

# Quantum-Correlated $D^0\bar{D}^0$ systems in LHCb Run 3 data

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08 April 2026



# Introduction

- From neutral meson oscillations, we know that neutral  $D$  mesons exist in  $CP$  eigenstates which are a superposition of flavour eigenstates (neglecting CPV):

$$|D_+\rangle = \frac{1}{\sqrt{2}}(|D^0\rangle + |\bar{D}^0\rangle)$$

$$|D_-\rangle = \frac{1}{\sqrt{2}}(|D^0\rangle - |\bar{D}^0\rangle)$$

- Quantum-correlated  $D^0\bar{D}^0$  systems refer to pairs of neutral  $D$  mesons which exist in eigenstates of  $C$  and  $P$  defined by (in absence of charm mixing):

$$C\text{-even} : \frac{|D^0\bar{D}^0\rangle + |\bar{D}^0D^0\rangle}{\sqrt{2}} = \frac{|D_+D_+\rangle - |D_-D_-\rangle}{\sqrt{2}}$$

and

$$C\text{-odd} : \frac{|D^0\bar{D}^0\rangle - |\bar{D}^0D^0\rangle}{\sqrt{2}} = \frac{|D_-D_+\rangle - |D_+D_-\rangle}{\sqrt{2}}$$

- Produced in decays of particles with definite  $J^{PC}$  quantum numbers to certain final states...e.g.  $\psi(3770)$ ,  $\chi_{c1}(3872)$ ...etc
- Reconstructing both neutral  $D$  mesons in  $CP$ -definite final states, e.g.  $K^+K^-$ ,  $\pi^+\pi^-$  enhances/suppresses contributions depending on the parent particle  $J^{PC}$

# Physics Motivation

- To date, QC  $D^0\bar{D}^0$  systems have been studied at  $e^+e^-$  experiments to obtain:
  - ▶ Time-integrated measurements of charm mixing parameters
  - ▶  $D^0$  decay strong phases (input to charm mixing and CKM  $\gamma$  measurement)
- May also be able to exploit them for other measurements (JHEP 03 (2023) 038):
  - ▶ Ruling out quantum numbers in hadron spectroscopy
  - ▶  $T/CPT$  symmetry measurements in neutral charm
- Need a large sample of correlated  $D^0\bar{D}^0$  pairs — charmonia(-like) states such as  $\chi_{c1}(3872)$  provide an opportunity to do so at LHCb

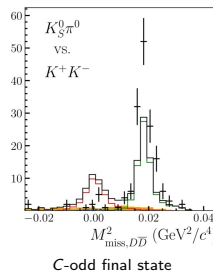
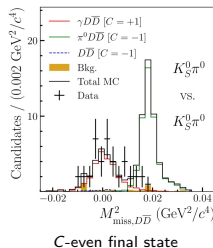


Figure from BESIII: [Phys. Rev. D 112 \(2025\), 072006](#)

## Quantum Correlations in $\chi_{c1}(3872)$ decays

- $\chi_{c1}(3872)$  quantum numbers were measured to be  $J^{PC} = 1^{++}$  in [Phys. Rev. Lett. 1101, 222001](#)
- $\chi_{c1}(3872)$  can decay to both the  $(D^0\bar{D}^0)_{C+\pi^0}$  and  $(D^0\bar{D}^0)_{C-\gamma}$  final states
- This analysis focuses mainly on  $D^0\bar{D}^0$  pairs produced by the decay  $\chi_{c1}(3872) \rightarrow D^0\bar{D}^0\pi^0/\gamma$  where the neutral particle is not reconstructed
- $\mathcal{O}(10^3)$  partially reconstructed  $\chi_{c1}(3872) \rightarrow D^0\bar{D}^0 X$  decays were seen by LHCb during run 1+2 ([JHEP 1907 \(2019\) 035](#) ( $9 \text{ fb}^{-1}$ ))
- Expect enhancement of  $D^0\bar{D}^0\pi^0$ , suppression of  $D^0\bar{D}^0\gamma$  components when both  $D$  mesons are reconstructed in the CP-even  $K^-K^+/\pi^-\pi^+$  final states

