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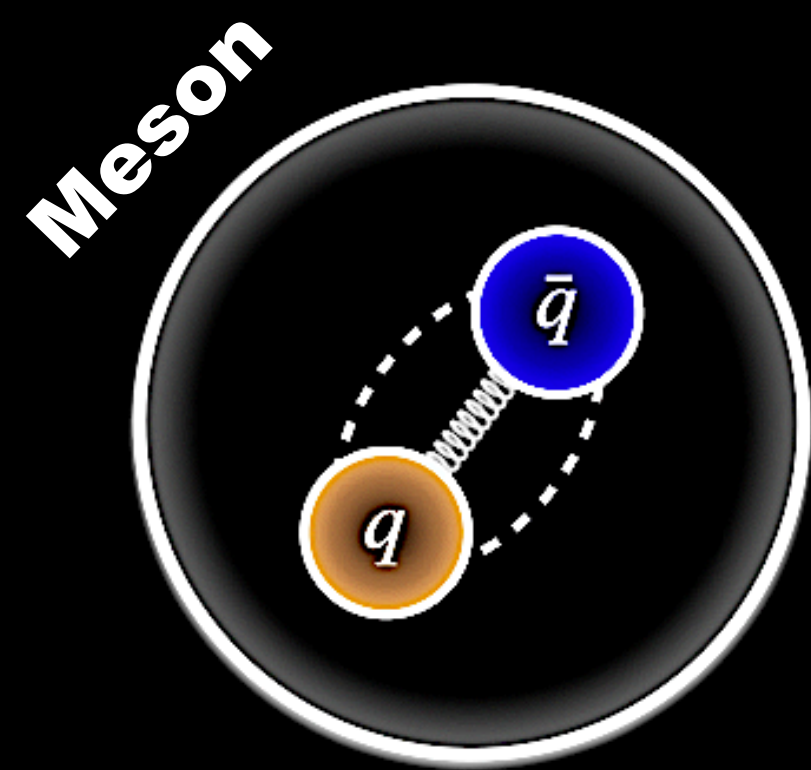
Exploring Photoproduced $\eta^{(\prime)} \pi^0$ Systems in the Search for Exotic Hadrons at GlueX

2026 Exotic Hadron Spectroscopy Workshop

The University of Cambridge

Zachary Baldwin | April 1, 2026

on behalf of the  collaboration



Total angular momentum | $J = 0, 1, 2, \dots$

Parity | $P = (-1)^{L+1}$

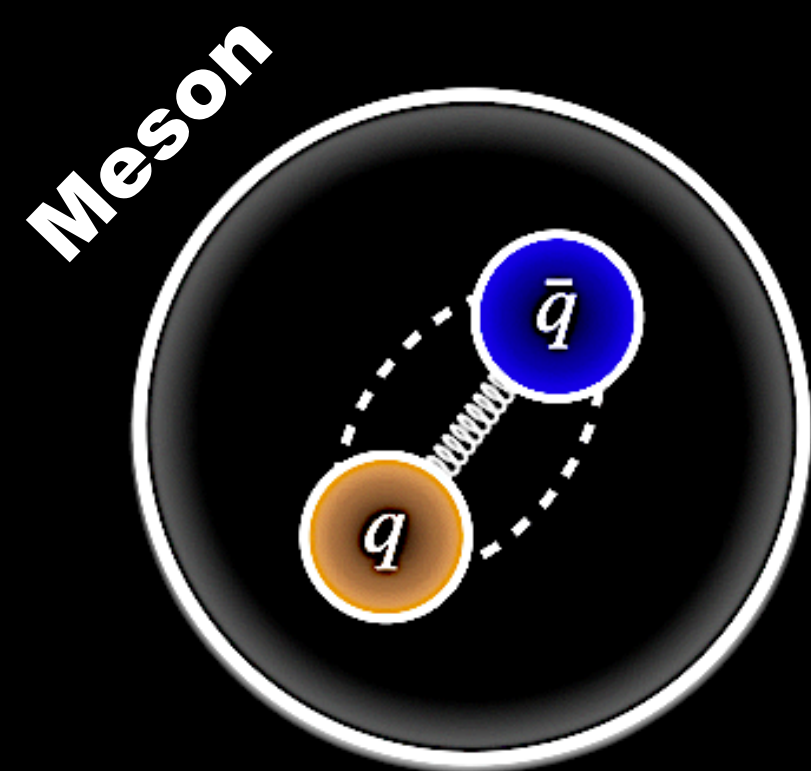
Charge Conjugation | $C = (-1)^{L+S}$

L is the relative orbital angular momentum of the q and \bar{q}

S is the total intrinsic spin of the $q\bar{q}$ pairs

Allowed J^{PC} quantum numbers

L	S	J^{PC}	L	S	J^{PC}	L	S	J^{PC}
0	0	0^{-+}	1	0	1^{+-}	2	0	2^{-+}
0	1	1^{--}	1	1	0^{++}	2	1	1^{--}
			1	1	1^{++}	2	1	2^{--}
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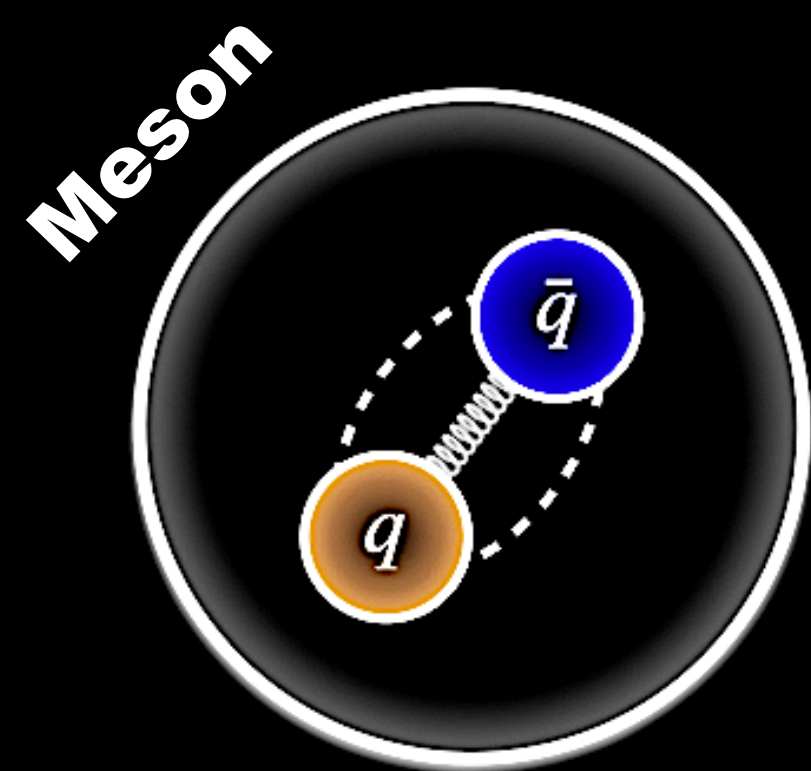
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Observation of any system with quantum numbers forbidden in the constituent quark model, provides direct evidence for a non- $q\bar{q}$ configuration

Forbidden J^{PC} quantum numbers

$0^{--}, 0^{+-}, 1^{-+}, 2^{+-}$



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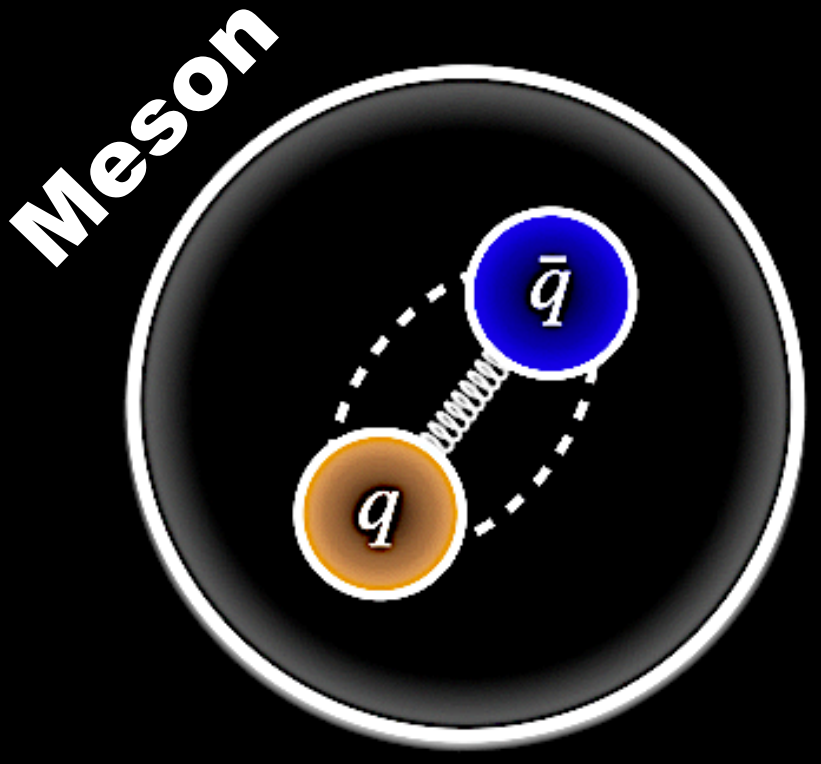
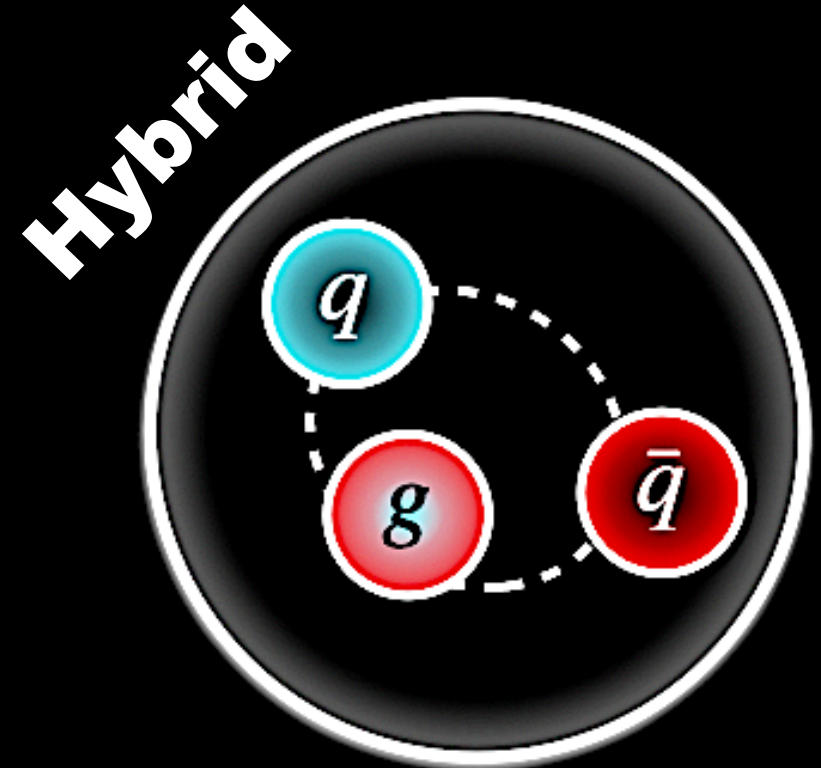
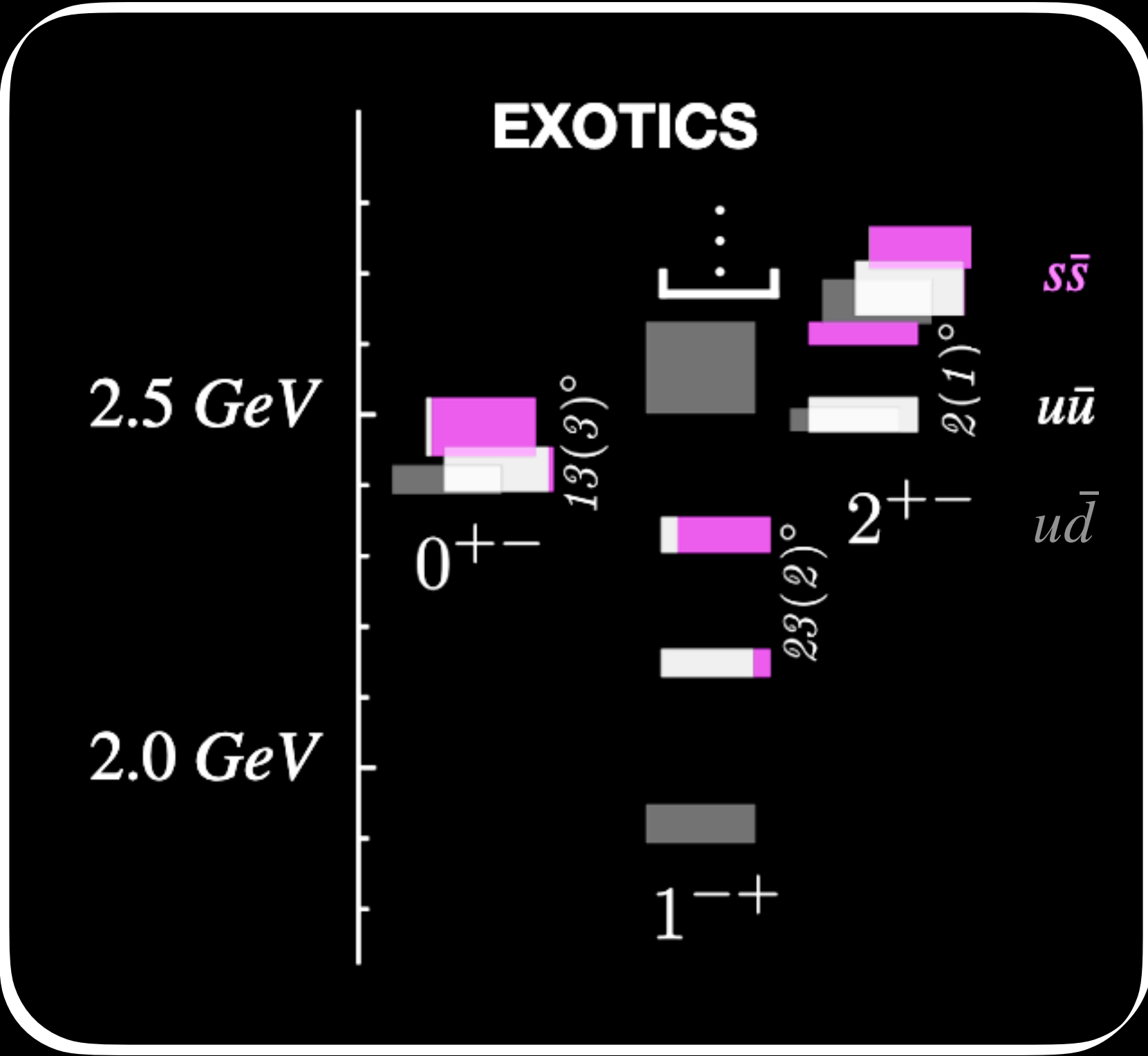
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Do gluons play a larger role within the structure of hadrons?

Are observed *exotics* genuinely non- $q\bar{q}$ states or kinematic effects?

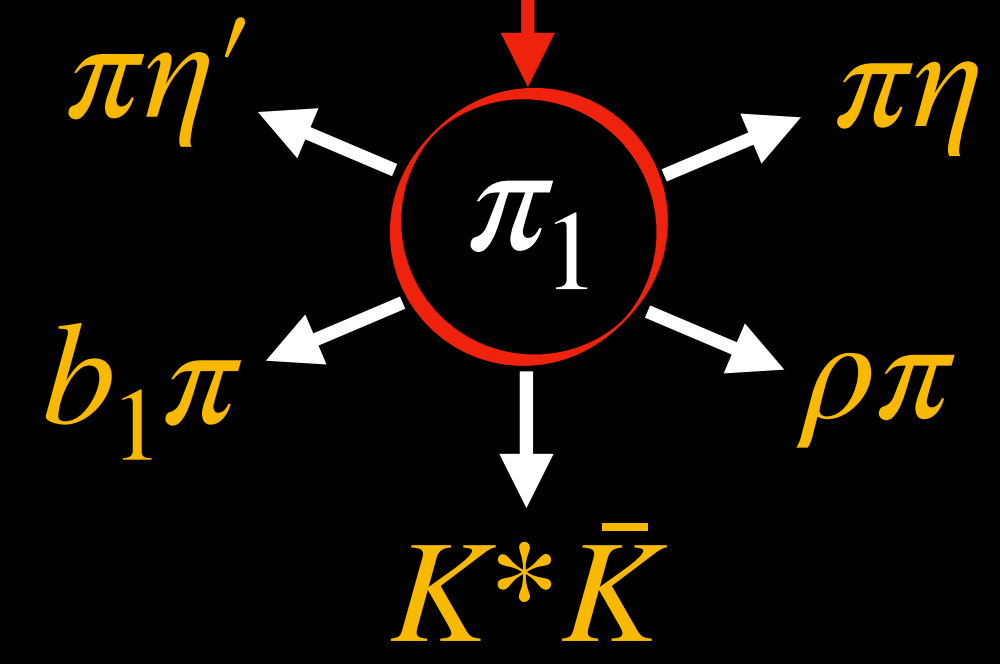
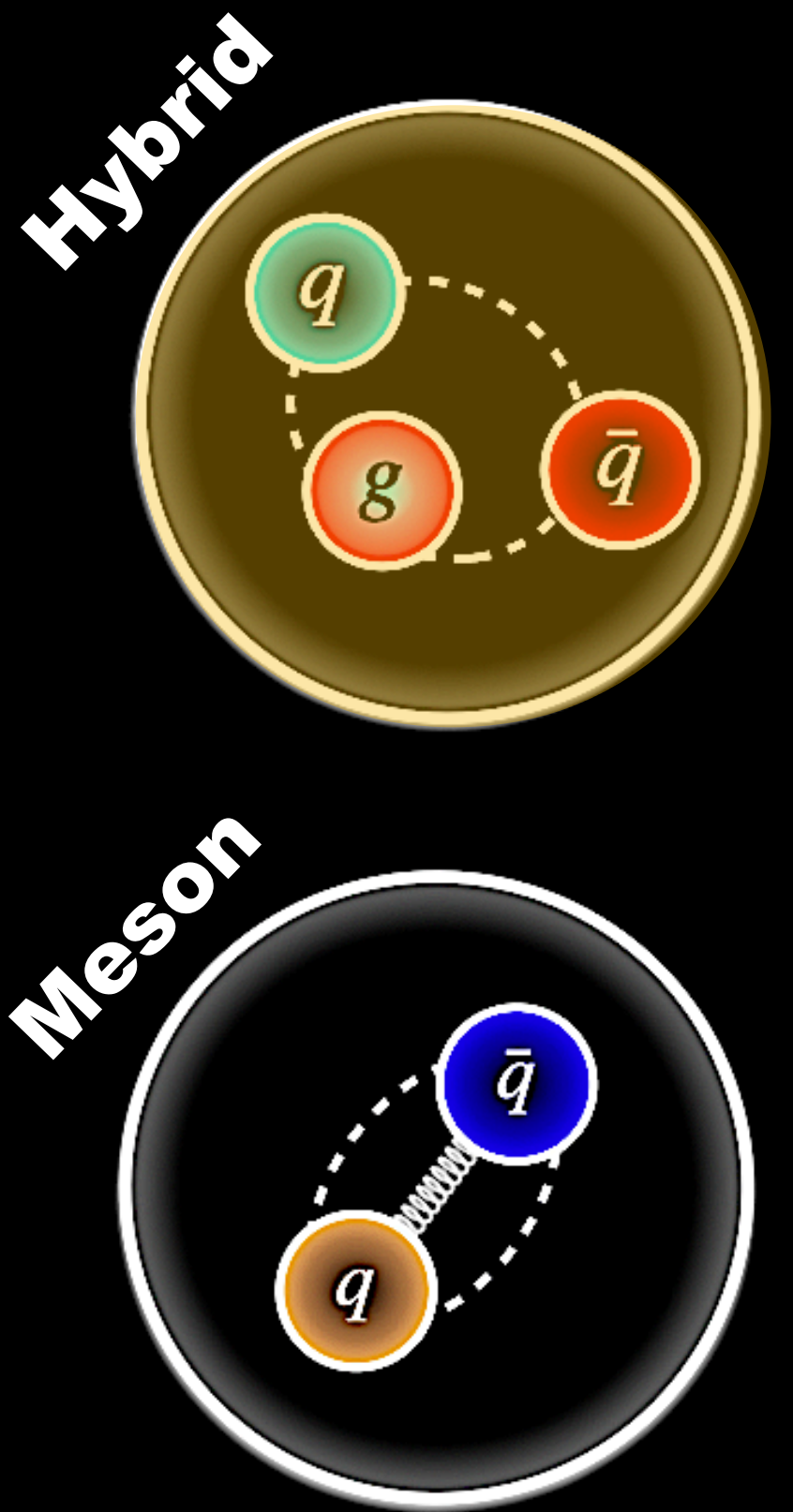
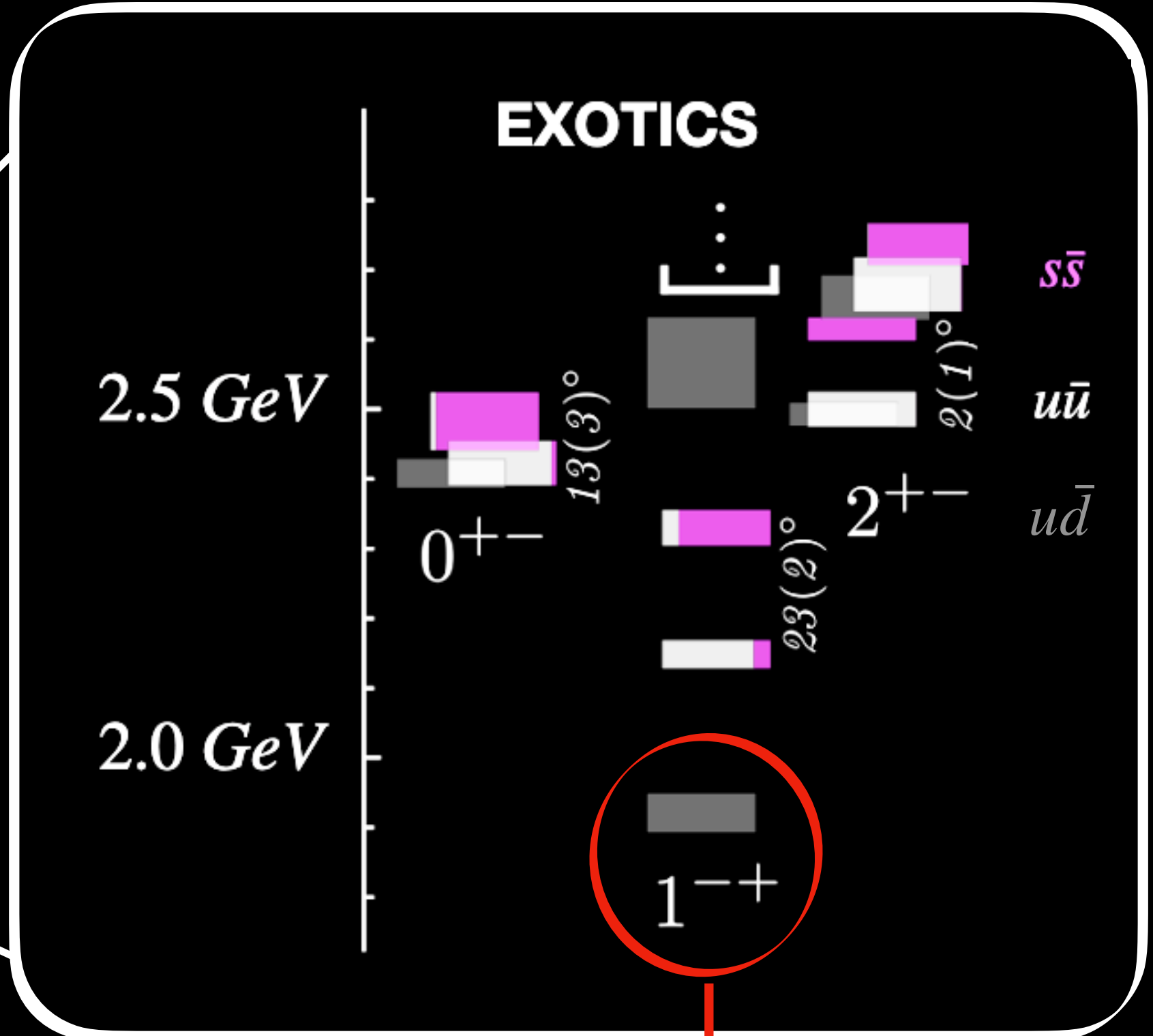
Lattice QCD predicts the existence of mesons with explicit gluonic excitations commonly referred to as spin-exotic hybrid mesons

J. Dudek et al. [Hadron Spectrum Collab], Phys. Rev. D 83, 111502 (2011)



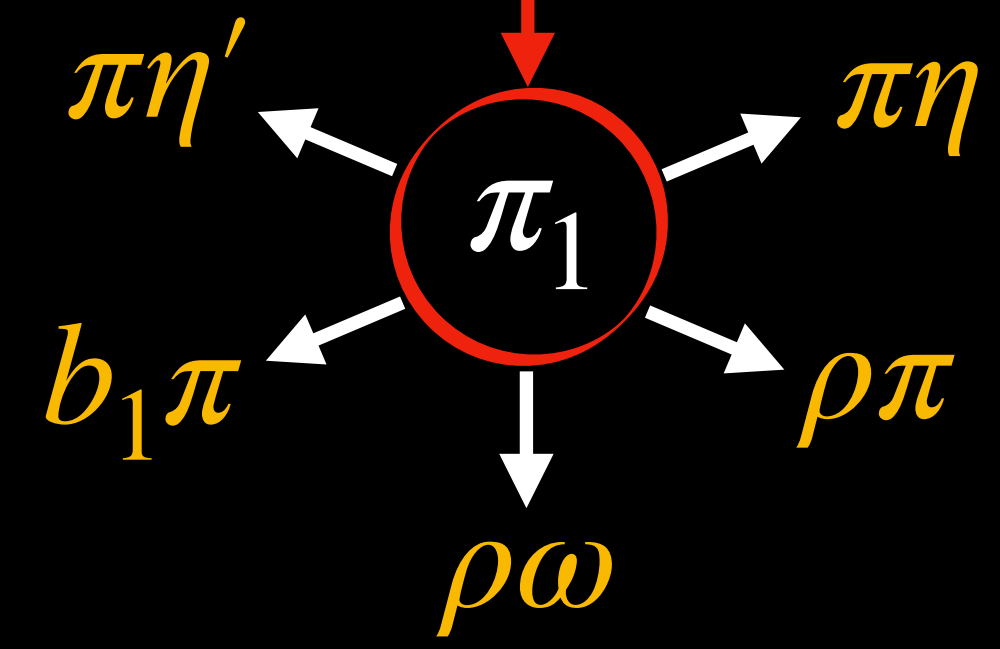
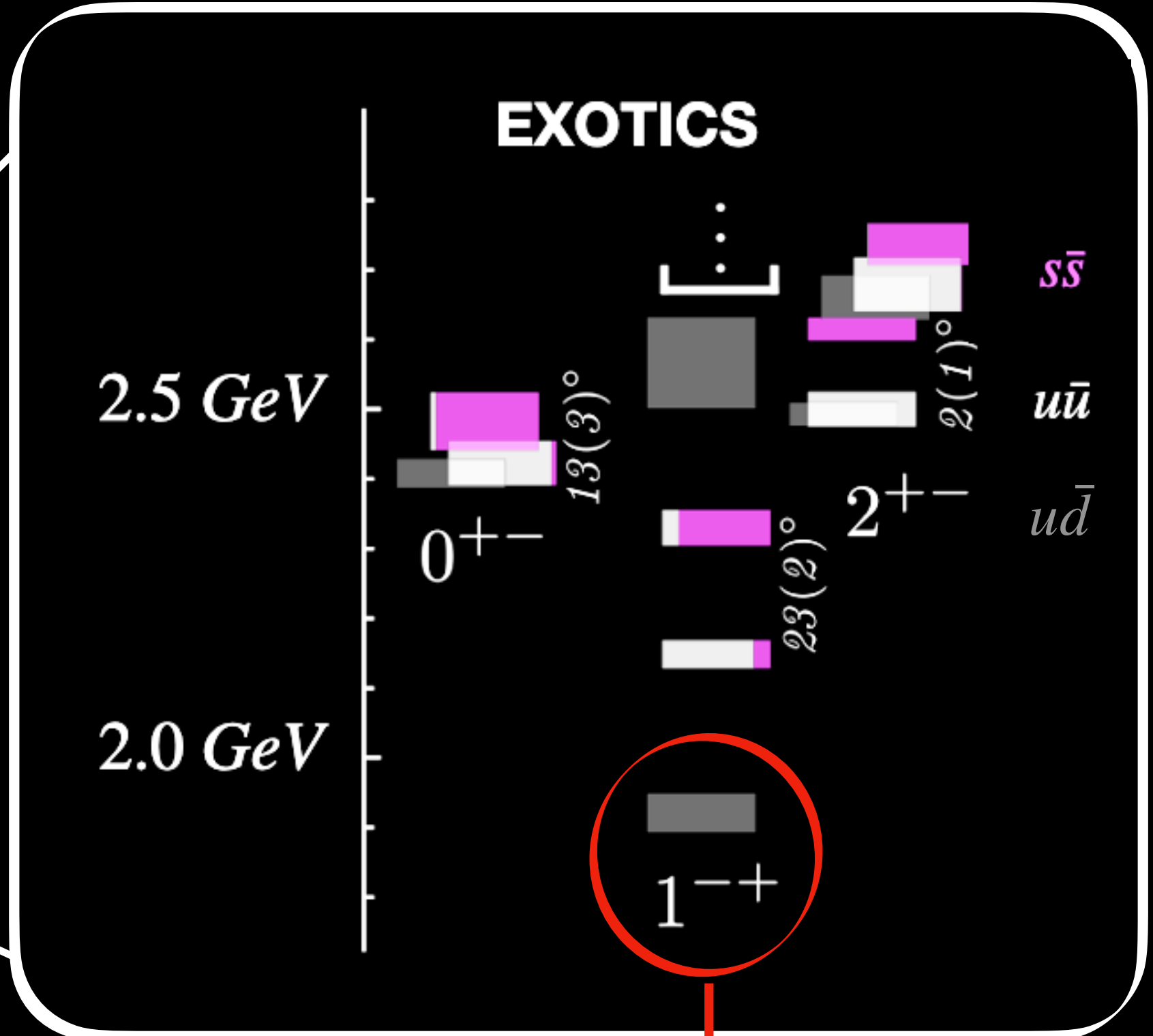
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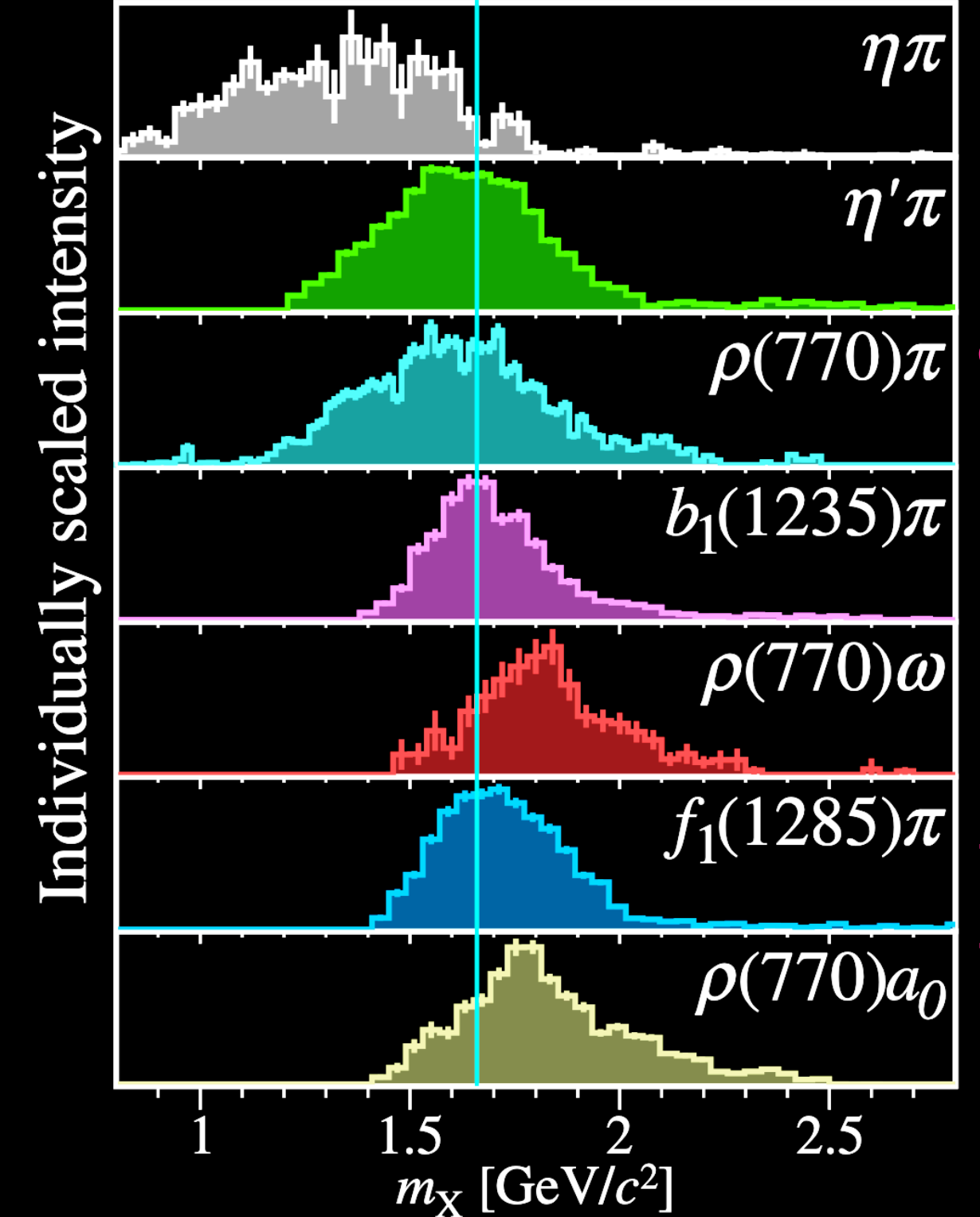


