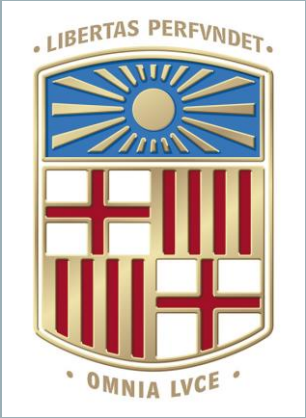




A UNITARY KAON- DEUTERON INTERACTION FOR FEMTOSCOPY



Juan M. Torres-Rincon

Universitat de Barcelona & Institut de Ciències del Cosmos



Sevilla, September 9, 2026,
*15th International Workshop on
the Physics of Excited Nucleons*
NSTAR 2026



Kaon-Deuteron interaction

Meson-baryon chiral interactions

E. Oset, A. Ramos, *Nucl.Phys.A* 635 (1998) 99-120

J.A. Oller, U.G. Meißner, *Phys.Lett.B* 500 (2001) 263-272

T. Hyodo, D. Jido, *Prog. Part. Nucl. Phys.* 67(1):55-98 (2012)

$$\mathcal{L}_{\Phi B} = \text{Tr} \left[\bar{B} (i \not{D} - M_0) B + \frac{D}{2} (\bar{B} \gamma^\mu \gamma_5 \{u_\mu, B\}) + \frac{F}{2} (\bar{B} \gamma^\mu \gamma_5 [u_\mu, B]) \right]$$



2-body unitarized amplitudes

$$t = [1 - V G_2]^{-1} V$$

OR model (E. Oset, A. Ramos, *Nucl.Phys.A* 635 (1998) 99; D. Jido, E. Oset, A. Ramos, *Phys.Rev.C* 66, 055203 (2002))

Lowest order + KN scattering data

BCN model (A. Feijoo, V. Magas, A. Ramos, *Phys.Rev.C* 99, 035211 (2019))

NLO + KN scattering data + kaonic hydrogen $1s$ state

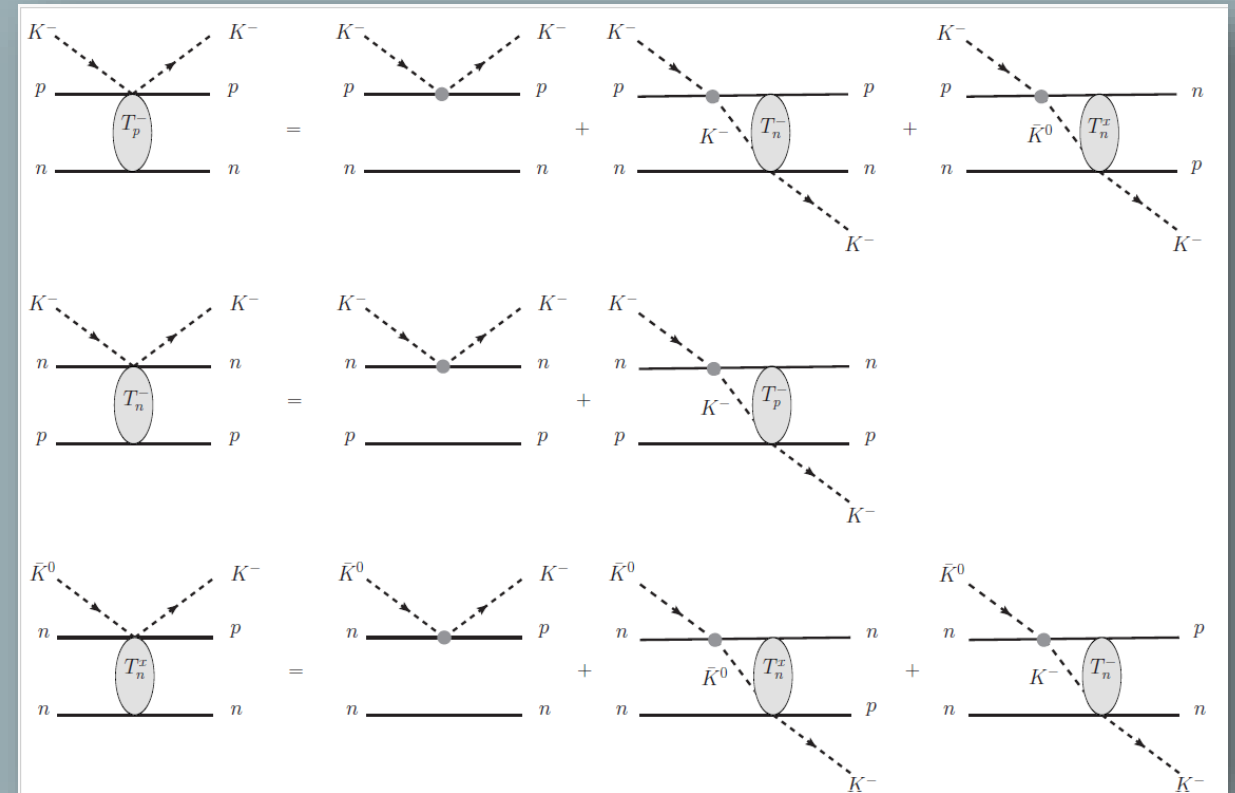
Kaon-Deuteron interaction

Kaon-deuteron interaction based on
2-body meson-baryon chiral interactions

**Faddeev equations in the
Fixed Center Approximation (FCA)**

$$T_{K^-d}^{\text{FCA}} = T_p + T_n$$

$$\begin{aligned} T_p &= t_p + t_p G T_n - t_p^x G T_n^x \\ T_n &= t_n + t_n G T_p \\ T_n^x &= t_n^x - t_n^0 G T_n^x + t_n^x G T_n \end{aligned}$$



S.S. Kamalov, E. Oset, A. Ramos, *Nucl.Phys.A* 690 (2001) 494

U. Meißner, U. Raha, A. Rusetsky, *Eur.Phys.J.C* (2006) 473

E. Oset et al., *Nucl.Phys.A* 881 (2012) 127 (beyond threshold amplitudes)

A. Martinez-Torres et al. *Phys.Rev.D* 78 (2008) 074031 & A. Martinez-Torres, Ph.D. thesis, IFIC (2009)

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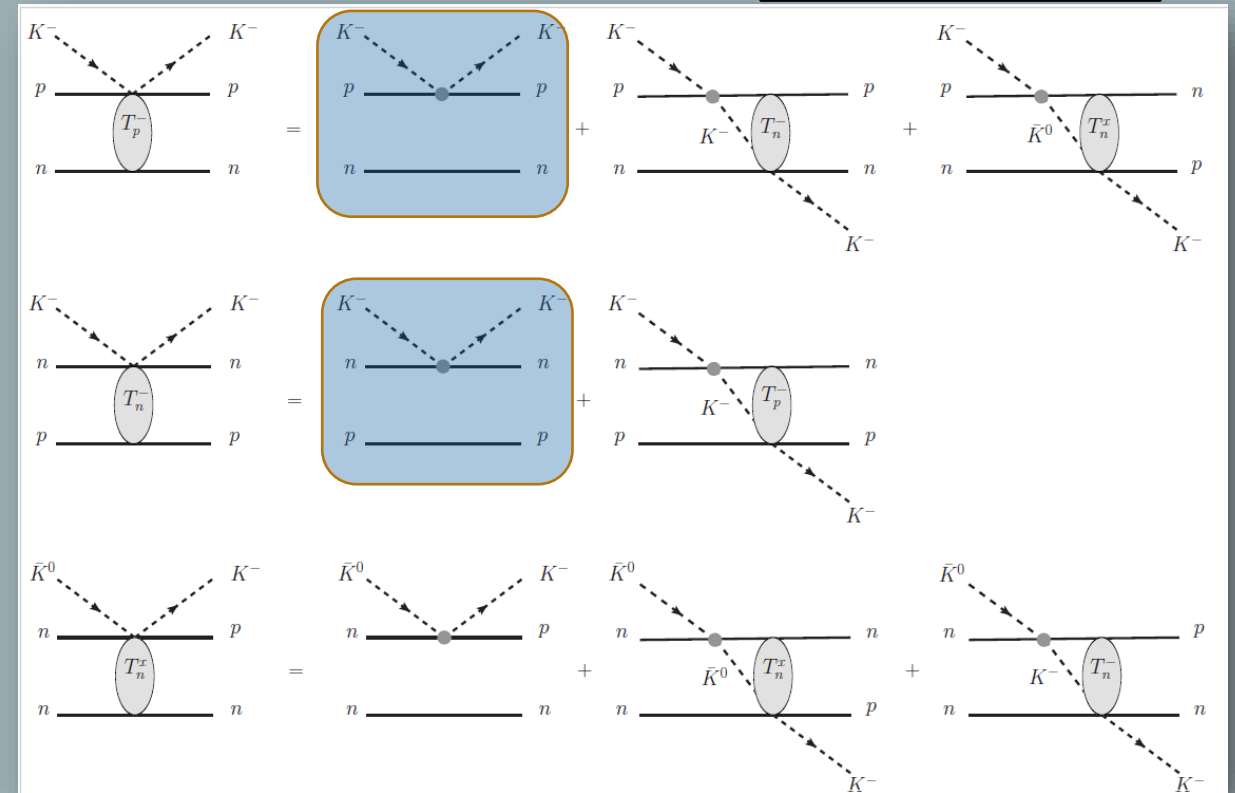
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Impulse Approximation

$$T_{K^-d}^{\text{IA}} = t_p + t_n$$



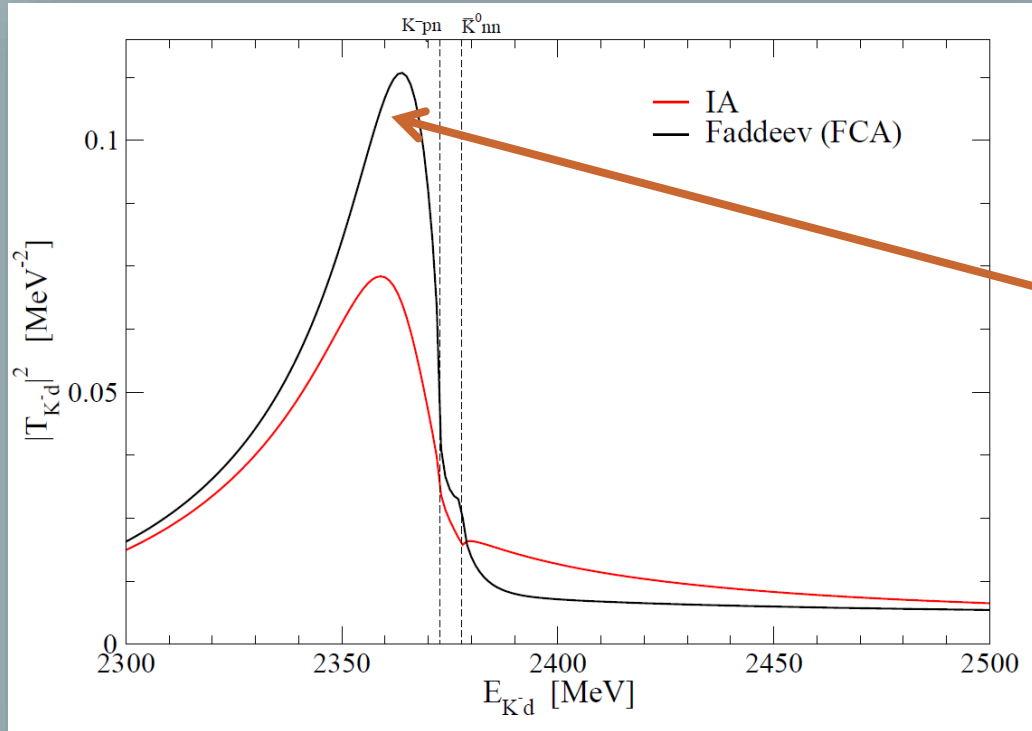
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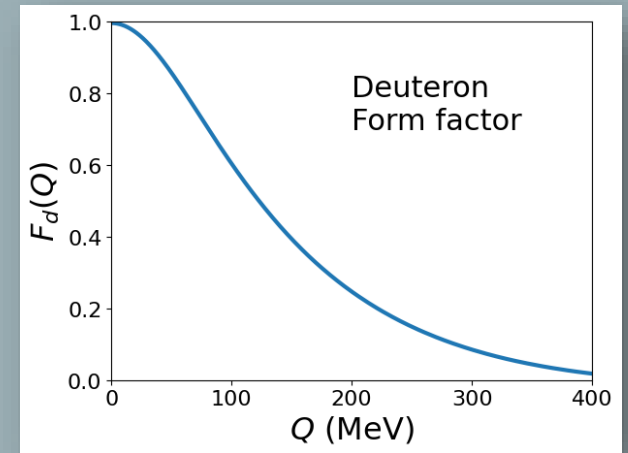
T-matrix of K^- d scattering

