

Phokhara @ NLO+

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Workshop on Radiative Corrections and Monte Carlo Simulations at Electron-Positron Colliders
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Università di Torino

LEVERHULME
TRUST

In Pisa [May 8, 2025], we discussed the ultimate plans of Phokhara

From Graziano's talk

Desiderable:

CODE	mmg	ppg	Comments (matrix element, FSC)
Phokhara	NNLO	NNLO	exponentiation, FxsQED, GVMD,FsQED

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Desiderable:

Today's talk:
current status of Phokhara

CODE	mmg	ppg	Comments (matrix element, FSC)
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Outline

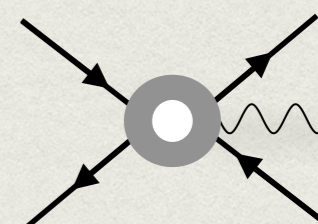
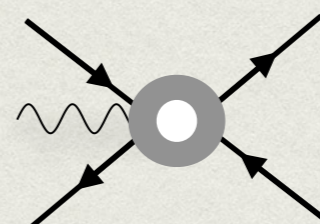
- Fixed order NLO + soft photon resummation
- GVMD (NLO) and $F \times$ sQED (NNLO) within Phokhara
- Fixed order NNLO :: in the making \rightarrow first look at Initial State Content

$$e^+e^- \rightarrow F^+F^-\gamma @ NLO$$

☑ Phokhara is a **Monte Carlo event generator** for low energy e^+e^- colliders, **with +20 years of development** (<https://looptreeduality.csic.es/phokhara/>)

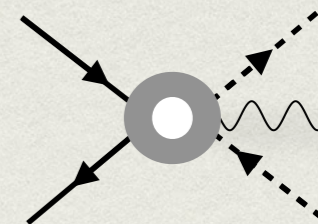
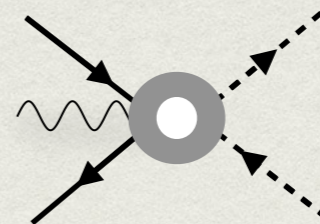
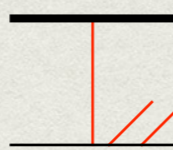
📌 Anatomy @ LO

- Born matrix element tree-level & n-pt process



📌 Anatomy @ NLO

- Real contribution tree-level $(n+1)$ -particles
- Virtual Contribution one-loop $(n+1)$ -particles



Use of **FXsQED**

☑ Phokhara10 does great for radiative return processes @ NLO accuracy

☑ Several cross checks against other codes:



☑ Workhorse :: efficient evaluation of *multi-loop scattering amplitudes*

$$e^+e^- \rightarrow F^+F^-\gamma \text{ @ NLO}$$

- **Phokhara10** calculates interference $\mathcal{M}^{(1)} = 2\Re(A^{(1)}A^{(0)*})$ from polarised amplitudes — let's exploit amplitudes!
- We construct efficient scattering amplitudes regardless of the loop order through a **four-dimensional tensor decomposition** -> immediately leads to independent **polarised/helicity amplitudes**

[Chen 2019; Peraro, Tancredi (2019)]

Tom's talk

- Analytic evaluation of polarised amplitudes
 $e^+(p_1)e^-(p_2) \rightarrow \pi^+(p_3)\gamma(p_4)\pi^-(p_5)$
 Understanding of Dirac structure in four dimensions.
 (same study carried out for $e^+e^- \rightarrow \mu^+\mu^-\gamma$)

$$\begin{aligned} \mathcal{T}_1 &= \bar{v}(p_2, m_e) \not{p}_3 u(p_1, m_e) (p_1 \cdot \varepsilon_4), \\ \mathcal{T}_2 &= \bar{v}(p_2, m_e) \not{p}_3 u(p_1, m_e) (p_3 \cdot \varepsilon_4), \\ \mathcal{T}_3 &= \bar{v}(p_2, m_e) \not{p}_5 u(p_1, m_e) (p_1 \cdot \varepsilon_4), \\ \mathcal{T}_4 &= \bar{v}(p_2, m_e) \not{p}_5 u(p_1, m_e) (p_3 \cdot \varepsilon_4), \\ \mathcal{T}_5 &= \bar{v}(p_2, m_e) u(p_1, m_e) (p_1 \cdot \varepsilon_4), \\ \mathcal{T}_6 &= \bar{v}(p_2, m_e) u(p_1, m_e) (p_3 \cdot \varepsilon_4), \\ \mathcal{T}_7 &= \bar{v}(p_2, m_e) \not{p}_3 \not{p}_5 u(p_1, m_e) (p_1 \cdot \varepsilon_4), \\ \mathcal{T}_8 &= \bar{v}(p_2, m_e) \not{p}_3 \not{p}_5 u(p_1, m_e) (p_3 \cdot \varepsilon_4). \end{aligned}$$