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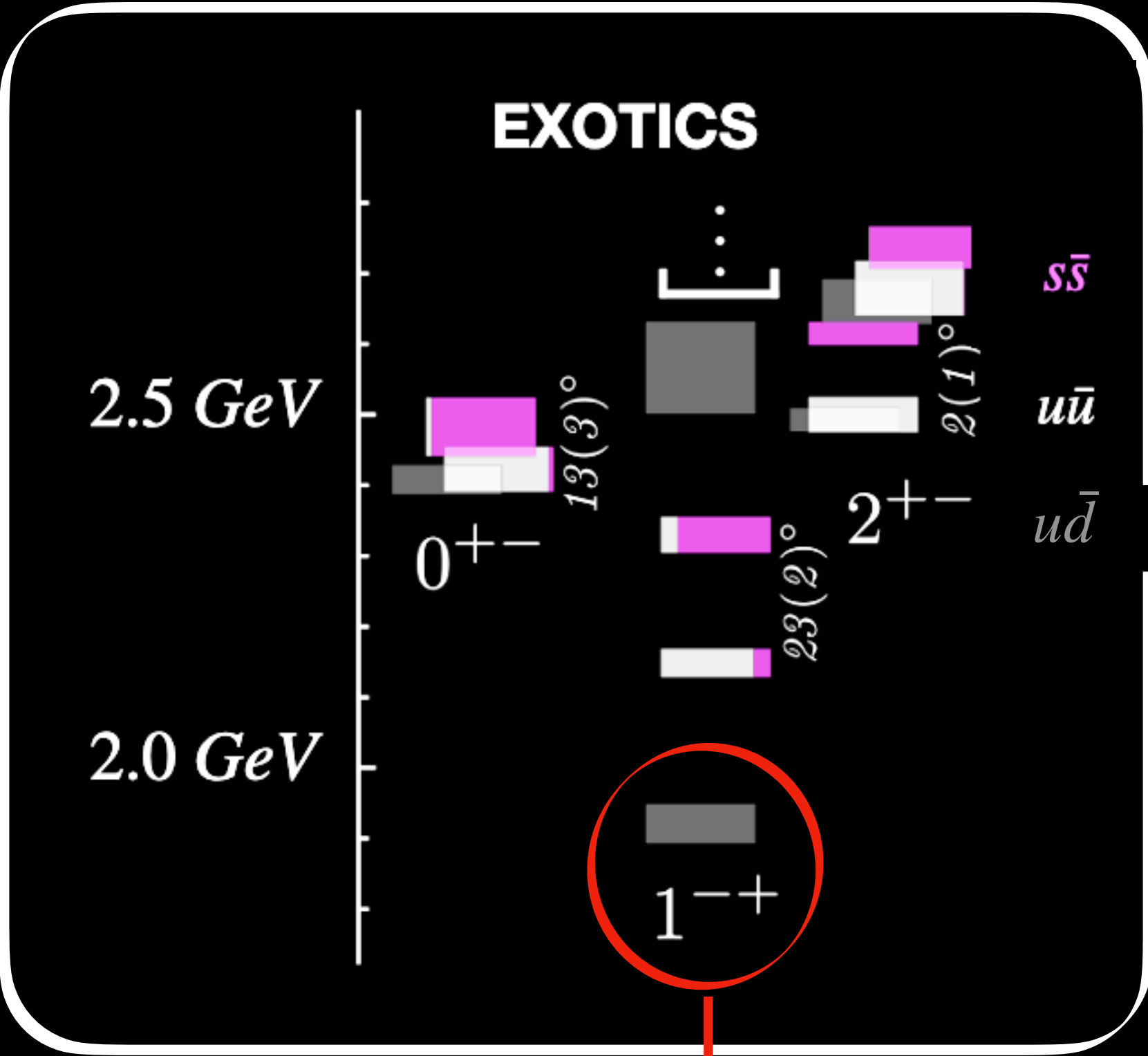
Spin-exotic $J^{PC}=1^{-+}$ waves at COMPASS preliminary
Nominal $\pi_1(1600)$ position



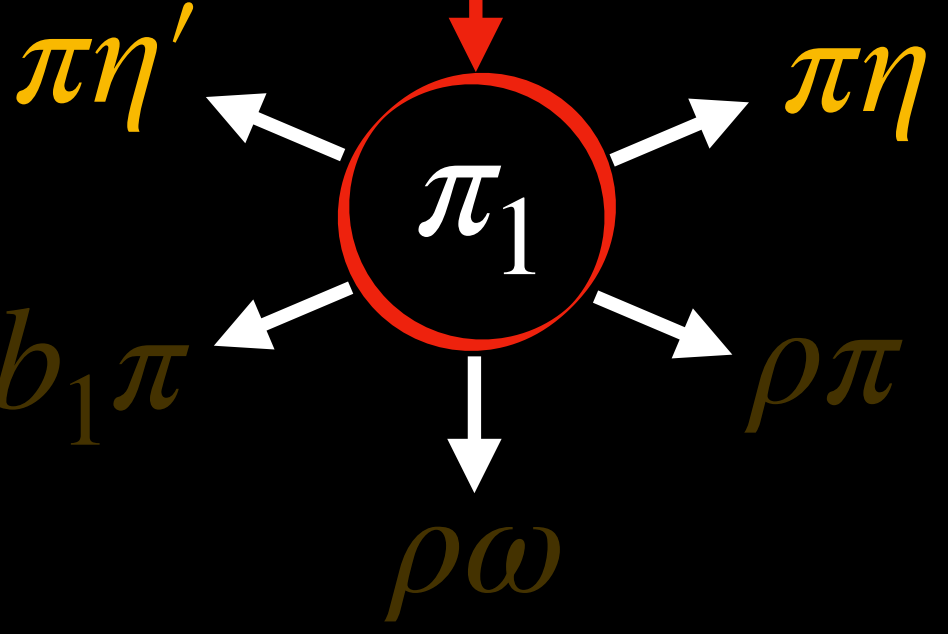
D. Spilbeck [on behalf of the COMPASS Collab], Proceedings of Science, HADRON (2026)

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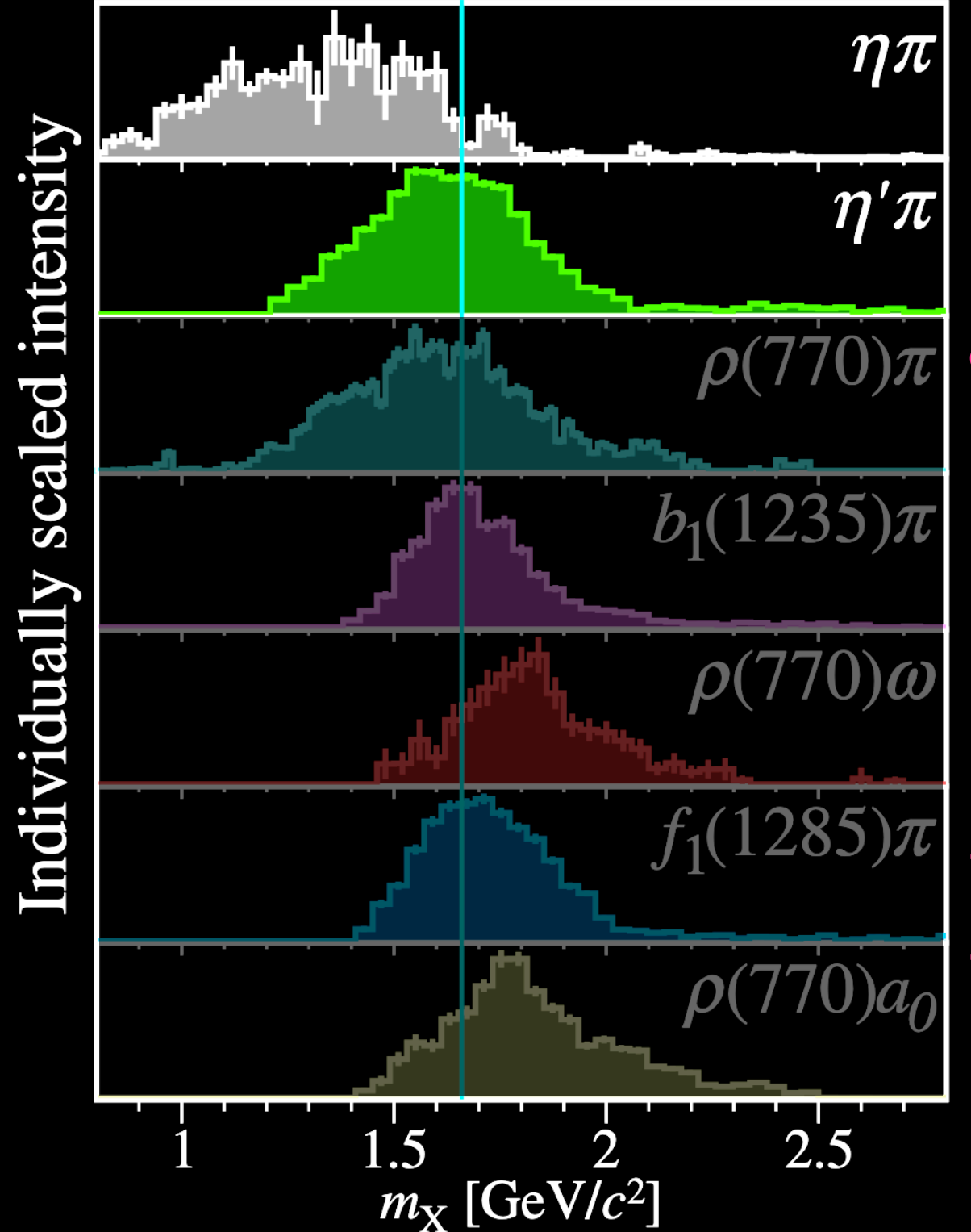
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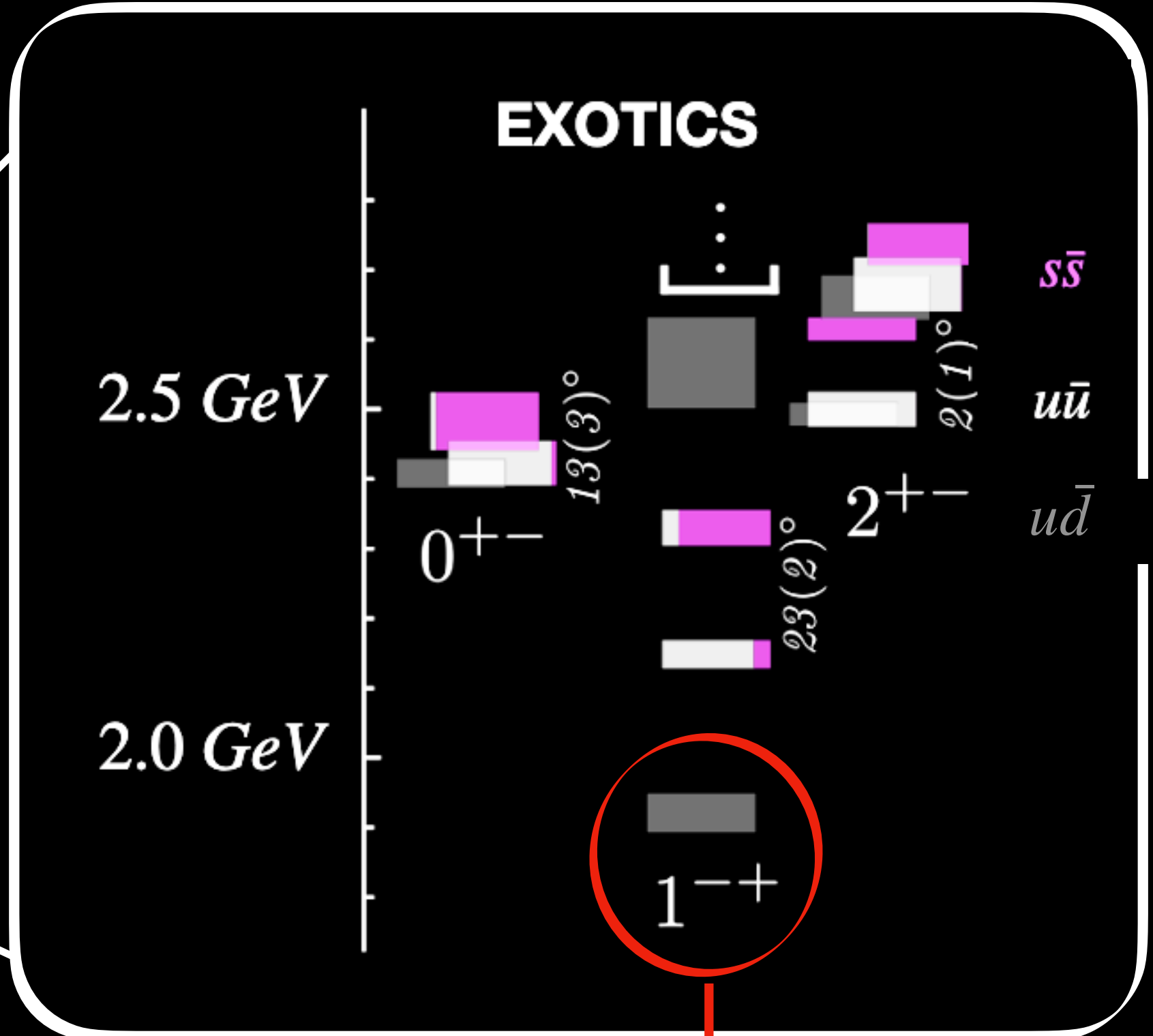
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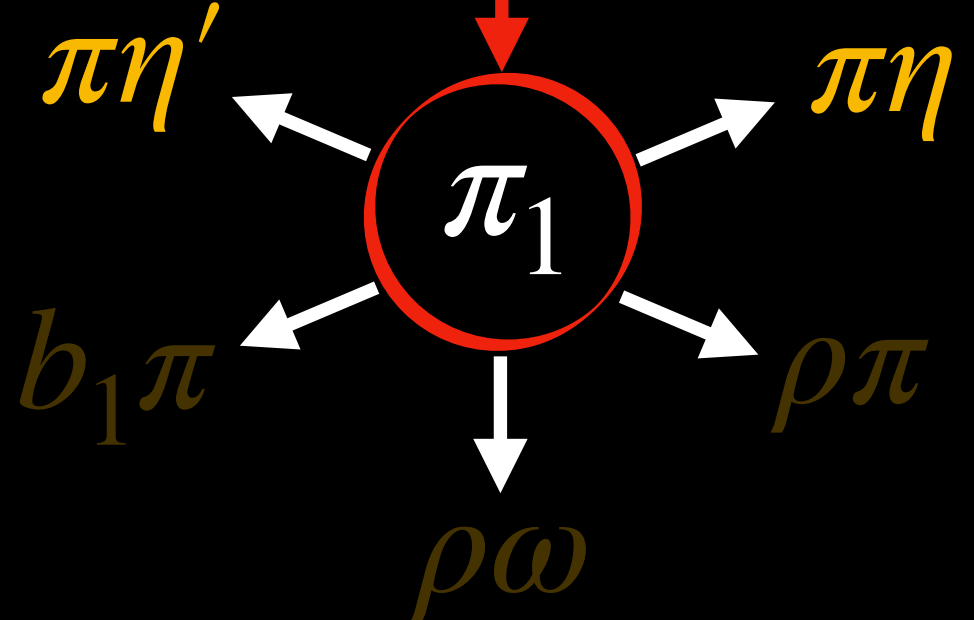
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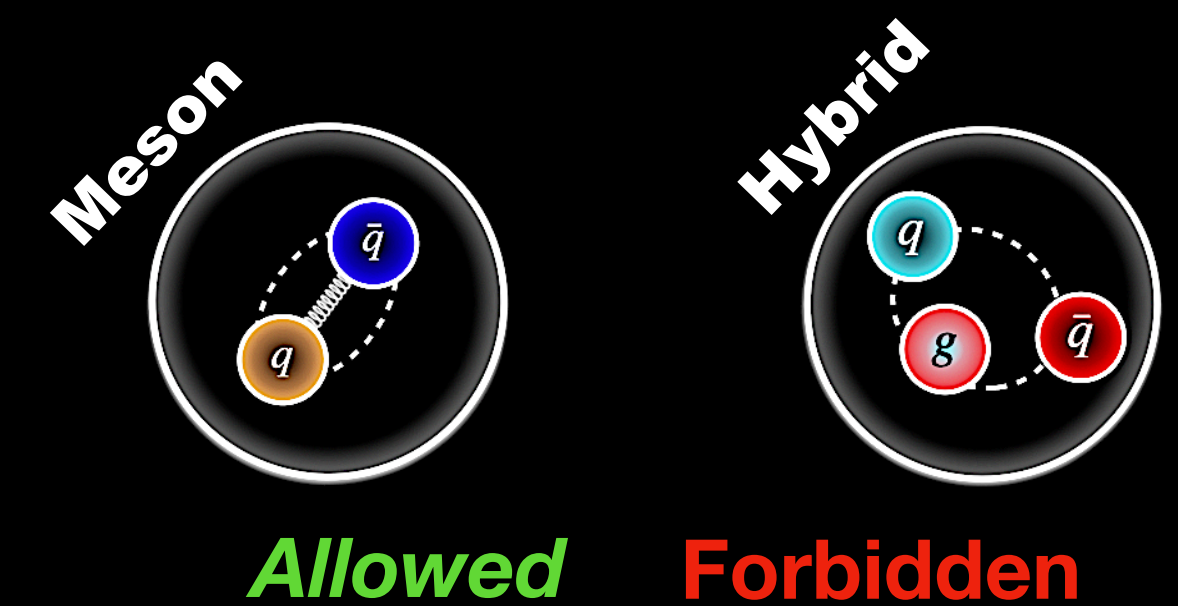


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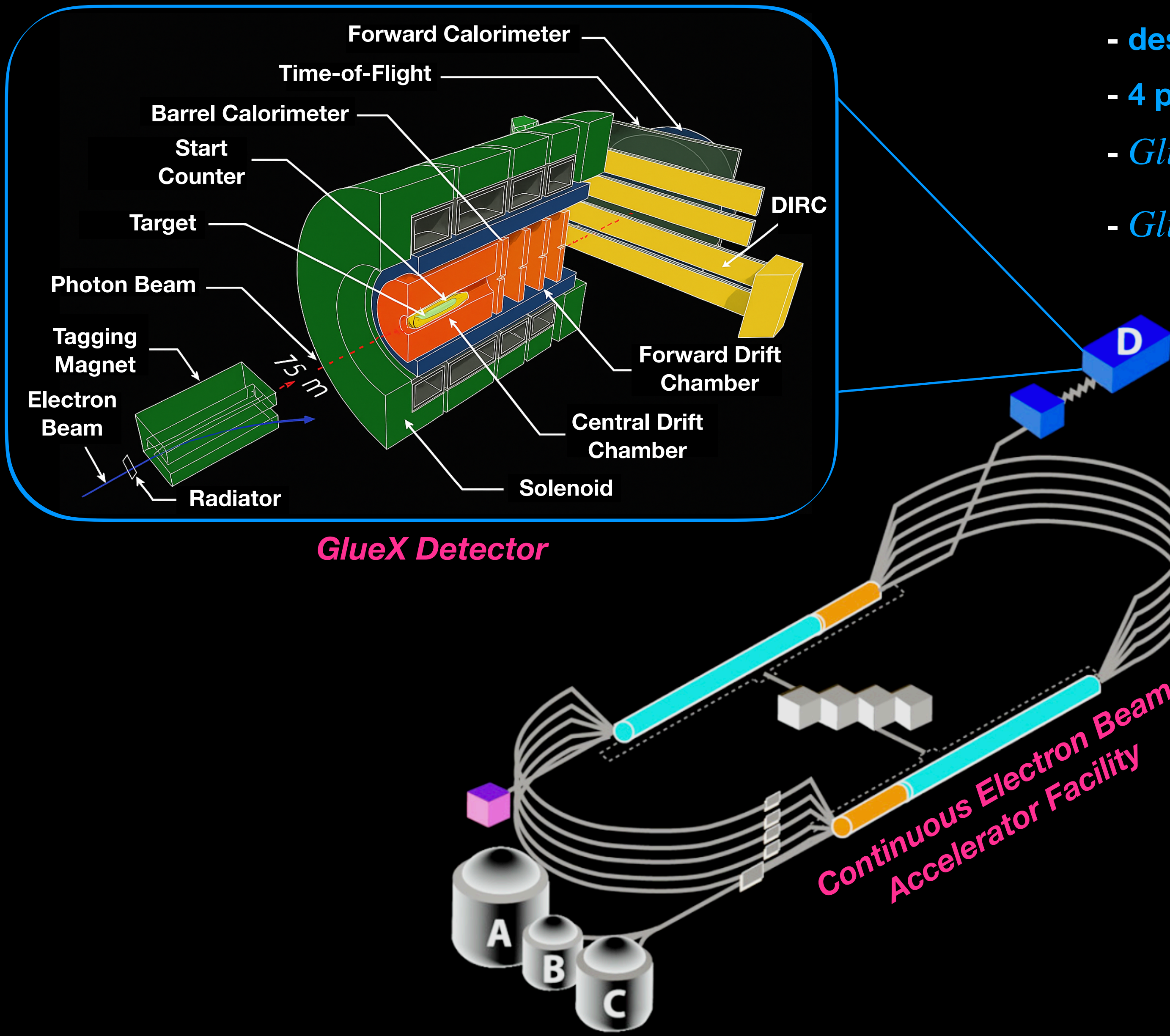


The $\eta\pi$ and $\eta'\pi$ channels are ideal for searches of spin-exotic hybrids

- all odd- L waves in $\eta^{(\prime)}\pi$ provide access to exotic quantum numbers
- simplistic 2-body final states
- historically, consistent observations of exotic resonances signal observed



L	S	P	D	F	...
J^{PC}	0^{++}	1^{-+}	2^{++}	3^{-+}	...



