

Bunched-beam transverse coherent direct space charge modes

Space Charge 2026

7th **ICFA** Beam Dynamics Mini-Workshop

Workshop Chairs:

Adrian Oeftiger, Oxford
Shinji Machida, RAL

30 September - 2 October

Lindemann Theatre, Physics Department
University of Oxford

Programme Committee:

Andreas Adelman, PSI
Michael Blaskiewicz, BNL
Nicholas Evans, ORNL
Giuliano Franchetti, GSI
Elias Métral, CERN
Hiroki Okuno, RIKEN
Thomas Planche, TRIUMF
Ji Qiang, LBNL
Kiersten Rulsard, ORNL
Yoichi Sato, JPARC/KEK
Guodong Shen, IMP/HIAF
Giulio Stancari, FNAL
Rob Williamson, RAL
Yaoshuo Yuan, CSNS

Accommodation at Trinity College



IOP Institute of Physics
Particle Accelerators
and Beams Group



Workshop Administration:

Julieta Estremadoyro-Vermejo
Kim Proudfoot



TMCI with space charge for a Gaussian bunch

<https://journals.aps.org/prab/pdf/10.1103/PhysRevAccelBeams.24.060101>

X. BUFFAT *et al.*

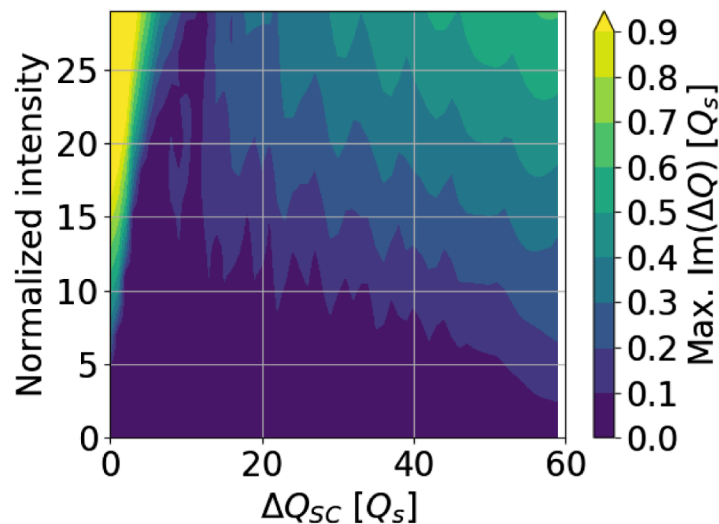


FIG. 8. Maximum imaginary tune shift for a Gaussian distribution obtained with 100 slices and 50 rings for different strengths of the wakefields and the space charge forces.

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Seminal paper from M. Blaskiewicz in 1998

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PHYSICAL REVIEW SPECIAL TOPICS - ACCELERATORS AND BEAMS, VOLUME 1, 044201 (1998)

Fast head-tail instability with space charge

M. Blaskiewicz

Alternating Gradient Synchrotron Department, Brookhaven National Laboratory, Upton, New York 11973-5000
(Received 29 April 1998; published 13 August 1998)

The fast head-tail instability with space charge is studied using series expansion techniques, numerical simulations, and a new formulation which allows for precise estimates of growth rates and thresholds. In regimes where they are reliable, all three techniques predict that space charge suppresses the fast head-tail instability. It is found that the series expansion techniques are unreliable for parameter regimes commonly employed in hadron accelerators. The numerical techniques are less prone to error, but the computational requirements become severe as space charge tune shifts increase. The new model has neither of these problems, but it underestimates the benefits of chromaticity, at least in its simplest form.

Physical model

Physical model

$$\frac{d^2 \tau}{d^2 \theta} + Q_s^2 \tau = 0, \quad (1)$$

$$\begin{aligned} \frac{d^2 x}{d^2 \theta} + Q_x^2 x = & C_{sc} \rho(\theta, \tau) [x - \langle x(\theta, \tau) \rangle] \\ & + \int_{-T}^{\tau} d\tau' W(\tau - \tau') \rho(\theta, \tau') \\ & \times \langle x(\theta, \tau') \rangle d\tau'. \end{aligned} \quad (2)$$

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Without wake field

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Without wake field

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Without wake field

Quadrupolar force

Dipolar force

Physical model

$$\frac{d^2 \tau}{d^2 \theta} + Q_s^2 \tau = 0,$$

Variations in transverse beam size are neglected
=> Linear space charge forces in the transverse direction (like KV distribution)

$$\begin{aligned} \frac{d^2 x}{d^2 \theta} + Q_x^2 x = & C_{sc} \rho(\theta, \tau) [x - \langle x(\theta, \tau) \rangle] \\ & + \int_{-T}^{\tau} d\tau' W(\tau - \tau') \rho(\theta, \tau') \\ & \times \langle x(\theta, \tau') \rangle d\tau'. \end{aligned} \quad (2)$$

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General eigenvalue system

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- ◆ The general eigenvalue system to solve is given by

$$\Delta Q_x a_{p,q} = \sum_{n,m} T_{p,q,n,m} a_{n,m} . \quad (11)$$

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- ◆ N.B.: it has the same form as the one to be solved for instance in my GALACTIC (impedance) Vlasov solver

$$\frac{\Delta Q}{Q_s} a_{kl} = H a_{ij} = \sum_{i,j=-\infty}^{+\infty} H_{kl,ij} a_{ij}$$

Longitudinal distribution

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* **AIR-BAG** bunch in a Square potential well (**ABS**)

*With a simple
analytical model*

Longitudinal distribution

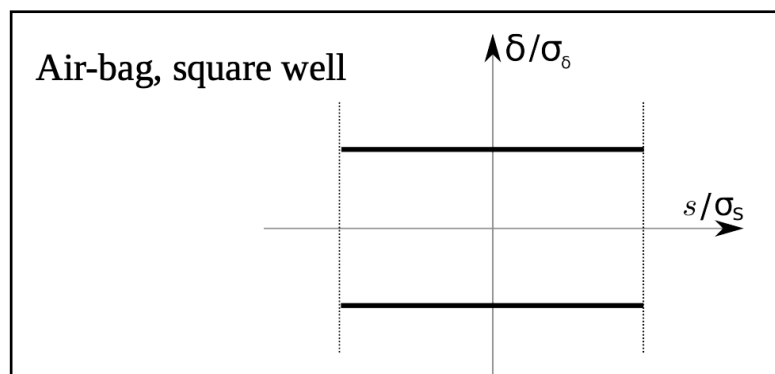
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Longitudinal distribution

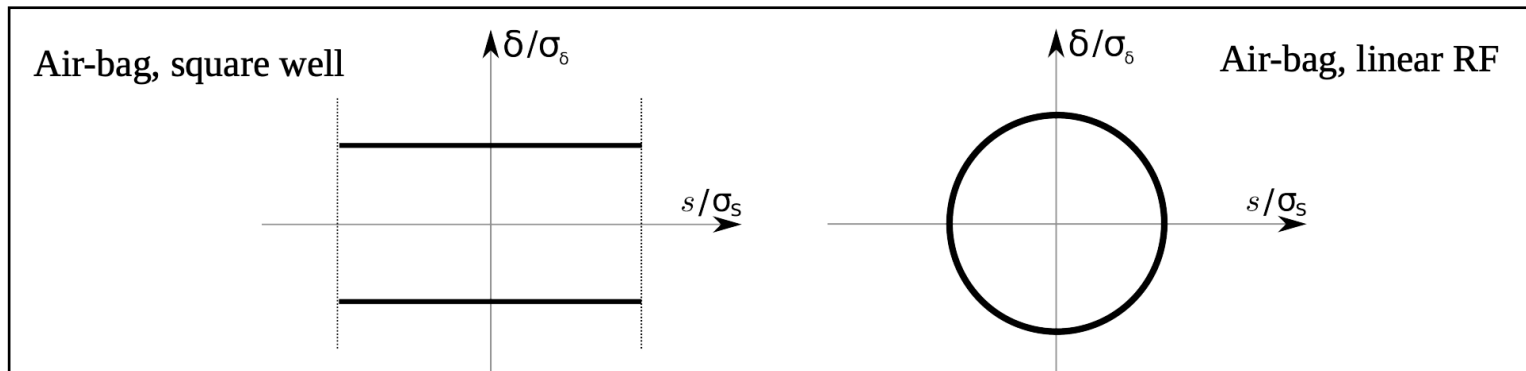
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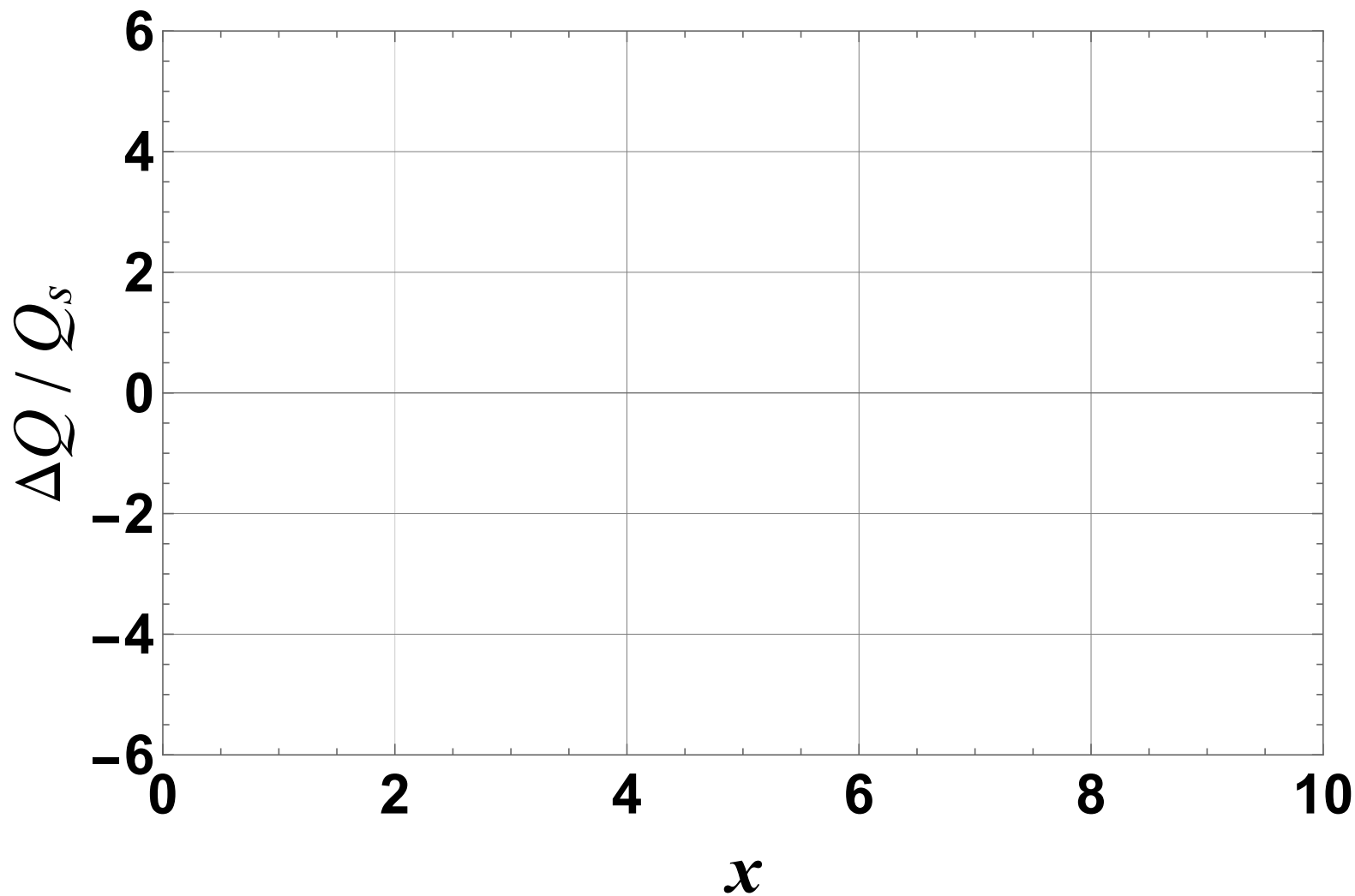
* **GAUSSIAN** line density

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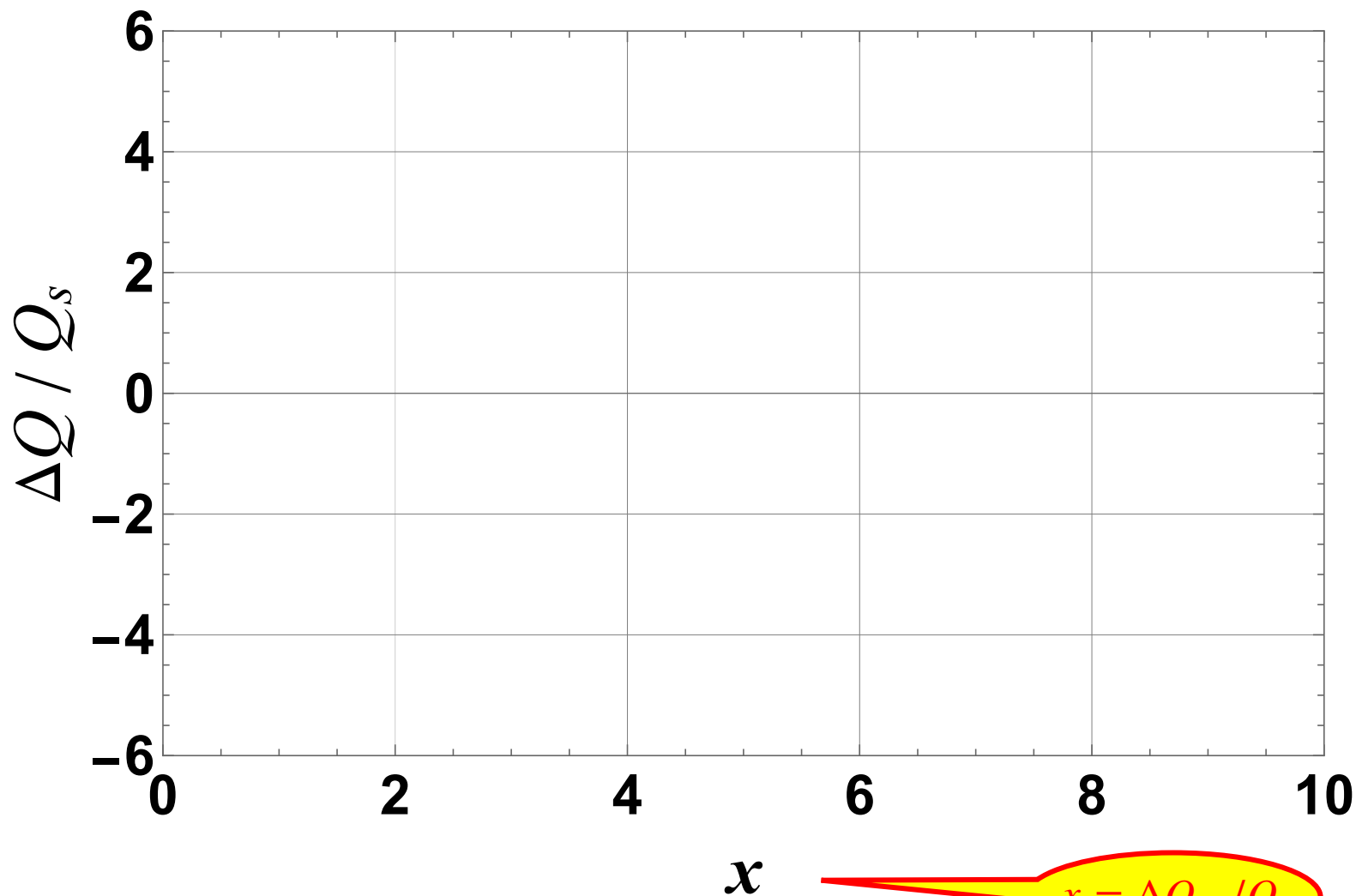
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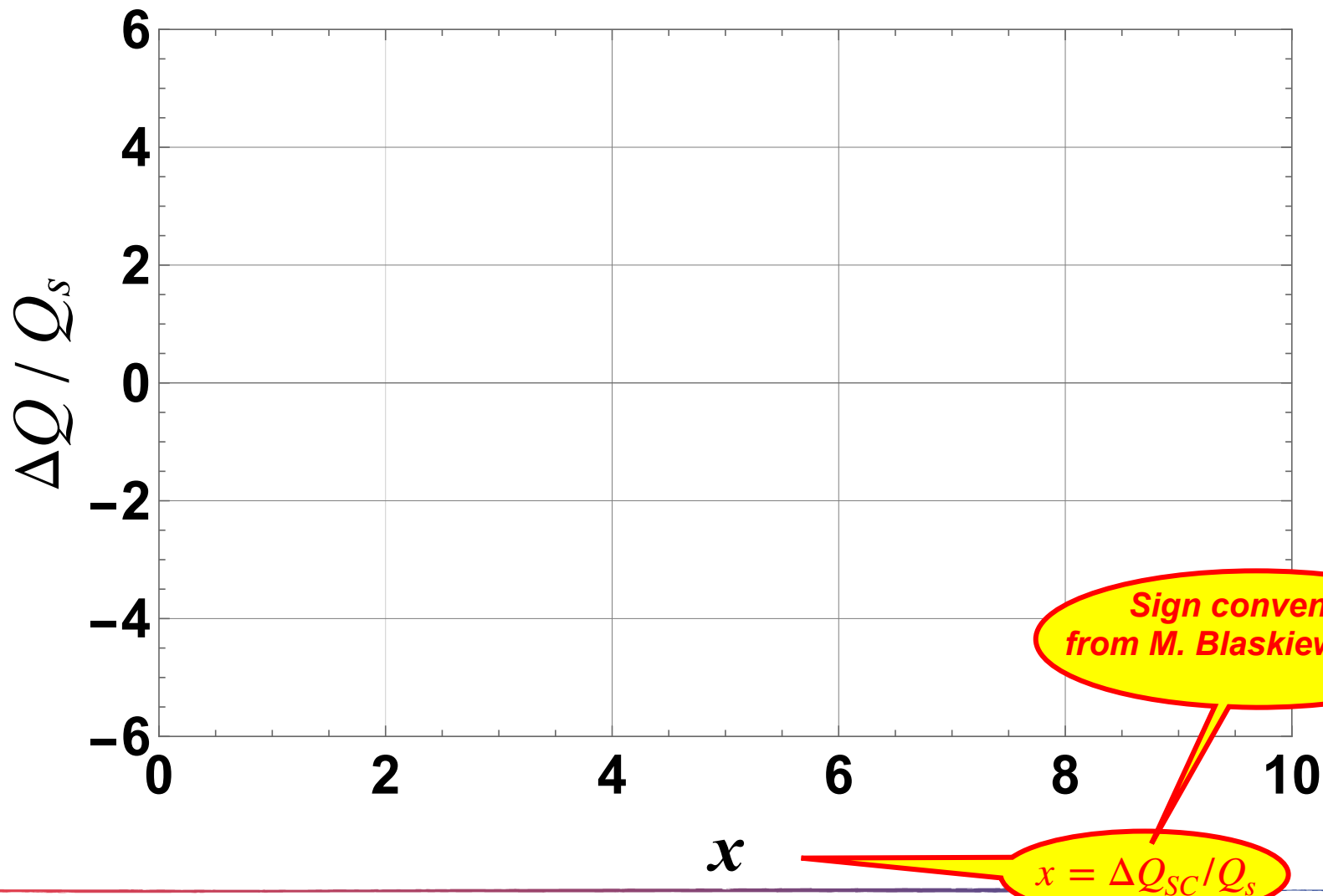
ABS



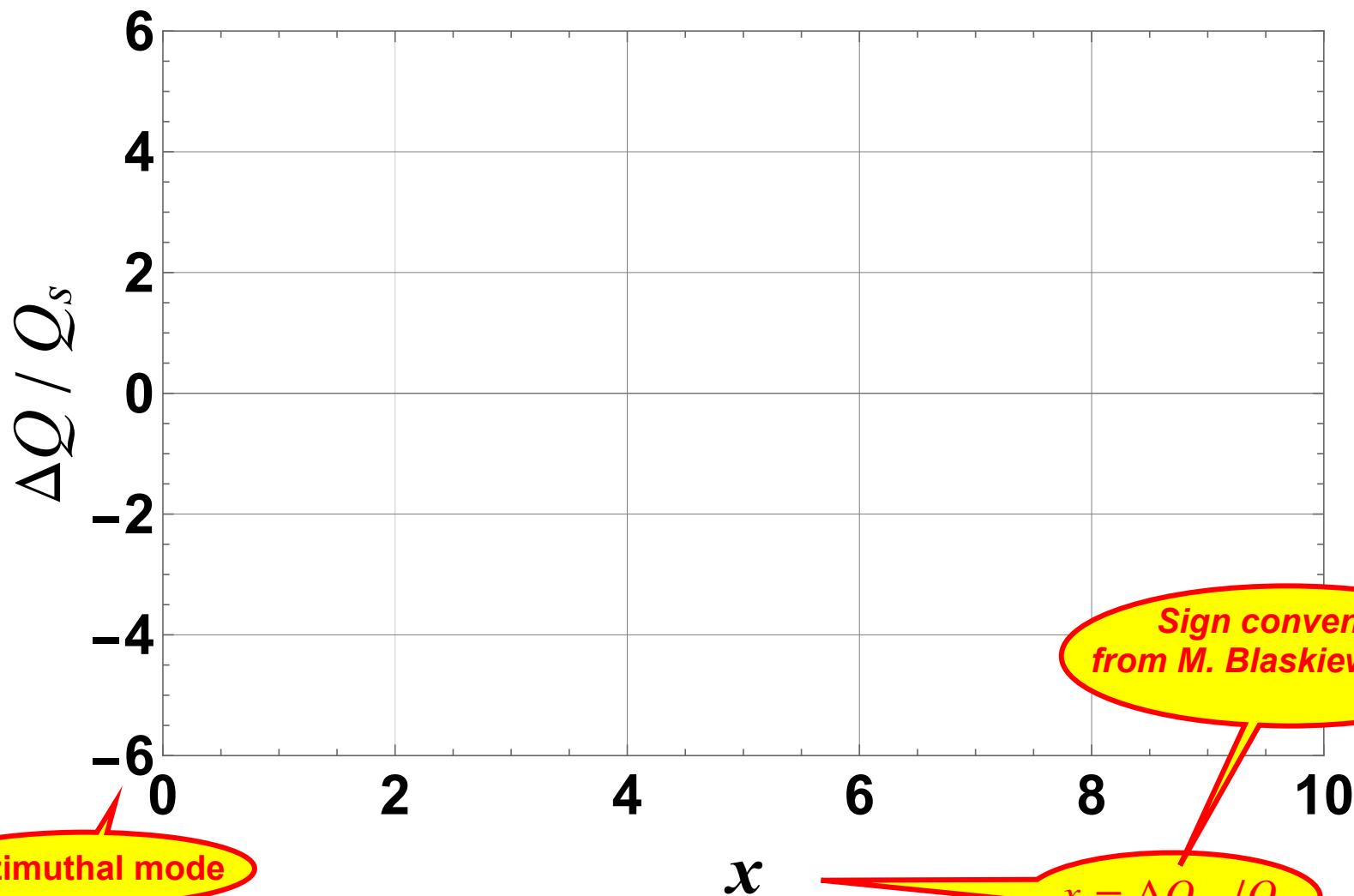
ABS



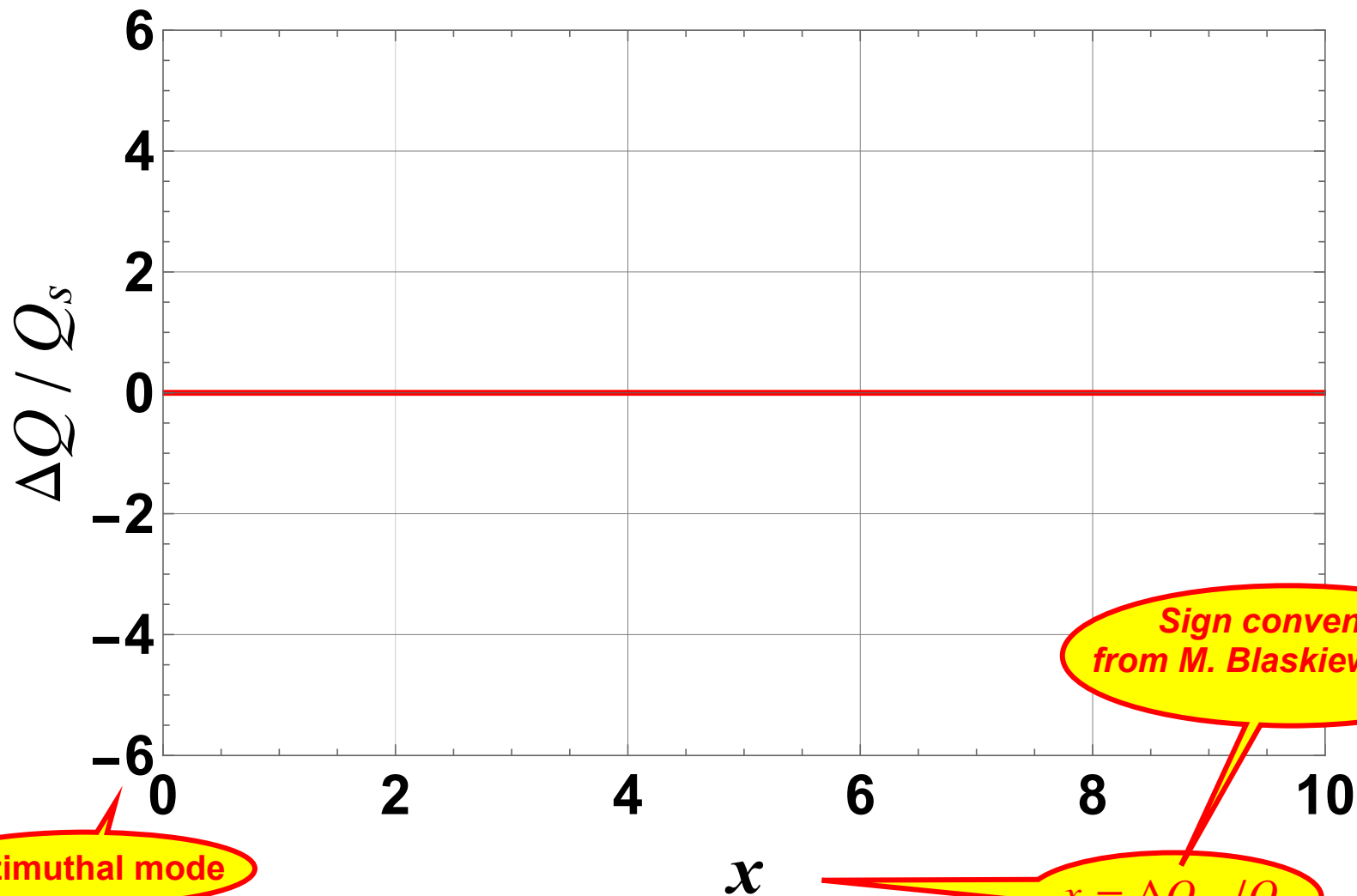
ABS



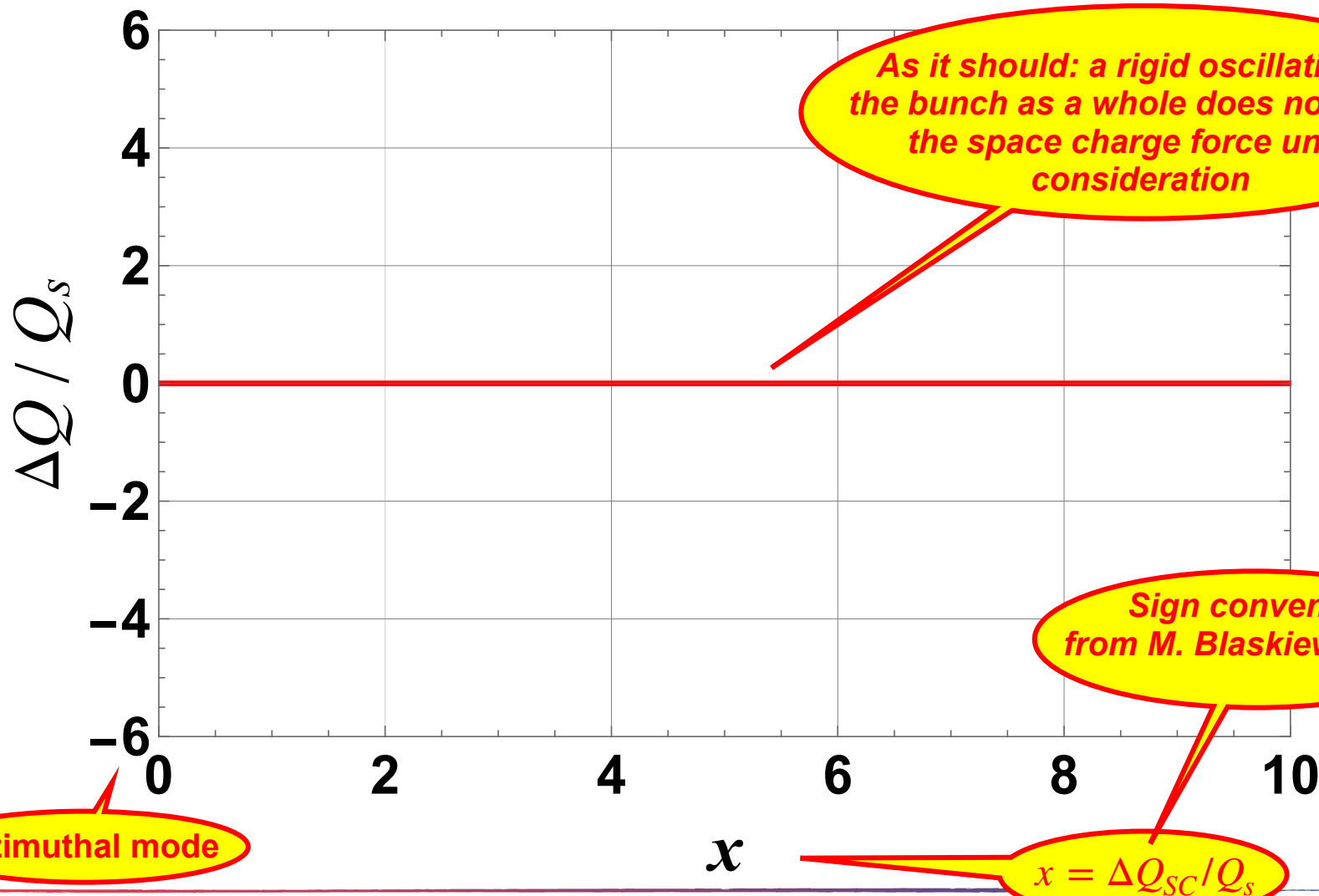
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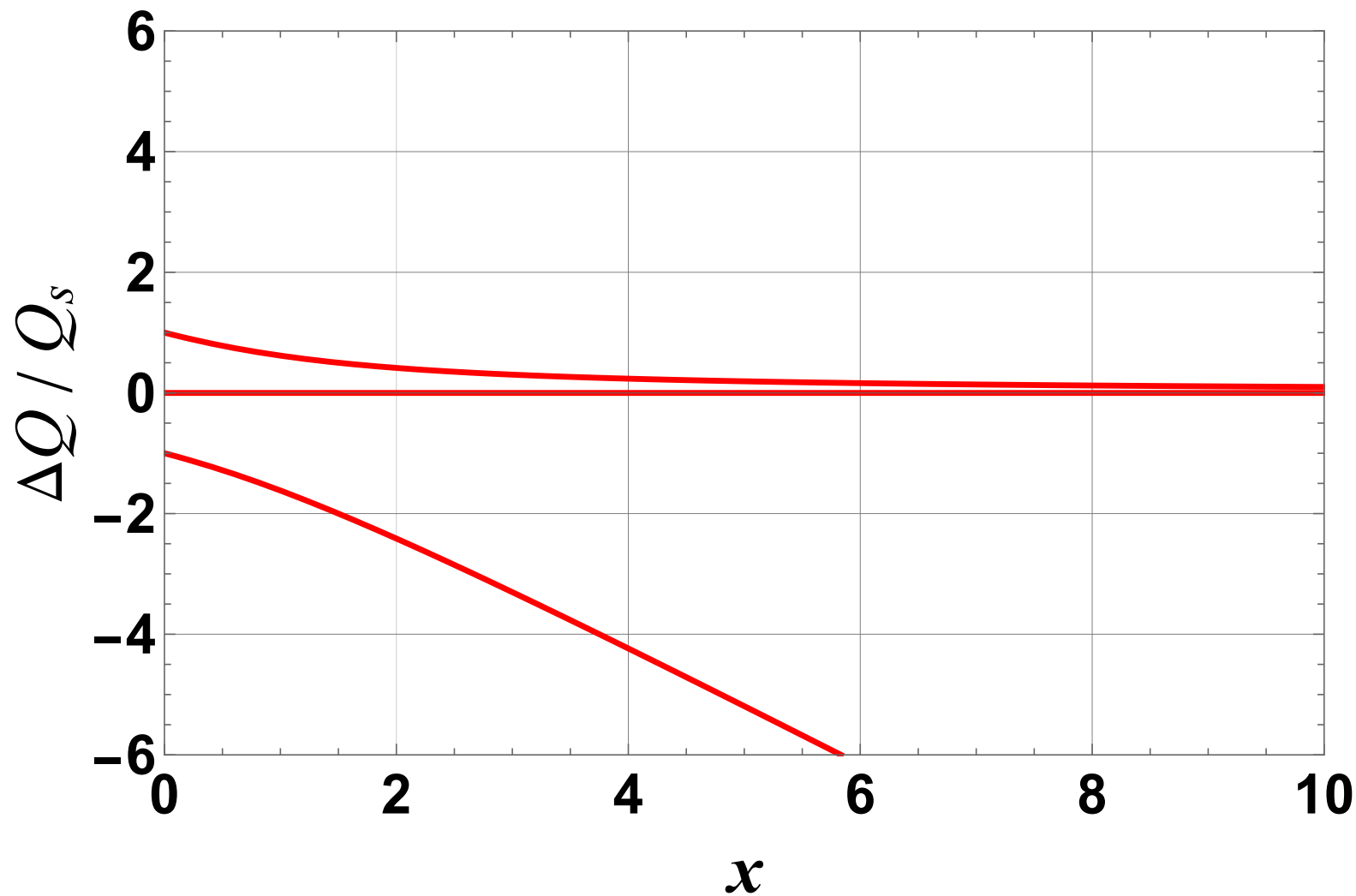
ABS



