

Not so minimal Warm Inflation

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Not so Minimal Warm Inflation

Cold inflation/Warm inflation

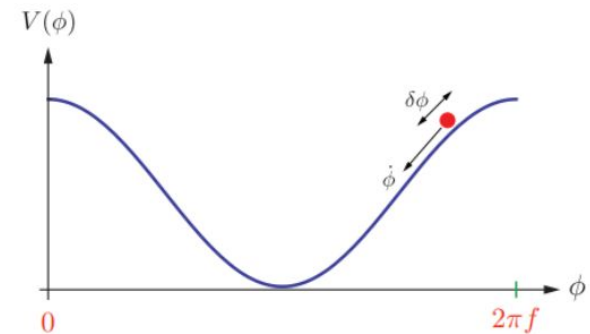
Minimal WI with natural potential

$$L = \dots - g^2 \frac{\phi}{64 \pi^2 f_a} \text{Tr} G_{\mu\nu} \tilde{G}^{\mu\nu} \quad \longrightarrow \quad Y \sim \frac{T^3}{f_a^2}$$

Model building constrains:

$$f_a \ll f_b$$

$$V = \Lambda^4 (1 - \cos \phi / f_b)$$



Expanding Universe

Flatness problem

$$\Omega_T = 1 \quad \longrightarrow \quad \Omega_T(t_{\text{nucl}}) - 1 \approx 10^{-16}$$

Horizon problem

The observable Universe was larger than the **particle horizon** at LSS

Inflation

Early period of accelerated expansion

$$\ddot{a} > 0: \quad P < -\rho/3$$

Primordial spectrum

Too small sub-horizon (**causal**) perturbations

Unwanted relics...

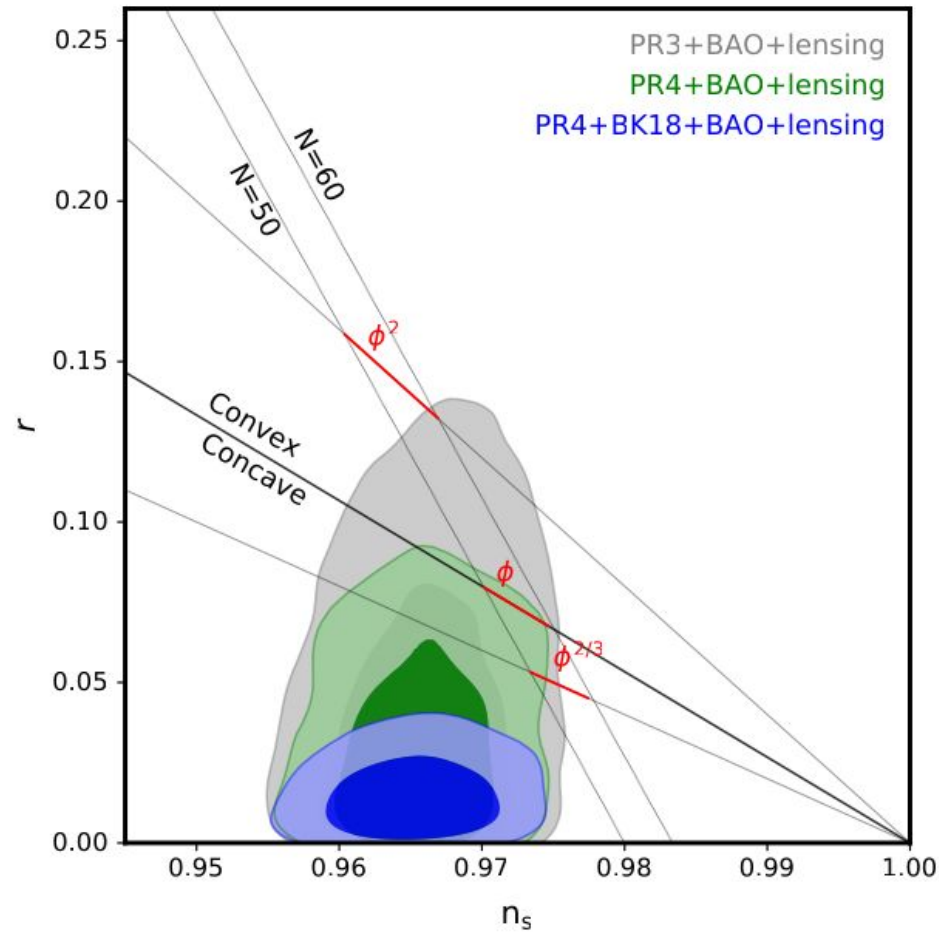
monopoles, moduli, gravitinos,...

[Starobinsky '80; Guth '81; Sato '81; Albrecht, Steinhardt '82; Linde '82]

Primordial spectrum: ~adiabatic, ~scale-invariant, gaussian?, tensors?