$$\mathcal{A}^{(L)} = \sum_{i=1}^8 \mathcal{F}_i^{(L)} \mathcal{T}_i$$

[Dave, Paltrinieri, Petit Rosas, WJT (2026)]

$$e^+e^- \rightarrow F^+F^-\gamma \text{ @ NLO}$$

○ **Phokhara10** provides fixed NLO theoretical predictions — **we need NLO+**

○ **Soft photon resummation** — YFS exponentiation

[Yennie, Frauschi, Suura (1961)]

$$\sigma_{\text{NLO}}^{\text{resum}} = \sum_{n_\gamma} \int_{\Omega} d\text{LIPS}_{n_\gamma+2} \frac{e^{\alpha Y_{\text{FF}}}}{n_\gamma!} \sum_{\text{hel}} \left| \sum_{\phi \in [I,F]^{n_\gamma}} \left(\prod_{i=1}^{n_\gamma} S_i(k_i) \Theta(k_i, \Omega) \right) \left(\beta_0 + \sum_{i=1}^{n_\gamma} \frac{\beta_1(k_i)}{S(k_i)} + \sum_{\substack{i,j=1 \\ j < k}}^{n_\gamma} \frac{\beta_1(k_i, k_j)}{S(k_i) S(k_j)} \right) \right|^2$$

Partitions (ISR/FSR)
Eikonal factors
Resolved domain
Regularised amplitudes

- * Virtual & real amplitudes inside 's
- * Control on ISR and FSR contributions!
- * Need momenta mappings to map N -photon events to i -photon events for β_i .

○ Warm up exercise :: Scan Mode NLO resummed, CMD scenario

Jeremy's talk

○ Work in progress :: Radiative return NLO resummed, KLOE scenarios

$e^+e^- \rightarrow \pi\pi\gamma @ 1L$

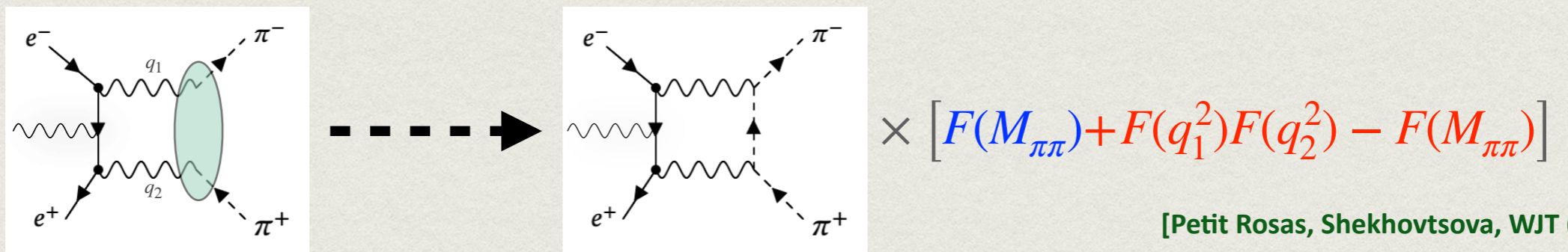
- **Phokhara10** employs $F \times$ sQED to describe the interaction between pions and photons

- Improve description between photon-pion interaction
 —> generalised vector-meson dominance (GVMD)

[Lee, Ignatov (2022)]

