All-Charm TetraQuarks at Hadron Colliders
A High-Precision Fragmentation Perspective

Francesco Giovanni Celiberto, UAH Madrid



Resummation, Evolution, Factorization 2026

🇪🇸 Santiago de Compostela - October 22nd 🇪🇸

F. G. C. [Phys. Rev. D Letters 111 (2025) 11, L111501 [arXiv:2504.03949]]
F. G. C. et al. [Phys. Rev. D 113 (2026) 5, 054014 [arXiv:2511.18984]]
F. G. C. [Phys. Rev. D 113 (2026) 11, 114037 [arXiv:2604.11646]]
J. V. Bulhões, F. G. C., V. P. Gonçalves [arXiv:2610:27396] [TQ4Qhf1.0]

Madrid UAH

talento cm Programa de atracción de talento investigador Comunidad de Madrid

EUGLOH

QCD Hydrogen Atoms

I

¿ Why studying (Heavy) Exotics ?

JETHAD
v0.5.5

Precision & Exploration

¿ What is an Exotic Hadron ?

- Baryons & Mesons → composite states → $|q_1 q_2 q_3\rangle$ & $|q_1 \bar{q}_2\rangle$
- Ordinary Hadrons ← simplest valence-quark color-singlet states
- QCD color neutrality → TetraQuarks, PentaQuarks, Hybrids allowed

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- First discovery of a (crypto-)exotic at Belle (🇯🇵) in 2003 → $X(3872)$

Interpretation of Exotic States

Compact MultiQuark



Interpretation of Exotic States

Compact MultiQuark





Hybrid

Interpretation of Exotic States

diQuark + AntidiQuark



Compact MultiQuark





Hybrid

Interpretation of Exotic States

Compact MultiQuark



diQuark + AntidiQuark





Hybrid



Molecule of Mesons

Interpretation of Exotic States

Compact MultiQuark



diQuark + AntidiQuark



HadroQuarkonium





Hybrid



Molecule of Mesons

All-Charm TetraQuarks: Experimental Highlights

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH (CERN)

CERN-EP-2025-118
2025/06/10

CMS-BPH-24-002

Determination of the spin and parity of all-charm tetraquarks

The CMS Collaboration*

arXiv:2506.07944v1 [hep-ex] 9 Jun 2025

Abstract

The traditional quark model accounts for the existence of baryons, such as protons and neutrons, which consist of three quarks, as well as mesons, composed of a quark-antiquark pair. Only recently has substantial evidence started to accumulate for exotic states composed of four or five quarks and antiquarks. The exact nature of their internal structure remains uncertain. This paper reports the first measurement of quantum numbers of the recently discovered family of three all-charm tetraquarks, using data collected by the CMS experiment at the Large Hadron Collider from 2016 to 2018. The angular analysis techniques developed for the discovery and characterization of the Higgs boson have been applied to the new exotic states. The quantum numbers for parity P and charge conjugation C symmetries are found to be $+1$. The spin J of these exotic states is consistent with $2\hbar$, while $0\hbar$ and $1\hbar$ are excluded at 95% and 99% confidence level, respectively. The $J^{PC} = 2^{++}$ assignment implies particular configurations of constituent spins and orbital angular momenta, which constrain the possible internal structure of these tetraquarks.

Submitted to Nature



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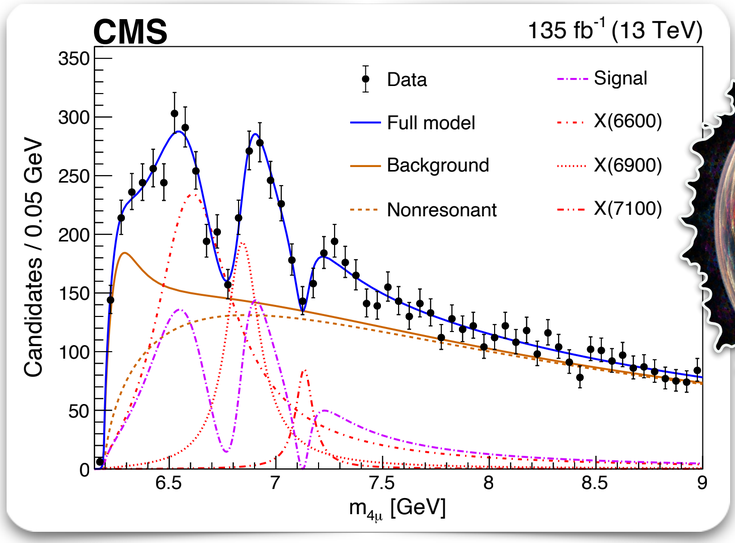

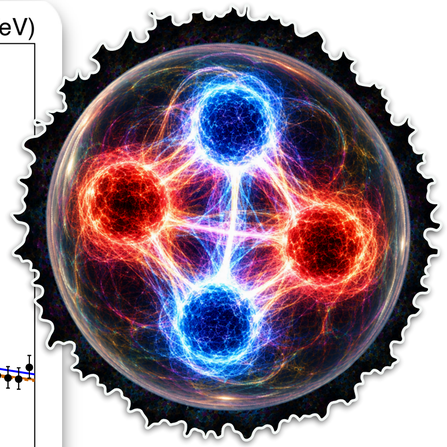

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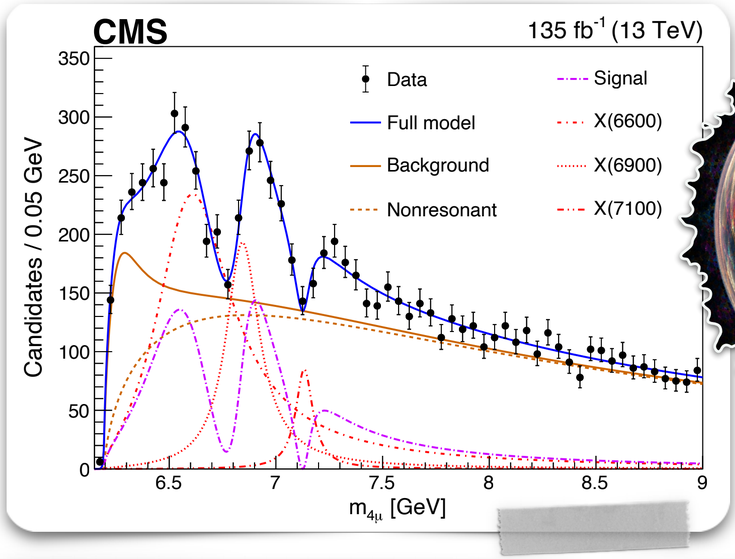

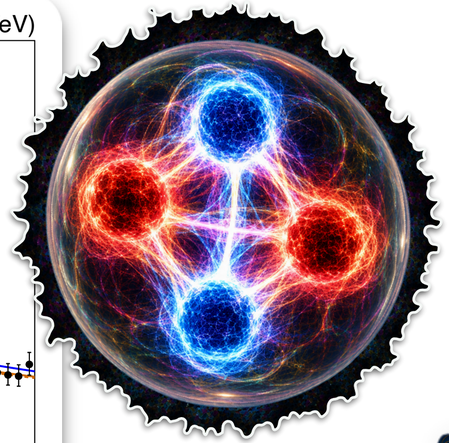



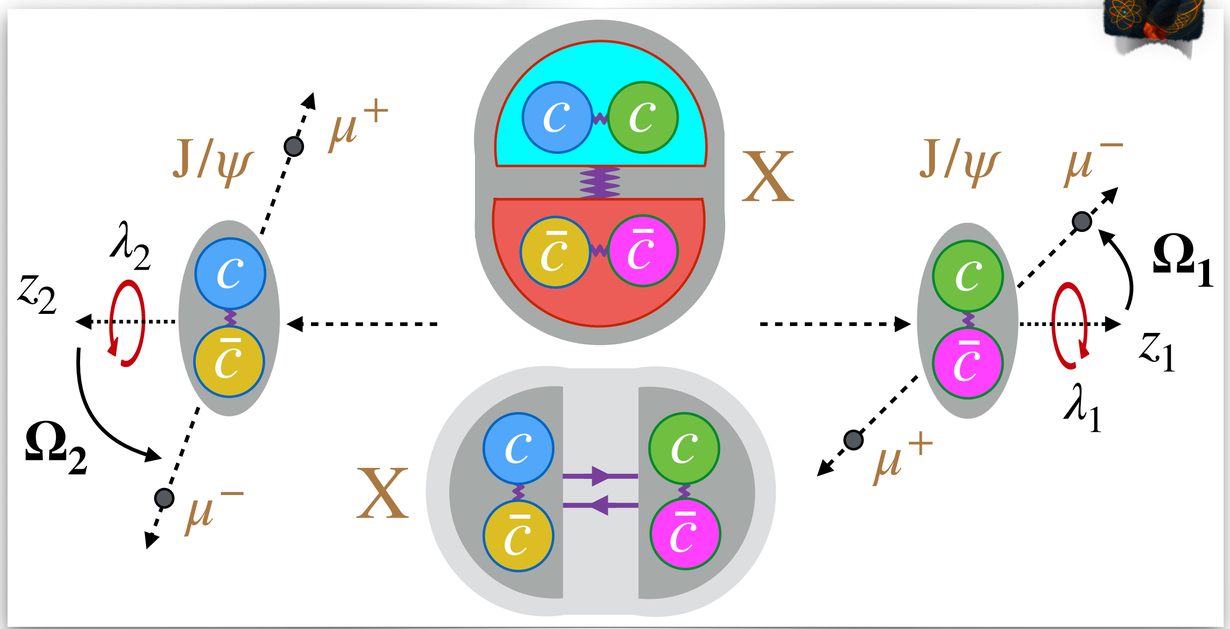

All-Charm TetraQuarks: Experimental Highlights

Article

Determination of the spin and parity of all-charm tetraquarks

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The CMS Collaboration*✉

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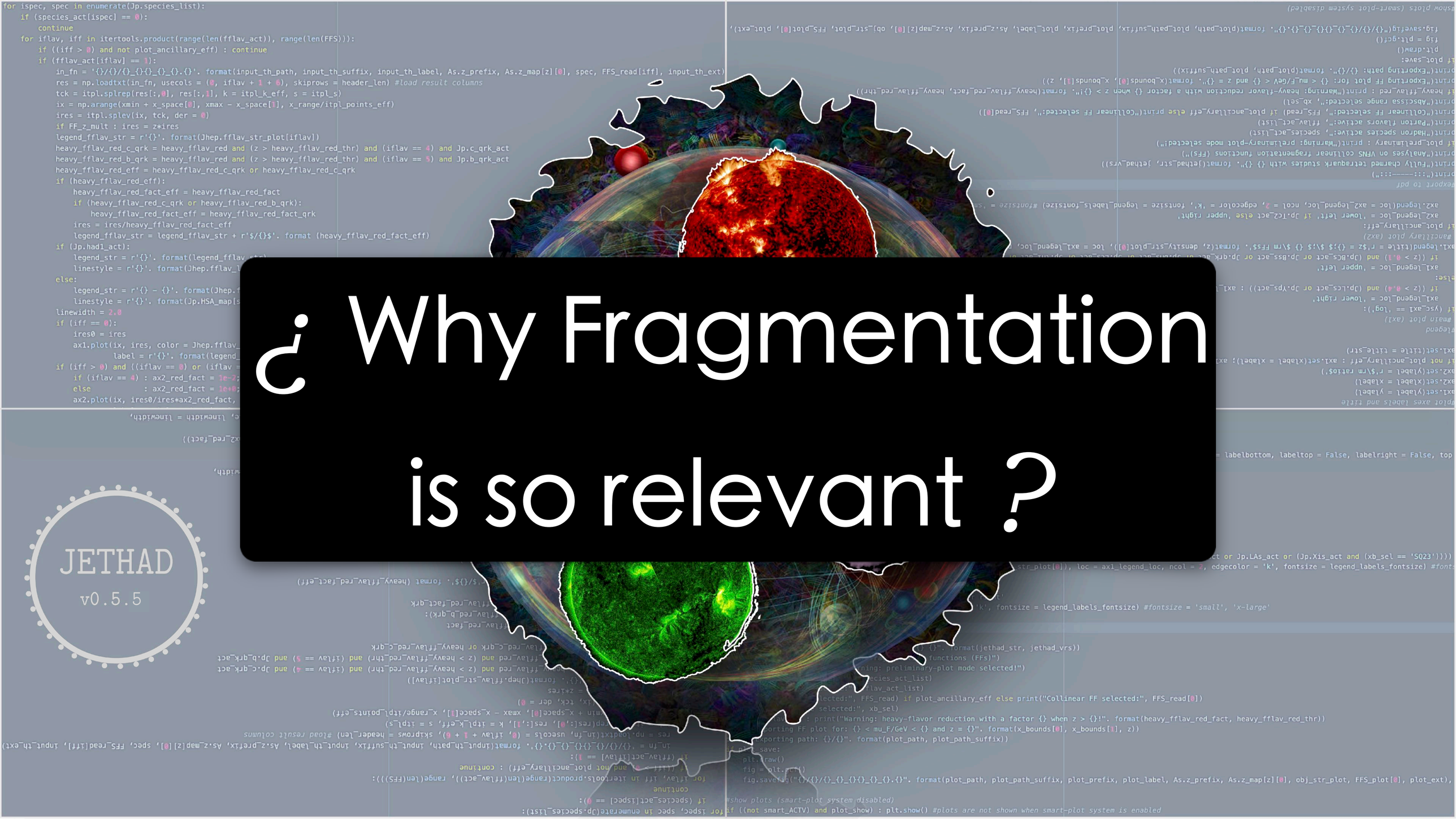
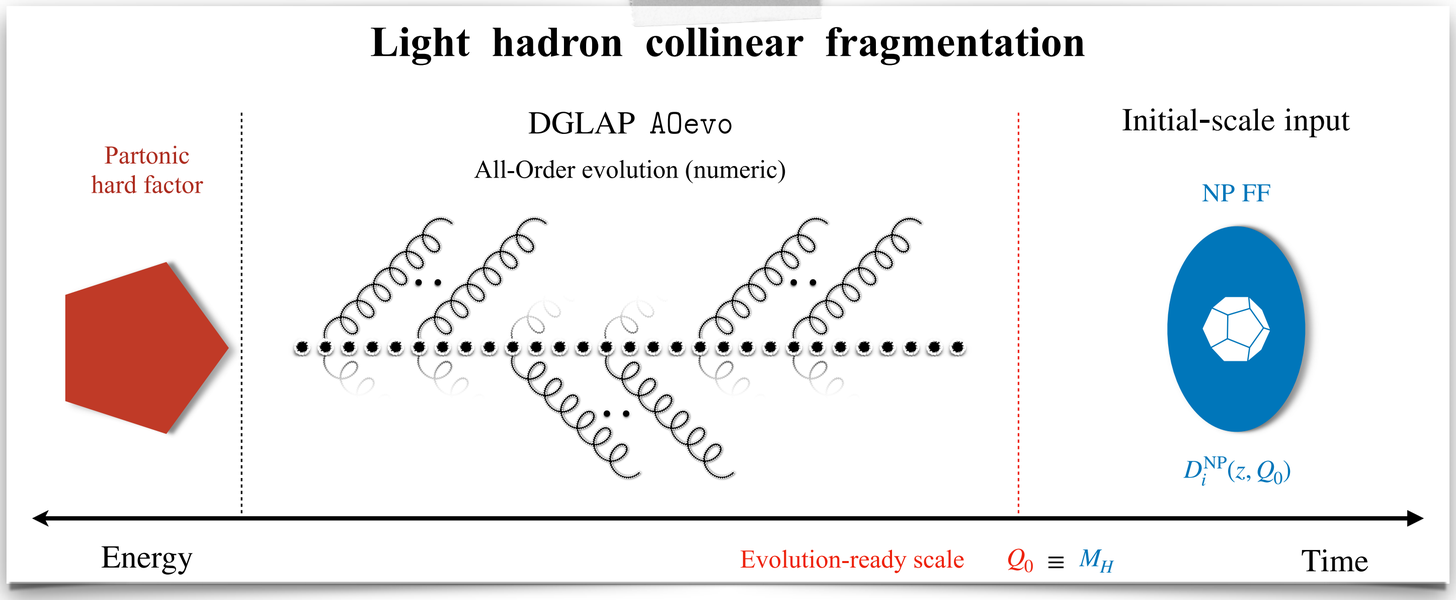

¿ Why Fragmentation
is so relevant ?
JETHAD
v0.5.5

Fragmentation in QCD (with Heavy Flavors)

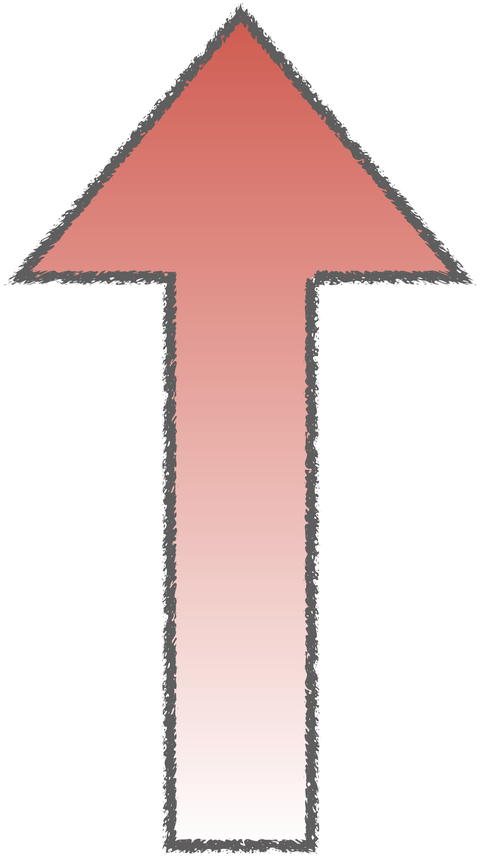
PRECISION CONTROL

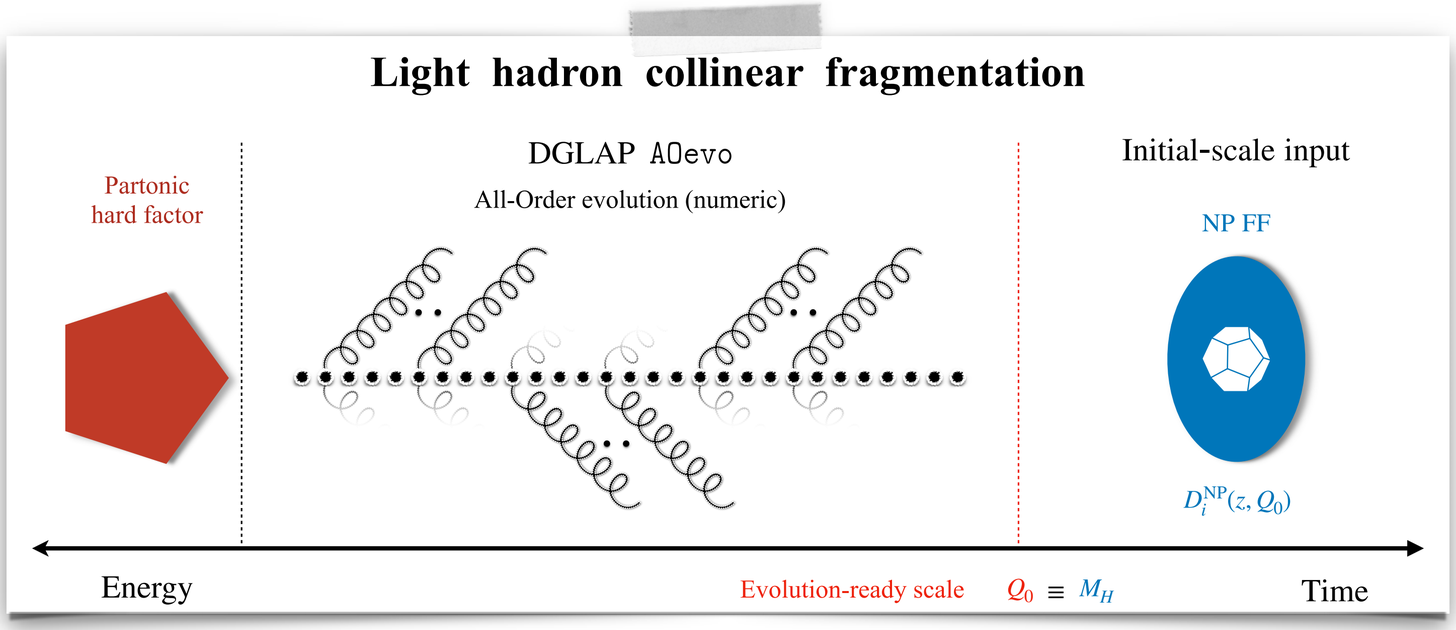


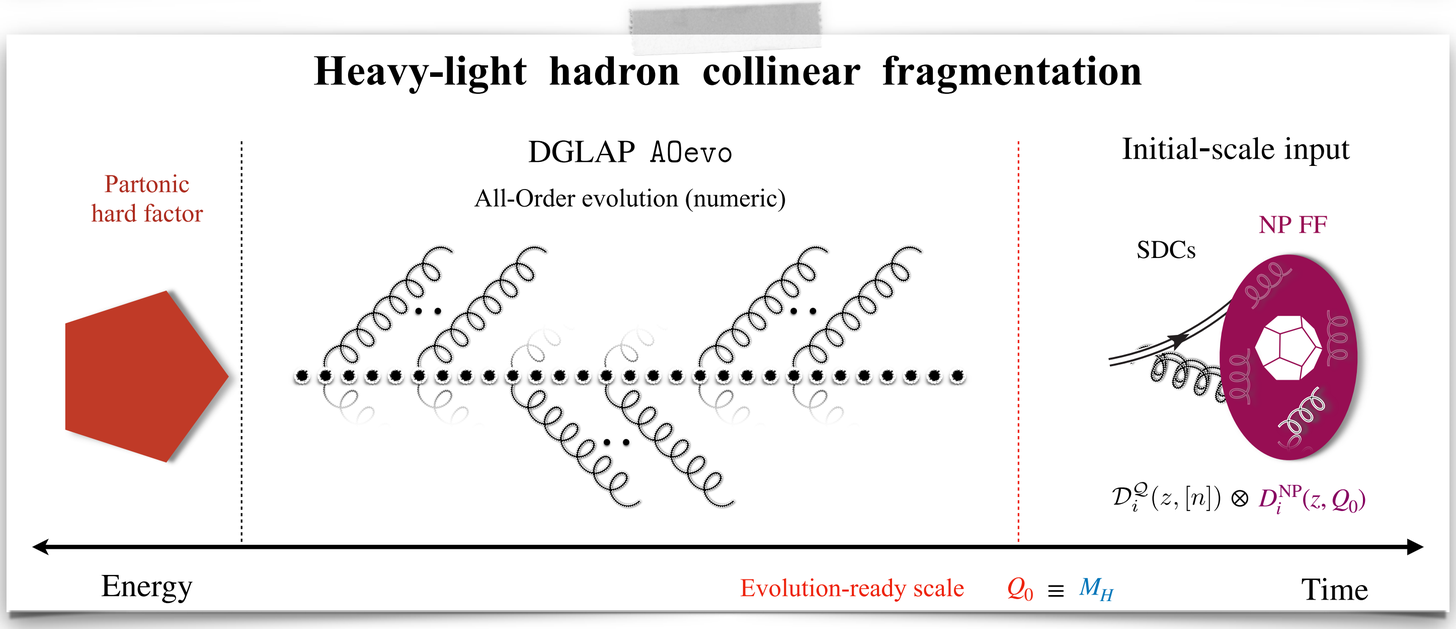
DISCOVERY POTENTIAL

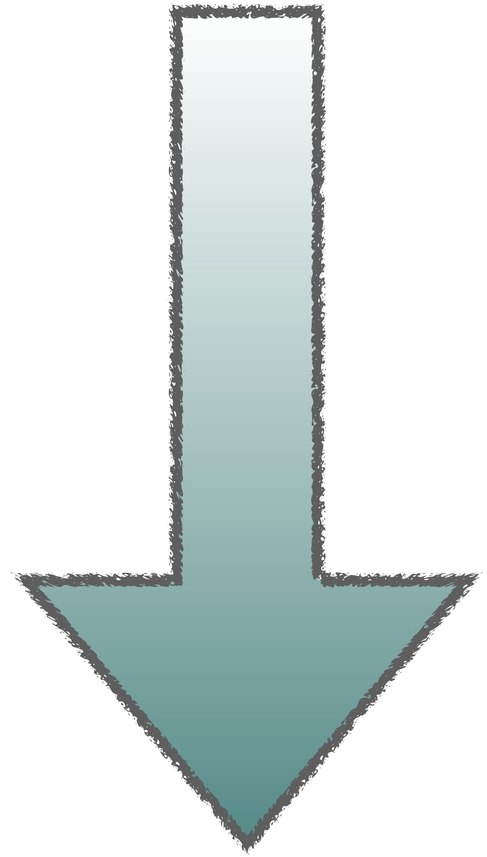
Fragmentation in QCD (with Heavy Flavors)

PRECISION CONTROL









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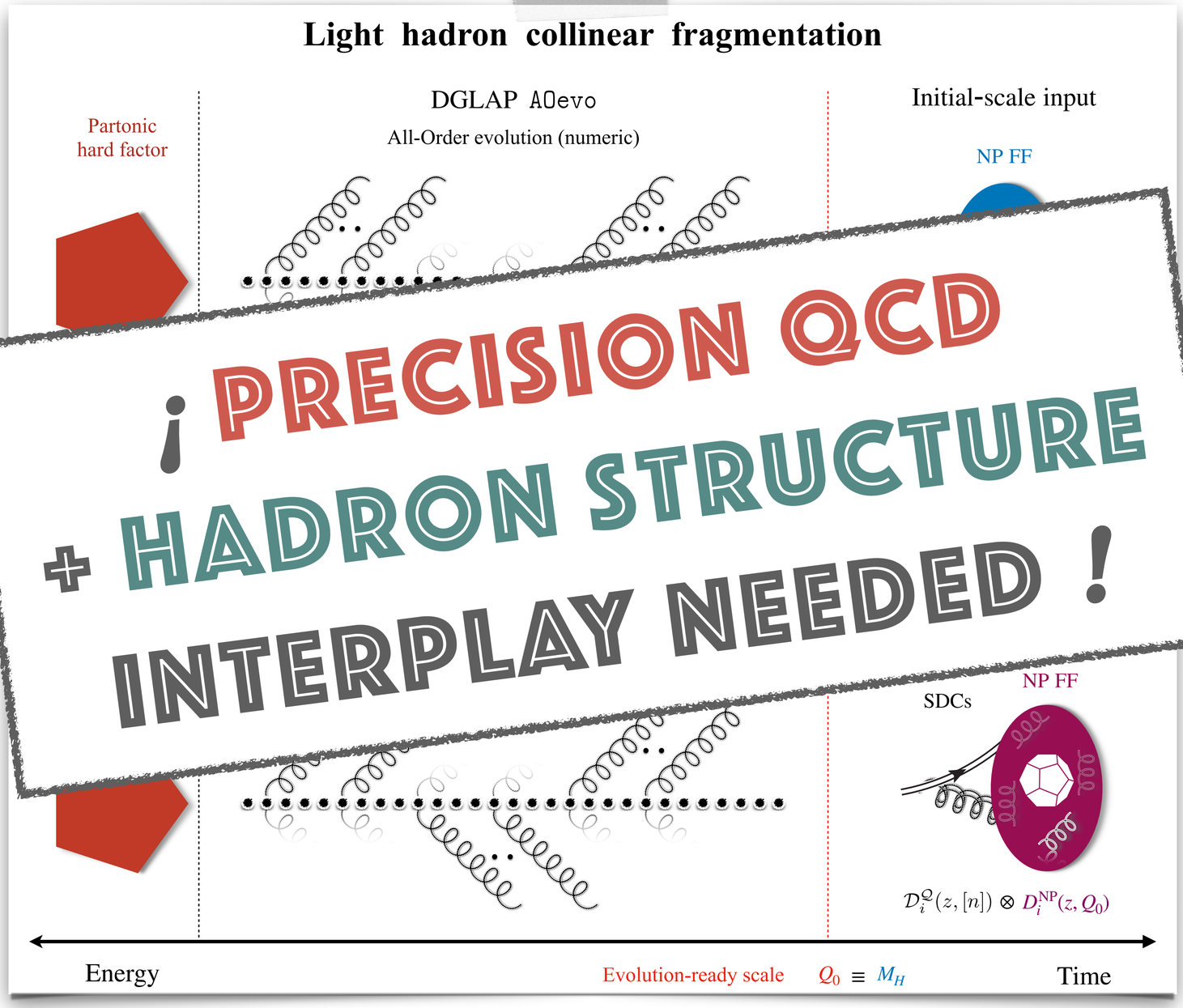
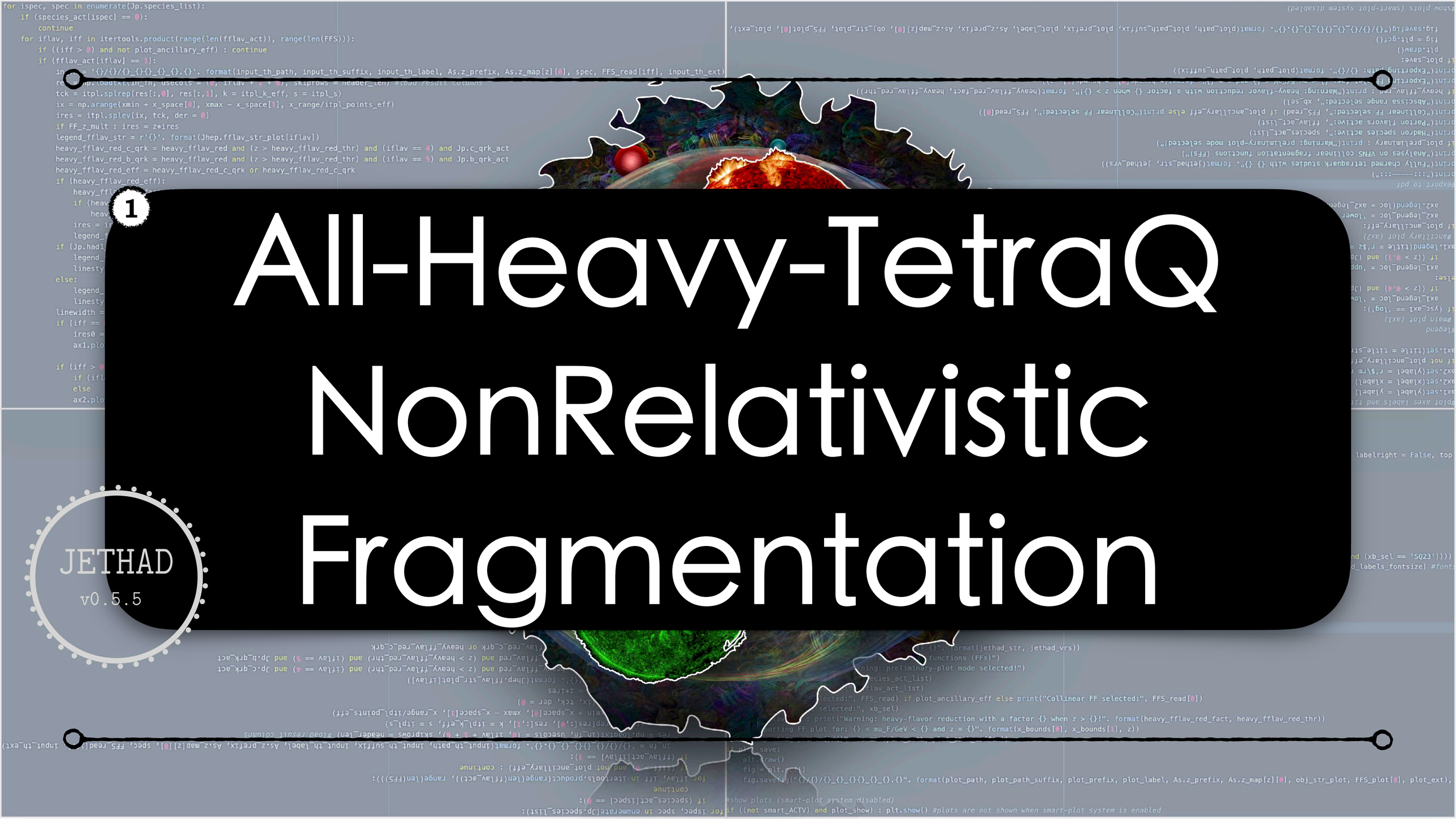


¡ PRECISION QCD + HADRON STRUCTURE INTERPLAY NEEDED !

DISCOVERY POTENTIAL


1
All-Heavy-TetraQ
NonRelativistic
Fragmentation
JETHAD
v0.5.5

Heavy-Flavor Non-Relativistic evolution

HADRONIC STRUCTURE

QUARKONIUM THEORY

Heavy-Flavor Non-Relativistic evolution

HADRONIC STRUCTURE

QUARKONIUM THEORY

PRECISION QCD

COLLINEAR FACTORIZATION

Heavy-Flavor Non-Relativistic evolution

HADRONIC STRUCTURE

QUARKONIUM THEORY

PRECISION QCD

COLLINEAR FACTORIZATION

Guiding principle ⇒ ¡ Use the Best of the Two Worlds as much as we can !

Interpretation

Heavy-Flavor Non-Relativistic evolution

HADRONIC STRUCTURE

QUARKONIUM THEORY

PRECISION QCD

COLLINEAR FACTORIZATION

Guiding principle ⇒ ¡ Use the Best of the Two Worlds as much as we can !

Interpretation

Evolution

Heavy-Flavor Non-Relativistic evolution

HADRONIC STRUCTURE

QUARKONIUM THEORY

PRECISION QCD

COLLINEAR FACTORIZATION

Guiding principle ⇒ ¡ Use the Best of the Two Worlds as much as we can !

