

Beyond the Starobinsky model after ACT

Dhong Yeon Cheong Min Gi Park[†] Seong Chan Park

Department of Physics Yonsei University Seoul 03722, Republic of Korea

[†]Speaker | Published in *Phys. Lett. B* **873** (2026) 140217 | arXiv:2509.04105

Abstract

Recent ACT observations show a mild tension ($\sim 2\sigma$) between the Starobinsky R^2 inflation prediction ($n_s = 0.965$) and the observed value $n_s = 0.975 \pm 0.005$. We revisit this model by introducing a small cubic curvature correction $f(R) = R + \frac{\beta}{2}R^2 + \frac{\gamma}{3}R^3$. A tiny deformation parameter $\delta = \gamma/\beta^2 \simeq -10^{-4}$ restores consistency with ACT data and predicts a reheating temperature $T_{\text{re}} \sim 5 \times 10^9 \text{ GeV}$ for a matter-dominated reheating phase. This work demonstrates that minimal higher-order corrections can reconcile inflationary theory with the latest CMB observations.

Motivation

Cosmic inflation explains the flatness, homogeneity, and origin of primordial fluctuations in the Universe. Among many models, the **Starobinsky model** — an R^2 correction to Einstein gravity — has been highly successful in matching *Planck* data. However, recent ACT-DR6 observations indicate a slightly larger scalar spectral index, $n_s = 0.975 \pm 0.005$, compared to the Starobinsky prediction $n_s = 0.965$, leading to a mild ($\sim 2\sigma$) tension.

This motivates revisiting higher-order curvature corrections, such as an R^3 term, within the minimal $f(R)$ framework to test whether such small deformations can naturally reconcile the theory with the new CMB data while preserving its predictive power.

Model

We consider a general form of the gravitational action,

$$S = \frac{1}{2} \int d^4x \sqrt{-g} f(R), \quad f(R) = R + \frac{\beta}{2}R^2 + \frac{\gamma}{3}R^3.$$

The first term corresponds to the Einstein–Hilbert action, the quadratic term (R^2) generates Starobinsky inflation, and **the cubic term** (R^3) represents a small higher-order deformation. Defining a dimensionless parameter

$$\delta = \frac{\gamma}{\beta^2} \ll 1,$$

we treat δ as a perturbative correction controlling the deviation from the pure R^2 case.

To make contact with the scalar-field picture, we move to the Einstein frame via a conformal transformation with $f'(R)$ as the rescaling factor:

$$s = \sqrt{\frac{3}{2}} \ln f'(R), \quad V_E(s) = \frac{R f'(R) - f(R)}{2 f'(R)^2}.$$

In this form, the model becomes dynamically equivalent to a scalar–tensor theory with a canonical field s driving inflation through the potential $V_E(s)$. The small cubic term slightly tilts this potential and hence modifies (n_s, r) .

