Kinetic Isocurvature Perturbations

w/K.J. Bae, J.-O. Gong, K. Harigaya, C.S. Shin

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Dhong Yeon (D.Y.) Cheong

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THE UNIVERSITY OF CHICAGO
THE UNIVERSITY OF CHICAGO LEINWEBER INSTITUTE FOR THEORETICAL PHYSICS
ENRICO FERMI INSTITUTE
KICP

Motivation

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Q : A new type of isocurvature perturbation?

Q : Any unique observational probe?

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- Kinetic Isocurvature Perturbations
- Adiabatic Number + Non-Adiabatic Momentum

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- Modulation of the Free Streaming Length
- Modulation of the Matter Power Spectrum

DM Energy Density Perturbations

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The energy density of DM

$p_{\rm DM} \equiv |\mathbf{p}_{\rm DM}|$

$$\rho_{\rm DM} = E_{\rm DM}\, n_{\rm DM}, \quad E_{\rm DM} \equiv \sqrt{m_{\rm DM}^2 + p_{\rm DM}^2}$$

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The energy density perturbation decomposes to

$$\frac{\delta\rho_{\rm DM}}{\rho_{\rm DM}} = \frac{\delta n_{\rm DM}}{n_{\rm DM}} + \frac{p_{\rm DM}^2}{E_{\rm DM}^2}\frac{\delta p_{\rm DM}}{p_{\rm DM}}$$

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Normally the dominant source

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Adiabatic $\lesssim 10^{-6}$

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For (warm) DM, as DM becomes non-relativistic

$$v_{\rm DM} = \frac{p_{\rm DM}}{E_{\rm DM}} \propto \frac{1}{a} \qquad \frac{\delta\rho_{\rm DM}}{\rho_{\rm DM}} \propto \frac{1}{a^2}$$

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Free-streaming length perturbation governed by initial p distribution

$$\delta_\lambda \equiv \frac{\delta\lambda_{\rm FS}}{\lambda_{\rm FS}} \propto \frac{\delta p_{\rm DM}}{p_{\rm DM}}$$

$$\lambda_{\rm FS} \equiv \int_{t_i}^{t_{\rm eq}} dt \frac{v(t)}{a(t)} = \int_{a_i}^{a_{\rm eq}} \frac{da}{a^2 H(a)} \frac{p_{\rm DM}(a)}{E_{\rm DM}(a)}$$

Kinetic Isocuvature Perturbation

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Setup to realize kinetic isocurvature perturbations?

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Setup

Key : adiabatic number perturbations + modulated kinetic distribution

Decaying massive scalar field : ϕ (subdominant component)

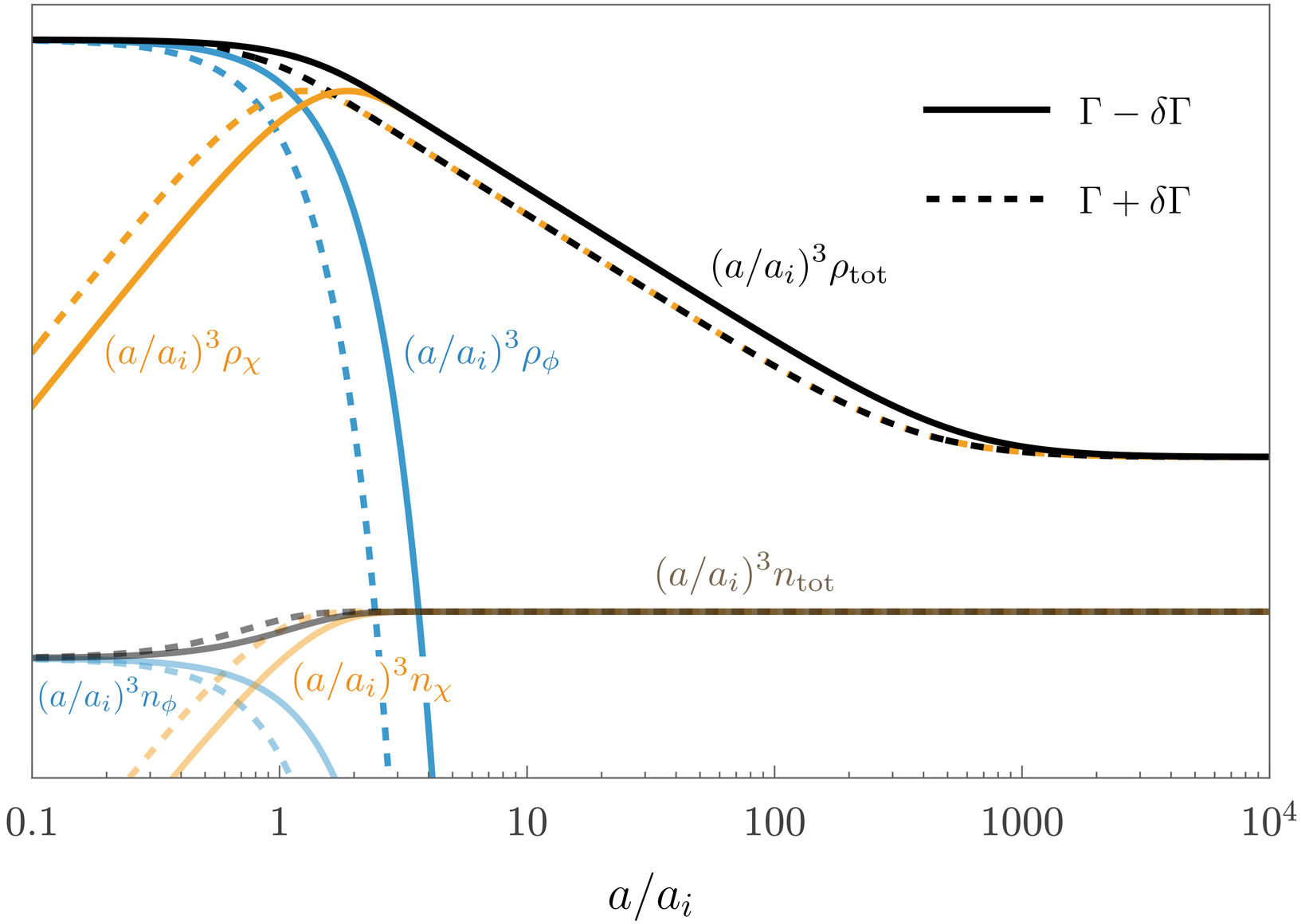
Light DM particle : χ / Modulating field σ (superhorizon fluctuations)