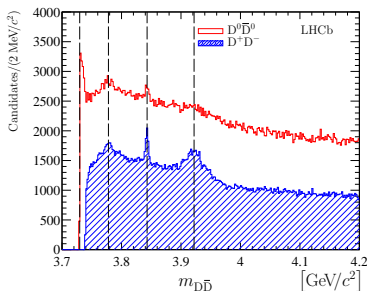
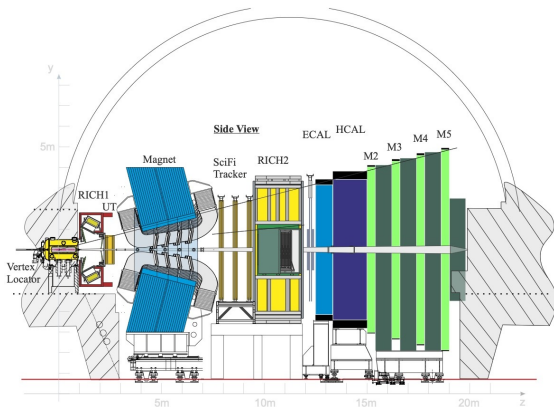


Figure from [JHEP 1907 \(2019\) 035](#)

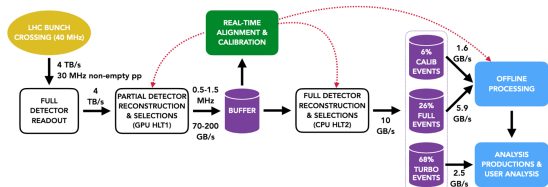
# The LHCb Detector

- Single-arm forward spectrometer built to study decays of  $b$  and  $c$  hadrons produced in the forward region
- Fully instrumented within pseudorapidity range of  $2 < \eta < 5$
- This analysis looks at  $5.4 \text{ fb}^{-1}$  and  $11.4 \text{ fb}^{-1}$  of  $pp$  collision data taken in 2024 and 2025 respectively



# LHCb in Run 3

- During LS2, LHCb underwent a major upgrade, allowing it to run at five times the luminosity of previous runs
- Key features of the upgrade include:
  - ▶ Overhaul of tracking and PID systems and associated readout electronics to accommodate full detector readout at 40 MHz
  - ▶ Replacement of hardware trigger with a software-based trigger  $\implies$  Improved efficiency especially for fully hadronic final states



- Changes in online selections for  $DD$  pair reconstruction in 2025 also brings a factor of 2 improvement in signal yield per lumi compared to 2024
- Makes it viable to study Cabibbo-suppressed  $CP$  double-tagged  $DD$  pairs despite low branching fraction ( $\mathcal{B}(D^0 \rightarrow K^- K^+) \sim 1/10$  of  $\mathcal{B}(D^0 \rightarrow K^- \pi^+)$ )

# Analysis Outline

- Measure the efficiency-corrected yield relative to the number of  $D\bar{D} \rightarrow K^-\pi^+$  vs  $K^+\pi^-$  decays:

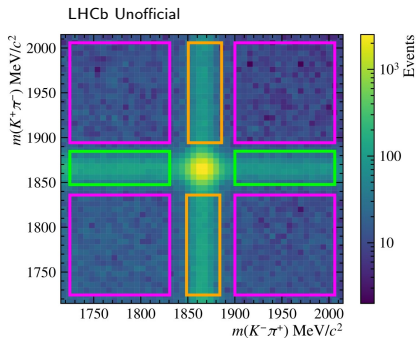
$$\begin{aligned}\kappa_{ij} &= \frac{\text{Yield in channel } ij}{\text{Yield in channel } ij \text{ without correlations}} \\ &= \frac{N_{X \rightarrow D^0 \bar{D}^0, D^0 \rightarrow i, \bar{D}^0 \rightarrow j}}{N_{X \rightarrow D^0 \bar{D}^0, D^0 \rightarrow K^-\pi^+, \bar{D}^0 \rightarrow K^+\pi^-}} \frac{\mathcal{B}(D^0 \rightarrow K^-\pi^+) \times \mathcal{B}(\bar{D}^0 \rightarrow K^+\pi^-)}{\mathcal{B}(D^0 \rightarrow i) \times \mathcal{B}(\bar{D}^0 \rightarrow j)} \frac{\epsilon_{K^-\pi^+/K^+\pi^-}}{\epsilon_{ij}}\end{aligned}$$

where  $i, j$  correspond to CP-even final states (i.e.  $K^-K^+, \pi^-\pi^+$ )

