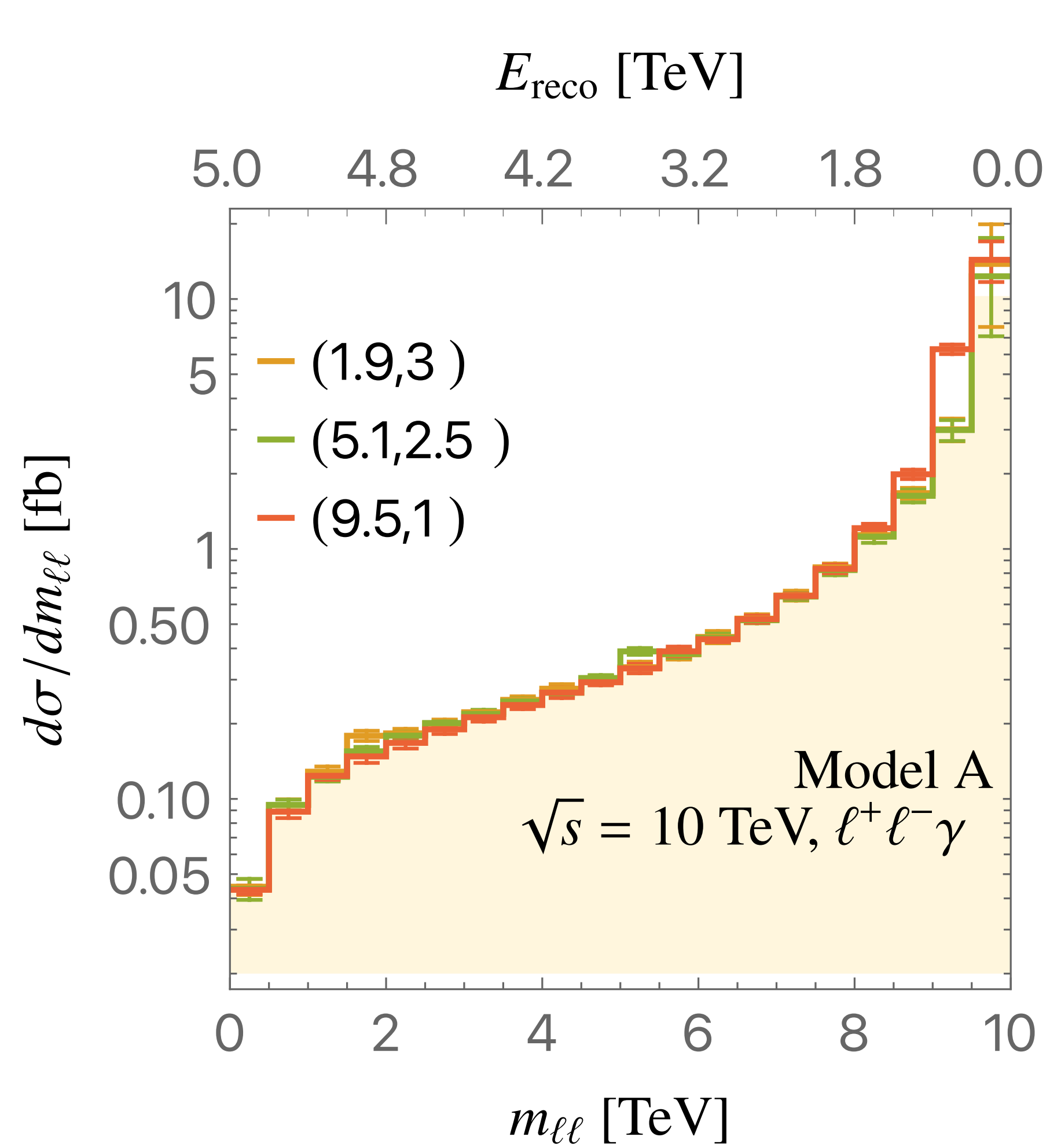


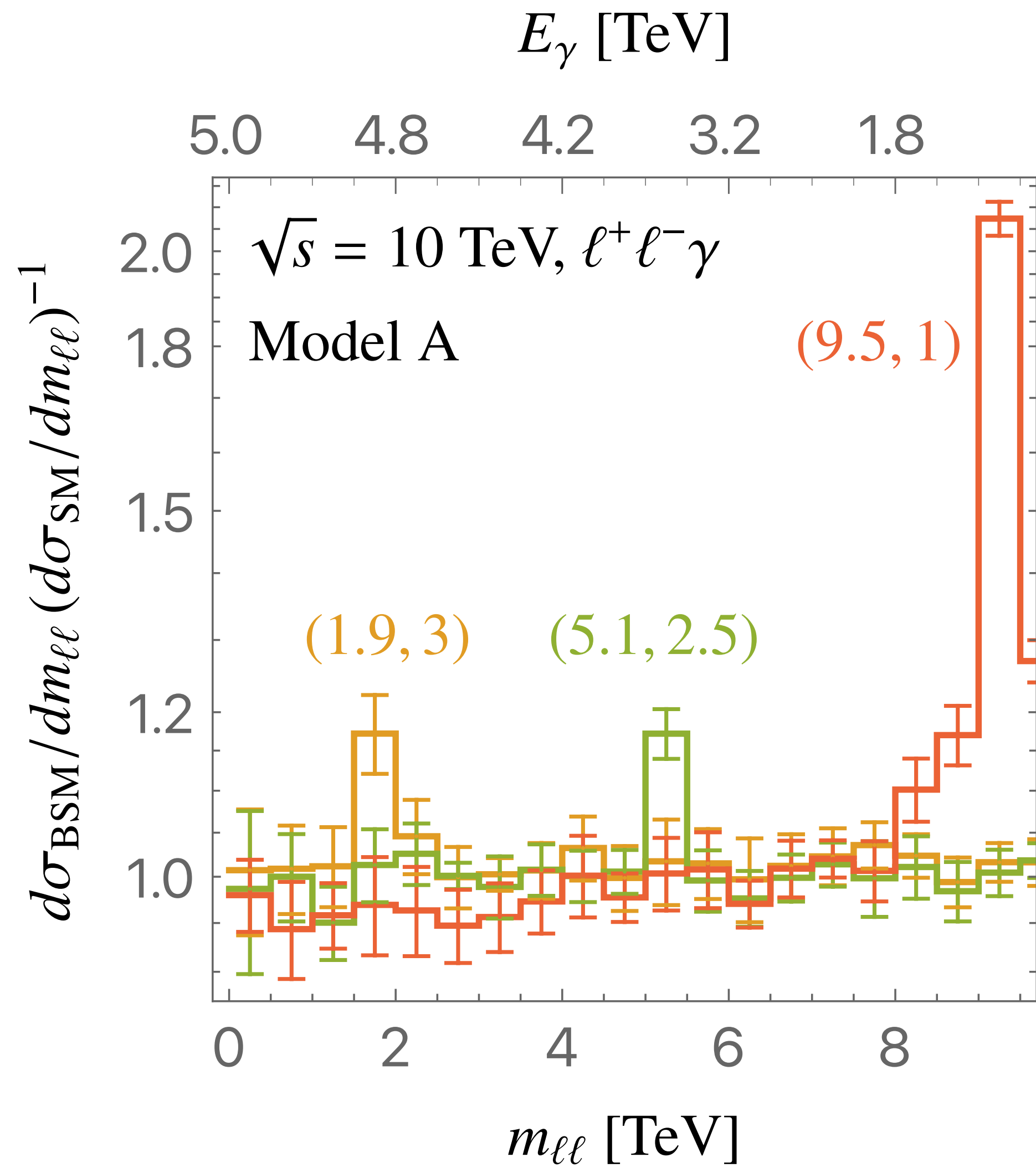
