

Unified and higher spin field theories from Hamiltonian cotangent bundle geometry

Deriving the Einstein-Maxwell System from a Single Phase-Space Scalar Field

Joint work with J. Relancio

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Dr. Christian Pfeifer
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1. Unified field theories
2. A scalar field equation for point particle Hamiltonians
3. From the scalar equation on phase space to the Einstein-Maxwell equations
4. Conclusion

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2. A scalar field equation for point particle Hamiltonians

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“a theory joining the gravitational and the electromagnetic field into one single hyper-field whose equations represent the conditions imposed on the geometrical structure of the universe”

[Tonnalet 1995, H. Goenner 2004]



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Gravity and Electromagnetism I
Extended spacetime geometry



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Not successful to reproduce
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Metric affine gravity

Teleparallel gravity

Poincaré-gauge theory



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All emerge from a fundamental electroweak gauge group

At low energies the gauge group decays into the weak and electromagnetic sector



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- Grand unified theory (GUT)
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Mostly applied as modified gravity
not unification

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No consensus how to consistently reproduce the standard model yet

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This Talk

Gravity and Electromagnetism Unification using the geometry of phase space

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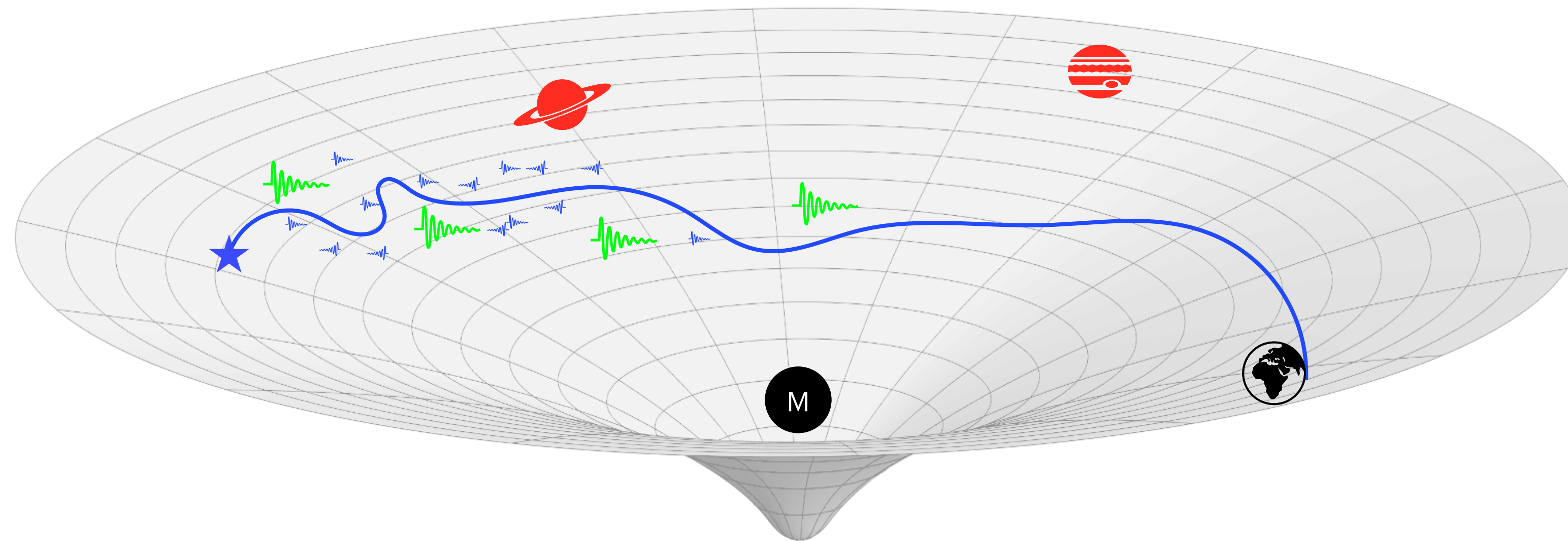
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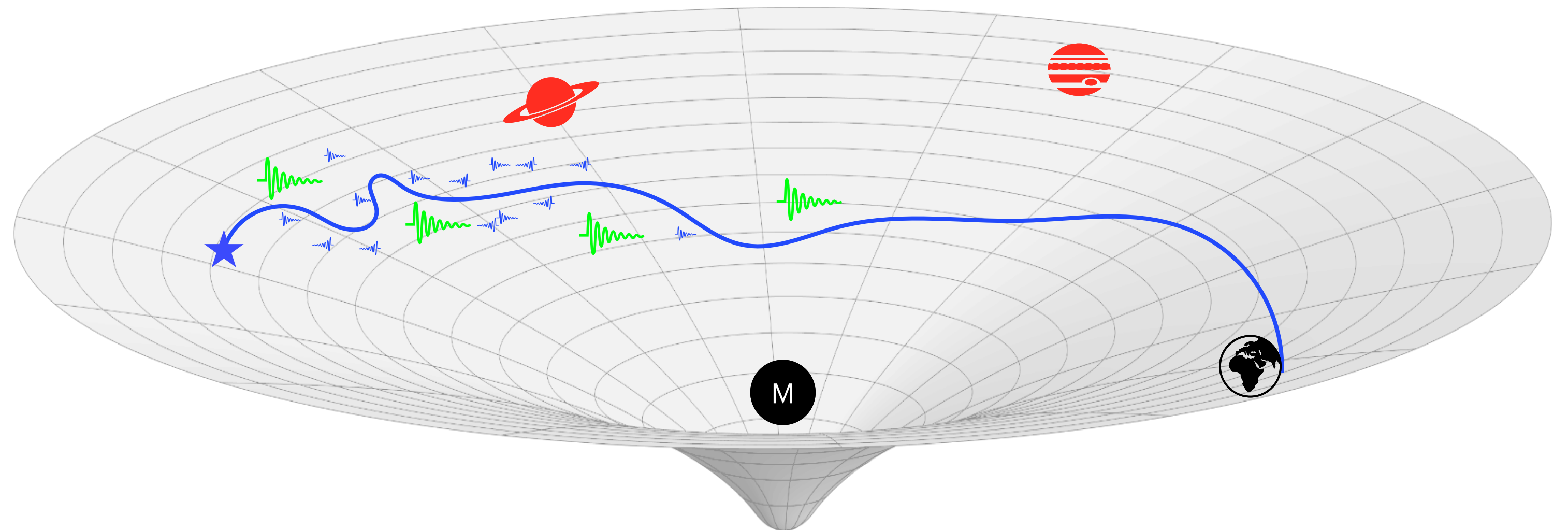
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Classical particles