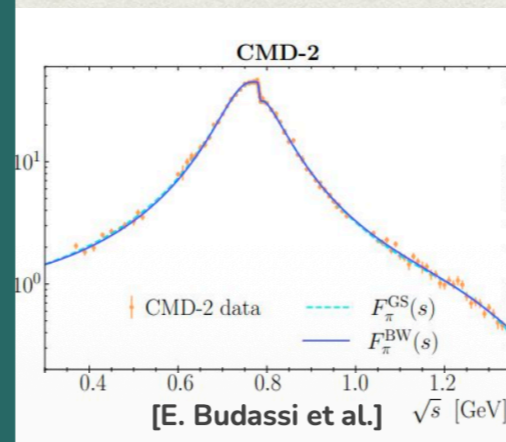
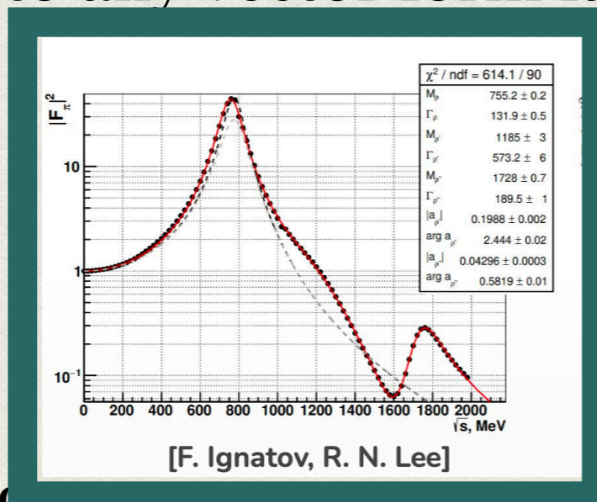
[Colangelo, Hoferichter, Monnard, Ruiz de Elvira (2022)]

- Insertion of GVMD in radiative return processes



[Petit Rosas, Shekhovtsova, WJT (2026)]

- Suitable to any vector form factor parametrisation



Validations

- ✓ Kloe scenarios: SA & LA
- ✓ Forward-Backward asymmetry in Kloe

$$e^+e^- \rightarrow F^+F^-\gamma @ NNLO$$

Anatomy @ NNLO

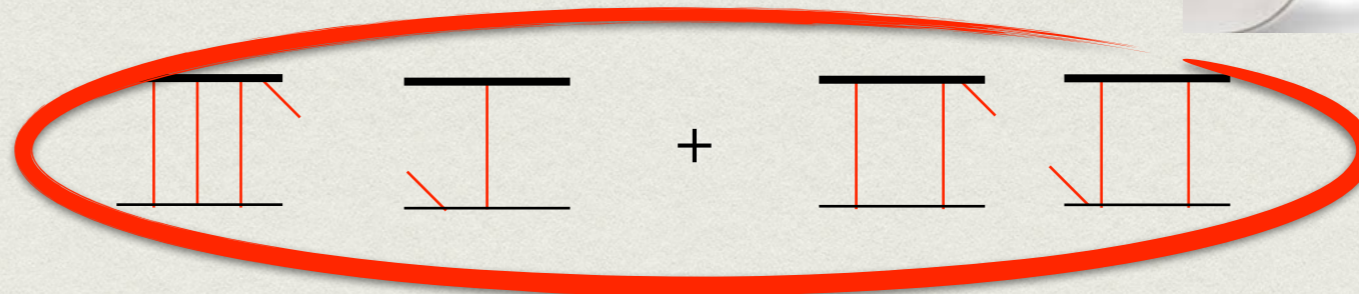
- Real-Real contribution
Tree-level $(n+2)$ -particles



- Real-Virtual Contribution
one-loop $(n+1)$ -particles



- Virtual-Virtual Contribution
two-loop n -particles



- Harder (but doable) phase-space integration
- Extend numerical evaluation of one-loop Feynman integrals
- Basis of two-loop Feynman integral not known



Kinematic invariant $(s_{12}, s_{23}, s_{34}, s_{45}, s_{51}, m_e^2, m_F^2)$

one-loop polarised amplitudes to higher orders in ϵ

Pronounced hierarchy of scales,

[Dave, Paltrinieri, Petit Rosas, WJT (2026)]

$$\frac{m_e^2}{s} \sim 10^{-7} - 10^{-8}$$

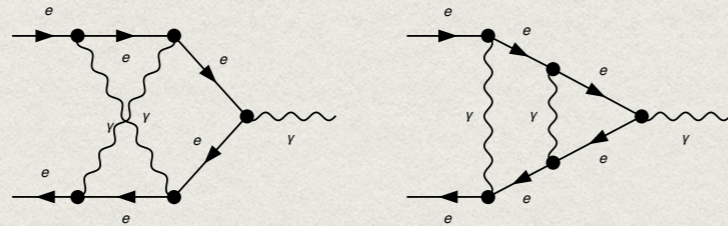
$e^+e^- \rightarrow \pi\pi\gamma @ 2L$

Amplitude generation

Two-loop gauge invariant pieces

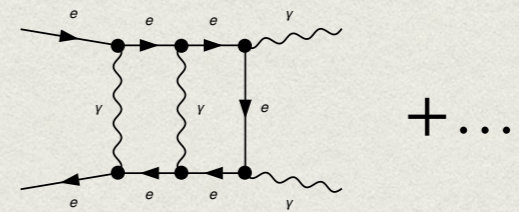
Algebraic decomposition

$$f^+ f^- \rightarrow \gamma^* \rightarrow F^+ F^- + \gamma$$



✓ Easy (m_f^2, s)