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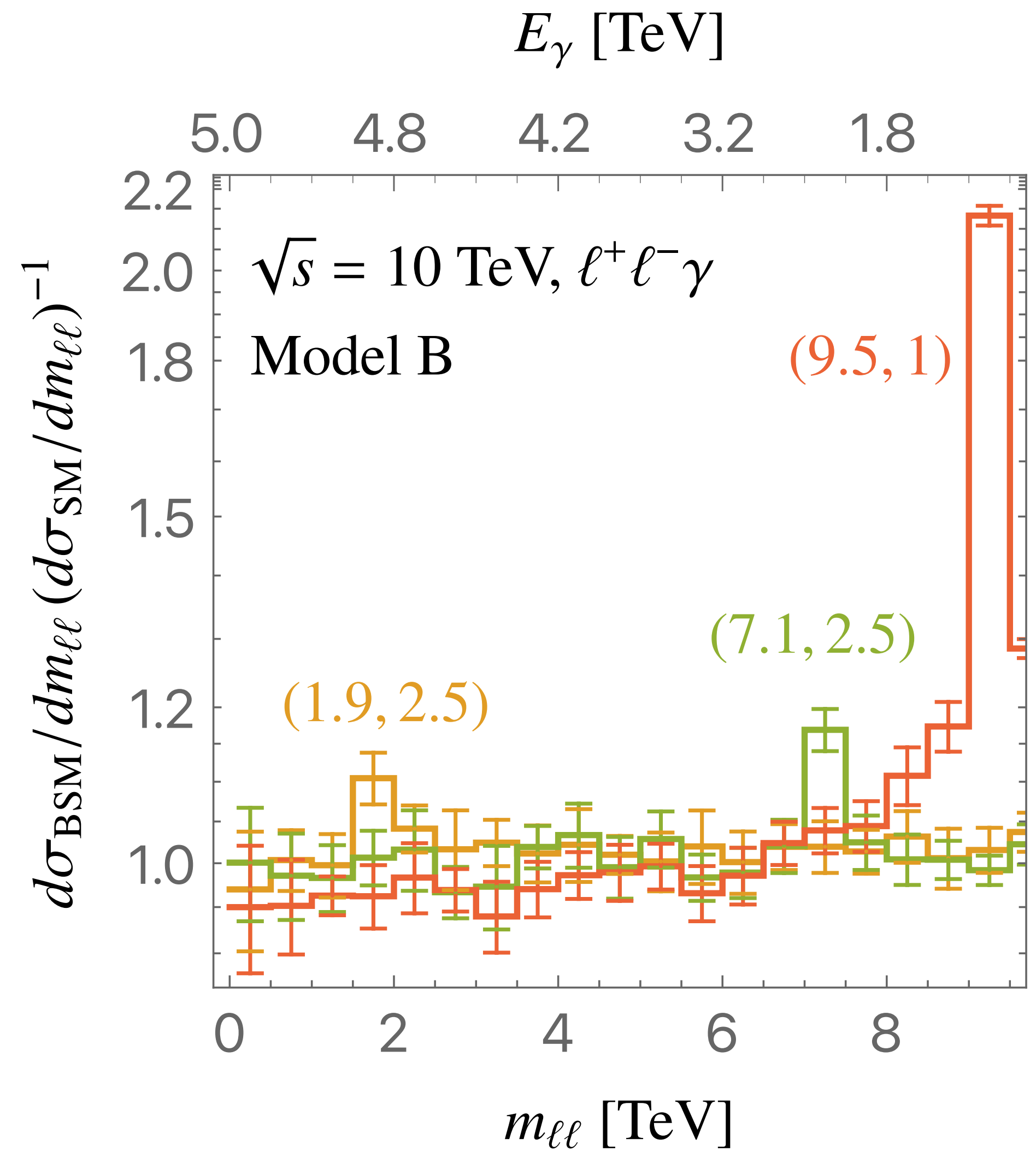
