Statistical Methods in High Energy Physics -1

Outline

- Probability
- Random variables, probability densities
- Probability densities in HEP
- Parameter estimation, hypothesis tests
- Maximum likelihood and least squares

Ref: Detection and Estimation Theory by Van Trees; Statistical data analysis by Cowan

Theory ↔ Statistics ↔ Experiment

Theory (model, hypothesis):

Experiment (observation):



Data analysis in particle physics



Observe events of a certain type

Measure characteristics of each event (particle momenta, number of muons, energy of jets,...) Theories (e.g. SM) predict distributions of these properties up to free parameters, e.g., α , $G_{\rm F}$, $M_{\rm Z}$, $\alpha_{\rm s}$, $m_{\rm H}$, ... Some tasks of data analysis:

Estimate (measure) the parameters;

Quantify the uncertainty of the parameter estimates;

Test the extent to which the predictions of a theory are in agreement with the data (\rightarrow presence of New Physics?)

Definition of probability
Consider a set S with subsets A, B, ...
For all
$$A \subset S, P(A) \ge 0$$

 $P(S) = 1$
If $A \cap B = \emptyset, P(A \cup B) = P(A) + P(B)$



Kolmogorov axioms (1933)

From these axioms we can derive further properties, e.g.

$$P(\overline{A}) = 1 - P(A)$$

$$P(A \cup \overline{A}) = 1$$

$$P(\emptyset) = 0$$

if $A \subset B$, then $P(A) \le P(B)$

$$P(A \cup B) = P(A) + P(B) - P(A \cap B)$$

Conditional probability, independence

Also define conditional probability of *A* given *B* (with $P(B) \neq 0$):

$$P(A|B) = \frac{P(A \cap B)}{P(B)}$$

E.g. rolling dice: $P(n < 3 | n \text{ even}) = \frac{P((n < 3) \cap n \text{ even})}{P(\text{even})} = \frac{1/6}{3/6} = \frac{1}{3}$

Subsets A, B independent if: $P(A \cap B) = P(A)P(B)$

If A, B independent,
$$P(A|B) = \frac{P(A)P(B)}{P(B)} = P(A)$$

do not confuse with disjoint subsets, i.e., $A \cap B = \emptyset$

Interpretation of probability

I. Relative frequency

A, B, ... are outcomes of a repeatable experiment

 $P(A) = \lim_{n \to \infty} \frac{\text{times outcome is } A}{n}$

quantum mechanics, particle scattering, radioactive decay...

II. Subjective probability (Bayesian)

A, B, ... are hypotheses (statements that are true or false)

P(A) =degree of belief that A is true

- Both interpretations consistent with Kolmogorov axioms.
- In particle physics frequency interpretation often most useful, but subjective probability can provide more natural treatment of non-repeatable phenomena.

Bayes' theorem

From the definition of conditional probability we have,

$$P(A|B) = \frac{P(A \cap B)}{P(B)}$$
 and $P(B|A) = \frac{P(B \cap A)}{P(A)}$

but $P(A \cap B) = P(B \cap A)$, so

$$P(A|B) = \frac{P(B|A)P(A)}{P(B)}$$

First published (posthumously) by the Reverend Thomas Bayes (1702–1761)

An essay towards solving a problem in the doctrine of chances, Philos. Trans. R. Soc. **53** (1763) 370; reprinted in Biometrika, **45** (1958) 293.

Bayes' theorem





Bayes' theorem becomes

$$P(A|B) = \frac{P(B|A)P(A)}{\sum_{i} P(B|A_i)P(A_i)}$$

Bayesian Statistics – general philosophy

In Bayesian statistics, use subjective probability for hypotheses:

probability of the data assuming hypothesis *H* (the likelihood) prior probability, i.e., before seeing the data $P(H|\vec{x}) = \frac{P(\vec{x}|H)\pi(H)}{\int P(\vec{x}|H)\pi(H) dH}$ posterior probability, i.e., after seeing the data over all possible hypotheses

Bayes' theorem has an "if-then" character: If your prior probabilities were $\pi(H)$, then it says how these probabilities should change in the light of the data.

