

Space Charge Modelling of Long Bunches and Data Driven Approaches

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Outline

- Two and a half dimensional space-charge solver for long bunches
- Neural network surrogate model for space-charge effects
- Differentiable space-charge model
- This session's talks

Space-Charge Model: from Maxwell Equations to Poisson Equations

Maxwell Equation with Lorenz Gauge

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{\partial^2 \phi}{\partial s^2} - \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} = -\frac{\rho(x, y, s, t)}{\epsilon_0}$$

$$\frac{\partial^2 A_s}{\partial x^2} + \frac{\partial^2 A_s}{\partial y^2} + \frac{\partial^2 A_s}{\partial s^2} - \frac{1}{c^2} \frac{\partial^2 A_x}{\partial t^2} = -\frac{J_s(x, y, s, t)}{c^2 \epsilon_0}$$

$$\rho = \rho(x, y, s - v_0 t), \quad \mathbf{J} = \mathbf{J}(x, y, s - v_0 t), \quad z = s - v_0 t,$$

Space-Charge Model: Poisson-like Equation

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2} + \frac{1}{\gamma^2} \frac{\partial^2 \phi}{\partial z^2} = -\frac{\rho(x, y, z)}{\epsilon_0}$$

$$\frac{\partial^2 A_s}{\partial x^2} + \frac{\partial^2 A_s}{\partial y^2} + \frac{1}{\gamma^2} \frac{\partial^2 A_s}{\partial z^2} = -\frac{J_s(x, y, z)}{c^2 \epsilon_0}$$

$$A_s(x, y, z) = \frac{\beta}{c} \phi(x, y, z)$$

2.5D Space-Charge Model for Long Bunch

long bunch: $\frac{\partial^2 \phi}{\partial x^2} \gg \frac{1}{\gamma^2} \frac{\partial^2 \phi}{\partial z^2}$ and $\frac{\partial^2 \phi}{\partial y^2} \gg \frac{1}{\gamma^2} \frac{\partial^2 \phi}{\partial z^2}$

zeroth-order solution:

$$\frac{\partial^2 \phi_0}{\partial x^2} + \frac{\partial^2 \phi_0}{\partial y^2} = -\frac{\rho(x, y, z)}{\epsilon_0} \quad \rho(x, y, z) = n(x, y)\lambda(z).$$

$$\frac{\partial^2 \phi_0}{\partial x^2} + \frac{\partial^2 \phi_0}{\partial y^2} = -\frac{n(x, y)\lambda(z)}{\epsilon}$$

for a transverse Gaussian density distribution, the electric potential:

$$\phi_0(x, y, z) = (\varphi(x, y) + \varphi_{00})\lambda(z)$$

$$\varphi(x, y) = \frac{1}{4\pi\epsilon_0} \int_0^\infty dt \frac{\exp\left[-\frac{x^2}{2(\sigma_x^2+t)} - \frac{y^2}{2(\sigma_y^2+t)}\right] - 1}{\sqrt{(\sigma_x^2+t)(\sigma_y^2+t)}}$$

$$\varphi_{00} = \frac{-1}{4\pi\epsilon_0} \iint \log((x^2 + y^2)/r_0^2) n(x, y) dx dy \quad \varphi_{00} = \frac{-1}{4\pi\epsilon_0} (\log 2 - \gamma_{em} + 2 \log(\frac{\sigma_x + \sigma_y}{2r_0}))$$

r_0 is reference length



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Reference Length

$$D(r_0) = \int_0^\infty (E_z^{2.5D}(0, 0, z; r_0) - E_z^{3D}(0, 0, z))^2 dz.$$

$$E_z^{2.5D}(0, 0, z; r_0) = \varphi_{00}(r_0) \frac{Qz}{\sqrt{2\pi}\sigma_z^3} \exp\left(-\frac{z^2}{2\sigma_z^2}\right),$$

$$E_z^{3D}(0, 0, z) = \frac{Qz}{4\pi\epsilon_0\sqrt{2\pi}} \int_0^\infty dt \frac{\exp\left[-\frac{z^2}{2(\sigma_z^2+t/\gamma^2)}\right]}{(\sigma_z^2+t/\gamma^2)\sqrt{(\sigma_x^2+t)(\sigma_y^2+t)(\sigma_z^2+t/\gamma^2)}}.$$

After minimizing $D(r_0)$ with respect to r_0 , one obtains

$$\log(r_0) = \log\left(\frac{\sqrt{2}(\sigma_x + \sigma_y)}{2}\right) - \frac{1}{2}\gamma_{em} + \sqrt{2}\sigma_z^3\gamma^3 \int_0^\infty dt \frac{1}{(2\sigma_z^2\gamma^2+t)^{3/2}\sqrt{(\sigma_x^2+t)(\sigma_y^2+t)}}.$$

Transverse open:

$$r_0 = 4e^{(-0.5\gamma_{em}-1)}\gamma\sigma_z \sim 1.1\gamma\sigma_z.$$

Transverse conducting pipe:

$$r_0 = a.$$

$$\varphi_{00} \approx \frac{-1}{4\pi\epsilon_0} \left(\log 2 - \gamma_{em} + 2 \log\left(\frac{\sigma}{a}\right) \right)$$

Symplectic 2.5D Space-Charge Model for a Transverse Gaussian Density Function

$$p_x(\tau) = p_x(0) + \tau K \frac{\lambda(z)}{4\pi\epsilon_0} \int_0^\infty dt \frac{x \exp\left[-\frac{x^2}{2(\sigma_x^2+t)} - \frac{y^2}{2(\sigma_y^2+t)}\right]}{(\sigma_x^2+t)\sqrt{(\sigma_x^2+t)(\sigma_y^2+t)}}$$

