



Precision measurement of the hyperfine constant for ${}^9\text{Be}^+$ in a linear Paul trap

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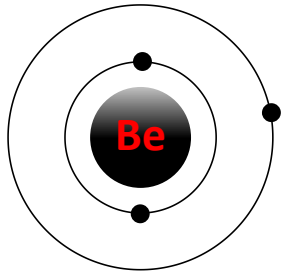
Outline

1. Background and Motivation

2. Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

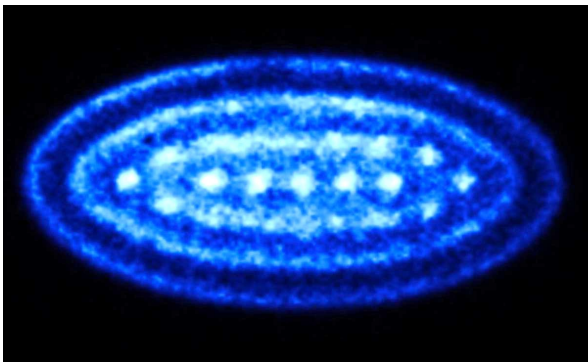
3. Conclusion and outlook

Background and Motivation



${}^9\text{Be}^+$

${}^9\text{Be}^+$ Coulomb crystal



- A few-body three-electron system, allowing for high precision QED calculation
- Forming Coulomb crystals after laser cooling
- Big charge mass ratio, good sympathetic coolants
- With $3/2$ nuclear spin

Background and Motivation

□ Hyperfine constant A and nuclear magnetic moment

- A stands for magnetic dipole hyperfine constant

$$A = \frac{2\mu_0}{3\pi} \mu_B \mu_I |\psi(0)|^2$$

nuclear magnetic moment

- Combining experiment measured A with theoretical calculations, the nuclear magnetic moment can be derived
- $|\psi(0)|^2$, Probability density of valence electrons at the nucleus. Substituting the nuclear magnetic moment can be used to test theoretical calculations

Background and Motivation

□ Hyperfine constant A and Zemach Radius

- Zemach Radius, characterizing the spatial correlation between nuclear charge and magnetic moment distributions

$$r_Z = \int d^3r d^3r' \rho_e(r) \rho_m(r') |r - r'|$$

nuclear charge density

magnetic moment density

- Comparing the experiment measured A with the QED calculated value based on point-nucleus model enables determining Zemach Radius

$$r_Z = \frac{a_0}{-2Z} \frac{A_{exp} - A_{theo}}{A_{exp}}$$

Background and Motivation

□ Current state of ground state hyperfine constant A of ${}^9\text{Be}^+$

*after diamagnetic correction

	Group	Year	M fields	A values / MHz
Strong magnetic field (Penning Trap)	J. J. Bollinger	1983	0.82 T	-625.008 837 048(10)
	M. Wada	2002	0.47 T	-625.008 835 23(75)
	J. J. Bollinger	2011	4.4609 T	-625.008 837 371(11)
				-625.008 837 044(12)*
Weak magnetic field	J.Vetter (Hot steam)	1976	5×10^{-8} T	-625.009(25)
	M. Wada (Paul Trap)	1998	0.7 mT	-625.008 82(26)
	W. Nörtershäuser (Beam)	2009	Earth field	-624.97(4)