Primordial spectrum: $P_R = P_R(k_0)(k/k_0)^{n_s-1}$ $k_0 = 0.05 \text{ Mpc}^{-1}$

Tensor-to-scalar Ratio: $r = P_T/P_R$



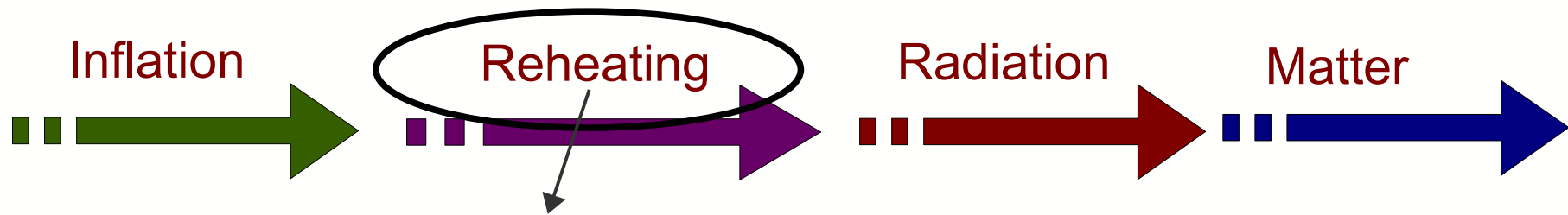
$$P_R = 2.2 \times 10^{-9}$$

$$r < 0.036$$

$$n_s = 0.9649 \pm 0.0042$$

$$\alpha_s = -0.013 \pm 0.009$$

[Tristam et al.: 2112.07961]

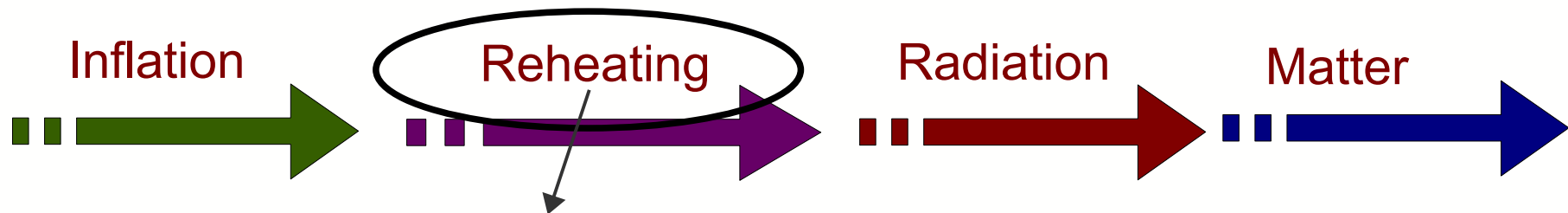


Inflaton interacts with other particles

"Cold" inflation:	Interactions negligible during Inflation	Reheating	→	Radiation
"Warm" inflation:	Inflaton "decay" into light dof	Dissipation	→	Radiation

$[T > H, \Gamma_\chi > H]$ (through a mediator)

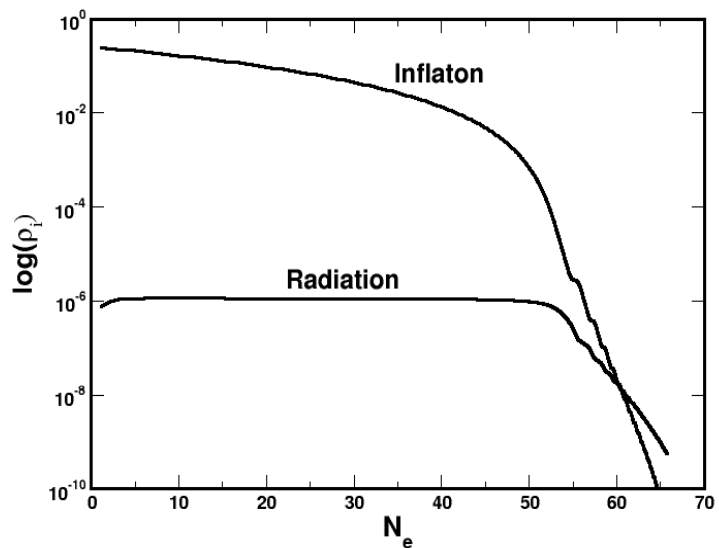
[Berera PRD55 '97; Berera, Moss & Ramos RPP72 '07]



Inflaton interacts with other particles

"Cold" inflation: Interactions negligible during Inflation $\xrightarrow{\text{Reheating}}$ Radiation

"Warm" inflation: Inflaton "decay" into light dof
 $[T > H, \Gamma_\chi > H]$ (through a mediator) $\xrightarrow{\text{Dissipation}}$ Radiation



$\swarrow$ extra friction

$$\ddot{\phi} + (3H + Y)\dot{\phi} + V_\phi = 0$$

A (small) fraction of the vacuum energy is converted into radiation during inflation

$$\dot{\rho}_r + 4H\rho_r = Y\dot{\phi}^2$$

Extra friction term: $Y(T, \phi)$, $Q = Y/3H$

Slow-roll : $3H(1+Q)\dot{\phi} \simeq -V_{\phi}(\phi, T)$, $\rho_r \simeq \frac{3}{4}Q\dot{\phi}^2$

