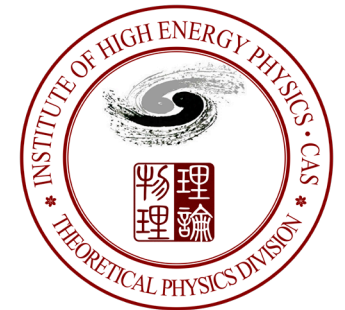


Loop Quantum Gravity Summer School 2026

Quantum Information and Quantum Entanglement

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(中国科学院高能物理研究所)

05/14/2026, Yangzhou University



Outlines

- **Lecture I: Background: quantum gravity meets quantum information**
- **Lecture II: Basic concepts in quantum information**
- **Lecture III: Quantum entanglement in tensor networks**
- **Lecture IV: Quantum entanglement and the structure of spacetime**

Acknowledgement to collaborators:

Peng Liu, Yu-Xuan Liu, Chao Niu, Jianpin Wu, Meng-he Wu, Zhuoyu Xian, Yikang Xiao

Lecture IV

Quantum entanglement and the structure of spacetime

- Entanglement creates geometric connection
- Tensor networks and AdS space
- Emergence of space from entanglement by spin networks

Entanglement creates geometric connection

Gedanken experiment

2013, Maldacena and Susskind *Cool horizons for entangled black holes*

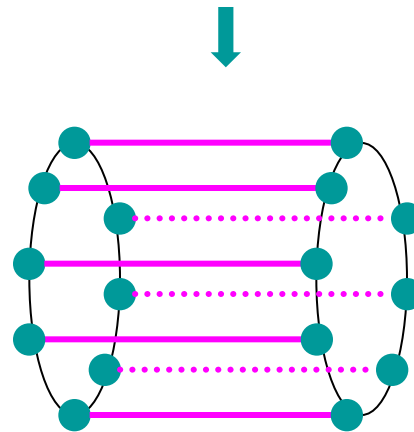
arXiv:1306.0533

EPR state

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|\uparrow_A \downarrow_B\rangle + |\downarrow_A \uparrow_B\rangle)$$

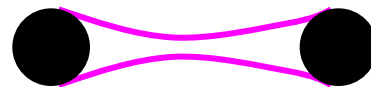


N EPR pairs

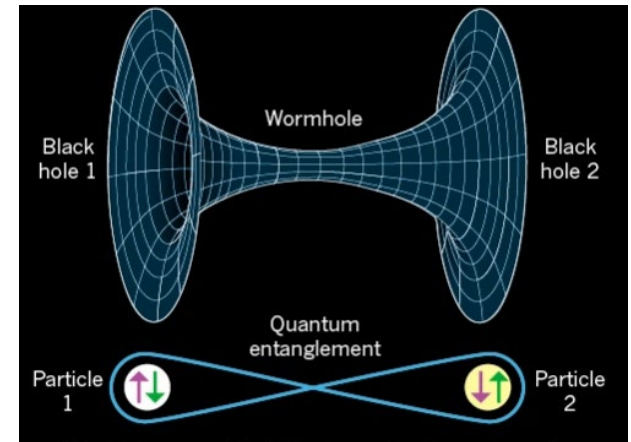


Increase N to form two black holes

ER bridge



ER=EPR



Entanglement creates geometric connection

2001, Maldacena

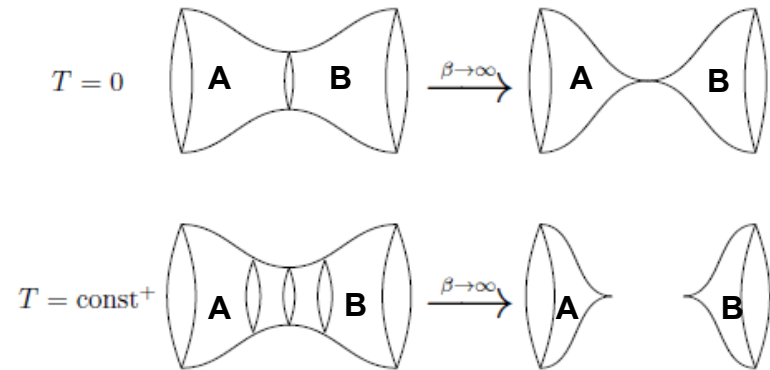
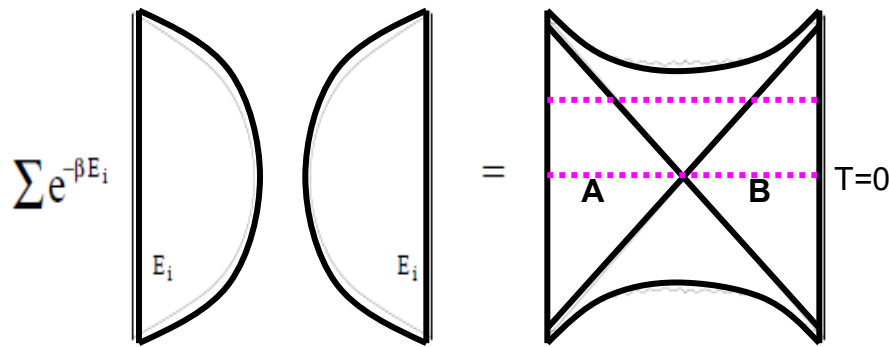
Eternal black holes in anti-de Sitter

arXiv:hep-th/0106112

2010, Van Raamsdonk

Building up spacetime with quantum entanglement

arXiv:1005.3035



TFD state

$$|\psi\rangle = \frac{1}{\sqrt{Z}} \sum_i e^{-\frac{\beta E_i}{2}} |\psi_i\rangle_A \otimes |\psi_i\rangle_B$$

$$\langle\psi|\psi\rangle = 1 \quad \rho = |\psi\rangle\langle\psi|$$

$$\begin{aligned} \rho_A &= \text{Tr}_B \rho_{AB} = \langle\psi_i|\psi\rangle\langle\psi|\psi_i\rangle_B \\ &= \frac{1}{Z} \sum_i e^{-\beta E_i} |\psi_i\rangle_A \langle\psi_i| = \rho_B \end{aligned}$$

$$|\psi\rangle \xrightarrow{\beta \rightarrow \infty} |\Phi\rangle$$

$$|\Phi\rangle = |\psi_0\rangle_A \otimes |\psi_0\rangle_B$$

Spacetime perhaps is just a geometric picture for the entanglement style of matter in a quantum system.

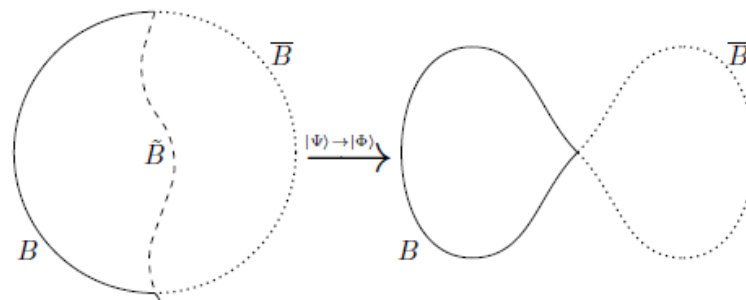
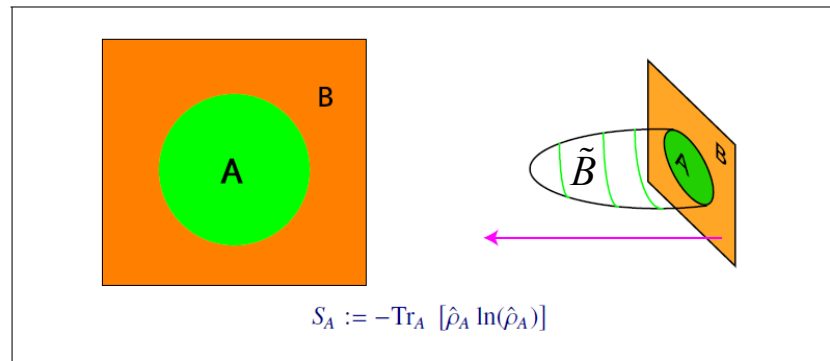
Entanglement creates geometric connection

2006, Ryu and Takayanagi

Holographic derivation of entanglement entropy from AdS/CFT

arXiv:hep-th/0603001

$$S_E = \frac{\text{Area}(\tilde{B})}{4G_N}$$



$$|\psi\rangle = \sum_{i,j} p_{i,j} |\phi_i^B\rangle \otimes |\phi_j^{\bar{B}}\rangle \longrightarrow |\Phi\rangle = \left(\sum_i c_i |\phi_i^B\rangle\right) \otimes \left(\sum_j d_j |\phi_j^{\bar{B}}\rangle\right)$$

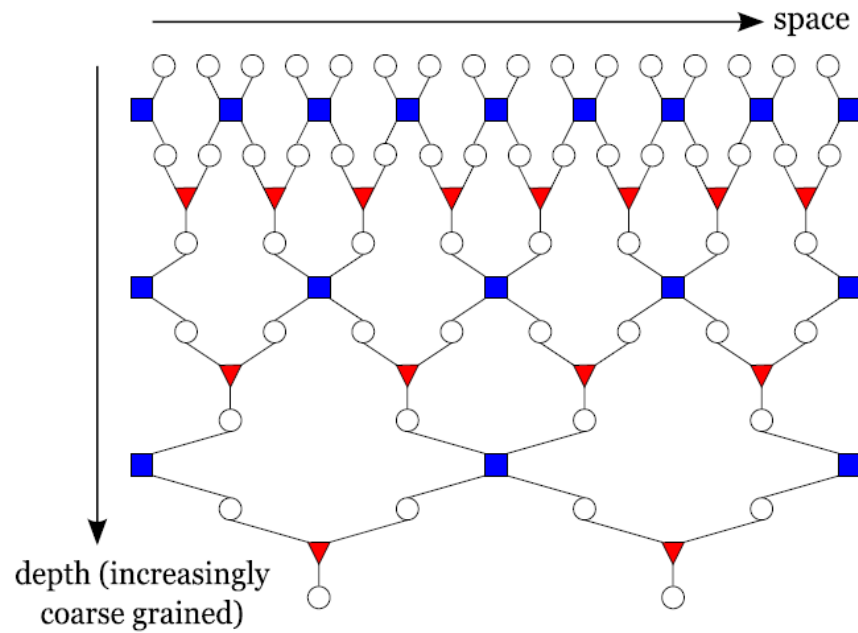
arXiv:1711.10854

Entanglement creates geometric connection

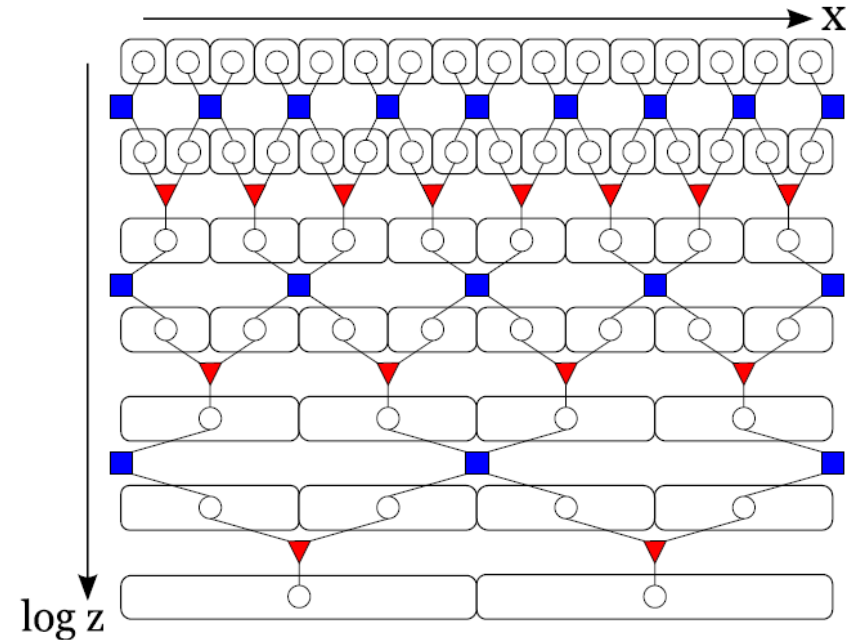
2009, Swingle

Entanglement Renormalization and Holography

arXiv:0905.1317



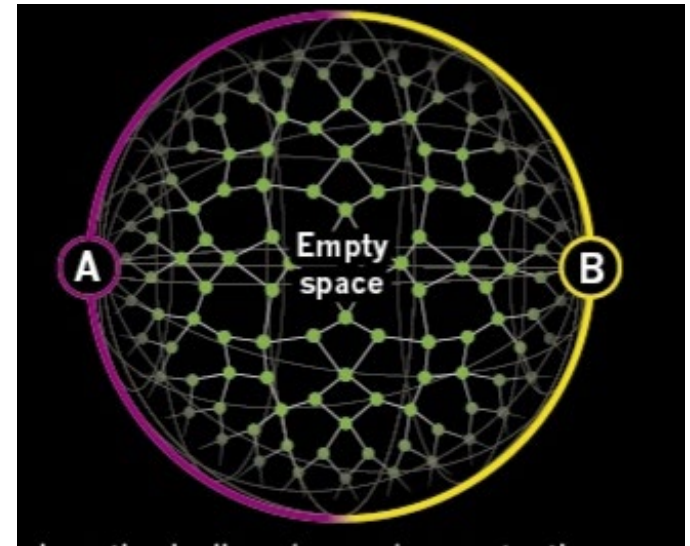
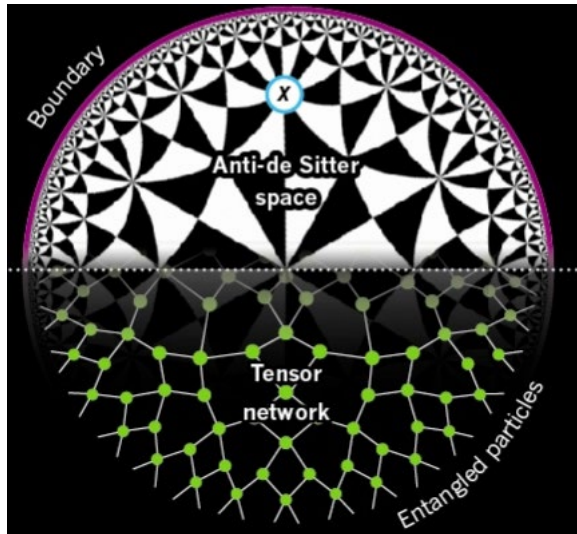
$$dS \sim \frac{L^{d_s-1}}{z^{d_s-1}} \frac{dz}{z}$$



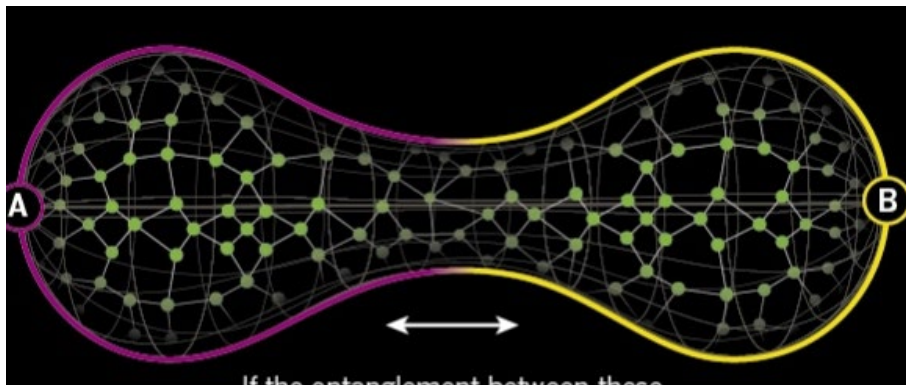
$$ds^2 = R^2 \left(\frac{dz^2 + dx^2}{z^2} \right) = R^2 (dw^2 + e^{-2w} dx^2)$$

