

# **Quantum mechanism angular momentum & graphical calculus**

**II**

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$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

## II. Introduction to:

1. Elementary variables
2. Quantum configuration space
3. A natural measure on the quantum configuration space
3. The Hilbert space
4. Spin network states
5. Elementary operators

## References:

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- [3] T. Thiemann, *Modern Canonical Quantum General Relativity* (Cambridge University Press, Cambridge, England, 2007). <https://doi.org/10.1017/CBO9780511755682>. arXiv:gr-qc/0110034
- [4] C. Rovelli, *Quantum Gravity* (Cambridge University Press, Cambridge, England, 2004). <https://doi.org/10.1017/CBO9780511755804>
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# 1. Elementary variables: holonomy and flux

The phase space: canonical variables  $(A_a^i, \tilde{E}_i^a)$  on  $\Sigma$

The only non-trivial Poisson bracket:

$$\{A_a^i(x), \tilde{E}_j^b(y)\} = \kappa\beta\delta_a^b\delta_j^i\delta^3(x, y), \text{ where } \kappa = 8\pi G$$

The **holonomy**  $h_e(A)$  of connection  $A_a^i$  along an edge  $e : [0,1] \rightarrow \Sigma$

$$A(e(t)) := \dot{e}^a(t)A_a^i(e(t))\tau_i$$

$$h_e(A) := \mathcal{P}\exp\left(\int_e A\right) = \mathbb{I}_2 + \sum_{n=1}^{\infty} \int_0^1 dt_1 \int_{t_1}^1 dt_2 \cdots \int_{t_{n-1}}^1 dt_n A(e(t_1)) \cdots A(e(t_n)) \in \text{SU}(2)$$

as the unique solution  $h_e(A) \equiv h_{e([0,t=1])}(A)$  of

$$\frac{dh_{e([0,t])}(A)}{dt} = h_{e([0,t])}(A) A(e(t)), \quad h_{e([0,0])}(A) = \mathbb{I}_2$$

Properties of holonomy

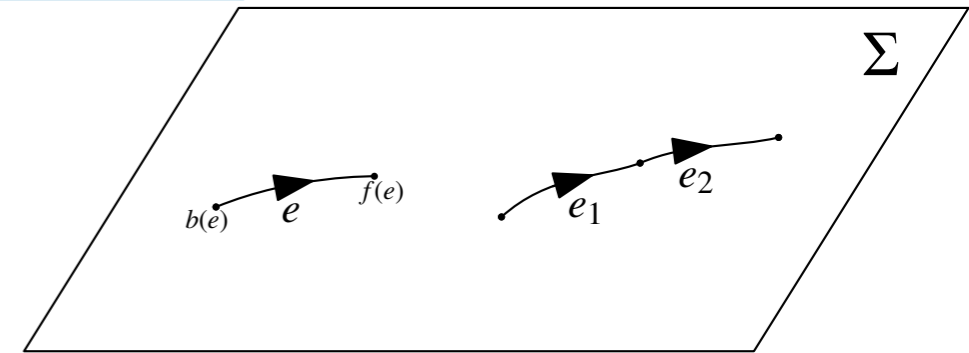
$$h_{e_1 \circ e_2}(A) = h_{e_1}(A)h_{e_2}(A)$$

$$h_{e^{-1}}(A) = h_e(A)^{-1}$$

The behavior of holonomy under gauge transformation

$$h_e(A^g) = g(b(e))h_e(A)g(f(e))^{-1}$$

$$A \mapsto A^g = -(dg)g^{-1} + gAg^{-1}$$



The **flux**  $\tilde{E}_i(S)$  of densitized triad  $\tilde{E}_i^a$  through a 2-surface  $S \subset \Sigma$  is defined by

$$\tilde{E}_i(S) := \int_S \left( * \tilde{E}_i \right)_{ab}$$

The elementary variables are defined in a background-independent way !

# 1. Elementary variables: holonomy and flux

$$e \cap S = b(e)/f(e)$$

$$\{A_a^i(x), \tilde{E}_j^b(y)\} = \kappa\beta\delta_a^b\delta_j^i\delta^3(x,y)$$

$$L^{(\tau_i)} g^A_B = (g\tau_i)^A_B \Leftrightarrow L^{(\tau_i)} g = g\tau_i$$

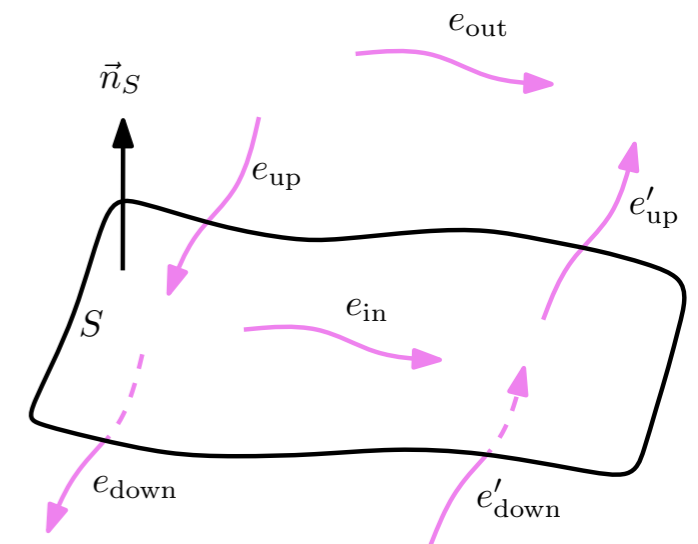
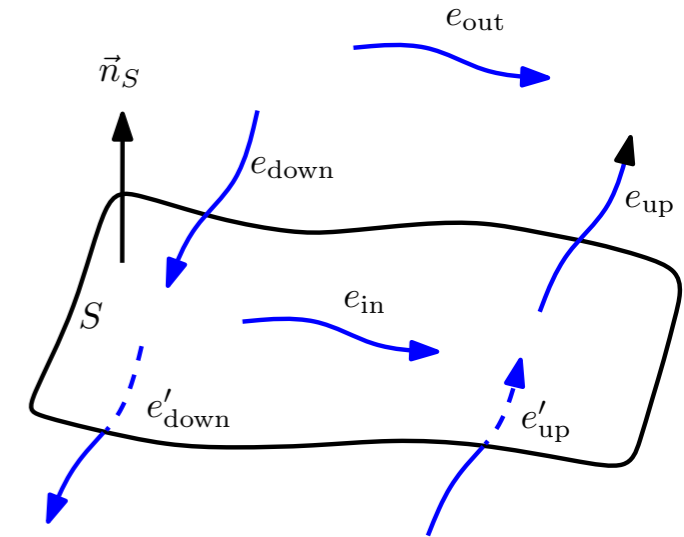
$$R^{(\tau_i)} g^A_B = -(\tau_i g)^A_B \Leftrightarrow R^{(\tau_i)} g = -\tau_i g$$