$|\eta_{\phi}| < (1+Q)$, $\epsilon_{\phi} < (1+Q)$, $\delta_T = T V_{T\phi}/V_{\phi} < 1$  Thermal corrections

•“Heavy” ($m > T$) mediator: [low-T regime] $Y \simeq C_{\phi} \frac{T^3}{\phi^2}$

Thermal corrections under control (inflaton coupled to heavy fields)

[BG, Berera, Ramos & Rosa 2012; BG, Berera & Kronberg 2015; R. Ayra et al. 2018]

•“Light” ($m < T$) mediator: [high-T regime] $Y \simeq C_{\phi} T$, C_{ϕ}/T

Thermal corrections may spoil inflation! $\Delta V_T = -\frac{\pi^2}{90} g_R T^4 + \frac{g^2 \phi^2}{12} T^2 + \dots$

[Berera, Gleiser & Ramos PRD'98; Yokoyama & Linde PRD '98]

Solution: use symmetries to protect the inflaton mass

- “Little Warm Inflation” :

Inflaton a PNCB of a broken U(1) symmetry + pair of fermions/scalars + exchange symmetry

$$\Delta V_T = -\frac{\pi^2}{90} g_R T^4 + \underbrace{\frac{g^2 M^2}{12} T^2 + \frac{g^4(\varphi) M^4}{16 \pi^2} \left(\log \frac{\mu^2}{T^2} - c_f\right)}_{\text{Fermion (scalars) masses are bounded}}$$

Fermion (scalars) masses are bounded $\Rightarrow$ No thermal mass for the inflaton

- “Minimal Warm Inflation” :

[Berghaus, Graham, Kaplan 1910.07525]

Axion-like inflaton+ gauge production [SU(N)]

$$\mathcal{L} = \dots - g^2 \frac{\phi}{64 \pi^2 f_a} \text{Tr} G_{\mu\nu} \tilde{G}^{\mu\nu}$$

[Mishra, Mohanty & Nautiyal, 1106.3039; Visinelli 1107.3523]

[Ferrerira & Notari 1711.07483; Kamali 1901.01897]

- Axion mass protected by shift symmetry at perturbative level; non-perturbative corrections axion mass $\sim$ negligible

- Dissipative coefficient $\Upsilon \sim \frac{\Gamma_{\text{sph}}}{2 T f_a^2}$

Sphaleron rate $\Gamma_{\text{sph}} \simeq N_c^5 \alpha^5 T^4$

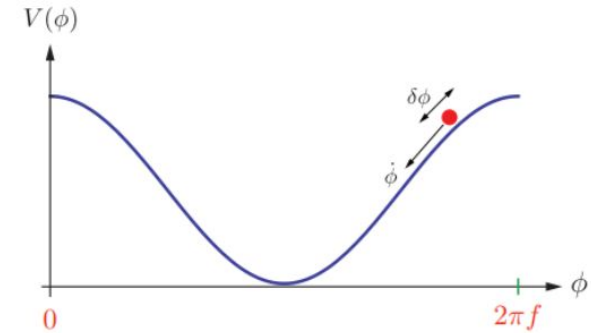
[Kuzmin, Rubakov & shaposhnikov, PLB '85]

- MWI + hybrid inflation : SDR ($Q \gg 1$) at horizon crossing

Our setup

- ◆ Natural potential: $V = \Lambda^4 \left(1 - \cos \frac{\phi}{f_b}\right)$

[Freese, Frieman & Olinto, PRL '90]



- ◆ Sphalerons induce (constant) dissipation even with $T=0$ (vacuum decay $Y \sim \text{Constant}$)

$$T=0 \rightarrow T > T_c \rightarrow T > H$$

[Laine & Procacci 2102.09913]

$$Y \simeq \frac{d_A \alpha^2}{f_a^2} \left(\underbrace{\kappa (\alpha N_c T)^3 F[\omega/T]}_{\text{Thermal}} + \underbrace{(1 + n_B(\omega)) \frac{\pi \omega^3}{(4\pi)^4}}_{\text{Vacuum}} \right) \quad [\omega \simeq m]$$

Thermal $[\omega \simeq m < \alpha^2 T]$

Vacuum $[\omega \simeq m \gg \alpha^2 T]$

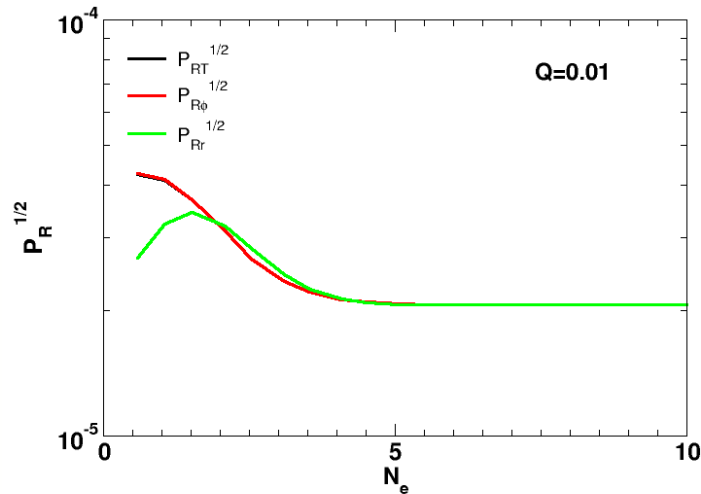
$$Y_{\text{IR}} \sim \frac{\alpha^5 T^3}{f_a^2}$$

