

A COMPLEX LPA FOR THE CRITICAL AND MULTICRITICAL LEE-YANG FIXED POINTS

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Based on:

- Dario Benedetti, FE, Omar Zanusso [arXiv:2601.15087]



Why study non-unitary theories?

Recent conjecture: [Zambelli, Zanusso, 16'] [Lencsés, Miscioscia, Mussardo, Takcs, 24'] [Katsevich, Klebanov, Sun, Tarnopolsky, 25'] [...]

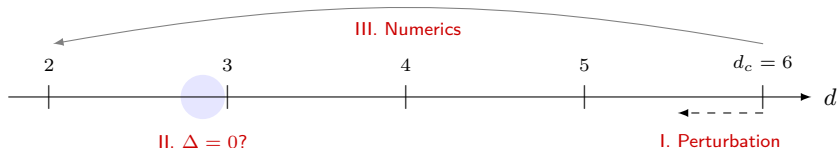
In $d = 2$, the non unitary minimal models $\mathcal{M}(2, 2n + 3)$ corresponds to the IR fixed point of

$$S = \int d^2x \frac{1}{2} \partial_\mu \varphi(x) \partial^\mu \varphi(x) + g \int d^2x i\varphi^{2n+1}(x)$$

- $\mathcal{PT}$ -symmetry $\rightarrow$ real, bounded from below spectrum
- Lee-Yang theory $\leftrightarrow \mathcal{M}(2, 5)$ [Fisher 78'] [Cardy, 85']

The idea:

Use the FRG approach to continue from the upper critical dimension to $d = 2$



Functional Renormalization Group

Wetterich equation [Wetterich, 1993]

$$k\partial_k \Gamma_k = \frac{1}{2} \text{Tr} \left[k\partial_k \mathcal{R}_k \left(\Gamma_k^{(2)} + \mathcal{R}_k \right)^{-1} \right]$$

Litim cutoff:

$$R_k(q) = Z_k \left(k^2 - q^2 \right) \Theta \left(k^2 - q^2 \right)$$


Anomalous dimension

$$\eta = -\frac{1}{Z_k} k\partial_k Z_k$$

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Local Potential Approximation

$$\Gamma_k[\phi] = \int d^d x \left(\frac{1}{2} Z_k (\partial\phi)^2 + V_k(\phi) \right)$$

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- $Z_k = 1$: LPA

$$\rightarrow \eta = 0, \Delta = \frac{d-2}{2}$$

- $Z_k \neq 1$: LPA'

$$\rightarrow \eta \neq 0, \Delta = \frac{d-2+\eta}{2}$$

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Fixed point equation

$$d v(\varphi) - \Delta \varphi v'(\varphi) = \left(1 - \frac{\eta}{d+2} \right) \frac{c_d}{1 + v''(\varphi)}$$

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$$\eta = 0 \quad \text{LPA}$$

$$\eta = \frac{c_d v'''(\varphi_0)^2}{(1 + v''(\varphi_0))^4} \quad \text{LPA'}$$

Perturbative treatment

Goal: Construct an ε -expansion in $d = d_c - \varepsilon$

→ Linearize by expanding around the Gaussian solution, $v(\varphi) = v_0 + \tilde{v}(\varphi)$

$$d v(\varphi) - \Delta \varphi v'(\varphi) - c_d \left(1 - \frac{\eta}{d+2}\right) \frac{1}{1 + v''(\varphi)} = 0$$

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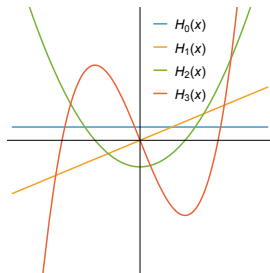
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→ Polynomially bounded solution: Hermite polynomials

$$\tilde{v}(\varphi) \simeq \lambda H_n(\varphi), \quad n = \frac{2d}{d-2}$$



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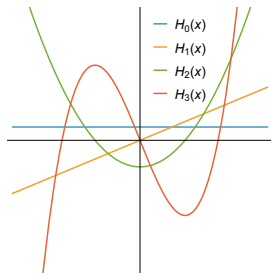
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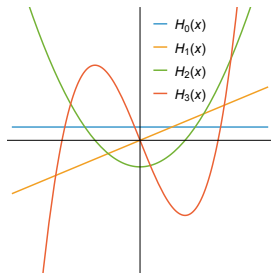
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Perturbative potential ($n = 1, d = 6 - \varepsilon$)

$$v(\varphi) = v_0 + \varepsilon^{1/2} \lambda H_3(\varphi) + \varepsilon \lambda^2 \sum_{q=0}^{\infty} \alpha_q H_q(\varphi) + \mathcal{O}(\varepsilon^{3/2})$$

$$\simeq 1.4 \cdot 10^{-5} + i \varepsilon^{1/2} \left(6.1 \varphi^3 - 7.6 \cdot 10^{-4} \varphi \right) - \varepsilon \left(5.6 \cdot 10^{-2} \varphi^2 - 3.1 \cdot 10^{-5} \right)$$

Complex potential

Coupled fixed point equations

Given $v(\varphi) = u(\varphi) + ih(\varphi)$,

$$u - \frac{\Delta}{d} \varphi u' = \frac{1 + u''}{(1 + u'')^2 + (h'')^2}, \quad h - \frac{\Delta}{d} \varphi h' = -\frac{h''}{(1 + u'')^2 + (h'')^2}$$