No unique prescription for priors (subjective!)

An example using Bayes' theorem

Suppose the probability (for anyone) to have a disease D is:

P(D)	=	0.001	← prior probabilities, i.e.,
P(no D)	=	0.999	before any test carried out

Consider a test for the disease: result is + or -

P(+|D) = 0.98 P(-|D) = 0.02 ← probabilities to (in)correctly identify a person with the disease

$$P(+|\text{no D}) = 0.03$$
 \leftarrow probabilities to (in)correctly
 $P(-|\text{no D}) = 0.97$ \leftarrow probabilities to (in)correctly

Suppose your result is +. How worried should you be?

Bayes' theorem example (cont.)

The probability to have the disease given a + result is

$$p(\mathbf{D}|+) = \frac{P(+|\mathbf{D})P(\mathbf{D})}{P(+|\mathbf{D})P(\mathbf{D}) + P(+|\mathrm{no}|\mathbf{D})P(\mathrm{no}|\mathbf{D})}$$

$= \frac{0.98 \times 0.001}{0.98 \times 0.001 + 0.03 \times 0.999}$

$$= 0.032 \leftarrow \text{posterior probability}$$

i.e. you're probably OK!

Your viewpoint: my degree of belief that I have the disease is 3.2%. Your doctor's viewpoint: 3.2% of people like this have the disease.

Random variables and probability density functions

A random variable is a numerical characteristic assigned to an element of the sample space; can be discrete or continuous.

Suppose outcome of experiment is continuous value *x*

P(x found in [x, x + dx]) = f(x) dx

 \rightarrow *f*(*x*) = probability density function (pdf)

 $\int_{-\infty}^{\infty} f(x) \, dx = 1 \qquad x \text{ must be somewhere}$

Or for discrete outcome x_i with e.g. i = 1, 2, ... we have

$$P(x_i) = p_i$$
probability mass function $\sum_i P(x_i) = 1$ x must take on one of its possible values

Cumulative distribution function

Probability to have outcome less than or equal to x is

 $\int_{-\infty}^{x} f(x') dx' \equiv F(x)$ cumulative distribution function



Other types of probability densities

Outcome of experiment characterized by several values, e.g. an *n*-component vector, $(x_1, \dots x_n)$

 \rightarrow joint pdf $f(x_1,\ldots,x_n)$

Sometimes we want only pdf of some (or one) of the components \rightarrow marginal pdf $f_1(x_1) = \int \cdots \int f(x_1, \dots, x_n) dx_2 \dots dx_n$ x_1, x_2 independent if $f(x_1, x_2) = f_1(x_1) f_2(x_2)$

Sometimes we want to consider some components as constant

$$\rightarrow$$
 conditional pdf $g(x_1|x_2) = \frac{f(x_1, x_2)}{f_2(x_2)}$

Expectation values

Consider continuous r.v. x with pdf f(x). Define expectation (mean) value as $E[x] = \int x f(x) dx$ Notation (often): $E[x] = \mu \sim$ "centre of gravity" of pdf. For a function y(x) with pdf g(y),

$$E[y] = \int y g(y) dy = \int y(x) f(x) dx$$
 (equivalent)

Variance: $V[x] = E[x^2] - \mu^2 = E[(x - \mu)^2]$

Notation: $V[x] = \sigma^2$ Standard deviation: $\sigma = \sqrt{\sigma^2}$

 σ ~ width of pdf, same units as *x*.