Interpretation

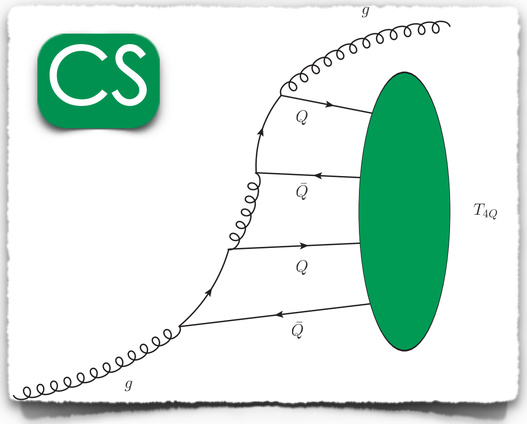
Evolution

Uncertainties

TetraQuark Fragmentation from NRQCD & symJETHAD

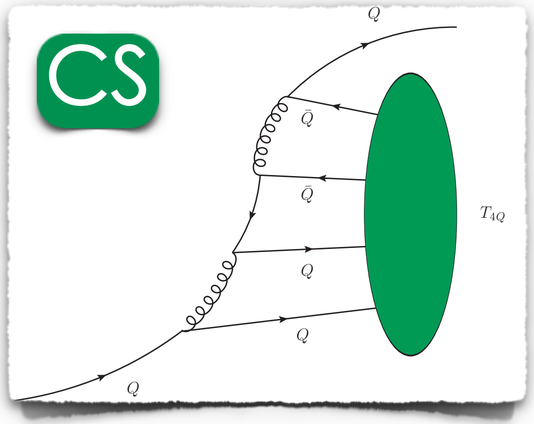
LO

Gluon



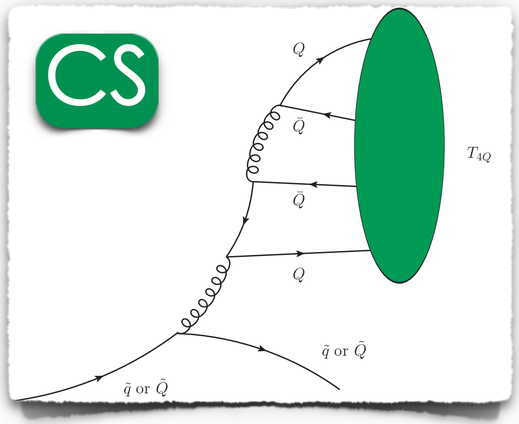
[F. Feng, Y. Huang et al., Phys. Rev. D 106 (2022) 11, 114029]

Constituent heavy quark



[X.-W. Bai, F. Feng et al., JHEP 09 (2024) 002

Non-constituent heavy or light quark



[X.-W. Bai, Y. Huang et al., Phys. Rev. D 111 (2025) 5, 054006]

$J^{PC} = 0^{++}$

$J^{PC} = 1^{+-}$

$J^{PC} = 2^{++}$

TetraQuark Fragmentation from NRQCD & symJETHAD

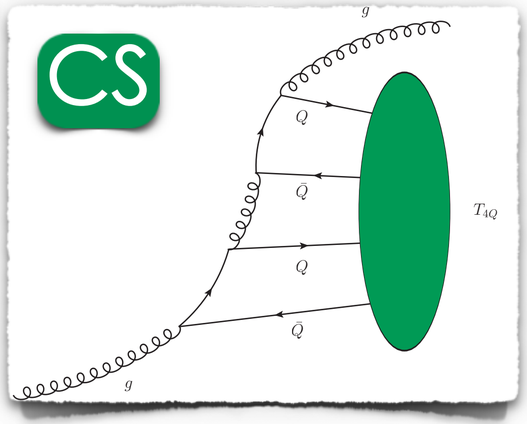
LO

2024

TQ4Q1.0

Gluon

CS



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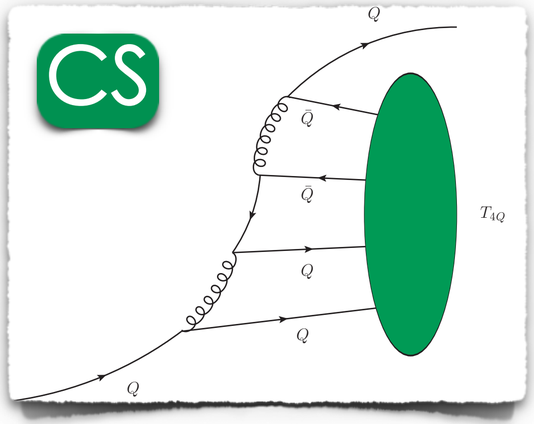

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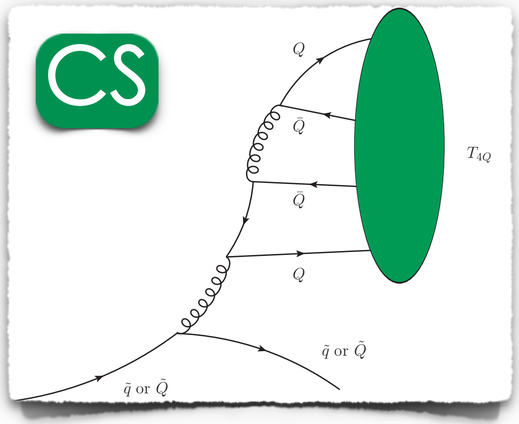

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	Gluon	Constituent heavy quark	Non-constituent heavy or light quark
$J^{PC} = 0^{++}$	TQ4Q1.0+ ✓	¡! Suzuki [TQ4Q1.0]	¡! Not implemented
$J^{PC} = 1^{+-}$	Landau-Yang LO: evolution only ✗	¡! Not implemented	C-parity LO: evolution only ✗
$J^{PC} = 2^{++}$	TQ4Q1.0+ ✓	¡! Suzuki [TQ4Q1.0]	¡! Not implemented

TetraQuark Fragmentation from NRQCD & symJETHAD

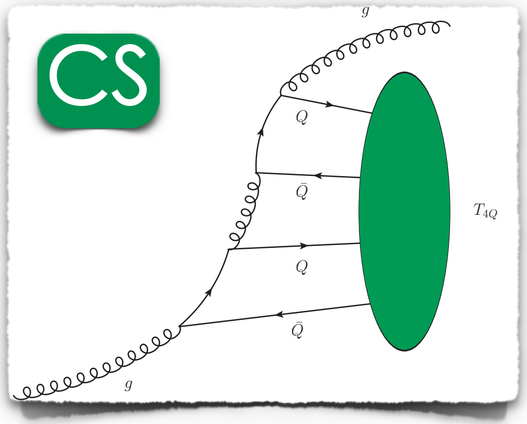
LO

2024 → 2025

TQ4Q1.1

Gluon

CS



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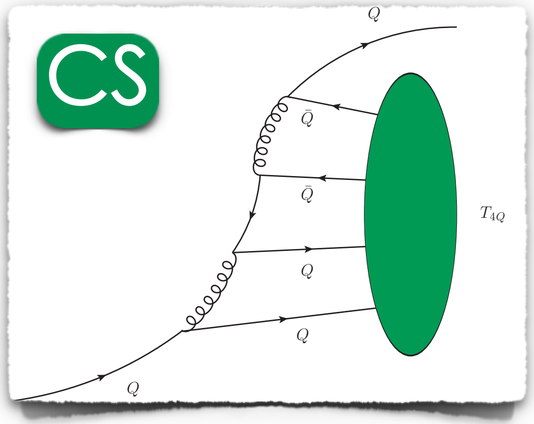
$J^{PC} = 0^{++}$ TQ4Q1.0+ ✓

$J^{PC} = 1^{+-}$ Landau-Yang LO: evolution only ✗

$J^{PC} = 2^{++}$ TQ4Q1.0+ ✓

Constituent heavy quark

CS



[X.-W. Bai, F. Feng et al., JHEP 09 (2024) 002]

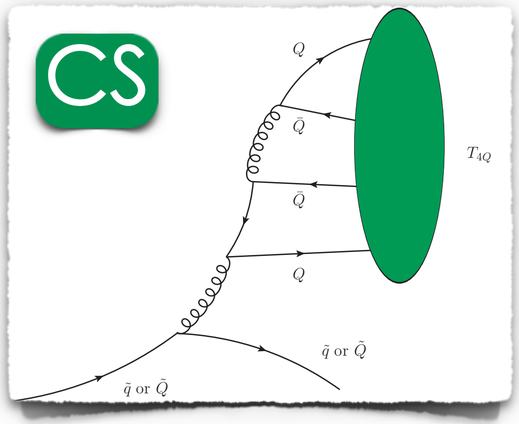
$J^{PC} = 0^{++}$ TQ4Q1.1+ ✓

$J^{PC} = 1^{+-}$ TQ4Q1.1+ ✓

$J^{PC} = 2^{++}$ TQ4Q1.1+ ✓

Non-constituent heavy or light quark

CS



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TetraQuark Fragmentation from NRQCD & symJETHAD

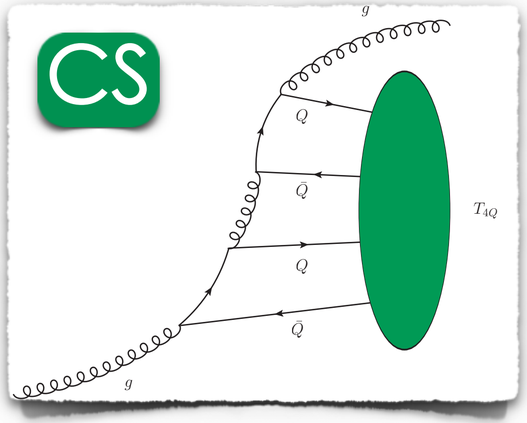
LO

2024 → 2025 → 2026

NEW

TQ4Q2.0

Gluon



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$J^{PC} = 0^{++}$ TQ4Q1.0+ ✓

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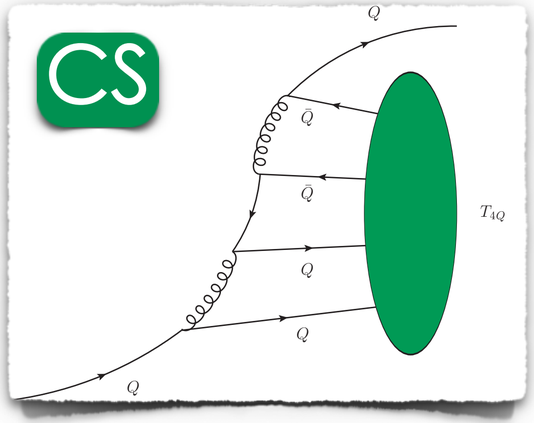
Constituent heavy quark



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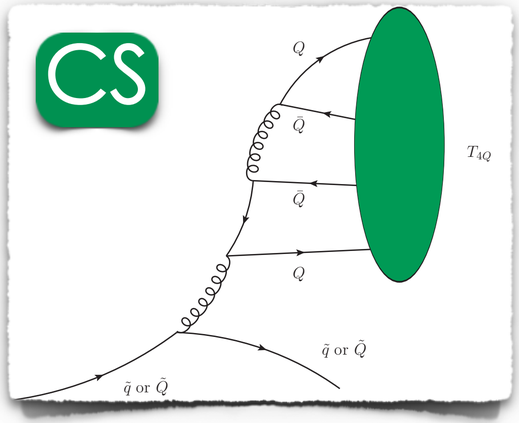
TQ4Q1.1+ ✓

TQ4Q1.1+ ✓

TQ4Q1.1+ ✓

Non-constituent heavy or light quark



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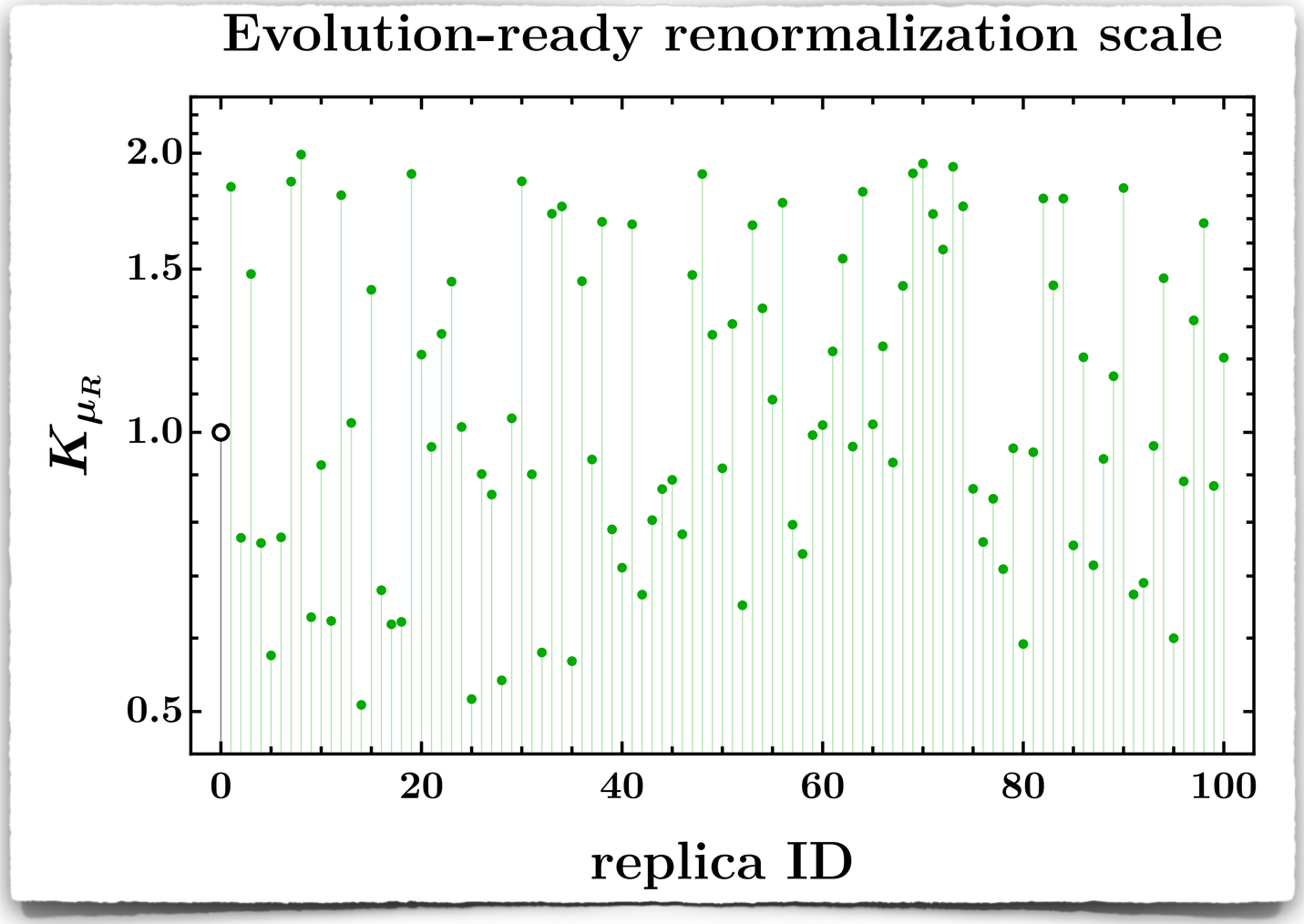
NEW TQ4Q2.0+ ✓

C-parity LO: evolution only ✗

NEW TQ4Q2.0+ ✓

TQ4Q2.0: F-MHOU Replica Analysis

T_{4Q}



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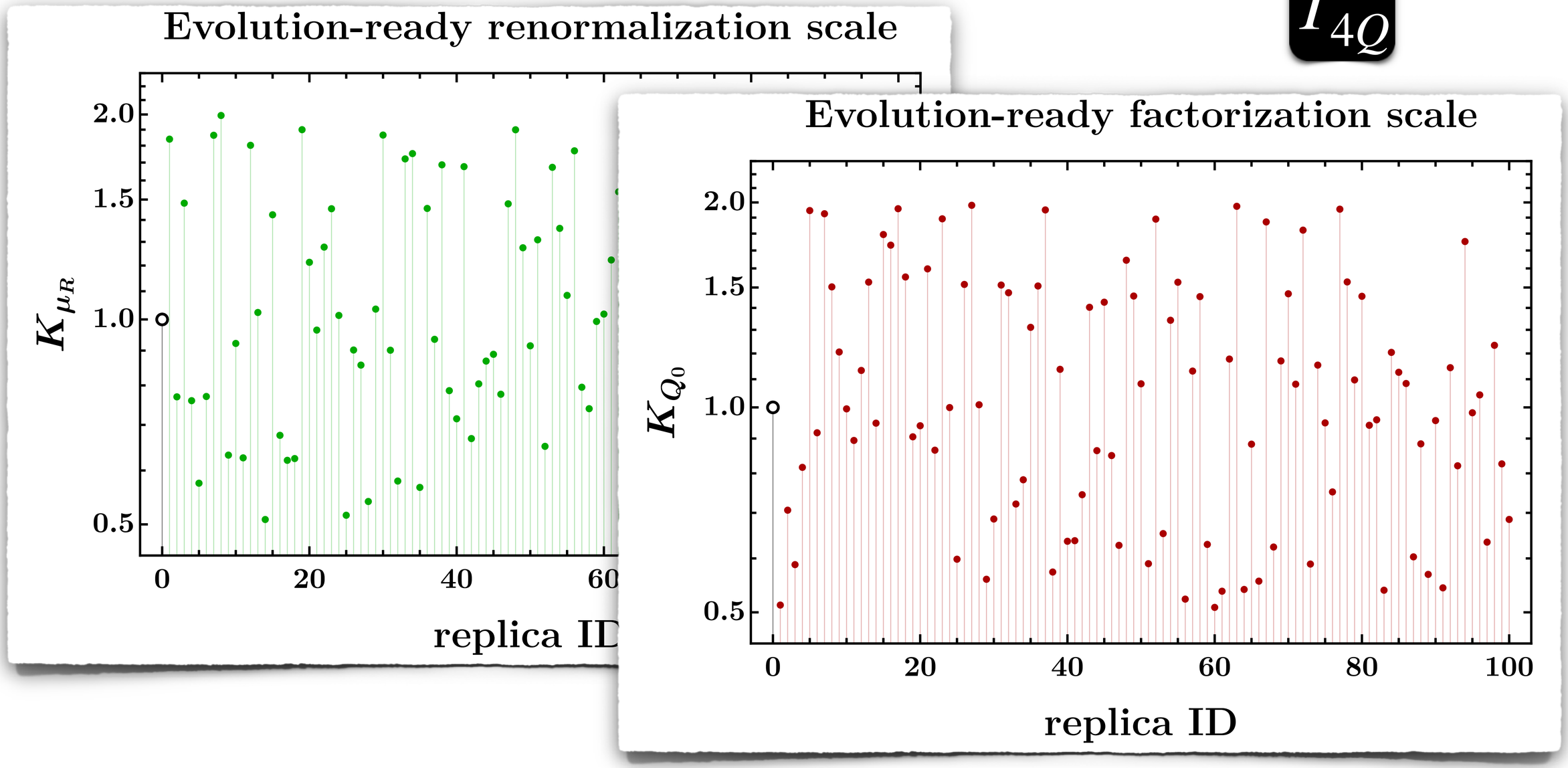
TQ4Q2.0: F-MHOU Replica Analysis

T_{4Q}



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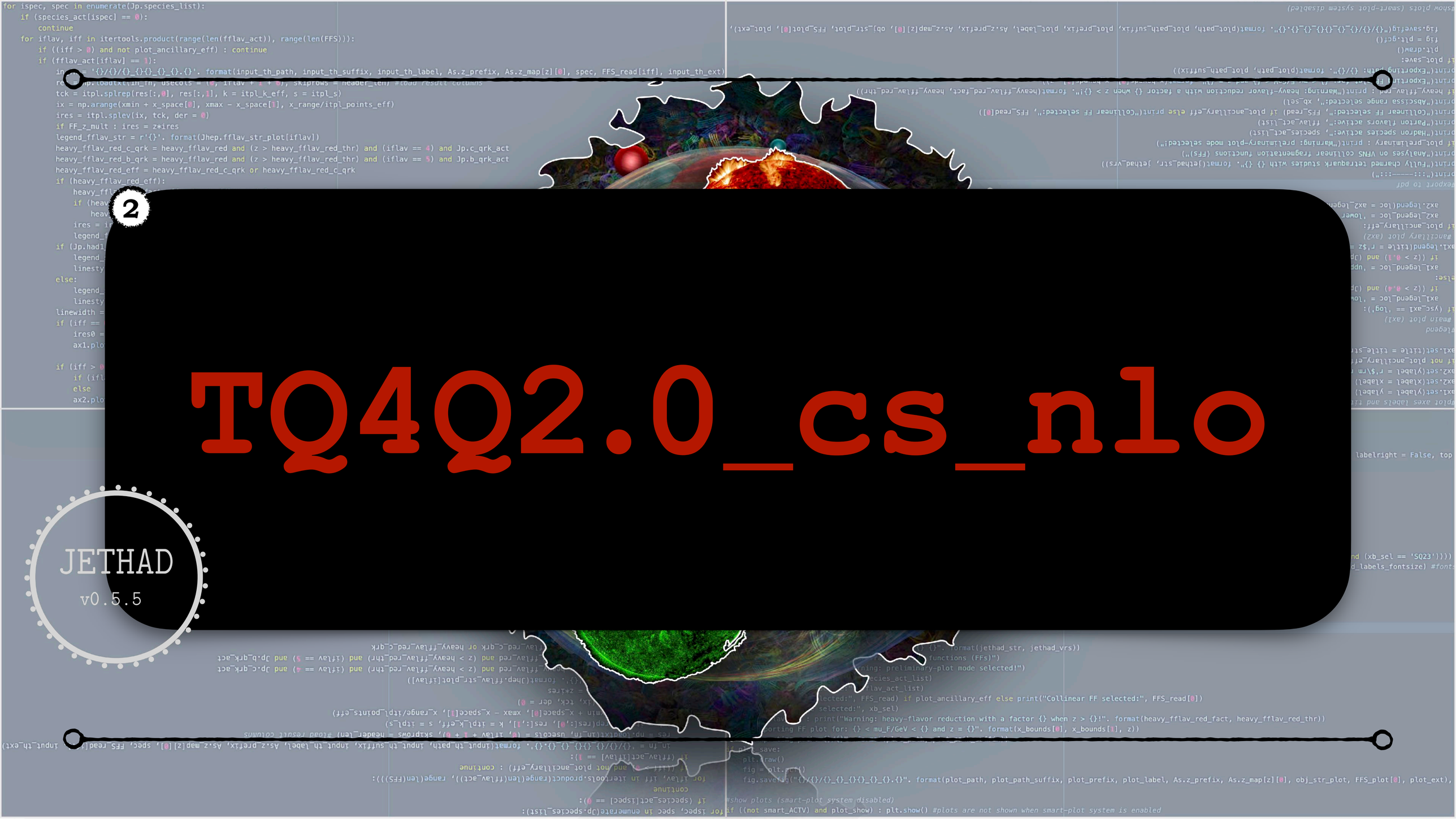
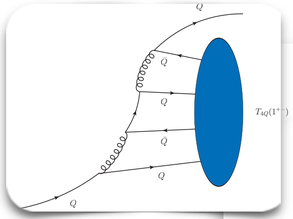
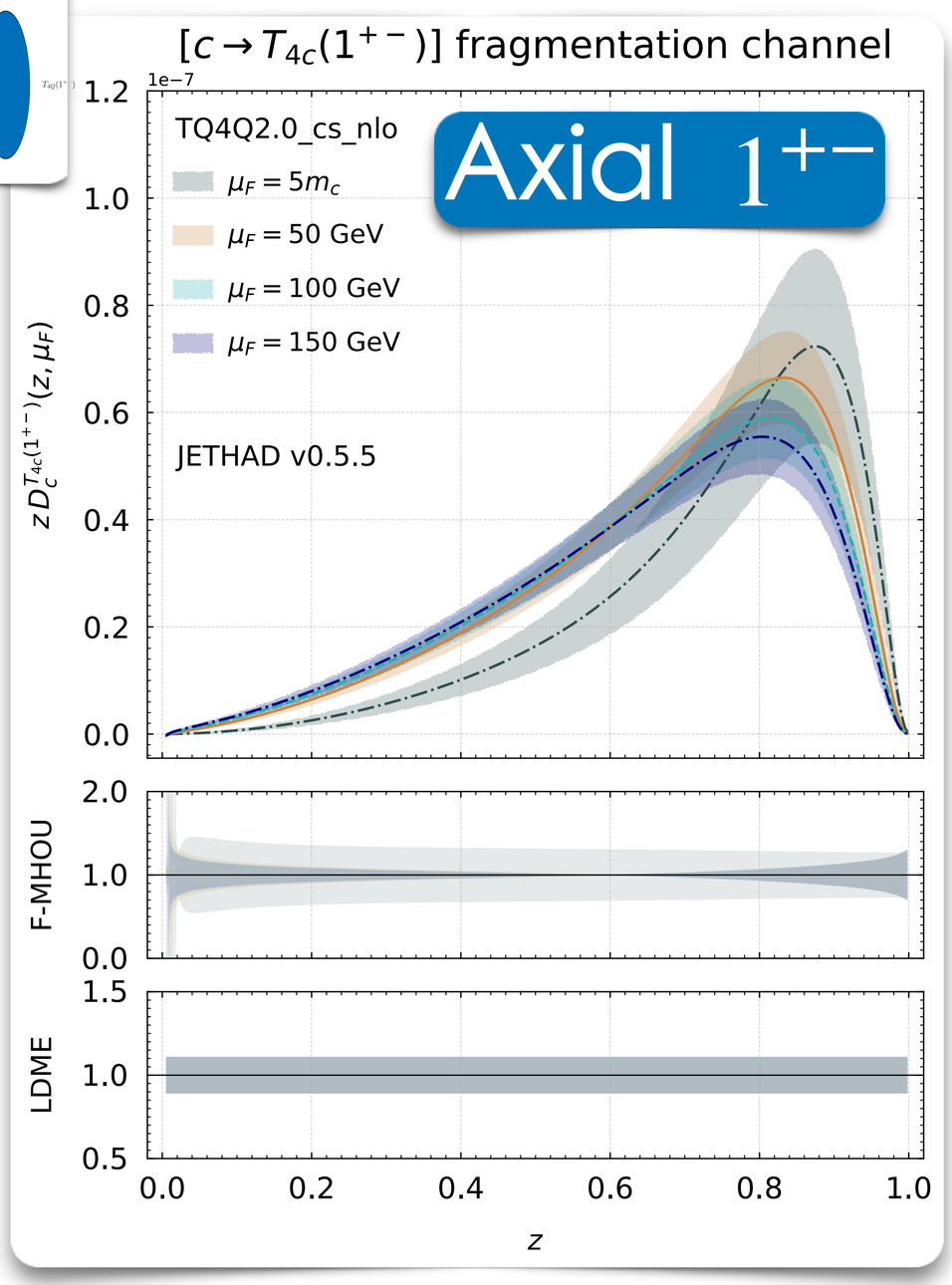
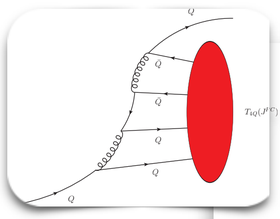
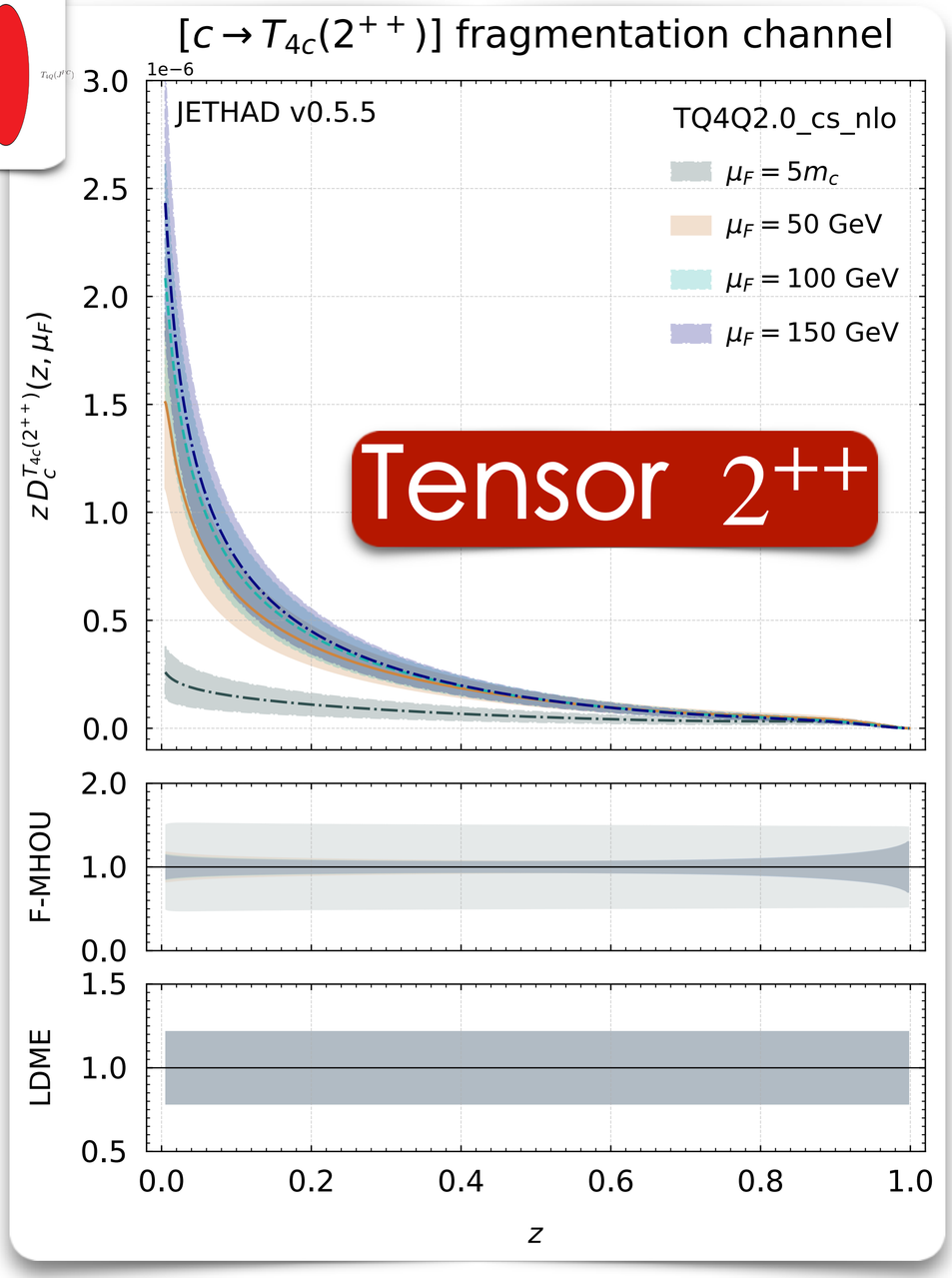