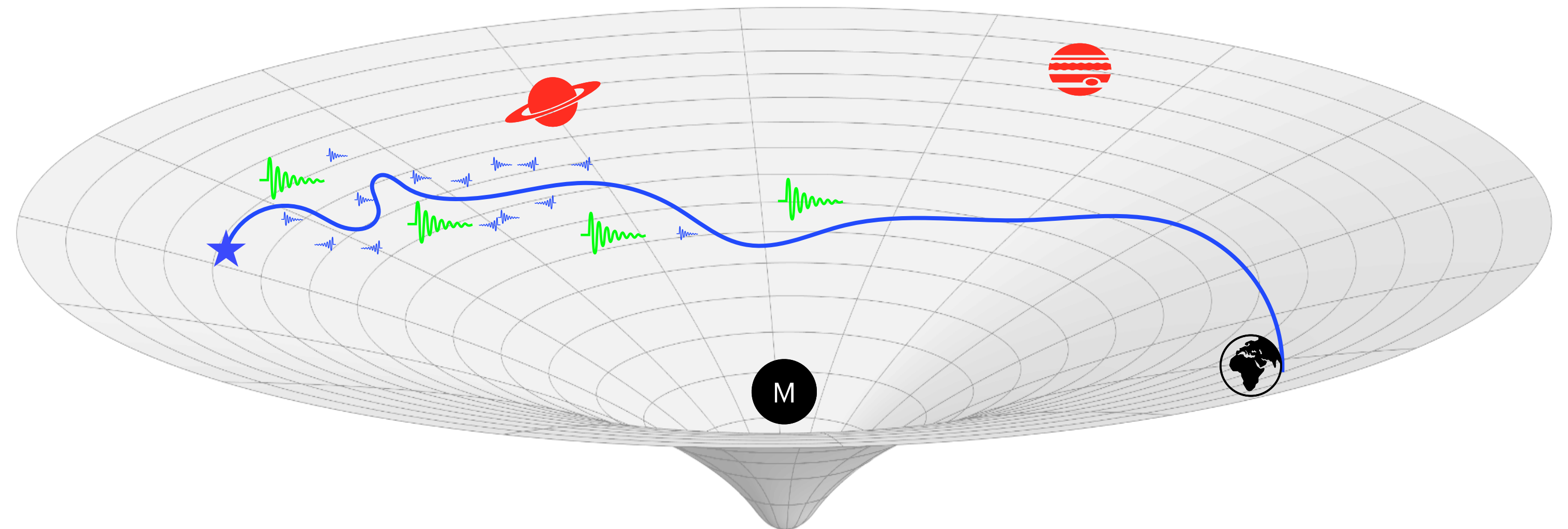


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- Particle motion from Hamilton fct.

$$H(x, p)$$





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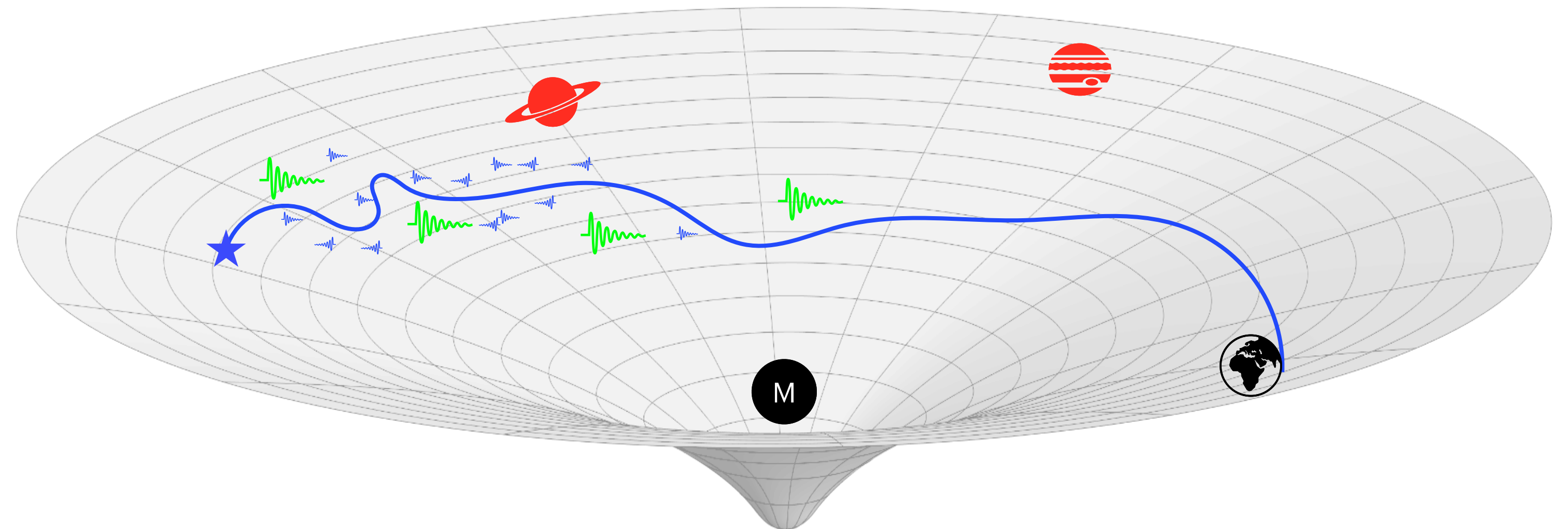
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- determining dispersion relation