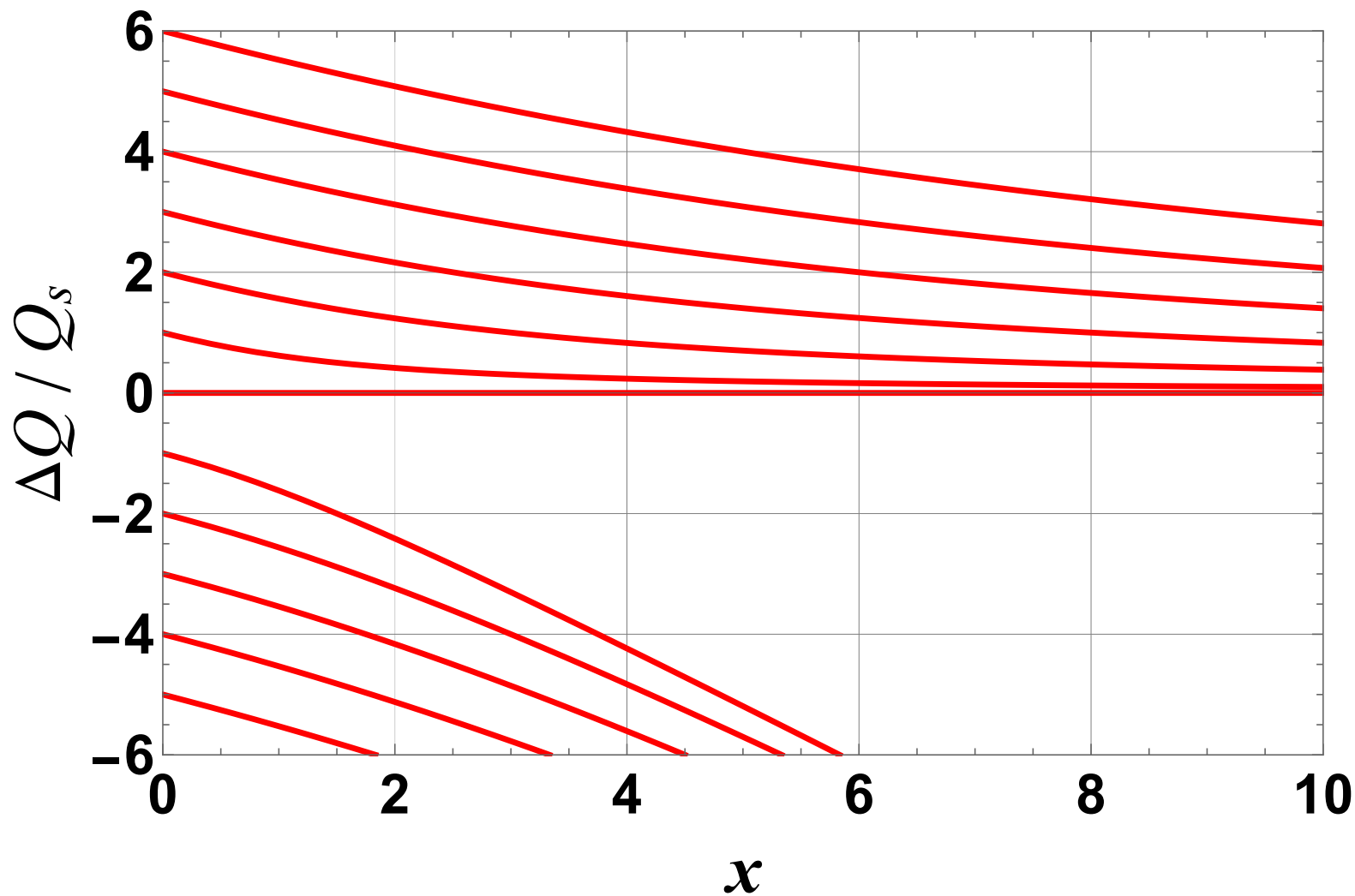
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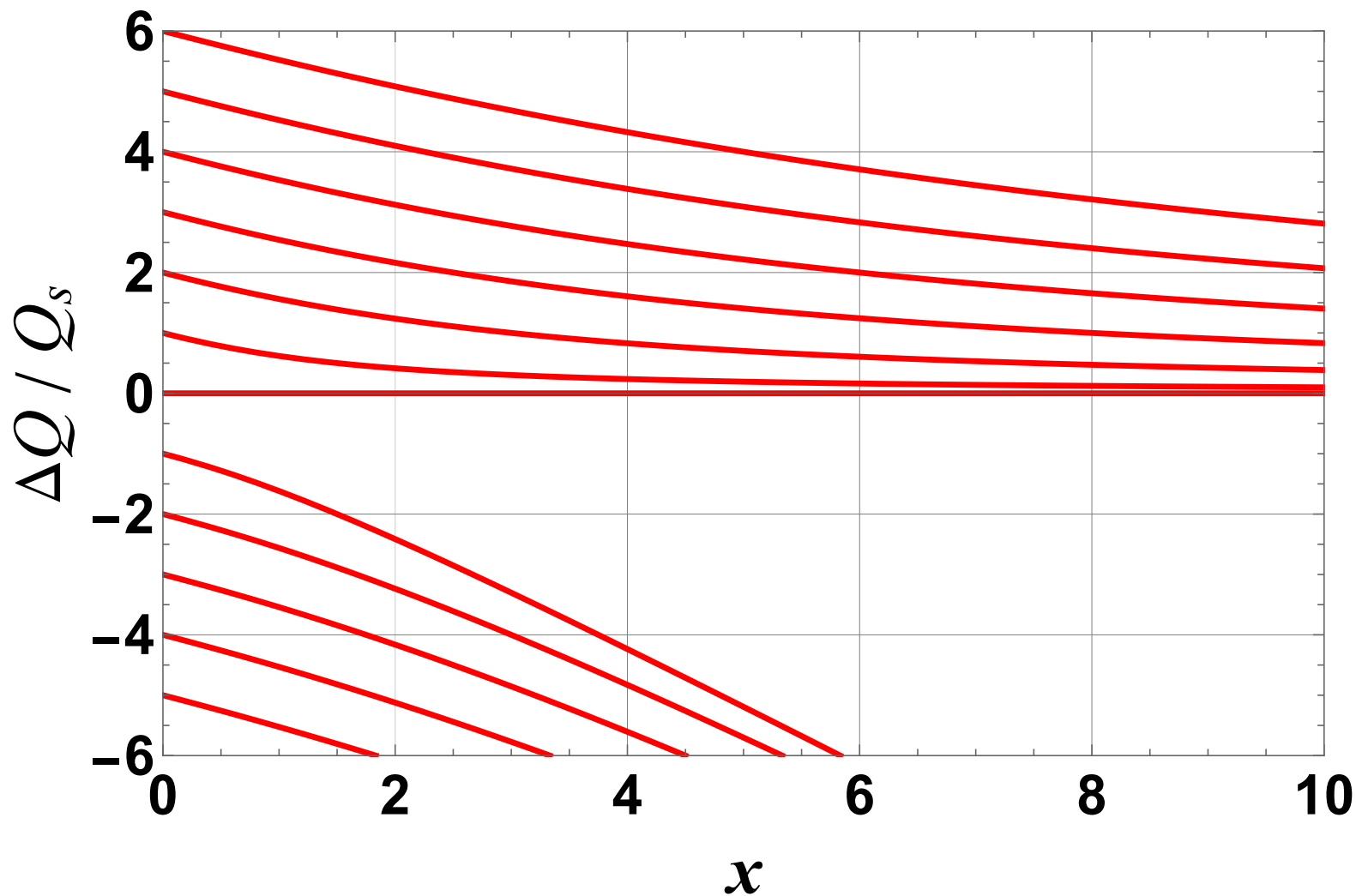
ABS



ABS

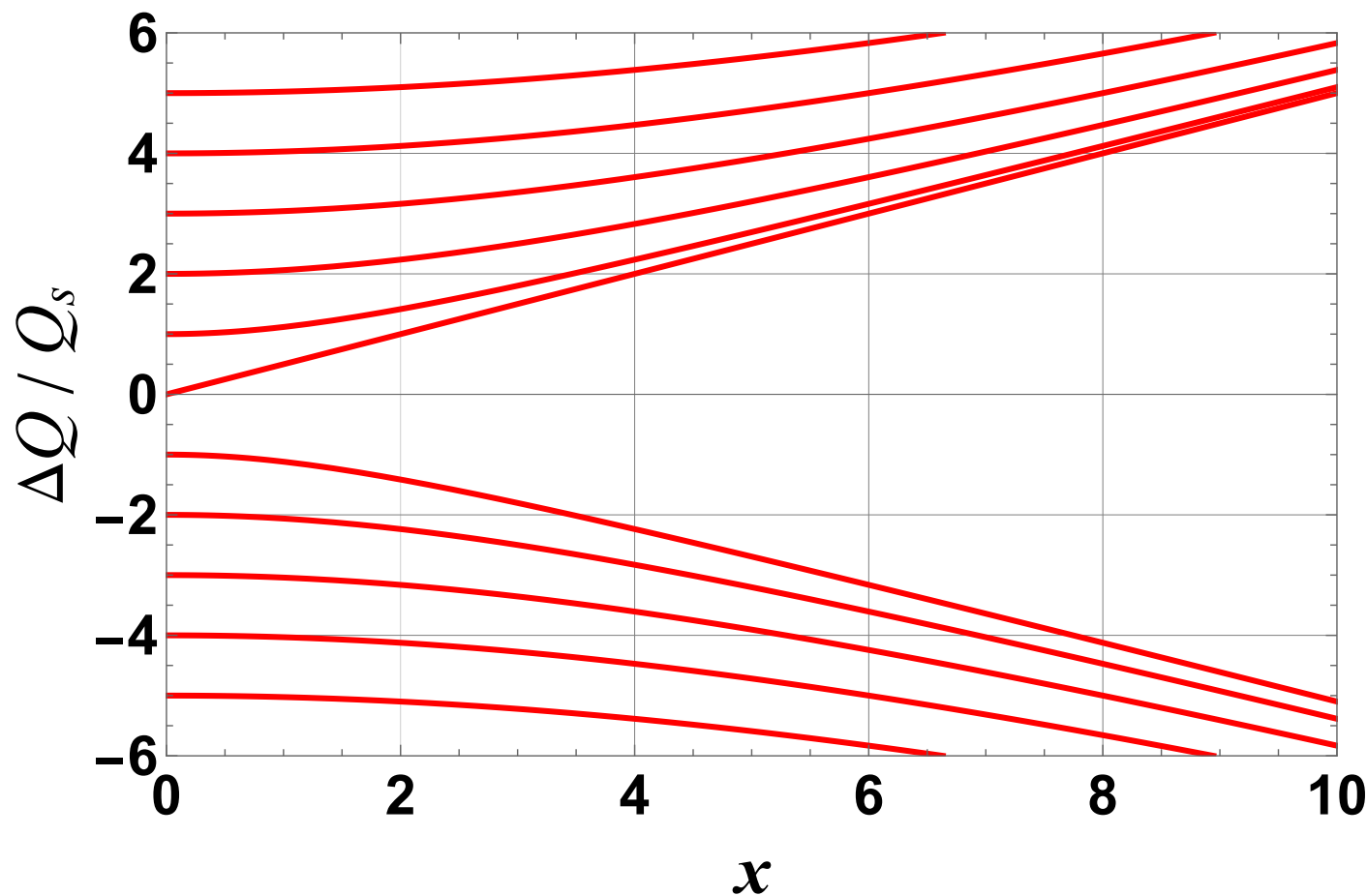


ABS $\Rightarrow$ No radial modes



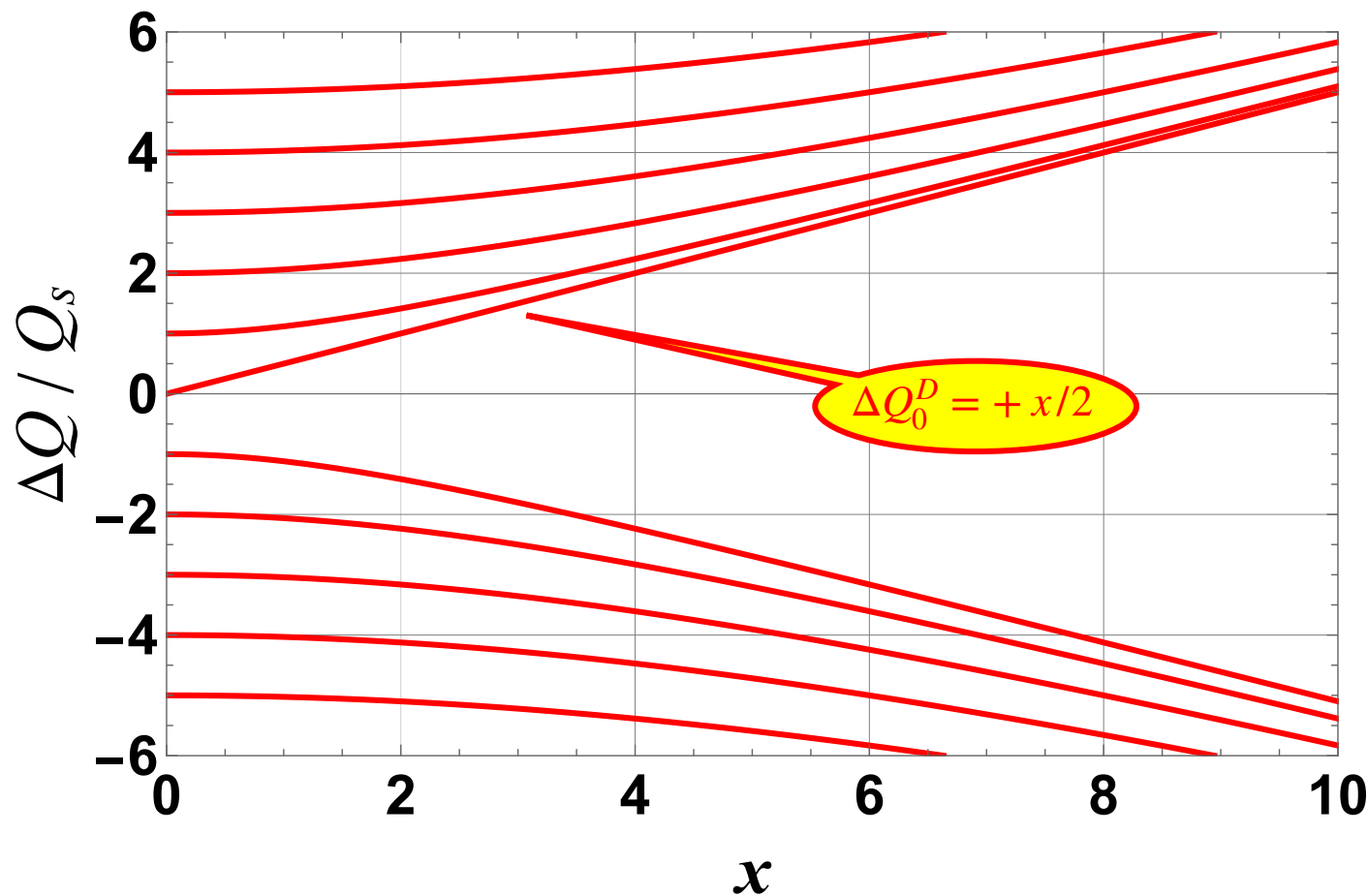
ABS $\Rightarrow$ No radial modes

◆ Dipolar force only



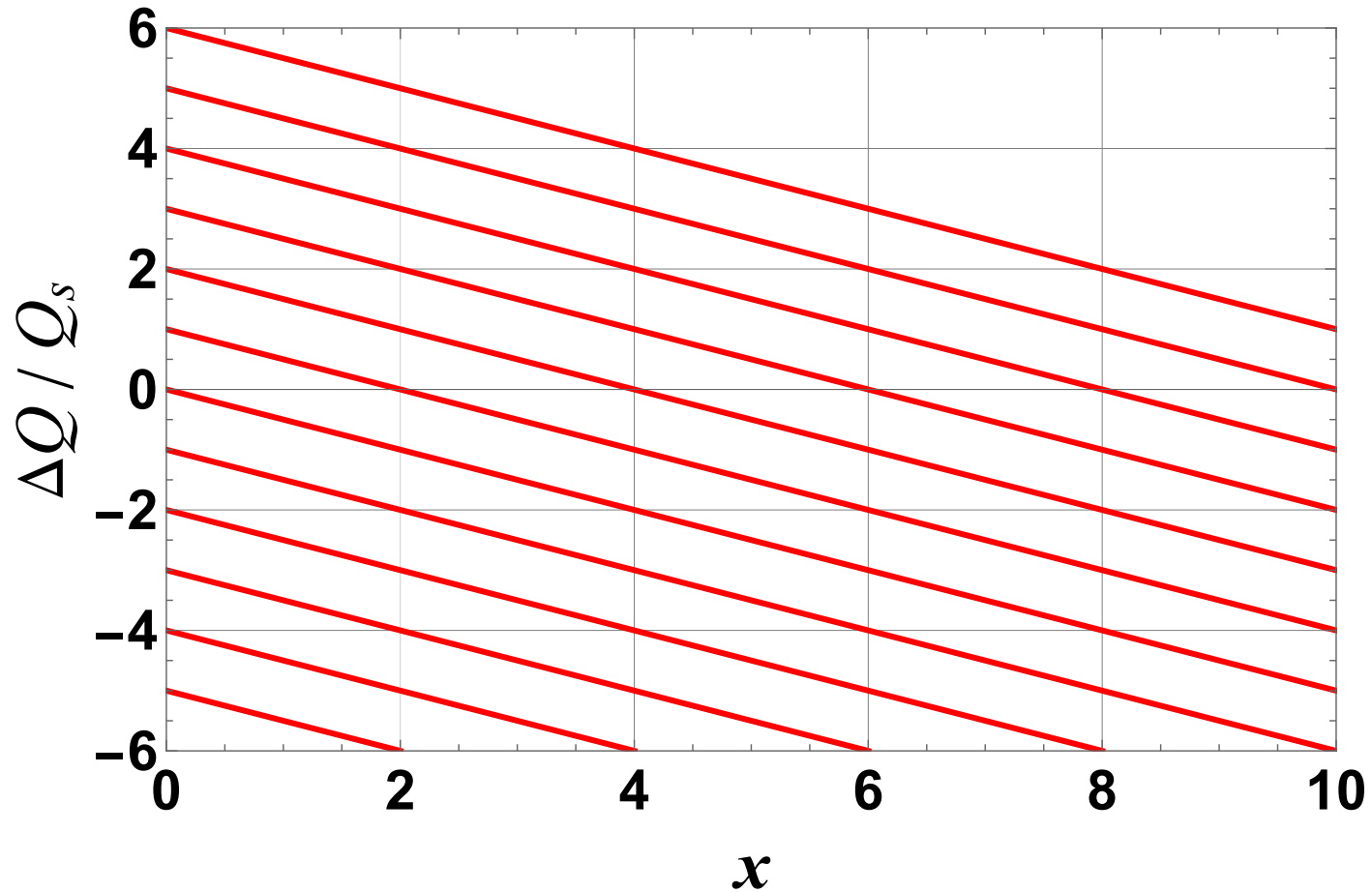
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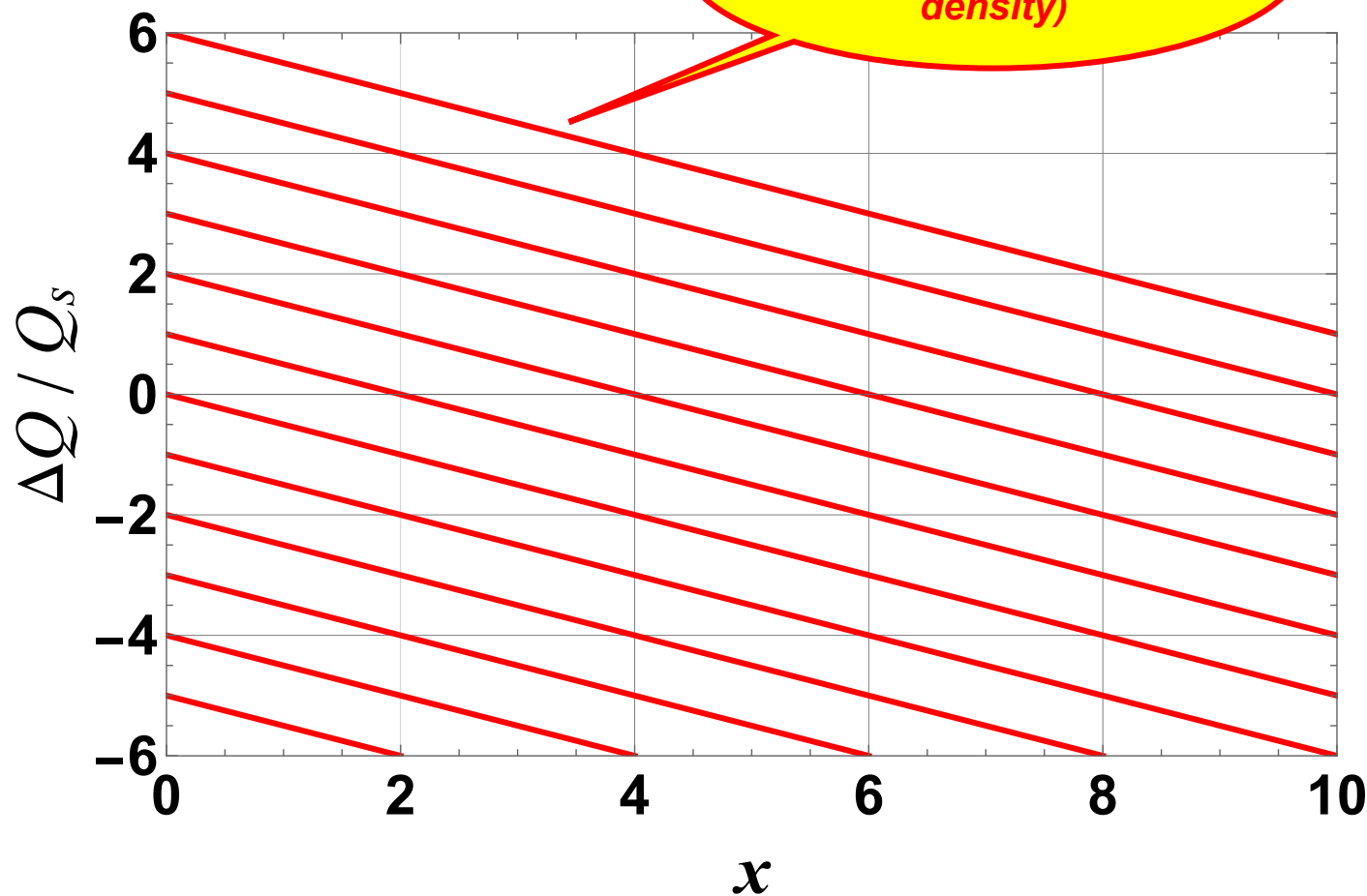
ABS => No radial modes

◆ Quadrupolar force only



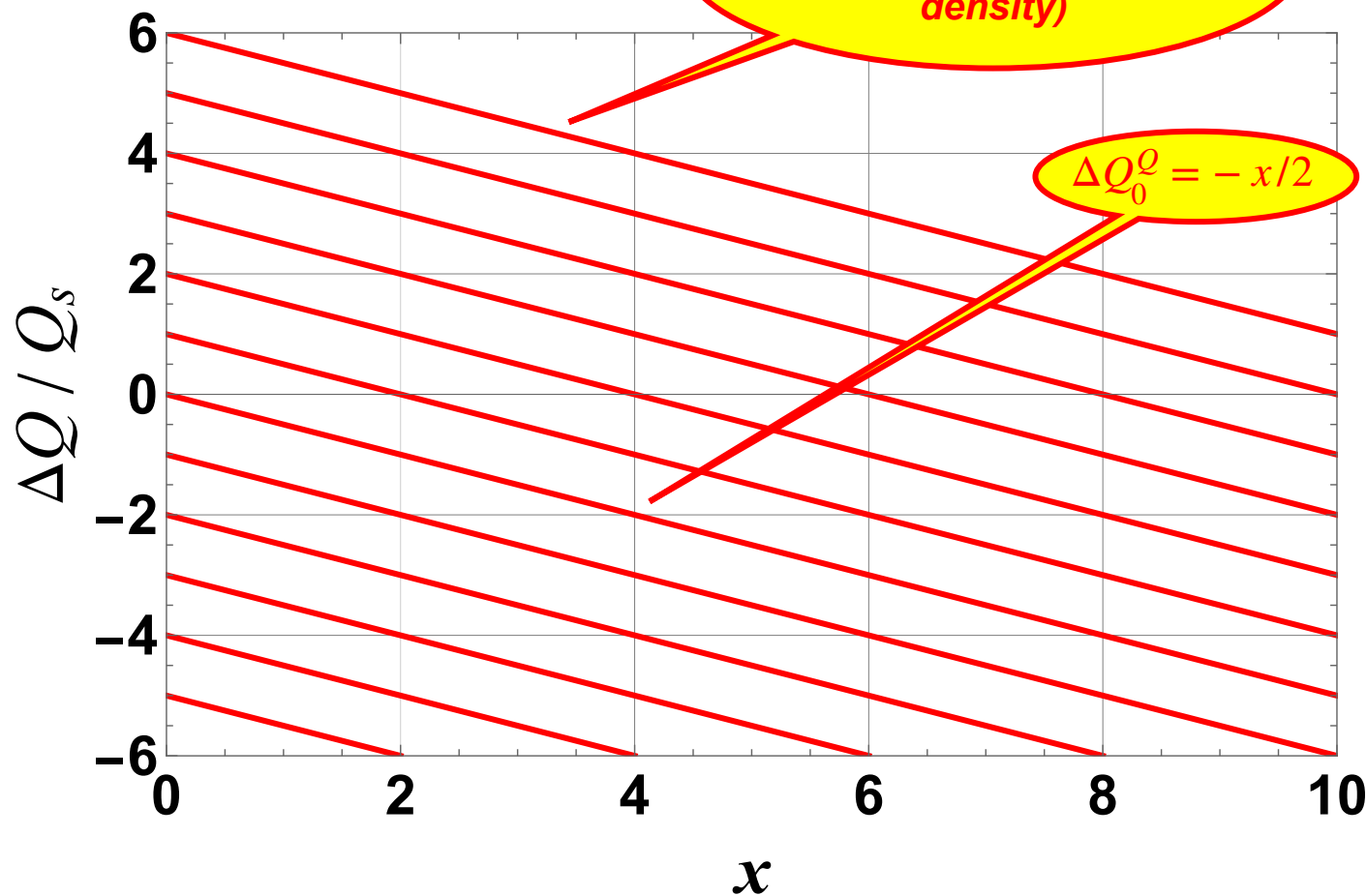
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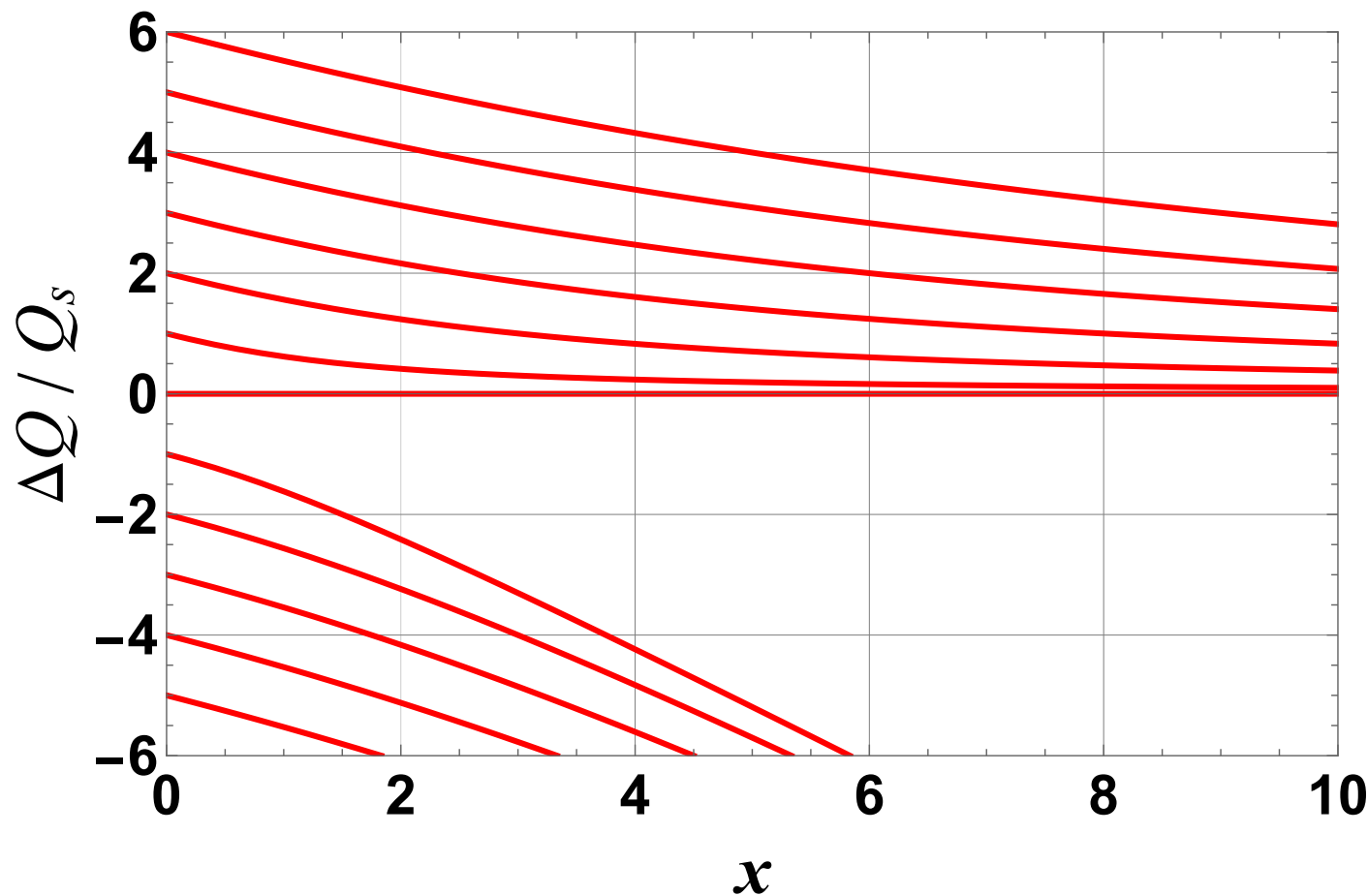
ABS $\Rightarrow$ No radial modes

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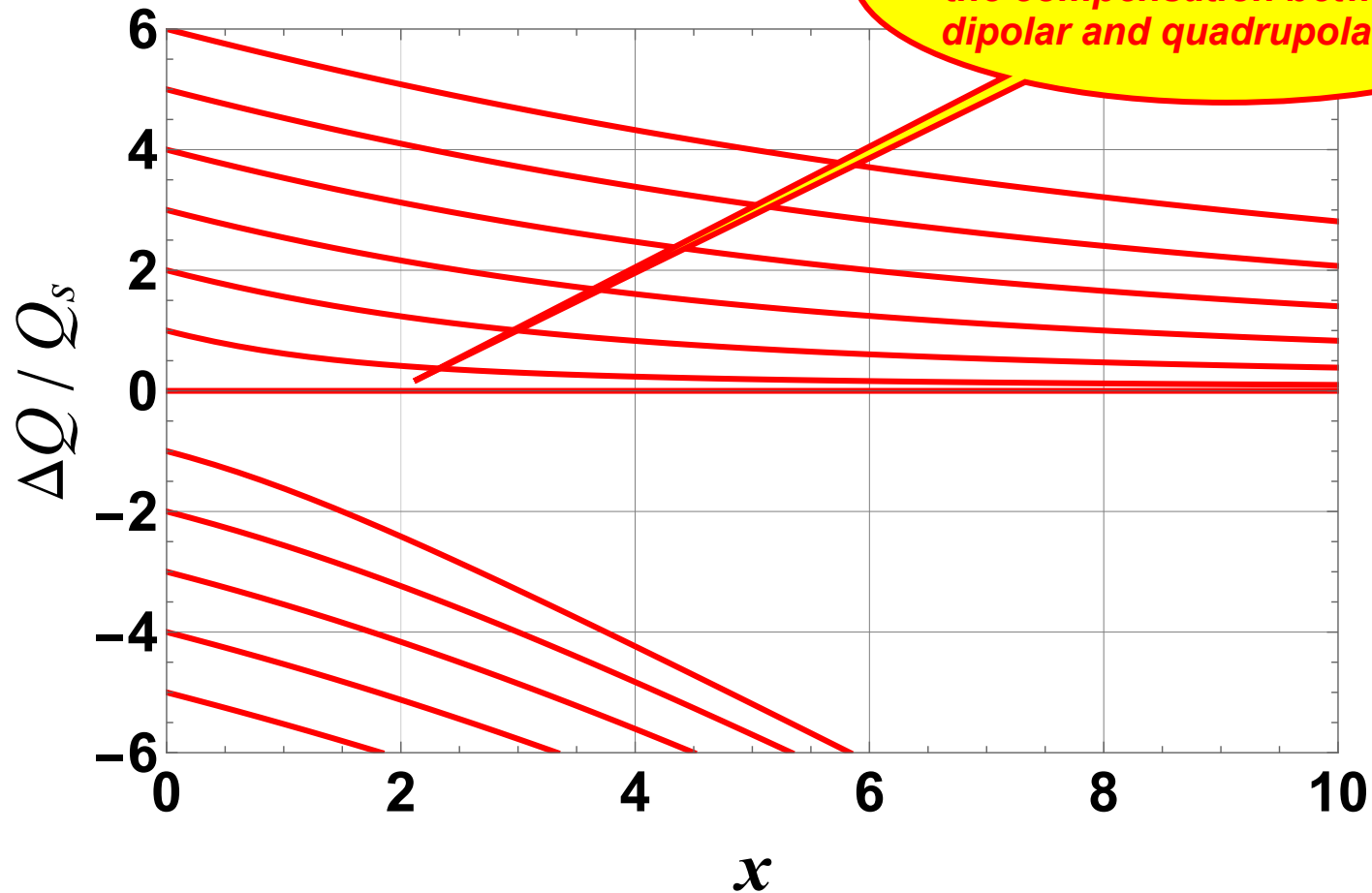
ABS => No radial modes

