

Asymptotically safe quantum gravity: functional and lattice perspectives

ERG 2026, University of Sussex

Marc Schiffer, Radboud University Nijmegen

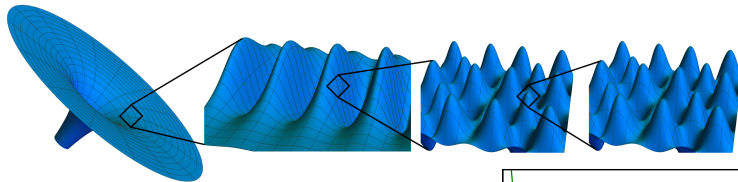
September 2, 2026

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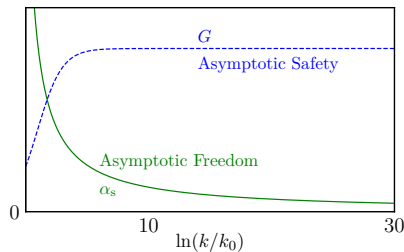


Asymptotically Safe Quantum Gravity

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- Perturbative quantum gravity:
loss of predictivity
- Key idea of asymptotic safety:
Quantum realization of
scale symmetry
 - ▶ imposes infinitely many conditions on
theory



Asymptotic freedom:

free fixed point, $\alpha_{s,*} = 0$

Asymptotic safety:

interacting fixed point, $G_* \neq 0$

Motivation: Exploring AS with functional and lattice methods

Perturbative methods: see talk by Y. Kluth

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- Lattice methods:

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 - ▶ extract scale dependence of couplings/operators
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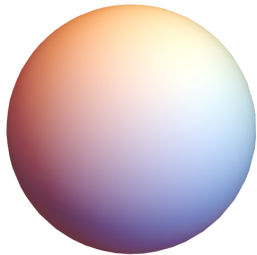
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Use different methods in concerted fashion to extract evidence for and physical features of asymptotically safe quantum gravity.

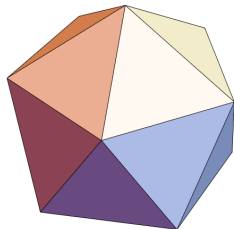
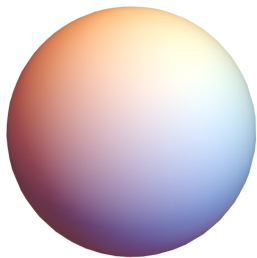
Lattice quantum gravity via dynamical triangulations

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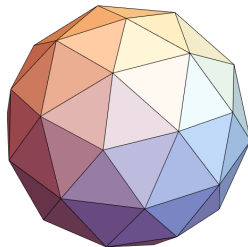
- Discretization of spacetime in terms of triangulations

Lattice quantum gravity via dynamical triangulations



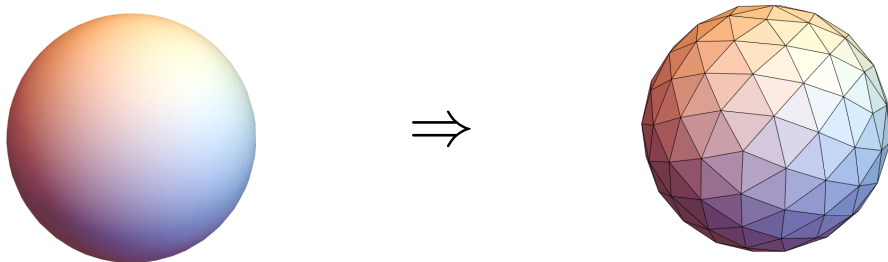
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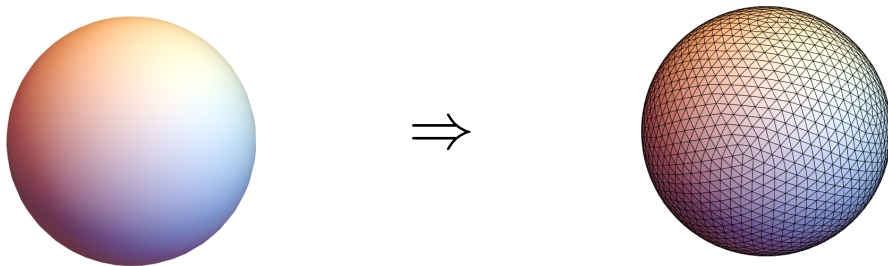
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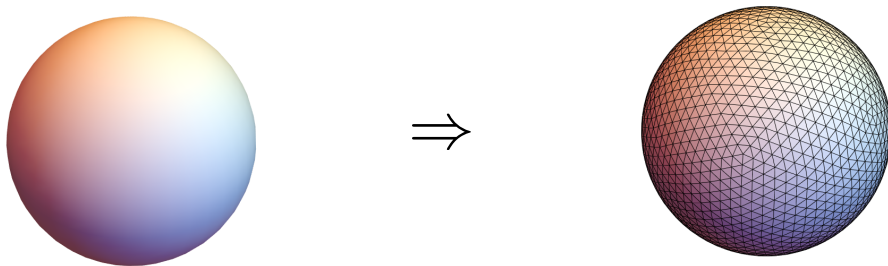
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- Discretization of spacetime in terms of triangulations

