

Randall-Sundrum Dark Matter Portal Model

Symmetries, Amplitudes and Phenomenology

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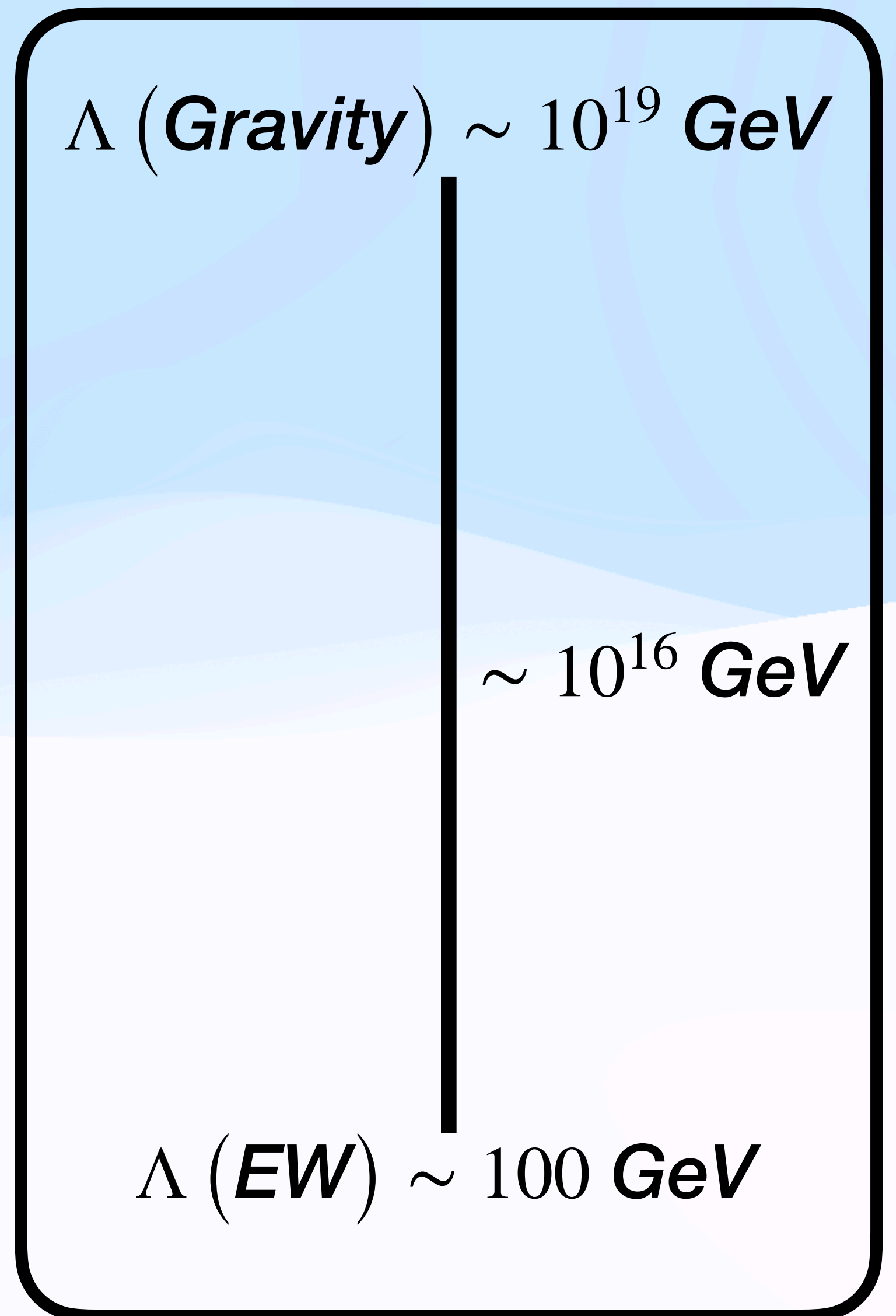
Randall-Sundrum Context

Some problems with SM:

- Hierarchy problem
- Highly fine tuned - Higgs mass
- Dark matters/dark energy unknowns

Randall-Sundrum models can answer* these and:

- Generate $\mathcal{O}(s)$ massive spin-2 EFT
- Be phenomenological probes for dark matter



General Relativity as Massless Spin-2

- **Perturbative** expansion about flat background by **spin-2 field**

$$g_{\mu\nu} = \eta_{\mu\nu} + \kappa_4 h_{\mu\nu}$$

$$\kappa_4 = \frac{2}{M_{\text{Pl}}} \sim 10^{-19} \frac{\text{GeV}}{c^2}$$

- Einstein's GR is precisely massless spin-2 theory - ghost free to all orders!

$$S_{\text{EH}} = \frac{1}{\kappa_4^2} \int d^4x \sqrt{|\det g|} R = \int d^4x \left[-\frac{1}{4} h^{\mu\nu} \zeta_{\mu\nu}^{\alpha\beta} h_{\alpha\beta} + \mathcal{O}(h^3) \right]$$

- Effective Field Theory with cutoff at $M_{\text{Pl}} \sim 10^{19}$ GeV
- Diffeomorphism invariance built in

Massive Gravity - Fierz-Pauli

$$S_{\text{FP}} = \int d^4x \left\{ -\frac{1}{4} h^{\mu\nu} \zeta_{\mu\nu}^{\alpha\beta} h_{\alpha\beta} - \frac{1}{4} m^2 \left(h_{\mu\nu} h^{\mu\nu} - h^2 \right) - \frac{\kappa_4^2}{2} h_{\mu\nu} T^{\mu\nu} \right\}$$

