

multi-loop methods and application: a status report

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thanks to my collaborators:

Becchetti, Brancaccio, Czakon, Hartanto, Peraro,
Poncelet, Ripani, Zoia



Radiative Corrections and Monte
Carlo Simulations at Electron-
Positron Colliders

Turin, 4th June 2026



Goal of this talk

Look ahead to see how we might reach the precision needs of current and future e^+e^- colliders
- low and high energy

Technology for underlying amplitude computations is the same independent of the collider

Different phase spaces can have a big effect on practical applications for precision phenomenology

Let's look at a few state-of-the art examples as motivation for discussion

**other talks this week:
Rocco, Pozzoli, Dave**

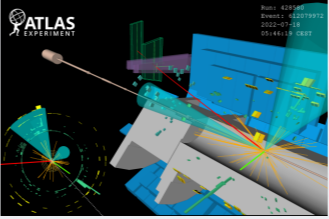
IBPs

DEQs

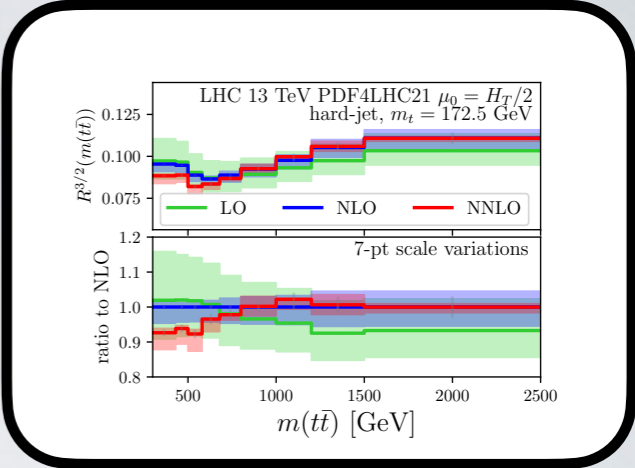
modular arithmetic

geometry/
mathematical structure

$$A = C.I$$



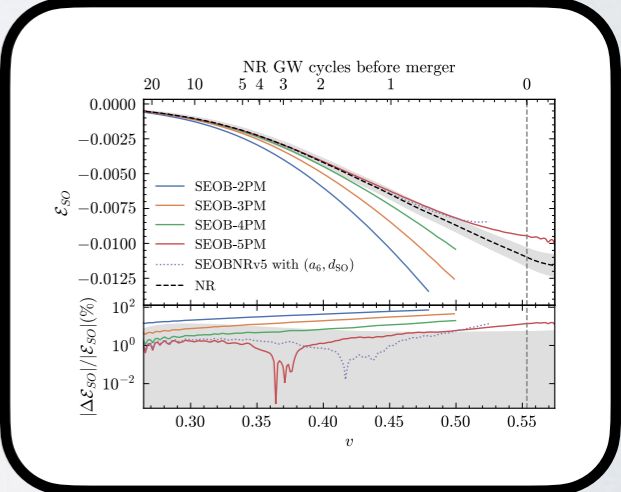
tt event



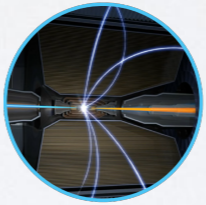
2511.11431



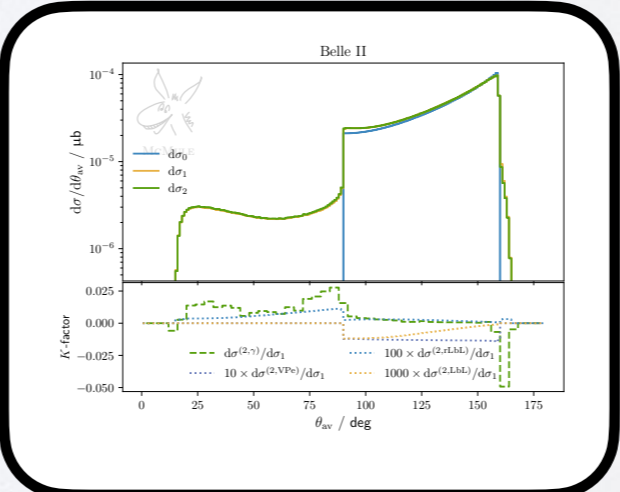
Ligo BH merger



2405.19181



Belle 2



2512.22929

perturbative precision

one-loop/NLO fully automated!*

current frontier



better measure of complexity would be number of scales

internal massive propagators: complex underlying integral geometries lead to rapid growth in complexity (elliptic, K3, Calabi Yau)