$$\tau_i = -\frac{i}{2}\sigma_i \text{ with } \sigma_i \text{ being the Pauli matrices.}$$

$$\begin{aligned} \{h_e(A), \tilde{E}_i(S)\} &= \frac{\kappa\beta}{2}\epsilon(e, S)\left(\delta_{e \cap S, b(e)}\tau_i h_e(A) + \delta_{e \cap S, f(e)}h_e(A)\tau_i\right) \\ &= \frac{\kappa\beta}{2}\kappa(e, S)\left(\delta_{e \cap S, b(e)}\tau_i h_e(A) - \delta_{e \cap S, f(e)}h_e(A)\tau_i\right) \\ &= -\frac{\kappa\beta}{2}\kappa(e, S)\left(\delta_{e \cap S, b(e)}R_e^{(\tau_i)} + \delta_{e \cap S, f(e)}L_e^{(\tau_i)}\right)h_e(A) \end{aligned}$$

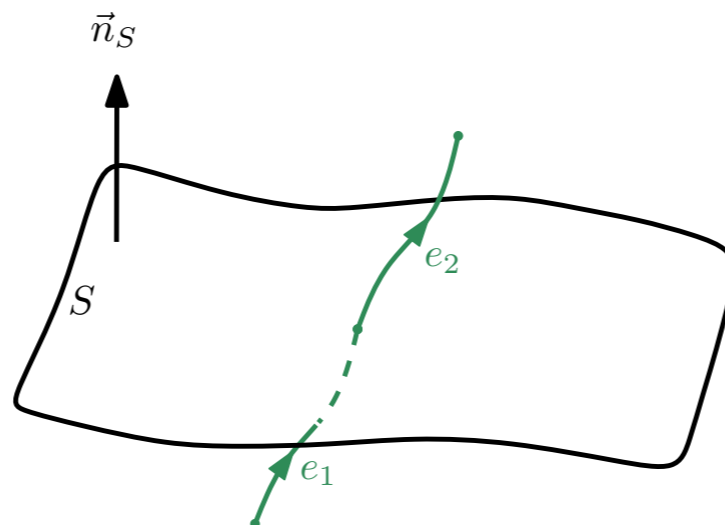
Direction factor:  $\epsilon(e, S) := \begin{cases} +1 & e \text{ is the 'up' type w.r.t. } S \\ -1 & e \text{ is the 'down' type w.r.t. } S \\ 0 & e \text{ is the 'inside/outside' type w.r.t. } S \end{cases} \left(n_a \dot{e}^a|_{e \cap S} > 0\right)$

Steric factor:  $\kappa(e, S) := \begin{cases} +1 & e \text{ is the 'up' type w.r.t. } S \\ -1 & e \text{ is the 'down' type w.r.t. } S \\ 0 & e \text{ is the 'inside/outside' type w.r.t. } S \end{cases}$



$$e = e_1 \circ e_2, \quad e_1 \cap e_2 = e \cap S$$

$$\{h_e(A), \tilde{E}_i(S)\} = \kappa\beta h_{e_1}(A) \tau_i h_{e_2}(A)$$



## 2. Quantum configuration space $\bar{\mathcal{A}}$

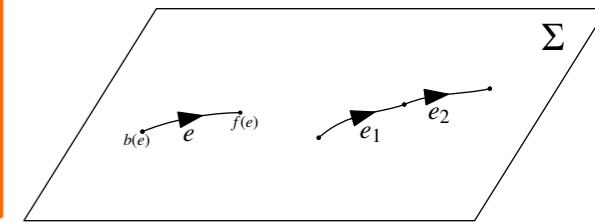
The **classical configuration space**  $\mathcal{A}$  is the set of **smooth** connections  $A$  on  $\Sigma$

The holonomy  $h_e(A)$  of connection  $A_a^i$  along the edge  $e : [0,1] \rightarrow \Sigma$

$$A(e) \equiv h_e(A) = \mathcal{P}\exp\left(\int_e A\right) \in \text{SU}(2)$$

**Properties:**

$$A(e_1 \circ e_2) = A(e_1)A(e_2), \quad A(e^{-1}) = A(e)^{-1}$$



**Def. (Set of paths).** A *piecewise analytic path* is a composition of edges. The set of paths is denoted by  $\mathcal{P}$ .

$\mathcal{A} = \text{Hom}(\mathcal{P}, \text{SU}(2))$ , be a **continue** homomorphism map

The **quantum configuration space**  $\bar{\mathcal{A}}$  as a specific extension of  $\mathcal{A}$