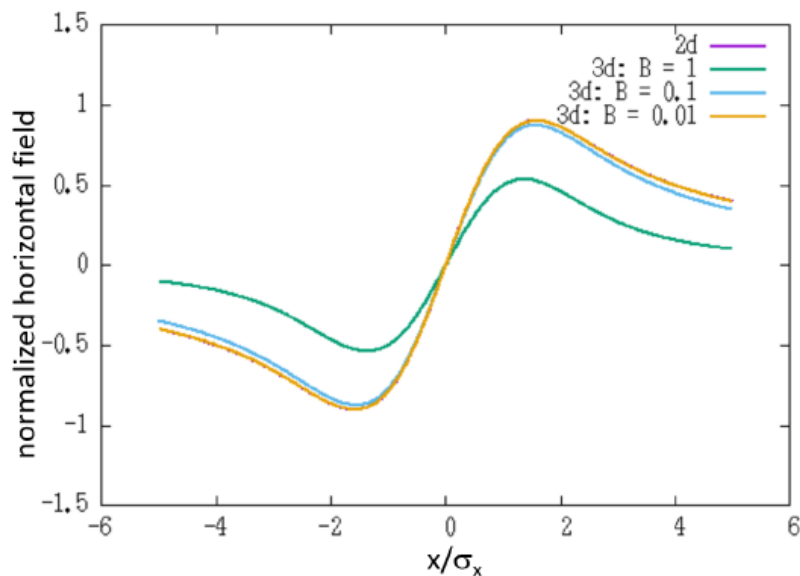
$$p_y(\tau) = p_y(0) + \tau K \frac{\lambda(z)}{4\pi\epsilon_0} \int_0^\infty dt \frac{y \exp\left[-\frac{x^2}{2(\sigma_x^2+t)} - \frac{y^2}{2(\sigma_y^2+t)}\right]}{(\sigma_y^2+t)\sqrt{(\sigma_x^2+t)(\sigma_y^2+t)}}$$

$$p_z(\tau) = p_z(0) - \tau K (\varphi(x, y) + \varphi_{00}) \frac{\partial \lambda(z)}{\partial z}$$

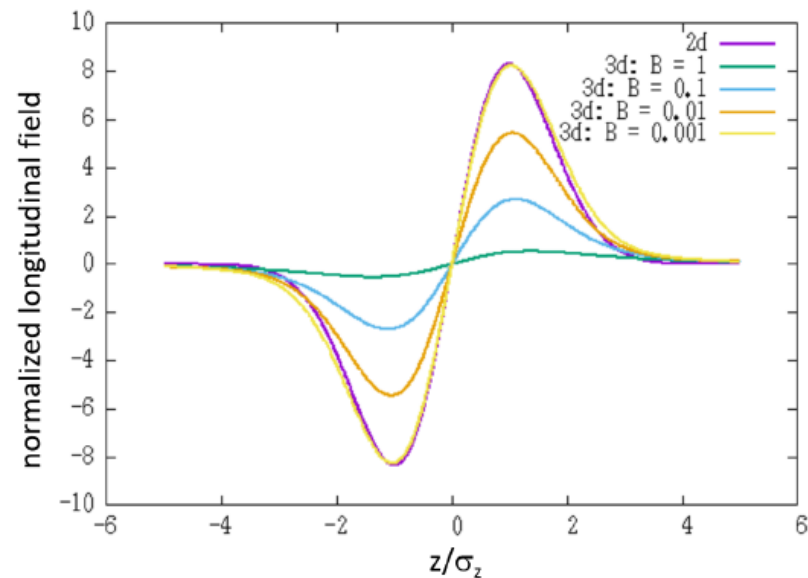
- Symplectic 2.5D space-charge model differs from previous 2.5D model by including the longitudinal space-charge kick

