

Dispersive and model-independent extraction of πN poles from revised total cross-section data

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[arXiv:2608.20458](https://arxiv.org/abs/2608.20458) [hep-ph]



NSTAR 2026, Seville, September 7th 2026

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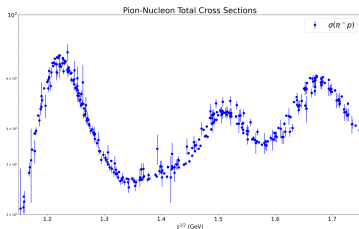
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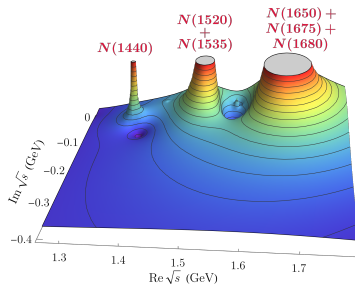
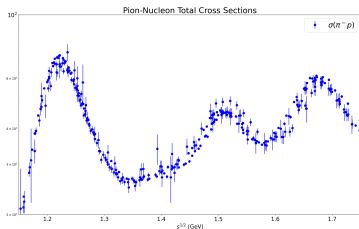
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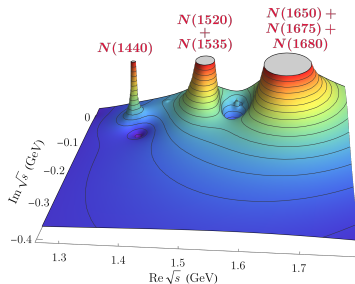
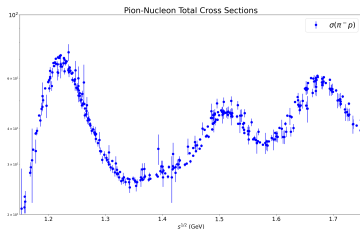
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2. Can we estimate the **isospin-breaking effects** in the $\Delta(1232)$ sector?

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- ⇒ **Dispersive analyses (FDRs) and a robust analytic continuation**

Why **dispersive**?

- * Consequence of **fundamental principles**, generally not imposed when extracting resonances
- * Integrals decrease the **parametrization dependence**

1. Motivation

Why Forward Dispersion Relations (FDRs)?

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- * They hold up to ∞ (Roy-like DR do not)
- * Through the **optical theorem** ($t=0$), the imaginary part can be directly related to **total cross-section data**, easier to measure!
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Why do we need a robust **analytic continuation** method?

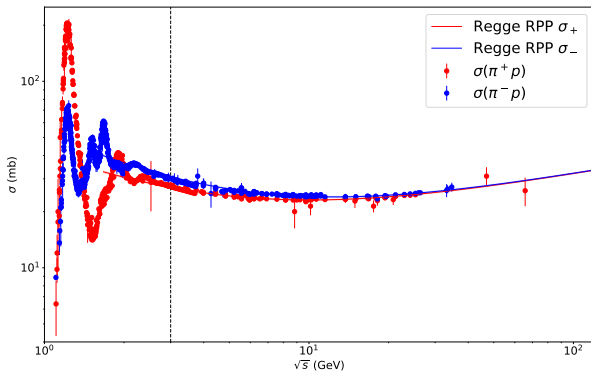
- * We only know exactly how to access the second Riemann sheet
- * Also reduces the **parametrization dependence**

1. Motivation

* No coupled-channel analysis $\Rightarrow$ just focus on πN

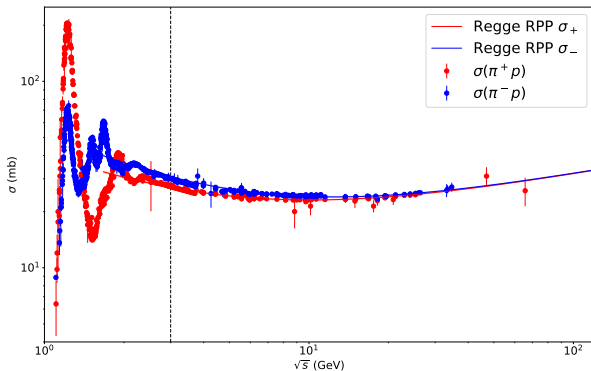
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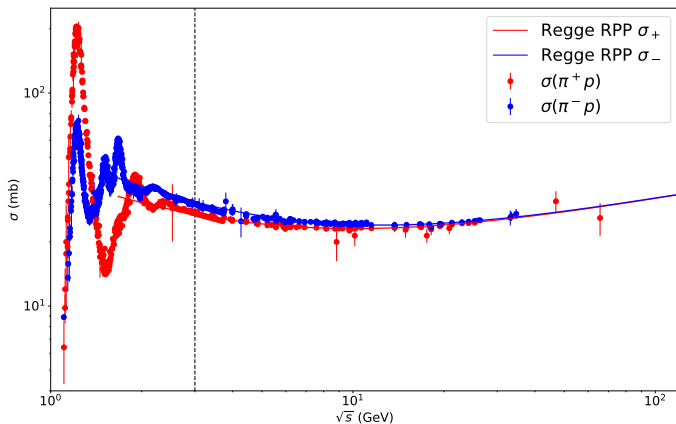