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Distributions



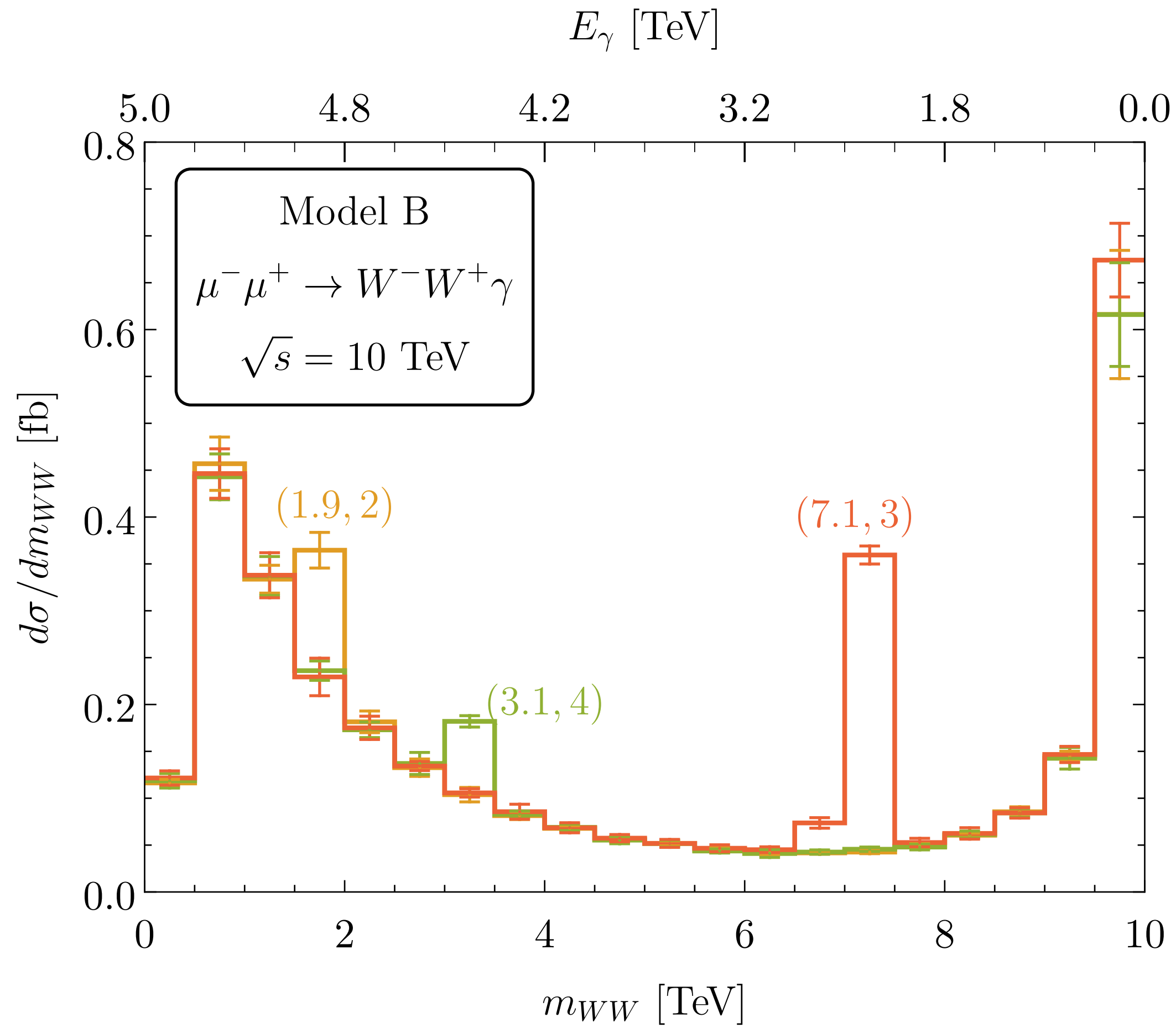