Comparison of Electric Fields from the 3D Gaussian and the 2.5D Gaussian Solvers

horizontal electric field

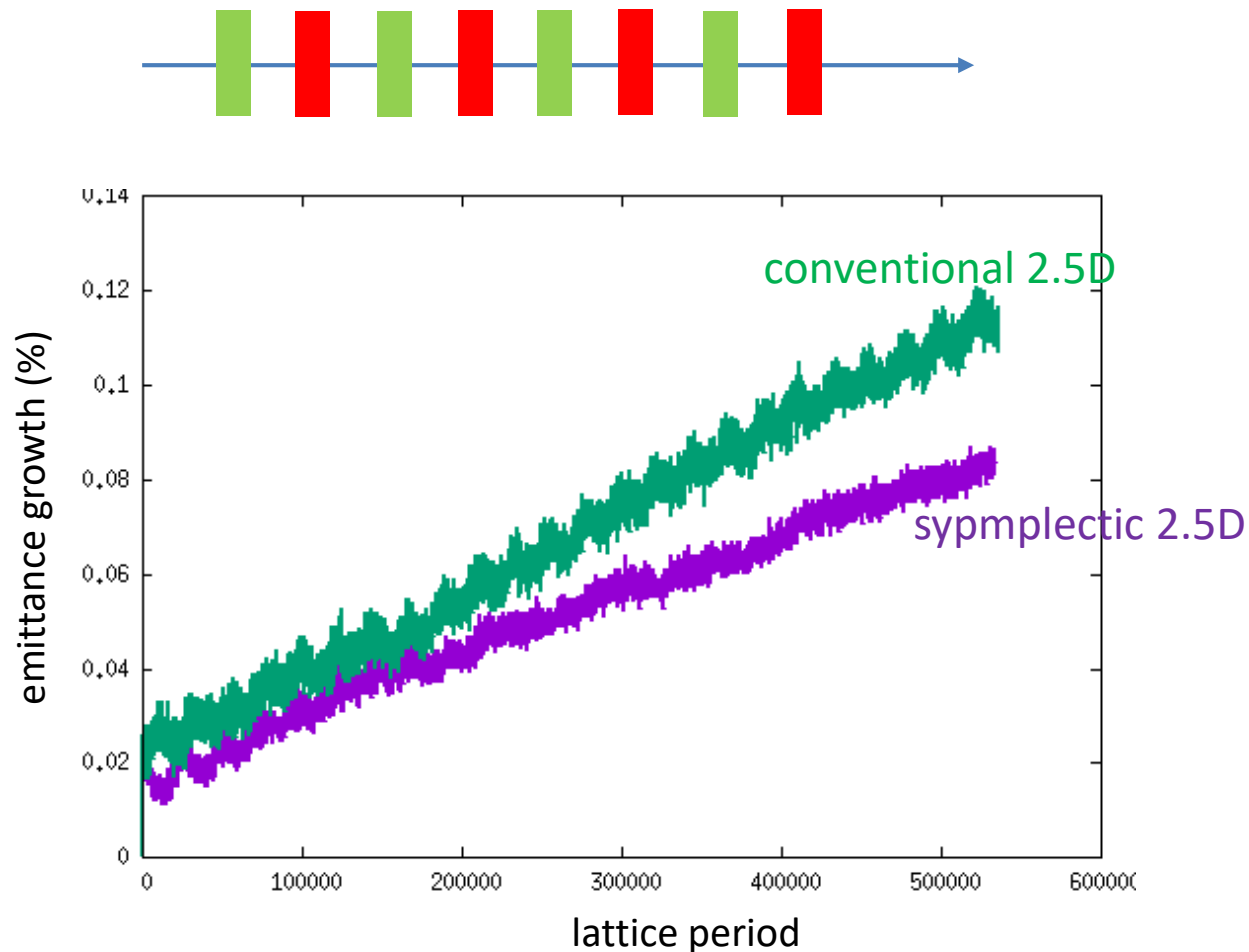


longitudinal electric field



- For a short bunch length (large transverse to longitudinal ratio), there exists significant discrepancy between the 3D and the 2.5D solver
- For a long bunch (small transverse to longitudinal ratio), the 2.5D solver agrees with the 3D solver quite well

Four Dimensional Emittance Growth from the symplectic 2.5D and conventional 2.5D Space-Charge Gaussian Model



- Symplectic 2.5D space-charge solver shows slower emittance growth.

2.5D Symplectic Space-Charge Model with Arbitrary Open Transverse Density Function

$$\varphi(x, y) = \frac{1}{2\pi\epsilon_0} \int \int G(x - x', y - y') n(x', y') dx' dy' \quad G(x, y) = -\frac{1}{2} \log(x^2 + y^2)$$

$$\bar{G}(x_i - x_{i'}, y_j - y_{j'}) = \int_{x_{i'} - h_x/2}^{x_{i'} + h_x/2} dx' \int_{y_{j'} - h_y/2}^{y_{j'} + h_y/2} dy' G(x_i - x', y_j - y')$$

$$\varphi(x_i, y_j) = \frac{1}{2\pi\epsilon_0} \sum_{i'=1}^{N_x} \sum_{j'=1}^{N_y} \bar{G}(x_i - x_{i'}, y_j - y_{j'}) n(x_{i'}, y_{j'})$$

$$p_x(\tau) = p_x(0) - \tau K \lambda(z) \sum_i \sum_j \frac{\partial S(x_i - x)}{\partial x} S(y_j - y) \varphi(x_i, y_j)$$

$$p_y(\tau) = p_y(0) - \tau K \lambda(z) \sum_i \sum_j S(x_i - x) \frac{\partial S(y_j - y)}{\partial y} \varphi(x_i, y_j)$$

$$p_z(\tau) = p_z(0) - \tau K \frac{\partial \lambda(z)}{\partial z} \sum_i \sum_j S(x_i - x) S(y_j - y) \varphi(x_i, y_j)$$

2.5D Symplectic Space-Charge Model with Arbitrary Transverse Density Function in Rectangular Pipe

$$p_x(\tau) = p_x(0) - \tau K \lambda(z) \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \varphi^{lm} \alpha_l \cos(\alpha_l x) \sin(\beta_m y)$$

$$p_y(\tau) = p_y(0) - \tau K \lambda(z) \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \varphi^{lm} \beta_m \sin(\alpha_l x) \cos(\beta_m y)$$

$$p_z(\tau) = p_z(0) - \tau K \frac{\partial \lambda(z)}{\partial z} \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \varphi^{lm} \sin(\alpha_l x) \sin(\beta_m y)$$

$$\varphi^{lm} = \frac{n^{lm}}{\epsilon_0 \gamma_{lm}^2}$$

$$n^{lm} = \frac{4}{ab} \int_0^a \int_0^b n(x, y) \sin(\alpha_l x) \sin(\beta_m y) dx dy$$

$$\text{where } \gamma_{lm}^2 = \alpha_l^2 + \beta_m^2$$

2.5D Symplectic Space-Charge Model with Arbitrary Transverse Density Function in Round Pipe

$$p_x(\tau) = p_x(0) - \tau K \lambda(z) \sum_{m=-N_m/2}^{N_m/2-1} \sum_{l=1}^{N_l} \varphi^{lm} (\gamma_{lm} J'_m(\gamma_{lm} r)) \frac{x}{r} + i m \frac{y}{r^2} J_m(\gamma_{lm} r) \exp(-im\theta),$$

$$p_y(\tau) = p_y(0) - \tau K \lambda(z) \sum_{m=-N_m/2}^{N_m/2-1} \sum_{l=1}^{N_l} \varphi^{lm} (\gamma_{lm} J'_m(\gamma_{lm} r)) \frac{y}{r} - i m \frac{x}{r^2} J_m(\gamma_{lm} r) \exp(-im\theta),$$

$$p_z(\tau) = p_z(0) - \tau K \frac{\partial \lambda(z)}{\partial z} \sum_{m=-N_m/2}^{N_m/2-1} \sum_{l=1}^{N_l} \varphi^{lm} J_m(\gamma_{lm} r) \exp(-im\theta)$$

$$\varphi^{lm} = \frac{n^{lm}}{\epsilon_0 \gamma_{lm}^2}$$

$$n^{lm} = \frac{1}{\pi J_m'^2(\gamma_{lm})} \int_0^{2\pi} \int_0^a n(r, \theta) \exp(im\theta) r J_m(\gamma_{lm} r) dr d\theta$$

$$J_m(\gamma_{lm} a) = 0.$$

2.5D Symplectic Space-Charge Model in FS Coordinate

$$\frac{\partial^2 \phi}{\partial x^2} + \frac{h}{1+hx} \frac{\partial \phi}{\partial x} + \frac{\partial^2 \phi}{\partial y^2} + \frac{1}{\gamma^2(1+hx)^2} \frac{\partial^2 \phi}{\partial z^2} = -\frac{\rho}{\epsilon_0}$$