Complex potential

Coupled fixed point equations

Assume $u(\varphi) = 0$,

$\mathcal{PT}$ -symmetry $v^*(\varphi) = v(-\varphi)$

$$h - \gamma \varphi h' = -\frac{h''}{1 + (h'')^2}, \quad \gamma = \frac{\Delta}{d}, \quad h(0) = 0, \quad h'''(0) = g \in \mathbb{R}$$

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Large field $\varphi \rightarrow \pm\infty$ behaviour

- $0 < \Delta < d/2$: Two (likely) isolated solutions
 - Gaussian solution (isolated) $h = 0$
 - nontrivial solution (one parameter family)

$$h(\varphi) \sim B \varphi^{\frac{1}{\gamma}} - \frac{\gamma^2}{2B(1-\gamma)^2} \varphi^{2-\frac{1}{\gamma}} + \mathcal{O}\left(\varphi^{4-\frac{2}{\gamma}}\right)$$

- $\Delta < 0$: Continuous (two parameters) family

$$h(\varphi) \sim b_1 \varphi^{\frac{1}{\gamma}} \left(1 + \mathcal{O}(\varphi^{-2})\right) + b_2 e^{\frac{\gamma}{2}\varphi^2} \varphi^{-\frac{1}{\gamma}-1} \left(1 + \mathcal{O}\left(\varphi^{-2}, e^{\gamma\varphi^2} \varphi^{-\frac{2}{\gamma}-2}\right)\right)$$

Newton equation

$$h - \gamma \varphi h' = -\frac{h''}{1 + (h'')^2}, \quad h(0) = 0, \quad h'''(0) = g$$

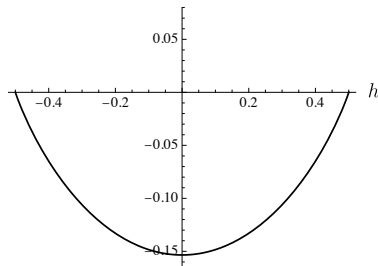
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$$h'' = -\frac{d}{dh}U(h), \quad h(0) = 0, \quad h'''(0) = -h'(0) = g$$

$$U(h) = \frac{\ln(1 - \sqrt{1 - 4h^2}) + \sqrt{1 - 4h^2}}{2}$$

→ Energy conservation:

$$\varphi = \frac{1}{\sqrt{2}} \int_0^h \frac{dz}{\sqrt{E - U(z)}}$$



Vanishing Δ

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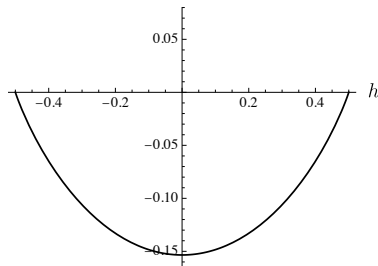
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LPA' upper bound

For any d , there exists $g = g_0(d)$ such that

$$\Delta = \left. \frac{d - 2 + \eta(d, g)}{2} \right|_{g=g_0(d)} = 0$$

→ At $d = d_0$, $E = 0 \implies g_0(d_0) = \pm\sqrt{1 - \ln 2}$



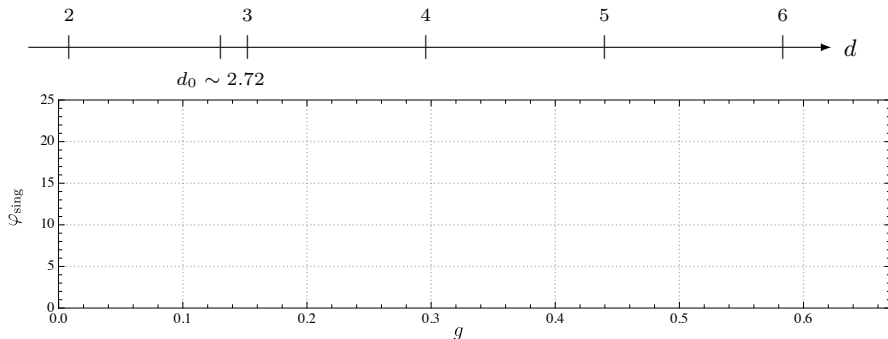
$$d_0 = 2 \frac{3 - \ln 2}{1 + \ln 2} \sim 2.7249$$

$$h - \frac{\Delta}{d} \varphi h' = -\frac{h''}{1 + (h'')^2}, \quad h(0) = 0, \quad h'''(0) = g \in \mathbb{R}$$

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- LPA' : $\eta = \frac{d+2}{2} \left(1 - \sqrt{1 + \frac{4dg^2}{(d+2)}} \right) \rightarrow \Delta = \frac{d-2+\eta}{2}$

