

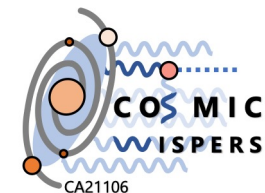
# Axion Physics and Detection

Andreas Ringwald  
PASCOS 2026  
University of Sheffield  
United Kingdom  
22-26 June 2026

**HELMHOLTZ** RESEARCH FOR  
GRAND CHALLENGES

**CLUSTER OF EXCELLENCE**  
QUANTUM UNIVERSE

 **cost**  
EUROPEAN COOPERATION  
IN SCIENCE & TECHNOLOGY



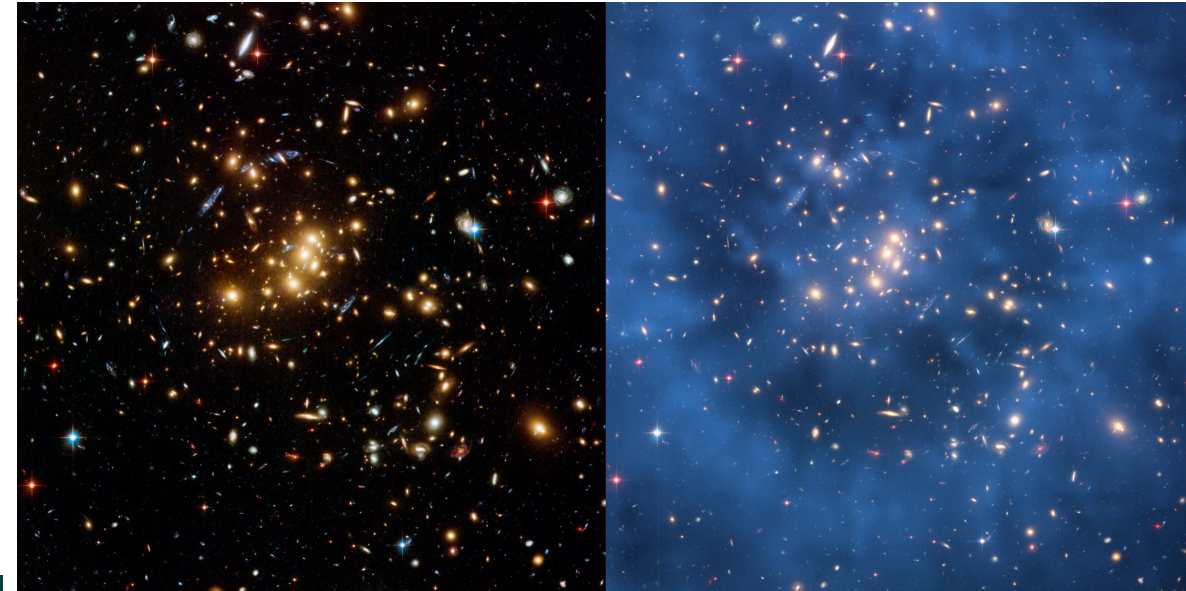
# The Axion

## Solving two puzzles in one go

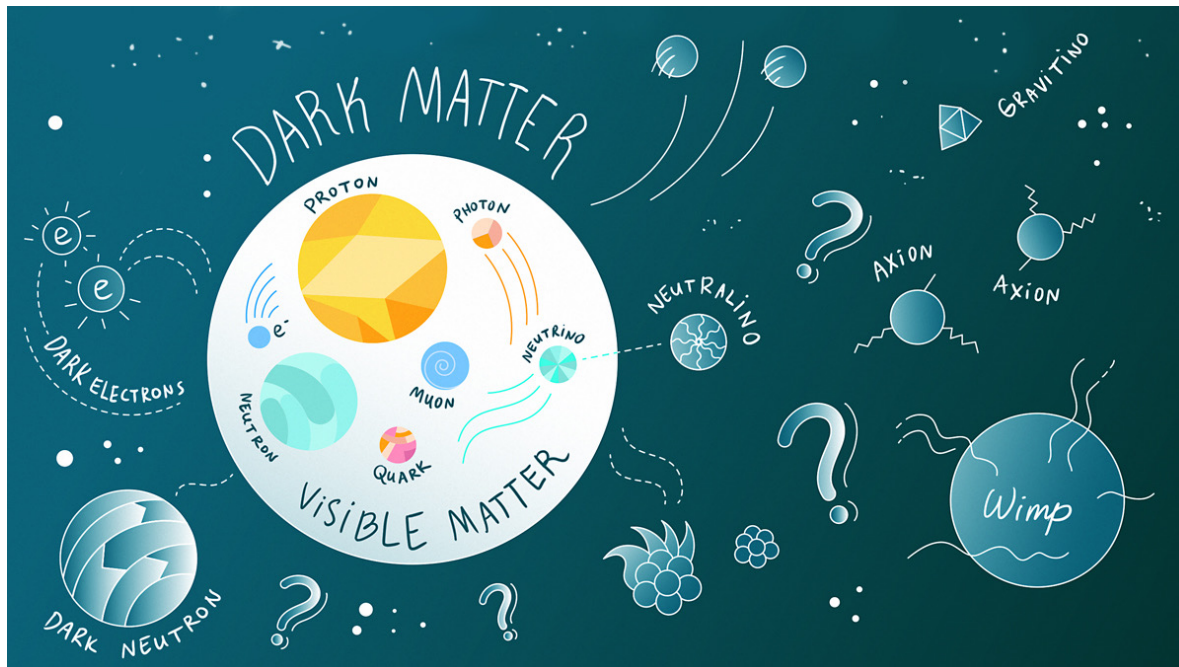
- Extensions of the Standard Model (SM) featuring the axion answer two fundamental questions in one go:

### 1. What is the nature of dark matter (DM)?

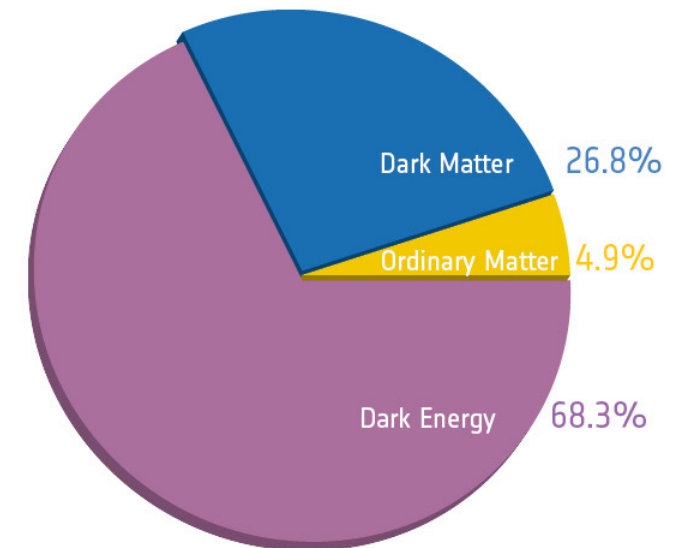
[Preskill,Wise,Wilczek 83; Abbott,Sikivie 83; Dine,Fischler 83]



<https://www.nasa.gov/content/discoveries-highlights-shining-a-light-on-dark-matter>



[Symmetry Magazine]



[PLANCK]