Entanglement creates geometric connection

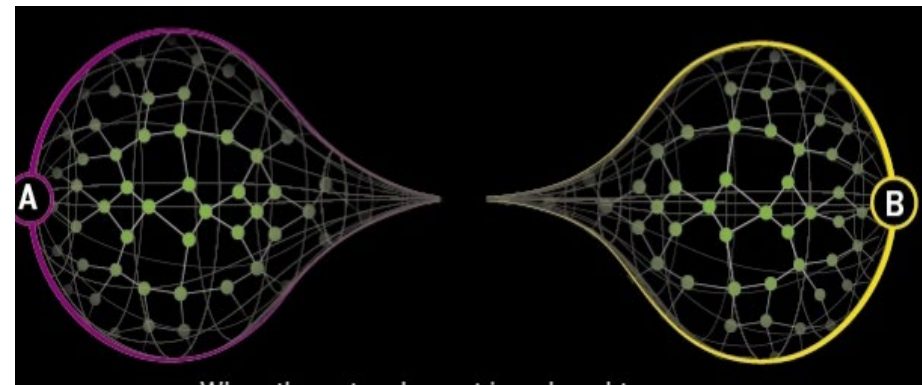
- The microscopic structure of space time by tensor networks



Space is precisely constructed through quantum entanglement via tensor networks



If the entanglement between these



When the entanglement is reduced to zero

Entanglement creates geometric connection

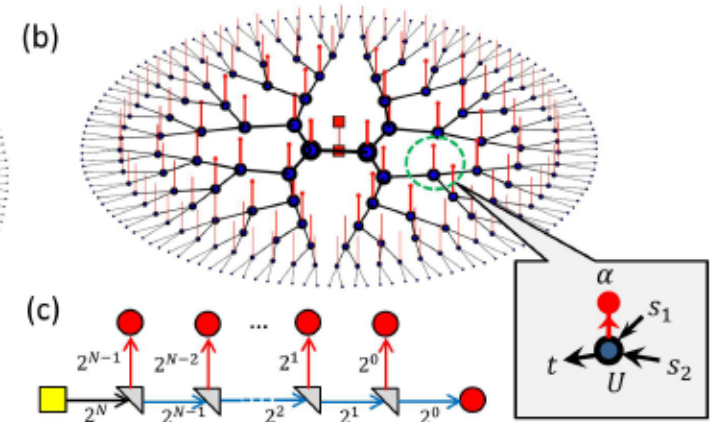
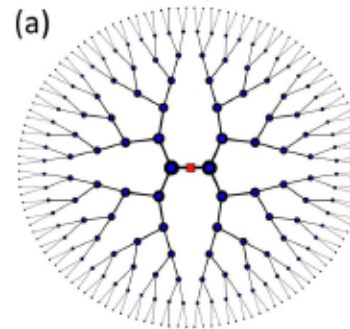
2013, X. Qi

Exact holographic mapping and emergent space-time geometry

arXiv:1309.6282

$$d_{(x,t_1)(y,t_2)} = -\xi \log \frac{\langle O_x(t_1) O_y(t_2) \rangle}{C_0}$$

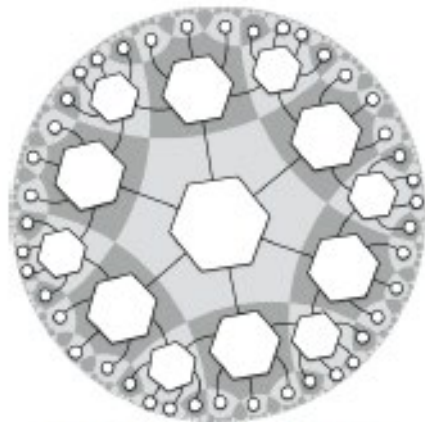
$$d_{xy} = -\xi \log \frac{I_{xy}}{I_0}$$



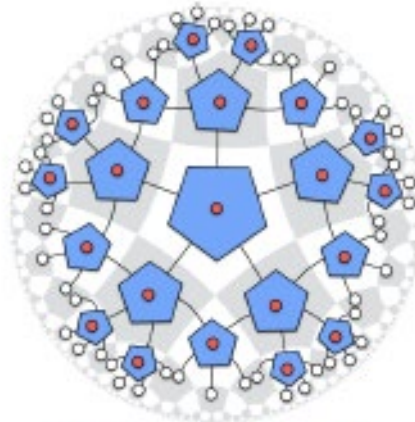
2015, Harlow and Preskill

*Holographic quantum error-correcting codes:
Toy models for the bulk/boundary correspondence*

arXiv:1503.06237



(a) Holographic hexagon state



(b) Holographic pentagon code

Holographic duality itself is a QEC

Perfect tensor (PT)

Entanglement creates geometric connection

2015, Miyaji et.al.

cMERA as Surface/State Correspondence in AdS/CFT

arXiv:1506.01353

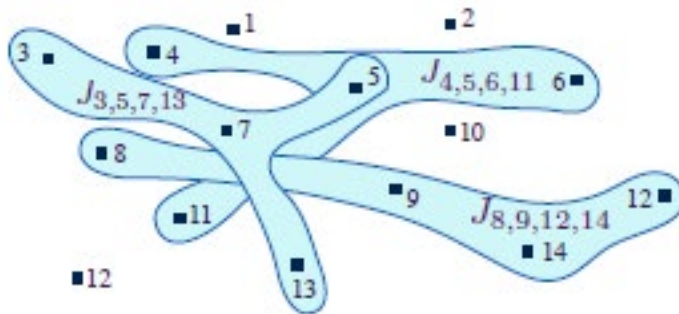
2015, Kitaev, Sachdev

A simple model of quantum holography - 2015

arXiv:1506.05111

Bekenstein-Hawking Entropy and Strange Metals

Sachdev-Ye-Kitaev Model



Einstein-Maxwell theory
+ cosmological constant

Horizon area \mathcal{A}_h ;
 $\text{AdS}_2 \times R^d$
 $ds^2 = (d\zeta^2 - dt^2)/\zeta^2 + d\vec{x}^2$
Gauge field: $A = (\mathcal{E}/\zeta)dt$

Boundary
area \mathcal{A}_b ;
charge
density \mathcal{Q}

$\zeta = \infty$

ζ

\vec{x}

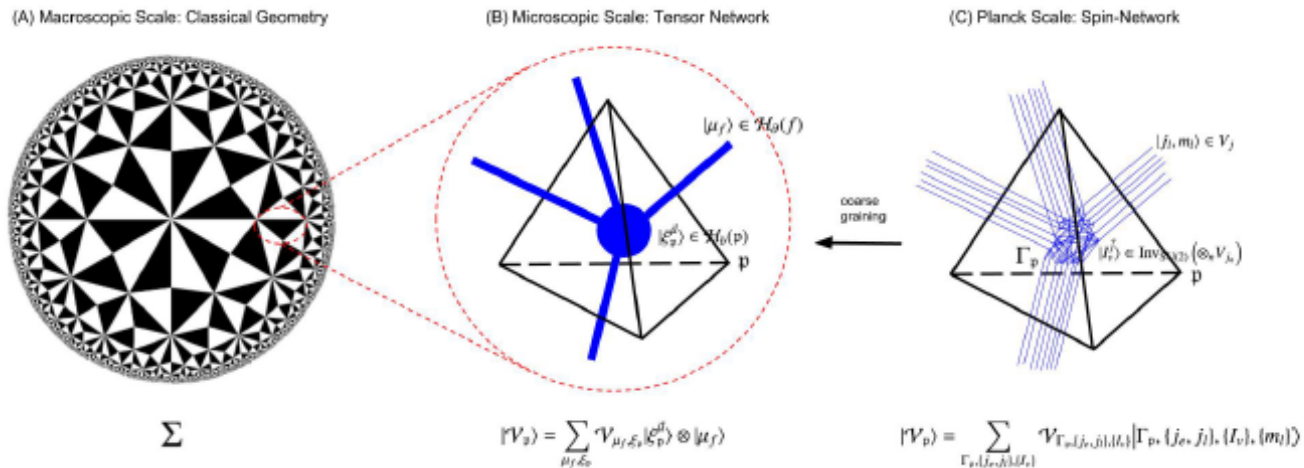
$$\mathcal{L} = \bar{\psi} \Gamma^\alpha D_\alpha \psi + m \bar{\psi} \psi$$

Entanglement creates geometric connection

2016, Han and Hung

*Loop Quantum Gravity, Exact Holographic Mapping,
and Holographic Entanglement Entropy*

arXiv:1610.02134



2019, A. Almheiri, N. Engelhardt,
D. Marolf, H. Maxfield

*The entropy of bulk quantum fields and the entanglement wedge
of an evaporating black hole*

arXiv: 1905.08762

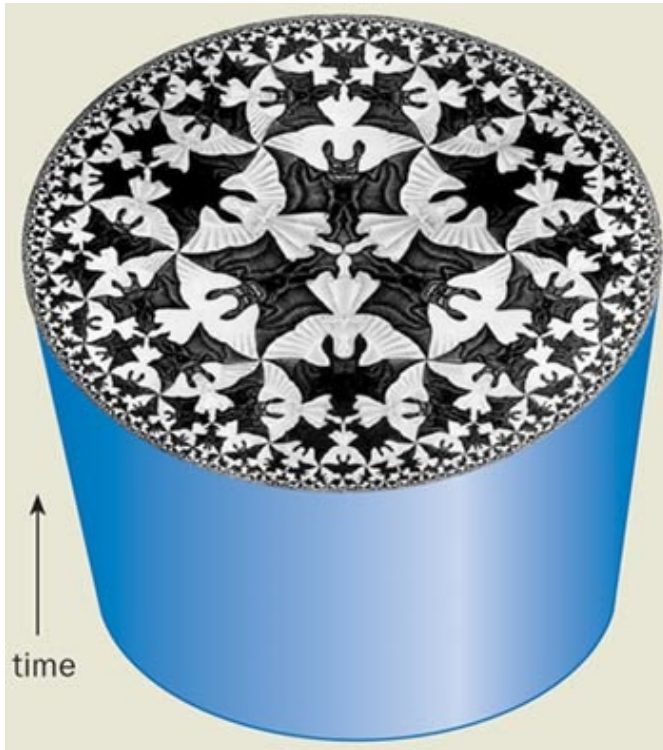
2019, G. Penington

*Entanglement Wedge Reconstruction and the Information
Paradox*

arXiv: 1905.08255

Tensor Networks and AdS space

■ Poincare patch of AdS space



Angels & Devils by M.C. Escher

Hyperbolic structure of H_2 space

$$ds^2 = L^2 \left(\frac{dz^2 + dx^2}{z^2} \right)$$

$$\xi = x + iz \quad L = 1$$

$$ds^2 = - \frac{d\xi d\xi^*}{(\xi - \xi^*)^2}$$

Isometry: $SL(2, R)$

$$\xi' = \frac{\alpha\xi + \beta}{\gamma\xi + \delta} \quad \alpha\delta - \beta\gamma = 1$$

Tensor Networks and AdS space

■ Two remarkable holographic features of AdS space

1. For the vacuum in AdS/CFT correspondence, Renyi entropy satisfies Cardy-Calabrese formula and a **non-flat ES** is inherent.

$$S_n = \left(\frac{n+1}{n}\right) \frac{A_{\min}}{8G}$$

$$S_n = \frac{1}{1-n} \ln \frac{\text{Tr} \rho_A^n}{(\text{Tr} \rho_A)^n}$$

Flat ES

$$\rho_A^2 = a \rho_A$$



$$S_n = \ln \frac{\text{Tr} \rho}{a}$$

2. A local operator in the bulk can be reconstructed in a subsystem A on the boundary. It can be viewed as the accomplishment of **QEC**.

Tensor Networks and AdS space

What kind of **tensor networks** could capture the holographic features of AdS?

1. **MERA** breaks the isometry $SL(2, \mathbb{R})$ and has a **preferred direction**, implying that QEC can not be realized along all directions.
2. **Perfect tensor** networks has a **flat** ES and trivial connected correlation functions, which is not a reflection of the holographic property of AdS.

Tensor Networks and AdS space

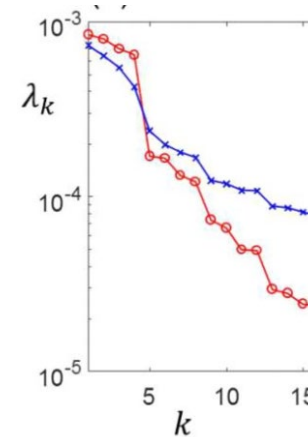
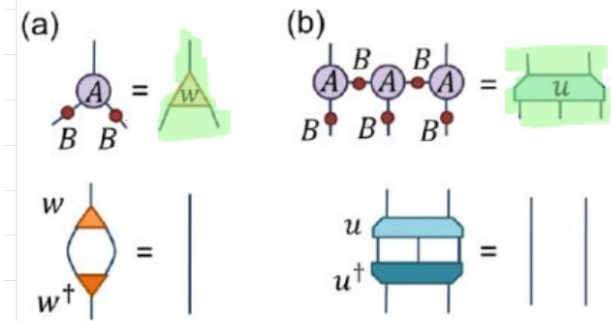
arXiv:1704.04229

■ Hyperinvariant Tensor Networks (HTN):

2017, Evenly

Hyperinvariant Tensor Networks and holography

(3,7) HTN isometries

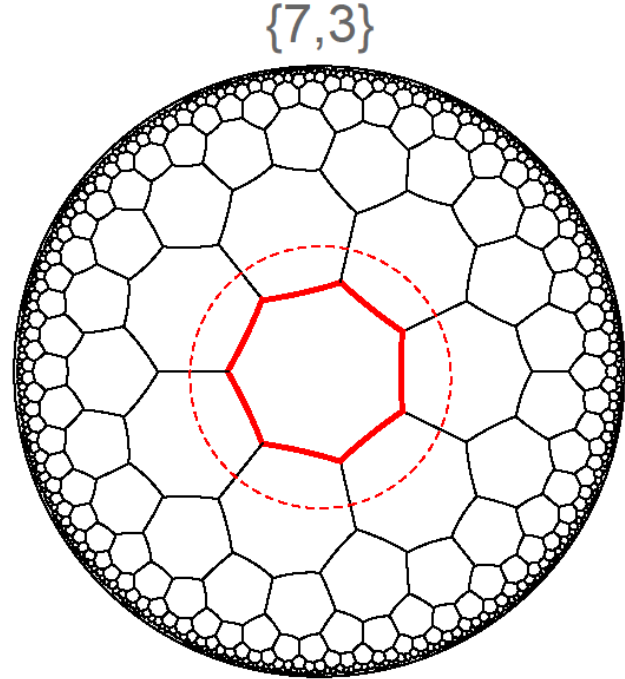
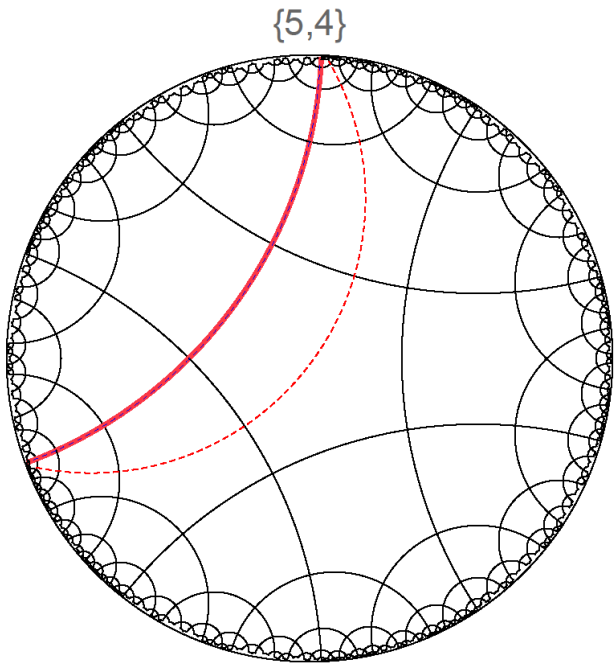


1. What kind of multi-tensor constraints could endow desirable features of AdS spacetime to a given tensor network?
2. Is there any criteria to justify the ability of QEC and the non-flatness of ES for a tensor network with given constraints?