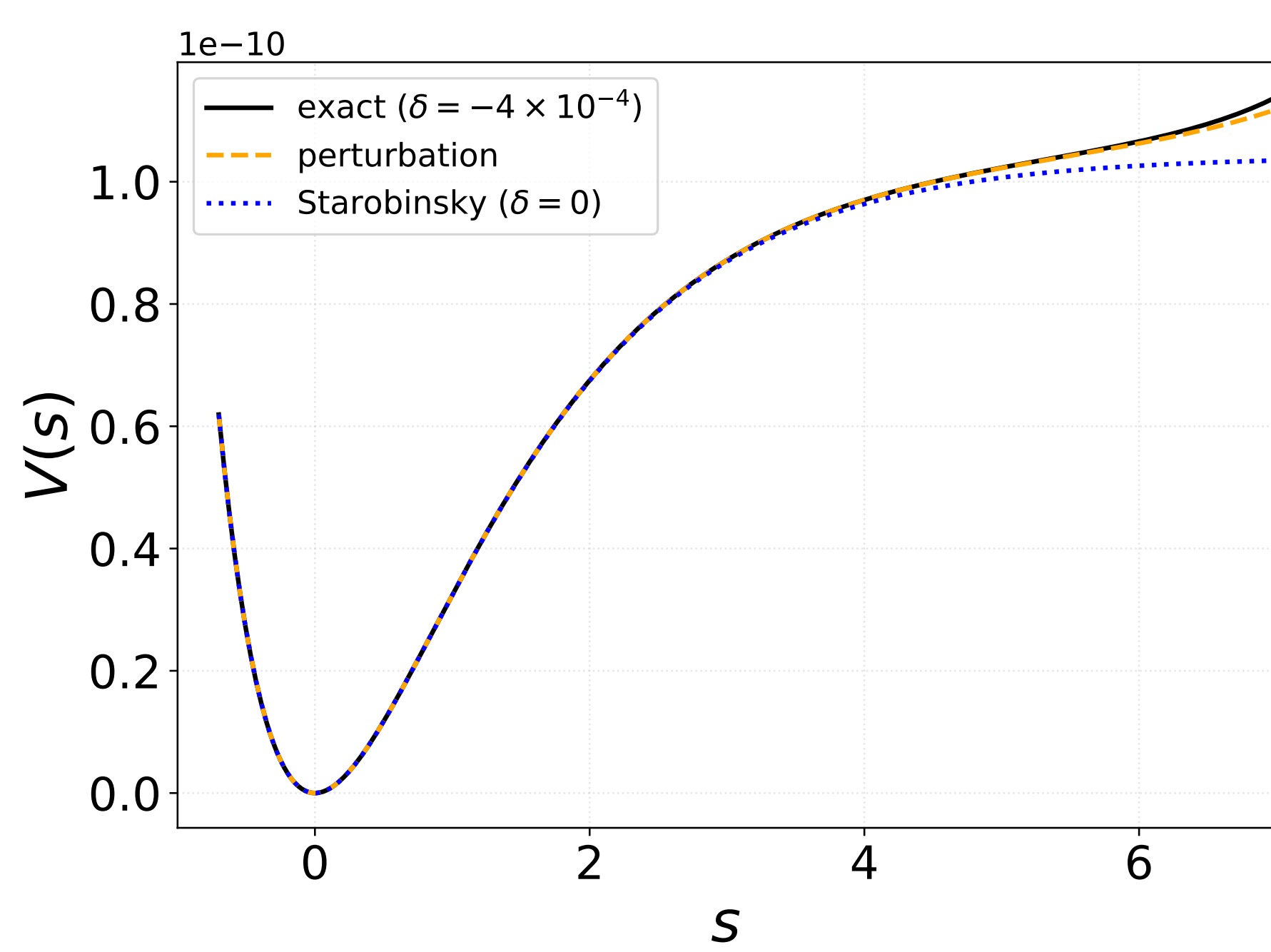


Figure: Einstein-frame potential $V_E(s)$ for the pure R^2 model (black) and the $R^2 + R^3$ case with $\delta = -10^{-4}$ (blue). The cubic correction slightly steepens the potential near the plateau, raising n_s while keeping r small.

Inflationary Observables

Inflationary predictions are characterized by the slow-roll parameters

$$\epsilon = \frac{1}{2} \left(\frac{V'_E}{V_E} \right)^2, \quad \eta = \frac{V''_E}{V_E}.$$

Expanding perturbatively in the small deformation parameter δ , the scalar spectral index n_s and tensor-to-scalar ratio r become

$$n_s \simeq 1 - \frac{2}{N'_e} - \frac{9}{2N_e'^2} - \delta \frac{128}{81} N'_e \quad (1)$$

$$r \simeq \frac{12}{N_e'^2} - \delta \frac{256}{27} \quad (2)$$

where $N'_e = N_e + \frac{3}{4} \ln(\frac{4}{3} N_e)$.

For the Starobinsky limit ($\delta = 0$), one finds $n_s \simeq 0.965$ and $r \simeq 3 \times 10^{-3}$ for $N_e \approx 55$. A small negative deformation ($\delta \simeq -10^{-4}$) slightly raises n_s to $\simeq 0.975$ without significantly changing r , bringing the prediction into excellent agreement with the ACT-DR6 data.