- * Determine resonance parameters in a different way (just total cross-section data, no partial waves)

2. πN total cross-section data

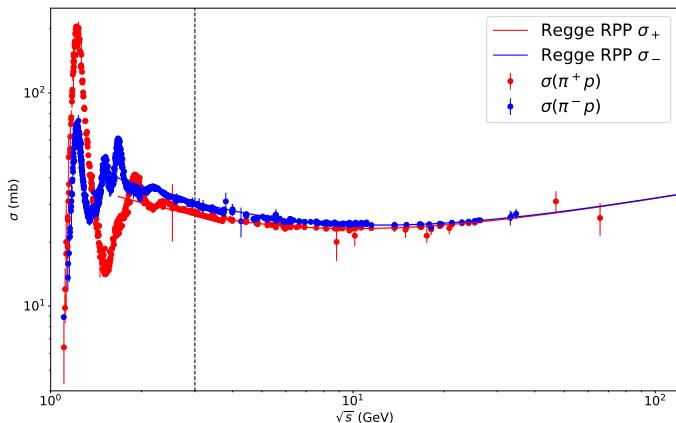
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Strategy:

$$\sigma_{\text{tot}}(\pi^\pm p) \Rightarrow \text{Im}D^\pm \Rightarrow \text{Re}D^\pm \Rightarrow \text{Resonances}$$

2. πN amplitudes and $\sigma_{\text{tot}}(\pi^\pm p)$

For $\pi^a(q) + N(p) \rightarrow \pi^b(q') + N(p')$

$$T_{fi}^{ba}(s, t) = \delta^{ba} T_{fi}^+(s, t) + i\epsilon^{bac} \tau^c T_{fi}^-(s, t)$$

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The relation between these amplitudes and the total cross-section data is the following (**optical theorem**):

$$\text{Im} D_\pm(s, 0) = k_L \sigma_{\text{tot}}(\pi^\pm p) \equiv k_L \sigma_\pm$$

with

$$D^+ = \frac{1}{2}(D_+ + D_-), \quad D^- = \frac{1}{2}(D_- - D_+),$$

and k_L the pion momentum in the laboratory frame

2. FDRs for D amplitudes

With these relations it is straightforward to obtain **once-subtracted FDRs!**
($\bar{D}^\pm$ are the pseudovector-Born-term-subtracted amplitudes)

$$\text{Re}\bar{D}^+(k_L) = 4\pi \left(1 + \frac{M_\pi}{m_N}\right) \bar{a}_{0+}^+ + \frac{k_L^2}{\pi} \mathcal{P} \int_0^\infty \frac{\sigma_+(k'_L) + \sigma_-(k'_L)}{k_L'^2 - k_L^2} dk'_L$$

$$\text{Re}\bar{D}^-(k_L) = \frac{g^2 \omega_L}{2m_N^2} + \frac{\omega_L}{\pi} \mathcal{P} \int_0^\infty \frac{k_L'^2}{\omega_L'} \frac{\sigma_-(k'_L) - \sigma_+(k'_L)}{k_L'^2 - k_L^2} dk'_L$$

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with

$$M_\pi \equiv \text{pion mass}, \quad m_N \equiv \text{proton mass}, \quad \omega_L = \sqrt{k_L^2 + M_\pi^2},$$

$$g^2 \equiv \text{pseudoscalar coupling constant},$$

and $\bar{a}_{0+}^+$ is related to even isospin S-wave scattering length, precisely extracted from **pionic-atoms data** Baru et al. PLB 694, 473 (2011), Baru et al. Nucl.Phys.A 872, 69 (2011)