- Perform measurements over the known  $DD$  resonances in the near-threshold region:  $\chi_{c1}(3872)$ ,  $\psi(3770)$ ,  $X(3842)$ ,  $\chi_{c2}(3930)$
- Expect to see  $\kappa \approx 2$  and  $\kappa = 0$  for  $C$ -even and  $C$ -odd resonances respectively
- Measurement of relative enhancement in  $D\bar{D} \rightarrow K^-\pi^+$  vs  $K^-K^+/\pi^-\pi^+$  also provides sensitivity to  $D^0 \rightarrow K^-\pi^+$  strong phase  $r_{K\pi}^D \cos \delta_{K\pi}^D$
- This talk covers some of the work done so far on the  $K^-\pi^+$  vs  $K^+\pi^-$  channel

# Backgrounds

- Combinatoric backgrounds forming fake  $D^0$  candidates:
  - Real  $D^0$  + Fake  $\bar{D}^0$
  - Fake  $D^0$  + Real  $\bar{D}^0$
  - Fake  $D^0$  + Fake  $\bar{D}^0$
- Real  $D^0$  faking prompt  $D^0\bar{D}^0$  candidates:
  - $D^0$  from different PV
  - $D^0$  from PV +  $D^0$  from  $B$ -decay
  - $B \rightarrow D^0\bar{D}^0 X$  decays
- Partially reconstructed  $D^+$ ,  $D_s^+$  decays
- $D^0$  decays with misidentified daughters
- Peaking backgrounds:
  - Partially reconstructed  $T_{cc}^+ \rightarrow D^0 D^0 \pi^+$  in suppressed  $D^0$  decay modes
  - $D^0$  from prompt  $D^{*+} \rightarrow D^0 \pi^+$

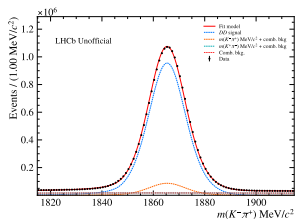


# $m(D^0)$ Mass Fits

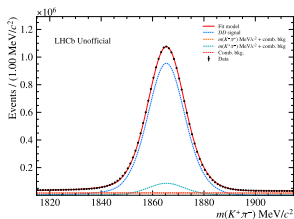
- Signal component is sum of a Double-sided Crystal Ball and Gaussian
- Combinatoric background components  $B_{1,2}(x)$  consists of a product of an exponential and a 2nd order polynomial
- For channels where both  $D^0$  final states are identical or charge conjugates of each other, require PDF to be symmetric

The fitted PDF can be written as:

$$\begin{aligned} PDF(x, y) = & N_{D, \bar{D}} G(x) G(y) \\ & + N_{D, bkg} G(x) B(y) + N_{\bar{D}, bkg} B(x) G(y) \\ & + N_{bkg, bkg} B(x) B(y) \end{aligned}$$



$$D^0 \rightarrow K^- \pi^+$$

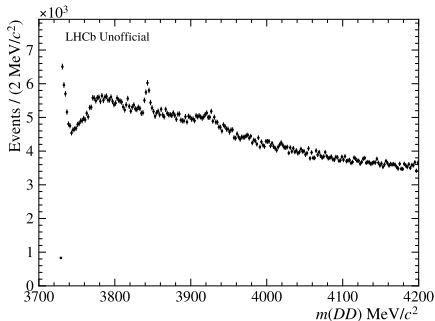


$$D^0 \rightarrow K^+ \pi^-$$

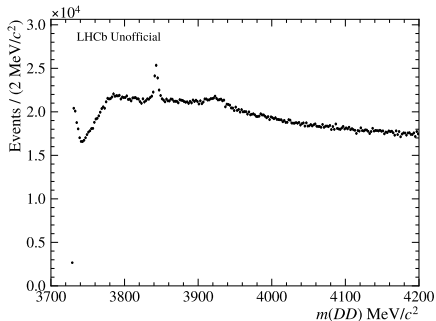
Partial 2025 dataset (5.7  $\text{fb}^{-1}$ )

## $m(DD)$ Mass Spectrum

- Using results of the  $D^0$  mass fit, we calculate per-event  $sWeights$  to subtract the  $D^0 \rightarrow hh$  combinatoric background
- Substantial difference between 2024, 2025  $m(DD)$  spectrum at open charm threshold — due to changes in trigger selections at start of 2025
- Clear peaking structures seen for the  $\chi_{c1}(3872)$ ,  $X(3842)$  and  $\chi_{c2}(3930)$
- $\psi(3770)$  less clear; wide resonance lying on top of  $D\bar{D}$  threshold