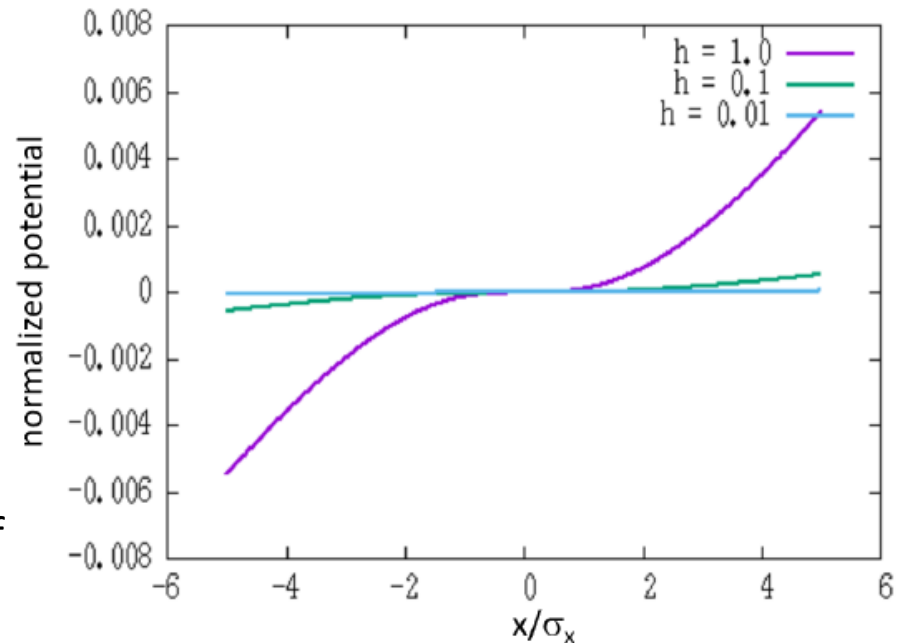
$$\frac{h}{1+hx} \frac{\partial \phi}{\partial x} \ll \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2}, \text{ and } \frac{1}{\gamma^2(1+hx)^2} \frac{\partial^2 \phi}{\partial z^2} \ll \frac{\partial^2 \phi}{\partial x^2} + \frac{\partial^2 \phi}{\partial y^2}$$

$$\frac{\partial^2 \phi_0}{\partial x^2} + \frac{\partial^2 \phi_0}{\partial y^2} = -\frac{\rho(x, y, z)}{\epsilon_0}$$

$$\frac{\partial^2 \phi_1}{\partial x^2} + \frac{\partial^2 \phi_1}{\partial y^2} = -\frac{\rho(x, y, z)}{\epsilon_0} - h \frac{\partial \phi_0}{\partial x}$$

- For a ring with radius > 100 m, the effect of curvature might be negligible.

first order horizontal electric potential



Neural Network Surrogate Space-Charge Model

$$H_2 = 4\pi \frac{K}{ab} \frac{1}{N_p} \sum_{i=1}^{N_p} \sum_{j=1}^{N_p} \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} \frac{1}{\gamma_{lm}^2} \sin(\alpha_l x_j) \sin(\beta_m y_j) \sin(\alpha_l x_i) \sin(\beta_m y_i)$$

particle distribution in x-y space

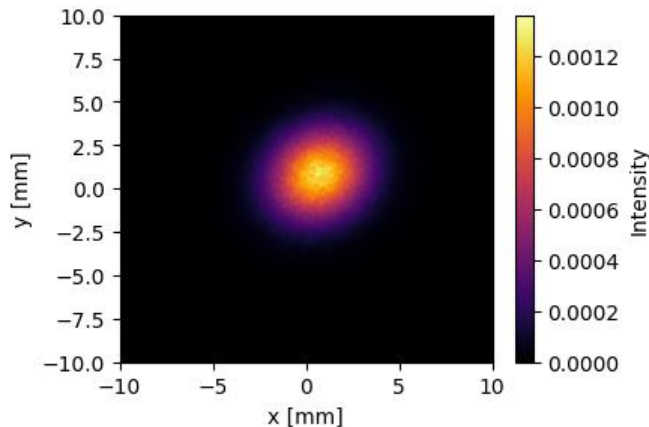
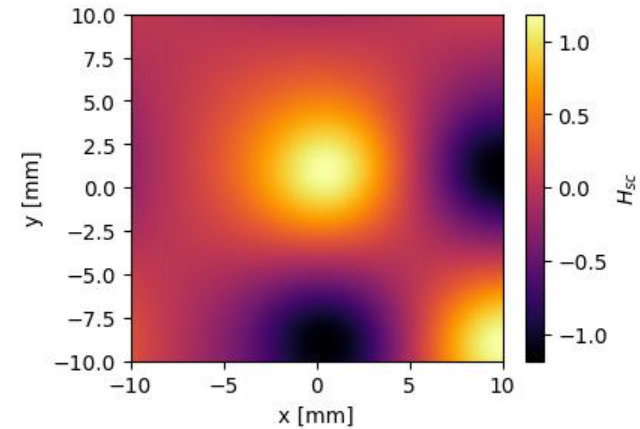