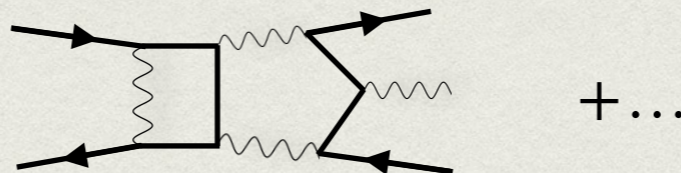
$$f^+ f^- \rightarrow \gamma \gamma^* \rightarrow F^+ F^-$$



□ "Normal" (s, t, m_e^2, q^2)

Loop integral evaluation

$$f^+ f^- \rightarrow F^+ F^- \gamma$$



□ Hard ($s_{12}, s_{23}, s_{34}, s_{45}, s_{51}, m_f^2, m_F^2$)

MC input

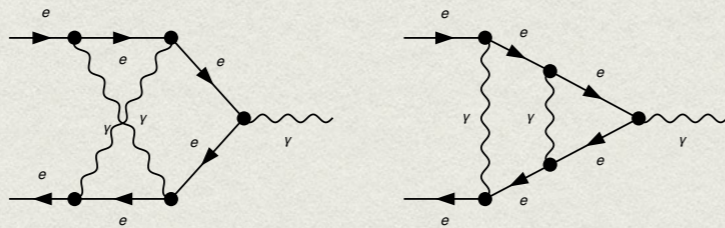
$e^+e^- \rightarrow \pi\pi\gamma @ 2L$

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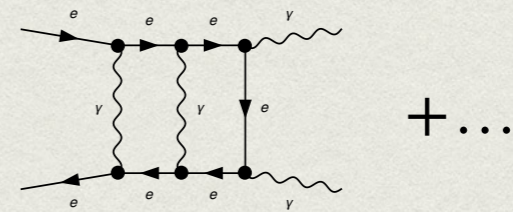
Algebraic decomposition

$$f^+ f^- \rightarrow \gamma^* \rightarrow F^+ F^- + \gamma$$



✓ Easy (m_f^2, s)

$$f^+ f^- \rightarrow \gamma \gamma^* \rightarrow F^+ F^-$$

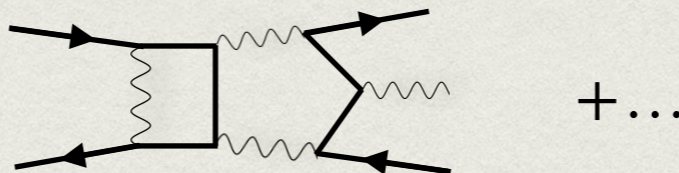


□ "Normal" (s, t, m_e^2, q^2)

Progress on these Feynman integrals

Loop integral evaluation

$$f^+ f^- \rightarrow F^+ F^- \gamma$$



□ Hard ($s_{12}, s_{23}, s_{34}, s_{45}, s_{51}, m_f^2, m_F^2$)