2
TQ4Q2.0_cs_nlo
JETHAD
v0.5.5

TQ4Q2.0: Charm Fragmentation to Charmoniumlike

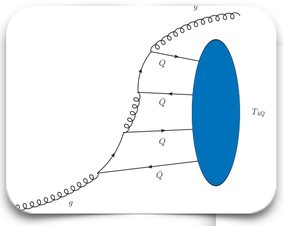


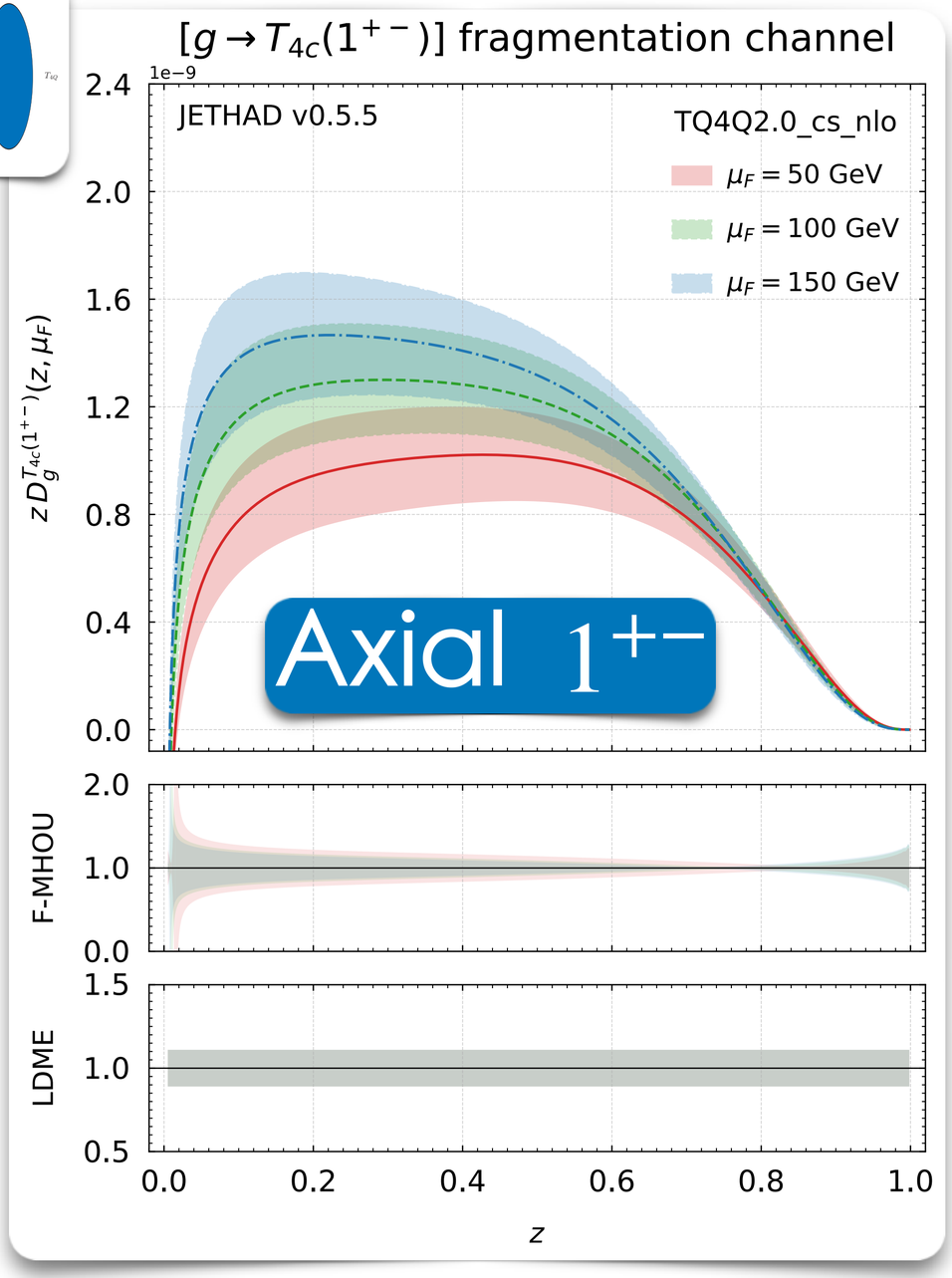


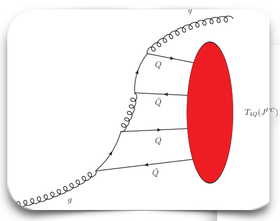


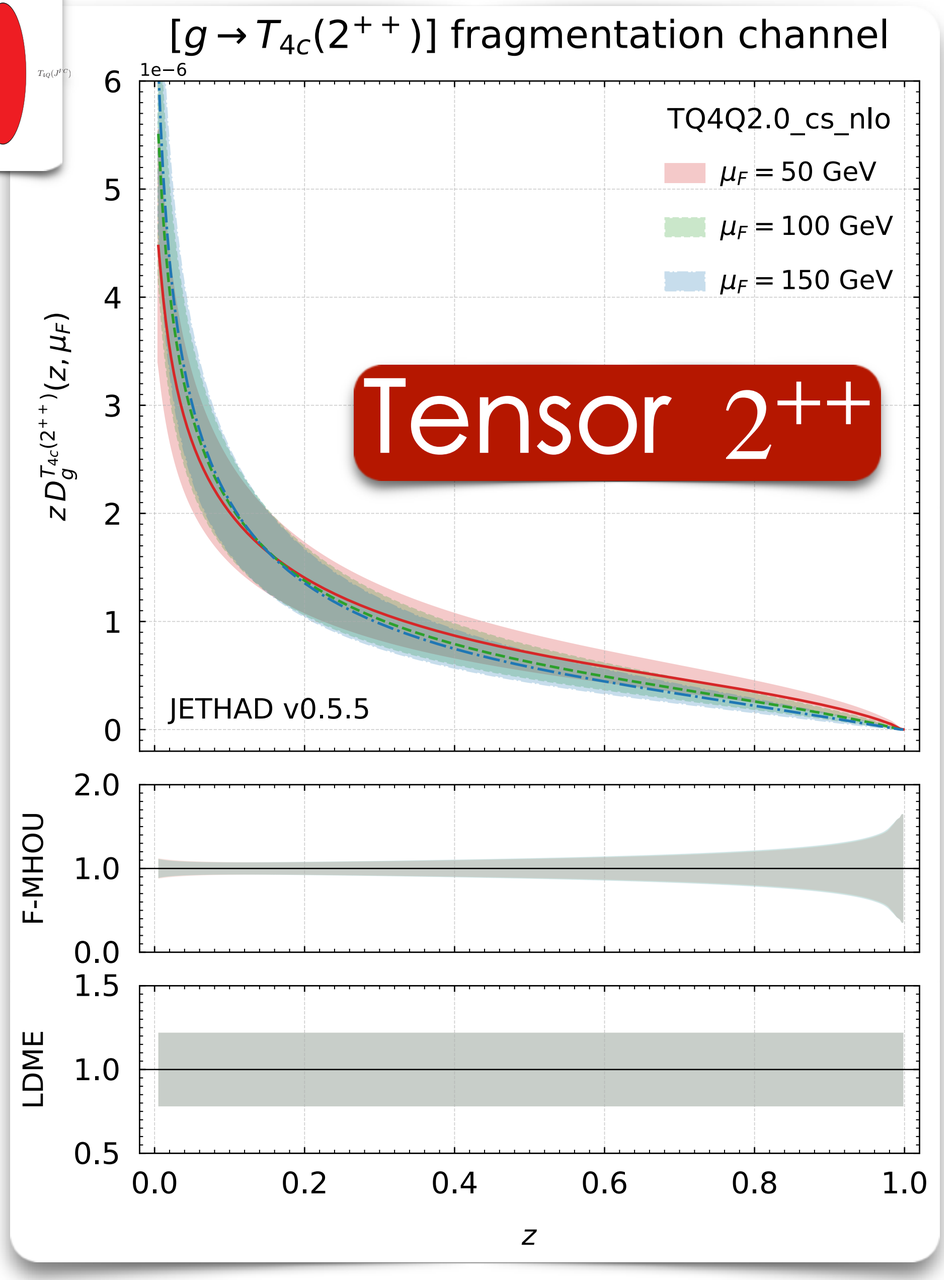


TQ4Q2.0: Gluon Fragmentation to Charmoniumlike

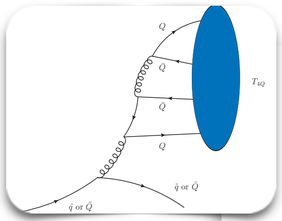


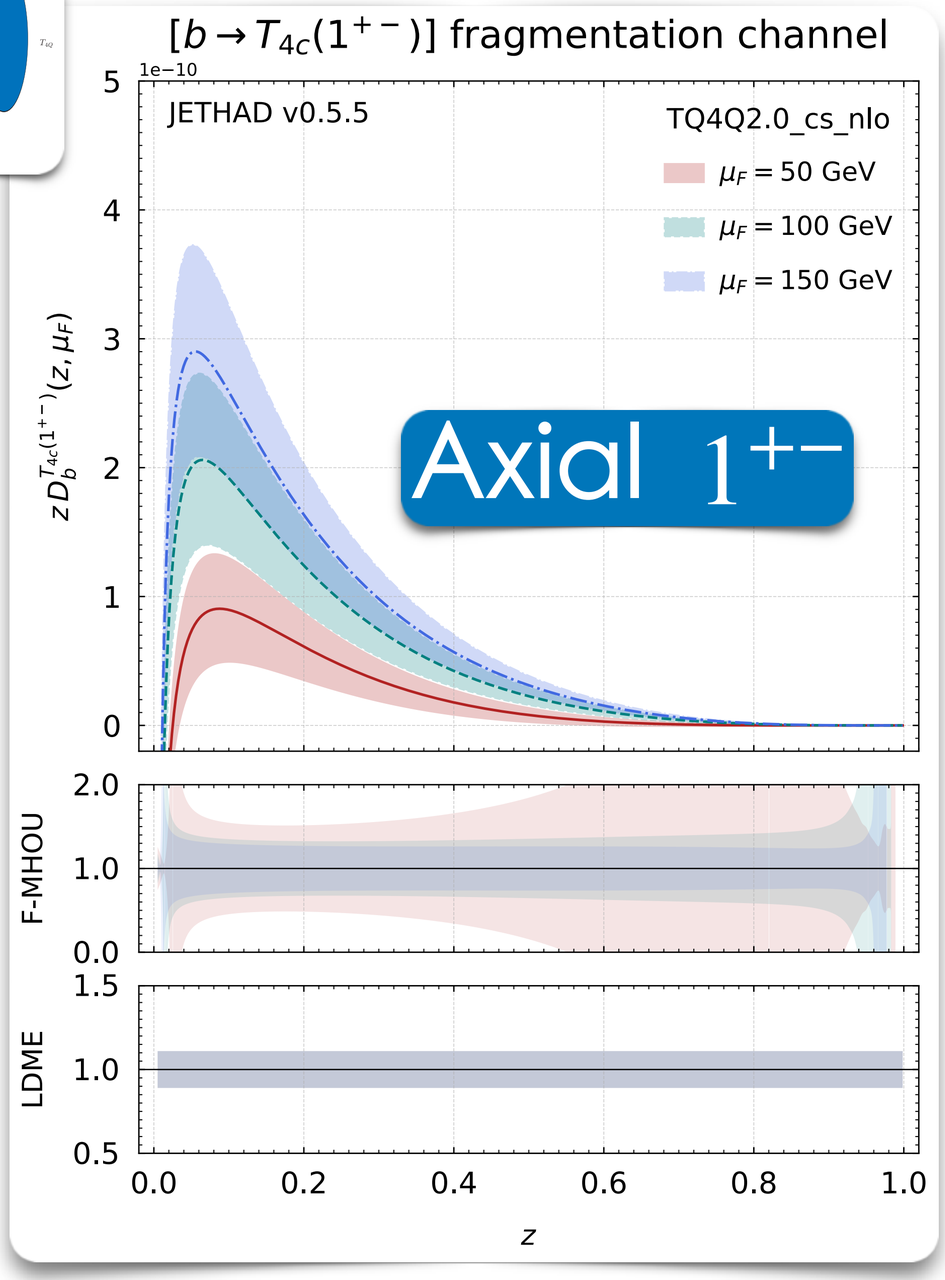


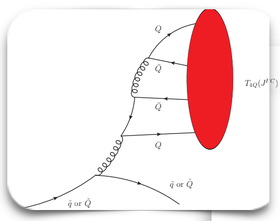


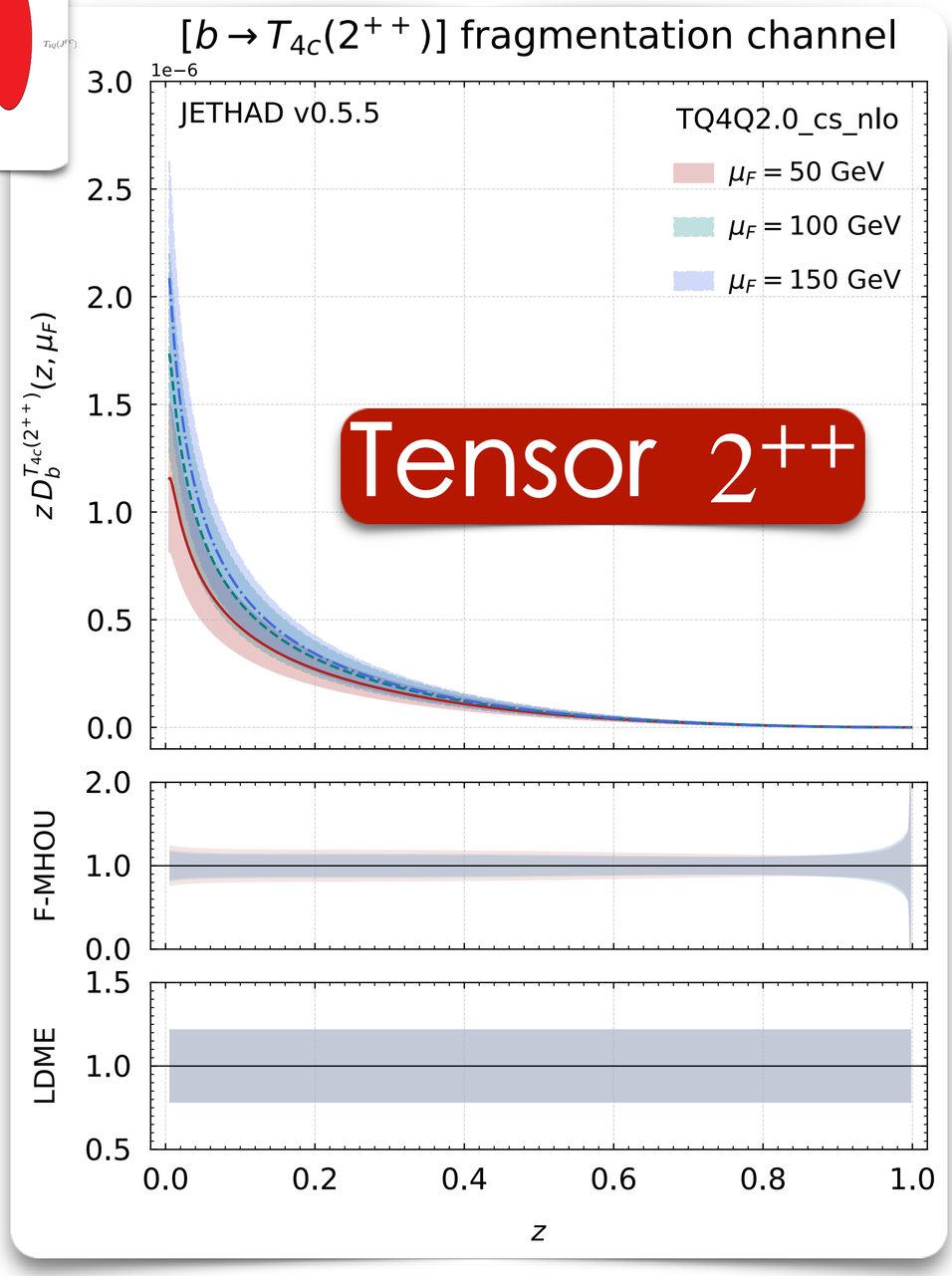


TQ4Q2.0: Bottom Fragmentation to Charmoniumlike

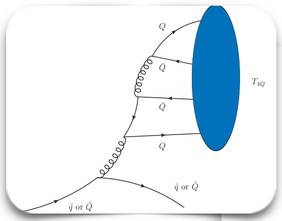


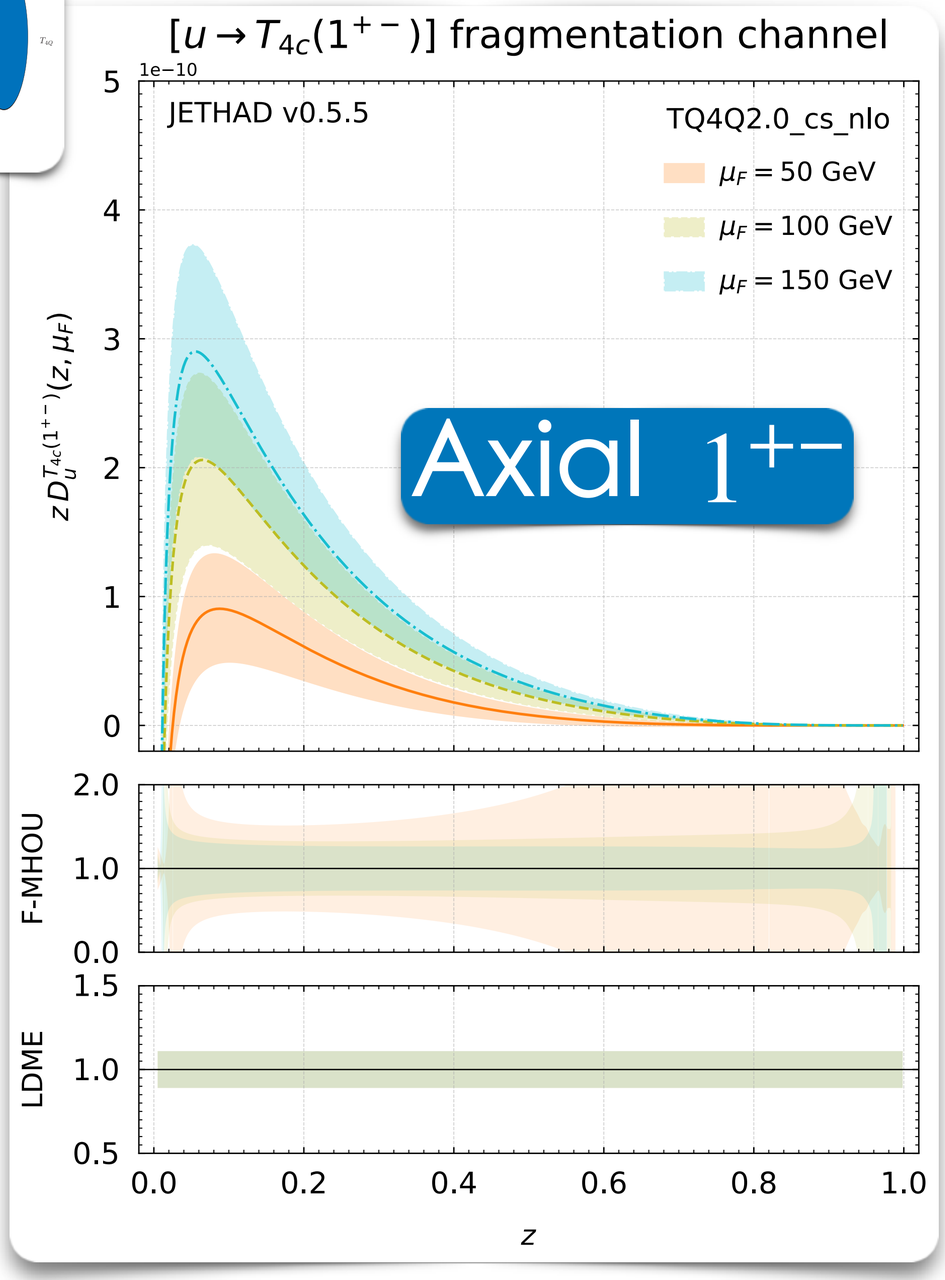


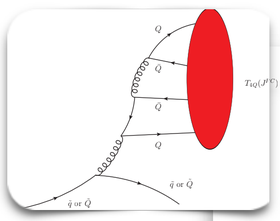


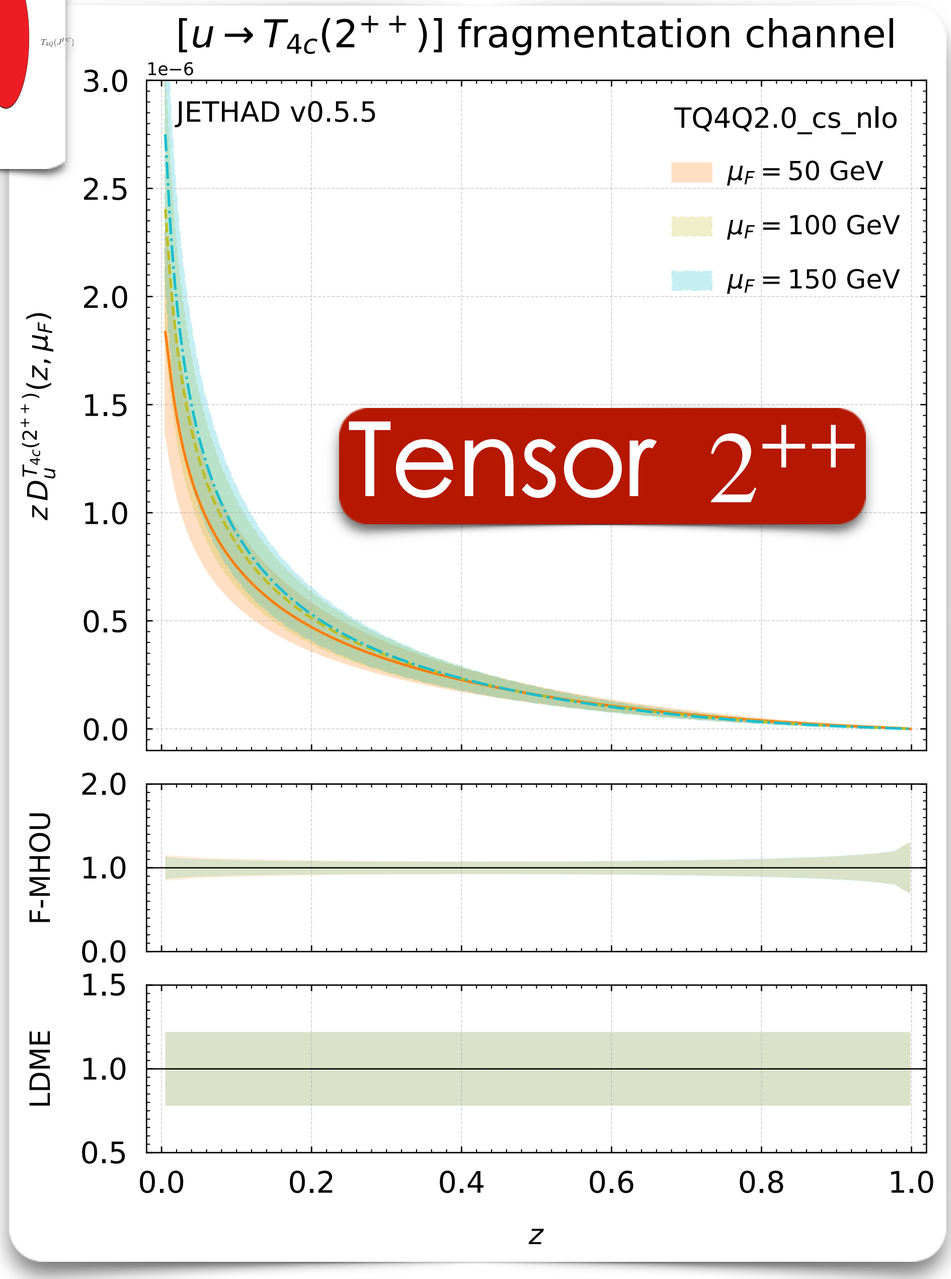


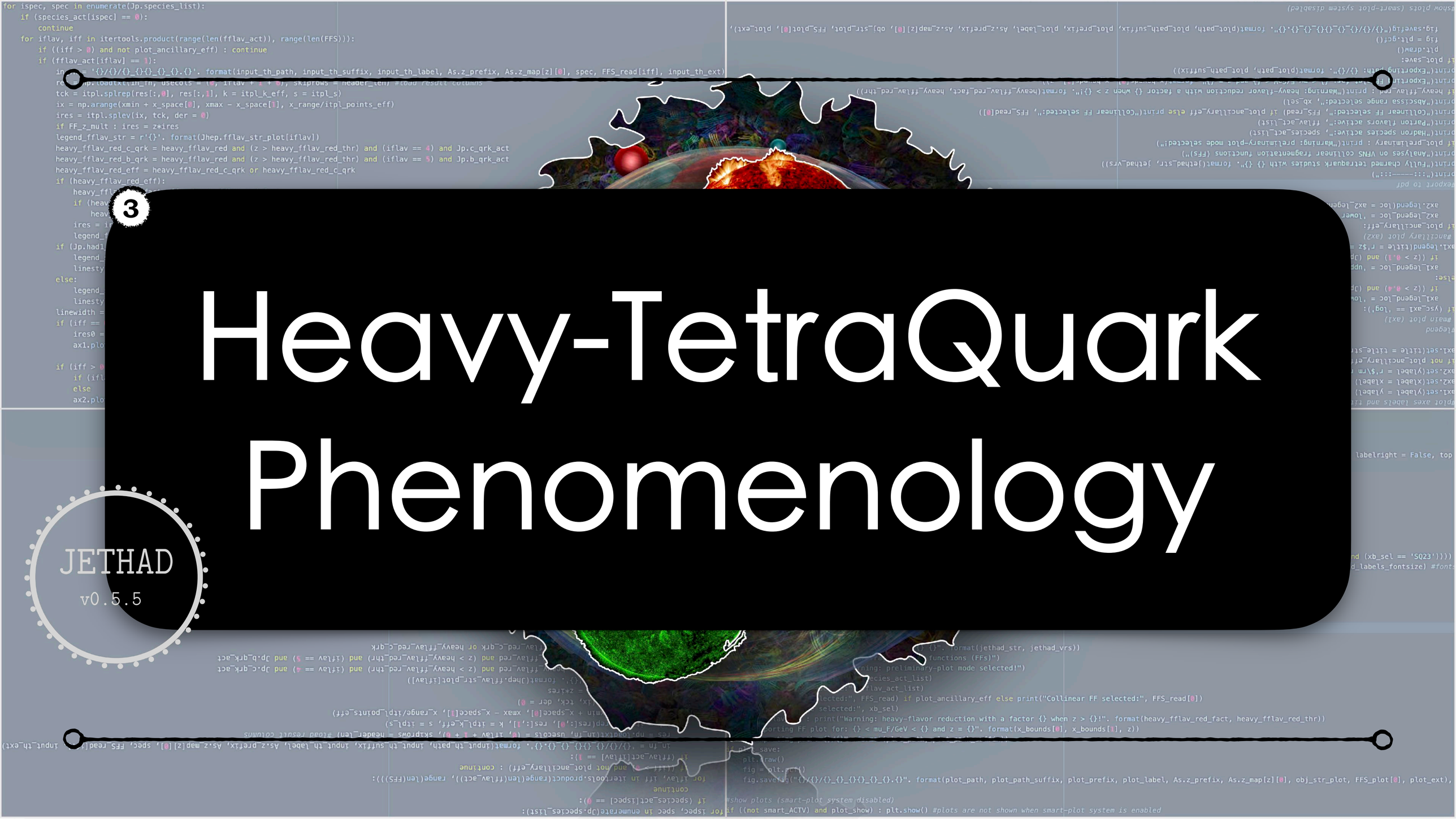
TQ4Q2.0: Light-Quark Fragmentation to Charmoniumlike










3
Heavy-TetraQuark
Phenomenology
JETHAD
v0.5.5

TetraQuark-Jet Systems with Resummation: Event Yields

Expected events yields [13 TeV Run 2]

T_{4Q}	J^{PC}	TQ4Q2.0	TQ4Q1.1	Δ (%)
T_{4c}	0^{++}	2016031 ± 399168	1768764 ± 309078	+13.98%
T_{4c}	1^{+-}	8619 ± 1770	8619 ± 1612	+0%
T_{4c}	2^{++}	3071882 ± 602910	2721362 ± 480942	+12.88%
T_{4b}	0^{++}	6046 ± 1168	5151 ± 945	+17.38%
T_{4b}	1^{+-}	17 ± 3	17 ± 2	+0%
T_{4b}	2^{++}	8095 ± 1552	7108 ± 828	+13.88%

TABLE IV: Expected event yields for all-heavy tetraquark production at $\sqrt{s} = 13$ TeV, obtained by integrating the NLL-resummed ΔY distributions over the rapidity range $|\Delta Y| < 6.5$. The yields assume an integrated luminosity of $\mathcal{L}^{(\mathrm{CMS})} = 138.6\ \mathrm{fb}^{-1}$, corresponding to the total dataset collected by CMS during Run 2 [308, 309]. The last column shows the relative variation $\Delta = (R - 1) \times 100$, with $R =$ TQ4Q2.0/TQ4Q1.1. Uncertainties are propagated from the ΔY distributions and reflect the combined effect of scale variations, FF evolution, and LDME inputs.

Expected events yields [100 TeV FCC]

T_{4Q}	J^{PC}	TQ4Q2.0	TQ4Q1.1	Δ (%)
T_{4c}	0^{++}	22140238 ± 4433814	19465674 ± 3373524	+13.74%
T_{4c}	1^{+-}	23463 ± 1094	23463 ± 1053	+0%
T_{4c}	2^{++}	32784967 ± 6421338	29144313 ± 5467493	+12.49%
T_{4b}	0^{++}	23424 ± 1982	20237 ± 1552	+15.75%
T_{4b}	1^{+-}	50 ± 2	50 ± 3	+0%
T_{4b}	2^{++}	27894 ± 1972	24493 ± 1323	+13.89%

TABLE V: Expected event yields for all-heavy tetraquark production at $\sqrt{s} = 100$ TeV, obtained by integrating the NLL-resummed ΔY distributions over the rapidity range $|\Delta Y| < 6.5$. The yields assume an integrated luminosity of $\mathcal{L}^{(\mathrm{CMS})} = 138.6\ \mathrm{fb}^{-1}$, corresponding to the full Run 2 CMS dataset [308, 309]. This value is conservatively adopted at $\sqrt{s} = 100$ TeV to enable a baseline comparison of energy scaling, independently of specific FCC projections. The last column shows the relative variation $\Delta = (R - 1) \times 100$, with $R =$ TQ4Q2.0/TQ4Q1.1. Uncertainties are propagated from the ΔY distributions and reflect the combined effect of scale variations, FF evolution, and LDME inputs.

JETHAD v0.5.6

Pile-up effects can degrade Experimental discrimination

TetraQuark-Jet Systems with Resummation: Event Yields

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TABLE IV: Expected event yields for a… tetraquark production at $\sqrt{s} = 13$ TeV, obta… integrating the NLL-resummed ΔY distributio… the rapidity range $|\Delta Y| < 6.5$. The yields assume an integrated luminosity of $\mathcal{L}^{(\mathrm{CMS})} = 138.6\ \mathrm{fb}^{-1}$, corresponding to the total dataset collected by CMS during Run 2 [308, 309]. The last column shows the relative variation $\Delta = (R - 1) \times 100$, with $R =$ TQ4Q2.0/TQ4Q1.1. Uncertainties are propagated from the ΔY distributions and reflect the combined effect of scale variations, FF evolution, and LDME inputs.

Δ(%)

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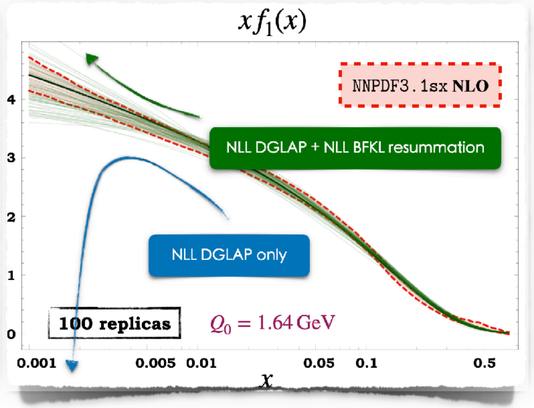
Δ(%)

JETHAD v0.5.6

Pile-up effects can degrade Experimental discrimination

Quakonium Studies: Challenges

Forward Onia ⇒ Gluon PDF & TMD Positivity



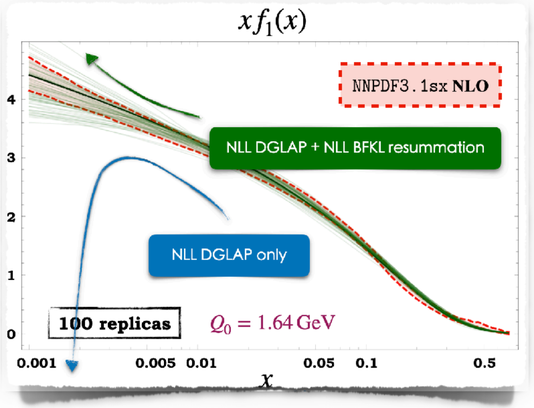
Low x and low Q^2

- [G. Altarelli et al. (1998)]
- [A. Candido et al. (2020)]
- [J. Collins et al. (2022)]
- [A. Candido et al. (2023)]

All-Charm TetraQuarks & Intrinsic Charm

Quakonium Studies: Challenges

Forward Onia ⇒ Gluon PDF & TMD Positivity



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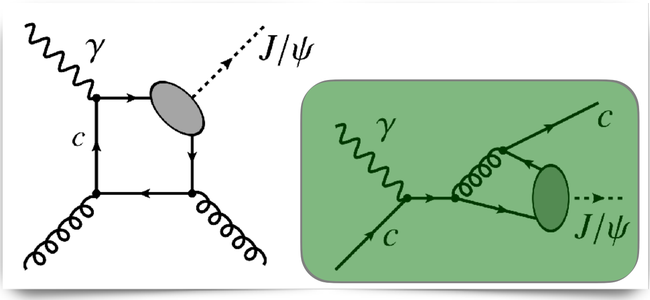
$\gamma + g \rightarrow J/\psi + c$ ⇒ Valence Intrinsic Charm @EIC

Large x
Moderate to large p_T

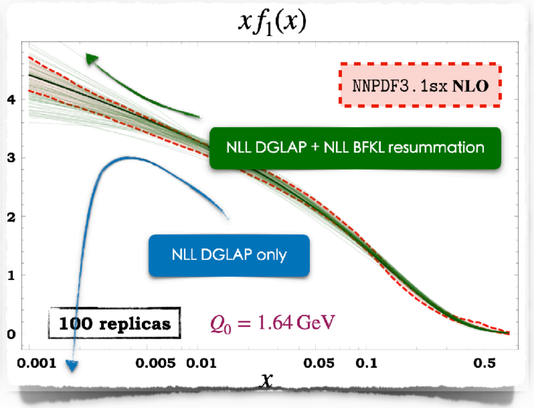


(J/ψ @EIC) 🔗 [C. Flore, J.-P. Lansberg, H.-S. Shao, Y. Yedelkina (2020)]
(Intrinsic Charm + valence) studies 🔗🔗 [NNPDF Collaboration (2022, 2023)]
🔗🔗 [CT Collaboration (2018, 2023)]

All-Charm TetraQuarks & Intrinsic Charm

Quakonium Studies: Challenges

Forward Onia ⇒ Gluon PDF & TMD Positivity



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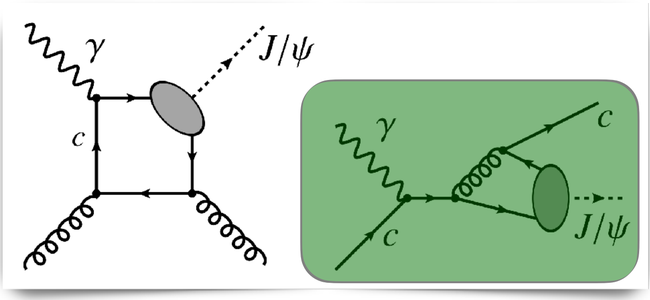
$\gamma + g \to J/\psi + c$ ⇒ Valence Intrinsic Charm @EIC

Large x
Moderate to large p_T



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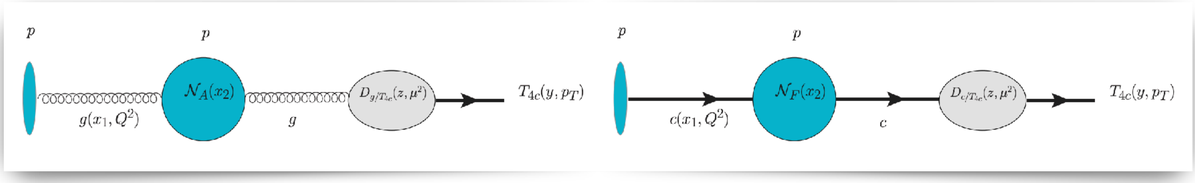
TetraQuark Explorations: Assets

$\{T_{4c}\} \ni X(6900)$ ⇒ Charm PDF @HL-LHC, FCC

Forward 0^{++}, 1^{+-}, 2^{++} within the CGC

🔗 [F.G. C, A.V. Giannini, V.P. Gonçalves, Y.N. Lima (2025)]

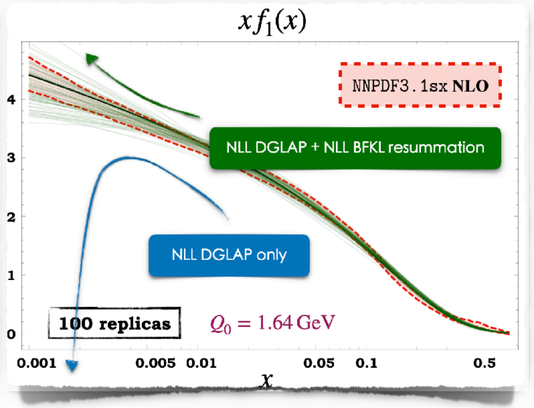


$$\sigma(pp \to T_{4c}X) \propto g(x_1, Q^2) \otimes \mathcal{N}_A(x_2) \otimes D_{g/T_{4c}} + c(x_1, Q^2) \otimes \mathcal{N}_F(x_2) \otimes D_{c/T_{4c}}$$

All-Charm TetraQuarks & Intrinsic Charm

Quakonium Studies: Challenges

Forward Onia ⇒ Gluon PDF & TMD Positivity



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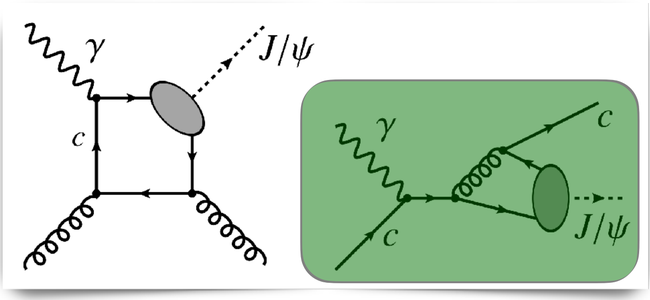
$\gamma + g \to J/\psi + c$ ⇒ Valence Intrinsic Charm @EIC

Large x
Moderate to large p_T



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🔗🔗 [CT Collaboration (2018, 2023)]

TetraQuark Explorations: Assets

$\{T_{4c}\} \ni X(6900)$ ⇒ Charm PDF @HL-LHC, FCC

Forward 0^{++}, 1^{+-}, 2^{++} within the CGC

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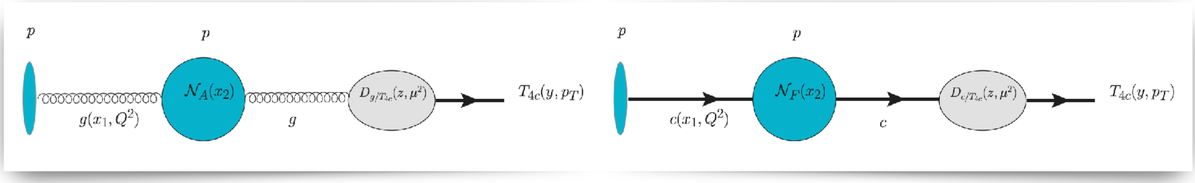
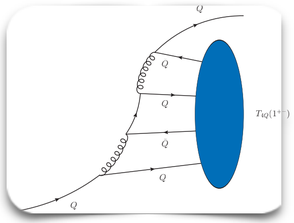


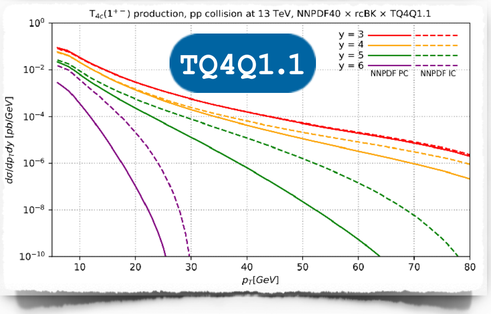
$$\sigma(pp \to T_{4c}X) \propto g(x_1, Q^2) \otimes \mathcal{N}_A(x_2) \otimes D_{g/T_{4c}} + c(x_1, Q^2) \otimes \mathcal{N}_F(x_2) \otimes D_{c/T_{4c}}$$

Forward $T_{4c}(1^{+-})$ ⇒ Intrinsic Charm @pp, pA



1^{+-} FFs

dilute-dense



A Multimodal Perspective

Final Remarks

JETHAD
v0.5.5

A Portal to Exotic Matter

Exotic Matter State-Driven MultiModal Fragmentation

Quantum Superposition



$+ e^{-i\vartheta_1 \Delta t}$



$+ e^{-i\vartheta_2 \Delta t}$



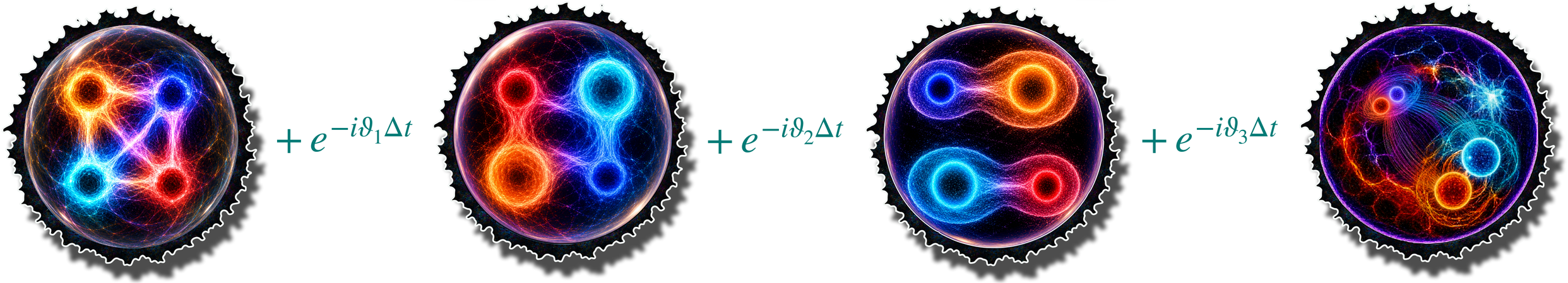
$+ e^{-i\vartheta_3 \Delta t}$



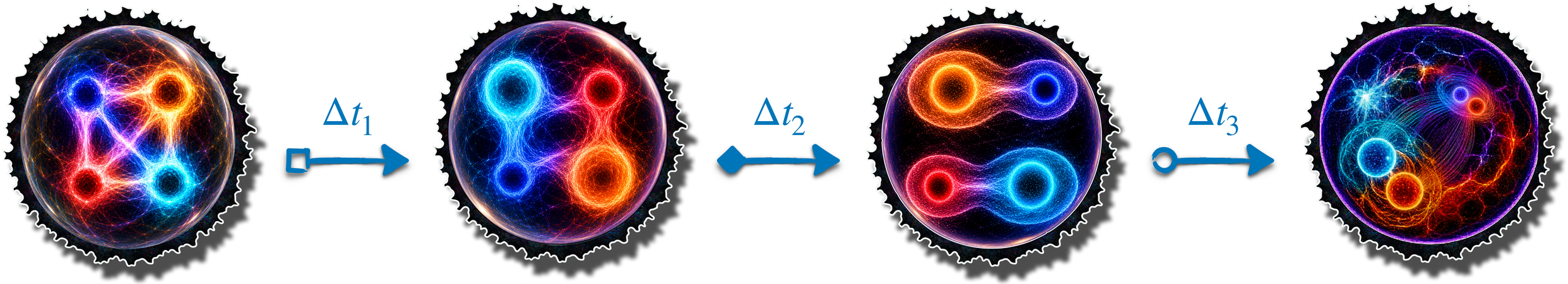
Compact | Diquark | Molecule | Hadro-onium