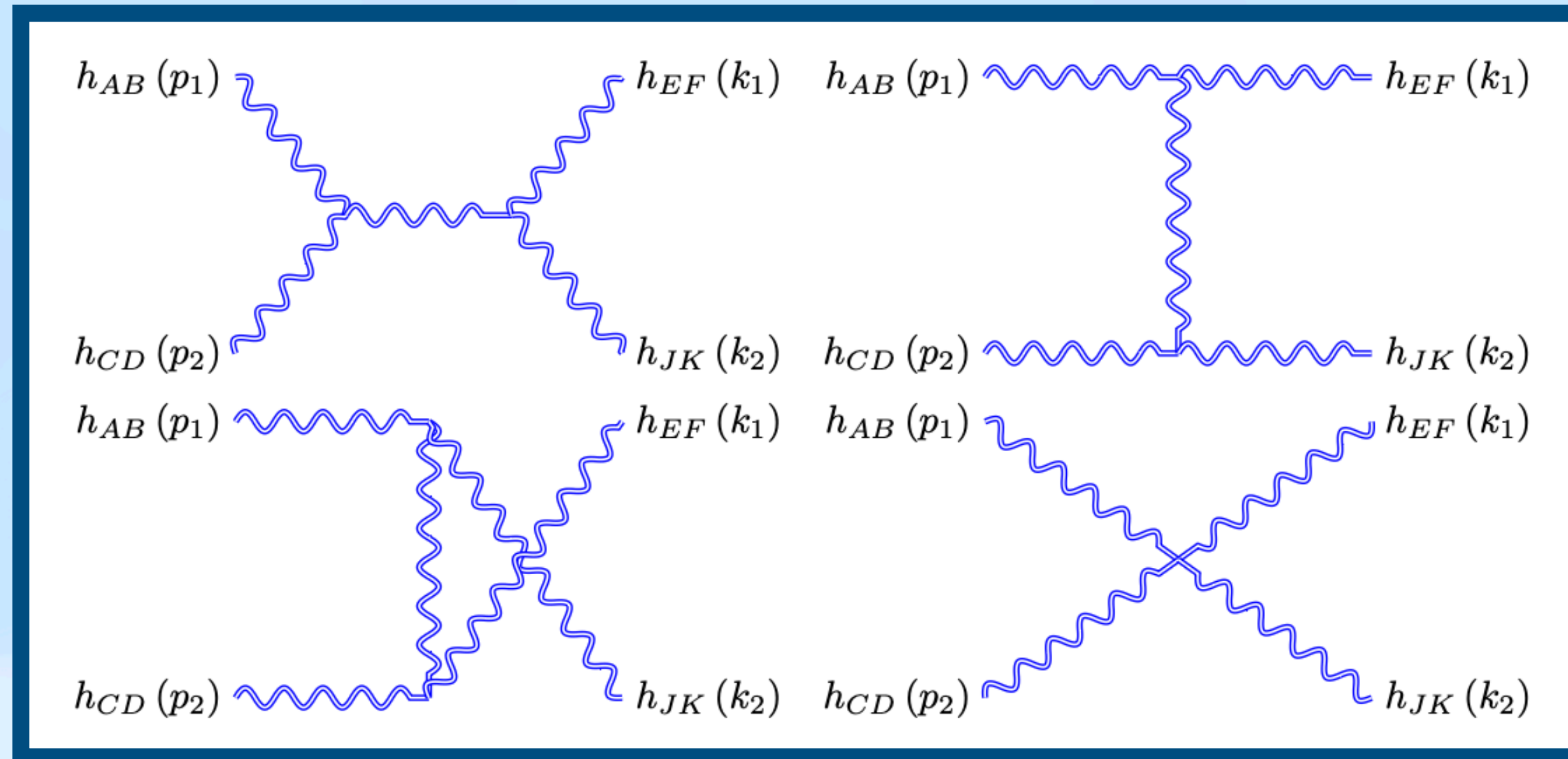
- Mass term explicitly breaks diffeomorphism invariance
- Expect GR for $m \rightarrow 0$

Massive Gravity - Fierz-Pauli

$$S_{\text{FP}} = \int d^4x \left\{ -\frac{1}{4} h^{\mu\nu} \zeta_{\mu\nu}^{\alpha\beta} h_{\alpha\beta} - \frac{1}{4} m^2 \left(h_{\mu\nu} h^{\mu\nu} - h^2 \right) - \frac{\kappa_4^2}{2} h_{\mu\nu} T^{\mu\nu} \right\}$$

- **Mass term** explicitly breaks diffeomorphism invariance
- Expect GR for $m \rightarrow 0$ - **doesn't happen**
- Now have five degrees of freedom
- Longitudinal mode couples with trace of stress-energy tensor
- **van Dam-Veltman-Zakharov discontinuity**