[Ambjørn and Jurkiewicz, 1992], [Agishtein and Migdal, 1992], . . .

$$\int \mathcal{D}g e^{-S[g]} \rightarrow \sum_{\mathcal{T}} \frac{1}{C_{\mathcal{T}}} e^{-S_{\text{ER}}}$$

with Euclidean Einstein-Regge action $S_{\text{ER}} = -\kappa_{d-2}N_{d-2} + \kappa_d N_d$ [Regge, 1961]

The Einstein-Regge action

- starting point: Einstein-Hilbert action

$$S_{\text{EH}} = -\frac{1}{16\pi G_{\text{N}}} \int d^4x \sqrt{g} (R - 2\Lambda)$$

The Einstein-Regge action

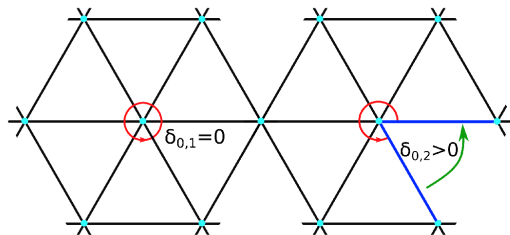
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- Discretize volume term:

$$\int d^4x \sqrt{g} \rightarrow \sum_{s_4} V_{s_4}$$

The Einstein-Regge action

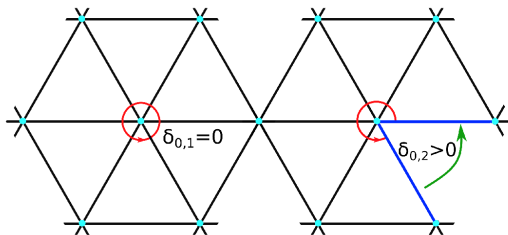


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- Curvature on the lattice: deficit angle δ_{s_2}

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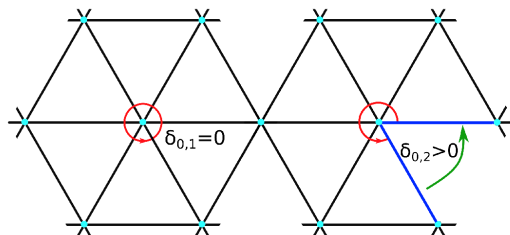
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$$\int d^4x \sqrt{g} R \rightarrow \sum_{s_2} V_{s_2} \delta_{s_2}$$