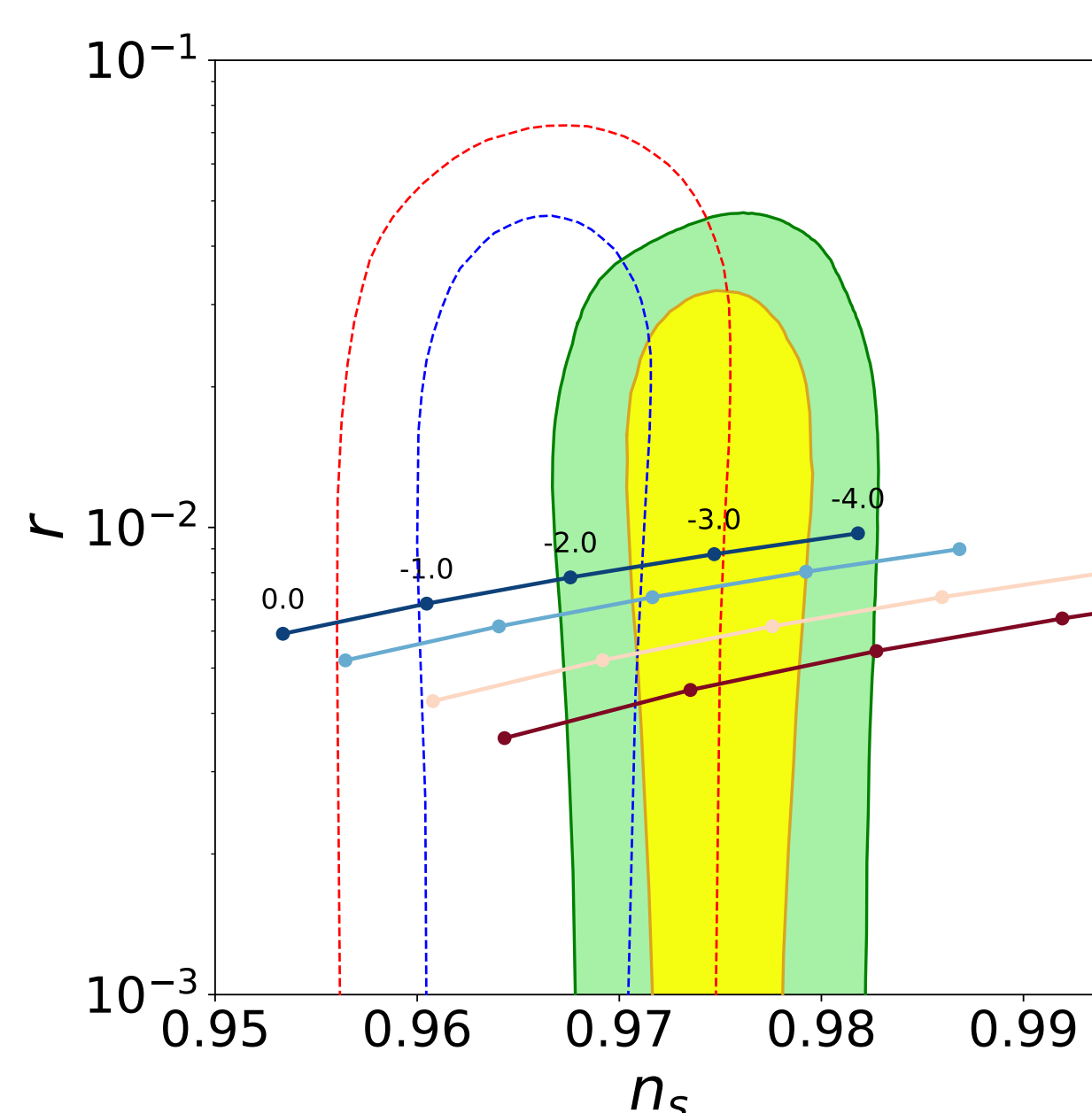


Figure: Predicted (n_s, r) values for different δ . The black line corresponds to $\delta = 0$ (Starobinsky), while the blue curve with $\delta = -10^{-4}$ lies within the ACT 1σ region (shaded).

Reheating Dynamics

After inflation ends, the Universe undergoes reheating, converting the inflaton's energy into radiation. The number of e -folds N_e is related to the reheating temperature T_{re} and the effective equation of state w_{eff} through

$$N_e = 61.4 + \frac{3w_{\text{eff}} - 1}{12(1 + w_{\text{eff}})} \ln \left(\frac{45 V_*}{\pi^2 g_* T_{\text{re}}^4} \right) - \ln \left(\frac{V_*^{1/4}}{H_*} \right).$$

This relation allows one to constrain reheating properties once (n_s, r) and δ are fixed. For a matter-dominated reheating phase ($w_{\text{eff}} = 0$), the model predicts

$$T_{\text{re}} \simeq 5.1 \times 10^9 \text{ GeV}, \quad \delta \simeq -1.6 \times 10^{-4}.$$

Positive δ values are disfavored, while mild negative δ naturally fit both the inflationary observables and the reheating constraints. This shows that a tiny cubic correction not only adjusts n_s but also yields a realistic reheating temperature compatible with the ACT-preferred region.

