



UNIVERSITÀ
DI TORINO



Andrea Signori

University of Turin and INFN

Tomography of hadrons in 3D

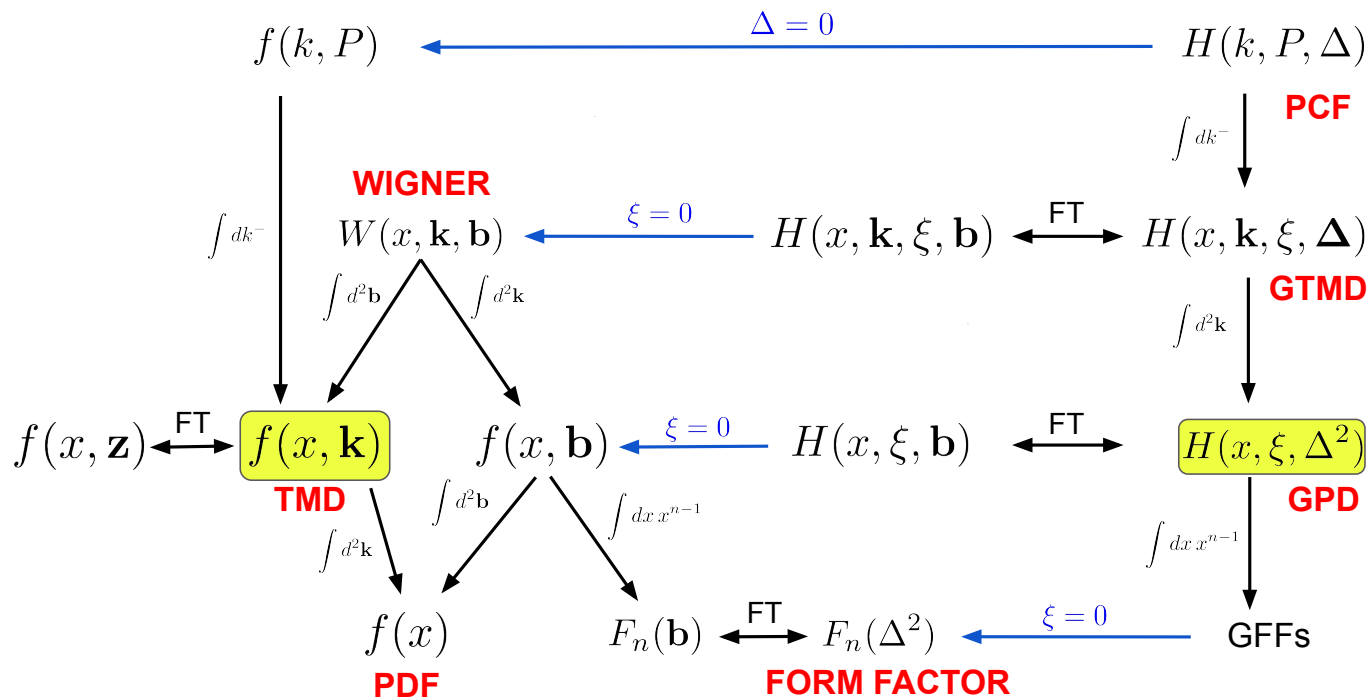


Outline

1. **Introduction: TMDs and GPDs**
2. **TMD formalism**
3. **Challenges and open questions**

1. Introduction

The hadron structure landscape



See <https://inspirehep.net/literature/3148063>

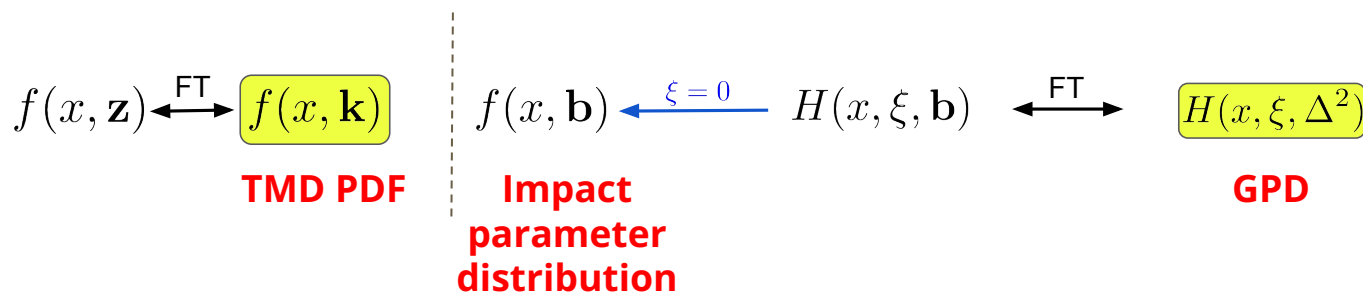
TMDs and GPDs



$$x = k^+ / P^+$$

$$\xi = -\Delta^+ / 2P^+$$

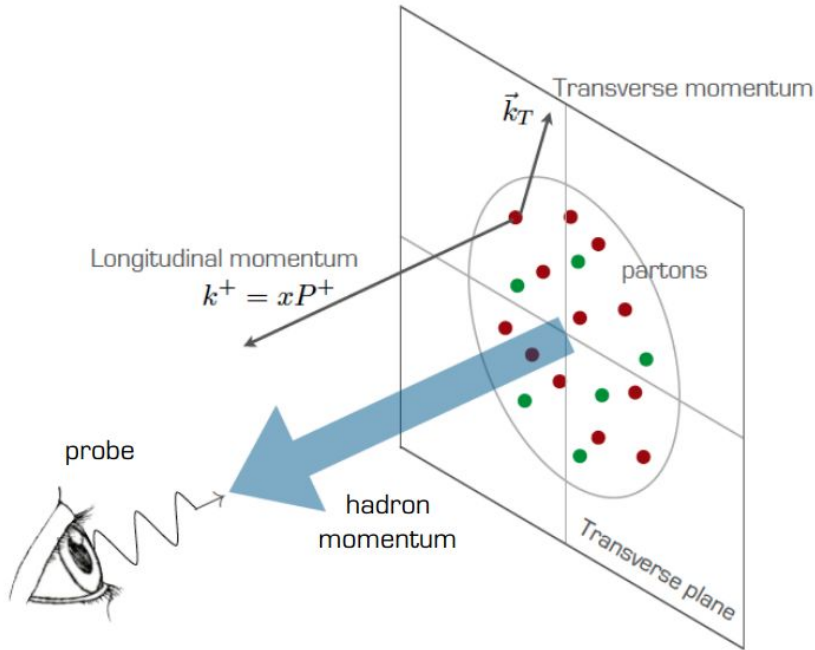
- **TMDs:** 3D picture in **momentum** space
- **GPDs:** picture in 1D **momentum** + 2D **position** space (**impact parameter distributions**)
- The transverse momentum **k** and transverse position **b** are **NOT related** to each other by a Fourier transform (**k** vs **z**)



For more details on GPDs, see C. Mezrag's talk and the associated contributions

Parton distribution functions

“Maps” of hadron structure in momentum space



$$f_1(x)$$

1D structure
in momentum space
("collinear")

$$f_1(x, k_T^2)$$

3D structure
in momentum space
("transverse momentum
dependent")

Credit picture: A. Bacchetta

Distributions with transverse hadron spin

$$f(x, \mathbf{b}, \mathbf{s}) = H(x, \mathbf{b}^2) + \frac{\epsilon^{ij} \mathbf{b}_i \mathbf{s}_j}{M} \frac{\partial}{\partial \mathbf{b}^2} E(x, \mathbf{b}^2)$$

Unpolarized GPD (I.P.D.)