Fierz-Pauli - Self-Scattering Amplitude

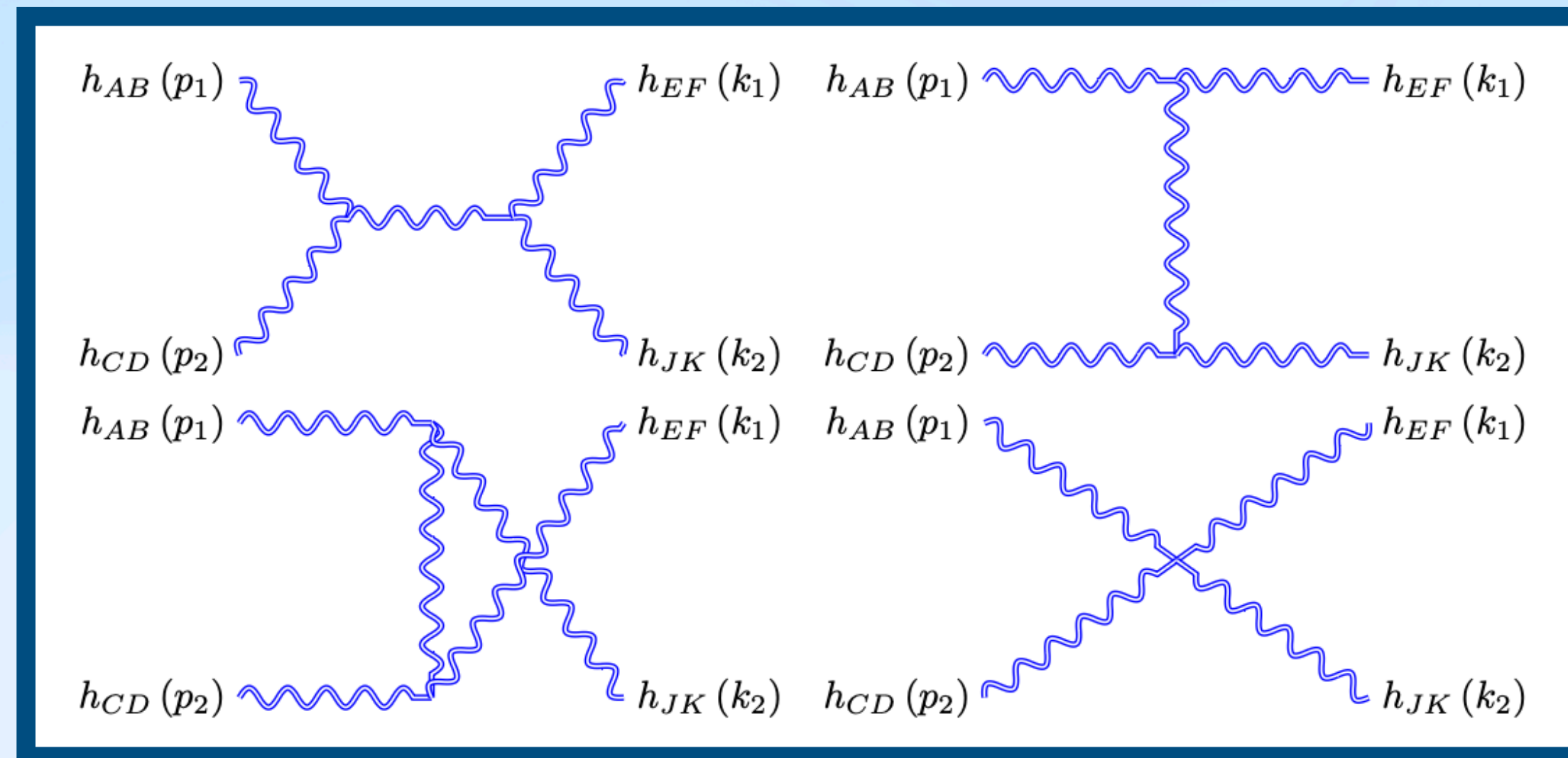


Amplitude Scales as $\mathcal{M} \sim \frac{s^5}{M_{\text{Pl}}^2 m^8}$

Effective Cutoff @ $\Lambda_{FP} = \left(M_{\text{Pl}}^2 m^8 \right)^{1/5}$

dRGT - Self-Scattering Amplitude

Partial resolution: de Rham-Gabadadze-Tolley (dRGT) massive gravity model

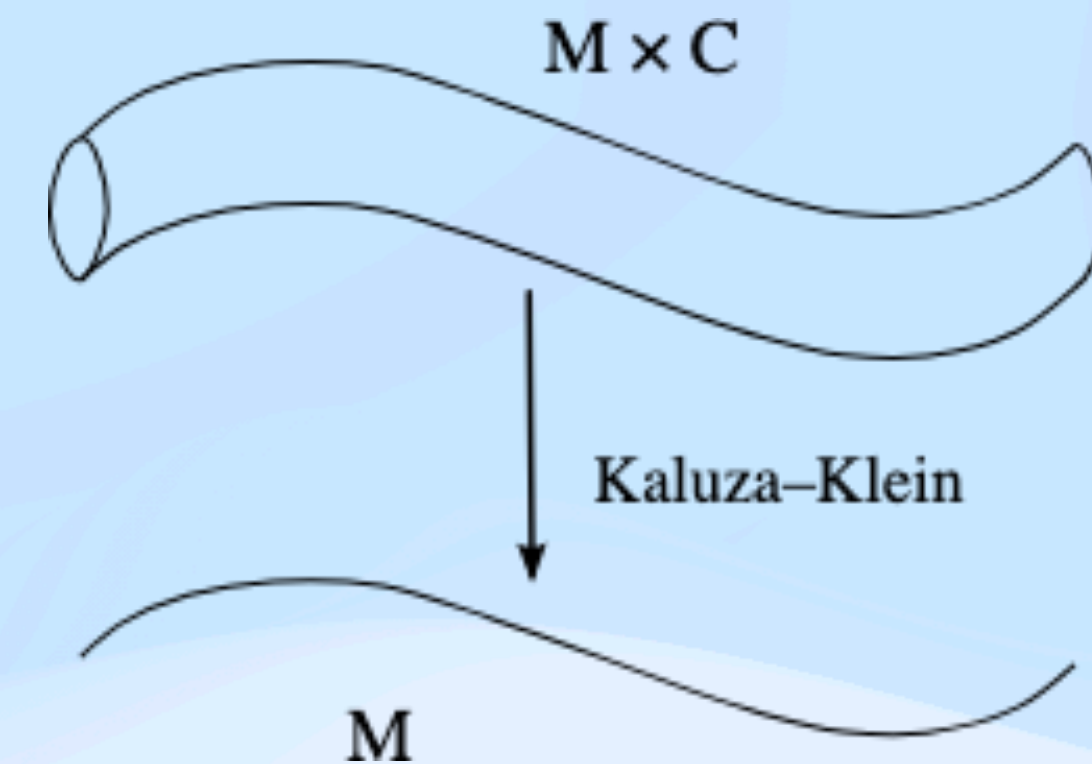


Amplitude Scales as $\mathcal{M} \sim \frac{s^3}{M_{\text{Pl}}^2 m^4}$

Effective Cutoff @ $\Lambda_{dRGT} = \left(M_{\text{Pl}}^2 m^4 \right)^{1/3}$

Compactified Extra-Dimensional Model of Gravity

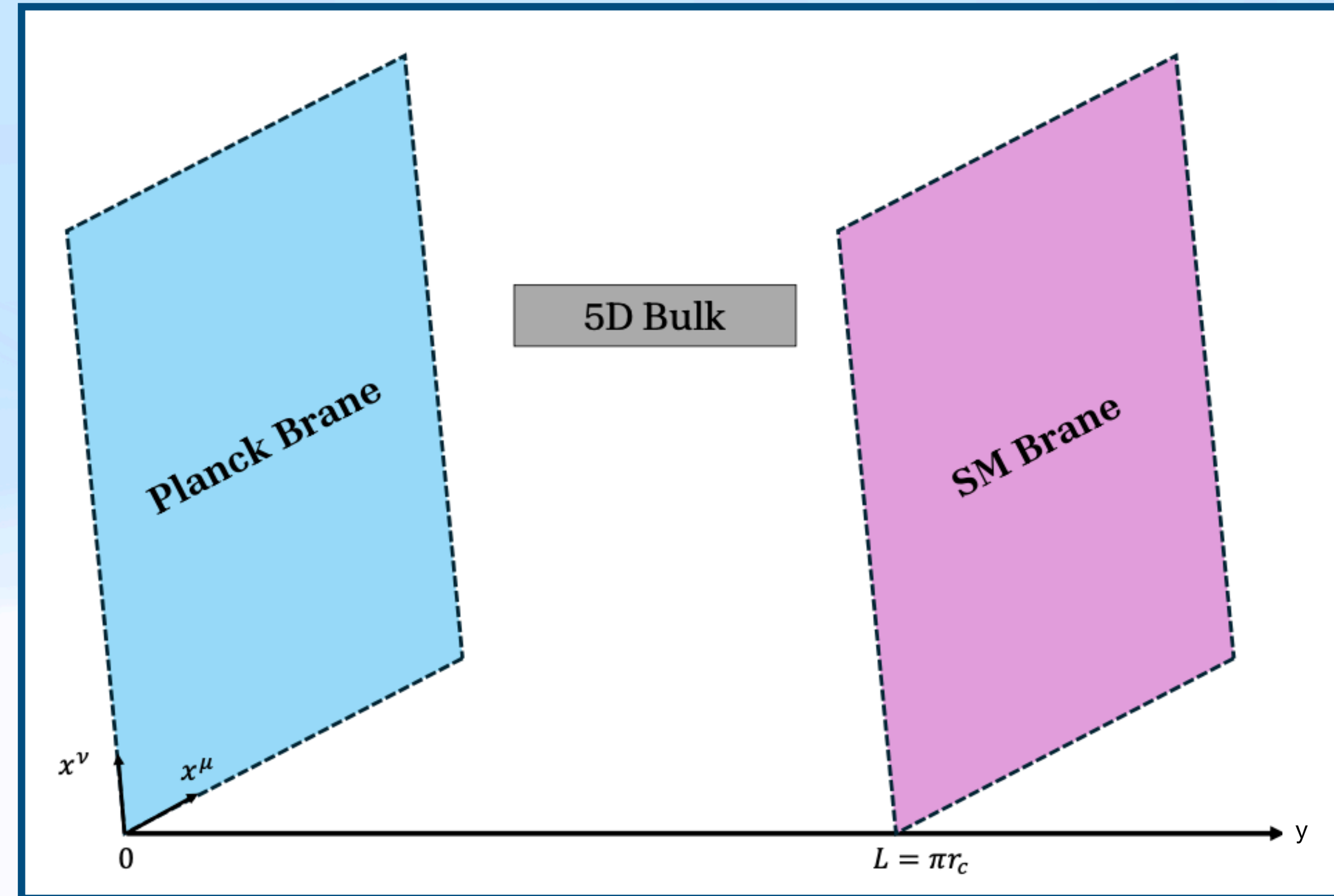
- **Kaluza-Klein** (1910/1920s) - unify **gravity** / **EM**
- **Introduces radion** - massless scalar field
- Extra-dim has **very small** extent $L \sim 10^{-30}$ cm
- **ONE Massless spin-2** state (the graviton - gravity)
- An **infinite tower of massive spin-2** states
- Also generate infinite tower of spin-1 and spin-0 fields
- **All spin-1 and spin-0** (except **massless spin-0 - radion**) can be gauged
- Independent of underlying geometry - choose a phenomenologically relevant one



**Geometric Higgs
Mechanism**

Randall-Sundrum (RS) Models

- Warped 4D, flat extra-dim
- Two 4D branes @ $y = 0, \pi r_c$
- Fundamental 5D Planck mass M_5
- Bulk Λ , brane-localised λ_i cosmological constants
- Constraints on Λ, λ_i fixed by $\Lambda_{4D} = 0$
- Additional - Goldberger-Wise mechanism
 - Length fixed by VEVs of bulk scalar
 - Radion becomes massive



$$ds^2 = e^{-2A(y)} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$$

$$S = -\frac{M_5^3}{4} \int d^4x \int_0^{\pi r_c} dy \sqrt{|\det G|} (\tilde{R} - 2\Lambda) + \int d^4x \sqrt{|\det G_i|} [\lambda_i]_{y=y_i}$$