$$H(x, p) = m^2$$

- and Hamilton equations of motion

$$\dot{x}^\mu = \partial_{p_\mu} H, \quad \dot{p}_\mu = -\partial_{x^\mu} H$$





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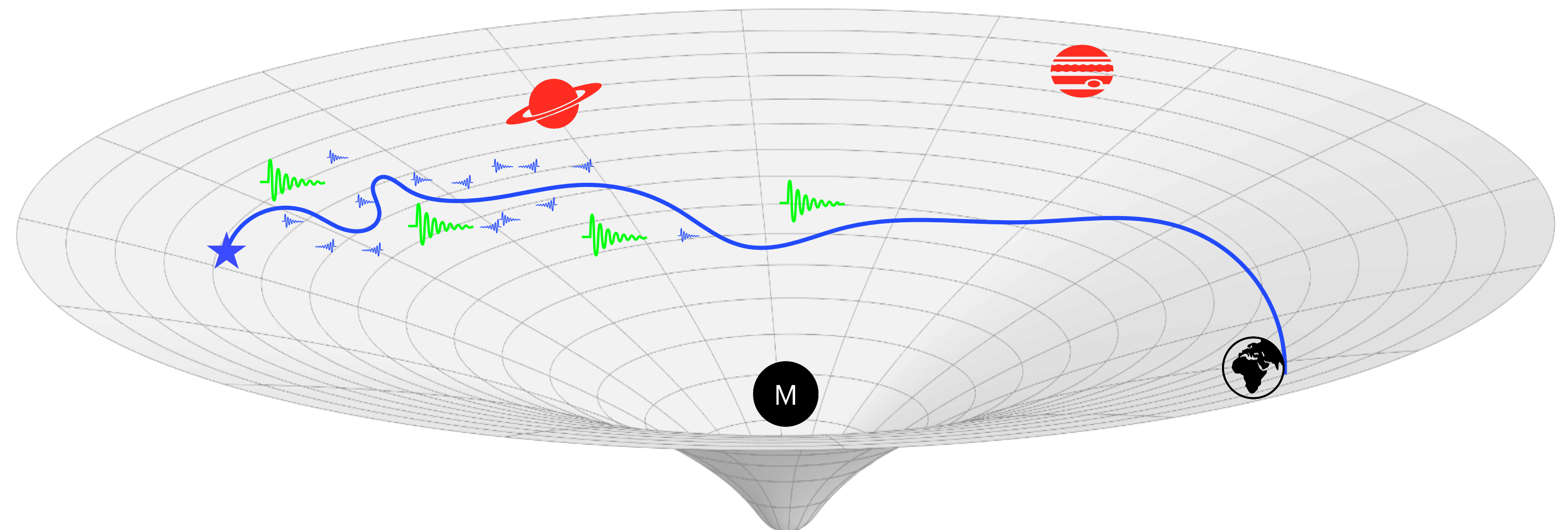
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- Hamiltonian main ingredient in Schroedinger equation

$$\frac{d}{dt}\Psi = \hat{H}\Psi$$





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- Further terms - QG Phenomenology

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[Addazzi 2022, Läanemets 2024]



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[Rubilar 2007; Hehl, Obukhov 2003]



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[Synge 1960; Tsupko 2025]



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$$\dot{x}^\mu = \partial_{p_\mu} H, \quad \dot{p}_\mu = -\partial_{x^\mu} H$$

Quantum particles

- Hamiltonian main ingredient in Schroedinger equation

$$\frac{d}{dt} \Psi = \hat{H} \Psi$$

Which Hamiltonian?

