

Space-time structure and phenomenology of Cartan Khronon theory

Lucy Zheng

Relevant publications “khronon theory” “Lorentz gauge theory” “Spin(4) gauge theory”, ...etc.

- 2018 • **Spacetime and dark matter from spontaneous breaking of Lorentz symmetry**
Złońnik, Urban, Marzola, Koivisto e-Print: 1807.01100 [gr-qc] DOI: 10.1088/1361-6382/aaea96
- 2019 • **The General Linear Cartan Khronon**
Koivisto, Hohmann, Złońnik e-Print: 1905.02967 [gr-qc] DOI: 10.3390/universe5070168
- 2021 • **The Λ and the CDM as Integration Constants**
Gallagher, Koivisto e-Print: 2103.05435 [gr-qc] DOI: 10.3390/sym13112076
- 2022 • **Pregeometric first order Yang-Mills theory**
Gallagher, Koivisto, Marzola e-Print: 2202.05657 [hep-th] DOI: 10.1103/PhysRevD.105.125010
- 2023 • **Paths to gravitation via the gauging of parametrized field theories**
Koivisto, Tom Zlosnik e-Print: 2212.04562 [gr-qc] DOI: 10.1103/PhysRevD.107.124013
- **Cosmology in the Lorentz gauge theory**
Koivisto e-Print: 2306.00963 [gr-qc] DOI: 10.1142/S0219887824500403
- **Hamiltonian formulation of gravity as a spontaneously-broken gauge theory of the Lorentz group**
Nikjoo, Zlosnik e-Print: 2308.01108 [gr-qc] DOI: 10.1088/1361-6382/ad1c84
- **Consistent first-order action functional for gauge theories** e-Print: 2311.07464 [hep-th]
Gallagher, Koivisto, Marzola, Varrin, Zlosnik DOI: 10.1103/PhysRevD.109.L061503
- 2024 • **Canonical aspects of pregeometric vector-based first order gauge theory**
Gallagher e-Print: 2403.02578 [hep-th] DOI: 10.1103/PhysRevD.110.085010
- **Black holes in Lorentz gauge theory**
Koivisto, Zheng e-Print: 2408.10100 [gr-qc] DOI: 10.1103/PhysRevD.111.064008
- 2025 • **Spin(4) gauge theory of space, time, gravitation, matter, and dark matter**
Koivisto, Zheng, Zlosnik e-Print: 2507.00968 [gr-qc] DOI: 10.1103/xx3x-s5wj

Relevant theses (2026)

(At the Corner of Quantum and Gravity)

Ludovic Varrin [arXiv: 2603.21941]

On the internal gauge theory analogy to the Cartan Khronon theory of gravity

Ernest Michael Priidik Gallagher [University of Tartu Press]

The Gravitating Universe

Mehraveh Nikjoo [TBA]

Cartan Khronon – Real space-time structure –

Lucy Zheng [TBA]

$(+, +, +, +)$

“But we live in a Lorentzian spacetime”...?

Lorentzian	Euclidean
Gravity	Quantum field theory
Cosmology	Black hole thermodynamics
⋮	Statistical field theory
Causal dynamics	

The necessity of Euclidean time

Path integral in quantum theory

$$\langle g_1, \phi_1, S_1 | g_2, \phi_2, S_2 | g_1, \phi_1, S_1 | g_2, \phi_2, S_2 \rangle = \int D[g, \phi] \exp(iI[g, \phi]),$$

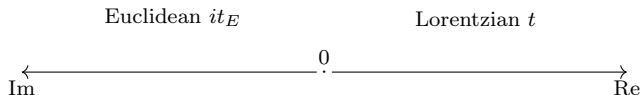
Black hole thermodynamics

$$Z = \text{Tr}\{\exp(-\beta H)\} = \int D[g, \phi] \exp(-I_E[g, \phi]),$$

(Hawking temperature $T = 1/\beta$, intrinsic entropy $S = \beta \frac{\partial I_E}{\partial \beta} - I_E$)