Tensor Networks and AdS space

$$\frac{1}{a} + \frac{1}{b} < \frac{1}{2}$$

- The (b,a) tiling of hyperbolic space



$$T_{ijkl} = T_{jkli} = T_{klij} = T_{lijk}$$

$$T_{ijk} = T_{jki} = T_{kij}$$

Tensor Networks and AdS space

Isometry

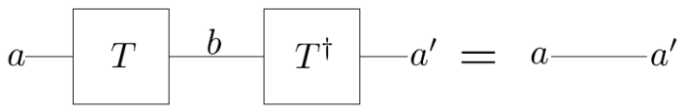


Figure 1. Diagrammatic tensor notation, here showing that T is an isometry.

$$\sum_b T_{ab} T^\dagger_{ba'} = \delta_{aa'}$$

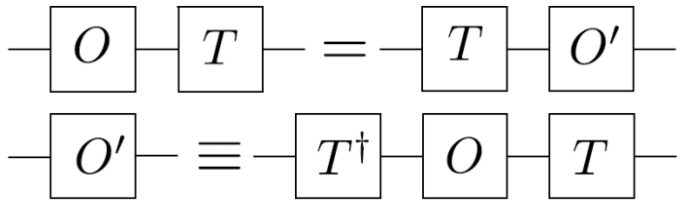
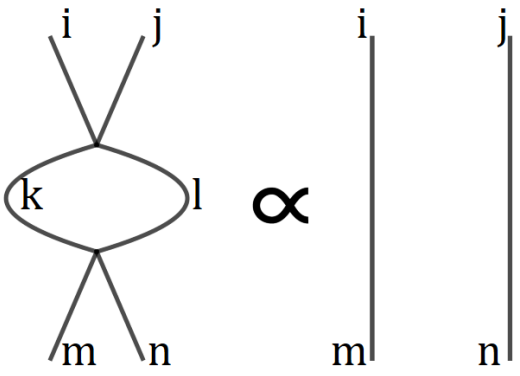


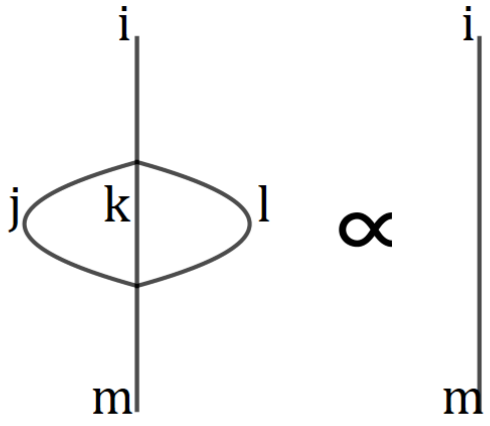
Figure 2. Operator pushing through an isometric tensor.

$$OT = TT^\dagger OT = T(T^\dagger OT) = TO'$$

Tensor constraints



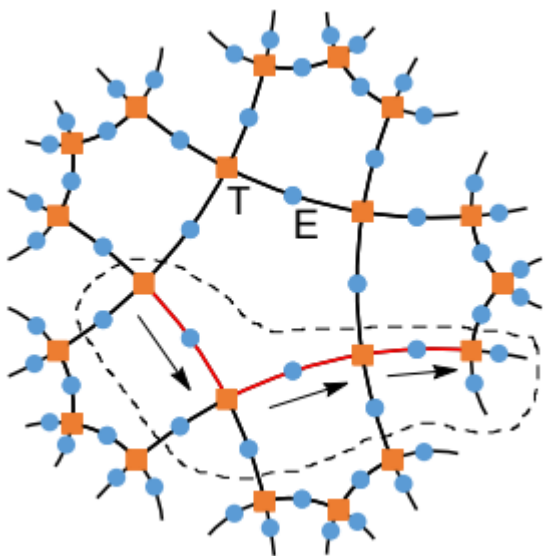
$$\sum_{kl} T_{ijkl} T^*_{mnkl} \propto \delta_{im} \delta_{jn}$$



$$\sum_{jkl} T_{ijkl} T^*_{mjkl} \propto \delta_{im}$$

Tensor Networks and AdS space

- The (b,a) tiling of hyperbolic space



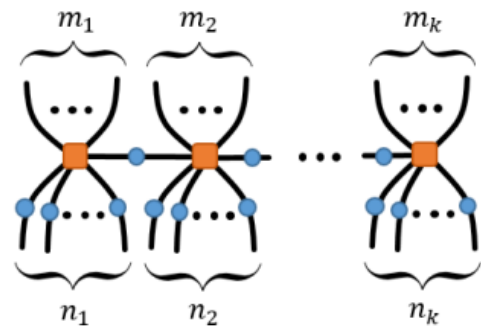
{5, 4}

$$T_{ijkl} = T_{jkli} = T_{klij} = T_{lijk}$$

$$E_{ij} = E_{ji}$$

$$\frac{1}{a} + \frac{1}{b} < \frac{1}{2}$$

- Tensor Chains



$$m_i + n_i = a - 2 + \delta_{i1} + \delta_{ik}$$

$k = 4$

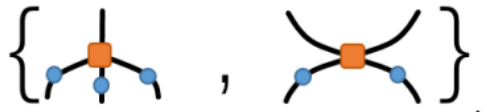
$(m_1, m_2, m_3, m_4) = (2, 0, 1, 1)$

$(n_1, n_2, n_3, n_4) = (1, 2, 1, 2)$

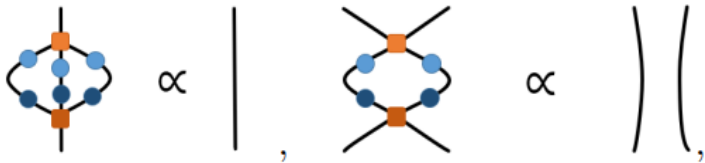
- Average induced interior angle

$$\kappa = \frac{1}{k} \left(\sum_{i=1}^k m_i + k - 1 \right)$$

Tensor Networks and AdS space



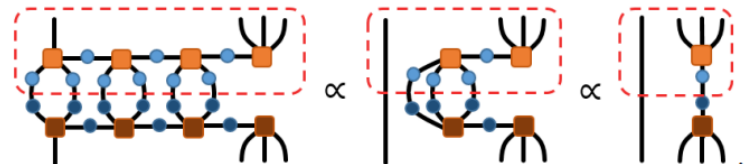
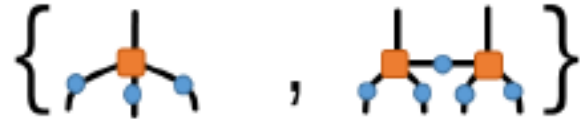
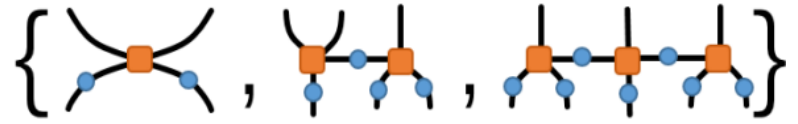
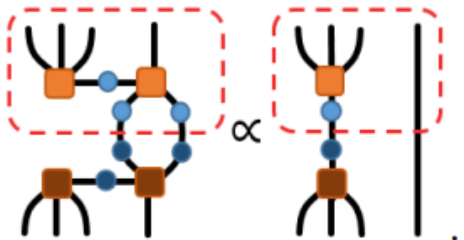
- Tensor constraints



- A set of isometry

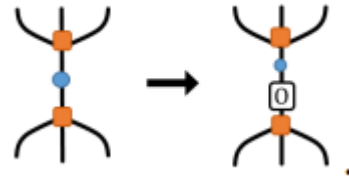


- Greedy algorithm

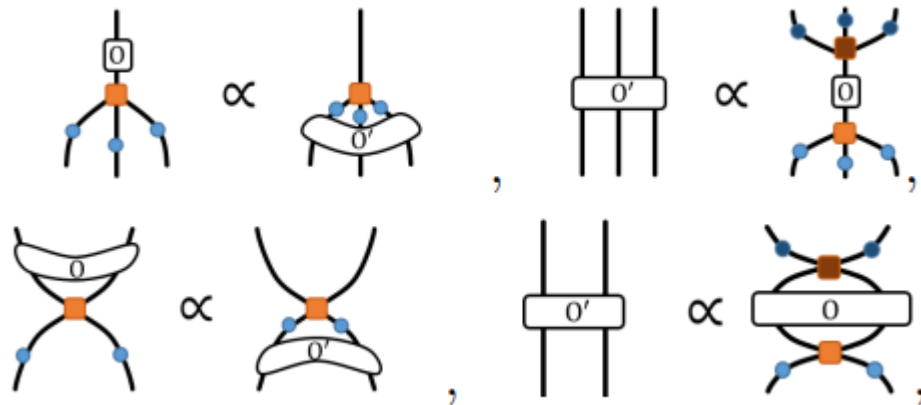


Tensor Networks and AdS space

■ QEC in a tensor network

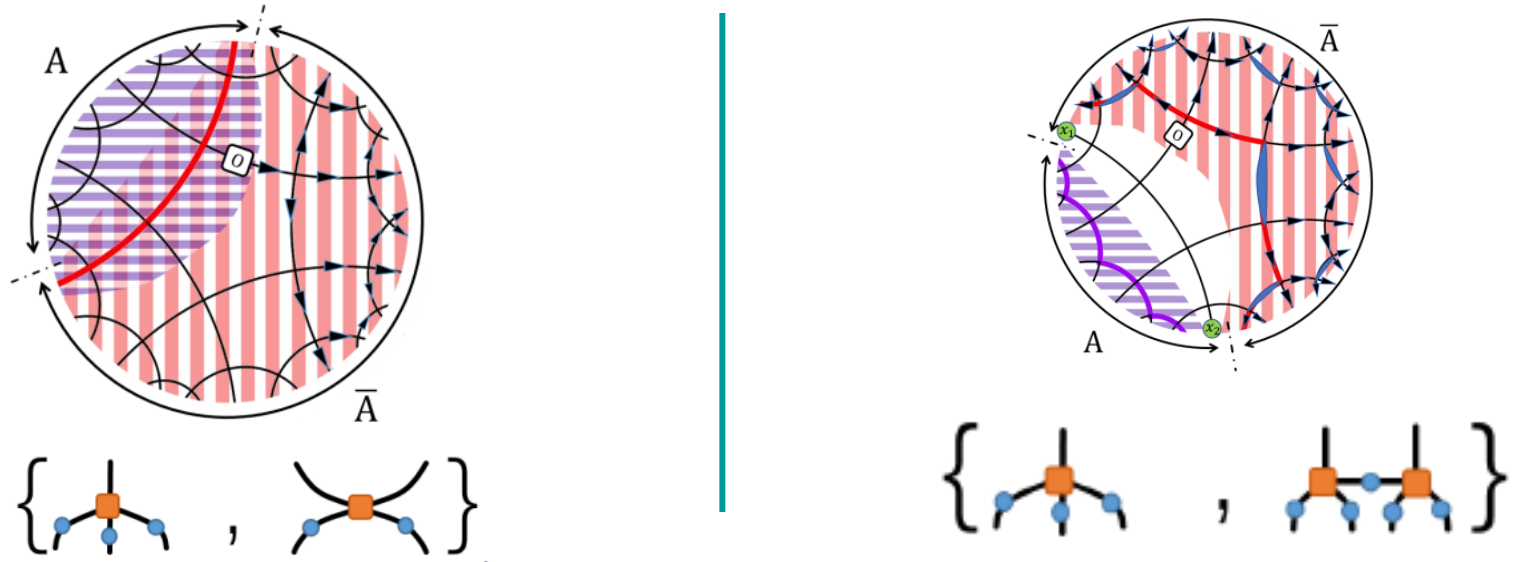


• Tensor pushing



Tensor Networks and AdS space

- QEC in networks



- The failure of QEC in a network

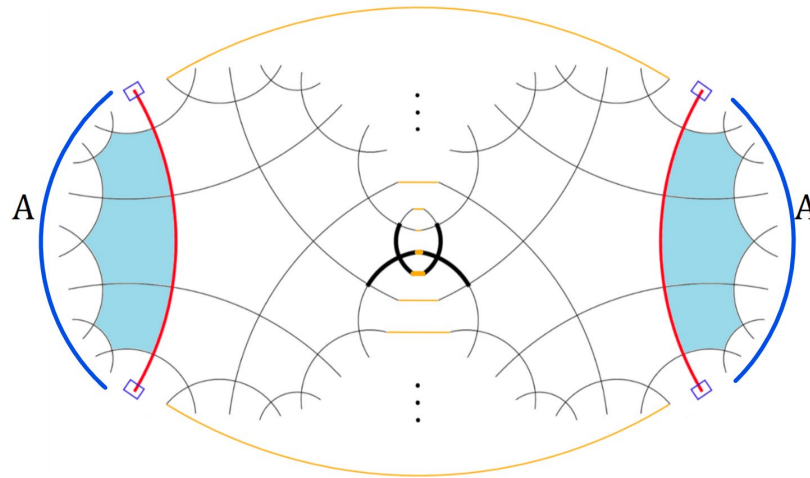


Tensor Networks and AdS space

■ The entanglement spectrum of a tensor network

- The reduced density matrix for a network

$$\rho_A = \text{Tr}_{\bar{A}} |\psi\rangle\langle\psi| = \psi\psi^\dagger$$



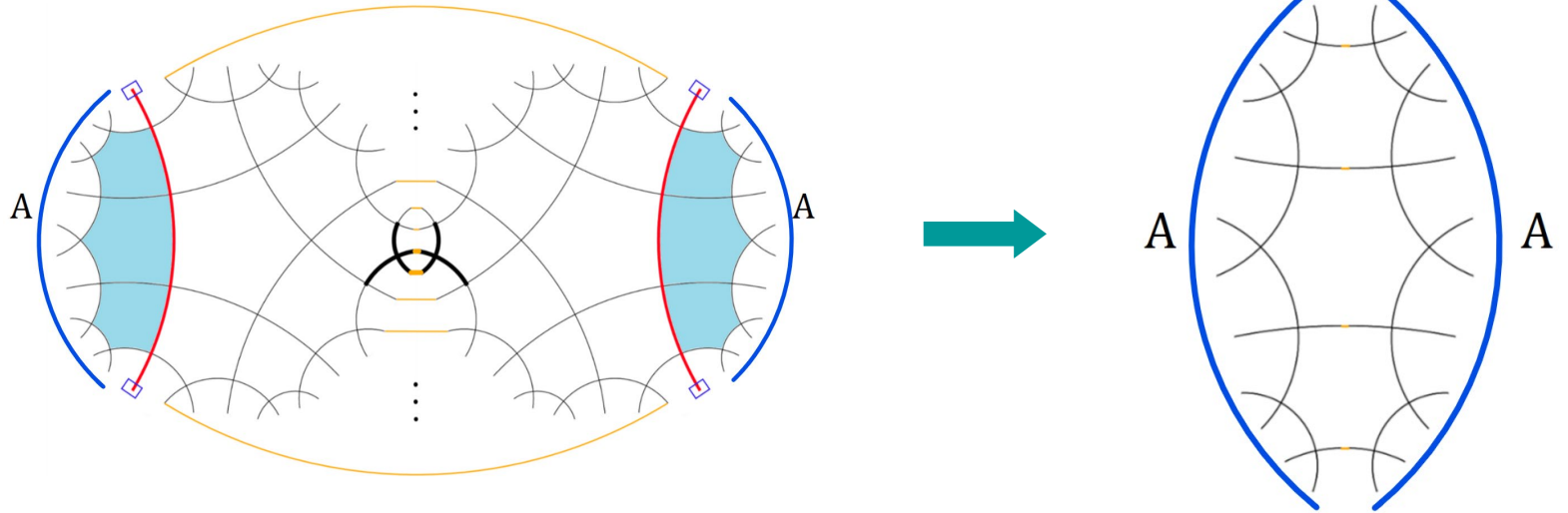
The **flatness** of ES is equivalent to

$$\boxed{\rho_A^2 = a\rho_A} \quad \longrightarrow \quad \rho_A\rho_A = \psi\psi^\dagger\psi\psi^\dagger \propto \psi\psi^\dagger = \rho_A$$