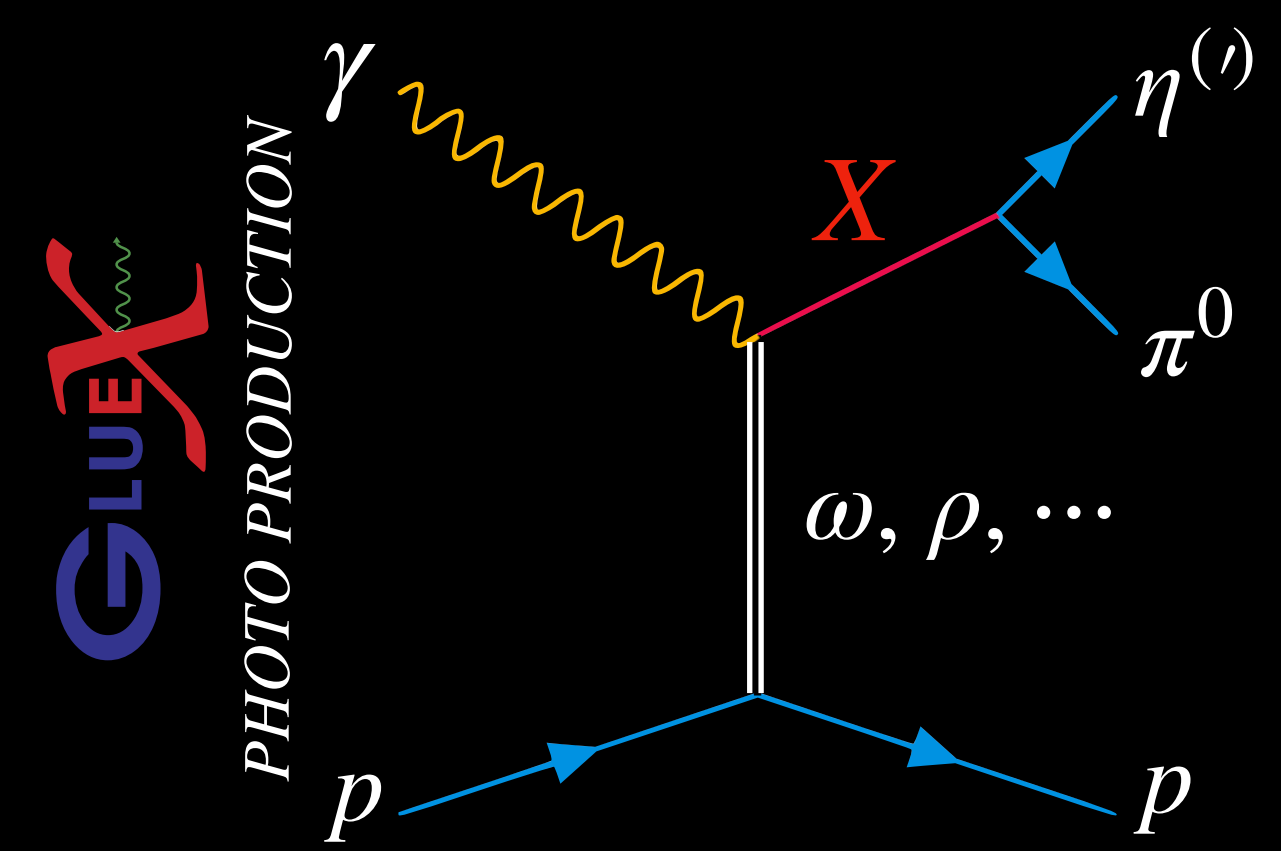
- designed to reconstruct final state particles from $\gamma p \rightarrow pM$
- 4 polarization orientations
- *GlueX-I* collected $\int L = 125 \text{ pb}^{-1}$ in coherent peak
- *GlueX-II* ~ 3-4 times more (currently ongoing)

Photoproduction \Rightarrow extremely versatile

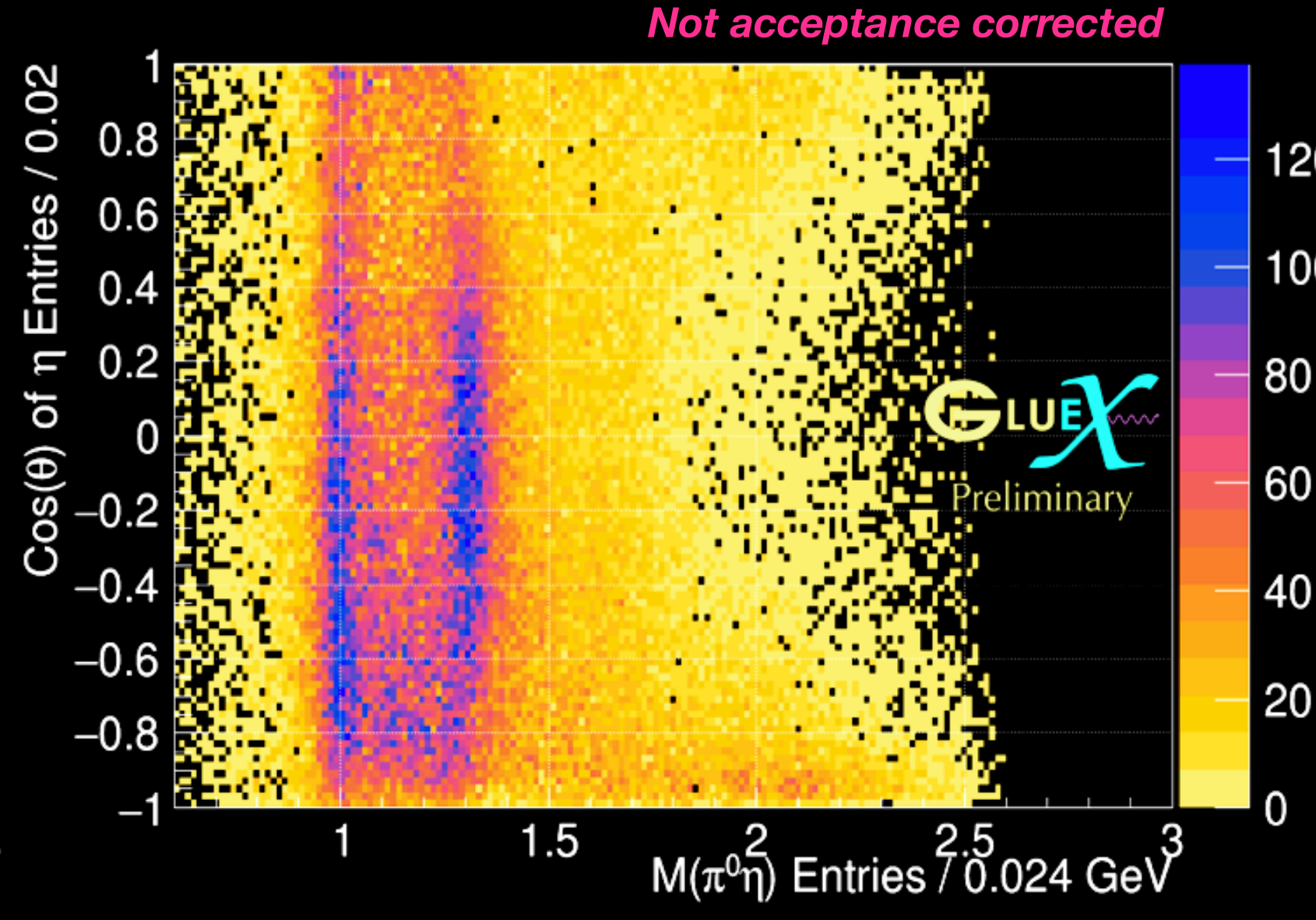
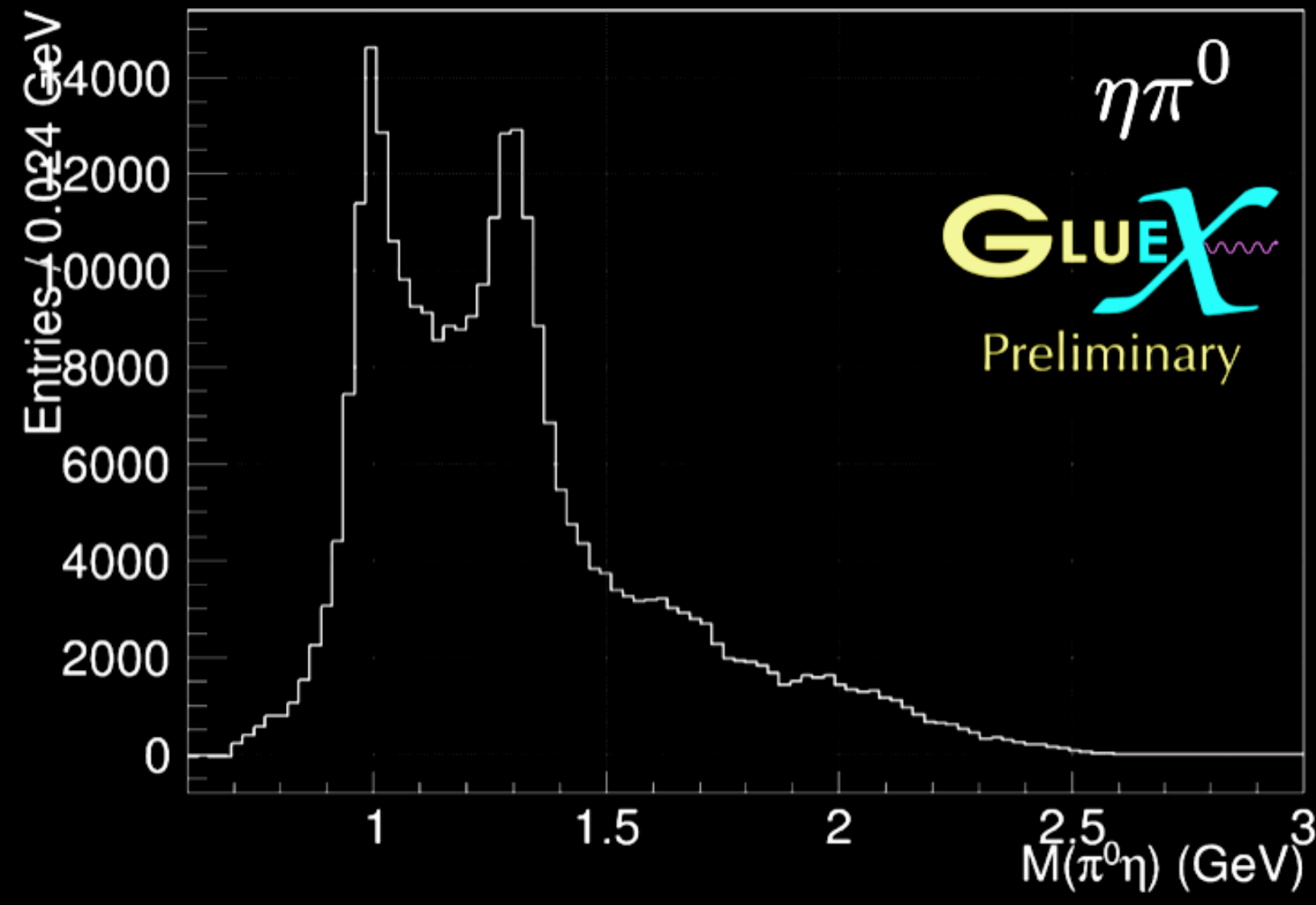
- access to large range of resonances
- complementary to hadroproduction (COMPASS, etc.)

Start simple to understand production of less complex hadrons

- study background, acceptance, non-resonant contributions, etc.

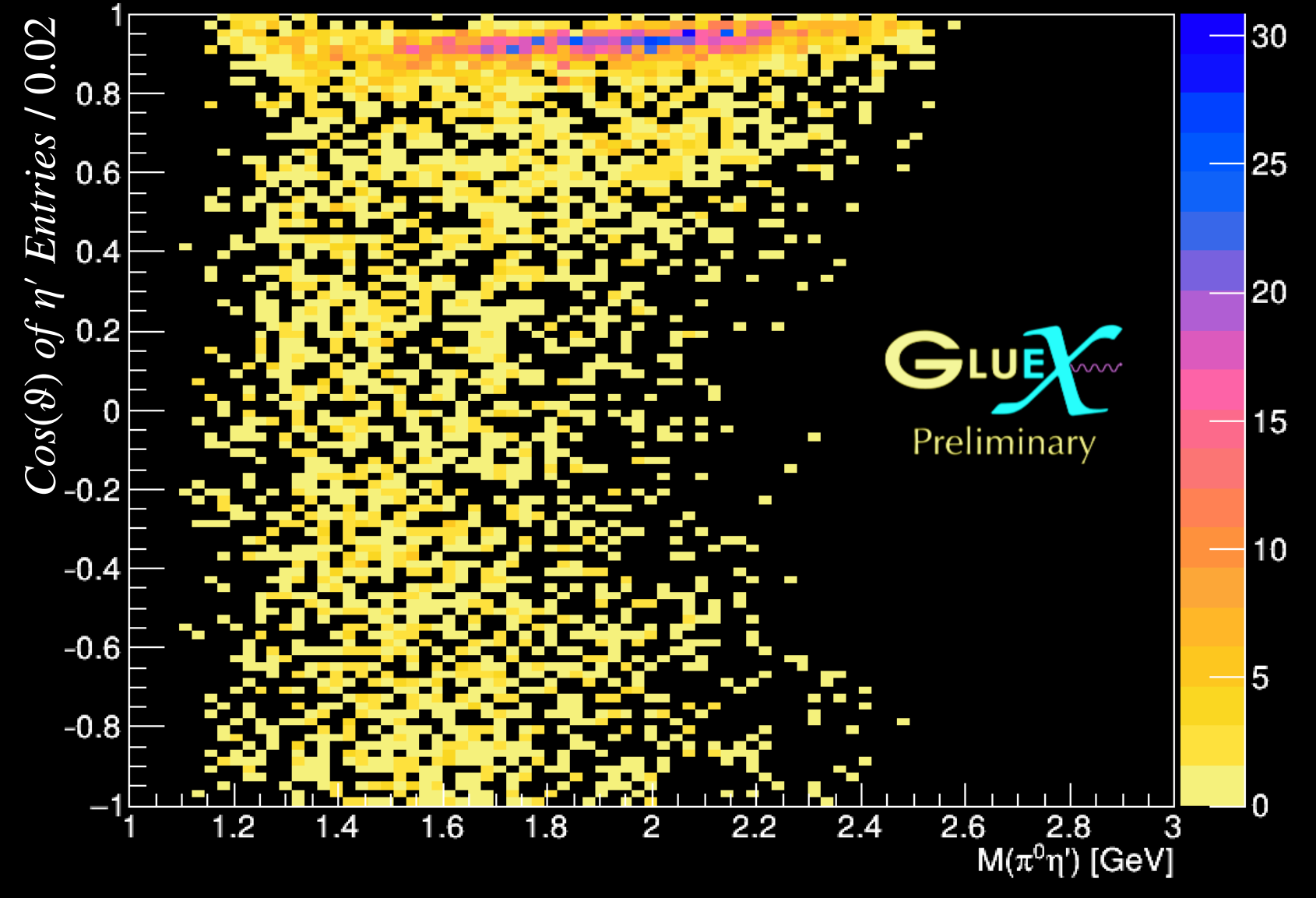
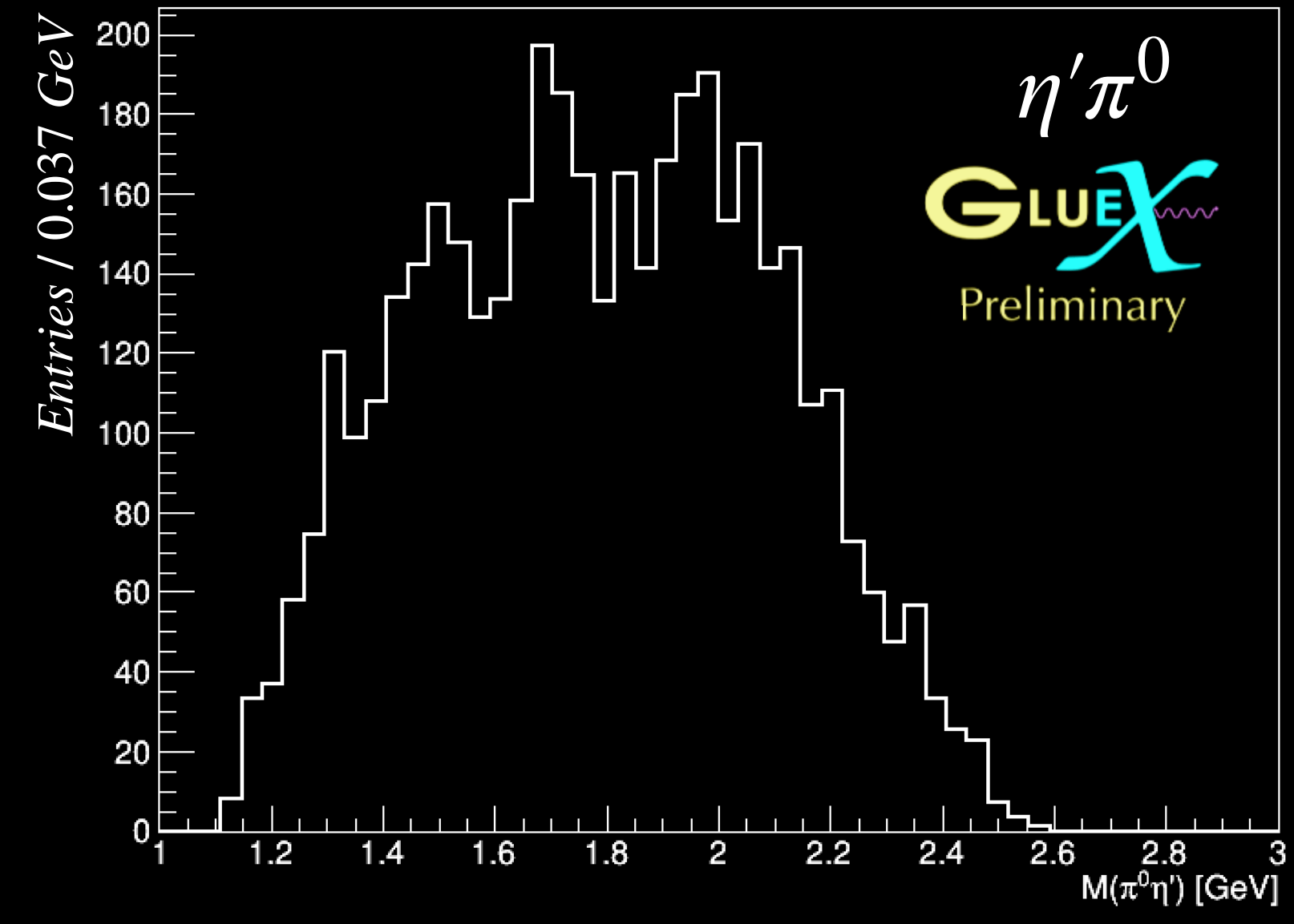
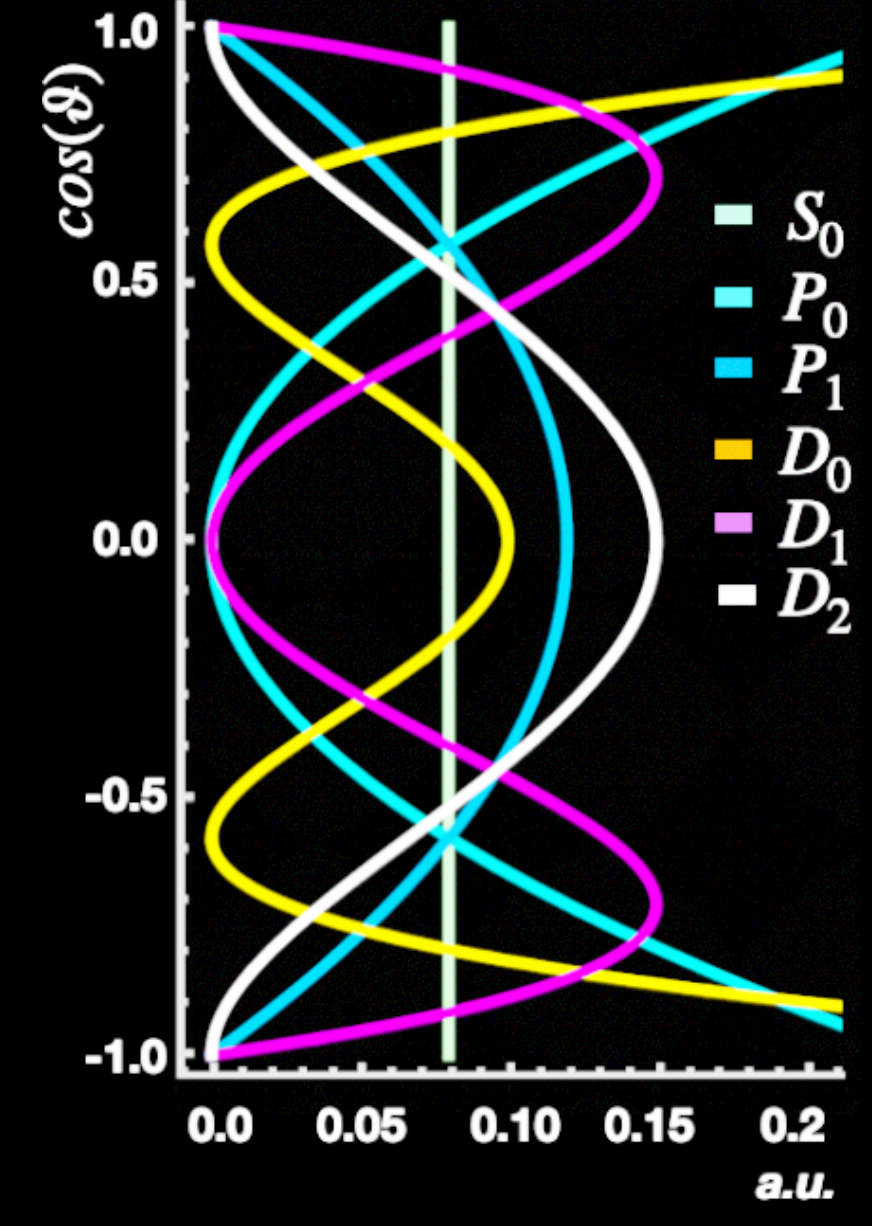


$0.1 < -t [GeV^2] < 0.3$



Neutral decay modes

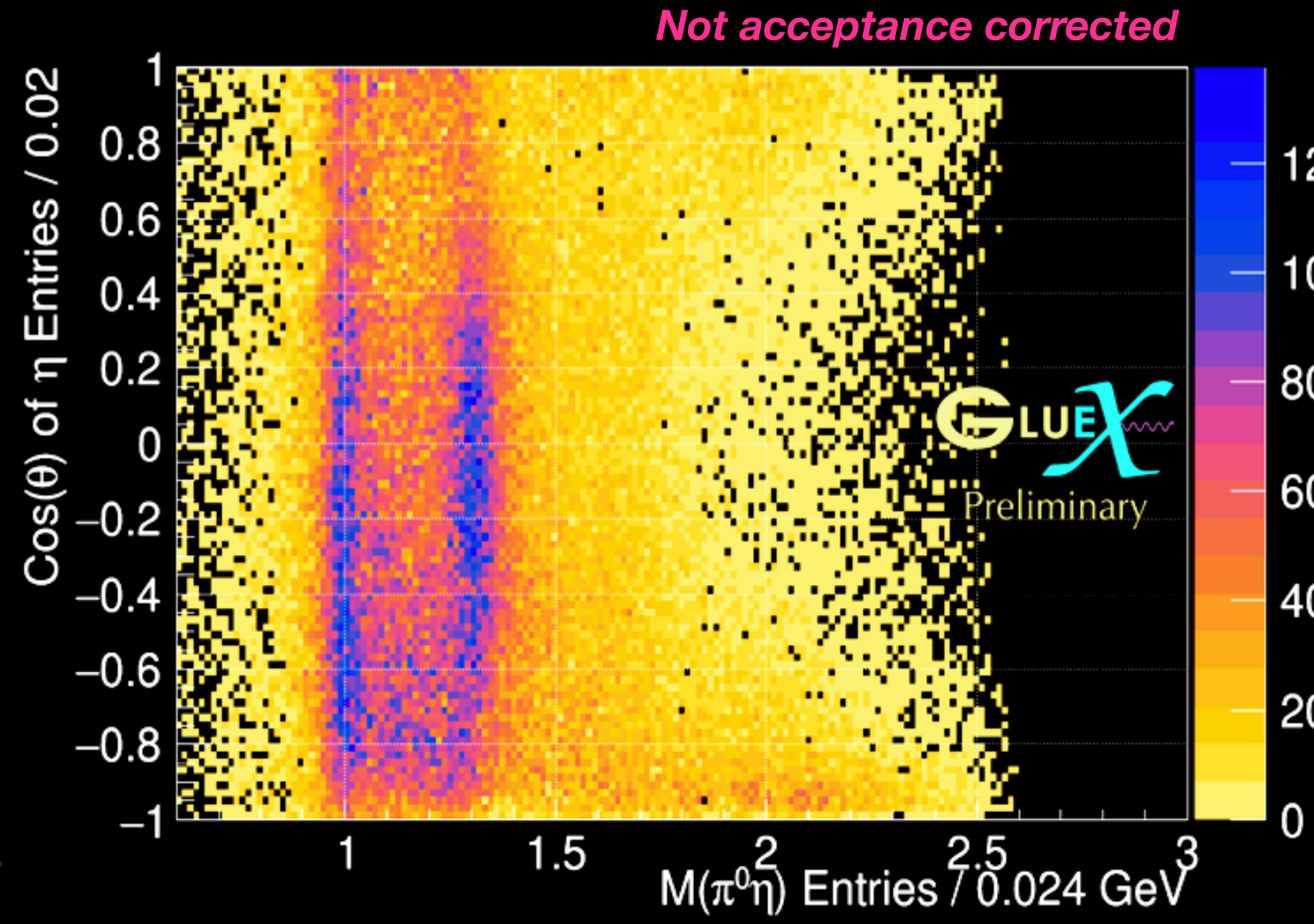
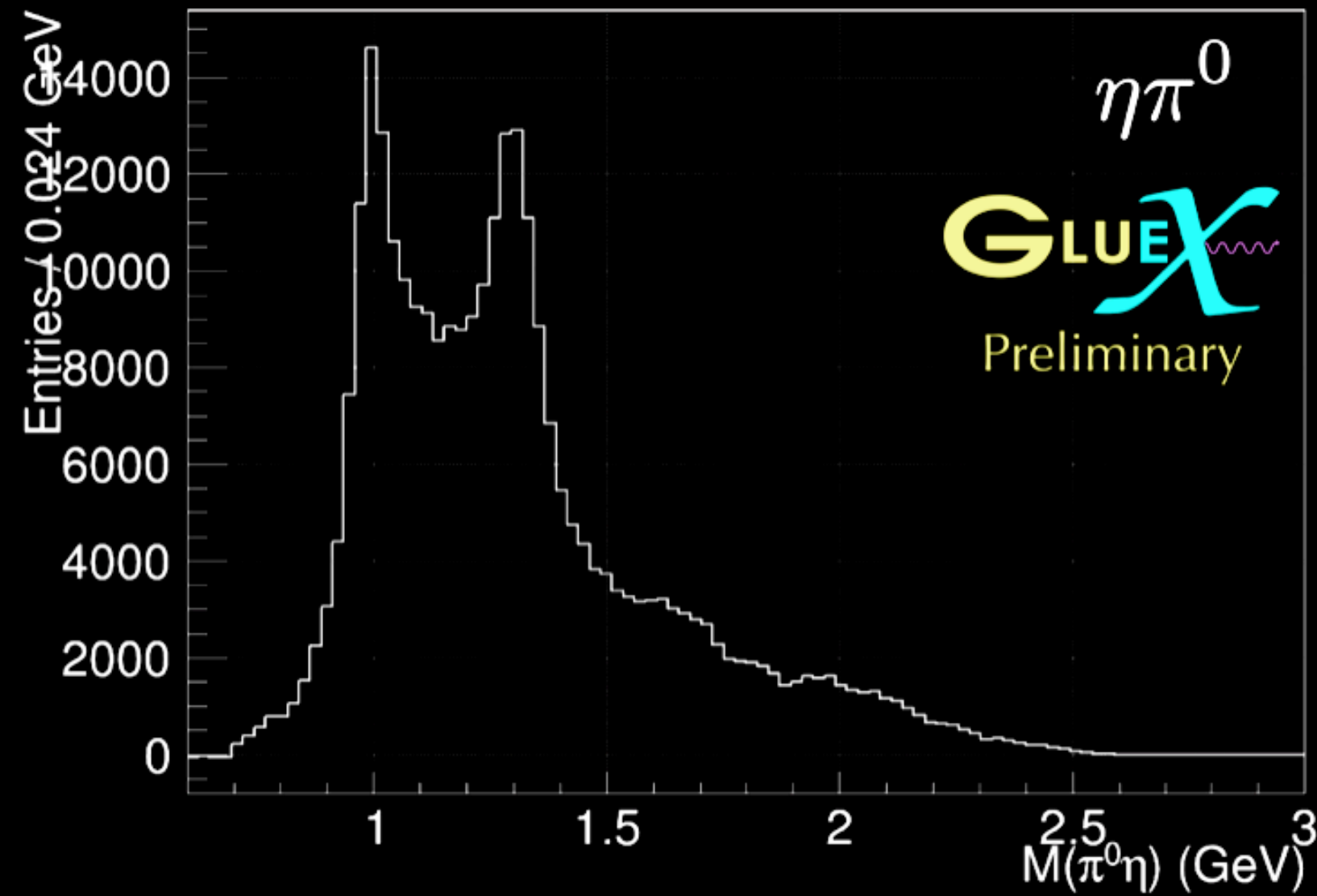
- $\eta\pi^0 \rightarrow 4\gamma$
- $\eta'\pi^0 \rightarrow 4\gamma\pi^+\pi^-$



Charged decay modes also being analyzed ...