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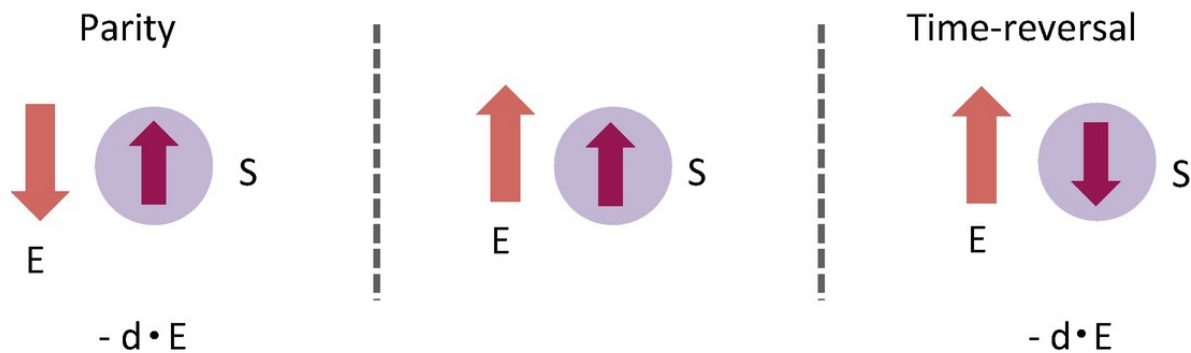
- Extensions of the Standard Model (SM) featuring the axion answer two fundamental questions in one go:
  1. **What is the nature of dark matter (DM)?**
  2. **Why do strong interactions conserve CP so accurately?**  
[Peccei,Quinn '77; Weinberg '78; Wilczek '78]

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    - Why is the electric dipole moment of the neutron so tiny?
      - Experiment:  $|d_n| < 1.8 \times 10^{-26} e \text{ cm}$  [Abel et al. 20]

Talk by Katia Michielsen on Thu 26/06 at 10:10 in Particles: Beyond the Standard Model & Flavour Physics parallel session

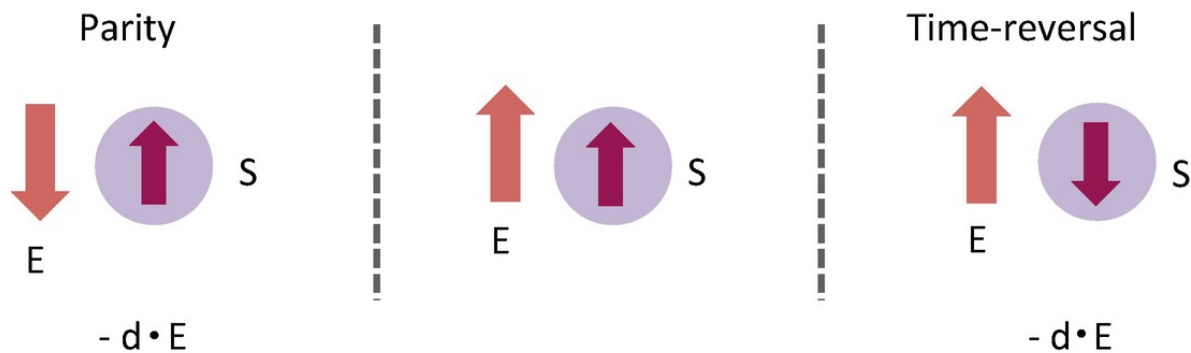


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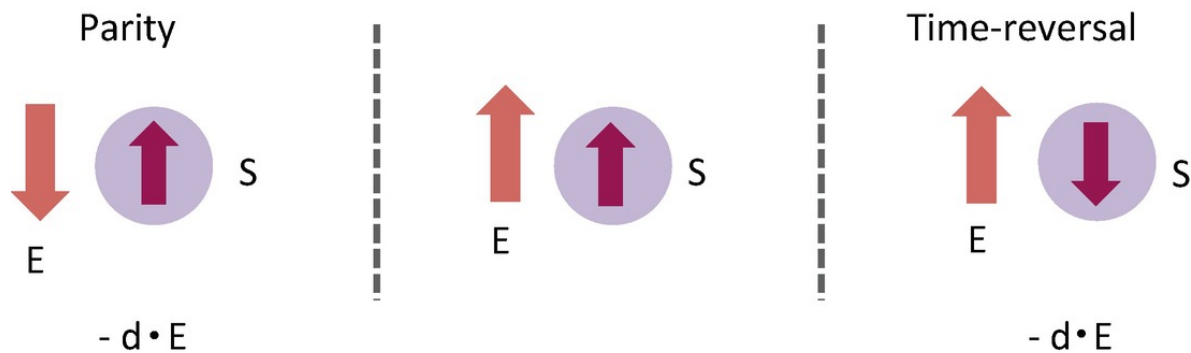
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QCD Lagrangian contains a term which violates P and T:

[Belavin et al. '75; 't Hooft 76; Callan et al. '76; Jackiw, Rebbi '76]

$$\mathcal{L}_{\text{QCD}} \supset -\theta \frac{\alpha_s}{8\pi} G_{\mu\nu}^b \tilde{G}^{b,\mu\nu} \equiv -\theta \frac{\alpha_s}{2\pi} \mathbf{E}^b \cdot \mathbf{B}^b$$



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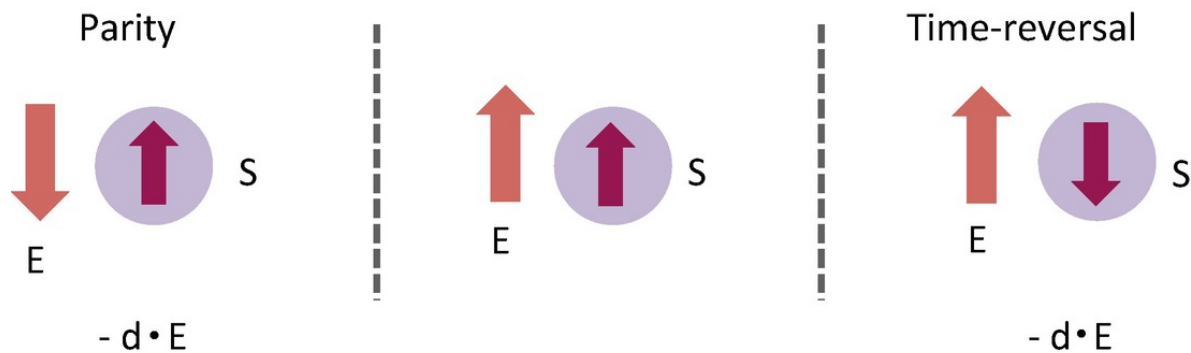
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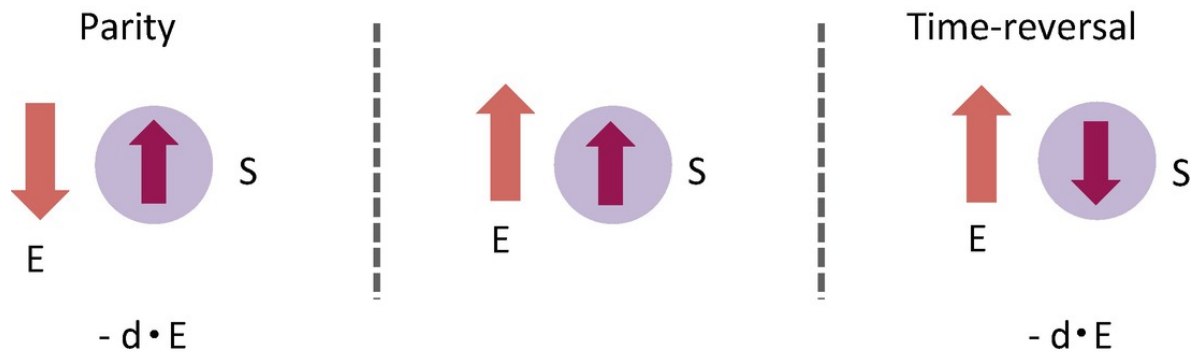
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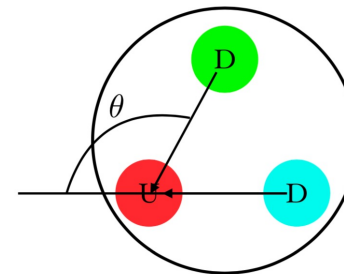
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- predicts neutron electric dipole moment (nEDM):  
[Crewther, Di Vecchia, Veneziano, Witten 79; ...; Pospelov, Ritz 00]

$$d_n \sim \bar{\theta} e \frac{m_u}{m_n^2} \sim 10^{-16} \bar{\theta} e \text{ cm}$$