Covariance and correlation

Define covariance cov[x,y] (also use matrix notation V_{xy}) as

$$\operatorname{cov}[x,y] = E[xy] - \mu_x \mu_y = E[(x - \mu_x)(y - \mu_y)]$$

Correlation coefficient (dimensionless) defined as

$$\rho_{xy} = \frac{\operatorname{COV}[x, y]}{\sigma_x \sigma_y}$$

If x, y, independent, i.e., $f(x, y) = f_x(x)f_y(y)$, then

$$E[xy] = \int \int xy f(x,y) \, dx \, dy = \mu_x \mu_y$$

 \rightarrow cov[x, y] = 0 x and y, 'uncorrelated'

- converse not always true.

Correlation (cont.)



Covariance matrix

Suppose we have a set of *n* random variables, say, $x_1, ..., x_n$. We can write the covariance of each pair as an *n* x *n* matrix:

$$V_{ij} = \operatorname{cov}[x_i, x_j] = \rho_{ij}\sigma_i\sigma_j$$

$$V = \begin{pmatrix} \sigma_1^2 & \rho_{12}\sigma_1\sigma_2 & \dots & \rho_{1n}\sigma_1\sigma_n \\ \rho_{21}\sigma_2\sigma_1 & \sigma_2^2 & \dots & \rho_{2n}\sigma_2\sigma_n \\ \vdots & & \ddots & \vdots \\ \rho_{n1}\sigma_n\sigma_1 & \rho_{n2}\sigma_n\sigma_2 & \dots & \sigma_n^2 \end{pmatrix}$$

Covariance matrix is: symmetric, diagonal = variances, positive semi-definite: $z^T V z \ge 0$ for all $z \in \mathbb{R}^n$

Correlation matrix

Closely related to the covariance matrix is the *n* x *n* matrix of correlation coefficients:

$$\rho_{ij} = \frac{\operatorname{cov}[x_i, x_j]}{\sigma_i \sigma_j}$$

$$\rho = \begin{pmatrix} 1 & \rho_{12} & \dots & \rho_{1n} \\ \rho_{21} & 1 & \dots & \rho_{2n} \\ \vdots & \ddots & \vdots \\ \rho_{n1} & \rho_{n2} & \dots & 1 \end{pmatrix}$$
By construction, diagonal elements are $\rho_{ii} = 1$

Distribution/pdf	Use in HEP
Binomial	Branching ratio
Multinomial	Histogram with fixed N
Poisson	Number of events
Uniform	Monte Carlo method
Exponential	Decay time
Gaussian	Measurement error
Chi-square	Goodness-of-fit
Cauchy	Mass of resonance
Landau	Ionization energy loss

Binomial distribution

Consider *N* independent experiments (Bernoulli trials): outcome of each is 'success' or 'failure', probability of success on any given trial is *p*. Define discrete r.v. n = number of successes ($0 \le n \le N$).

Probability of a specific outcome (in order), e.g. 'ssfsf' is

$$pp(1-p)p(1-p) = p^n(1-p)^{N-n}$$

N!

But order not important; there are

$$\overline{n!(N-n)!}$$

ways (permutations) to get *n* successes in *N* trials, total probability for *n* is sum of probabilities for each permutation.

Binomial distribution (2)

The binomial distribution is therefore

$$f(n; N, p) = \frac{N!}{n!(N-n)!}p^n(1-p)^{N-n}$$
random parameters
variable

For the expectation value and variance we find:

$$E[n] = \sum_{n=0}^{N} nf(n; N, p) = Np$$
$$V[n] = E[n^{2}] - (E[n])^{2} = Np(1-p)$$

Binomial distribution (3)

Binomial distribution for several values of the parameters:



Example: observe *N* decays of W^{\pm} , the number *n* of which are $W \rightarrow \mu \nu$ is a binomial r.v., *p* = branching ratio.

Multinomial distribution

Like binomial but now *m* outcomes instead of two, probabilities are

$$\vec{p} = (p_1, \dots, p_m)$$
, with $\sum_{i=1}^m p_i = 1$.

For *N* trials we want the probability to obtain:

 n_1 of outcome 1, n_2 of outcome 2,

 n_m of outcome m.