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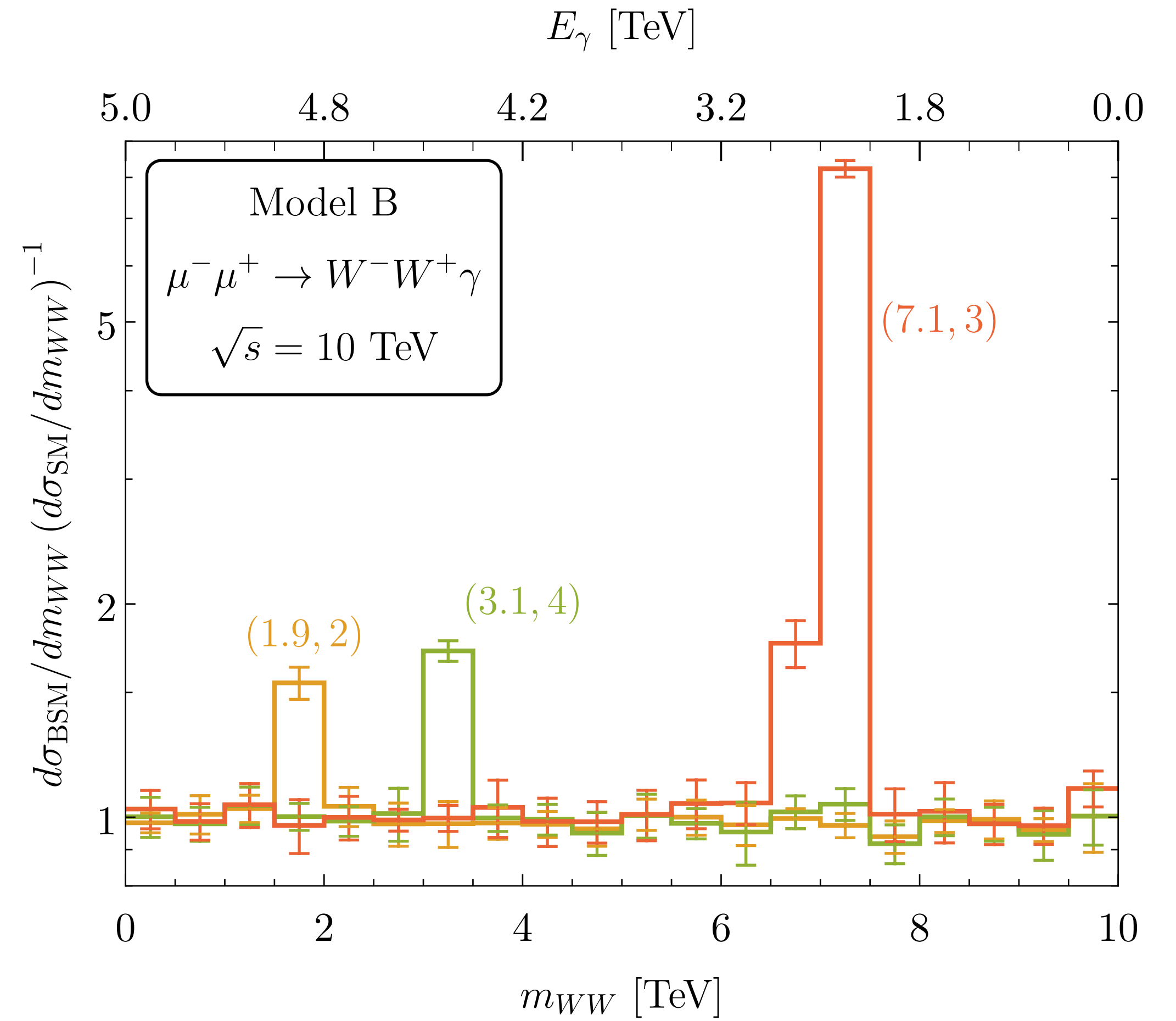


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