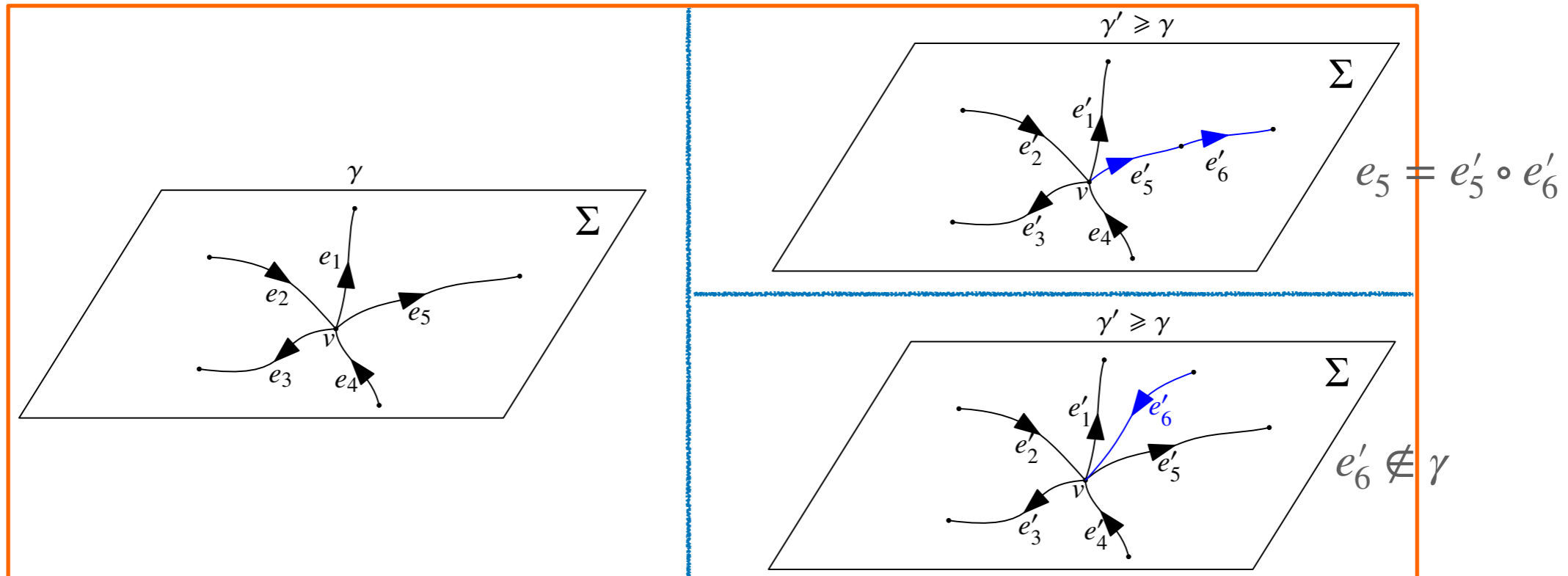
$\bar{\mathcal{A}} := \text{Hom}(\mathcal{P}, \text{SU}(2))$ , can be an arbitrary **discontinue** homomorphism map

**Relation:** in a natural topology (due to Gel'fand),  $\mathcal{A}$  is densely embedded in  $\bar{\mathcal{A}}$ , and  $\bar{\mathcal{A}}$  is compact in a natural (Gel'fand) topology.

$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

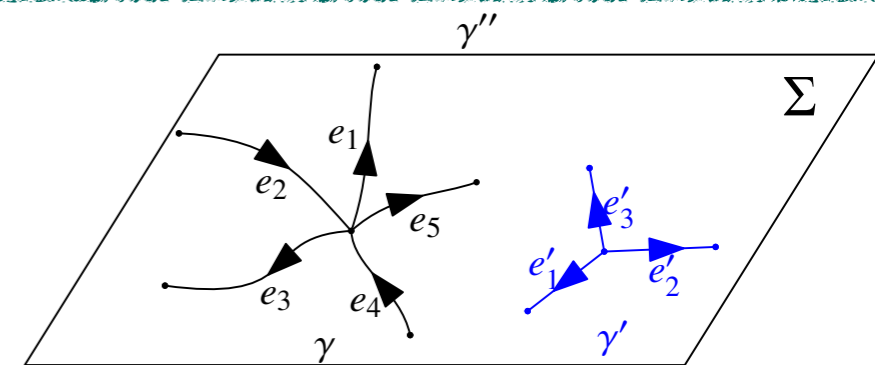
### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

**Def. (Graph).** A graph  $\gamma$  is a collection of edges and vertices, in which the edges can only intersect each other at their endpoints. A graph  $\gamma'$  is said to be large than another graph  $\gamma$  (or contain  $\gamma$ ),  $\gamma' \geq \gamma$ , if every  $e$  of  $\gamma$  can be written as  $e = e'_1 \circ e'_2 \cdots \circ e'_k$  for some edges  $e', \dots, e'_k$  of  $\gamma'$ .



#### Remarks:

- (i) Not all two graphs,  $\gamma$  and  $\gamma'$ , need to have a containment relation.
- (ii) For any two graph  $\gamma, \gamma'$ , there exists  $\gamma''$  such that  $\gamma'' \geq \gamma, \gamma'$ , where  $\gamma''$  contain  $\gamma \cup \gamma'$ .



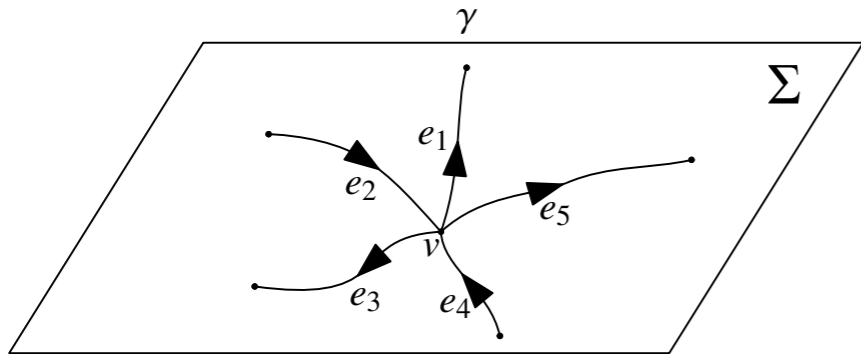
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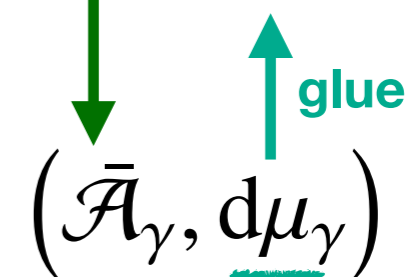
**Remarks:**

- (i) Not all two graphs,  $\gamma$  and  $\gamma'$ , need to have a containment relation.
- (ii) For any two graph  $\gamma, \gamma'$ , there exists  $\gamma''$  such that  $\gamma'' \geq \gamma, \gamma'$ , where  $\gamma''$  contain  $\gamma \cup \gamma'$ .