$$\mathcal{L} \supset \begin{cases} A\phi\chi^2 & \chi : \text{ scalar} \\ y\phi\bar{\chi}\chi & \chi : \text{ fermion.} \end{cases}$$

$$\Gamma = \begin{cases} \dfrac{A^2}{16\pi m_\phi}\left(1 - \dfrac{4m_\chi^2}{m_\phi^2}\right)^{1/2} \\ \dfrac{y^2 m_\phi}{8\pi}\left(1 - \dfrac{4m_\chi^2}{m_\phi^2}\right)^{3/2} \end{cases}$$

$$\boxed{\delta\sigma \to (\delta A, \delta y) \to \delta\Gamma}$$

Energy density modulation



Decay yield

$$n_\chi\left(a_i\right) = 2n_\phi\left(a_i\right).$$

Fixed, no δn

Different decay times t_i

Momentum redshift

$$p(a) = p_i\left(a_i/a\right)$$

Induced $\delta\rho_\chi$

Relation between $\delta\rho$ and δ_λ

Within our region of interest,

$u_i = p_i/m_\chi$

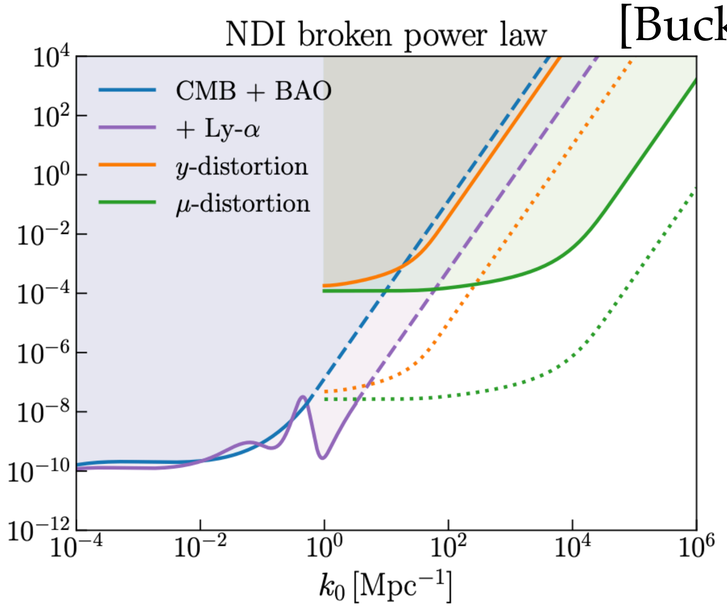
$$\lambda_{\rm FS} \simeq 1{\rm Mpc}\left(\frac{u_i^2 t_i}{10^6{\rm sec}}\right)^{1/2}\left[1 - 0.07\ln\left(\frac{u_i^2 t_i}{10^6{\rm sec}}\right)\right]$$

$$|\delta_\lambda| \simeq \left|\frac{\delta\rho_\chi}{\rho_\chi}\right|\left[1 + 1.56 g_*^{-1/2}\left(\frac{1{\rm Mpc}}{\lambda_{\rm FS}}\right)^2\left(1 + 0.27\ln\left(\frac{1{\rm Mpc}}{\lambda_{\rm FS}}\right)\right)\left(\frac{6.7{\rm Mpc}^{-1}}{k}\right)^2\right]$$