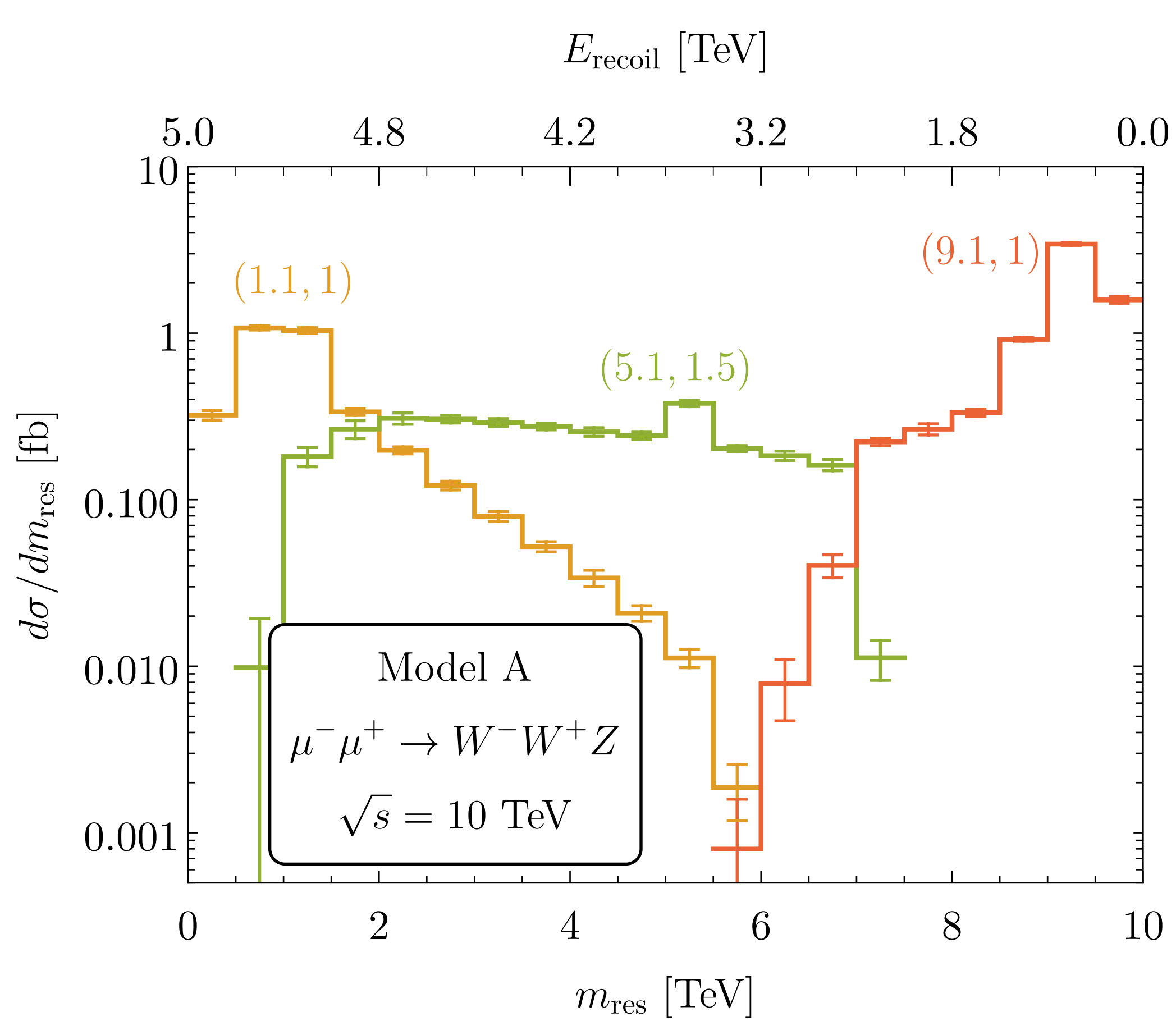


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