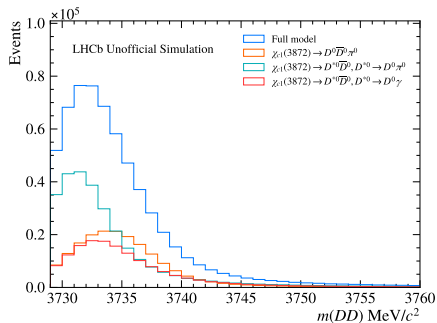
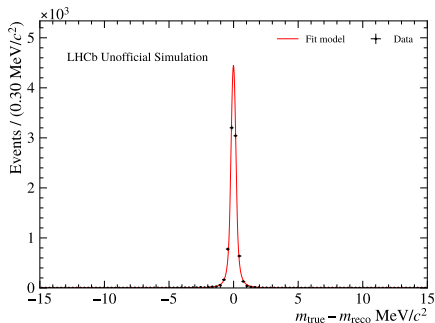
Full 2024 dataset ( $5.4 \text{ fb}^{-1}$ )



Full 2025 dataset ( $11.4 \text{ fb}^{-1}$ )

## $\chi_{c1}(3872)$ Signal Model

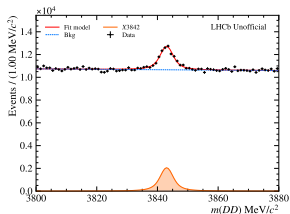
- Threshold peak from partially reconstructed  $\chi_{c1}(3872) \rightarrow D^0 \bar{D}^0 \pi^0 / \gamma$  decay consists of both C-even and C-odd contributions  $\implies$  Needs different treatment from  $c\bar{c} \rightarrow D^0 \bar{D}^0$  peaks
- Generate MC of the various  $\chi_{c1}(3872) \rightarrow D^0 \bar{D}^0 X$  components:
  - ▶  $\chi_{c1}(3872) \rightarrow D^0 \bar{D}^0 \pi^0$
  - ▶  $\chi_{c1}(3872) \rightarrow D^{*0} \bar{D}^0, D^{*0} \rightarrow D^0 \pi^0$
  - ▶  $\chi_{c1}(3872) \rightarrow D^{*0} \bar{D}^0, D^{*0} \rightarrow D^0 \gamma$
- Fix ratio of  $D^* D$  components using known branching fractions
- Convolve with resolution function obtained from fit to  $m(DD)_{\text{true}} - m(DD)_{\text{reco}}$  in simulation



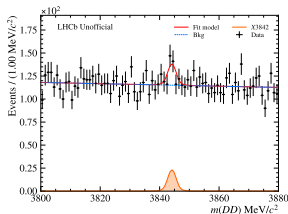
# Uncorrelated Normalisation

$$\kappa_{ij} = \frac{\text{Yield in channel } ij}{\text{Yield in channel } ij \text{ without correlations}}$$

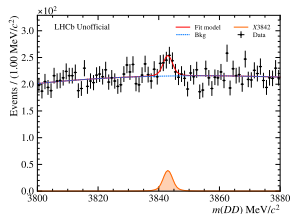
- Obtain uncorrelated normalisation for each  $m(DD)$  peak from the  $K^-\pi^+/K^+\pi^-$  channel
- Perform fit to  $m(DD)$  and generate toys with statistics scaled by branching fraction and efficiency ratios
- Based on toy fits, expect to observe  $X(3842)$  with a significance  $> 5\sigma$  under the no-correlation hypothesis only when including All 3  $CP$  double-tag  $D \rightarrow hh$  channels ( $K^-K^+$  vs  $K^-K^+$ ,  $\pi^-\pi^+$  vs  $\pi^-\pi^+$  and  $K^-K^+$  vs  $\pi^-\pi^+$ )



2025  $K^-\pi^+/K^+\pi^-$  Data



Expected 2025 toy ( $K^-K^+/K^-K^+$ )



Expected 2025 toy (All  $CP$  double-tag channels)