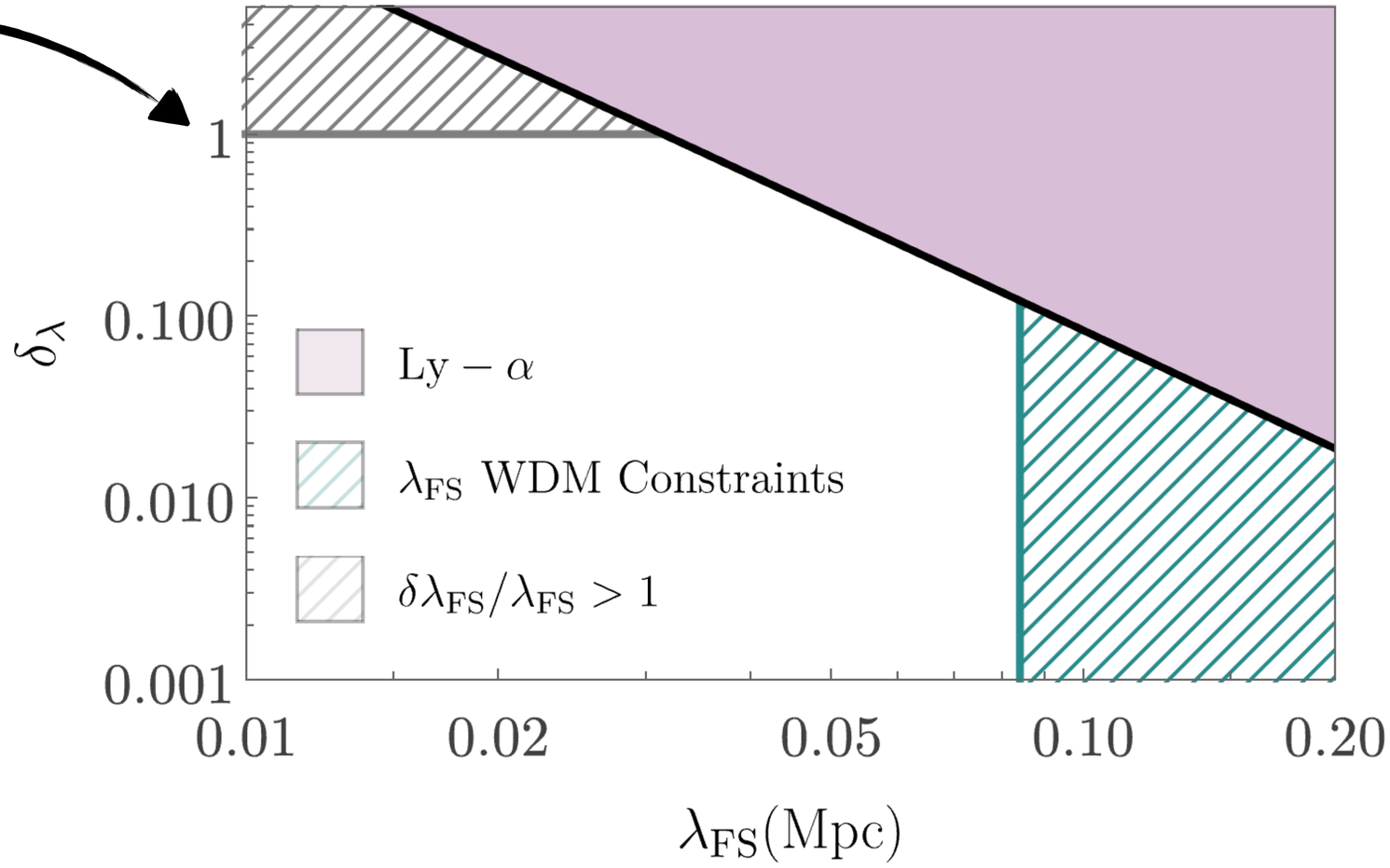
Any constraints on $\delta\rho_\chi/\rho_\chi \leftrightarrow \delta_\lambda$?