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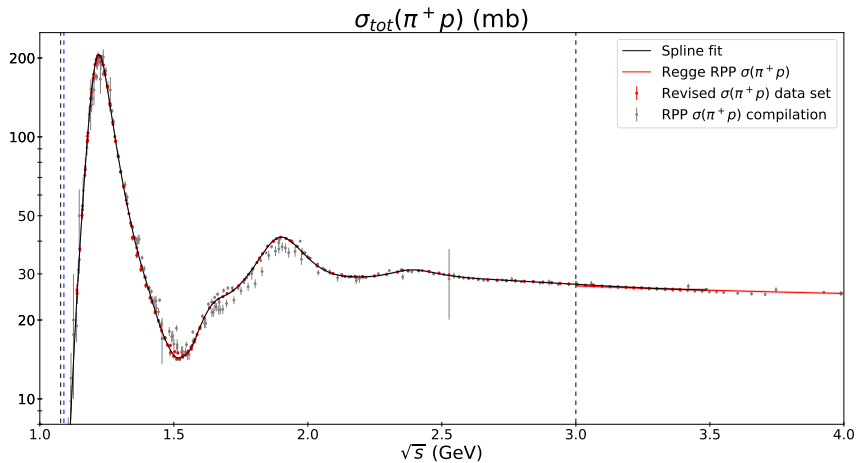
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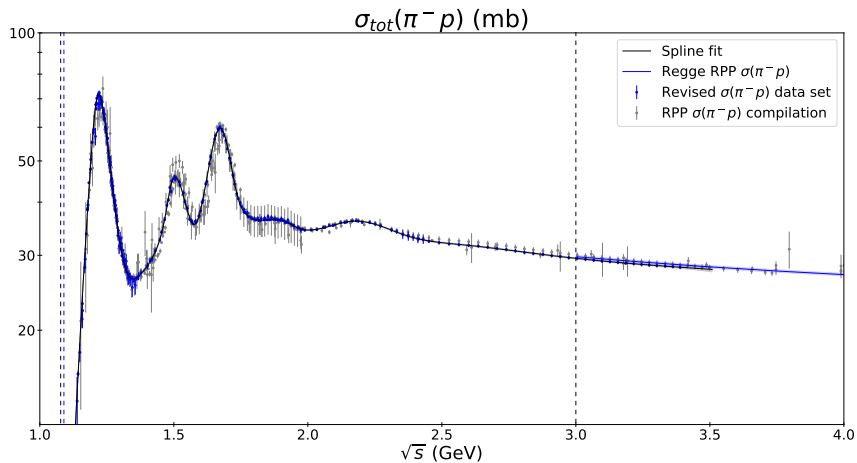
- * Fit to Roy-Steiner analysis near threshold (tiny region) to stabilize the fit
Hoferichter et al., Phys. Rept 625 (2016)