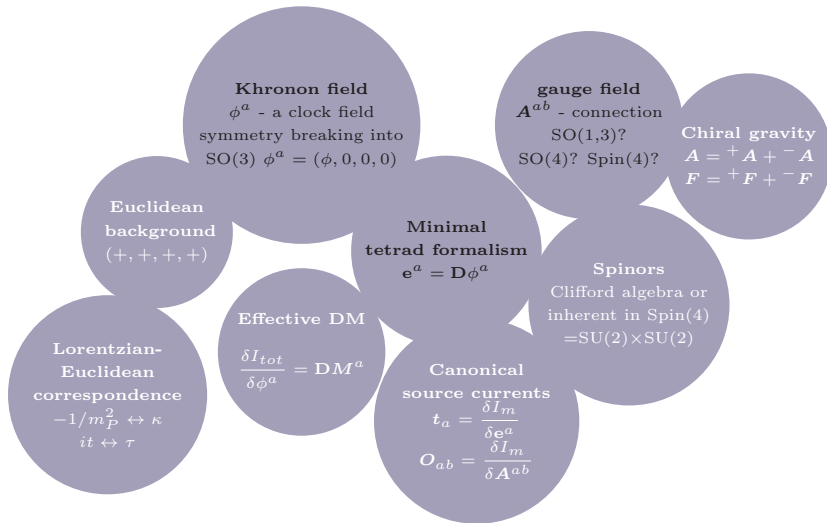
Is Euclidean spacetime the fundamental framework to define physics?

Standard Wick rotation

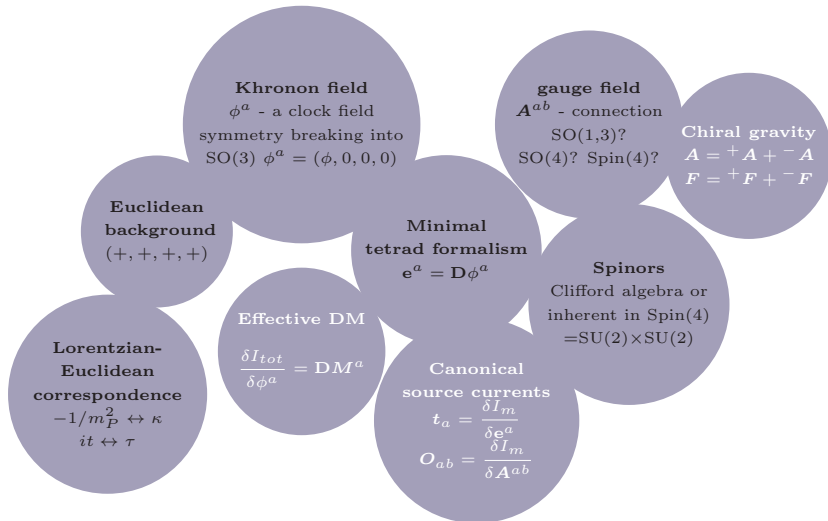


- ▶ does not apply to certain spacetime metrics
- ▶ index rotation of the gamma matrices with $\gamma^4 = -i\gamma^0$ fails to
 - establish Euclidean symmetry in Dirac action
 - maintain hermiticity $\psi_{\pm}^{\dagger} = \psi_{\mp}$
 - maintain self-conjugacy of Majorana spinor (components become complex)

Cartan Khronon framework



Khronon + Spin(4) as a solution



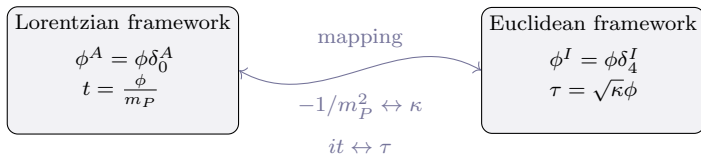
Works for Kerr, Cosmological background & perturbations, etc.