Kinetic Isocurvature Constraints

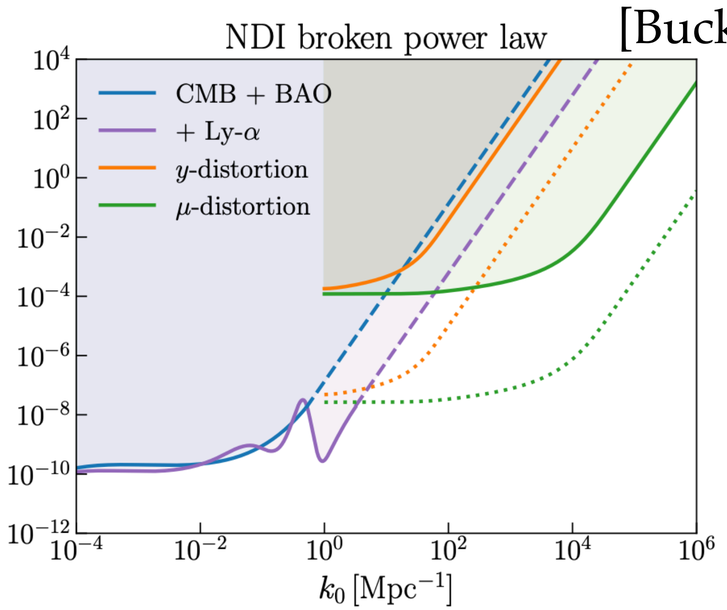
[Buckley, 2025]



$$P^{(2)}_{\delta_\lambda}(k) = A_{\delta_\lambda}$$



Kinetic Isocurvature Constraints



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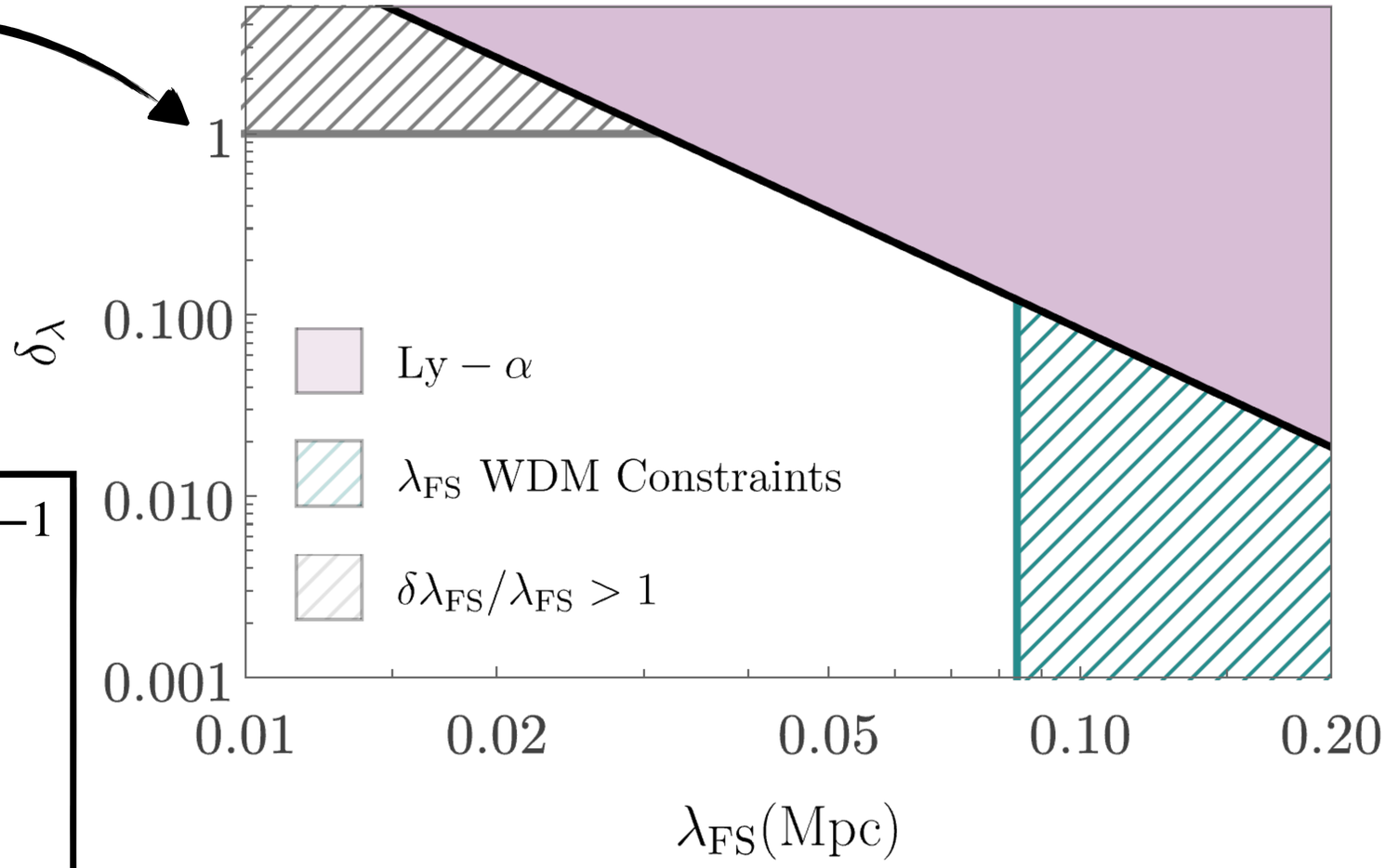
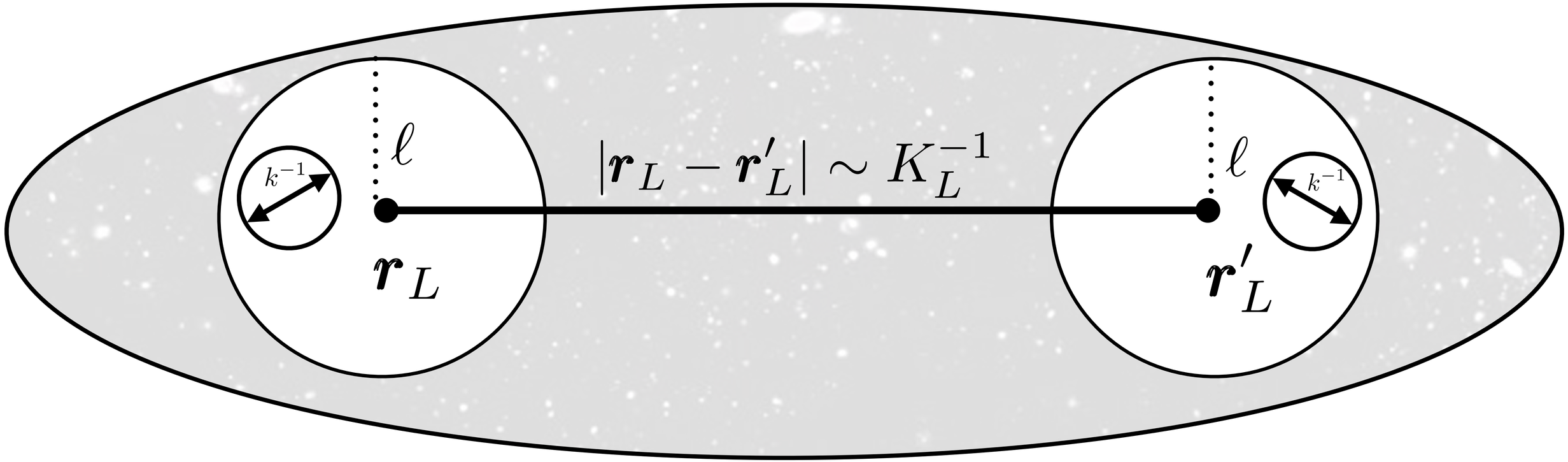


