

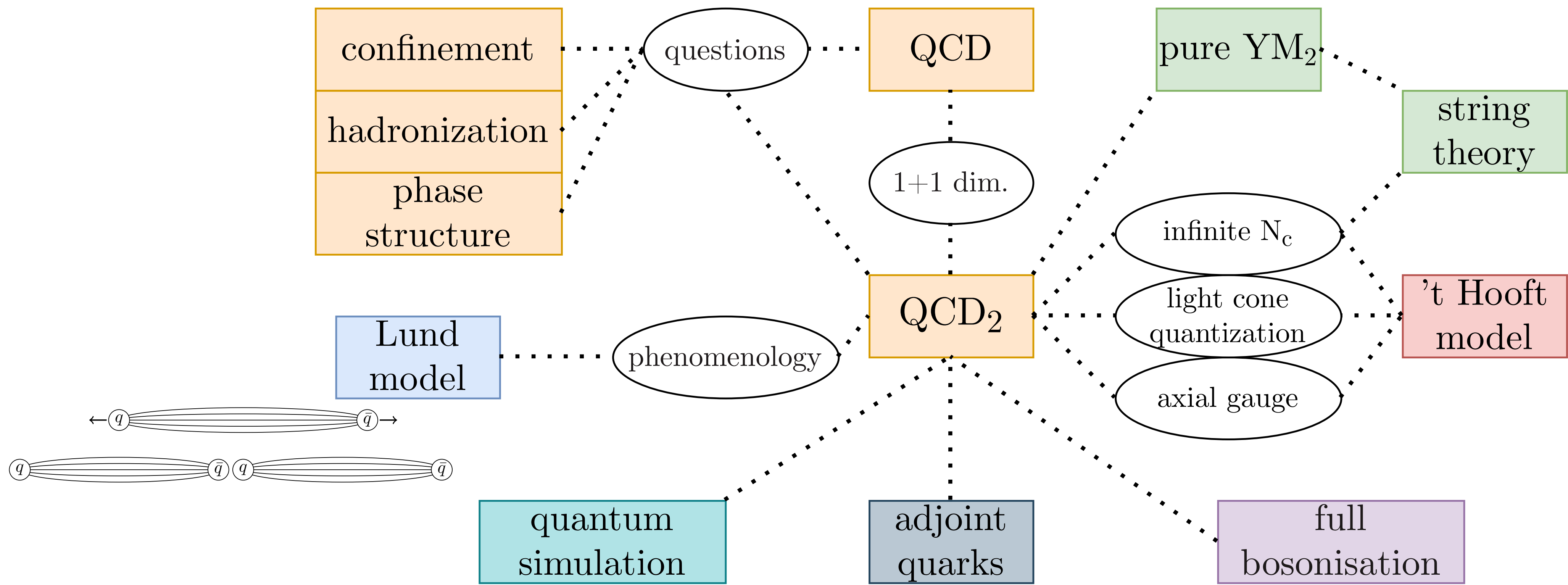
Image of Jena: [Entdecke-Deutschland.de](https://www.entdecke-deutschland.de), 2024

Functional Renormalization of QCD in $1+1$ dimensions

Eric OEVERMANN
Friedrich Schiller University Jena

ERG 2026
02 September 2026

Contextualization



For image and rererences see: EO, Adrian Koenigstein, Stefan Floerchinger, Phys.Rev.D 111 (2025)

Contents

1. Four-fermion interactions from quark-gluon dynamics

- QCD in $1+1$
- Regulator properties
- Fierz completeness

2. Non-relativistic theory

- Access full two-body quark-antiquark bound state and scattering spectrum

Functional renormalization of QCD in $1+1$ dimensions: four-fermion interactions from quark-gluon dynamics

EO, Koenigstein, Floerchinger, Phys.Rev.D 111 (2025) : first FRG paper on QCD in $1+1$