Euclidean-Lorentzian correspondence – the dimension

For a dimensionless symmetry breaking scalar field ϕ^a , $\mathbf{e}^a = \mathbf{D}\phi^a \sim [E]$.
Introducing dimensionful $[E]$ parameters m_P and $1/\sqrt{\kappa}$ in respective framework,

$$t = \phi/m_P \text{ and } \tau = \sqrt{\kappa}\phi$$

$\mathbf{e}^0 = \mathbf{d}t$ and $\mathbf{e}^4 = \mathbf{d}\tau$ both become dimensionless.



All the dimensionful parameters require mapping, e.g., spin parameter $a = J/M$

Euclidean-Lorentzian correspondence – spinors

A unitary rotation $U = \exp\left(\gamma^4\gamma^5\frac{\theta}{2}\right)$ by $\theta = \pi/2$ maps spinors between the two frameworks without violating each group symmetry

$$\gamma_E^i = \gamma_L^i, \quad \gamma_E^4 = \gamma_L^5, \quad \gamma_E^5 = \gamma_E^1\gamma_E^2\gamma_E^3\gamma_E^4 = -\gamma^4,$$

which satisfy the corresponding hermiticity relations.

$$\left\{\gamma_E^I, \gamma_E^J\right\} = 2\delta^{IJ}\mathbb{1}, \quad (\gamma_E^I)^\dagger = -\gamma_E^5\gamma_E^I\gamma_E^5 = \gamma_E^I, \quad (\gamma_E^5)^2 = \mathbb{1}.$$

This new system can be understood as a reinterpretation of γ^4 and γ^5 .

Phenomenology in the purely right-(left-)handed limit

The full ambidextrous action

$$I = \int \mathbf{D}\phi^I \wedge \mathbf{D}\phi^J \wedge (g_+{}^+ \mathbf{F}_{IJ} - g_-{}^- \mathbf{F}_{IJ}) - \int \star\lambda - \int \mathbf{D}\phi^I \wedge \mathbf{t}_I - \int \mathbf{A}^{IJ} \wedge \mathbf{O}_{IJ},$$

In the $(g_+, g_-) = (1, 0), (0, 1)$ limit, when $\mathbf{O}^{IJ} = 0$, we recover...

- ▶ FLRW background + $\Lambda (= -m_P^{-2}\lambda)$ + effective ideal dust \mathbf{M}^I
(vanishing pressure and sound speed for \mathbf{M}^I)
- ▶ Standard GR black hole solutions
- ▶ Gravitational wave propagating at the speed of light

Phenomenology in the general chiral configuration

The full ambidextrous action

$$I = \int \mathbf{D}\phi^I \wedge \mathbf{D}\phi^J \wedge (g_+{}^+ \mathbf{F}_{IJ} - g_-{}^- \mathbf{F}_{IJ}) - \int \star\lambda - \int \mathbf{D}\phi^I \wedge \mathbf{t}_I - \int \mathbf{A}^{IJ} \wedge \mathbf{O}_{IJ},$$

For the general $g_+g_- \neq 0$, when $\mathbf{O}^{IJ} = 0$, we obtain...

- ▶ FLRW background with \mathbf{M}^I -effect included in the vacuum + $\Lambda (= -m_P^{-2}\lambda)$
(vanishing sound speed for \mathbf{M}^I)
- ▶ No black hole solutions have been found
- ▶ Gravitational wave propagating at the speed of $|\gamma| \equiv \left| \frac{g_+ + g_-}{g_+ - g_-} \right|$

multi-messenger data from the neutron star merger GW170817 :
 $|g_{\mp}/g_{\pm}| \lesssim 10^{-15}$ at present

Phenomenology in the chirally symmetric limit

The full ambidextrous action

$$I = \int \mathbf{D}\phi^I \wedge \mathbf{D}\phi^J \wedge (g_+{}^+ \mathbf{F}_{IJ} - g_-{}^- \mathbf{F}_{IJ}) - \int \star\lambda - \int \mathbf{D}\phi^I \wedge \mathbf{t}_I - \int \mathbf{A}^{IJ} \wedge \mathbf{O}_{IJ},$$