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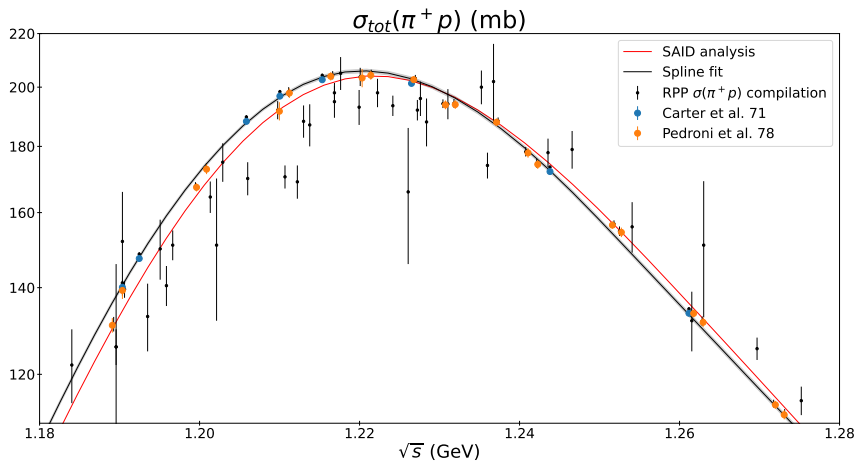
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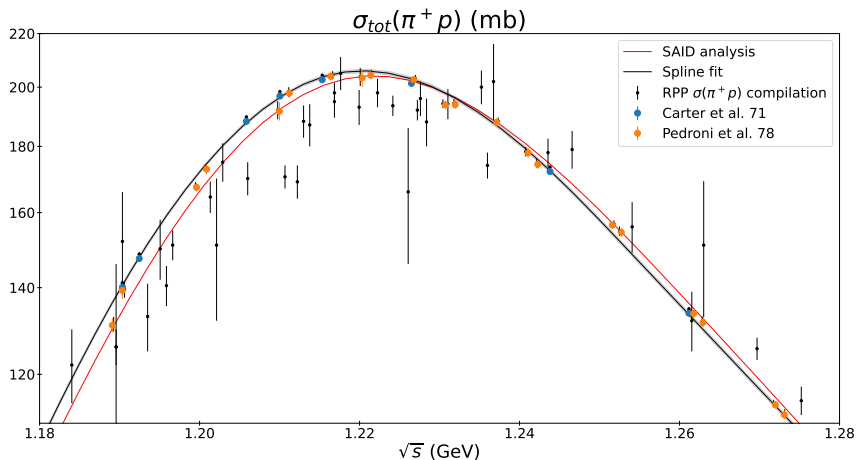
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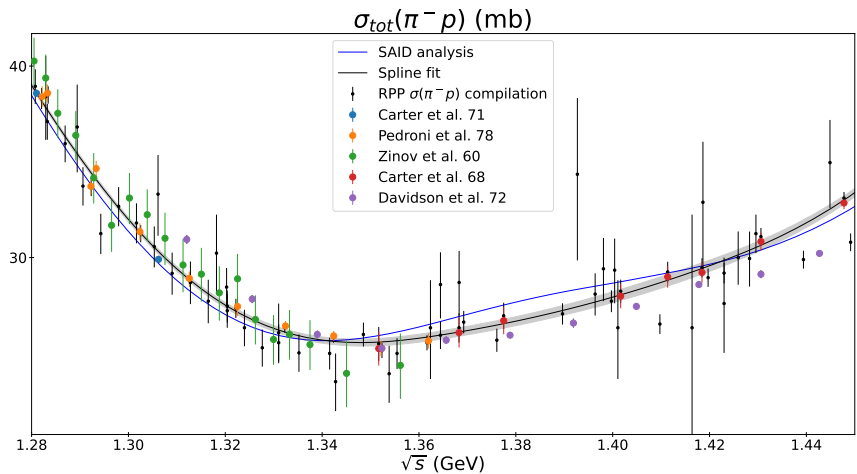
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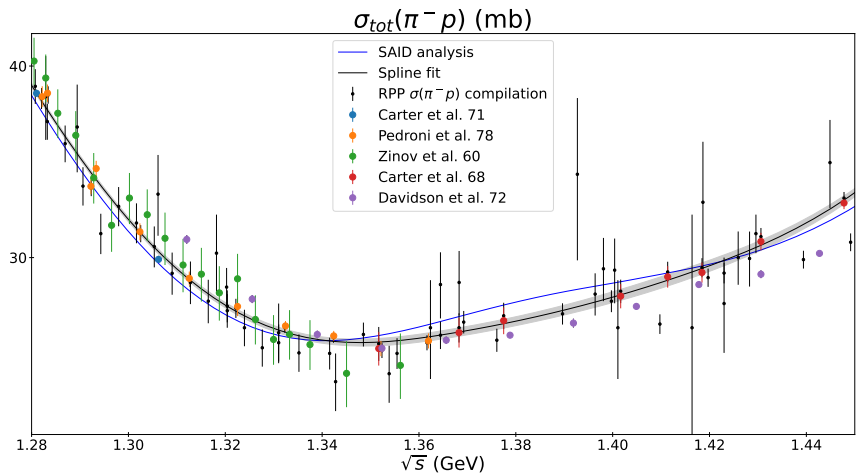
* In the resonance region, **fair agreement with SAID** partial-wave analysis (different data)

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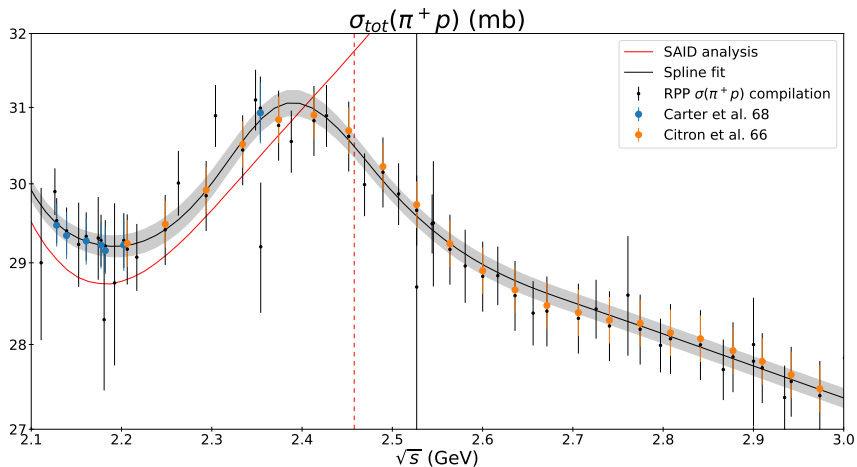