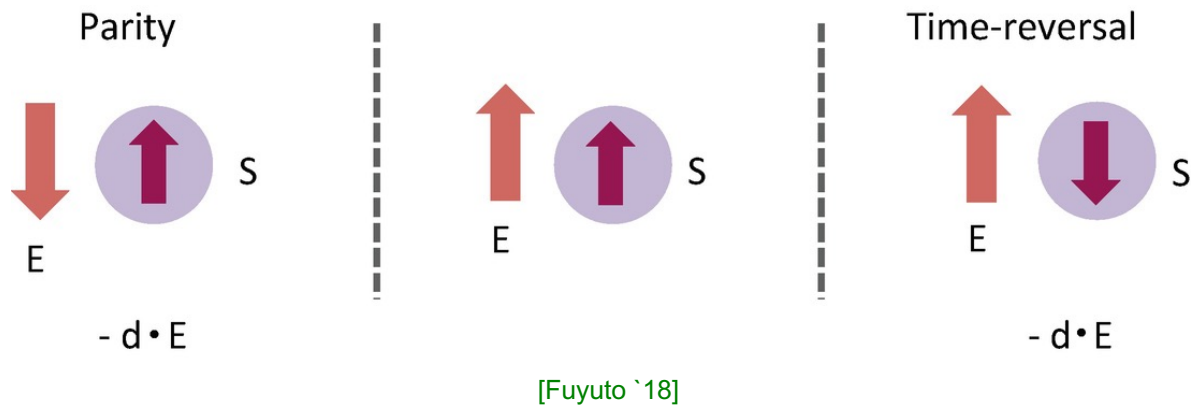
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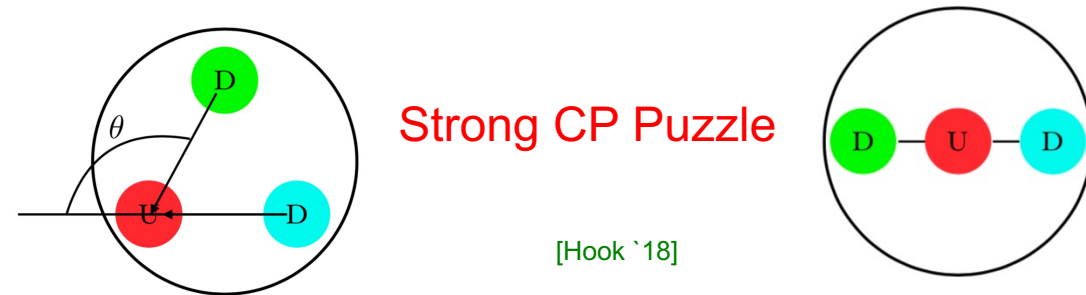
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- Comparison with experiment:  $\Rightarrow |\bar{\theta}| \lesssim 10^{-10}$



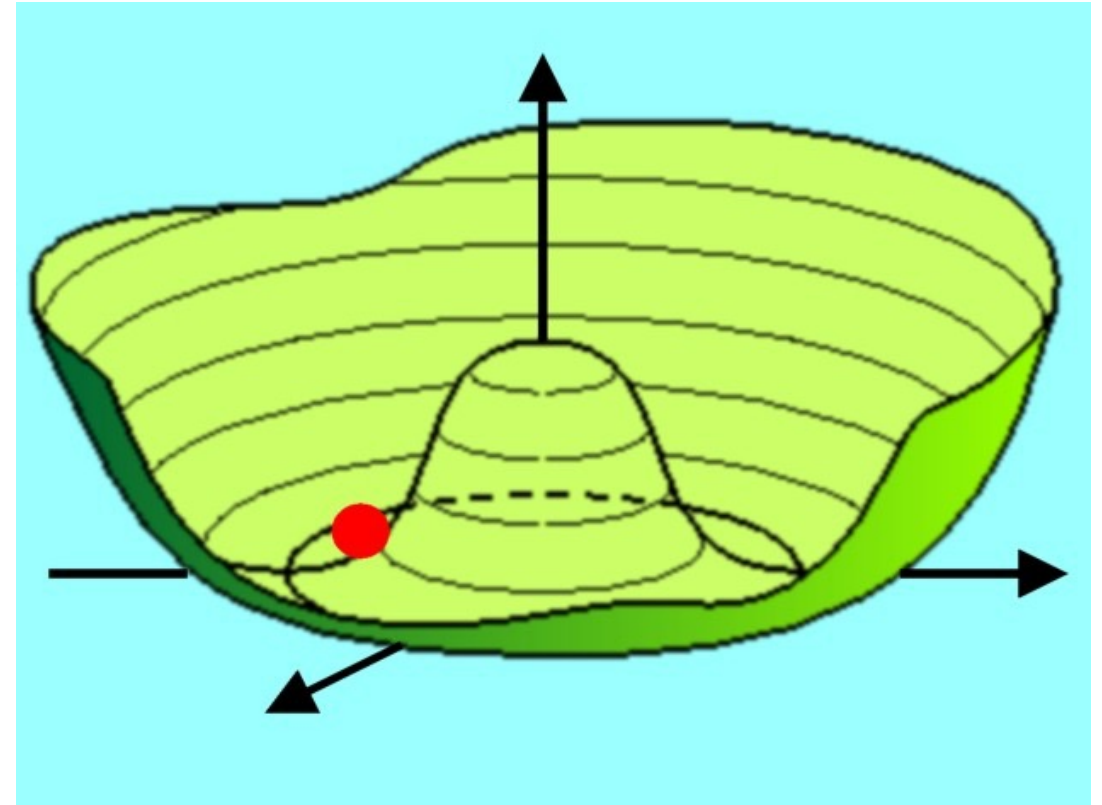
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## Minimal SM extension solving strong CP puzzle

[Kim 79; Shifman, Vainshtein, Zakharov 80]

- Add to SM a singlet complex scalar field  $\sigma$ , featuring a spontaneously broken global axial  $U(1)_{PQ}$  symmetry, and an exotic quark  $Q$  charged under it:

$$\mathcal{L} \supset -\lambda_\sigma \left( |\sigma|^2 - \frac{v_{PQ}^2}{2} \right)^2 + y \sigma \bar{Q}_L Q_R$$



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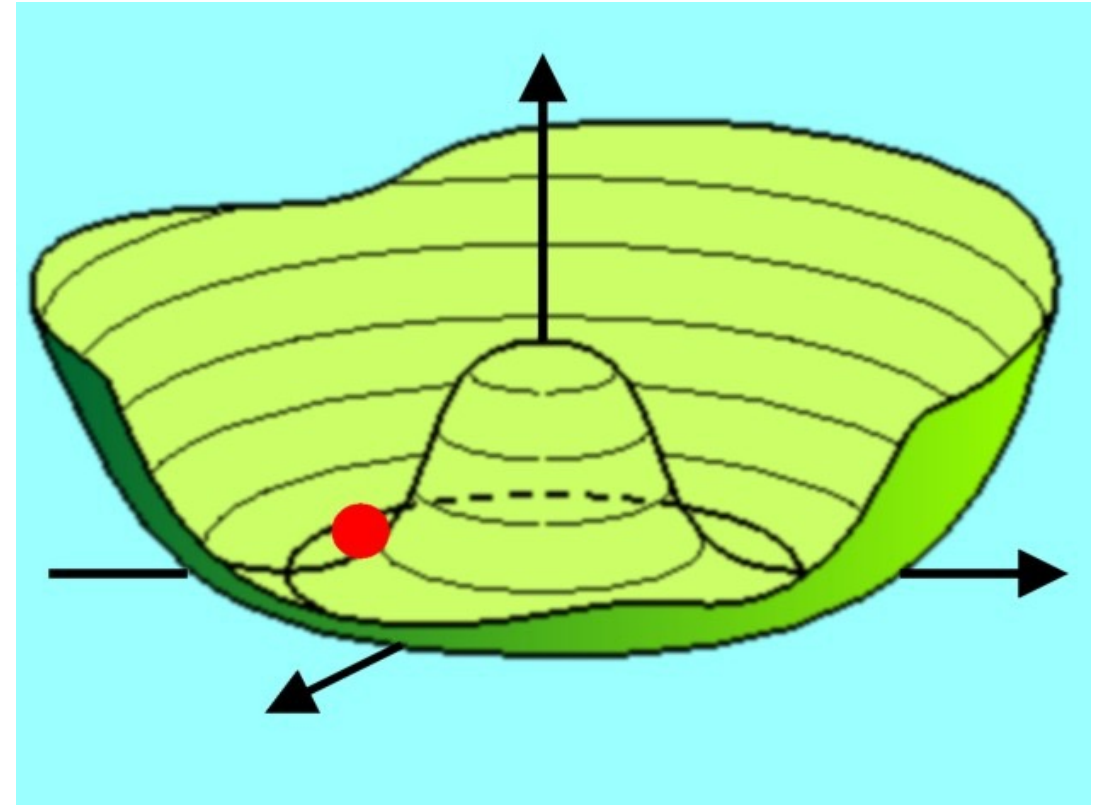
- Features three particles beyond SM:

1. Excitation of angular field  $a$ : axion (NG-boson)

2. Excitation of modulus field  $\rho$ :  $m_\rho = \sqrt{2\lambda_\sigma} v_{PQ}$

3. Exotic quark:  $m_Q = \frac{y}{\sqrt{2}} v_{PQ}$

$$\left[ \sigma(x) = \frac{1}{\sqrt{2}} (v_{PQ} + \rho(x)) e^{ia(x)/v_{PQ}} \right]$$



[Raffelt]

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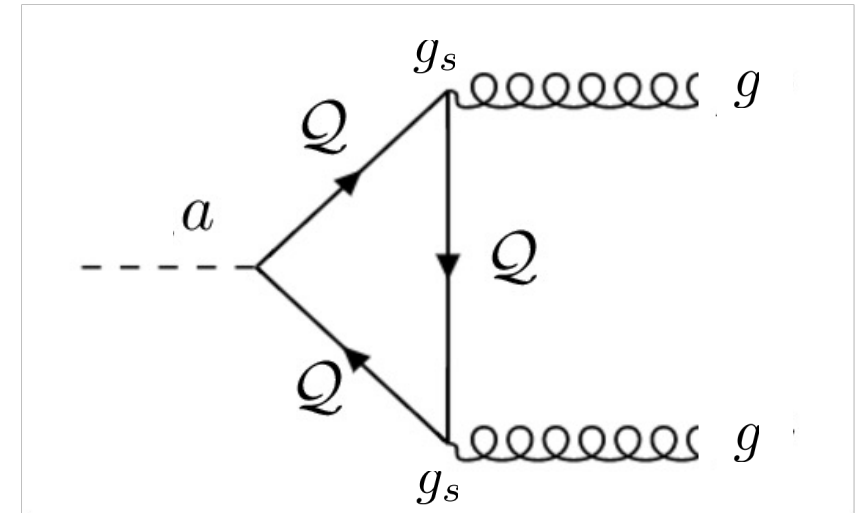
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- Integrating out the modulus and the exotic quark



[Chadha-Day, Ellis, Marsh, Science Advances 22]

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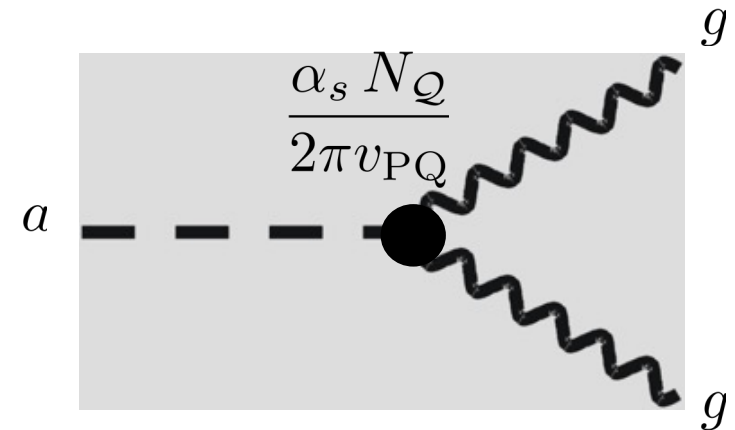
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$$\mathcal{L}_{\text{eff}} \supset \frac{1}{2} \partial^\mu a \partial_\mu a + \frac{\alpha_s}{8\pi} \frac{a}{f_a} G_{\mu\nu}^a \tilde{G}^{a\mu\nu} \quad \text{with} \quad f_a = v_{PQ}/N_Q$$



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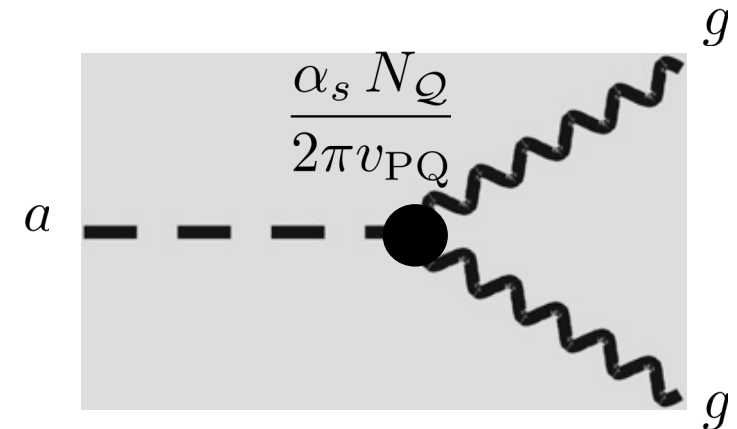
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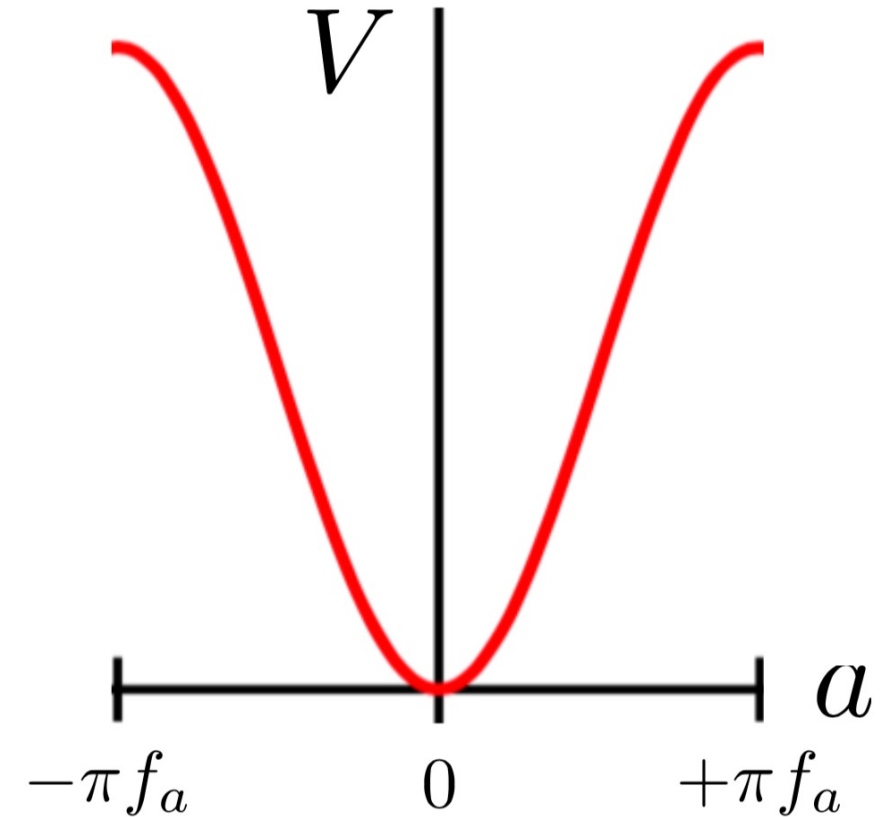
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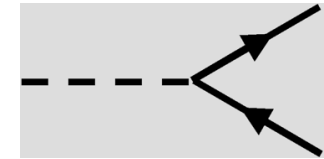
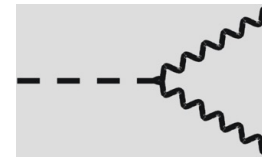
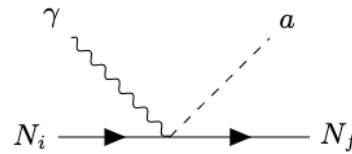
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 with  $f_a = v_{PQ}/N_Q$
- $\theta$ -parameter can be eliminated:  $a(x) \rightarrow a(x) + \bar{\theta} f_a$
- Shifted field (the axion) acquires mass from mixing with neutral pion,  $m_a \sim m_\pi f_\pi / f_a$ :  $\langle a \rangle = 0$ , CP conserved