$\bar{\mathcal{A}}_\gamma := \bar{\mathcal{A}}|_\gamma$ , the restriction of the domain of  $\bar{\mathcal{A}}$  on  $\gamma$  with  $n$  edges  $e_1, \dots, e_n$



$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$



### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

$\downarrow$   
 $(\bar{\mathcal{A}}_\gamma, d\mu_\gamma)$

$\uparrow$  glue

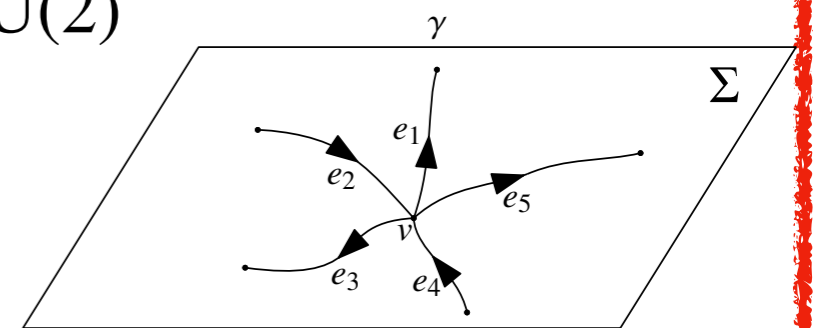
$\bar{\mathcal{A}}_\gamma := \bar{\mathcal{A}}|_\gamma$ , the restriction of the domain of  $\bar{\mathcal{A}}$  on  $\gamma$  with  $n$  edges  $e_1, \dots, e_n$

There exists a bijection

$$\mathcal{I}_E : \bar{\mathcal{A}}_\gamma \rightarrow \text{SU}(2)^n \equiv \text{SU}(2) \times \dots \times \text{SU}(2)$$

$$A \mapsto (A(e_1), \dots, A(e_n))$$

$|E(\gamma)|$  : the number of edges of  $\gamma$



which helps us to

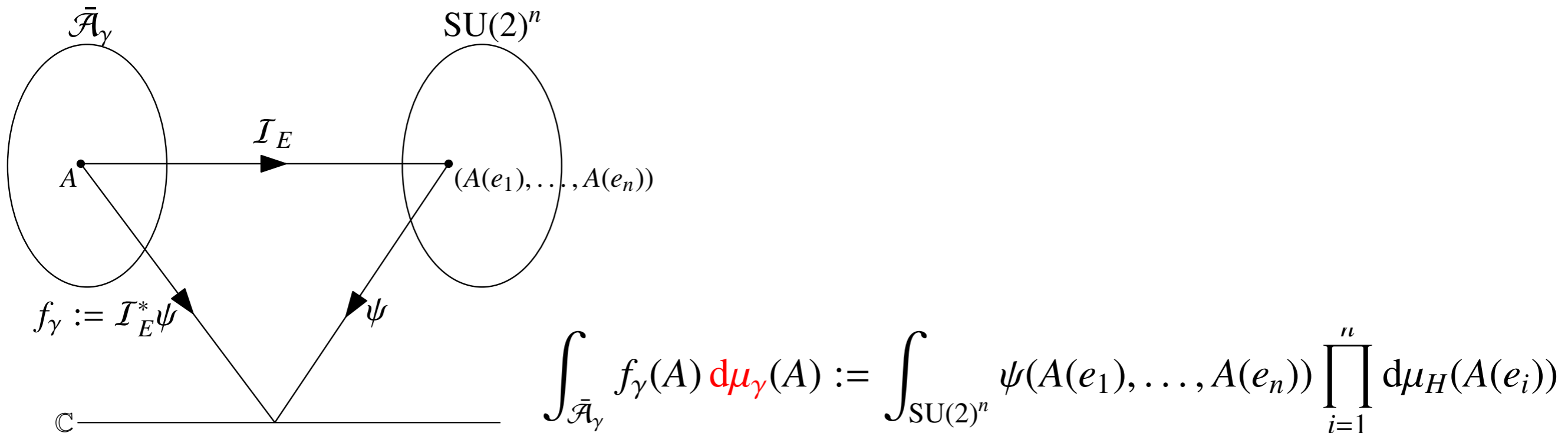
- (i) equip  $\bar{\mathcal{A}}_\gamma$  with a topology induced from that of  $\text{SU}(2)^n$ .
- (ii) define a nature measure  $d\mu_\gamma$  on  $\bar{\mathcal{A}}_\gamma$  induced from the Haar measure  $d\mu_H$  on  $\text{SU}(2)$ .

### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

Step I: a natural measure  $d\mu_\gamma$  on  $\bar{\mathcal{A}}_\gamma$

There exists a bijection

$$\begin{aligned} \mathcal{I}_E : \bar{\mathcal{A}}_\gamma &\rightarrow \text{SU}(2)^n \equiv \text{SU}(2) \times \cdots \times \text{SU}(2) \\ A &\mapsto (A(e_1), \dots, A(e_n)) \end{aligned}$$



#### Remarks:

(i) For any given  $\gamma$  and for each  $\bar{A} \in \bar{\mathcal{A}}$ , there exists a smooth  $A \in \mathcal{A}$  such that