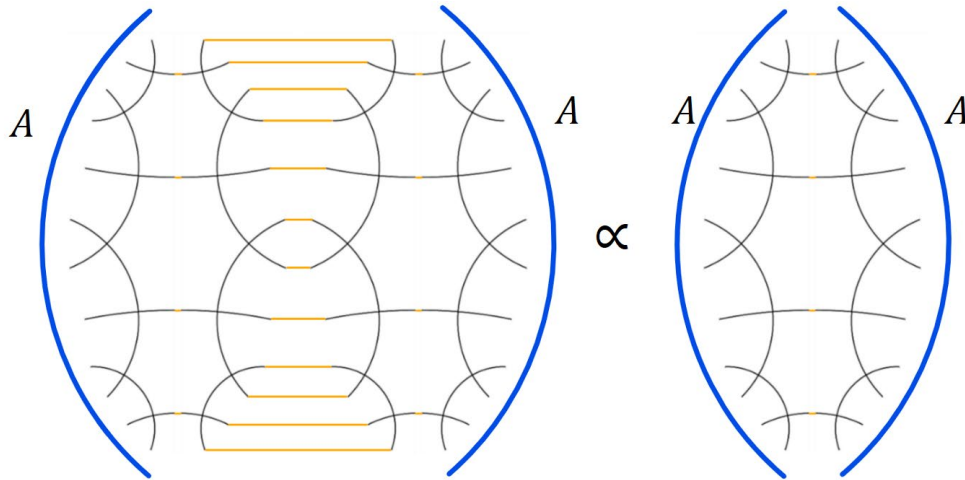
Tensor Networks and AdS space

- Single tensor constraint leads to flat ES

$$\rho_A$$



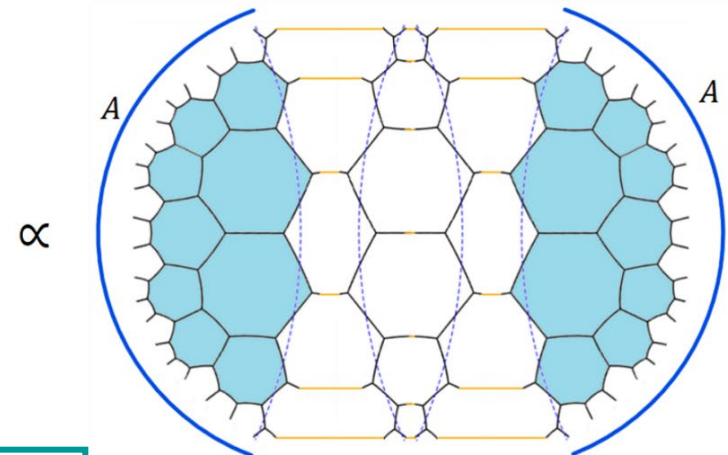
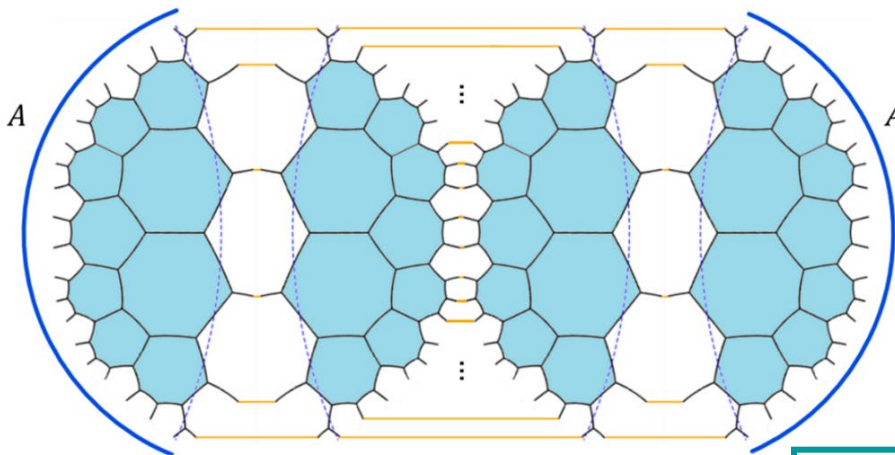
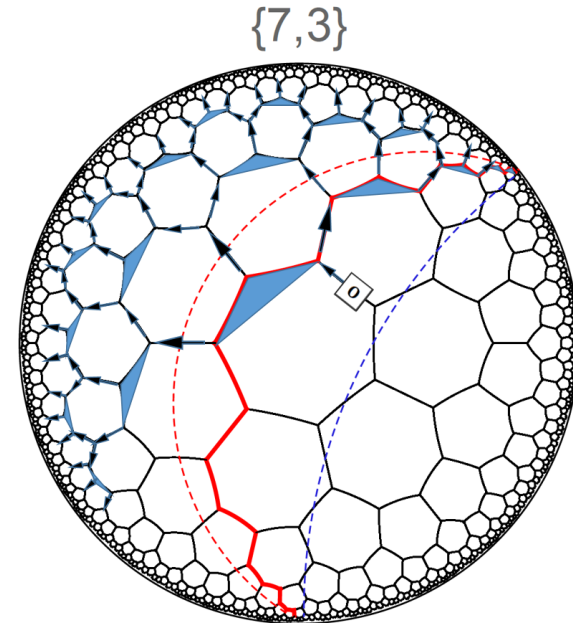
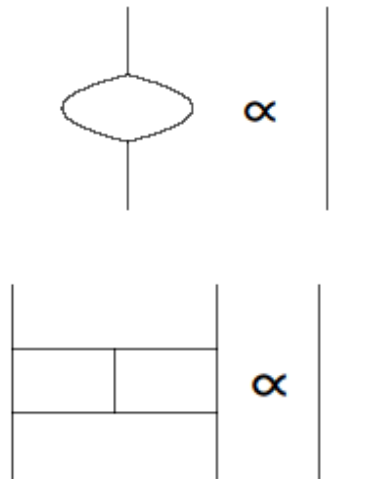
$$\rho_A^2$$



$$\rho_A^2 \propto \rho_A$$

Tensor Networks and AdS space

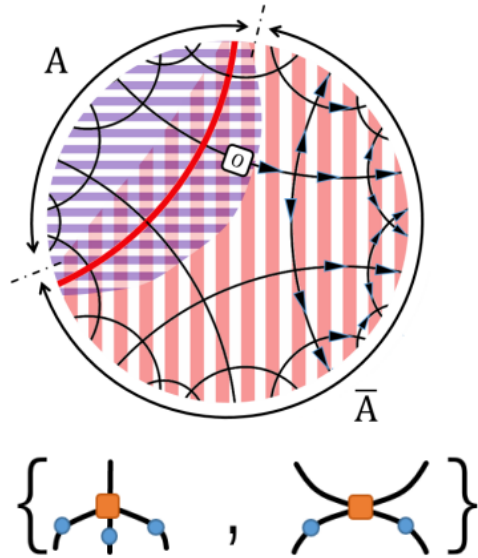
- Multi-tensor constraints and non-flat ES



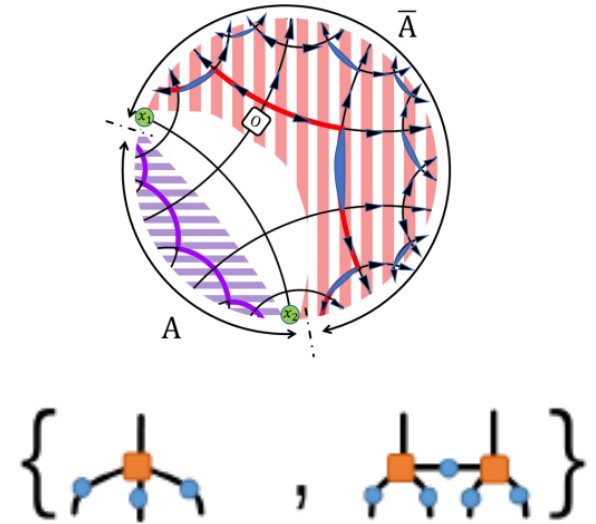
$$\rho_A^2 \neq \rho_A$$

Tensor Networks and AdS space

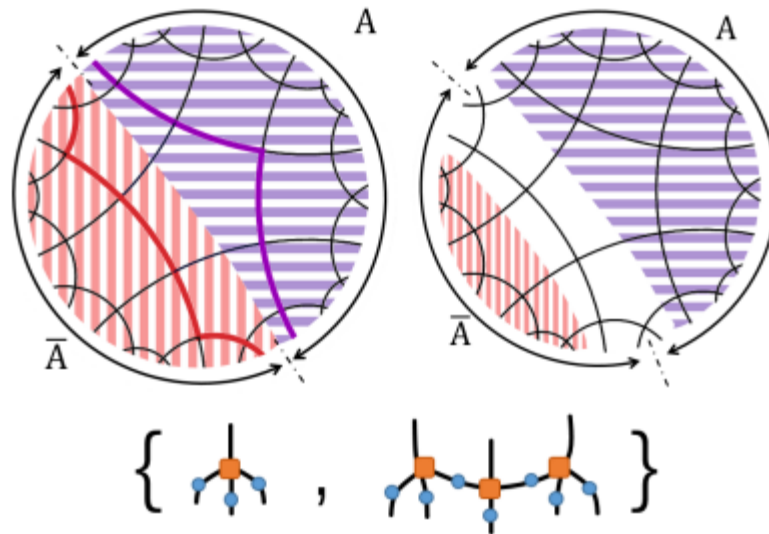
- A flat ES



- A non-flat ES

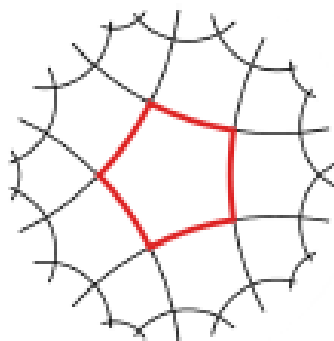


- A mixed ES

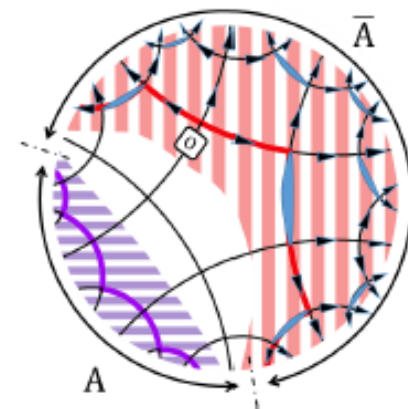
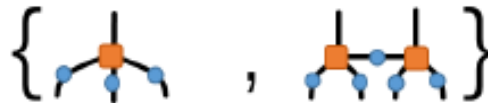


Tensor Networks and AdS space

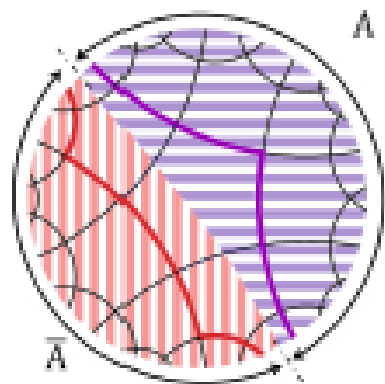
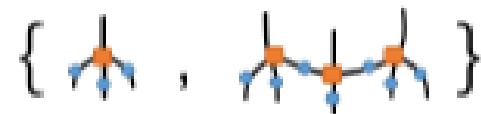
- κ_c of CP tensor chain



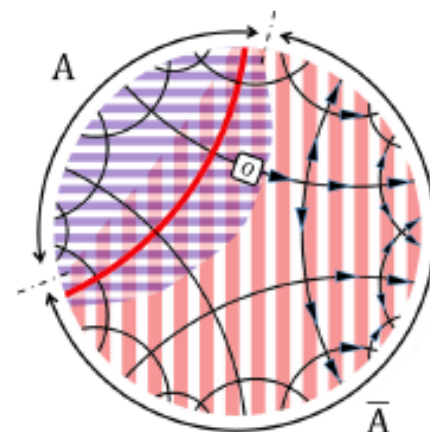
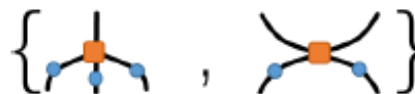
$$\kappa_c = 1$$



$$\kappa_c = 3/2$$



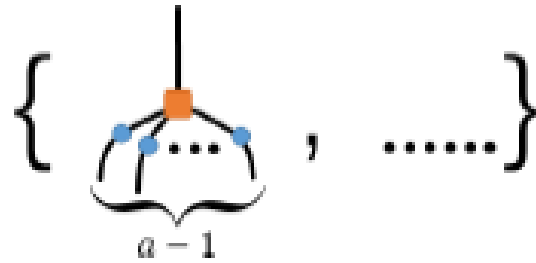
$$\kappa_c = 5/3$$



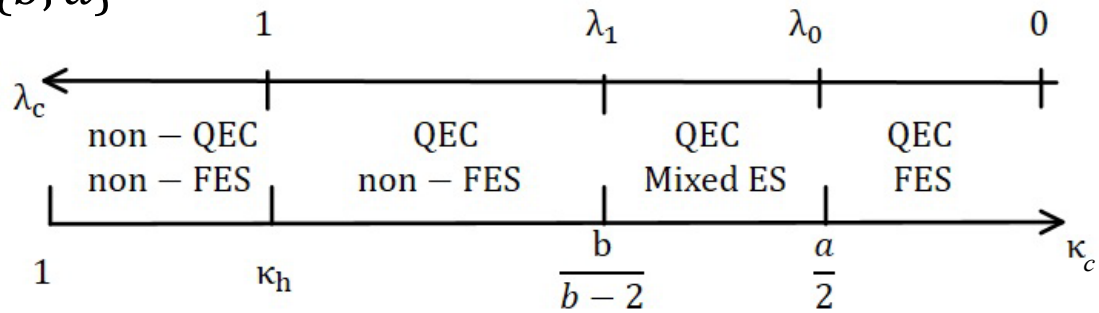
$$\kappa_c = 2$$

Tensor Networks and AdS space

- General constraints $\{b, a\}$



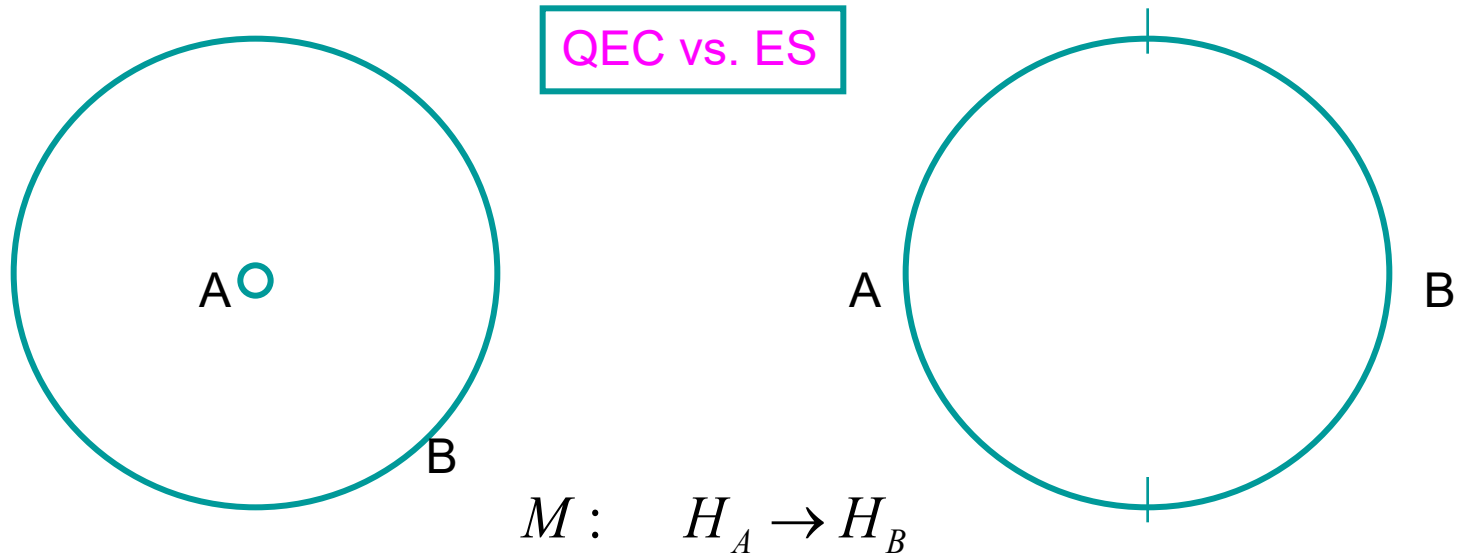
Classification $\{b, a\}$



$$\kappa_h = \frac{a}{2} - \frac{a-2}{2} \sqrt{\frac{ab - 2a - 2b}{ab - 2a - 2b + 4}}$$