# The Axion

## Effective field theory below QCD scale

$$\mathcal{L} \supset -\frac{1}{2}m_a^2 a^2 - \frac{i}{2} \frac{e C_{\text{NEDM}}}{f_a} a \bar{\psi}_N \sigma_{\mu\nu} \gamma_5 \psi_N F^{\mu\nu} + C_{a\gamma} \frac{\alpha}{8\pi} \frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu} + \frac{1}{2} C_{af} \frac{\partial_\mu a}{f_a} \bar{\psi}_f \gamma^\mu \gamma_5 \psi_f$$



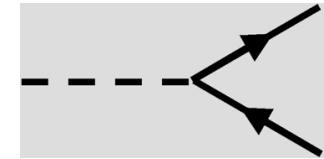
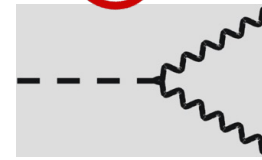
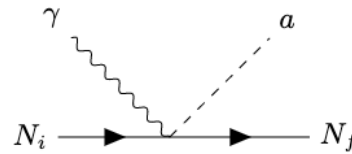
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$$m_a \simeq \frac{\sqrt{z}}{1+z} \frac{m_\pi f_\pi}{f_a} \approx 6 \text{ meV} \left( \frac{10^9 \text{ GeV}}{f_a} \right)$$

$$z = m_u/m_d \simeq 1/2$$



- Suppression by inverse power of symmetry breaking scale  $f_a$  leads to tiny mass and feeble couplings to SM

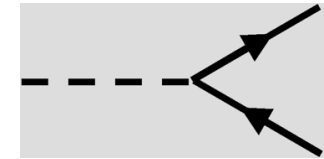
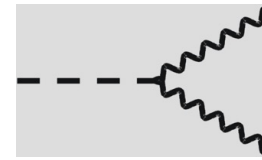
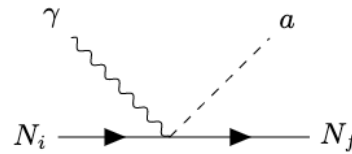
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- Suppression by inverse power of symmetry breaking scale  $f_a$  leads to tiny mass and feeble couplings to SM
- Size of coefficients in vanilla axion models:

- EDM coupling:

$$C_{\text{NEDM}} = 2.4(1.0) \times 10^{-16} \text{ cm}$$

- Photon coupling:

$$C_{a\gamma} = \frac{E}{N} - \frac{2}{3} \frac{4+z}{1+z}$$

- Nucleon couplings:

$$C_{aN} = \mathcal{O}(1)$$

[Pospelov, Ritz '00]

[Kaplan 85; Srednicki '85; Grilli di Cortona et al. '16]

[Grilli di Cortona et al. '16]

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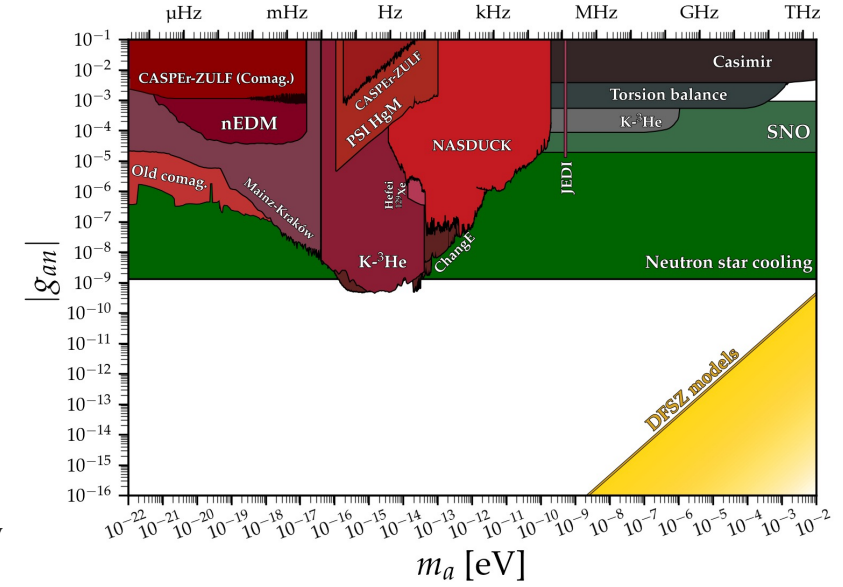
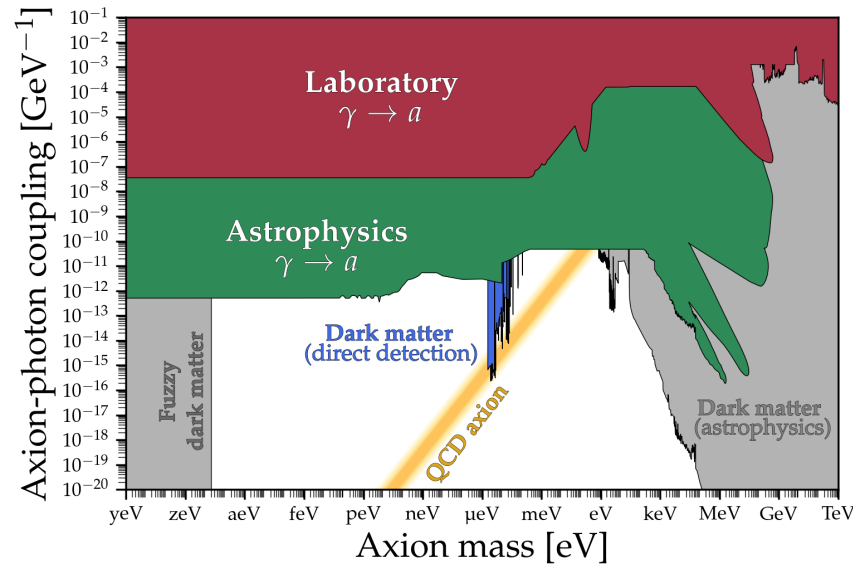
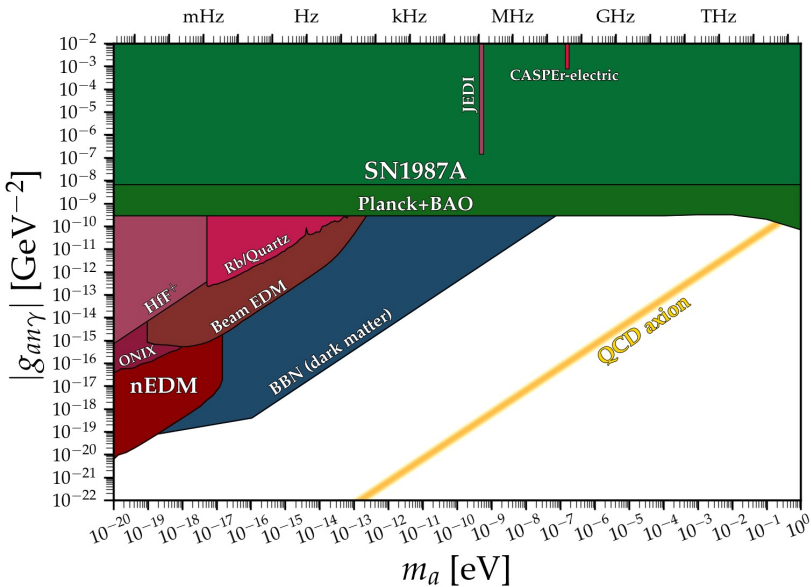
## Expectations for axion couplings to SM particles