$$Y_{\text{UV}} \sim \frac{\alpha^2 \omega^3}{f_a^2} \quad [\omega \simeq m]$$

$$F[\omega/T] \simeq \frac{1 + \frac{\omega^2}{(c_{\text{IR}} \alpha^2 N_c^2 T)^2}}{1 + \frac{\omega^2}{(c_{\text{M}} \alpha N_c T)^2}}, \quad d_A = N_c^2 - 1, \quad \kappa \simeq 1.5, \quad c_{\text{IR}} \simeq 106, \quad c_{\text{M}} \simeq 5.1$$

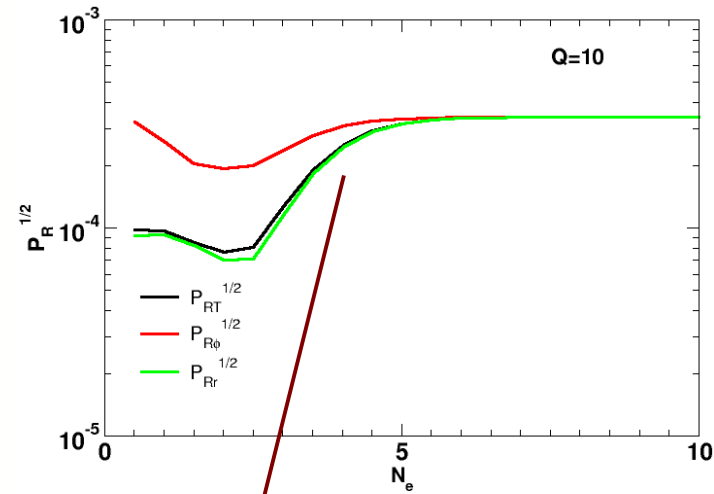
Fluctuations & primordial spectrum: coupled system

• Weak dissipative regime ($Q=Y/H \ll 1$)



$$G[Q] \simeq 1$$

• Strong dissipative regime ($Q=Y/H > 1$)



Growing mode!

Numerical integration

$$P_R \simeq \left(\frac{H^2}{2\pi\dot{\phi}} \right)^2 \left(1 + 2N + \frac{T}{H} \frac{2\pi Q}{\sqrt{1+4\pi Q/3}} \right) G[Q]$$

$$G[Q] \sim Q^\alpha$$

Dissipative processes may maintain a non-trivial distribution of inflaton particles: $N \simeq n_{BE} = (e^{k/aT} - 1)^{-1}$

[Ramos, da Silva, 1302.3544]

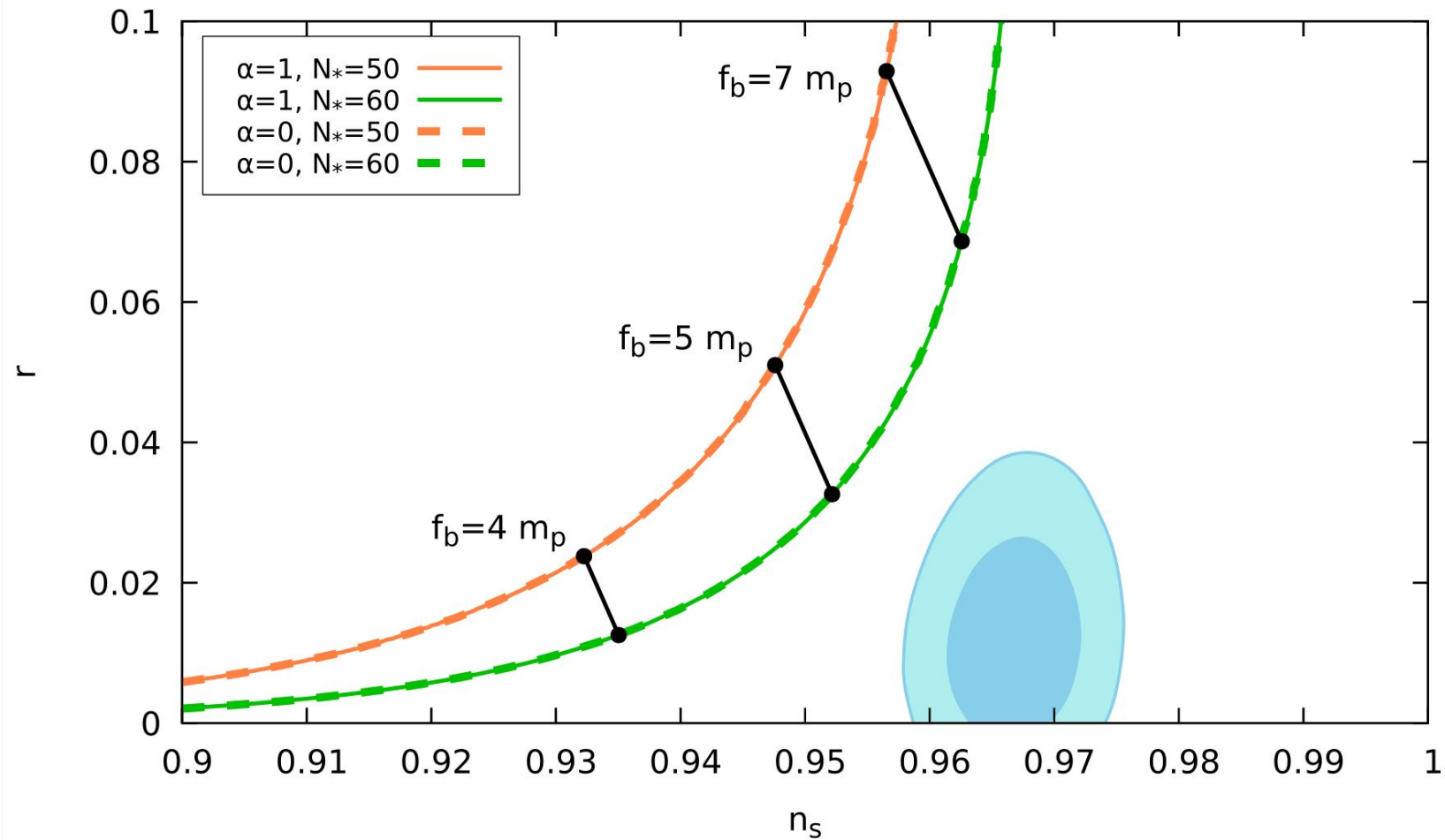
[BG, Berera, Moss & Ramos, 1401.1149]