À. Ramos, JMT-R, A. de Fagoaga, E. Cabré,
Phys.Rev.D 113, 036020 (2026)

3-body bound state originated from the $\Lambda(1405)$

See also

E. Oset et al., *Nucl.Phys.A*
881 (2012) 127



$$T_{K^-d}^{\text{FCA}} = \frac{t_p + t_n + (2t_p t_n - \tilde{t}_x^2) G - 2t_n \tilde{t}_x^2 G^2}{1 - t_p t_n G^2 + t_n \tilde{t}_x^2 G^3}$$

$$\tilde{t}_x^2 = t_x^2 / (1 + t_0 G)$$

$$G(q^0) = \int \frac{d^3 q}{(2\pi)^3} \frac{F_d(\mathbf{q})}{(q^0)^2 - \mathbf{q}^2 - m_K^2 + i\eta}$$

Kaon-Deuteron interaction

Scattering Length

$$A_{Kd} = -\frac{1}{4\pi} \frac{M_d}{m_K + M_d} T_{Kd}$$



From K-d femtoscopy in heavy-ion collisions
ALICE Coll (2602.22833 [nucl-ex])

À. Ramos, JMT-R, A. de Fagoaga, E. Cabré, *Phys.Rev.D* 113, 036020 (2026)

(OR model)

	A_{K-d} [fm]
Impulse Approx.	$-0.59 + i2.15$
Fixed-Center Approx.	$-2.06 + i1.77$

$$\text{Re } A_{K-d} = -1.44 \pm 0.15 \text{ (stat)}^{+0.10}_{-0.10} \text{ (syst) fm}$$

$$\text{Im } A_{K-d} = +1.34 \pm 0.33 \text{ (stat)}^{+0.21}_{-0.15} \text{ (syst) fm}$$

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Energy shift and width of the $1s$ atomic level in kaonic deuterium

$$\varepsilon_{1s}^d - i\frac{\Gamma_{1s}^d}{2} = -\frac{2\alpha^3 \mu_r^2 A_{K-d}}{1 + 2\alpha \mu_r A_{K-d} (\ln \alpha - 1)}$$

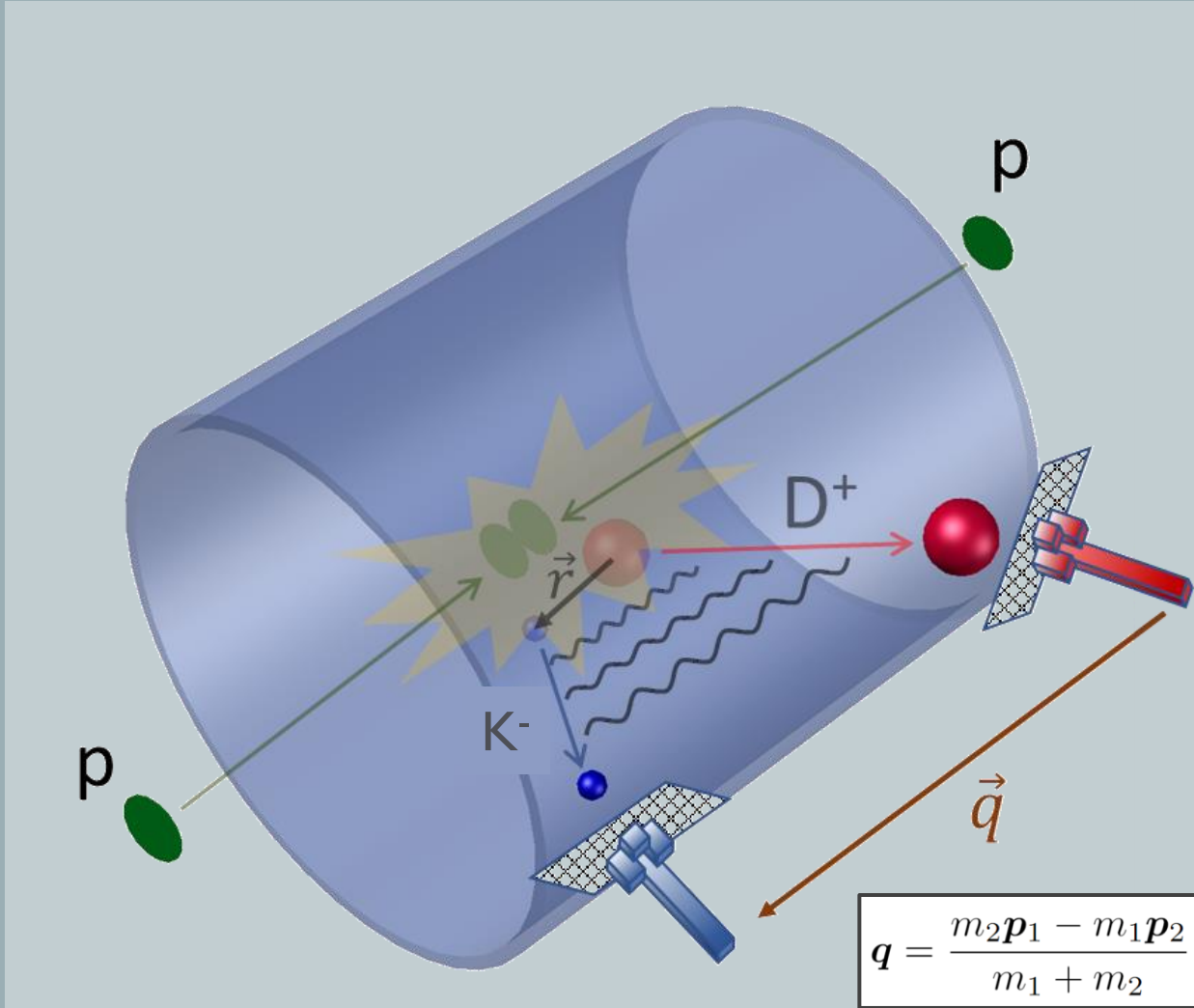
	ε_{1s}^d [eV]	Γ_{1s}^d [eV]
Impulse Approx.	791	2063
Fixed-Center Approx.	1040	1106



Experimental value from SIDDHARTA-2 experiment
(2608.10749 [nucl-ex])

$$\begin{aligned} \varepsilon_{1s}^d &= 810.9 \pm 24.5 \text{ (stat)} \pm 2.1 \text{ (syst) eV} \\ \Gamma_{1s}^d &= 812 \pm 97 \text{ (stat)} \pm 33 \text{ (syst) eV} \end{aligned}$$

Femtoscscopy in RHICs



U. Heinz, B. Jacak, *Ann.Rev.Nucl.Part.Sci.* 49 (1999) 529

M. Lisa, S. Pratt, U. Wiedemann, *Ann.Rev.Nucl.Part.Sci.* 55 (2005) 357

Pair Correlation Function

$$C(\mathbf{q}) = \mathcal{N} \frac{N_{\text{same}}(\mathbf{q})}{N_{\text{mixed}}(\mathbf{q})}$$

L. Fabbietti, V. Mantovani Sarti, O. Vazquez Doce,
Ann. Rev. Nucl. Part. Sci. 71, 377 (2021)

$C(\mathbf{q}) > 1$: attraction

$C(\mathbf{q}) < 1$: repulsion

Koonin-Pratt formula

Single-channel KP formula

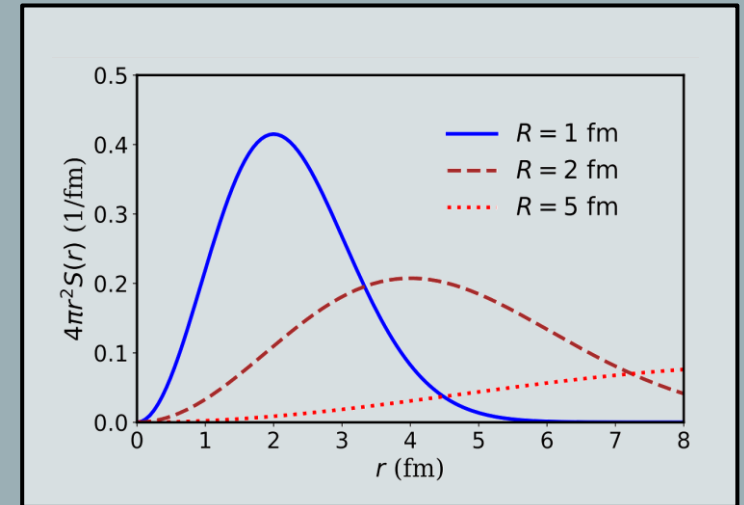
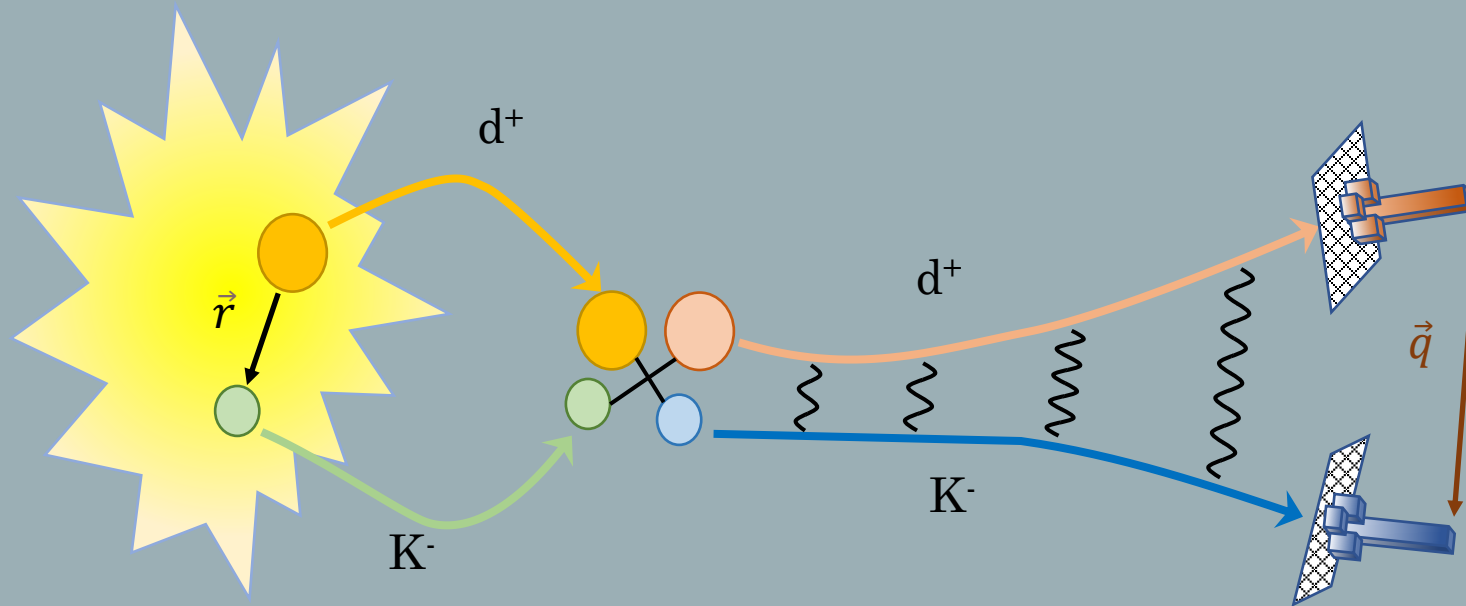
$$C(\mathbf{q}) = \int d^3r S(\mathbf{r}) |\Psi(\mathbf{q}; \mathbf{r})|^2$$

Koonin, *Phys.Lett.B*, 70, 43 (1977)

Pratt, Csorgo, Zimanyi, *Phys.Rev.C*, 42, 2646(1990)