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* Special case: $N(1440)$ Roper resonance $\Rightarrow$ Clear differences

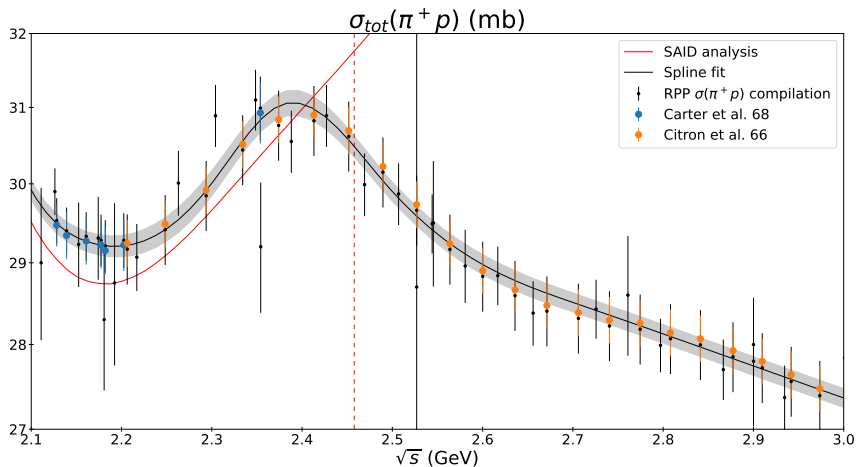
$\Rightarrow$ The Roper effect is imperceptible in the $\sigma_{\text{tot}}(\pi^{\pm} p)$ data

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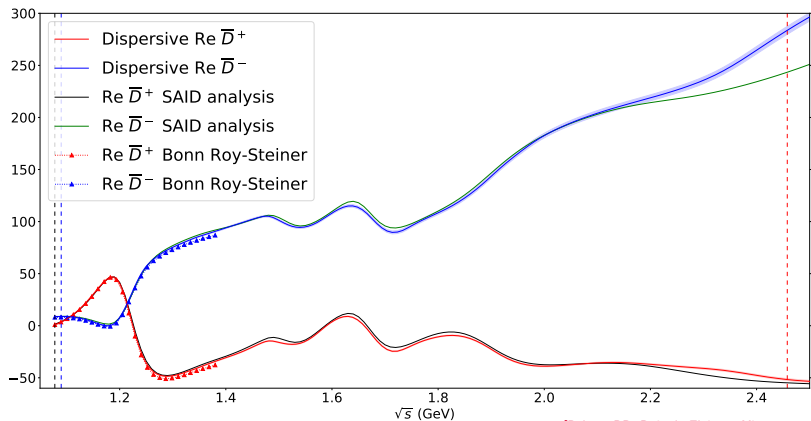
* Huge differences at higher energies! (Above SAID upper validity limit)

3. πN FDRs' output

* Coulomb distortion correction with Tromborg method Tromborg et al. PRD 15, 725 (1976)
and isospin $\Delta(1232)$ correction from SAID database (thanks Strakovsky and Döring!) to have purely hadronic amplitudes for the FDRs!

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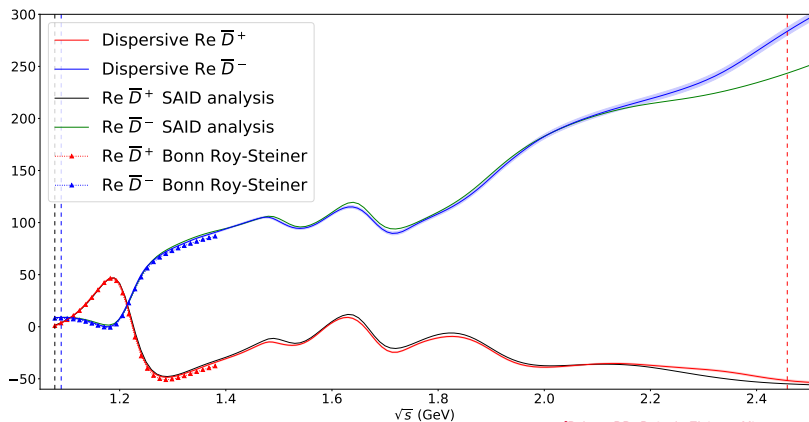
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- * **Good agreement** with both analyses up to 2.2 GeV

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FDR_{C_N} : “Interpolate” the FDR output with this functional form in a real s segment, which contains **poles** ($\sim N/2$) (already tested in $\pi\pi$)

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- * **We only keep stable poles under N , segments and fit parameters!**

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- * We now have the FDR outputs $\Rightarrow$ We can apply the FDR_{C_N} method!