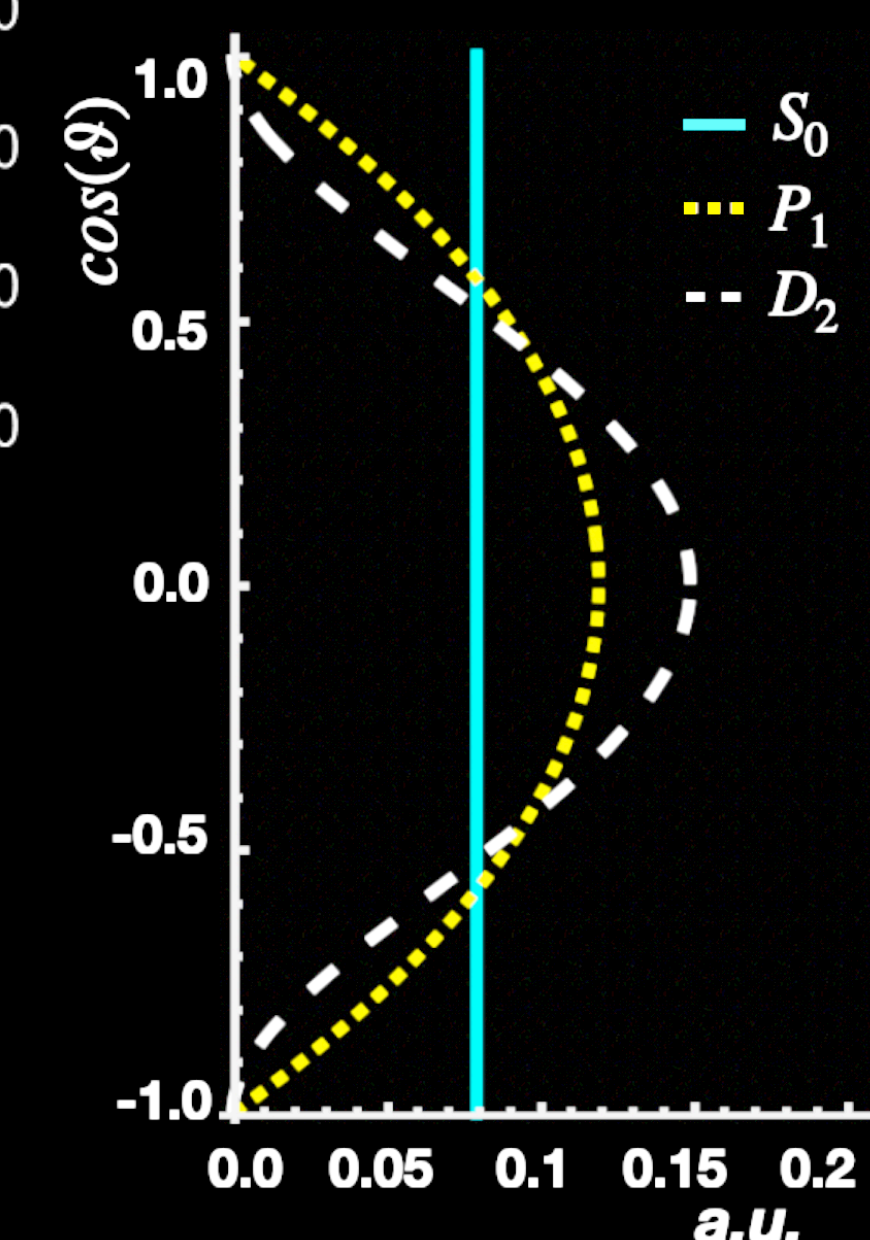
- $\eta^{(\prime)}\pi^-\Delta^{++} \mid \Delta^{++} \rightarrow \pi^+p$
 $\eta' \rightarrow \eta\pi^+\pi^-$

$0.1 < -t [GeV^2] < 0.3$

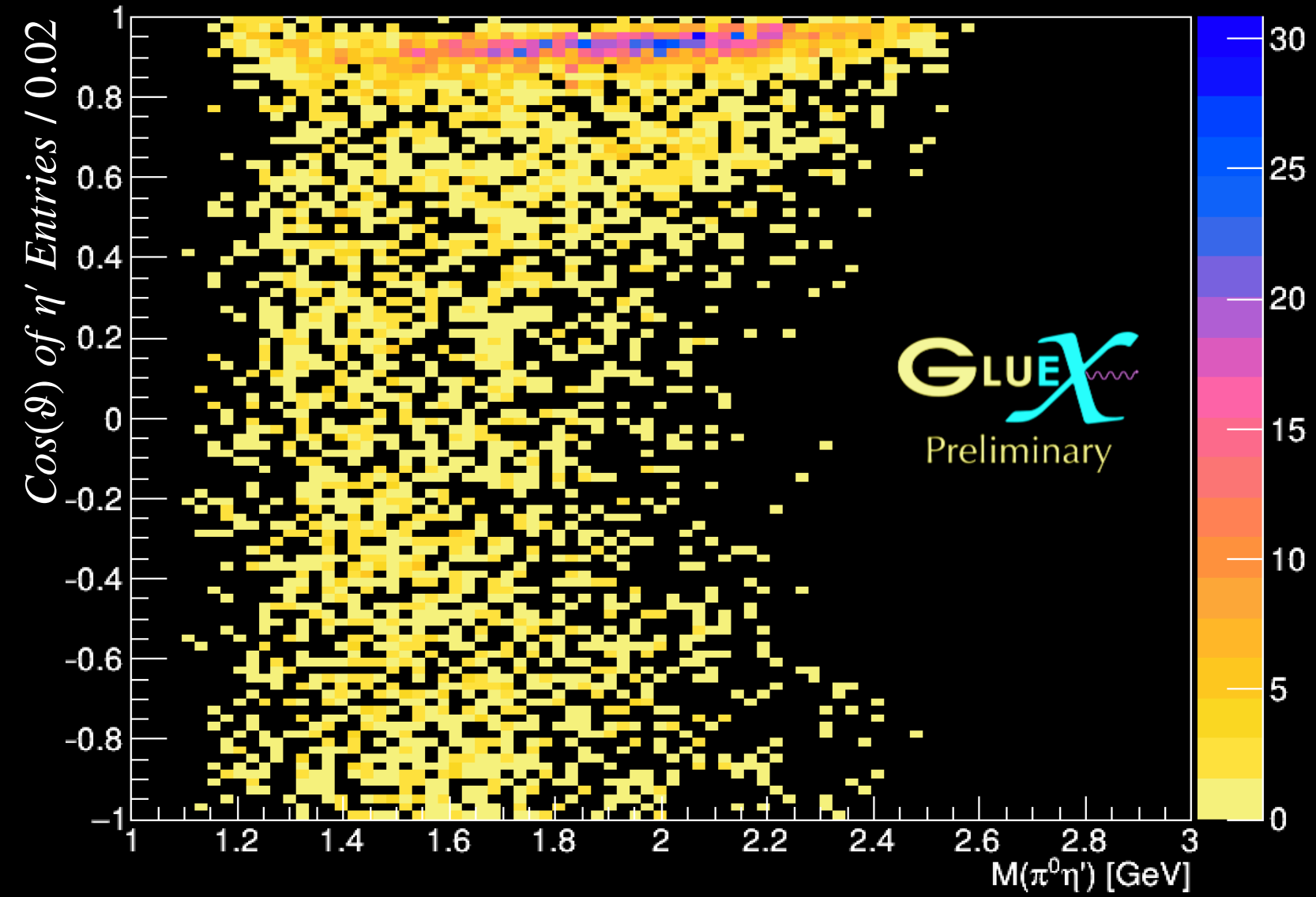
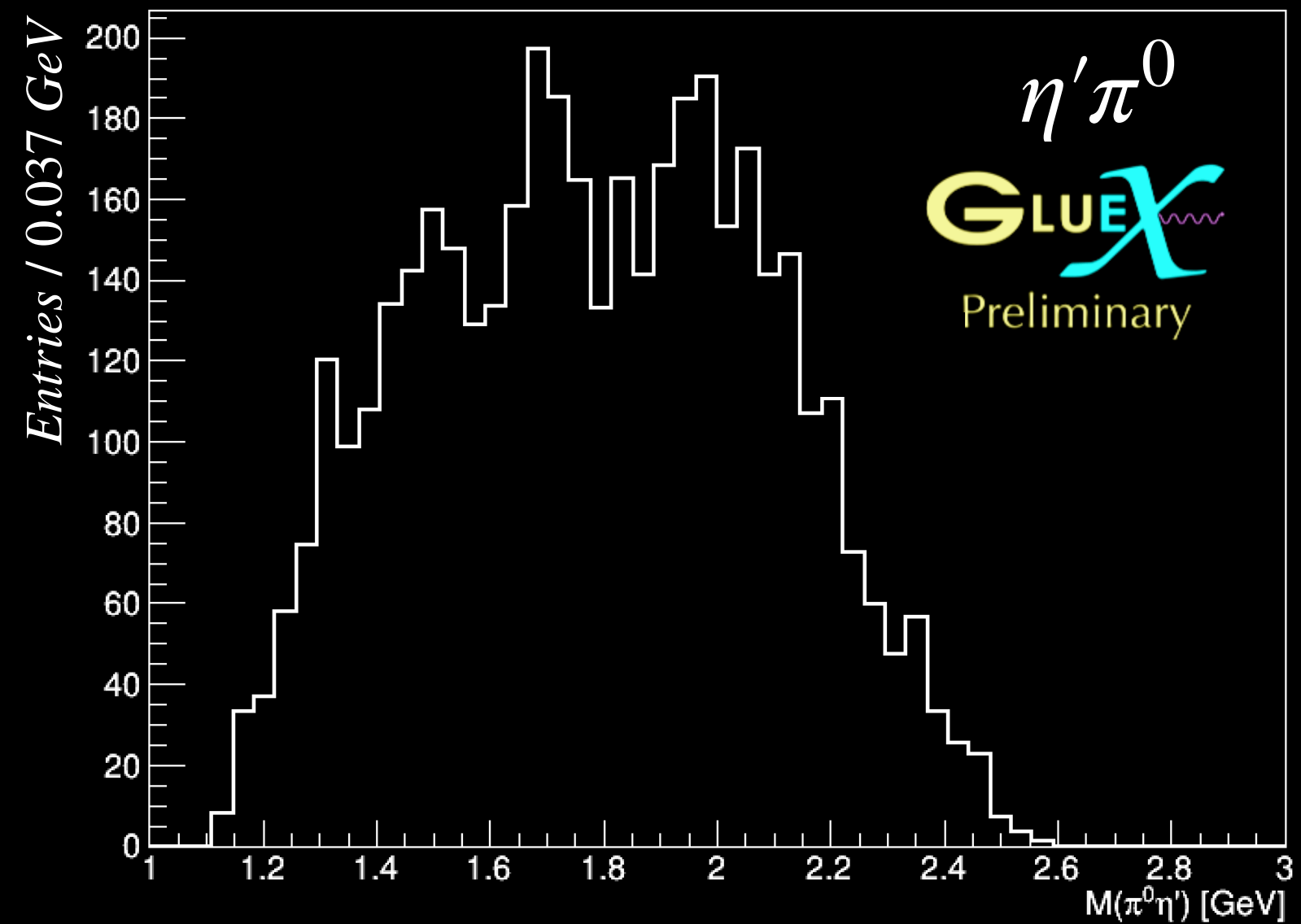


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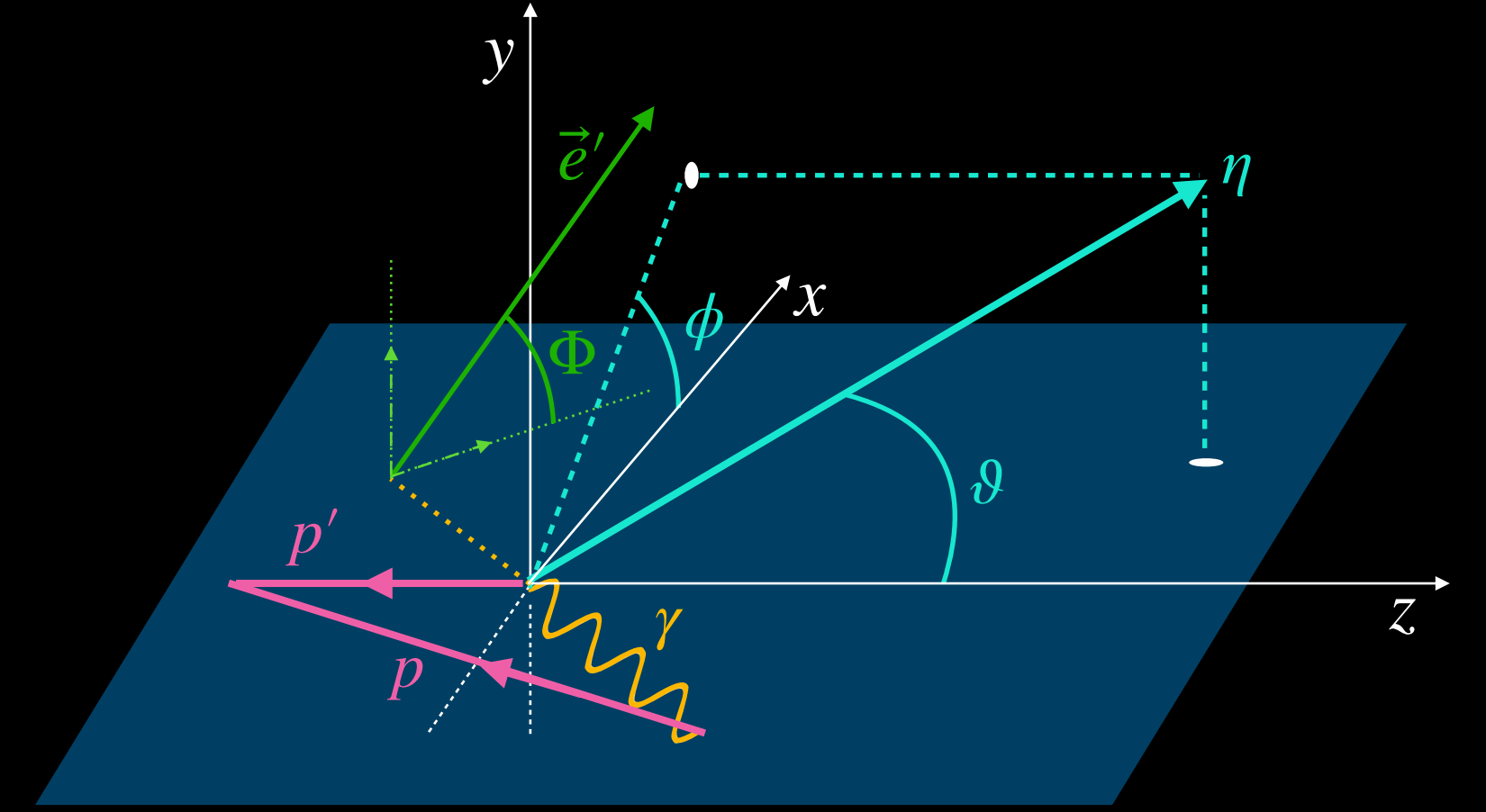
Dominant S_0 and D_2 contributions observed



Describes all two-pseudoscalar systems (i.e. all $\eta^{(\prime)} \pi$)

- basis $\rightarrow Z_l^m(\Omega, \Phi) = Y_l^m(\Omega) e^{-i\Phi}$

- described by 3 angles: $\left. \begin{matrix} \cos \vartheta_{\eta^{(\prime)}} \\ \phi_{\eta^{(\prime)}} \\ \Phi \end{matrix} \right\}$ in the resonance frame
 $\Phi \rightarrow$ between the polarization and production plane



V. Mathieu et al. [JPAC], PRD 100, 054017 (2019)

- reflectivity corresponds to exchange being natural (+) and unnatural (-) parity

- 4x more amplitudes than hadroproduction

$$\Rightarrow \mathcal{I}(\Omega, \Phi) = 2\kappa \sum_k \left\{ (1 - P_\gamma) \left| \sum_{l,m} [l]_m^{(-)} \mathcal{R}e[Z_l^m(\Omega, \Phi)] \right|^2 + (1 - P_\gamma) \left| \sum_{l,m} [l]_m^{(+)} \mathcal{I}m[Z_l^m(\Omega, \Phi)] \right|^2 + \right. \\ \left. (1 + P_\gamma) \left| \sum_{l,m} [l]_m^{(+)} \mathcal{R}e[Z_l^m(\Omega, \Phi)] \right|^2 + (1 + P_\gamma) \left| \sum_{l,m} [l]_m^{(-)} \mathcal{I}m[Z_l^m(\Omega, \Phi)] \right|^2 \right\}$$

Assume $a_2(1320)$ and $a_2(1700)$ are text book Breit-Wigner resonances

- share only 1 common phase parameter for each in the D_{waves}

S_{wave} contributions more complicated

- define mass independent piecewise parameterization

Individual fit results across $-t$

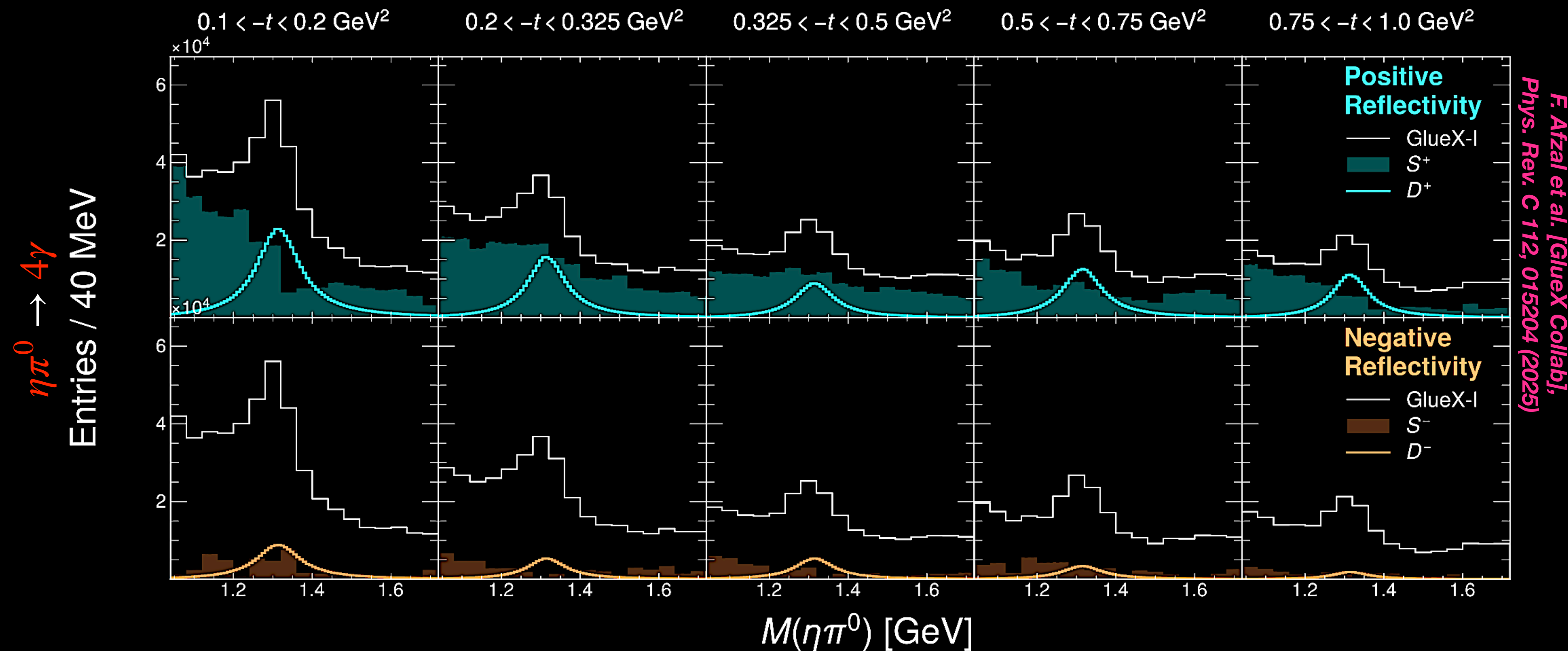
- incoherent sums of (+) and (-) reflectivities

$$\eta = P(-1)^J$$

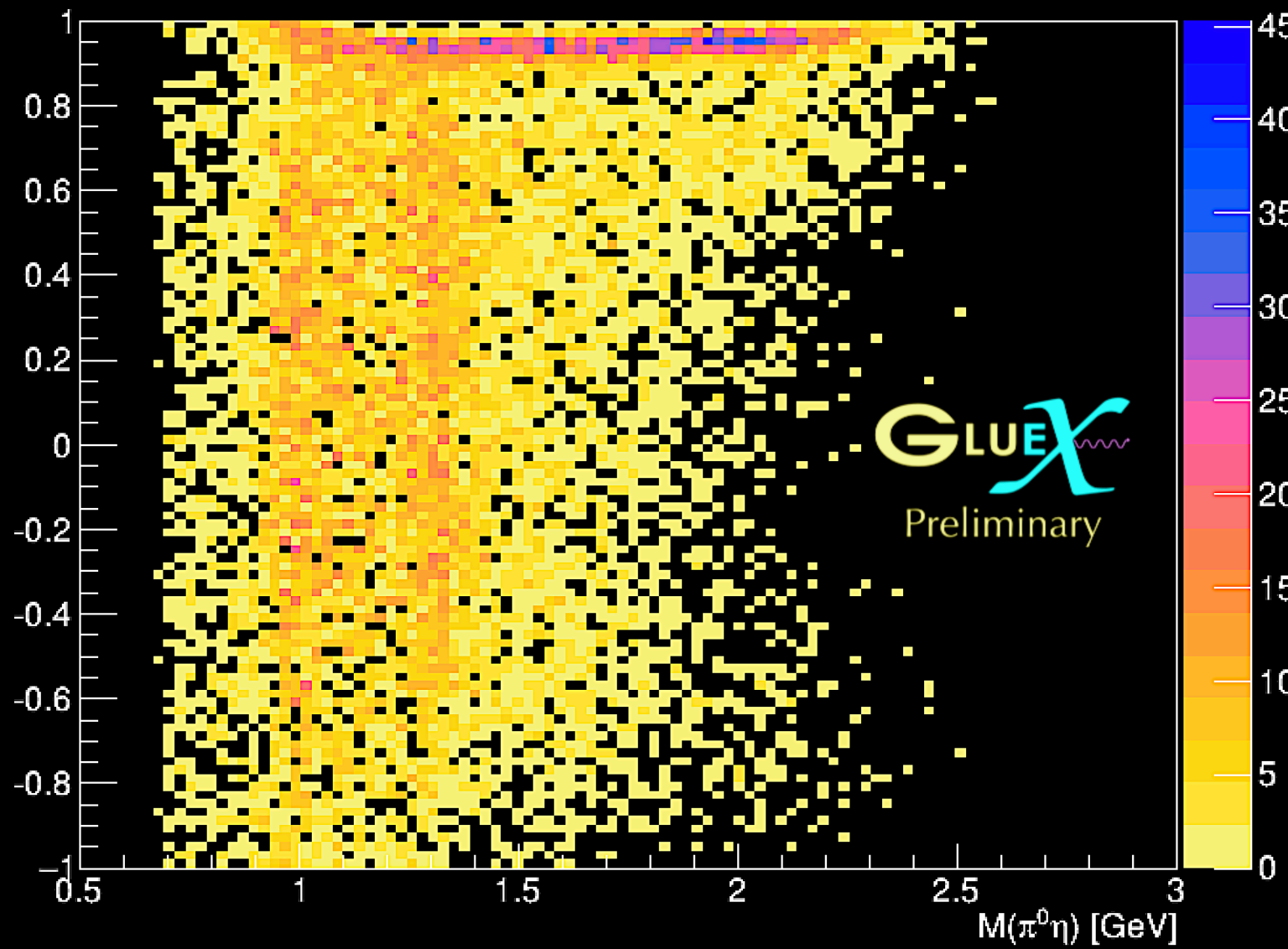
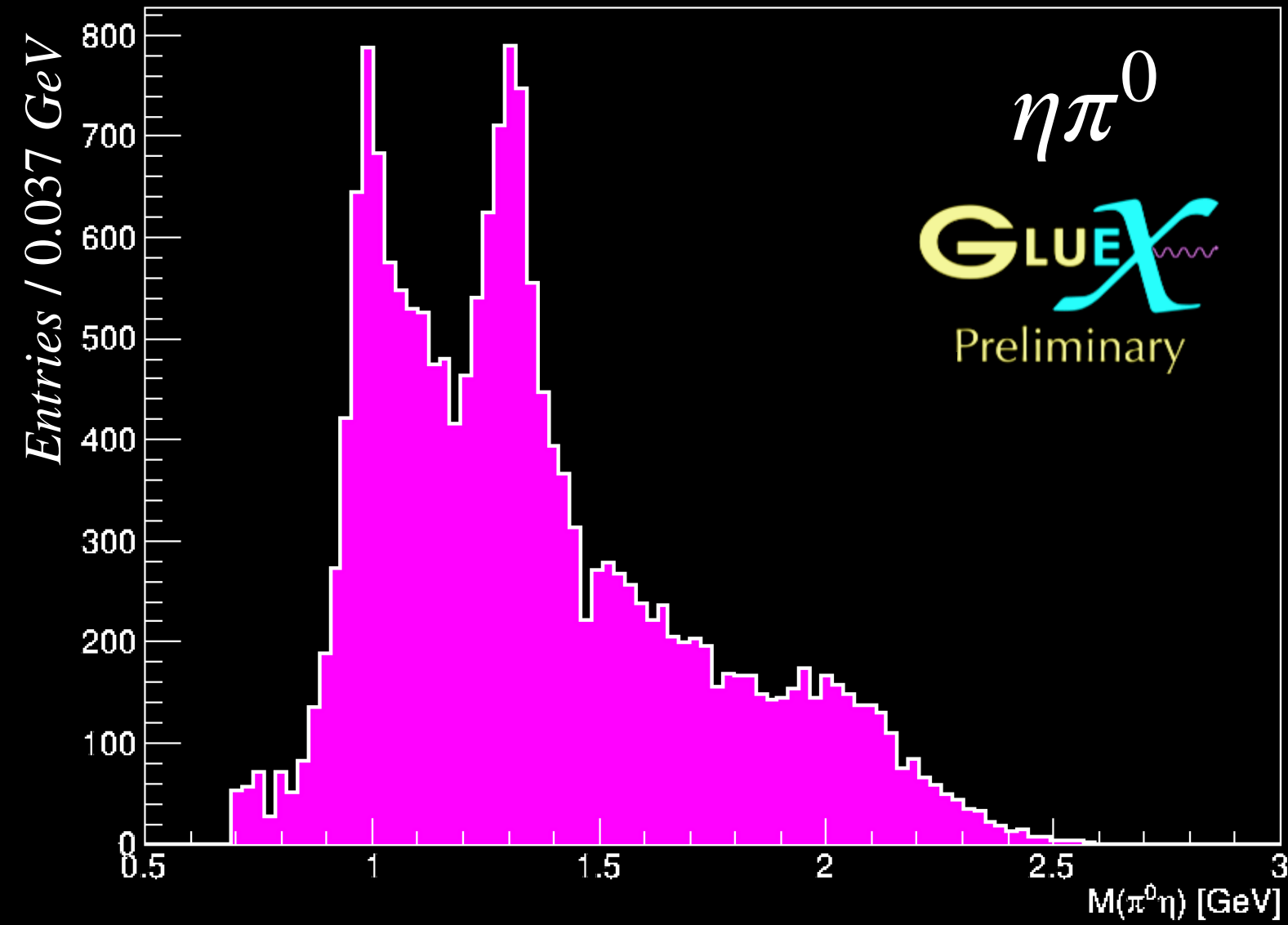
Why ?

Dominant contribution in the $\eta\pi^0$ channel

- reasonably isolated
- limited P_{Wave} contribution predicted
- use as reference for search of exotic π_1 in $\eta'\pi$