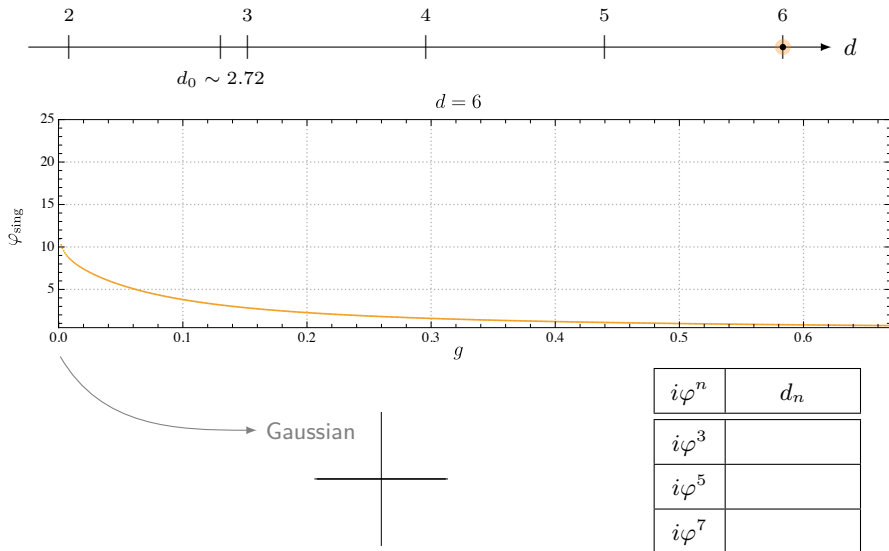
Numerical resolution



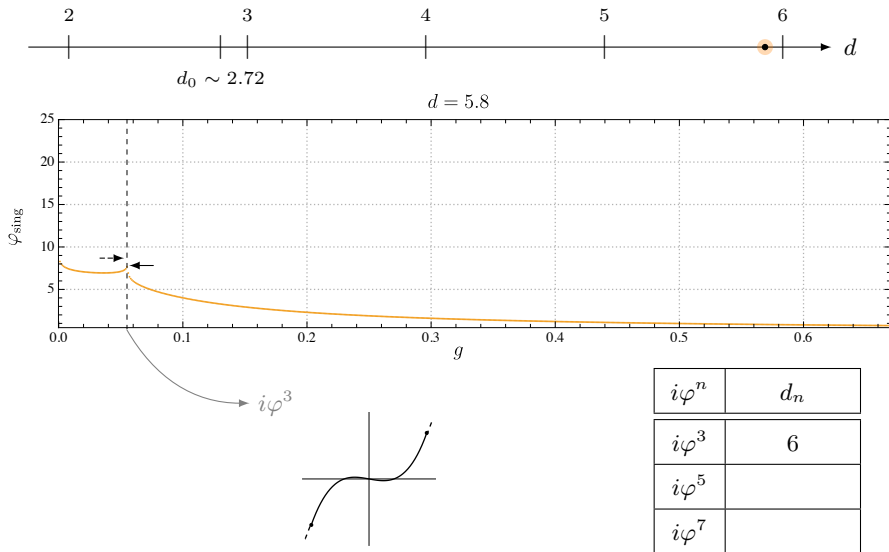
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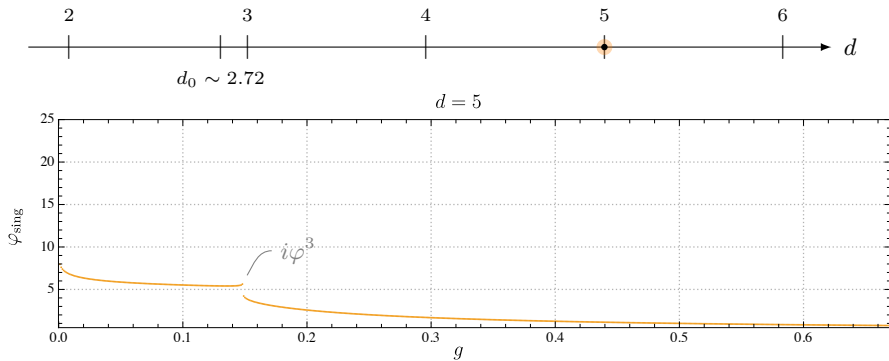
Spike plots LPA



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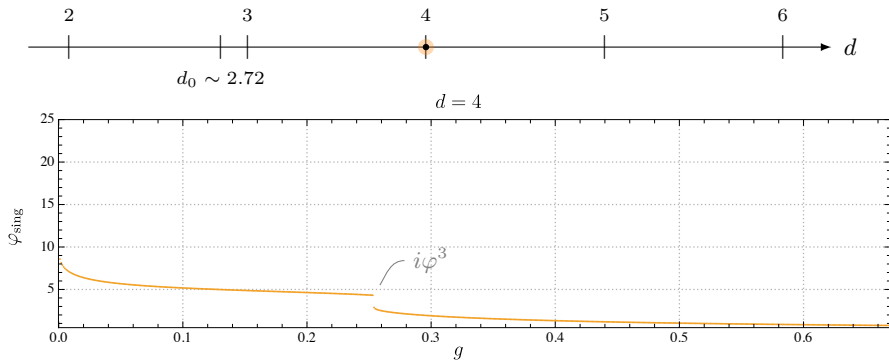