MC input

$$e^+e^- \rightarrow \pi\pi\gamma @ 2L$$

○ Initial state contributions



Amplitude generation



Algebraic decomposition

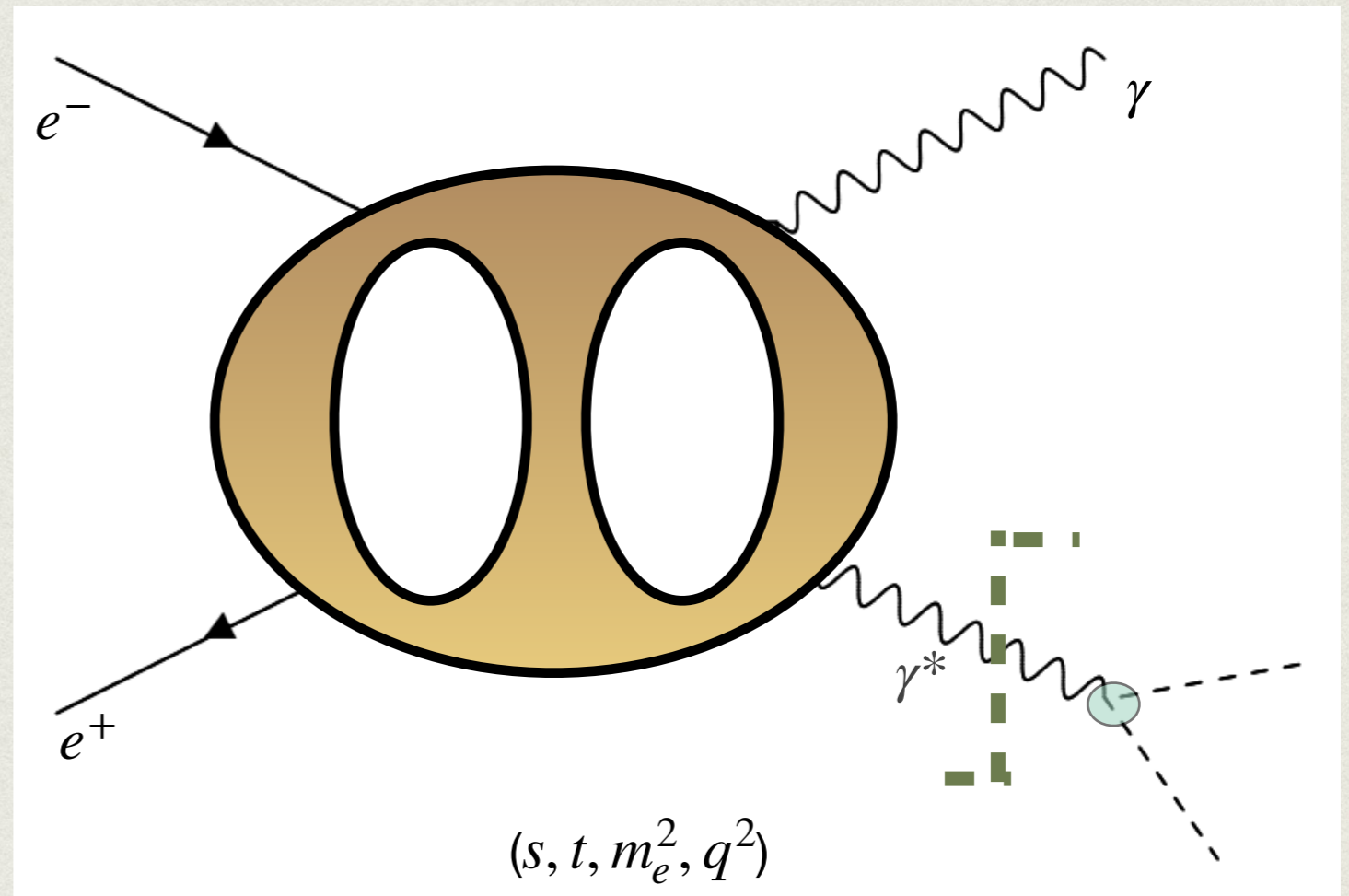


Loop integral evaluation



MC input

$$e^+e^- \rightarrow \gamma\gamma^* \rightarrow \pi\pi$$



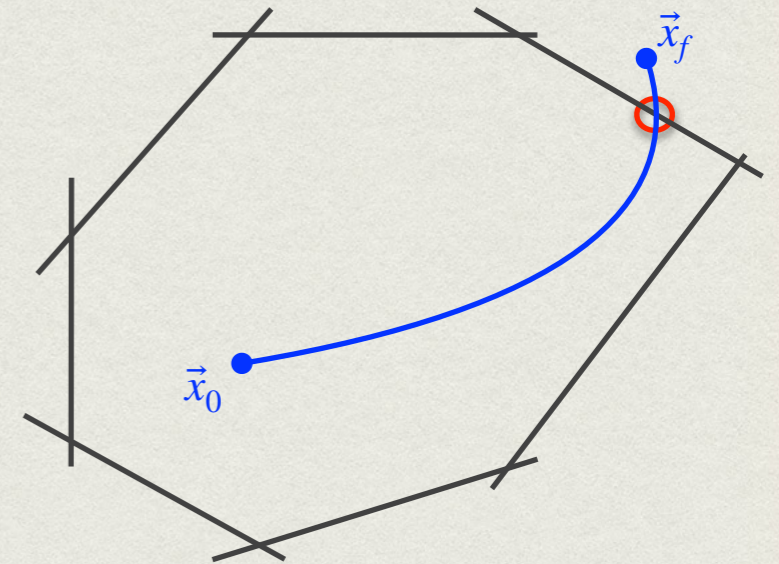
Evaluation of Feynman integrals

- Evaluation of Feynman integrals by the method of differential equations

$$\partial_x \vec{I}(\vec{x}; \epsilon) = A_x(\vec{x}; \epsilon) \vec{I}(\vec{x}; \epsilon)$$