QCD in 1+1 dimensions

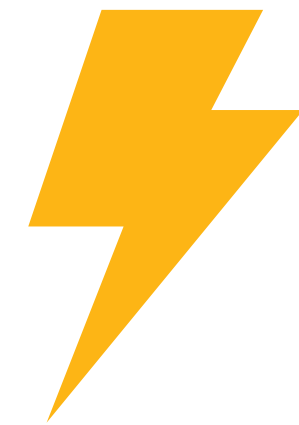
Microscopic action

$$S = \int d^2x \left[\bar{\psi} [\gamma^\mu (\partial_\mu - i A_\mu) + m] \psi + \frac{1}{2g^2} \text{tr}(F_{\mu\nu} F^{\mu\nu}) \right]$$



- A. Gauge coupling g of positive energy dimension: **super-renormalizable theory**
- B. **No propagating gauge DoF**
- C. Quarks: 2 components, **no spin/angular DoF**

Need **gauge fixing** and **IR regularization**



Technical steps on gauge fixing:

- Background & fluctuation field splitting
 $A = \bar{A} + a$
- Gauge fixing term & ghosts

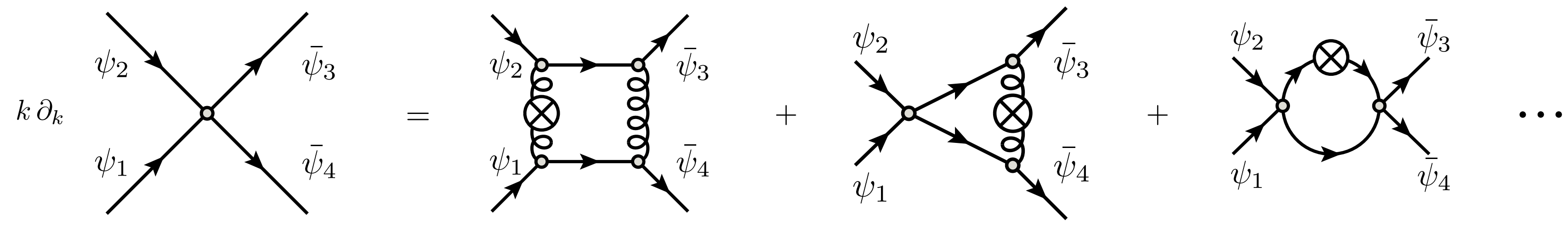
Technical steps on regularization:

- Callan-Symanzik type / masslike / momentum-independent regulators
- Preserve Osterwalder-Schrader axioms

Fluctuation field gauge variance and background independence are encoded in identities that are unmodified by the regulator presence!

Local four-fermion interactions

Truncation of Γ_k : **Fierz complete** set of local four-fermion interactions



$$k \partial_k \lambda_i = \# \left(\frac{g}{k} \right)^4 + \sum_{j=1}^3 \#_j \left(\frac{g}{k} \right)^2 + \sum_{j,l=1}^3 \#_{jl} \lambda_j \lambda_l$$

Flow equation can be used to
determine a Fierz complete basis

[procedure similarly described in Braun, Geißel, Pawłowski, Sattler and Wink (2025) 'Juggling with tensor bases in functional approaches']

3 interaction terms needed

1. Write down over complete set of interactions
2. Take 4 functional derivatives of flow eq.
3. Project the equations / contract indices with projectors