. . .

This is the multinomial distribution for $\vec{n} = (n_1, \dots, n_m)$

$$f(\vec{n}; N, \vec{p}) = \frac{N!}{n_1! n_2! \cdots n_m!} p_1^{n_1} p_2^{n_2} \cdots p_m^{n_m}$$

Multinomial distribution (2)

Now consider outcome *i* as 'success', all others as 'failure'.

 \rightarrow all n_i individually binomial with parameters N, p_i

$$E[n_i] = Np_i, \quad V[n_i] = Np_i(1-p_i) \quad \text{for all } i$$

One can also find the covariance to be

$$V_{ij} = Np_i(\delta_{ij} - p_j)$$

Example: $\vec{n} = (n_1, \dots, n_m)$ represents a histogram with *m* bins, *N* total entries, all entries independent.

Poisson distribution

Consider binomial *n* in the limit

$$N \to \infty, \qquad p \to 0, \qquad E[n] = Np \to \nu.$$

 \rightarrow *n* follows the Poisson distribution:

$$f(n;\nu) = \frac{\nu^n}{n!} e^{-\nu} \quad (n \ge 0)$$

$$E[n] = \nu$$
, $V[n] = \nu$.

Example: number of scattering events *n* with cross section σ found for a fixed integrated luminosity, with $\nu = \sigma \int L dt$.



Uniform distribution

Consider a continuous r.v. x with $-\infty < x < \infty$. Uniform pdf is:



For any r.v. x with cumulative distribution F(x), y = F(x) is uniform in [0,1].

Example: for $\pi^0 \to \gamma\gamma$, E_{γ} is uniform in $[E_{\min}, E_{\max}]$, with $E_{\min} = \frac{1}{2} E_{\pi} (1 - \beta)$, $E_{\max} = \frac{1}{2} E_{\pi} (1 + \beta)$

Exponential distribution

The exponential pdf for the continuous r.v. *x* is defined by:



Example: proper decay time *t* of an unstable particle

 $f(t;\tau) = \frac{1}{\tau}e^{-t/\tau}$ (τ = mean lifetime)

Lack of memory (unique to exponential): $f(t - t_0 | t \ge t_0) = f(t)$

Gaussian distribution

The Gaussian (normal) pdf for a continuous r.v. *x* is defined by:

$$f(x; \mu, \sigma) = \frac{1}{\sqrt{2\pi\sigma}} e^{-(x-\mu)^2/2\sigma^2}$$

$$E[x] = \mu$$

$$F[x] = \mu$$

$$F[x] = \mu$$

$$F[x] = \mu$$

$$F[x] = \sigma^2$$

$$F[x]$$

x

Special case: $\mu = 0$, $\sigma^2 = 1$ ('standard Gaussian'):

$$\varphi(x) = \frac{1}{\sqrt{2\pi}} e^{-x^2/2} , \quad \Phi(x) = \int_{-\infty}^x \varphi(x') \, dx'$$

If $y \sim \text{Gaussian}$ with μ , σ^2 , then $x = (y - \mu) / \sigma$ follows $\varphi(x)$.

Gaussian pdf and the Central Limit Theorem

The Gaussian pdf is so useful because almost any random variable that is a sum of a large number of small contributions follows it. This follows from the Central Limit Theorem:

For *n* independent r.v.s x_i with finite variances σ_i^2 , otherwise arbitrary pdfs, consider the sum

$$y = \sum_{i=1}^{n} x_i$$

In the limit $n \to \infty$, y is a Gaussian r.v. with

$$E[y] = \sum_{i=1}^{n} \mu_i \qquad V[y] = \sum_{i=1}^{n} \sigma_i^2$$

Measurement errors are often the sum of many contributions, so frequently measured values can be treated as Gaussian r.v.s.