“Yellow bands” of vanilla axion couplings:

$$g_{an\gamma} = e \frac{C_{\text{NEDM}}}{f_a}$$

$$g_{a\gamma} = \frac{\alpha}{2\pi f_a} C_{a\gamma}$$

$$g_{an} = \frac{C_{an} m_n}{f_a}$$



[<https://github.com/cajohare/AxionLimits>]

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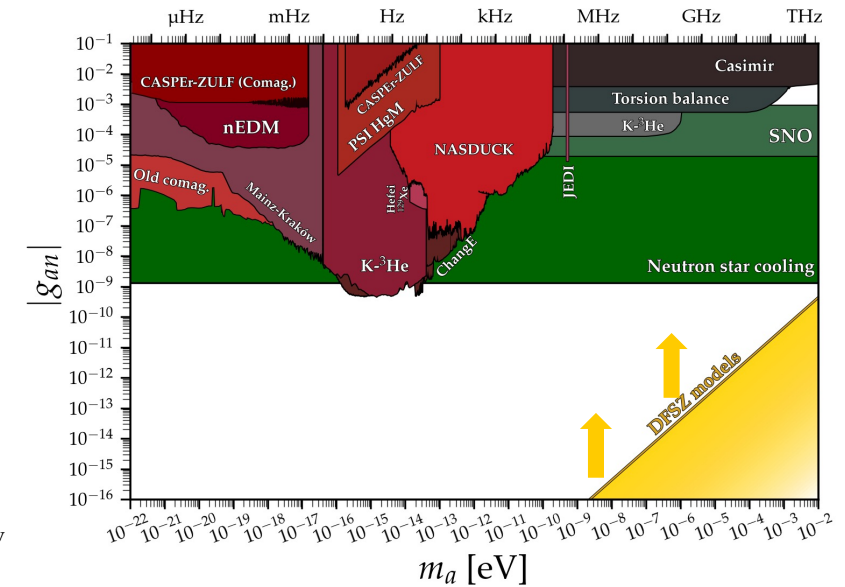
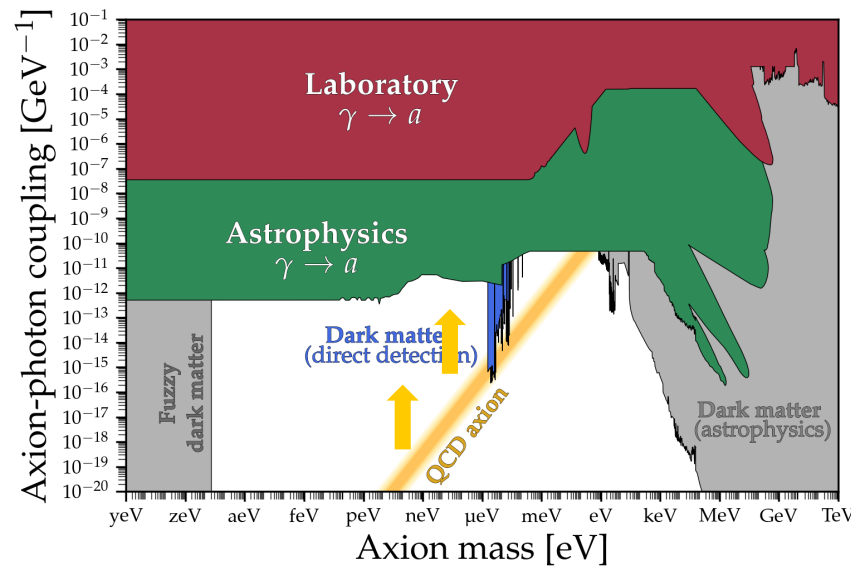
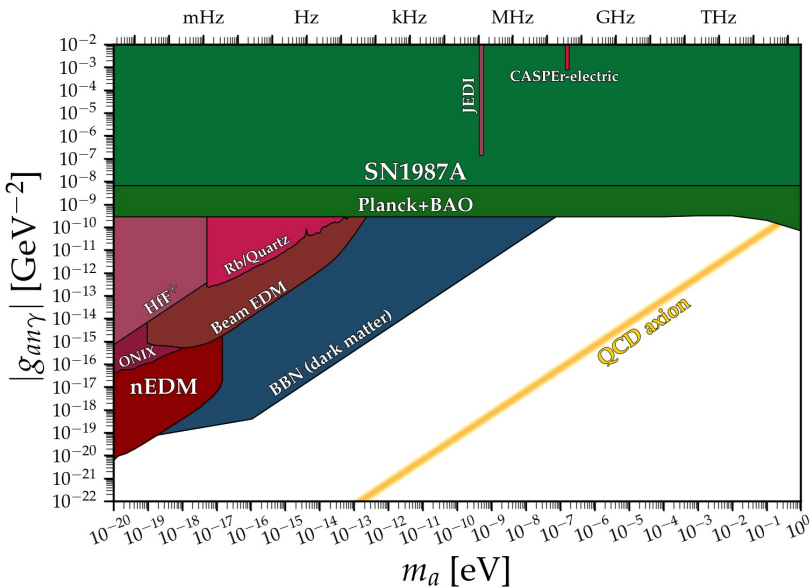
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There are axion models with enhanced  $C_{a\gamma}$  and/or  $C_{aN}$ :

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## Expectations for axion couplings to SM particles