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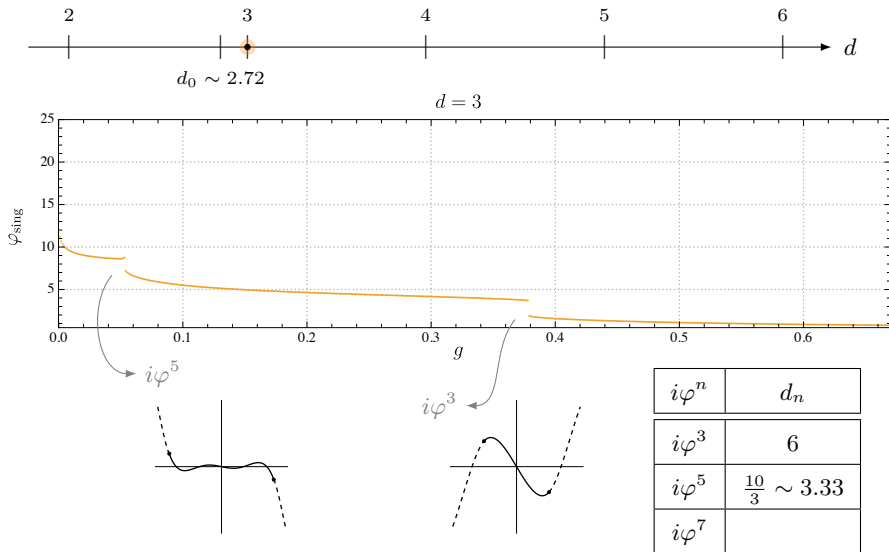
$i\varphi^n$	d_n
$i\varphi^3$	6
$i\varphi^5$	
$i\varphi^7$	

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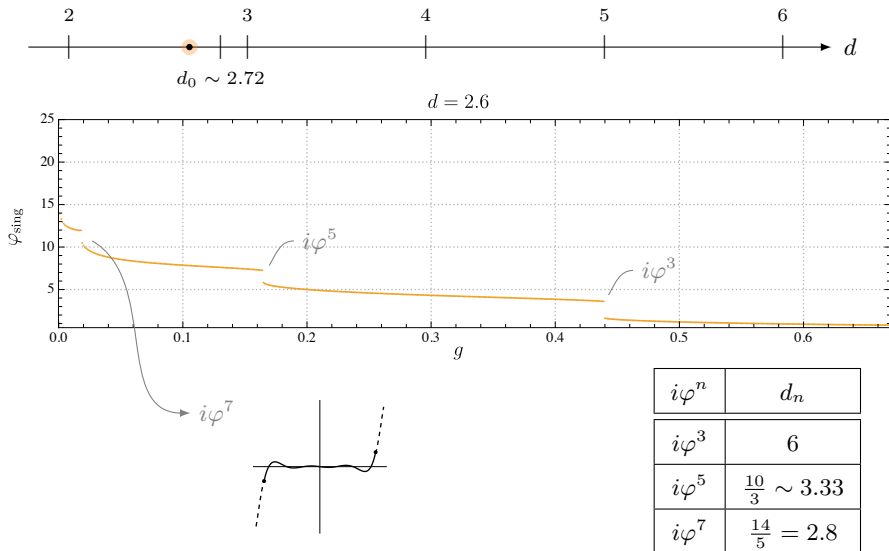


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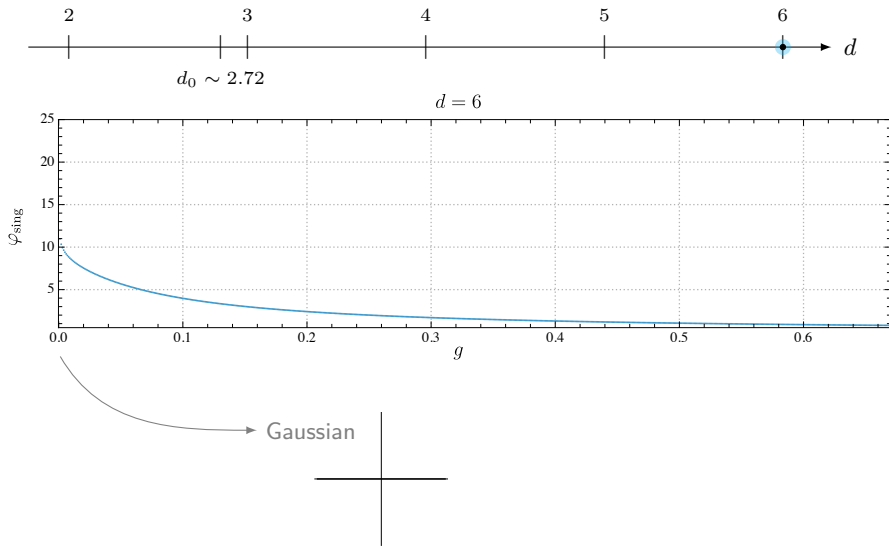
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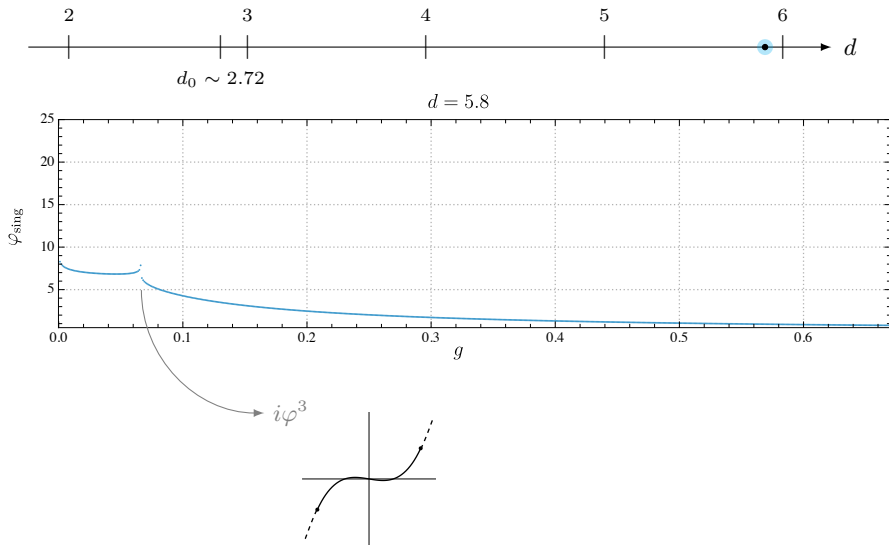
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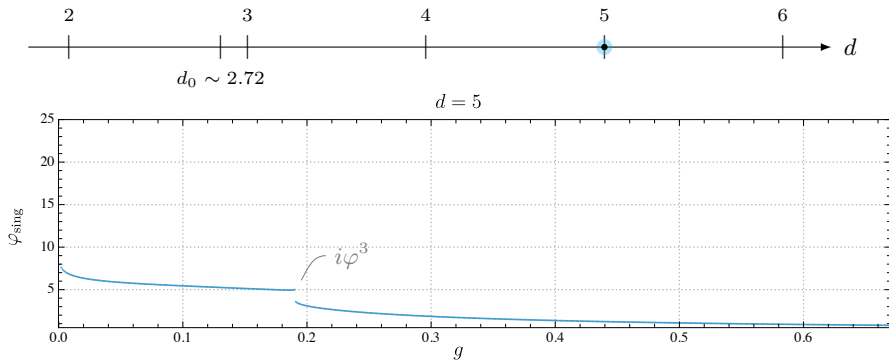
Spike plots LPA'