Ly−α bounds at $k \sim 1\,\mathrm{Mpc}^{-1}$

$\delta_\lambda = \mathcal{O}(0.1 - 1)$ allowed!

Any observational effect?

Position Dependent Power Spectrum



[Takada & Wu (2013), Chiang et.al. (2014)]

Position dep. power spectrum within sub volume V_s

$$P_\delta\left(\boldsymbol{k}, \boldsymbol{r}_L\right) \equiv \frac{1}{V_s}\left|\delta\left(\boldsymbol{k}, \boldsymbol{r}_L\right)\right|^2,$$

$$\delta\left(\boldsymbol{k}, \boldsymbol{r}_L\right) \equiv \int d^3r\delta(\boldsymbol{r})W_s\left(\boldsymbol{r}-\boldsymbol{r}_L\right)e^{-i\boldsymbol{k}\cdot\boldsymbol{r}} = \int \frac{d^3q}{(2\pi)^3}\delta^{(3)}(\boldsymbol{k}-\boldsymbol{q})W_s(\boldsymbol{q})e^{-i\boldsymbol{q}\cdot\boldsymbol{r}_L}$$

Position Dependent Power Spectrum

Modulation in $\lambda_{\rm FS}$

$$\lambda_{\rm FS}(\boldsymbol{x}) = \bar{\lambda}_{\rm FS} + \delta\lambda_{\rm FS}(\boldsymbol{x}).$$

Separate universe approach : $\lambda_{\rm FS}$ constant within a V_s patch.