Image to image



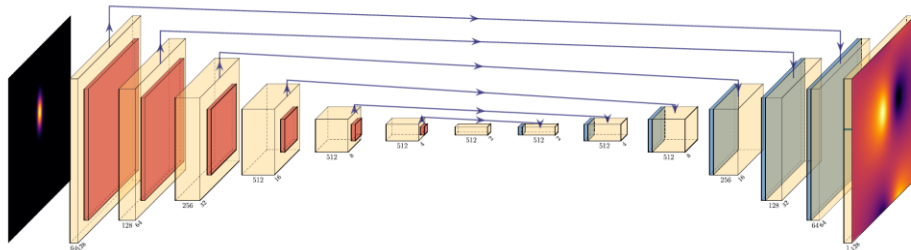
space-charge H_{sc}



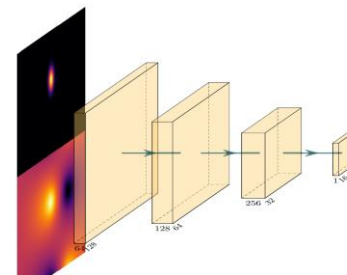
- Space-charge model computes the Hamiltonian associated with the space-charge potential by solving the Poisson equation
- A machine learning model can be used to speed up the calculation of H_{sc} w/o solving the Poisson equation directly

A neural network based ML model enables fast space-charge simulation

schematic diagram of the generator



schematic diagram of the discriminator

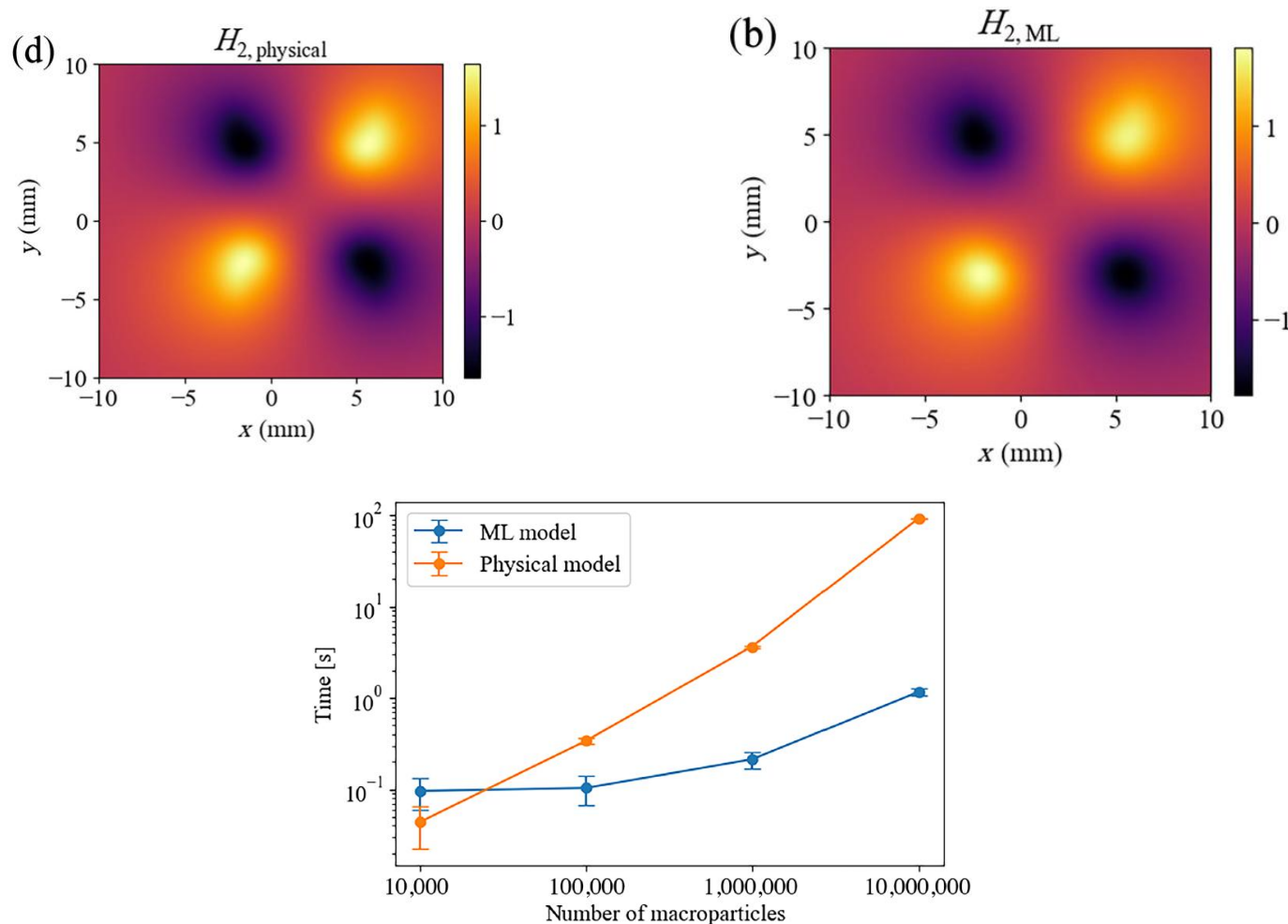


$$\mathcal{L}_G = -\frac{1}{N} \sum_{i=1}^N \{y_i \log[D(x_i, G(x_i))] + (1 - y_i) \log[1 - D(x_i, G(x_i))]\} + \lambda \frac{1}{N} \sum_{i=1}^N |G(x_i) - Y_i|$$

$$\mathcal{L}_D = -\frac{1}{N} \sum_{i=1}^N [y_i \log(D(x_i, y_i))] - \frac{1}{N} \sum_{i=1}^N \{(1 - y_i) \log[1 - D(x_i, G(x_i))]\}$$

- x_i is the input data of the i^{th} training sample, y_i represents its label (1 for real samples and 0 for fake samples), G represents the fake data generated by the generator based on the input x , D represents the discriminator's output, and Y represents the real target image.
- The first term in $\mathcal{L}_G$ is the adversarial loss that encourages the generator to produce images as realistic as possible to fool the discriminator, and the second term in $\mathcal{L}_G$ is the L1 loss measuring the pixelwise similarity between the target and generated images. Based on a systematic parameter search, the coefficient of the L1 loss λ is set to be 1. The first and second terms in $\mathcal{L}_D$ represent the loss function for evaluating the real and fake samples, which encourage the discriminator to output 1 for real images and 0 for fake images.