$$\bar{A}(e) = h_e(A), \quad \forall e \in \gamma$$

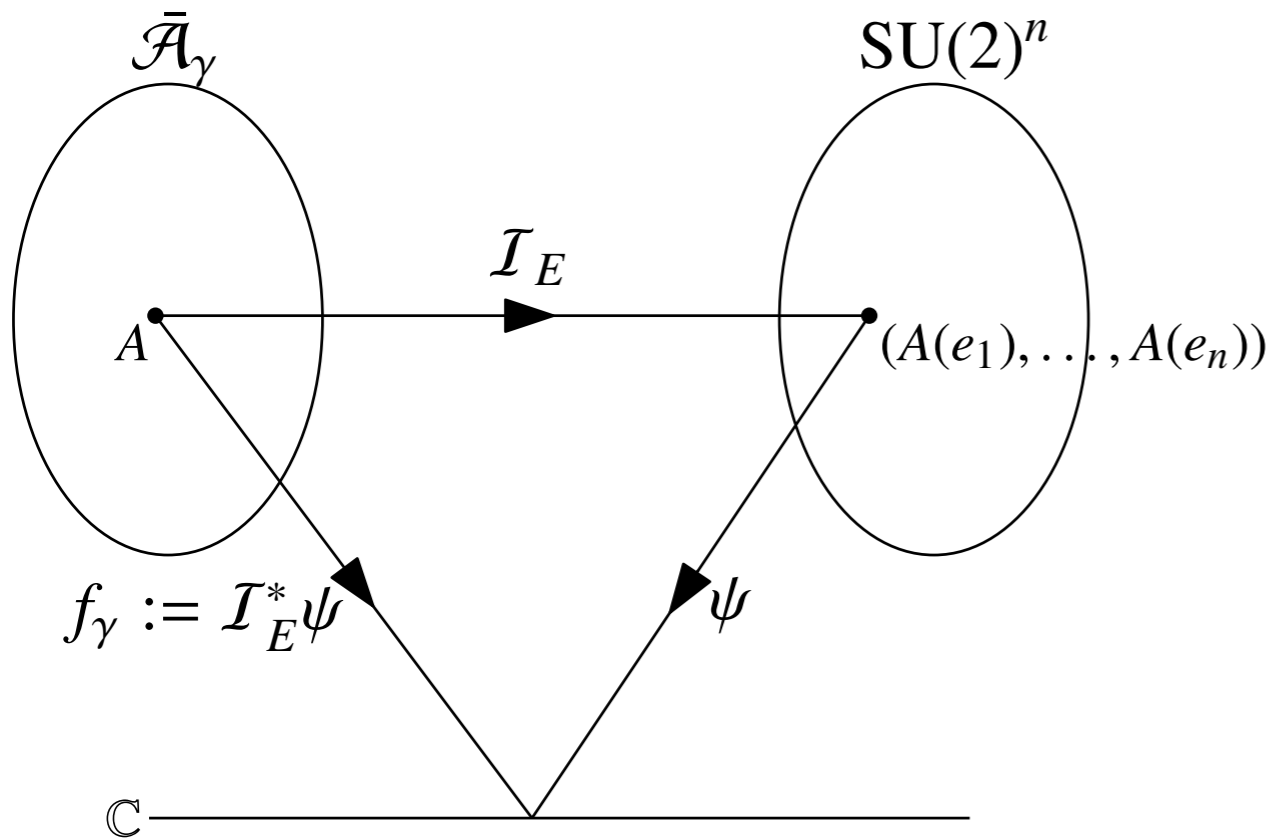
(ii)  $f_\gamma(A) = \psi(A(e_1), \dots, A(e_n)) = \psi(h_{e_1}(A), \dots, h_{e_n}(A)) \equiv f_\gamma(h_{e_1}(A), \dots, h_{e_n}(A))$

### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

Step I: a natural measure  $d\mu_\gamma$  on  $\bar{\mathcal{A}}_\gamma$

There exists a bijection

$$\begin{aligned} \mathcal{I}_E : \bar{\mathcal{A}}_\gamma &\rightarrow \text{SU}(2)^n \equiv \text{SU}(2) \times \cdots \times \text{SU}(2) \\ A &\mapsto (A(e_1), \dots, A(e_n)) \end{aligned}$$



$$\int_{\bar{\mathcal{A}}_\gamma} f_\gamma(A) d\mu_\gamma(A) := \int_{\text{SU}(2)^n} \psi(A(e_1), \dots, A(e_n)) \prod_{i=1}^n d\mu_H(A(e_i))$$