- Hamiltonian chosen depending on particle - background field interaction

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Can we determine the dynamical equations for background fields from a dynamical equation for H?

1. Unified field theories
- 2. A scalar field equation for point particle Hamiltonians**
3. From the scalar equation on phase space to the Einstein-Maxwell equations
4. Conclusion



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(spacetime coordinate invariant, ...)
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The answer: Hamilton geometry - Derive the geometry of phase space from H

[Miron; Barcaroli 2015; CP, J. Relancio 2026]

**The action**

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$



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$$N_{\mu\nu} = \frac{1}{4} \left(\{g_{\mu\nu}, H\} - g_{\mu\rho} \partial_\nu \bar{\partial}^\rho H - g_{\nu\rho} \partial_\mu \bar{\partial}^\rho H \right)$$

$$\delta_\mu H = 0 \Leftrightarrow H(x, \lambda p) = \lambda^n H(x, p)$$

- The canonical covariant derivative ∇_μ

$$\nabla_\mu Z^\nu = \delta_\mu Z^\nu + H^\nu{}_{\mu\sigma} Z^\sigma$$

$$H^\sigma{}_{\mu\nu} = \frac{1}{2} g^{\sigma\lambda} (\delta_\mu g_{\lambda\nu} + \delta_\nu g_{\lambda\mu} - \delta_\lambda g_{\mu\nu})$$

- The curvature

$$\mathcal{R}^\mu{}_{\nu\rho\sigma}(x, p) = \delta_\rho H^\mu{}_{\nu\sigma} - \delta_\sigma H^\mu{}_{\nu\rho} + H^\mu{}_{\lambda\rho} H^\lambda{}_{\nu\sigma} - H^\mu{}_{\lambda\sigma} H^\lambda{}_{\nu\rho}$$

$$\mathcal{R} = \mathcal{R}^\mu{}_{\nu\mu\sigma} g^{\sigma\nu} = \mathcal{R}_{\nu\sigma} g^{\sigma\nu}.$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$



$$\delta_H S = 0$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$

- The Hamilton function

$$H(x, p)$$

- x and p derivatives

$$\partial_\mu = \frac{\partial}{\partial x^\mu} \quad \bar{\partial}^\mu = \frac{\partial}{\partial p_\mu}$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = \frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu H$$

- The horizontal derivative

$$\delta_\mu = \partial_\mu + N_{\mu\nu} \bar{\partial}^\nu$$

$$N_{\mu\nu} = \frac{1}{4} \left(\{g_{\mu\nu}, H\} - g_{\mu\rho} \partial_\nu \bar{\partial}^\rho H - g_{\nu\rho} \partial_\mu \bar{\partial}^\rho H \right)$$

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$$\mathcal{R} = \mathcal{R}^\mu{}_{\nu\mu\sigma} g^{\sigma\nu} = \mathcal{R}_{\nu\sigma} g^{\sigma\nu}.$$

- The Berwald covariant derivative $\tilde{\nabla}_\mu$

$$\tilde{\nabla}_\mu Z^\nu = \delta_\mu Z^\nu + \bar{\partial}^\nu N_{\mu\sigma} Z^\sigma$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$



$$\delta_H S = 0$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + \mathcal{Q}^\rho{}_\rho g_{\mu\nu} - \mathcal{Q}_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$

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$$H(x, p)$$

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$$\partial_\mu = \frac{\partial}{\partial x^\mu} \quad \bar{\partial}^\mu = \frac{\partial}{\partial p_\mu}$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = \frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu H$$