RS Behaviour

- Motivated by Anti de-Sitter/Conformal Field Theory duality, choose $\Lambda > 0$

$$ds^2 = e^{-2A(y)} \eta_{\mu\nu} dx^\mu dx^\nu - dy^2$$

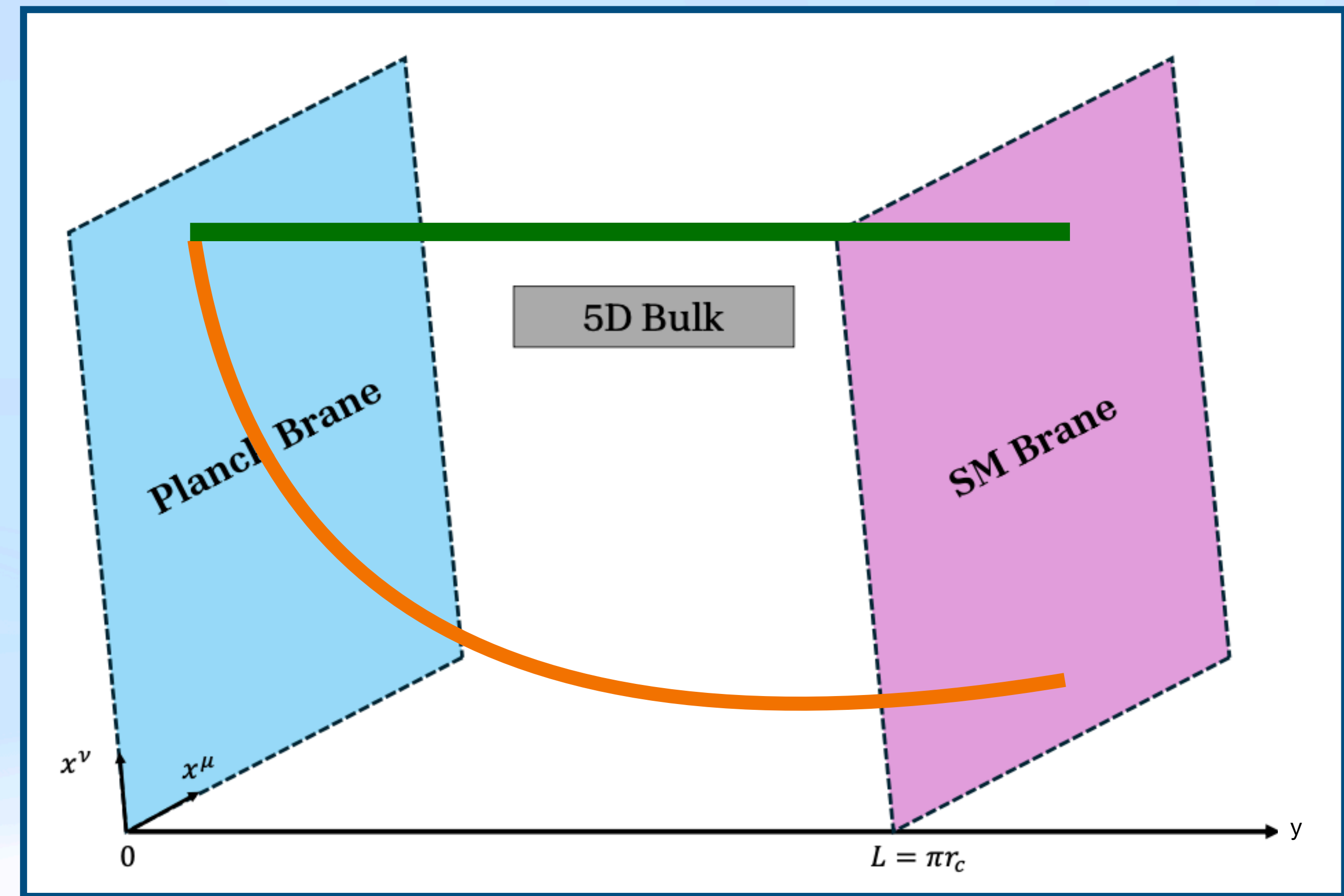
$$A(y) = ky, \quad k = \sqrt{\frac{\Lambda}{3}}$$

- Graviton has **universal coupling @ M_{Pl}**
- SM Higgs VEV** exponentially smaller

$$\frac{\langle \Phi \rangle_{\text{SM}}}{\langle \Phi \rangle_{\text{Pl}}} = e^{-k\pi r_c}$$