$$\mathcal{A} \rightarrow d\sigma$$

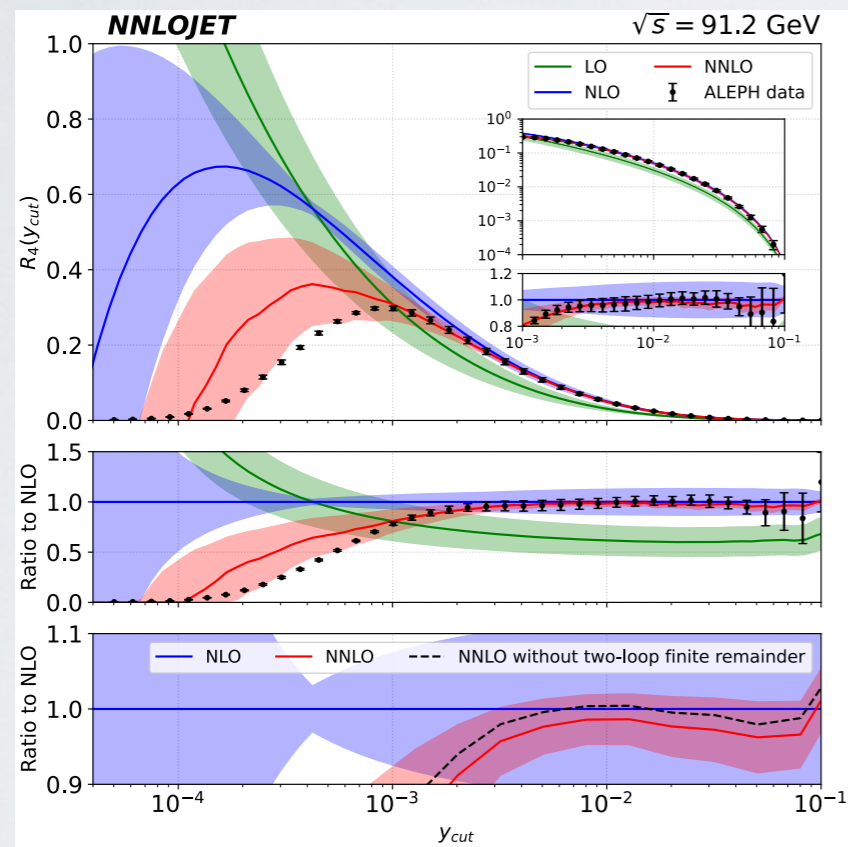
combining multi-loop multi-scale amplitudes with IR subtraction and challenging phase space

$e^+e^- \rightarrow \text{jets}$

new results from antenna subtraction w/ NNLOJET

$e^+e^- \rightarrow 4 \text{ jets @ NNLO}$

[Chen et al. 2602.18185](#)



RV and RR small in this case but dominate computation time

compact two-loop leading colour amplitudes

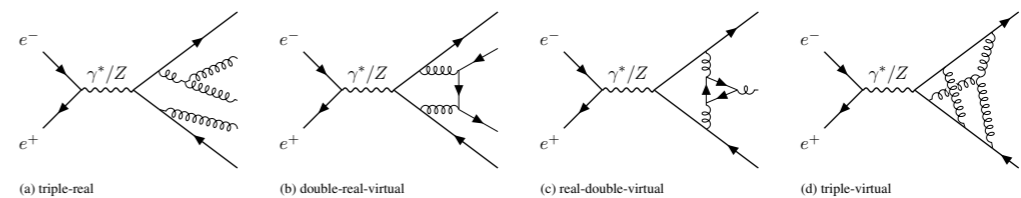
[De Laurentis et al. 2503.10595](#)

fast stable evaluation with (analytically continued) pentagon functions

[Abreu et al. 2306.15431](#)

$e^+e^- \rightarrow 2 \text{ jets @ N3LO}$

[Chen et al. 2505.10618](#)



The numerical evaluation of these matrix elements in double-unresolved configurations at one loop and single-unresolved limits at two loops required careful treatment to achieve the necessary accuracy. In double-unresolved limits of one-loop ma-

example amplitude toolchain

tensor integral input

Feynman diagrams

Projectors [in 4D Tancredi, Peraro]

Unitarity cuts

tensor reduction to master integrals

integration-by-parts identities [Tkachov, Chetykin] [Laporta]

solve over finite fields e.g. FiniteFlow, FireFly, FIRE, Kira

optimized systems (Blade, NeatIBP etc.)

master integral expansions

identity function basis via differential equations for master integrals

canonical 'ε' -factorized form (Henn)

$$A = C.I$$



$$\mathcal{F} = c.f$$

finite remainder

analytic reconstruction

extract rational coefficients through multiple finite field evaluations

(in parallel) [Thiele interpolation]

optimize with linear relations, partial fractioning etc.