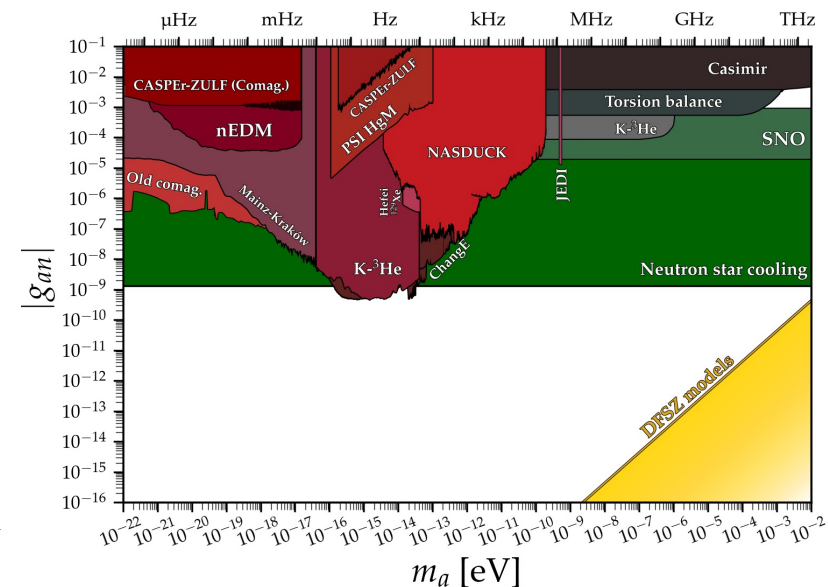
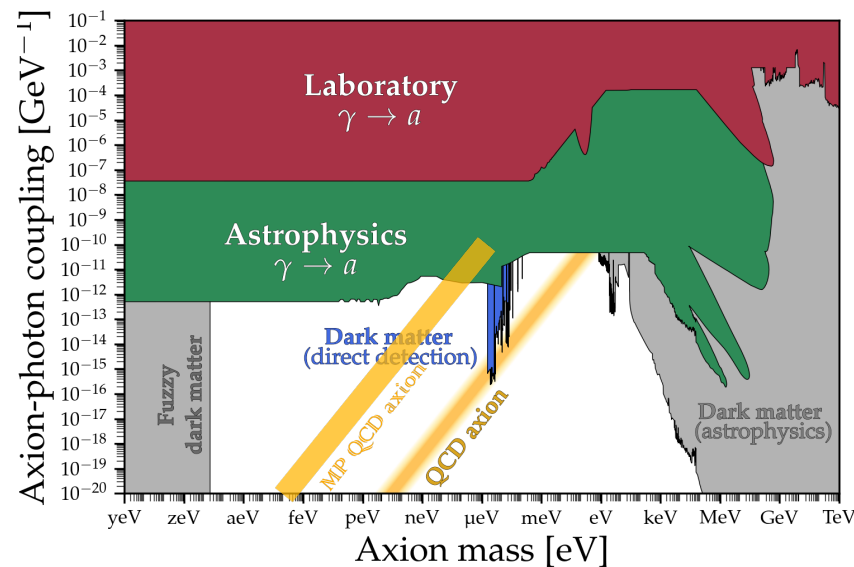
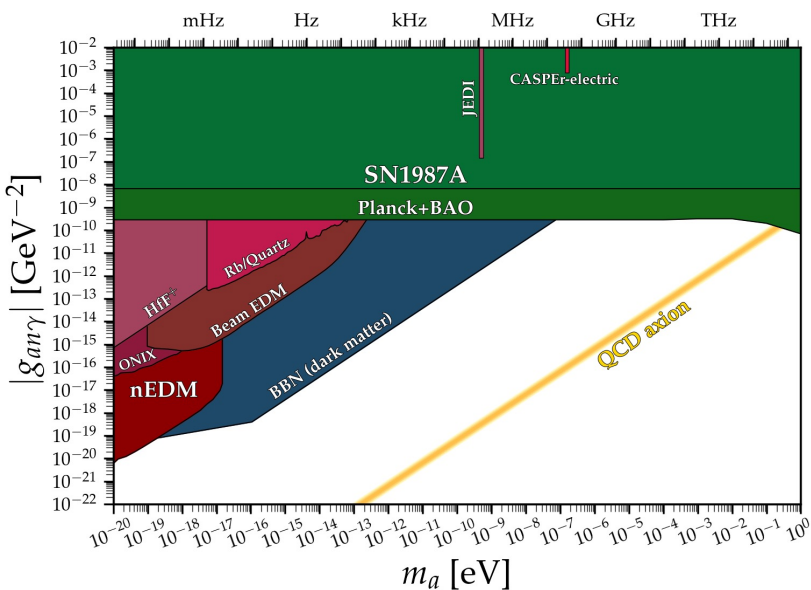
There are axion models with enhanced  $C_{a\gamma}$  and/or  $C_{aN}$ , e.g. Monopole Philic (MP) axion

[Sokolov,AR, 2104.02574; 2109.08503; 2205.02605; 2303.10170]

$$g_{an\gamma} = e \frac{C_{\text{NEDM}}}{f_a}$$

$$g_{a\gamma} = \frac{\alpha}{2\pi f_a} C_{a\gamma}$$

$$g_{an} = \frac{C_{an} m_n}{f_a}$$



[<https://github.com/cajohare/AxionLimits>]

# The Axion

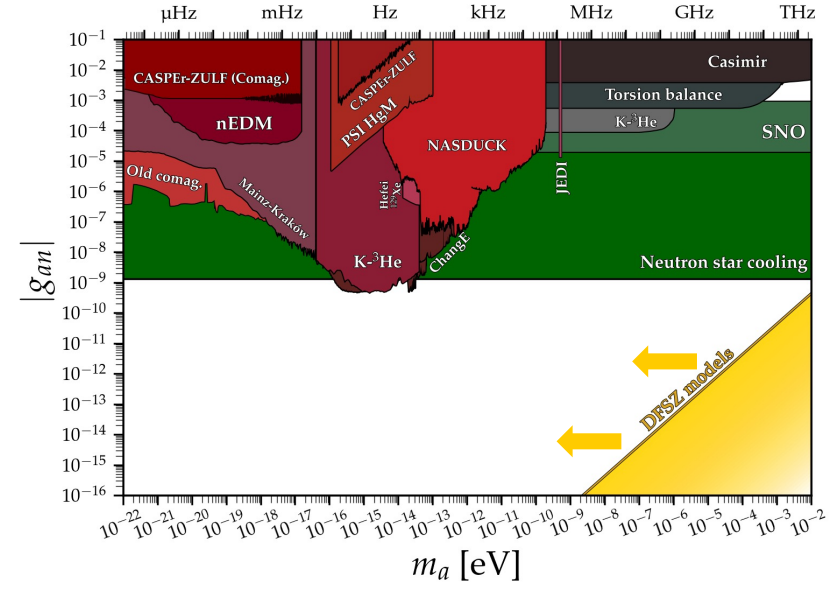
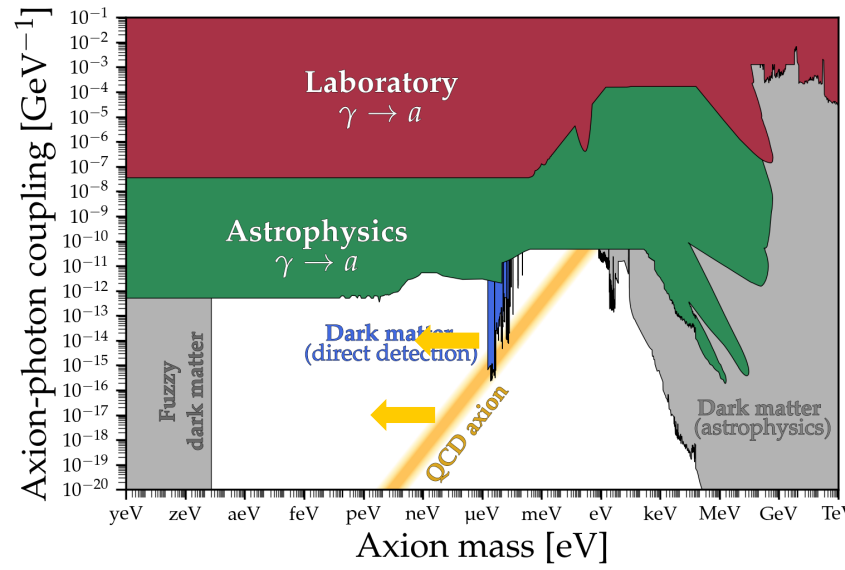
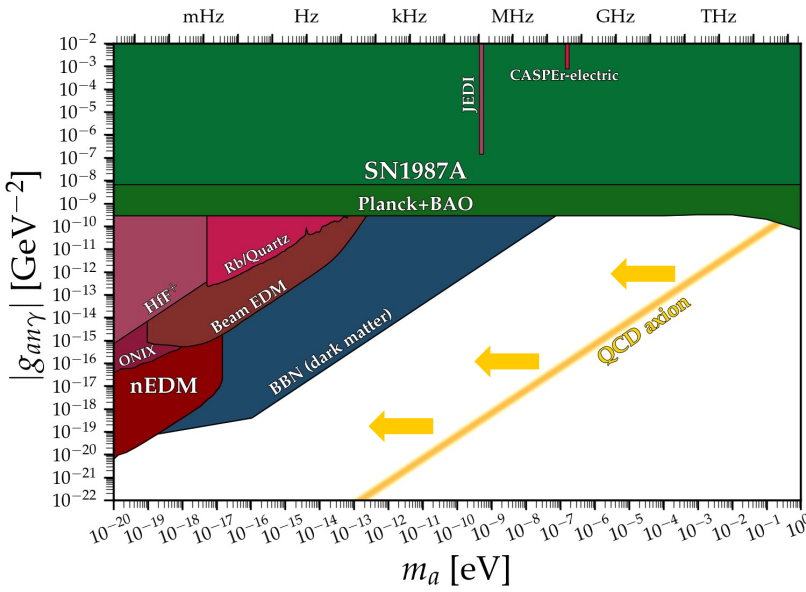
## Expectations for axion couplings to SM particles