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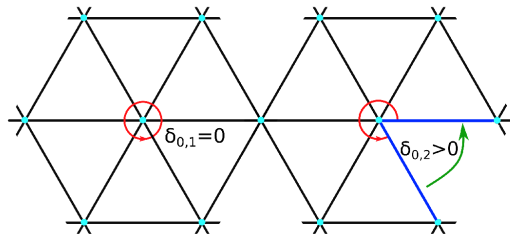
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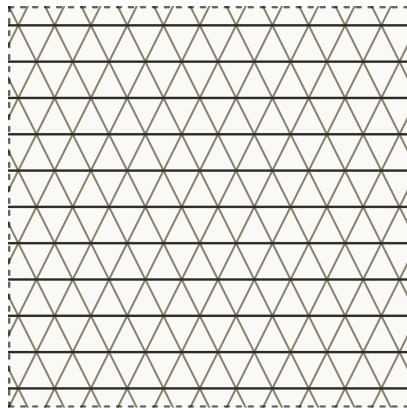
$$S_{\text{ER}} = -\frac{1}{8\pi G_{\text{N}}} \left(\sum_{s_2} V_{s_2} \delta_{s_2} - \Lambda \sum_{s_4} V_{s_4} \right) = \kappa_4 N_4 - \kappa_2 N_2$$

Dynamical Triangulations

- Discretization of spacetime in terms of triangulations

$$\int \mathcal{D}g e^{-S[g]} \rightarrow \sum_{\mathcal{T}} \frac{1}{C_{\mathcal{T}}} e^{-S_{\text{ER}}}$$

- Perform sum via Monte Carlo methods



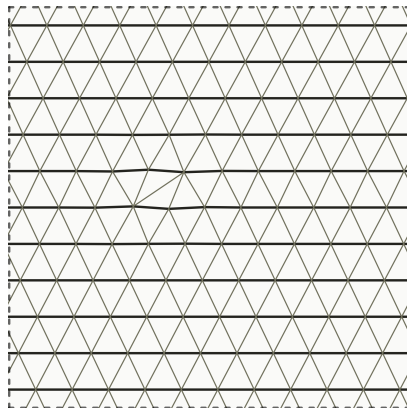
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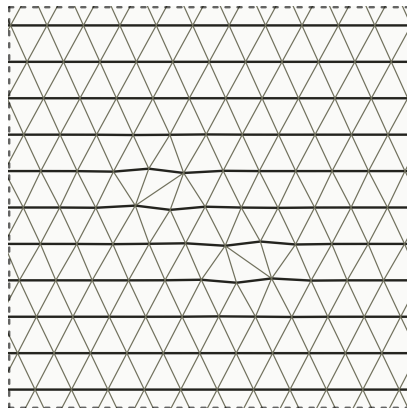
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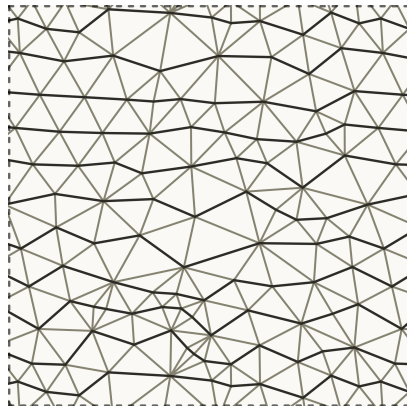
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- Perform sum via Monte Carlo methods
 - ▶ Pachner moves: change volume
 - ▶ Tune κ_4 automatically

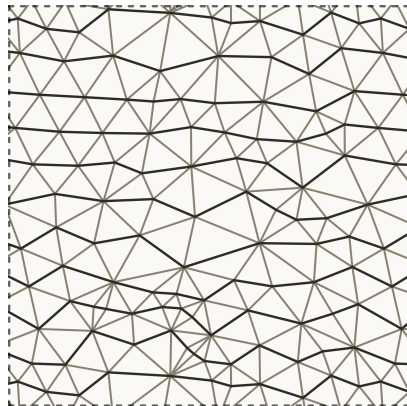


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- Perform sum via Monte Carlo methods
 - ▶ Pachner moves: change volume
 - ▶ Tune κ_4 automatically
 - ▶ Search for continuous phase transition



[Credits: J. van der Duin]

Lattice quantum gravity in $d = 4$

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- in $d = 4$: no physical phase, no indications for higher-order transition in κ_2 - κ_4 - space

[Ambjørn, Jain, Jurkiewicz, Kristjansen, 1993], [Bakker, Smit, 1994]

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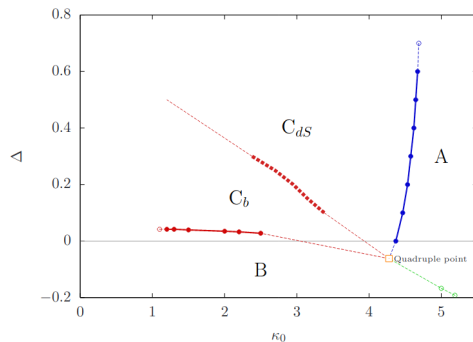
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- CDT: impose causal structure

[Ambjørn, Loll, 1998], [Ambjørn, Jurkiewicz, Loll, 2000], . . .