◆ Sum of the 2 forces



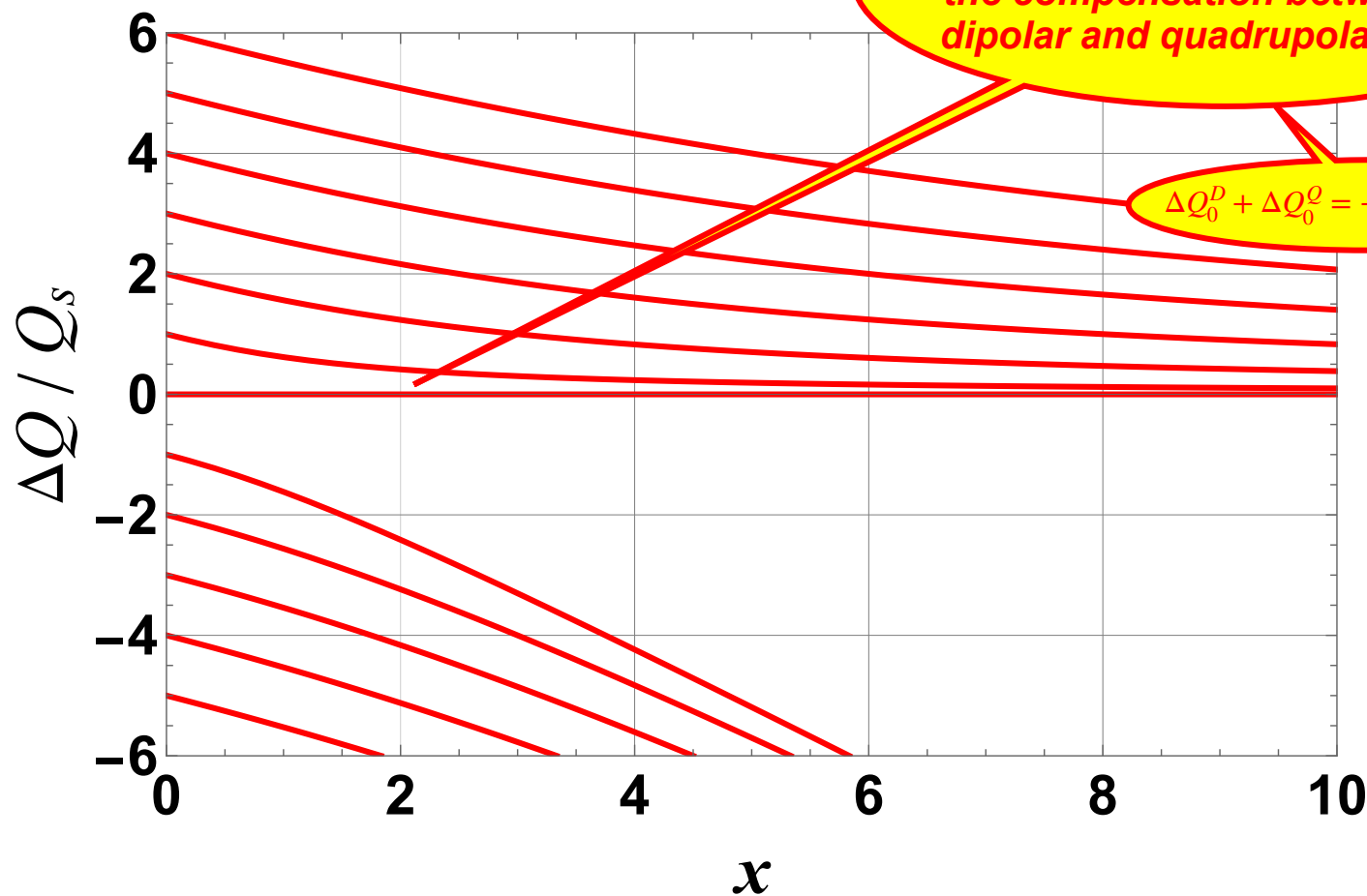
ABS => No radial modes

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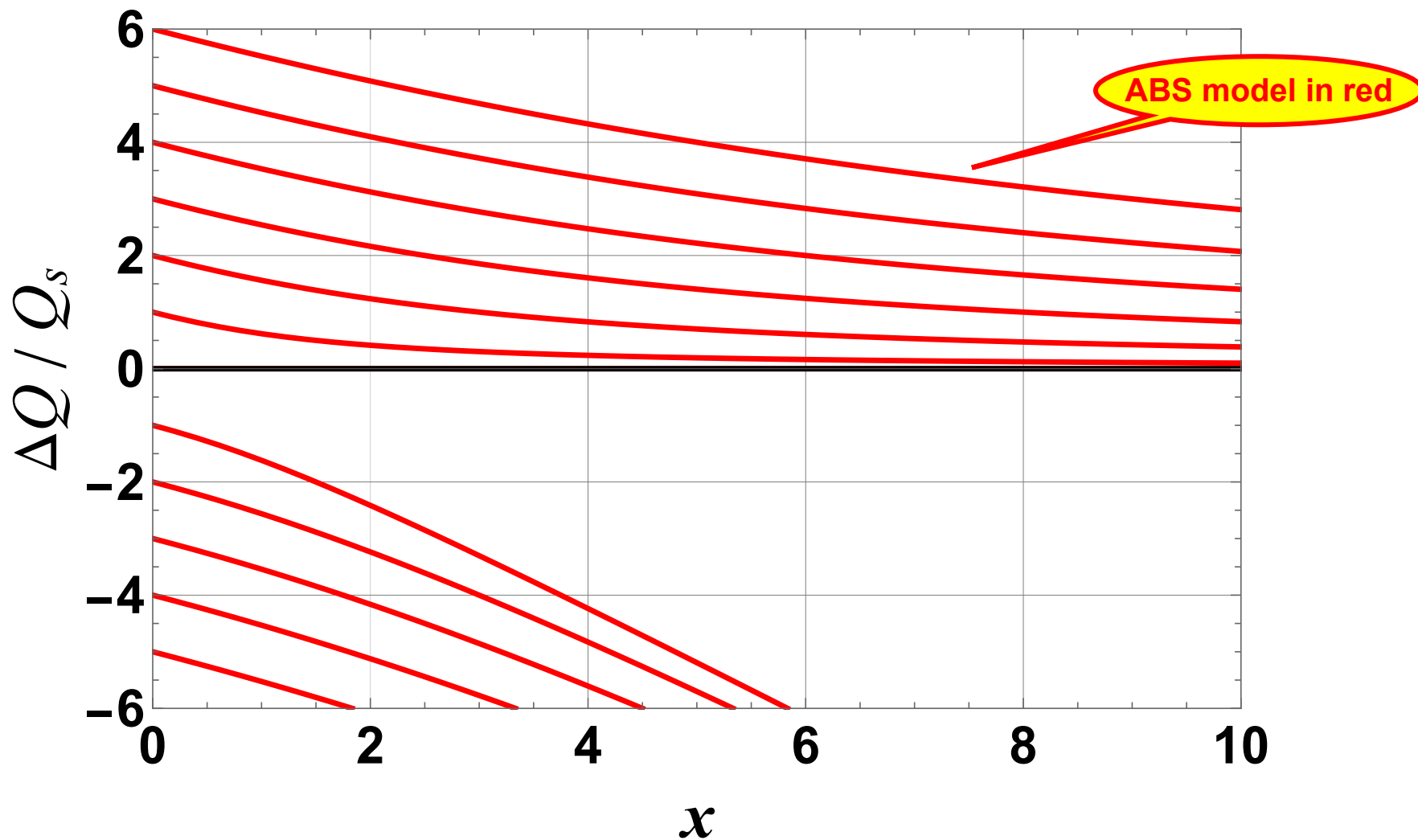