$\{5,4\}$ $\kappa_h = 1.42,$ $\frac{b}{b-2} = \frac{5}{3},$ $\frac{a}{2} = 2$

Summary:



QEC is better



ES more easily becomes flat!

Angles vs. Devils

AdS is fantastic!

Emergence of space from entanglement by spin networks



Roger Penrose

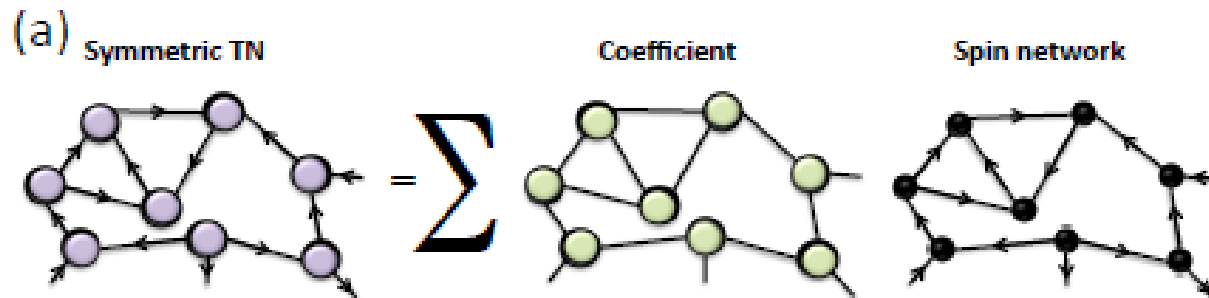


Guifre Vidal

2014 Orus

*Advances on Tensor Network Theory:
Symmetries, Fermions, Entanglement, and holography*

ArXiv:1407.6552



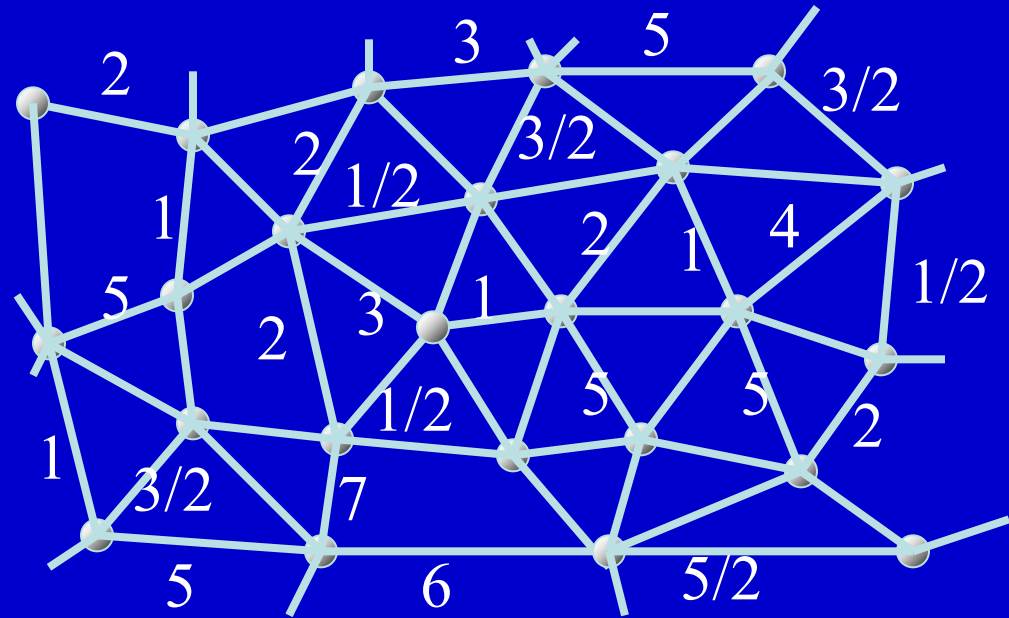
Spin networks and quantum geometry

Spin network states

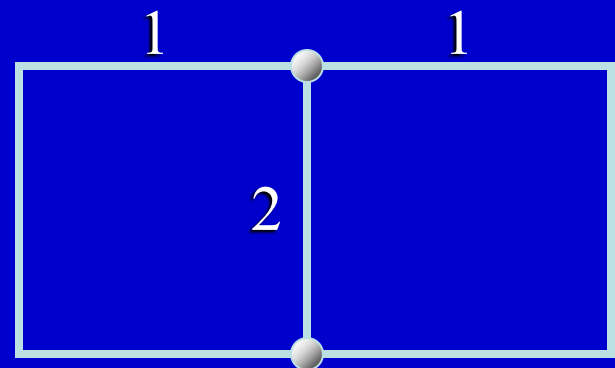
$$|\Gamma, j_m, I_v\rangle$$

$$j = \frac{1}{2}, 1, \frac{3}{2}, 2, \dots$$

$$|n, l, m\rangle$$



Γ

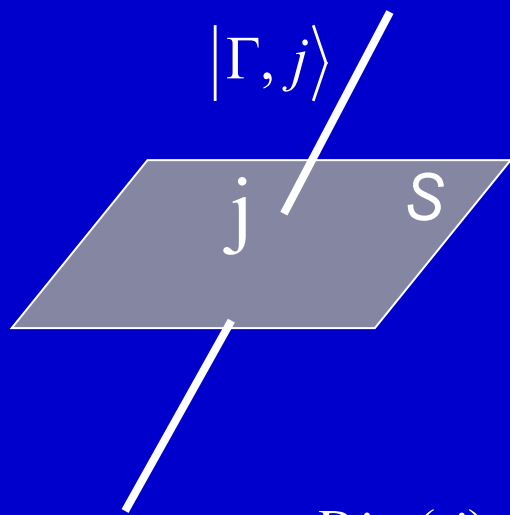


Spin networks and quantum geometry

■ Area operator

$$\begin{aligned} A(\mathbf{S}) &= \int_{\mathcal{S}} d\sigma^2 \sqrt{\det({}^2h)} \\ &= \int_{\mathcal{S}} d\sigma^2 \sqrt{\hat{E}_a^i \hat{E}^{aj} n_i(\sigma) n_j(\sigma)} \end{aligned}$$

$$\hat{A}(\mathbf{S}) |\Gamma, j\rangle = 8\pi\gamma l_p^2 \sqrt{j(j+1)} |\Gamma, j\rangle$$

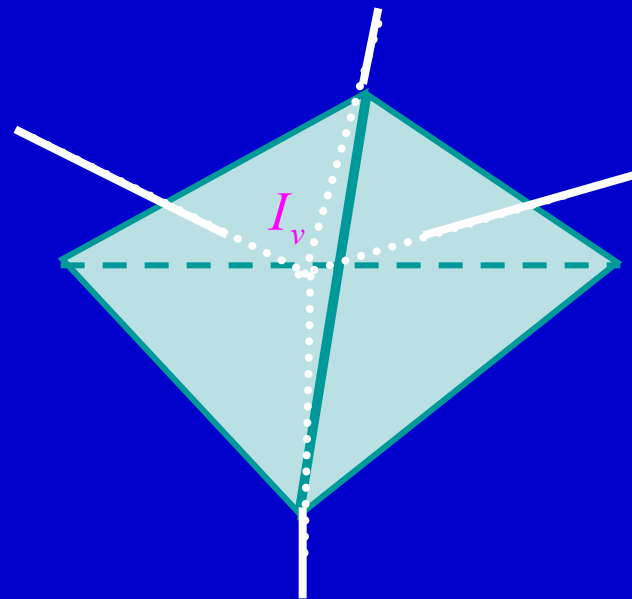


$$\text{Dim}(j) = 2j + 1$$

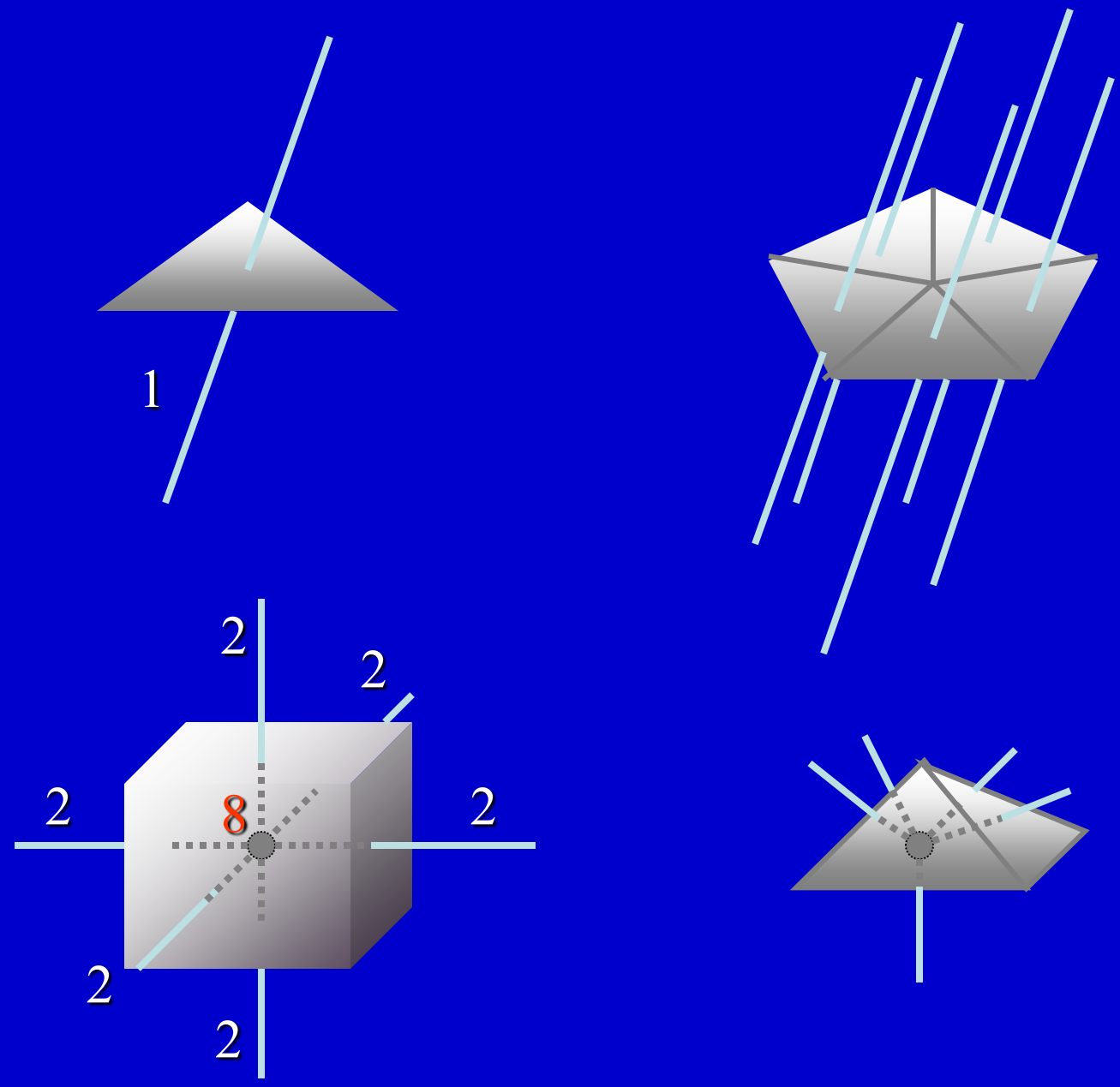
■ Volume operators

$$\begin{aligned} V(\mathbf{R}) &= \int_{\mathcal{R}} d^3x \sqrt{\det g} \\ &= \int_{\mathcal{R}} d^3x \sqrt{\frac{1}{3!} |\varepsilon_{abc} \varepsilon_{ijk} \hat{E}^{ai} \hat{E}^{bj} \hat{E}^{ck}|} \end{aligned}$$