RG on the lattice: [Ambjørn, Gizbert-Studnicki, Goerlich, Nemeth; 2024]



Taken from [Loll, 2020]

Lattice quantum gravity in $d = 4$

- Discretization of spacetime in terms of triangulations

$$\int \mathcal{D}g e^{-S[g]} \rightarrow \sum_{\mathcal{T}} \frac{1}{C_{\mathcal{T}}} \left[\prod_{j=1}^{N_2} \mathcal{O}(t_j)^{\beta} \right] e^{-S_{\text{ER}}}$$

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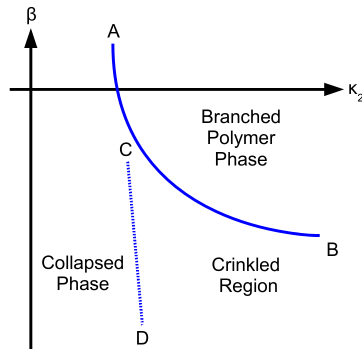
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- EDT: include additional operator

[Bruegmann, Marinari, 1993], [Bilke, Burda, Krzywicki, Petersson, Tabaczek, Thorleifsson, 1998], [Laiho, Coumbe, 2011], [Ambjørn, Glaser Goerlich, Jurkiewicz, 2013], . . .

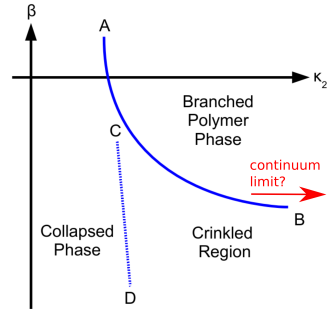


[Ambjørn, Glaser, Goerlich, Jurkiewicz, 2013]

[Coumbe, Laiho, 2014]

Phase diagram of EDT

- Idea:
follow AB-line towards large κ_2 and
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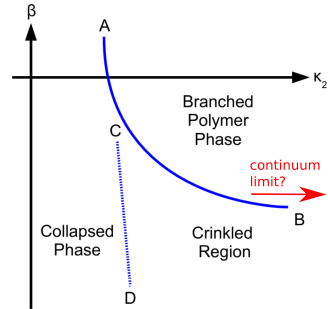


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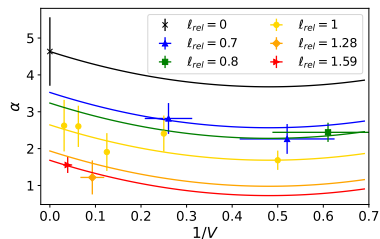
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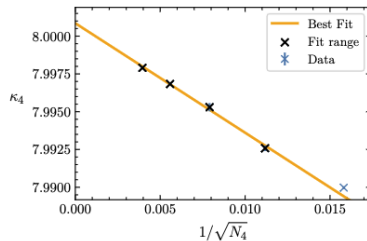
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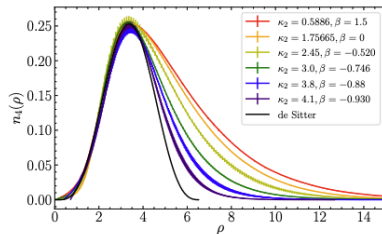
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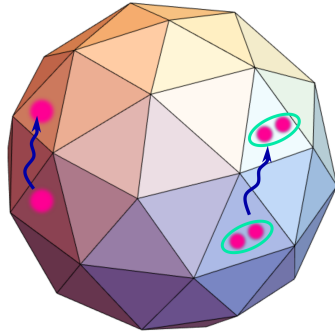
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 - ▶ Is there a continuum limit?
 - ▶ Decreasing lattice spacings?