Pair wave function with CM momentum $\mathbf{q}$

Source function



Gaussian source function

$$S(r) = \frac{1}{(2\sqrt{\pi}R)^3} \exp\left(-\frac{r^2}{4R^2}\right)$$

K^- d correlation function

À. Ramos, JMT-R, A. de Fagoaga, E. Cabré, *Phys.Rev.D* 113, 036020 (2026)

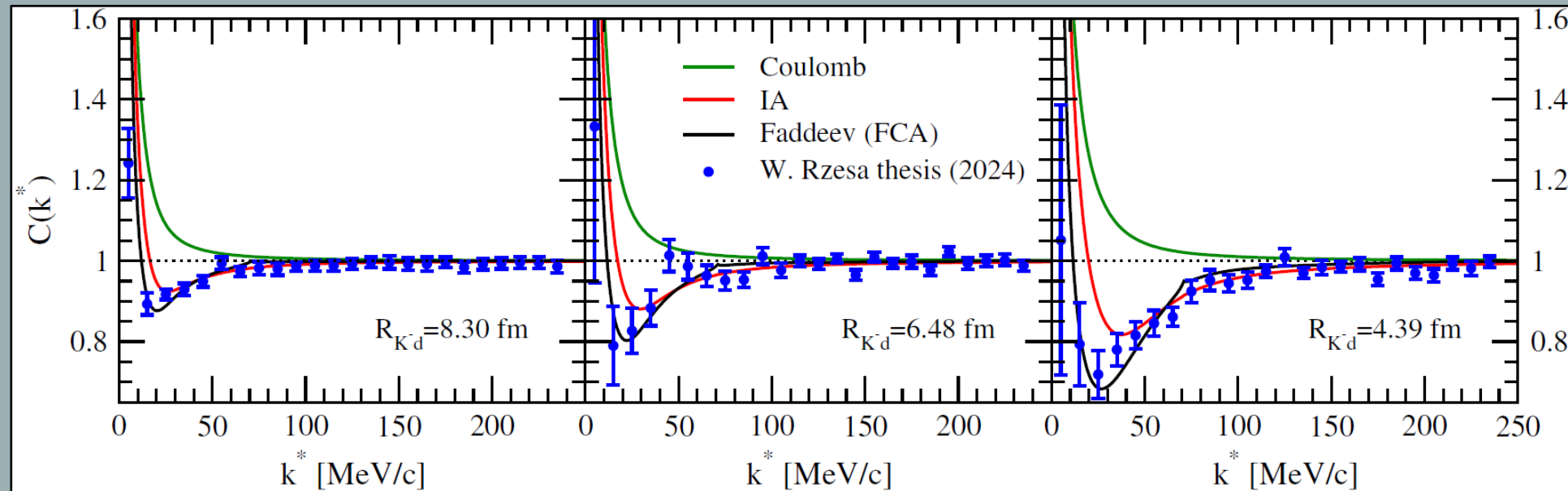
Pb-Pb collisions @ $s^{1/2}_{NN}=5.02$ TeV

W. Rzes, Ph.D. thesis,
Warsaw U. of Tech. (2024) (ALICE Coll.)

Centrality: 0-10%

10-30%

30-50%



Source radii only **indicative** from experimental fits

FCA works very well! No new parameters, only model assumptions

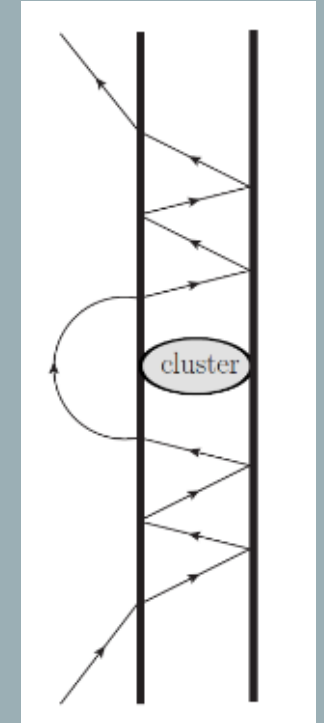
Kaon-Deuteron interaction

Lack of unitarity in the particle-cluster scattering concerns in FCA raised by

- P. Encarnación, A. Feijoo, and E. Oset, *Phys. Rev. D* 111, 114023 (2025),
- N. Ikeno and E. Oset, *Phys. Rev. D* 111 (2025) 114023

Solved by coherent propagation in

- B. Agatão, P. Brandão, A. Martínez Torres, K. P. Khemchandani, L. M. Abreu, and E. Oset, *Eur. Phys. J. C* 85, 1136 (2025)



Unitarity not realized at the particle-cluster threshold

Small effect for the deuteron case (binding energy of 2.2 MeV).

Tiny shift in nucleon masses adjust unitary branch point

*) À. Ramos, JMT-R, A. de Fagoaga, E. Cabré, *Phys. Rev. D* 113, 036020 (2026)

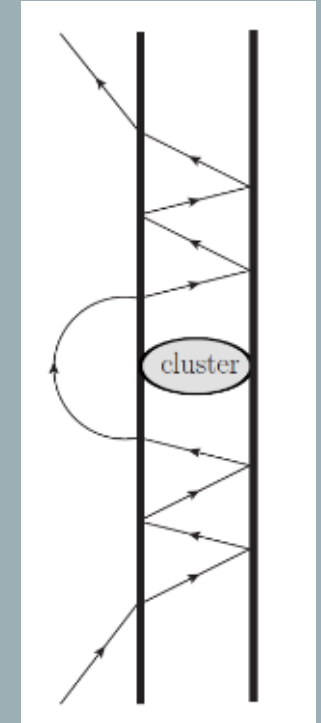
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Effective range expansion

$$-\frac{4\pi\sqrt{s}}{M_d}(T_{Kd})^{-1} = \frac{1}{A_{Kd}} + \frac{1}{2}r_0q_{\text{CM}}^2 - iq_{\text{CM}}$$

Unitary Factor should be 1!

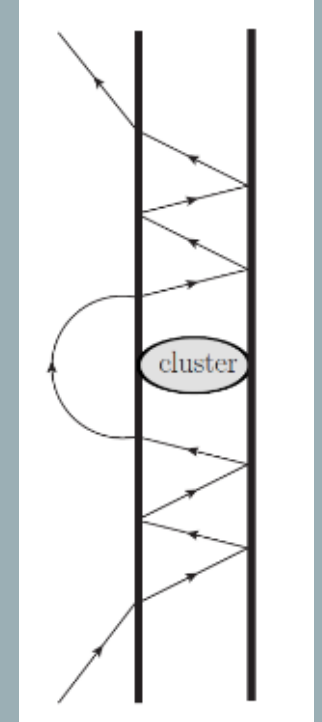
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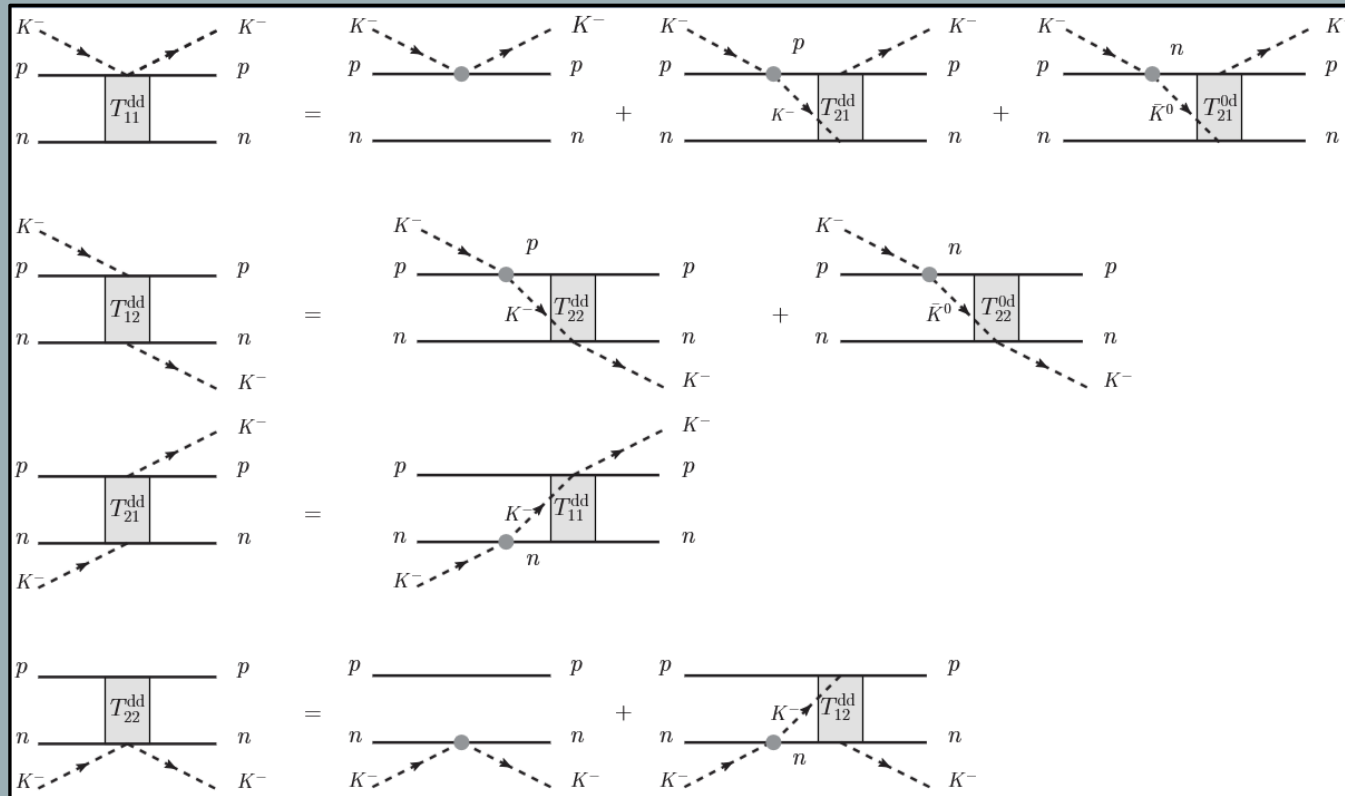
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Unitary Factor should be 1!

	FCA (OR model*)	FCA (BCN model)
Unitary Factor	0.69	0.62

Kaon-Deuteron interaction

A. Feijoo, À. Ramos, JMT-R, in preparation



Rewriting of Faddeev partitions

Keep track of

1) first and last kaon interaction with particle 1 or 2

2) Channel type

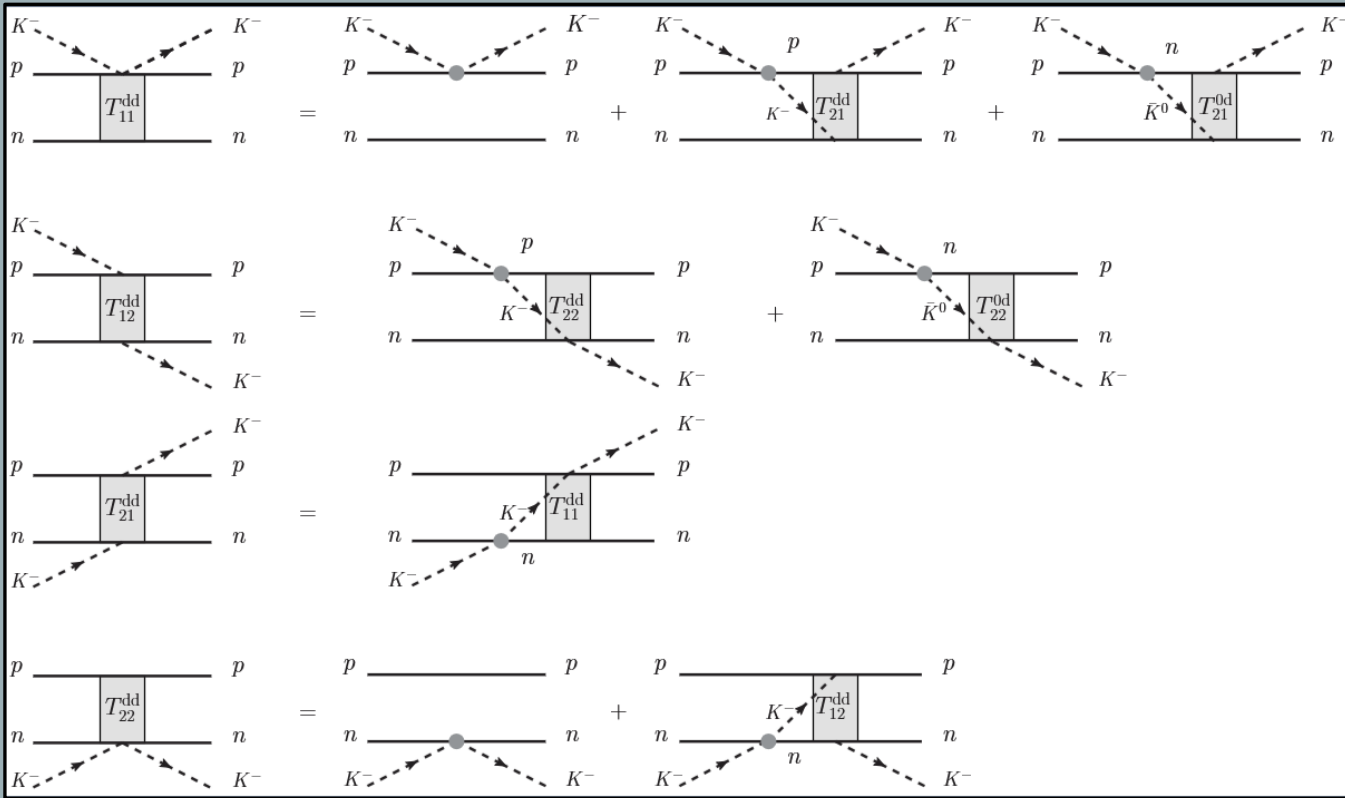
2.1) direct ($K^- p n$)