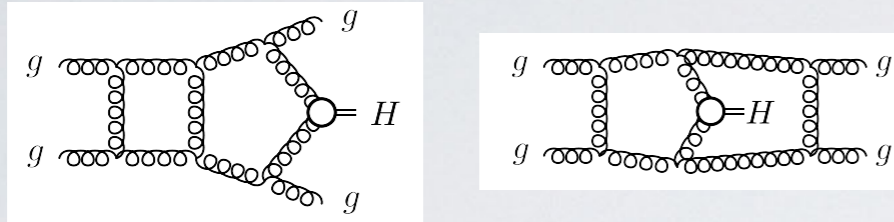
Successful program at 2 loops: $pp \rightarrow 3j$, $pp \rightarrow 3\gamma$, $pp \rightarrow \gamma\gamma j$,

$pp \rightarrow \gamma jj$, $pp \rightarrow W+2j$, $pp \rightarrow Wbb$, $pp \rightarrow Hbb$, $pp \rightarrow W\gamma\gamma$

these cases have massless propagators...

pp \rightarrow H+2 jets

De Laurentis et al. 2605.04009



leading colour, heavy top quark limit

numerical unitarity method over finite fields (Caravel), IBPs with surface term decomposition (syzygy)

analytic reconstruction possible using optimized ansatz in terms of spinor helicity variables and multivariate partial fractioning

see also pp \rightarrow bbH Hartanto, Poncelet 2603.29480

	$\mathcal{A}_{0 \rightarrow \bar{b}bggH}^{(2),N_c^2}$	$\mathcal{A}_{0 \rightarrow \bar{b}bggH}^{(2),N_c n_f}$	$\mathcal{A}_{0 \rightarrow \bar{b}bqgH}^{(2),N_c^2}$
$x_1 = 1$	256/251	166/161	145/140
linear relations	187/183	157/153	118/114
denominator matching	187/0	157/0	118/0
univariate partial fraction in x_5	53/52	50/49	41/38
factor matching	53/0	50/0	41/0
number of sample points	157798	122037	60482
required prime fields	3	2	3
evaluation time per point	2050s	190s	841s

Vector space	Univ. slices	Biv. slices	Random	Total
B_{4g}^{--}	60×171	85×85	20k	$\approx 55k$
B_{4g}^{+-}	55×138	62×62	25k	$\approx 45k$
B_{4g}^{++}	79×66	53×53	2k	$\approx 12.5k$
B_{2q2g}^-	55×109	69×69	10k	$\approx 22.5k$
B_{2q2g}^+	50×79	40×40	5k	$\approx 12.5k$
B_{4q}	50×79	49×49	3k	$\approx 10k$

Table 2: Number of points collected for reconstruction across each of the vector spaces, including a breakdown of the main sources.

special functions with massive particles

massive internal propagators
lead quickly to elliptic Feynman
integrals (or worse)

what is a basis for the function space?
numerical evaluation not clear
even (nested) square roots
can cause trouble for the
numerical evaluations

$$dI = M(\epsilon, x) \cdot I \quad [\text{Kotikov}][\text{Gehrmann, Remiddi}]$$

$$dI = \epsilon M(x) \cdot I \quad [\text{Henn}]$$

$$M(x) = \sum_w \text{dlog}(w)$$

leads to MPLs \Rightarrow good numerical performance

$$dI = \sum_{k=0}^G \epsilon^k M^{(k)}(x) \cdot I$$

can be achieved in general without elliptic/other objects
further rotation to ϵ factorised form can be done

numerical integration when MPLs are not possible...



integrate along path from
boundary point x to y

generalized series expansions

e.g. Liu, Ma, Wang [1711.09572] Moriello [1907.13234]

automated codes

DiffExp

AMFlow

LINE

SeaSyde

Runge-Kutta methods (e.g.
ODEINT Boost Library)