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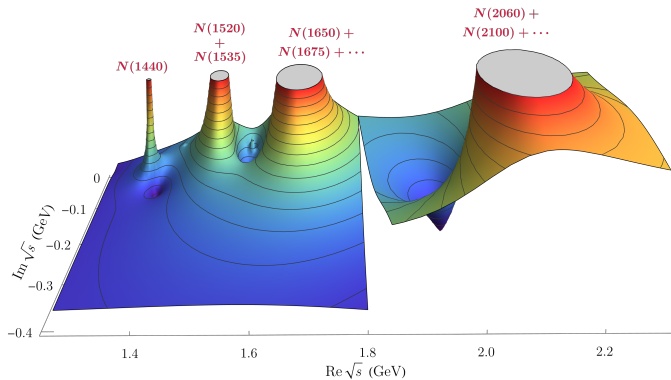
- * **But we can** distinguish the **isospin** of resonances!

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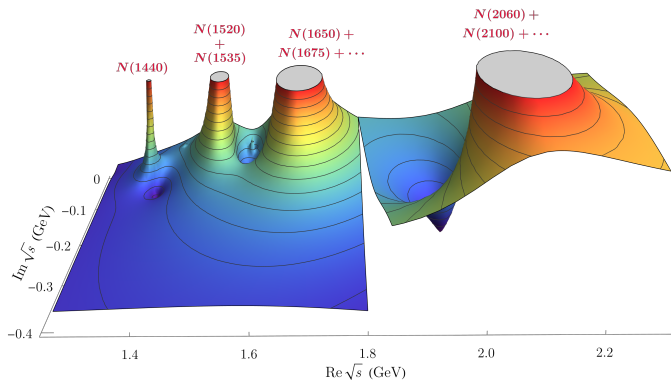
- * We can then extract isospin-definite **M** , **Γ** , and **residues**!

4. πN pole-resonance extraction $I = 1/2$



[Pelaez, PR, Ruiz de Elvira, arXiv:2608.20458]

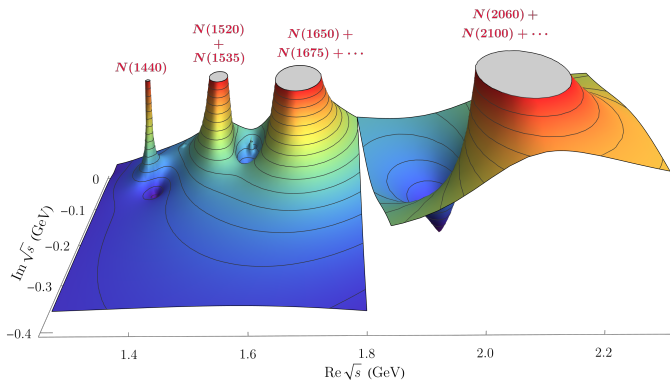
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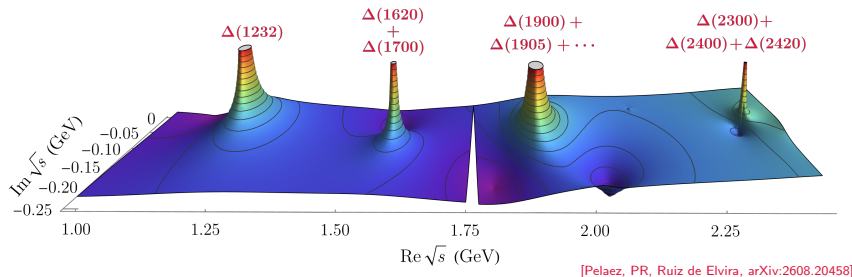
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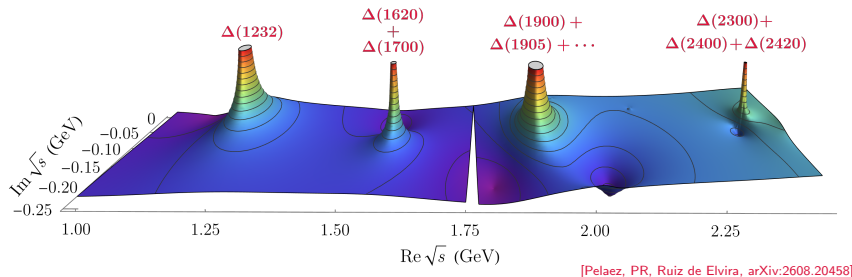
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- * We obtain 4 stable $I = \frac{1}{2}$ poles that we associate with these resonances
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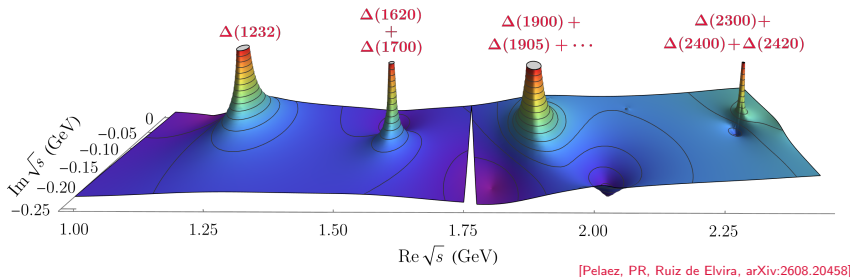