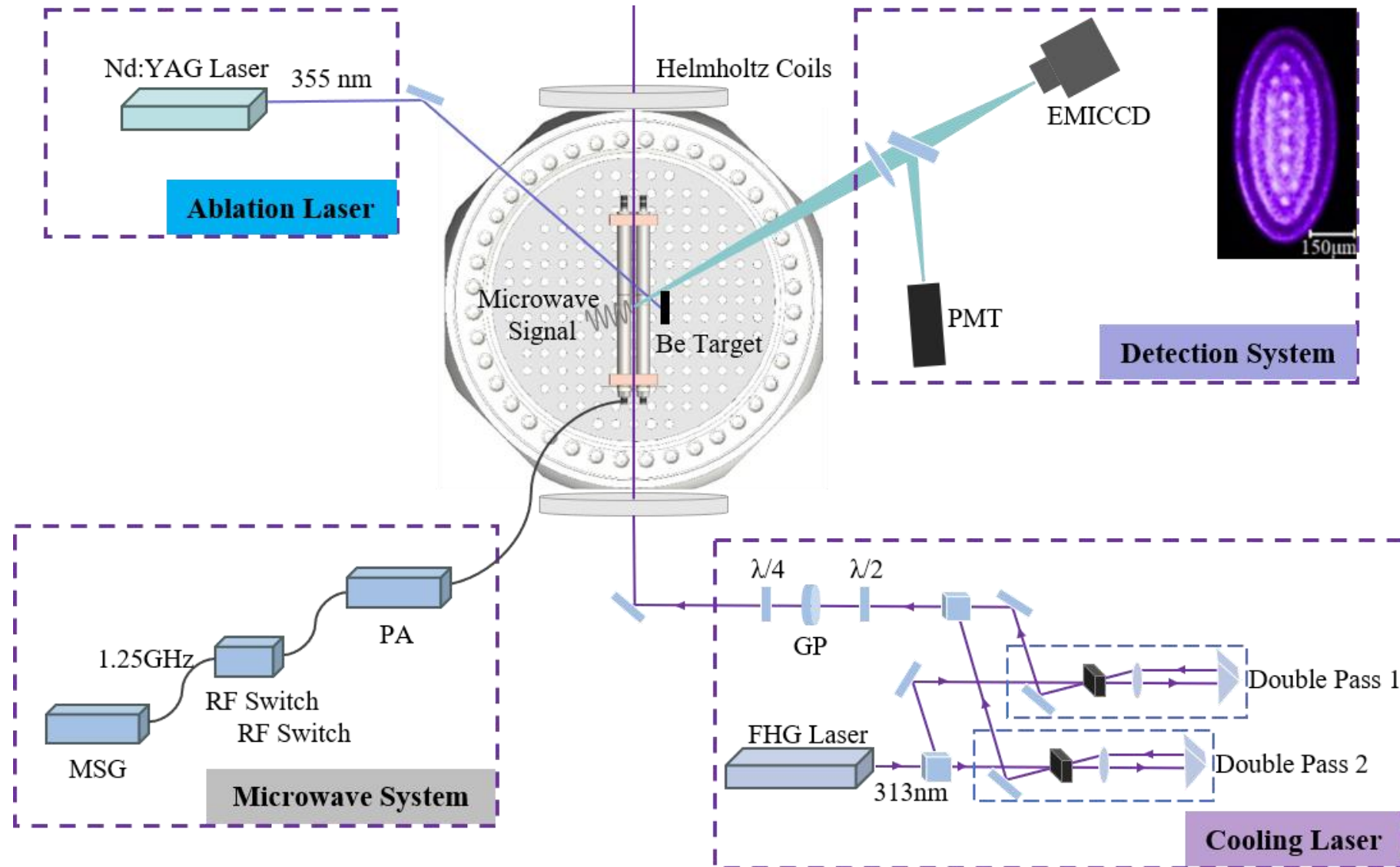
- Even considering the diamagnetic effect, there is still an error of 2.4σ
- Diamagnetic effect can be neglected under weak magnetic fields

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Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