[WI2easy: Rodrigues, Ramos 2504.17760]

[SWIM: Kumar & Das 2604.24654]

◆ Natural potential+ dissipation: $f_a = f_b$

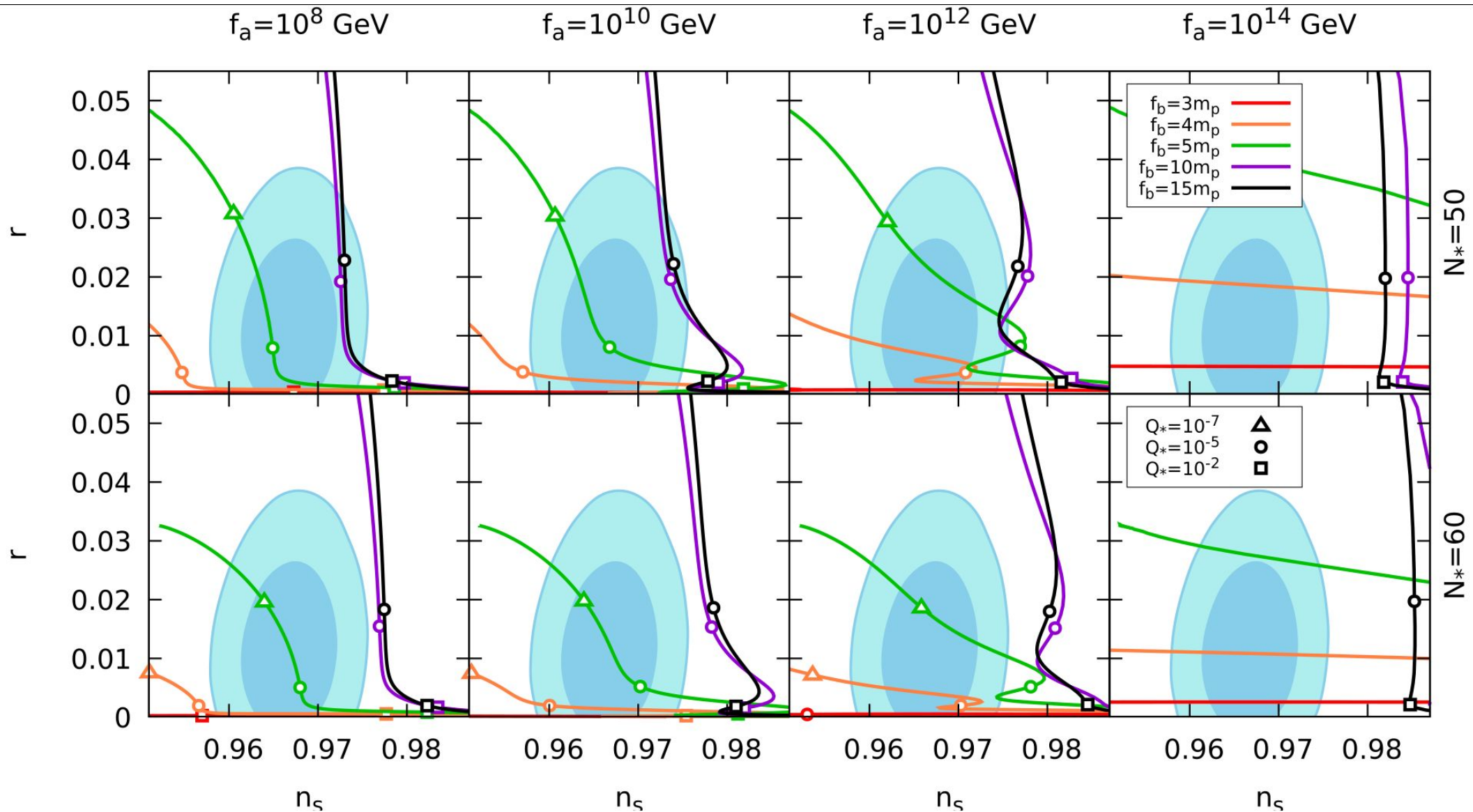
Dissipation suppressed during inflation, excluded by observations