ABS => No radial modes

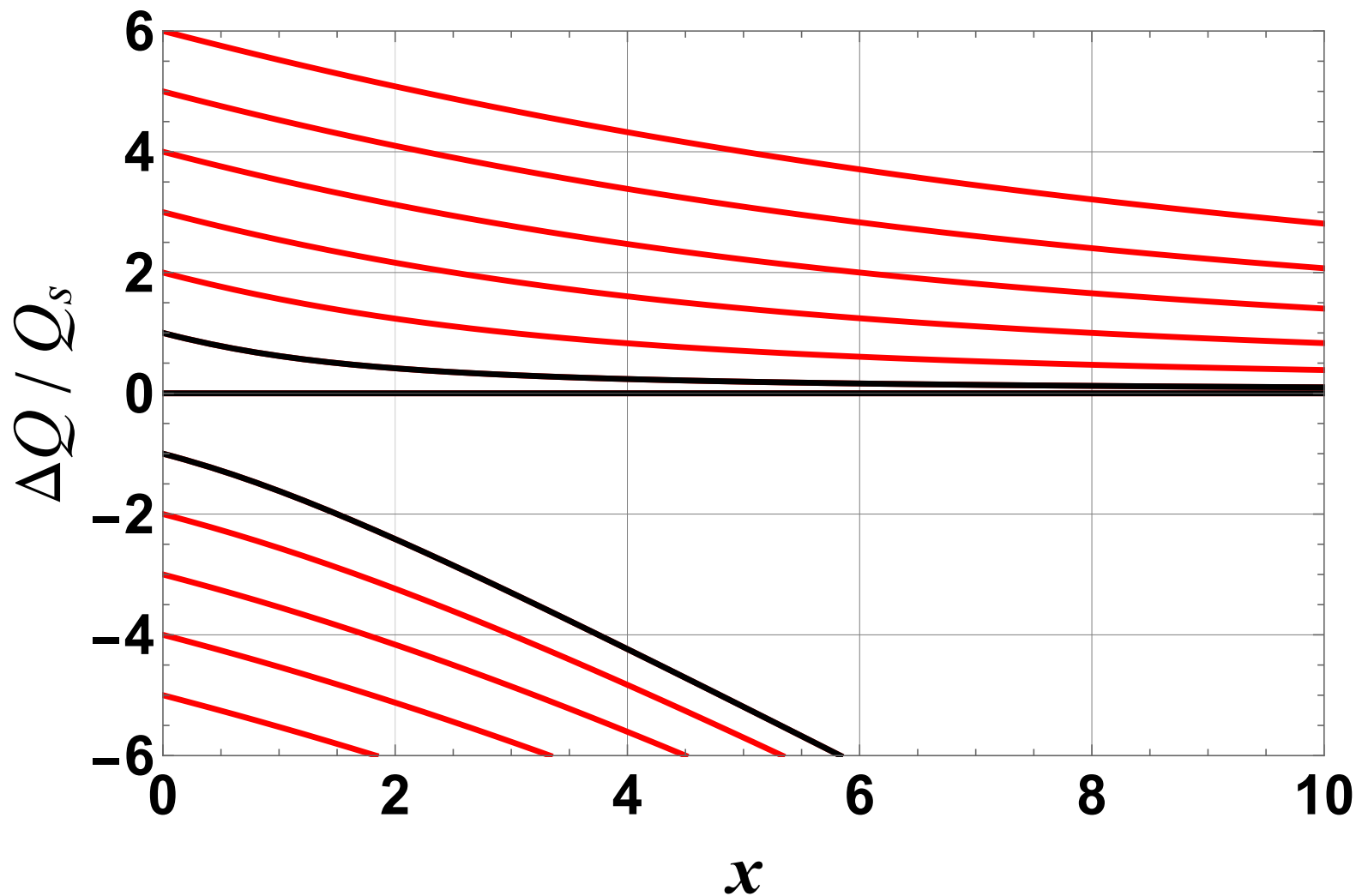
◆ Sum of the 2 forces



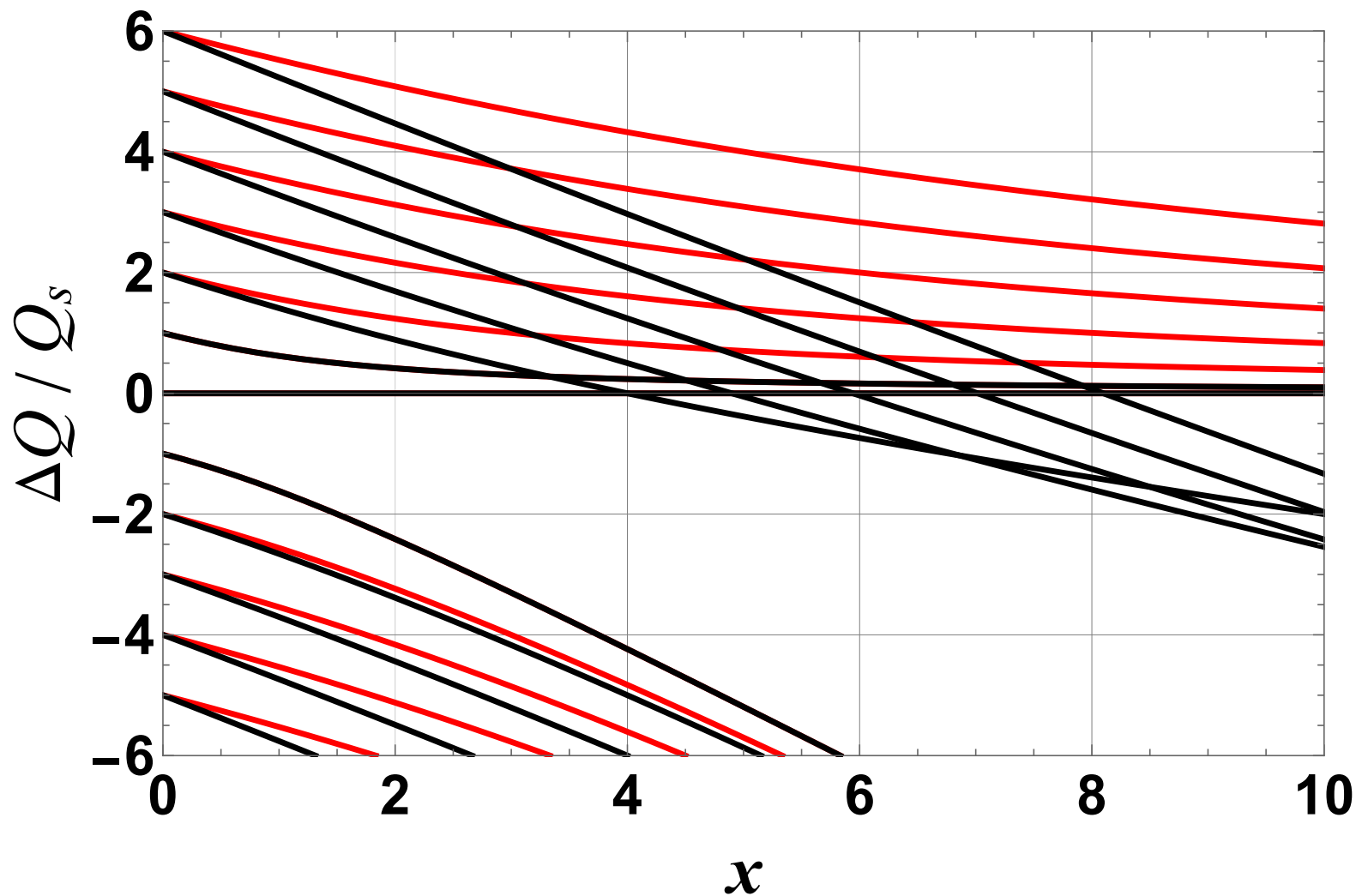
BOXCAR: radial mode = 0



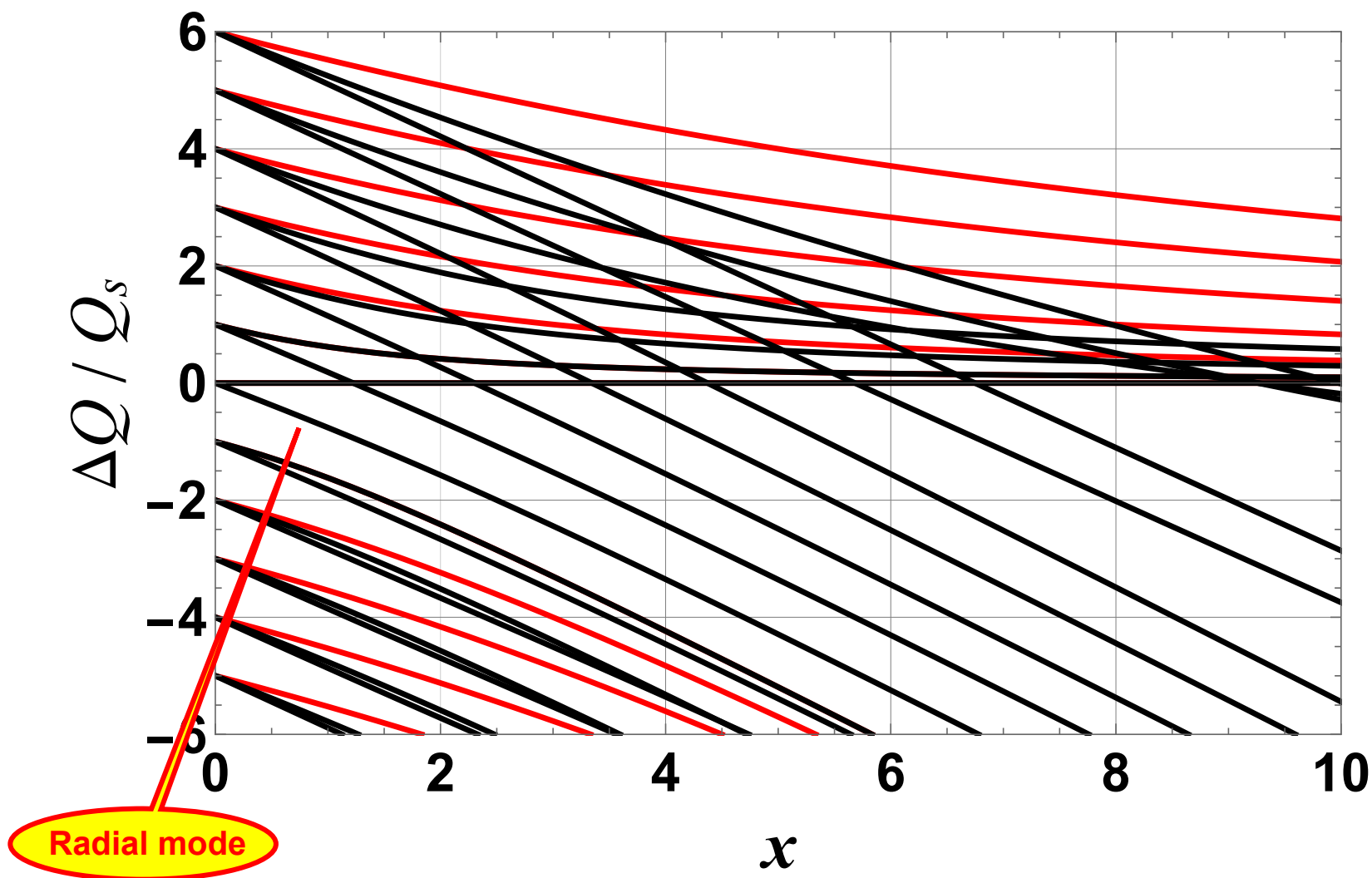
BOXCAR: radial mode = 0



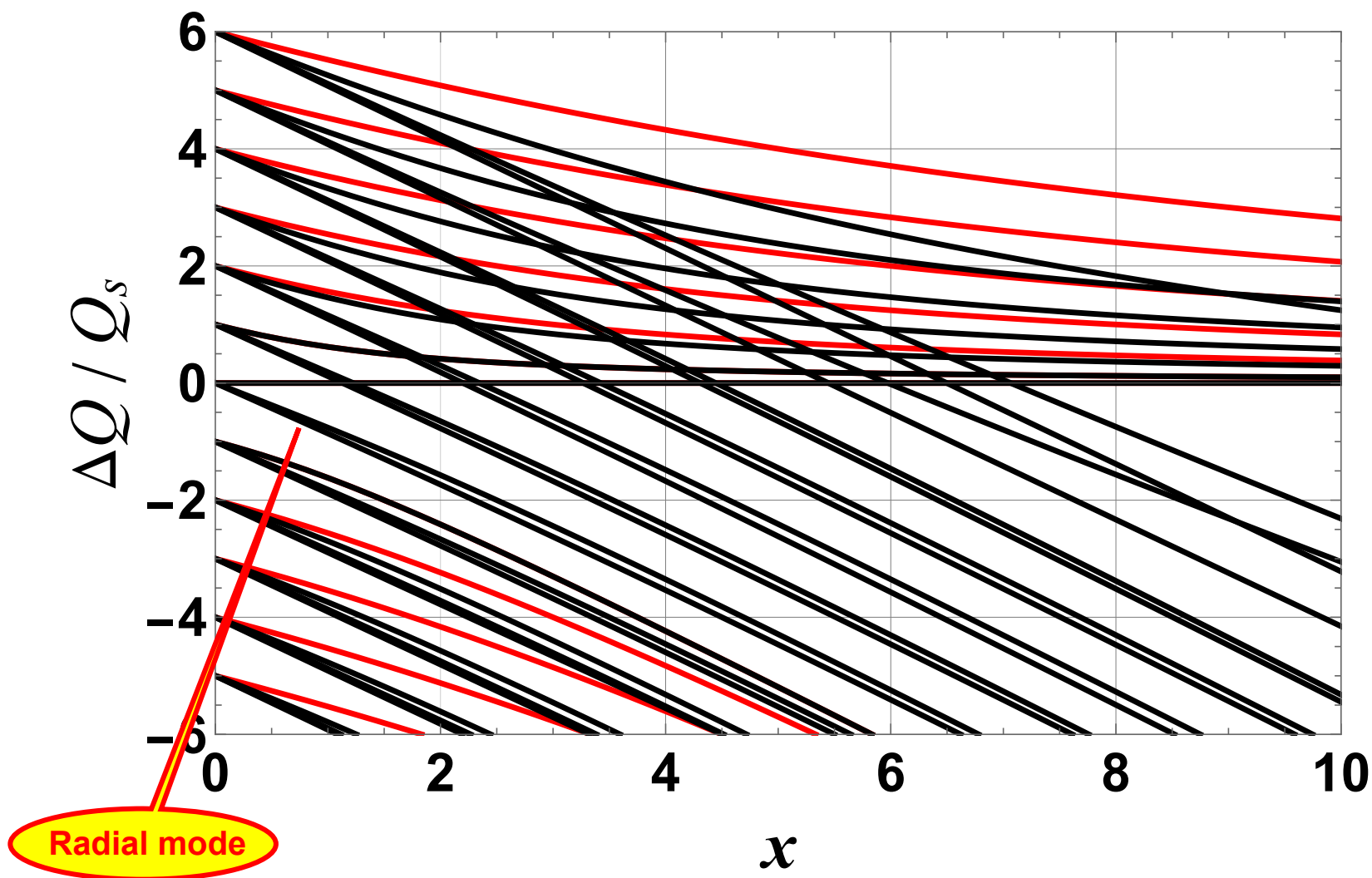
BOXCAR: radial mode = 0



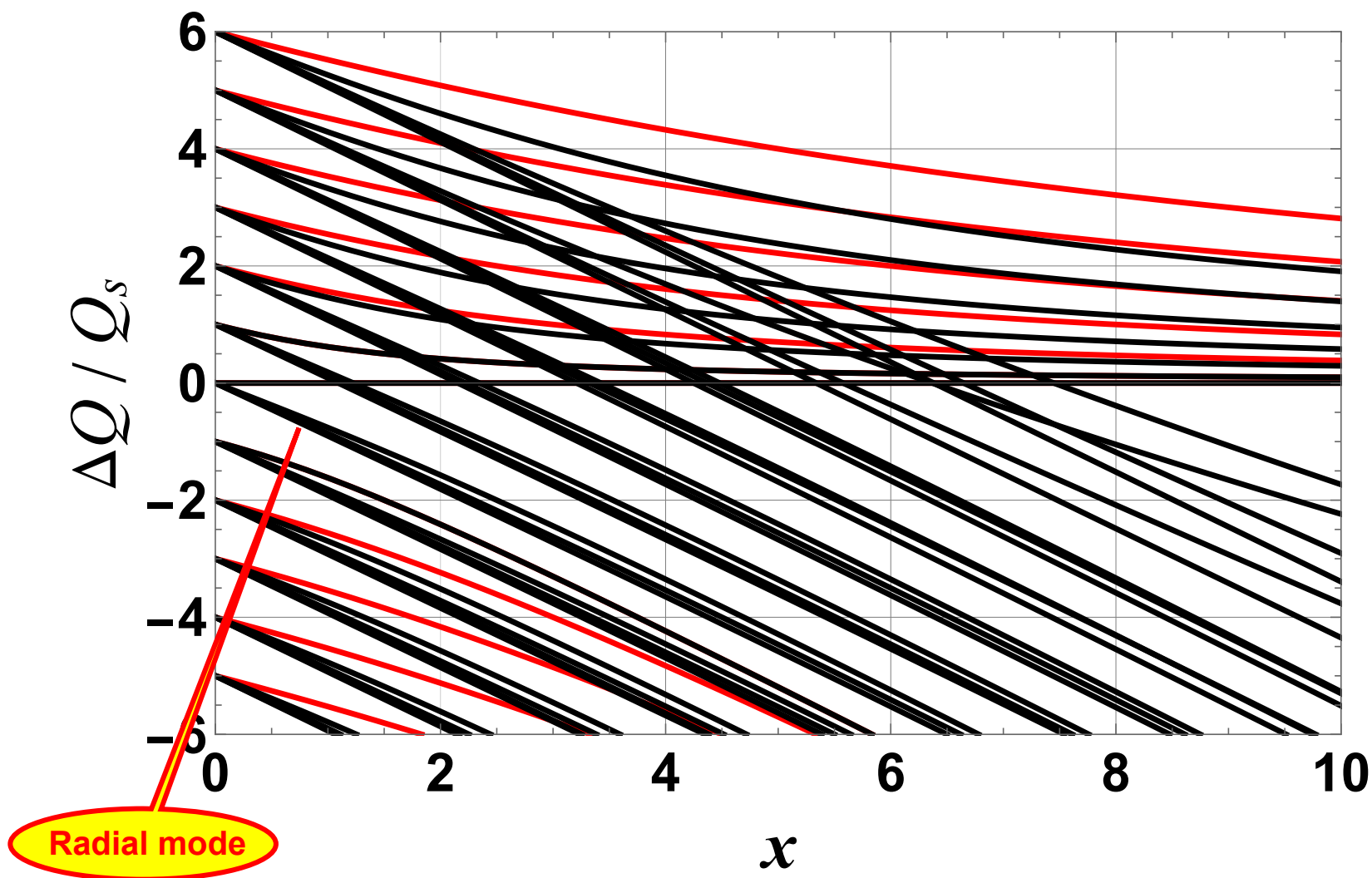
BOXCAR: radial mode = 1



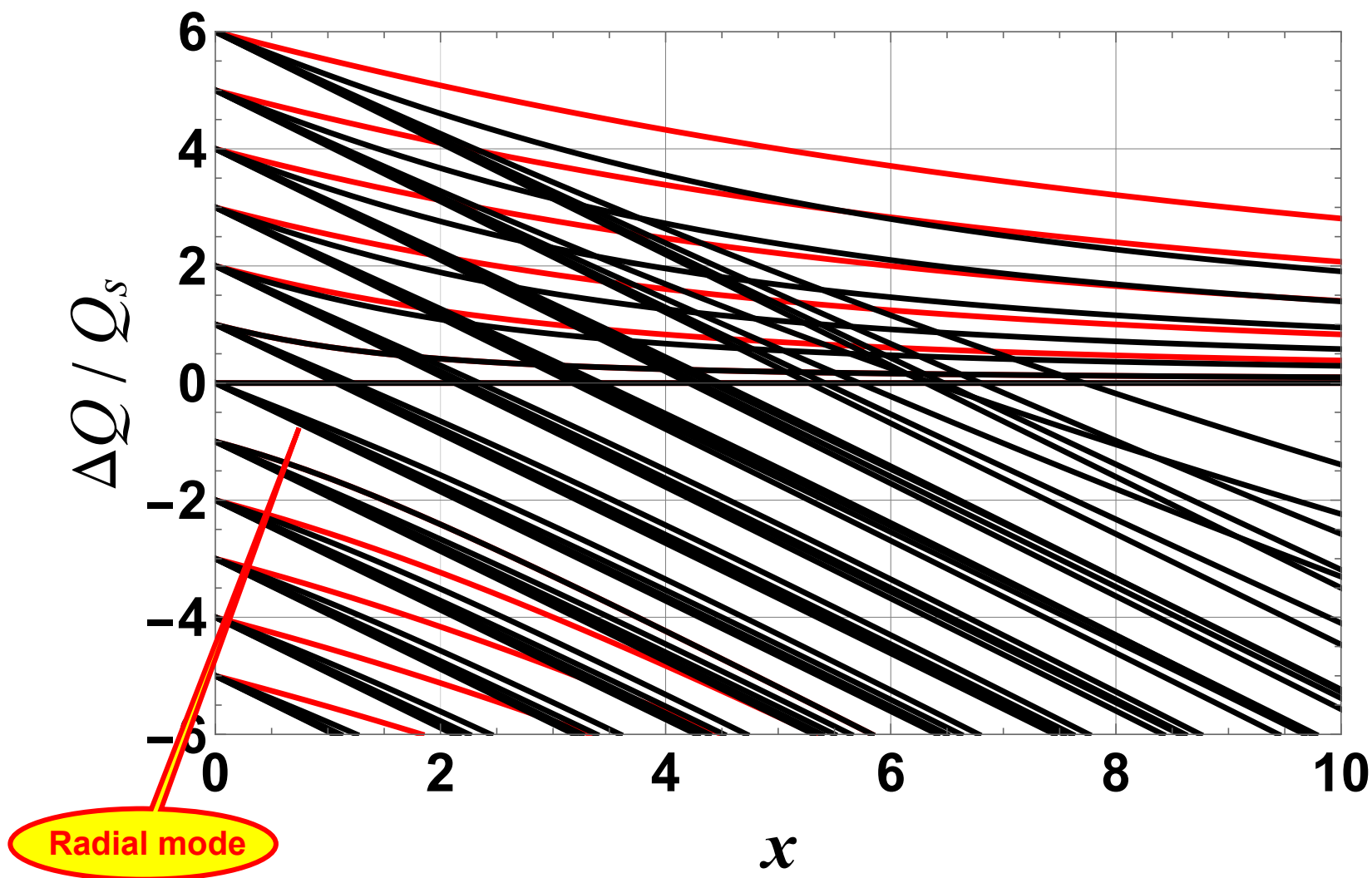
BOXCAR: radial mode = 2



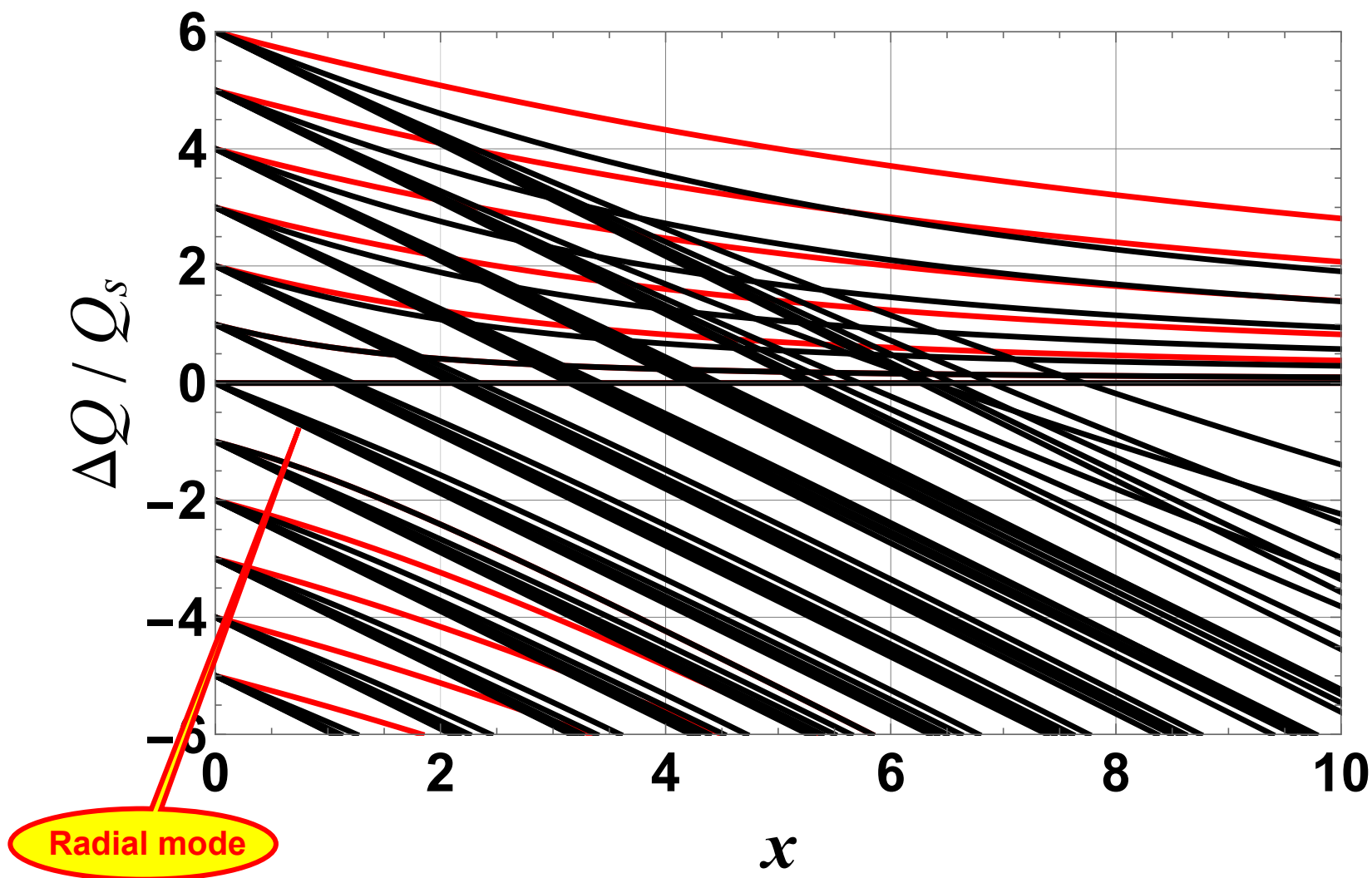
BOXCAR: radial mode = 3



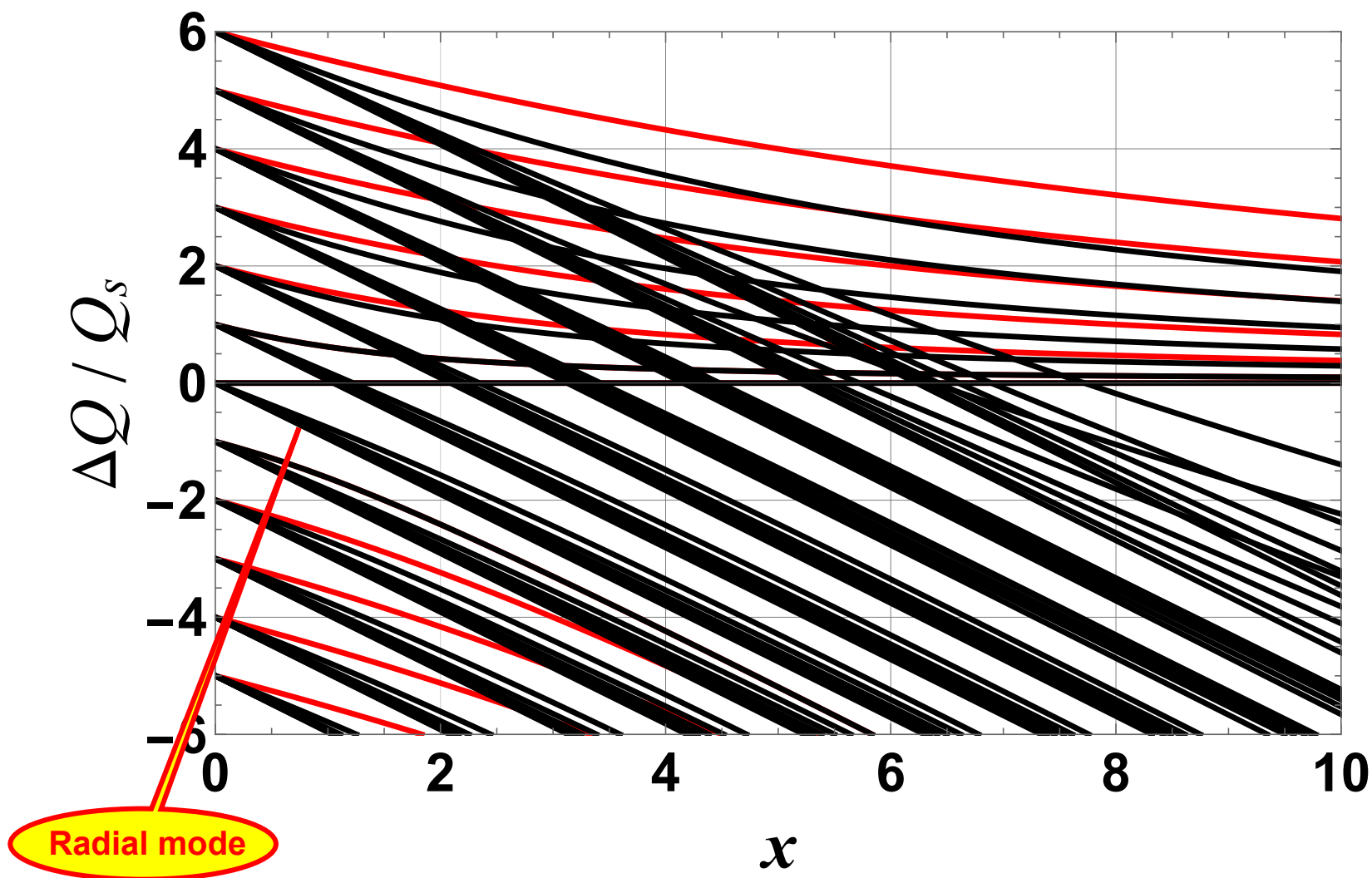
BOXCAR: radial mode = 4



BOXCAR: radial mode = 5

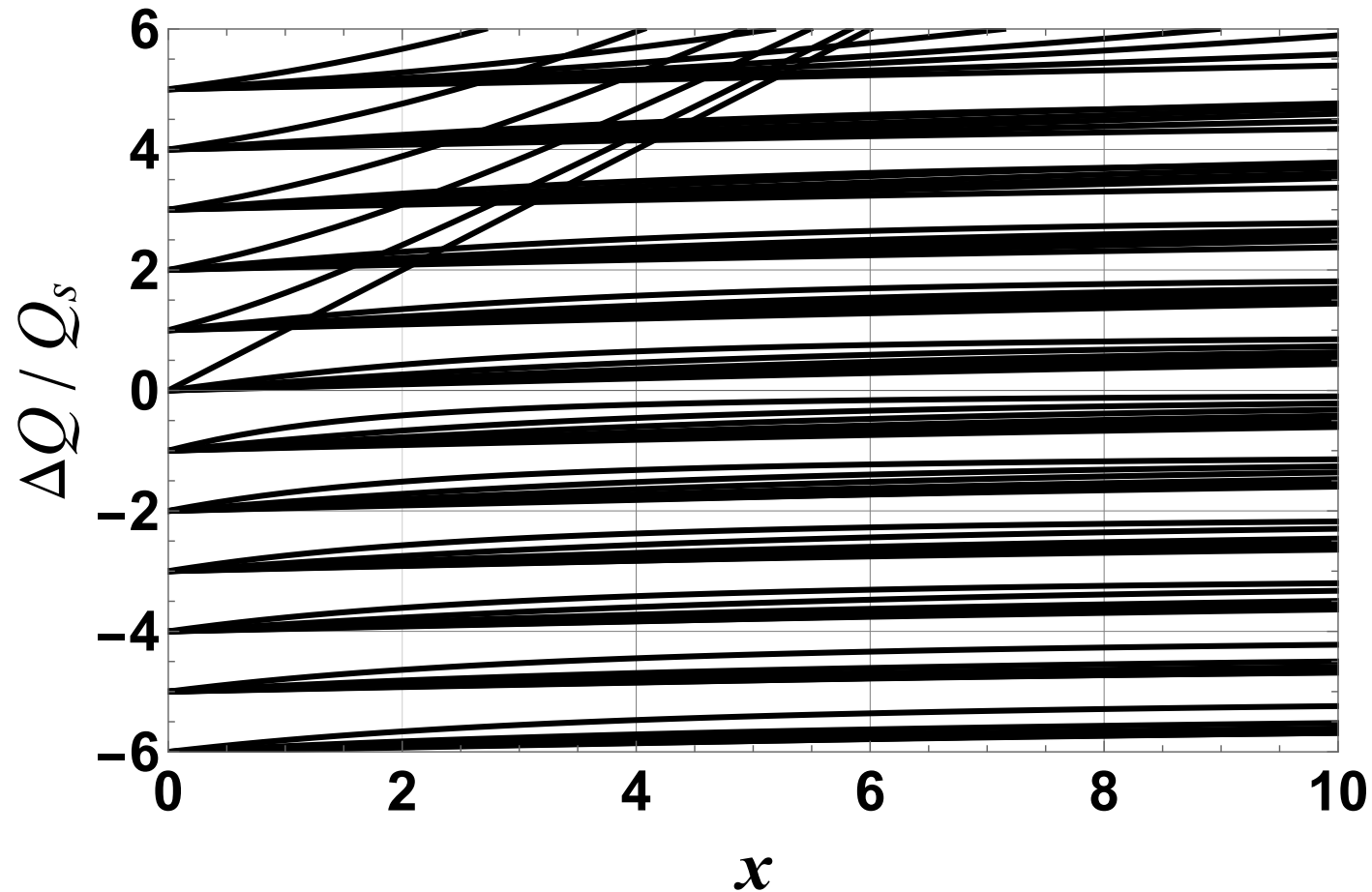


BOXCAR: radial mode = 6



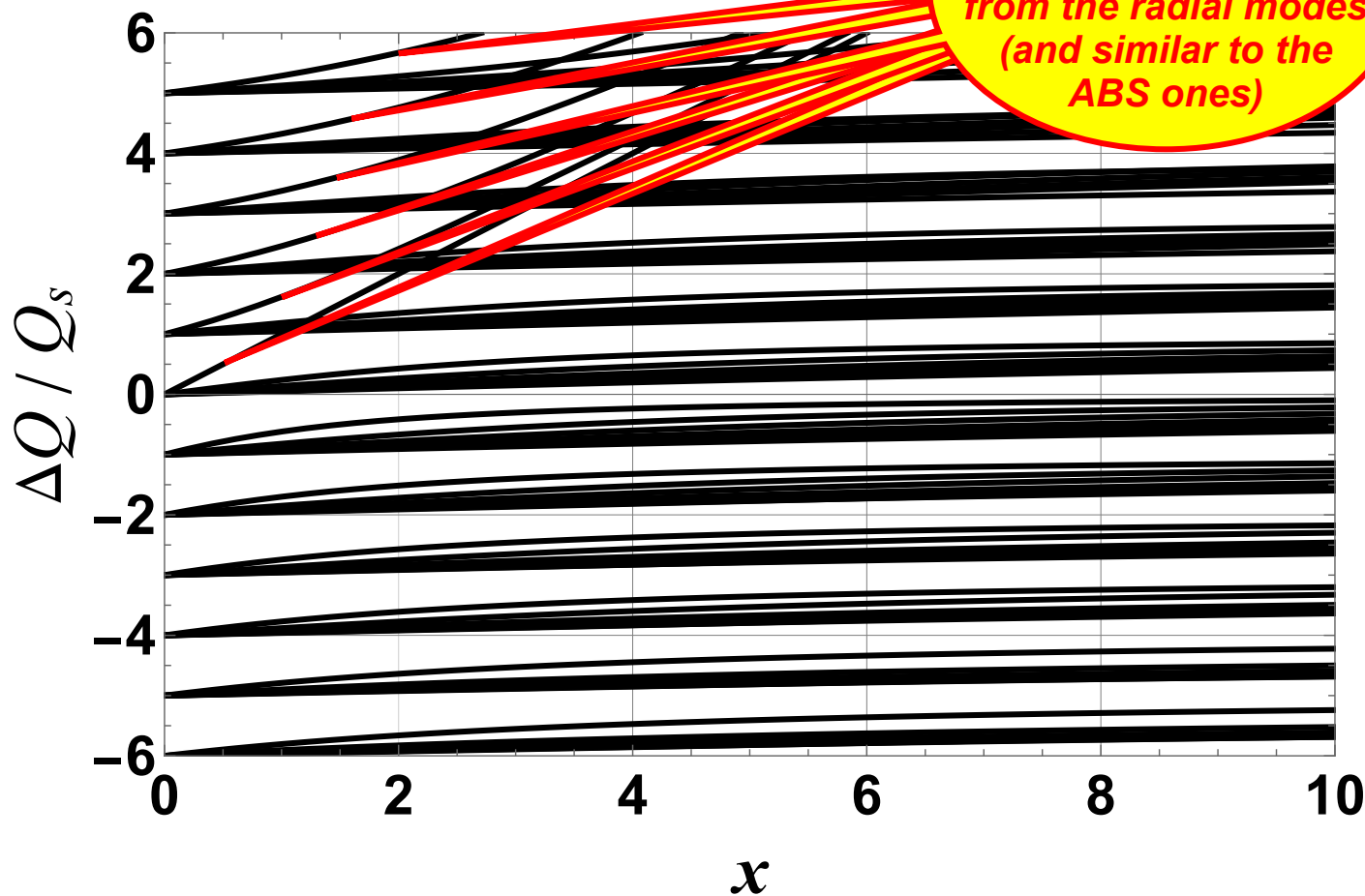
BOXCAR: radial mode = 6

◆ Dipolar force only



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◆ Dipolar force only



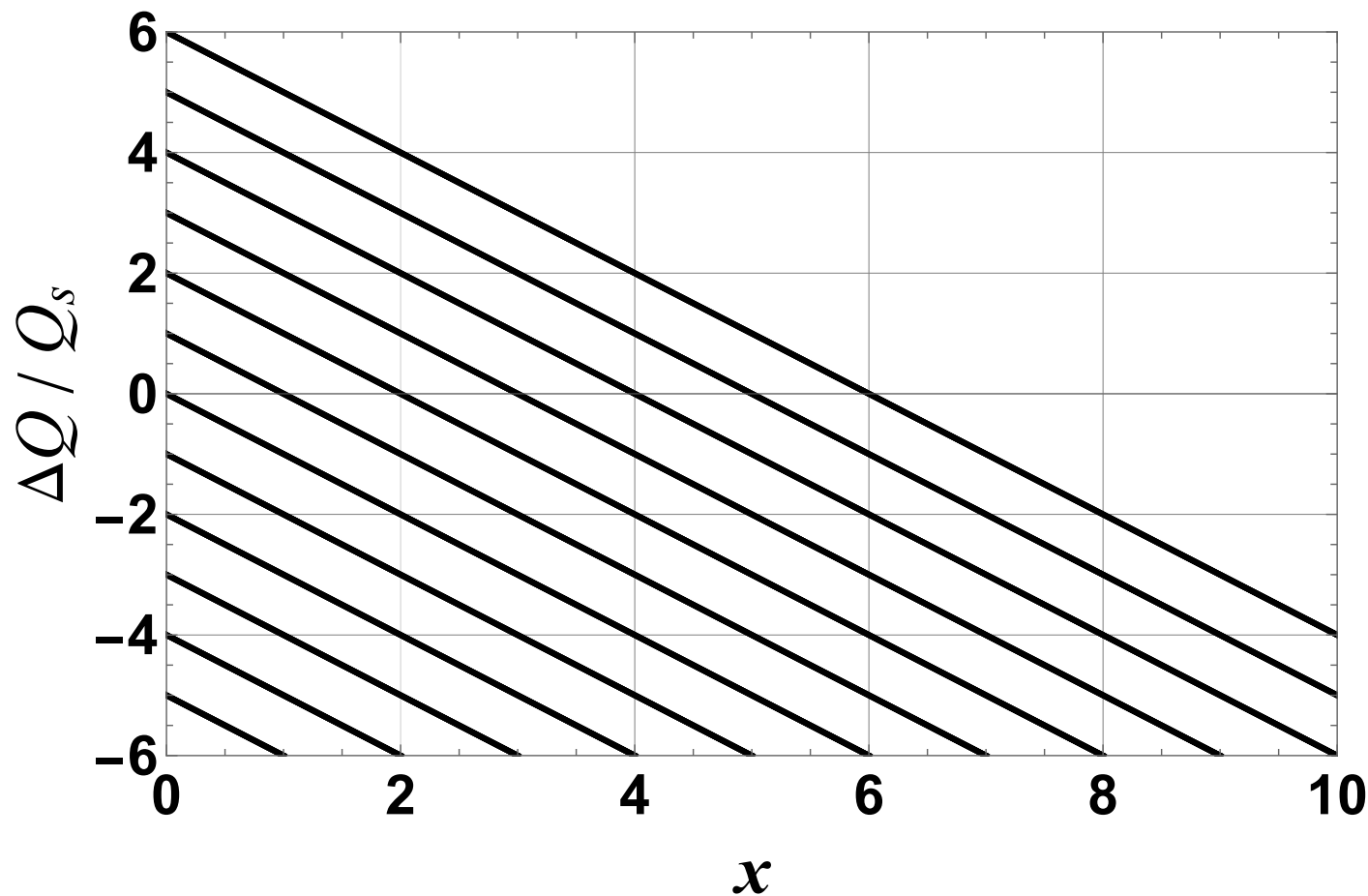
BOXCAR: radial mode = 6

◆ Dipolar force only



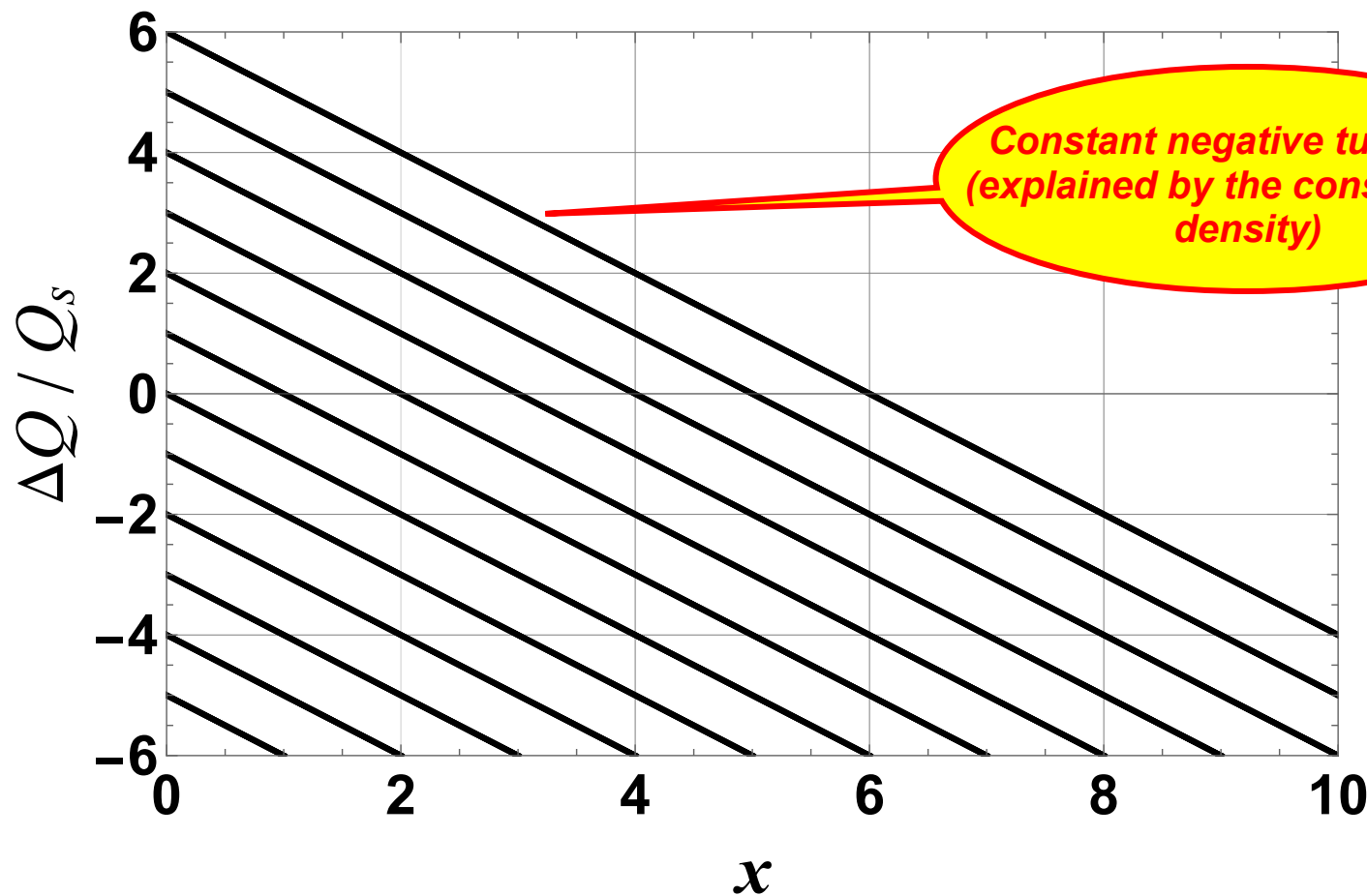
BOXCAR: radial mode = 6

◆ Quadrupolar force only



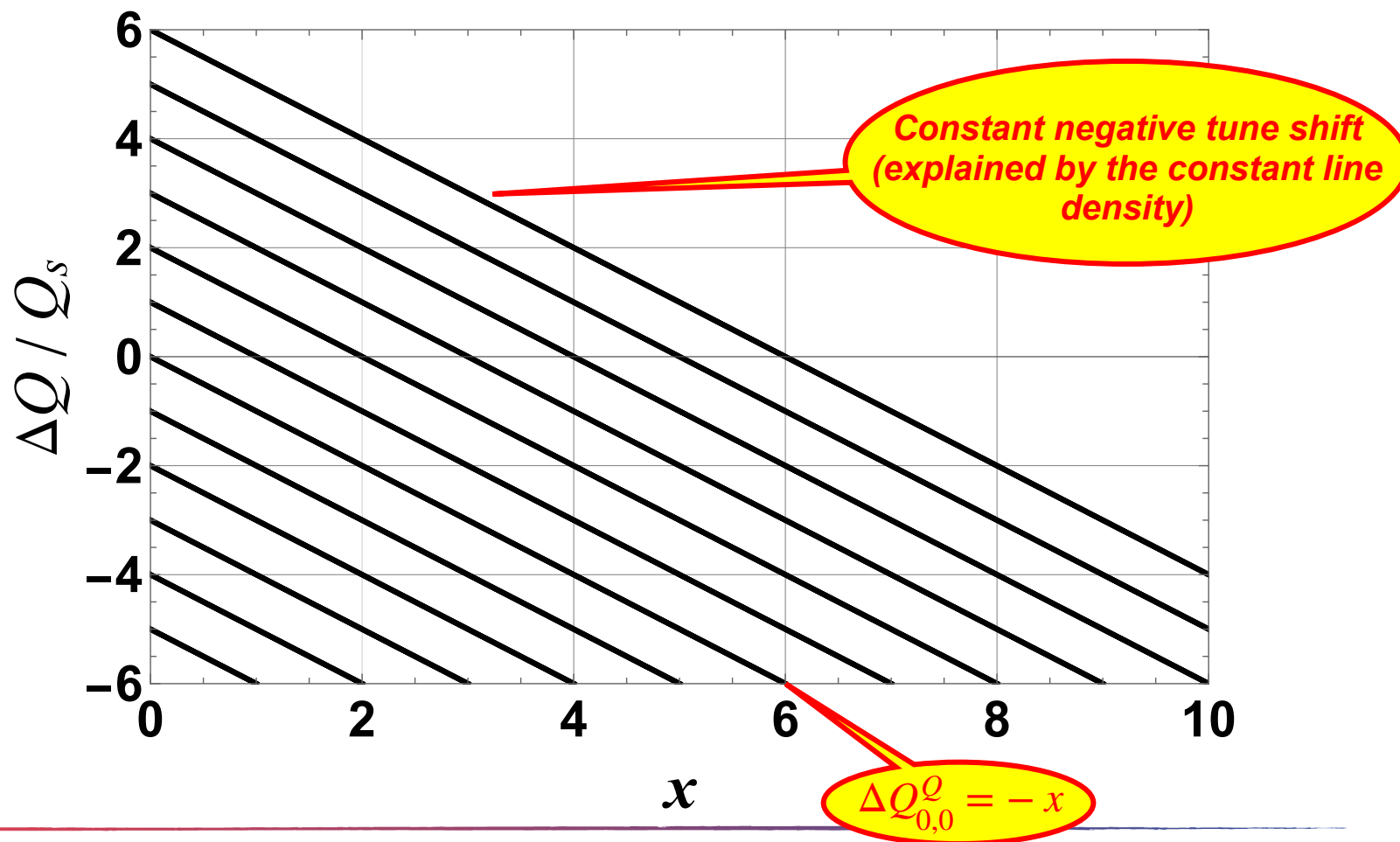
BOXCAR: radial mode = 6

◆ Quadrupolar force only



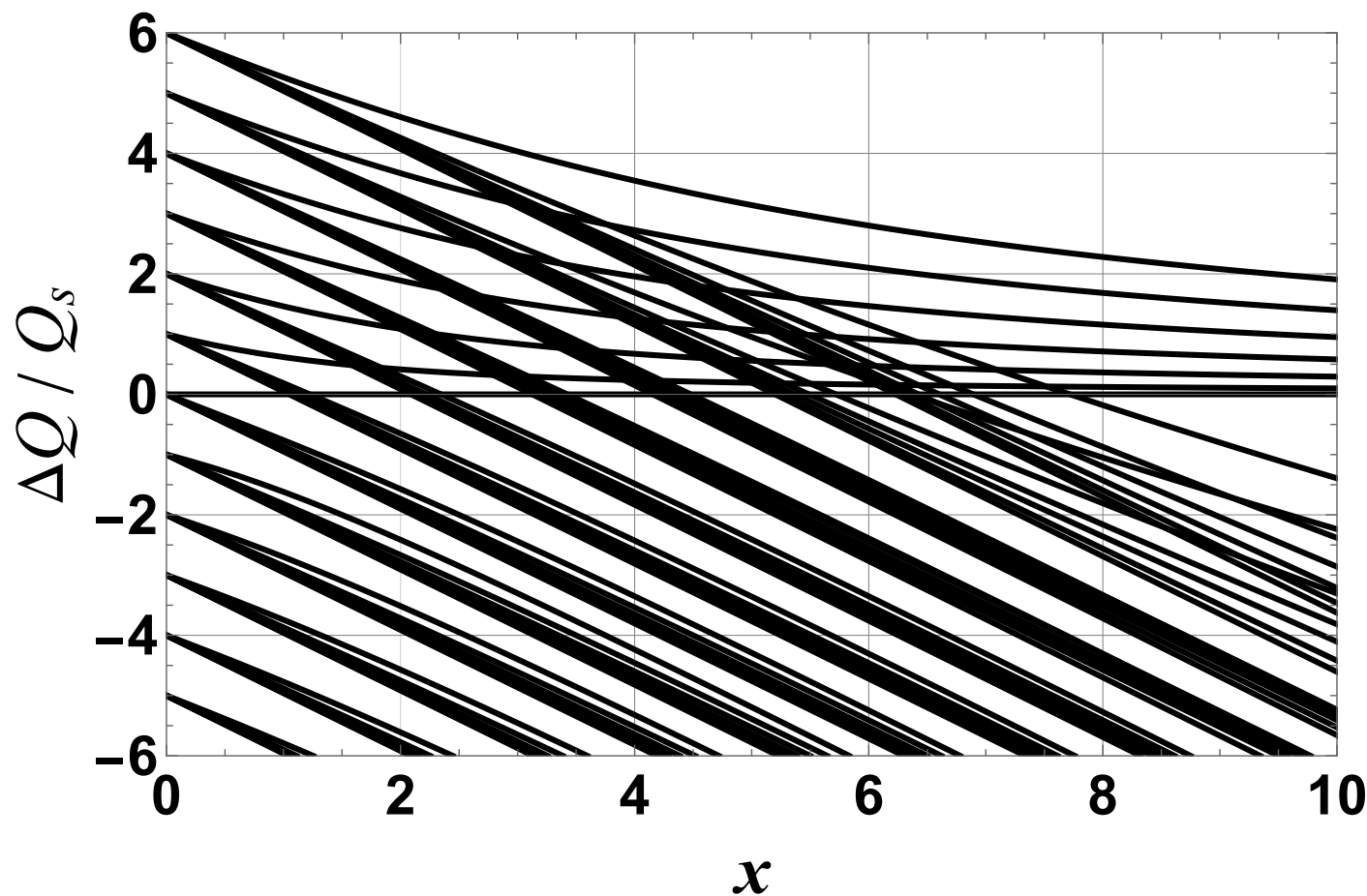
BOXCAR: radial mode = 6

◆ Quadrupolar force only



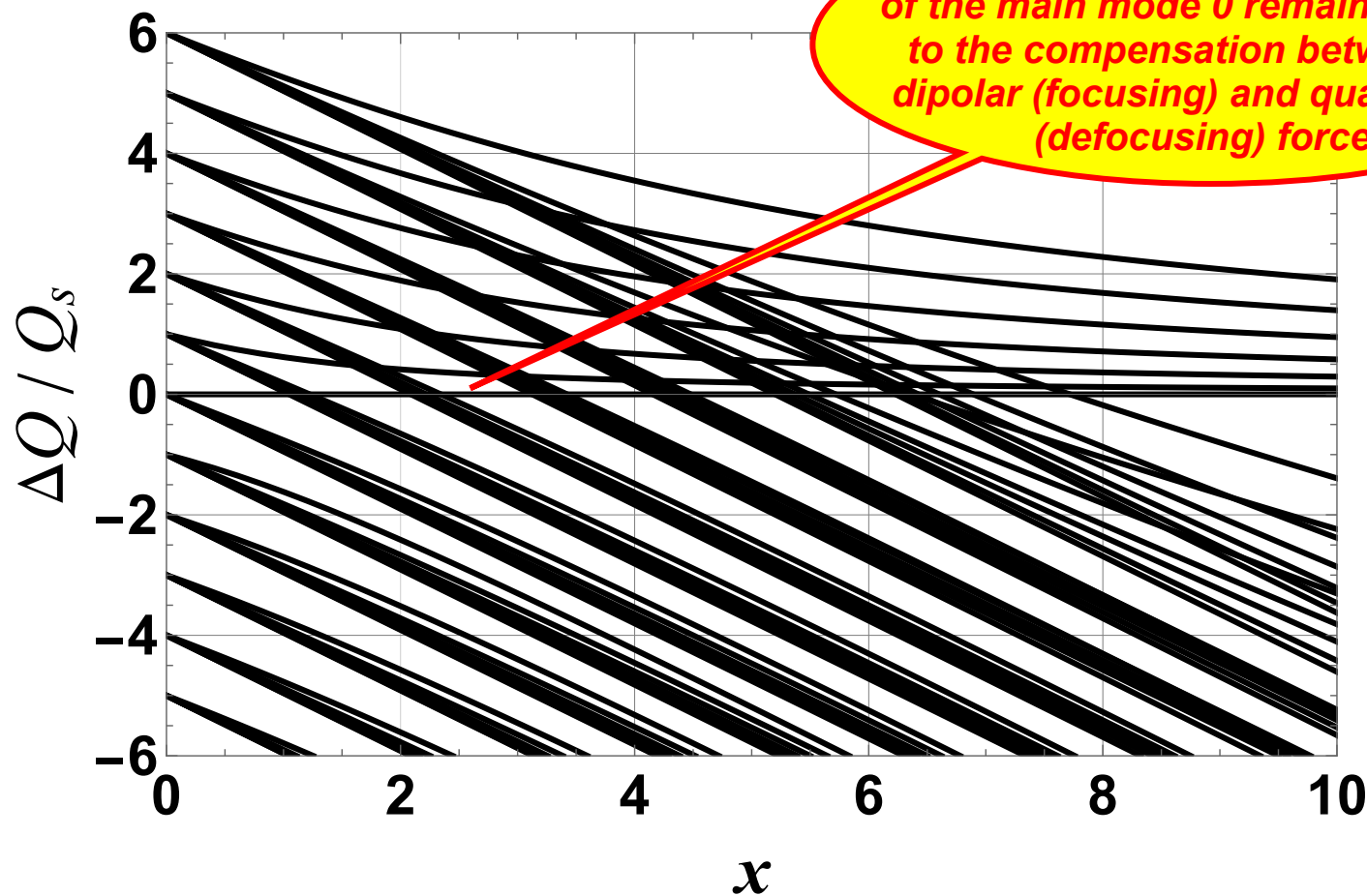
BOXCAR: radial mode = 6

◆ Sum of the 2 forces



BOXCAR: radial mode = 6

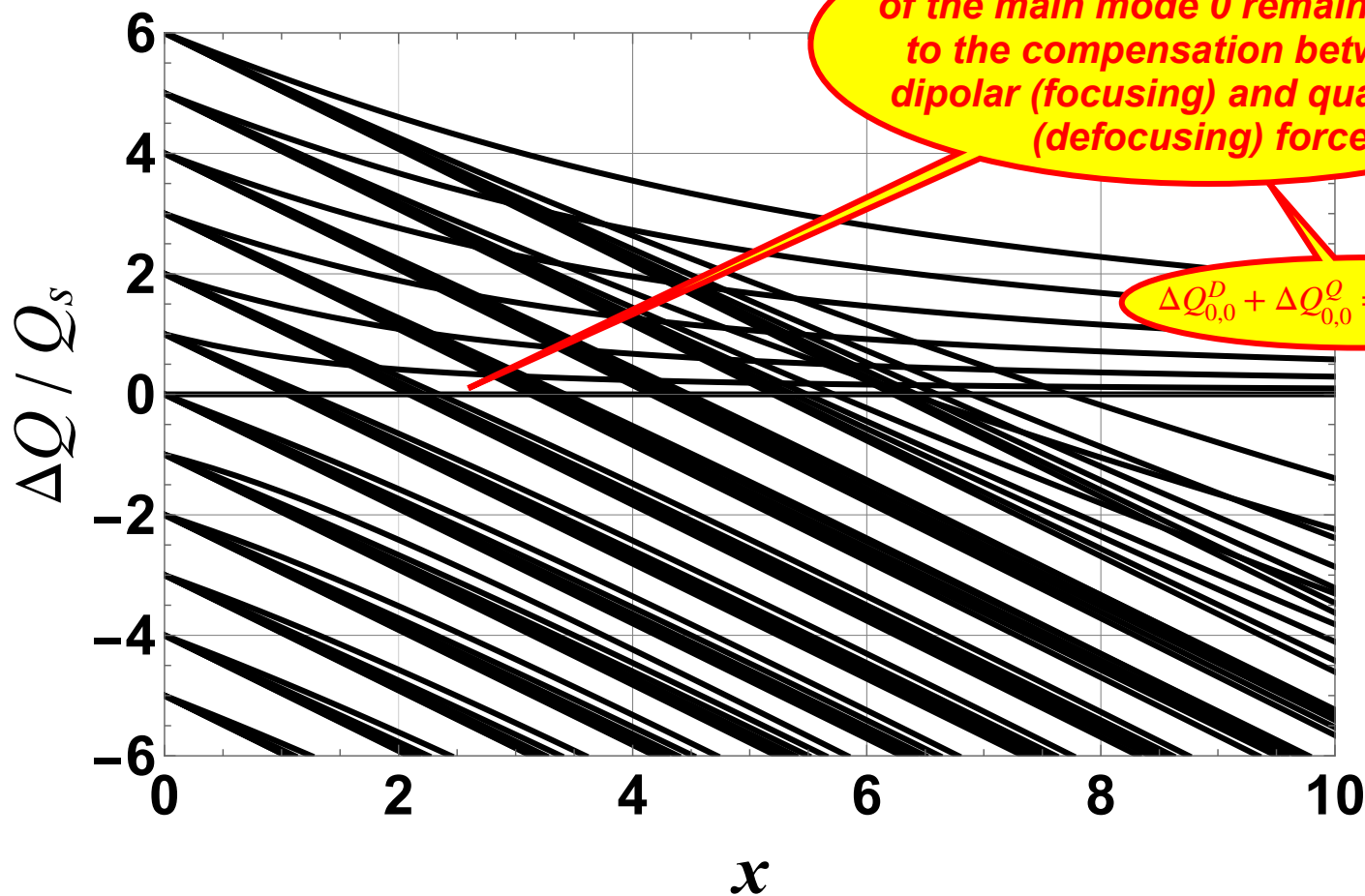
◆ Sum of the 2 forces



The fact that the tune shift of the main mode 0 remains 0 is due to the compensation between the dipolar (focusing) and quadrupolar (defocusing) forces