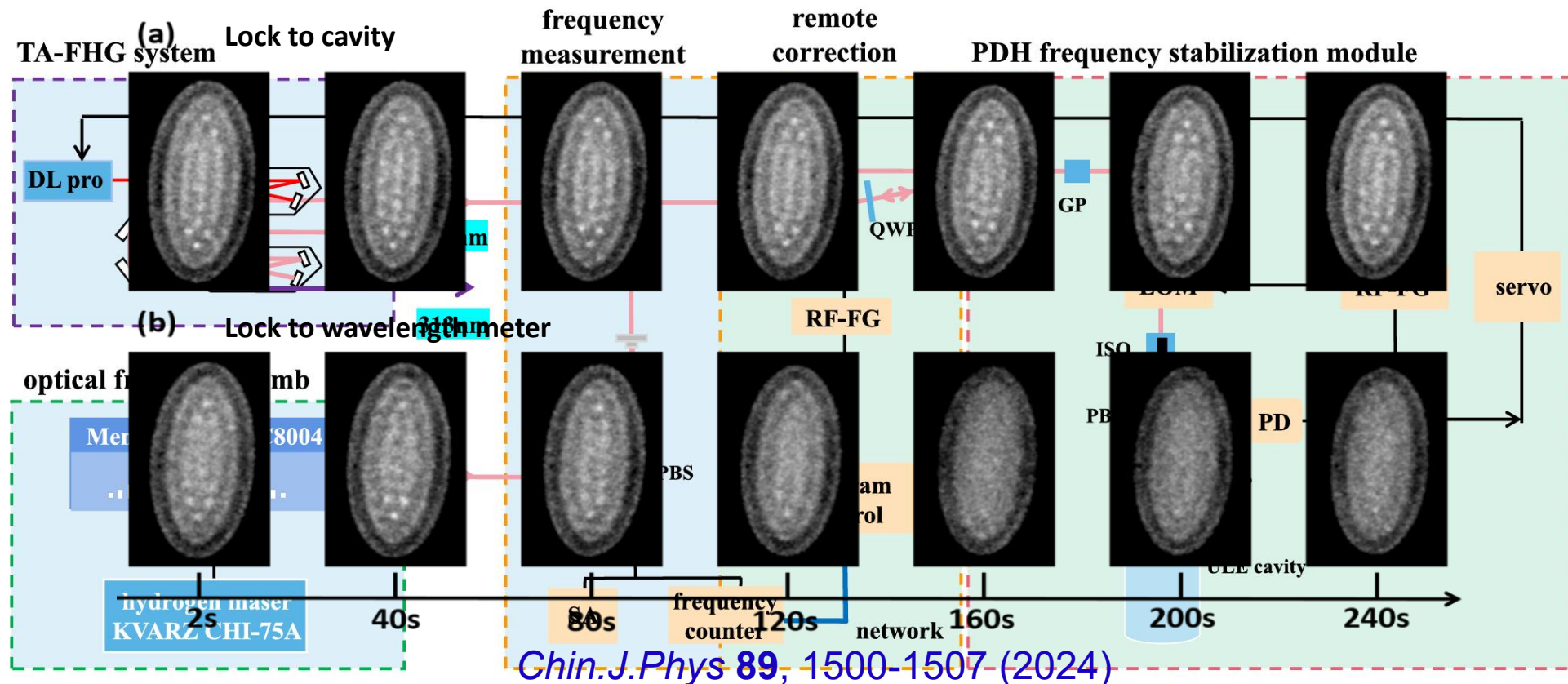
Experimental setup



Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

Preparation of stable ${}^9\text{Be}^+$ Coulomb crystal

- 626 nm laser is locked to a ultra-stable cavity, resulting frequency instability of $1.8\text{E-}12@1\text{s}$ and $3.8\text{E-}15@1024$
- ${}^9\text{Be}^+$ Coulomb crystal maintains single-ion resolution in 240s



Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

□ Hyperfine constant A and involved physical quantities

- Based on **Breit-Rabi parameterized framework**, hyperfine splitting under magnetic field \mathbf{B} for ground state can be expressed as:

$$\Delta E(x) = \Delta E_{\text{HFS}} \left[-\mathbf{g}'_I m_F x \pm \frac{1}{2} \sqrt{1 + \frac{4m_F}{2I+1} (\mathbf{g}_J + \mathbf{g}'_I)x + (\mathbf{g}_J + \mathbf{g}'_I)^2 x^2} \right]$$

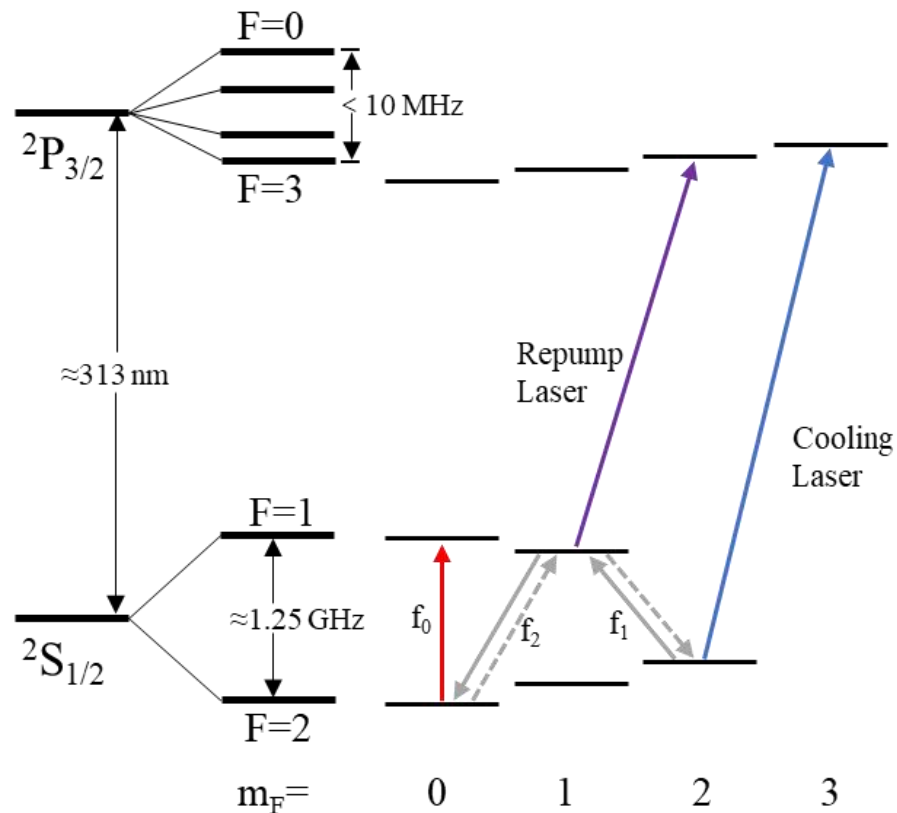
where $\Delta E_{\text{HFS}} = hA(I + \frac{1}{2})$, A stands for magnetic dipole hyperfine constant;