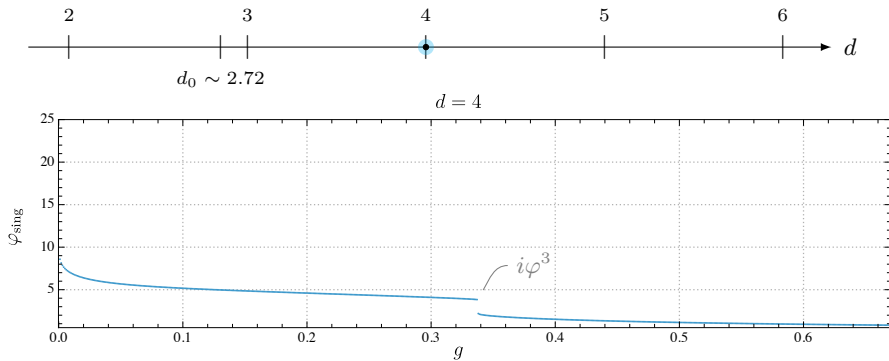
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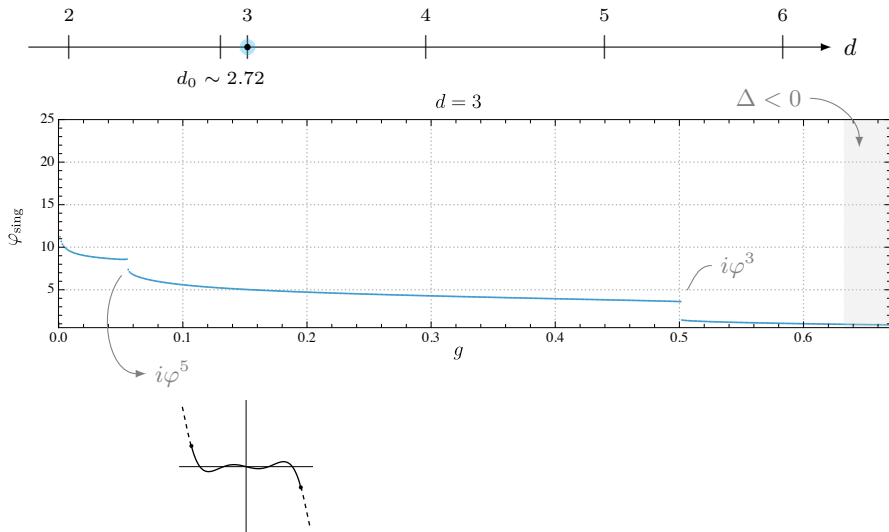
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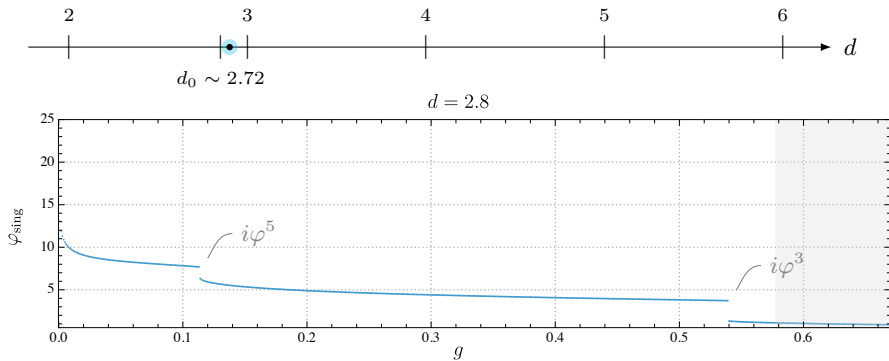
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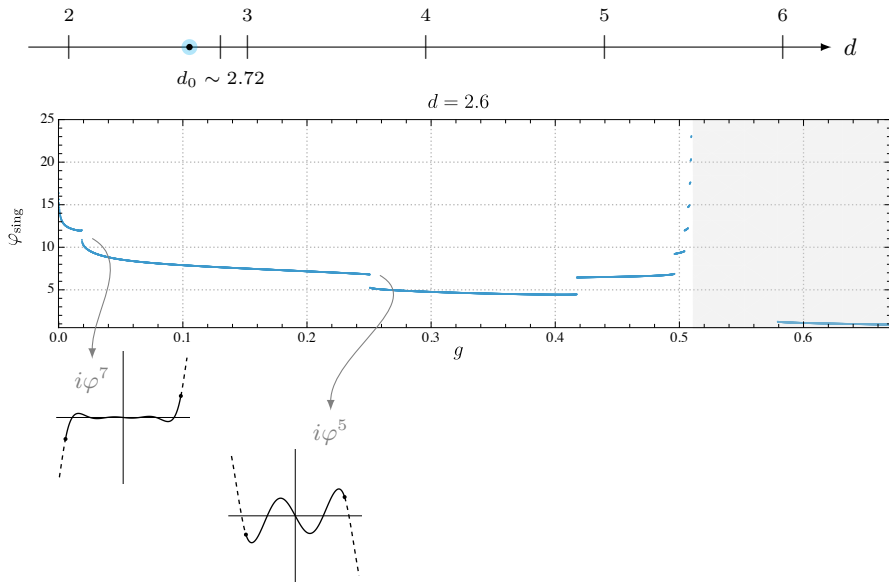
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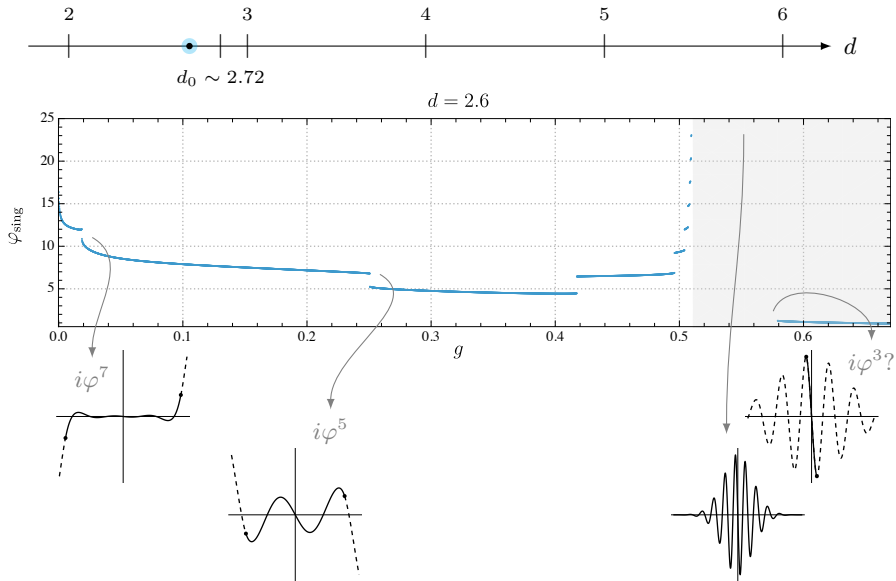