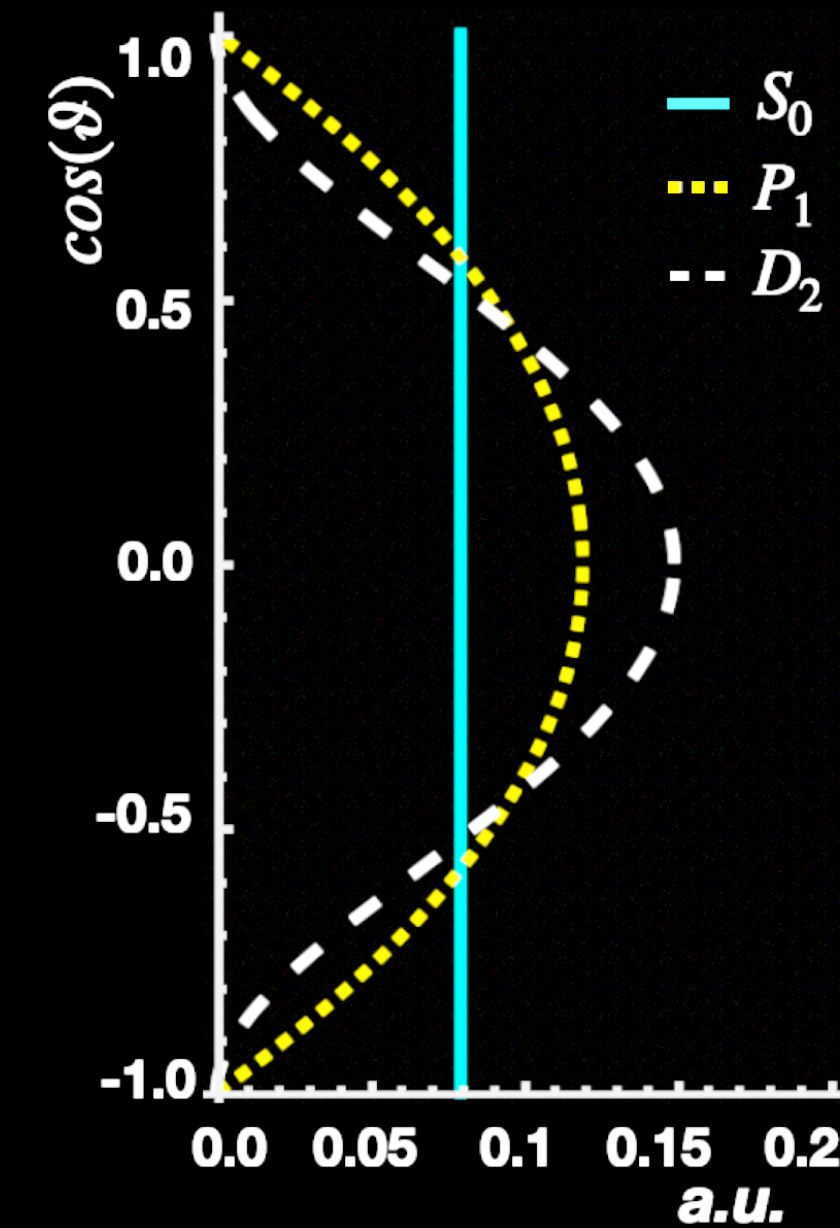


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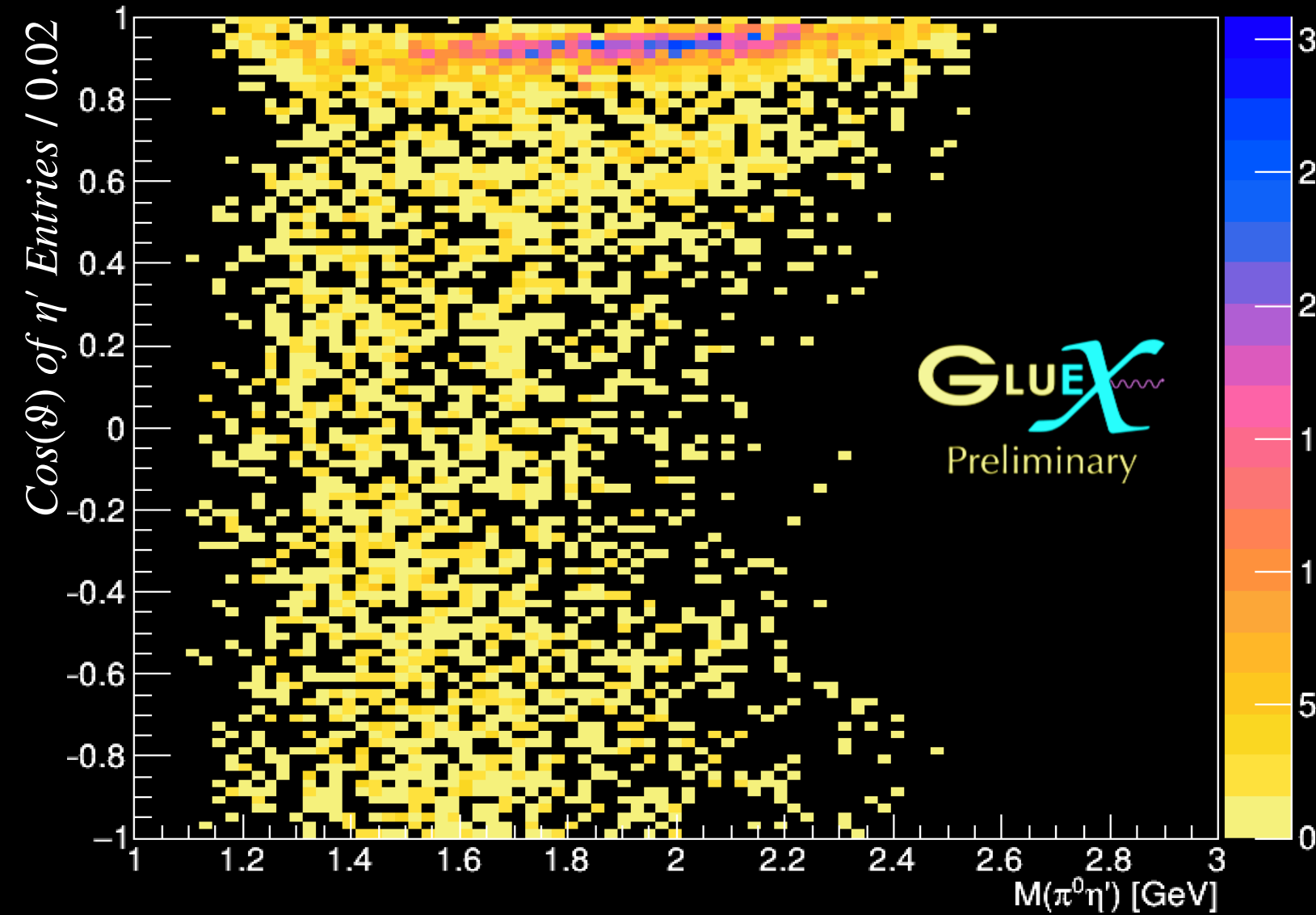
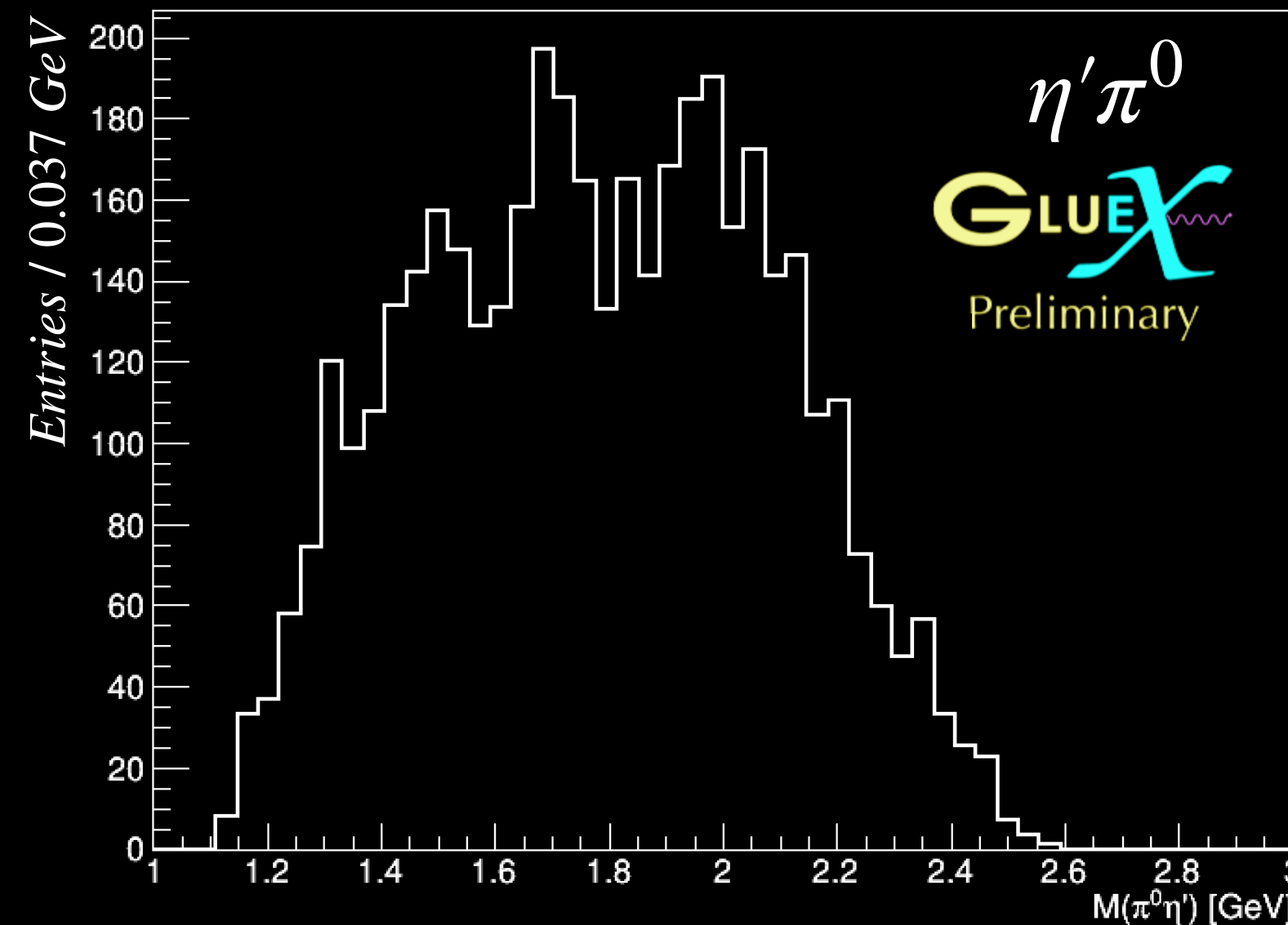
Neutral decay modes

- $\eta\pi^0 \rightarrow 4\gamma\pi^+\pi^-$
 - $\eta'\pi^0 \rightarrow 4\gamma\pi^+\pi^-$
- } Same final state



Analyzing a different final state for the same channel

- cross-validation
- completeness of amplitude solutions
- channel dependent backgrounds



Nature of strong interactions:

- governed by non-perturbative QCD
- allows resonance formation and decay across multiple channels

Spectra are not observables of resonance structure

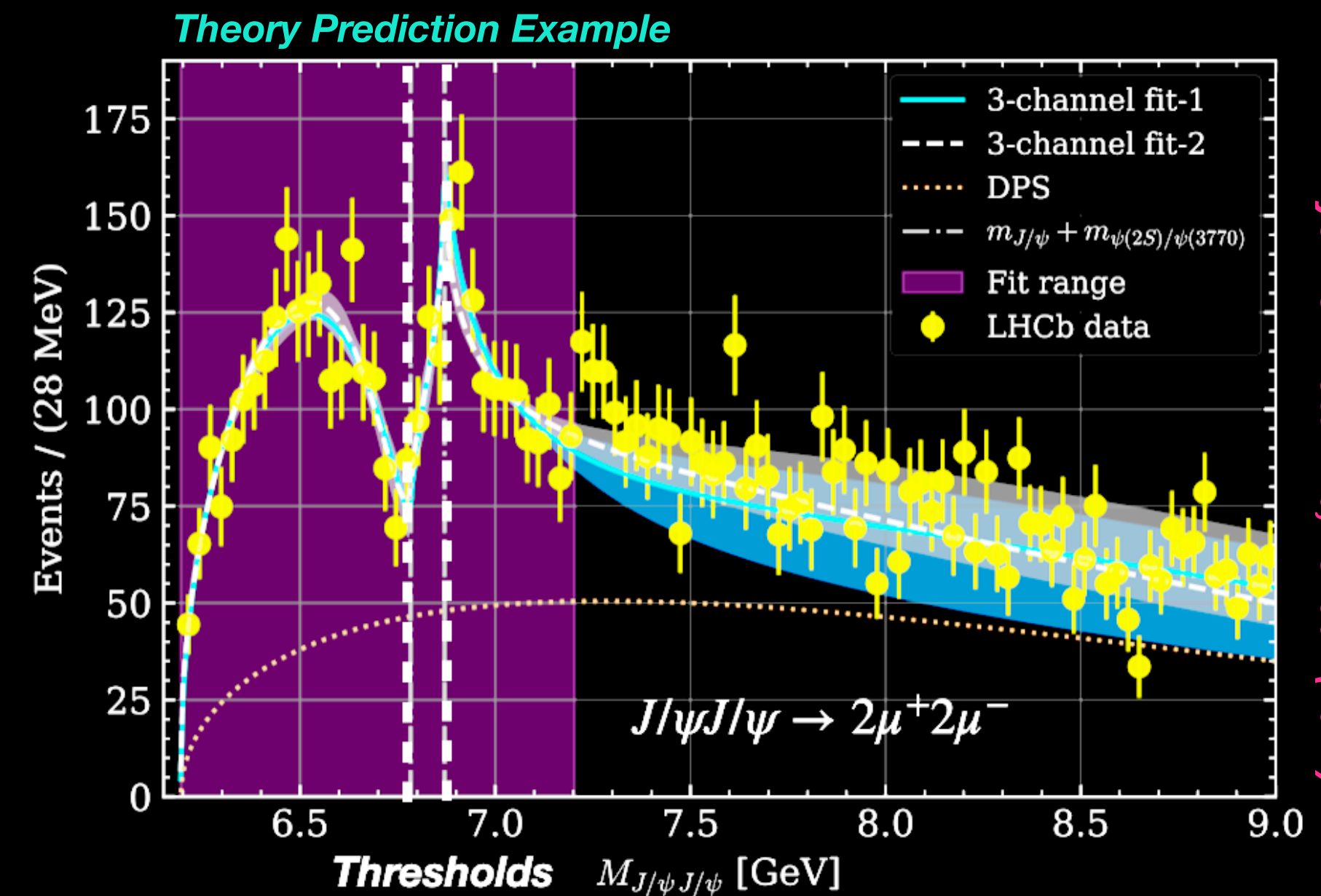
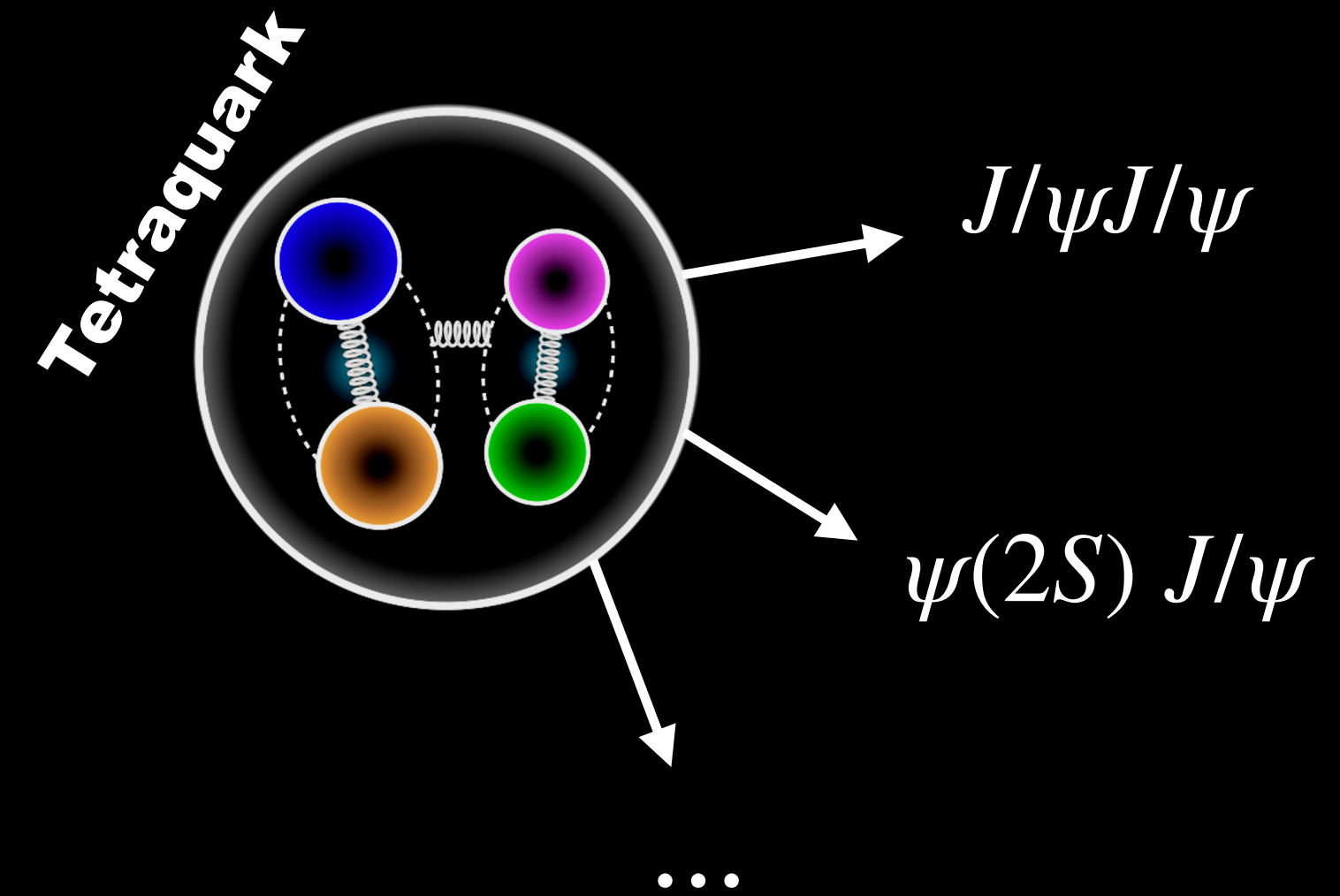
- measured distributions depend on production dynamics
- a single decay channel samples only a projection of the full S -matrix

highly populated spectrums
= overlapping & interfering resonances

Overall, single channel analyses **cannot** fully disentangle complex interference or threshold behavior

pole structure is process-independent
(the resonance's nature is fixed)

∴ production and decay can shape the spectrum,
but not the resonance itself



Xiang-Kun Dong et al.,
Phys. Rev. Lett. 126, 132001 (2021)

JPAC analysis utilizing COMPASS data $\rightarrow N(s)/D(s)$

- coupled channel fit to both $\eta^{(\prime)}\pi$ systems
- describes dominate a_2 resonances and the π_1
- poles identified via analytic continuation (zeros of $D(s)$), not peak positions

Angular Momentum Barrier Factors

Production Term (smooth, real functions)

$$a_i^J(s) = q^{J-1} p_i^J \sum_k n_k^J(s) [D^J(s)^{-1}]_{ki}$$

$$D_{ki}^J(s) = [K^J(s)^{-1}]_{ki} - \frac{s}{\pi} \int_{s_k}^{\infty} ds' \frac{\rho N_{ki}^J(s')}{s'(s' - s - i\epsilon)}$$

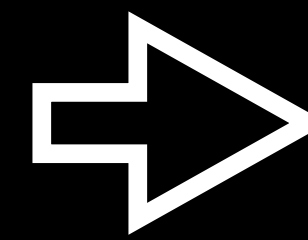
Analytic Denominator Matrix



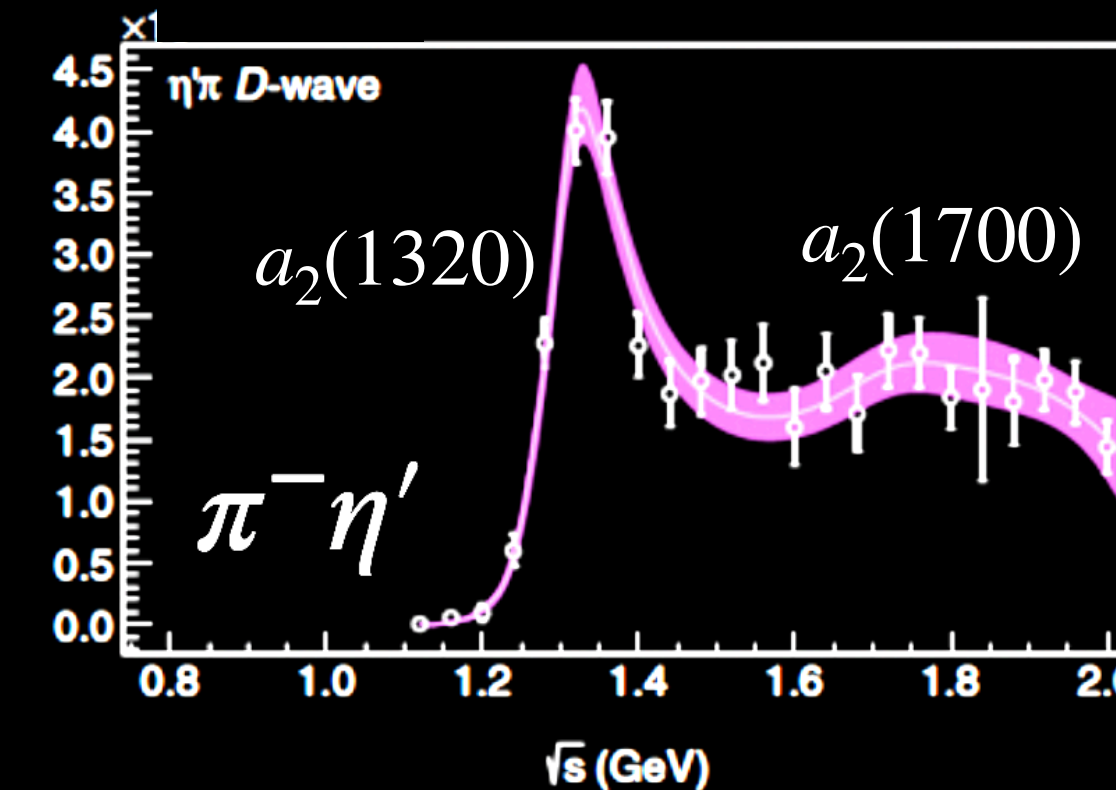
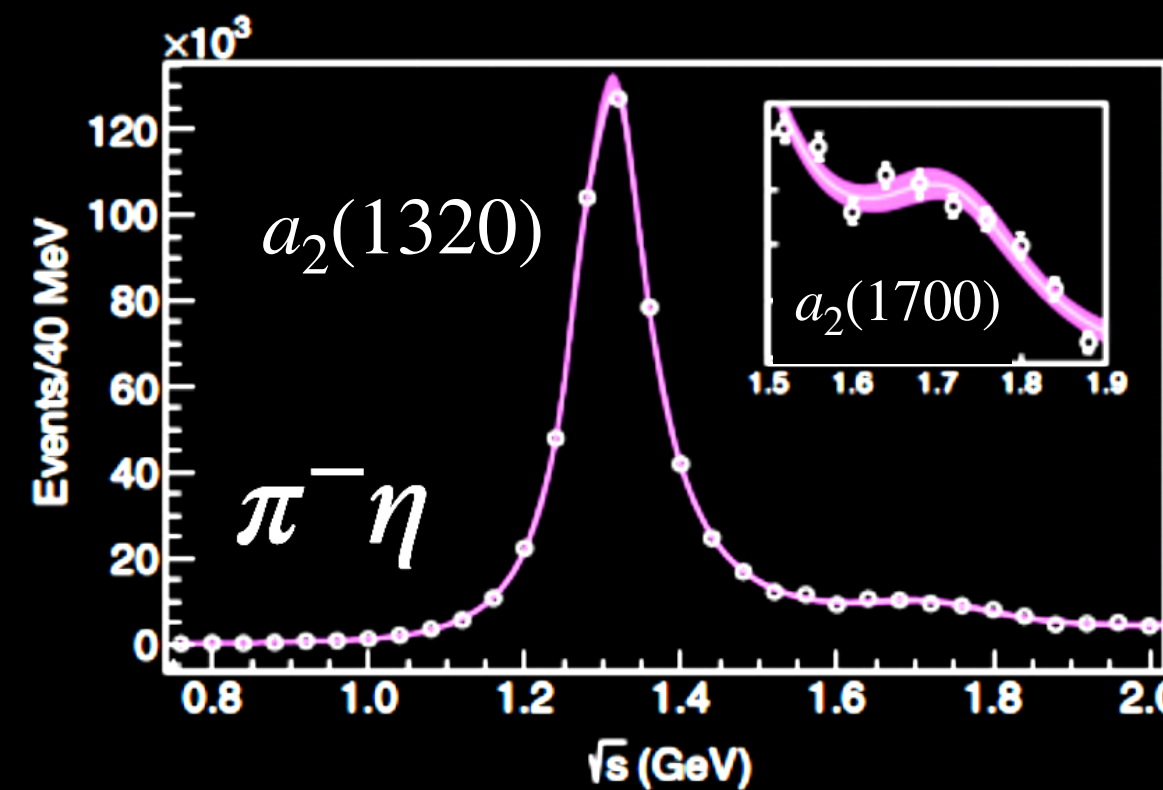
A. Rodas et al. [Joint Physics Analysis Center], PRL 122, 042002 (2019)

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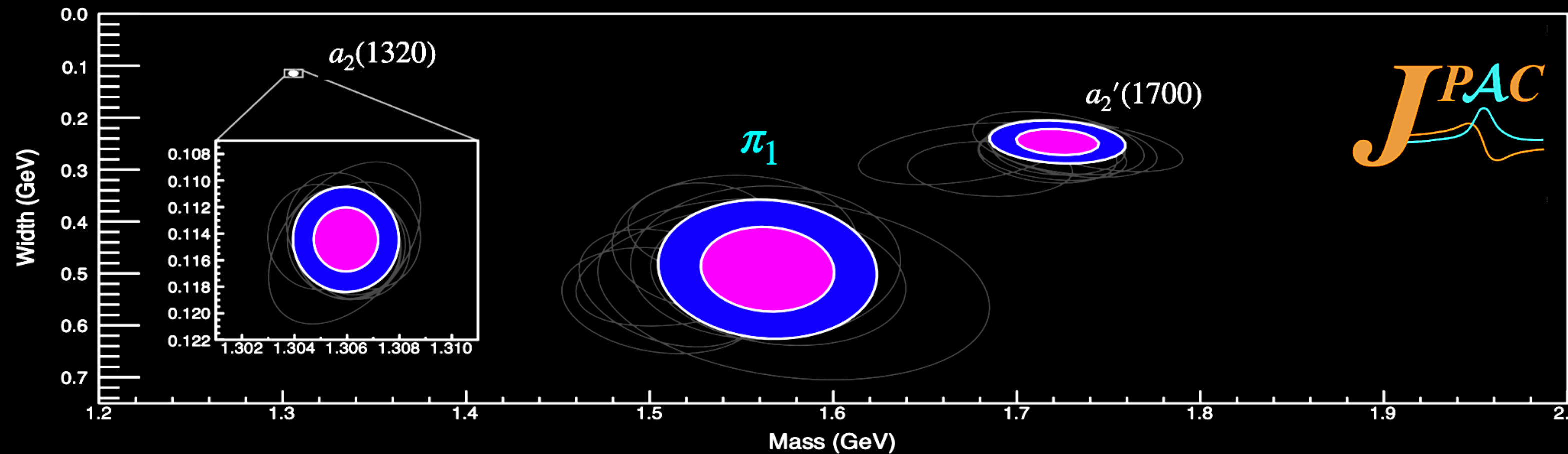
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A single pole explains structures in multiple channels with shifted peaks

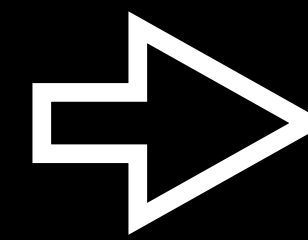


A. Rodas et al. [Joint Physics Analysis Center], PRL 122, 042002 (2019)

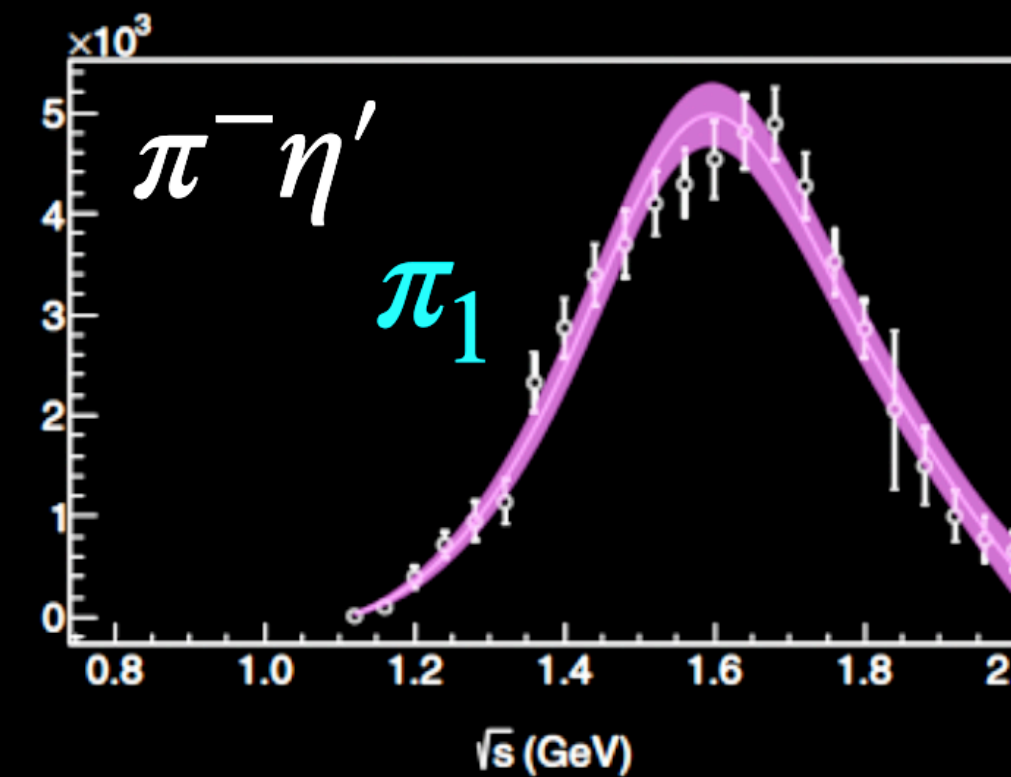
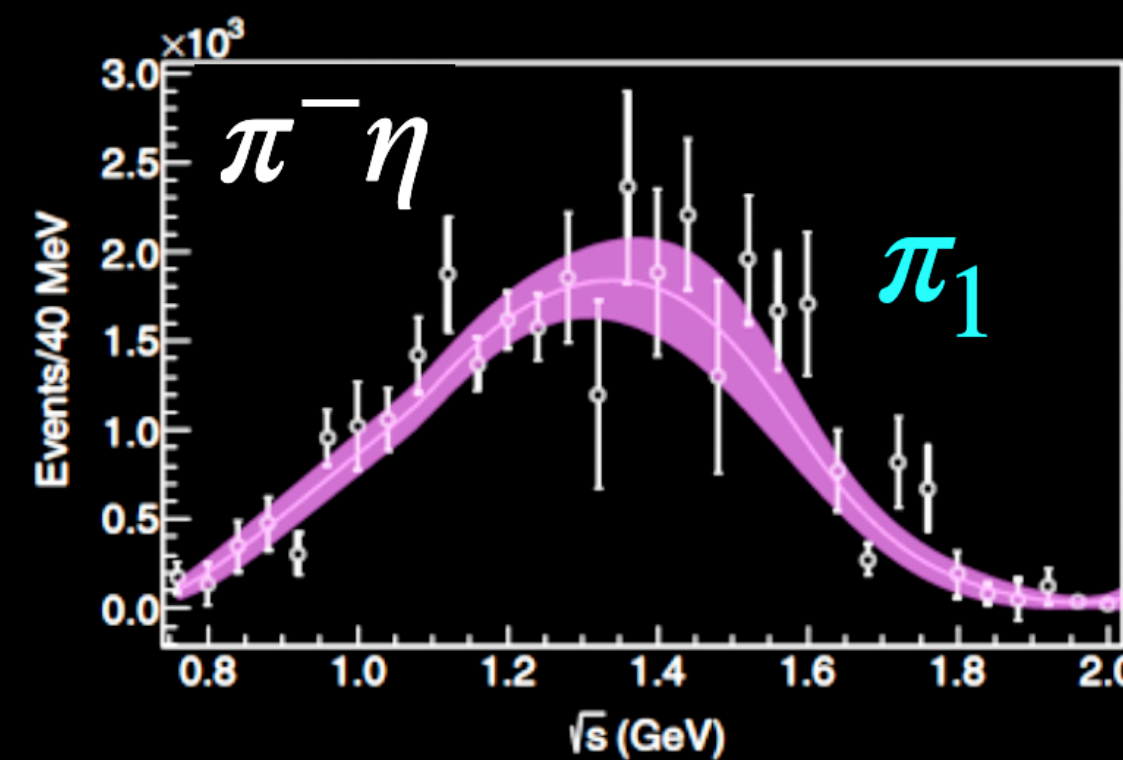