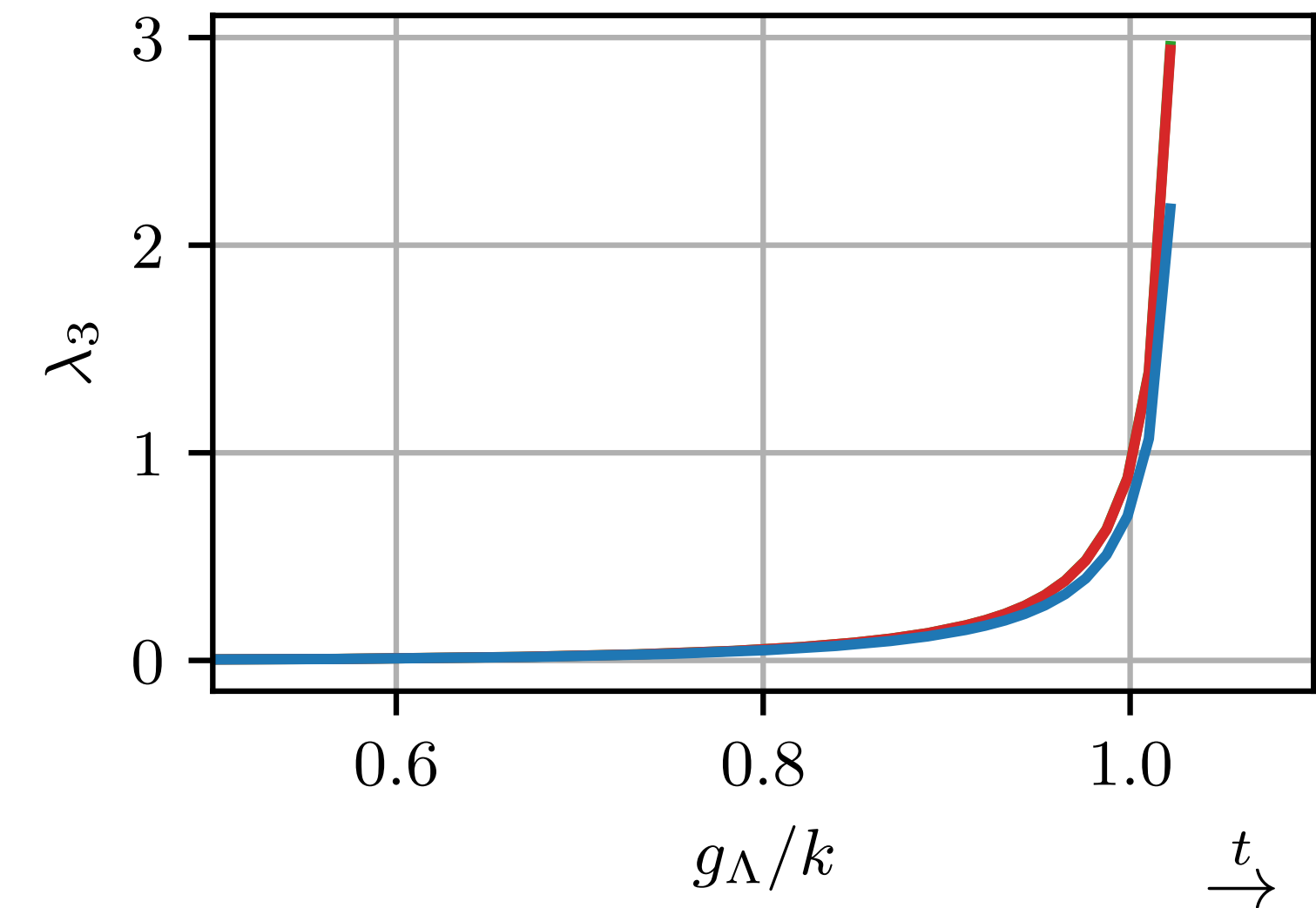
$$M_{ij}(k \partial_k \lambda_j) = b_i$$