Bulirsch-Stoer shows
promising performance

Czakon, Niggetiedt 2001.03008

Czakon, Harlander, Klappert, Niggetiedt 2105.04436

Rosàs, Bobadilla 2507.12548

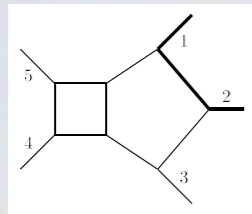
this week! → Baur, Duhr 2606.02744

SB, Becchetti, Brancaccio, Czakon,
Hartanto, Poncelet, Zoia

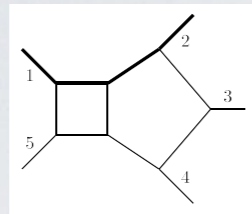
$$pp \rightarrow ttj$$

2404.12325 2412.13876

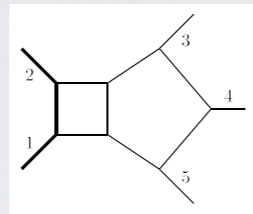
2511.11424 2511.11431



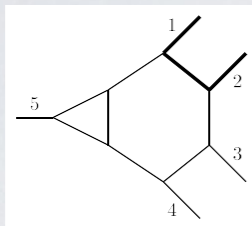
(a) Family PB_A .



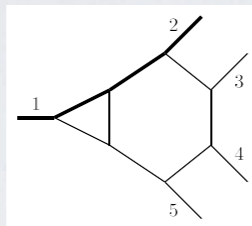
(b) Family PB_B .



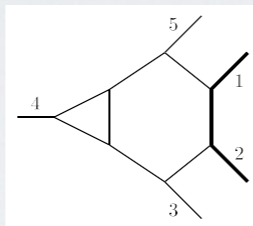
(c) Family PB_C .



(d) Family HT_A .



(e) Family HT_B .



(f) Family HT_C .

physical projectors, NeatIBP system

float type boundary	requested $ \Delta Q $	achieved $ \Delta Q $	maximal $ \Delta f_i/f_i $	number of steps	number of calls	call time	time
double \vec{d}_{01}	10^{-10}	8.2×10^{-10}	1.4×10^{-6}	17	716	0.77 ms	0.56 s
double \vec{d}_{02}	10^{-10}	8.2×10^{-10}	1.4×10^{-6}	28	1184	0.73 ms	0.89 s
dd_real \vec{d}_{01}	10^{-10}	7.7×10^{-11}	1.3×10^{-6}	16	709	9.3 ms	6.7 s
dd_real \vec{d}_{02}	10^{-10}	7.7×10^{-11}	1.3×10^{-6}	27	1166	9.3 ms	11 s
dd_real \vec{d}_{01}	10^{-20}	3.8×10^{-21}	1.2×10^{-16}	42	3002	9.0 ms	27 s
dd_real \vec{d}_{02}	10^{-20}	3.8×10^{-21}	1.2×10^{-16}	68	4873	9.0 ms	44 s

analytic reconstruction

	1	2	3	4
master integral coefficients of mass-renormalised amplitude (full ϵ dependence)	404/393	398/389	411/402	421/411
special function coefficients of finite remainder	314/303	305/296	321/312	326/317
linear relations	291/280	287/278	299/293	304/299
denominator matching #1	291/0	287/0	299/0	304/0
partial fraction decomposition in x_{5123}	44/40	55/51	57/54	58/54
denominator matching #2	44/0	54/0	54/0	56/0
number of sample points (1 prime field)	137076	89624	161482	179838

stable enough BUT rational coefficients
are extremely complicated

quadruple precision often necessary for the
coefficients - basis functions much quicker

$d\sigma$ computed using STRIPPER IR subtraction
is leading colour enough?

Channel	Functions [s]	Coefficients [s]	Assembly [s]	total [s]
$gg \rightarrow \bar{t}tj$	2.69	30.90	1.58	35.17
$q\bar{q} \rightarrow \bar{t}t\bar{q}$	2.16	9.40	0.18	11.74
$qg \rightarrow \bar{t}tq$	2.50	9.62	0.21	12.33
$q\bar{q} \rightarrow \bar{t}tg$	2.12	9.30	0.18	11.60



Becchetti, Canko, Chen, Chestnov,
Delto, Ditsch, Kallweit, Peraro, Pozzoli,
Savoini, Tancredi, Zoia, Grazzini