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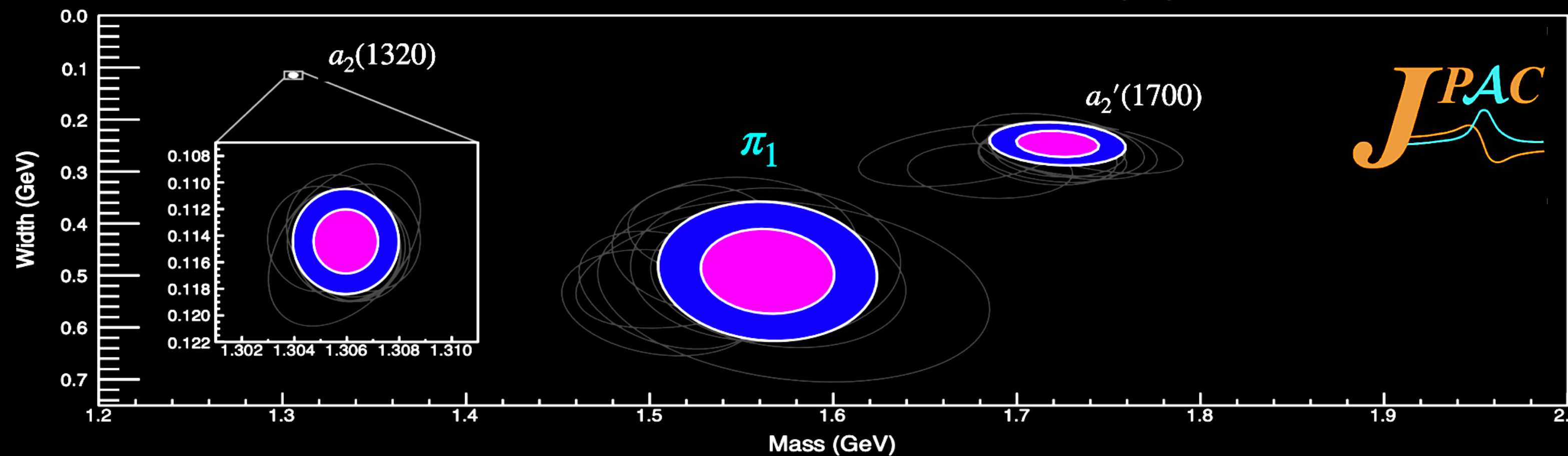
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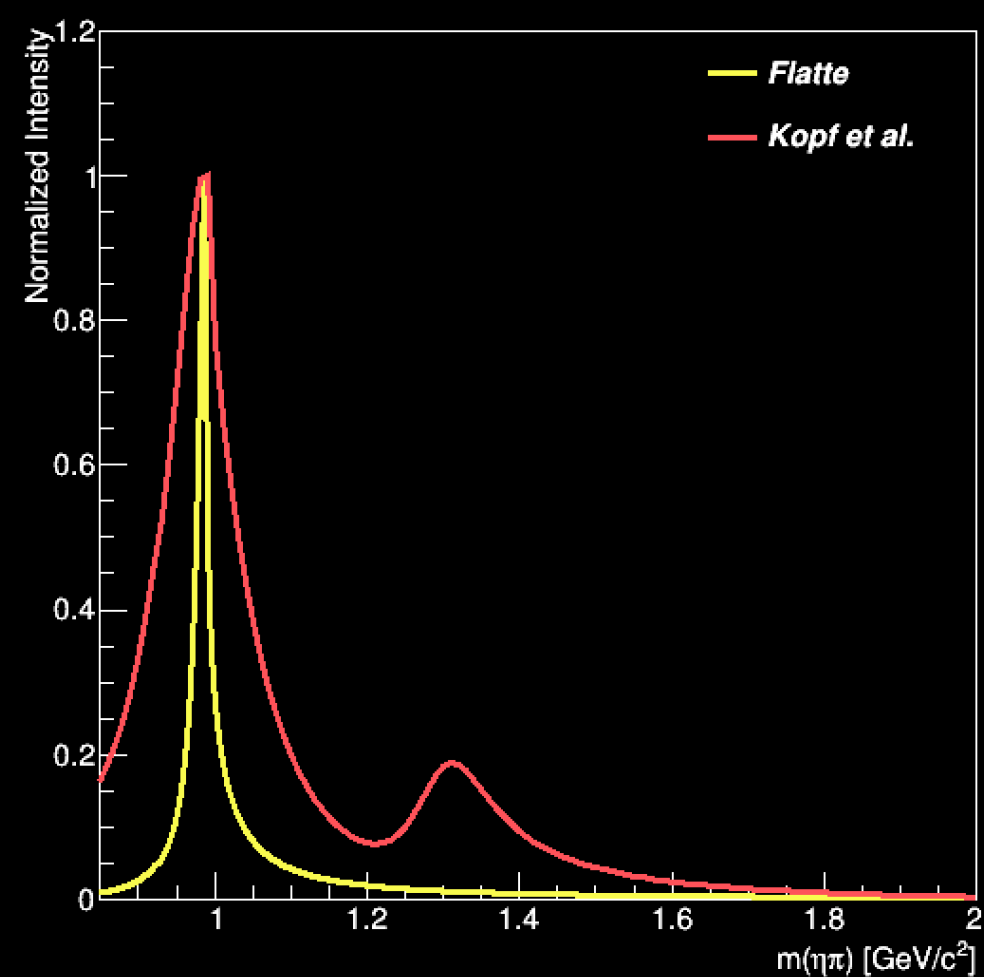
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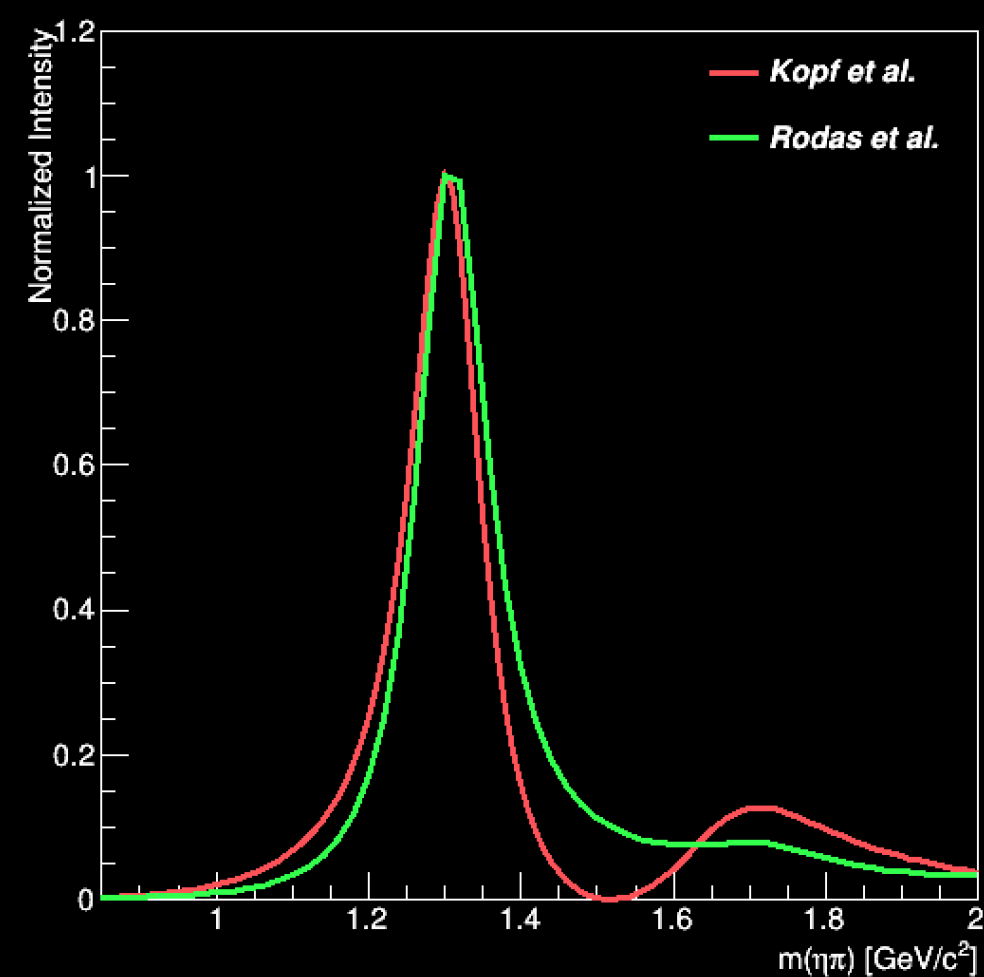
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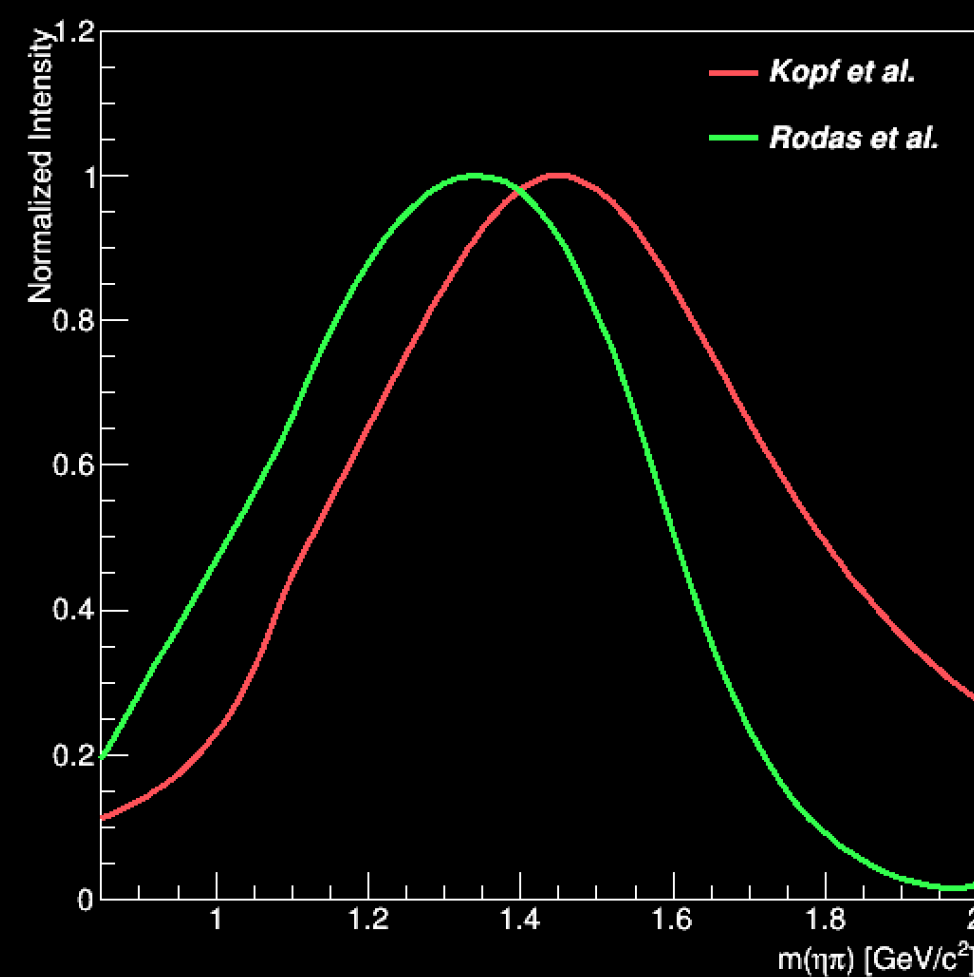
S_{wave}



D_{wave}



P_{wave}

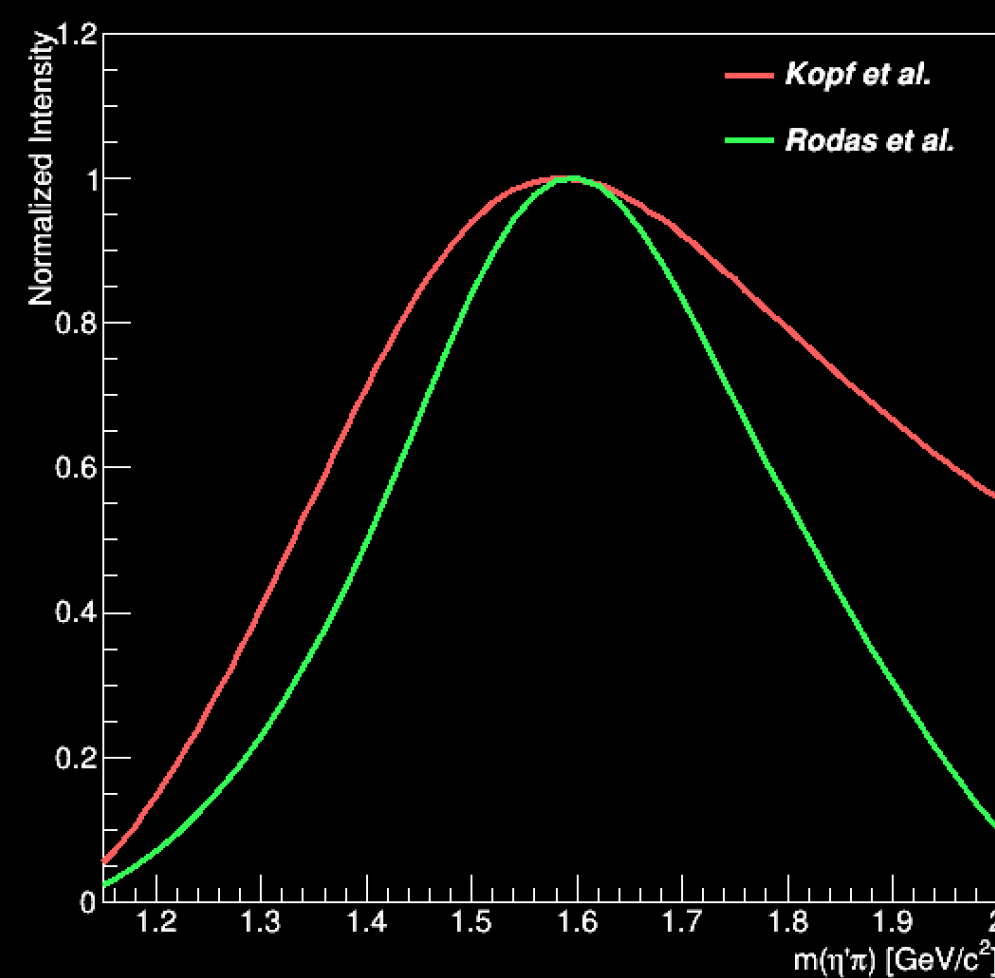
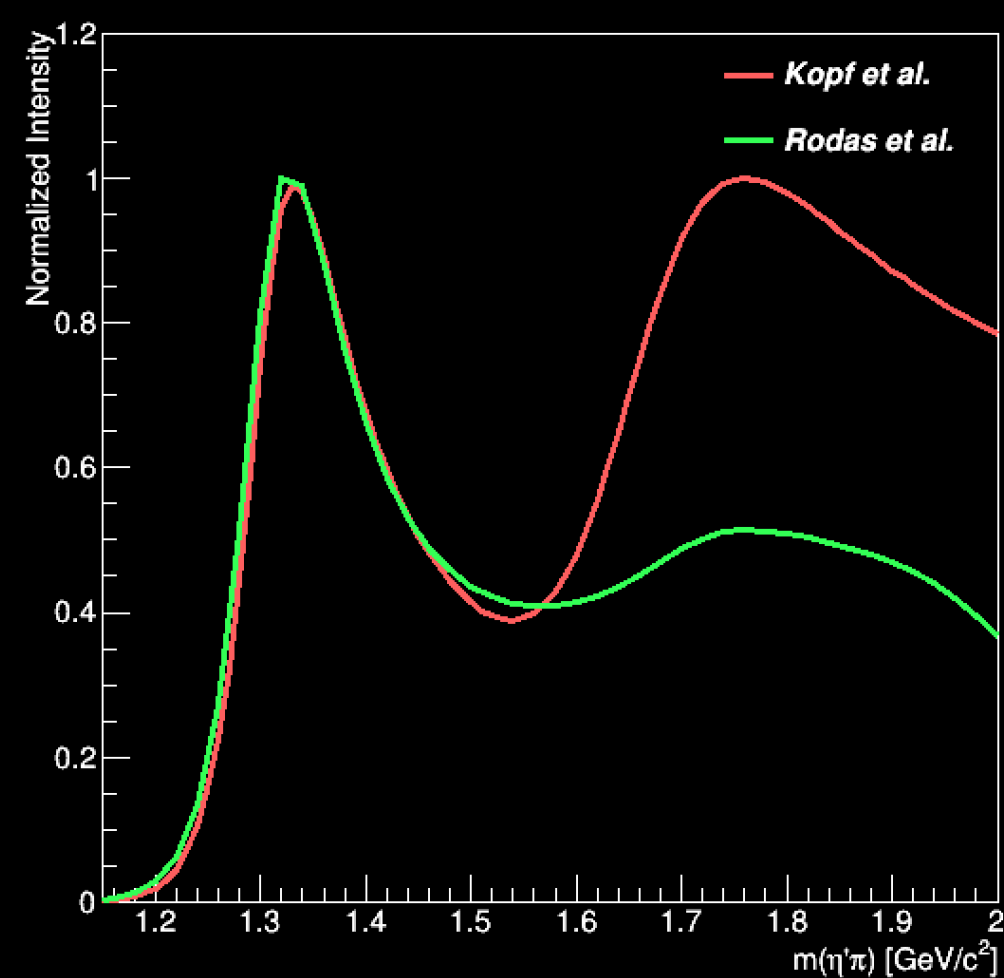


Numerically evaluation of amplitudes from Rodas, *et al.* & other coupled-channel models to GlueX data

Scattering and pole structures are universal

- can use published results to fit photoproduced $\eta\pi$ and $\eta'\pi$!

No coupled-channel parameterization of S_{wave} in $\eta'\pi$



Other methods:

Flatté

S. M. Flatté, Phys. Lett. B 63, 224 (1976)

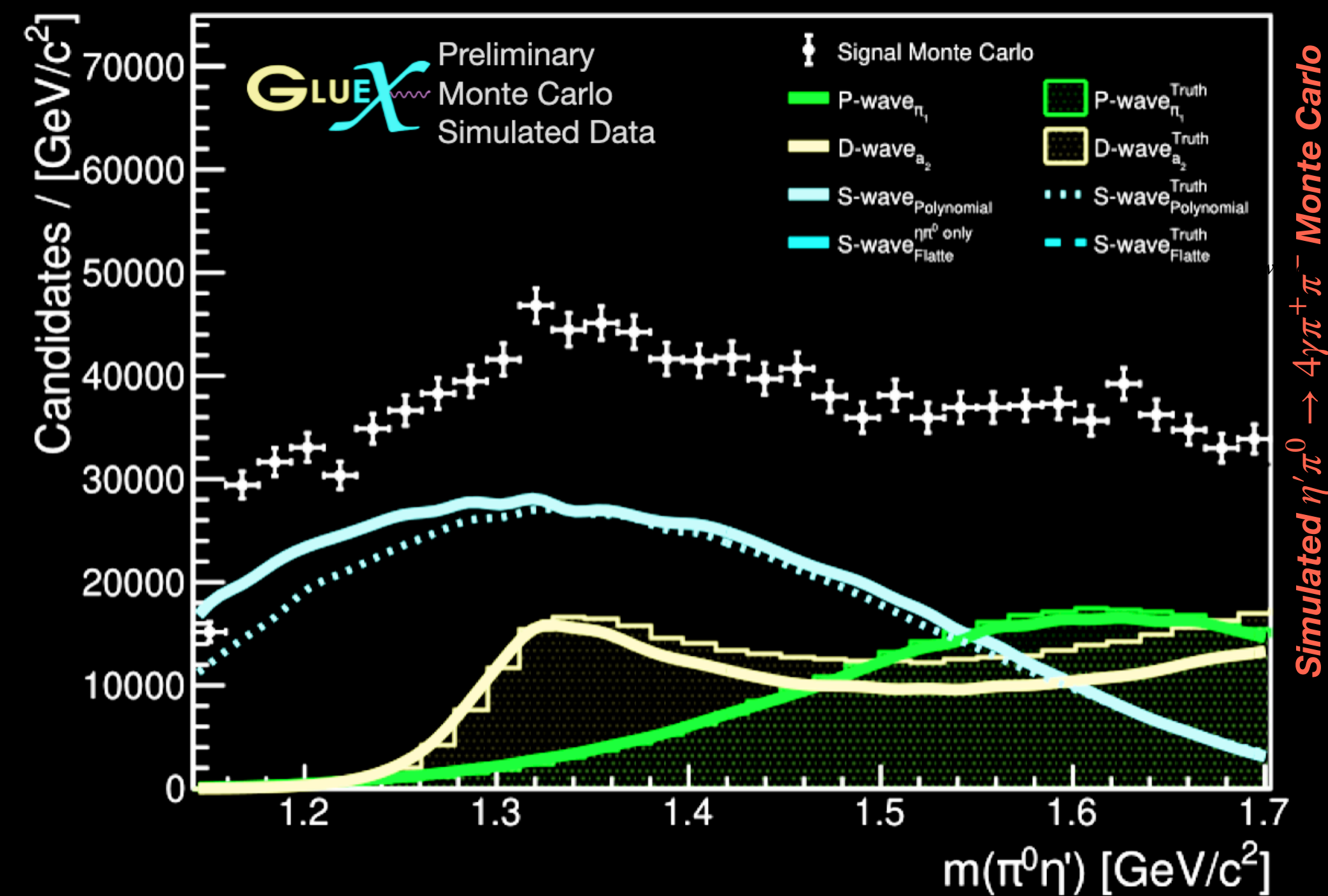
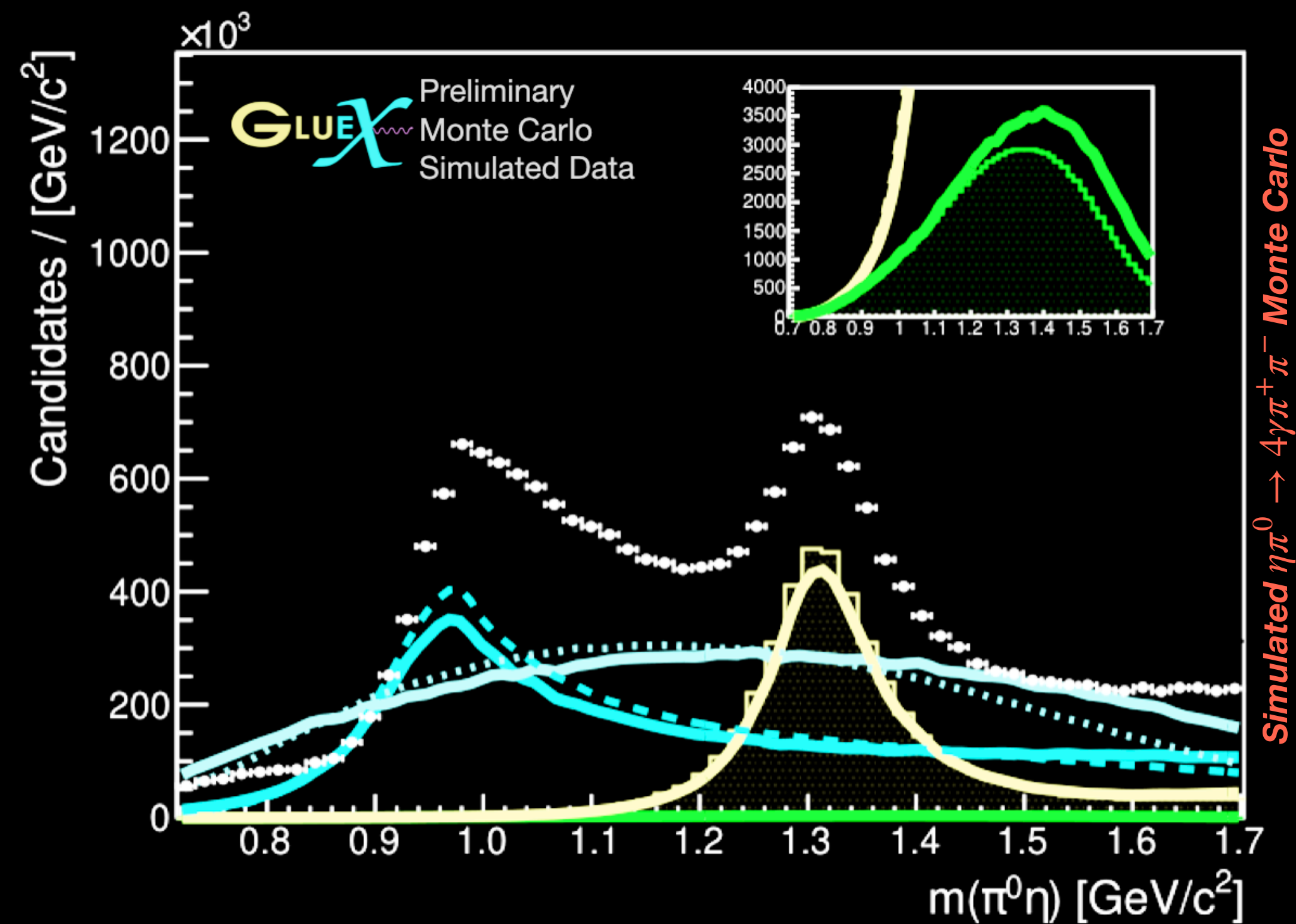
Kopf, *et al.*

B. Kopf et al., Eur. Phys. J. C 81, 1056 (2021)

Note: for applications to photoproduction, models need tweaking

Input-Output check with generated signal Monte Carlo for photoproduction kinematics using coupled channel approach

- lines-shapes remain consistent to the truth
- extremely small P_{wave} in $\eta\pi^0$ channel compared to $\eta'\pi^0$
- S_{wave} contributions = greater difficulty (especially in $\eta'\pi$)



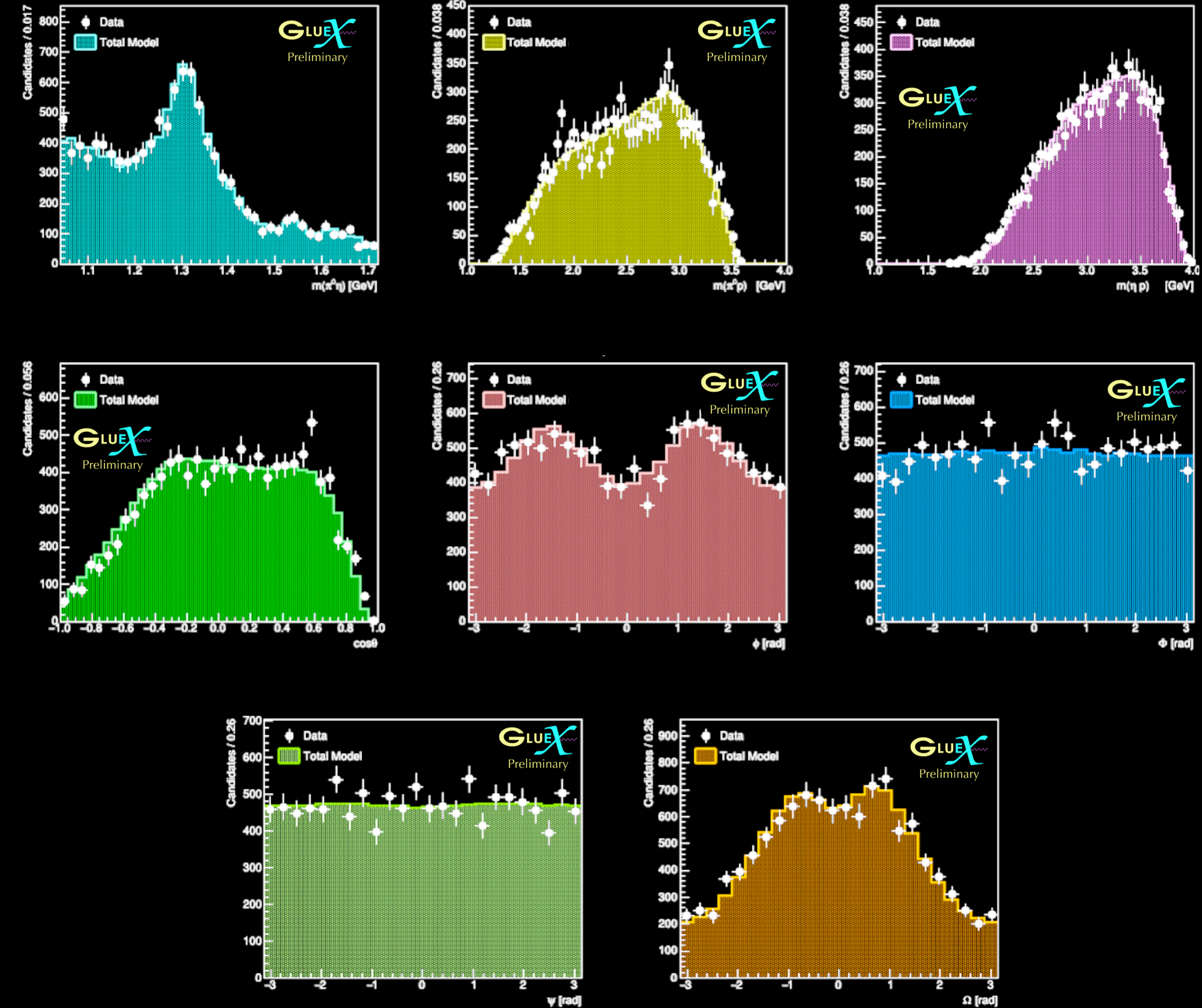
Decompose angular components into spherical harmonics $Y_l^m(\Omega)$

- moment decomposition is *unique*
- determine sensitivity to exotic contributions

No direct access to partial wave amplitudes

- BUT can calculate moments from partial waves

EXAMPLE: PWA fit results projected onto masses and angles



Decompose angular components into spherical harmonics $Y_l^m(\Omega)$

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- determine sensitivity to exotic contributions