Helicity-flip GPD (I.P.D.)

?

$$f^{[U]}(x, \mathbf{k}, \mathbf{s}) = f_1^{[U]}(x, \mathbf{k}^2) - \frac{\epsilon^{ij} \mathbf{k}_i \mathbf{s}_j}{M} f_{1T}^{\perp[U]}(x, \mathbf{k}^2)$$

Unpolarized TMD PDF

Sivers TMD PDF

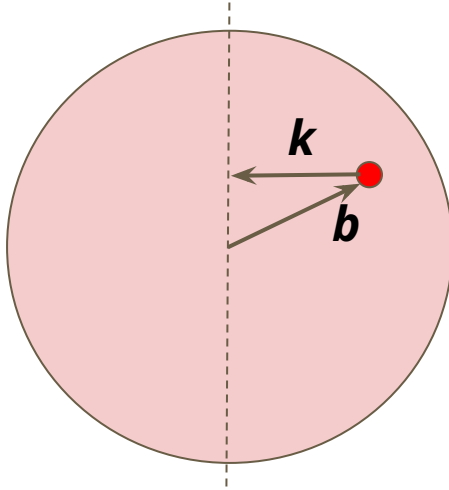
Process dependence!

The lensing mechanism

$$f(x, \mathbf{b}, \mathbf{s}) = H(x, \mathbf{b}^2) + \frac{\epsilon^{ij} \mathbf{b}_i \mathbf{s}_j}{M} \frac{\partial}{\partial \mathbf{b}^2} E(x, \mathbf{b}^2)$$

Symmetric distribution
in transverse plane

Distorted distribution
in transverse plane

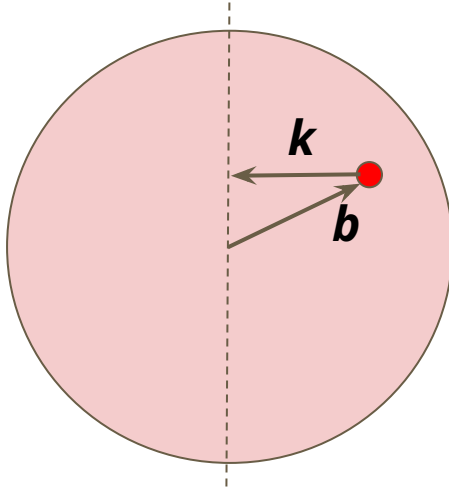


An attractive **color force** between the parton and the spectator turns the transverse position into a transverse momentum

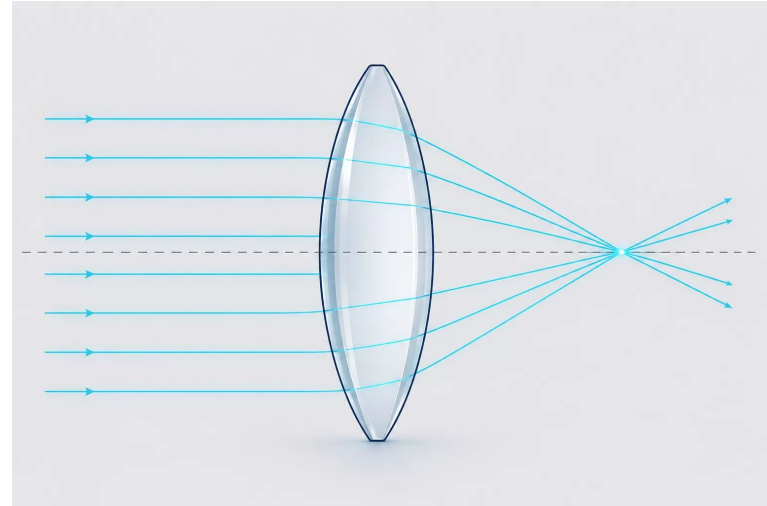
The lensing mechanism

$$f(x, \mathbf{b}, \mathbf{s}) = H(x, \mathbf{b}^2) + \frac{\epsilon^{ij} \mathbf{b}_i \mathbf{s}_j}{M} \frac{\partial}{\partial \mathbf{b}^2} E(x, \mathbf{b}^2)$$

Symmetric distribution
in transverse plane



Distorted distribution
in transverse plane

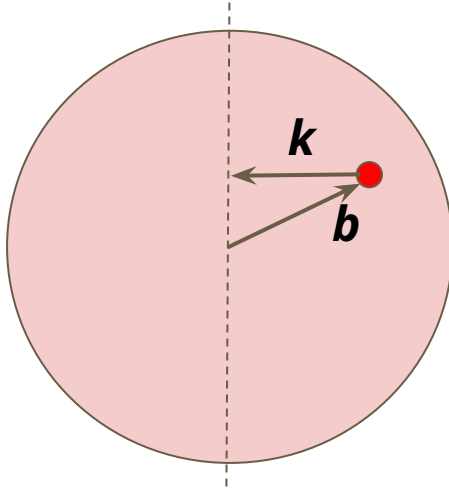


The lensing mechanism

$$f(x, \mathbf{b}, \mathbf{s}) = H(x, \mathbf{b}^2) + \frac{\epsilon^{ij} \mathbf{b}_i \mathbf{s}_j}{M} \frac{\partial}{\partial \mathbf{b}^2} E(x, \mathbf{b}^2)$$

Symmetric distribution
in transverse plane

Distorted distribution
in transverse plane



Lensing function $\mathbf{I}_\perp(\mathbf{b})$

$$\mathcal{P}(x, \mathbf{k}) = \int d^2\mathbf{b} \, q(x, \mathbf{b}) \, \delta^{(2)}(\mathbf{k} - \mathbf{I}_\perp(\mathbf{b}))$$

Lensing: a function of the **confining potential**

GPDs and angular momentum

Ji's sum rule

$$J^q(\mu) = \frac{1}{2} \int_{-1}^1 dx x \left[H^q(x, \xi, \Delta^2; \mu) + E^q(x, \xi, \Delta^2; \mu) \right]_{\Delta^2=0}$$

↑
Total angular momentum of (anti)quarks in
proton polarized along the z direction