- The horizontal derivative

$$\delta_\mu = \partial_\mu + N_{\mu\nu} \bar{\partial}^\nu$$

$$N_{\mu\nu} = \frac{1}{4} \left(\{g_{\mu\nu}, H\} - g_{\mu\rho} \partial_\nu \bar{\partial}^\rho H - g_{\nu\rho} \partial_\mu \bar{\partial}^\rho H \right)$$

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$$\nabla_\mu Z^\nu = \delta_\mu Z^\nu + H^\nu{}_{\mu\sigma} Z^\sigma$$

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- The curvature

$$\mathcal{R}^\mu{}_{\nu\rho\sigma}(x, p) = \delta_\rho H^\mu{}_{\nu\sigma} - \delta_\sigma H^\mu{}_{\nu\rho} + H^\mu{}_{\lambda\rho} H^\lambda{}_{\nu\sigma} - H^\mu{}_{\lambda\sigma} H^\lambda{}_{\nu\rho}$$

$$\mathcal{R} = \mathcal{R}^\mu{}_{\nu\mu\sigma} g^{\sigma\nu} = \mathcal{R}_{\nu\sigma} g^{\sigma\nu}.$$

- The Berwald covariant derivative $\tilde{\nabla}_\mu$

$$\tilde{\nabla}_\mu Z^\nu = \delta_\mu Z^\nu + \bar{\partial}^\nu N_{\mu\sigma} Z^\sigma$$

- The Q tensor

$$Q_{\lambda\nu} = (-\nabla_\nu + S^\rho{}_{\rho\nu}) \left(\frac{1}{H^2} \delta_\lambda H + \frac{1}{H} S^\sigma{}_{\sigma\lambda} \right)$$

1. Unified field theories
2. A scalar field equation for point particle Hamiltonians
- 3. From the scalar equation on phase space to the Einstein-Maxwell equations**
4. Conclusion



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$

- The Hamilton function

$$H(x, p) = g^{\mu\nu}(x) p_\mu p_\nu$$

- x and p derivatives

$$\partial_\mu = \frac{\partial}{\partial x^\mu} \quad \bar{\partial}^\mu = \frac{\partial}{\partial p_\mu}$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = 0, \quad Q_{\mu\nu} = 0$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



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$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

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$$\delta_\mu H = 0, \quad Q_{\mu\nu} = 0$$

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} \right) - \frac{\mathcal{R}}{H^2} \right) = 0$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



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$$\partial_\mu = \frac{\partial}{\partial x^\mu} \quad \bar{\partial}^\mu = \frac{\partial}{\partial p_\mu}$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = 0, \quad Q_{\mu\nu} = 0$$

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} \right) - \frac{\mathcal{R}}{H^2} \right) = 0$$

\Leftrightarrow

$$-\frac{2}{H^2} \dot{R} + \frac{4}{H^3} \dot{R}^{\mu\nu} p_\mu p_\nu = 0$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



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- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = 0, \quad Q_{\mu\nu} = 0$$

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} \right) - \frac{\mathcal{R}}{H^2} \right) = 0$$

\Leftrightarrow

$$-\frac{2}{H^2} \dot{R} + \frac{4}{H^3} \dot{R}^{\mu\nu} p_\mu p_\nu = 0$$

\Leftrightarrow

$$4 \left(\frac{1}{2} g^{\mu\nu} \dot{R} - \dot{R}^{\mu\nu} \right) p_\mu p_\nu = 0$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



- The Hamilton function

$$H(x, p) = g^{\mu\nu}(x) p_\mu p_\nu$$

- x and p derivatives

$$\partial_\mu = \frac{\partial}{\partial x^\mu} \quad \bar{\partial}^\mu = \frac{\partial}{\partial p_\mu}$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = 0, \quad Q_{\mu\nu} = 0$$

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} \right) - \frac{\mathcal{R}}{H^2} \right) = 0$$

\Leftrightarrow

$$-\frac{2}{H^2} \overset{\circ}{R} + \frac{4}{H^3} \overset{\circ}{R}{}^{\mu\nu} p_\mu p_\nu = 0$$