ML Based Surrogate Model Reproduces the Physical Model Well but Is More Computational Efficient



Differentiable Space-Charge Modeling

- The differentiable simulation is a simulation that can automatically compute derivatives of the simulation result with respect to its input parameters.
- Differentiable space-charge modeling is to include space-charge effects in the differentiable simulation.
- Differentiable simulation can be used to study:
 - sensitivity of target physical quantities w.r.t. design or initial parameters
 - included in fast gradient-based optimizer

J. Qiang, “Differentiable self-consistent space-charge simulation for accelerator design,” Phys. Rev. Accel. Beams 26, 024601 (2023).



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Forward Mode Differentiable Simulation Using DA/TPSA

- Conventional graphic representation of a complex function/code might not be easy.
- DA/TPSA has been used to calculate high-order transfer maps in accelerator beam dynamics.
- The same library can be used to calculate derivatives w.r.t. design parameters.
- DA/TPSA changes the calculation derivatives of a function w.r.p. variables into the calculation of function evaluation of DA vector like variables.

Forward AD Based on First Derivatives of DA/TPSA

<https://github.com/qianglbi/mAD>

Given a function $f(x_1, x_2, \dots, x_n)$

find its derivatives w.r.p. variable $x_1, x_2, \dots, x_n$

Define a DA/TPSA vector like variable $F = (f, f_{x_1}, f_{x_2}, \dots, f_{x_n})$

Addition Rule: $F + G = (f + g, f_{x_1} + g_{x_1}, f_{x_2} + g_{x_2}, \dots, f_{x_n} + g_{x_n})$

Multiplication Rule: $FG = (fg, gf_{x_1} + fg_{x_1}, gf_{x_2} + fg_{x_2}, \dots, gf_{x_n} + fg_{x_n})$

Division Rule: $\frac{F}{G} = (\frac{f}{g}, \frac{f_{x_1}g - g_{x_1}f}{g^2}, \frac{f_{x_2}g - g_{x_2}f}{g^2} \dots, \frac{f_{x_n}g - g_{x_n}f}{g^2})$

Function Rule: $h(F) = (h(f), h_f f_{x_1}, h_f f_{x_2}, \dots, h_f f_{x_n})$

Emittance and Its Derivatives w.r.p. Seven Lattice Param. from the Use of a TPSA Library and the mAD Module

Emittance and its 7 derivatives from the use of a TPSA library

Emx 1.061687927800431E-006 9.404478867221511E-008 -1.812609047369974E-009 -4.761253438860015E-007
-1.763704718594396E-008 2.345272532436834E-011 -1.977599625054010E-008 -1.196832248156261E-008

Emy 1.009205402120444E-006 6.725695553868752E-008 1.304519894794864E-009 4.491663448734034E-007
7.826349966999831E-008 2.017384176315016E-010 -4.701173761686457E-009 3.276319957786300E-008

computing time: 858.41 seconds

Emittance and its 7 derivatives from the use of the mAD module

Emx 1.061687927800431E-006 9.404478867221511E-008 -1.812609047369819E-009 -4.761253438860811E-007
-1.763704718594396E-008 2.345272532436833E-011 -1.977599625054010E-008 -1.196832248156509E-008

Emy 1.009205402120444E-006 6.725695553868883E-008 1.304519894794905E-009 4.491663448734086E-007
7.826349967000353E-008 2.017384176313790E-010 -4.701173761675994E-009 3.276319957786431E-008

computing time: 4.16 seconds

More than a factor of 200 reduction of computing time by using the fast mAD module!

Differentiable Space-Charge Simulation through a FODO Lattice



$$\mathcal{M}_1(\tau) = \begin{pmatrix} \cos(\sqrt{Dk}D\tau) & \frac{1}{\sqrt{Dk}} \sin(\sqrt{Dk}D\tau) \\ -\sqrt{Dk} \sin(\sqrt{Dk}D\tau) & \cos(\sqrt{Dk}D\tau) \end{pmatrix}$$

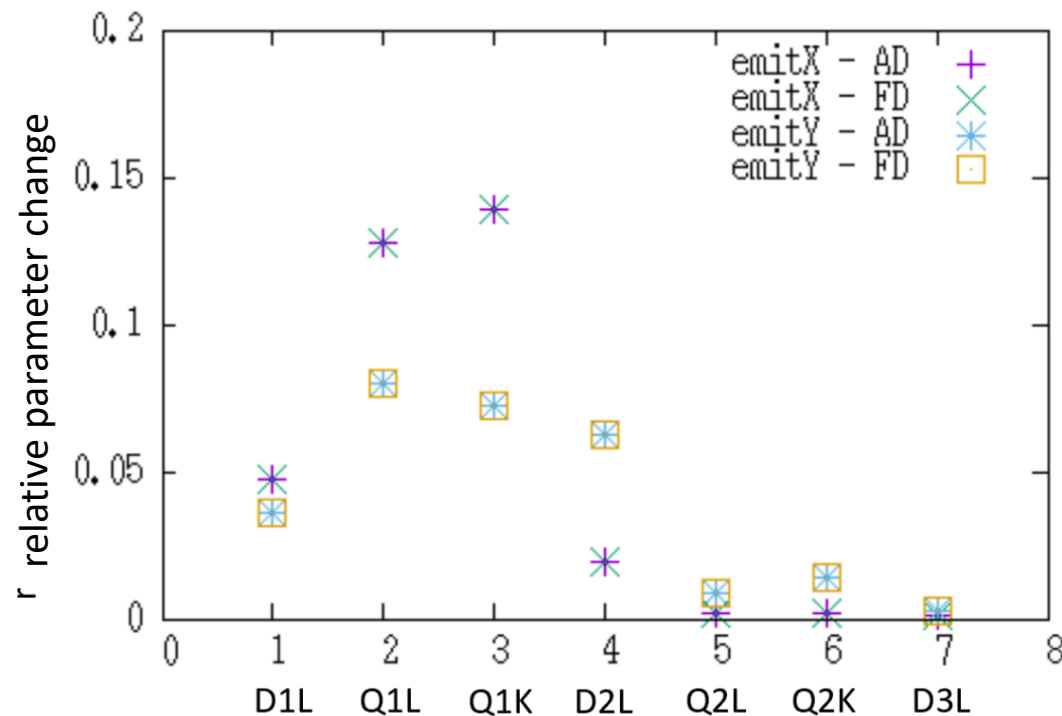
$$Dp_{xi}(\tau) = Dp_{xi}(0) - D\tau \frac{K}{2} \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} D\phi^{lm} \alpha_l \cos(\alpha_l Dx_i) \sin(\beta_m Dy_i)$$