No direct access to partial wave amplitudes

- BUT can calculate moments from partial waves

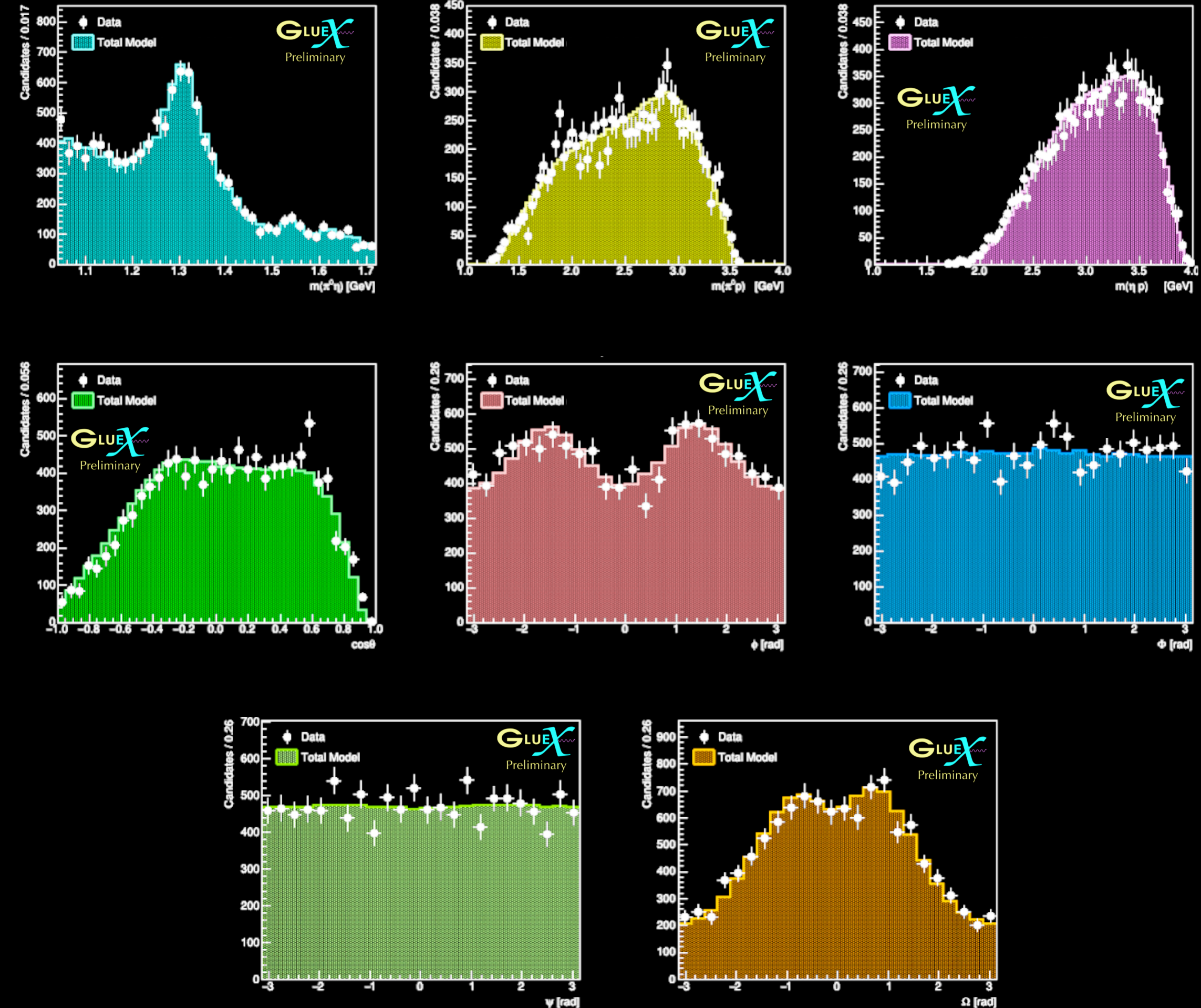
PRO

- model independence
- computationally simple

CON

- abstract - does not provide direct access to physical amplitudes
- interpretation = ambiguous with out theory

EXAMPLE: PWA fit results projected onto masses and angles



Decompose angular components into spherical harmonics $Y_l^m(\Omega)$

- moment decomposition is *unique*
- determine sensitivity to exotic contributions

No direct access to partial wave amplitudes

- BUT can calculate moments from partial waves

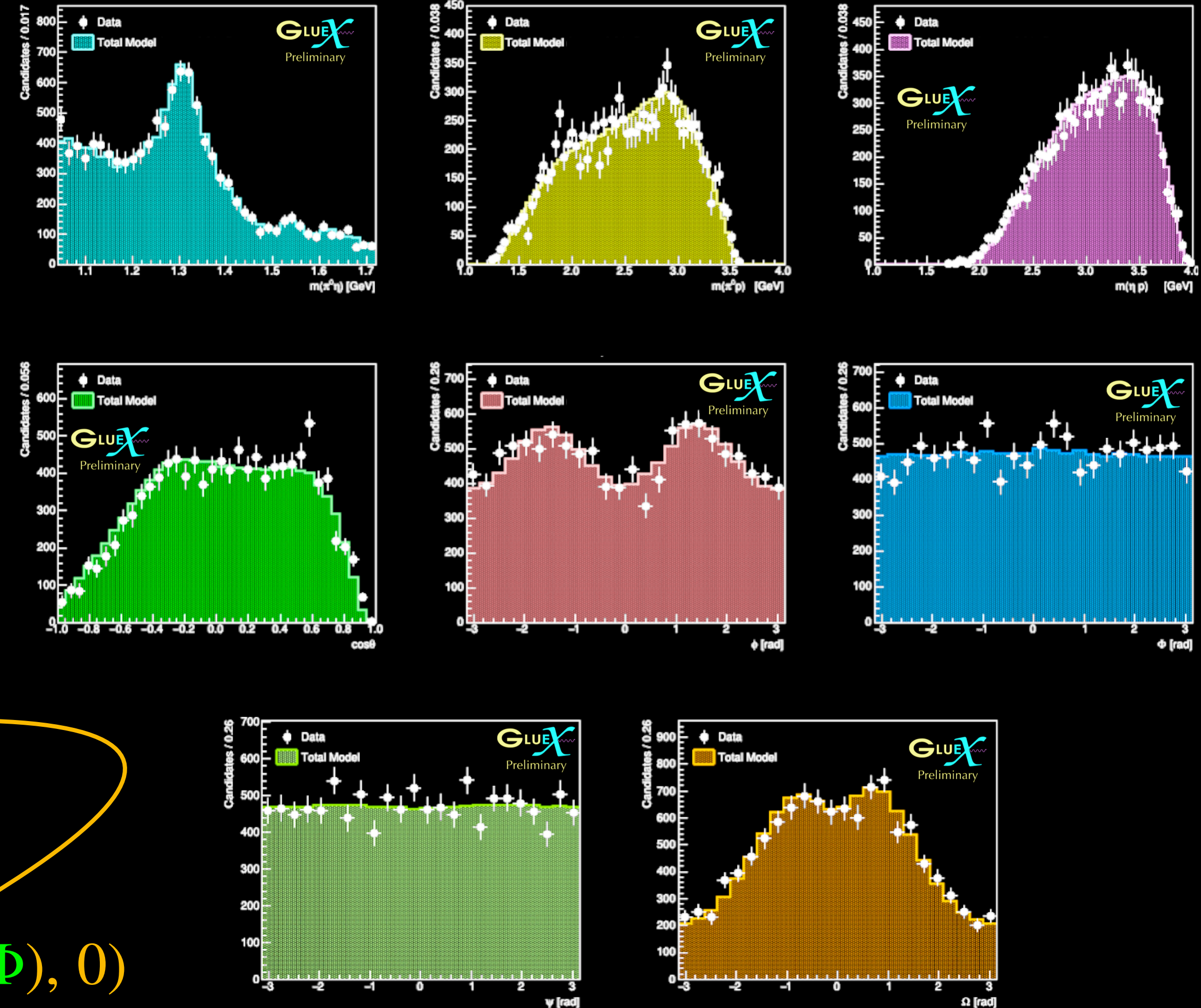
PRO

- model independence
- computationally simple

CON

- abstract - does not provide direct access to physical amplitudes
- interpretation = ambiguous with out theory

EXAMPLE: PWA fit results projected onto masses and angles



Linearly polarized γ

$$\mathcal{I}(\Omega, \Phi) = \sum_{\alpha=0} \mathcal{I}_{\alpha}(\Omega) P_{\gamma}^{\alpha}(\Phi)$$

$$P_{\gamma}^{\alpha}(\Phi) = (1, -P_{\gamma} \cos(2\Phi), -P_{\gamma} \sin(2\Phi), 0)$$

$$\mathcal{I}(\Omega, \Phi) = \mathcal{I}_0(\Omega) - \mathcal{I}_1(\Omega) P_{\gamma} \cos(2\Phi) - \mathcal{I}_2(\Omega) P_{\gamma} \sin(2\Phi)$$

V. Mathieu et al. [JPAC], PRD 100, 054017 (2019)

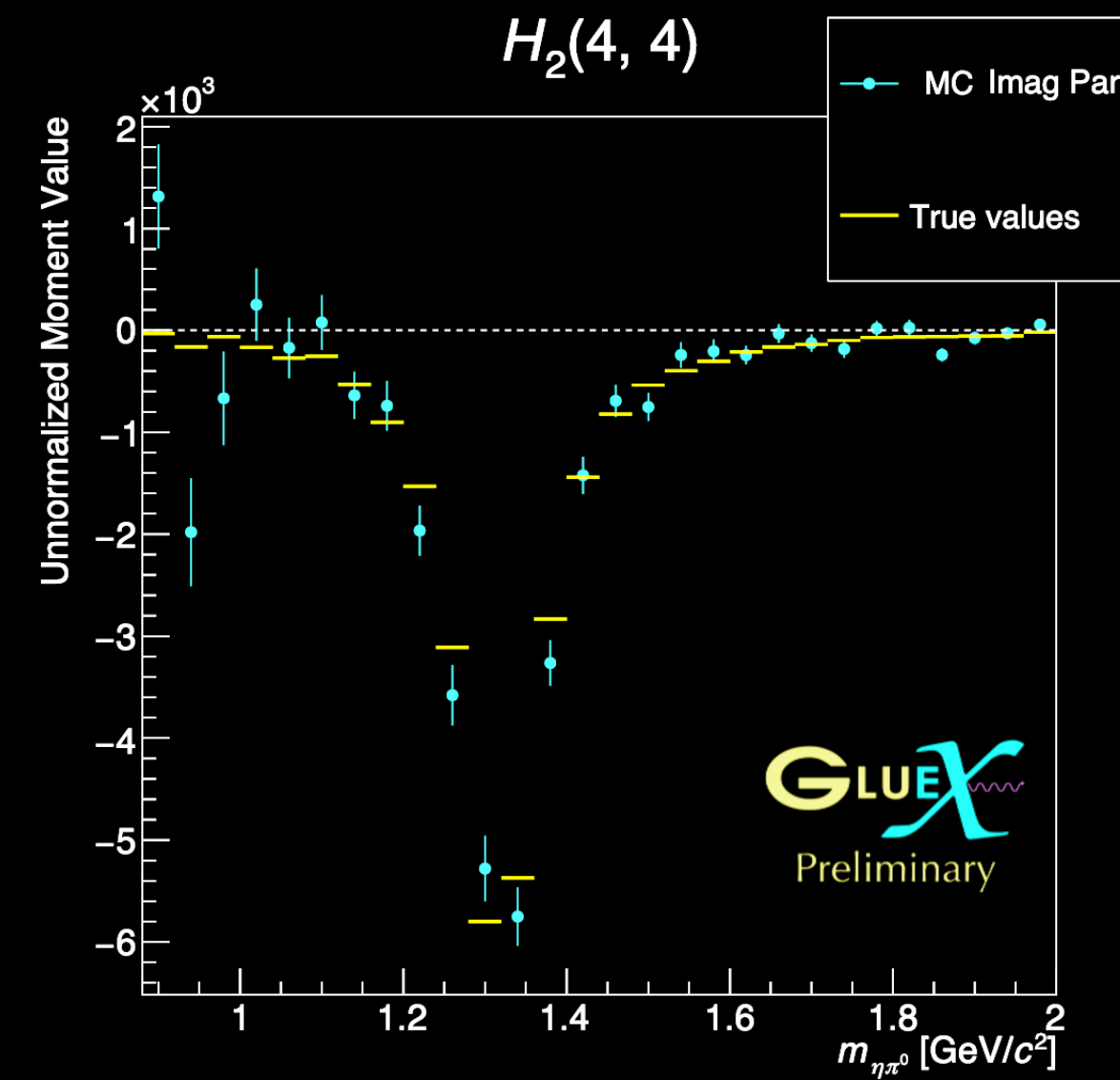
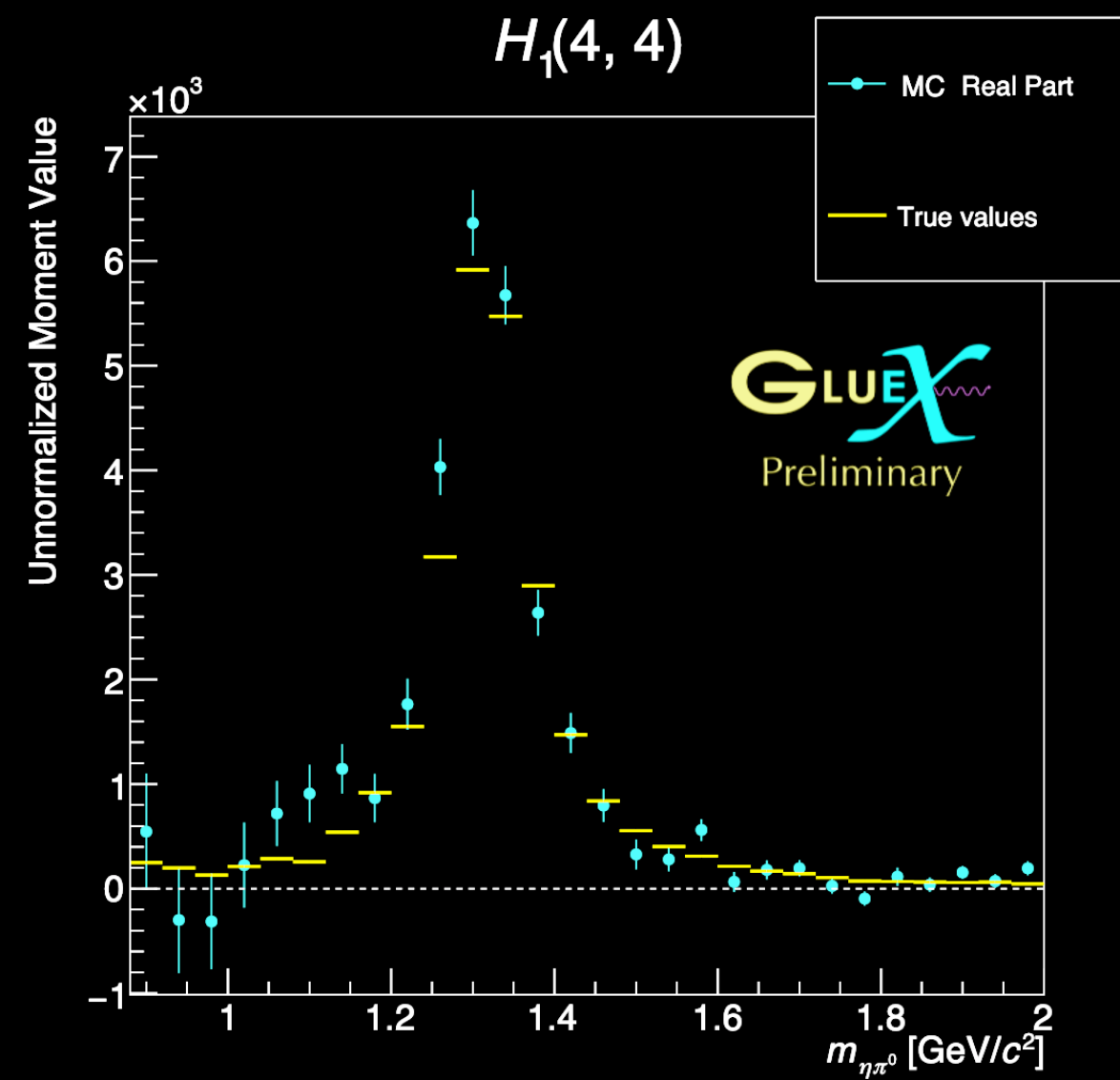
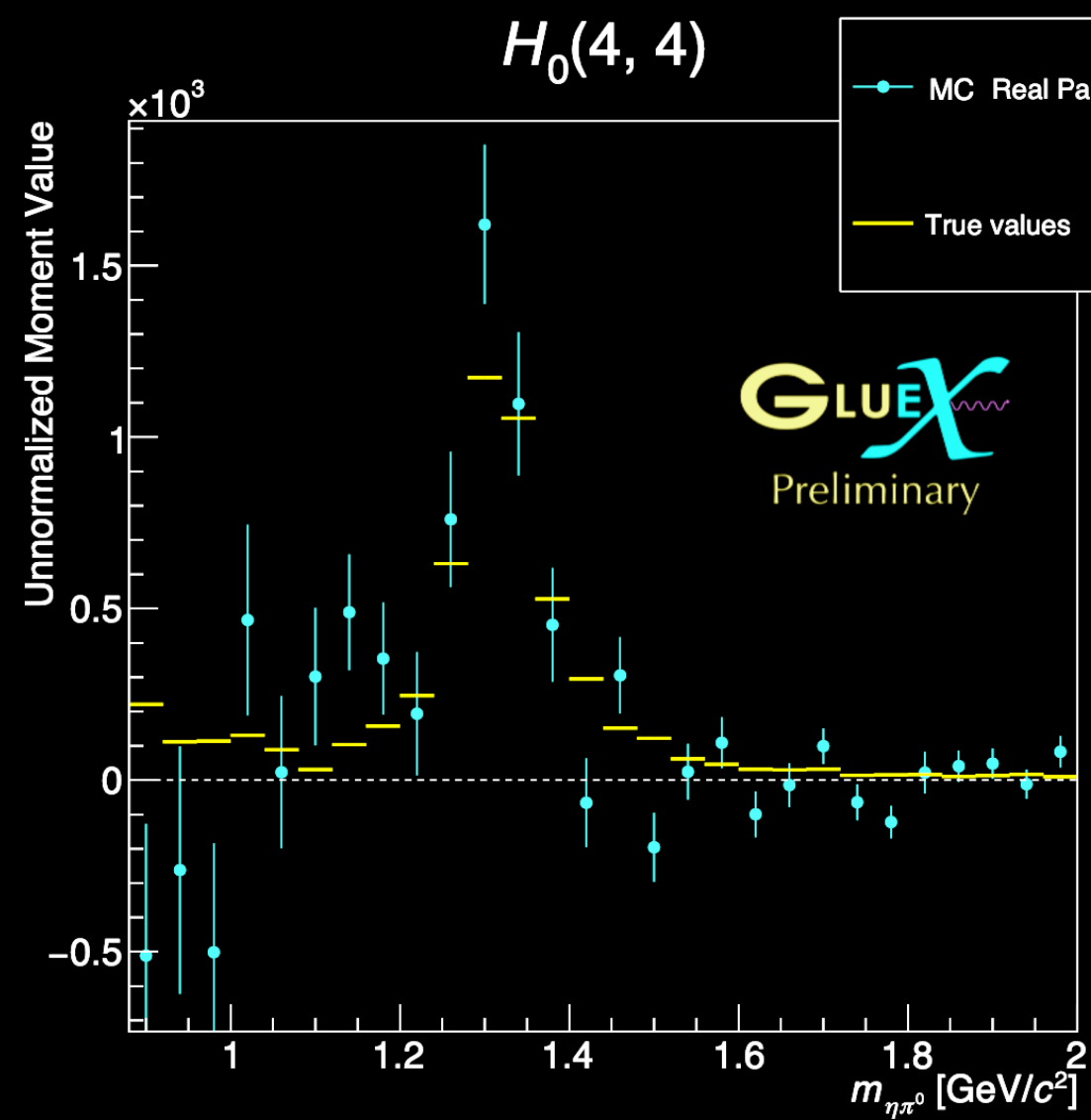
$$\mathcal{I}_0(\Omega) = \sum_{L,M} \sqrt{\frac{2L+1}{4\pi}} H_0(L,M) Y_l^m(\Omega)$$

$$\mathcal{I}_{1,2}(\Omega) = - \sum_{L,M} \sqrt{\frac{2L+1}{4\pi}} H_{1,2}(L,M) Y_l^m(\Omega)$$

$H_0(L,M) \rightarrow$ intensity-like (phase-insensitive)

$H_1(L,M) \rightarrow$ polarization dependent cosine interference

$H_2(L,M) \rightarrow$ polarization dependent sine interference (phase-sensitive)



EXAMPLE: Monte Carlo I/O with simulated $a_2(1320)$

$H_i(4,4)$ contains clear $a_2(1320)$ signal

- partial wave description of the amplitude
Model Dependent

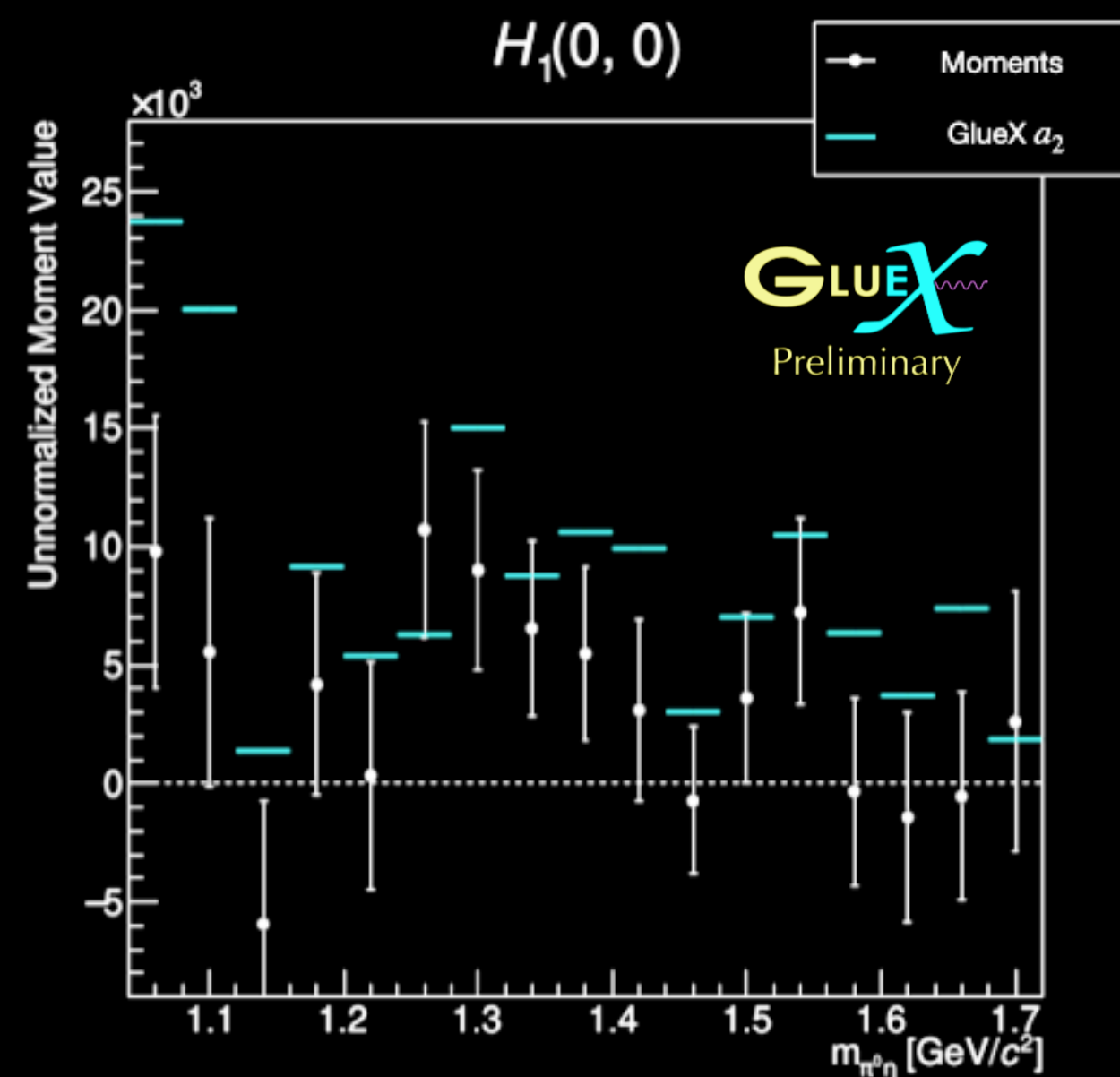
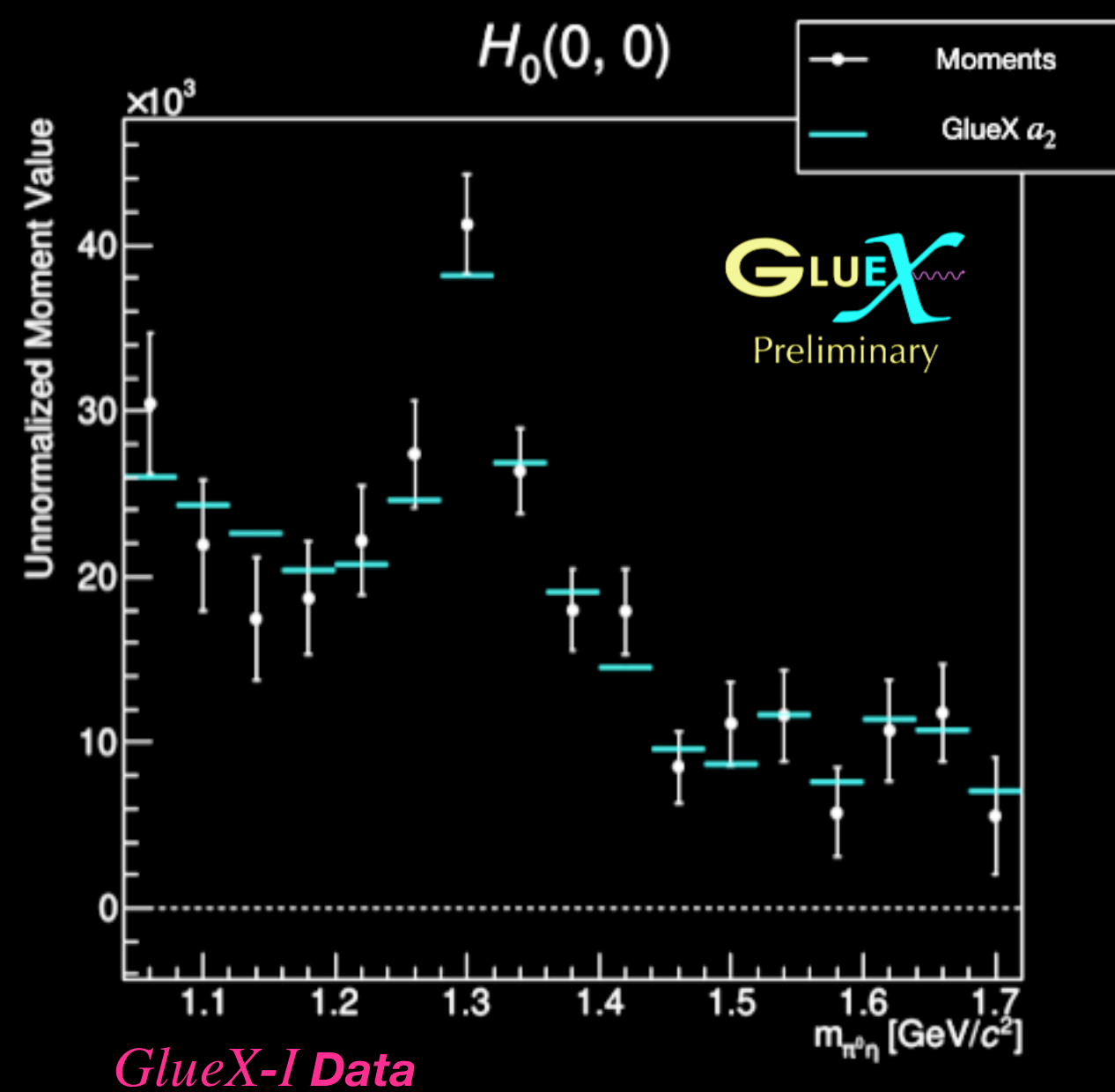
- moment description of the angular intensity
Model Independent

Utilizing the same model used in the first measurement of $a_2(1320)$
polarized photoproduction cross section

F. Afzal et al. [GlueX Collab],
Phys. Rev. C 112, 015204 (2025)

A piecewise ← **GlueX $a_2(1320)$ paper** → $\mathcal{A}_{a_2}^{BW} + \mathcal{A}_{a'_2}^{BW}$

$$\mathcal{A} = \mathcal{A}_{S_{wave}} + \cancel{\mathcal{A}_{P_{wave}}} + \mathcal{A}_{D_{wave}} + \dots$$