How large the hierarchy between scales?

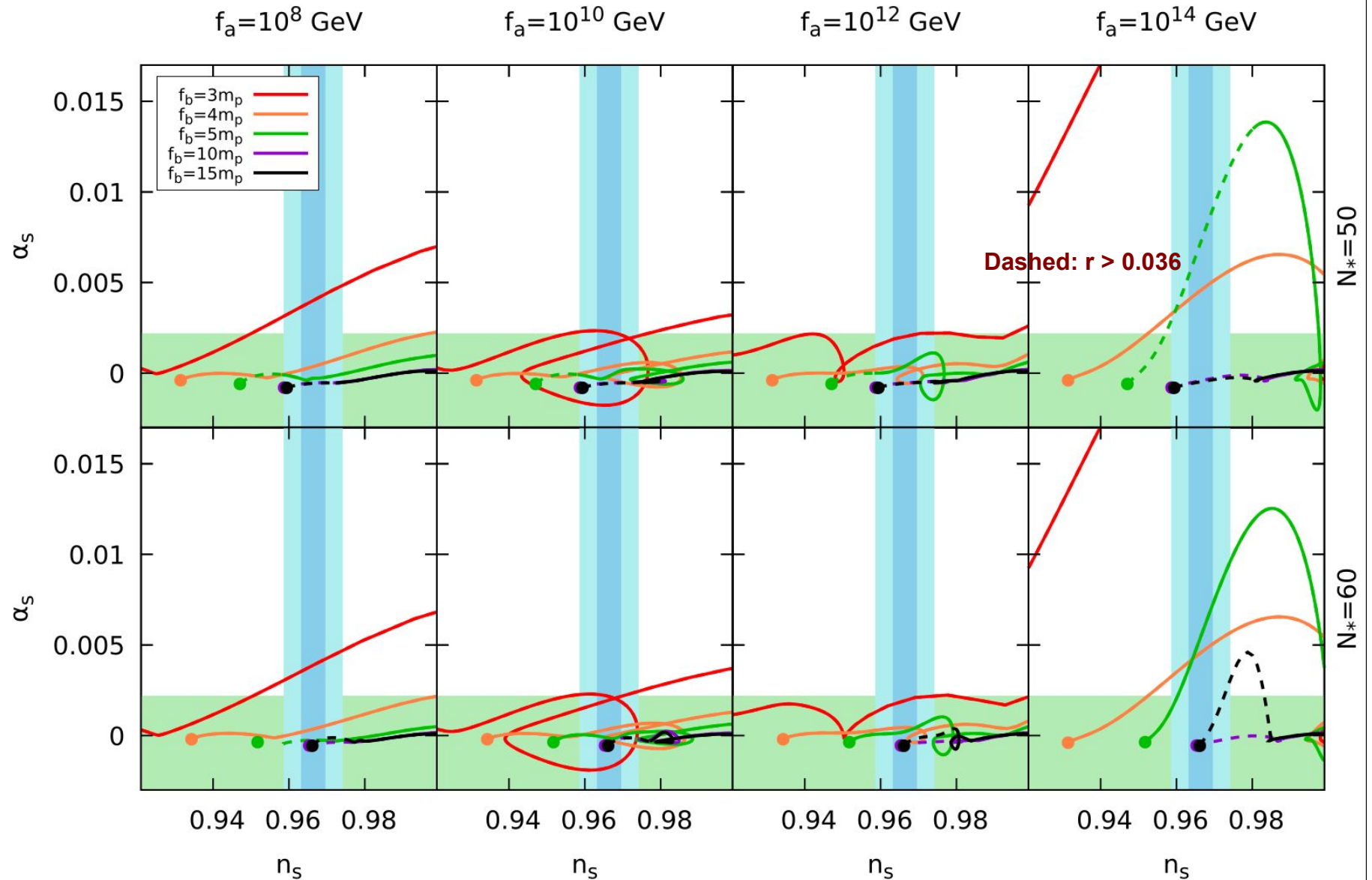
- ◆ Natural potential+ dissipation: $f_a \ll f_b$, weak dissipation, $f_b > m_P$

$$r \simeq 16\epsilon_H / (1+n_{BE}) < 16\epsilon_H$$



Running of the spectral index

- ◆ Natural potential+ dissipation: $f_a \ll f_b$, weak dissipation, $f_b \gtrsim 3m_P$, $f_b/f_a > 10^6$



How large the hierarchy between scales?