↑
E flips the helicity of the target,
while preserving the helicity of the quark:
it is connected to **orbital angular momentum**

$$f_{1T}^{\perp(1)}(x) \sim \int d^2\mathbf{b} \, \mathbf{I}_{\perp}(x, \mathbf{b}_{\perp}) \frac{\partial}{\partial \mathbf{b}^2} E(x, \mathbf{b}^2)$$

Lensing of **E** GPD - **Sivers** TMD PDF → **OAM**

CAVEAT: this is a strictly **model-dependent** relation:
S. Rodini et al. - <https://inspirehep.net/literature/1744103>

2. TMD formalism

TMD and collinear PDFs for quarks in nucleon

		quark pol.		
		U	L	T
nucleon pol.	U	f_1		$h_1^\perp$
	L		g_{1L}	$h_{1L}^\perp$
	T	$f_{1T}^\perp$	g_{1T}	$h_1, h_{1T}^\perp$

$$\Phi_{ij}(k, P) = \text{F.T.} \langle P | \bar{\psi}_j(0) U \psi_i(\xi) | P \rangle$$

At leading twist: 8 TMD PDFs

(similar classification for gluons and for fragmentation functions)

- **Black**: time-reversal even AND collinear
- **Blue**: time-reversal even
- **Red**: time-reversal odd (*process dependence*)

The **symmetries of QCD** play a crucial role in this classification

Why should we study these distributions ?

f_1

- Test **factorization** and **universality**
- *Precise* knowledge: impact on **HEP**, e.g. **mW** determination

h_1

- Tensor charge of the nucleon: **CP** violation and access to **BSM** physics

$f_{1T}^\perp, h_1^\perp$

- **(TMD only)** Test the **symmetries** of QCD

e

- Quark-gluon correlations and quark contribution to **hadron mass**

E

- Quark-gluon correlations and **dynamical** generation of **quark mass**

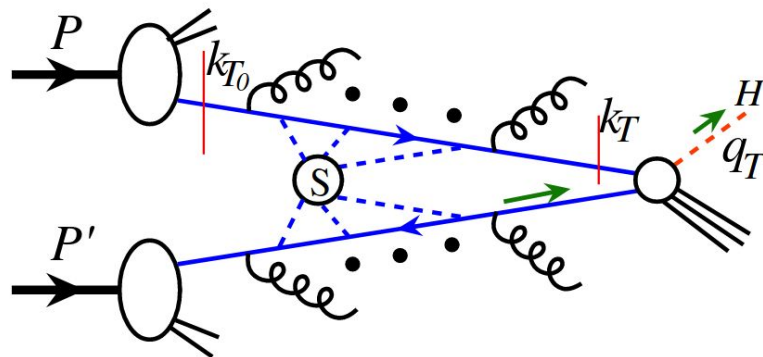
TMD factorization

$$q_T \ll Q$$

$$pp \longrightarrow \gamma^* / Z \longrightarrow l \bar{l} + X$$

$$\frac{d\sigma}{dq_T} \sim \mathcal{H} f_1(x_a, k_{Ta}, Q, Q^2) f_1(x_b, k_{Tb}, Q, Q^2) \delta^{(2)}(q_T - k_{Ta} - k_{Tb}) + \mathcal{O}(q_T/Q) + \mathcal{O}(\Lambda/Q)$$

- TMDs & partonic cross section:
same **IR poles** = same non-perturbative physics
- **observed transverse momentum q_T** :
transverse momenta of **quarks**
- quark transverse momentum :
radiative (perturbative) and **intrinsic**
(non-perturbative) components
- Renormalization = **evolution** equations tell us
how to distinguish between the two



Quarks

Drell-Yan / Z / W production (hh)

Semi-Inclusive DIS (eh)

2h-inclusive
e+e- annihilation

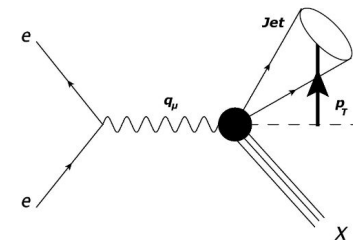
TMD
factorization

Gluons

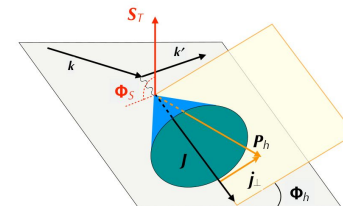
Higgs production
in hadronic collisions

Quarkonium production (e.g. $\eta_{b,c}$,
double J/Psi in hadronic collisions)

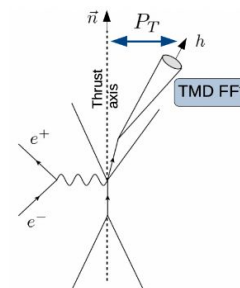
Jets:
(e.g. jet SIDIS,
di-jet SIDIS)



Hadron “in jet”:
(eh, hh, e+e-)



Y “+ jet”:
(e.g. Y = γ, h)



Quarks

Drell-Yan / Z / W production (hh)

Semi-Inclusive DIS (eh)

2h-inclusive
e+e- annihilation

Global fits (unpolarized)

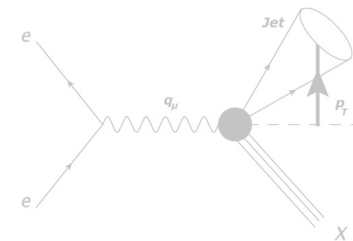
Gluons

Higgs production
in hadronic collisions

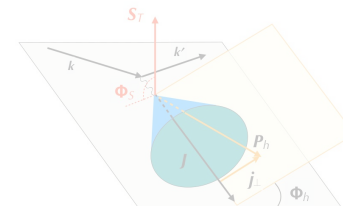
Quarkonium production (e.g.
 $\eta_{b,c}$, double J/Psi in hadronic
collisions)

**Not enough data
(or not at all)
for the other
processes**

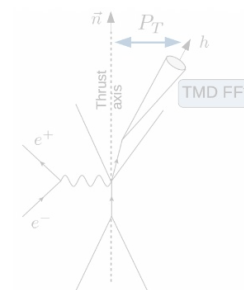
Jets:
(e.g. jet SIDIS,
di-jet SIDIS)



Hadron "in jet":
(eh, hh, e+e-)



Y "+ jet":
(e.g. Y = γ, h)



Approximations

$$pp \longrightarrow \gamma^* / Z \longrightarrow l \bar{l} + X$$

$$\frac{d\sigma}{dq_T} \sim \mathcal{H} f_1(x_a, k_{Ta}, Q, Q^2) f_1(x_b, k_{Tb}, Q, Q^2) \delta^{(2)}(q_T - k_{Ta} - k_{Tb})$$

[TMD region, $q_T \ll Q$]

$$+ \mathcal{O}(q_T/Q) + \mathcal{O}(\Lambda/Q) \quad [\text{large } q_T \text{ and low } Q \text{ corrections}]$$