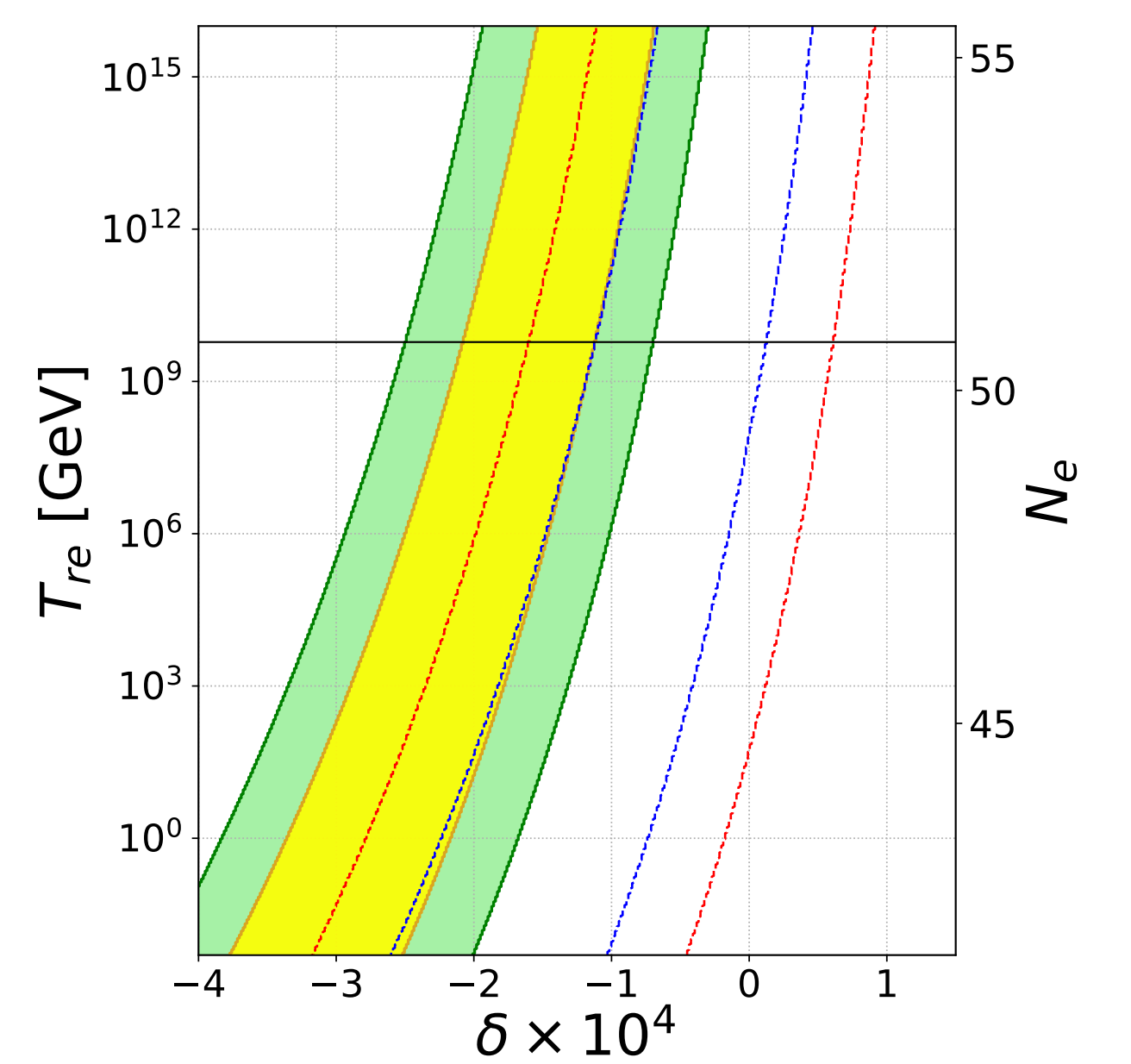


Figure: Allowed region of reheating parameters in the (δ, T_{re}) plane for $w_{\text{eff}} = 0$. The shaded band denotes the range consistent with ACT and Planck constraints on (n_s, r) .

Parameter Constraints

Scanning over the deformation parameter δ and the reheating equation of state w_{eff} , we identify the regions consistent with both the inflationary observables (n_s, r) and the reheating equation.

Our numerical analysis shows:

- Positive δ values are excluded by the ACT-DR6 spectral tilt.
- Mild negative values ($\delta \simeq -10^{-4}$) provide an excellent fit to ACT-P data.
- The allowed (δ, w_{eff}) region forms a narrow diagonal band, implying a stable one-to-one mapping between model deformation and inflationary predictions.

This means that the cubic term not only improves the fit to n_s but also keeps the inflationary operator perturbatively well-behaved, preserving predictivity within the minimal $f(R)$ framework.