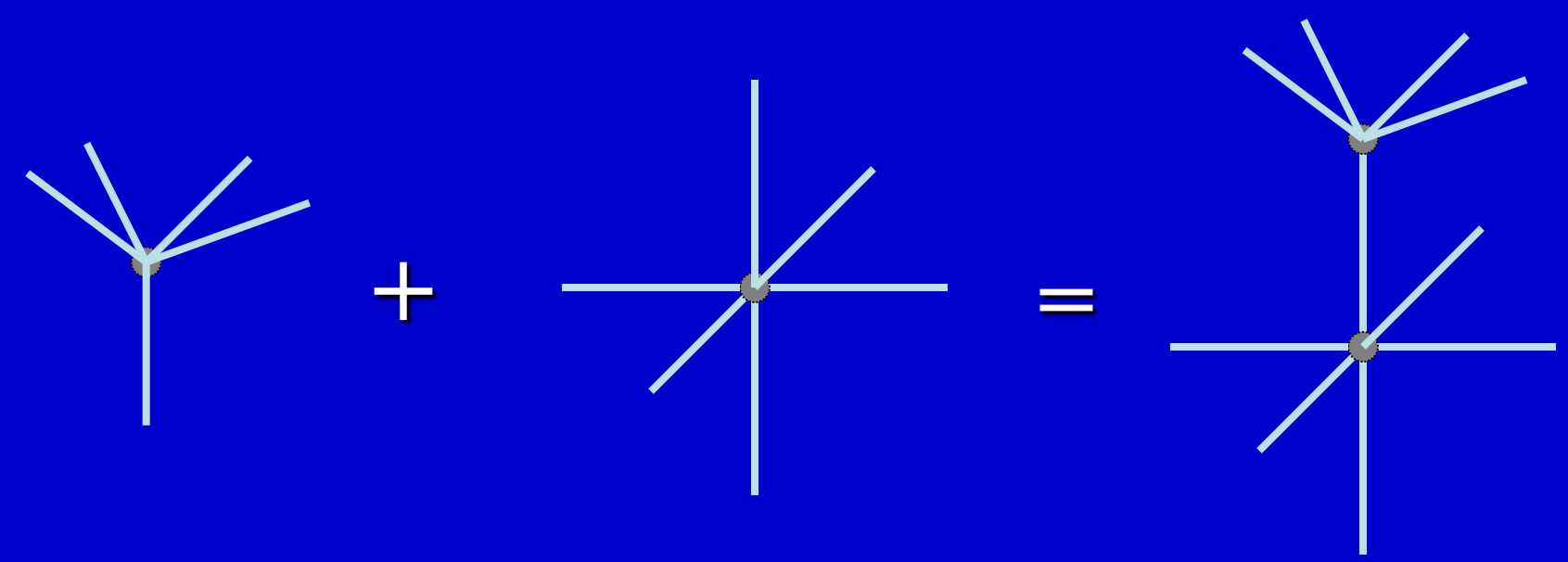
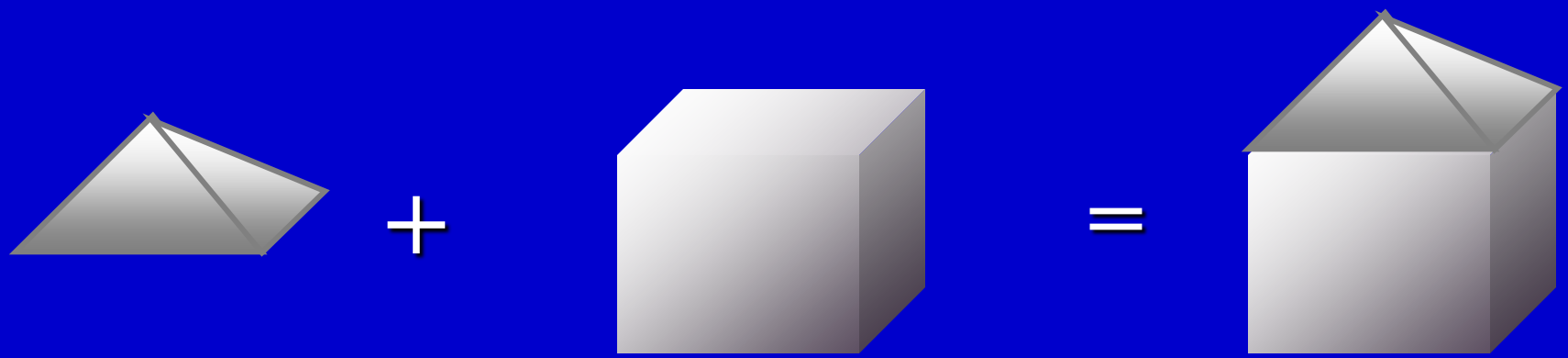
$$\hat{V}(\mathbf{R}) |\Gamma, j_m, I_v\rangle = \sum V(I_v) |\Gamma, j_m, I_v\rangle$$



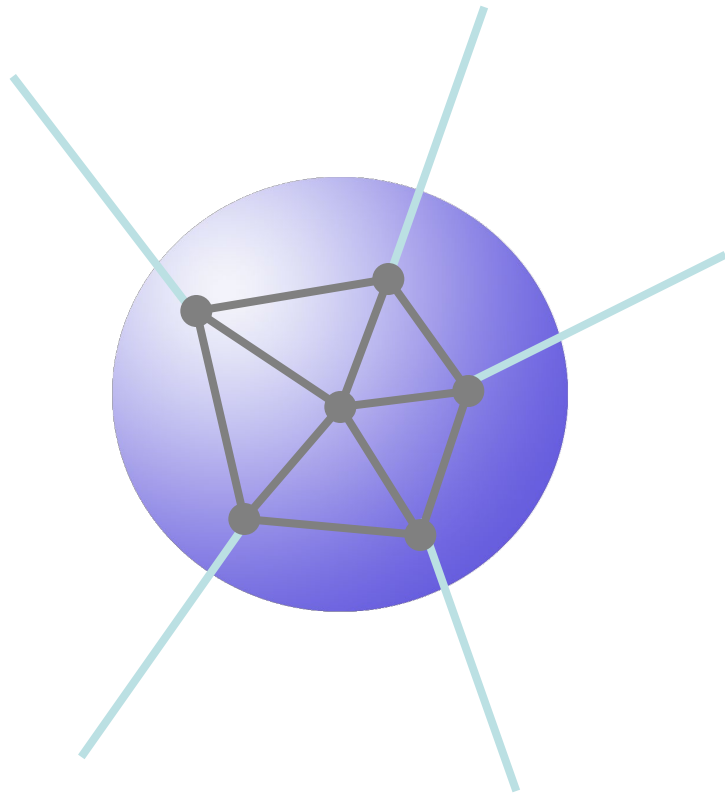
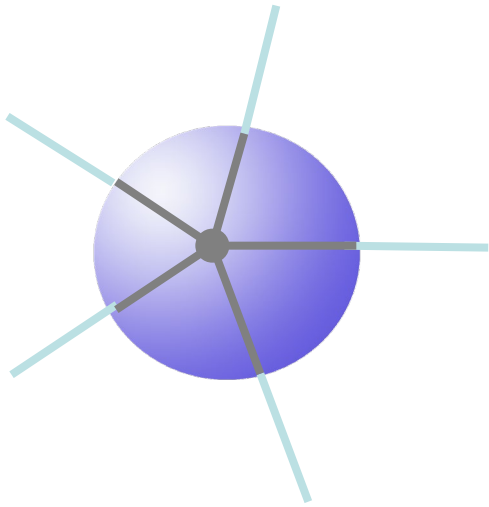
Spin networks and quantum geometry



Spin networks and quantum geometry



Spin networks and quantum geometry



Spin networks with boundary

- Spin network states with boundary

$$|\Gamma, \{j_e, j_o\}, \{I_v\}, \{M_o\}\rangle = \otimes |I_v^{\{j_e, j_o\}}\rangle \otimes |j_o, M_o\rangle \in H_\Gamma \otimes H^\partial$$

$$\langle \{h_e\}, \{h_o\} | \Gamma, \{j_e, j_o\}, \{I_v\}, \{M_o\} \rangle = \sum_{m_e, n_e, m_o} \prod_e U_{m_e n_e}^{j_e}(h_e) \prod_v (I_v)^{\{j_e, j_o\}}_{\{m_e, n_e, m_o\}} \prod_o U_{m_o n_o}^{j_o}(h_o)$$

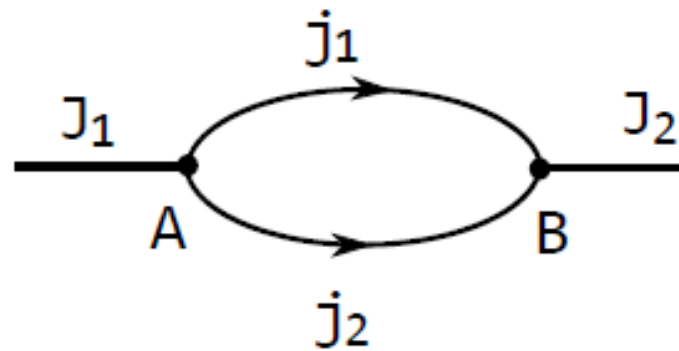
- The boundary states

$$|\tilde{\psi}[\Gamma, \{j_e, j_o\}, \{I_v\}, \{h_o\}]\rangle \in H^\partial$$

$$\otimes \langle j_o, M_o | \tilde{\psi}[\Gamma, \{j_e, j_o\}, \{I_v\}, \{h_o\}] \rangle = \langle \Gamma, \{j_e, j_o\}, \{I_v\}, \{M_o\} | \{h_e\}, \{h_o\} \rangle$$

Spin networks with boundary

- The boundary states in simple SN



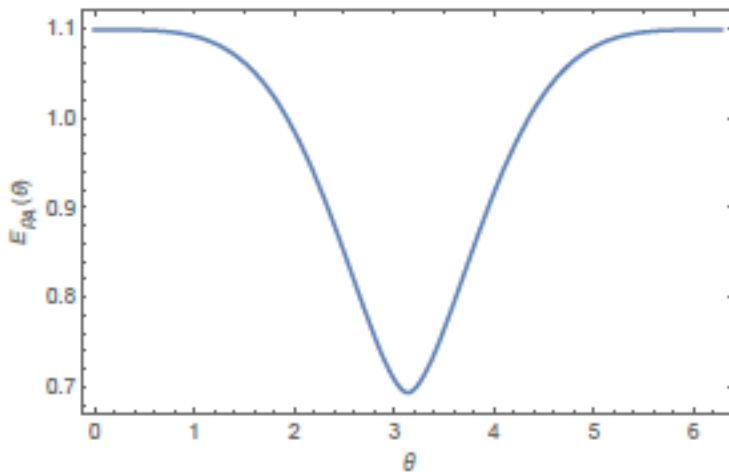
$$J_1 = J_2 = 1$$

$$|\psi_{J_1; J_2}\rangle = \sum_{M_1 M_2} N C_{M_1}^{J_1 j_1 j_2} D^{m_1 n_1} D^{m_2 n_2} U_{n_1}^{k_1*} [g(\theta)] U_{n_2}^{k_2*} [g(0)] C_{M_2}^{J_2 j_1 j_2} |M_1 M_2\rangle$$

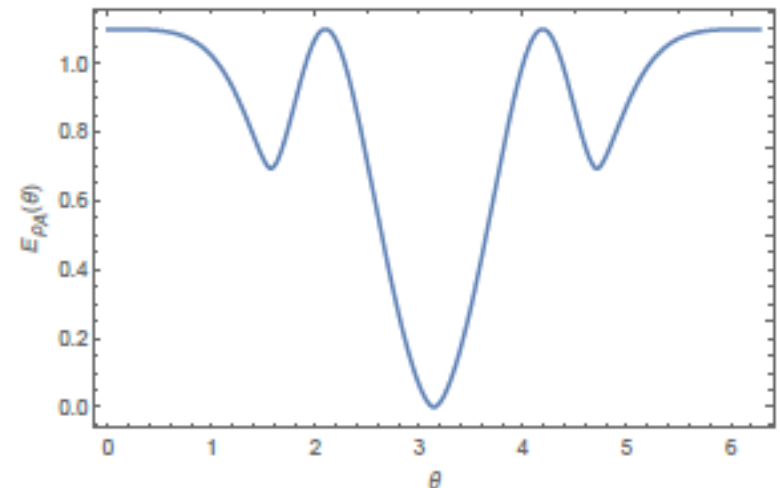
$$C_M^{J j_1 j_2} = \langle j_1 m_1; j_2 m_2 | J, M \rangle$$

$$D^{mn} = (-1)^{j-m} \delta^{m,-n}$$

$$U_{n_1}^{k_1} [g(\theta)] = e^{-ik_1 \theta} \delta_{n_1}^{k_1}$$



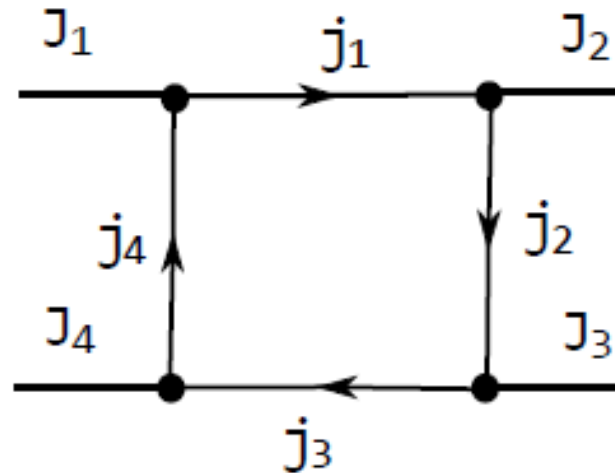
$$j_1 = j_2 = 1/2$$



$$j_1 = j_2 = 1$$

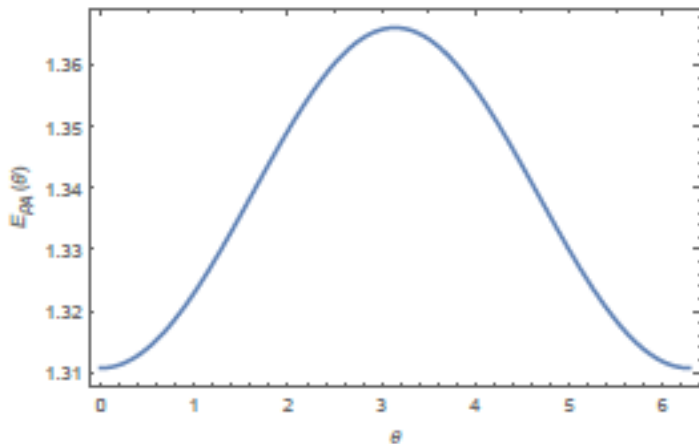
Spin networks with boundary

- The boundary states in simple SN

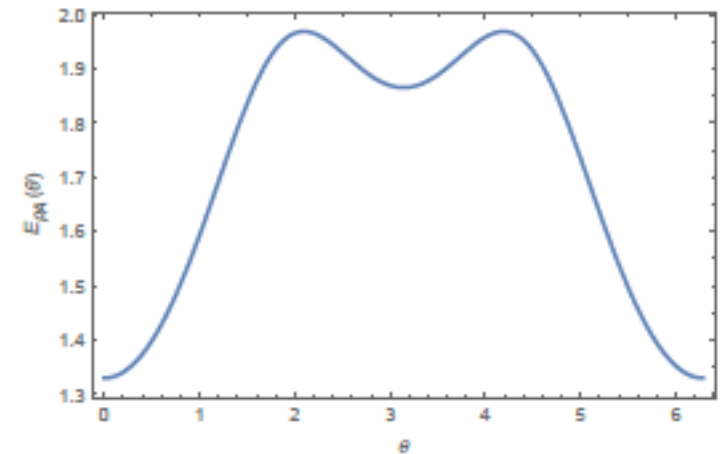


$$J_1 = J_2 = J_3 = J_4 = 1$$

$$|\psi_{J_1 J_4; J_2 J_3}\rangle = \sum_{M_l} N C_{M_1 m_1 k_4}^{J_1 j_1 j_4} D^{m_1 n_1} U_{n_1}^{k_1*} [g(\theta)] C_{M_2 m_2 k_1}^{J_1 j_2 j_1} D^{m_2 n_2} U_{n_2}^{k_2*} [g(0)] \\ \times C_{M_3 m_3 k_2}^{J_3 j_3 j_2} D^{m_3 n_3} U_{n_3}^{k_3*} [g(0)] C_{M_4 m_4 k_3}^{J_4 j_4 j_3} D^{m_4 n_4} U_{n_4}^{k_4*} [g(0)] |M_1 M_2 M_3 M_4\rangle$$