Exotic Matter State-Driven MultiModal Fragmentation

Quantum Superposition



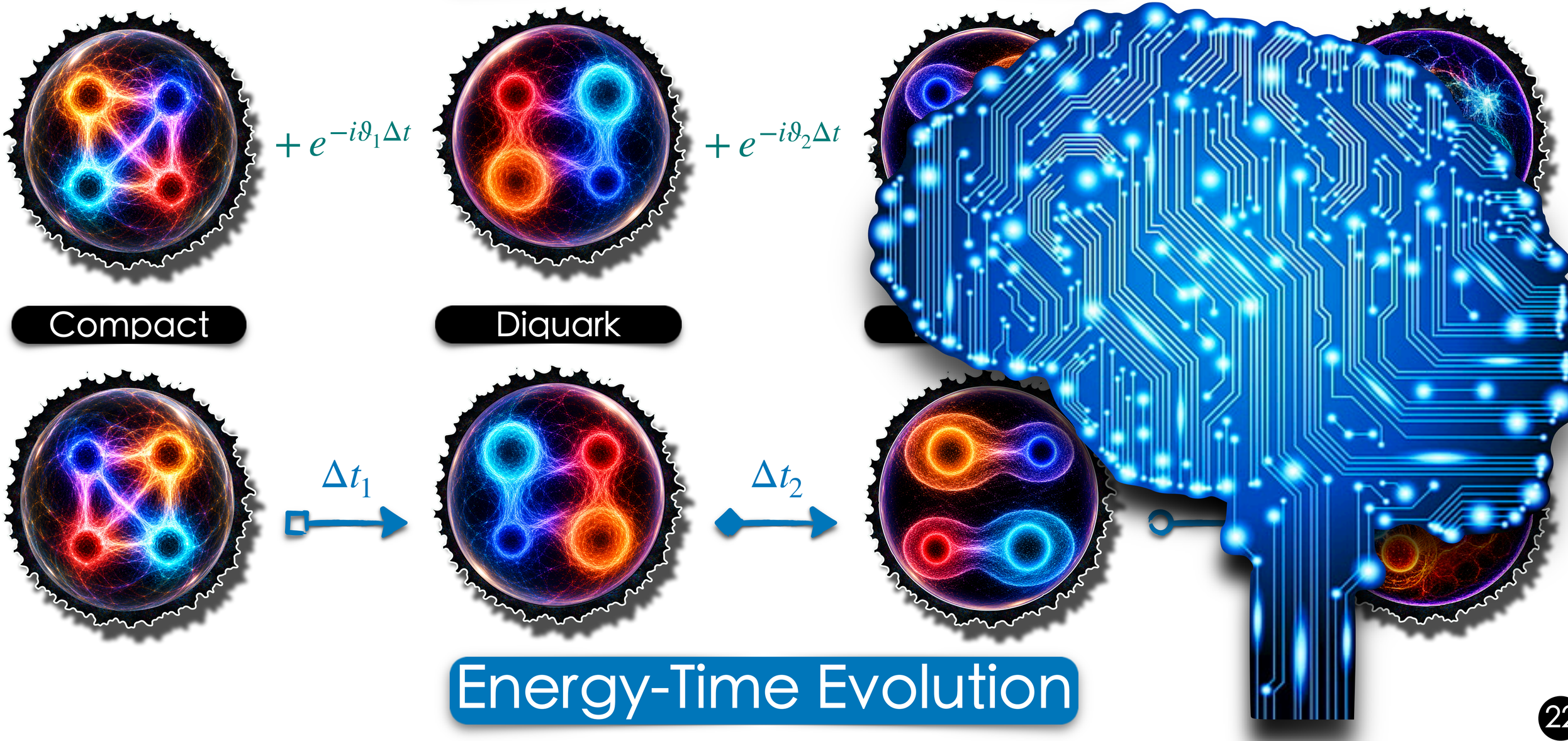
Compact | Diquark | Molecule | Hadro-onium



Energy-Time Evolution

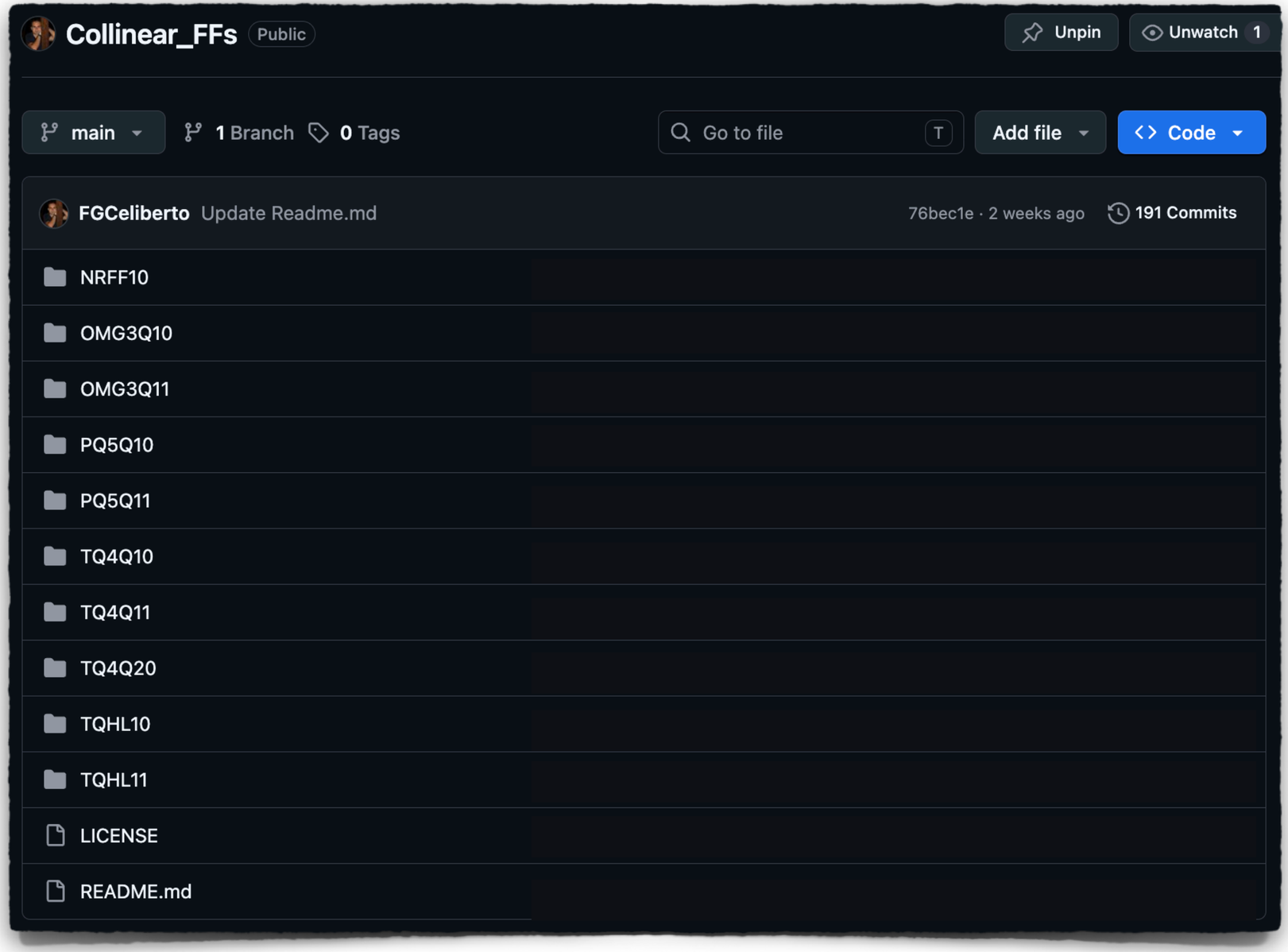
Exotic Matter State-Driven MultiModal Fragmentation

Quantum Superposition



Energy-Time Evolution

Heavy-Flavor Fragmentation from HF-NRevo



Heavy-Flavor Fragmentation from **HF-NRevo**

Collinear_FFs Public

main | 1 Branch | 0 Tags | Go to file

FGCeliberto Update Readme.md

- NRFF10
- OMG3Q10
- OMG3Q11
- PQ5Q10
- PQ5Q11
- TQ4Q10
- TQ4Q11
- TQ4Q20
- TQHL10
- TQHL11
- LICENSE
- README.md

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Collinear_FFs

Collections of collinear fragmentation functions (FFs) given as LHPDF grids.

Available FF families to date:

1. NRFFx.y (NonRelativistic Fragmentation Functions)

   NRFF1.0 (NLO, 2 sets, central value only, LDME uncertainties)

2. OMG3Qx.y (OMeGa baryons with 3 heavy Quarks VFNS FFs)

   OMG3Q1.0 (NLO, 2 sets, central value only, diquark-like)

   OMG3Q1.1 (NLO, 2 sets, F-MHOU + F-NPWF uncertainties, replica 0->8, diquark-like)

3. TQHLx.y (TetraQuarks with Heavy and Light flavors)

   TQHL1.0 (NLO, 4 sets, central value only)

   TQHL1.1 (NLO, 4 sets, central value only)

4. TQ4Qx.y (TetraQuarks with 4 heavy Quarks)

   TQ4Q1.0 (NLO, 4 sets, central value only)

   TQ4Q1.1 (NLO, 6 main + 4 supplementary sets, replica-like MHOU uncertainties in main sets)

   TQ4Q2.0 (NLO, 6 sets, replica-like MHOU uncertainties)

5. PQ5Qx.y (PentaQuarks with 5 heavy Quarks)

   PQ5Q1.0 (NLO, 2 sets, replica 0 = direct multiquark, replica 1 = diquark-like)

   PQ5Q1.1 (NLO, 1 set, F-MHOU + F-NPWF uncertainties, replica 0->8 = direct multicharm, replica 9->17 = diquark-like)

🔗 https://github.com/ FGCeliberto/Collinear_FFs/

Heavy-Flavor Fragmentation from **HF-NRevo**

Collinear_FFs Public

main | 1 Branch | 0 Tags | Go to file

FGCeliberto Update Readme.md

- NRFF10
- OMG3Q10
- OMG3Q11
- PQ5Q10
- PQ5Q11
- TQ4Q10
- TQ4Q11
- TQ4Q20
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Available FF families to date:

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   NRFF1.0 (NLO, 2 sets, central value only, LDME uncertainties)

2. OMG3Qx.y (OMeGa baryons with 3 heavy Quarks VFNS FFs)

   OMG3Q1.0 (NLO, 2 sets, central value only, diquark-like)

   OMG3Q1.1 (NLO, 2 sets, F-MHOU + F-NPWF uncertainties, replica 0->8, diquark-like)

3. TQHLx.y (TetraQuarks with Heavy and Light flavors)

   TQHL1.0 (NLO, 4 sets, central value only)

   TQHL1.1 (NLO, 4 sets, central value only)

4. TQ4Qx.y (TetraQuarks with 4 heavy Quarks)

   TQ4Q1.0 (NLO, 4 sets, central value only)

   TQ4Q1.1 (NLO, 6 main + 4 supplementary sets, replica-like MHOU uncertainties in main sets)

   TQ4Q2.0 (NLO, 6 sets, replica-like MHOU uncertainties)

5. PQ5Qx.y (PentaQuarks with 5 heavy Quarks)

   PQ5Q1.0 (NLO, 2 sets, replica 0 = direct multiquark, replica 1 = diquark-like)

   PQ5Q1.1 (NLO, 1 set, F-MHOU + F-NPWF uncertainties, replica 0->8 = direct multicharm, replica 9->17 = diquark-like)

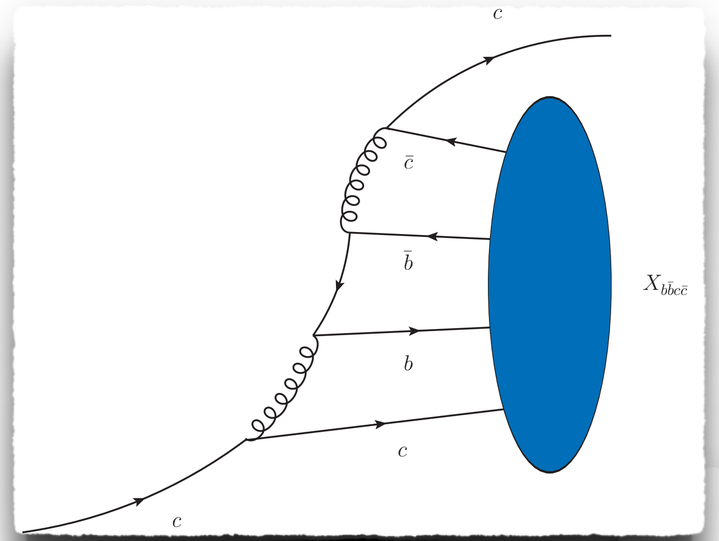
NEW

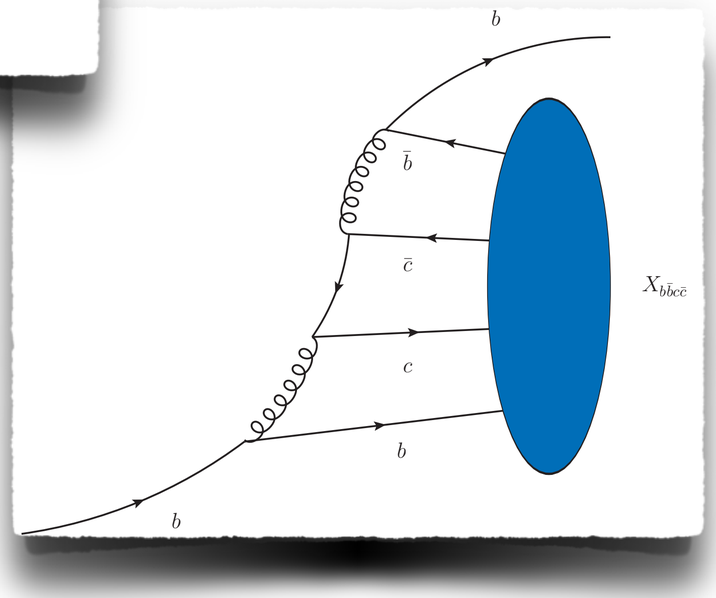
TQ4Qhf1.0 ¡released this week!

🔗 https://github.com/ FGCeliberto/Collinear_FFs/

MultiModal Fragmentation to TetraQuarks

Mode 1
Direct MultiQuark
Scalar S-wave $X_{b\bar{b}c\bar{c}}$





Initial-scale input: 🔗🔗 Suzuki model

TQ4Qhf1.0 FFs: $(Q \to X_{b\bar{b}c\bar{c}}) \otimes$ **APFEL++**

F-MHOUs [P] + F-NPWF [NP]

TQHL1.0 (Hidden-charm) 🔗 [F. G. C., A. Papa, Phys. Lett. B 848 (2024) 138406, arXiv:2308.01539]

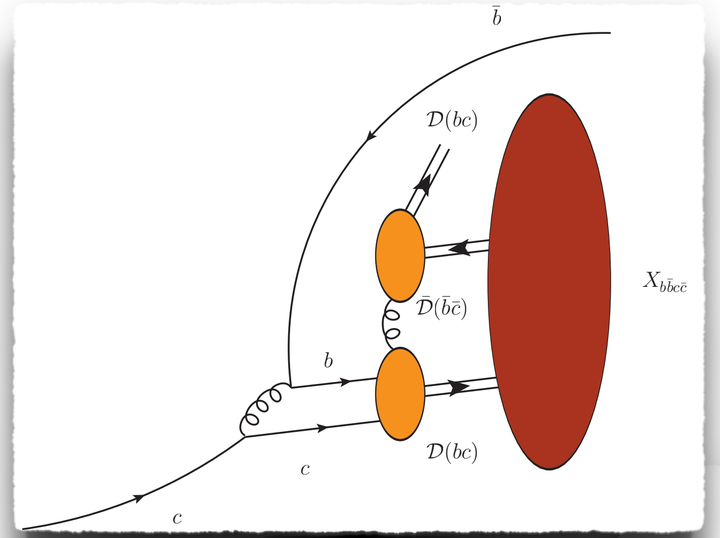
TQHL1.0 (Hidden-bottom, Best Paper Award) 🔗 [F. G. C., Symmetry 16 (2024) 5, 550, arXiv:2403.15639]

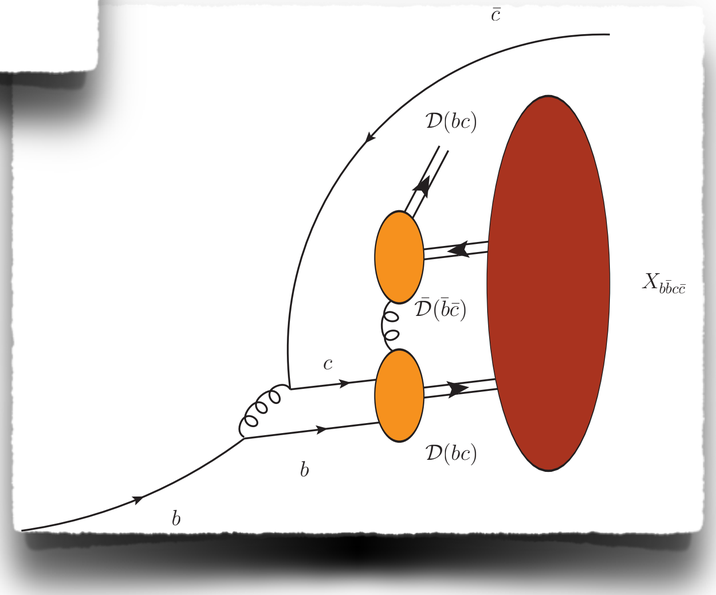
TQHL1.1 (Hidden-heavy) 🔗 [F. G. C., G. Gatto, Phys. Rev. D. 111 (2025) 3, 034037, arXiv:2412.10549]

T4Qhf1.0 (All-heavy mixed flavor, uncertainties, observability) 🔗 [J. V. Bulhões, F. G. C., V. P. Gonçalves, arXiv:2610.27396]

MultiModal Fragmentation to TetraQuarks

Mode 2
Scalar DiQuark
Scalar S-wave $X_{b\bar{b}c\bar{c}}$





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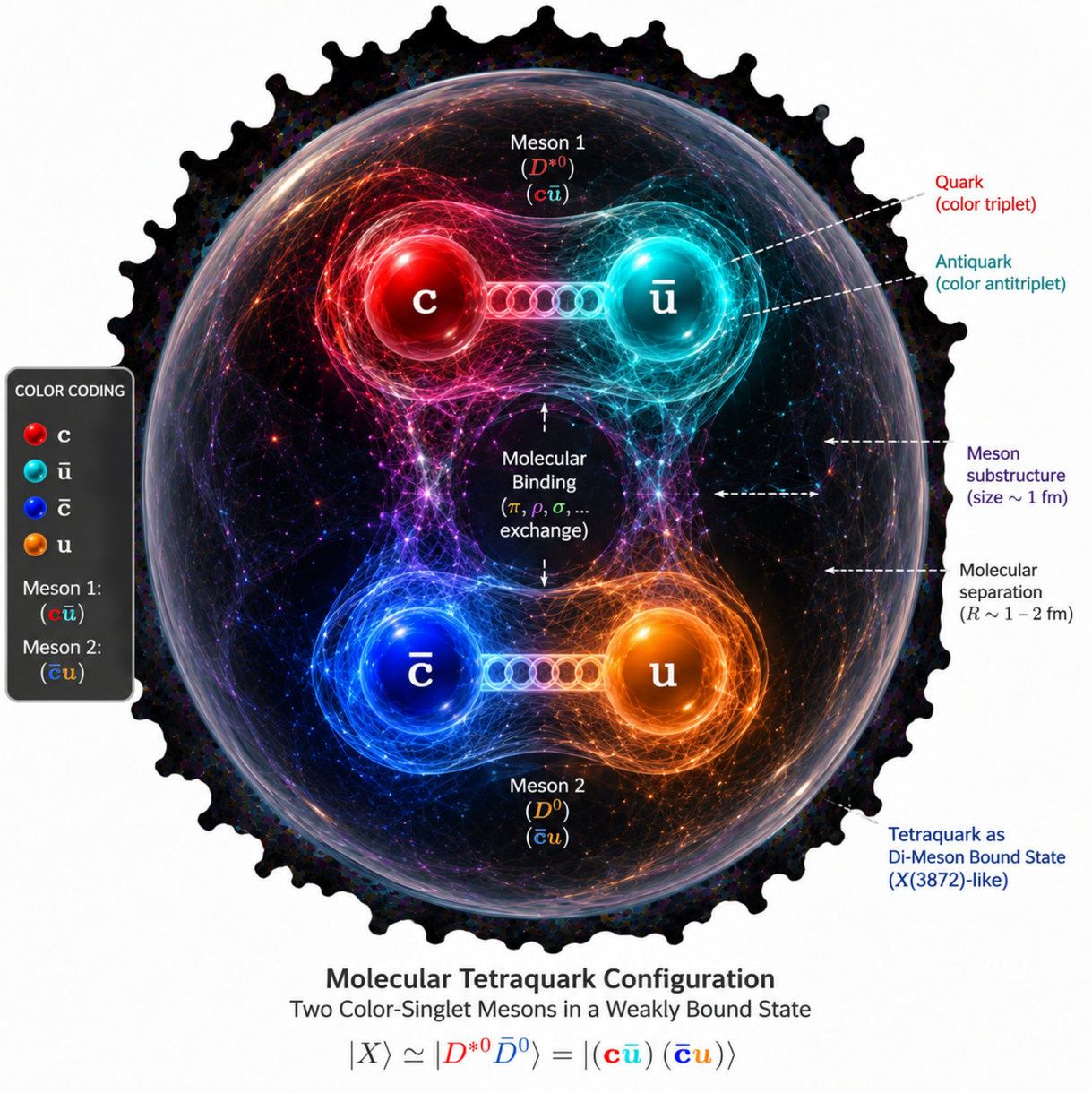
¿ What's next ?

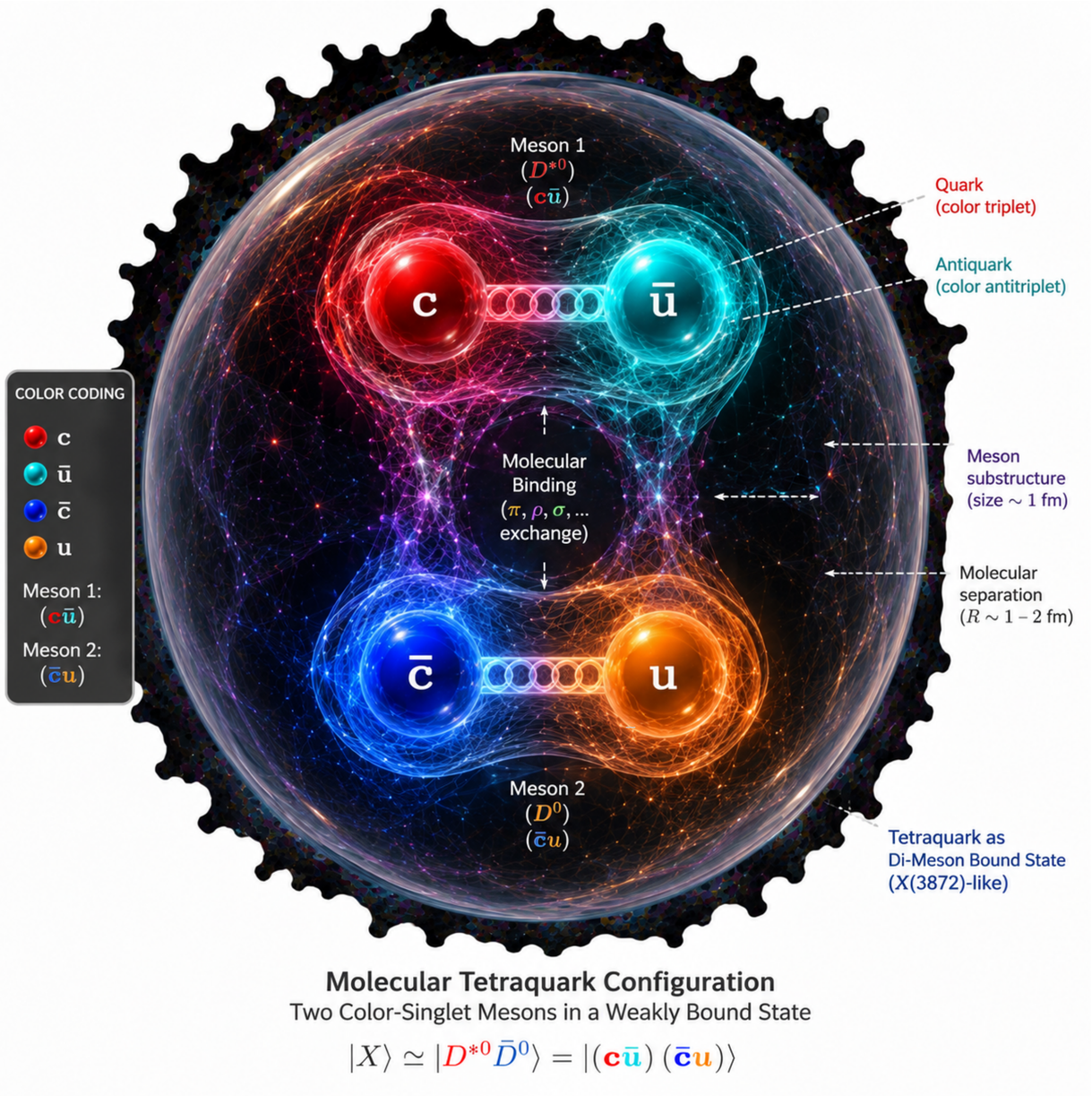


Molecular Tetraquark Configuration

Two Color-Singlet Mesons in a Weakly Bound State

$$|X\rangle \simeq |D^{*0}\bar{D}^0\rangle = |(c\bar{u})\,(\bar{c}u)\rangle$$

¿ What's next ?



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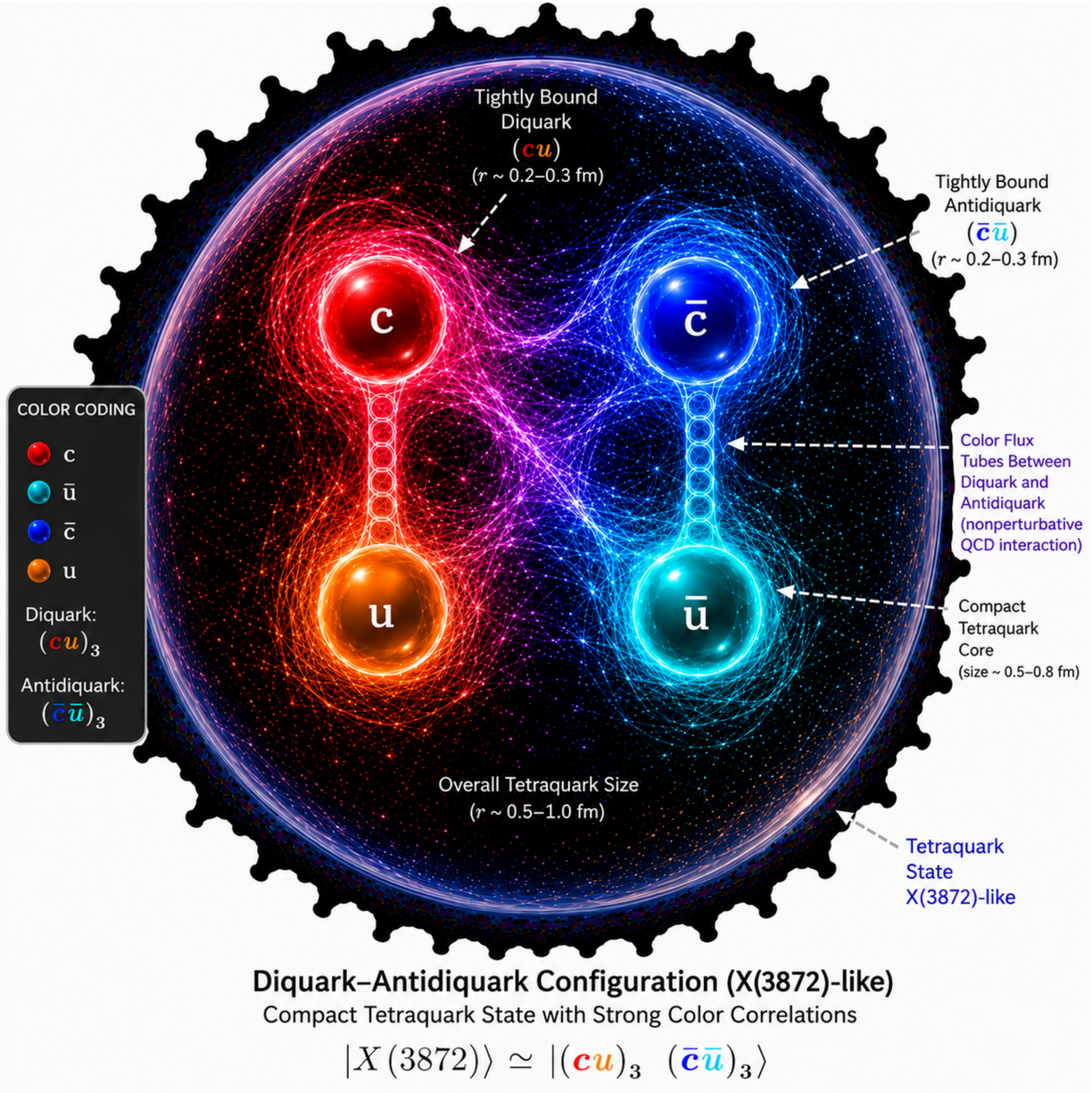
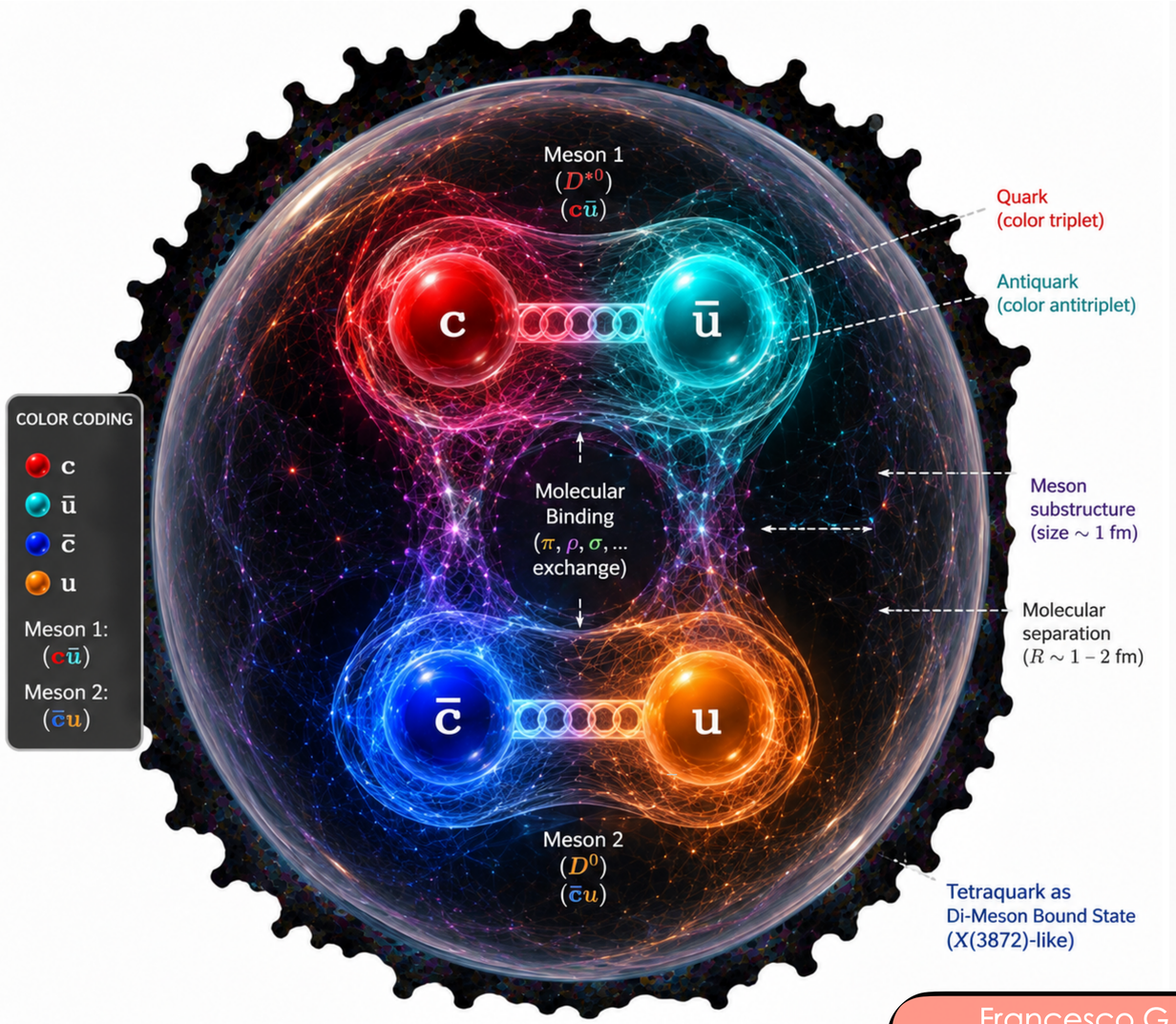


Diquark–Antidiquark Configuration (X(3872)-like)

Compact Tetraquark State with Strong Color Correlations

$$|X(3872)\rangle \simeq |(cu)_3\ \ (\bar{c}\bar{u})_3\rangle$$

¿ What's next ?