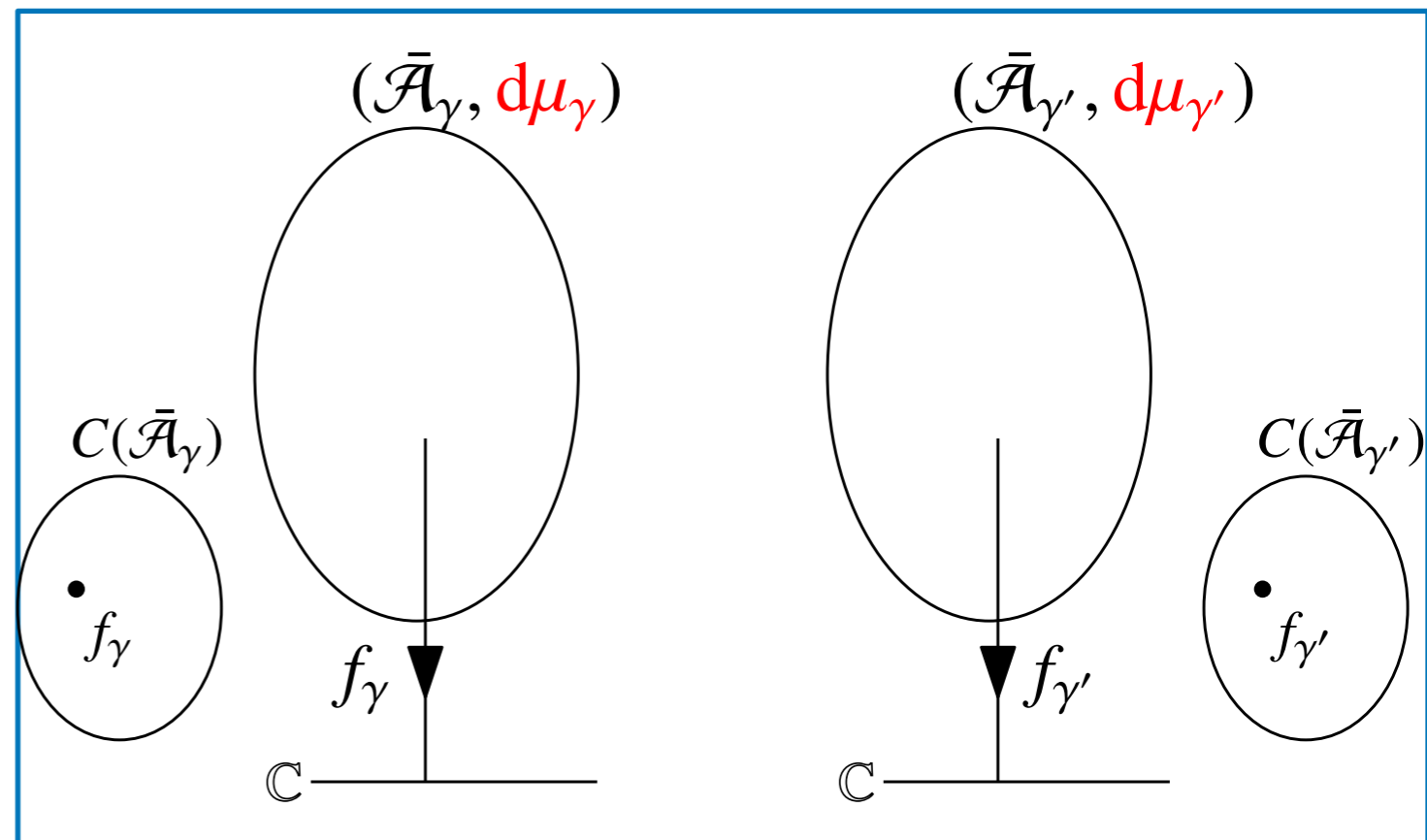
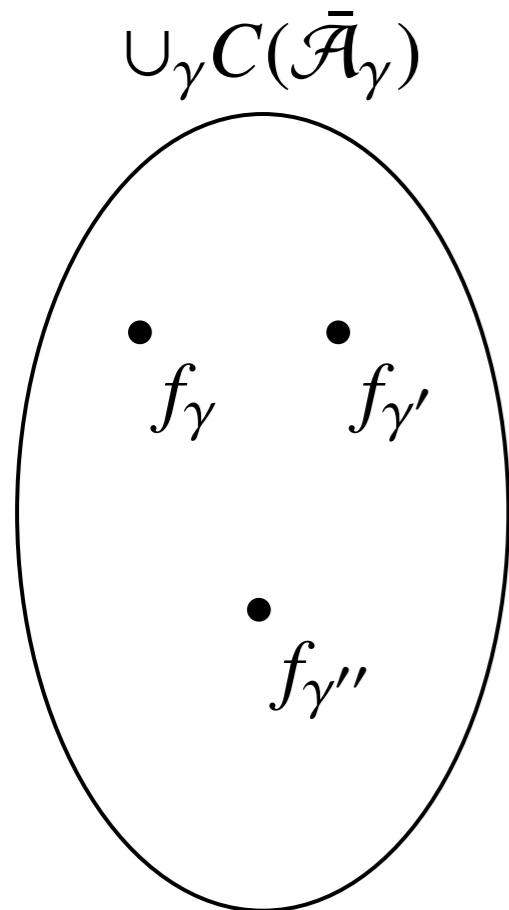
$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

glue  
 $(\bar{\mathcal{A}}_\gamma, d\mu_\gamma)$

### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

Step II: a natural measure  $d\mu_{AL}$  on  $\bar{\mathcal{A}}$  defined by  $d\mu_\gamma$  on  $\bar{\mathcal{A}}_\gamma$

$C(\bar{\mathcal{A}}_\gamma)$ : the set of all continuous complex functions on  $\bar{\mathcal{A}}_\gamma$



$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

$\downarrow$  ✓  
 $(\bar{\mathcal{A}}_\gamma, d\mu_\gamma)$  ✓  
 $\uparrow$  glue

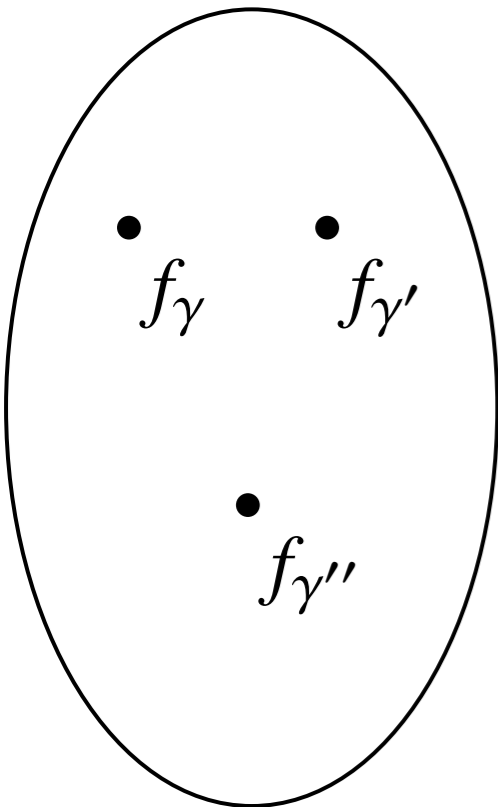
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Step II: a natural measure  $d\mu_{AL}$  on  $\bar{\mathcal{A}}$  defined by  $d\mu_\gamma$  on  $\bar{\mathcal{A}}_\gamma$

For  $\gamma' \geq \gamma$ , there exists a surjective map (projective map)

$$P_{\gamma'\gamma} : \bar{\mathcal{A}}_{\gamma'} \rightarrow \bar{\mathcal{A}}_\gamma, \quad \text{restricting the domain from } \gamma' \text{ to } \gamma$$

$\cup_\gamma C(\bar{\mathcal{A}}_\gamma)$



$$\forall \gamma'' \geq \gamma, \gamma', \quad P_{\gamma''\gamma'}^* f_{\gamma'} = P_{\gamma''\gamma}^* f_\gamma \quad \longrightarrow \quad \int_{\bar{\mathcal{A}}_{\gamma'}} f_{\gamma'}(A) d\mu_{\gamma'}(A) = \int_{\bar{\mathcal{A}}_\gamma} f_\gamma(A) d\mu_\gamma(A)$$