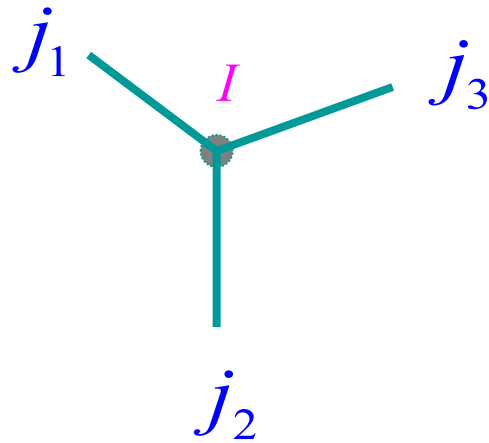
$$j_1 = j_2 = j_3 = j_4 = 1/2$$



$$j_1 = j_2 = j_3 = j_4 = 1$$

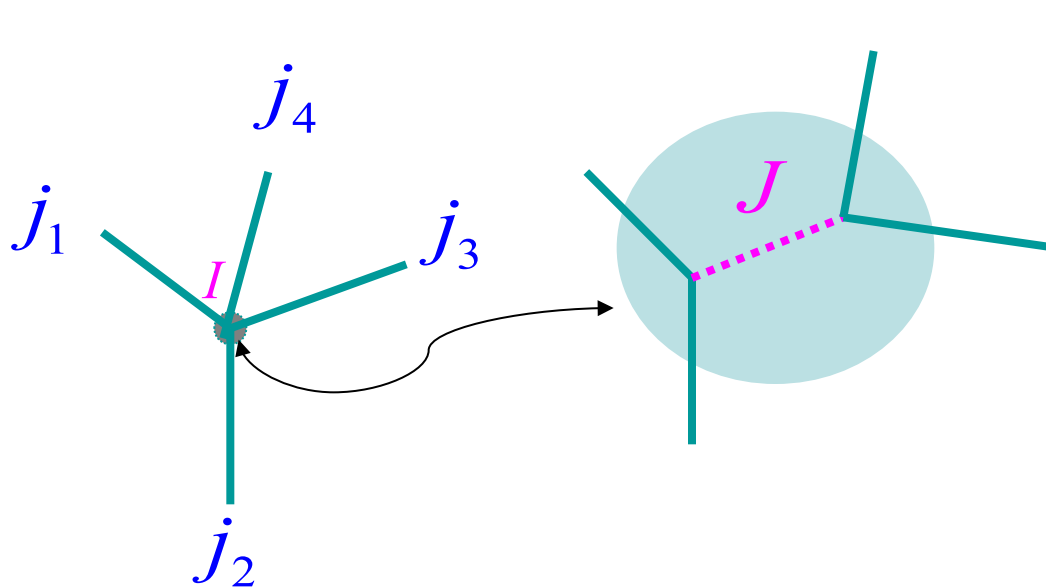
Quantum entanglement and quantum geometry

- Intertwiner (编织结)



$$\dim(I) = 1$$

$$V(I) = 0$$



$$j_1 = j_2 = j_3 = j_4 = \frac{1}{2}$$

$$J = 0, 1$$

$$\dim(I) = 2$$

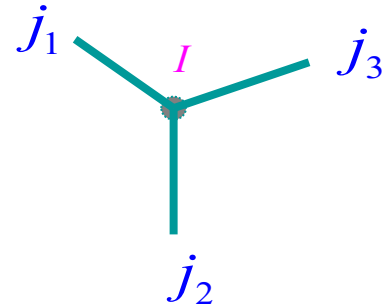
Quantum entanglement and quantum geometry

- Volume operator and eigenvalues

$$\hat{V} = l_p^3 \sqrt{|\hat{W}|}$$

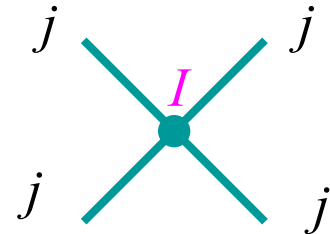
$$W_{\nu\lambda} = \langle I_\nu | \hat{W} | I_\lambda \rangle$$

$$\hat{W} |I_3\rangle = 0$$



$$\hat{W} = \frac{1}{8} \varepsilon^{ijk} (J_i^{(1)} J_j^{(2)} J_k^{(3)} - J_i^{(1)} J_j^{(2)} J_k^{(4)} + J_i^{(1)} J_j^{(3)} J_k^{(4)} - J_i^{(2)} J_j^{(3)} J_k^{(4)})$$

$$\hat{W} |I_4\rangle \neq 0$$



$$\hat{f}_i^{(p)} = \frac{1}{2} \int_{S_2^p} \varepsilon_{abc} \hat{E}_i^a dx^b dx^c$$

Quantum entanglement and quantum geometry

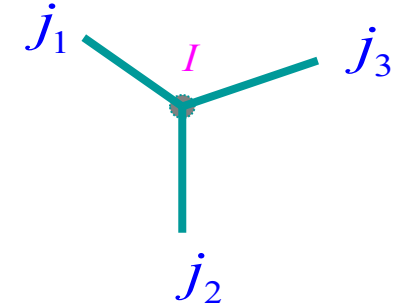
Y. Li, M. Han, M. Grassl and B. Zeng, (2017), ArXiv:1612.04504.

- Tri-valent SU(2) invariant perfect tensors

$$|\psi_3\rangle = \sum_{m_1 m_2 m_3} \begin{pmatrix} j_1 & j_2 & j_3 \\ m_1 & m_2 & m_3 \end{pmatrix} |m_1 m_2 m_3\rangle$$

Perfect tensors

$$S_{tot} = S_1 + S_2 + S_3 = \ln(2j_1 + 1) + \ln(2j_2 + 1) + \ln(2j_3 + 1)$$

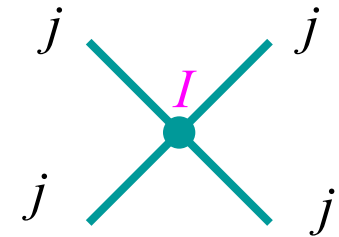


- No Four-valent SU(2) invariant perfect tensors

Perfect tensors $S_{12} = S_{13} = S_{14} = \ln 4$ $(j = \frac{1}{2})$

$$S_{tot} = S_{12} + S_{13} + S_{14} = 3 \ln 4$$

$$\begin{aligned} |\psi_4\rangle &= \sum_{J=0,1} \alpha(J) |I(J)\rangle \\ &= \sum_{J=0,1} \frac{\alpha(J)}{\sqrt{2J+1}} C_{m_1 m_2 M}^{j j J} C_{m_3 m_4 N}^{j j J} (-1)^{J-M} \delta^{M,-N} |m_1 m_2 m_3 m_4\rangle \end{aligned}$$



Quantum entanglement and quantum geometry

- Conjecture

The emergence of the space with **non-zero volume** in three dimensions is the reflection of the **non-perfectness** of SU(2)-invariant tensors.

- The measure of the non-perfectness

Maximally entangled SU(2)-invariant states S_{tot}^{\max}

$$\Delta S = S_{tot}^{pt} - S_{tot}^{\max}$$

$$\delta = \Delta S / S_{tot}^{pt}$$

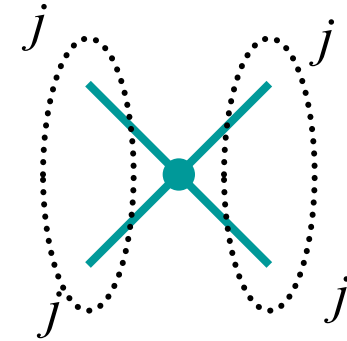
Quantum entanglement and quantum geometry

- Entanglement of boundary states

$$\begin{aligned}
 |\psi_4\rangle &= \sum_{J=0,\dots,2j} \alpha(J) |I(J)\rangle \\
 &= \sum_{J=0,\dots,2j} \frac{\alpha(J)}{\sqrt{2J+1}} C_{m_1 m_2 M}^{j j J} C_{m_3 m_4 N}^{j j J} (-1)^{J-M} \delta^{M,-N} |m_1 m_2 m_3 m_4\rangle
 \end{aligned}$$

$$\rho_{ik} = \text{Tr}_{ik} \frac{|\psi_4\rangle\langle\psi_4|}{\langle\psi_4|\psi_4\rangle}$$

$$S_{ik} = -\text{Tr}(\rho_{ik} \ln \rho_{ik})$$



Y. Ling, Y.Xiao, M.Wu. [arXiv:1907.01215](https://arxiv.org/abs/1907.01215)

- SU(2)-invariant maximally entangled states

Determine $2j+1$ complex numbers in intertwiner:

$$\alpha(J), \quad J = 0, \dots, 2j$$

$$j_i = \frac{1}{2}, i = 1 \dots 4$$

$$\alpha(0) = \frac{1}{\sqrt{2}}, \alpha(1) = \pm \frac{i}{\sqrt{2}}$$

$$j_i = 1, i = 1 \dots 4$$

$$\alpha(0) = \frac{\sqrt{2}}{3}, \alpha(1) = \pm \frac{i}{\sqrt{2}}, \alpha(2) = -\frac{1}{3} \sqrt{\frac{5}{2}}$$

Quantum entanglement and quantum geometry

- E.g.:

$$j_i = \frac{1}{2}, i = 1 \dots 4$$

Perfect tensor

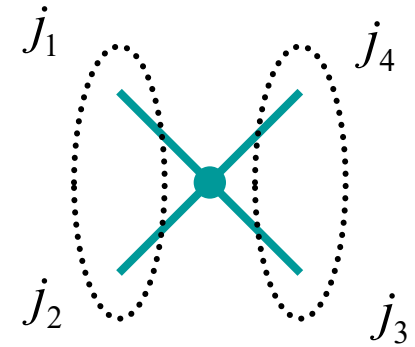
$$S_{12} = S_{13} = S_{14} = \ln 4$$

$$S_{tot}^{pt} = S_{12} + S_{13} + S_{14} = 3 \ln 4$$

Maximally entangled states $S_{tot}^{\max} = S_{12} + S_{13} + S_{14} = 3 \ln(2\sqrt{3})$

$$\Delta S \cong 0.43$$

$$\delta \cong 0.104$$



Y. Ling, Y.Xiao, M.Wu. [arXiv:1907.01215](https://arxiv.org/abs/1907.01215)

$$j_i = 1, i = 1 \dots 4$$

Perfect tensor

$$S_{12} = S_{13} = S_{14} = 2 \ln 3$$

$$S_{tot}^{pt} = S_{12} + S_{13} + S_{14} = 6 \ln 3$$

$$S_{tot}^{\max} = S_{12} + S_{13} + S_{14} = \frac{5}{3} \ln 2 + \frac{9}{2} \ln 3$$

Maximally entangled states

$$S_{tot}^{\max} = S_{12} + S_{13} + S_{14} = \frac{5}{3} \ln 2 + \frac{9}{2} \ln 3$$

$$\Delta S \cong 0.49$$

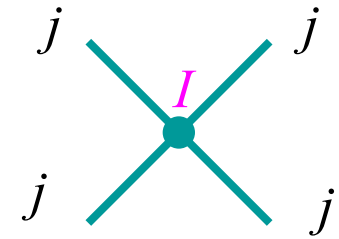
$$\delta \cong 0.075$$

Quantum entanglement and quantum geometry

- Eigenstates and Eigenvalues of volume operator

$$\hat{V} = l_p^3 \sqrt{|\hat{W}|}$$

$$W_{\nu\lambda} = \langle I_\nu | \hat{W} | I_\lambda \rangle$$



$$j_1 = j_2 = j_3 = j_4 = 1/2$$



$$W = \begin{pmatrix} 0 & -\frac{\sqrt{3}}{8}i \\ \frac{\sqrt{3}}{8}i & 0 \end{pmatrix}$$



$$W_{\pm} = \pm \frac{\sqrt{3}}{8}$$

$$V = l_p^3 \sqrt{\frac{\sqrt{3}}{8}}$$

$$|\pm\rangle = \frac{1}{\sqrt{2}}|J=0\rangle \pm \frac{i}{\sqrt{2}}|J=1\rangle$$

Maximally entangled states have definite orientation!

$$j_1 = j_2 = j_3 = j_4 = 1$$



$$W = \begin{pmatrix} 0 & -\frac{i}{\sqrt{3}} & 0 \\ \frac{i}{\sqrt{3}} & 0 & -\frac{i}{2}\sqrt{\frac{5}{3}} \\ 0 & 0 & \frac{i}{2}\sqrt{\frac{5}{3}} \end{pmatrix}$$



$$W_{\pm} = \pm \frac{\sqrt{3}}{2}, W_3 = 0 \quad V = l_p^3 \sqrt{\frac{\sqrt{3}}{2}}, 0$$

$$|\pm\rangle = \frac{\sqrt{2}}{3}|J=0\rangle \pm \frac{i}{\sqrt{2}}|J=1\rangle - \frac{1}{3}\sqrt{\frac{5}{2}}|J=2\rangle$$

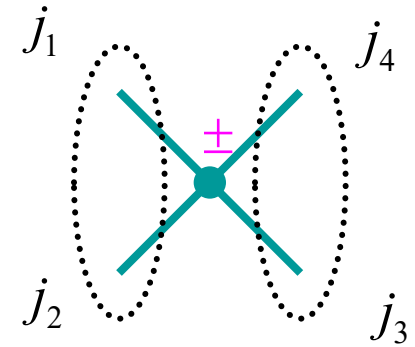
Quantum entanglement and quantum geometry

- Numerical Check

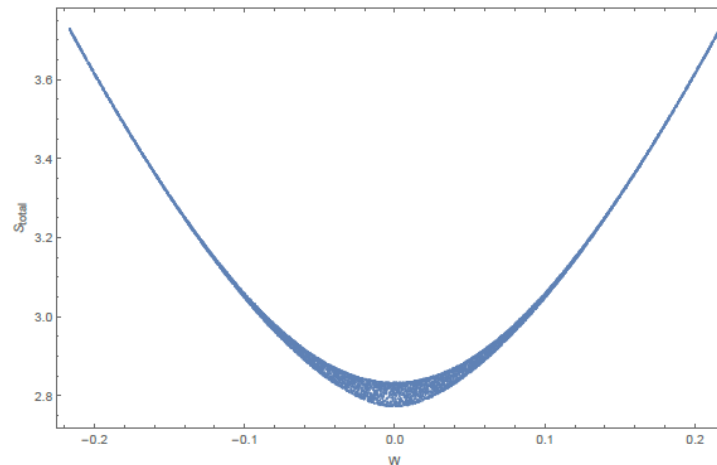
$$j_i = \frac{1}{2}, i = 1 \dots 4$$

$$\begin{aligned} |\psi_4\rangle &= \sum_{J=0,1} \alpha(J) C_{m_1 m_2 M}^{j j J} C_{m_3 m_4 N}^{j j J} (-1)^{J-M} \delta^{M,-N} |m_1 m_2 m_3 m_4\rangle \\ &\equiv C_+ |+\rangle + C_- |-\rangle \end{aligned}$$

$$\langle \widehat{W} \rangle = \frac{\langle \psi_4 | \widehat{W} | \psi_4 \rangle}{\langle \psi_4 | \psi_4 \rangle} = \frac{\sqrt{3} |C_+|^2 - |C_-|^2}{8 |C_+|^2 + |C_-|^2} \quad \langle \widehat{V} \rangle = \sqrt{\frac{\sqrt{3}}{8}}$$