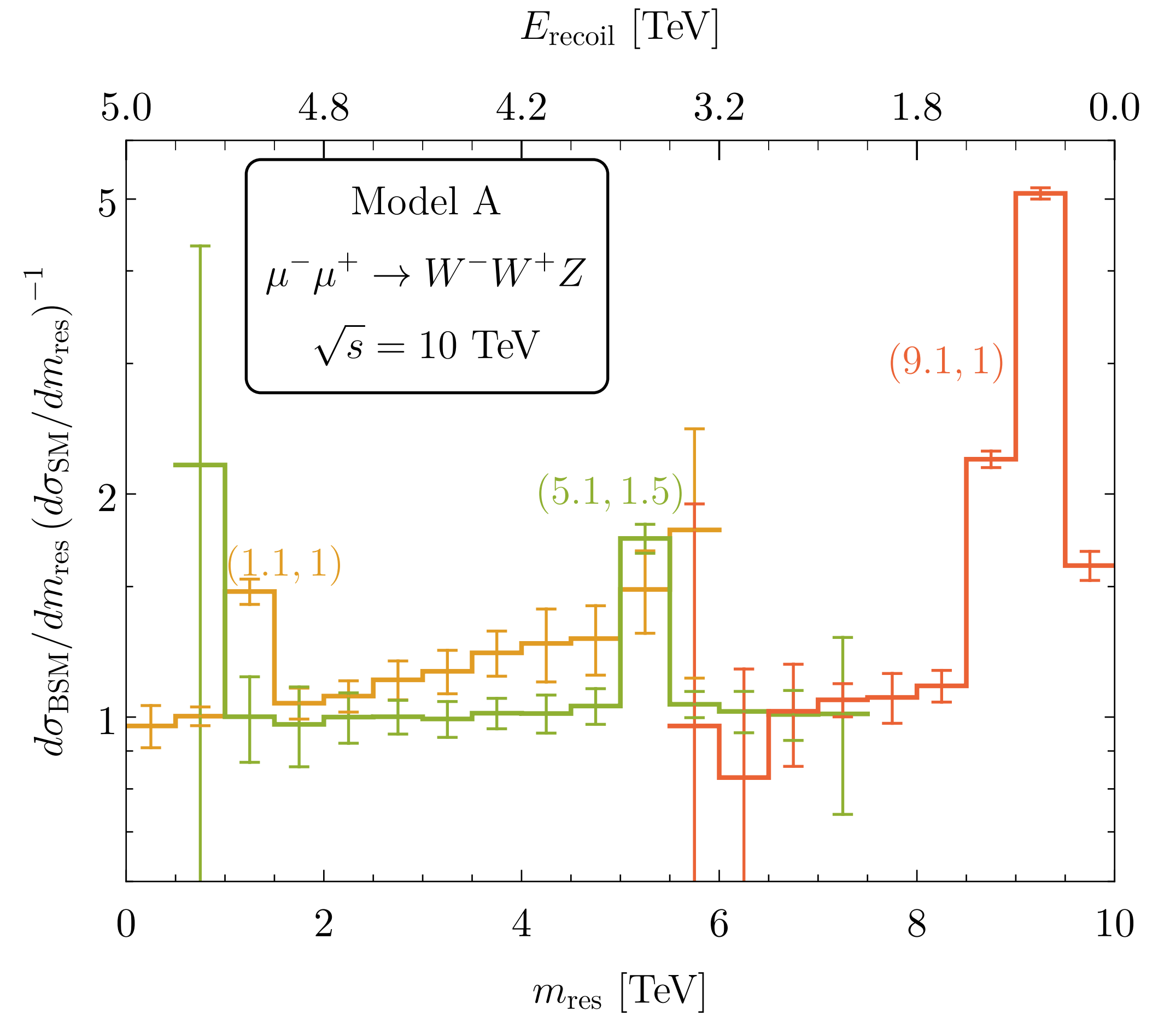


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