\Leftrightarrow

$$4 \left(\frac{1}{2} g^{\mu\nu} \overset{\circ}{R} - \overset{\circ}{R}{}^{\mu\nu} \right) p_\mu p_\nu = 0$$

The Einstein vacuum equations from a scalar field equation on the tangent bundle



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$

- The Hamilton function

$$\begin{aligned} H(x, p) &= g^{\mu\nu}(x) (p_\mu - eA_\mu) (p_\nu - eA_\nu) \\ &= g^{\mu\nu}(x) \tilde{p}_\mu \tilde{p}_\nu \end{aligned}$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = -e F_{\mu\nu} \tilde{p}^\nu$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



- The Hamilton function

$$H(x, p) = g^{\mu\nu}(x) (p_\mu - eA_\mu) (p_\nu - eA_\nu) \\ = g^{\mu\nu}(x) \tilde{p}_\mu \tilde{p}_\nu$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = -eF_{\mu\nu} \tilde{p}^\nu$$

$$0 = -\frac{2\alpha_1(H\dot{R} - 2\dot{R}_{\mu\nu}\tilde{p}^\mu\tilde{p}^\nu)}{H^3} \\ + \frac{\alpha_2\kappa e(eF_{\mu\nu}F^{\mu\nu}H + 3\tilde{p}^\mu(-2eF_\mu{}^\sigma F_{\nu\sigma}\tilde{p}^\nu + H\dot{\nabla}_\nu F_\mu{}^\nu))}{H^3} \\ + \frac{\alpha_3\kappa e(-3eF_{\mu\nu}F^{\mu\nu}H + 8\tilde{p}^\mu(2eF_\mu{}^\sigma F_{\nu\sigma}\tilde{p}^\nu - H\dot{\nabla}_\nu F_\mu{}^\nu))}{H^3}$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



- The Hamilton function

$$H(x, p) = g^{\mu\nu}(x) (p_\mu - eA_\mu) (p_\nu - eA_\nu)$$

$$= g^{\mu\nu}(x) \tilde{p}_\mu \tilde{p}_\nu$$

- the Hamilton metric

$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

$$\delta_\mu H = -e F_{\mu\nu} \tilde{p}^\nu$$

$$0 = \left[-2\alpha_1 (g^{\mu\nu} \dot{R} - 2\dot{R}^{\mu\nu}) + \kappa e^2 \left(2(8\alpha_3 - 3\alpha_2) F^{\mu\sigma} F^\nu{}_\sigma + (\alpha_2 - 3\alpha_3) F_{\sigma\rho} F^{\sigma\rho} g^{\mu\nu} \right) \right] \tilde{p}_\mu \tilde{p}_\nu$$

$$+ \left[\kappa e (3\alpha_2 - 8\alpha_3) g^{\mu\nu} \dot{\nabla}_\rho F^{\sigma\rho} \right] \tilde{p}_\mu \tilde{p}_\nu \tilde{p}_\sigma$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

The scalar field equation

$$\alpha_1 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\mathcal{R}_{\mu\nu}}{H} + Q^\rho{}_\rho g_{\mu\nu} - Q_{\mu\nu} \right) - \frac{\mathcal{R}}{H^2} \right) + \kappa \left\{ \alpha_2 \left(\frac{1}{2} \bar{\partial}^\mu \bar{\partial}^\nu \left(\frac{\delta_\mu H \delta_\nu H}{H^2} \right) - 2 \frac{g^{\mu\nu} \delta_\mu H \delta_\nu H}{H^3} - 2 \tilde{\nabla}_\mu \left(\frac{g^{\mu\nu} \delta_\nu H}{H^2} \right) \right) + \alpha_3 \left(-\frac{1}{H^2} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) + 2 \tilde{\nabla}_\mu \bar{\partial}^\rho \left(\frac{\bar{\partial}^\mu \delta_\rho H}{H} \right) \right) \right\} + \frac{1}{4} X = 0$$