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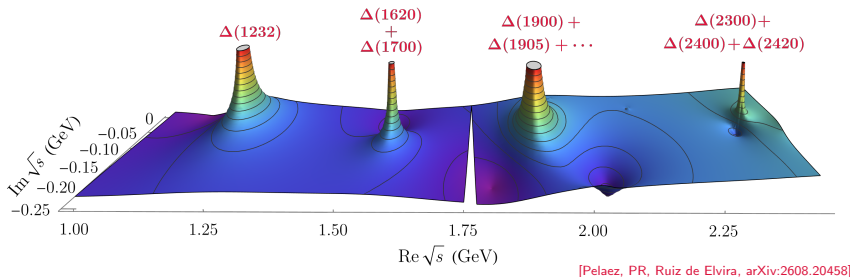
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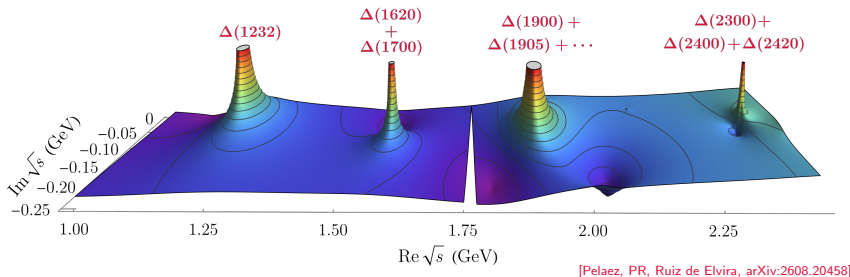
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- * Total cross-section data have not enough resolution to distinguish two or more overlapping resonances $\Rightarrow$ Some poles correspond to several resonances!
- * However, we have more information: residue of the poles

4. πN pole-resonance extraction $I = 3/2$



- * We obtain 4 stable $I = \frac{3}{2}$ poles that we associate with these resonances
- * Total cross-section data have not enough resolution to distinguish two or more **overlapping resonances** $\Rightarrow$ Some poles correspond to several resonances!
- * However, we have more information: **residue of the poles**
 $\Rightarrow$ By doing a Cauchy integral enclosing several j resonances

$$R_{\text{FDR}} = \sum_j R_j$$

4. FDR_{CN} : Resonance identification from the residues

* **Pole residues** help identify the contributing resonances!

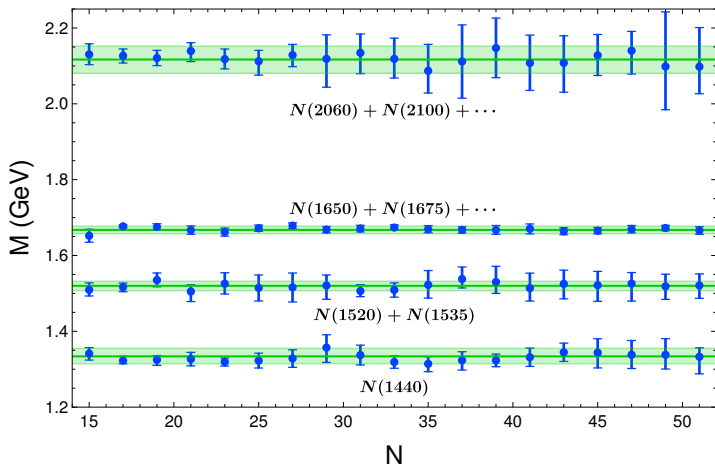
Pole	$ R_{FDR} $ (GeV)	$ R_{PDG} $ (GeV)
$N(1440)$	6_{-3}^{+4}	7.8 ± 0.7
$N(1520) + N(1535)$	18_{-8}^{+9}	12.8 ± 1.6
$N(1650) + \dots + N(1720)$	32 ± 13	34 ± 6
$N(2060) + \dots + N(2250)$	133_{-47}^{+51}	77 ± 20
$\Delta(1232)$	25 ± 10	18.3 ± 0.7
$\Delta(1620) + \Delta(1700)$	10 ± 4	7 ± 4
$\Delta(1905) + \Delta(1910) + \Delta(1950)$	32 ± 6	36 ± 7
$\Delta(2420)$	14 ± 3	17 ± 8