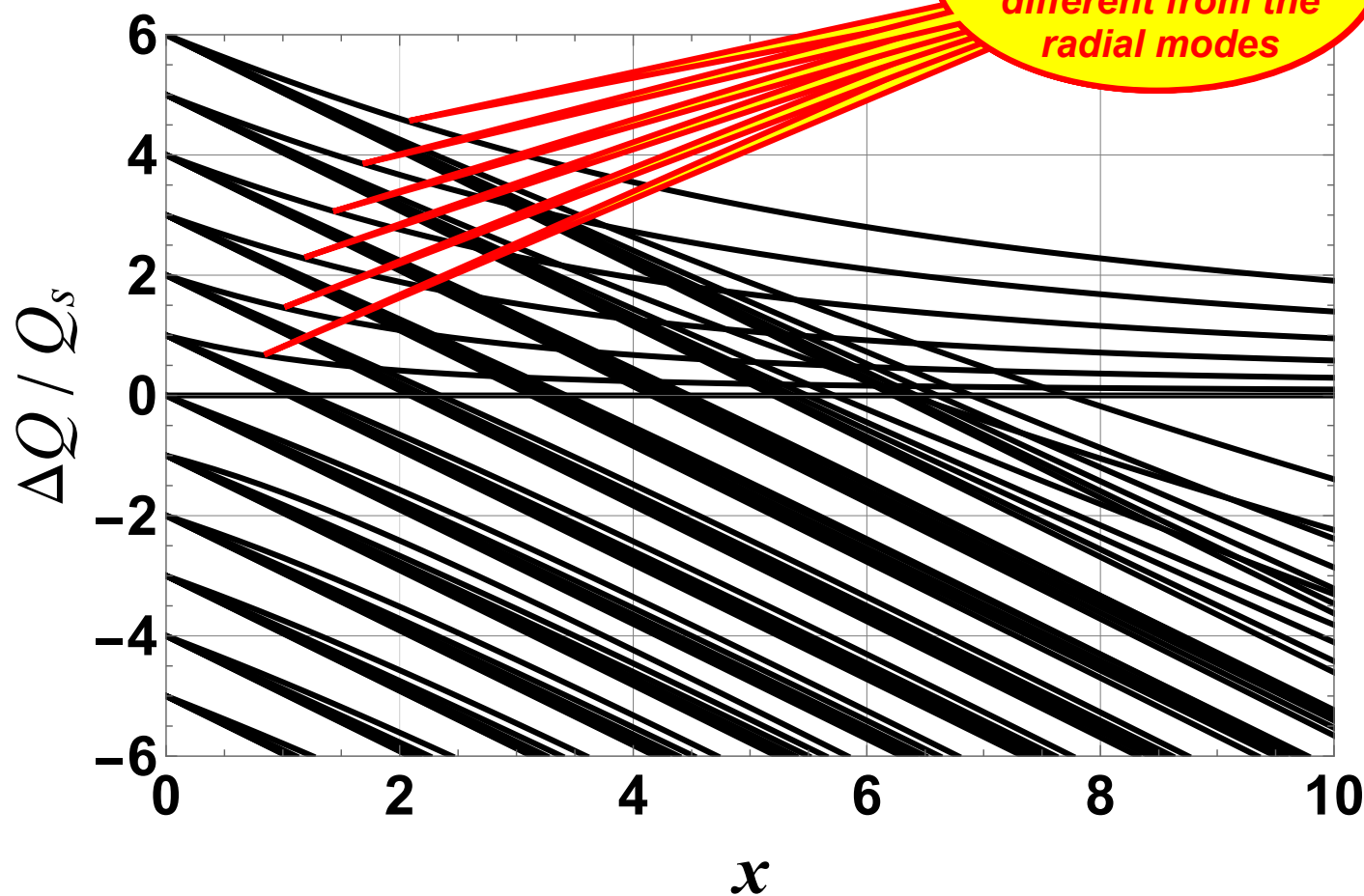
BOXCAR: radial mode = 6

◆ Sum of the 2 forces



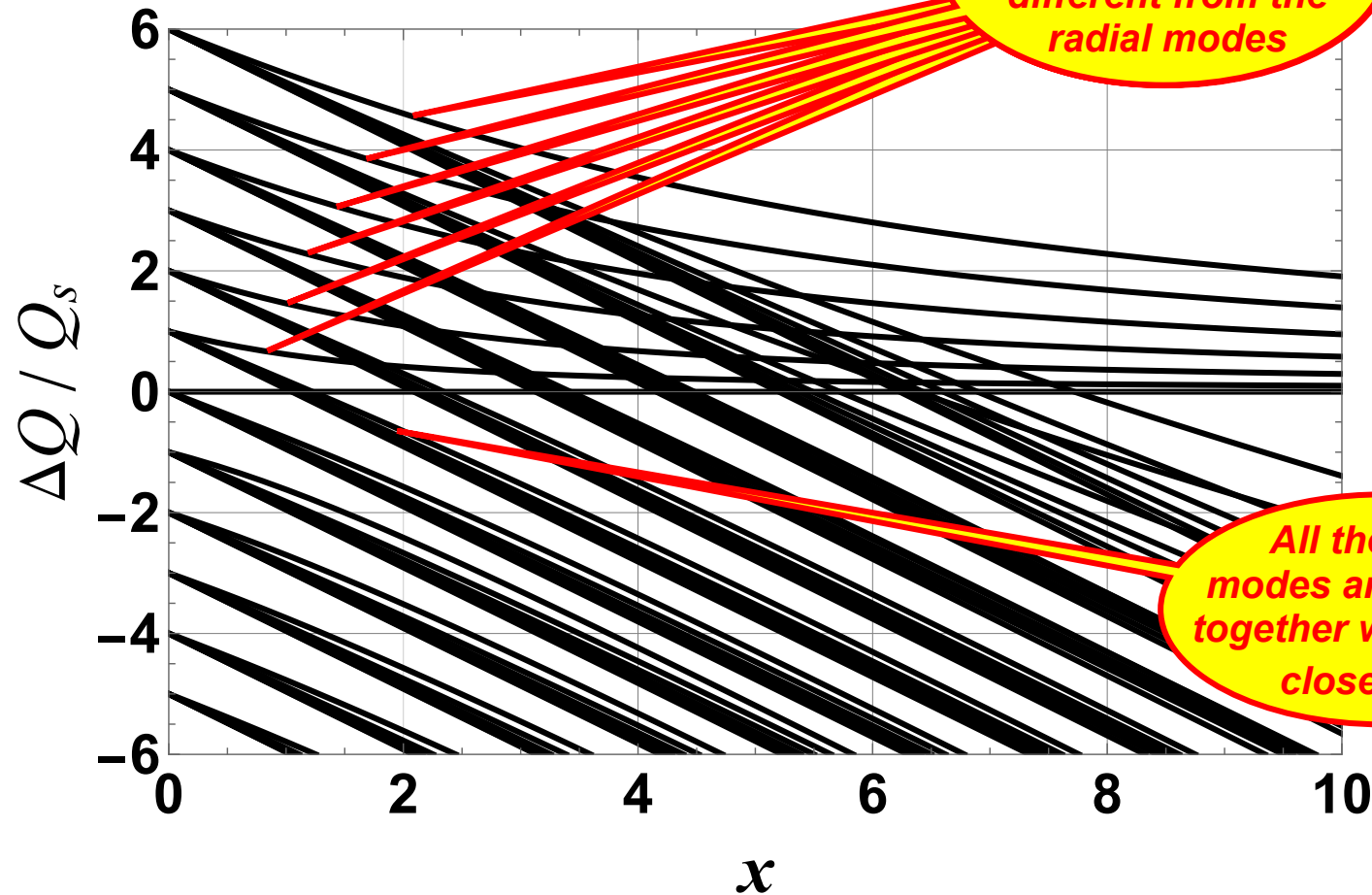
BOXCAR: radial mode = 6

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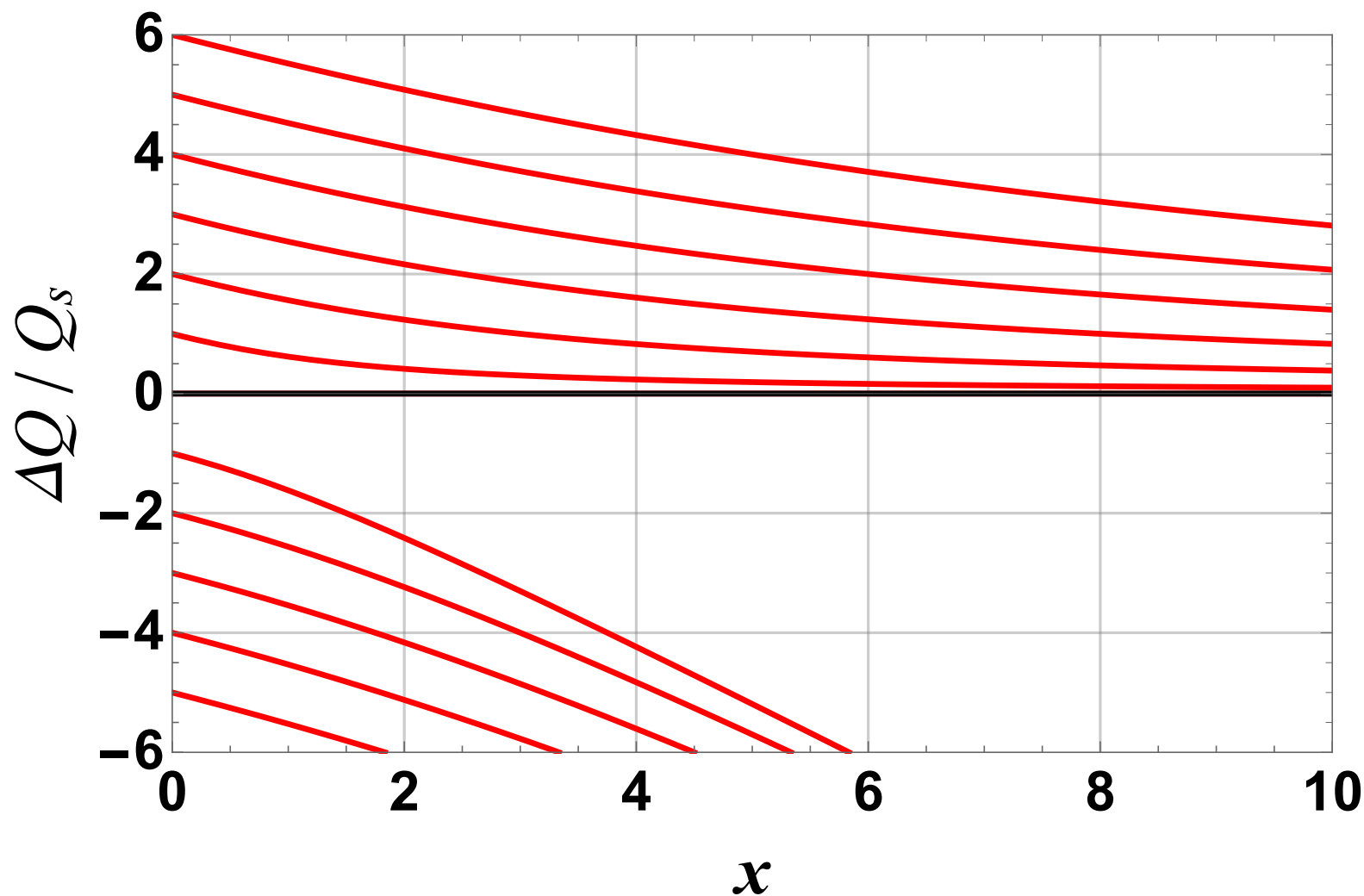
◆ Sum of the 2 forces



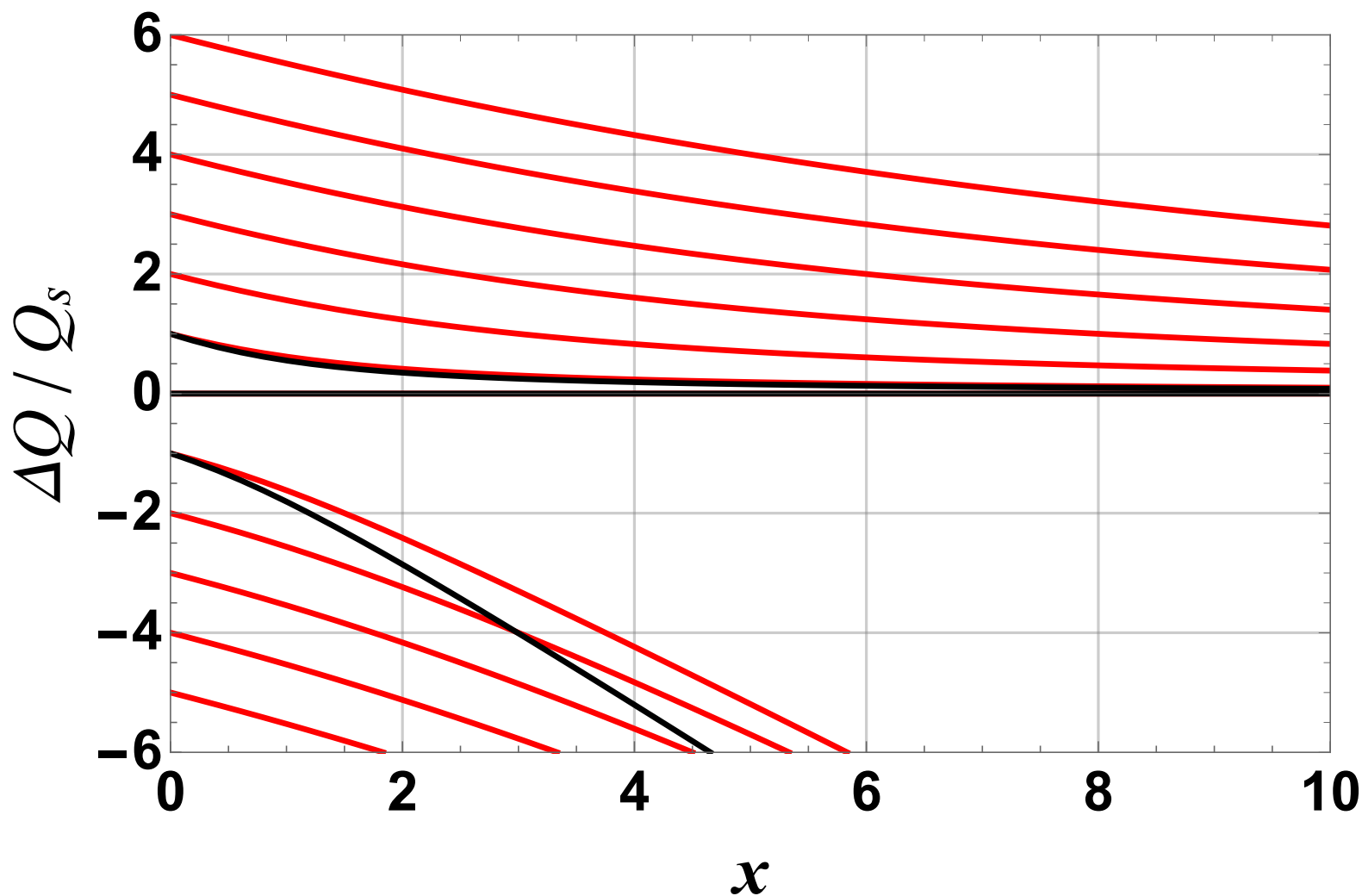
The main positive azimuthal modes are different from the radial modes