- The Hamilton function

$$H(x, p) = g^{\mu\nu}(x) (p_\mu - eA_\mu) (p_\nu - eA_\nu)$$

$$= g^{\mu\nu}(x) \tilde{p}_\mu \tilde{p}_\nu$$

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$$g^{\mu\nu}(x, p) = g^{\mu\nu}(x)$$

- Properties

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$$+ \left[\kappa e (3\alpha_2 - 8\alpha_3) g^{\mu\nu} \dot{\nabla}_\rho F^{\sigma\rho} \right] \tilde{p}_\mu \tilde{p}_\nu \tilde{p}_\sigma$$



The action

$$S[H] = \int_{T^*M} d^4x d^4p \left(\alpha_1 \frac{\mathcal{R}}{H} + \kappa \left(\alpha_2 \frac{1}{H^2} g^{\mu\nu} \delta_\mu H \delta_\nu H + \alpha_3 \frac{1}{H} (\bar{\partial}^\rho \delta_\mu H) (\bar{\partial}^\mu \delta_\rho H) \right) \right)$$

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\Leftrightarrow

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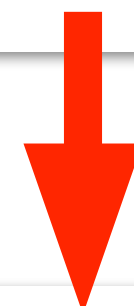


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The Einstein Maxwell equations from a scalar field equation on the tangent bundle

1. Unified field theories
2. A scalar field equation for point particle Hamiltonians
3. From the scalar equation on phase space to the Einstein-Maxwell equations
4. **Conclusion**



**“a theory joining the gravitational and the electromagnetic field into one single hyperfield
whose equations represent the conditions imposed on the geometrical structure of the universe”**
[Tonnalet 1995, H. Goenner 2004]



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The idea

- Identify the **Hamiltonian** as the **hyperfield** that joins (all) bosons
- Construct dynamical equations that determine the Hamiltonian **as scalar field on phase space**
- Break the equation down to **tensorial field equations on spacetime** for bosons tensor field building up the Hamiltonian

Hamilton Geometry of Phase Space determines the geometric structure of the universe.



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The realisation

- An action principle for a scalar on the tangents bundle

$$S[H] = \int_{T^*M} \Sigma \mathcal{L}(H, \partial_x H, \partial_p H, \partial_x \partial_x H, \partial_p \partial_x H, \partial_p \partial_p H, \dots)$$

Variation yields non linear, geometric scalar field equation for H

- Decompose field equation into n-th order powers of p

$$\sum_{i=0}^n G^{\mu_1 \dots \mu_n}(x) p_{\mu_1} p_{\mu_n} = 0$$

- Each power is induced tensorial field equation on spacetime

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Einstein Vacuum Equations

- $H = g^{\mu\nu}(x) p_\mu p_\nu$
- ($n = 2$) $G^{\mu\nu}(x) = 0 \Leftrightarrow \overset{\circ}{R}{}^{\mu\nu} = 0$

Einstein Maxwell Equations

- $H = g^{\mu\nu}(x) (p_\mu - eA_\mu)(p_\nu - eA_\nu)$
- ($n = 3$) $G^{\mu\nu\rho}(x) = 0 \Leftrightarrow \nabla_\mu F^{\mu\nu} = 0$

$$(n = 2) \quad G^{\mu\nu} = 0 \Leftrightarrow$$

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The next step: General polynomial Hamiltonians

$$H = V(x) + A^\mu(x) p_\mu + g^{\mu\nu}(x) p_\mu p_\nu + \sum_{i=3}^n \Psi^{\mu_1 \dots \mu_i}(x) p_{\mu_1} \dots p_{\mu_i}$$



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Scalar field equation on phase space yield coupled non-linear tensorial equations for higher spin fields



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Thank you for your attention!

Aitäh :)



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