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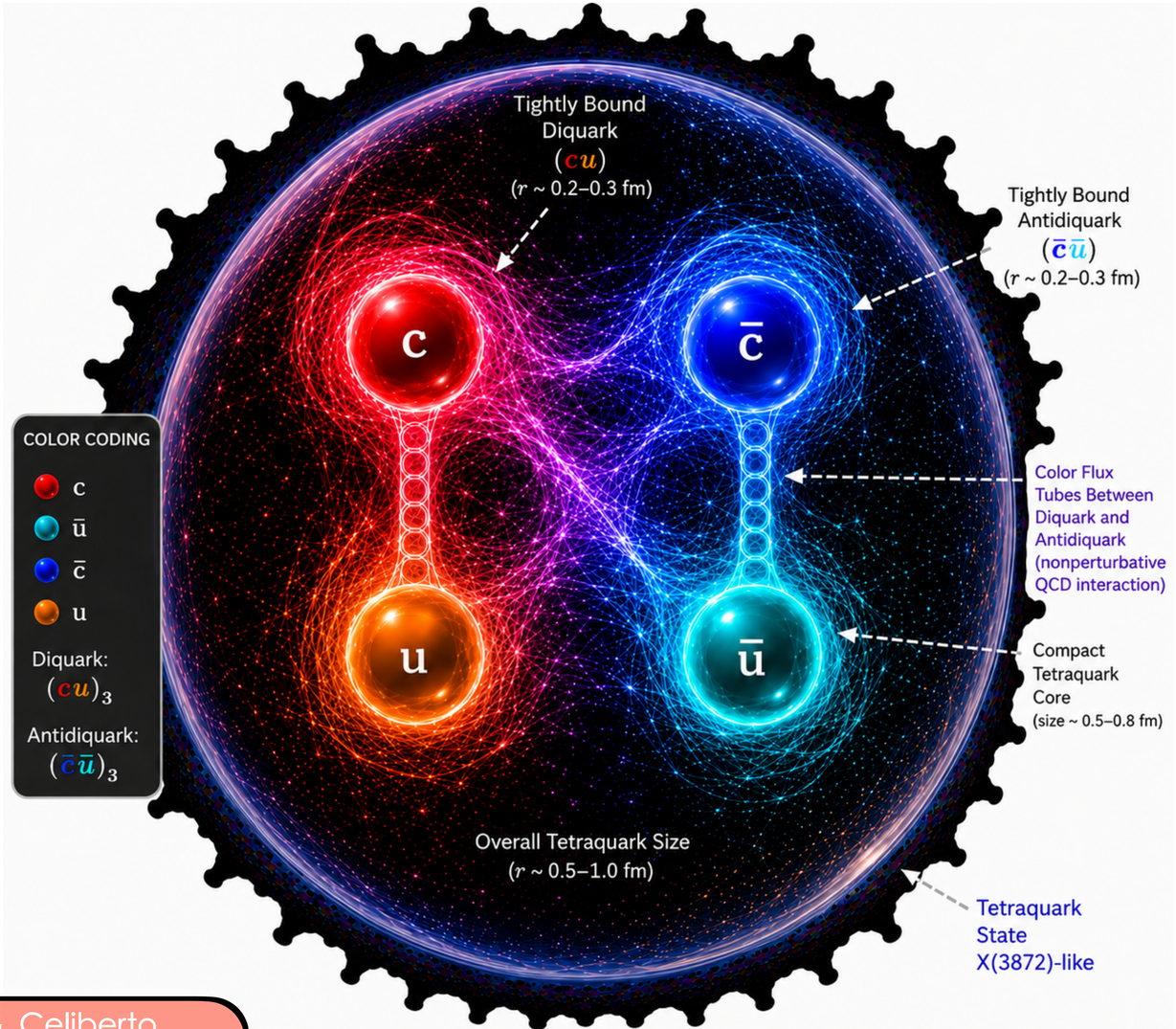

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Francesco G. Celiberto
Carlo L. Guerrero Conteras
Francesca Lonigro
Julio Bohigas Couto
Gonzalo Tudanca Munsuri

Toward **ExoFF**

Toward ExoFF



Exotic hadrons ⇒ precision QCD & production mechanisms

Toward **ExoFF**



Exotic hadrons ⇒ precision QCD & production mechanisms



HF-NRevo ⇒ consistent evolution, uncertainty-aware FFs

Toward **ExoFF**

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- **HF-NRevo** ⇒ consistent evolution, uncertainty-aware FFs
- AI-driven extraction & global fits of heavy-hadron FFs

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- From theory-driven benchmarks to state-oriented fragmentation

Toward **ExoFF**

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- **HF-NRevo** ⇒ consistent evolution, uncertainty-aware FFs
- AI-driven extraction & global fits of heavy-hadron FFs
- From theory-driven benchmarks to state-oriented fragmentation
- Road to a MULTIMODAL **ExoFF** framework

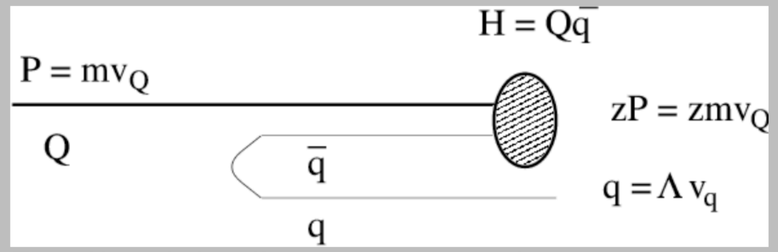
EXTRAS

Highlights on Heavy-Flavor Fragmentation

We observe hadrons, whose momentum is different from that of the parent heavy quark. We usually write this as

$$\text{FF(hadron)} = \text{FF(quark)} \otimes \text{FF(quark}\rightarrow\text{hadron)}$$

What are the characteristics of FF(quark→hadron)?



The heavy quark has momentum $P = mv_Q$, the light quarks have momentum $q = \Lambda v_q$, with Λ a hadronic mass scale

For the binding we need $v_Q \simeq v_q = v$. We have then

$$P = zP + q \qquad mv = zmv + \Lambda v$$

and therefore

$$\langle z \rangle \simeq 1 - \frac{\Lambda}{m}$$

Hence, FF(quark→hadron) is peaked near $z = 1$ and binding effects scale linearly in the heavy quark mass (Bjorken/Suzuki ($\sim$ 1977))

(NB. more rigorous derivations of this result are available)

Highlights of Non-Relativistic QCD

NRQCD → effective field theory → nonrelativistic (NR) $Q\bar{Q}$ system

Heavy-quark spinor fields as NR DOFs in the NQRCD Lagrangian

$$\mathcal{L}_{\mathrm{NRQCD}} = \psi^\dagger \left(iD_0 + \frac{\vec{D}^2}{2m_Q} \right) \psi + \chi^\dagger \left(iD_0 - \frac{\vec{D}^2}{2m_Q} \right) \chi + \mathcal{L}_{\mathrm{light}} + \delta\mathcal{L}$$

Highlights of Non-Relativistic QCD

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[cross section] = [short-distance coeffs.] ⊗ [long-distance matrix elements]

$$d\sigma(H + X) = \sum_n d\hat{\sigma}(Q\overline{Q}\,[n] + X) \left\langle \mathcal{O}^H\,[n] \right\rangle$$

Highlights of Non-Relativistic QCD

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$$d\sigma(H+X) = \sum_n d\hat{\sigma}(Q\overline{Q}\,[n] + X)\,\langle \mathcal{O}^H\,[n] \rangle$$

$$|\mathcal{Q}\rangle = \mathcal{O}(1)\,|Q\bar{Q}[^3S_1^{(1)}]\rangle + \mathcal{O}(v)\,|Q\bar{Q}[^3P_J^{(8)}g]\rangle + \mathcal{O}(v^2)\,|Q\bar{Q}[^1S_0^{(8)}g]\rangle + \mathcal{O}(v^2)\,|Q\bar{Q}[^3S_1^{(1,8)}gg]\rangle + \mathcal{O}(v^2)\,|Q\bar{Q}[^3D_J^{(1,8)}gg]\rangle + \ldots .$$

Heavy-Flavor Non-Relativistic evolution

HF-NRevo

Interpretation | Evolution | Uncertainties

- NRQCD ⇒ proxy for initial-scale FF inputs
- A heavy-flavor-scheme based vision
- Low vs large p_T ⇒ FFNS vs VFNS fragmentation

Heavy-Flavor Non-Relativistic evolution

HF-NRevo

Interpretation | Evolution | Uncertainties

- Proper DGLAP collinear evolution
- Thresholds at work for all partons
- Two-step strategy $\Rightarrow$ `EDevo` $\otimes$ `AOevo`

Heavy-Flavor Non-Relativistic evolution

HF-NRevo

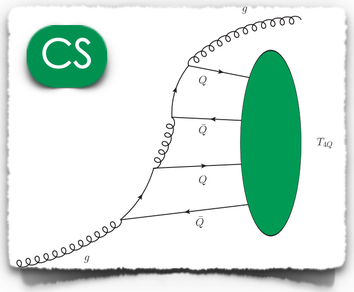
Interpretation

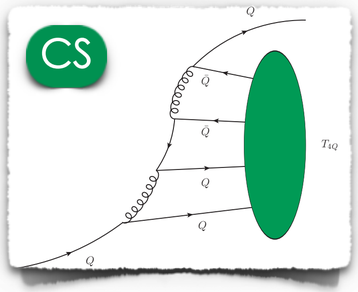
Evolution

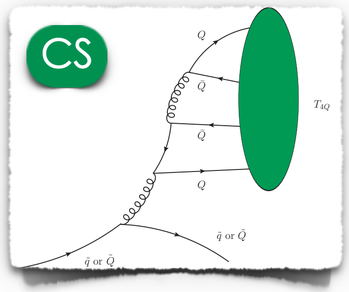
Uncertainties

- Initial-scale MHOUs ⇒ first systematic analysis
- Idea ⇔ `ThCovM` 🔗🔗 & `MCscales` 🔗 MHOUs
- Future ⇐ extraction of NP LDMEs from data

HF-NRevo: Initial-Scale NRCQD Inputs





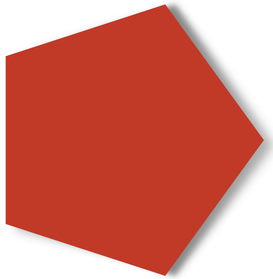


Scalar $[J^{PC} = 0^{++}]$

Tensor $[J^{PC} = 2^{++}]$

HF-NRevo for $T_{4c}(0^{++}, 2^{++})$ collinear fragmentation

Partonic hard factor



DGLAP AOevo

All-Order evolution (numeric)

DGLAP EDevo

Expanded Decoupled evolution

NRQCD

Color-composite LDMEs

SDCs

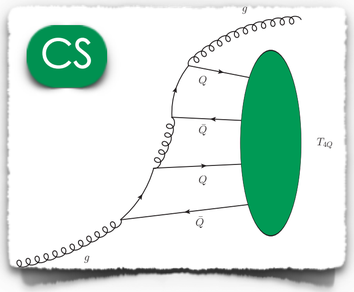
$\mathcal{D}_i^{\mathcal{Q}}(z,[n]) \otimes \langle \mathcal{O}^{\mathcal{Q}}([n]) \rangle$

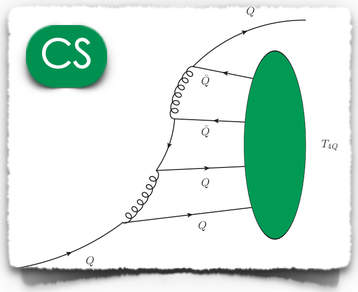
Energy

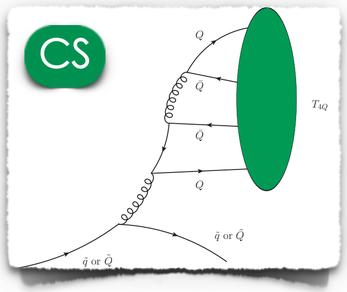
Evolution-ready scale $Q_0 \equiv m_b + 4m_c \quad 5m_c \quad m_{\tilde{q}} + 4m_c$

Time

HF-NRevo: Initial-Scale NRCQD Inputs







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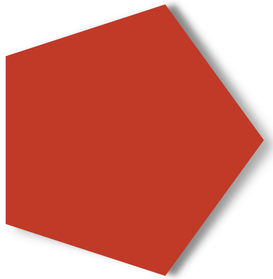

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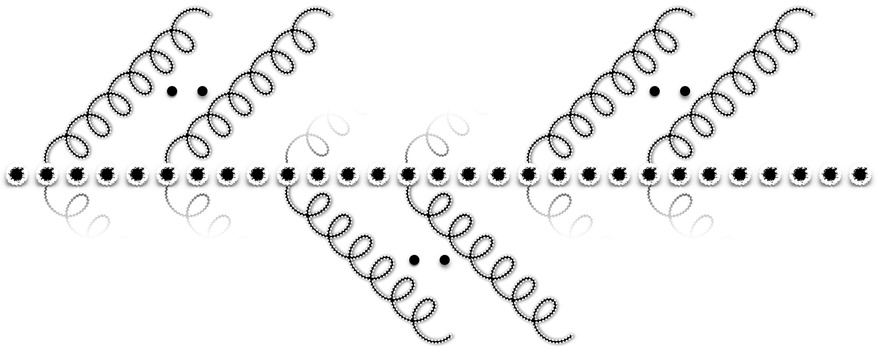
All-Order evolution (numeric)



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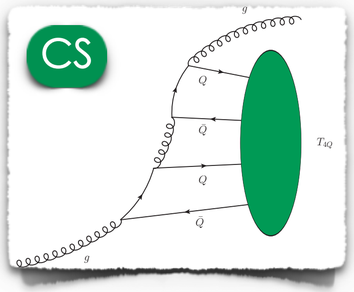

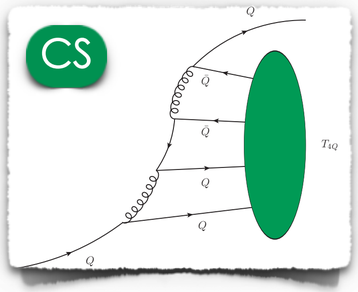

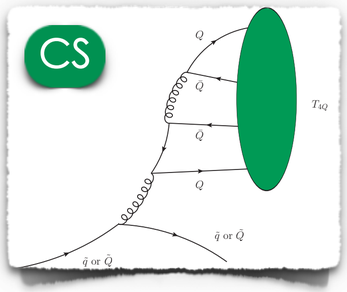

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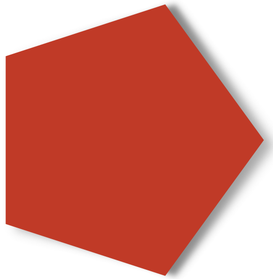

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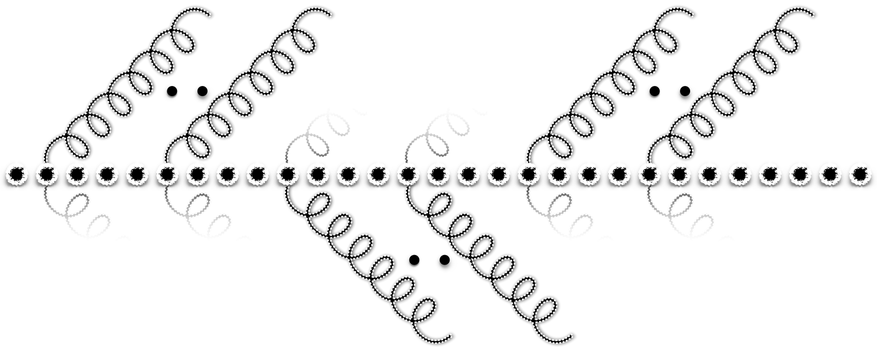

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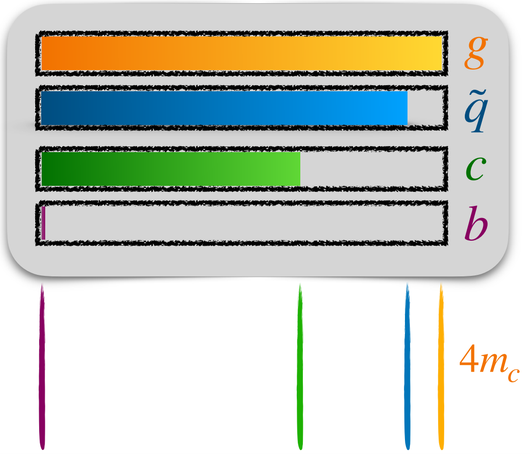

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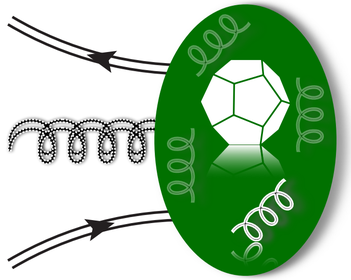

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Time

ALL-HEAVY PENTAQUARK FRAGMENTATION

MultiModal Fragmentation to PentaQuarks

Mode 1
Direct MultiQuark
Scalar S-wave P_{5c}

🔗 [F. G. C., Eur. Phys. J. C 85 (2025) 12, 1395]; 🔗 [F. G. C., Phys. Rev. D. (to appear), arXiv:2605.01539]

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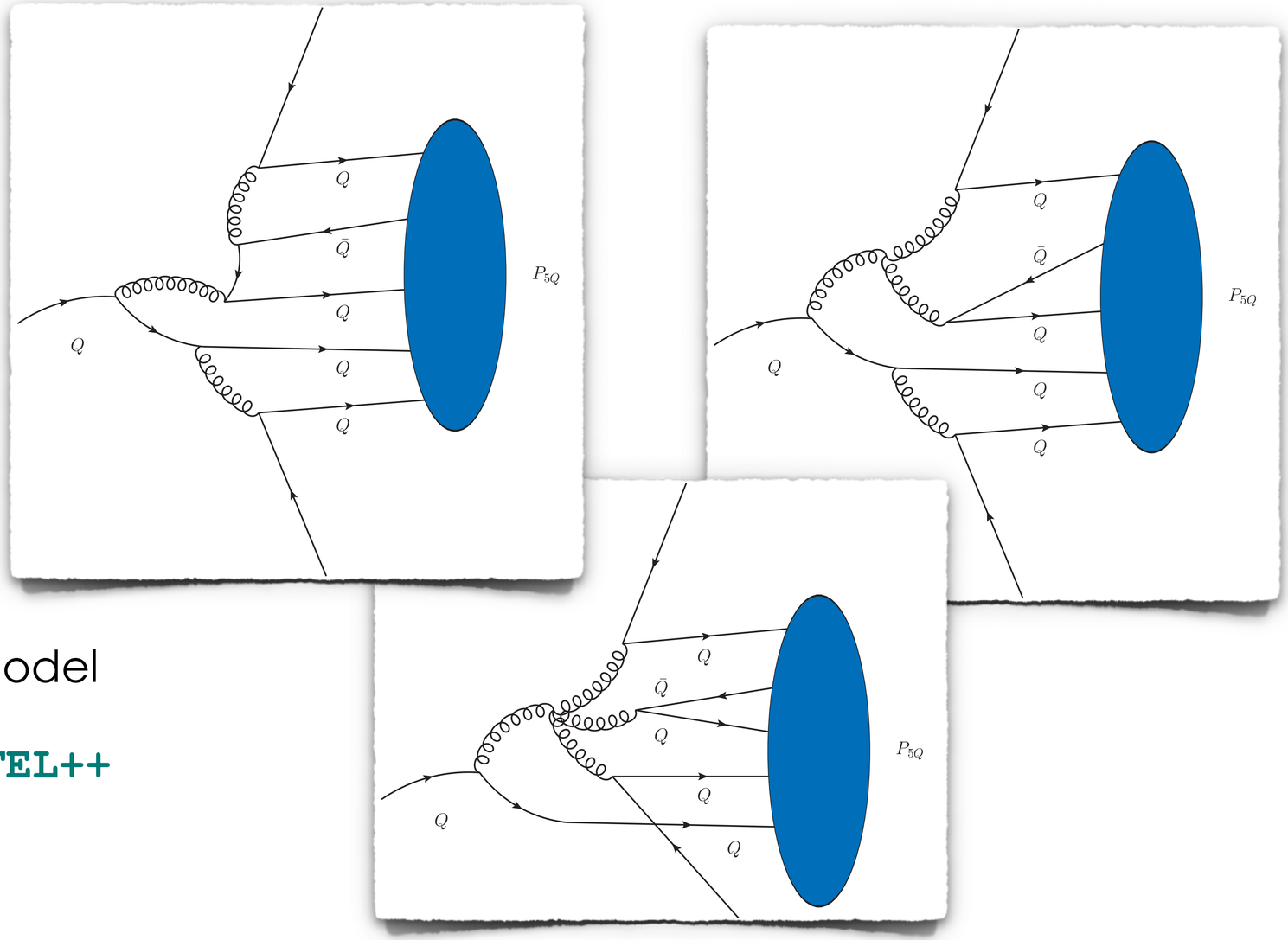
Scalar S-wave P_{5c}



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MultiModal Fragmentation to PentaQuarks

Mode 2
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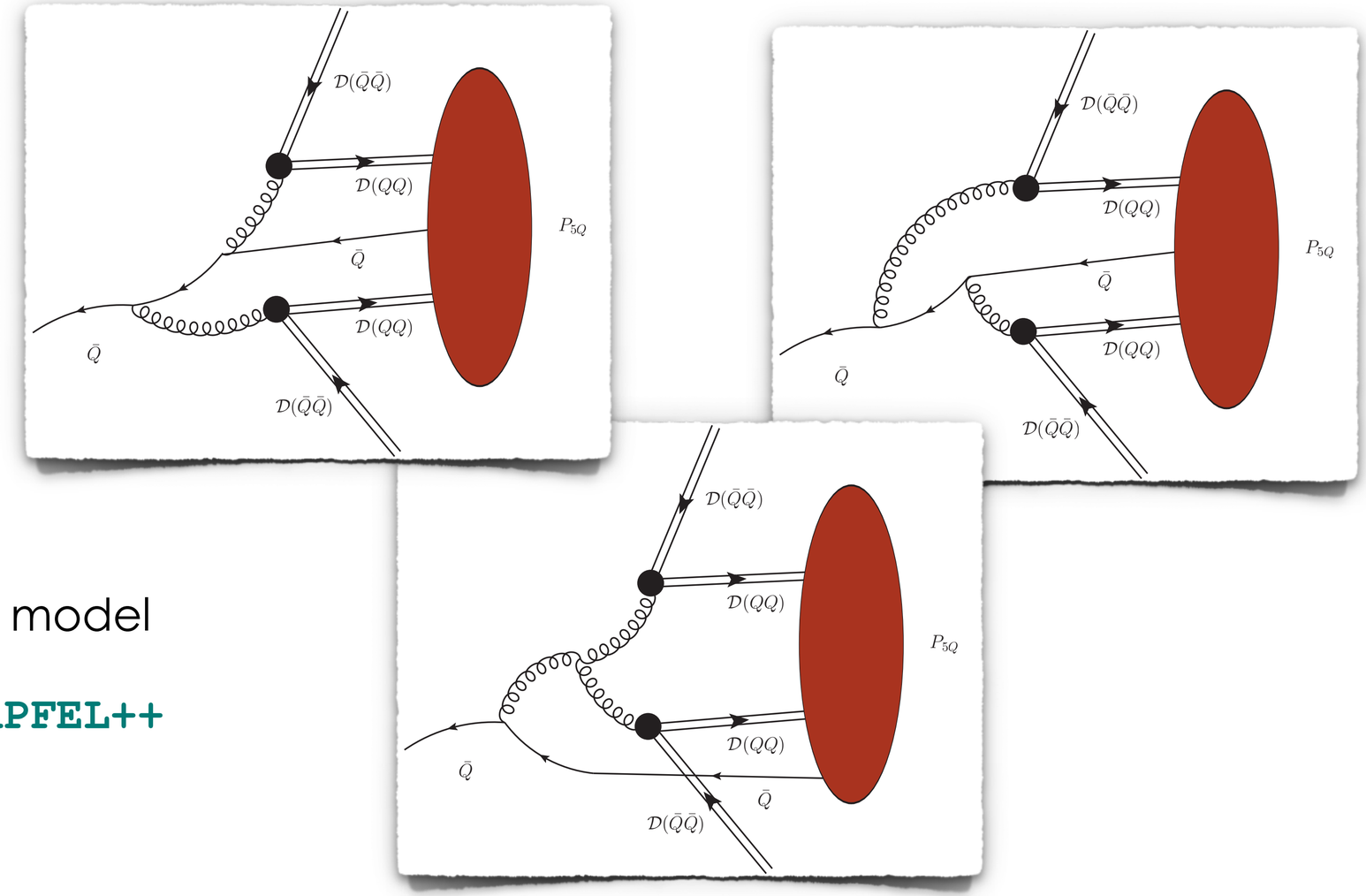
Scalar S-wave P_{5c}



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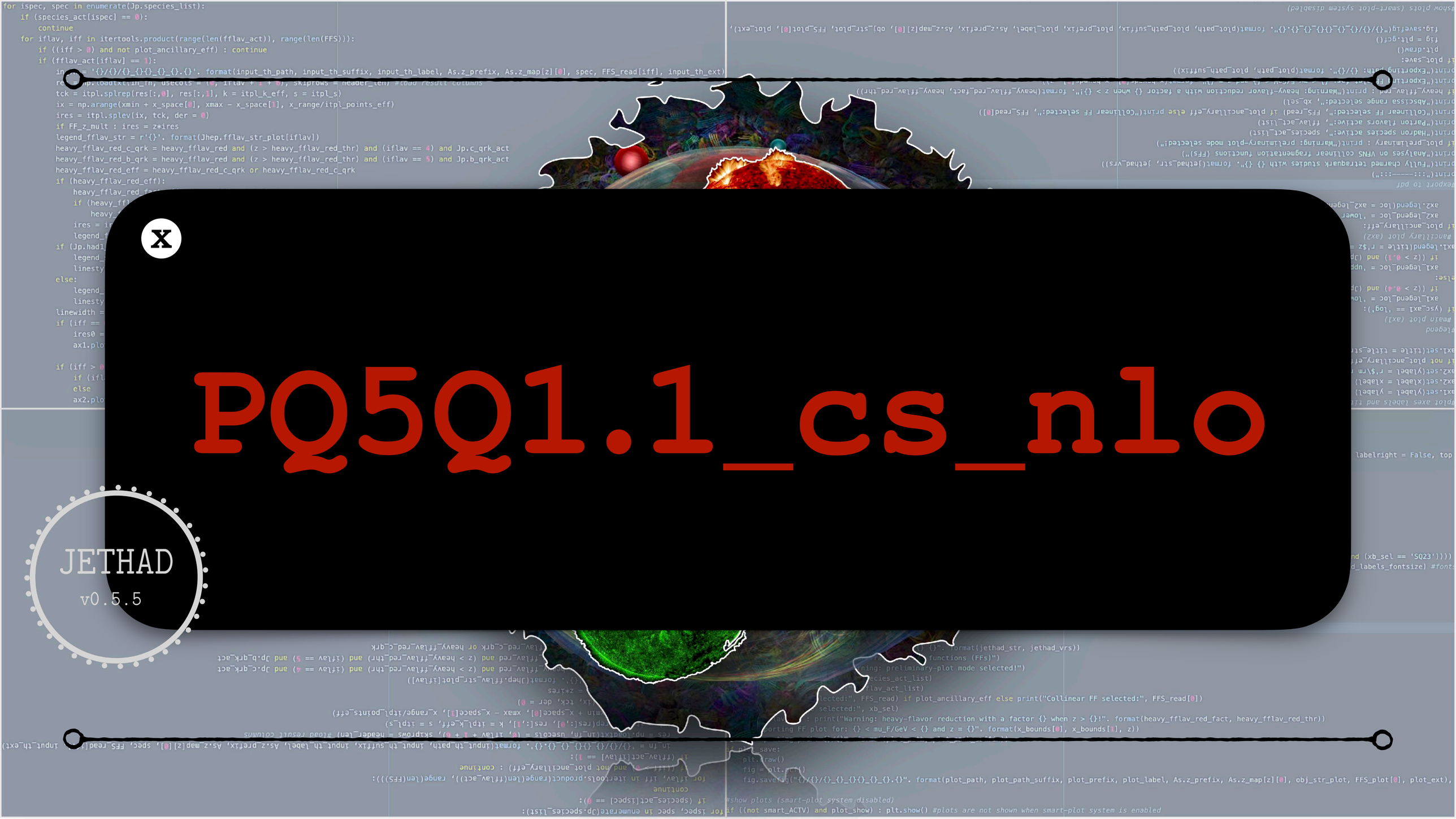
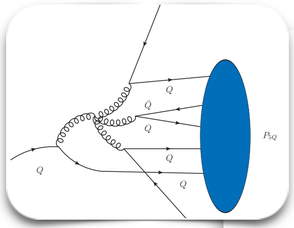
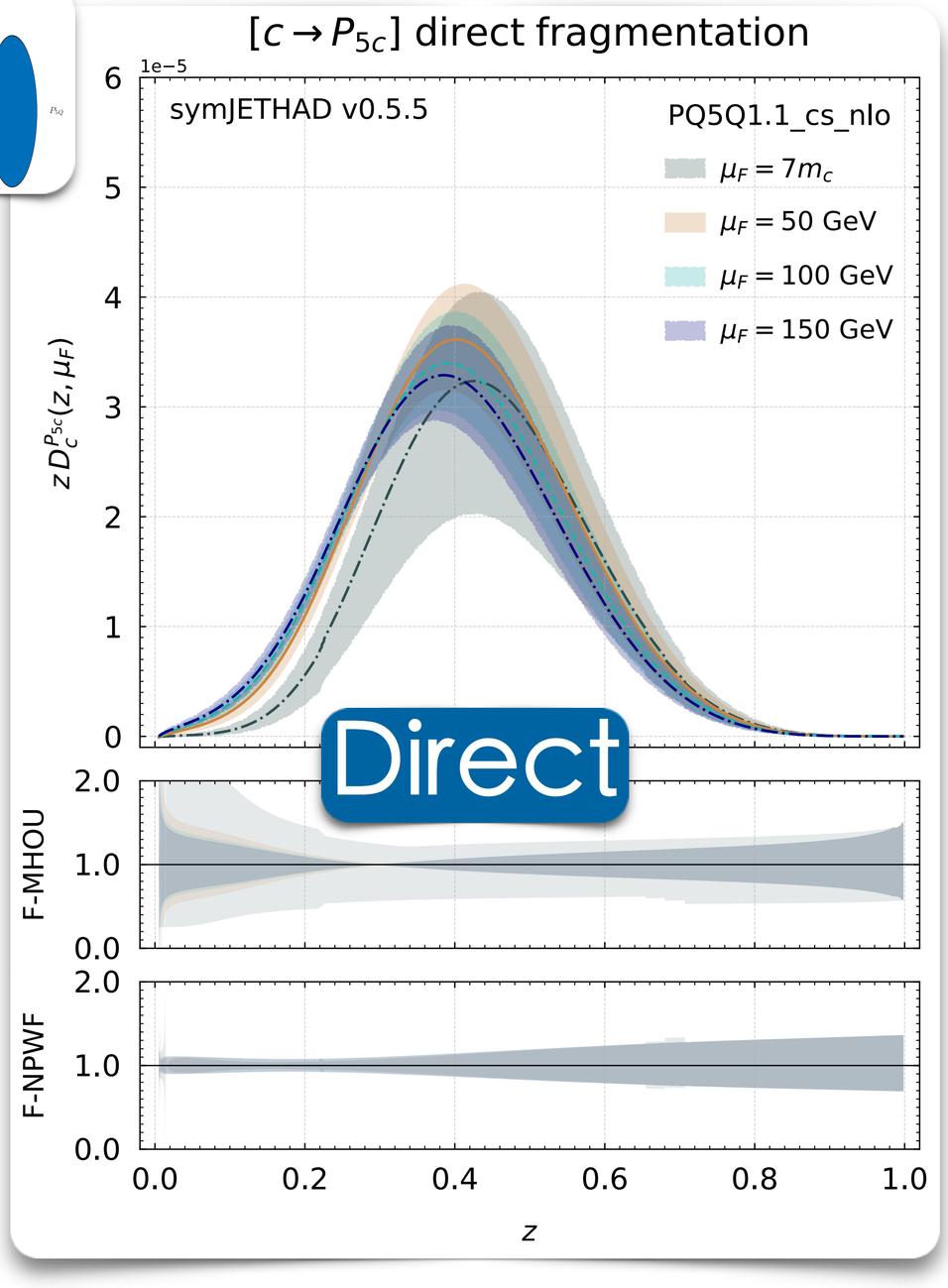
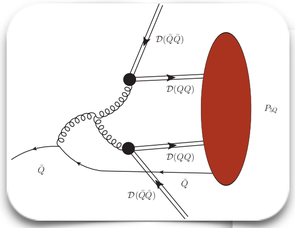
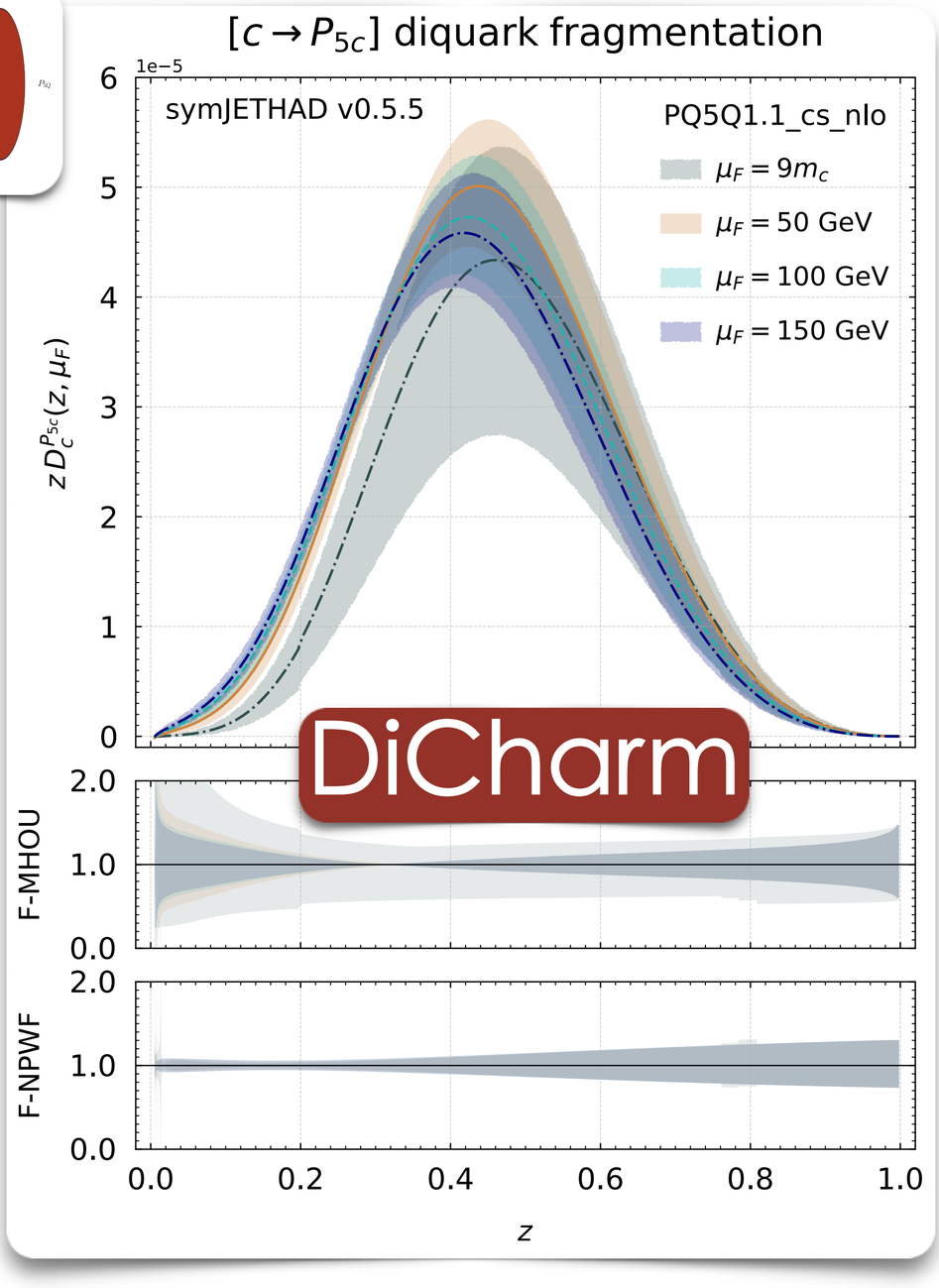