κ_2	G_a	a_{rel}	ℓ_{rel}	$\chi^2/\text{d.o.f}$	p-value
2.45	0.477(27)(57)	1.125(92)	1.45(12)	0.61	0.61
3.0	0.604(38)(35)	1.00	1.00	0.47	0.71
3.4	0.805(81)(50)	0.866(67)	0.618(49)	0.39	0.76
3.8	0.98(23)(06)	0.78(10)	0.386(51)	1.55	0.20

Newtonian binding energy from EDT

- Use matter as probe for gravity



Newtonian binding energy from EDT

- Use matter as probe for gravity
- Study scalar fields on EDT
(quenched approximation)

$$S_{\text{M}}^{\text{lattice}} = \sum_{x,y} \phi_x L_{xy} \phi_y ,$$

$$\text{with } L_{xy} = (D_x + m_0^2)\delta_{xy} - A_{xy} .$$

Newtonian binding energy from EDT

- Use matter as probe for gravity
- Study scalar fields on EDT
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- Extract binding energy E_b and renormalized mass m from one- and two-particle two-point correlators

[de Bakker, Smit, 1996]

$$G_s(r) \sim L^{-1},$$
$$\propto \frac{e^{-m r}}{r^p} \quad (\text{for large } r).$$

$$G_s^{(2)}(r) \sim (L^{-1})^2,$$
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[\[de Bakker, Smit, 1996\]](#)
- Fit $E_b = A m^\alpha$
- Continuum, non-relativistic limit:
 $E_b = \frac{G m^5}{4}$ (in $d = 4$)

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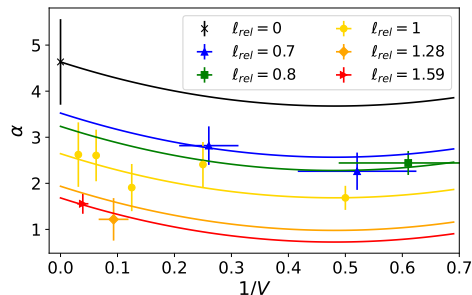
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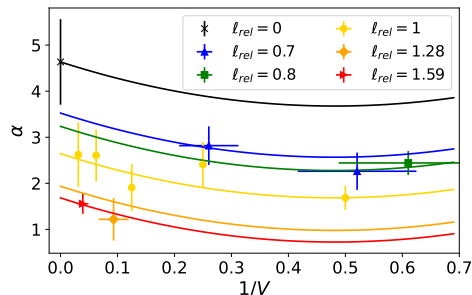
- EDT fit: $\alpha = 4.6 \pm 0.9$
 $\Rightarrow d = 3.9 \pm 0.2$

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[Dai, Laiho, MS, Unmuth-Yockey; 2021]

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EDT might feature well-behaved non-relativistic limit! [Dai, Laiho, MS, Unmuth-Yockey; 2021]

The Saddle point approximation in EDT II

- Saddle point approximation:

[Ambjorn, Goerlich, Jurkiewicz, Loll, 2012]

$$\langle N_4 \rangle \simeq \frac{A^2}{4(\kappa_4 - \kappa_4^c)^2} \Rightarrow A \sim |\kappa_4 - \kappa_4^c| \sqrt{N_4}.$$

Use finite-volume scaling of κ_4 to test recovery of semi-classical limit

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Use finite-volume scaling of κ_4 to test recovery of semi-classical limit

- Match lattice saddle point approximation with continuum calculation:

$$Z(\kappa_2, \kappa_4) \approx \exp\left(\frac{A^2(\kappa_2)}{4(\kappa_4 - \kappa_4^c)}\right)$$

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Use finite-volume scaling of κ_4 to test recovery of semi-classical limit

- Match lattice saddle point approximation with continuum calculation:

$$Z(\kappa_2, \kappa_4) \approx \exp\left(\frac{A^2(\kappa_2)}{4(\kappa_4 - \kappa_4^c)}\right) = \exp\left(\frac{3\pi}{G\Lambda}\right)$$