4. Reduce system until it becomes invertible

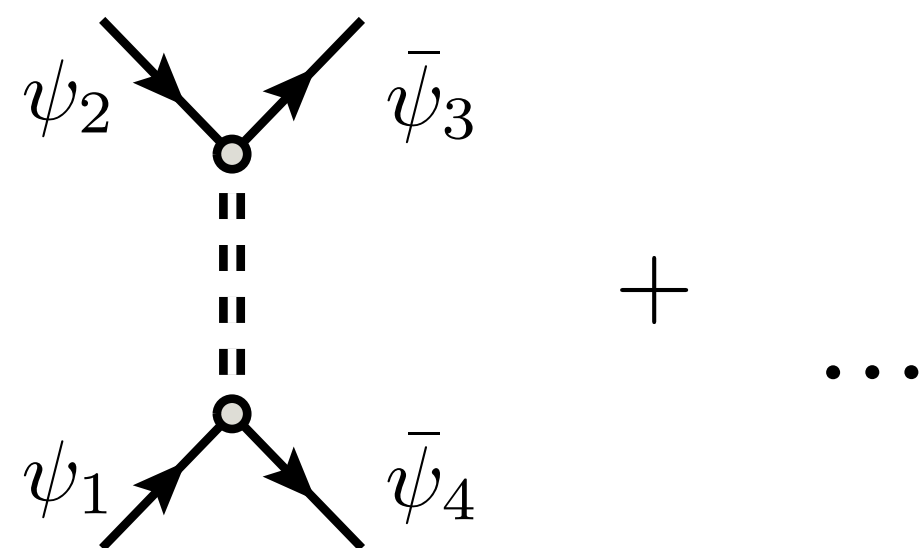
(Some) Results

- Scale at which **divergence occurs given by g_Λ**
- Divergence of four-fermion coupling manifest as soon as non-zero (dimensionless couplings)
- Our paper: further discussions, e.g. on large N

[Cresswell-Hogg and Litim (2025) with agreements]



- Next: **resolve momentum dependence** of four-fermion vertex, include bound states



Non-relativistic QCD in $1+1$ dimensions

EO, Steller, Floerchinger (work in progress)

Example for fully calculable two-body quark-antiquark bound state spectrum via a Hubbard-Stratonovich transformation

Non-relativistic action

$$S = \int_{\mathbb{R}^2} dt dx \left[-\frac{1}{4g^2} F_{\mu\nu}^z F^{\mu\nu z} + \chi_c^\dagger i (\partial_t - iA_0^z T^z) \chi + \frac{1}{2m} \chi_c^\dagger (\partial_x - iA_1^z T^z)^2 \chi \right. \\ \left. + \tilde{\chi}_c^\dagger i (\partial_t - iA_0^z T^z) \tilde{\chi} + \frac{1}{2m} \tilde{\chi}_c^\dagger (\partial_x - iA_1^z T^z)^2 \tilde{\chi} \right]$$

EoM are type
Schrödinger eqs.

- Quark field χ and anti-quarks $\tilde{\chi}$ transform under $U \in \text{SU}(N)$ as

$$\chi \mapsto U \chi, \quad \tilde{\chi} \mapsto U^* \tilde{\chi}$$

- Independent global $U(1)$ symmetry for quarks and anti-quarks

➡ Time dependent $U(1)$ symmetry: add a potential to temporal derivative

- **Galilei invariance** of all terms in the action!

Gauge fixing and IR regularisation

Spatial axial gauge

$$A_1^z = 0$$

- Static gluons
- No gluon self-interactions

$$S = \int_{\mathbb{R}^2} dt dx \left[\frac{1}{2g^2} A^{0z} \partial_x^2 A_0^z + (\chi^\dagger T^z \chi - \tilde{\chi}^\dagger (T^z)^T \tilde{\chi}) A^{0z} + \chi_c^\dagger \left(i \partial_t + \frac{1}{2m} \partial_x^2 \right) \chi^c + \tilde{\chi}_c^\dagger \left(i \partial_t + \frac{1}{2m} \partial_x^2 \right) \tilde{\chi}^c \right]$$