$$\Delta P_m\left(\boldsymbol{k}, \lambda_{\rm FS}\left(\boldsymbol{r}_L\right)\right) \simeq \lambda_{\rm FS} \left.\frac{dP_m}{d\lambda_{\rm FS}}\right|_{\lambda_{\rm FS}=\bar{\lambda}_{\rm FS}} \delta_\lambda\left(\boldsymbol{r}_L\right)$$

$$\Delta P_m\left(\boldsymbol{k}, \lambda_{\rm FS}\left(\boldsymbol{r}_L\right)\right) \equiv P_m\left(\boldsymbol{k}, \lambda_{\rm FS}\left(\boldsymbol{r}_L\right)\right) - P_m\left(\boldsymbol{k}, \bar{\lambda}_{\rm FS}\right)$$

Position Dependent Power Spectrum

Correlation of ΔP_m governed by P_{δ_λ}

$$P_{PP}\left(\boldsymbol{k}, \boldsymbol{K}_L\right) = \lambda_{\mathrm{FS}}^2 \left(\frac{dP_m}{d\lambda_{\mathrm{FS}}}\right)^2 P_{\delta_\lambda}\left(\boldsymbol{K}_L\right)$$

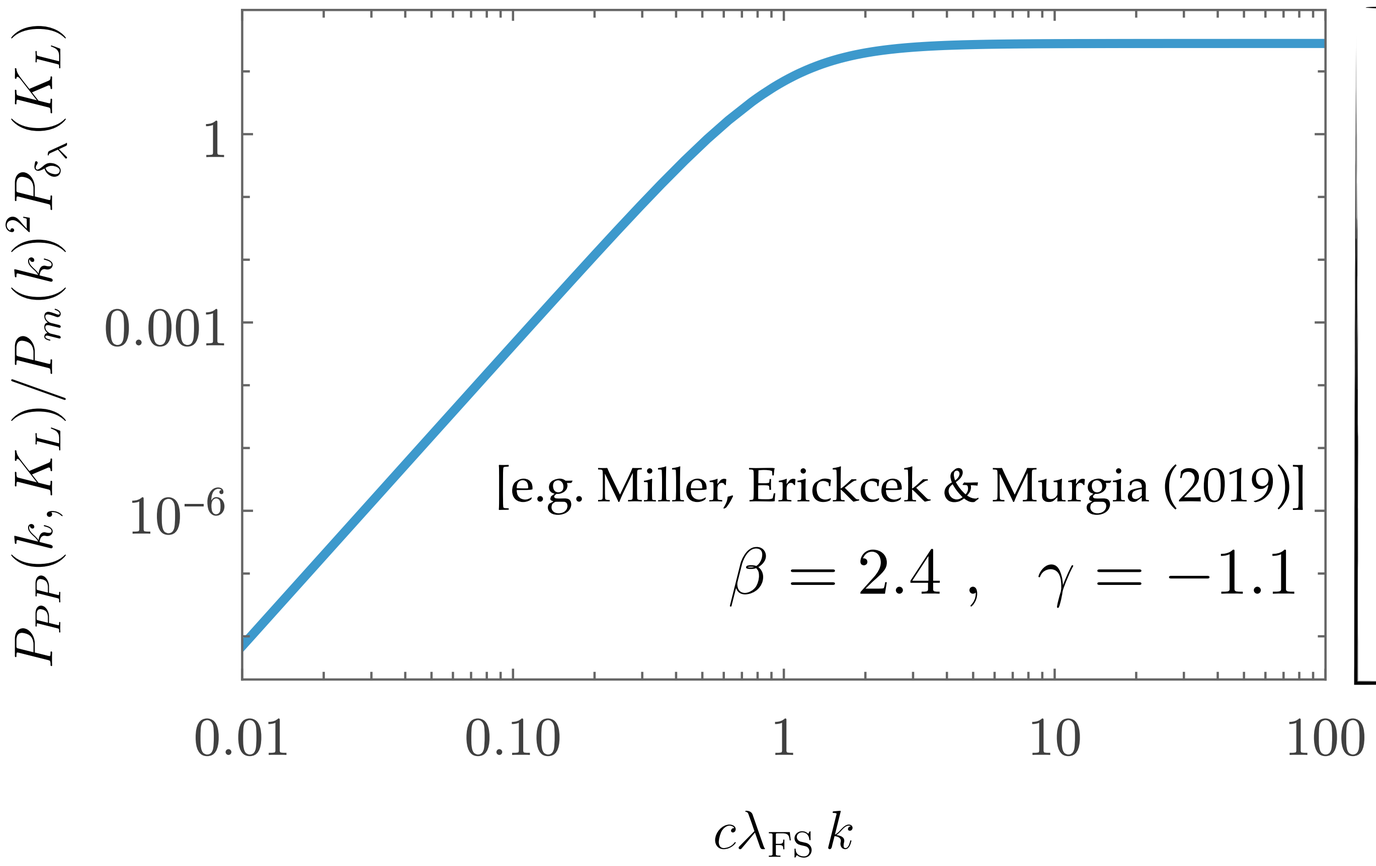
$$\langle \Delta P_m(\boldsymbol{k}, \lambda_{\mathrm{FS}}; \boldsymbol{K}_L) \Delta P_m(\boldsymbol{k}, \lambda_{\mathrm{FS}}; \boldsymbol{K}'_L) \rangle \equiv (2\pi)^3 \delta^{(3)}(\boldsymbol{K}_L + \boldsymbol{K}'_L) P_{PP}(\boldsymbol{k}, \boldsymbol{K}_L)$$