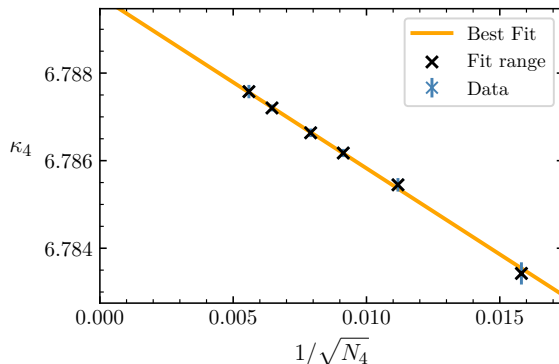
Assumption: Continuum is dominated by de Sitter instanton [Hawking, Moss, 1987]

Extract G from lattice data:

$$\frac{G}{\ell_{\text{fid}}^2} \sim \left(\frac{a}{\ell}\right)^2 \frac{\ell_{\text{rel}}^2}{|s|}, \quad s : \text{slope of } \kappa_4(N_4),$$

Numerical result: finite volume scaling

Example at $\beta = -0.6$, $\kappa_2 = 2.245$:



Extract slope s : $\kappa_4(N_4) = A + s \frac{1}{\sqrt{N_4}}$

Fit result: $s = -0.393 \pm 0.022$

Recovery of semi-classical scaling!

[Bassler, Laiho, MS, Unmuth-Yockey; 2021]

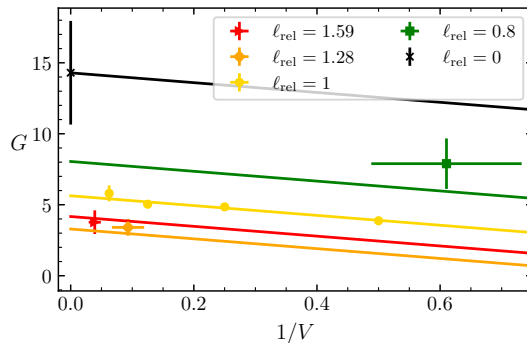
Numerical result: the Newton coupling

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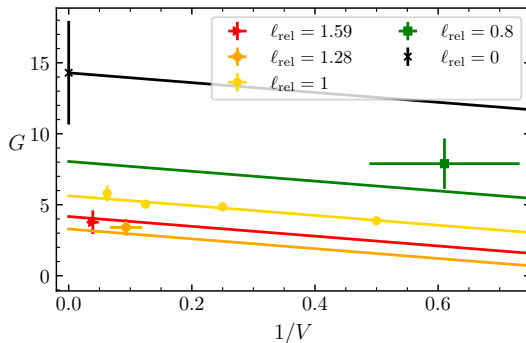
$$G = \frac{H_G}{V} + I_G \ell_{\text{rel}}^2 + J_G \ell_{\text{rel}}^4 + K_G,$$



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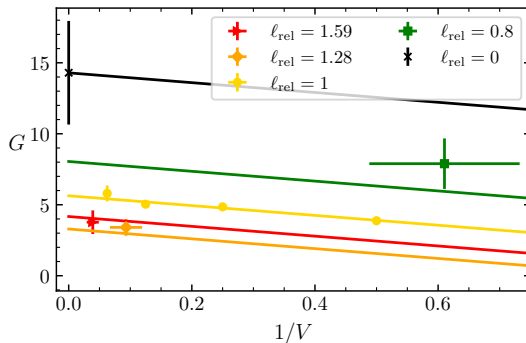
$$\Rightarrow \ell_{\text{fid}} \sim 0.26 \ell_{\text{Pl}}$$

[Bassler, Laiho, MS, Unmuth-Yockey; 2021]