Galilei invariance preserved

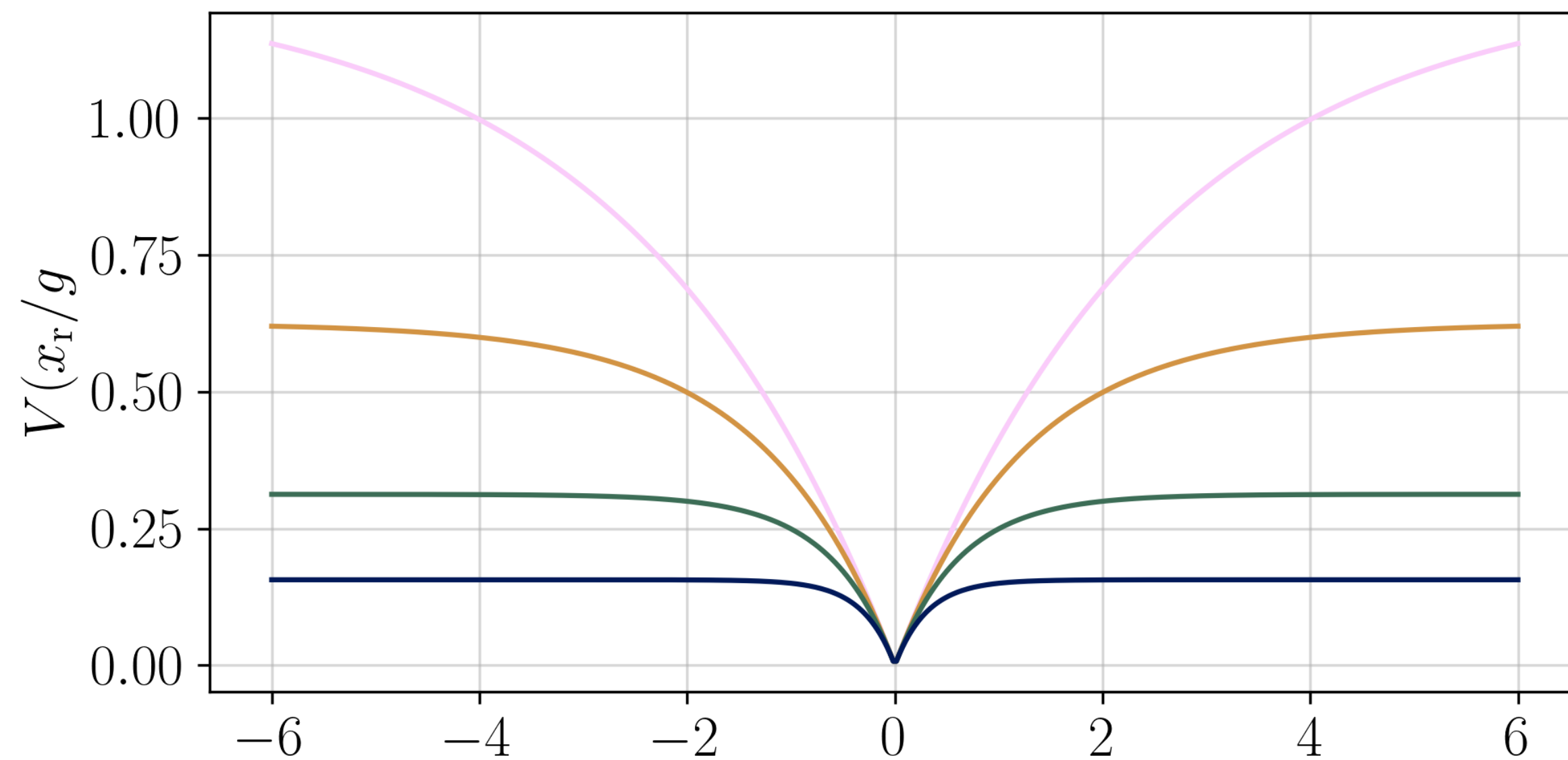
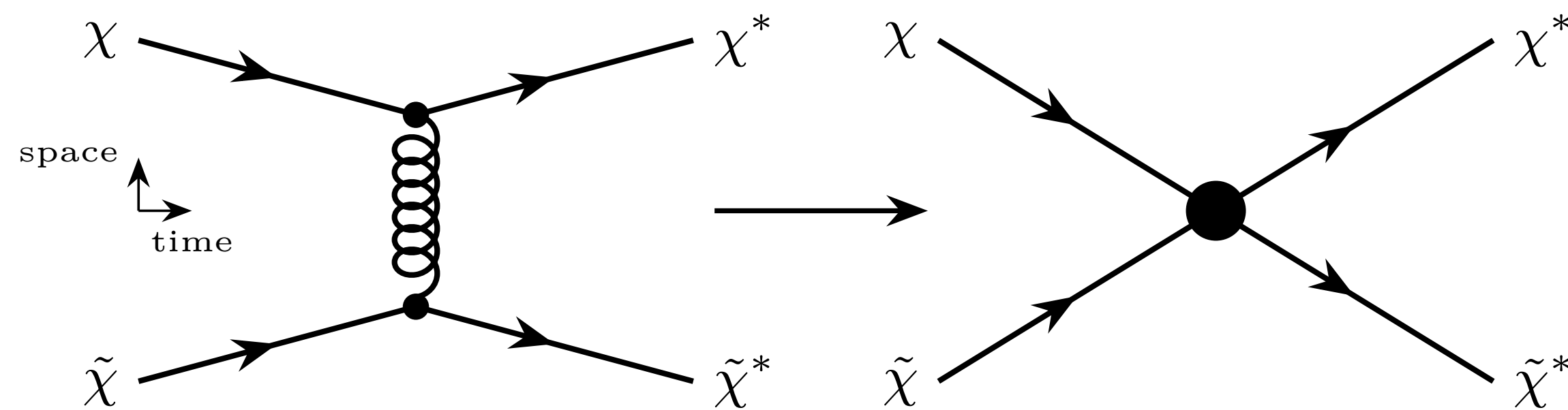
Adding a gluon mass term [c.f. regularization in Goldman, Nefediev and Wagebrunn '25]