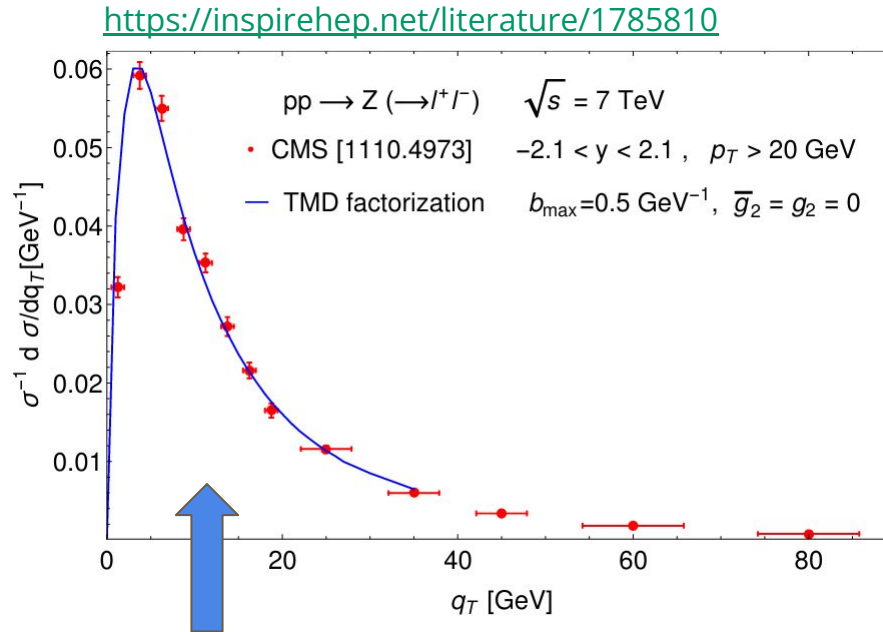
Large transverse
momentum corrections:
**"Matching" to
collinear factorization
at large q_T**



Higher twist corrections:
Multiparton correlations (?)

TMD region: low transverse momentum

$$q_T \ll Q$$

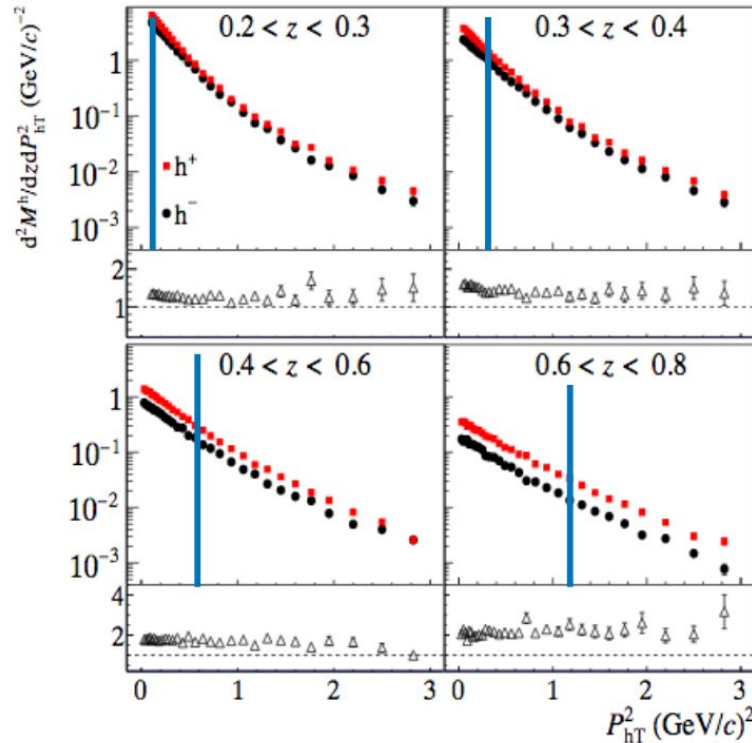


Hadronic collisions

$$q_T / Q < 0.3$$

TMD region: low transverse momentum

$$q_T \ll Q$$



SIDIS - TMD region

$$P_{hT}^2/z^2 \ll Q^2$$

Let's highlight

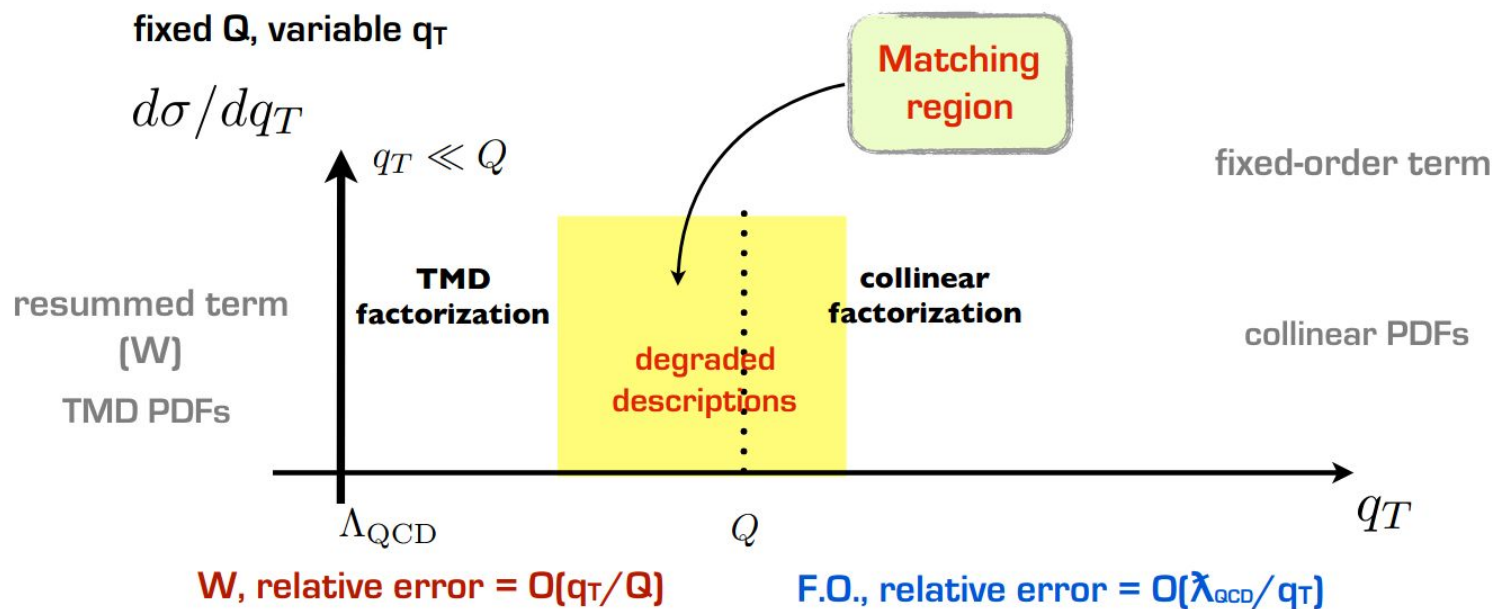
$$P_{hT}^2/z^2 \sim 0.25 Q^2$$

One of the bins with highest Q :

$$\langle Q^2 \rangle = 9.78 \text{ GeV}^2$$

$$\langle x \rangle = 0.149$$

Matching TMD and collinear factorization



Evolution equations

$$F_a(x, b_T^2; \mu, \zeta) = \int d^2 \mathbf{k}_T e^{i \mathbf{k}_T \cdot \mathbf{b}_T} F_a(x, k_T^2; \mu, \zeta)$$