♦ “Clockwork” : $V = \Lambda^4 \left(1 - \cos \frac{\phi}{j f_a} \right)$, $j < f_b / \Lambda$, $\Lambda \leq f_a$

[Choi & Im 1511.00132; Kaplan & Rattazzi 1511.01827]

♦ Partial waves unitarity bounds (scattering of gauge fields and inflaton) :

$$\Lambda < 32 \pi^{3/2} \frac{f_a}{\alpha}$$

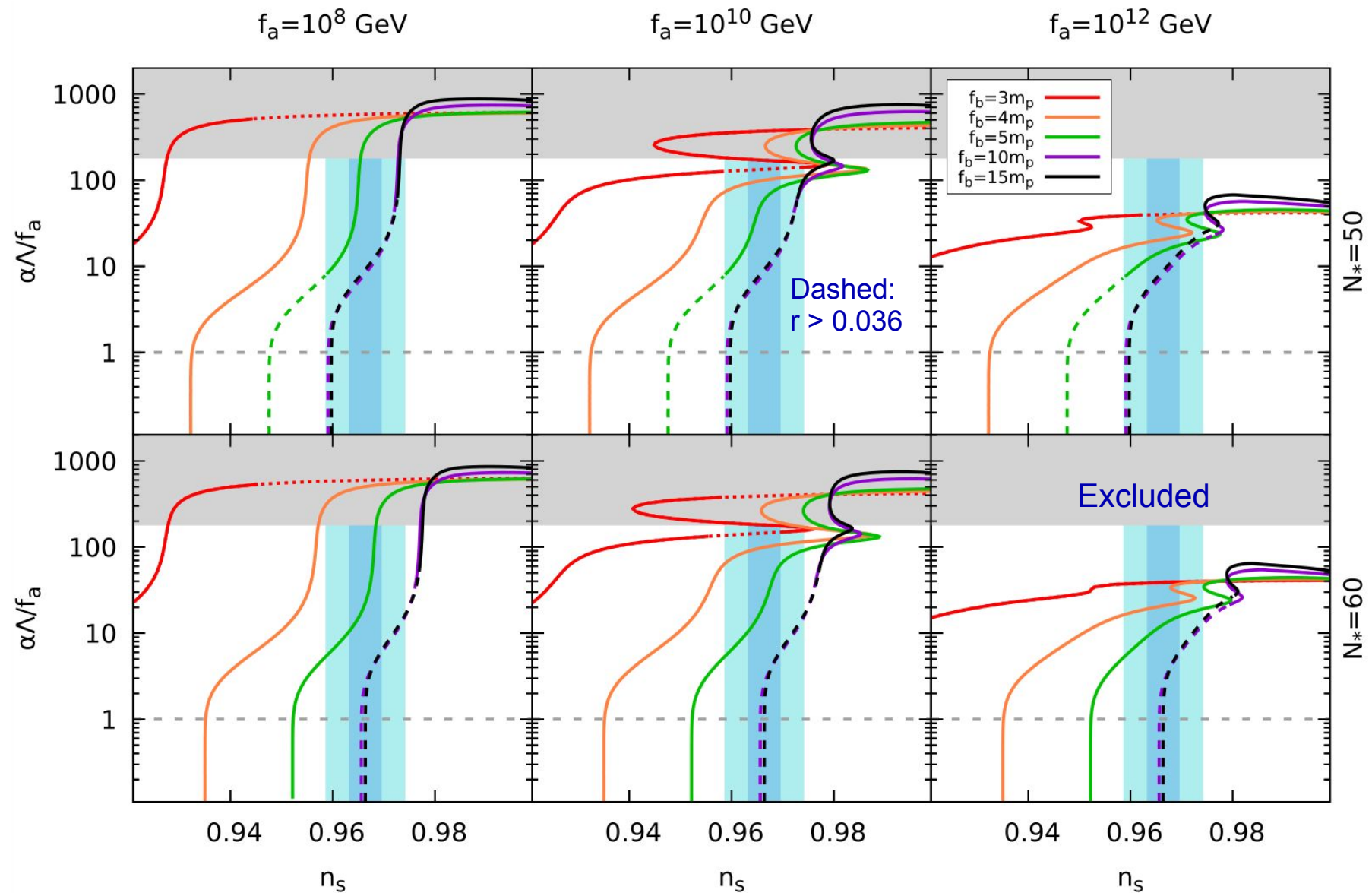
[Bagherian, Reece & Xu 2207.11262]

The value of Λ is set from the amplitude of the primordial spectrum

How large the hierarchy between scales?