NNLO **total** cross sections using
approximate 2-loop amplitudes

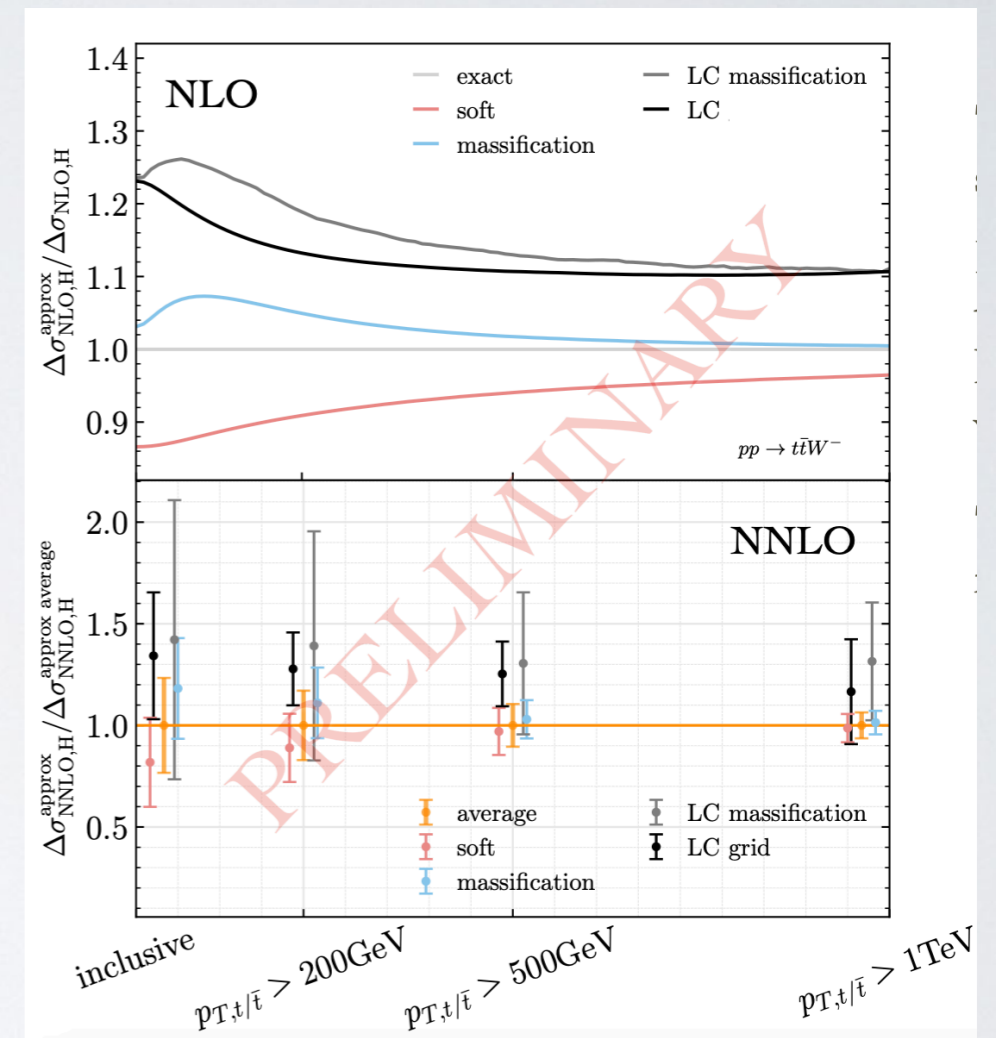
- soft boson limit
- massification [Penin][Moch, Mitov][Becher, Melnikov]
[Engel, Gnendiger, Signer, Ulrich][Wang et a.]

NEW! Leading colour amplitudes

challenging amplitude computation
and numerical evaluation

analytic reconstruction not possible

200K points, interpolation
grid, integrals via AMFlow



Grazzini, Top XS WG, 29/05/26

outlook

mass effects particularly relevant for low energy QED computations

multi-loop computations with full mass dependence are developing fast

are we at the limit of analytic workflow with reconstruction over finite fields?

NLO solved? e.g. analytic 1L $2 \rightarrow 4$ for $pp \rightarrow \gamma\gamma$ @ N3LO [Czakon et al. 2604.12613]

Simulations are non-trivial:

- can we share/interface to MC in a way useable to experiments?
- full mass effects come at considerable cost - CPU and human hours!
- fixed order will never be enough for all phase space region
- approximations can be much more efficient for event generation - can we guarantee reliable results when using them?