PQ5Q1.1_cs_nlo
JETHAD
v0.5.5

PQ5Q1.1: Charm Fragmentation to PentaCharm

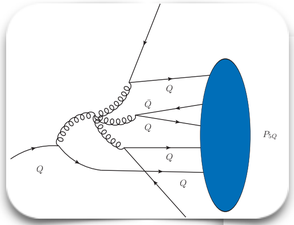


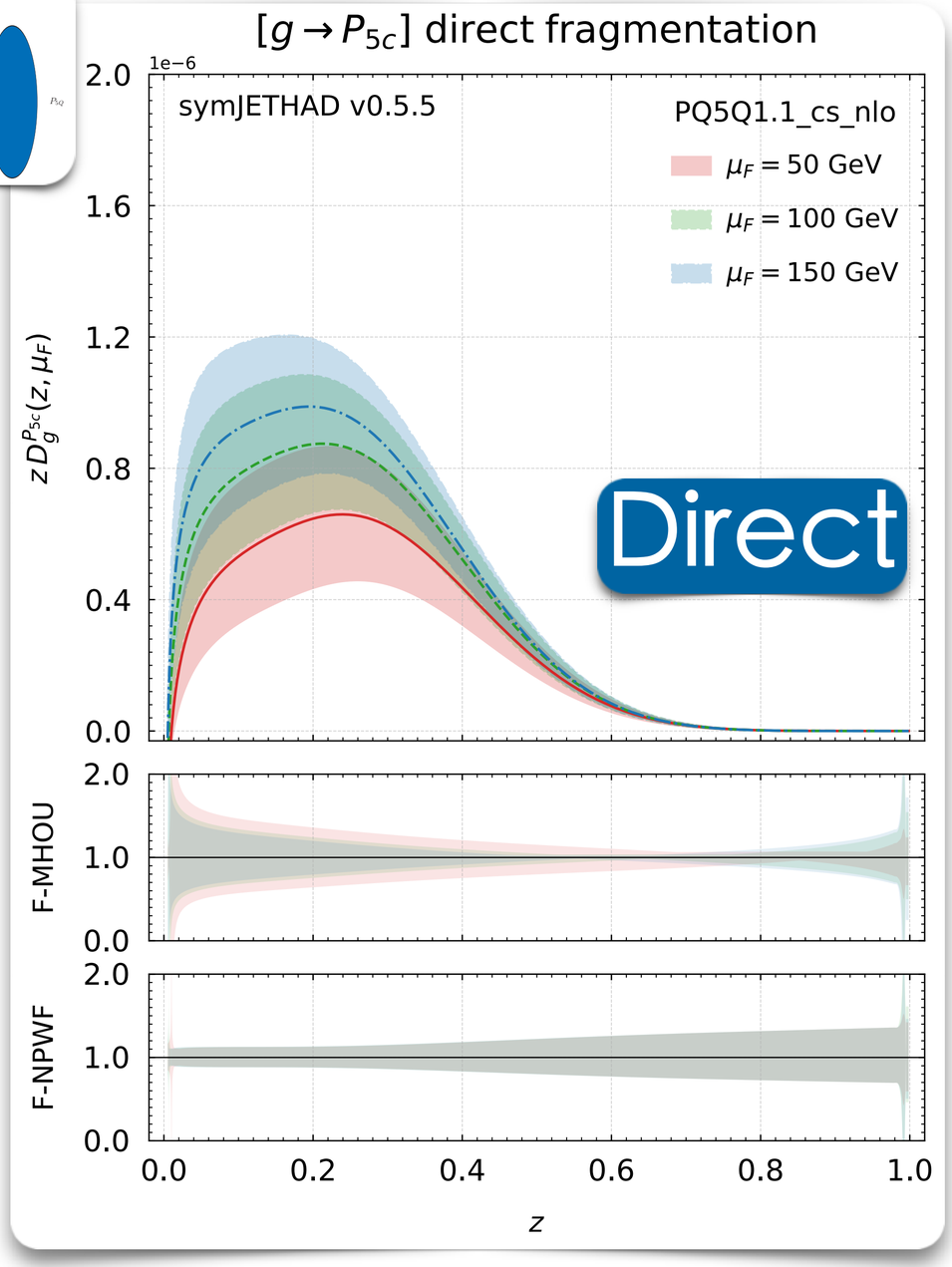


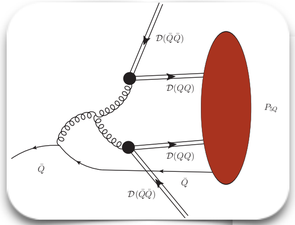


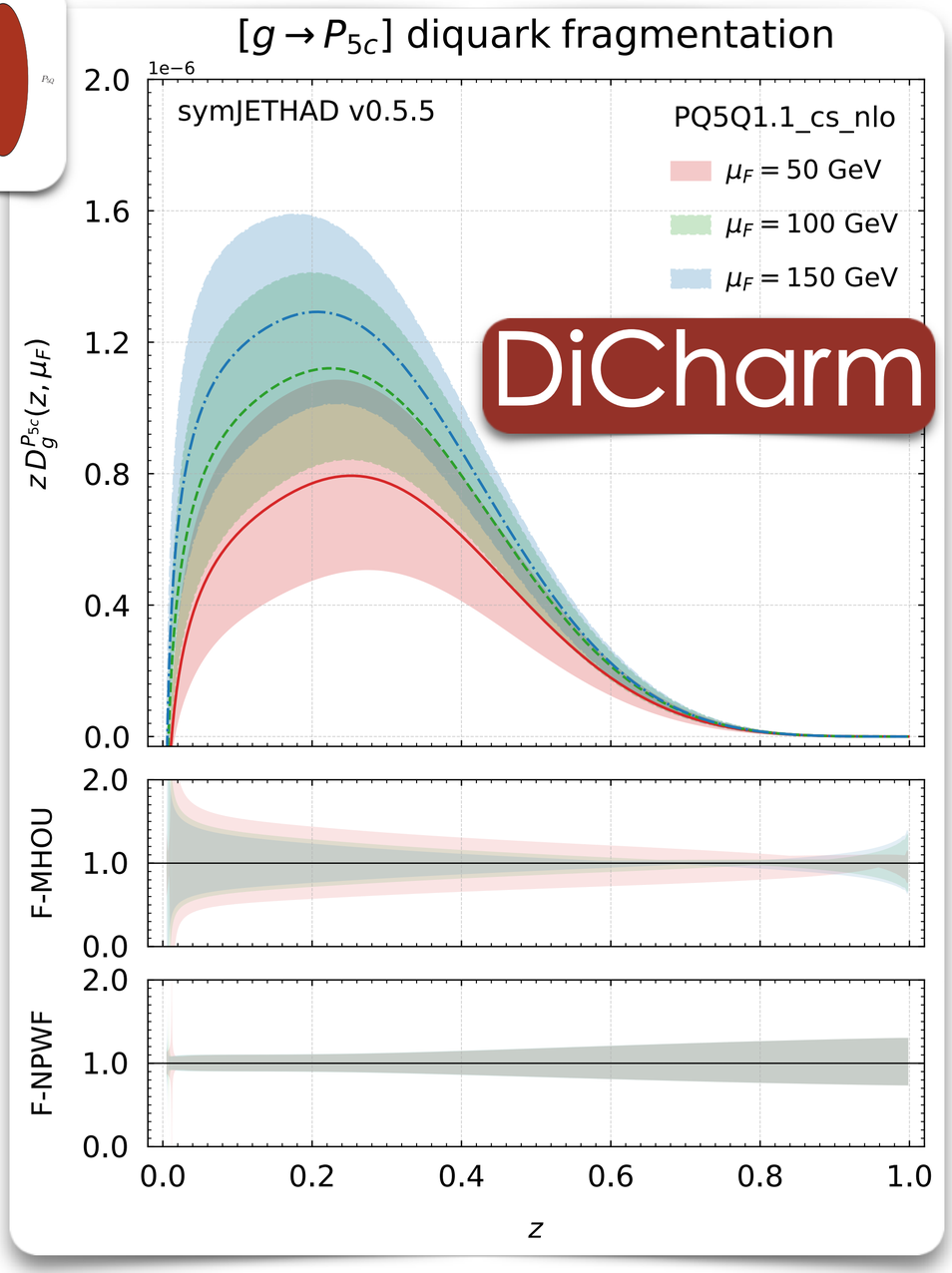


PQ5Q1.1: Gluon Fragmentation to PentaCharm

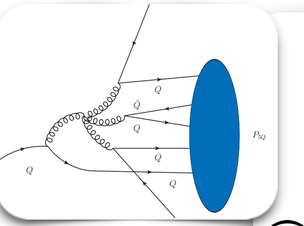


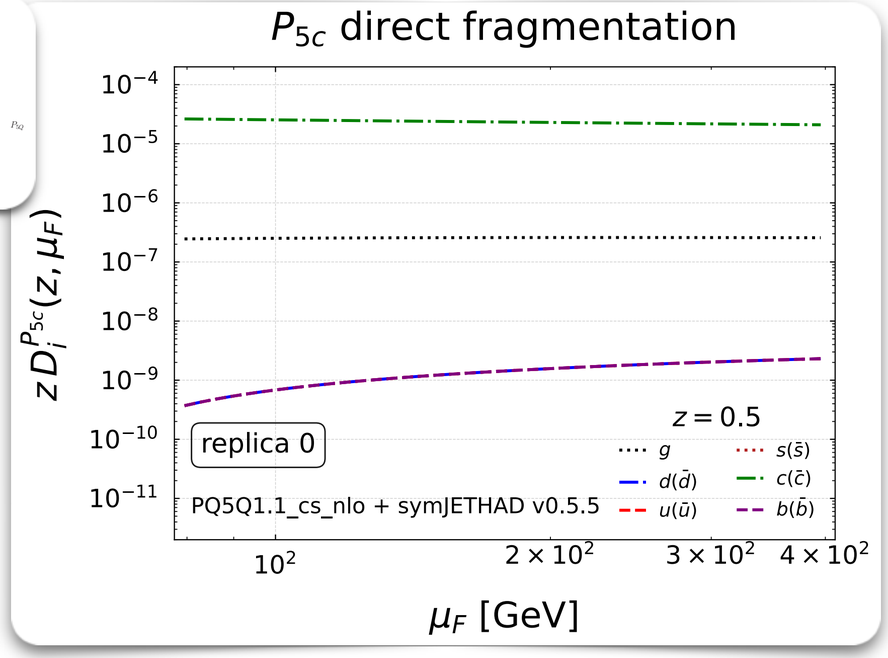


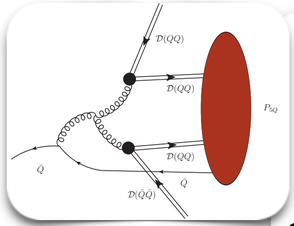


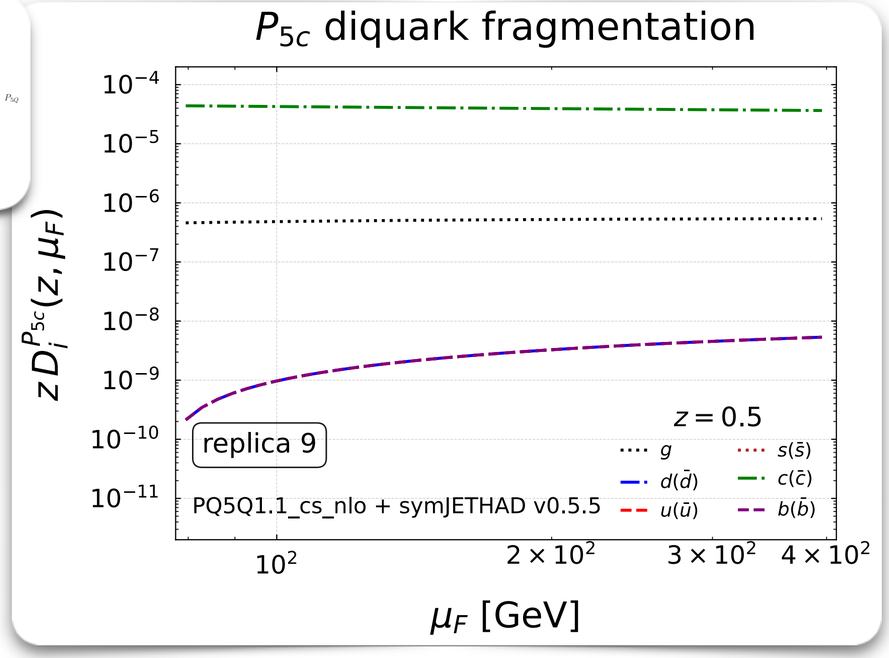


PQ5Q1.1: Energy Spectrum





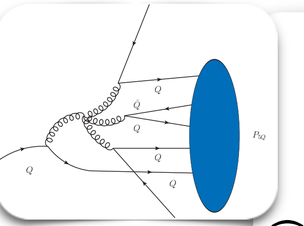


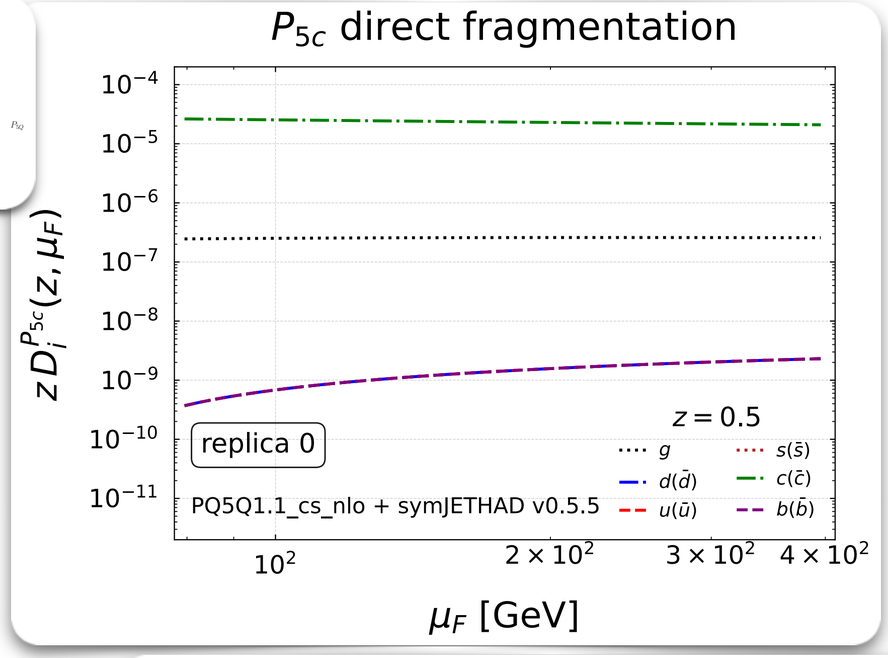


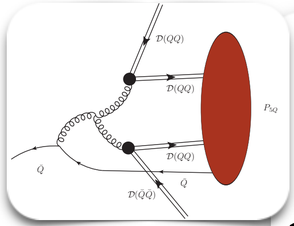
DiCharm

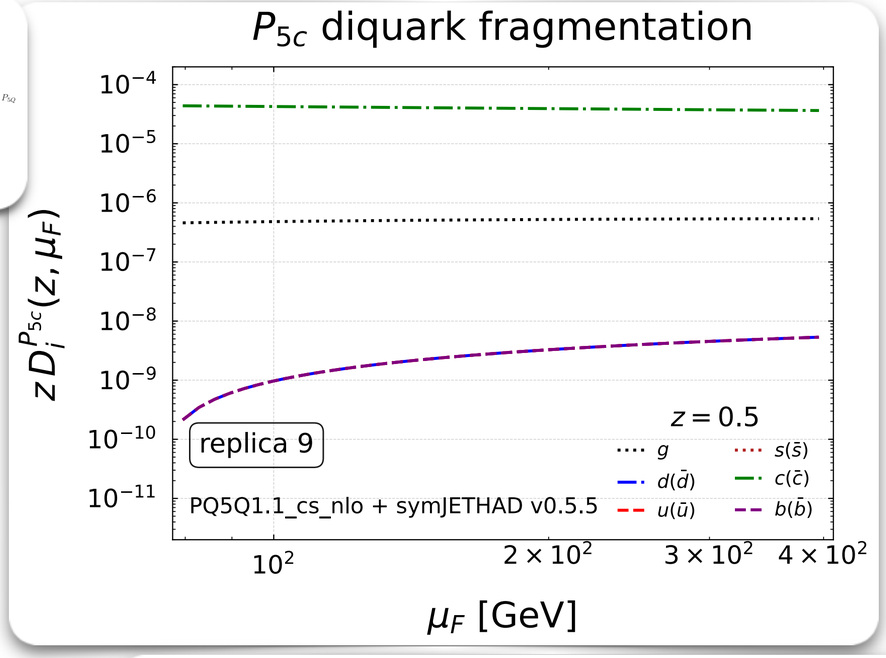
Direct

PQ5Q1.1: Energy Spectrum



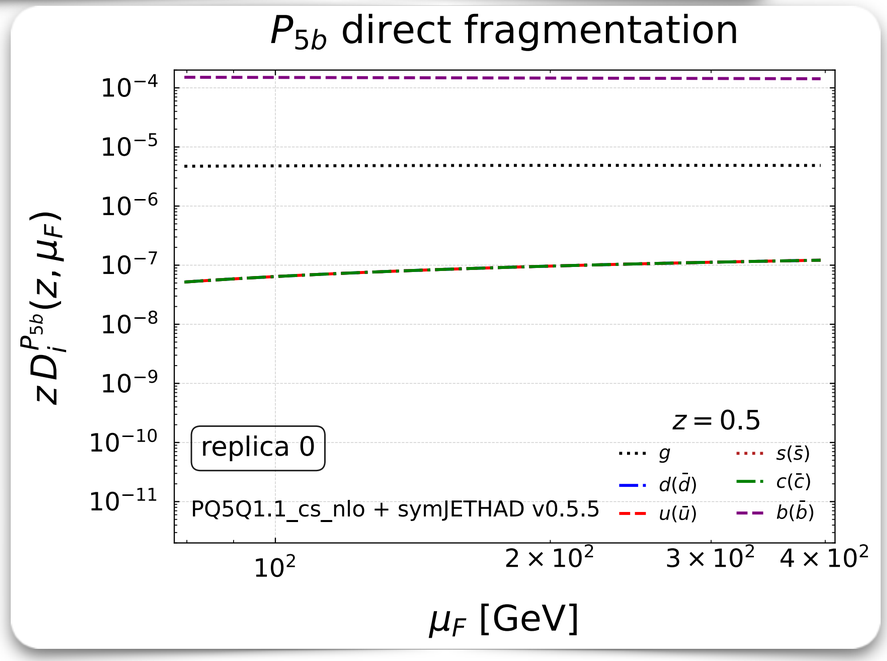


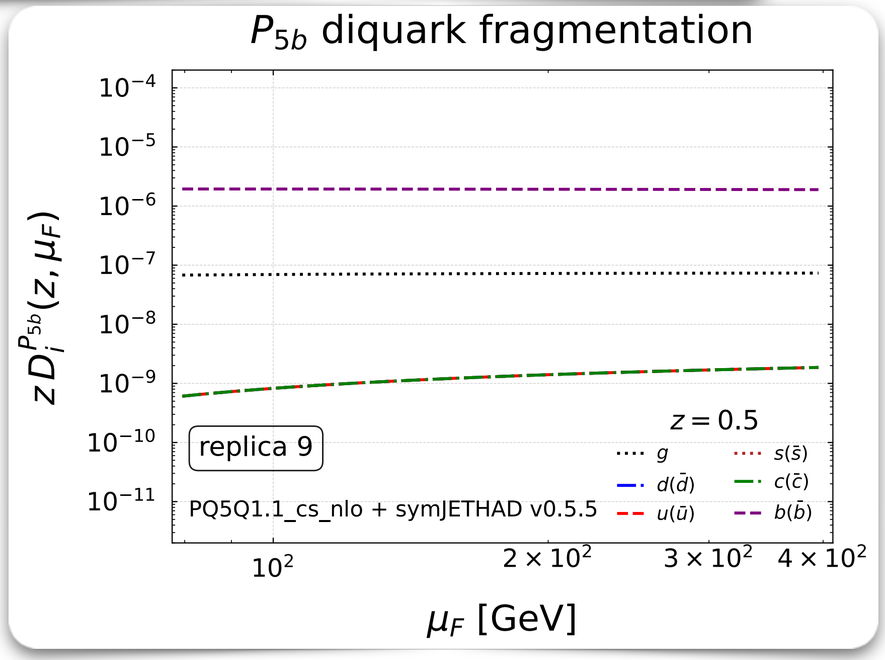




DiCharm

DiRect



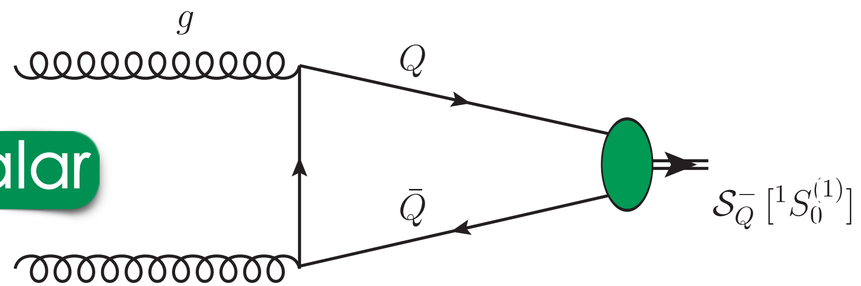


QUARKONIUM FRAGMENTATION

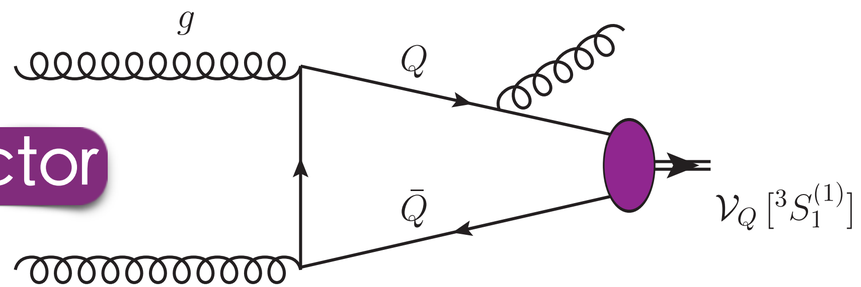
From low energies to high energies

Low p_T $\Rightarrow$ two-parton FFs
Genuine higher-twist FFNS FFs

scalar

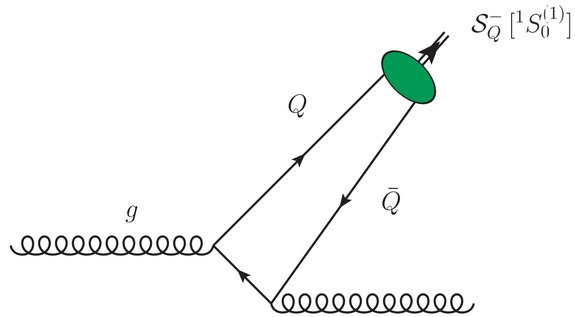


vector



High p_T $\Rightarrow$ single-parton FFs
Twist-2 VFNS FFs + power corrections

scalar



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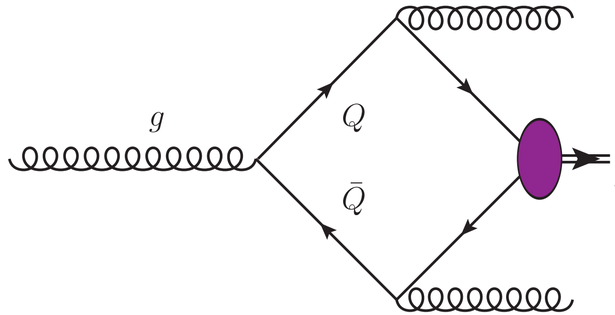
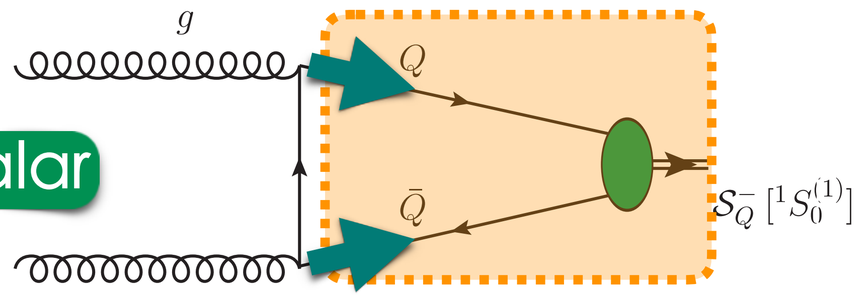


From low energies to high energies

Low p_T ⇒ two-parton FFs
Genuine higher-twist FFNS FFs

scalar



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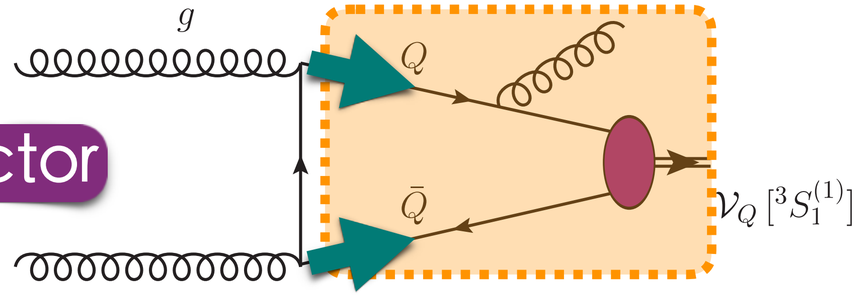
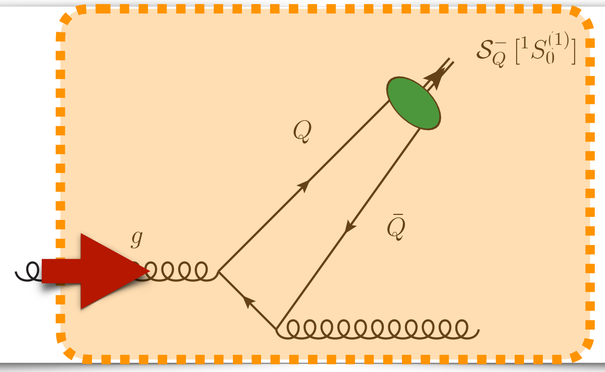


High p_T ⇒ single-parton FFs
Twist-2 VFNS FFs + power corrections

scalar



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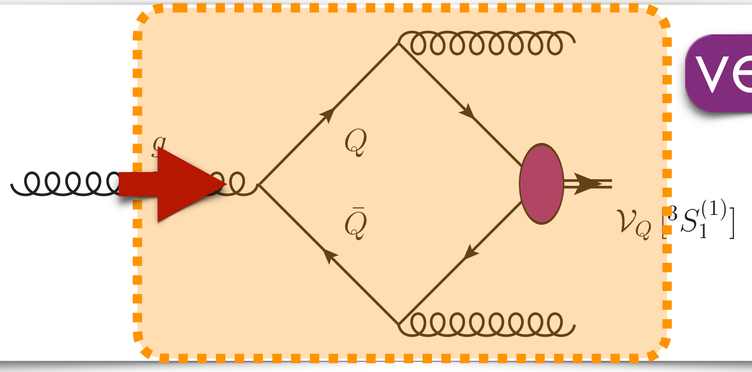


Collinear Fragmentation Correlator (CFC)

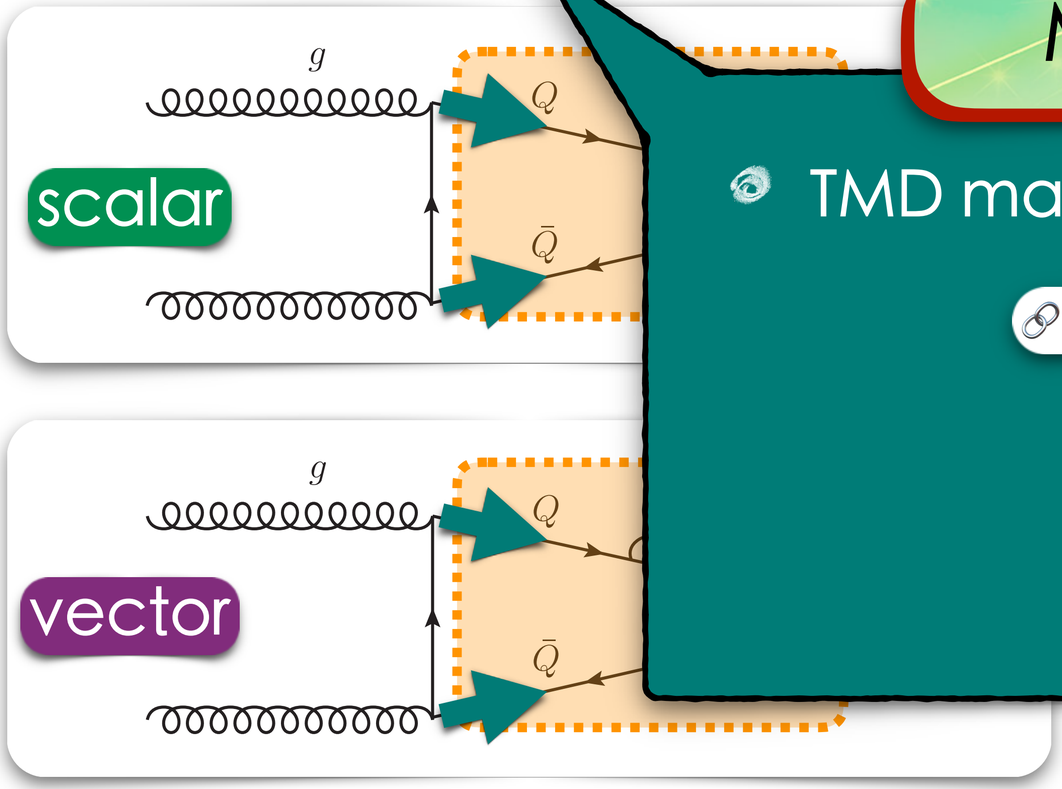
From low energies to high energies

Low p_T $\Rightarrow$ two-parton FFs

Genuine higher-twist FFNS FFs

High p_T $\Rightarrow$ single-parton FFs

Twist-2 VFNS FFs + power corrections



Matching low & high p_T

- TMD matching tails $\Rightarrow$ different singularities

[D. Boer, J. Bor, L. Maxia, C. Pisano, F. Yuan (2023)]

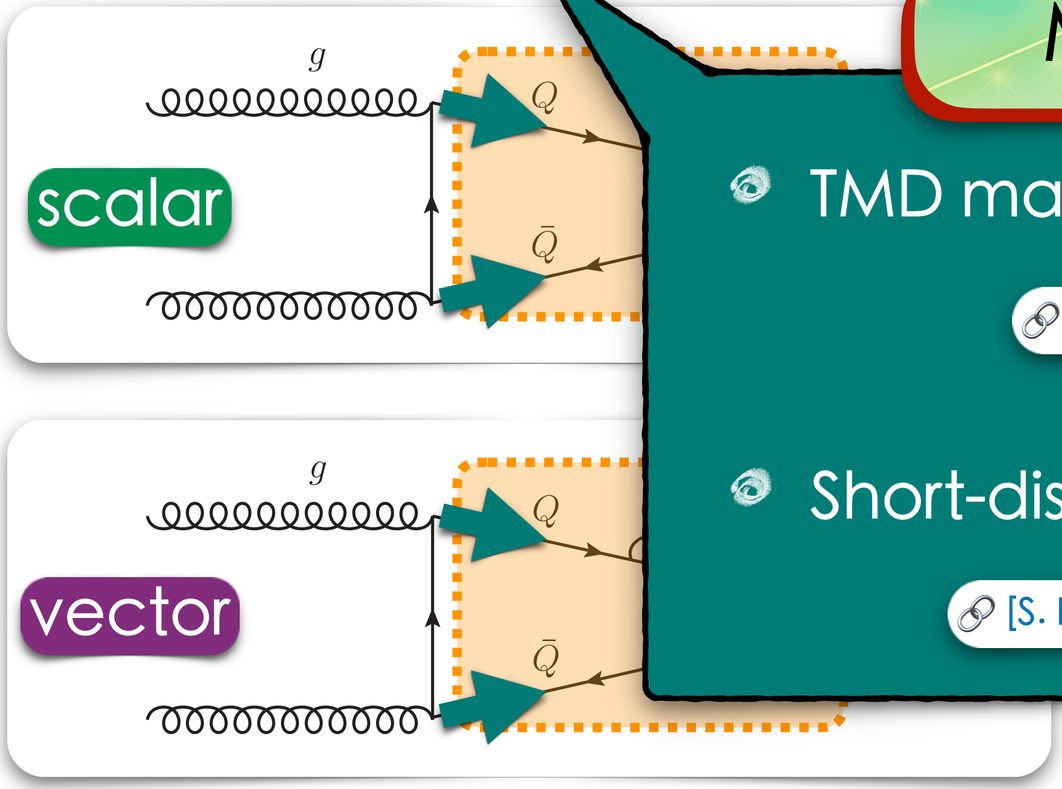
From low energies to high energies

Low p_T $\Rightarrow$ two-parton FFs

Genuine higher-twist FFNS FFs

High p_T $\Rightarrow$ single-parton FFs

Twist-2 VFNS FFs + power corrections



Matching low & high p_T

- TMD matching tails $\Rightarrow$ different singularities

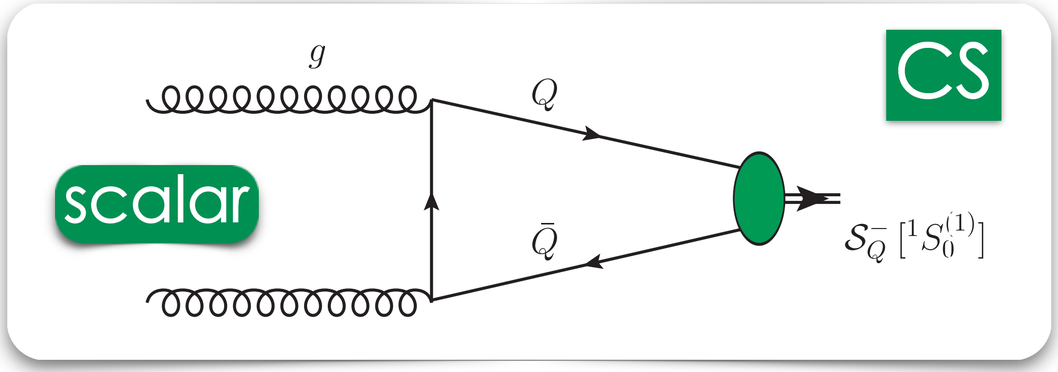
[D. Boer, J. Bor, L. Maxia, C. Pisano, F. Yuan (2023)]

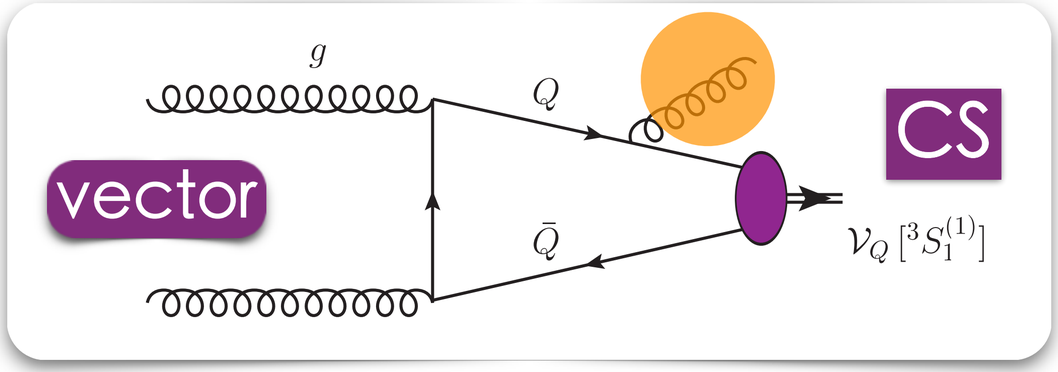
- Short-distance $\Leftrightarrow$ double-parton FFs (SCET)

[S. Fleming, A.K. Leibovich, T. Mehen, I.Z. Rothstein (2012)]

From low energies to high energies

Low energy, low p_T
Short-distance production



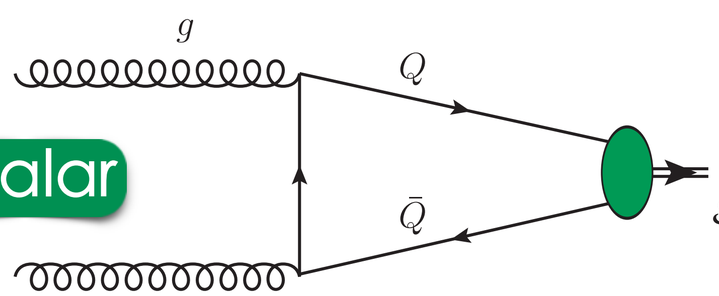


Extra gluon due to Landau-Yang selection rule

From low energies to high energies

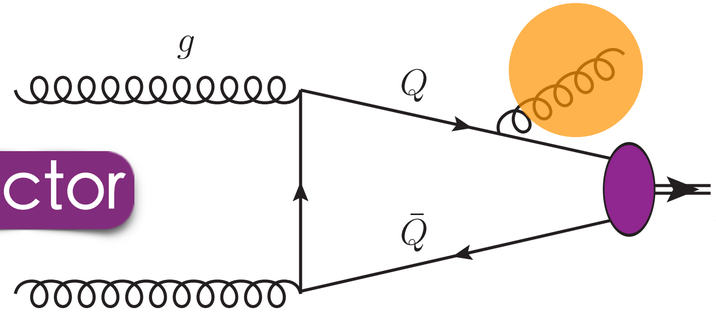
Low energy, low p_T
Short-distance production

scalar



CS

vector

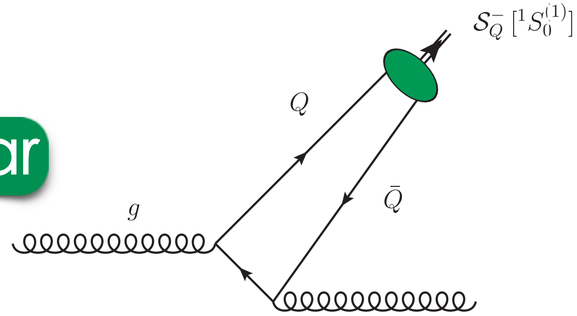


CS

$p_T \gtrsim 10 \div 15$ GeV
(charmonia)

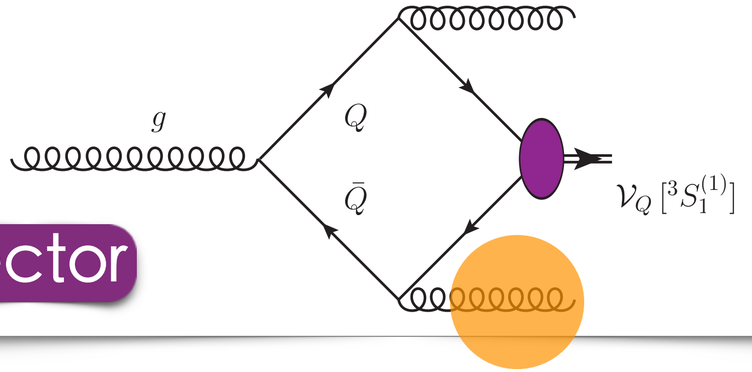
High energy, high p_T
Single-parton fragmentation

scalar



CS

vector



CS

Extra gluon due to Landau-Yang selection rule

HF-NRevo: Initial-scale NRCQD inputs

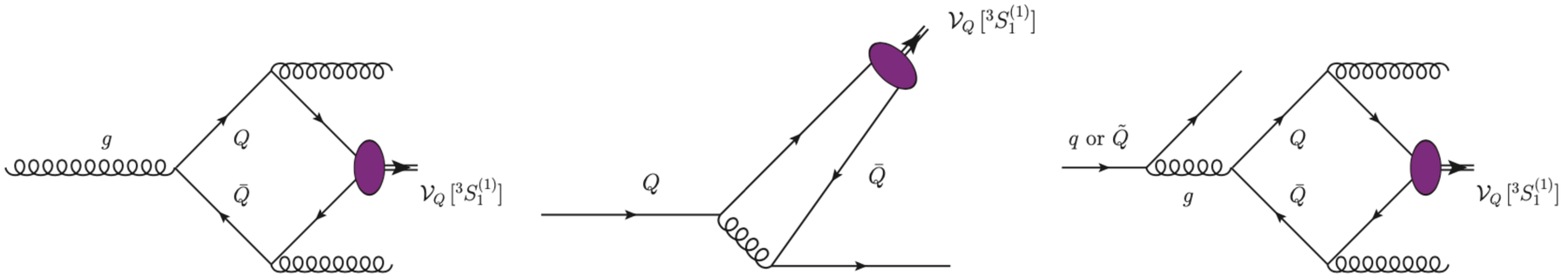


Fig. 1 Left panel: a leading diagram for the fragmentation of a gluon to $\mathcal{V}_Q\,[^3S_1^{(1)}]$ at $\mathcal{O}(\alpha_s^3)$. Central panel: one of the leading diagrams for the fragmentation of a constituent heavy quark (Q) to a S-wave color singlet vector quarkonium $\mathcal{V}_Q\,[^3S_1^{(1)}]$ at $\mathcal{O}(\alpha_s^2)$. Right panel: a leading diagram for the fragmentation of a nonconstituent light (q) or heavy ($\tilde{Q}$) quark to $\mathcal{V}_Q\,[^3S_1^{(1)}]$ at $\mathcal{O}(\alpha_s^4)$. Violet blobs refer to the $\mathcal{V}_Q\,[^3S_1^{(1)}]$ nonperturbative NRQCD LDME.