2.2) exchange ($K^- n p$)

2.3) neutral ($\bar{K}^0 n n$)

Kaon-Deuteron interaction

A. Feijoo, À. Ramos, JMT-R, in preparation



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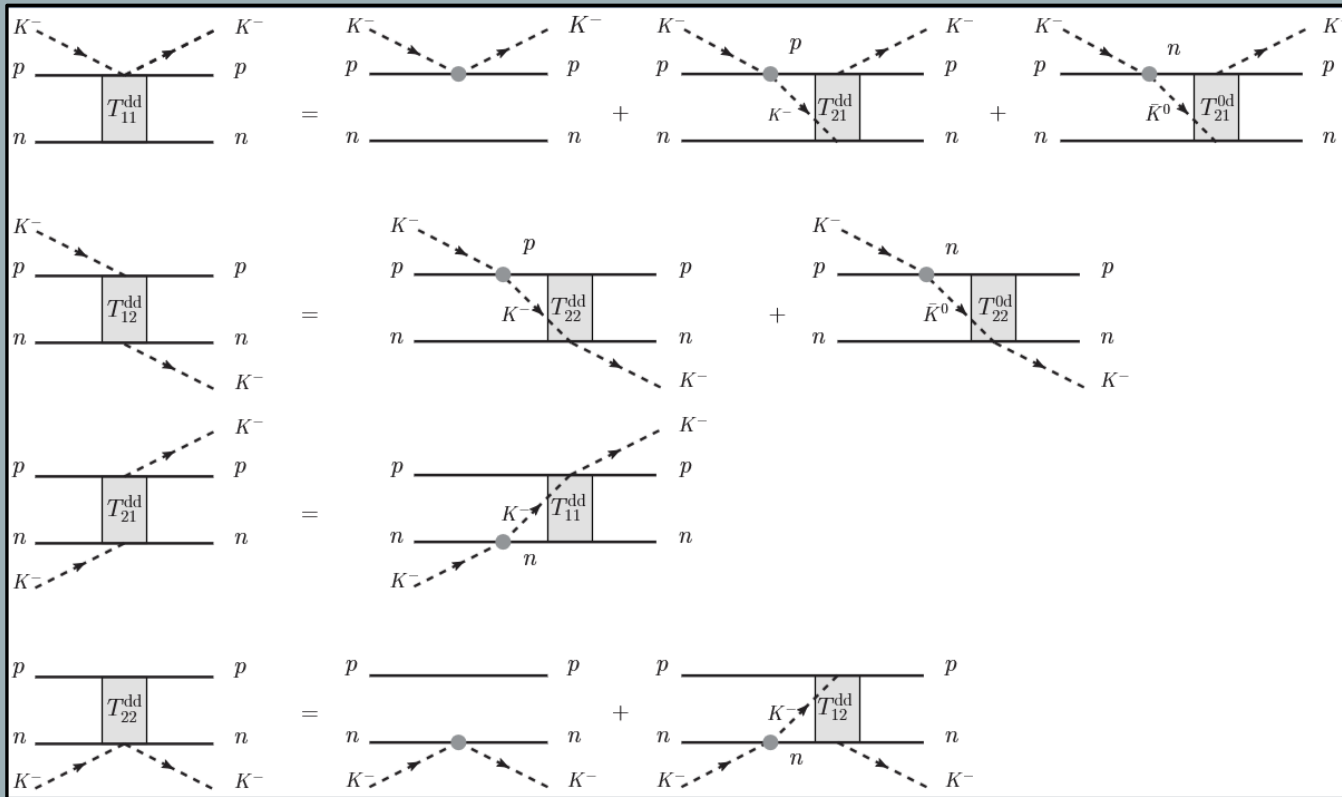
2.2) exchange (K^- n p)

2.3) neutral ($\bar{K}^0$ n n)

$$\mathbf{T} = \begin{pmatrix} T_{11}^{\text{dd}} & T_{12}^{\text{dd}} & T_{11}^{\text{de}} & T_{12}^{\text{de}} & T_{11}^{\text{d0}} & T_{12}^{\text{d0}} \\ T_{21}^{\text{dd}} & T_{22}^{\text{dd}} & T_{21}^{\text{de}} & T_{22}^{\text{de}} & T_{21}^{\text{d0}} & T_{22}^{\text{d0}} \\ T_{11}^{\text{red}} & T_{12}^{\text{red}} & T_{11}^{\text{ee}} & T_{12}^{\text{ee}} & T_{11}^{\text{e0}} & T_{12}^{\text{e0}} \\ T_{21}^{\text{red}} & T_{22}^{\text{red}} & T_{21}^{\text{ee}} & T_{22}^{\text{ee}} & T_{21}^{\text{e0}} & T_{22}^{\text{e0}} \\ T_{11}^{\text{0d}} & T_{12}^{\text{0d}} & T_{11}^{\text{0e}} & T_{12}^{\text{0e}} & T_{11}^{\text{00}} & T_{12}^{\text{00}} \\ T_{21}^{\text{0d}} & T_{22}^{\text{0d}} & T_{21}^{\text{0e}} & T_{22}^{\text{0e}} & T_{21}^{\text{00}} & T_{22}^{\text{00}} \end{pmatrix}$$

Kaon-Deuteron interaction

A. Feijoo, À. Ramos, JMT-R, in preparation



$$T_{K^-d}^{\text{FCA}} = T^p + T^n = \sum_{i,j=1}^2 (T_{ij}^{\text{dd}} - T_{ij}^{\text{de}})$$

This coupled system gives the previous FCA

Rewriting of Faddeev partitions

Keep track of

1) first and last kaon interaction with particle 1 or 2

2) Channel type

2.1) direct ($K^- p n$)

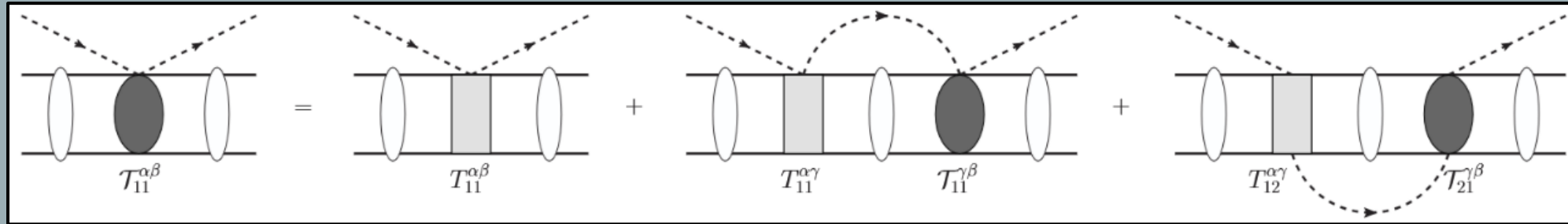
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Kaon-Deuteron interaction

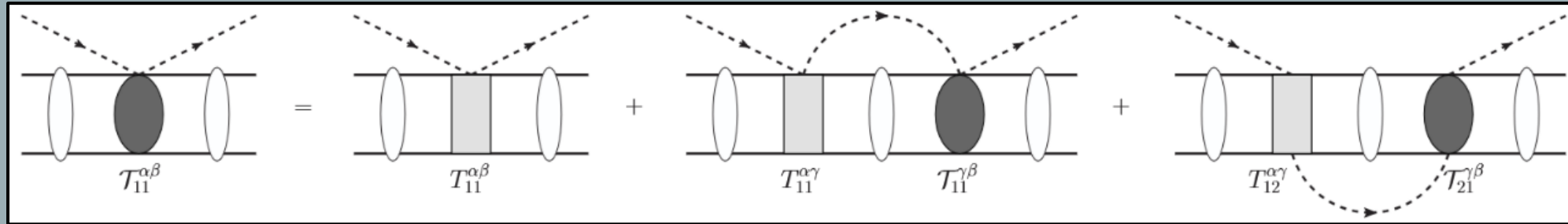
A. Feijoo, À. Ramos, JMT-R, in preparation



$$\mathbf{T}^{-1} = \begin{pmatrix} \frac{t_0}{t_0 t_1 - t_x^2} & -G_- & 0 & 0 & -\frac{t_x}{t_0 t_1 - t_x^2} & 0 \\ -G_- & t_2^{-1} & 0 & 0 & 0 & 0 \\ 0 & 0 & t_2^{-1} & -G_- & 0 & 0 \\ 0 & 0 & -G_- & \frac{t_0}{t_0 t_1 - t_x^2} & 0 & -\frac{t_x}{t_0 t_1 - t_x^2} \\ -\frac{t_x}{t_0 t_1 - t_x^2} & 0 & 0 & 0 & \frac{t_1}{t_0 t_1 - t_x^2} & -G_0 \\ 0 & 0 & 0 & -\frac{t_x}{t_0 t_1 - t_x^2} & -G_0 & \frac{t_1}{t_0 t_1 - t_x^2} \end{pmatrix}$$

Kaon-Deuteron interaction

A. Feijoo, À. Ramos, JMT-R, in preparation



Unitary Amplitude

$$\mathcal{T}^{-1} = \mathbf{T}^{-1} - \mathbf{G}_c$$

$$\mathbf{G}_c = \text{diag} (G_-^{(1)}, G_-^{(2)}, G_-^{(1)}, G_-^{(2)}, G_0^{(1)}, G_0^{(2)})$$

$$\mathbf{T}^{-1} = \begin{pmatrix} \frac{t_0}{t_0 t_1 - t_x^2} & -G_- & 0 & 0 & -\frac{t_x}{t_0 t_1 - t_x^2} & 0 \\ -G_- & t_2^{-1} & 0 & 0 & 0 & 0 \\ 0 & 0 & t_2^{-1} & -G_- & 0 & 0 \\ 0 & 0 & -G_- & \frac{t_0}{t_0 t_1 - t_x^2} & 0 & -\frac{t_x}{t_0 t_1 - t_x^2} \\ -\frac{t_x}{t_0 t_1 - t_x^2} & 0 & 0 & 0 & \frac{t_1}{t_0 t_1 - t_x^2} & -G_0 \\ 0 & 0 & 0 & -\frac{t_x}{t_0 t_1 - t_x^2} & -G_0 & \frac{t_1}{t_0 t_1 - t_x^2} \end{pmatrix}$$

$$G_m^{(i)} = \int \frac{d^3 q}{(2\pi)^3} \frac{\left(F_d^{(i)}(\mathbf{q})\right)^2}{2\omega_m(\mathbf{q}) (\sqrt{s_{Kd}} - \omega_m(\mathbf{q}) - E_m(\mathbf{q}) + i\eta)}$$

$$T_{K^-d}^{\text{uni}} = \sum_{i,j=1}^2 (\mathcal{T}_{ij}^{\text{dd}} - \mathcal{T}_{ij}^{\text{de}})$$