Size of effect governed by P_m dependence on λ_{FS}, through the T_{FS} !

Modulating Power Spectrum

$$P_m(k) = T_{\rm FS}(k)^2 P_{\mathcal{R}}(k)\,|W_s(k)|^2$$

Representative form $T_{\rm FS}\left(k, \lambda_{\rm FS}\right) = \left[1 + \left(c\lambda_{\rm FS} k\right)^{\beta}\right]^{\gamma}$



$$\frac{P_{PP}\left(k_{\rm FS}, K_L\right)}{P_m\left(k_{\rm FS}\right)^2} \simeq 10 P_{\delta_\lambda}\left(K_L\right)$$

$\mathcal{O}(1)$ fluctuations in normalized correlation

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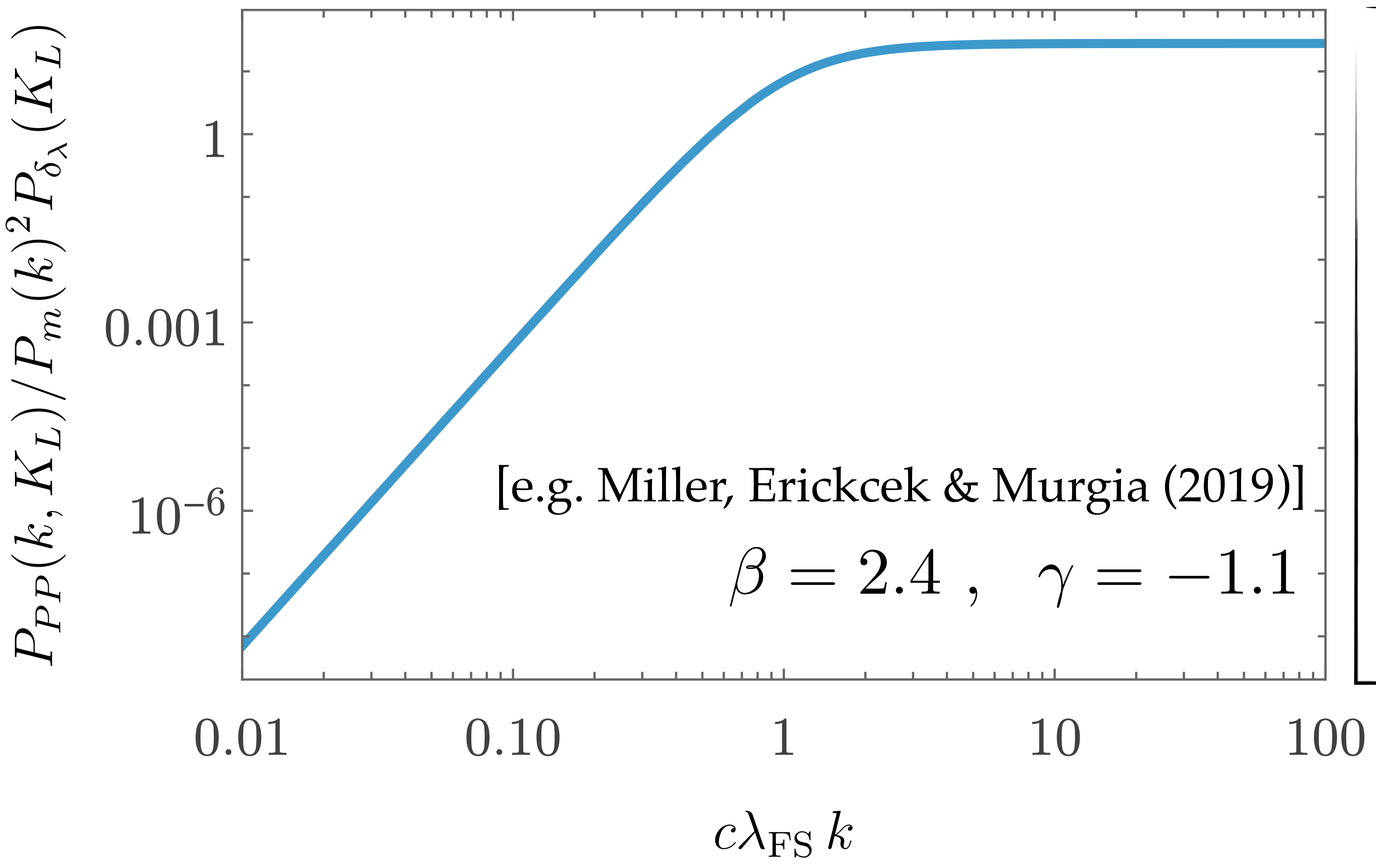

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$\mathcal{O}(1)$ fluctuations in normalized correlation

Modulations in galaxy surveys, small scale observations?