- Find DEQs of the form

$$d A^{(\mathbf{X})}(\vec{x}; \epsilon) = \sum_{k=0}^2 \epsilon^k \left[\sum_{\alpha} c_{k\alpha}^{(\mathbf{X})} d \log (W_{\alpha}(\vec{x})) + \sum_{\beta} d_{k\beta}^{(\mathbf{X})} \omega_{\beta}(\vec{x}) \right]$$



$$\omega_{\beta}(\vec{x}) = \omega_{\beta}^{(s)}(\vec{x}) ds + \omega_{\beta}^{(t)}(\vec{x}) dt + \omega_{\beta}^{(m^2)}(\vec{x}) dm^2 + \omega_{\beta}^{(q^2)}(\vec{x}) dq^2$$

Move **difficult** integrals to very late stages

Available Mathematica implementations (DiffExp, ...) but slow for MC evaluations and costly to generate grids

- Our numerical approach $\mathcal{O}(10 - 100\text{ms})$ per point

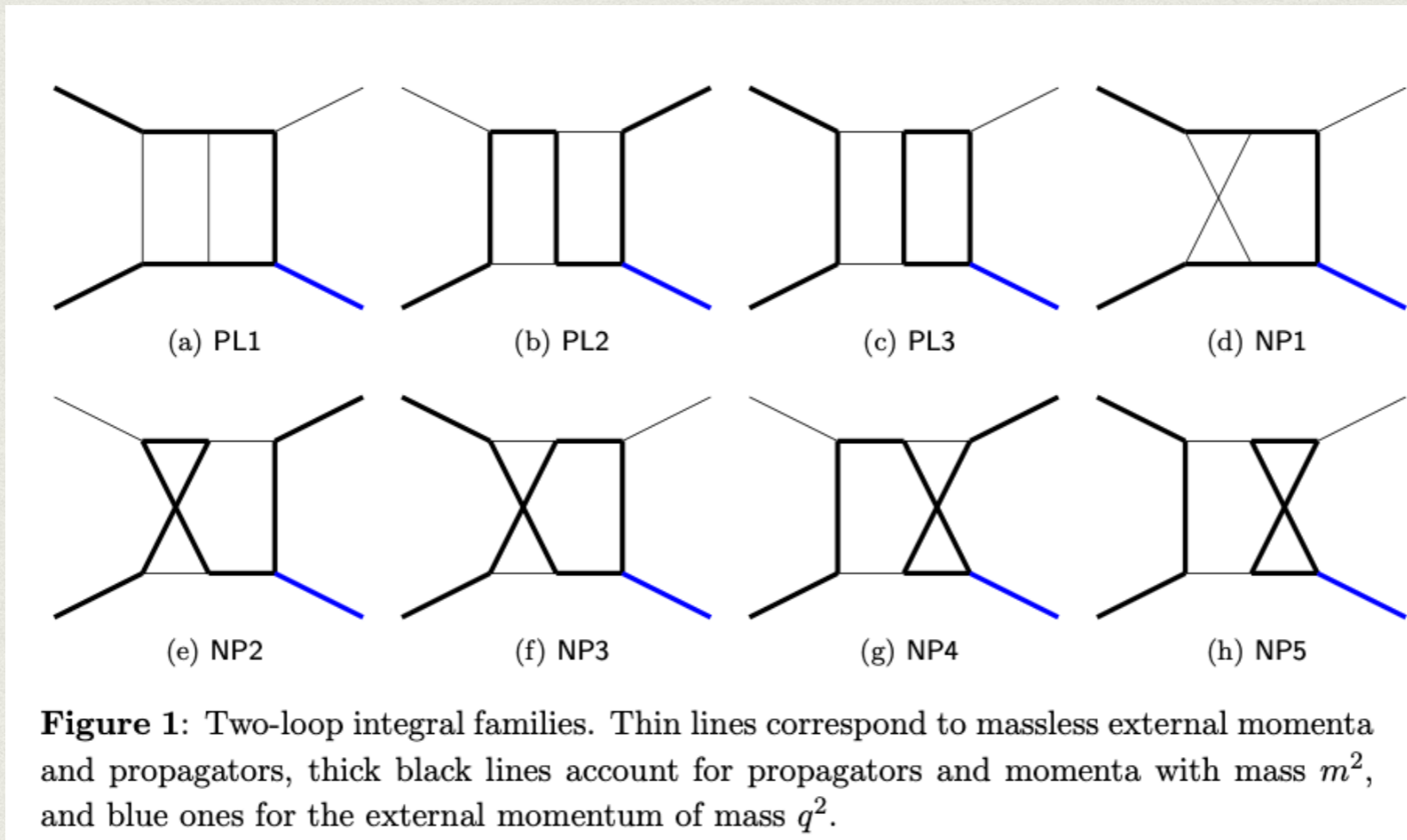
[Petit Rosas, WJT (2025)]