$x = \frac{\mu_B B}{\Delta E_{\text{HFS}}}$, implied another physical constant μ_B

Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

□ Preparation and detection of magnetic field insensitive transition

• Experimental scheme



• Sensitivity to magnetic field

	$1\mu\text{T}$	$10\mu\text{T}$	$100\mu\text{T}$
$F=2, m_F=2$ to $F=1, m_F=1$	$\sim 21\text{ kHz}/\mu\text{T}$	$\sim 21\text{ kHz}/\mu\text{T}$	$\sim 21\text{ kHz}/\mu\text{T}$
$F=1, m_F=1$ to $F=1, m_F=0$	$\sim 7\text{ kHz}/\mu\text{T}$	$\sim 7\text{ kHz}/\mu\text{T}$	$\sim 7\text{ kHz}/\mu\text{T}$
$F=1, m_F=0$ to $F=1, m_F=0$	$\sim 0.6\text{ Hz}/\mu\text{T}$	$\sim 6\text{ Hz}/\mu\text{T}$	$\sim 60\text{ Hz}/\mu\text{T}$

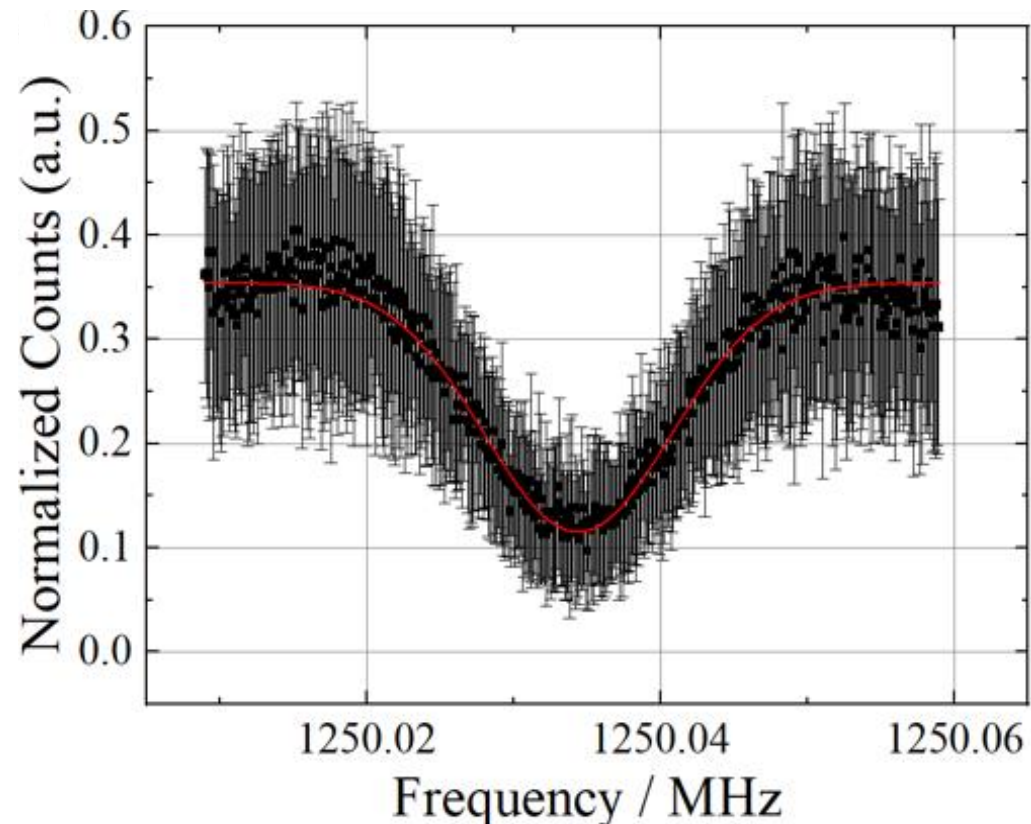
• Experimental sequence

- Optically pump to $F=2, m_F=2$
- Two pulse transfer to $F=2, m_F=0$
- Drive target transition
- Transfer the remaining ions back to $F=2, m_F=2$
- Detecting by measuring the fluorescence reduction

Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

□ Spectrum of magnetic field insensitive transition

- Measuring the relative change in the population by varying the frequency of the microwave