All the radial modes are grouped together with a slope close to $-x$

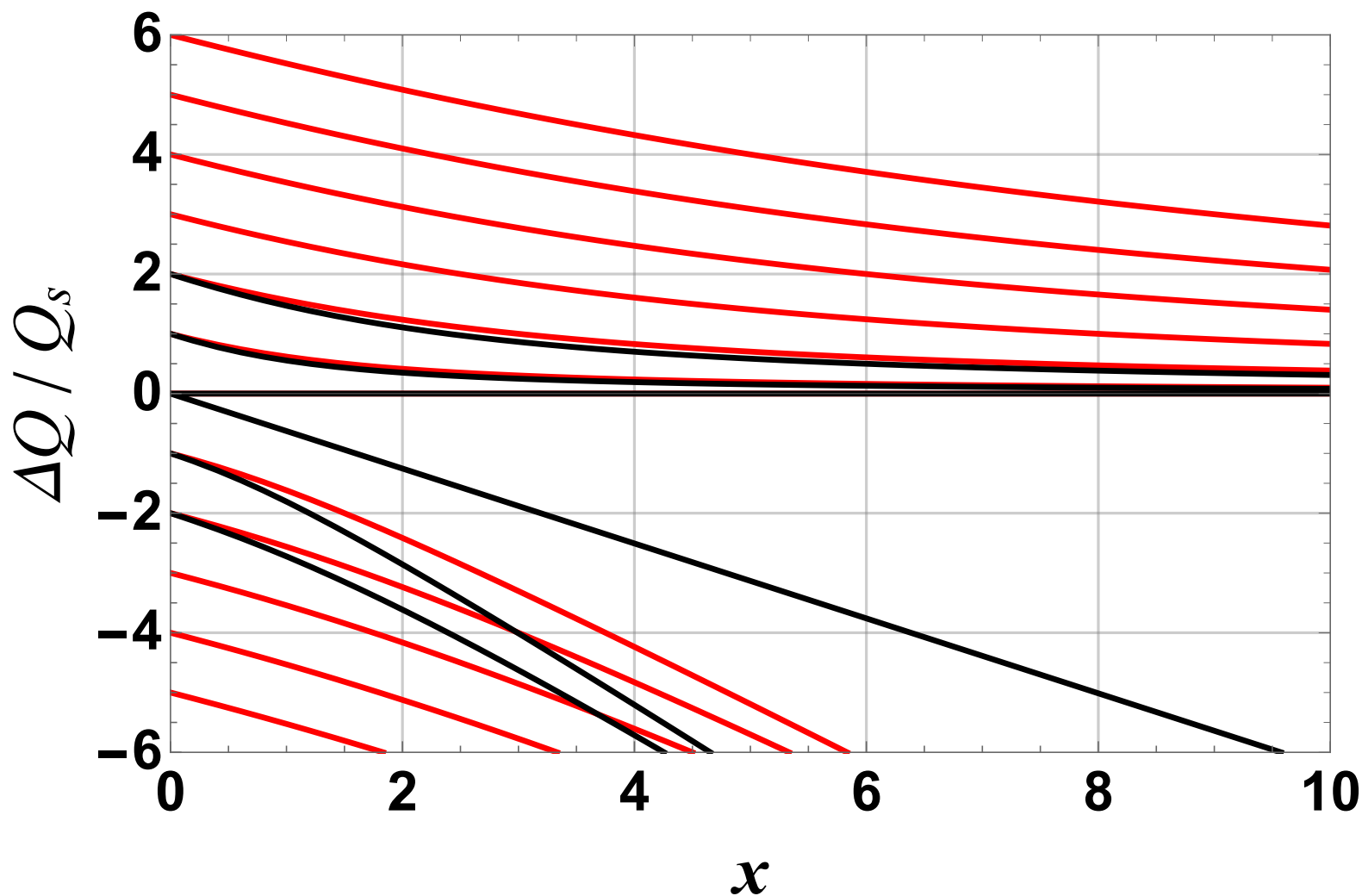
GAUSSIAN: $m_{max} = 0$



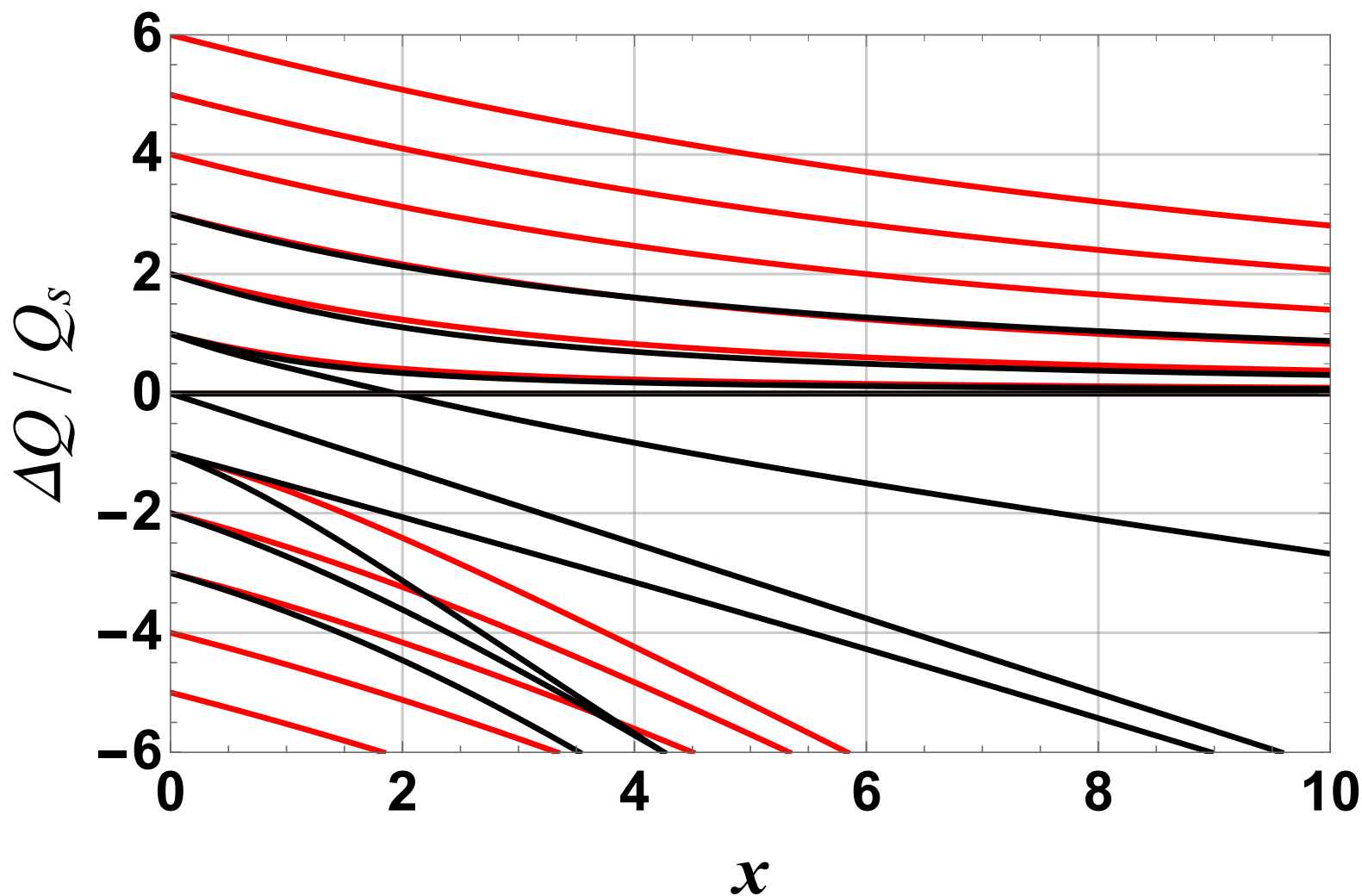
GAUSSIAN: $m_{max} = 1$



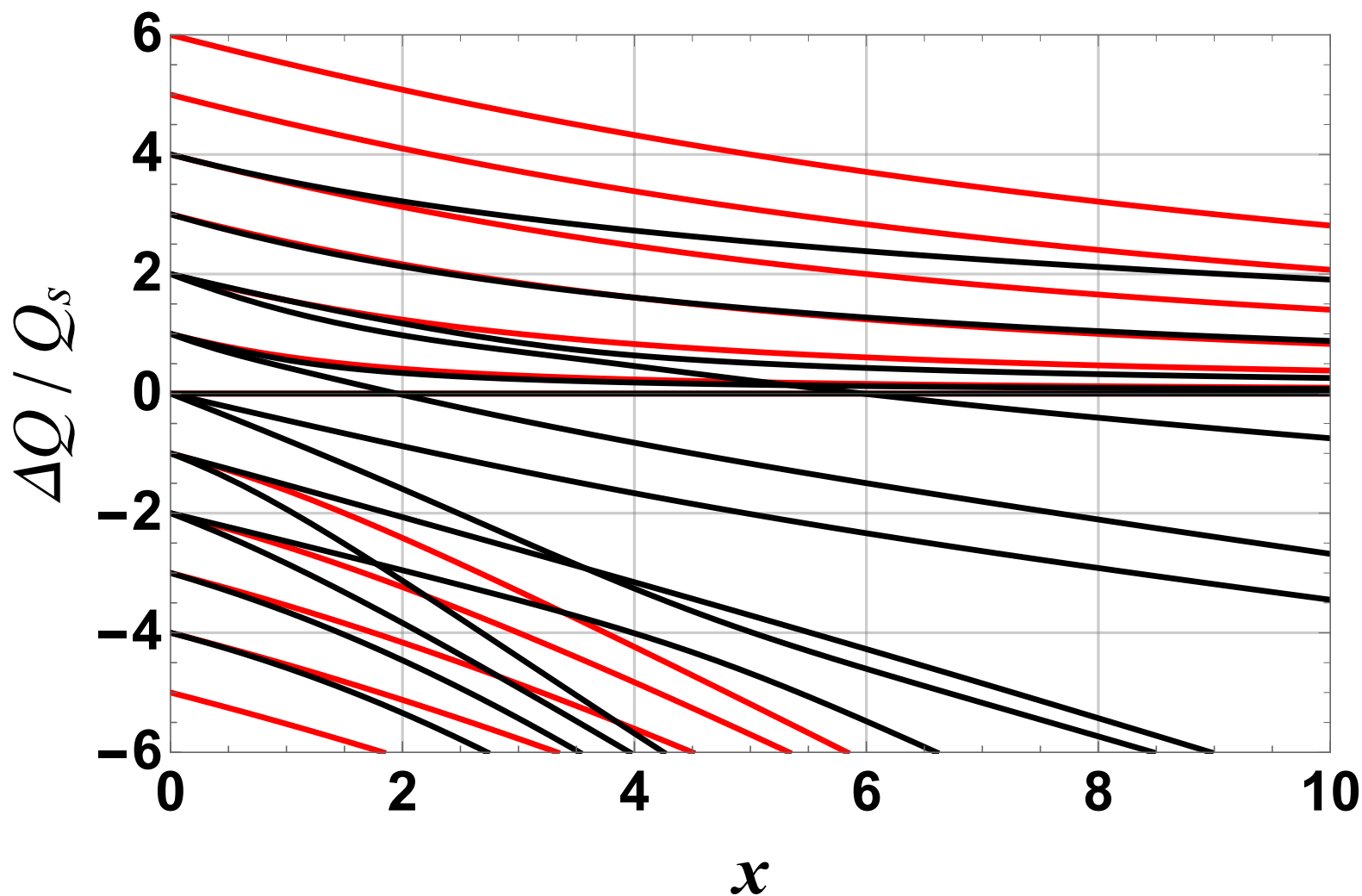
GAUSSIAN: $m_{max} = 2$



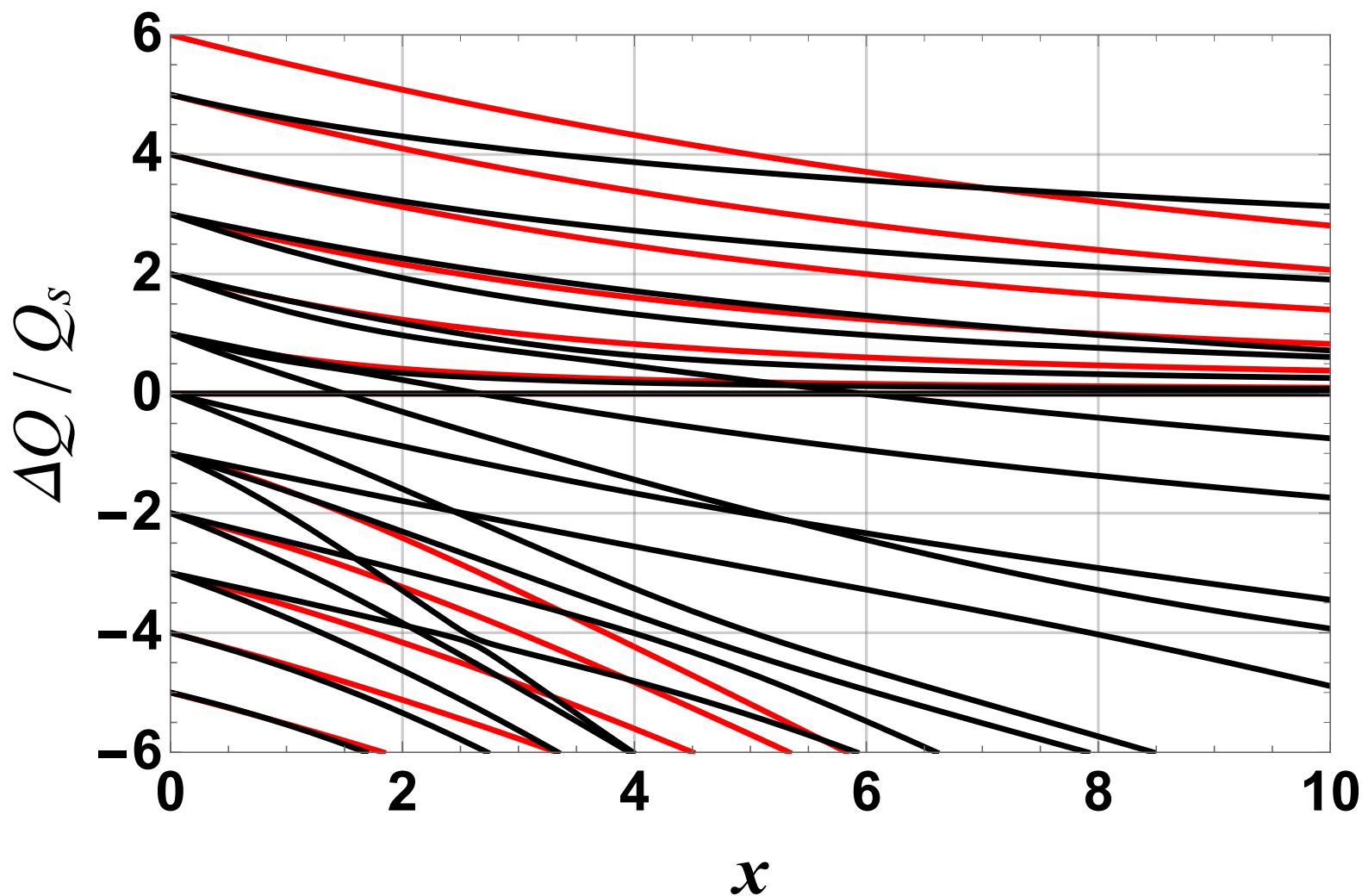
GAUSSIAN: $m_{max} = 3$



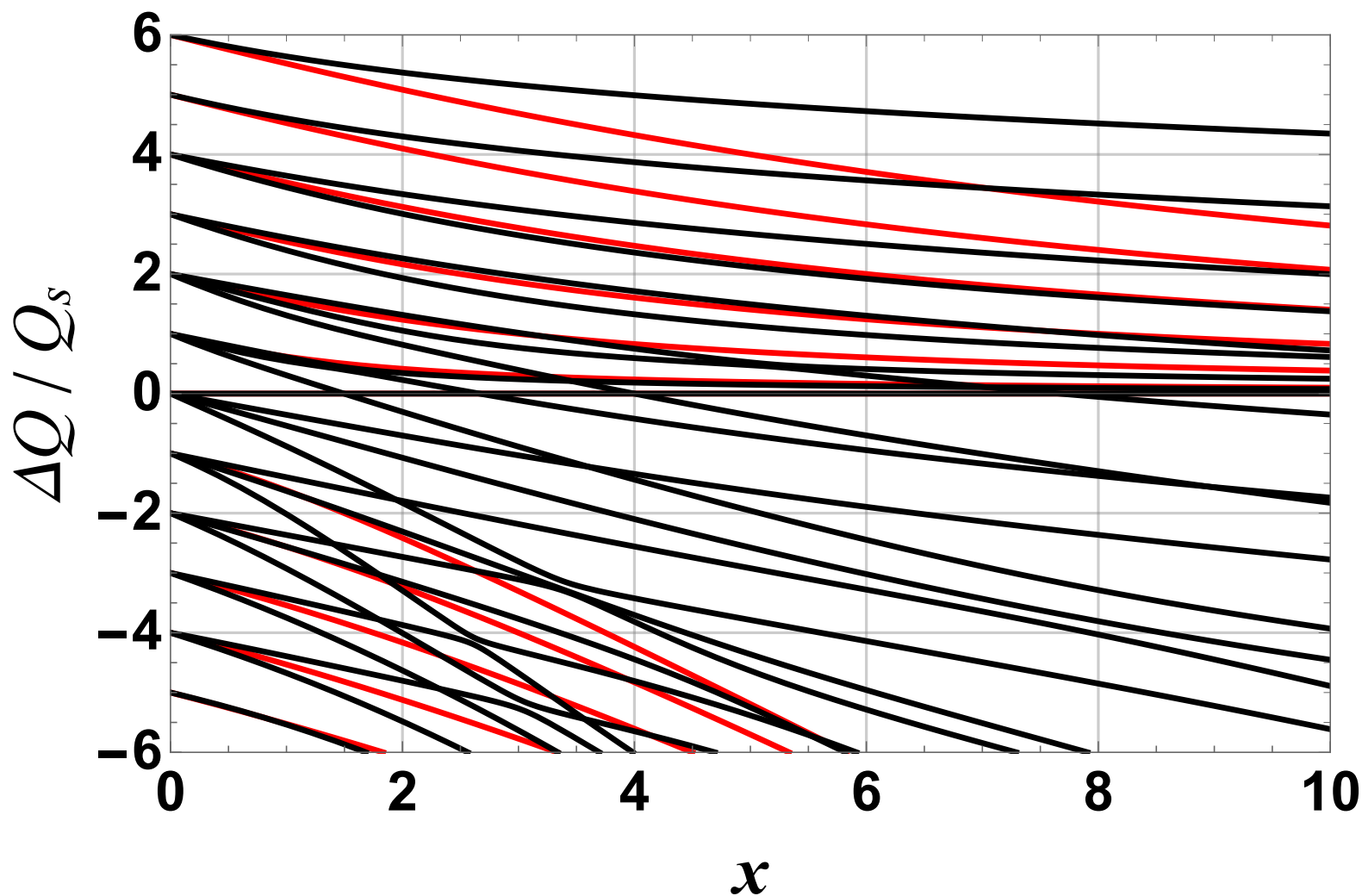
GAUSSIAN: $m_{max} = 4$



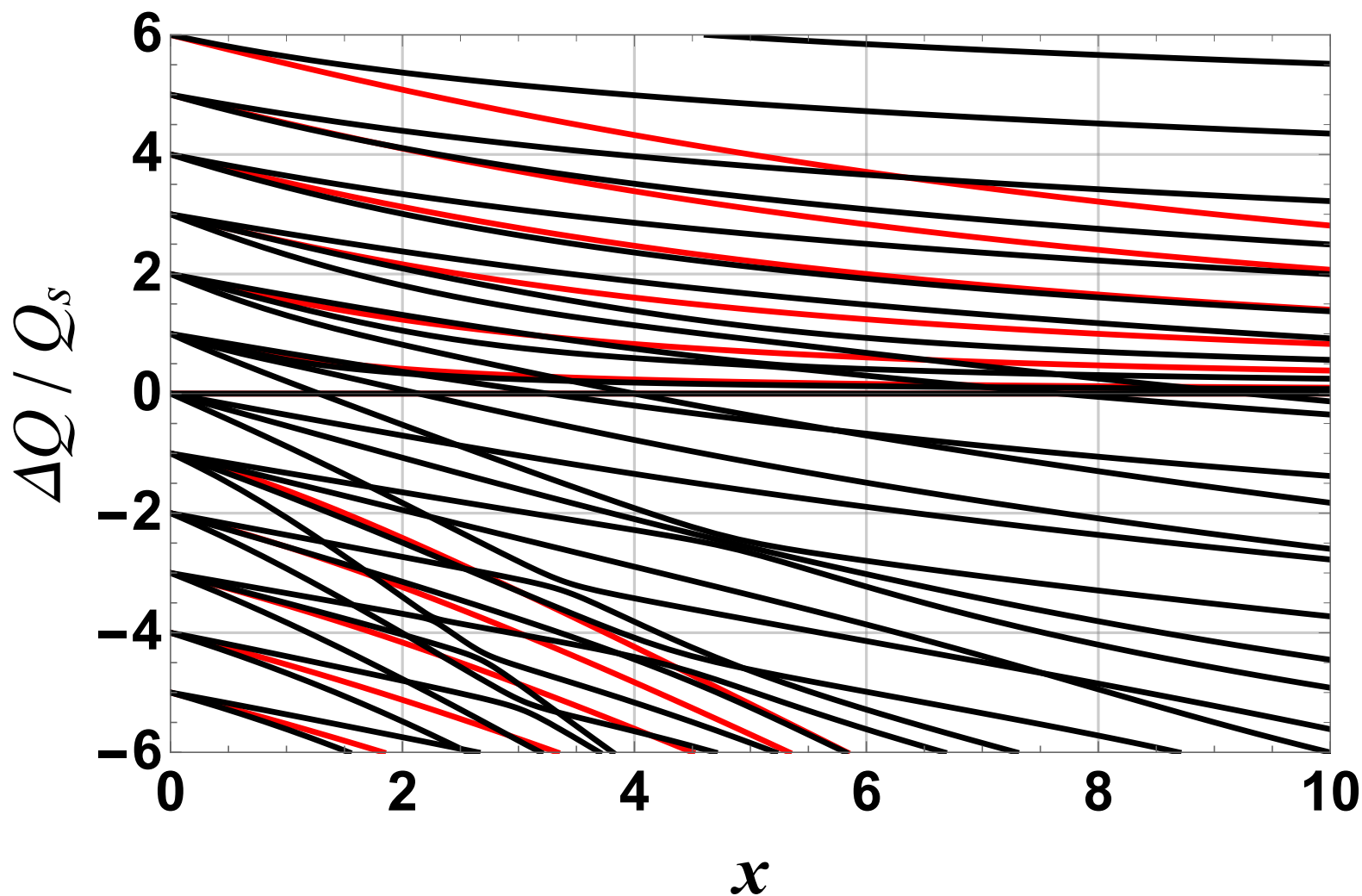
GAUSSIAN: $m_{max} = 5$



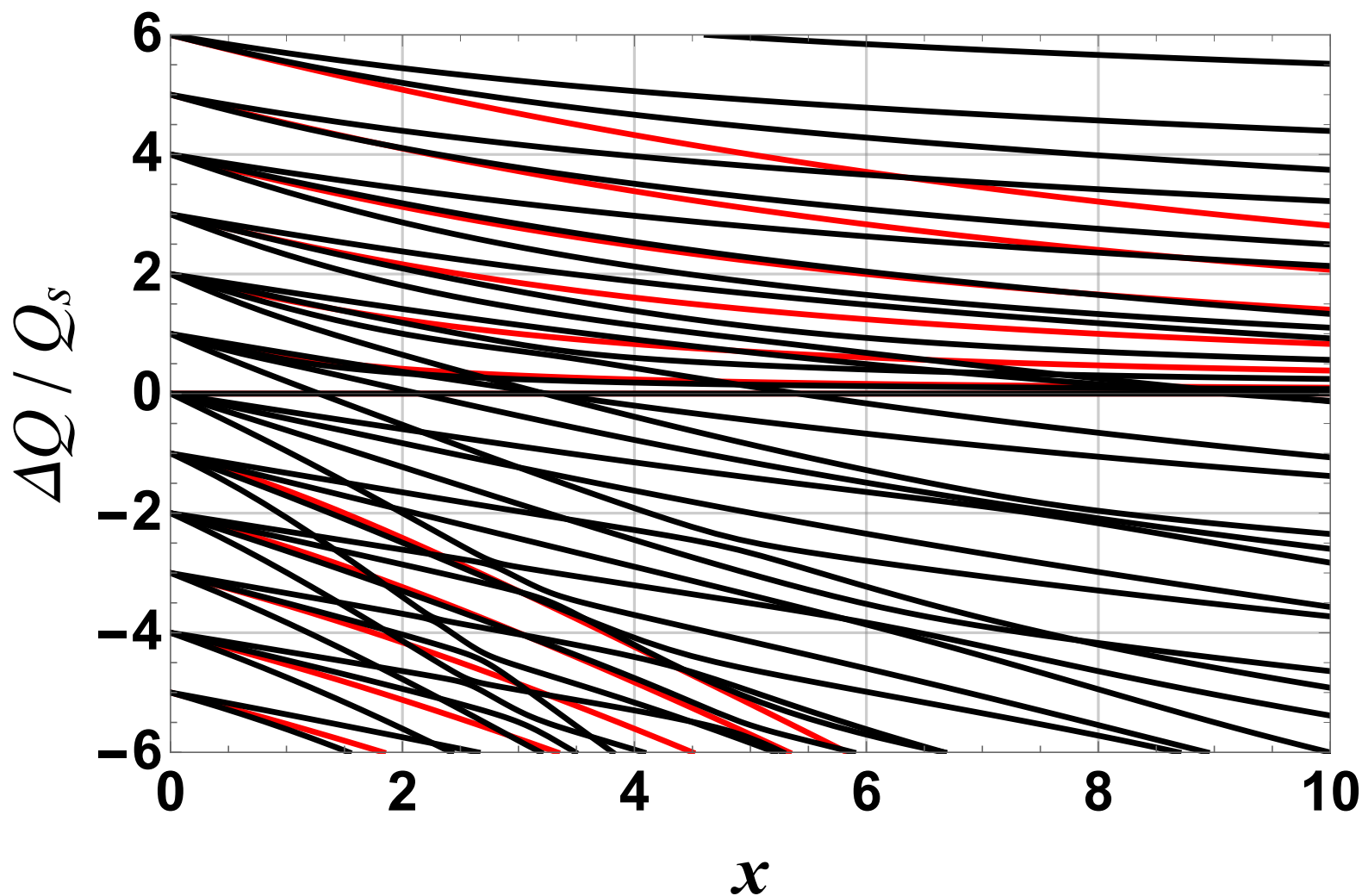
GAUSSIAN: $m_{max} = 6$



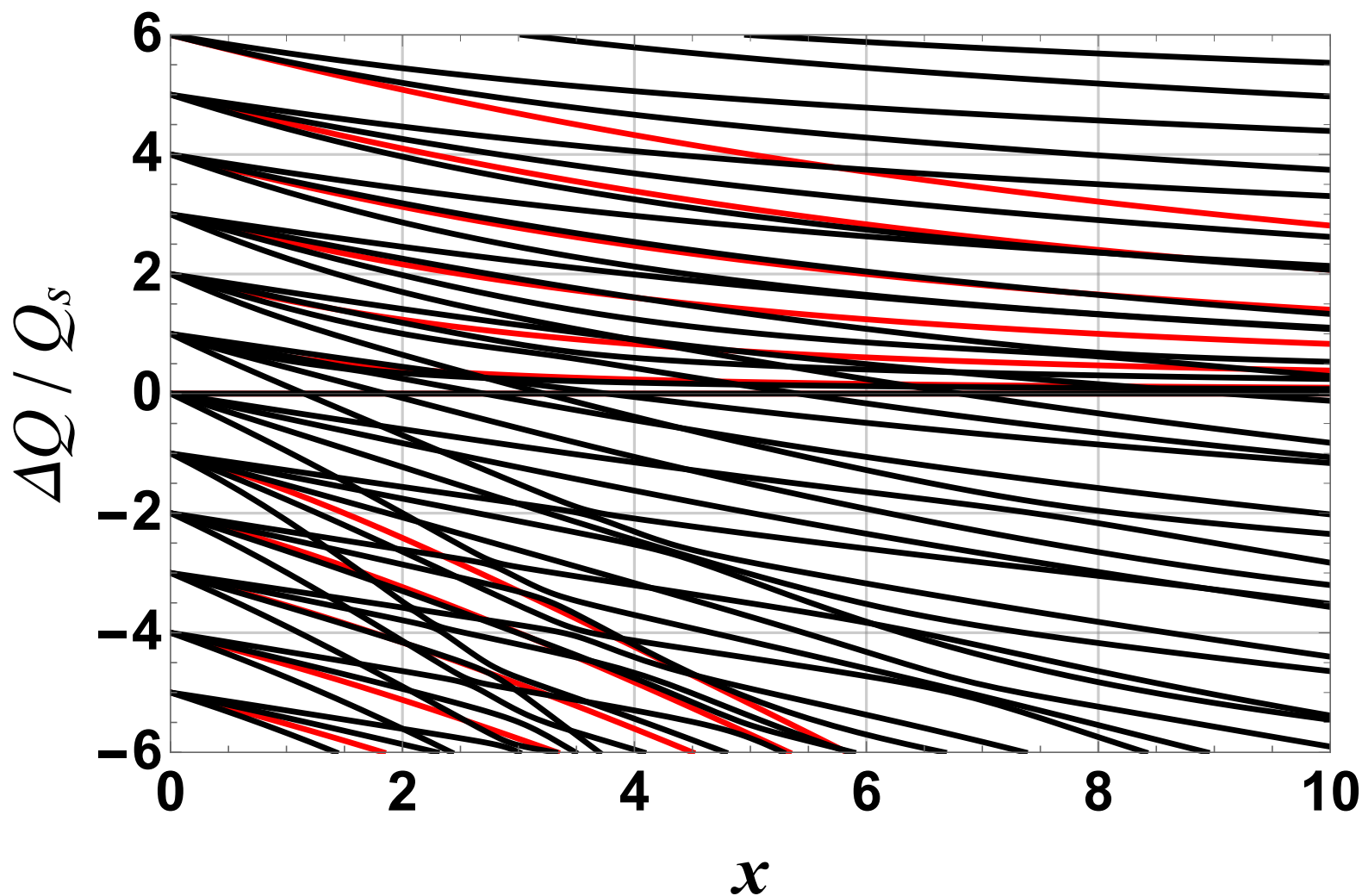
GAUSSIAN: $m_{max} = 7$



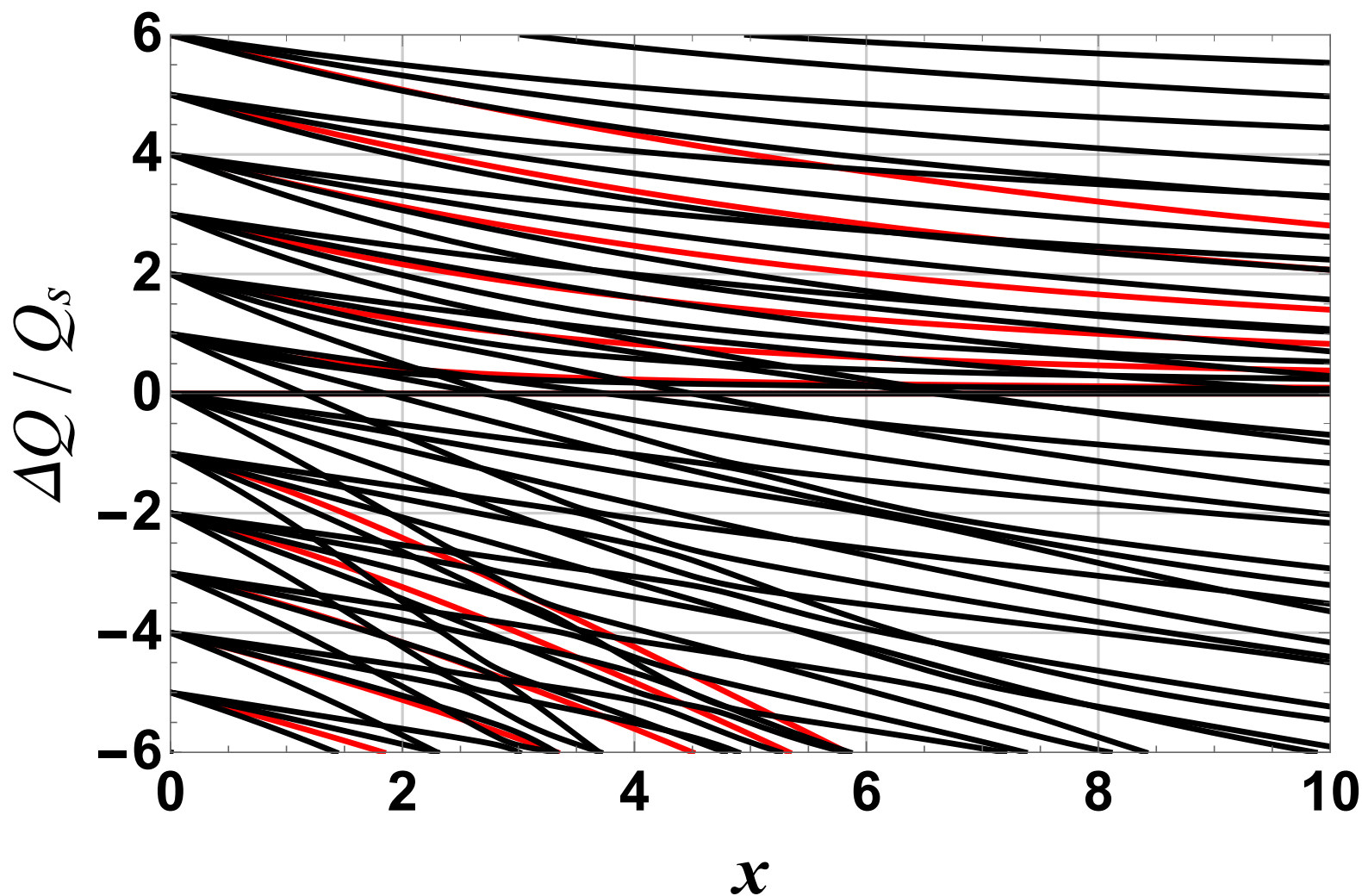
GAUSSIAN: $m_{max} = 8$



GAUSSIAN: $m_{max} = 9$

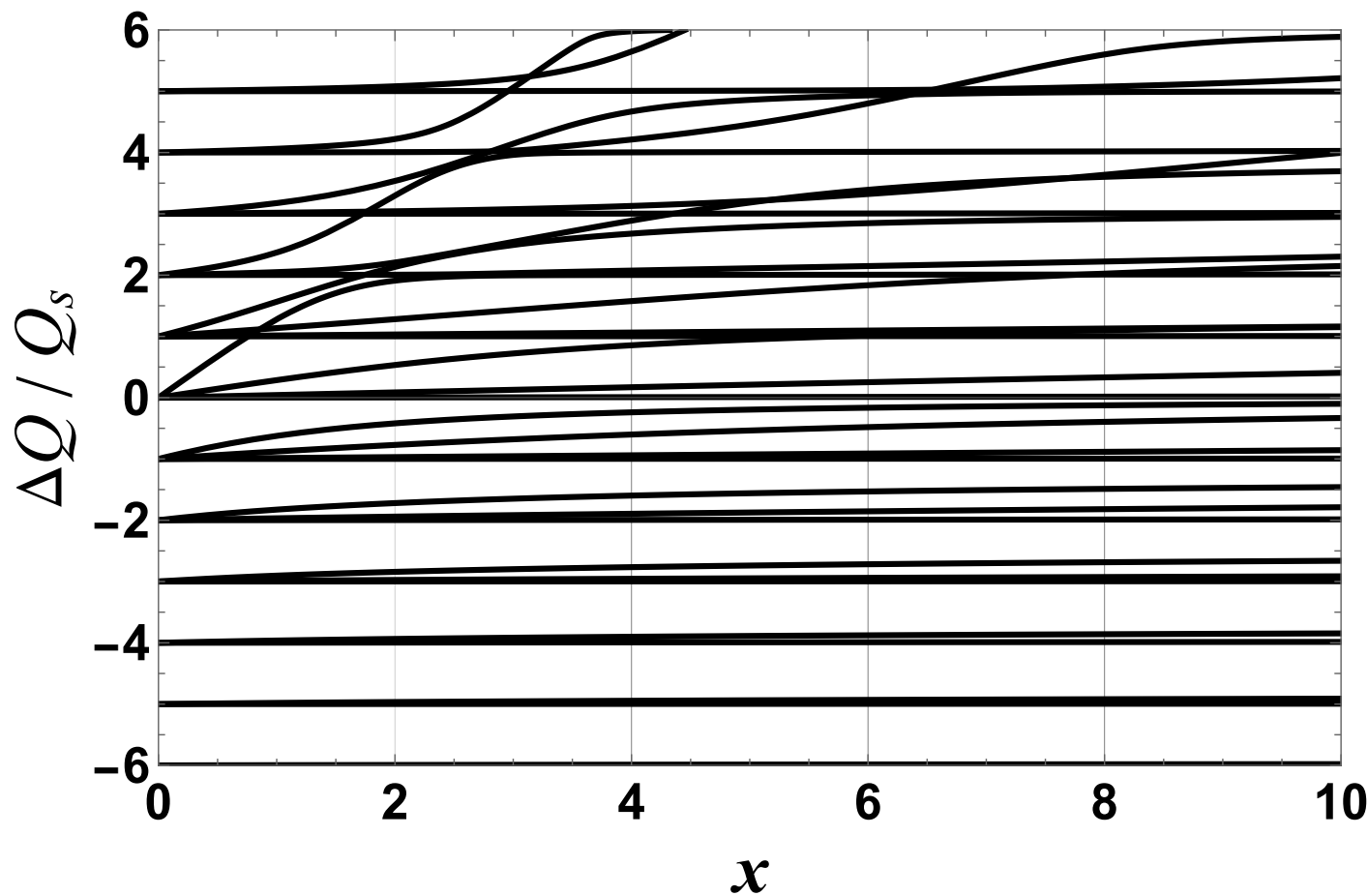


GAUSSIAN: $m_{max} = 10$



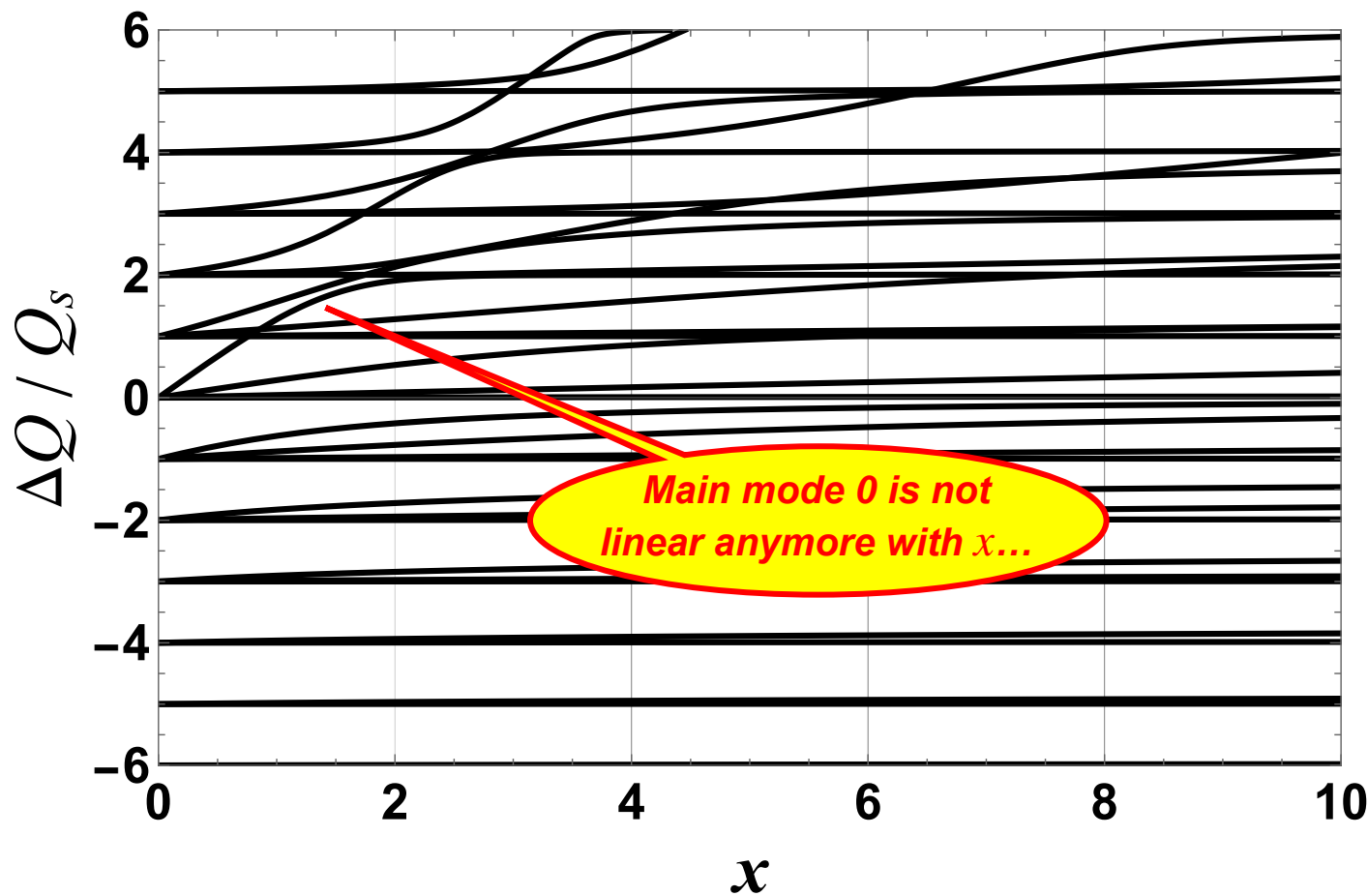
GAUSSIAN: $m_{max} = 7$

◆ Dipolar force only



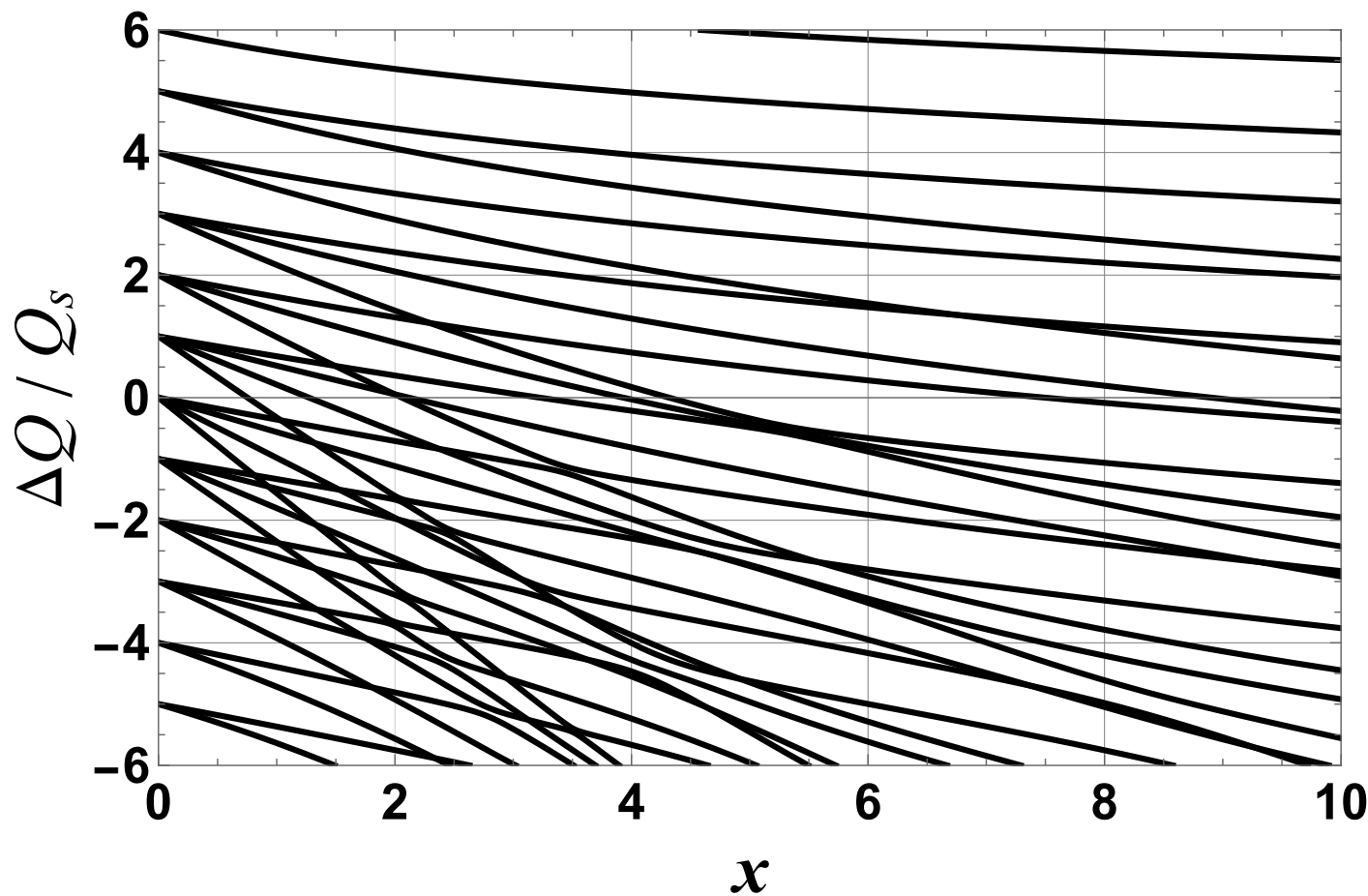
GAUSSIAN: $m_{max} = 7$

◆ Dipolar force only



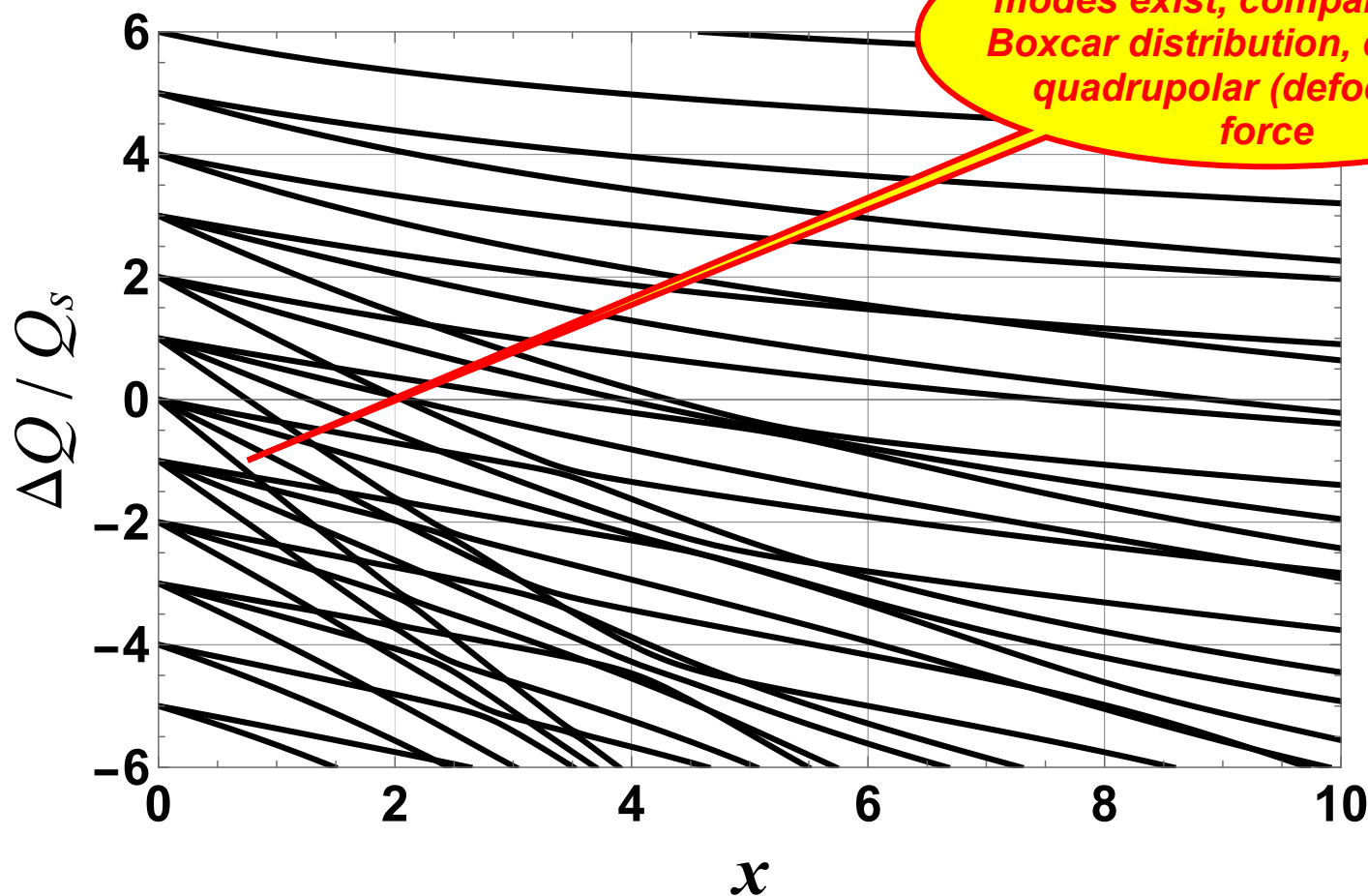
GAUSSIAN: $m_{max} = 7$

◆ Quadrupolar force only



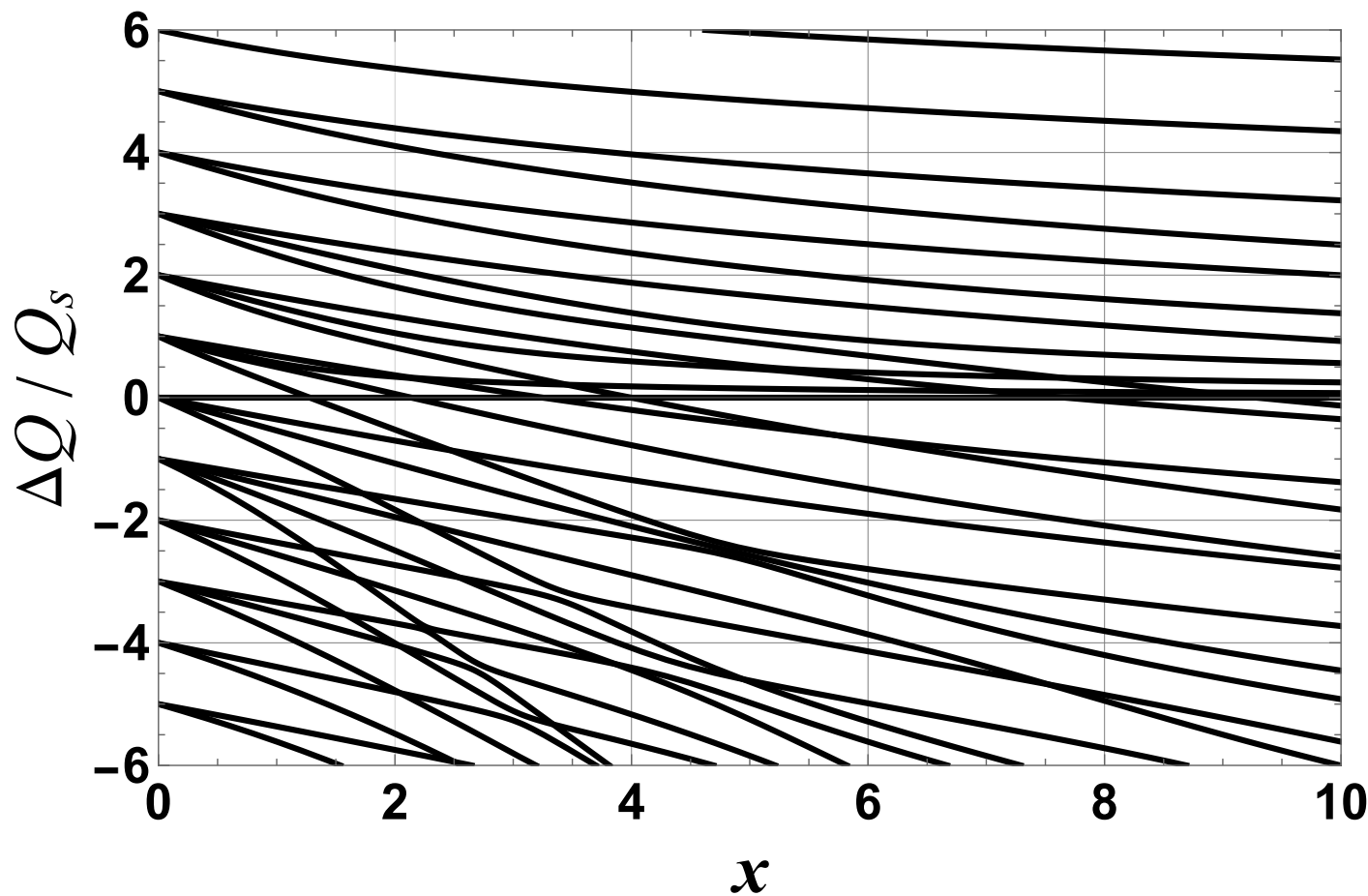
GAUSSIAN: $m_{max} = 7$

◆ Quadrupolar force only



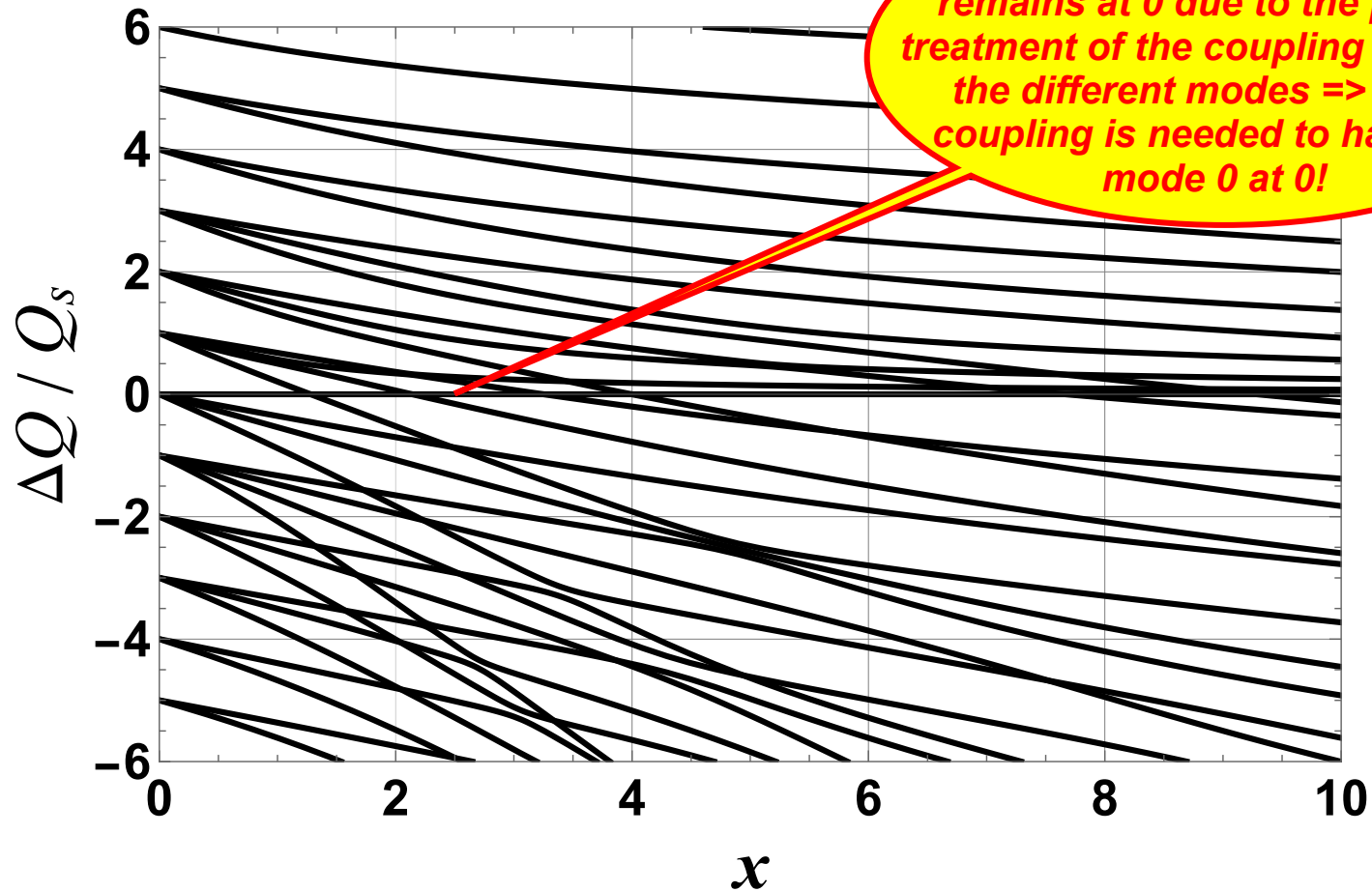
GAUSSIAN: $m_{max} = 7$

◆ Sum of the 2 forces



GAUSSIAN: $m_{max} = 7$

◆ Sum of the 2 forces



... but the main mode 0 still remains at 0 due to the proper treatment of the coupling between the different modes => Mode coupling is needed to have the mode 0 at 0!