- ✓ Fast integrator in a low-level language
- ✓ Precise, but no need for 50 significant figures!
- ✓ Ideally, fast enough to not use grids

Validations

- ✓ $e^+e^- \rightarrow \pi^+\pi^-\gamma$ up $\mathcal{O}(\epsilon^2)$
- ✓ $pp \rightarrow ttj$ @ 2L
- ✓ $e^+e^- \rightarrow \gamma\gamma^*$ @ 2L

Evaluation of Feynman integrals for $e^+e^- \rightarrow \gamma\gamma^*$ @ 2L



Mattia's talk

- ✎ Presence of Elliptic and hyper elliptic integrals
- ✓ Planar integrals (PL1, PL2, PL3) completed
- 🔄 Non planar integrals in the making

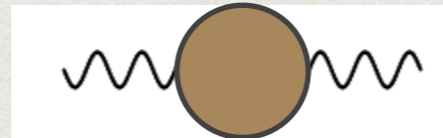
[Pozzoli, WJT (to appear)]

$$e^+e^- \rightarrow \gamma\gamma^* @ 2L$$

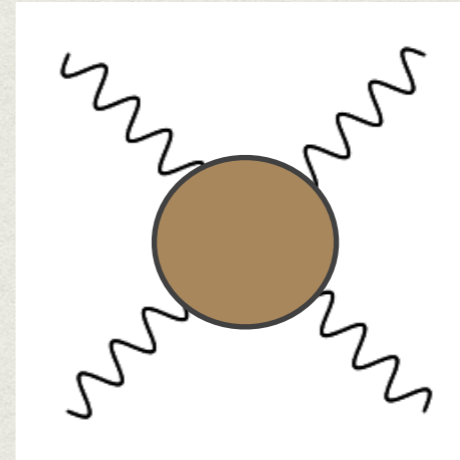
[Dave, Petit Rosas, Pozzoli, WJT (in preparation)]

Decomposition of the two-loop amplitudes (from leptonic to hadronic corrections)

$$\mathcal{A}^{(2)} = A^{(2)} + N_F B^{(2)}$$

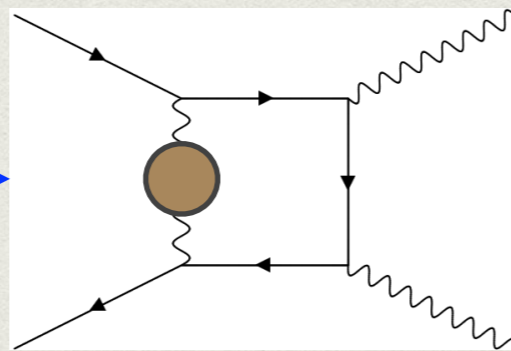
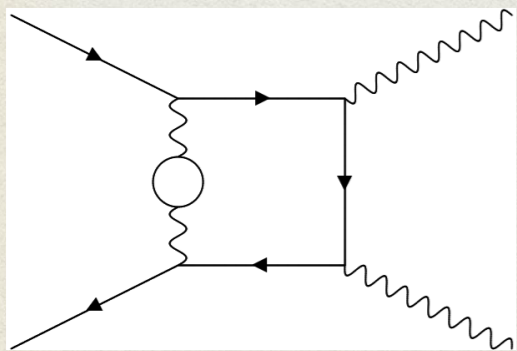


Photon correction



Light-by-light

$$\Pi^{(\text{had})}(q^2) = -\frac{q^2}{\pi} \int_0^\infty \frac{d\Lambda}{\Lambda^2} \frac{\tilde{\Pi}^{(\text{had})}(\Lambda^2)}{q^2 - \Lambda^2 + i0}$$

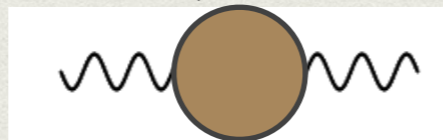


$$-\frac{1}{\pi} \int_0^\infty \frac{d\Lambda}{\Lambda^2} \tilde{\Pi}^{(\text{had})}(\Lambda^2) \times [1L]$$