$$Dp_{yi}(\tau) = Dp_{yi}(0) - D\tau \frac{K}{2} \sum_{l=1}^{N_l} \sum_{m=1}^{N_m} D\phi^{lm} \beta_m \sin(\alpha_l Dx_i) \cos(\beta_m Dy_i)$$

$$D\phi^{lm} = 4\pi \frac{4}{ab} \frac{1}{N_p} \sum_{j=1}^{N_p} \frac{1}{\gamma_{lm}^2} \sin(\alpha_l Dx_j) \sin(\beta_m Dy_j)$$

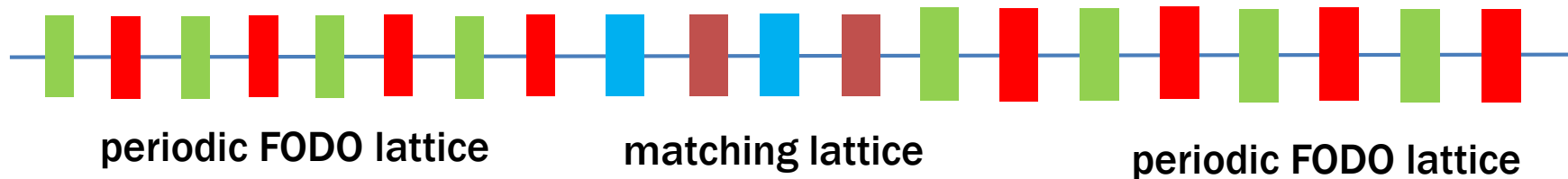
$$D\epsilon_x = \sqrt{D \langle x^2 \rangle D \langle p_x^2 \rangle - (D \langle xp_x \rangle)^2}$$

Relative emittance changes w.r.p lattice parameters from a single AD simulation and from the FD approximation using 14 simulations



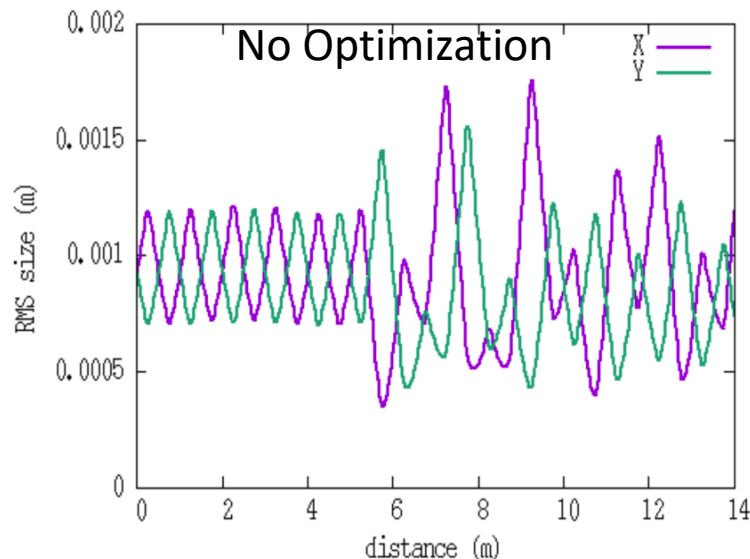
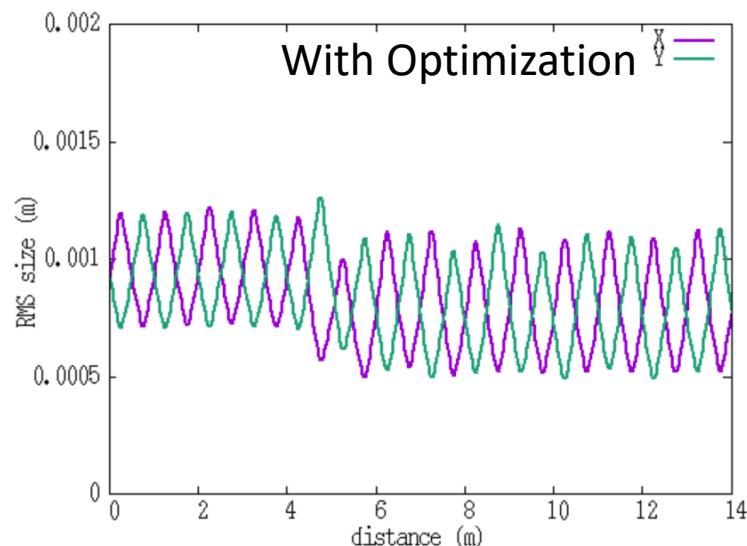
- Excellent agreement between the AD and the FD approximation.

Evolution of the transverse RMS beam size through the FODO lattice w/o optimization (left) and with gradient-based optimization using the auto-differentiable parallel PIC simulation



$$f(\mathbf{k}) = \frac{(\beta_x(\mathbf{k}) - \beta_{xt})^2}{\beta_{xt}^2} + (\alpha_x(\mathbf{k}) - \alpha_{xt})^2 + \frac{(\beta_y(\mathbf{k}) - \beta_{yt})^2}{\beta_{yt}^2} + (\alpha_y(\mathbf{k}) - \alpha_{yt})^2$$

- 4 control knobs (quadrupoles) in the matching lattice section



- Optimized matching solution yields more regular envelope oscillation including space-charge effects.

