

Linearity in the renormalization group equations

Denis Tolkachev

In collaboration with D. Kazakov, A. Mukhaeva and R. Iakhibbaev



Bogoliubov Laboratory of Theoretical Physics
Joint Institute for Nuclear Research



Single operator model

The Callan-Symanzik Equation

$$\left(\mu^2 \frac{d}{d\mu^2} + \beta(\alpha) \frac{d}{d\alpha} - \gamma(\alpha) \right) \Gamma \left(\frac{Q^2}{\mu^2}, \alpha \right) = 0$$

Running coupling

$$\frac{\partial}{\partial L} \bar{\alpha}[\alpha, L] = \beta(\bar{\alpha}) = \beta_0 \bar{\alpha}^2 + \beta_1 \bar{\alpha}^3 + \beta_2 \bar{\alpha}^4 + \beta_3 \bar{\alpha}^5 + O(\bar{\alpha}^6)$$

$$L = \log(Q^2/\mu^2)$$

1-loop approximation	Initial conditions	Solutions	Exact solutions for higher orders of the running coupling are not known
$\frac{\partial}{\partial L} \alpha_1 = \beta_0 \alpha_1^2$	$\alpha_1[\alpha, 0] = \alpha$	$\alpha_1 = \frac{\alpha}{1 - \beta_0 \alpha L}$	
2-loop approximation		The solution contains a non-elementary Lambert function	?
$\frac{\partial}{\partial L} \alpha_2 = \beta_0 \alpha_2^2 + \beta_1 \alpha_2^3$	$\alpha_2[\alpha, 0] = 0$	$\alpha_2 = \frac{1}{1 + W(\beta_0, \beta_1, L)}$	$\alpha_3, \alpha_4 \dots$

Analysis of coefficients in perturbative expansion

Let us consider the n-loop approximation and the coefficients of the corresponding logarithms

$$f_1(\beta_0)\log(x)^N + f_2(\beta_0, \beta_1)\log(x)^{N-1} + f_3(\beta_0, \beta_1, \beta_2)\log(x)^{N-2} + \dots \quad x = Q^2/\mu^2$$
$$\frac{\partial}{\partial L}\alpha_1 = \beta_0\alpha_1^2 \quad \frac{\partial}{\partial L}\alpha_2 = \beta_0\alpha_2^2 + \beta_1\alpha_2^3 \quad \frac{\partial}{\partial L}\alpha_3 = \beta_0\alpha_3^2 + \beta_1\alpha_3^3 + \beta_2\alpha_3^4$$

However, it turns out that you can extract the third power of the logarithm from the equation in a subleading approximation

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$$\frac{\partial}{\partial L}\alpha_1 = \beta_0\alpha_1^2$$

$$\frac{\partial}{\partial L}\alpha_2 = \beta_0\alpha_2^2 + \beta_1\alpha_2^3$$

$$\frac{\partial}{\partial L}\alpha_3 = \beta_0\alpha_3^2 + \beta_1\alpha_3^3 + \beta_2\alpha_3^4$$

However, it turns out that you can extract the third power of the logarithm from the equation in a subleading approximation

$$\alpha_2 = \alpha_2(\beta_0, \beta_1) + \cancel{\alpha_2(\beta_0, \beta_1, \beta_2, \beta_3 \dots)}$$

Running coupling

Running coupling is the sum of logarithms in all orders of the PT

$$\bar{\alpha}[\alpha, L] = \alpha \left(1 + \sum_{k=1}^{\infty} A_k \alpha^k L^k + \sum_{k=2}^{\infty} B_k \alpha^k L^{k-1} + \sum_{k=3}^{\infty} C_k \alpha^k L^{k-2} + \dots \right) \quad L = \log(Q^2/\mu^2)$$

$$\bar{\alpha}[\alpha, L] = \sum_{i=1}^{\infty} \alpha_i = \alpha_1 + \alpha_2 + \alpha_3 \dots$$

How can we extract the corresponding α_n ?

Kazakov, et al *Phys.Rev.D* 113 (2026) 12, 125021

Kazakov, et al *JHEP* 04 (2023) 128

In a series of our papers, we developed a method for summing logarithmic corrections based on the locality of Quantum field theory and the Bogoliubov-Parasiuk theorem.

1) A general method that is applicable regardless of the renormalizability of the model

2) the resulting equation is always linear

Running coupling

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1) A general method that is applicable regardless of the renormalizability of the model

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$$\alpha[\alpha, L] = \sum_{i=1}^{\infty} \alpha_i = \alpha_1 + \alpha_2 + \alpha_3 \dots$$

Kazakov, et al [arXiv:2604.04766](#) , Phys.Lett.B

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Let's make the substitution into the equation for the running coupling

$$\frac{\partial}{\partial L} \bar{\alpha} = \beta(\bar{\alpha}) = \beta_0 \bar{\alpha}^2 + \beta_1 \bar{\alpha}^3 + \beta_2 \bar{\alpha}^4 + \beta_3 \bar{\alpha}^5 + O(\bar{\alpha}^6)$$