Final result too large to be shown

Kaon-Deuteron interaction

Particular cases

Set all kaon-deuteron propagators equal

$$G_{-}^{(1)} = G_{-}^{(2)} = G_{-} \quad ; \quad G_0^{(1)} = G_0^{(2)} = G_0$$

Unitarized Impulse Approximation!

$$\left(T_{\text{uni}}^{K^-d}\right)^{-1} = \frac{1}{t_p + t_n} - G_{-}$$

Neglect charge-exchange process

$$t_x = 0$$

We recover previous result in
B. Agatão, P. Brandão, A. Martínez Torres,
K. P. Khemchandani, L. M. Abreu, and
E. Oset, *Eur. Phys. J. C* 85, 1136 (2025)

$$\left(T_{\text{uni}}^{K^-d}\right)^{-1} = \frac{t_p + t_n + (2G - G_{-}^{(1)} - G_{-}^{(2)})t_p t_n}{1 - G_{-}^{(1)}t_p - G_{-}^{(2)}t_n - \left(G^2 - G_{-}^{(1)}G_{-}^{(2)}\right)t_p t_n}$$

Kaon-Deuteron interaction

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	FCA (OR model*)	FCA-UNIT (OR model)	FCA-UNIT (BCN model)
Unitary Factor	0.69	1.00	1.00
Scattering Length [fm]	$-2.06 + i1.77$	$-1.04 + i0.90$	$-0.85 + i1.07$



From K-d femtoscopy in heavy-ion collisions
ALICE Coll (arXiv:2602.22833 [nucl-ex])

$$\begin{aligned}\text{Re } A_{K-d} &= -1.44 \pm 0.15 \text{ (stat)} \text{ }^{+0.10}_{-0.10} \text{ (syst) fm} \\ \text{Im } A_{K-d} &= +1.34 \pm 0.33 \text{ (stat)} \text{ }^{+0.21}_{-0.15} \text{ (syst) fm}\end{aligned}$$

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	FCA (OR model*)	FCA-UNIT (OR model)	FCA-UNIT (BCN model)
ε_{1s}^d [eV]	1040	583	525
Γ_{1s}^d [eV]	1106	767	958

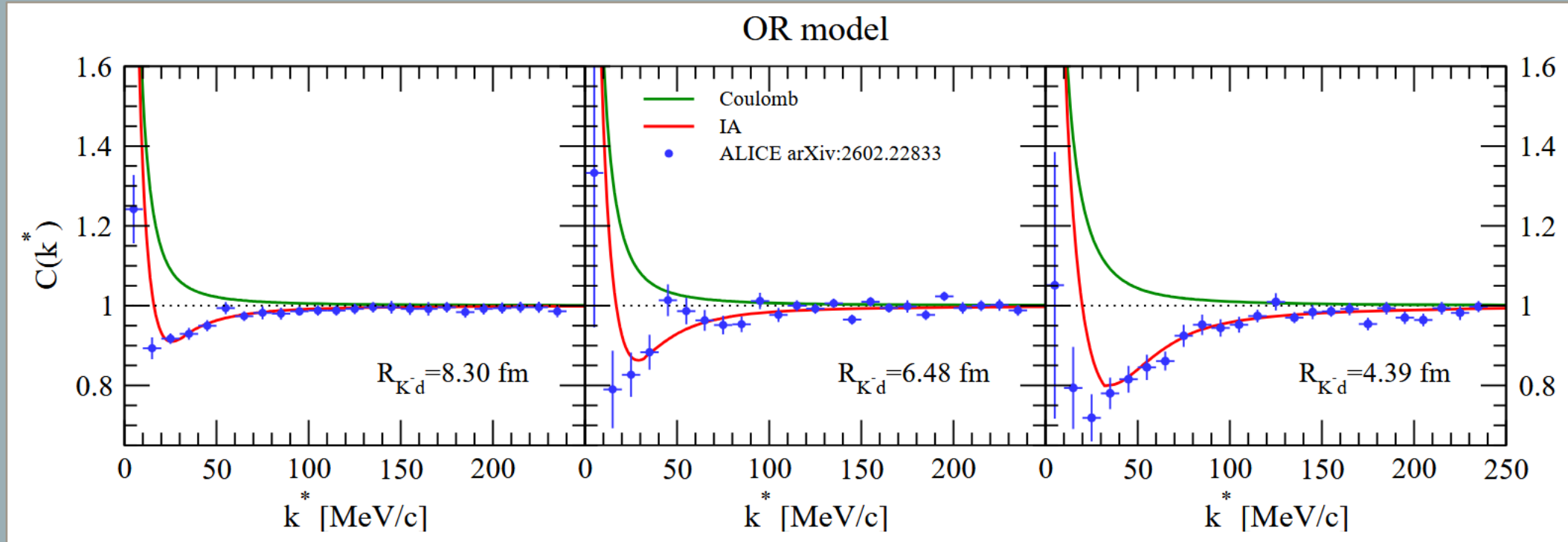


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$$\begin{aligned} \varepsilon_{1s}^d &= 810.9 \pm 24.5 \text{ (stat)} \pm 2.1 \text{ (syst) eV} \\ \Gamma_{1s}^d &= 812 \pm 97 \text{ (stat)} \pm 33 \text{ (syst) eV} \end{aligned}$$

Kaon-Deuteron interaction

Experimental from K^- d femtoscopy in heavy ions ALICE Coll (arXiv:2602.22833 [nucl-ex])

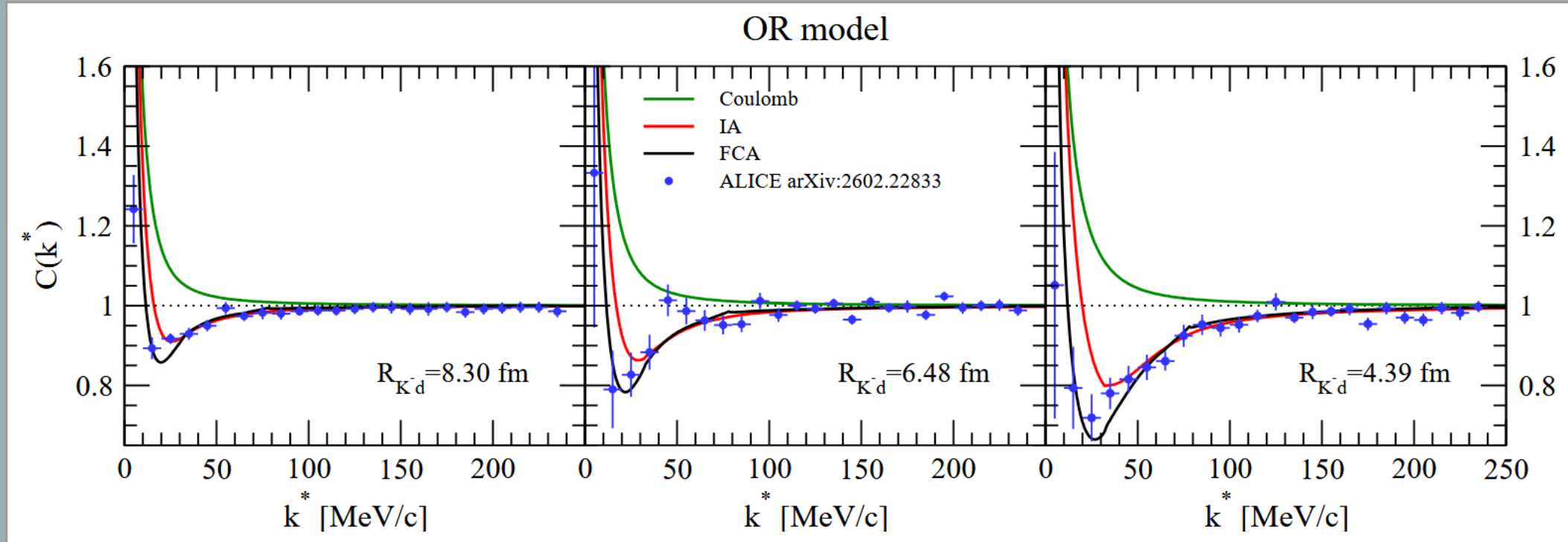


OR model

Source radii only **indicative** from experimental fits
(they are attached to a different interaction)

Kaon-Deuteron interaction

Experimental from K^- d femtoscopy in heavy ions ALICE Coll (arXiv:2602.22833 [nucl-ex])

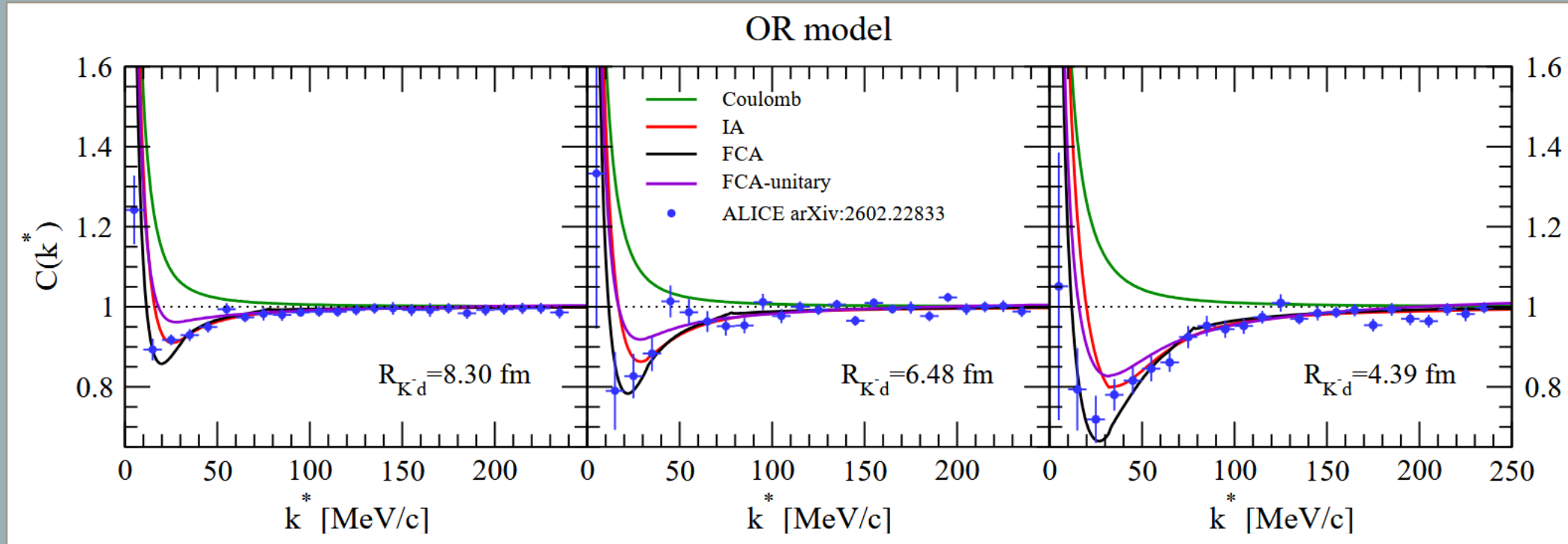


OR model

Source radii only **indicative** from experimental fits
(they are attached to a different interaction)

Kaon-Deuteron interaction

Experimental from K^- d femtoscopy in heavy ions ALICE Coll (arXiv:2602.22833 [nucl-ex])