Multivariate Gaussian distribution

Multivariate Gaussian pdf for the vector $\vec{x} = (x_1, \dots, x_n)$:

$$f(\vec{x};\vec{\mu},V) = \frac{1}{(2\pi)^{n/2}|V|^{1/2}} \exp\left[-\frac{1}{2}(\vec{x}-\vec{\mu})^T V^{-1}(\vec{x}-\vec{\mu})\right]$$

 $\vec{x}, \vec{\mu}$ are column vectors, $\vec{x}^T, \vec{\mu}^T$ are transpose (row) vectors,

$$E[x_i] = \mu_i, , \quad \text{cov}[x_i, x_j] = V_{ij}.$$

For n = 2 this is

$$f(x_1, x_2; \mu_1, \mu_2, \sigma_1, \sigma_2, \rho) = \frac{1}{2\pi\sigma_1\sigma_2\sqrt{1-\rho^2}} \\ \times \exp\left\{-\frac{1}{2(1-\rho^2)} \left[\left(\frac{x_1-\mu_1}{\sigma_1}\right)^2 + \left(\frac{x_2-\mu_2}{\sigma_2}\right)^2 - 2\rho\left(\frac{x_1-\mu_1}{\sigma_1}\right)\left(\frac{x_2-\mu_2}{\sigma_2}\right) \right] \right\}$$

where $\rho = \operatorname{cov}[x_1, x_2]/(\sigma_1 \sigma_2)$ is the correlation coefficient.

Chi-square (χ^2) distribution

The chi-square pdf for the continuous r.v. z ($z \ge 0$) is defined by

$$f(z;n) = \frac{1}{2^{n/2}\Gamma(n/2)} z^{n/2-1} e^{-z/2} \quad \underbrace{\underbrace{\underbrace{}}_{04}^{0.5}}_{n_4} \underbrace{\underbrace{}_{04}^{-n_{-1}}}_{n_{-1}} \underbrace{\underbrace{}_{04}^{-n$$

n = 1, 2, ... = number of 'degrees of freedom' (dof)

$$E[z] = n, \quad V[z] = 2n.$$



For independent Gaussian x_i , i = 1, ..., n, means μ_i , variances σ_i^2 ,

$$z = \sum_{i=1}^{n} \frac{(x_i - \mu_i)^2}{\sigma_i^2} \quad \text{follows } \chi^2 \text{ pdf with } n \text{ dof.}$$

Example: goodness-of-fit test variable especially in conjunction with method of least squares.

Landau distribution

For a charged particle with $\beta = v/c$ traversing a layer of matter of thickness *d*, the energy loss Δ follows the Landau pdf:



L. Landau, J. Phys. USSR **8** (1944) 201; see also W. Allison and J. Cobb, Ann. Rev. Nucl. Part. Sci. **30** (1980) 253.

Example: decay of an unstable particle

As an example that we'll use to illustrate several statistical methods, consider measuring the proper decay time of an unstable particle such as a B meson:

R. Barate et al. / Physics Letters B 492 (2000) 259-274



Measure flight distance d and momentum p of decay products of B meson with mass $m_{\rm B}$.

These are related to the proper decay time t_p (time in B rest frame) by

$$d = vt_{
m lab} = eta c imes \gamma t_{
m p} = rac{p_{
m B}}{m_{
m B}} t_{
m p}$$

so $t_{
m p} = rac{m_{
m B} d}{p_{
m B}}$

Exponential pdf for proper decay time

We can model *t* as following an exponential pdf:



We can show (exercise) that the mean and variance of *t* are:

$$E[t] = \int_0^\infty t f(t;\tau) \, dt = \tau \qquad \qquad V[t] = E[t^2] - (E[t])^2 = \tau^2$$

Frequentist hypothesis tests

Suppose a measurement produces data x; consider a hypothesis H_0 we want to test and alternative H_1

 H_0 , H_1 specify probability for \mathbf{x} : $P(\mathbf{x}|H_0)$, $P(\mathbf{x}|H_1)$

A test of H_0 is defined by specifying a critical region w of the data space such that there is no more than some (small) probability α , assuming H_0 is correct, to observe the data there, i.e.,

 $P(\mathbf{x} \in w \mid H_0) \le \alpha$

Need inequality if data are discrete.