$$\int_{\mathbb{R}^2} dt dx \left(-\frac{1}{4g^2} F_{\mu\nu}^z F^{\mu\nu z} - \frac{m_{\text{gl}}^2}{2g^2} A_\mu^z A^{\mu z} \right)$$

keeps Galilei invariance after gauge fixing

Integrate out gauge field

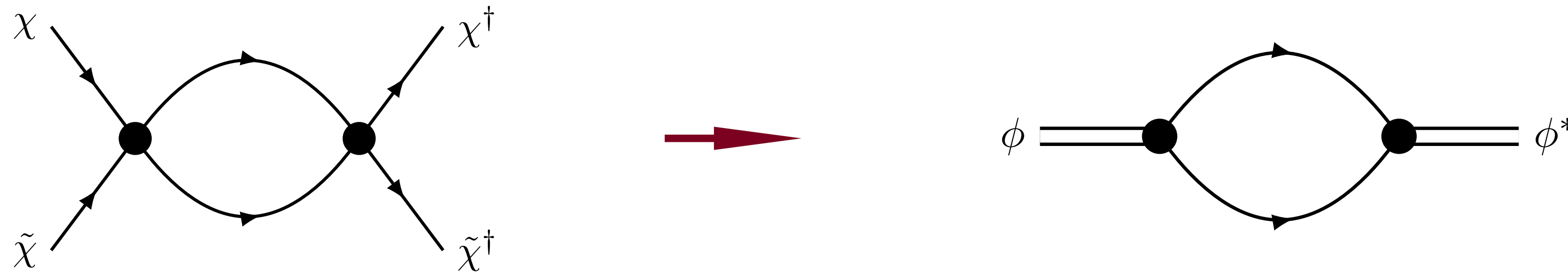
Transition to purely fermionic description [c.f. e.g. Alkhofer and Reinhardt '95]



[Floerchinger '10]

Use FRG to connect field description and two-body Schrödinger equation

I. Introduce bi-local field via a HS transformation



2. Solve flow equation for propagator $G_\phi = P_\phi^{-1}$ of ϕ

$$\Gamma_k^{(\phi,2)} = \int_{E,p,q,q'} \phi^*(E,p,q) P_{\phi,k}(E,p,q,q') \phi(E,p,q')$$

Non-relativistic limit:

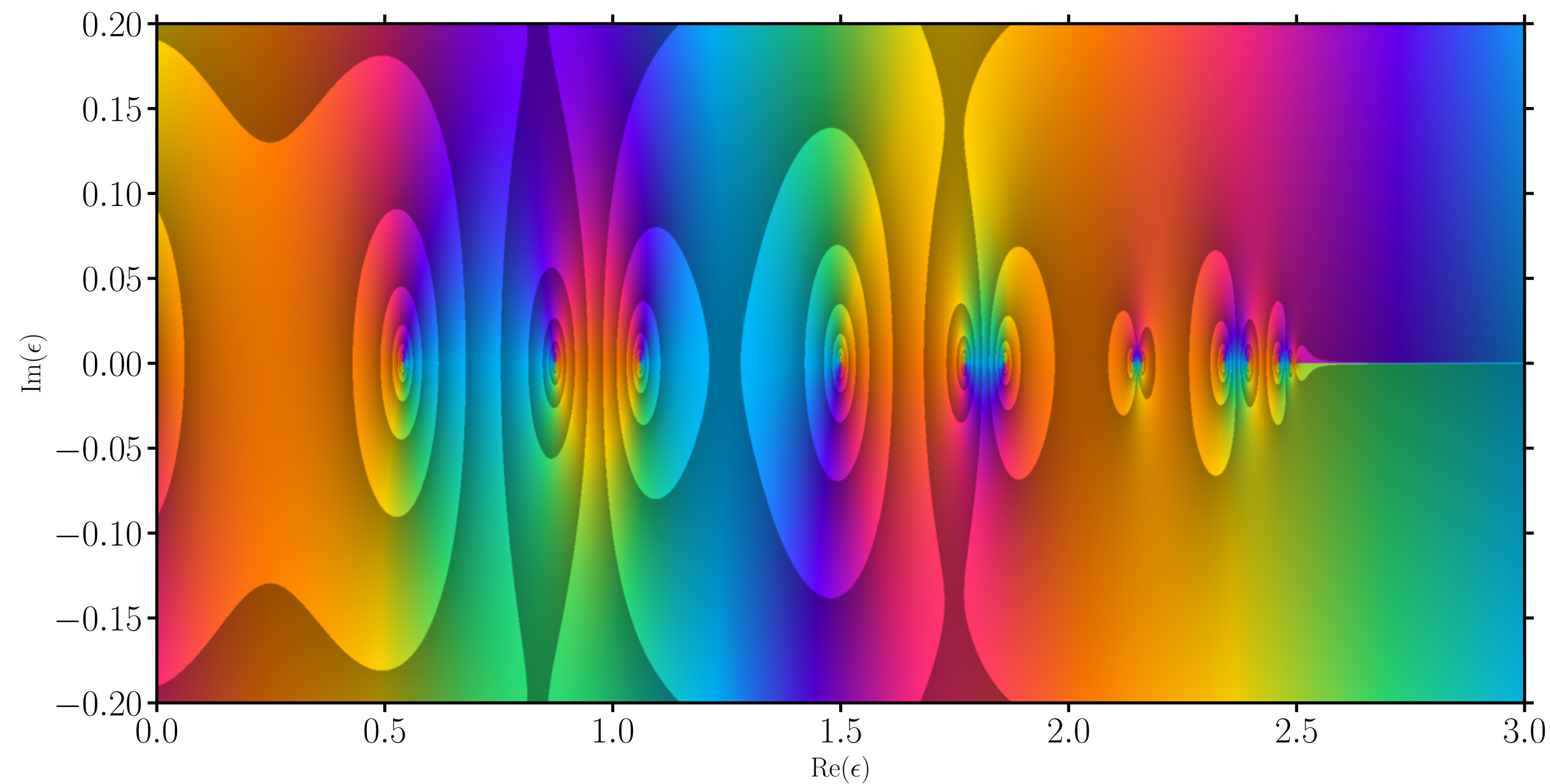
- Only non-tour diagrams
- N-body problems decouple

Spectrum

Poles of $G_\phi(E)$ are poles of $(E - H)^{-1}$ with

$$H = -\frac{1}{4m} \frac{\partial^2}{\partial x_c^2} - \frac{1}{m} \frac{\partial^2}{\partial x_r^2} + V_1(x_r)$$