BENCHMARK 1: Boxcar

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- ◆ M. Blaskiewicz's solution for the Boxcar line density is (without wake field)

$$T_{p,q,n,m} = (qQ_s - \Delta Q_{sc})\delta_q^m \delta_p^n + \Delta Q_{sc} B_n^m B_p^q \frac{2\delta_{|m|+2n}^{|q|+2p}}{2(|m| + 2n) + 1} - \Delta Q_w \frac{B_n^m B_p^q}{2(|m| + 2n) + 1} \left\{ \delta_{|q|+2p}^0 \delta_{|m|+2n}^0 + \frac{\delta_{|q|+2p}^{|m|+2n+1} - \delta_{|q|+2p}^{|m|+2n-1}}{2(|q| + 2p) + 1} \right\}, \quad (12)$$

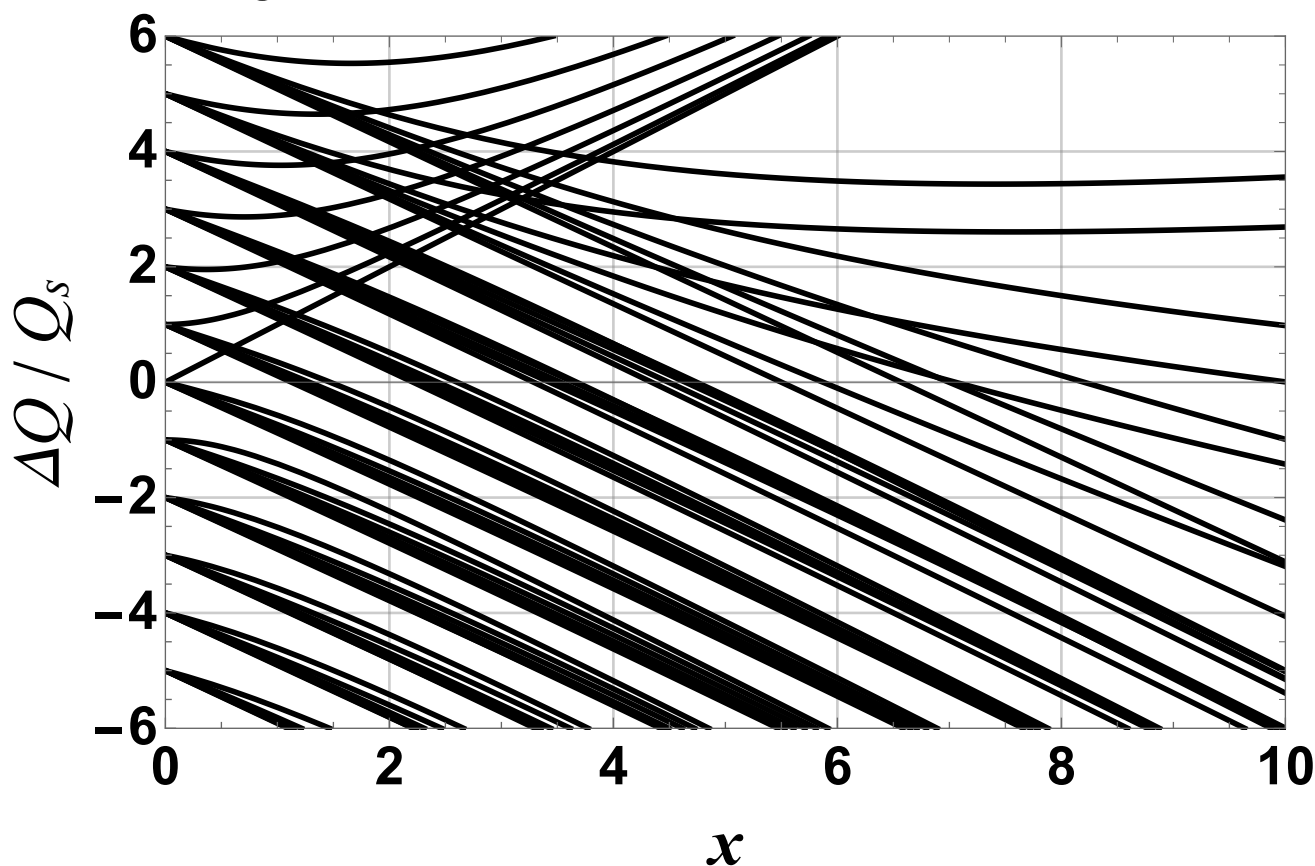
$$B_n^m = \left\{ \frac{\Gamma(n + 1/2)(2|m| + 1 + 4n)\Gamma(|m| + n + 1/2)}{\pi \Gamma(n + 1)\Gamma(|m| + n + 1)} \right\}^{1/2}$$

BENCHMARK 1: Boxcar

- ◆ But the problem is that when we plot the solutions of these equations, strange results are obtained...

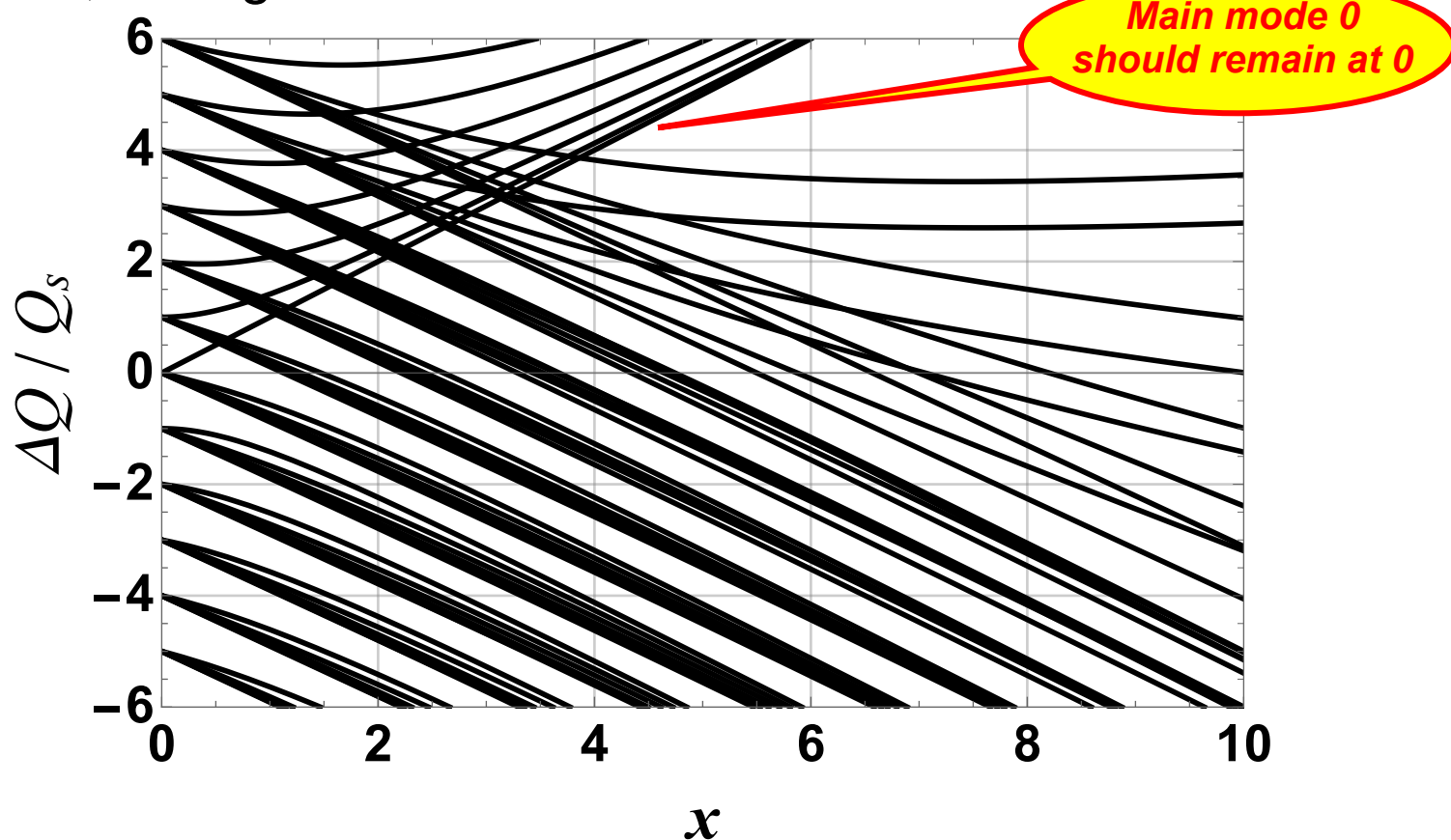
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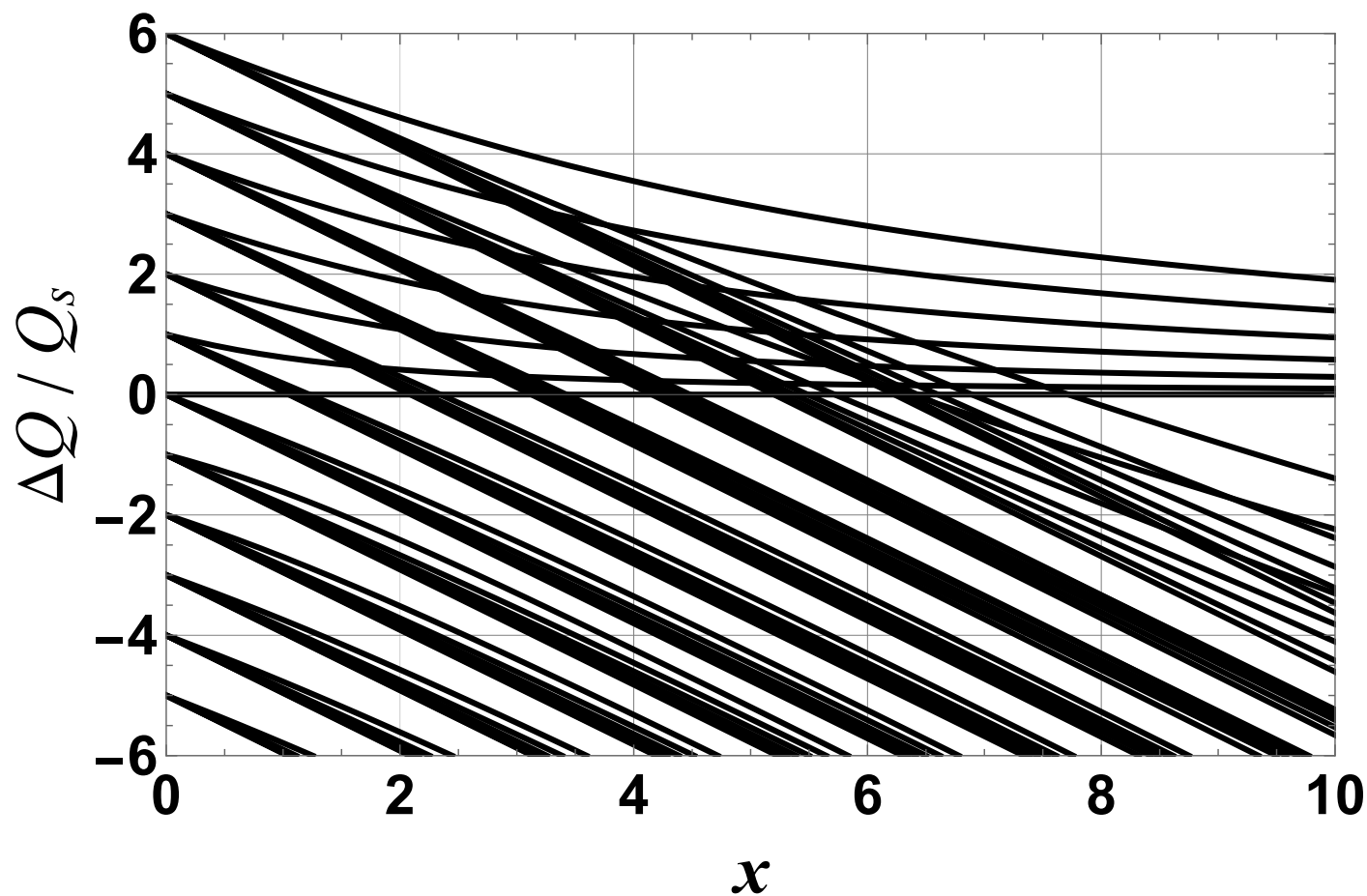
BENCHMARK 1: Boxcar

◆ Typo:

*A factor 1/2
is missing here*

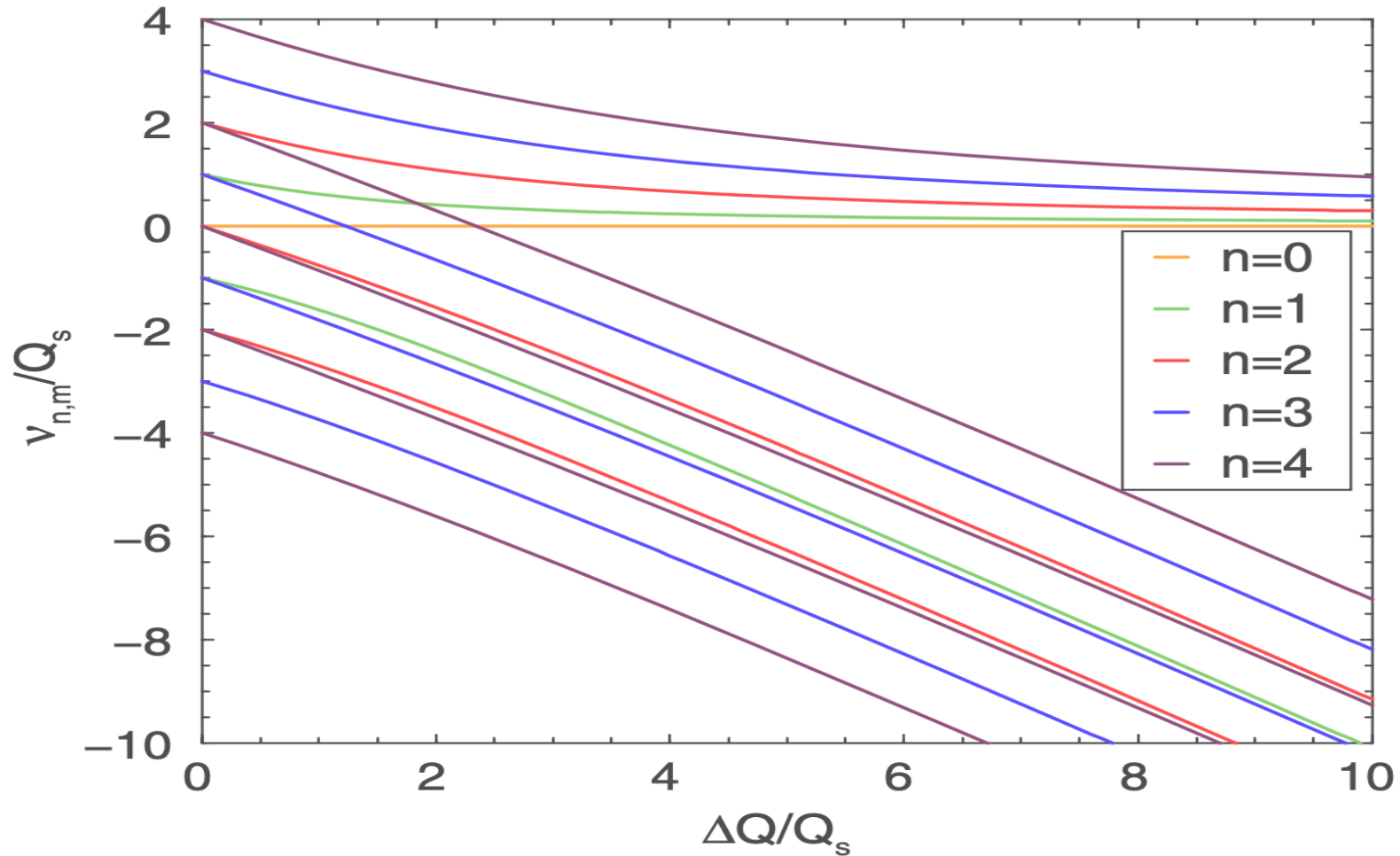
$$T_{p,q,n,m} = (qQ_s - \Delta Q_{sc})\delta_q^m \delta_p^n + \Delta Q_{sc} B_n^m B_p^q \frac{2\delta_{|m|+2n}^{|q|+2p}}{2(|m| + 2n) + 1} - \Delta Q_w \frac{B_n^m B_p^q}{2(|m| + 2n) + 1} \left\{ \delta_{|q|+2p}^0 \delta_{|m|+2n}^0 + \frac{+\delta_{|q|+2p}^{|m|+2n+1} - \delta_{|q|+2p}^{|m|+2n-1}}{2(|q| + 2p) + 1} \right\}, \quad (12)$$

BENCHMARK 1: Boxcar



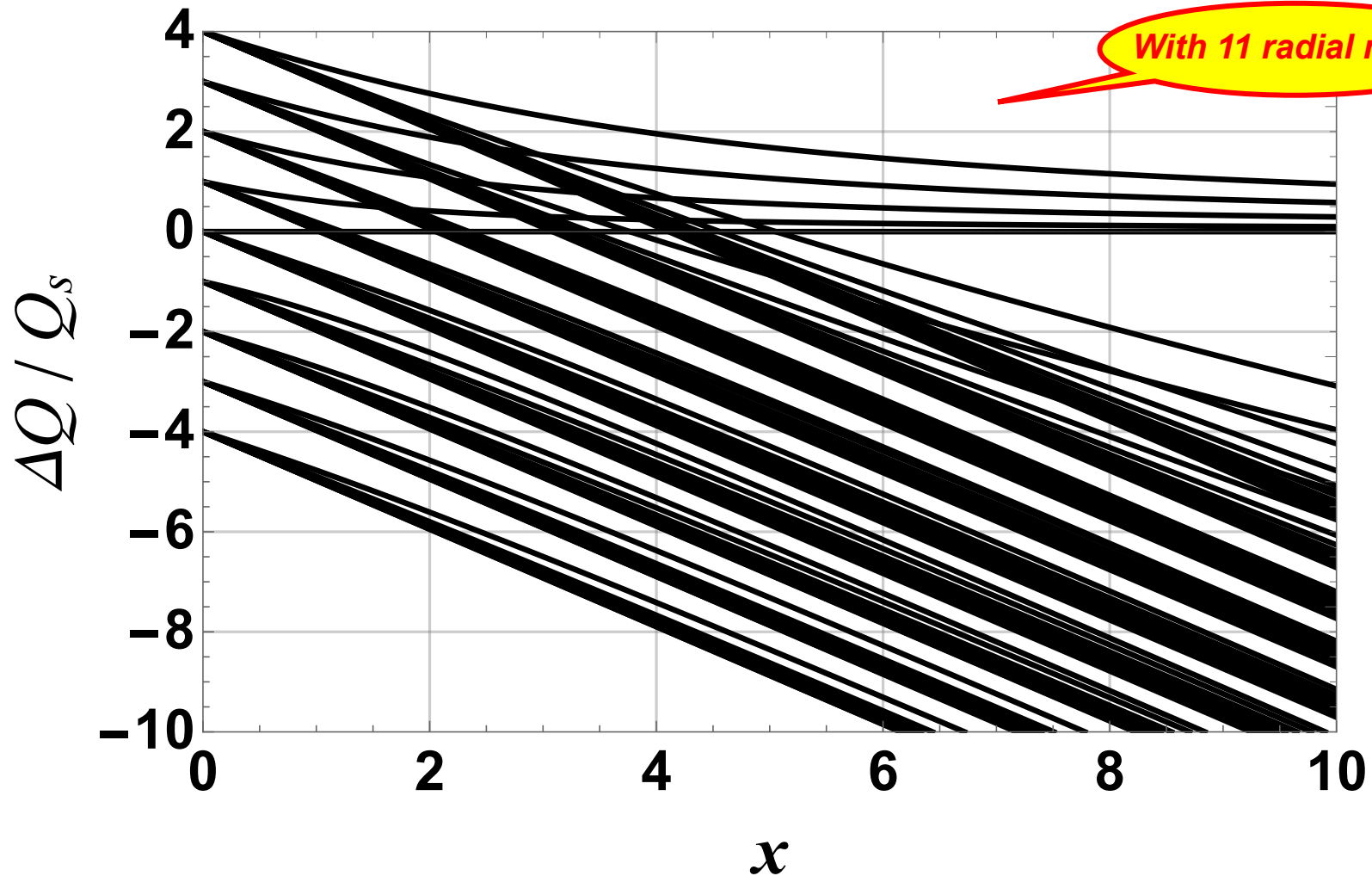
Comparison to the Boxcar model from V. Balbekov in 2017

(<https://journals.aps.org/prab/pdf/10.1103/PhysRevAccelBeams.20.034401>)



