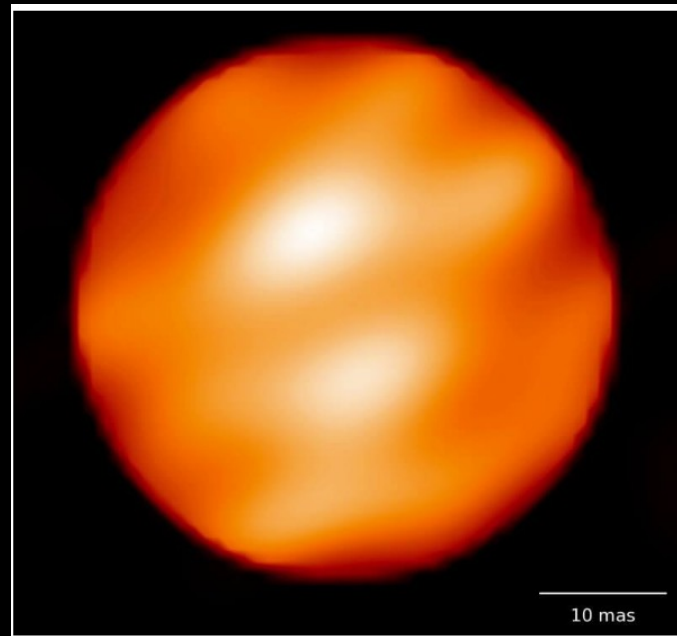


Astrophysical Consequences of an η_W Pseudoscalar

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Credit: Observatoire de Paris/Xavier Haubois et al

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θ angles

- An interaction allowed by all known symmetries: $\frac{\alpha_s \theta_{\text{QCD}}}{8\pi} G_{\mu\nu}^a G^{\tilde{a}\mu\nu}$
 - Breaks P and T , hence CP
 - CP is not a preserved symmetry of Nature (electroweak interactions)
- Neutron EDM bounds: $\theta_{\text{QCD}} \lesssim 10^{-10}$ (Why? Strong CP problem)
- One can rotate θ_{QCD} into the quark mass matrix M_q
 \Rightarrow physical parameter $\bar{\theta}_{\text{QCD}} \equiv \theta_{\text{QCD}} + \arg |M_q| \Rightarrow$ Why $\bar{\theta}_{\text{QCD}} \lesssim 10^{-10}$
- New physics solutions: *E.g.*, Peccei-Quinn (1977), Nelson (1984)-Barr (1984)
- Obvious solution: one massless quark would allow to rotate $\bar{\theta}_{\text{QCD}}$ away
 - Introduces a classical chiral symmetry
 - Not favored by low energy hadronic data (lattice, phenomenology)
- For $m_q \rightarrow 0$, theoretical arguments imply η' would dynamically remove $\bar{\theta}_{\text{QCD}}$

E.g., Witten, 1979

 - Peccei-Quinn axion can accomplish this, but requires new physics

The electroweak sector

- Analogous $SU(2)_L$ angle: $\frac{\alpha_2 \theta_w}{8\pi} W_{\mu\nu}^a \tilde{W}^{a\mu\nu}$ $a = 1, 2, 3; \alpha_2 = g_2^2/(4\pi)$

- A classically good symmetry $U(1)_{B+L}$ (good quality)

- However $B + L$ broken by an anomaly at quantum level $\partial_\mu J^{\mu B+L} \propto W_{\mu\nu}^a \tilde{W}^{a\mu\nu}$

- One can be removed θ_w away by rotating quark and lepton fields

$$Q \rightarrow e^{i\alpha/6} Q \quad \text{and} \quad \ell \rightarrow e^{i\alpha/2} \ell$$

See, e.g., Fileviez Pérez, Patel, 2014

$$\theta_w \rightarrow \theta_w + N_f \alpha \quad (N_f = 3 \text{ number of families})$$

- Akin to the massless quark case for QCD

The case for η_W

- Recent work has raised the question: is θ_W removed dynamically by an η_W ?
 - Analogy with η' for QCD with $|M_q| = 0$

Dvali, Kobakhidze, Sakhelashvili (DKS), 2408.07535, 2509.16043
- DKS argue η_W can emerge from the Standard Model (SM) electroweak dynamics
- Theoretical arguments for this require more careful examination
- No firm prediction for the properties (mass m_{η_W} , decay constant f_{η_W}) of η_W
- Nonetheless, the emergence of this state can have significant import
 - A deeper understanding of electroweak (EW) dynamics and phenomenology
- One may ask: can the typical expectations for η_W be consistent with data?
 - This can be complementary to formal investigations
 - The answers may challenge the underlying assumptions and arguments