Conclusions and Discussions

- Kinetic Isocurvature Perturbations as a new class of isocurvature!
- Governed by p fluctuation, $O(1)$ fluctuations compatible with isocurvature constraints.
- $O(1)$ modulations in the patch-by-patch free streaming length.

Open questions :

UV complete model, optimized observables?

Backup Material

A Concrete Model

Massive scalar field : ϕ (zero mode), scalar DM χ, modulating scalar σ.

$$\mathcal{L}_{\rm int} = g\sigma\phi\chi^2 \ , \quad \Gamma_\phi \simeq \frac{1}{8\pi}\frac{g^2\sigma_i^2}{m_\phi}$$

Temperature scales : $T_{\rm dec}$, $T_{\rm osc}$, $T_{\rm NR}$, assuming RD at $T_{\rm dec}$ and $T_{\rm osc}$.

$$T_{\rm dec} \simeq \sqrt{\Gamma M_P} \ , \quad T_{\rm osc} \simeq \left(\frac{90}{\pi^2 g_*}\right)^{1/4}\sqrt{m_\phi M_P} \ , \quad T_{\rm NR} \simeq T_{\rm dec}\frac{m_\chi}{m_\phi}$$

$$\frac{\rho_\chi}{s} \simeq \frac{m_\chi n_\chi}{s} \simeq m_\chi \frac{n_\phi}{s} \sim \frac{m_\chi m_\phi \phi_i^2}{g_* \left(m_\phi M_P\right)^{3/2}}$$

$$\phi_i \sim \left(\frac{g_* T_{\rm dec} T_{\rm eq}}{T_{\rm NR}}\right)^{1/2} \frac{M_P^{3/4}}{m_\phi^{1/4}} \sim 10^{15}{\rm GeV} \times \left(\frac{T_{\rm dec}}{10^9{\rm GeV}}\right)^{1/2} \left(\frac{10{\rm keV}}{T_{\rm NR}}\right)^{1/2} \left(\frac{10^8{\rm GeV}}{m_\phi}\right)^{1/4}$$

A Concrete Model

Requirements on m_σ

In order for large Γ modulation, σ should start oscillating at $T < T_{\rm dec}$

$$m_\sigma < \frac{T_{\rm dec}^2}{M_P} = 1{\rm GeV}\left(\frac{T_{\rm dec}}{10^9{\rm GeV}}\right)^2$$

In order for $\delta m_\sigma^{\text{1-loop}} \sim \frac{g}{4\pi}\phi < \frac{T_{\rm dec}^2}{M_P}$, $g \lesssim 10^{-8}\left(\frac{T_{\rm NR}}{10{\rm keV}}\right)^{1/2}\left(\frac{m_\phi}{10^8{\rm GeV}}\right)$

A Concrete Model

σ eventually starts oscillations and dissipated at $T_{\rm dis}$, needs suppressed

$$\mathcal{R}_\sigma \sim \frac{\delta\sigma}{\sigma_i}\frac{\sigma_i^2}{M_P^2}\frac{\sqrt{m_\sigma M_P}}{T_{\rm dis}}$$

$$\sim 10^{-6} \times \frac{\delta\sigma}{\sigma_i}\frac{m_\phi}{10^8{\rm GeV}}\left(\frac{T_{\rm dec}}{10^9{\rm GeV}}\right)^2\left(\frac{m_\sigma}{1{\rm GeV}}\right)^{1/2}\frac{10^3{\rm GeV}}{T_{\rm dis}}\left(\frac{10^{-8}}{g}\right)^2$$

$\mathcal{R}_\sigma$ suppressed for sufficiently large $T_{\rm dis}$

e.g. If $\mathcal{L} \supset A_H\sigma|H|^2$, σ is dissipated with a rate $\sim \dfrac{A_H^2}{8\pi T}$

$$T_{\rm dis} \sim 10^3{\rm GeV}\left(\frac{A_H}{10^{-3}{\rm GeV}}\right)^{2/3}$$

$m_\sigma \gg$ MeV for suppressing the $\Delta N_{\rm eff}$