We take only a linear part

And the same order for PT $[\alpha_n] = [\alpha_1^n]$

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$$\frac{d\alpha_1}{dL} = \beta_0 \alpha_1^2$$

$$\frac{d\alpha_2}{dL} = 2\beta_0 \alpha_1 \alpha_2 + \beta_1 \alpha_1^3$$

α_1 is excluded automatically

$$\frac{d\alpha_3}{dL} = 2\beta_0 \alpha_1 \alpha_3 + \beta_0 \alpha_2^2 + 3\beta_1 \alpha_1^2 \alpha_2 + \beta_2 \alpha_1^4,$$

$$\frac{d\alpha_4}{dL} = 2\beta_0 \alpha_1 \alpha_4 + 2\beta_0 \alpha_2 \alpha_3 + 3\beta_1 \alpha_1^2 \alpha_3 + 3\beta_1 \alpha_1 \alpha_2^2 + 4\beta_2 \alpha_1^3 \alpha_2 + \beta_3 \alpha_1^5$$

$$\frac{d\alpha_5}{dL} = 2\beta_0 \alpha_1 \alpha_5 + 2\beta_0 \alpha_2 \alpha_4 + \beta_0 \alpha_3^2 + \beta_1 \alpha_2^3 + 3\beta_1 \alpha_1^2 \alpha_4 + 6\beta_1 \alpha_1 \alpha_2 \alpha_3 +$$

$$+ 4\beta_2 \alpha_1^3 \alpha_3 + 6\beta_2 \alpha_1^2 \alpha_2^2 + 5\beta_3 \alpha_1^4 \alpha_2 + \beta_4 \alpha_1^6$$

$$\alpha[\alpha, L] = \sum_{i=1}^{\infty} \alpha_i = \alpha_1 + \alpha_2 + \alpha_3 \dots$$

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Solving equations

$$\alpha_1[\alpha, L] = \frac{\alpha}{1 - \beta_0 \alpha L}$$

$$\bar{\beta}_N = \beta_N / \beta_0$$

It is important that the solutions to these equations are easy to find and expressed in elementary functions.

$$\alpha_2[\alpha, L] = \bar{\beta}_1 \alpha_1^2 \log \left(\frac{\alpha_1}{\alpha} \right)$$

$$\alpha_3[\alpha, L] = \bar{\beta}_1^2 \alpha_1^3 \left(\log^2 \left(\frac{\alpha_1}{\alpha} \right) + \log \left(\frac{\alpha_1}{\alpha} \right) \right) + \alpha_1^2 (\alpha_1 - \alpha) (\bar{\beta}_2 - \bar{\beta}_1^2)$$

$$\alpha_4[\alpha, L] = \bar{\beta}_1^3 \alpha_1^4 \left(\log^3 \left(\frac{\alpha_1}{\alpha} \right) + \frac{5}{2} \log^2 \left(\frac{\alpha_1}{\alpha} \right) - \bar{\beta}_2 (3\alpha_1 - 2\alpha) \log \left(\frac{\alpha_1}{\alpha} \right) \right) + \text{et al}$$

$$\alpha_5[\alpha, L] = \bar{\beta}_1^4 \alpha_1^5 \left(\log^4 \left(\frac{\alpha_1}{\alpha} \right) + \frac{13}{3} \log^3 \left(\frac{\alpha_1}{\alpha} \right) - \frac{3}{2} \log^2 \left(\frac{\alpha_1}{\alpha} \right) - 4 \log \left(\frac{\alpha_1}{\alpha} \right) \right) +$$

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New approximation

$$\alpha_N[\text{new approximation for N order of PT}] = \sum_{i=1}^N \alpha_i(L) + \sum_{i=1}^N \alpha_1^i \alpha_i \left[\bar{\beta}_1 \alpha_1 \log \left(\frac{\alpha_1}{\alpha} \right) \right]$$

New approximation

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An object with anomalous dimensions

$$\left(\mu^2 \frac{d}{d\mu^2} + \beta(\alpha) \frac{d}{d\alpha} - \gamma(\alpha) \right) \Gamma \left(\frac{Q^2}{\mu^2}, \alpha \right) = 0$$

$$\frac{d \log \Gamma}{d \log \mu^2} = \gamma(\alpha) \quad \text{Formal solution} \quad \log \Gamma = \int_{\alpha}^{\bar{\alpha}} \frac{\gamma(x)}{\beta(x)} dx$$

$$\gamma(\alpha) = \gamma_0 \alpha + \gamma_1 \alpha^2 + \gamma_2 \alpha^3 + \dots$$

Linearisation

$$\log \Gamma = \log \Gamma_1 + \log \Gamma_2 + \log \Gamma_3 + \dots$$

Summary

- ✓ A method has been found for obtaining analytical solutions to the renormalisation group equations in the single-coupling model
- ✓ A new approximation formula for taking into account quantum corrections in a given order of perturbation theory is proposed

Thank you for attention!