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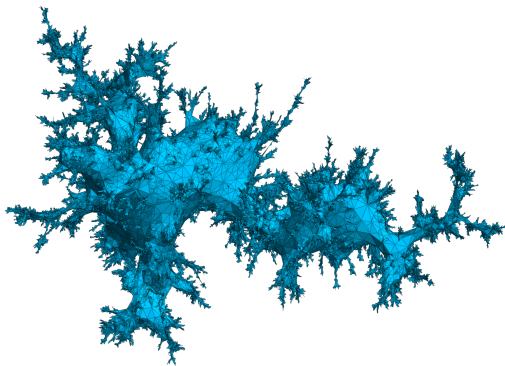
[Bassler, Laiho, MS, Unmuth-Yockey; 2021]

To date: $\ell_{\text{SF}} \sim 0.057 \ell_{\text{Pl}}$; $a_{\text{SF}} \sim 1.3 \ell_{\text{Pl}}$

[Dai, Freeman, Laiho, MS, Unmuth-Yockey; 2025]

Towards de Sitter geometries in EDT

- Shelling function $n_4(\rho)$: counts number of four-simplices at geodesic distance τ away from source-simplex.

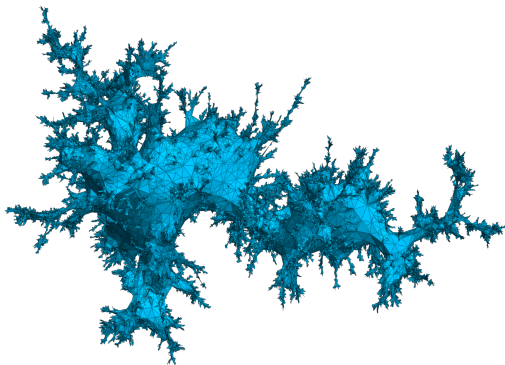


Taken from [van der Duin, Loll, MS, Silva; 2025]

Towards de Sitter geometries in EDT

- Shelling function $n_4(\rho)$: counts number of four-simplices at geodesic distance τ away from source-simplex.
- For de Sitter with $d_H = 4$:

$$n_4(\rho) \sim \cos^3(\rho)$$



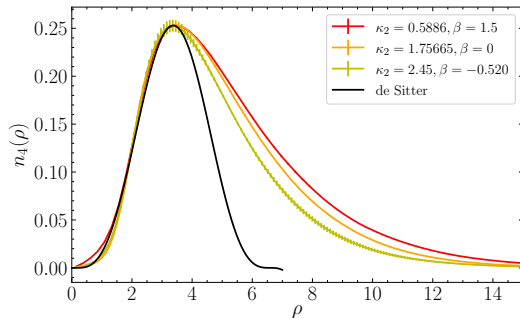
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- Peak-height: order parameter of AB-transition

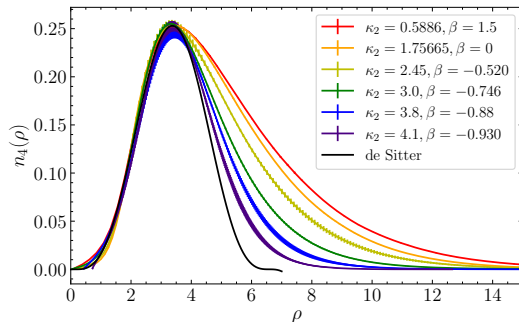


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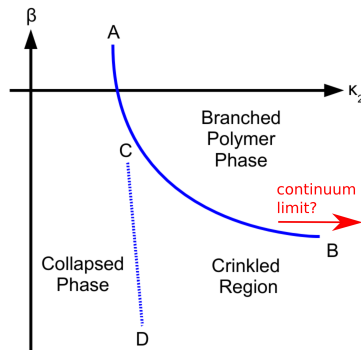


[Dai, Freeman, Laiho, MS, Unmuth-Yockey; 2023, 2024]

Lattice volume profiles: approximate de Sitter profile better for larger κ_2

Phase diagram of EDT

- Idea:
follow AB-line towards large κ_2 and negative β
- Key open questions:
 - ▶ Is there a physical phase?
 - ▶ Four-dimensional geometries
 - ▶ Semi-classical limit
 - ▶ Physical volume profiles
 - ▶ Is there a continuum limit?
 - ▶ Decreasing lattice spacings



[Ambjorn, Glaser, Goerlich, Jurkiewicz, 2013]
[Coulme, Laiho, 2014]

(Tentative) evidence for continuum limit and physical phase in $4d$ EDT.