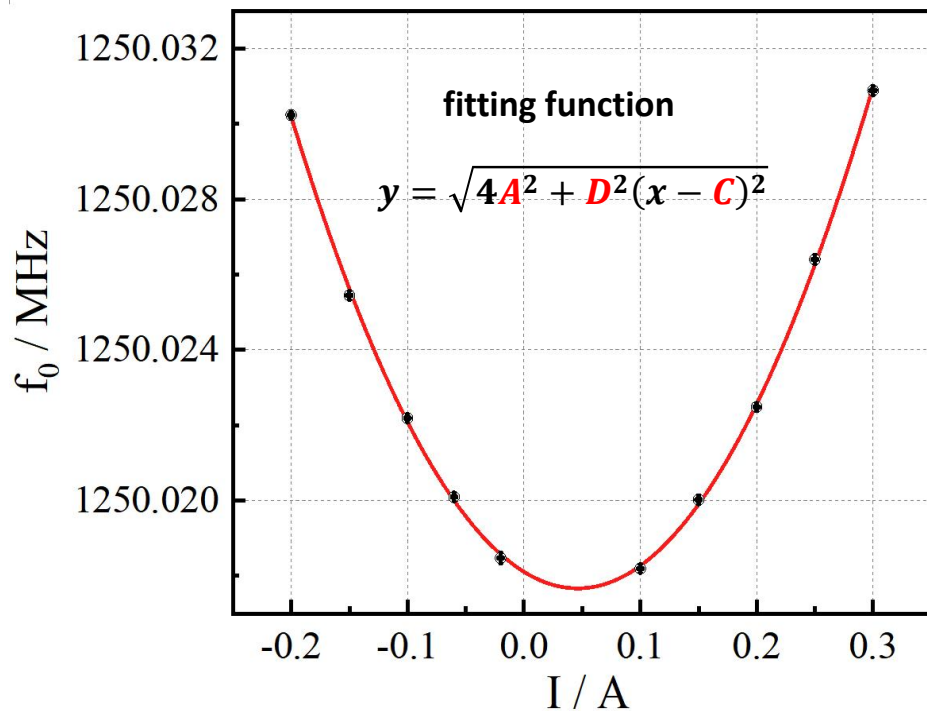


- $F=2, m_F=0 \rightarrow F=1, m_F=0$ @ 0.3A coil current ($\sim 190\mu\text{T}$)

Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

□ Determination of hyperfine constant A

$$f_{00} = \frac{\Delta E_{HFS}}{h} \sqrt{1 + \left[\frac{\mu_B (g_J - g_{I'})}{\Delta E_{HFS}} \right]^2 B^2} = \sqrt{4A^2 + \left[\frac{\mu_B (g_J - g_{I'})}{h} \right]^2 B^2}$$



- f_{00} : magnetic field insensitive transition frequencies
- A : hyperfine structure constant in near-zero field
- $D = \frac{\mu_B (g_J - g_{I'})}{h} k$, k current to magnetic field conversion coefficient
- C : current offset from zero magnetic field ($B = k(I - C)$)

$$A = -625.008833(34) \text{ MHz}$$

$$D = 22.8(1) \text{ MHz A}^{-1}$$

$$C = 0.0462(6) \text{ A}$$

Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

□ systematic error evaluation

- Determination of A : **$-625.008833(35)$ MHz** (5.6×10^{-8})
- Statistical error represents the maximum error due to the magnetic field drifts
- Errors arising from magnetic field inhomogeneity, microwave frequency uncertainty, power broadening and the electric field are negligible