# Summary

- LHCb Run 3 opens up a unique opportunity to study quantum-correlated  $D\bar{D}$  mesons at hadron colliders for the first time
- Ongoing analysis to demonstrate quantum correlations in  $D^0\bar{D}^0$  system produced at LHCb using run 3 dataset
- Procedures for event selection and mass fits are currently being finalised, and work on fitting the uncorrelated  $K^-\pi^+/K^+\pi^-$  channel near the open charm threshold is underway
- To-do: Apply selections to and fit mass spectrum in other  $D^0 \rightarrow hh$  channels, systematic uncertainties, measure yield ratios and  $r_{K\pi}^D \cos \delta_{K\pi}^D$

# Backup

## T/CPT symmetry measurements

- Can make direct measurements of  $T/CPT$  symmetry using flavour- $CP$  transitions (e.g.  $D_+ \rightarrow D^0$  with a  $CPT$  conjugate decay of  $\bar{D}^0 \rightarrow D_+$ ) by reconstructing  $D_1$  at time  $t_1$  in a  $CP$  state (e.g.  $K^-K^+$ ) and  $D_2$  in a  $CF$  final state

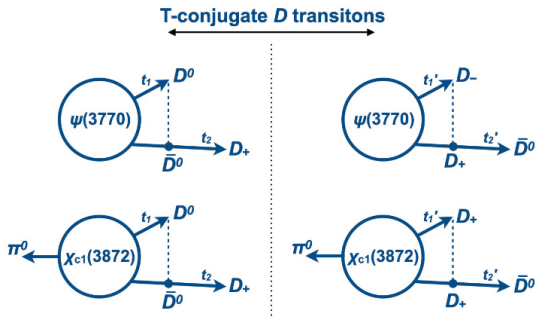


Figure from [JHEP 03 \(2023\) 038](#)

## Strong phase measurements

- The decay rate for  $D\bar{D} \rightarrow Y_1$  vs  $Y_2$  is given by:

$$\begin{aligned} \frac{\Gamma(D\bar{D} \rightarrow Y_1 \text{ vs. } Y_2)}{A_1^2 A_2^2} &= [r_1^2 + r_2^2 + 2CR_1R_2r_1r_2 \cos(\delta_1 - \delta_2)] \\ &\quad - (1 + C)y [R_1r_1 \cos \delta_1(1 + r_2^2) + R_2r_2 \cos \delta_2(1 + r_1^2)] \\ &\quad - (1 + C)x [R_1r_1 \sin \delta_1(1 - r_2^2) + R_2r_2 \sin \delta_2(1 - r_1^2)] \\ &\quad + \mathcal{O}(x^2, y^2), \end{aligned}$$

where  $Y_1, Y_2$  denotes the final state of the  $D, C$  the even/odd correlation of the  $D\bar{D}$  pair and  $x, y$  are charm mixing parameters

- For a correlated  $D\bar{D}$  pair decaying to a  $D\bar{D} \rightarrow K^-\pi^+$  vs  $K^-K^+/\pi^-\pi^+$  final state, the relative enhancement wrt the uncorrelated case is given by:

$$1 + \frac{2Cr_{K\pi}^D \cos \delta_{K\pi}^D - (1 + C)y}{1 + (r_{K\pi}^D)^2}$$

where  $\lambda$  is  $+1$  and  $-1$  for  $C$ -even and odd correlated  $D\bar{D}$  pairs respectively

# Fitted $m(D^0\bar{D}^0)$ spectrum — Run 1+2 analysis

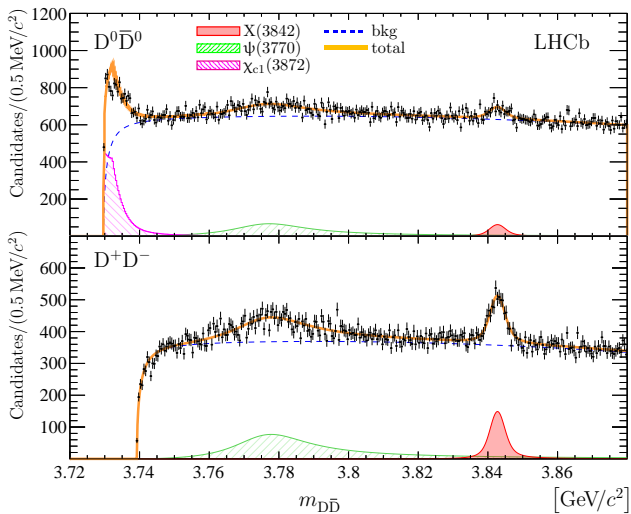


Figure from JHEP 1907 (2019) 035 ( $9 \text{ fb}^{-1}$ )