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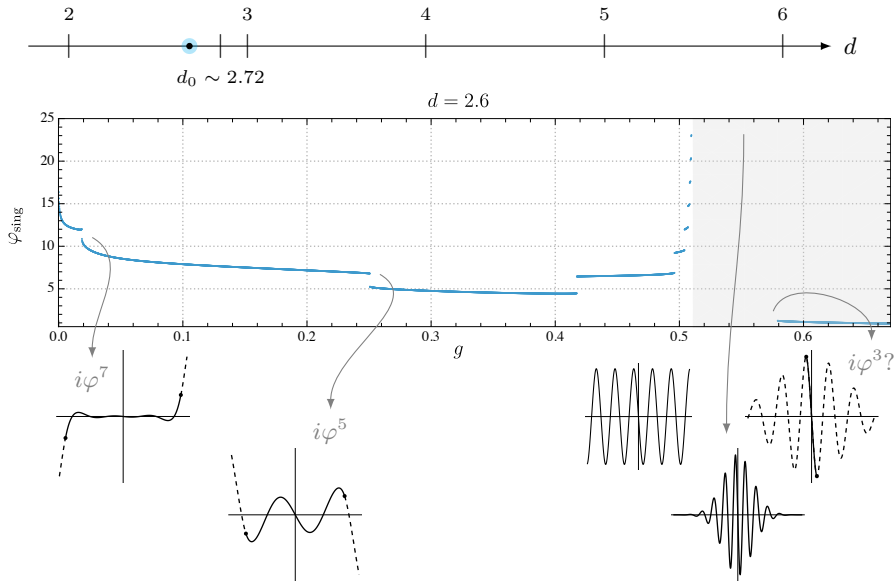
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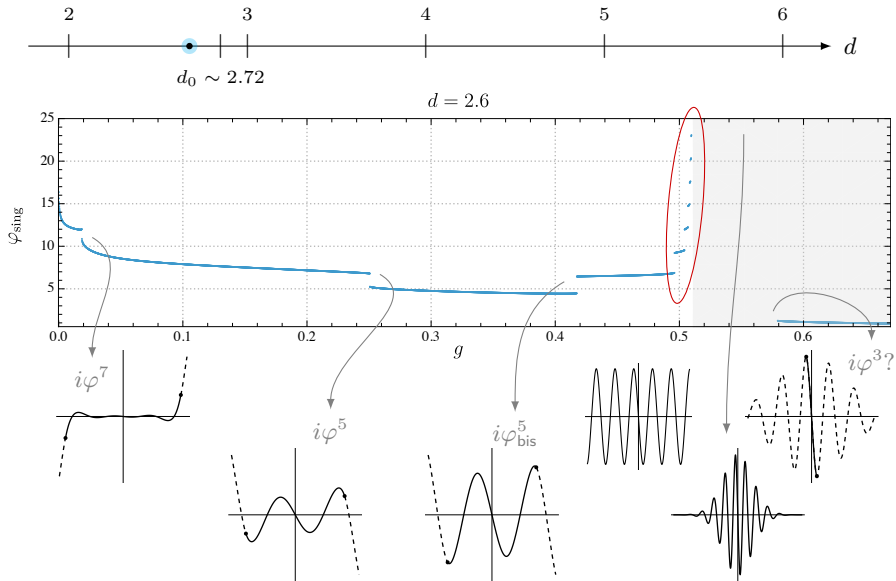
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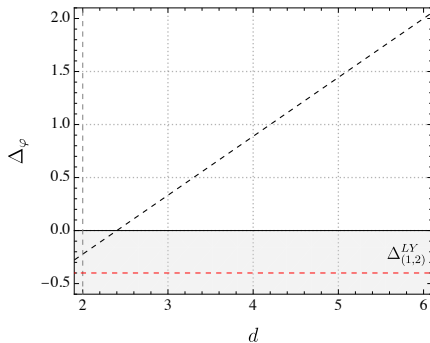
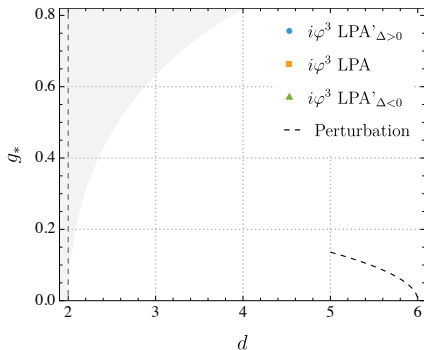