← Fourier transform of a TMD PDF
(bT conjugated to kT)

$$\frac{d \ln F_a(x, b_T^2; \mu, \zeta)}{d \ln \zeta} = -D(b_T \mu, \alpha_s(\mu)) , \quad \text{rapidity}$$

$$\frac{d \ln F_a(x, b_T^2; \mu, \zeta)}{d \ln \mu} = \gamma_F \left(\alpha_s(\mu), \frac{\zeta}{\mu^2} \right) , \quad \text{UV}$$

$$\frac{d D(b_T \mu, \alpha_s(\mu))}{d \ln \mu^2} = \frac{1}{2} \gamma_K(\alpha_s(\mu)) , \quad \text{UV}$$

$$\frac{d \gamma_F \left(\alpha_s(\mu), \frac{\zeta}{\mu^2} \right)}{d \ln \zeta} = -\gamma_K(\alpha_s(\mu)) . \quad \text{rapidity}$$



QCD evolution of a TMD PDF

$$F_a(x, b_T^2; \mu, \zeta) = F_a(x, b_T^2; \mu_0, \zeta_0) \quad \rightarrow \text{TMD distribution at initial scales}$$

$$\times \exp \left[\int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_F \left(\alpha_s(\mu'), \frac{\zeta}{\mu'^2} \right) \right] \quad \rightarrow \text{evolution in } \mu$$

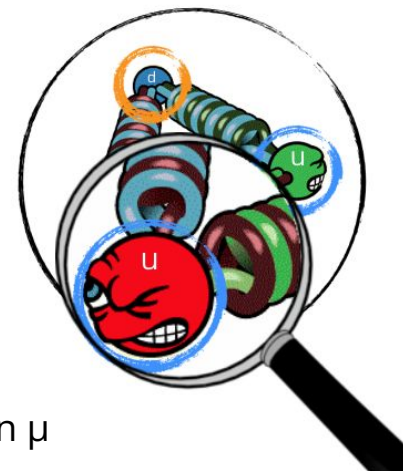
Calculable in pQCD

$$\times \left(\frac{\zeta}{\zeta_0} \right)^{-D(b_T \mu_0, \alpha_s(\mu_0)) + g_K(b_T; \lambda)} \quad \rightarrow \text{evolution in } \zeta$$

Non-pert. corrections (large b_T)

$$F_a(x, b_T^2; \mu_0, \zeta_0) = \sum_b \underbrace{C_{a/b}(x, b_T^2, \mu_0, \zeta_0)}_{\text{small } b_T \text{ (OPE)}} \otimes \underbrace{f_b(x, \mu_0)}_{\text{prior knowledge assumed (?)}} \underbrace{F_{NP}(b_T; \lambda)}_{\text{large } b_T}$$

Prior knowledge assumed (?)



Non-perturbative TMD parts

$$F_a(x, b_T^2; \mu, \zeta) = F_a(x, b_T^2; \mu_0, \zeta_0)$$

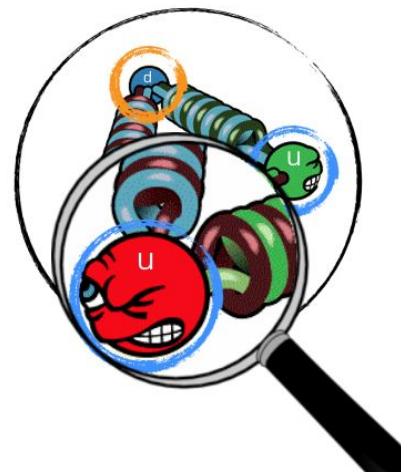
$$\exp \left[\int_{\mu_0}^{\mu} \frac{d\mu'}{\mu'} \gamma_F \left(\alpha_s(\mu'), \frac{\zeta}{\mu'^2} \right) \right]$$

$$\left(\frac{\zeta}{\zeta_0} \right)^{-D(b_T \mu_0, \alpha_s(\mu_0))} + g_K(b_T; \lambda)$$

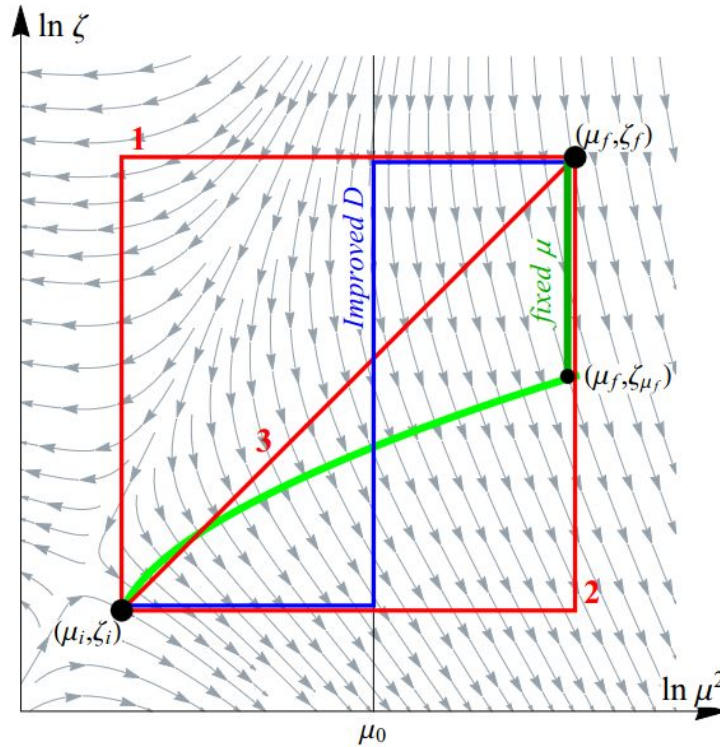
$$F_a(x, b_T^2; \mu_0, \zeta_0) = \sum_b C_{a/b}(x, b_T^2, \mu_0, \zeta_0) \otimes f_b(x, \mu_0) F_{NP}(b_T; \lambda)$$

**Non-pert. corrections
(large bT)**

**Intrinsic transverse
momentum,
potentially flavor
dependent**



A two-scale evolution



The combined evolution along the light green curve is absent.