Source of uncertainty	Shift (Hz)	Uncertainty (Hz)
Statistical error	0	34
Magnetic field inhomogeneity	0	0.02
Second-order Doppler	0.058	0
Microwave signal generator (MSG)	0	0.003
Lifetime	17.4*	1.4
Total	0	35

Precision determination of hyperfine constant A for ${}^9\text{Be}^+$

□ Zemach radius

- Bohr radius (CODATA2022)
- Experimental determination A value: $-625.008833(35)\text{MHz}$
- Theoretical value from Pachucki: $-625.3927(36)(16)\text{MHz}$

$$r_z = \frac{a_0}{-2Z} \frac{A_{exp} - A_{theo}}{A_{exp}}$$

- Zemach radius of ${}^9\text{Be}^+$ is determined to be **4.03(5) fm**

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Conclusion

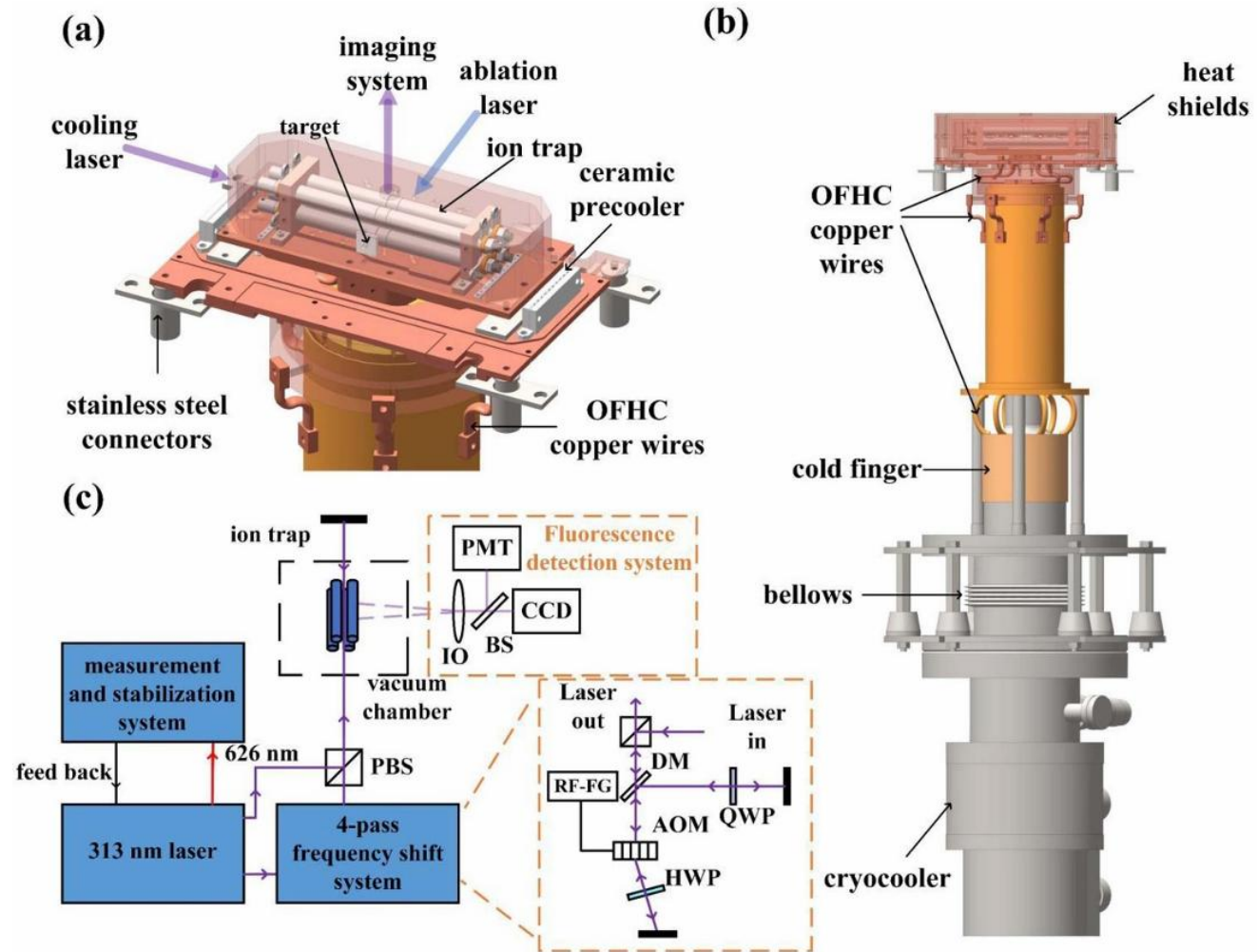
- Zero-field ground-state hyperfine constant: $-625.008833(35)$ MHz (5.6×10^{-8})
- Combining experimental and theoretical value, ${}^9\text{Be}$'s nuclear Zemach radius is calculated as **4.03(5) fm**.