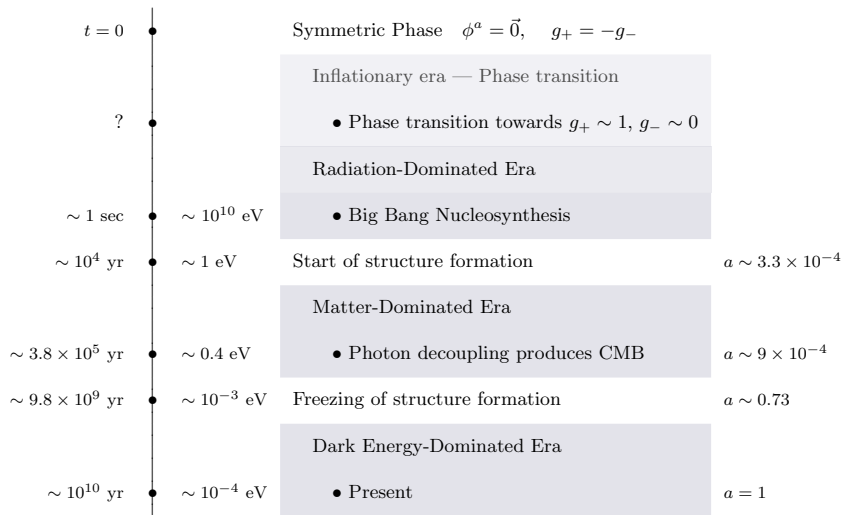
For the chirally symmetric $g_+ = -g_-$ case, when $\mathbf{O}^{IJ} = 0$, we obtain...

- ▶ khronon field acting as scale factor $a \sim \phi$
- ▶ radiation-dominated universe $H \sim t^{-1}$
- ▶ M^I : candidate for dark radiation...?

Likely scenario:

chirally symmetric gravity $g_+ = -g_-$ in early universe (before BBN)
transitions into the (almost) purely right-(left-)handed gravity

Cartan Khronon Cosmic Timeline



Phenomenology with spin current

The full ambidextrous action

$$I = \int \mathbf{D}\phi^I \wedge \mathbf{D}\phi^J \wedge (g_+^+ \mathbf{F}_{IJ} - g_-^- \mathbf{F}_{IJ}) - \int \star\lambda - \int \mathbf{D}\phi^I \wedge \mathbf{t}_I - \int \mathbf{A}^{IJ} \wedge \mathbf{O}_{IJ},$$

When $\pm \mathbf{O}^i = \sqrt{\kappa} \pm O \star \mathbf{e}^i \neq 0$, we obtain...

- Coupling between effective DM and material spin current

$$\dot{\hat{\rho}} + 3H \left(\hat{\rho} + \frac{2\beta}{\kappa\tau^2} \right) = \frac{6\beta}{\kappa\tau^3} - \frac{+O + -O}{2\gamma\tau^2}$$

- New effect in the Friedmann equations

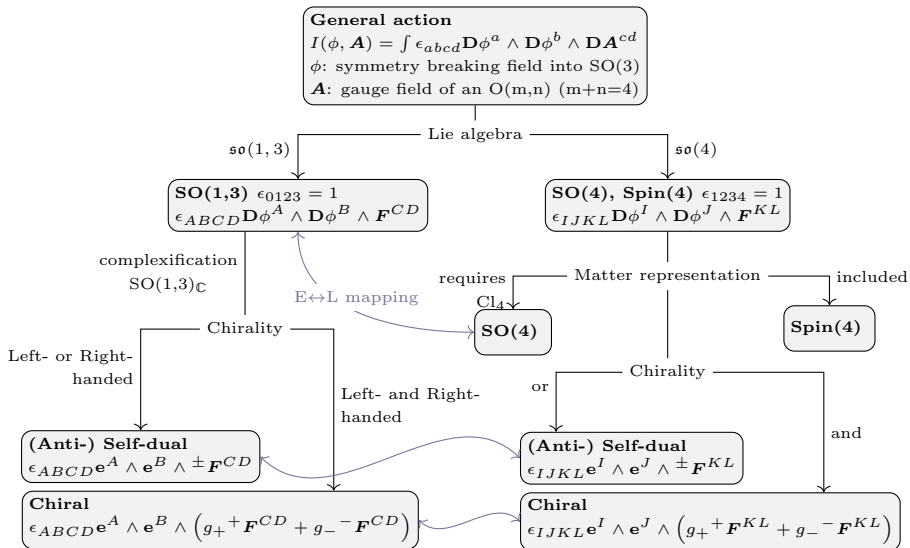
$$\Omega \equiv \frac{\kappa}{2\gamma} (+O + -O), \quad \Sigma \equiv \frac{\kappa^{3/2}}{\alpha\gamma^2\phi} (g_+^- O - g_-^+ O)$$