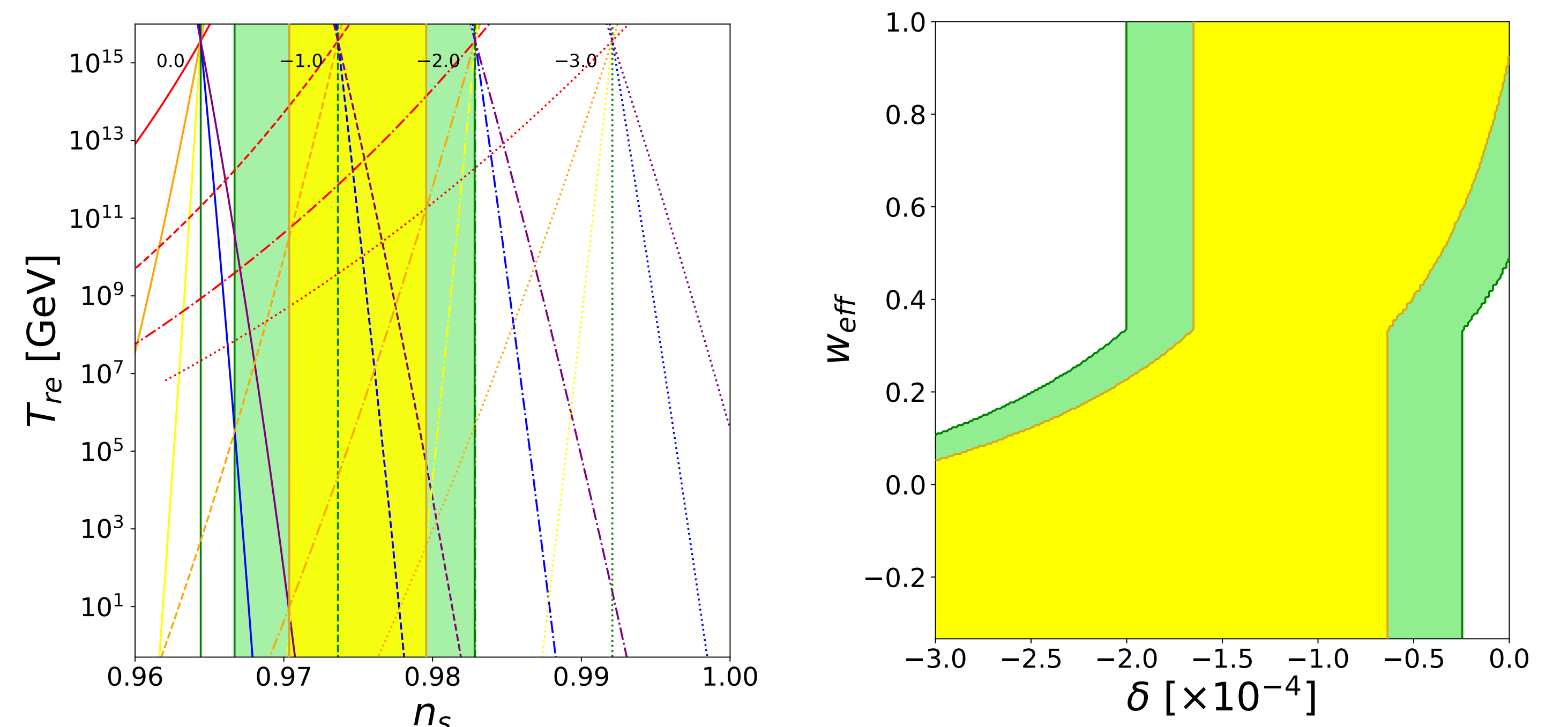


Figure: Combined parameter constraints from reheating and inflation observables. Both plots are adapted from Figs. 3–4 of the paper.

Conclusion

The original Starobinsky model ($\delta = 0$) is mildly disfavored by the recent *Atacama Cosmology Telescope* (ACT) DR6 data, showing a $\sim 2\sigma$ tension in n_s . Introducing a small cubic curvature correction,

$$\delta = \frac{\gamma}{\beta^2} \simeq -10^{-4},$$

successfully restores agreement with ACT-P results while keeping the tensor-to-scalar ratio low ($r \simeq 3 \times 10^{-3}$).

This minimal higher-order modification provides a unified picture of inflation and reheating:

- Explains the slight upward shift of n_s observed by ACT.
- Predicts a consistent reheating temperature $T_{\text{re}} \sim 5 \times 10^9 \text{ GeV}$ for $w_{\text{eff}} \simeq 0$.
- Preserves the predictive simplicity of the Starobinsky framework.

Future extensions may include additional higher-curvature terms (such as Gauss–Bonnet couplings) or multi-field generalizations to explore the robustness of these predictions across broader parameter spaces.