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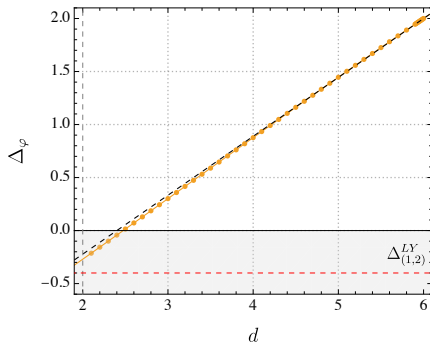
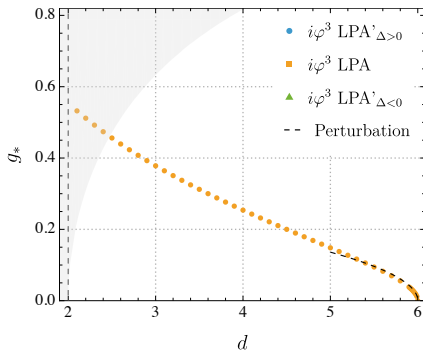


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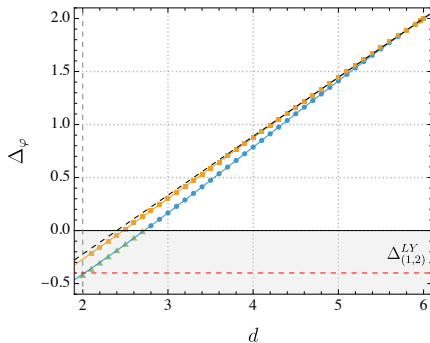
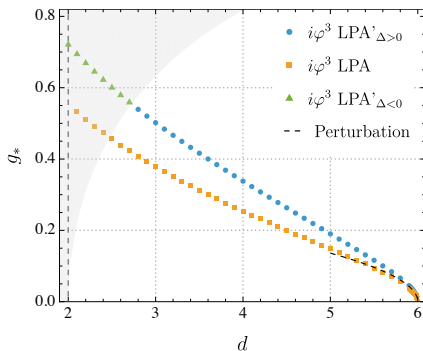




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5-loops [Borinsky et al., 21']	1.4245(10)	0.827(6)	0.215(10)	-0.390(15)
CFT $\mathcal{M}_{(2,5)}$				-0.4