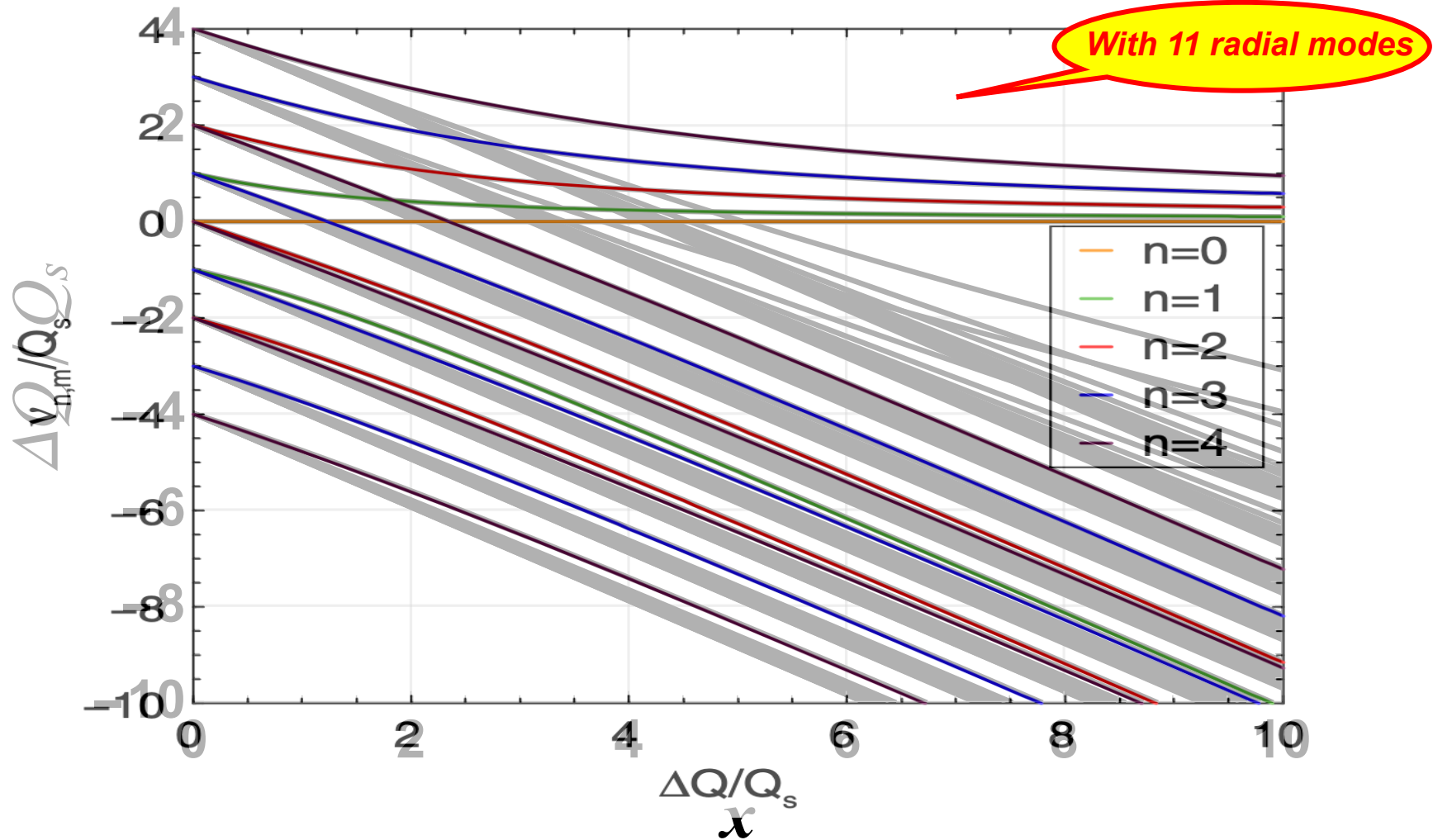
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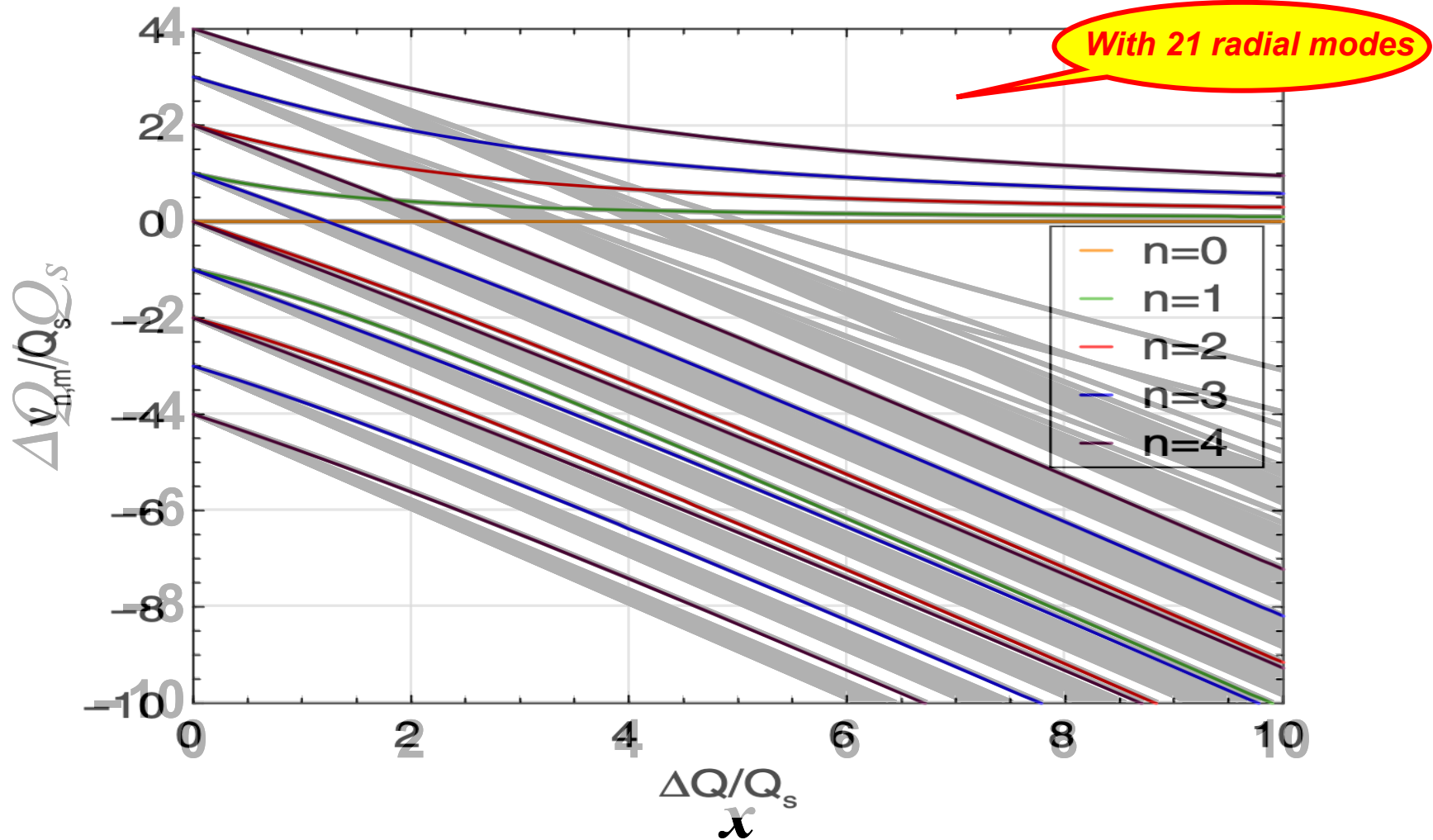
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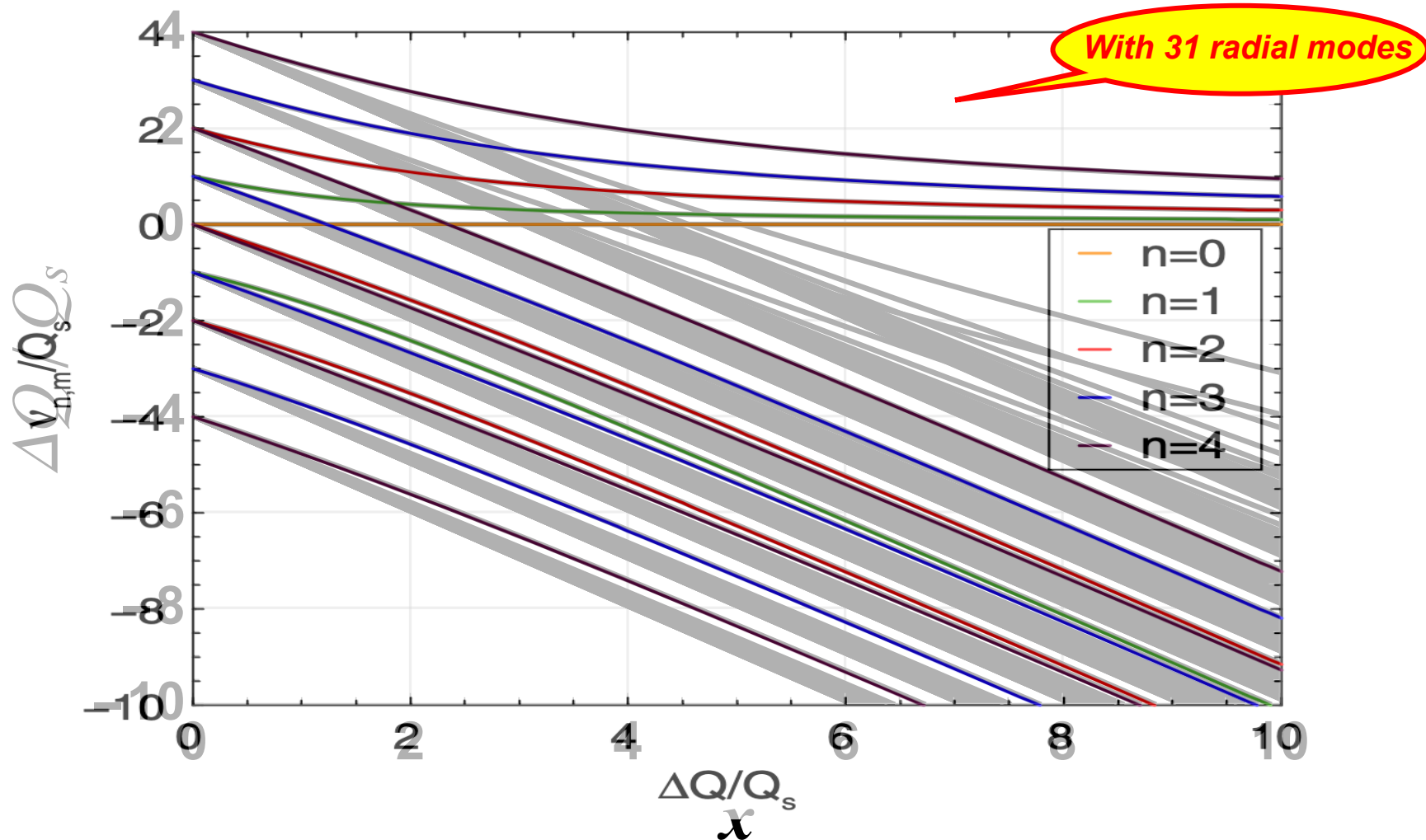
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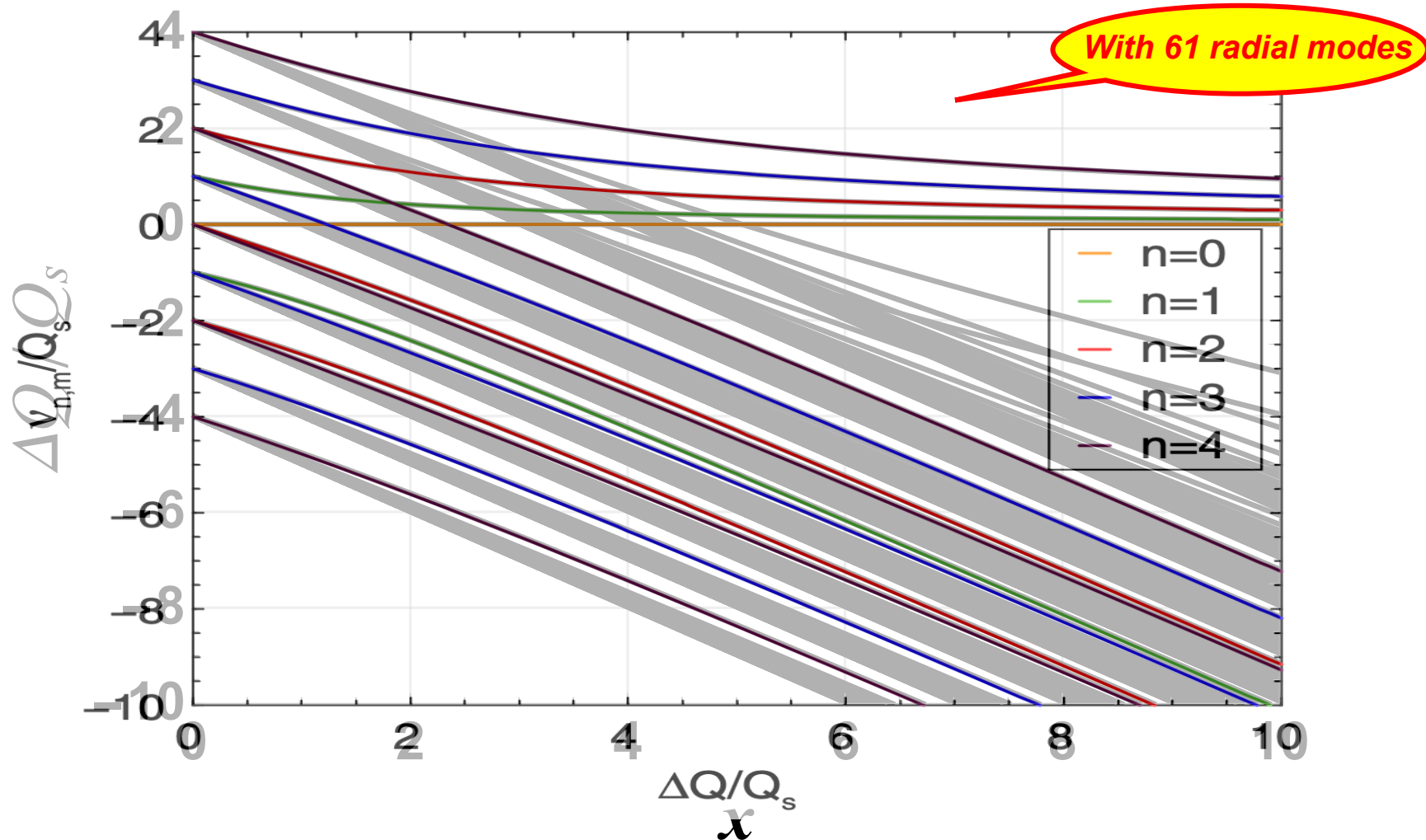
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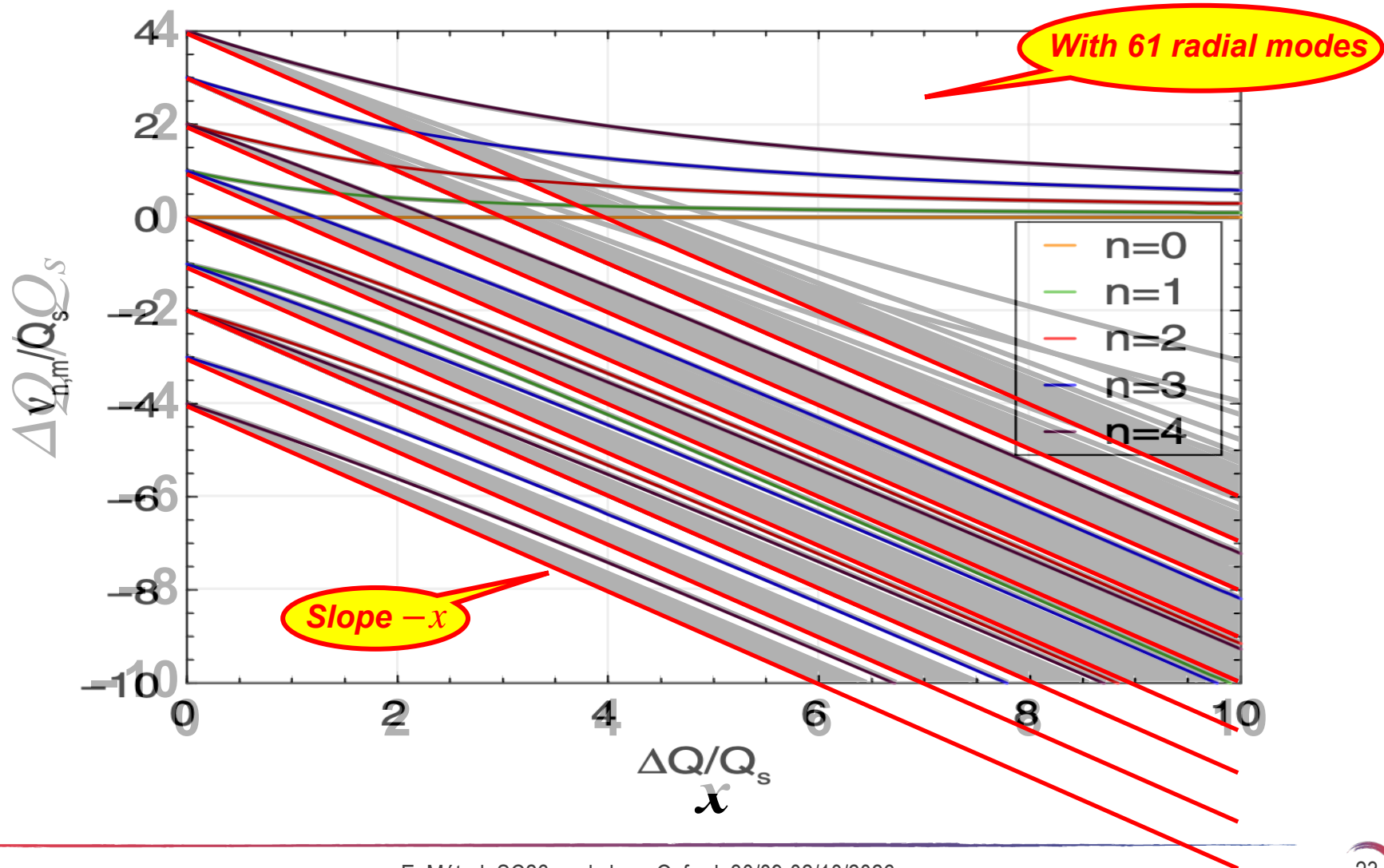
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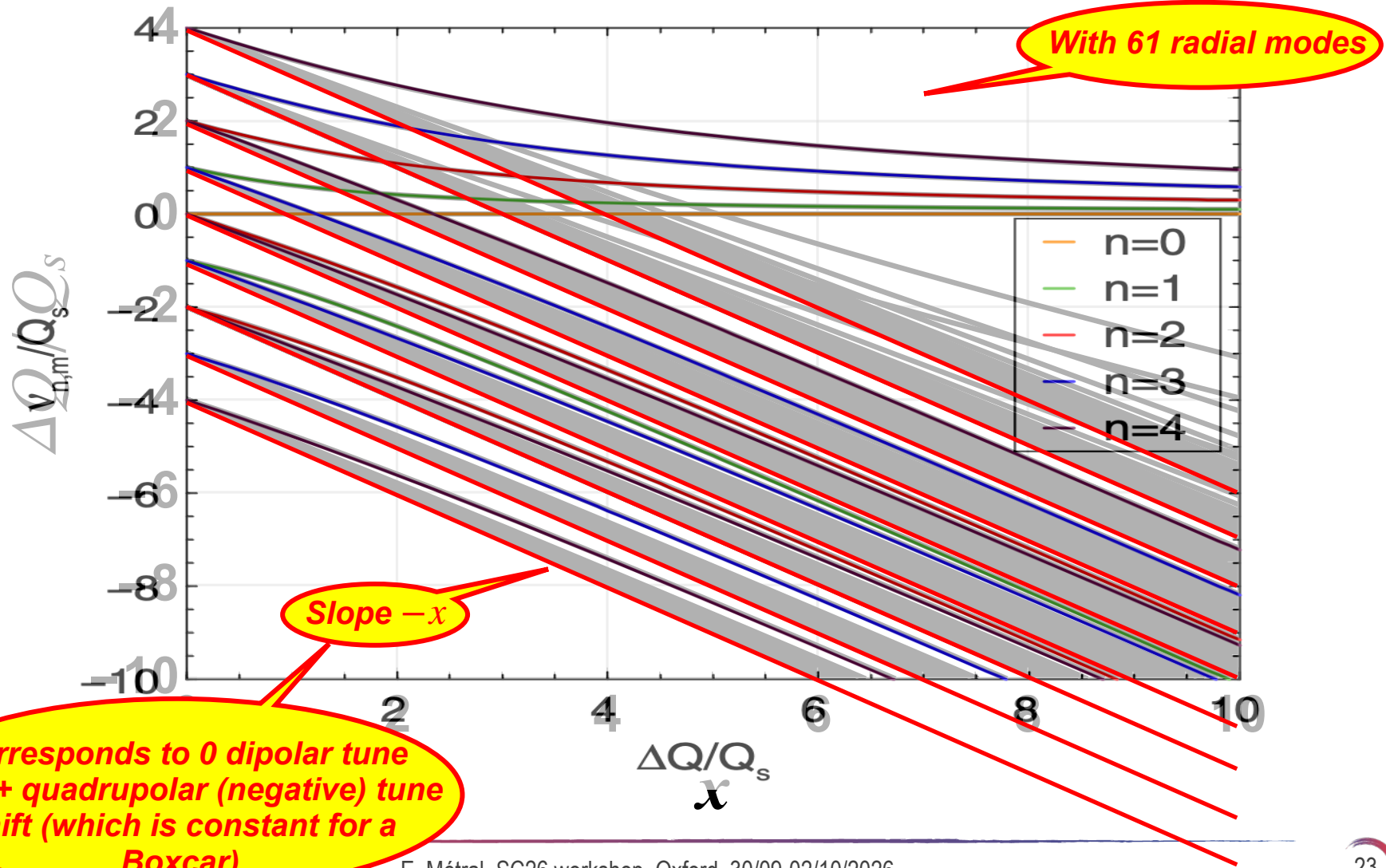
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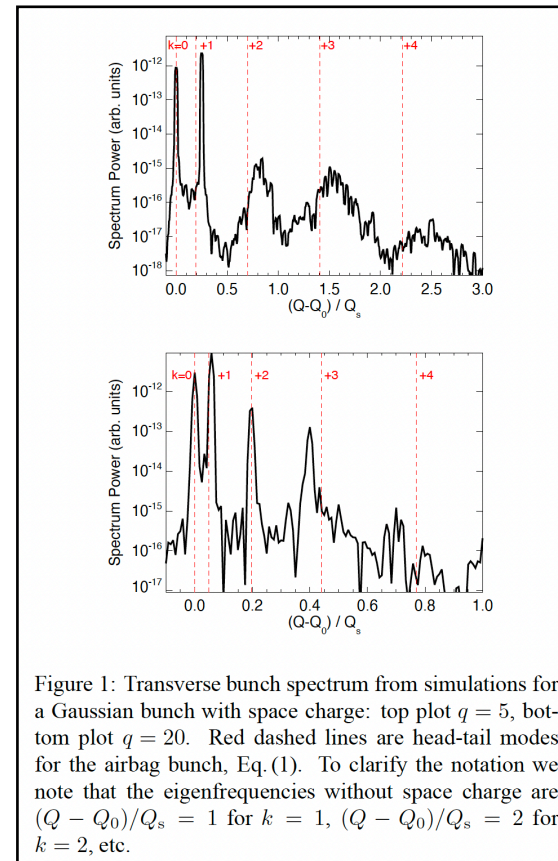
BENCHMARK 2: Gaussian

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- ◆ The main azimuthal modes 0, +1, +2, +3 and +4 were checked in 2010 by V. Kornilov and O. Boine-Frankenheim using PATRIC and HEADTAIL (<https://proceedings.jacow.org/HB2010/papers/weo1a02.pdf>)

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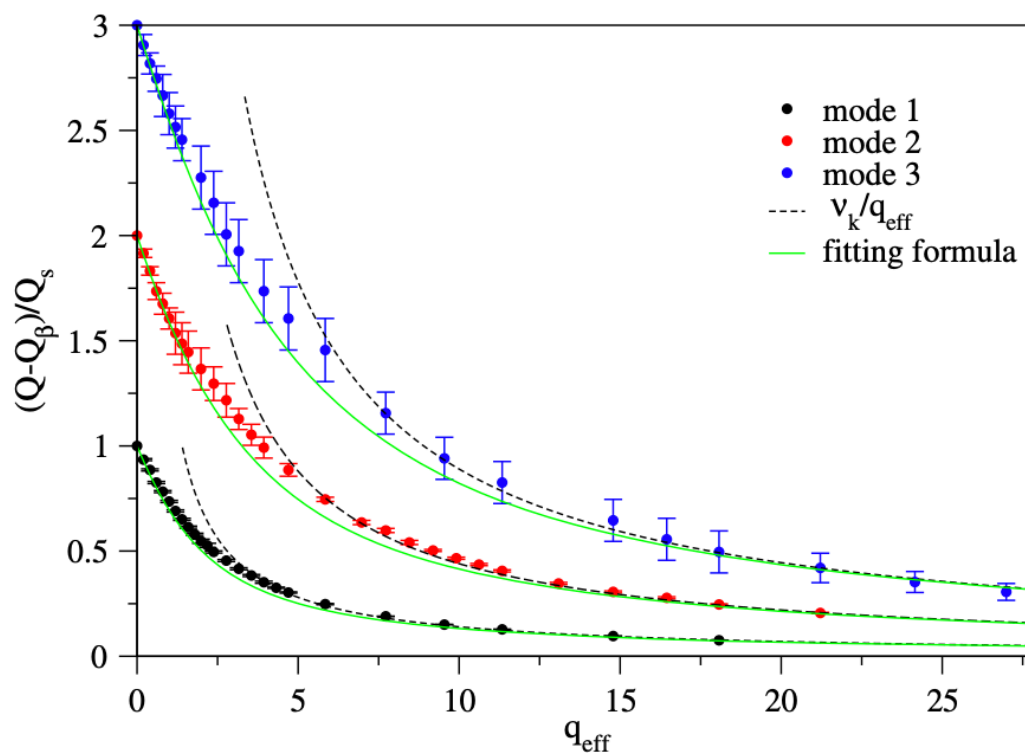
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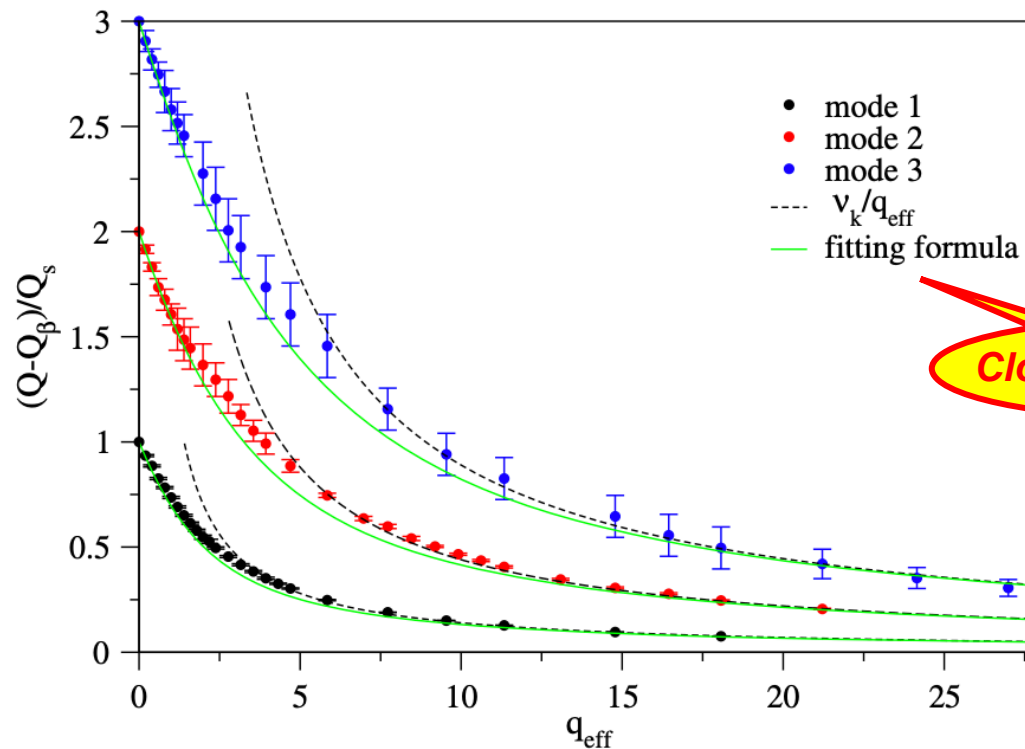
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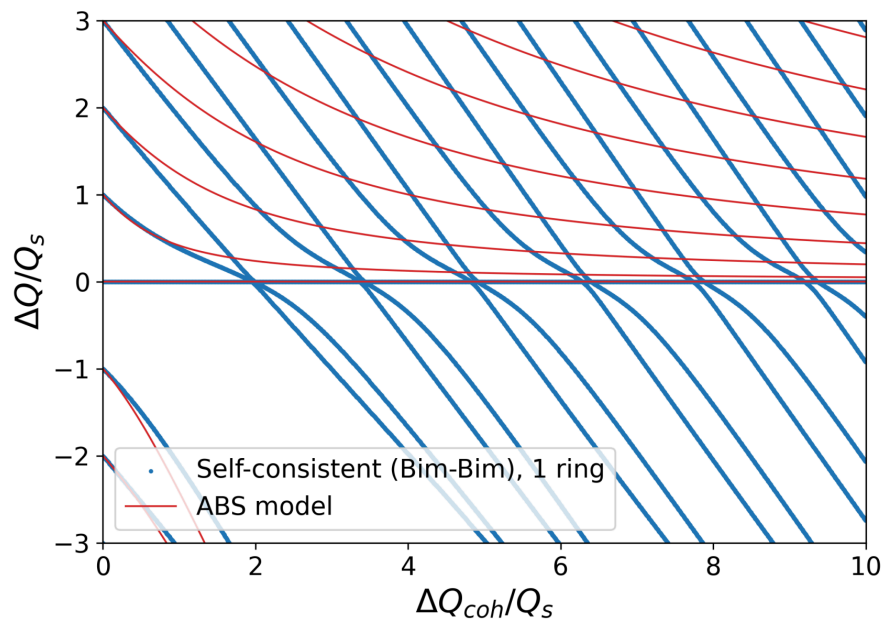
Close to ABS model

BENCHMARK 4: Gaussian

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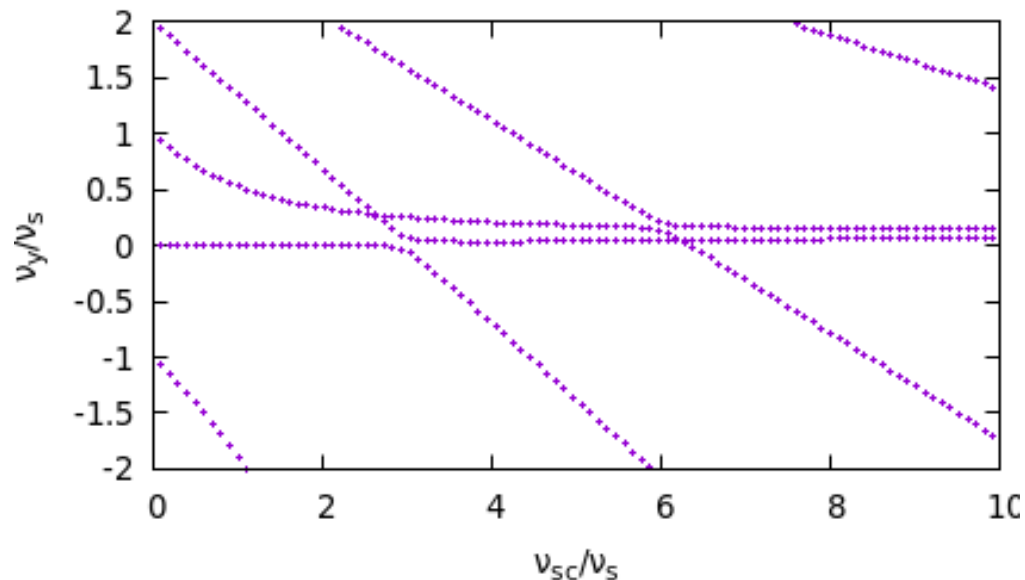
BimBim's (2021) result with 0 radial mode

(<https://accelconf.web.cern.ch/ipac2024/pdf/MOPS35.pdf>)



Ohmi's (2025) result with 0 radial mode

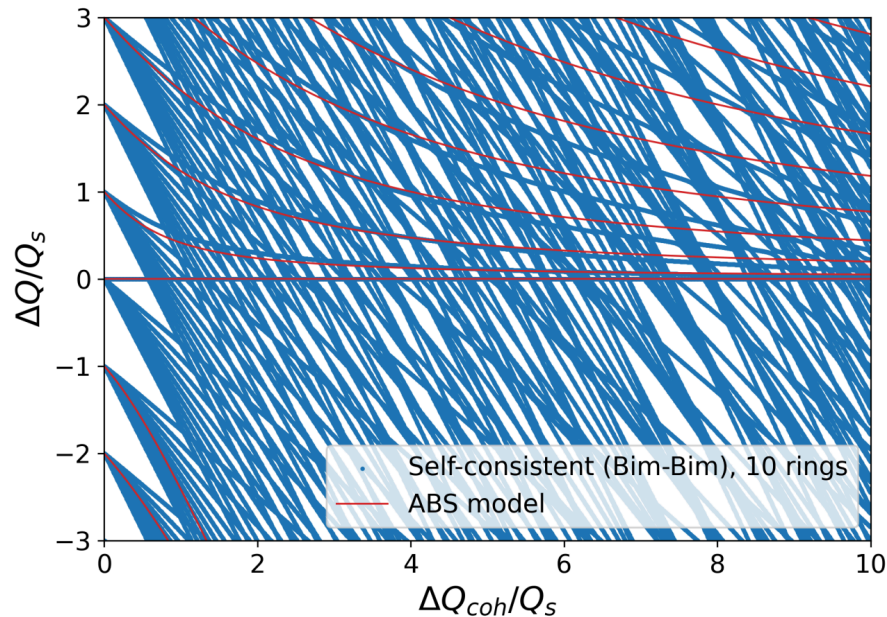
(https://indico.global/event/9305/contributions/90661/attachments/41536/77818/bbscmc_ohmi_slide.pdf)



BENCHMARK 4: Gaussian

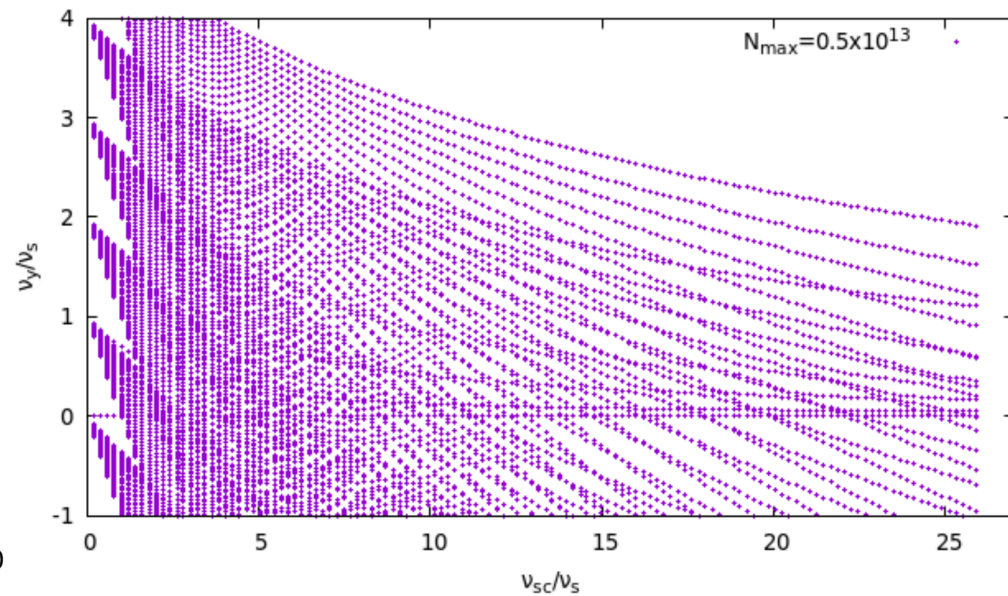
BimBim's (2021) result
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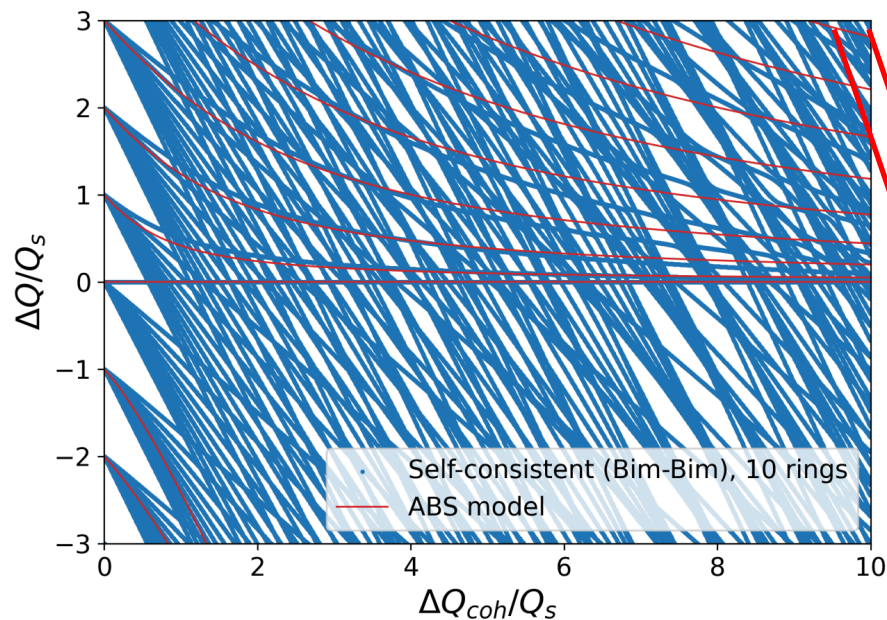
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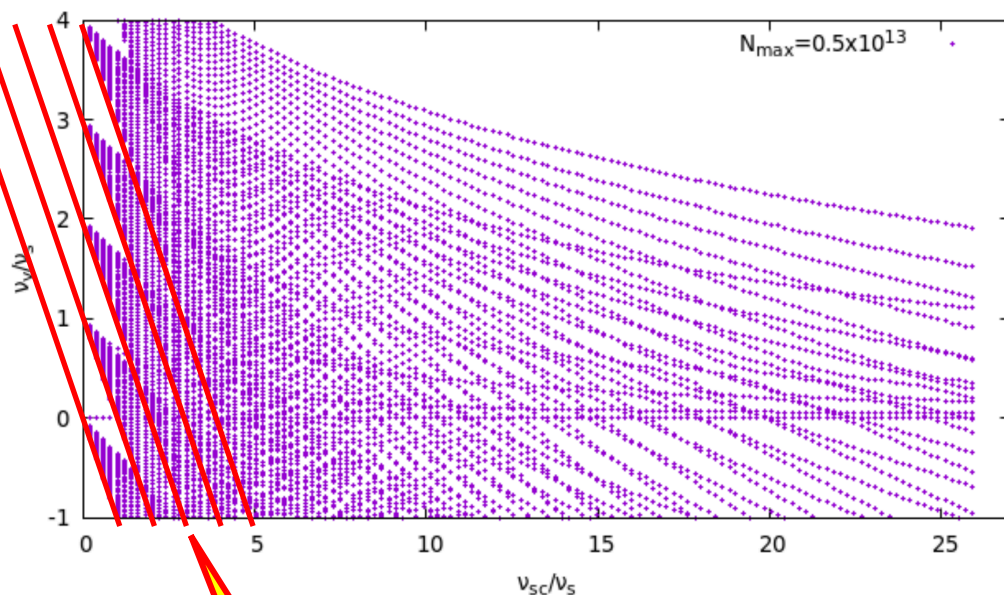
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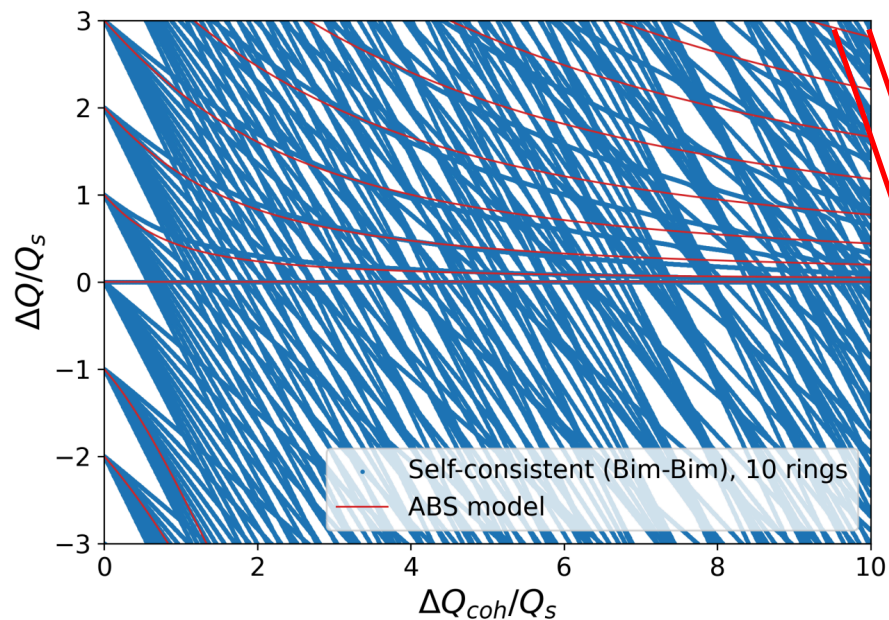
Slope $-x$

BENCHMARK 4: Gaussian

See also [D. Amorim et al., IPAC2025](#)

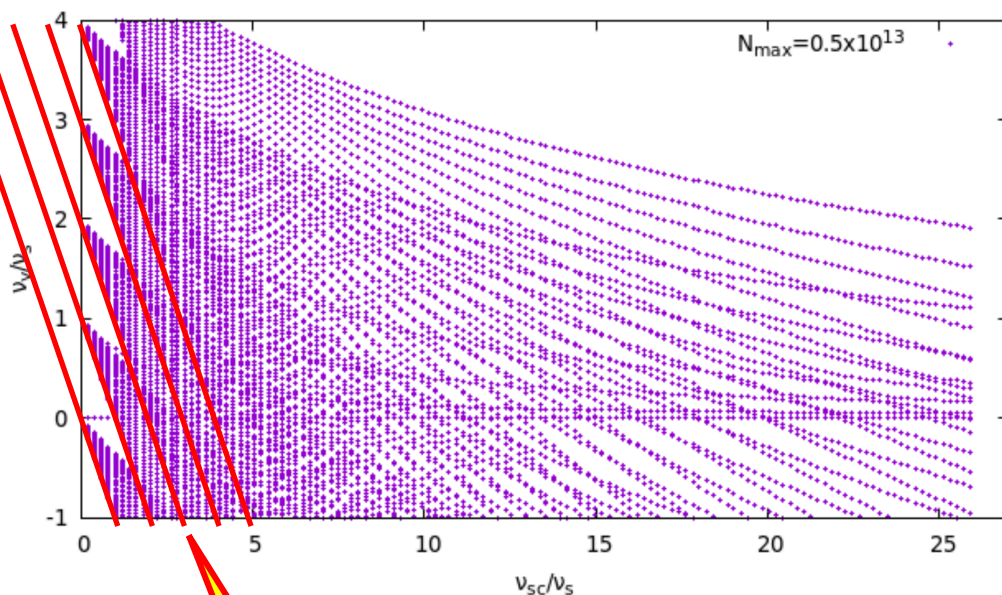
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 - * All the radial modes are grouped together with a slope close to $-x$

Conclusions

- ◆ GAUSSIAN distribution

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◆ GAUSSIAN distribution

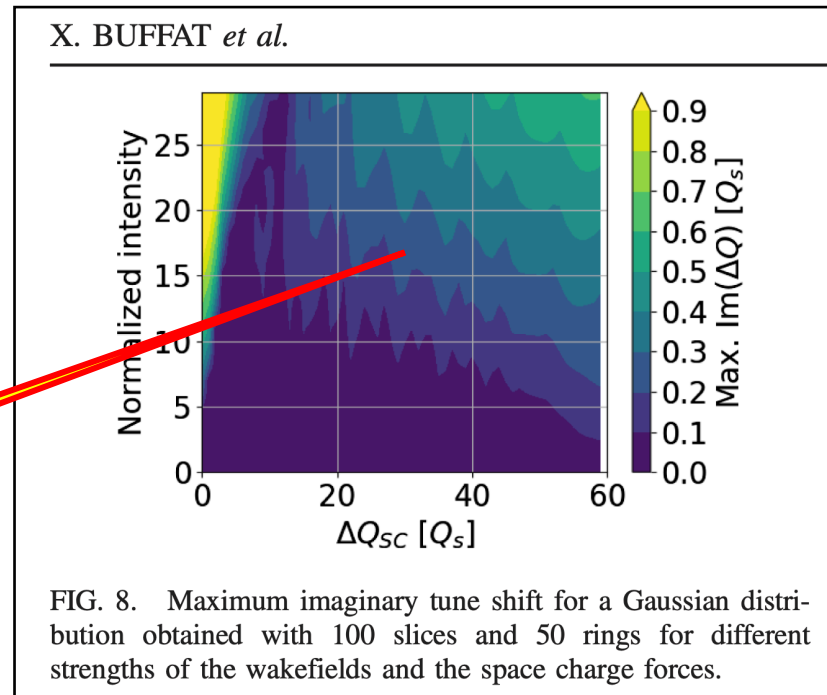
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Due to radial modes!



Thank you for your attention!



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