➡ Solve Schrödinger equation, Jost / scattering solutions



Full Solution

$$\phi^*(E, p, x_r) \bar{G}_\phi(E, p, x_r, x'_r) \phi(E, p, x'_r) = \sum_n \frac{\varphi_n^*(x_r) \varphi_n(x'_r)}{E - E_c - E_n} + \int_{4/9g^2/(2m_{gl})}^{\infty} \frac{dE_r}{2\pi} \frac{1}{E - E_c - E_r} \sum_{i \in \{L, R\}} \left(\overline{\psi_i(E_r, x'_r)} \psi_i(E_r, x_r) \Theta(x_r - x'_r) + \overline{\psi_i(E_r, x_r)} \psi_i(E_r, x'_r) \Theta(x'_r - x_r) \right)$$

- Can express all quantities explicitly and relate to Jost function/solution

IR limit:

- Linear potential
- Continuum removed from spectrum

UV limit:

- Delta potential
- One bound state of vanishing energy

Thank you! Questions?

Summary

- QCD in $1+1$ with the FRG: minimal truncation in glue and quark sector but with Fierz complete set of local four-fermion interactions
 - A. Discussed properties of Callen-Symanzik regulator
 - B. Fierz completeness determined by flow equation
 - C. Divergence of couplings signals formation of mesonic degrees of freedom
- Non-relativistic QCD in $1+1$ as example for fully calculable two-body quark-antiquark bound state spectrum, e.g., via a Hubbard-Stratonovich transformation

*Back-up

Equations

$$S = \int d^2x \left[\bar{\psi} (\gamma^\mu D_\mu + m) \psi + \frac{1}{2g^2} \text{tr}(F_{\mu\nu} F^{\mu\nu}) + \frac{1}{\xi g^2} \text{tr}((D_\mu[\bar{A}] a^\mu) (D_\nu[\bar{A}] a^\nu)) - 2 \text{tr}(\bar{c} D^\mu[\bar{A}] D_\mu[A] c) \right]$$

$$\Delta S_k = \int d^2x \left[\bar{\psi} k \psi + \frac{1}{g^2} \text{tr}(a_\mu k^2 a^\mu) + 2 \text{tr}(\bar{c} k^2 c) \right] \quad V_1(x_r) = -\frac{4}{9} \frac{g^2}{2m_{\text{gl}}} \left(1 - e^{-m_{\text{gl}}|x_r|} \right)$$

$$S = \int_{t,x} \left[\chi^\dagger \left(i \partial_t + \frac{1}{2m} \partial_x^2 \right) \chi + \tilde{\chi}^\dagger \left(i \partial_t + \frac{1}{2m} \partial_x^2 \right) \tilde{\chi} \right] - \int_{t,x_c,x_r} \left[(\tilde{\chi}^T \chi)^\dagger(t, x_c, x_r) V_1(x_r) (\tilde{\chi}^T \chi)(t, x_c, x_r) + \dots \right]$$

$$Z = \int \mathcal{D}\chi^\dagger \mathcal{D}\chi \mathcal{D}\tilde{\chi}^\dagger \mathcal{D}\tilde{\chi} e^{iS[\chi, \tilde{\chi}, \chi^\dagger, \tilde{\chi}^\dagger]} \propto \int \mathcal{D}\chi^\dagger \mathcal{D}\chi \mathcal{D}\tilde{\chi}^\dagger \mathcal{D}\tilde{\chi} \mathcal{D}\phi^* \mathcal{D}\phi e^{i(S+S_{\text{pb}})}$$

$$S_{\text{pb}} = \int_{t,x_c,x_r} \left[\phi^*(t, x_c, x_r) - (\tilde{\chi}^T \chi)^\dagger(t, x_c, x_r) \right] V_1(x_r) \left[\phi(t, x_c, x_r) - (\tilde{\chi}^T \chi)(t, x_c, x_r) \right]$$