$$3\alpha H^2 + 3H\Omega + \frac{3}{4} \frac{\Omega^2}{\alpha} = \kappa(\rho + \hat{\rho}) + 3\beta\tau^{-2},$$

$$3\alpha H^2 + 2\alpha\dot{H} + 2H\Omega + \frac{\Omega^2}{4\alpha} + \dot{\Omega} + \frac{\Sigma}{\sqrt{\kappa}} = -\kappa p + \beta\tau^{-2}.$$

Is Lorentzian spacetime – where we *observe* physics –
the optimal framework to *define* physics?

Cartan-khronon flowchart



Perturbations in GR limit $\beta = 0$: FLRW + Λ + ideal dust

► Standard cosmology

Gravitational wave

$$\ddot{h}_{ij} + 3\frac{\dot{a}}{a}\dot{h}_{ij} + \frac{1}{a^2}\nabla^2 h_{ij} = 0 \xrightarrow{\text{Wick}} -h''_{ij} - 3Hh'_{ij} + \frac{1}{a^2}\nabla^2 h_{ij} = 0 \quad (25)$$

No propagation for vector perturbation

$$\begin{aligned} \nabla^2 V_i &= 0 \\ V'_i + 3HV_i &= 0 \end{aligned} \quad (26)$$

Continuity equation under scalar perturbation

$$\delta\rho' + 3H(\delta\rho + \delta p) = (\rho + p) \left[\nabla^2 \left(\frac{v}{a} + c' - \frac{\Phi}{a^2} \right) - 3\Psi' \right] \quad (27a)$$

Euler equation

$$a(\rho + p)(v' + 4aHv) + (\rho + p)'v = \delta p - \frac{2}{3}\nabla^2\Pi \quad (27b)$$

► + effective ideal dust

Continuity equation

$$\delta\hat{\rho}' + 3H\delta\hat{\rho} = \Phi'\hat{\rho}' + 3 \left(H\Phi' - \Psi' + \frac{1}{3}\nabla^2 c' \right) \hat{\rho} \quad (28a)$$

Density evolution of CDM

($\hat{\delta} \equiv \delta\hat{\rho}/\hat{\rho}$)

$$\hat{\delta}'' + 2H\hat{\delta}' = \frac{\hat{\rho}}{2m_P^2}\hat{\delta} \quad (28b)$$

$g_+^+ F - g_-^- F$: nature of M^a

- ▶ effective pressure

$$\hat{p} = \beta \left(\frac{H}{t} - \frac{1}{t^2} \right)$$

- ▶ effective energy density sourced by its own pressure

$$\hat{\rho}' + 3H\hat{\rho} = -3\hat{p}/t$$

- ▶ acceleration driven by time lapse fluctuation

$$\hat{v}' \sim \frac{\partial\varphi}{\partial\phi}$$

- ▶ clustering

$$\delta\hat{\rho}' + 3H\delta\hat{\rho} = \frac{\nabla^2}{a^2} \left[2\beta m_P^2 (\Psi' - H\Phi') + \left(\hat{\rho} - 2\beta m_P^2 H' + \frac{2\beta m_P^2}{t^2} \right) \Phi \right]$$

- ▶ vanishing sound speed

$$\hat{c}_s \equiv \frac{\hat{p}'}{\hat{\rho}'} = 0$$