◆ Clockwork : $\Lambda < f_a$ ($\alpha < 1$)

◆ Unitarity bound : $\Lambda < 32\pi^{3/2} f_a / \alpha$

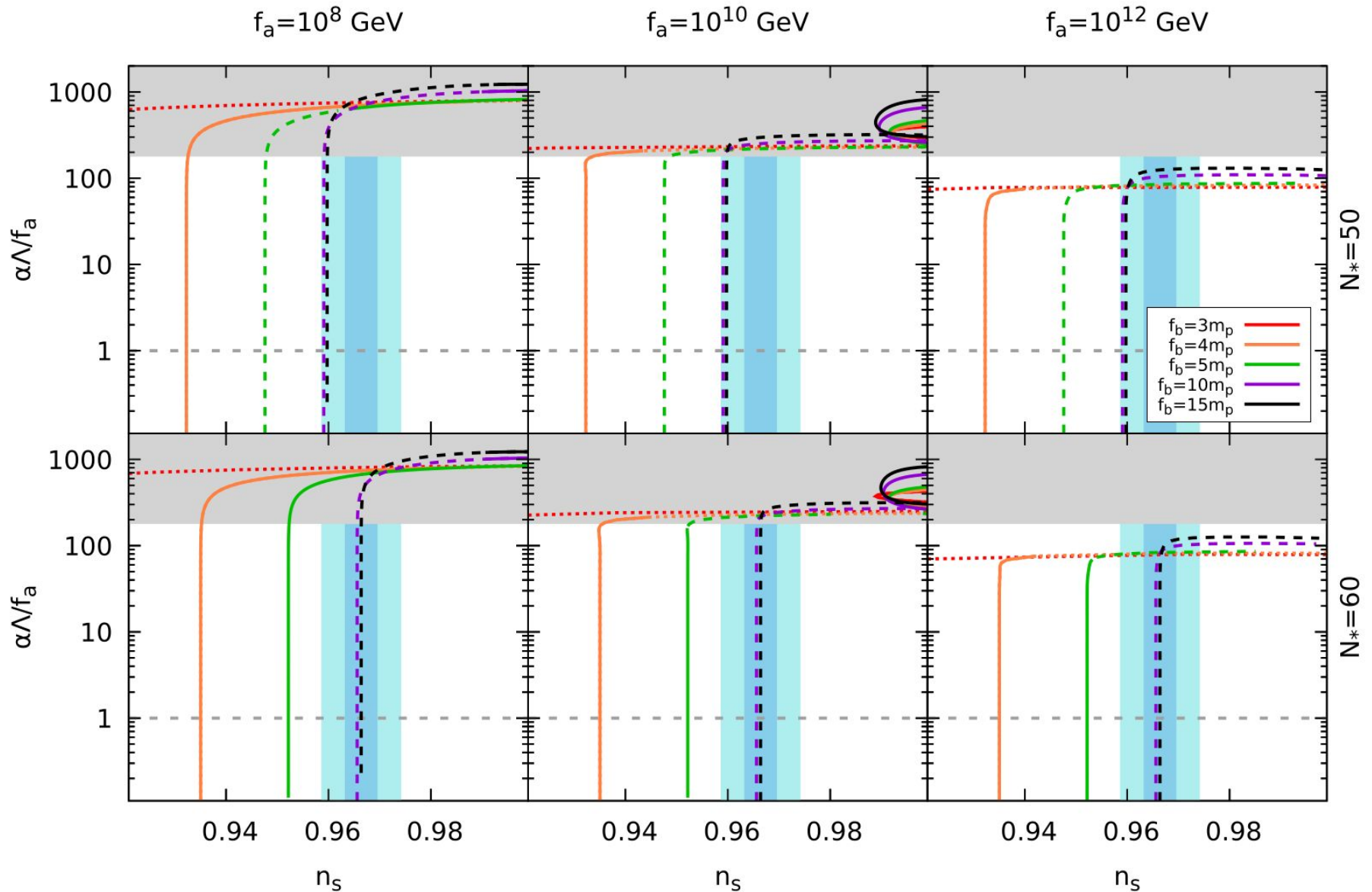


There are no values compatibles with “clockwork” mechanism to generate the hierarchy

Non-thermal inflaton

◆ Clockwork : $\Lambda < f_a$ ($\alpha < 1$)

◆ Unitarity bound : $\Lambda < 32\pi^{3/2}f_a/\alpha$



No values compatibles with “clockwork” & unitarity constraints

Summary

- Dissipative effects due to decaying fields can be relevant during inflation, and modify the inflationary predictions

- Thermal corrections to inflation potential under control with symmetries:

LWI: Inflaton a PNGB of a broken U(1) symmetry + pair of fermions/scalars + exchange sym.

MWI: Inflaton a PNGB + shift symmetry + gauge field production (non perturbative)

- Case of study: natural inflation + gauge dissipative coefficient

$$V = \Lambda^4 (1 - \cos \phi / f_b) \quad Y \sim \frac{T^3}{f_a^2} F_3[\omega] + F_0[\omega]$$

- Observations (CMB spectrum): only compatible in the WDR ($Q \ll 1$)

Large hierarchy required $f_b \sim O(m_P) \gg f_a$

Clockwork mechanism not enough to explain the hierarchy

Compatible with unitarity bounds, but only with a thermal inflaton

Thank you!