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	Our work (Paul Trap)	2025	<0.3 mT	-625.008 833(35)
Theory	V. A. Yerokhin	2008	\	-625.08(2)
				-625.11(3)
	K. Pachucki	2014	\	-625.3927(36)(16)

Outlook

□ Doing better

- Adding magnetic shielding to reduce the effects of magnetic field fluctuations and magnetic field gradients
- Increase the vacuum level to reduce the effect of ion lifetime

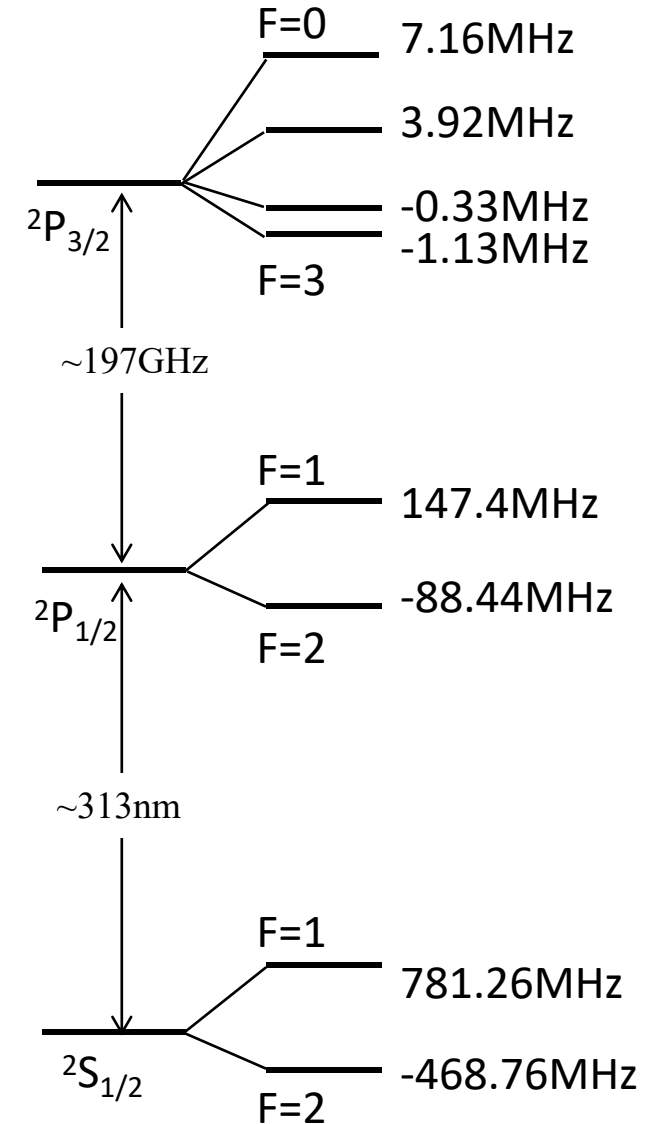


Outlook

□ Going further

- Determination the hyperfine constant $A_{P_{1/2}}$ for $^2P_{1/2}$ state
- Determination the hyperfine constant $A_{P_{3/2}}$ and $B_{P_{3/2}}$ for $^2P_{3/2}$ state

	Theory		Experiment
$A_{P_{1/2}}$	-117.919(4)	-117.932(3)	-117.92(4)
$A_{P_{3/2}}$	-1.015(3)	-1.026(3)	-1.01(2)
$B_{P_{3/2}}$		-2.29940(3)	-2.23(8)



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