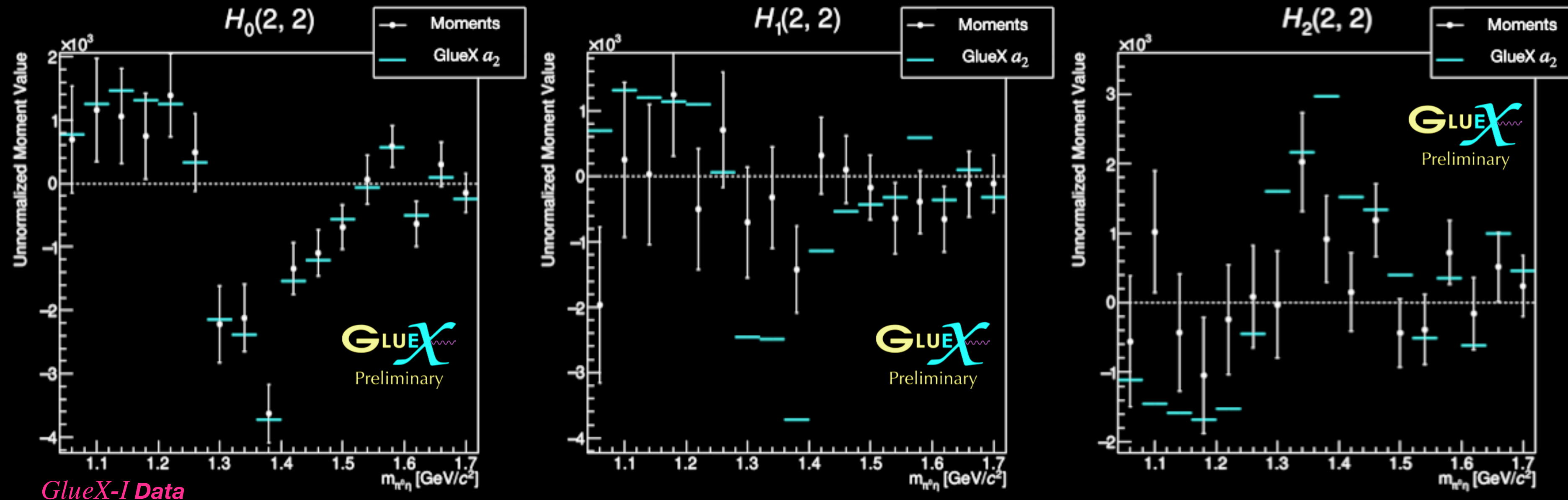
$H_2(0, 0) = 0$ due to symmetries

Utilizing the same model used in the first measurement of $a_2(1320)$
polarized photoproduction cross section

F. Afzal et al. [GlueX Collab],
Phys. Rev. C 112, 015204 (2025)

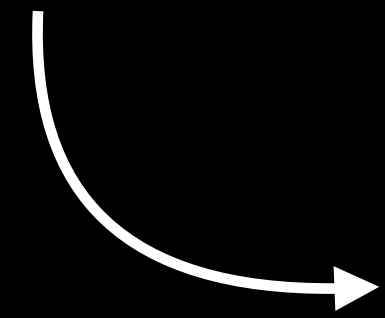
$A^{piecewise}$ \leftarrow **GlueX $a_2(1320)$ paper** $\rightarrow A_{a_2}^{BW} + A_{a'_2}^{BW}$

$$A = A_{S_{wave}} + \cancel{A_{P_{wave}}} + A_{D_{wave}} + \dots$$



GlueX-I Data

S_{wave}



Piecewise

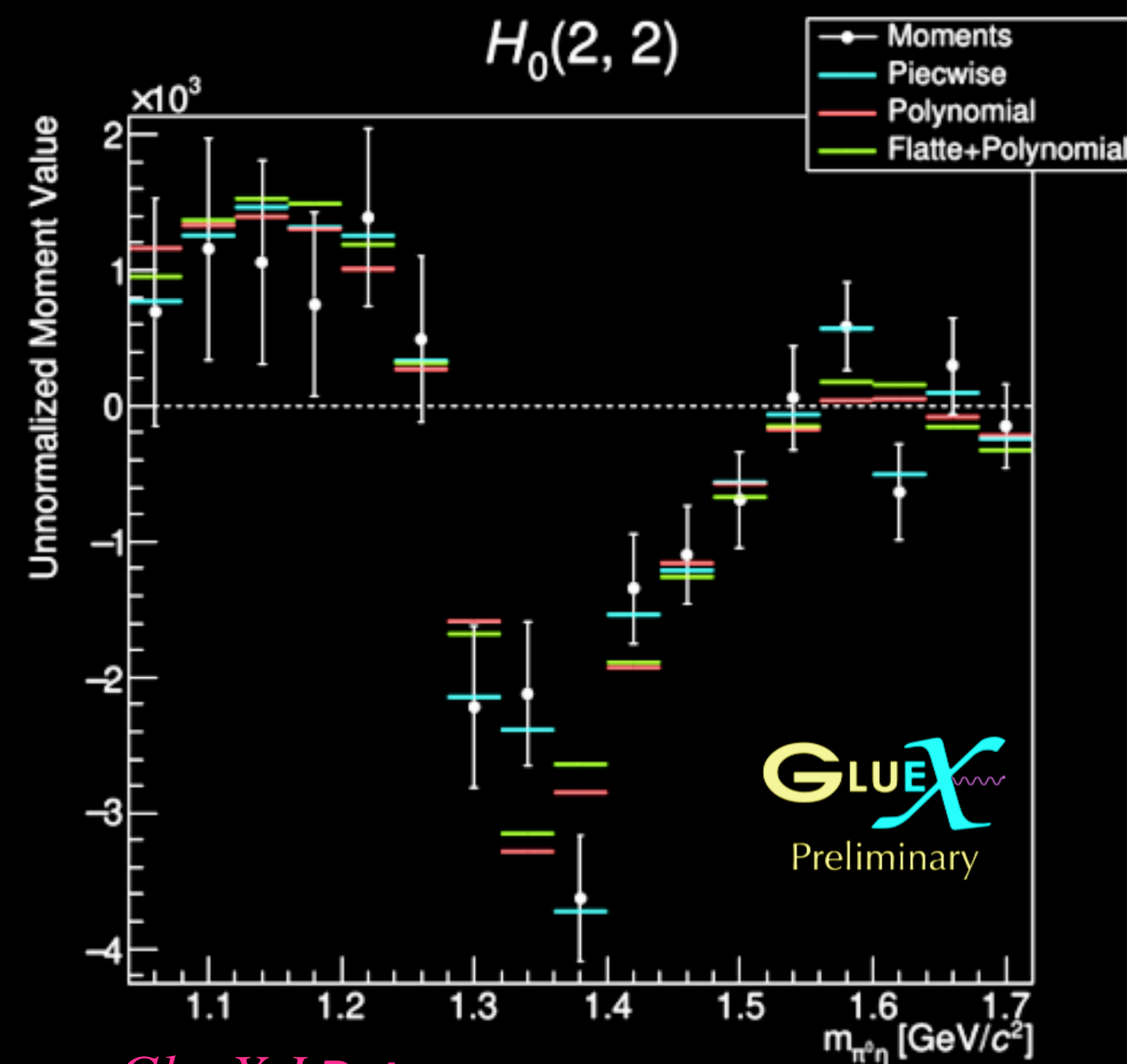
Polynomial

coupled-channel

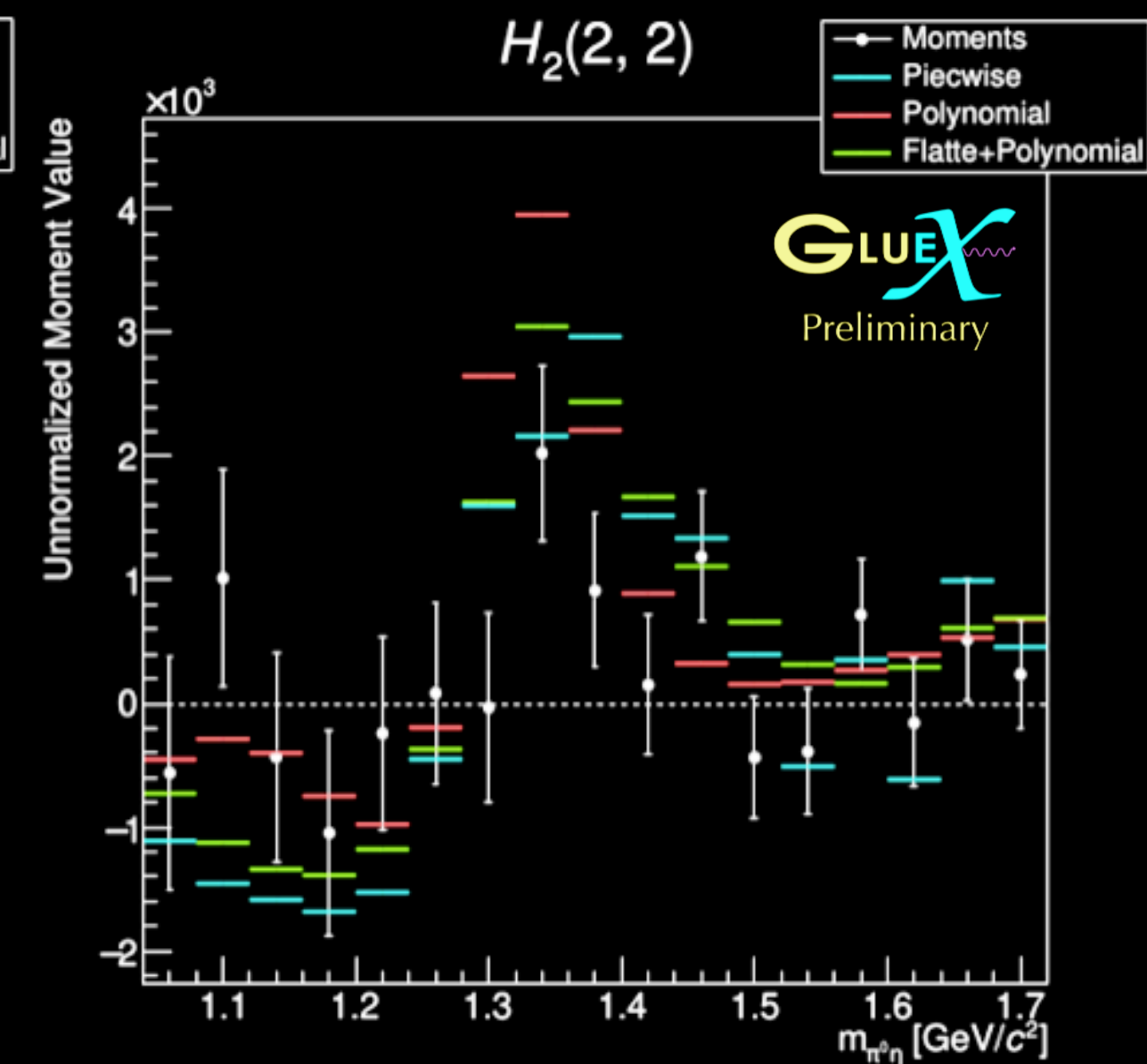
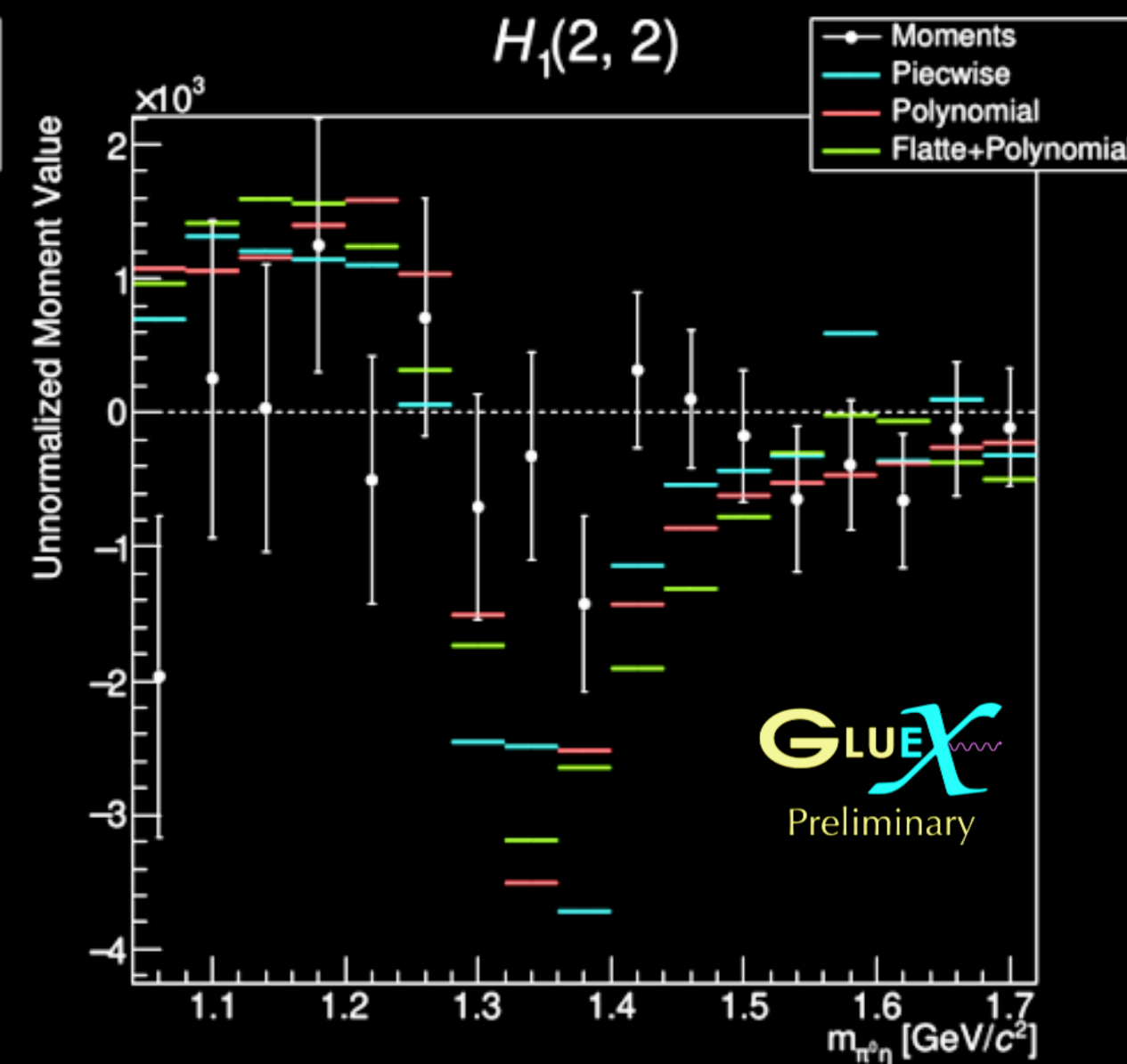
Flatté + Polynomial

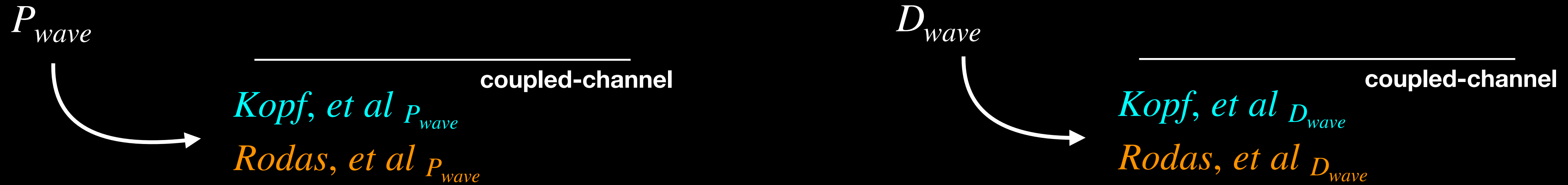
Overall, strong agreement across differing models

- global χ^2 across mass bins roughly equal

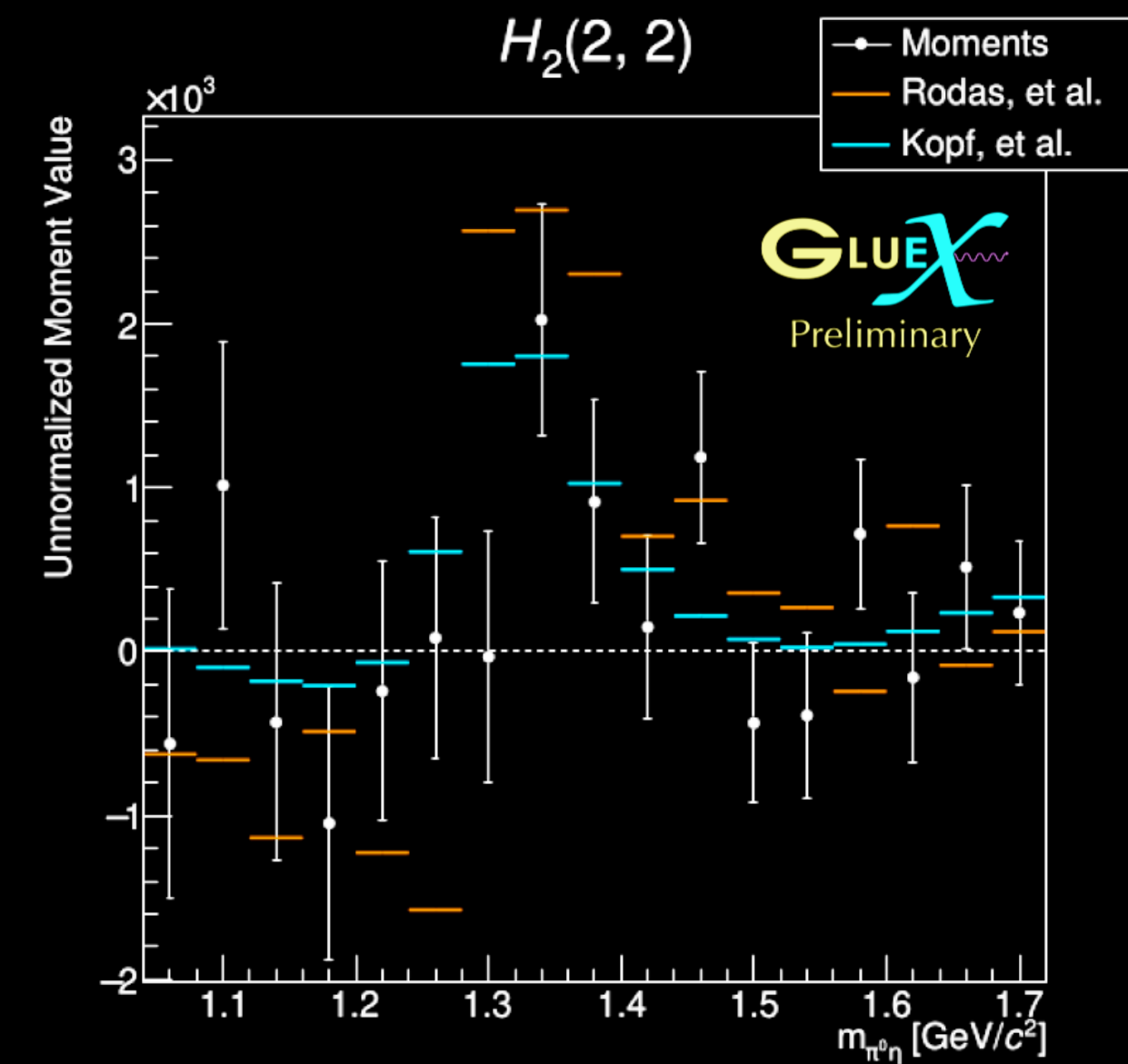
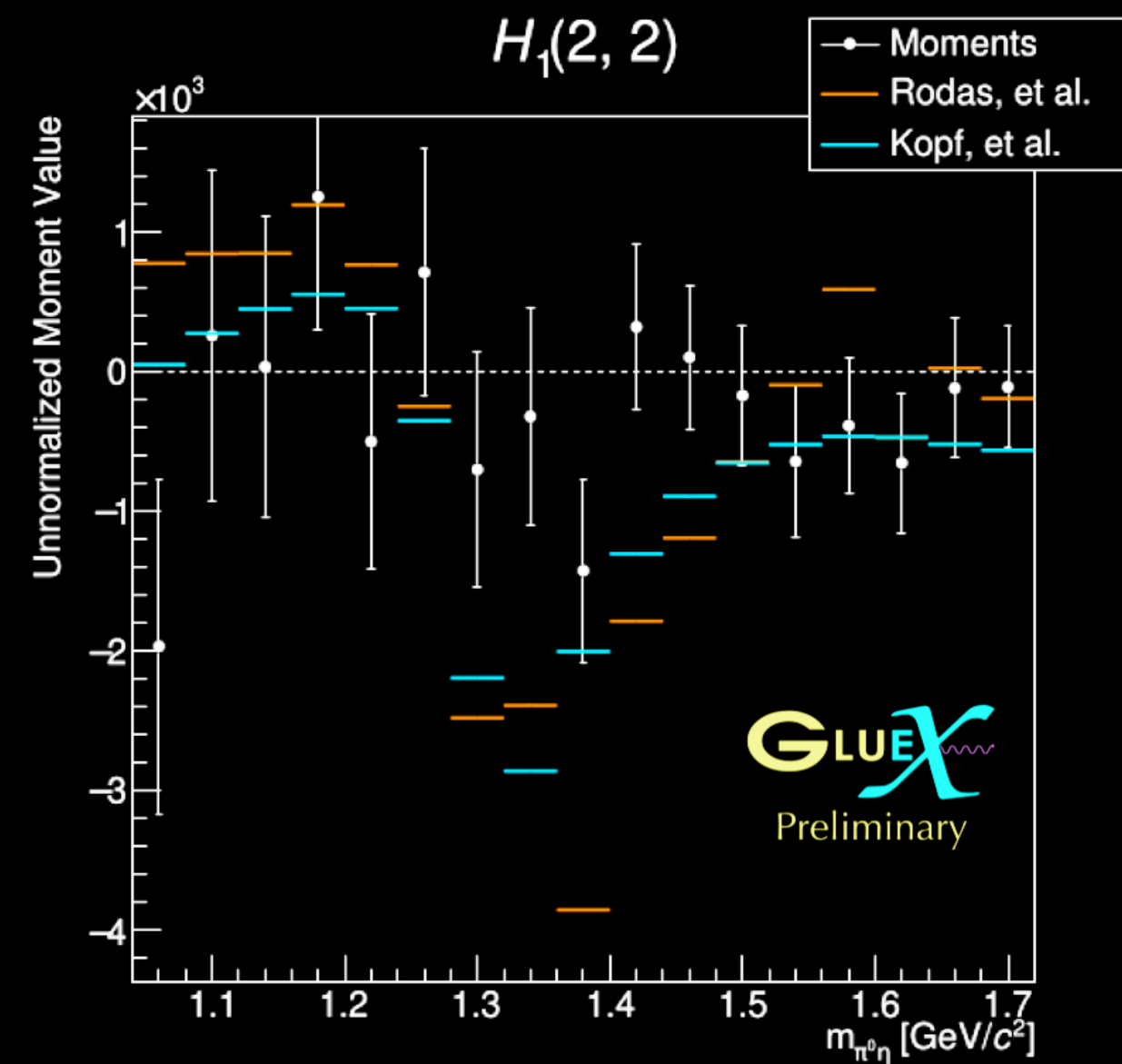
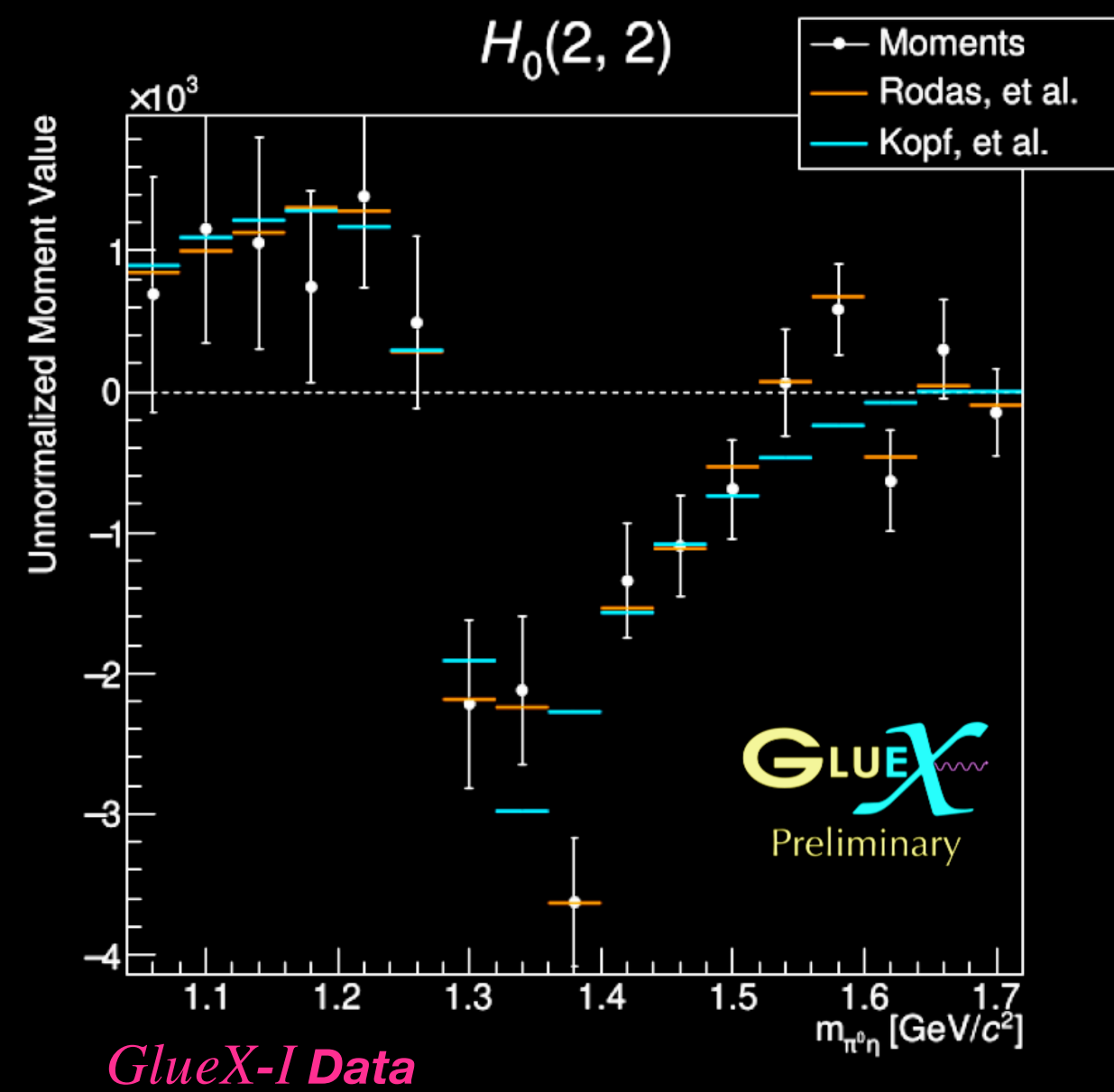


GlueX-I Data





Decent agreement still in $H_i(2, 2)$



GlueX has collected large quantity of photoproduced data

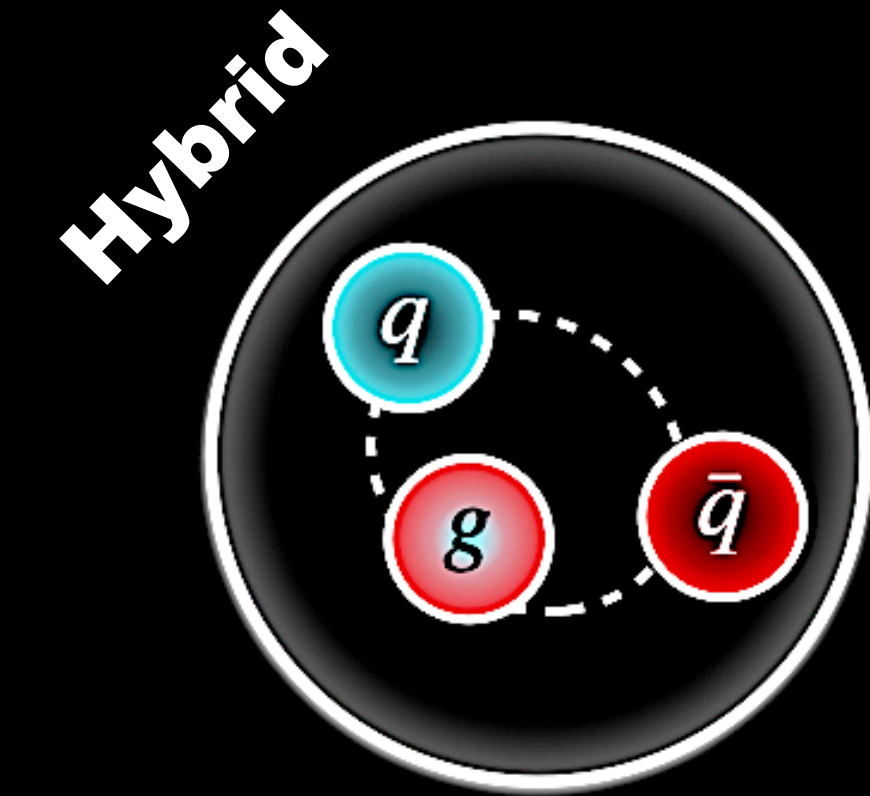
- recent results extracted a_2 cross section which will be used as a reference signal
- can analyze production mechanisms using polarization information
- strong effort to look for *exotic hybrid* π_1 meson in $\eta^{(\prime)}\pi$ systems using several different analysis methods

first look into utilizing coupled channel methods at GlueX

- extraction of moment decomposition from data

Next immediate steps:

- further analyze coupled-channel systematics
- perform model comparison studies



EXCITING TIMES
FOR
EXOTICS SEARCHES
AT
GLUEX

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glueX.org/thanks



Office of Science

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BACKUP SLIDES

Assume $a_2(1320)$ and $a_2(1700)$ are text book Breit-Wigner resonances

- share only 1 common phase parameter for each in the D_{waves}

S_{wave} contributions more complicated

- define *mass independent* piecewise parameterization

Individual fit results across $-t$

- incoherent sums of (+) and (-) reflectivities

$$\eta = P(-1)^J$$

Decent agreement between JPAC TMD model

- the first measurement of the $a_2(1320)$ polarized photoproduction cross section