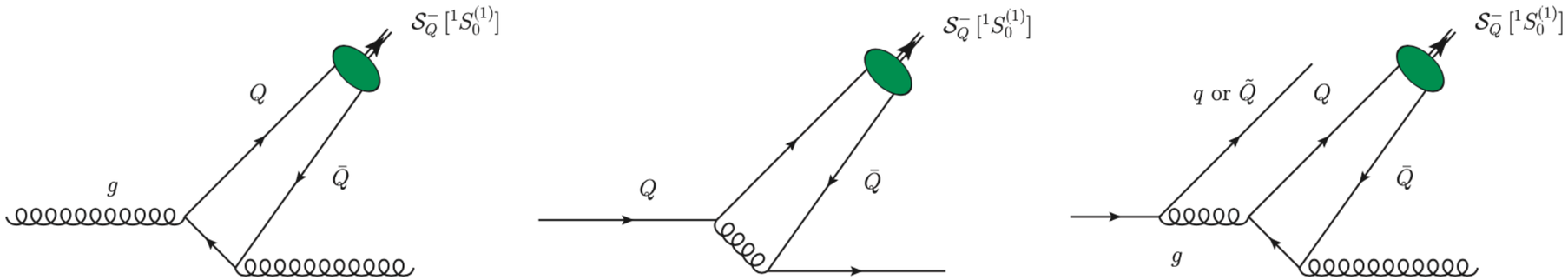
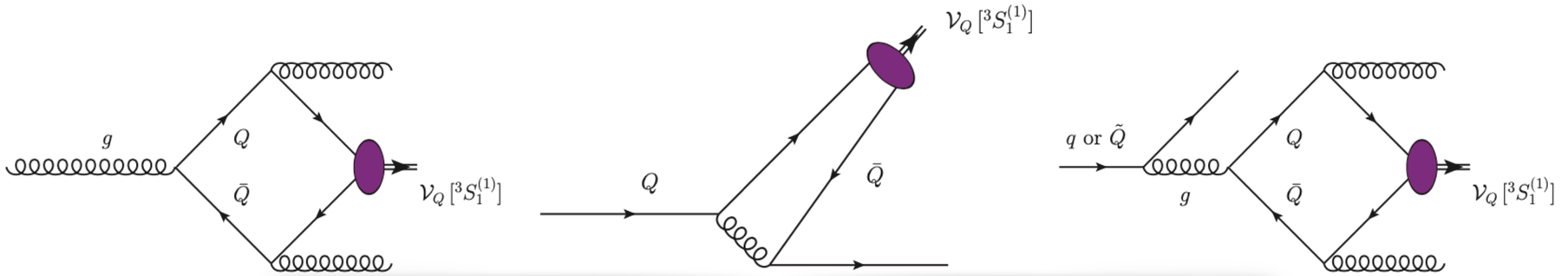


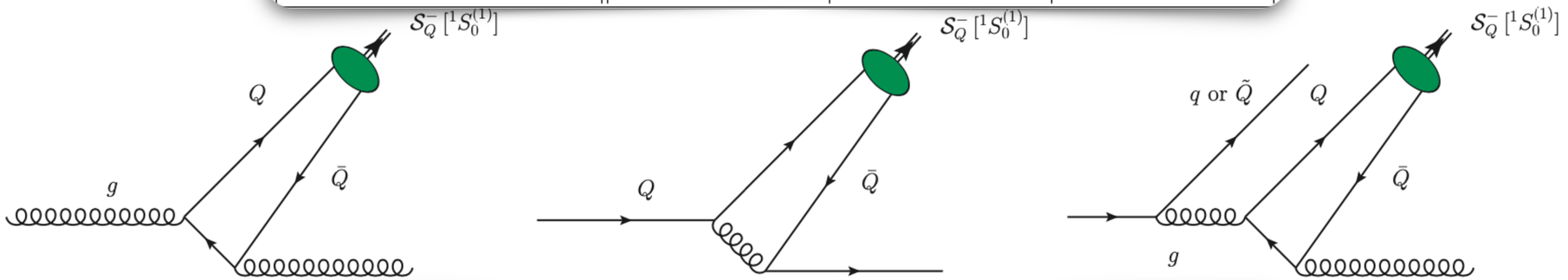
Fig. 2 Left panel: a leading diagram for the fragmentation of a gluon to $\mathcal{S}_Q^-\,[^1S_0^{(1)}]$ at $\mathcal{O}(\alpha_s^2)$. Central panel: one of the leading diagrams for the fragmentation of a constituent heavy quark (Q) to a S-wave color singlet pseudoscalar quarkonium $\mathcal{S}_Q^-\,[^1S_0^{(1)}]$ at $\mathcal{O}(\alpha_s^2)$. Right panel: a leading diagram for the fragmentation of a nonconstituent light (q) or heavy ($\tilde{Q}$) quark to $\mathcal{S}_Q^-\,[^1S_0^{(1)}]$ at $\mathcal{O}(\alpha_s^3)$. Green blobs refer to the $\mathcal{S}_Q^-\,[^1S_0^{(1)}]$ nonperturbative NRQCD LDME.

[F. G. C., in preparation]

HF-NRevo: Initial-scale NRCQD inputs



Quarkonium ($\mathcal{Q}$)	$g \to \mathcal{Q}$	$Q \to \mathcal{Q}$	q or $\tilde{Q} \to \mathcal{Q}$
$\mathcal{Q} \equiv \mathcal{V}_Q\,[^1S_0^{(1)}] \equiv J/\psi, \Upsilon$	–	$\mathcal{O}(\alpha_s^2) \equiv$ LO	–
	$\mathcal{O}(\alpha_s^3) \equiv$ NLO	$\mathcal{O}(\alpha_s^3) \equiv$ NLO	–
	$\mathcal{O}(\alpha_s^4) \equiv$ NNLO	$\mathcal{O}(\alpha_s^4) \equiv$ NNLO	$\mathcal{O}(\alpha_s^4) \equiv$ NNLO
	⋮	⋮	⋮

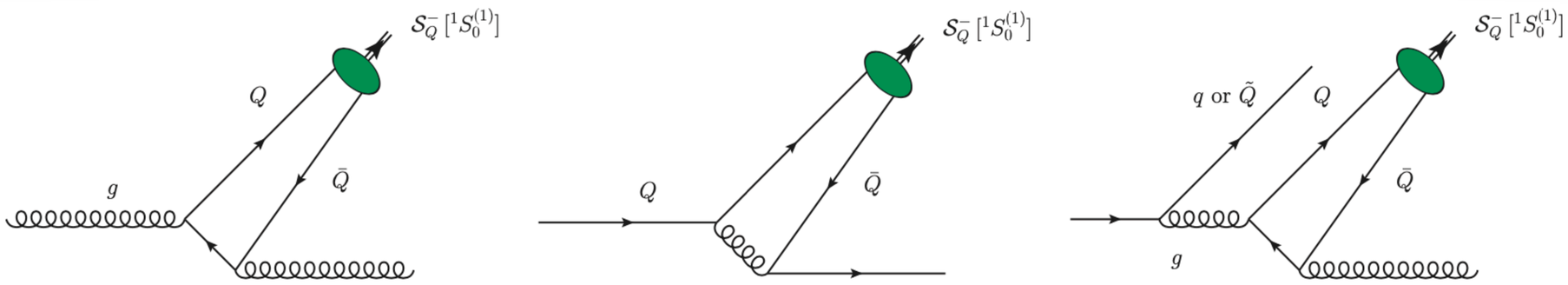


$\mathcal{Q} \equiv \mathcal{S}_Q^-\,[^3S_1^{(1)}] \equiv \eta_c, \eta_b$	$\mathcal{O}(\alpha_s^2) \equiv$ LO	$\mathcal{O}(\alpha_s^2) \equiv$ LO	–
	$\mathcal{O}(\alpha_s^3) \equiv$ NLO	$\mathcal{O}(\alpha_s^3) \equiv$ NLO	$\mathcal{O}(\alpha_s^3) \equiv$ NLO
	$\mathcal{O}(\alpha_s^4) \equiv$ NNLO	$\mathcal{O}(\alpha_s^4) \equiv$ NNLO	$\mathcal{O}(\alpha_s^4) \equiv$ NNLO
	⋮	⋮	⋮

[F. G. C., in preparation]

HF-NRevo: Initial-scale NRCQD inputs

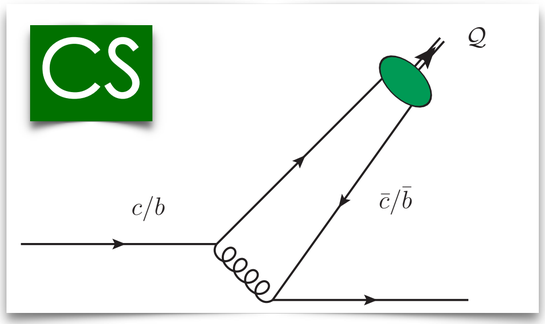
$\mathcal{S}_Q^-\,[^1S_0^{(1)}]$ quarkonium	$\mu_{F,0}(g \to \mathcal{S}_Q^-)$	$\mu_{F,0}(d,u,s \to \mathcal{S}_Q^-)$	$\mu_{F,0}(Q \to \mathcal{S}_Q^-)$	$\mu_{F,0}(\tilde{Q} \to \mathcal{S}_Q^-)$	$Q_0 \equiv \max(\{\mu_{F,0}\})$
η_c	$2m_c$	$m_{d,u,s} + 2m_c$	$3m_c$	$m_b + 2m_c$	$m_b + 2m_c$
η_b	$2m_b$	$m_{d,u,s} + 2m_b$	$3m_b$	$m_c + 2m_b$	$3m_b$



[F. G. C., in preparation]

Vector quarkonium from single-parton fragmentation

¡ *Let us consider* J/ψ *and* Υ *at large* p_T $\rightarrow$ *initial-scale* ***heavy-quark*** + ***gluon*** *from* ***NRQCD***!



$$D_{\mathcal{Q}}^{\mathcal{Q}}(z,\mu_F \equiv \mu_0) = D_{\mathcal{Q}}^{\mathcal{Q},\mathrm{LO}}(z) + \frac{\alpha_s^3(\mu_R)}{m_Q^3}\,|\mathcal{R}_{\mathcal{Q}}(0)|^2\,\Gamma^{\mathcal{Q},\mathrm{NLO}}(z)$$

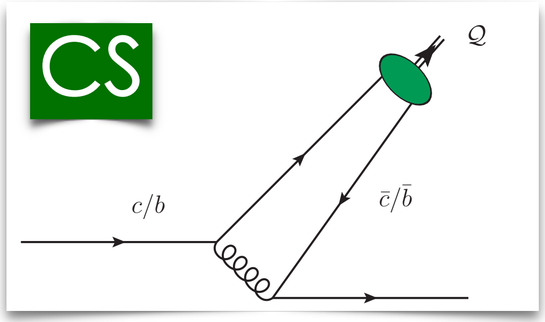
$(\mathcal{Q} \to \mathcal{Q}\,Q)$ at $\mu_0 = 3m_Q$

(LO) [E. Braaten et al., Phys. Rev. D 48 (1993) 4230-4235]
(NLO) [X. Zheng et al., Phys. Rev. D 100 (2019) 1, 014005]

[My talk at the Higgs Centre, Edinburgh (2023)]
[F. G. C., M. Fucilla, Eur. Phys. J. C 82 (2022) 10, 929]

Vector quarkonium from single-parton fragmentation

¡ *Let us consider* J/ψ *and* Υ *at large* p_T $\rightarrow$ *initial-scale* ***heavy-quark*** + ***gluon*** *from* ***NRQCD*** !



$$D_{\mathcal{Q}}^{\mathcal{Q}}(z,\mu_F \equiv \mu_0) = D_{\mathcal{Q}}^{\mathcal{Q},\mathrm{LO}}(z) + \frac{\alpha_s^3(\mu_R)}{m_Q^3}\,|\mathcal{R}_{\mathcal{Q}}(0)|^2\,\Gamma^{\mathcal{Q},\mathrm{NLO}}(z)$$

$(Q \rightarrow \mathcal{Q}\,Q)$ at $\mu_0 = 3m_Q$

(LO) 🔗 [E. Braaten et al., Phys. Rev. D 48 (1993) 4230-4235]

(NLO) 🔗 [X. Zheng et al., Phys. Rev. D 100 (2019) 1, 014005]

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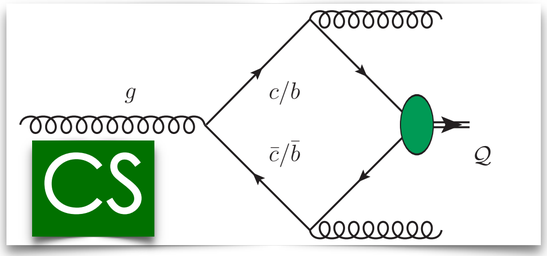


$$D_g^{\mathcal{Q}}(z,2m_Q) = \frac{5}{36(2\pi)^2}\alpha_s^3(2m_Q)\frac{|\mathcal{R}_{\mathcal{Q}}(0)|^2}{m_{\mathcal{Q}}^3}\int_0^z \mathrm{d}\xi \int_{(\xi+z^2)/2z}^{(1+\xi)/2} \mathrm{d}\tau \frac{1}{(1-\tau)^2(\tau-\xi)^2(\tau^2-\xi)^2}$$

$$\sum_{i=1}^{2} z^i \left[f_i^{(g)}(\xi,\tau) + g_i^{(g)}(\xi,\tau)\frac{1+\xi-2\tau}{2(\tau-\xi)\sqrt{\tau^2-\xi}} \ln\left(\frac{\tau-\xi+\sqrt{\tau^2-\xi}}{\tau-\xi-\sqrt{\tau^2-\xi}}\right)\right],$$

$(g \rightarrow \mathcal{Q}\,gg)$ at $\mu_0 = 2m_Q$

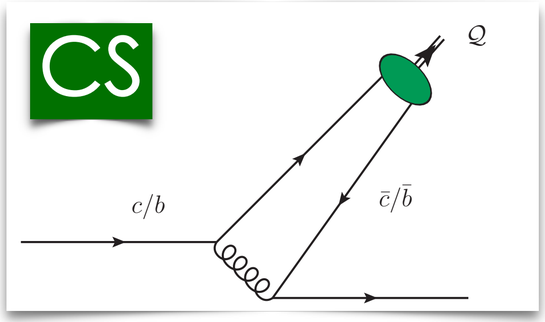
(LO) 🔗 [A. Braaten, T.C Yuan, Phys. Rev. Lett. 71 (1993), 1673]

🔗 [My talk at the Higgs Centre, Edinburgh (2023)]

🔗 [F. G. C., M. Fucilla, Eur. Phys. J. C 82 (2022) 10, 929]

Vector quarkonium from single-parton fragmentation

¡ *Let us consider* J/ψ *and* Υ *at large* p_T $\rightarrow$ *initial-scale* ***heavy-quark*** + ***gluon*** *from* ***NRQCD*** !

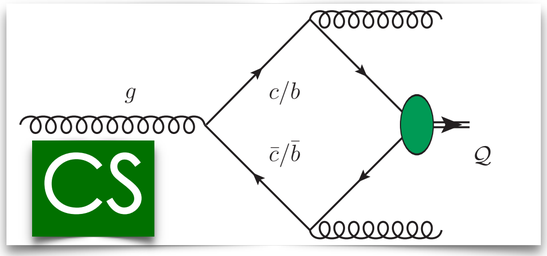


$$D_{\mathcal{Q}}^{\mathcal{Q}}(z,\mu_F \equiv \mu_0) = D_{\mathcal{Q}}^{\mathcal{Q},\mathrm{LO}}(z) + \frac{\alpha_s^3(\mu_R)}{m_Q^3}\,|\mathcal{R}_{\mathcal{Q}}(0)|^2\,\Gamma^{\mathcal{Q},\mathrm{NLO}}(z)$$

$(Q \to \mathcal{Q}\,Q)$ at $\mu_0 = 3m_Q$

(LO) 🔗 [E. Braaten et al., Phys. Rev. D 48 (1993) 4230-4235]
(NLO) 🔗 [X. Zheng et al., Phys. Rev. D 100 (2019) 1, 014005]

+



$$D_g^{\mathcal{Q}}(z,2m_Q) = \frac{5}{36(2\pi)^2}\alpha_s^3(2m_Q)\frac{|\mathcal{R}_{\mathcal{Q}}(0)|^2}{m_{\mathcal{Q}}^3}\int_0^z \mathrm{d}\xi \int_{(\xi+z^2)/2z}^{(1+\xi)/2} \mathrm{d}\tau \frac{1}{(1-\tau)^2(\tau-\xi)^2(\tau^2-\xi)^2}$$

$$\sum_{i=1}^{2} z^i \left[f_i^{(g)}(\xi,\tau) + g_i^{(g)}(\xi,\tau)\frac{1+\xi-2\tau}{2(\tau-\xi)\sqrt{\tau^2-\xi}} \ln\left(\frac{\tau-\xi+\sqrt{\tau^2-\xi}}{\tau-\xi-\sqrt{\tau^2-\xi}}\right)\right],$$

$(g \to \mathcal{Q}\,gg)$ at $\mu_0 = 2m_Q$

$\otimes$

(LO) 🔗 [A. Braaten, T.C Yuan, Phys. Rev. Lett. 71 (1993), 1673]

🔗 [My talk at the Higgs Centre, Edinburgh (2023)]
🔗 [F. G. C., M. Fucilla, Eur. Phys. J. C 82 (2022) 10, 929]

NRFF1.0
quarkonium FFs

=

APFEL++/EKO

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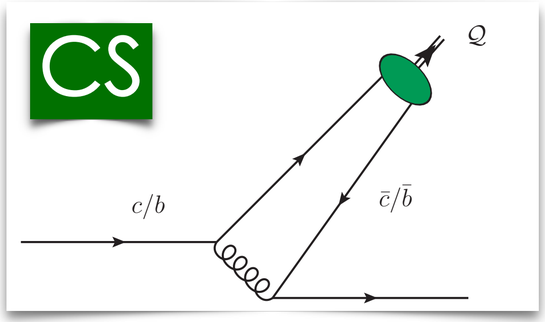

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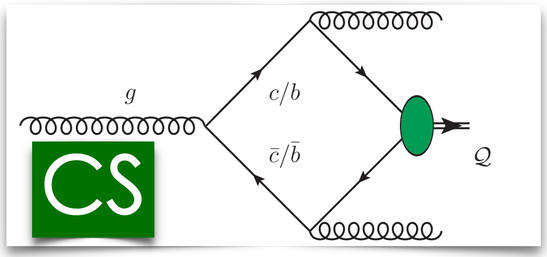
+



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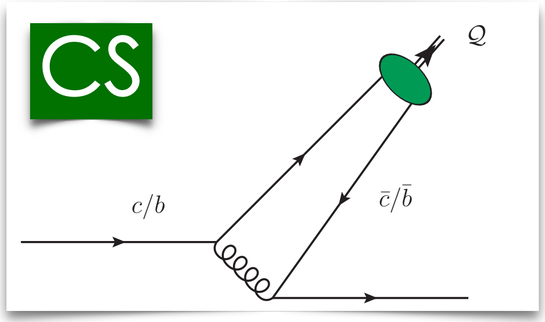

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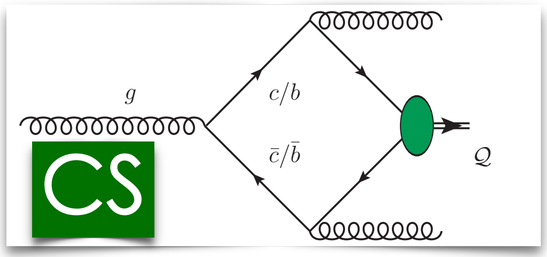

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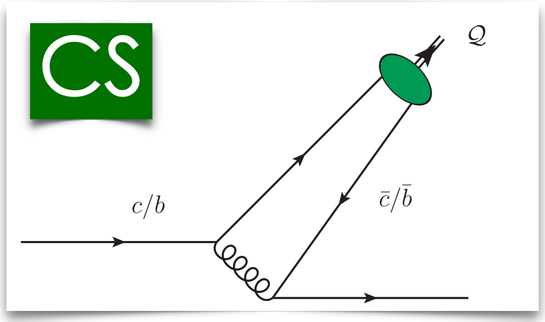

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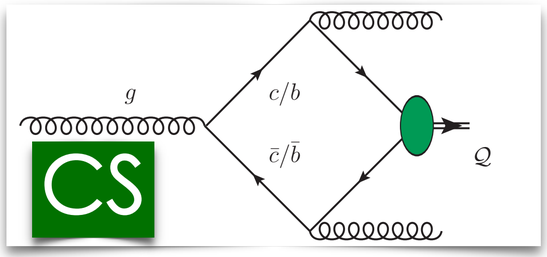

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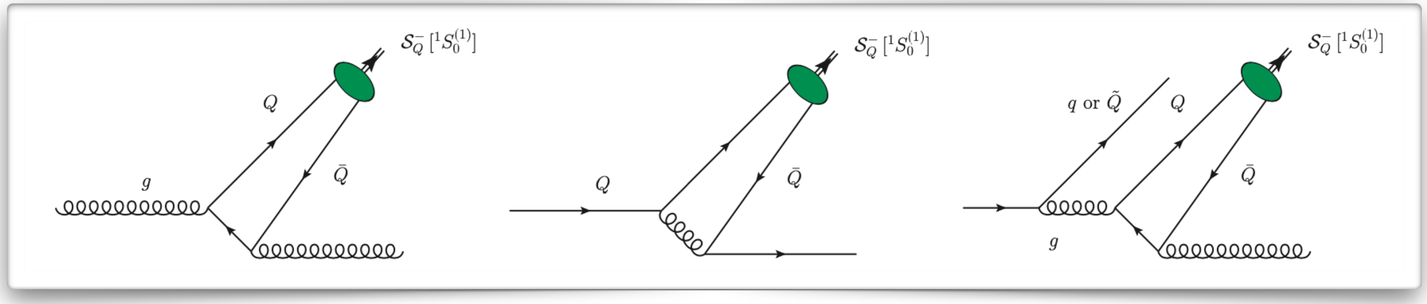
HF-NRevo: Initial-scale NRCQD inputs



Scalar η_c

HF-NRevo for $\eta_c[^1S_0^{(1)}]$ collinear fragmentation

Partonic hard factor

DGLAP AOevo

All-Order evolution (numeric)

DGLAP EDevo

Expanded Decoupled evolution

NRQCD

LDMEs

SDCs

g
$\tilde{q}$
c
b
$2m_c$

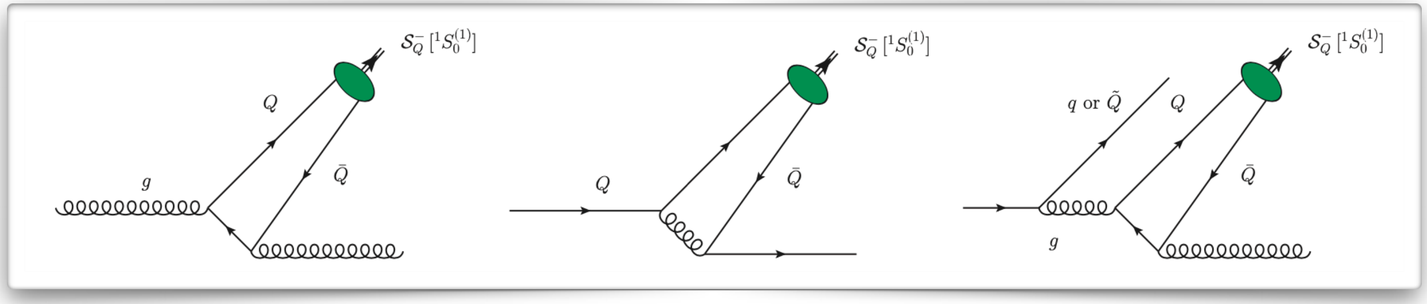
$\mathcal{D}_i^{\mathcal{Q}}(z,[n]) \otimes \langle \mathcal{O}^{\mathcal{Q}}([n]) \rangle$

Energy

Evolution-ready scale $Q_0 \equiv m_b + 2m_c \quad 3m_c \quad m_{\tilde{q}} + 2m_c$

Time

HF-NRevo: Initial-scale NRCQD inputs



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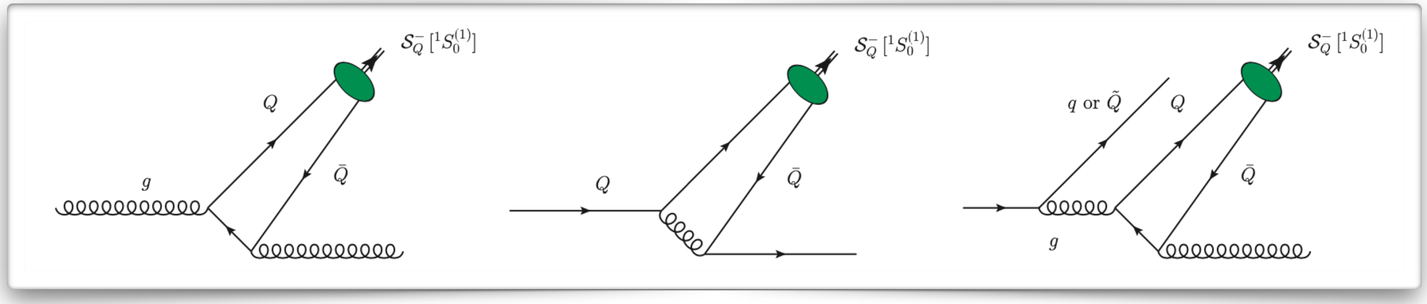

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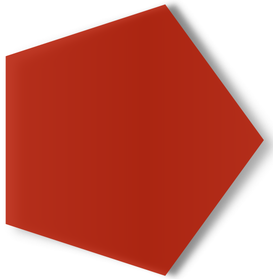
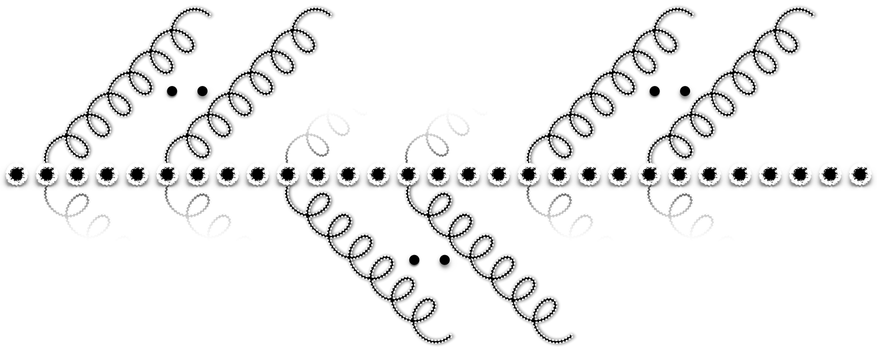
Partonic hard factor



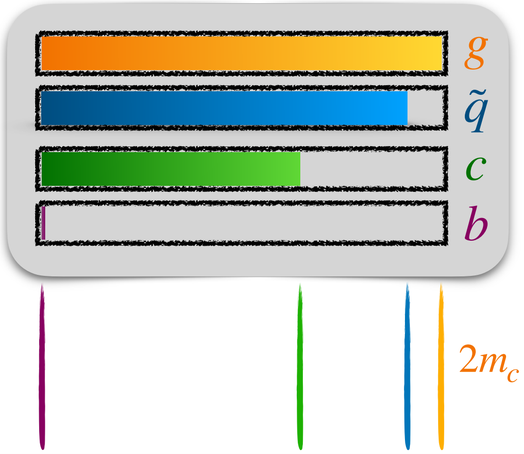
DGLAP AOevo

All-Order evolution (numeric)



DGLAP EDevo

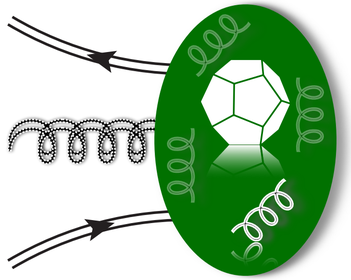
Expanded Decoupled evolution



NRQCD

LDMEs

SDCs



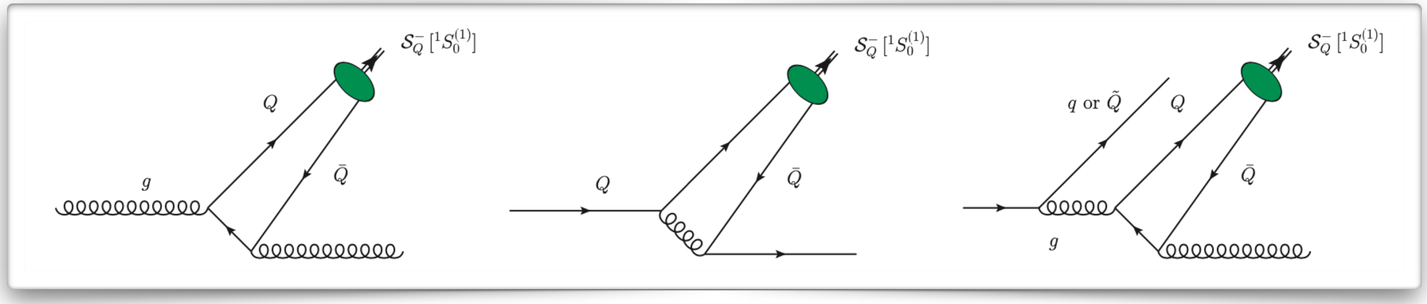
$\mathcal{D}_i^{\mathcal{Q}}(z,[n]) \otimes \langle \mathcal{O}^{\mathcal{Q}}([n]) \rangle$

Energy

Evolution-ready scale $Q_0 \equiv m_b + 2m_c$ $3m_c$ $m_{\tilde{q}} + 2m_c$

Time

HF-NRevo: Initial-scale NRCQD inputs



Scalar η_b

HF-NRevo for $\eta_b[^1S_0^{(1)}]$ collinear fragmentation

Partonic hard factor

DGLAP AOevo

All-Order evolution (numeric)

DGLAP EDevo

Expanded Decoupled evolution

g
$\tilde{q}$
c
b
$2m_b$

NRQCD

SDCs

LDMEs

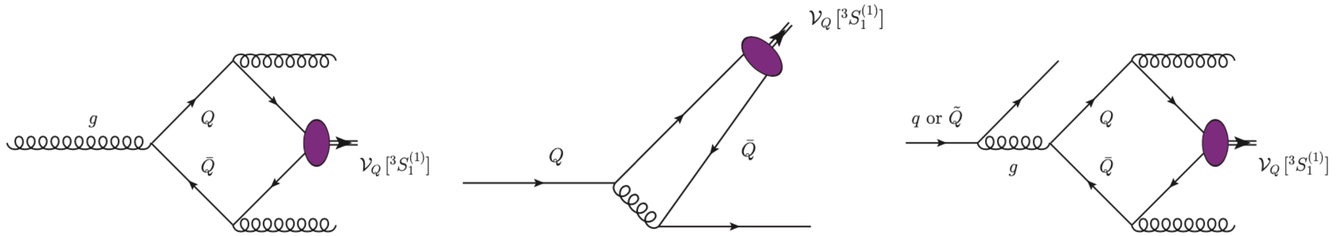
$\mathcal{D}_i^{\mathcal{Q}}(z,[n]) \otimes \langle \mathcal{O}^{\mathcal{Q}}([n]) \rangle$