Motivation: Exploring AS with functional and lattice methods

- Functional methods:

- ▶ extract scale dependence of couplings/operators
- ▶ AS: $k\partial_k\vec{g} = 0$
- ▶ semi-analytical treatment, inclusion of matter, connection to low-energy physics, ...
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- Lattice methods:

- ▶ extract lattice correlators, order parameters, ...
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Concrete example: critical exponent of scalar mass

Θ_{m^2} from CDT

Based on [van der Duin, Loll, MS; WIP]

- Use quenched approximation
- Invert L to compute propagator

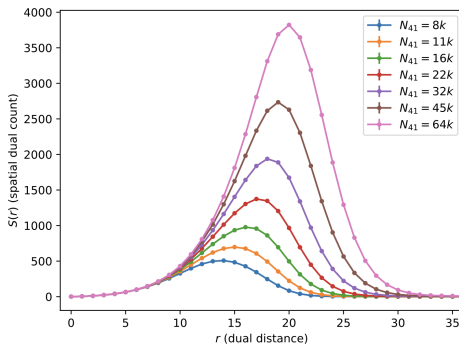
$$S_M^{\text{lattice}} = \sum_{x,y} \phi_x L_{xy} \phi_y ,$$

$$\text{with } L_{xy} = (D_x + m_0^2)\delta_{xy} - A_{xy} .$$

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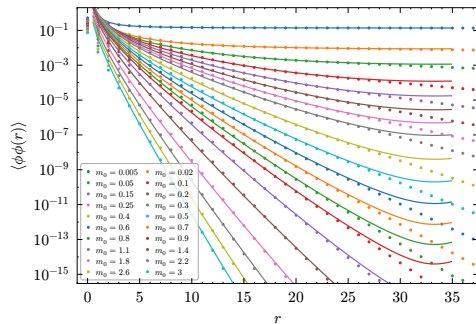
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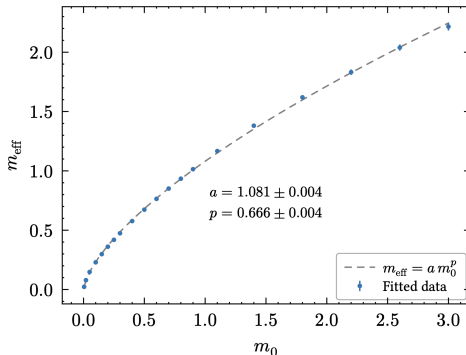


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Fit radius from volume profiles
- Extract Θ_{m^2} from (volume)
scaling of $m_{\text{eff}}(m_0)$

[Jha, Laiho, Unmuth-Yockey; 2018]



Based on [Gyftopoulos, Portas-Chiesa, Reichert, MS; WIP]

- Θ_{m^2} computed extensively
see e.g., [Narain, Percacci; 2009], [Narain, Rahmede; 2010], [Percacci, Vacca; 2015], [Labus, Percacci, Vacca; 2015], [Pawlowski, Reichert, Wetterich, Yamada; 2018], [Wetterich, Yamada; 2019], [Pastor-Gutierrez, Pawlowski, Reichert; 2023], . . .
- Here: take all operators into account, which contribute to Θ_{m^2} at $\mathcal{O}(G^2)$ (NLO).
See also [de Brito, Reichert, MS; 2025]

Θ_{m^2} from FRG

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- Here: take all operators into account, which contribute to Θ_{m^2} at $\mathcal{O}(G^2)$ (NLO).
See also [de Brito, Reichert, MS; 2025]
- These include: $R\phi^2$, $R^2\phi^2$, $C^2\phi^2$, $(\phi\Box\phi)\phi^2$, $F^2\phi^2$, $(\bar{\psi}\not{D}\psi)\phi^2$
- Results (preliminary):
 $\Theta_{m^2} = 0.4$ ($N_S = 1$); $\Theta_{m^2} = 2.04$ (SM matter)
 $R\phi^2$ is irrelevant
- Ahead: stability analysis, estimate of systematic uncertainties

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Thank you for your attention!