Hence, these two functions are said to be equivalent

$$f_\gamma \sim f_{\gamma'}$$

### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

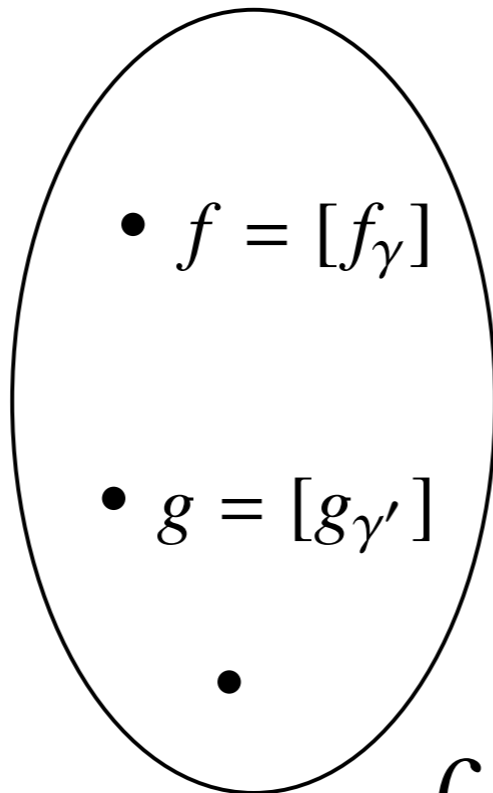
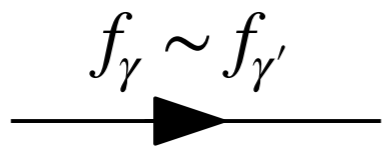
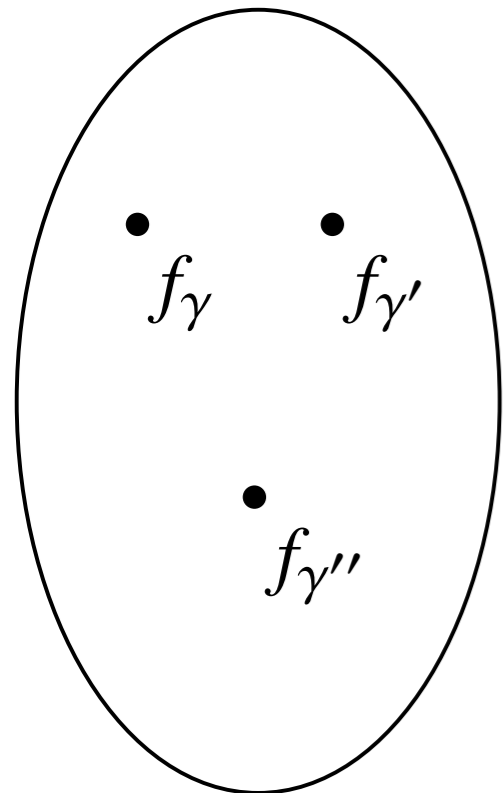
Step II: a natural measure  $d\mu_{AL}$  on  $\bar{\mathcal{A}}$  defined by  $d\mu_\gamma$  on  $\bar{\mathcal{A}}_\gamma$

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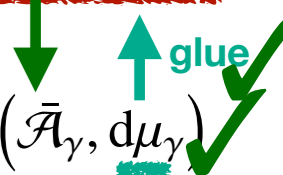
$$\text{Cyl}(\bar{\mathcal{A}}) := \cup_\gamma C(\bar{\mathcal{A}}_\gamma) / \sim$$



$$\int_{\bar{\mathcal{A}}} f(A) d\mu_{AL}(A) := \int_{\bar{\mathcal{A}}_\gamma} f_\gamma(A) d\mu_\gamma(A)$$

$$\langle g, f \rangle_{\text{Cyl}(\bar{\mathcal{A}})} := \int_{\bar{\mathcal{A}}} \overline{g(A)} f(A) d\mu_{AL}(A) := \int_{\bar{\mathcal{A}}_{\gamma''}} \overline{(P_{\gamma''\gamma'}^* g_{\gamma'})}(A) (P_{\gamma''\gamma}^* f_\gamma)(A) d\mu_{\gamma''}(A)$$

$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H) \xrightarrow{?} \mathcal{H}_{\text{kin}} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$



### 3. A natural measure $d\mu_{AL}$ on $\bar{\mathcal{A}}$

$d\mu_{AL}$  : The Ashtekar-Lewandowski measure.

**Thm.** *The Ashtekar-Lewandowski measure is invariant under internal gauge transformations  $g(x)$  and spatial diffeomorphisms  $\varphi$ , i.e.,  $\forall f \in \text{Cyl}(\bar{\mathcal{A}})$ ,*

$$\int_{\bar{\mathcal{A}}} g \circ f(A) d\mu_{AL}(A) = \int_{\bar{\mathcal{A}}} f(A) d\mu_{AL}(A),$$
$$\int_{\bar{\mathcal{A}}} \varphi \circ f(A) d\mu_{AL}(A) = \int_{\bar{\mathcal{A}}} f(A) d\mu_{AL}(A).$$

# 4. The kinematical Hilbert space $\mathcal{H}_{\text{kin}}$

$$\mathcal{H}_{\text{kin}} := \overline{\text{Cyl}(\bar{\mathcal{A}})} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

Spin network states: the basis of  $\mathcal{H}_{\text{kin}}$

$$\mathcal{H}_{\text{kin}} := \overline{\text{Cyl}(\bar{\mathcal{A}})} = L^2(\bar{\mathcal{A}}, d\mu_{AL})$$

Spin network states

$$T_{\gamma, \vec{j}, \vec{m}, \vec{n}}(A) := \prod_{e \in E(\gamma)} \sqrt{d_{j_e}} [\pi_{j_e}(A(e))]^{m_e}_{n_e}$$

$$\langle T_{\gamma', \vec{j}', \vec{m}', \vec{n}'}, T_{\gamma, \vec{j}, \vec{m}, \vec{n}} \rangle_{\mathcal{H}_{\text{kin}}} = \delta_{\gamma, \gamma'} \delta_{\vec{j}', \vec{j}} \delta_{\vec{m}', \vec{m}} \delta_{\vec{n}', \vec{n}}$$