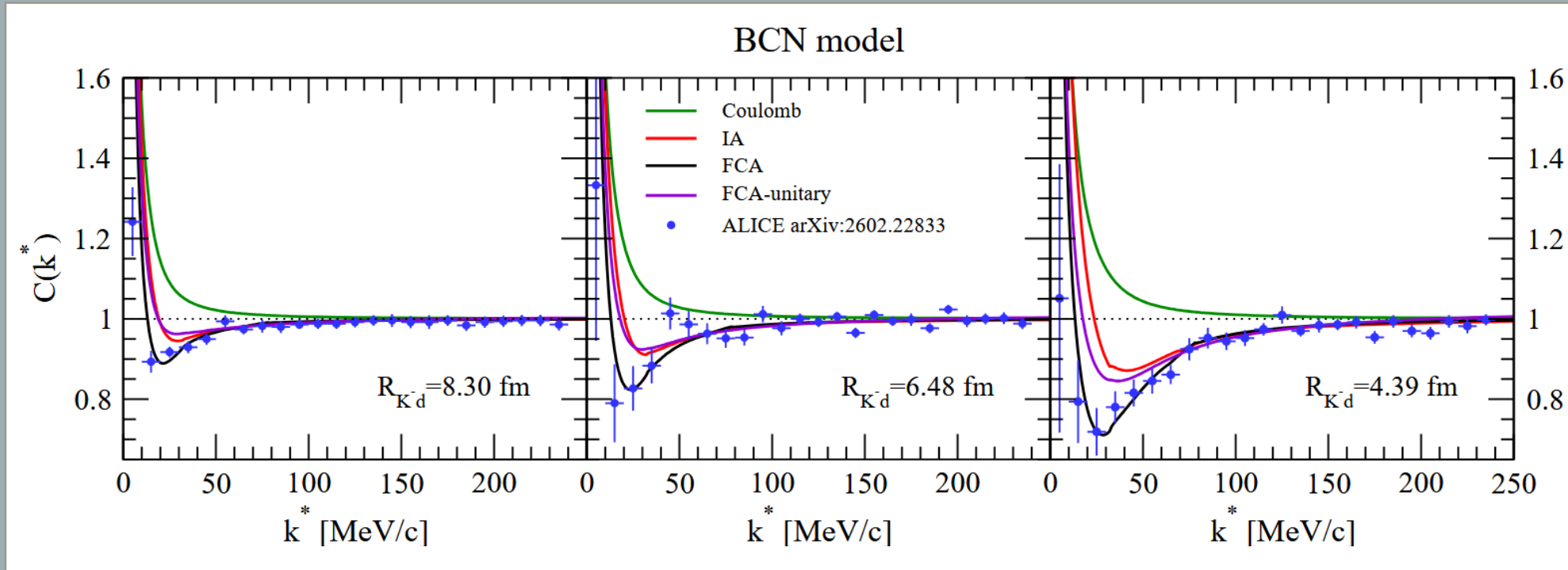


OR model

Source radii only **indicative** from experimental fits
(they are attached to a different interaction)

Kaon-Deuteron interaction

Experimental from K^- d femtoscopy in heavy ions ALICE Coll (arXiv:2602.22833 [nucl-ex])



BCN model

Source radii only **indicative** from experimental fits
(they are attached to a different interaction)

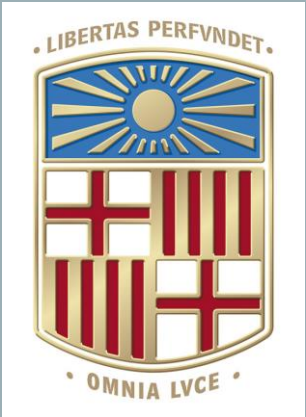
FCA with exact unitarity deviates from data

K^- d correlation function

Summary

Kaon-deuteron interaction in FCA provide a good description of ALICE femtoscopy data

- Unitarity has been restored! Scattering length and kaonic deuterium energy shift and width have improved from previous results.
- Femtoscopy correlation function with untarized FCA has deviated from data. Radii need to be modified/fitted. However, interaction model might need also to be assessed.
- OUTLOOK: Calibrate break-up, recombination, rearrangement processes in Faddeev equations.



A UNITARY KAON- DEUTERON INTERACTION FOR FEMTOSCOPY



Juan M. Torres-Rincon

Universitat de Barcelona & Institut de Ciències del Cosmos



Sevilla, September 9, 2026,
*15th International Workshop on
the Physics of Excited Nucleons*
NSTAR 2026



Meson-baryon chiral effective theory

$$\mathcal{L}_{\Phi B} = \text{Tr} \left[\bar{B} (i \not{D} - M_0) B + \frac{D}{2} (\bar{B} \gamma^\mu \gamma_5 \{u_\mu, B\}) + \frac{F}{2} (\bar{B} \gamma^\mu \gamma_5 [u_\mu, B]) \right]$$

$$D_\mu B = \partial_\mu B + [\Gamma_\mu, B] \quad \Gamma_\mu = \frac{1}{2} [u^\dagger, \partial_\mu u] = \frac{1}{2} (u^\dagger \partial_\mu u + u \partial_\mu u^\dagger) \quad u_\mu = \frac{i}{2} \{u^\dagger, \partial_\mu u\} = \frac{i}{2} (u^\dagger \partial_\mu u - u \partial_\mu u^\dagger)$$

$$B = \begin{pmatrix} \frac{1}{\sqrt{2}}\Sigma^0 + \frac{1}{\sqrt{6}}\Lambda & \Sigma^+ & p \\ \Sigma^- & -\frac{1}{\sqrt{2}}\Sigma^0 + \frac{1}{\sqrt{6}}\Lambda & n \\ \Xi^- & \Xi^0 & -\frac{2}{\sqrt{6}}\Lambda \end{pmatrix} \quad \Phi = \begin{pmatrix} \frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & \pi^+ & K^+ \\ \pi^- & -\frac{1}{\sqrt{2}}\pi^0 + \frac{1}{\sqrt{6}}\eta & K^0 \\ K^- & \bar{K}^0 & -\frac{2}{\sqrt{6}}\eta \end{pmatrix}$$

Ecker, *Prog. Part. Nucl. Phys.*, 35:1-80 (1995), Bernard, Kaiser, Meißner, *Int. J. Mod. Phys. E*, 4:193-346 (1995), Oset, Ramos, *Nucl. Phys. A* 635 (1998) 99-120, Hyodo, Jido, *Prog. Part. Nucl. Phys.* 67(1):55-98 (2012)

Meson-baryon chiral effective theory

$$\mathcal{L}_{\text{WT}} = \frac{1}{4f^2} \text{Tr} [\bar{B} i \gamma^\mu [\Phi \partial_\mu \Phi - (\partial_\mu \Phi) \Phi, B]]$$



$$V_{ij} = \frac{-C_{ij}}{4f_i f_j} \bar{u}(j) \gamma_\mu u(i) (k_i + k_j)^\mu$$

$$Q = -1$$

$$K^- n, \pi^0 \Sigma^-, \pi^- \Sigma^0, \pi^- \Lambda, \eta \Sigma^-, K^0 \Xi^-$$

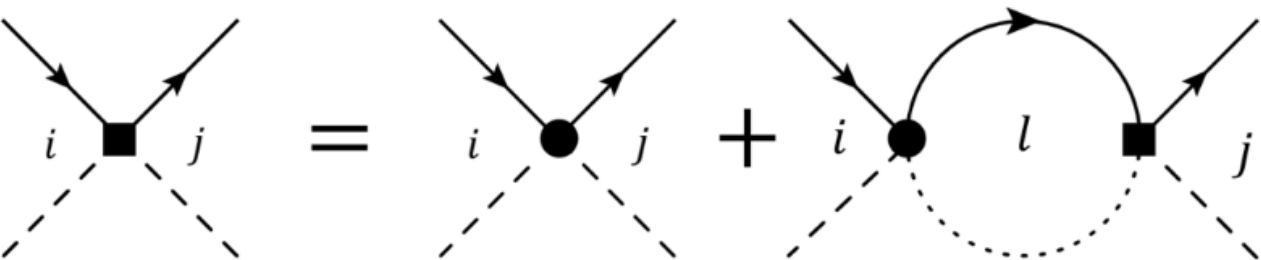
$$Q = 0$$

$$K^- p, \bar{K}^0 n, \pi^0 \Lambda, \pi^0 \Sigma^0, \pi^+ \Sigma^-, \pi^- \Sigma^+, \eta \Lambda, \eta \Sigma^0, K^+ \Sigma^-, K^0 \Sigma^0$$

$$V_{ij}^{s\text{-wave, on-shell}} = -\frac{C_{ij}}{4f_i f_j} \sqrt{\frac{E_j + M_j}{2M_j}} \sqrt{\frac{E_i + M_i}{2M_i}} (2\sqrt{s} - M_i - M_j)$$

$$\text{Isospin limit: } f_i = f_j = 1.123 f_\pi$$

T-matrix equation with on-shell factorization

$$t_{ij} = V_{ij} + \sum_l V_{il} G_l t_{lj}$$


$$G_l(P) = i \int \frac{d^4 q}{(2\pi)^4} \frac{2M_l}{(P - q)^2 - M_l^2 + i\epsilon} \frac{1}{q^2 - m_l^2 + i\epsilon}$$

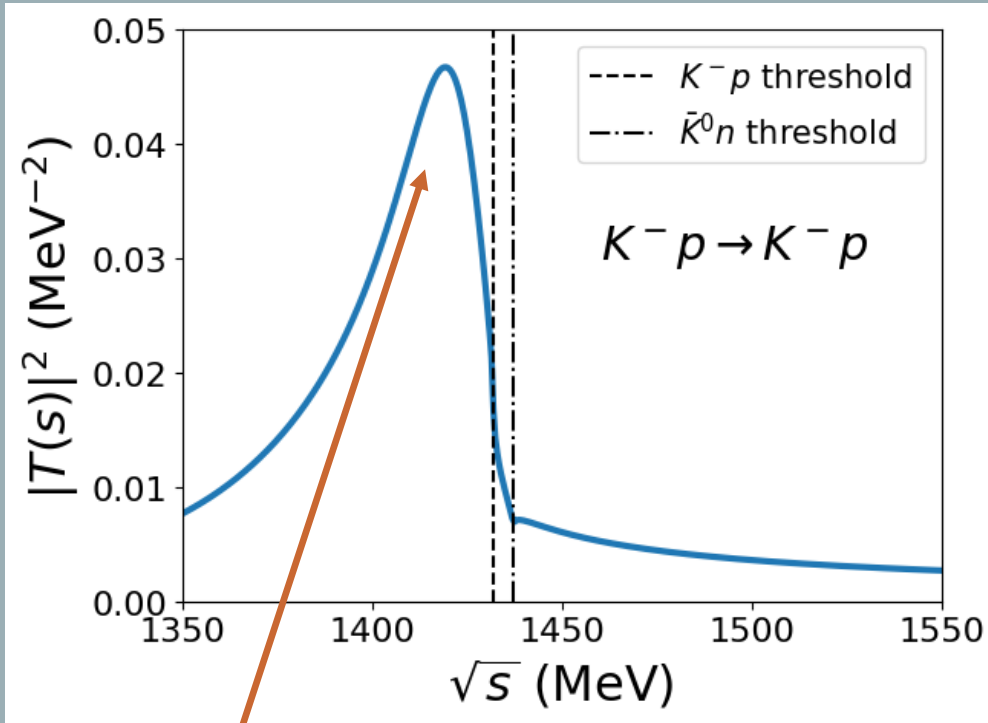
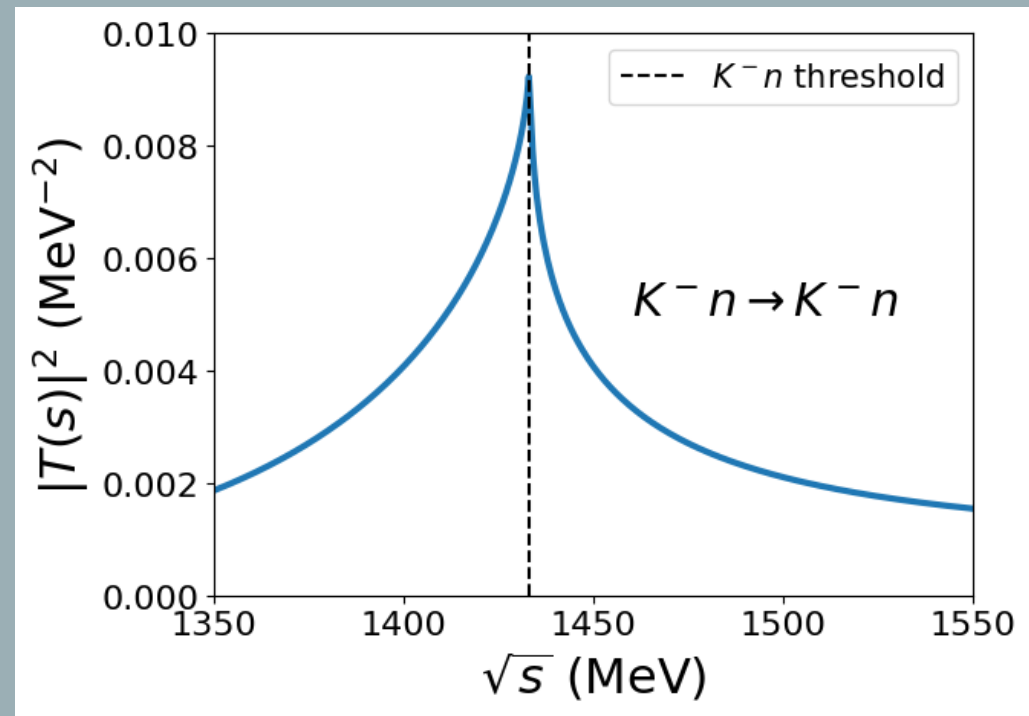
$$t = [1 - VG]^{-1} V$$