 α is called the size or significance level of the test.

If x is observed in the critical region, reject H_0 .



Definition of a test (2)

But in general there are an infinite number of possible critical regions that give the same size α .

Use the alternative hypothesis H_1 to motivate where to place the critical region.

Roughly speaking, place the critical region where there is a low probability (α) to be found if H_0 is true, but high if H_1 is true:



Example of a test for classification



For each event in a mixture of signal (s) and background (b) test

 H_0 : event is of type b

using a critical region W of the form: $W = \{x : x \le x_c\}$, where x_c is a constant that we choose to give a test with the desired size α .

Classification example (2)

Suppose we want $\alpha = 10^{-4}$. Require:

$$\alpha = P(x \le x_{c}|b) = \int_{0}^{x_{c}} f(x|b) dx = \frac{4x^{4}}{4} \Big|_{0}^{x_{c}} = x_{c}^{4}$$

and therefore $x_{\rm c} = \alpha^{1/4} = 0.1$

For this test (i.e. this critical region W), the power with respect to the signal hypothesis (s) is

$$M = P(x \le x_{\rm c}|{\rm s}) = \int_0^{x_{\rm c}} f(x|{\rm s}) \, dx = 2x_{\rm c} - x_{\rm c}^2 = 0.19$$

Note: the optimal size and power is a separate question that will depend on goals of the subsequent analysis.

Parameter estimation

The parameters of a pdf are constants that characterize its shape, e.g. $1 - \frac{1}{2}$



random variable parameter

Suppose we have a sample of observed values: $\vec{x} = (x_1, \ldots, x_n)$

We want to find some function of the data to estimate the parameter(s):

 $\hat{\theta}(\vec{x}) \leftarrow \text{estimator written with a hat}$

Sometimes we say 'estimator' for the function of $x_1, ..., x_n$; 'estimate' for the value of the estimator with a particular data set.

Maximum Likelihood Estimator The likelihood function for i.i.d.* data

* i.i.d. = independent and identically distributed

Consider *n* independent observations of *x*: $x_1, ..., x_n$, where *x* follows $f(x; \theta)$. The joint pdf for the whole data sample is:

$$f(x_1,\ldots,x_n;\theta) = \prod_{i=1}^n f(x_i;\theta)$$

In this case the likelihood function is

$$L(\vec{\theta}) = \prod_{i=1}^{n} f(x_i; \vec{\theta}) \qquad (x_i \text{ constant})$$

Properties of estimators

If we were to repeat the entire measurement, the estimates from each would follow a pdf:



We want small (or zero) bias (systematic error): $b = E[\hat{\theta}] - \theta$

 \rightarrow average of repeated measurements should tend to true value.

And we want a small variance (statistical error): $V[\hat{\theta}]$

 \rightarrow small bias & variance are in general conflicting criteria

Maximum Likelihood Estimators (MLEs)

We *define* the maximum likelihood estimators or MLEs to be the parameter values for which the likelihood is maximum.



Could have multiple maxima (take highest).

MLEs not guaranteed to have any 'optimal' properties, (but in practice they're very good).