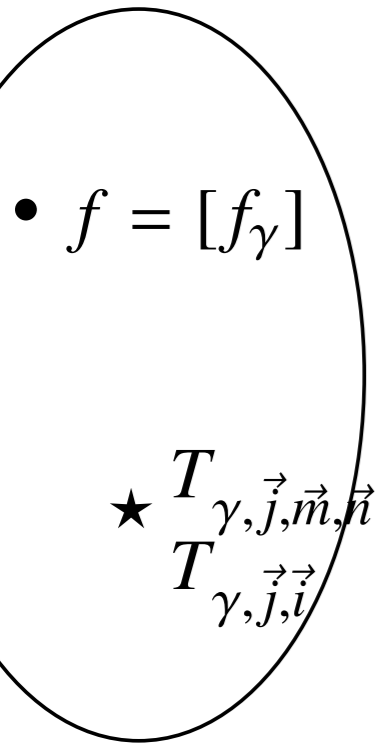
$$T_{\gamma, \vec{j}, \vec{i}}(A) = \bigotimes_{v \in V(\gamma)} i_v \cdot \bigotimes_{e \in E(\gamma)} \sqrt{d_{j_e}} \pi_{j_e}(A(e))$$

$$\langle T_{\gamma', \vec{j}', \vec{i}'}, T_{\gamma, \vec{j}, \vec{i}} \rangle_{\mathcal{H}_{\text{kin}}} = \delta_{\gamma, \gamma'} \delta_{\vec{j}', \vec{j}} \delta_{\vec{i}', \vec{i}}$$

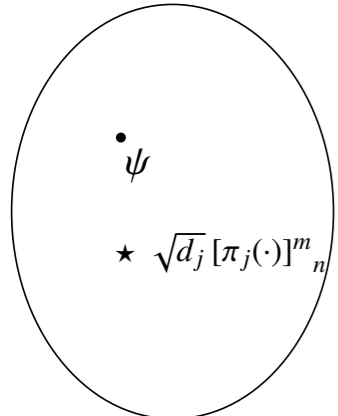
$$i_v \sim \left( i_{v; j_1 \dots j_n}^{J; \vec{a}} \right)_{m_1 \dots m_n}^M \equiv (-1)^{j_1 - \sum_{i=2}^n j_i - J} \sum_{k_2, \dots, k_{n-1}} \langle a_2 k_2 | j_1 m_1 j_2 m_2 \rangle \dots \langle JM | a_{n-1} k_{n-1} j_n m_n \rangle$$

Gauge variant:  $i_v \sim \left( i_{v; j_1 \dots j_n}^{J \neq 0; \vec{a}} \right)_{m_1 \dots m_n}^{M \neq 0}$

Gauge invariant:  $i_v \sim \left( i_{v; j_1 \dots j_n}^{J=0; \vec{a}} \right)_{m_1 \dots m_n}^{M=0} \equiv \left( i_{v; j_1 \dots j_n}^{J=0; \vec{a}} \right)_{m_1 \dots m_n}$



$$\mathcal{H}_{\text{SU}(2)} = L^2(\text{SU}(2), d\mu_H)$$



## 5. The elementary operators on $\mathcal{H}_{\text{kin}}$

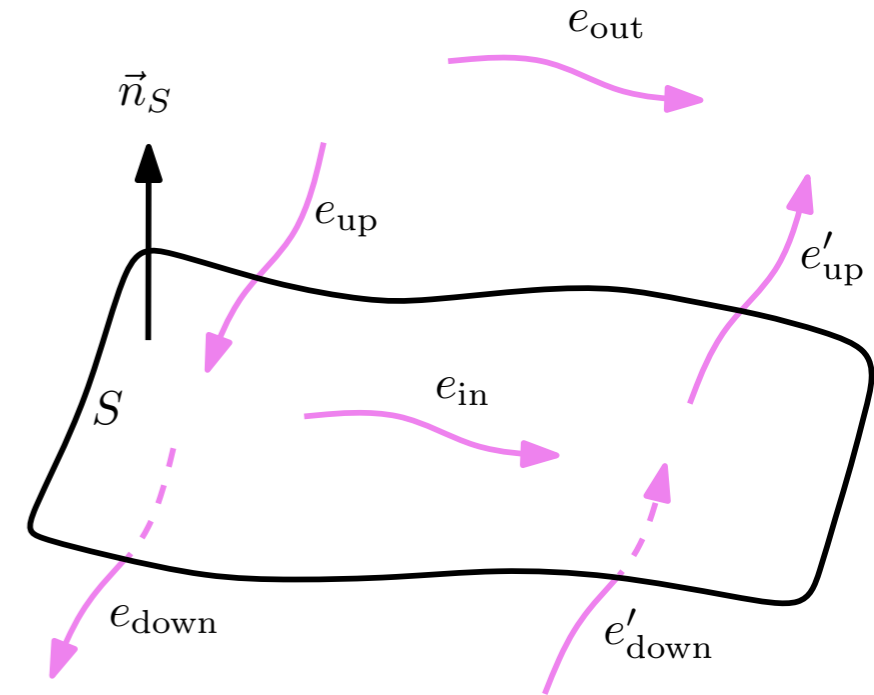
$$\widehat{A(e)} \cdot f_\gamma(A) := A(e) f_\gamma(A)$$

$$\begin{aligned} \widehat{\tilde{E}_i(S)} \cdot f_\gamma(A) &:= -i\hbar \{f_\gamma(A), \tilde{E}_i(S)\} \\ &= \frac{\hbar\kappa\beta}{2} \sum_{v \in V(\gamma)} \left( \sum_{b(e)=v} \kappa(e, S) J_{e,R}^i + \sum_{f(e)=v} \kappa(e, S) J_{e,L}^i \right) f_\gamma(A) \end{aligned}$$

$$J_{e,R}^i := iR_e^{(\tau_i)}, \quad J_{e,L}^i := iL_e^{(\tau_i)}$$

$$\{h_e(A), \tilde{E}_i(S)\} = -\frac{\kappa\beta}{2} \kappa(e, S) \left( \delta_{e \cap S, b(e)} R_e^{(\tau_i)} + \delta_{e \cap S, f(e)} L_e^{(\tau_i)} \right) h_e(A)$$

$$\text{Steric factor: } \kappa(e, S) := \begin{cases} +1 & e \text{ is the 'up' type w.r.t. } S \\ -1 & e \text{ is the 'down' type w.r.t. } S \\ 0 & e \text{ is the 'inside/outside' type w.r.t. } S \end{cases}$$



*Thanks for your attention!*