Dimensional regularization for loop functions at $\mu=630$ MeV

E. Oset, A. Ramos, *Nucl.Phys.A* 635 (1998) 99-120

J.A. Oller, E. Oset, *Nucl.Phys.A* 620 (1997) 438-456, J.A. Oller, U.G. Meißner, *Phys.Lett.B* 500 (2001) 263-272

T matrices for (elastic) K^-p and K^-n

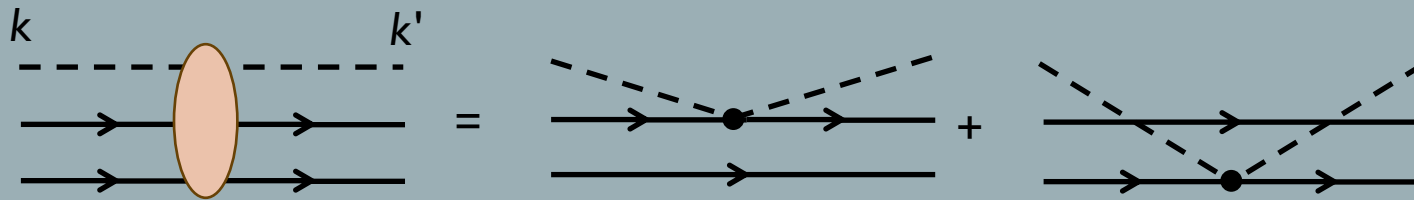
 t_p  t_n 

$\Lambda(1405)$ resonance below K^-p threshold (but above $\pi\Sigma$)

E. Oset, A. Ramos, *Nucl. Phys. A* 635 (1998) 99-120; J.A. Oller, U.G. Meißner, *Phys. Lett. B* 500 (2001) 263-272; S. Navas et al. (Particle Data Group), *Phys. Rev. D* 110, 030001 (2024)

Kaon-Deuteron interaction

Impulse Approximation

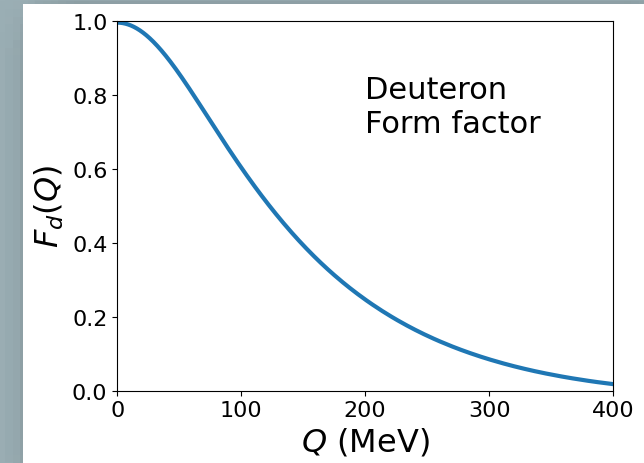


$$T_{Kd}(k', k) = [t_p(k', k) + t_n(k', k)] F_d(Q)$$

Chiral Kp interaction

Chiral Kn interaction

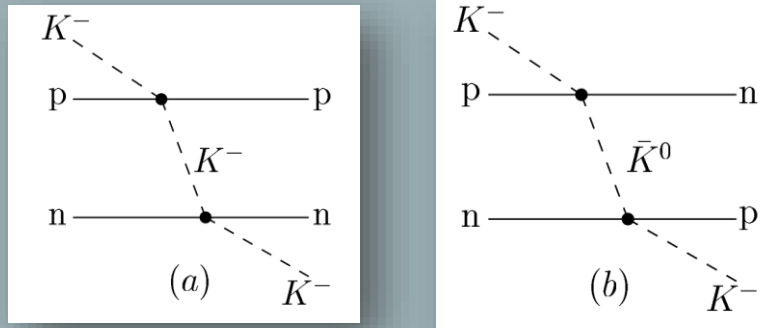
$$F_d(Q) = \int e^{-i\mathbf{Q}\cdot\mathbf{r}} |\phi_d(\mathbf{r})|^2 d\mathbf{r}$$



Argonne AV18 wave functions

Kaon-Deuteron interaction

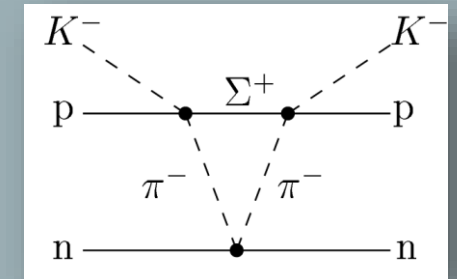
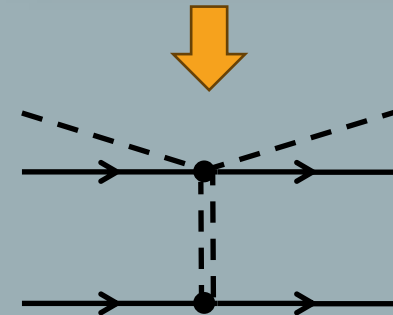
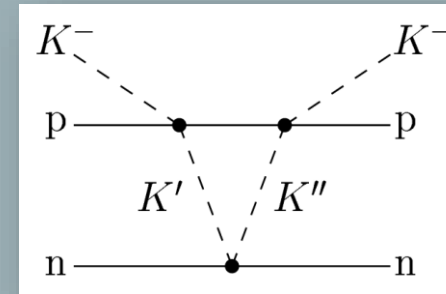
Double scattering approximation



Cancellation of off-shell 2-body T operators with chiral 3-body contact terms

Fixed-Center Approximation

Triple scattering approximation



Only 1% contribution to the K^-d amplitude

S.S. Kamalov, E. Oset, A. Ramos, *Nucl.Phys.A* 690 (2001) 494-508;

E. Oset et al., *Nucl.Phys.A* 881 (2012) 127-140;

A. Martinez-Torres et al. *Phys.Rev.D* 78 (2008) 074031; A. Martinez-Torres, Ph.D. thesis, IFIC (2009)

Deuterons in RHICs

- Deuteron production in HICs still under debate (coalescence versus thermal models)

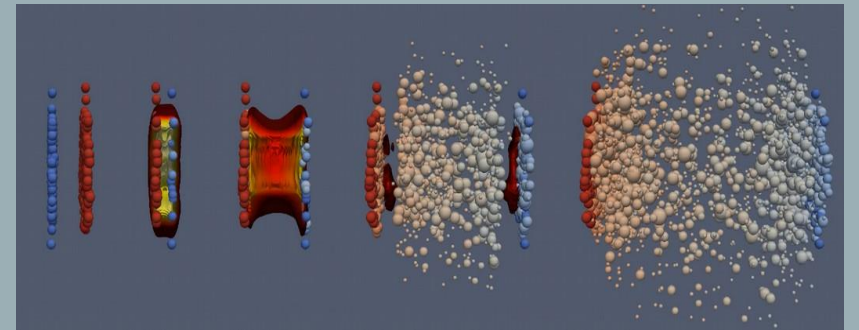
L. P. Csernai and J. I. Kapusta, *Phys. Rept.* 131, 223 (1986); P. Danielewicz and G. F. Bertsch, *Nucl. Phys. A* 533, 712 (1991)., R. Scheibl and U.W. Heinz, *Phys. Rev. C* 59, 1585 (1999), A. Andronic, P. Braun-Munzinger, and J. Stachel, *Nucl. Phys. A* 772, 167 (2006), A. Andronic, P. Braun-Munzinger, K. Redlich, and J. Stachel, *Nature* 561, 321 (2018), V. Vovchenko et al. *Phys. Lett. B* 800, 135131 (2020); T. Neidig, et al. *Phys. Lett. B* 827, 136891 (2022); J. Rais, H. van Hees and C. Greiner, *Phys. Rev. C* 106 (2022), 064004...

- Deuteron **femtoscopy** might help to distinguish between them

S. Mrówczyński and P. Słoń, *Acta Phys. Polon. B* 51, 1739 (2020)

- Hints of deuteron early formation via deuteron-pion femtoscopy

O. Vázquez Doce, D. Mihaylov, L. Fabbietti, *Eur. Phys. J. A* 61 (2025) 3, 53



MADAI collaboration,
Hannah Petersen and Jonah Bernhard

Deuterons in RHICs

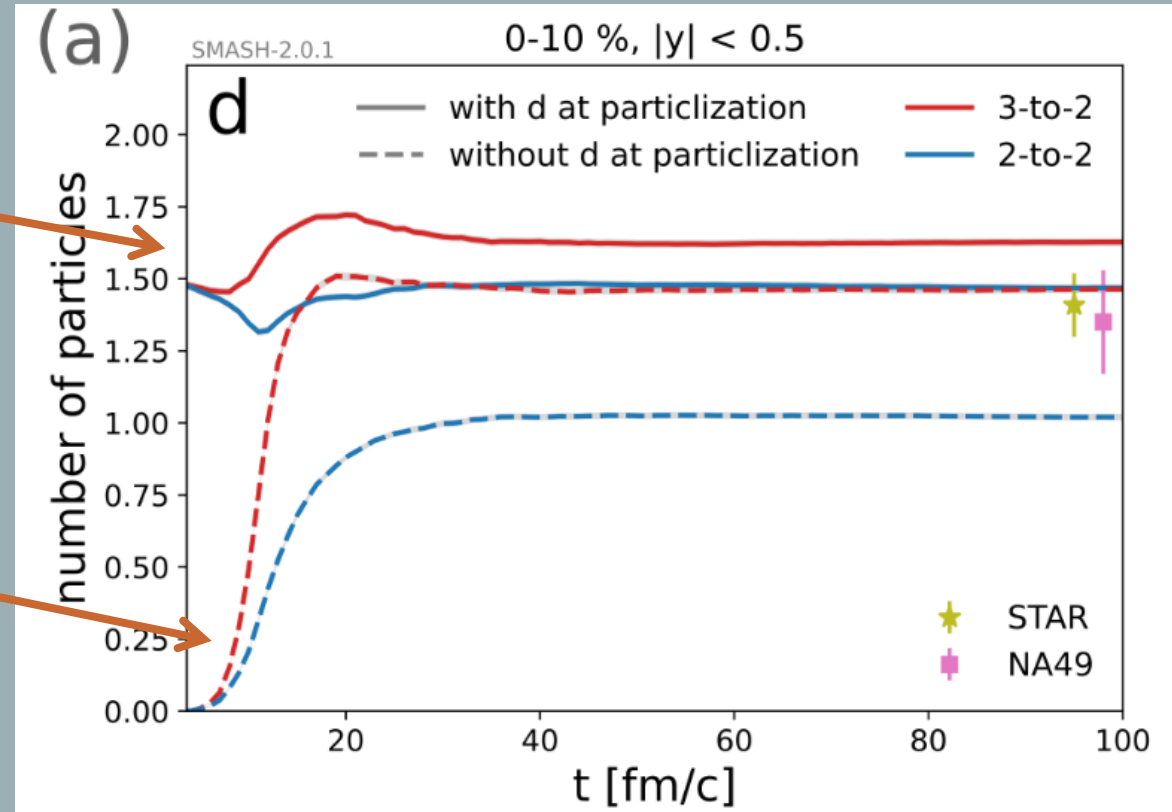
Time evolution of deuteron yield in Au Au collisions at $s^{1/2}_{NN} = 7.7$ GeV



Statistical Thermal Model:
Production at freeze out

Coalescence Model:
Dynamical production of
composed hadrons

“Stochastic rates” to account for boost-invariant, 3-to-2 reactions:

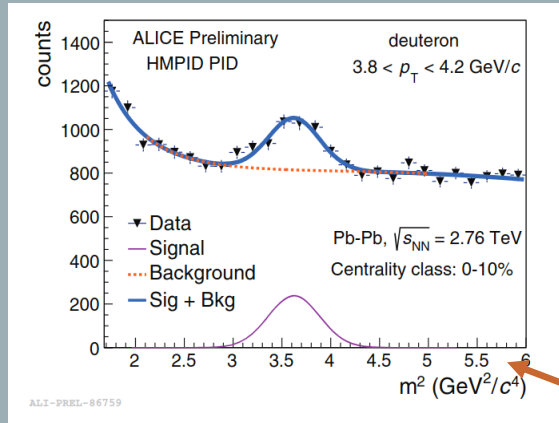


J. Staudenmaier, D. Oliinychenko, JMT-R and H. Elfner
Phys.Rev.C 104 (2021) 3, 034908

Femtoscscopy in RHICs

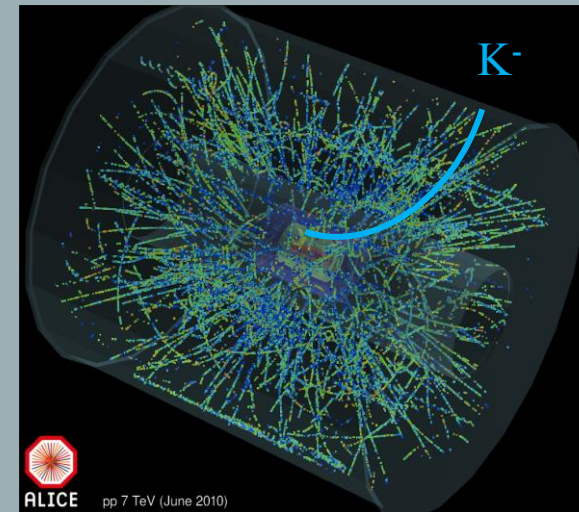
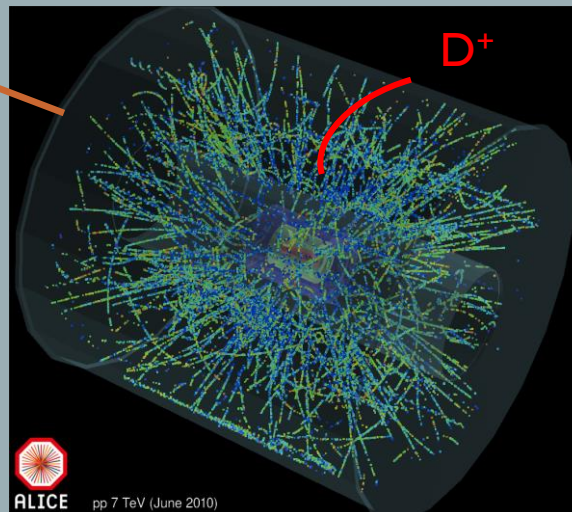
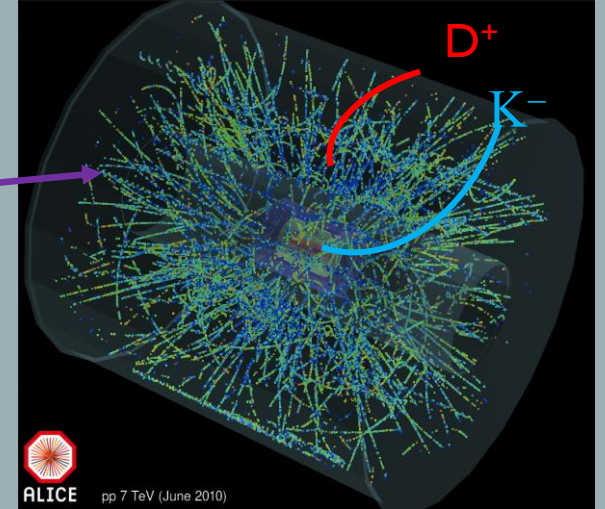
Pair Correlation Function

$$C(\mathbf{q}) = \mathcal{N} \frac{N_{\text{same}}(\mathbf{q})}{N_{\text{mixed}}(\mathbf{q})}$$



D^+ reconstruction

'Event mixing' technique



10^6 - 10^7
events

Correlation function

KP Correlation function with Coulomb interaction

$$C(q) = \int d^3r S(r) |\Phi^C(\mathbf{q}; \mathbf{r})|^2 + 4\pi \int r^2 dr S(r) \sum_{\ell} (2\ell+1) [|\varphi_{\ell}(q; r)|^2 - |\Phi_{\ell}^C(qr)|^2]$$

Complete Coulomb wave function

ℓ -wave strong + Coulomb wf

ℓ -wave Coulomb wf

As in 2-body, we keep only s-wave interactions

Two-body approach

S. Mrowczynski, 2501.08283

The three-body approach to hadron-deuteron correlations changes into the two-body approach if the three-particle hadron-deuteron wave function factorizes into the deuteron wave-function and the wave function of a hadron motion relative to the deuteron. The assumption of factorization is less restrictive for the kaon-deuteron than the proton-deuteron system, as there is no ambiguity which particle should be bound to the neutron.

Kaon-deuteron akin to having a faithful two-body description

We consider an effective two-body approach from Faddeev equations

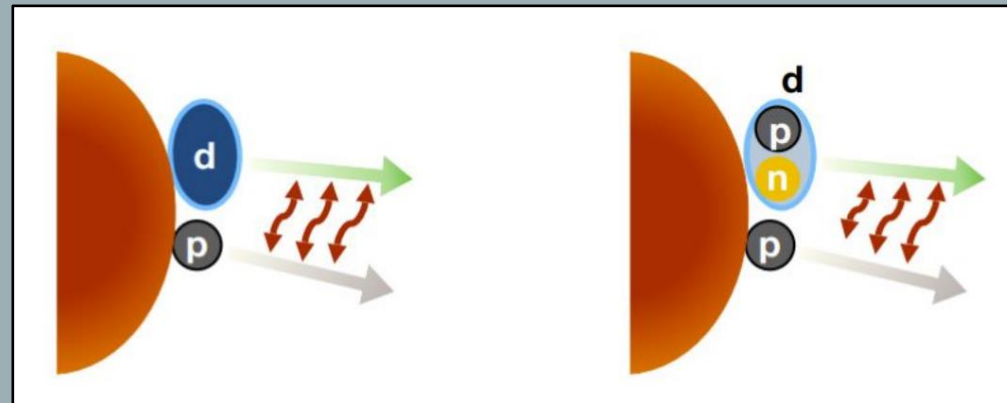


Image Credit:
ALICE coll.

Kaon-Deuteron Correlation Function

We use T-matrix to compute the wave function (Lippmann-Schwinger equation)

$$\varphi_i(q; r) = j_0(qr)\delta_{if} + \int_0^\infty \frac{4\pi q'^2 dq'}{(2\pi)^3} \frac{T_{if}(q', q; \sqrt{s}) j_0(q'r)}{2\omega_{H,i} 2\omega_{\phi,i}(\sqrt{s} - \omega_{H,i} - \omega_{\phi,i} + i\eta)}$$

$$T(q', q; \sqrt{s}) = T^S(\sqrt{s}) + T^C(q', q; \sqrt{s})$$

Total Coulomb wave function and $\ell=0$ partial wave

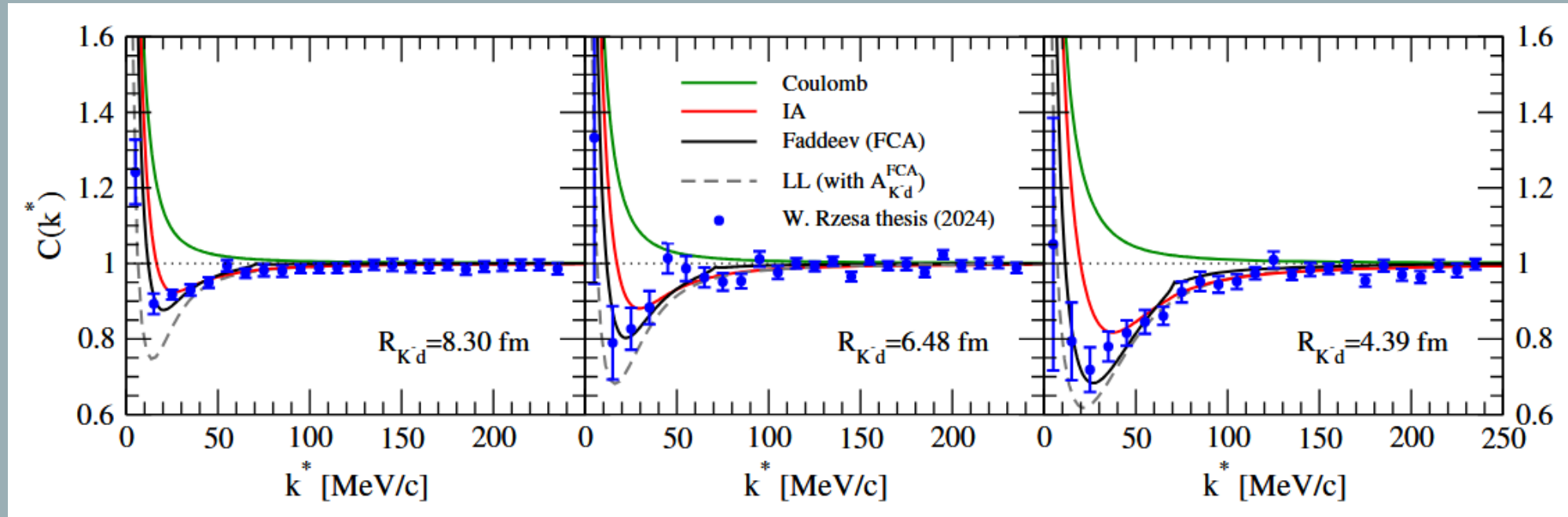
$$\Phi^C(r, z; k) = e^{-\pi\gamma/2} \Gamma(1 + i\gamma) e^{ikz} {}_1F_1(-i\gamma; 1; ik(r - z))$$

$$\gamma = Z_1 Z_2 \mu \alpha / k$$

$$\Phi_{0f}^C(qr) = (qr)^{-1} e^{i\sigma_0} F_0(\gamma; qr)$$

$$\sigma_l = \arg \Gamma(1 + l + i\gamma)$$

ALICE results in PbPb collisions



$$|\psi(r^*, k^*)| = \sqrt{A_C(\eta)} \left[\exp(-ik^*r^*) F(-i\eta, 1, i\xi) + f_c(k^*) \frac{G}{r^*} \right]$$

$$S(\vec{r}) \sim \exp\left(-\frac{r^2}{2R_{Kd}^2}\right)$$

Asymptotic part of the wave function

W. Rzes (ALICE Coll.) *PhD Thesis*

Effect of charge-exchange channel

A. Feijoo, À. Ramos, JMT-R, in preparation

	FCA-UNIT (OR model)	FCA-UNIT (BCN model)
Unitary Factor	1.00/1.00	1.00/1.00
Scattering Length [fm]	-1.04/-1.12 + i0.90/0.84	-0.85/-0.93 + i1.07/1.00

Inelastic channel removal $t_x = 0$

	FCA-UNIT (OR model)	FCA-UNIT (BCN model)
ε_{1s}^d [eV]	583/607	525/547
Γ_{1s}^d [eV]	767/701	958/876

Inelastic channel removal $t_x = 0$



2602.22833 [nucl-ex]

$$\begin{aligned} \text{Re } A_{K-d} &= -1.44 \pm 0.15 \text{ (stat)} \text{ }^{+0.10}_{-0.10} \text{ (syst) fm} \\ \text{Im } A_{K-d} &= +1.34 \pm 0.33 \text{ (stat)} \text{ }^{+0.21}_{-0.15} \text{ (syst) fm} \end{aligned}$$



2608.10749 [nucl-ex]

$$\begin{aligned} \varepsilon_{1s}^d &= 810.9 \pm 24.5 \text{ (stat)} \pm 2.1 \text{ (syst) eV} \\ \Gamma_{1s}^d &= 812 \pm 97 \text{ (stat)} \pm 33 \text{ (syst) eV} \end{aligned}$$