- Higher KK modes couple to SM as

$$\Lambda_{\text{TeV}} = M_{\text{Pl}} e^{-k\pi r_c}$$

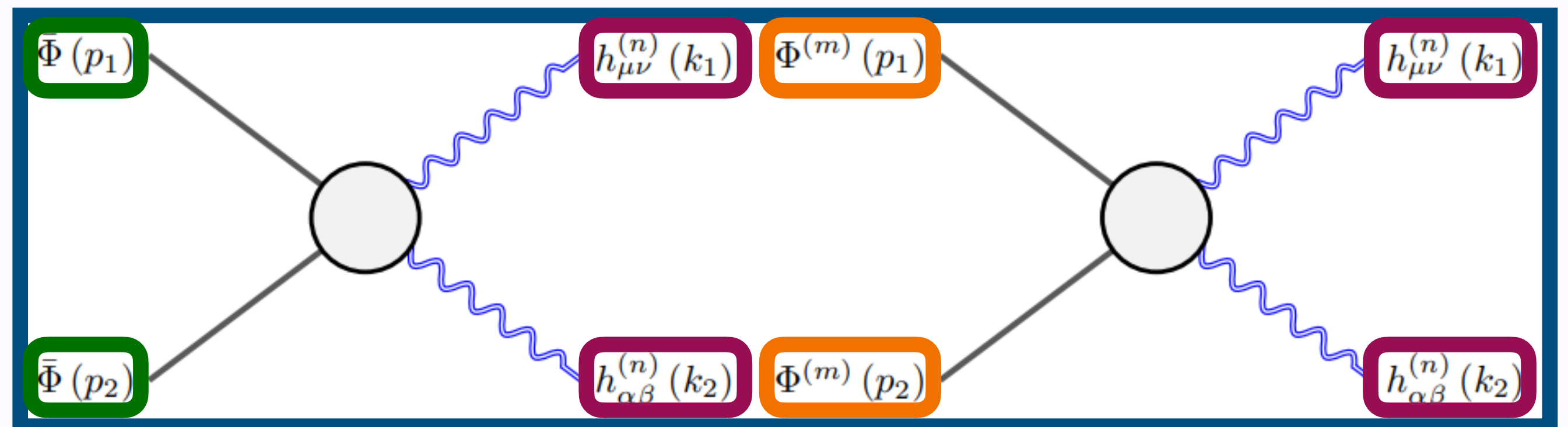


- Now want to investigate UV behaviour of this model

Scattering Amplitudes with Matter

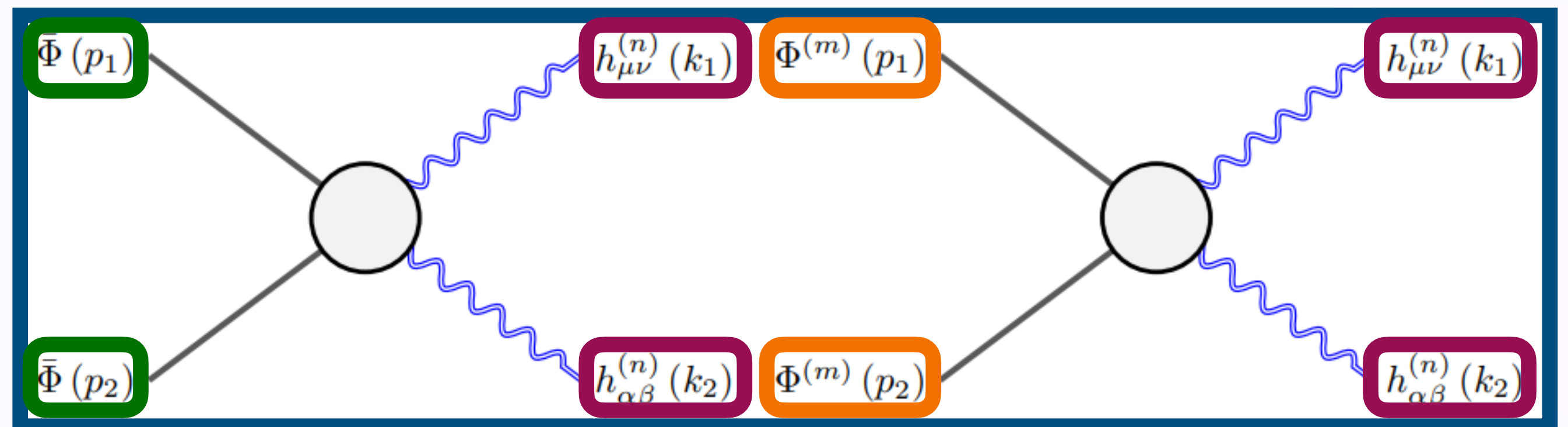
Phys. Rev. D 109, 015033 arXiv:2311.00770

- Produce **on-shell KK modes** from **brane-localised** or **bulk** matter
- Takeaways:
 - We define **sum-rules** as a series of rules that relate the summation of propagating KK modes to external on-shell KK modes.
 - For all (brane or bulk) matter species the combination of these sum-rules eliminate the bad high energy behaviour.
 - The **leading high energy behaviour** of the amplitude is always $\mathcal{M} \propto s/\Lambda$



Gauge Implications

- Previous result was in unitary gauge - try 't-Hooft Feynman-like gauge
- Generates unphysical vector and scalar KK modes - KK mass degeneracy
- $N = 2$ SUSY relationship between spin-2/spin-1/scalar KK
- High energy limit of amplitude appears as though scattering scalar KK modes, not spin-2 KK modes
- Behaviour resembles a Goldstone-Equivalence theorem...

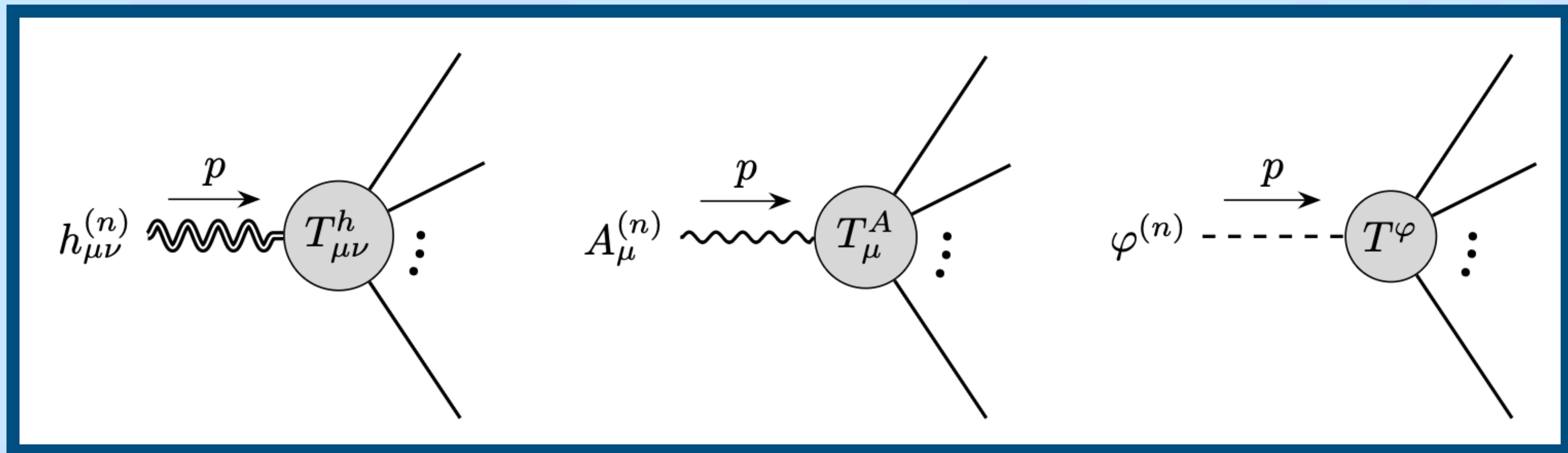


Symmetries and Equivalence theorems

Phys. Rev. D 109, 075016 arXiv:2312.08576

- The compactification is a ‘soft’ breaking of the 5D diffeomorphism
- Residual 4D diffeomorphism invariance \rightarrow equivalence theorem
- Exchange
 - (Unitary) **Spin-2 helicity-1** KK mode \rightarrow ('t-Hooft-Feynman) **spin-1 vector** KK mode
 - (Unitary) **Spin-2 helicity-0** KK mode \rightarrow ('t-Hooft-Feynman) **scalar** KK mode
- The **gauge choice and mass degeneracy** produce Ward identities
- Ward identities give transparent power counting for scattering processes

Example



$$\mathcal{M}^h = \epsilon^{\mu\nu}(p) T_{\mu\nu}^h$$

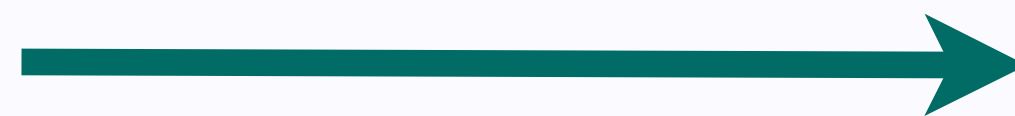
$$\mathcal{M}^A = \epsilon^{\mu}(p) T_{\mu}^A$$

$$\mathcal{M}^{\phi} = T^{\phi}$$

Mass degeneracy between helicity modes
ensures the following Ward Identities hold