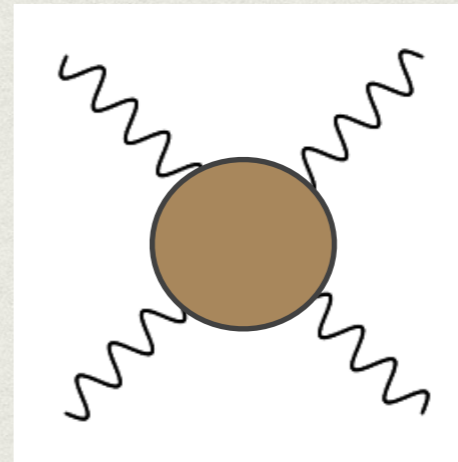
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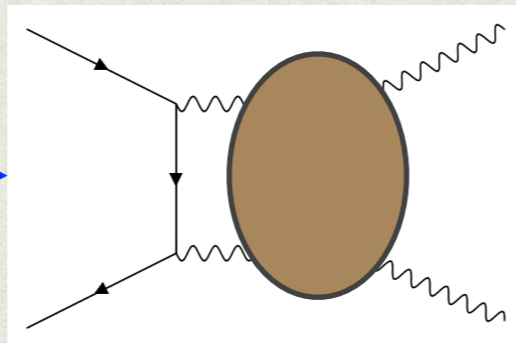
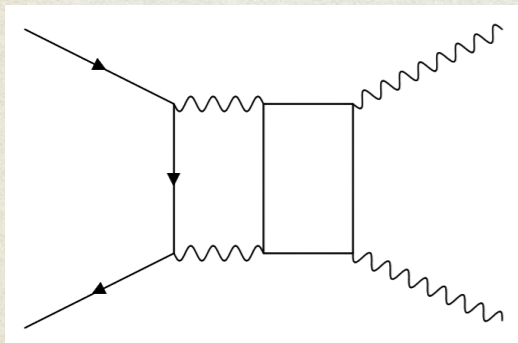


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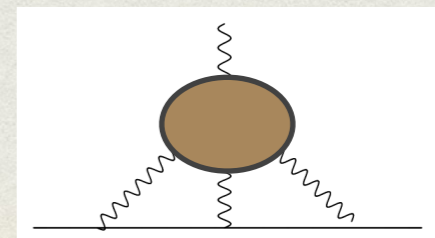


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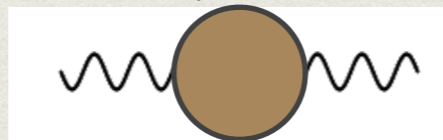
Can we borrow knowledge from $(g-2)_\mu$?



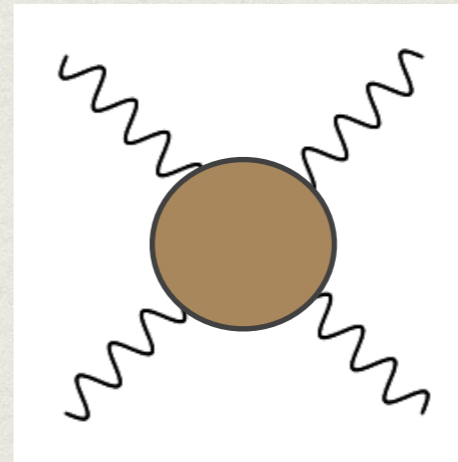
$$e^+e^- \rightarrow \gamma\gamma^* @ 2L$$

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Photon correction



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$$\Pi^{(\text{had})}(q^2) = -\frac{q^2}{\pi} \int_0^\infty \frac{d\Lambda}{\Lambda^2} \frac{\tilde{\Pi}^{(\text{had})}(\Lambda^2)}{q^2 - \Lambda^2 + i0}$$

Getting a dispersive representation in terms of one-loop amplitudes will immediately a check for $B^{(2)}$ calculation and show possible pitfalls in the calculation of $A^{(2)}$

Conclusions

• We have reached (and will be discussed in more details):

- Analytic study of radiative return processes
- First improvements and validation of the Phokhara generator
GVMD + Soft-photon resummation
- First look at the evaluation of two-loop Feynman integrals for $e^+e^- \rightarrow \gamma \gamma^*$

• Open questions & future directions

- Complete radiative return NLO resummed for KLOE scenarios
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