Different frameworks, same observable

$$q_T^{\text{res.}} \propto_{\text{PB}} e^{2S} [f_1 \otimes \mathcal{H} \otimes f_2]$$

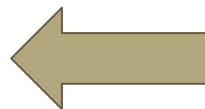
$$\left(\frac{d\sigma}{dq_T} \right)_{\text{res.}} \propto^{\text{TMD}} H \times F_1 \times F_2 + \mathcal{O} \left[\left(\frac{q_T}{Q} \right)^m \right]$$

$$\propto^{\text{SCET}} H \times B_1 \times B_2 \times S$$

$$\mathcal{H} = HC_1C_2$$

$$F_i = e^S C_i \otimes f_i$$

$$F_i = \sqrt{S} \times B_i$$

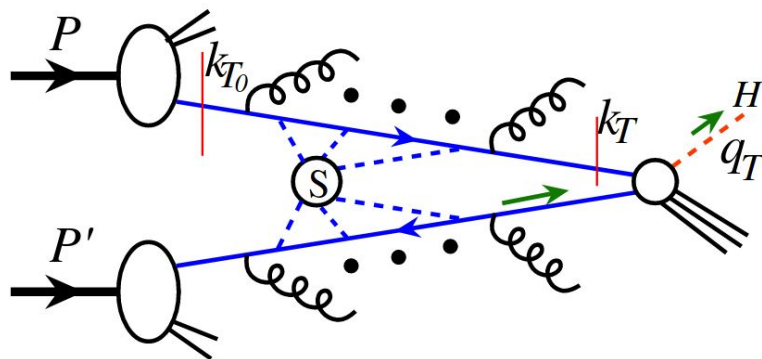


Dictionary to compare different factorization frameworks

“equivalent” to the extent of describing TMD physics

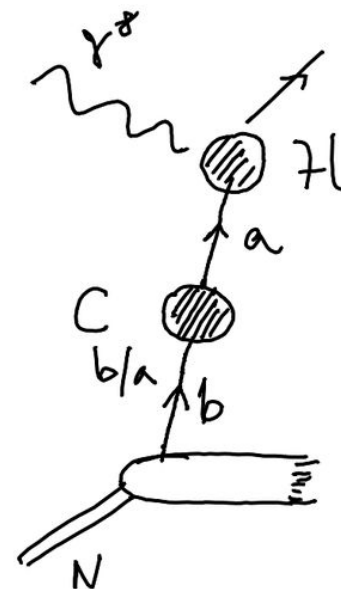
TMD evolution vs parton shower

- **Parton shower**: generates the radiation (**exclusively**) and then computes its recoil q_T



- **TMD evolution**: sums (**inclusively**) over the radiation and predicts the recoil q_T distribution.

$$F_a(x, b_T^2; \mu, \zeta)$$



3. Challenges

Challenges (TMD) as an outlook

- **Semi-inclusive DIS**: is the TMD factorization formalism under control?
(see talks by M. Cerutti, S. Piloneta)
- **Parametrization bias**: are we using the best methodology to extract information from data?
(see talk by M. Cerutti)
- **Factorization**: what can we say about its breaking? See also Gluon TMD-related observables
(see talk by C. Flore)
- **Higher twists**: we are getting an improved picture (theory and pheno developments)
(see talk by A. Vladimirov et al.)
- **Non-perturbative approaches**: what can we learn from LQCD and continuum methods?
(see talk by M. Ding et al.)
- **Hadronization**: it deserves separate presentations
(see talk by W. Brooks, F.G. Celiberto et al.)
- ...

Backup

TMD PDFs: traffic light notation

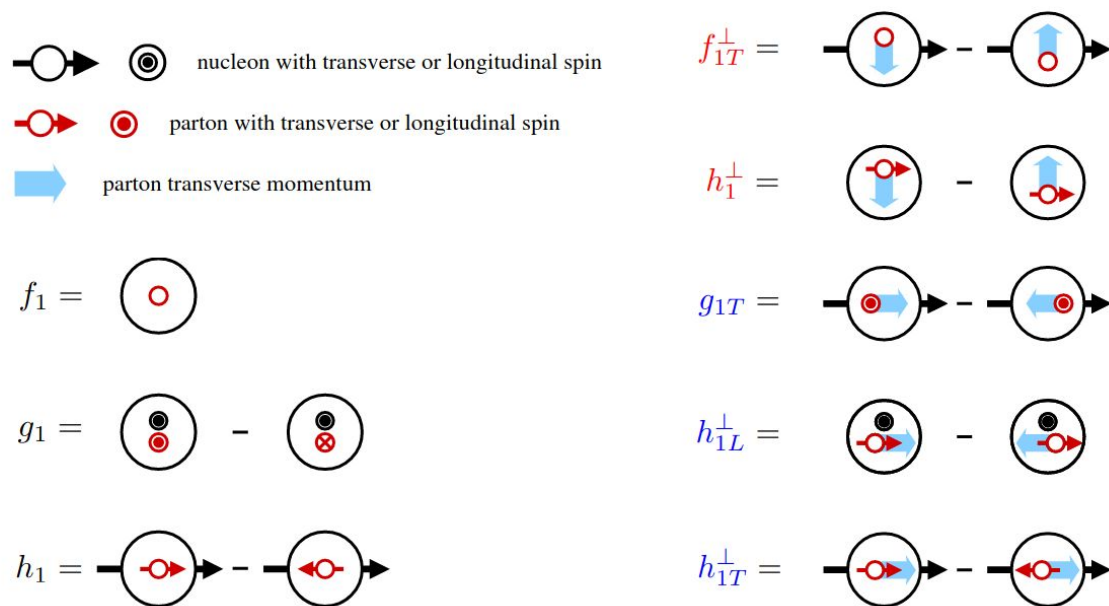


Figure 3.5: Probabilistic interpretation of twist-2 transverse-momentum-dependent distribution functions. To avoid ambiguities, it is necessary to indicate the directions of quark's transverse momentum, target spin and quark spin, and specify that the proton is moving out of the page, or alternatively the photon is moving into the page.

Non-TMD-factorizable processes

For $pp \rightarrow h1 h2 X$ TMD factorization is violated

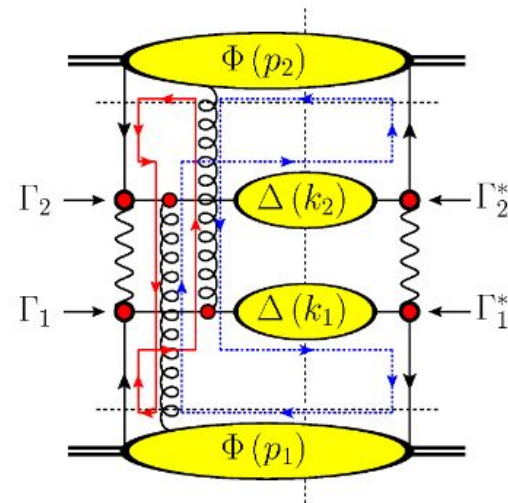
See :

- Collins, Qiu (2007) : <https://inspirehep.net/literature/750627>
- Rogers, Mulders (2010) : <https://inspirehep.net/literature/843028>
- Buffing (2016) : <https://inspirehep.net/literature/1391461> (see figure)

This **endangers** also other processes such as $pp \rightarrow h X$ (and similar)

Quantify factorization breaking effects ? See e.g. :

- Buffing, Kang, Lee, Liu (2018) : <https://inspirehep.net/literature/1709823>
- Aidala (2019) : <https://inspirehep.net/literature/1772224>
- LHCb collaboration (2021) : <https://inspirehep.net/literature/1901628>



Small ..?