$$ip^{\nu} T_{\mu\nu}^h - \frac{1}{\sqrt{2}} m_n T_{\mu}^A = 0$$

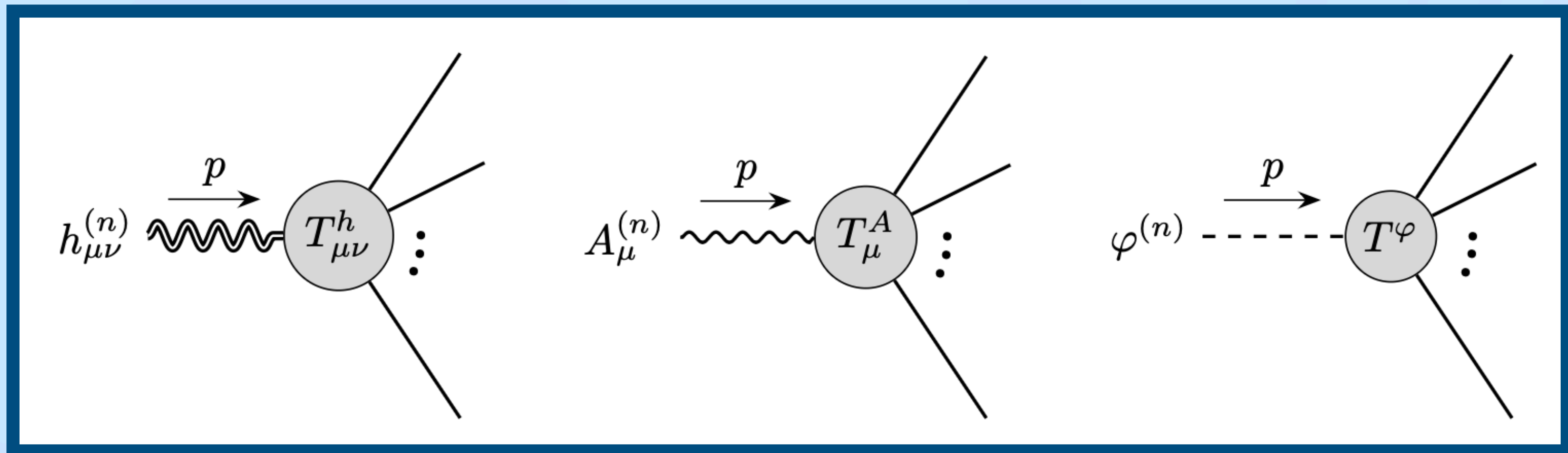
$$-\frac{1}{2} m_n \eta^{\mu\nu} T_{\mu\nu}^h + \frac{i}{\sqrt{2}} p^{\mu} T_{\mu}^A + \sqrt{\frac{3}{2}} m_n T^{\phi} = 0$$



$$\frac{1}{\sqrt{6}} T_{\mu\nu}^h \left(\eta^{\mu\nu} + \frac{2p^{\mu} p^{\nu}}{m} \right) = T^{\phi}$$

Helicity-2 sub-amplitude related to Goldstone scalar sub-amplitude

Example



$$\mathcal{M}^h = \epsilon^{\mu\nu}(p) T_{\mu\nu}^h$$

$$\mathcal{M}^A = \epsilon^{\mu}(p) T_{\mu}^A$$

$$\mathcal{M}^{\phi} = T^{\phi}$$

High energy scattering of massive longitudinal spin-2 KK mode has same behaviour as scalar field

$$T_{\mu\nu}^h \epsilon_0^{\mu\nu} = T^{\phi} + \mathcal{O}(E^0)$$

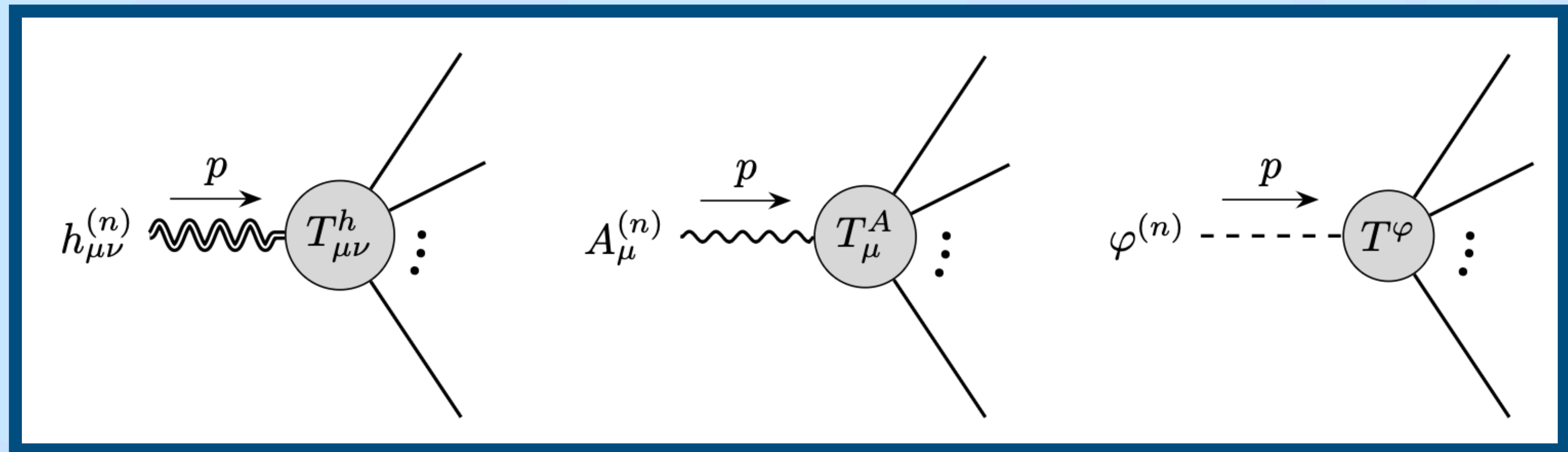
Massive spin-2 modes in unitary gauge are eaten Goldstones in 't-Hooft Feynman gauge

$$\underbrace{\epsilon_0^{\mu\nu}}$$

$$\frac{1}{\sqrt{6}} T_{\mu\nu}^h \left(\eta^{\mu\nu} + \frac{2p^{\mu} p^{\nu}}{m} \right) = T^{\phi}$$