MLE example: parameter of exponential pdf

Consider exponential pdf, $f(t; \tau) = \frac{1}{\tau}e^{-t/\tau}$

and suppose we have i.i.d. data, t_1, \ldots, t_n

The likelihood function is
$$L(\tau) = \prod_{i=1}^{n} \frac{1}{\tau} e^{-t_i/\tau}$$

The value of τ for which $L(\tau)$ is maximum also gives the maximum value of its logarithm (the log-likelihood function):

$$\ln L(\tau) = \sum_{i=1}^{n} \ln f(t_i; \tau) = \sum_{i=1}^{n} \left(\ln \frac{1}{\tau} - \frac{t_i}{\tau} \right)$$

MLE example: parameter of exponential pdf (2)

Find its maximum by setting $\frac{\partial \ln L(\tau)}{\partial \tau} = 0$,

$$\rightarrow \hat{\tau} = \frac{1}{n} \sum_{i=1}^{n} t_i$$

Monte Carlo test: generate 50 values using $\tau = 1$:

We find the ML estimate:

$$\hat{\tau} = 1.062$$



MLE example: parameter of exponential pdf (3) For the exponential distribution one has for mean, variance:

$$E[t] = \int_0^\infty t \, \frac{1}{\tau} e^{-t/\tau} \, dt = \tau$$

$$V[t] = \int_0^\infty (t-\tau)^2 \, \frac{1}{\tau} e^{-t/\tau} \, dt = \tau^2$$
For the MLE $\hat{\tau} = \frac{1}{n} \sum_{i=1}^n t_i$ we therefore find
$$E[\hat{\tau}] = E\left[\frac{1}{n} \sum_{i=1}^n t_i\right] = \frac{1}{n} \sum_{i=1}^n E[t_i] = \tau \implies b = E[\hat{\tau}] - \tau = 0$$

$$V[\hat{\tau}] = V\left[\frac{1}{n} \sum_{i=1}^n t_i\right] = \frac{1}{n^2} \sum_{i=1}^n V[t_i] = \frac{\tau^2}{n} \implies \sigma_{\hat{\tau}} = \frac{\tau}{\sqrt{n}}$$

Variance of estimators from information inequality

The information inequality (RCF) sets a lower bound on the variance of any estimator (not only ML):

Often the bias b is small, and equality either holds exactly or is a good approximation (e.g. large data sample limit). Then,

$$V[\widehat{\theta}] \approx -1 \left/ E\left[\frac{\partial^2 \ln L}{\partial \theta^2}\right] \right.$$

Estimate this using the 2nd derivative of $\ln L$ at its maximum:

$$\widehat{V}[\widehat{\theta}] = -\left(\frac{\partial^2 \ln L}{\partial \theta^2}\right)^{-1}\Big|_{\theta=\widehat{\theta}}$$

MVB for MLE of exponential parameter

Find MVB =
$$-\left(1 + \frac{\partial b}{\partial \tau}\right)^2 / E\left[\frac{\partial^2 \ln L}{\partial \tau^2}\right]$$

We found for the exponential parameter the MLE $\hat{\tau} = \frac{1}{n} \sum_{i=1}^{n} t_i$

and we showed b = 0, hence $\partial b / \partial \tau = 0$.

We find
$$\frac{\partial^2 \ln L}{\partial \tau^2} = \sum_{i=1}^n \left(\frac{1}{\tau^2} - \frac{2t_i}{\tau^3} \right)$$

and since
$$E[t_i] = \tau$$
 for all i , $E\left[\frac{\partial^2 \ln L}{\partial \tau^2}\right] = -\frac{n}{\tau^2}$,
and therefore MVB = $\frac{\tau^2}{n} = V[\hat{\tau}]$. (Here MLE is "efficient").

The Monte Carlo method

What it is: a numerical technique for calculating probabilities and related quantities using sequences of random numbers.

g(r)

 \mathbf{O}

The usual steps:

- (1) Generate sequence $r_1, r_2, ..., r_m$ uniform in [0, 1].
- Use this to produce another sequence x₁, x₂, ..., x_n distributed according to some pdf f (x) in which we're interested (x can be a vector).
- (3) Use the *x* values to estimate some property of f(x), e.g., fraction of *x* values with a < x < b gives $\int_a^b f(x) dx$.
 - \rightarrow MC calculation = integration (at least formally)
- MC generated values = 'simulated data'
 - \rightarrow use for testing statistical procedures