Energy

Evolution-ready scale $Q_0 \equiv 3m_b$ $m_c + 2m_b$ $m_{\tilde{q}} + 2m_b$

Time

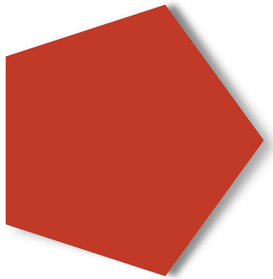
HF-NRevo: Initial-scale NRCQD inputs



Vector J/ψ

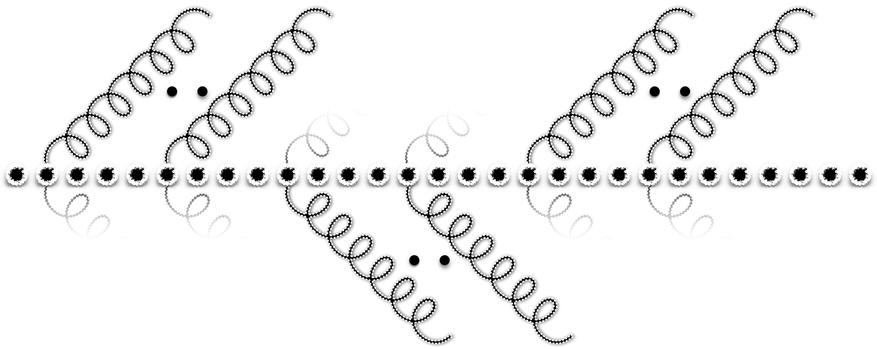
HF-NRevo for $J/\psi[^3S_1^{(1)}]$ collinear fragmentation

Partonic hard factor



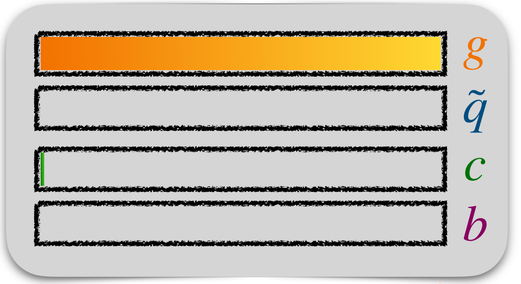
DGLAP AOevo

All-Order evolution (numeric)



DGLAP EDevo

Expanded Decoupled evolution

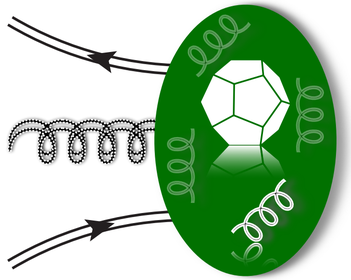


$2m_c$

NRQCD

LDMEs

SDCs



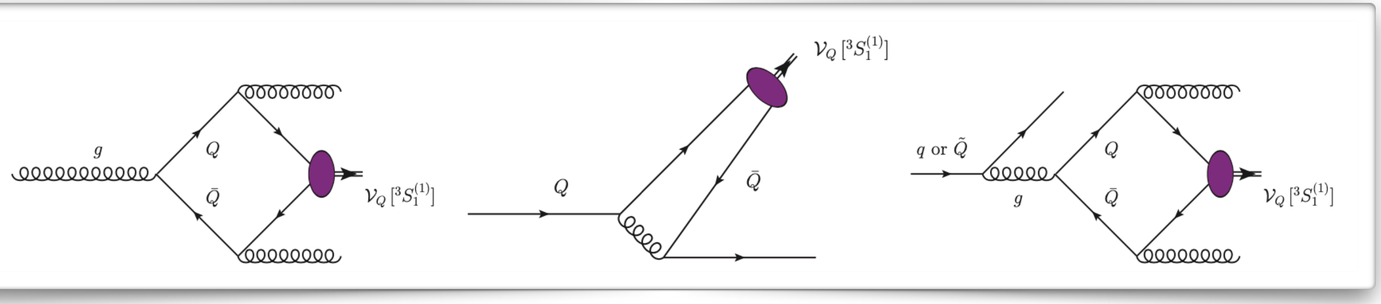
$\mathcal{D}_i^{\mathcal{Q}}(z,[n]) \otimes \langle \mathcal{O}^{\mathcal{Q}}([n]) \rangle$

Energy

Evolution-ready scale $Q_0 \equiv 3m_c$

Time

HF-NRevo: Initial-scale NRCQD inputs



Vector Υ

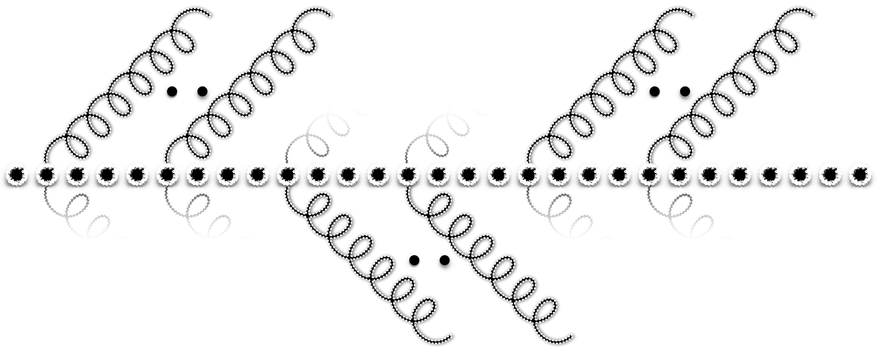
HF-NRevo for $\Upsilon[^3S_1^{(1)}]$ collinear fragmentation

Partonic hard factor



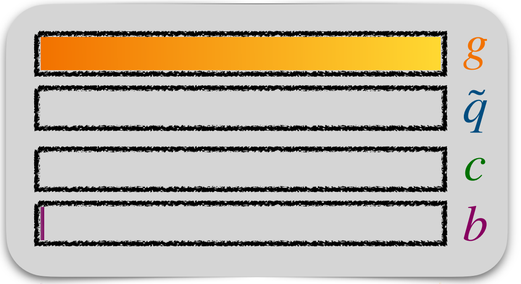
DGLAP AOevo

All-Order evolution (numeric)



DGLAP EDevo

Expanded Decoupled evolution

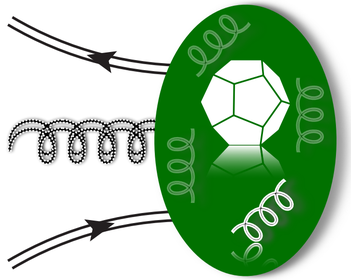


$2m_c$

NRQCD

LDMEs

SDCs



$\mathcal{D}_i^{\mathcal{Q}}(z,[n]) \otimes \langle \mathcal{O}^{\mathcal{Q}}([n]) \rangle$

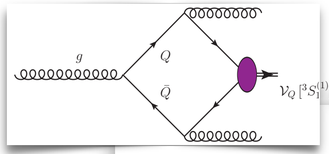
Energy

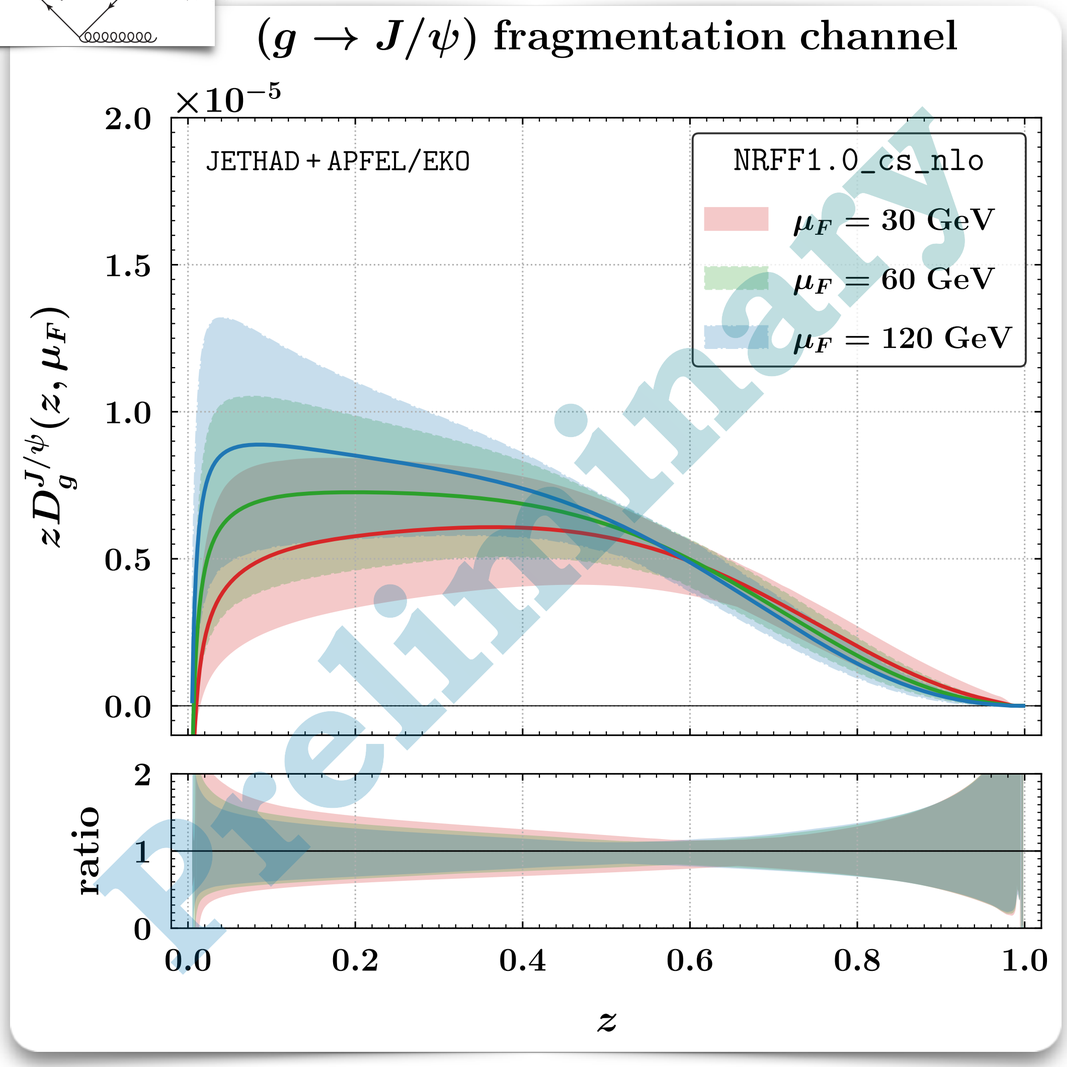
Evolution-ready scale $Q_0 \equiv 3m_b$

Time

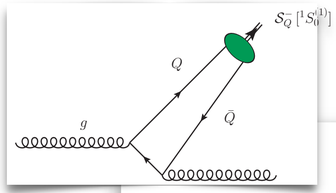
NRFF1.0: Gluon fragmentation to charmonia

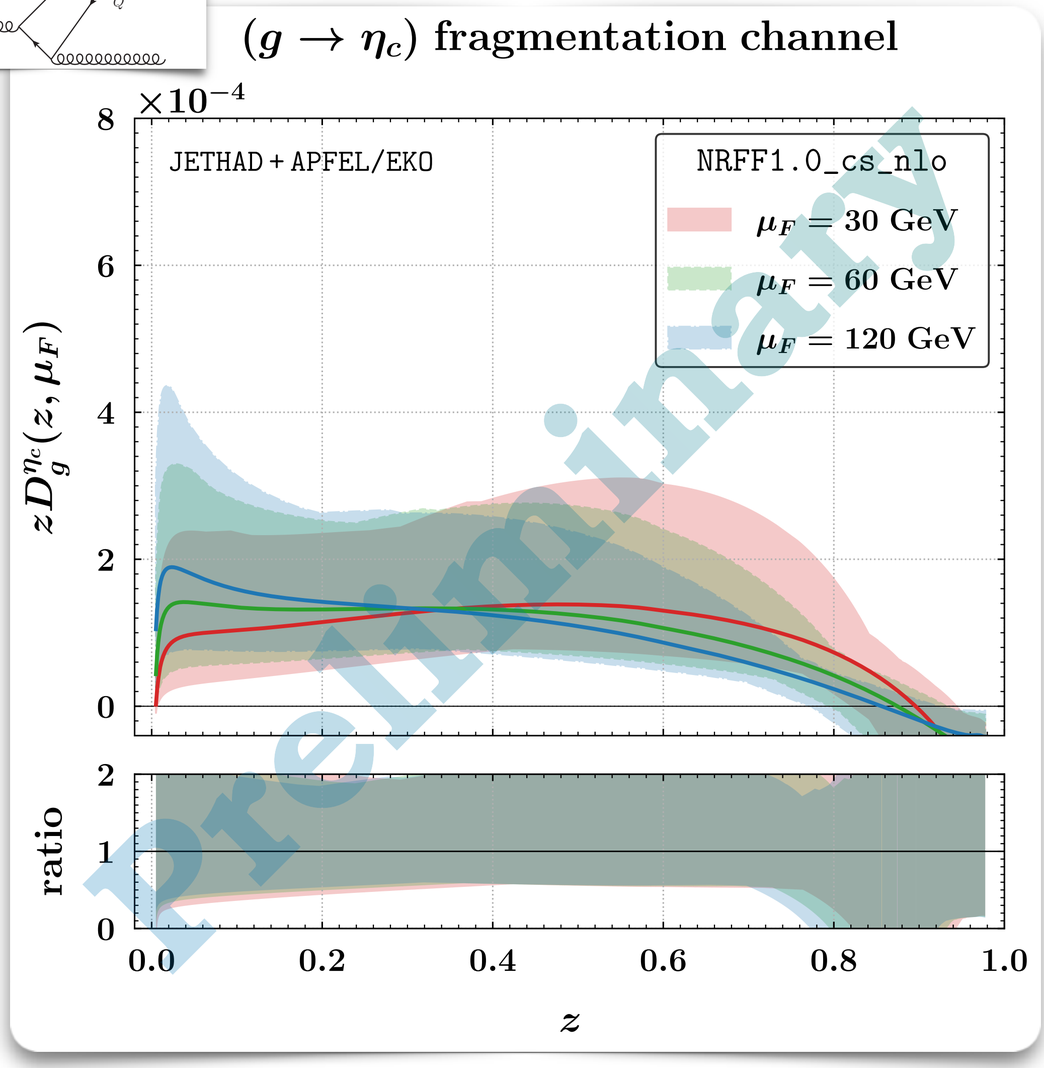
Vector J/ψ





Scalar η_c

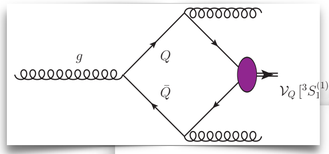


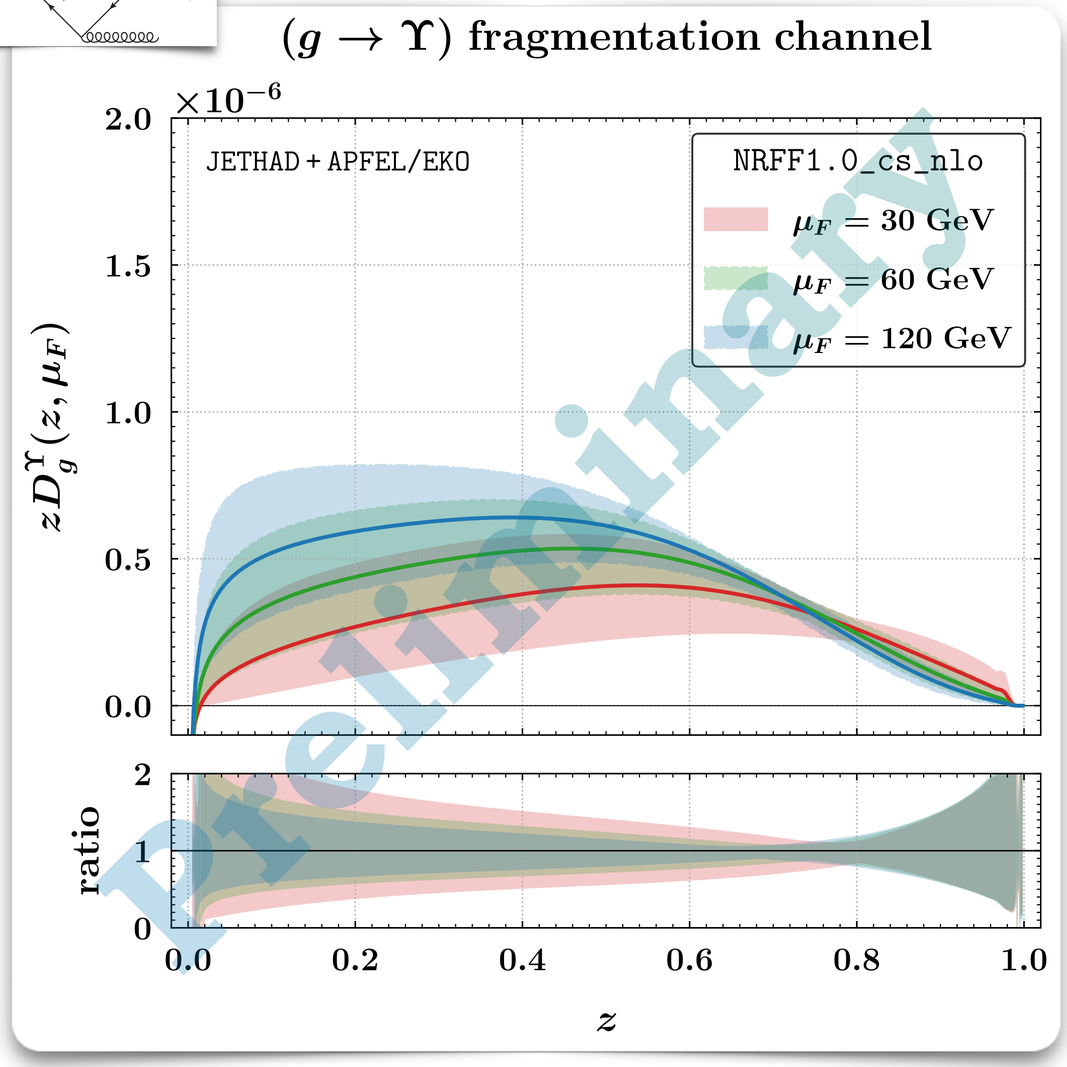


[F. G. C., in preparation]

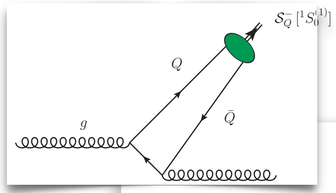
NRFF1.0: Gluon fragmentation to bottomonia

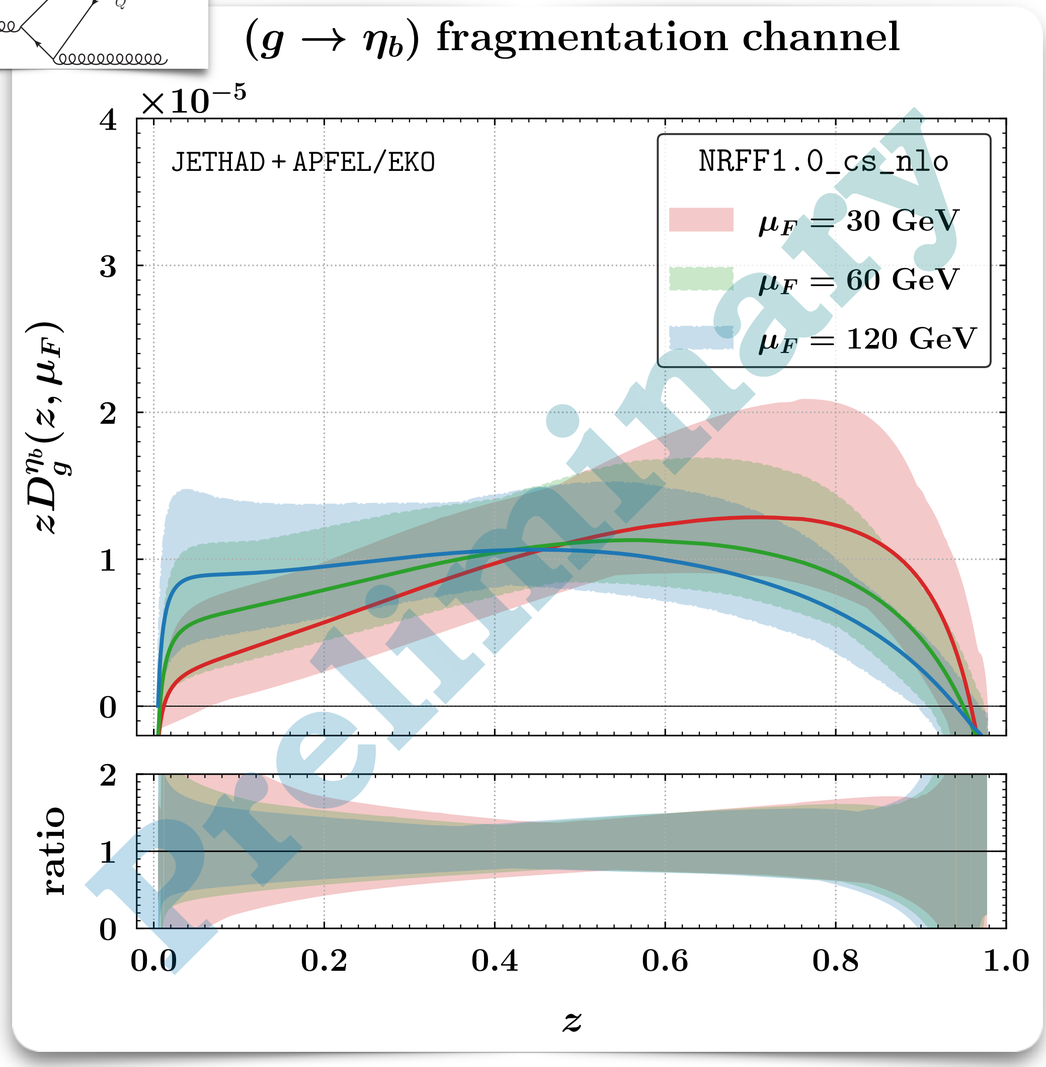
Vector Υ





Scalar η_b

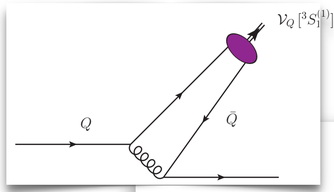


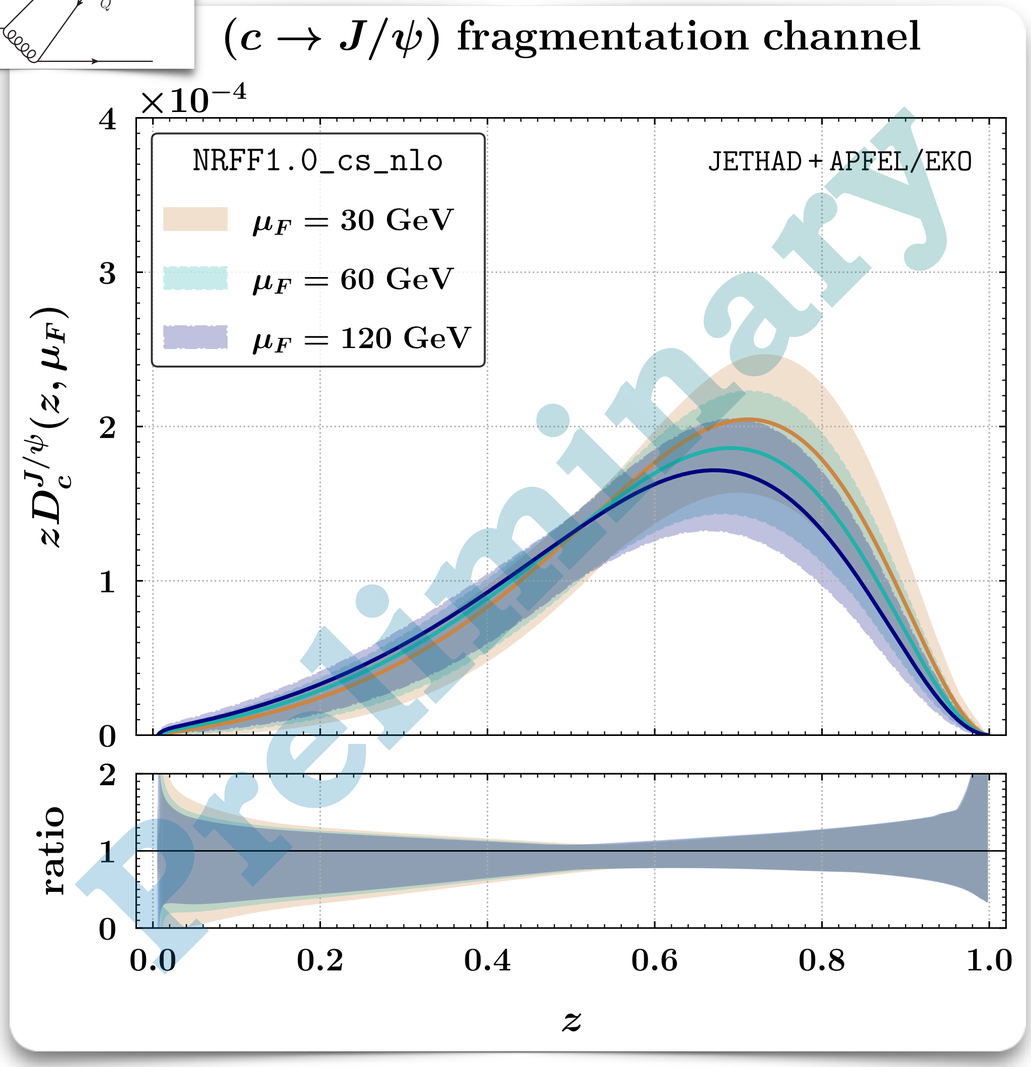


[F. G. C., in preparation]

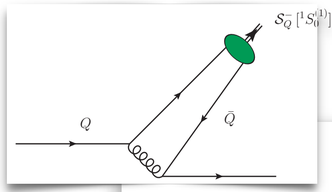
NRFF1.0: Charm fragmentation to charmonia

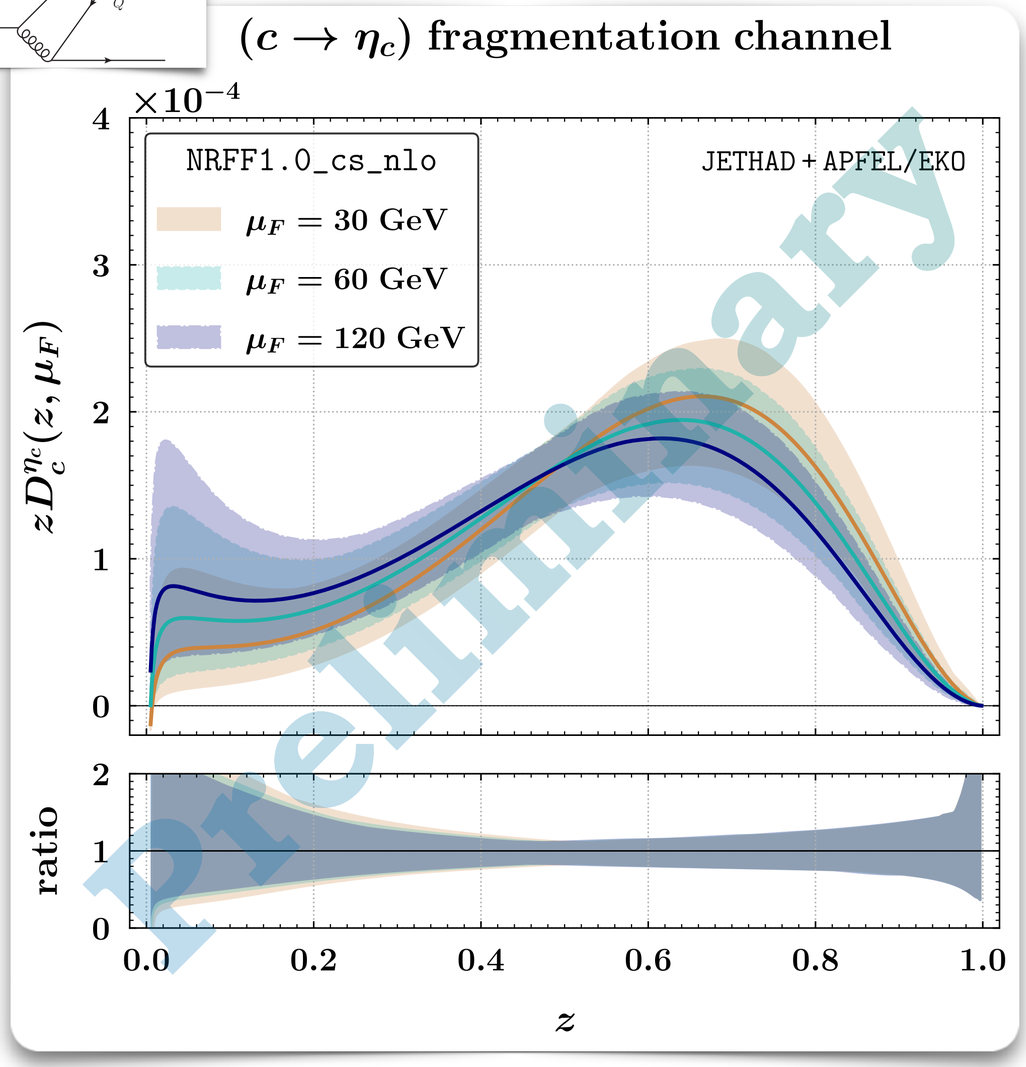
Vector J/ψ





Scalar η_c



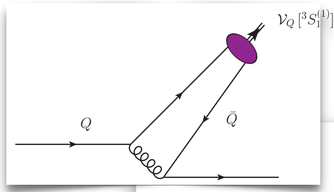


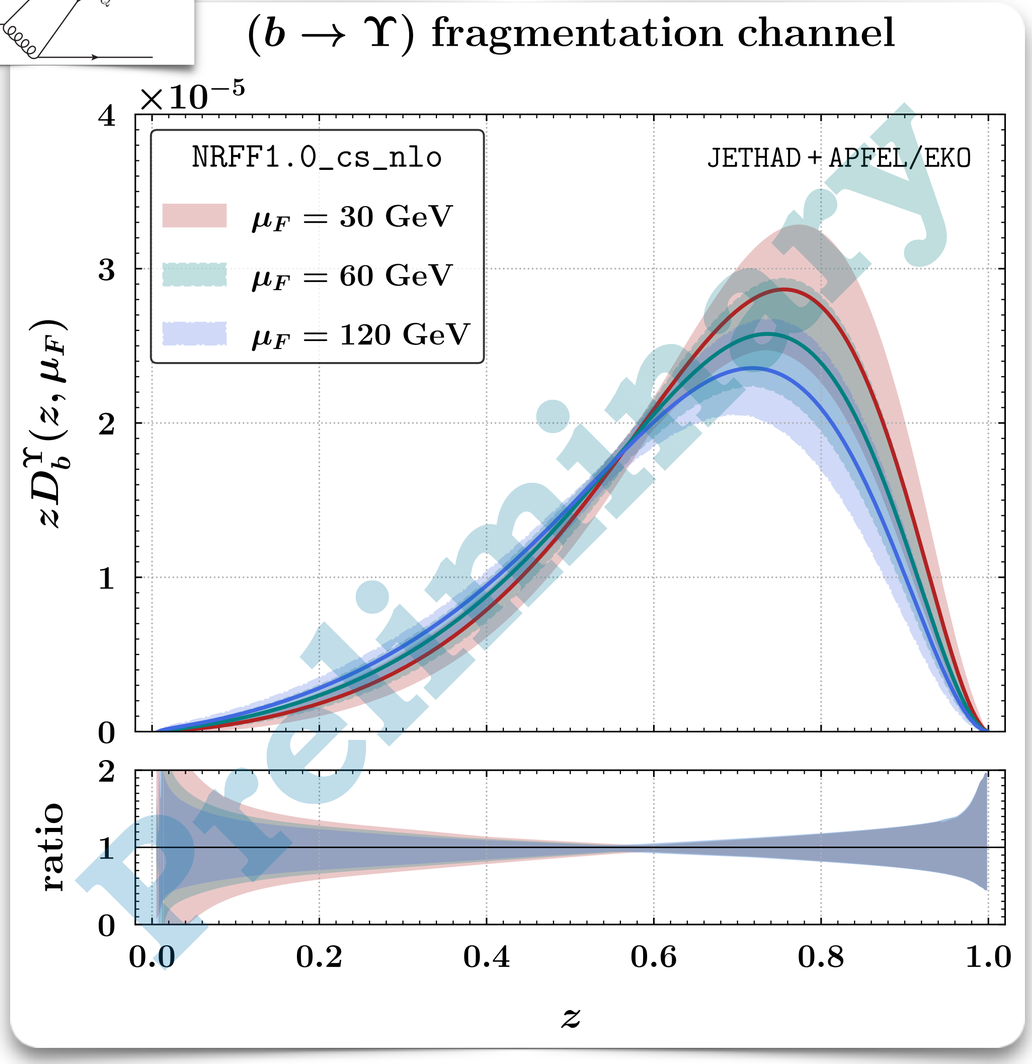
Preliminary

[F. G. C., in preparation]

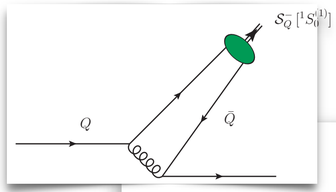
NRFF1.0: Bottom fragmentation to bottomonia

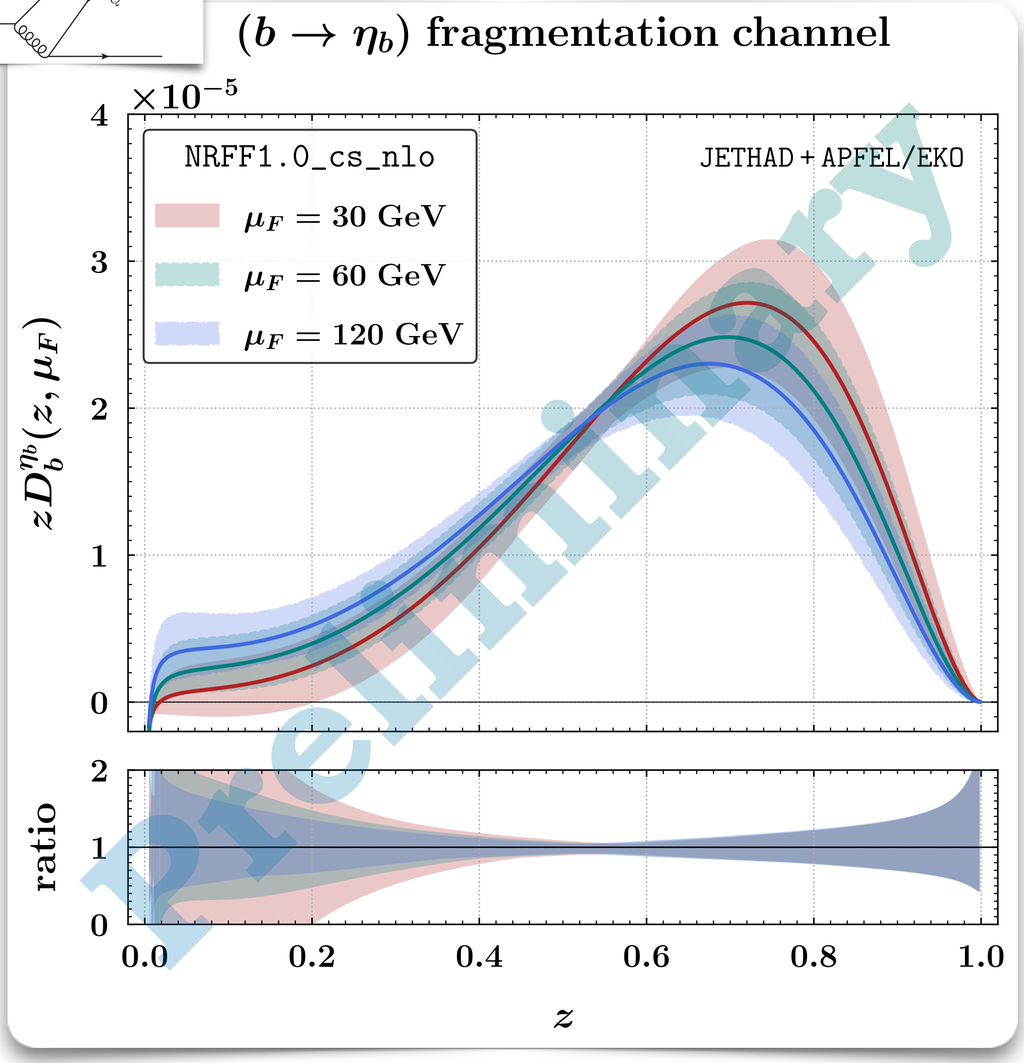
Vector Υ





Scalar η_b





Preliminary

[F. G. C., in preparation]