Helicity-2 sub-amplitude related to Goldstone scalar sub-amplitude

Results



$$\mathcal{M}^h = \epsilon^{\mu\nu}(p) T_{\mu\nu}^h$$

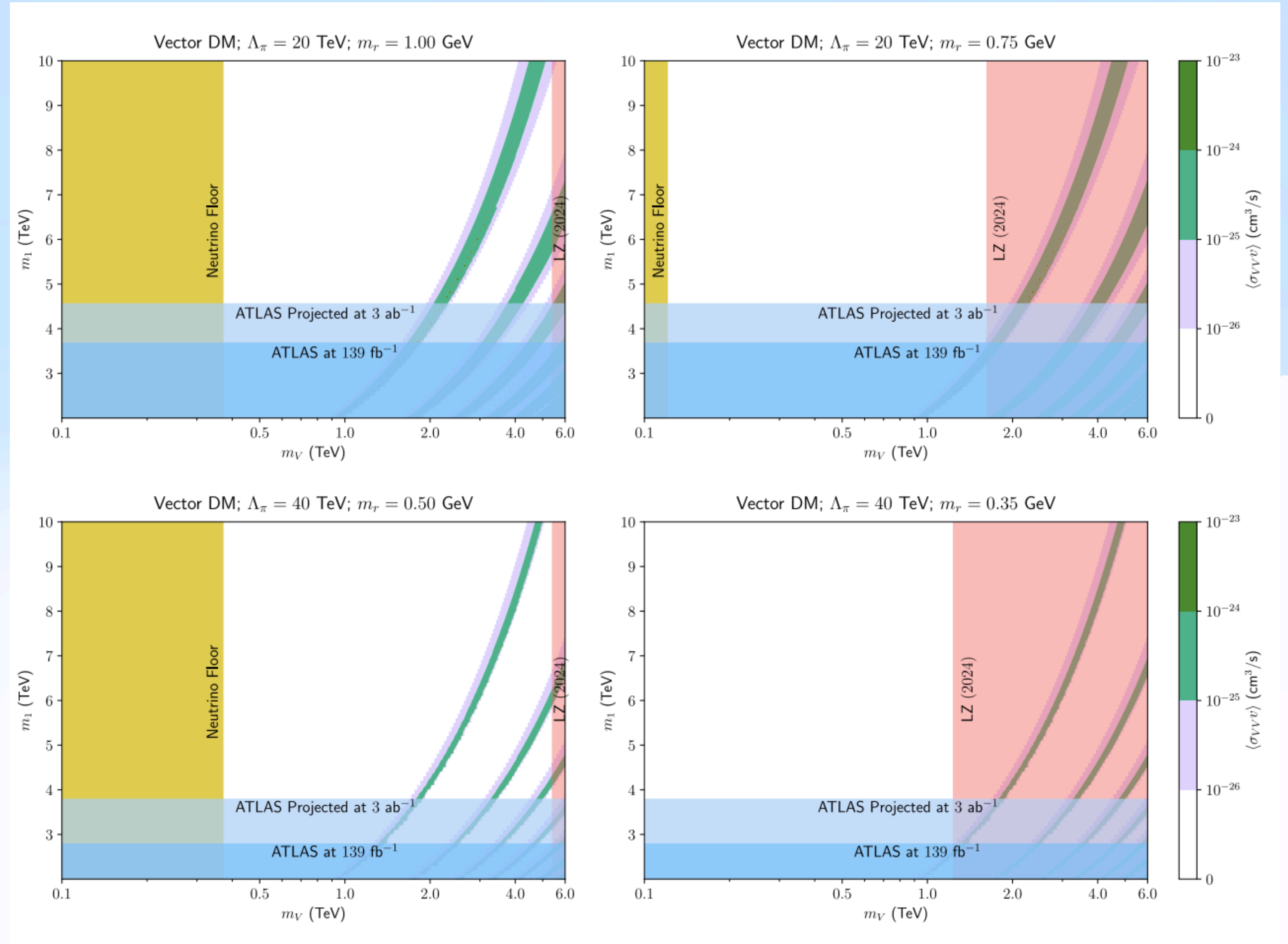
$$\mathcal{M}^A = \epsilon^{\mu}(p) T_{\mu}^A$$

$$\mathcal{M}^{\phi} = T^{\phi}$$

- Generalises to **any number** of **external spin-2** fields
- AND can include **propagators** in framework
- Confirms unitary gauge result (previous work)
- No need for complex sum-rules in this formalism