Y. Ling, Y.Xiao, M.Wu. [arXiv:1907.01215](https://arxiv.org/abs/1907.01215)

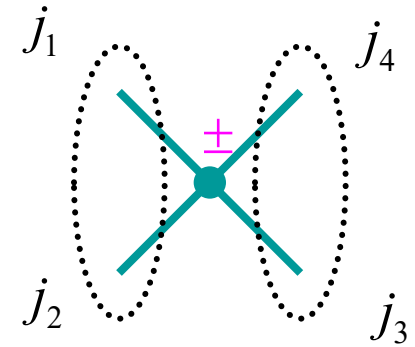


Quantum entanglement and quantum geometry

• Numerical Check

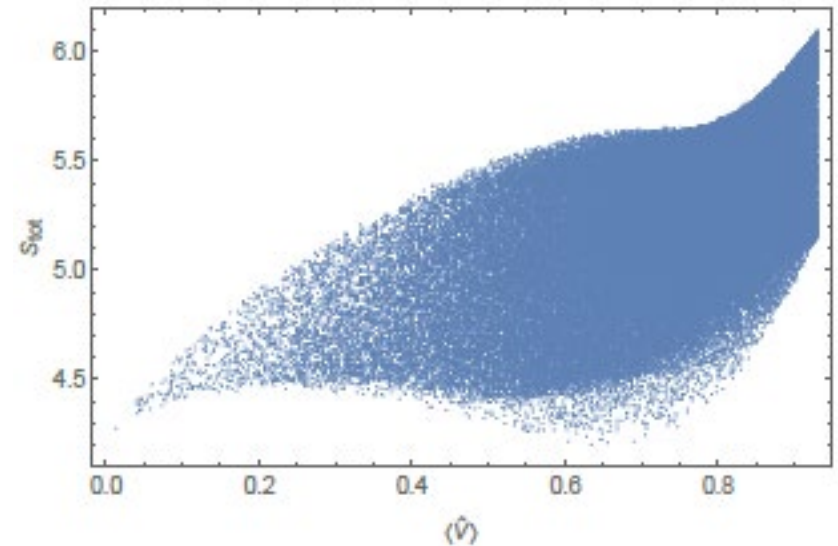
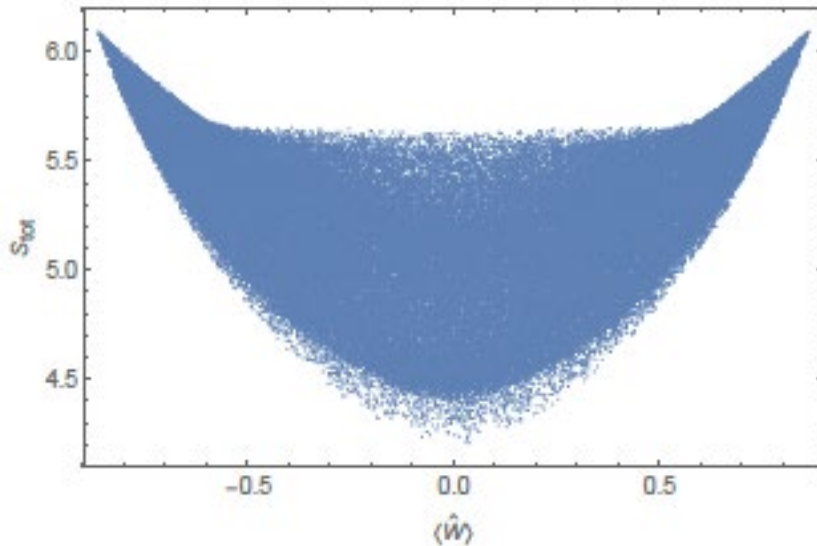
$$j_i = 1, i = 1 \dots 4$$

$$|\psi_4\rangle = \alpha_+ |\lambda_+\rangle + \alpha_0 |\lambda_0\rangle + \alpha_- |\lambda_-\rangle$$



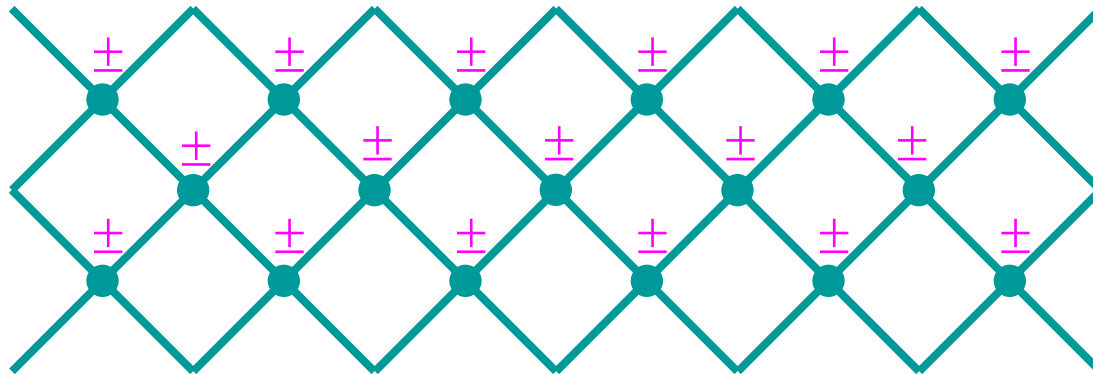
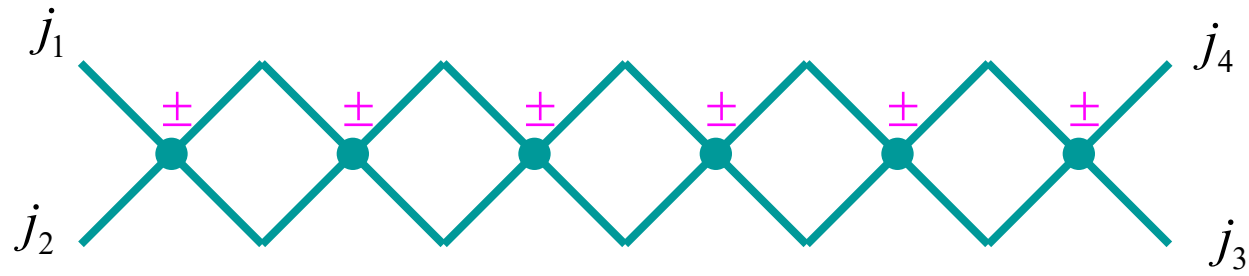
$$\langle \hat{W} \rangle = \frac{\sqrt{3}}{2} \frac{|\alpha_+|^2 - |\alpha_-|^2}{|\alpha_+|^2 + |\alpha_0|^2 + |\alpha_-|^2}$$

$$\langle \hat{V} \rangle = \sqrt{\frac{\sqrt{3}}{2}} \frac{|\alpha_+|^2 + |\alpha_-|^2}{|\alpha_+|^2 + |\alpha_0|^2 + |\alpha_-|^2}$$

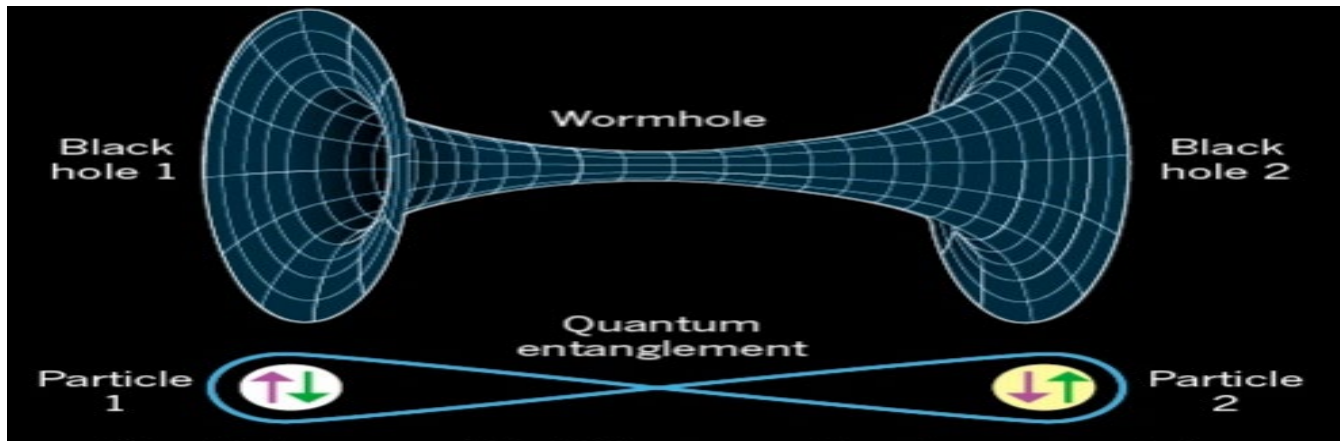
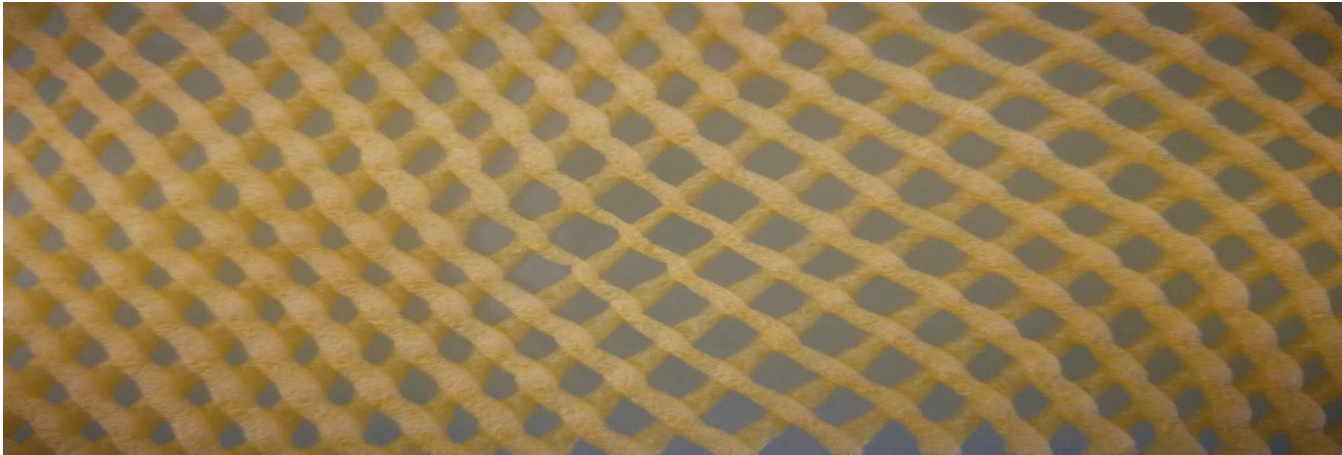


Quantum entanglement and quantum geometry

$$j_i = \frac{1}{2}, i = 1 \dots 4$$



Quantum entanglement and quantum geometry



Quantum entanglement and quantum geometry

■ AI for Quantum Gravity is on the road!

1. Suppose you knew nothing about LQG.
2. Feed DS4 with the ArXiv paper, for instance, [arXiv:1907.01215](#).
3. Input commands and ask the DS4 to design a Mathematica codes to compute the eigenvalues of the volume operator when acting on 4-valent vertex with $J = \frac{1}{2}, 1$ with detailed requirement.
4. Using this Mathematica codes to compute the eigenvalues of the volume operator with **higher** spins which do not appear in [arXiv:1907.01215](#).
5. Check the answers in [gr-qc/9602023!](#)

Quantum entanglement and quantum geometry

```
(* ===== *)
(* Mathematica program to compute volume eigenvalues *)
(* for a 4-valent vertex with edges j (configurable) *)
(* *)
(* Based on: Y. Ling, Y. Xiao, M. Wu, arXiv:1907.01215 *)
(* "Note on quantum entanglement and quantum geometry" *)
(* *)
(* Strategy: *)
(* 1. Build SU(2) generators as exact symbolic matrices *)
(* 2. Build tensor product operators via KroneckerProduct *)
(* 3. Compute W operator exactly (symbolic) *)
(* 4. Project onto intertwiner basis *)
(* 5. Diagonalize *)
(* ===== *)
```

```
ClearAll["Global`*"];
```

```
Print["====="];
Print[" Volume eigenvalues for 4-valent vertex (exact symbolic)"];
Print[" Based on Y. Ling, Y. Xiao, M. Wu, arXiv:1907.01215"];
Print["====="];
```

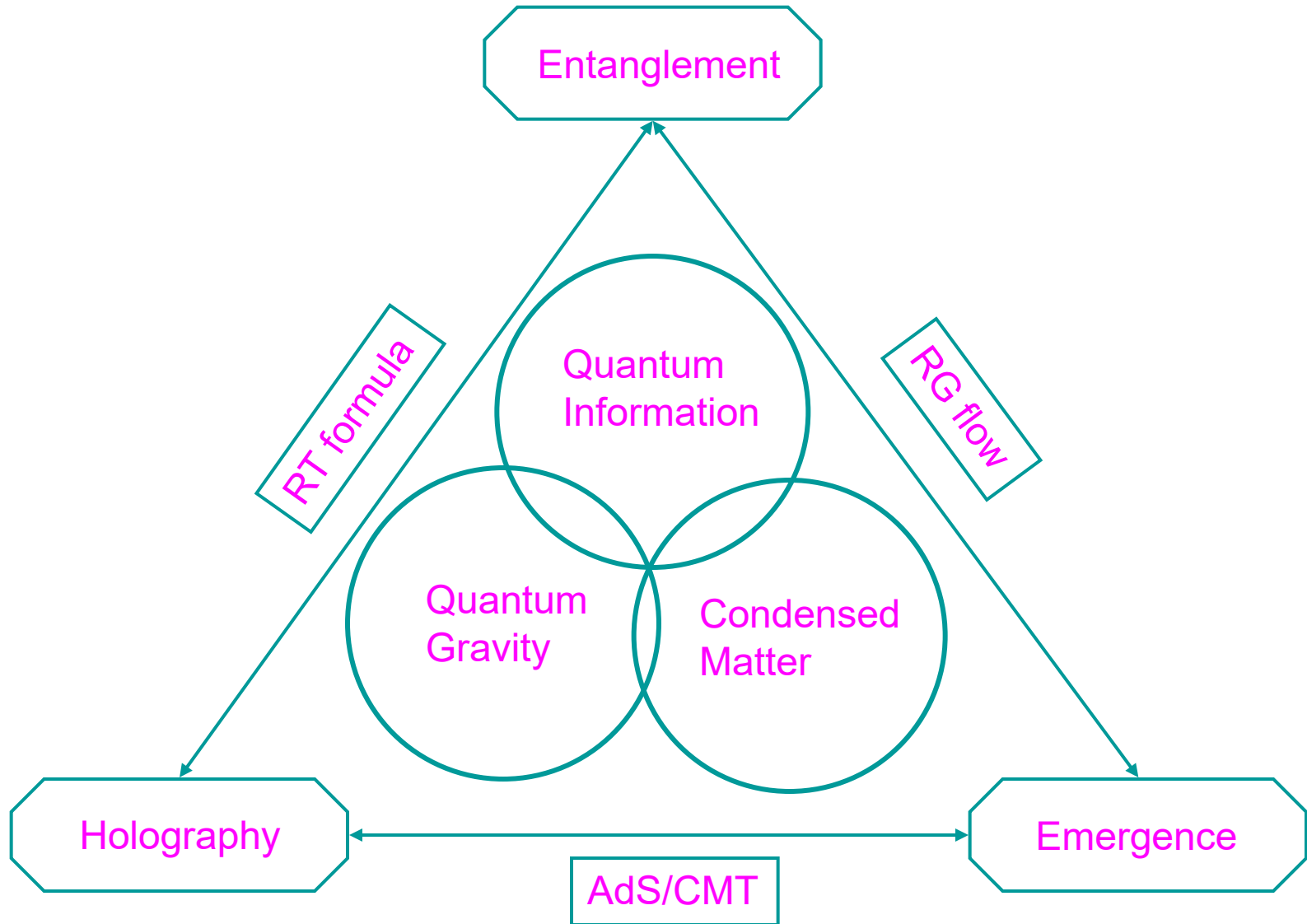
```
(* ===== *)
(* USER CONFIGURABLE PARAMETER *)
(* ===== *)
```

```
spin = 2; (* <---- CHANGE THIS VALUE AS NEEDED *)
```

```
Print["\nConfigured spin: j = ", spin];
```

```
(* ===== *)
(* SU(2) generators as EXACT symbolic matrices *)
(* Ordering:  $|j,j\rangle = \text{index } 1, \dots, |j,-j\rangle = d$  *)
(* ===== *)
```

Summary:



Thanks!