[Pelaez, PR, Ruiz de Elvira, arXiv:2608.20458]

And same for the residue phase!

4. FDR_{CN} : Stability of the analytic continuation

* **Stable** results for different N (e.g. for the $l = \frac{1}{2}$ poles)



[Pelaez, PR, Ruiz de Elvira, arXiv:2608.20458]

4. $\Delta(1232)$ isospin-breaking effects: Strategy

* Our final objective is to try to estimate the isospin-breaking effects in the $\Delta(1232)$ sector. Δ^0/Δ^{++} appear in $\sigma_{\text{tot}}(\pi^\pm p)$ data, respectively

4. $\Delta(1232)$ isospin-breaking effects: Strategy

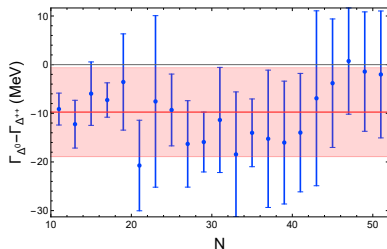
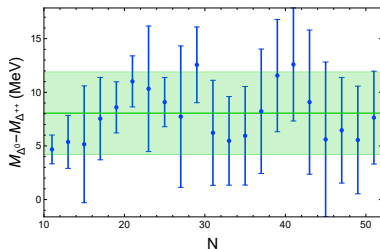
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 - $\Rightarrow$ **Problem:** FDRs involve both $\sigma_\pm \rightarrow$ non isospin invariant amplitude

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- * Compute FDRs without the isospin and electromagnetic corrections
 - $\Rightarrow$ **Problem:** FDRs involve both $\sigma_\pm \rightarrow$ **non isospin invariant amplitude**
- * FDRs for $D_\pm$ in the $\Delta(1232)$ region are dominated by the RHC, i.e. $\sigma_\pm$
 - $\Rightarrow$ **LHC contribution small** compared to our uncertainties

4. $\Delta(1232)$ isospin-breaking effects: Results

Results:



[Pelaez, PR, Ruiz de Elvira, arXiv:2608.20458]

$$M_{\Delta^0} - M_{\Delta^{++}} = 8 \pm 4 \text{ MeV}$$

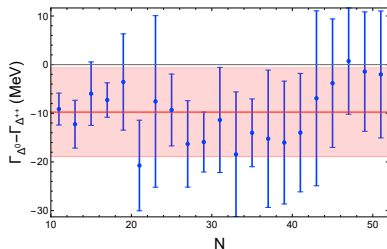
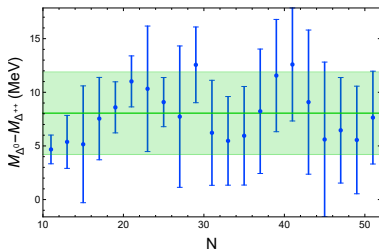
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$$|R_{\Delta^0}| - |R_{\Delta^{++}}|/3 = -0.8 \pm 1.2 \text{ GeV}$$

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* Large uncertainties, but clear hints of $M_{\Delta^0} > M_{\Delta^{++}}$ in agreement with PDG!

★ Consistent with the PDG estimate and Arndt et al. PRC 74, 045205 (2006) for isolated resonances (Roper is difficult to obtain from $\sigma_{\text{tot}}(\pi^\pm p)$ data!)

	$\sqrt{s_{\Delta(1232)}} \text{ (MeV)}$
FDR $_{C_N}$	$(1211 \pm 4) - i(52 \pm 4)$
PDG estimate	$(1210 \pm 1) - i(50 \pm 1)$
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★ IB effects in Δ^0/Δ^{++} estimated $\Rightarrow$ Clear indication of $M_{\Delta^0} > M_{\Delta^{++}}$

Thank you for your attention