General features of η_W

- We follow DKS in general description of the main points
- η_W as a pseudo-Nambu-Goldstone boson (pNGB): spontaneous $U(1)_{B+L}$ breaking
- Mass generation from EW $SU(2)$ anomaly of $B + L$ current

$$m_{\eta_W}^2 \lesssim e^{\frac{-2\pi}{\alpha_2(M)}} M^2$$

- M is the relevant scale

- We take $M \sim v$, with $v \approx 246$ GeV the Higgs vev $\Rightarrow e^{-2\pi/\alpha_2(v)} \sim 10^{-80}$

$$m_{\eta_W} \lesssim 10^{-29} \text{ eV}$$

- Too light for “fuzzy” dark matter (de Broglie wavelength \gg Mpc)

- DKS: spontaneous $U(1)_{B+L}$ breaking via the 't Hooft vertex of the EW instantons

$$\langle (QQQ\ell)^{N_f} \rangle \propto e^{\frac{-2\pi}{\alpha_2(v)}} v^{6N_f}$$

Quark Q and lepton ℓ doublets

- Analogy with QCD chiral condensate $\langle \bar{q}_L q_R \rangle$, setting the pion decay constant f_π .

- This suggests typical expectation (also consistent with DKS arguments)

$$f_{\eta_W} \lesssim v$$

Interactions of η_W with SM

- To constrain η_W , we need to find its couplings to matter
- Couplings to $SU(2)$ and $U(1)_Y$ gauge bosons via anomalies

$$\frac{\eta_W}{f_{\eta W}} \left(\frac{\alpha_2}{8\pi} W_{\mu\nu}^a \tilde{W}^{a\mu\nu} - \frac{\alpha_1}{8\pi} B_{\mu\nu} \tilde{B}^{\mu\nu} \right)$$

- Since $B + L$ is not anomalous under EM/QCD, no coupling to $\gamma\gamma/gg$
- However, there is a 1-loop coupling to fermions $i\xi m_\psi \frac{\eta_W}{f_{\eta W}} \bar{\psi} \gamma_5 \psi$, where

$$\xi = \frac{3\alpha^2}{16\pi^2} \left[\frac{3}{4 \sin^4 \theta_W} - \frac{Y_L^2 + Y_R^2}{\cos^4 \theta_W} \right] \ln \left(\frac{\Lambda^2}{m_W^2} \right)$$

See, e.g., Bauer, Neubert, Thamm, 2017

Fermion mass m_ψ , α EM fine structure constant, θ_W weak mixing angle, Y_L and Y_R hypercharges of the doublet and singlet ψ , respectively.

W^\pm mass $m_W \approx 80.4$ GeV, $\Lambda > m_W$ a cutoff scale, $\sin^2 \theta_W(m_Z) \approx 0.231$, Z mass $m_Z \approx 91.2$ GeV, $\alpha(m_Z^2) \approx 1/127.9$

Coupling to electrons

- We may derive a coupling electrons $y_e \eta_W \bar{e} \gamma_5 e$ based on the above
- One could set $\Lambda = 4\pi f_{\eta W}$ and get

$$y_e \approx 3 \times 10^{-5} \ln \left(\frac{16\pi^2 f_{\eta W}^2}{m_W^2} \right) \frac{m_e}{f_{\eta W}}$$

- Globular cluster red giant anomalous energy loss: $y_e \lesssim 10^{-13}$ [Caputo, Raffelt, 2024](#)

$$\Rightarrow f_{\eta W} \gtrsim 10^3 \text{ TeV}$$

Similar result, in a different context: [Cacciapaglia, Sannino, Turner, 2510.14104](#)

- The above lower bound in tension with an EW origin for η_W

Conclusions

- Recent suggestion for emergence of a light pseudoscalar η_W from EW dynamics warrants examination
- If present, it has significant implications for non-perturbative EW interactions
- One may ask whether the expected properties of such a state are accommodated by data
- We considered this question and found that red giant cooling bounds strongly disfavor values of couplings for η_W suggested by theoretical considerations
- This can be interpreted as a challenge to theory arguments for its emergence from SM interactions
- Further theoretical investigations regarding η_W seem warranted, given its potential impact on a fundamental description of the EW sector