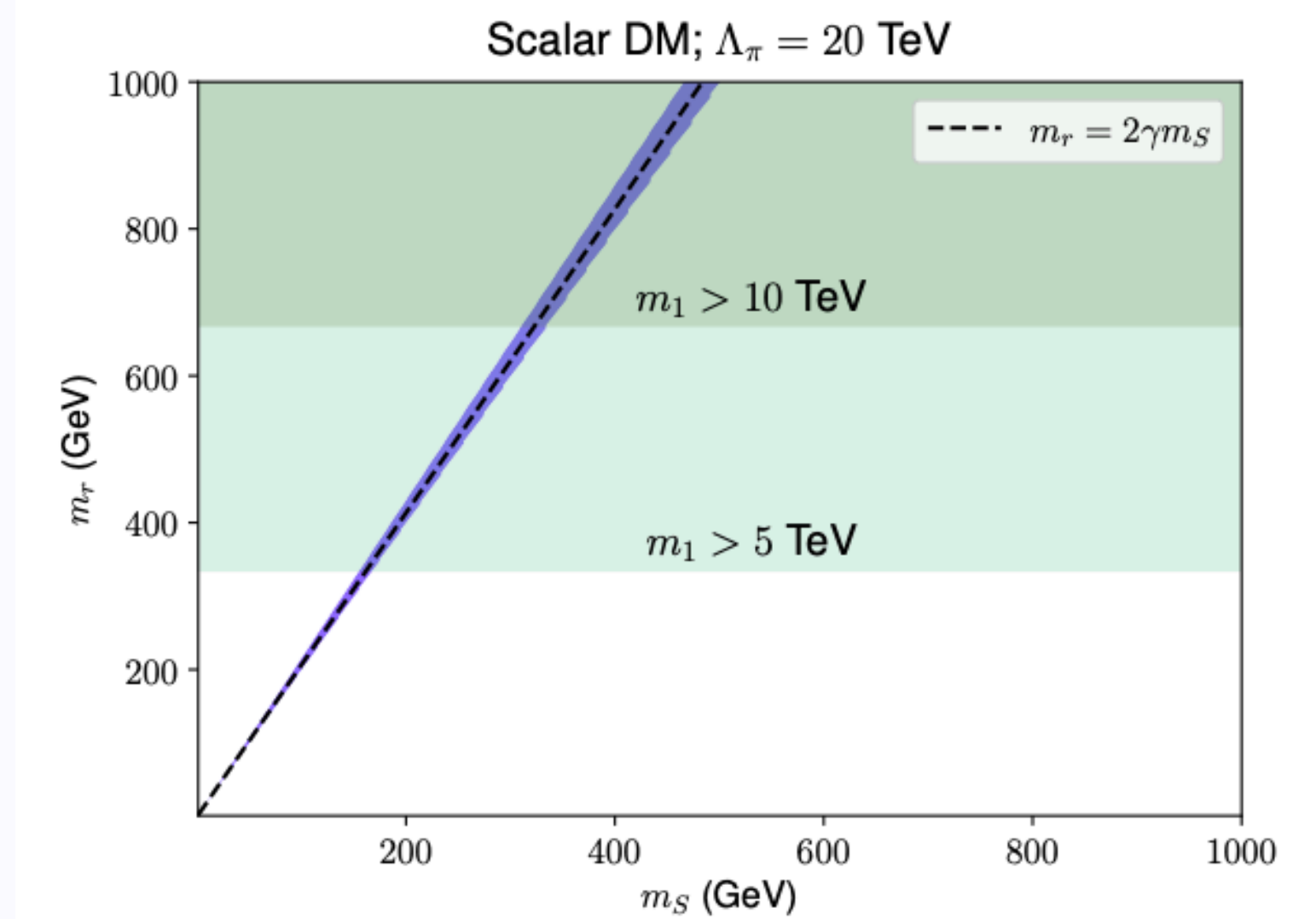
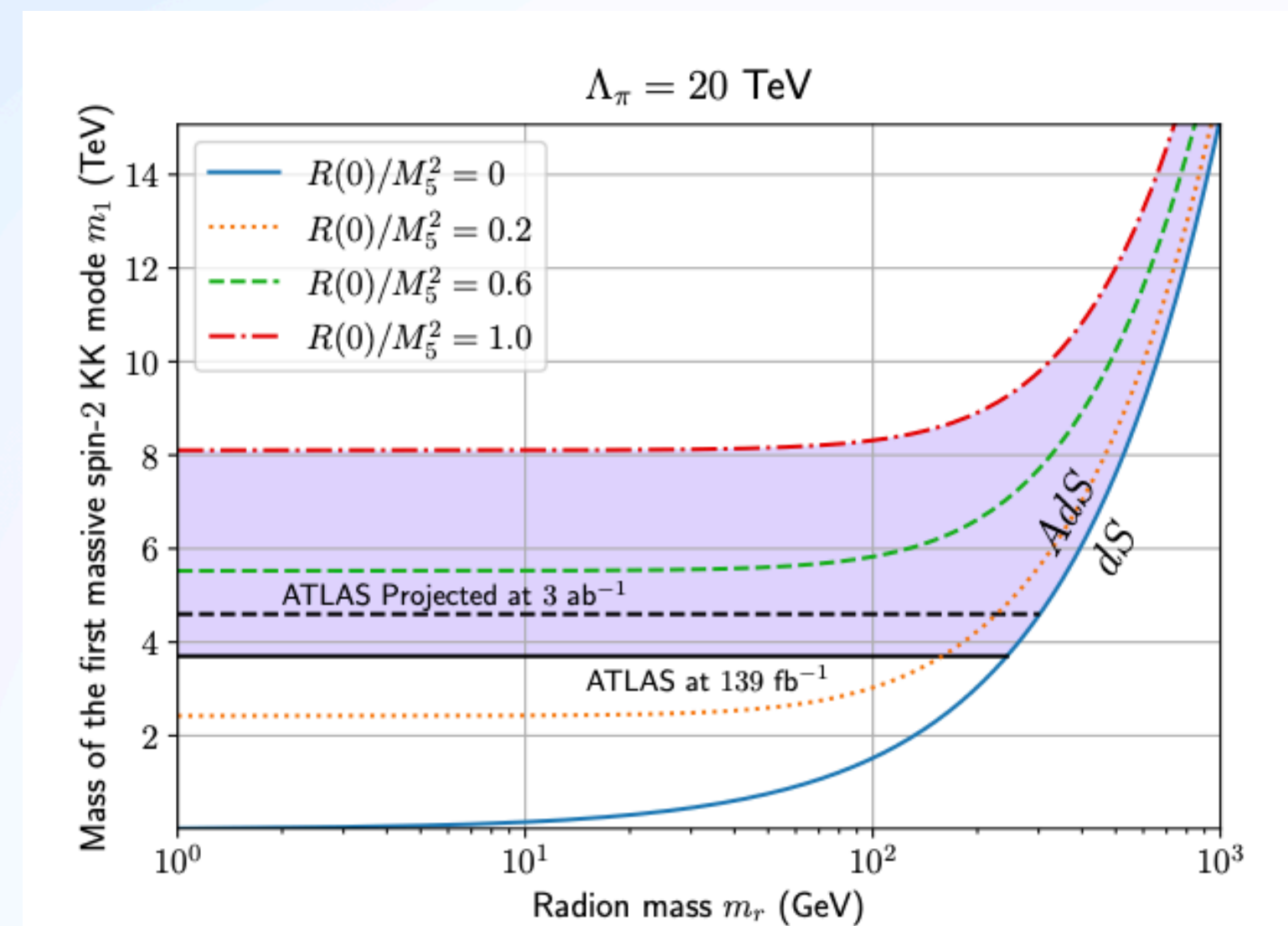
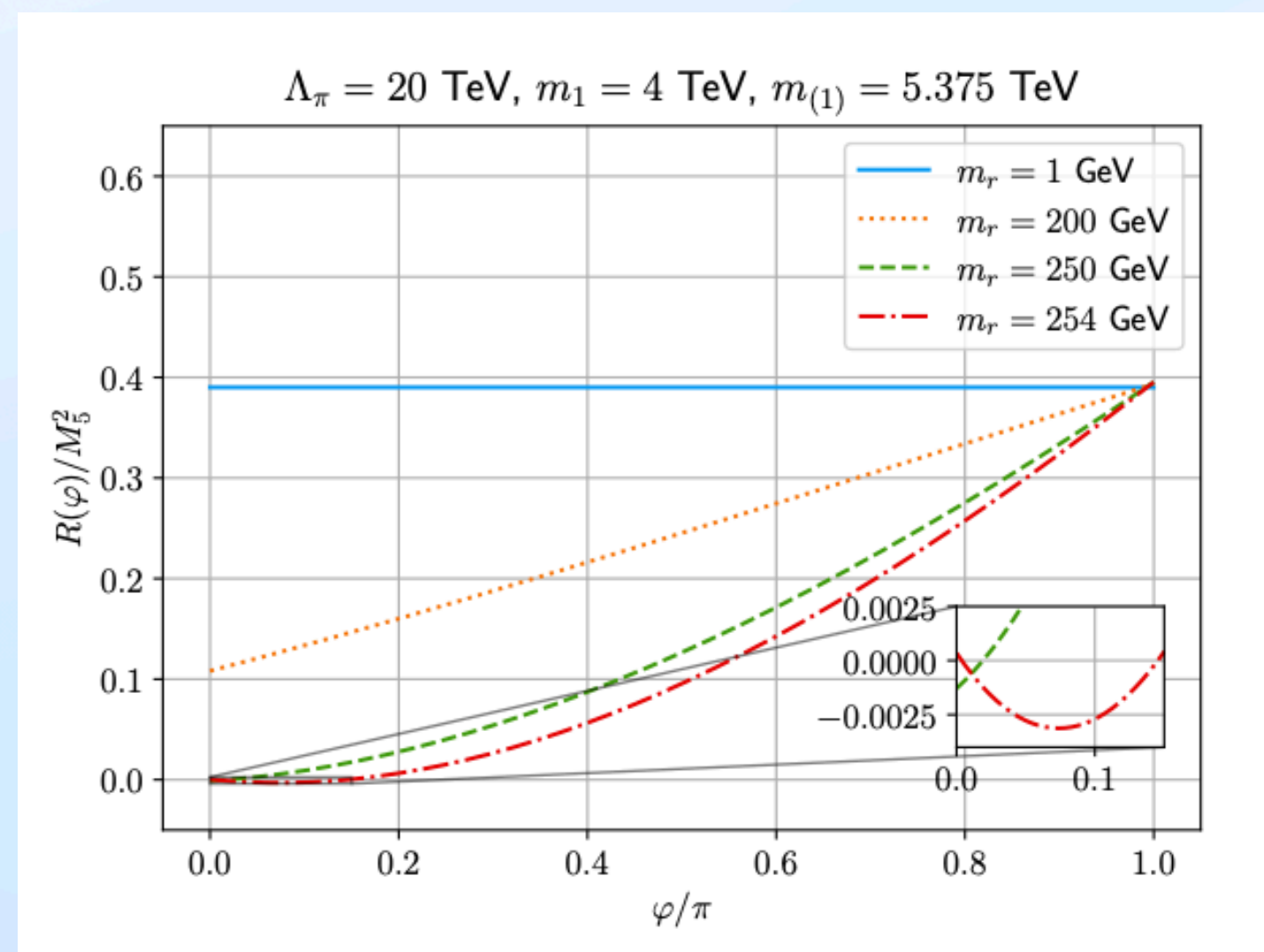
DM Phenomenology - Light Radion

- Localise DM on TeV brane
- Vector DM - viable gap in parameter space
- Lilac and green regions
 - Equals or exceeds requirement for correct DM relic density
- Low backreaction from GW
 - $m_r \sim \mathcal{O}(1 \text{ GeV})$



DM Phenomenology - Heavy Radion

- Localise DM on TeV brane, but high backreaction from GW $m_r \sim \mathcal{O}(100 \text{ GeV})$
- Change for AdS to dS - new constraint'
- Scalar and vector DM viable on resonance in a narrow band ($\Lambda = 20 \text{ TeV}$)
- Annihilation to radion (decays to b-quarks) rules out $m_{\text{DM}} \approx 5 - 80 \text{ GeV}$



Next Steps

Multibrane Model

- Introduce another hierarchy
- Changes to mass spectrum
- Investigate self scattering behaviour
- Preliminary findings:
 - There is no $\mathcal{O}(s^3)$ contribution to the amplitude
 - Leading order behaviour is of $\mathcal{O}(s^2)$