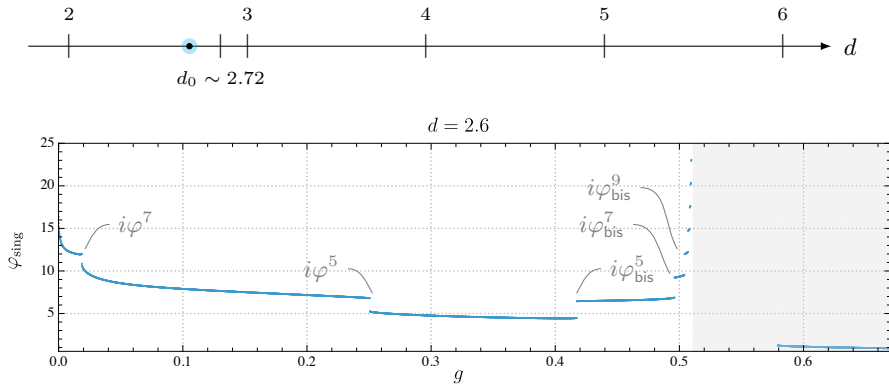


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LPA at $\varphi = 0$	1.44587	0.87621	0.30122	-0.27004
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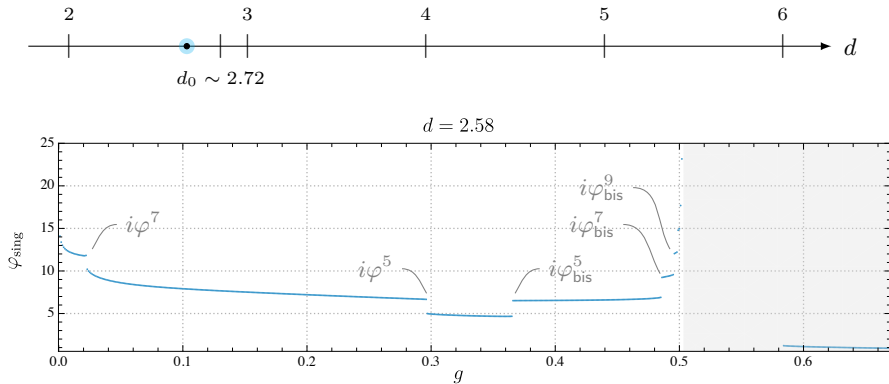


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LPA' $\Delta_{>0}$	1.41178	0.78675	0.16659	-0.41035
LPA' $\Delta_{<0}$				-0.42809
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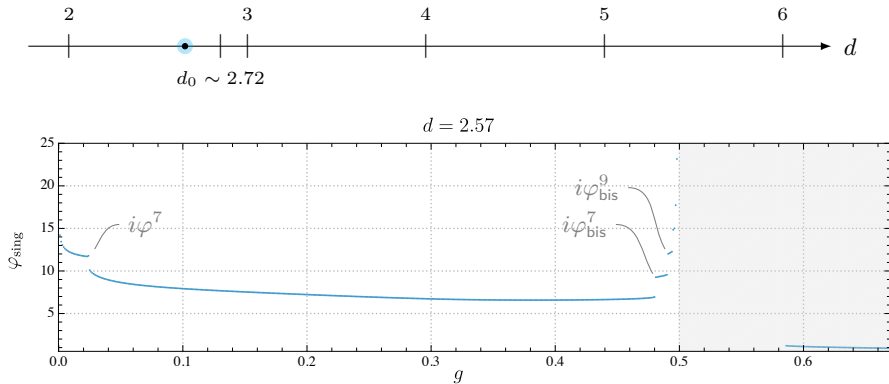
Back to spike plots (LPA')

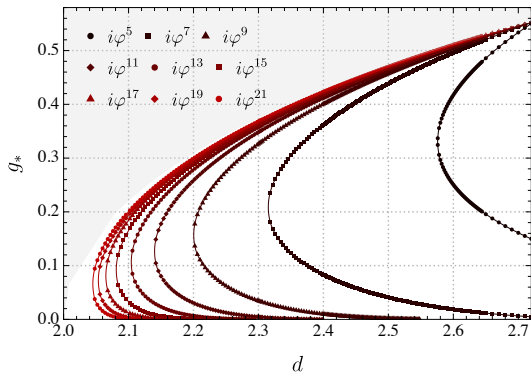


Back to spike plots (LPA')



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Fixed points $d < d_0$ 

$i\varphi^n$	d_c
$i\varphi^3$	/
$i\varphi^5$	2.5758
$i\varphi^7$	2.315
$i\varphi^9$	2.200
$i\varphi^{11}$	2.1399
$i\varphi^{13}$	2.104
$i\varphi^{15}$	2.080

$$h - \frac{\Delta}{d} \varphi h' = -\frac{h''}{1 + (h'')^2}, \quad h(0) = 0, \quad h'''(0) = g \in \mathbb{R}$$

- **Perturbation theory:** Identified weakly coupled $i\varphi^{2n+1}$ theories near $d = d_{2n+1} - \varepsilon$
- Classified **global solutions** of the flow equation
 - $\Delta > 0$: isolated global solutions
 - $\Delta < 0$: continuum of solutions
- **Numerical results:**
 - $\Delta_\varphi|_{d=2}$ with an error of 2.6% – 7% compared to $\mathcal{M}(2, 5)$ for $i\varphi^3$
 - $d = d_0$, the value at which $\Delta = 0$ for $i\varphi^3$
 - infinite sequence of non-perturbative fixed points annihilating with the $i\varphi^{2n+1}$, $n > 1$ fixed points

Outlook

- Improve the benchmark for $i\phi^3$ theory
- Understand the non-perturbative fixed points...