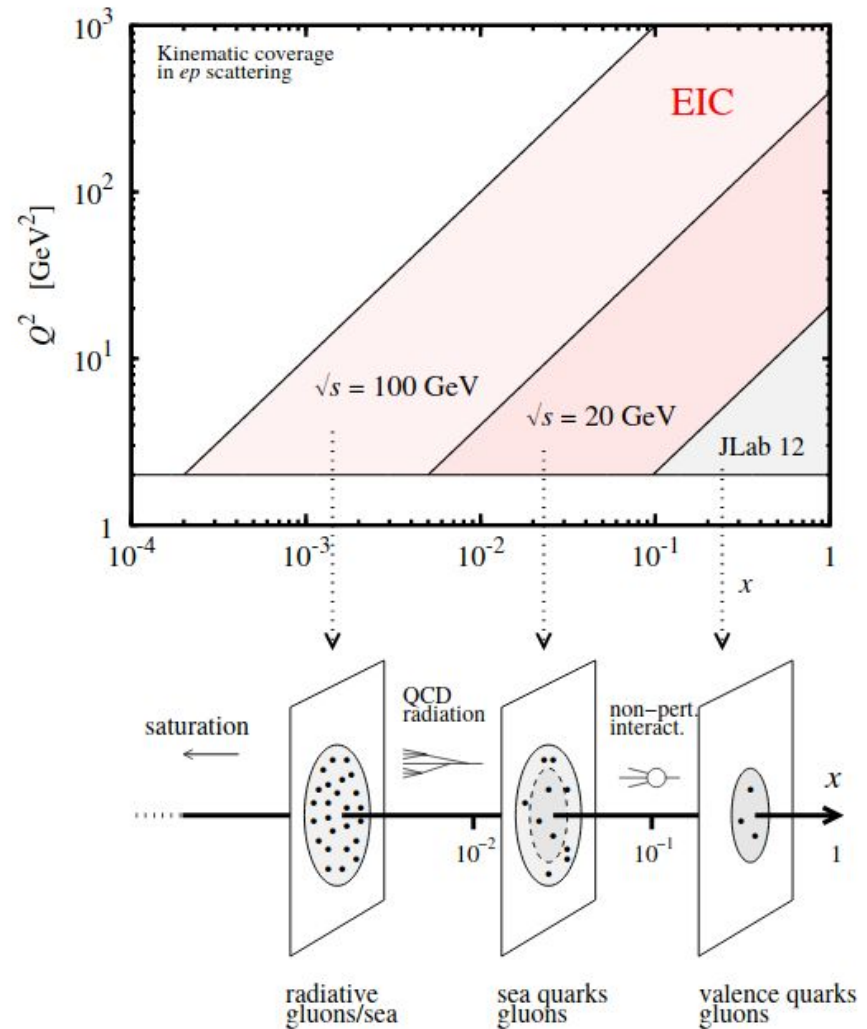
SIDIS coverage

Importance of
complementary experiments

from JLab 12 GeV, Hermes, Compass
to the EIC

zooming into hadron structure

Credit picture: C. Weiss



Approximations

See <https://inspirehep.net/literature/1732230>

Semi-Inclusive Deep-Inelastic Scattering: Nachtmann variable(s)

$$x_{\text{Bj}} = \frac{Q^2}{2P \cdot q}, \quad x_{\text{N}} = -\frac{q^+}{P^+} = \frac{2x_{\text{Bj}}}{1 + \sqrt{1 + \frac{4x_{\text{Bj}}^2 M^2}{Q^2}}},$$

Nachtmann and Bjorken- x are the same up to **target mass corrections**

$$z_{\text{N}} = \frac{P_{\text{B}}^-}{q^-}$$

$$z_{\text{N}} = \frac{Q^4 x_{\text{N}} z_{\text{h}} \left(1 \pm \sqrt{1 - \frac{4M^2 M_{\text{B}}^2 x_{\text{Bj}}^2 (Q^4 + x_{\text{N}}^2 M^2 \mathbf{q}_{\text{T}}^2)}{Q^8 z_{\text{h}}^2}} \right)}{2x_{\text{Bj}} (Q^4 + x_{\text{N}}^2 M^2 \mathbf{q}_{\text{T}}^2)} \stackrel{\text{Fixed } x_{\text{N}}, z_{\text{h}}, \mathbf{q}_{\text{T}}}{=} z_{\text{h}} \left(1 + O\left(\frac{m^4}{Q^4}\right) \right)$$

$\mathcal{O}(q_{\text{T}}^2/Q^2)$

And there are also q_{T} -dependent corrections!

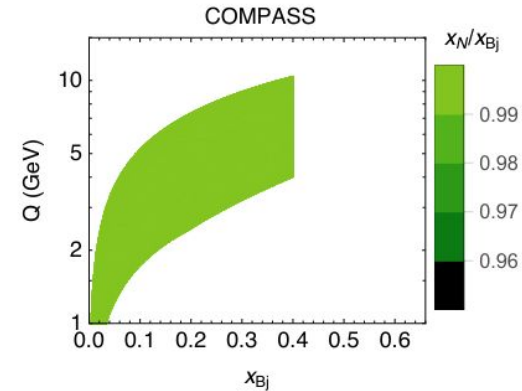
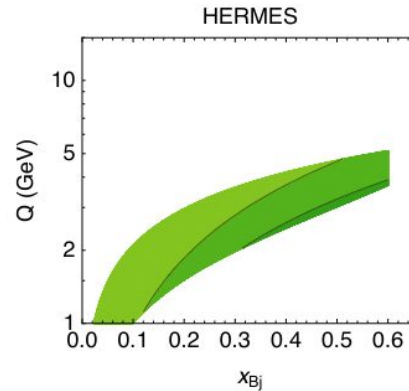
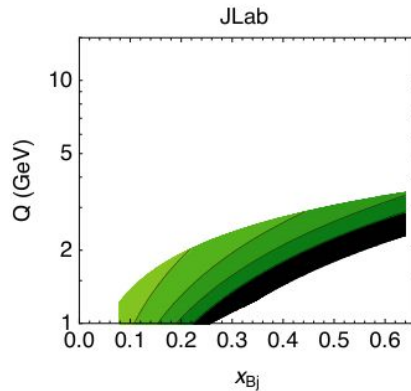
Approximations

See <https://inspirehep.net/literature/1732230>

Semi-Inclusive Deep-Inelastic Scattering: Nachtmann variable(s)