Study of the fully coupled 3D envelope instability using AD

3D Coupled Envelope Equation of 6x6 Covariance Matrix

$$\frac{d\Sigma}{ds} = F \Sigma + (F \Sigma)^T$$

$$F(s) = \begin{pmatrix} 0 & \frac{\delta}{lp_0} & 0 & 0 & 0 & 0 \\ F_{21}(s) & 0 & F_{23}(s) & 0 & F_{25}(s) & 0 \\ 0 & 0 & 0 & \frac{\delta}{lp_0} & 0 & 0 \\ F_{41}(s) & 0 & F_{43}(s) & 0 & F_{45}(s) & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{\delta}{lp_0} \\ F_{61}(s) & 0 & F_{63}(s) & 0 & F_{65}(s) & 0 \end{pmatrix}$$

$$\Sigma(s) = \Sigma_m(s) + \tilde{\Sigma}(s)$$

$$\xi = \left(\tilde{\Sigma}_{11}, \tilde{\Sigma}_{12}, \dots, \tilde{\Sigma}_{66} \right)^T$$

$$\frac{d\xi}{ds} = D(\Sigma_m, s) \xi(s)$$

$$\xi(s) = M(s) \xi(0)$$

$$\frac{dM(s)}{ds} = D(\Sigma_m, s) M(s) \longrightarrow 441 \text{ coupled ODEs}$$

- Reduction of the number of ODEs from 441 to 21 using AD!



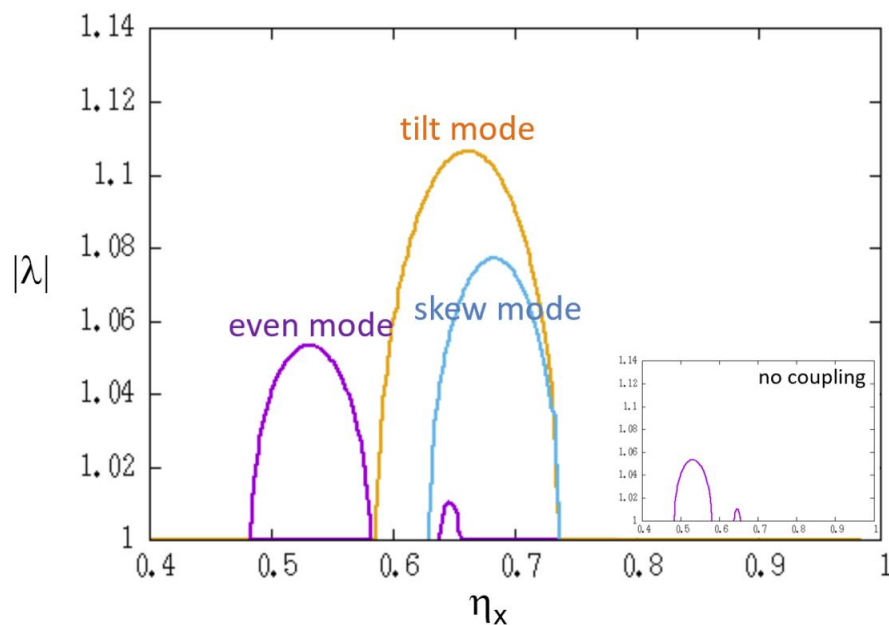
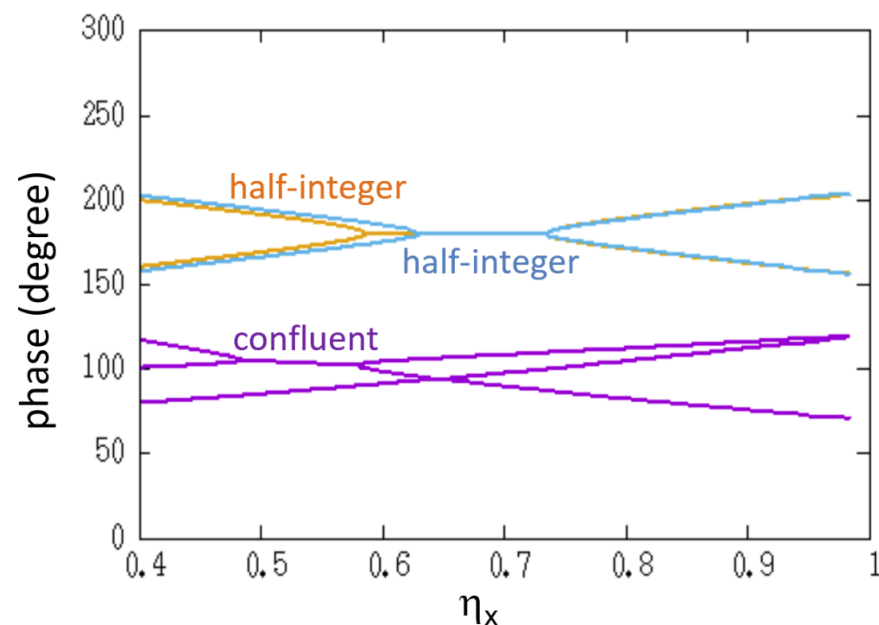
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Growth rate amplitudes and phases of unstable envelope modes as a function of the horizontal tune depression



- Including SC induced coupling in the envelope equation shows new instability stopbands

Introducing this session's talks

- Performance-Portable Space-Charge Methods in OPALX for Hadron and Lepton Accelerator Modelling – Michail Zampetakis
- Fast-Tracking Space Charge tool for Nonlinear Dynamics study in Realistic Collider Lattices – Helena Alamprese
- Longitudinal Phase-Space Reconstruction Using Differentiable Space-Charge Solvers – Sergei Sherstiuk
- Space Charge Physics in the Code ImpactX – Chad Mitchell

Thank You for Your Attention!