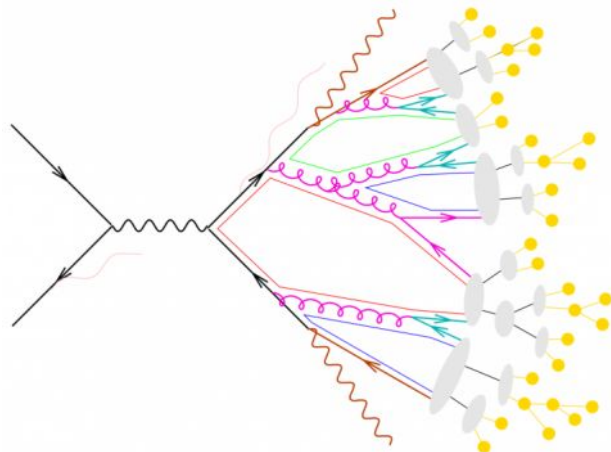
$$x_{\text{Bj}} = \frac{Q^2}{2P \cdot q}, \quad x_{\text{N}} = -\frac{q^+}{P^+} = \frac{2x_{\text{Bj}}}{1 + \sqrt{1 + \frac{4x_{\text{Bj}}^2 M^2}{Q^2}}},$$

Nachtmann and Bjorken- x are the same up to **target mass corrections**



Hadronization and fragmentation functions (FFs)

“Maps” of hadron formation in momentum space



$D_1^h(z)$ single-hadron collinear FF

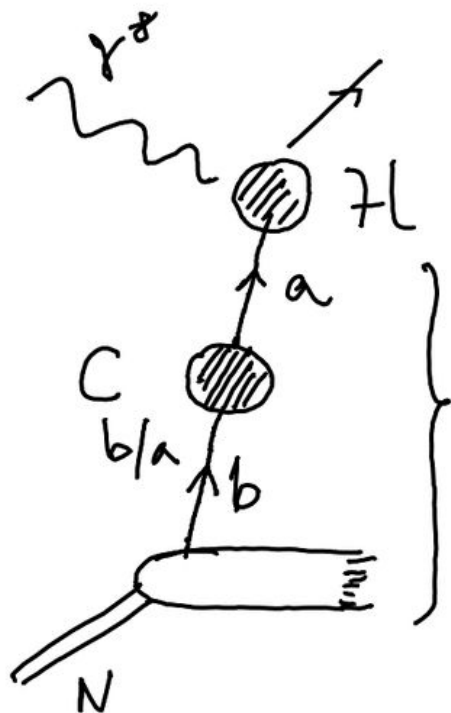
$D_1^h(z, P_T^2)$ single-hadron TMD FF

$D_1^{h_1 h_2}(z, \zeta)$ di-hadron FF

$J(s)$ inclusive jet FF

$\mathcal{G}^h(s, z)$ in-jet FF

Intrinsic transverse momentum



Provides fixed-order (perturbative)
transverse-momentum structure at the
matching scale
(**DGLAP** splitting and **logs** of b_T)

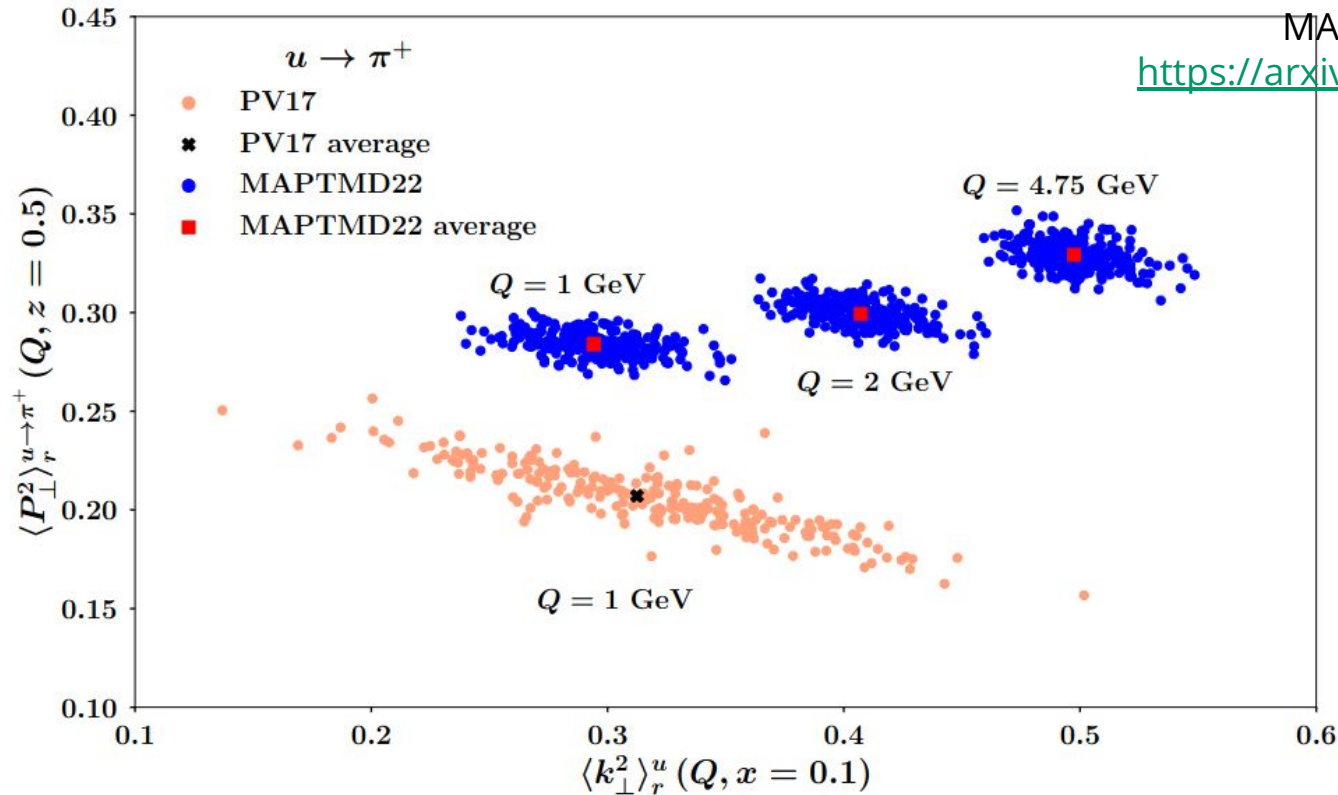
$$f_1(x, b_T)^{a/N} = \sum_b c_{b/a} \otimes f_1(x)^{b/N} + \mathcal{O}(b_T \Lambda)^{a/N}$$

$F_{NP}(b_T; \lambda)$

**For quark flavor a ,
not b**

Intrinsic transverse momentum

Kinematic dependence of
average intrinsic
transverse momenta



Intrinsic transverse momentum

Kinematic dependence of
average intrinsic
transverse momenta

MAP coll. 2024 -

<https://arxiv.org/pdf/2405.13833>