There are axion models with reduced  $(m_a f_a)/(m_\pi f_\pi)$

$$g_{an\gamma} = e \frac{C_{\text{NEDM}}}{f_a}$$

$$g_{a\gamma} = \frac{\alpha}{2\pi f_a} C_{a\gamma}$$

$$g_{an} = \frac{C_{an} m_n}{f_a}$$



[<https://github.com/cajohare/AxionLimits>]

# The Axion

## Expectations for axion couplings to SM particles

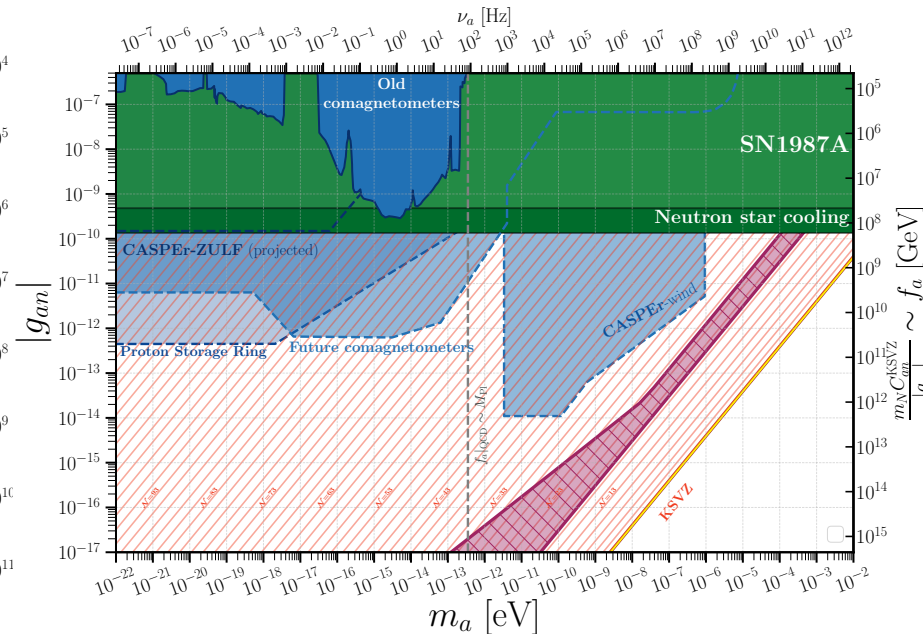
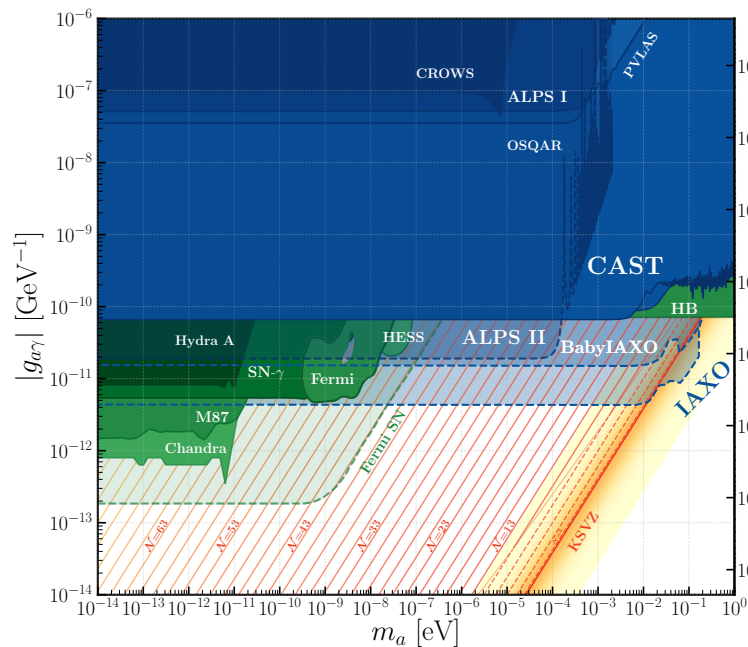
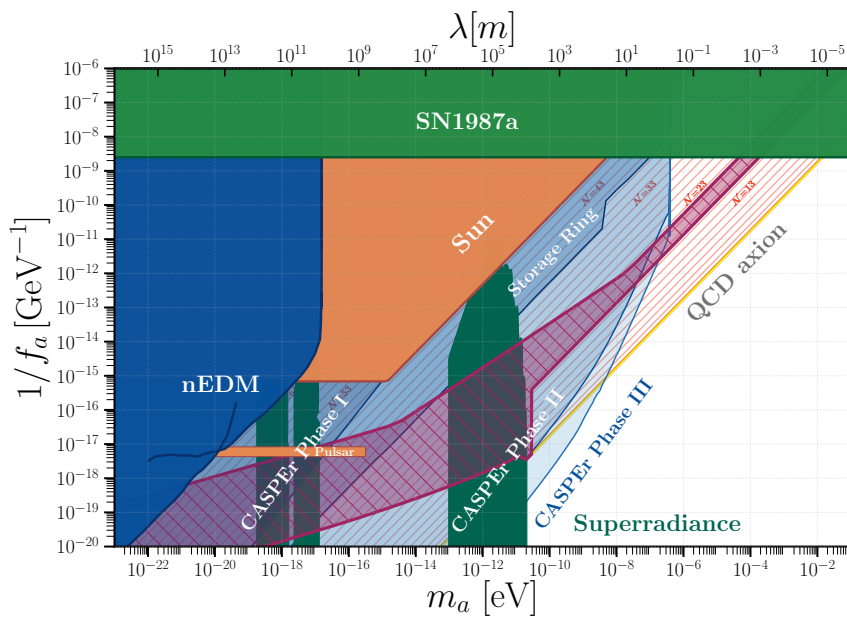
There are axion models with reduced  $(m_a f_a)/(m_\pi f_\pi)$ , e.g. the  $Z_N$  model, where  $(m_a f_a)/(m_\pi f_\pi) \sim (m_u/m_d)^N$

[Hook, 1802.10093; Di Luzio, Gavela, Quilez, AR, 2102.00012; 2102.01082]

$$g_{an\gamma} = e \frac{C_{\text{NEDM}}}{f_a}$$

$$g_{a\gamma} = \frac{\alpha}{2\pi f_a} C_{a\gamma}$$

$$g_{an} = \frac{C_{an} m_n}{f_a}$$



# Axion Experiments Not Relying on Axion Dark Matter

# Axion Experiments Not Relying on Axion Dark Matter

Searching for solar or home-made axions

# Axion Experiments Not Relying on Axion Dark Matter

Searching for solar or home-made axions

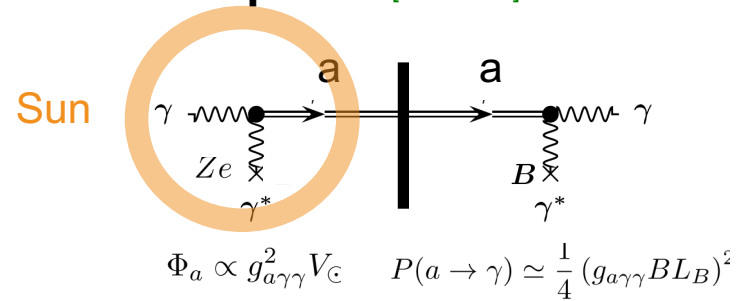
$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

# Axion Experiments Not Relying on Axion Dark Matter

## Searching for solar or home-made axions

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- **Helioscopes:** [Sikivie 83]

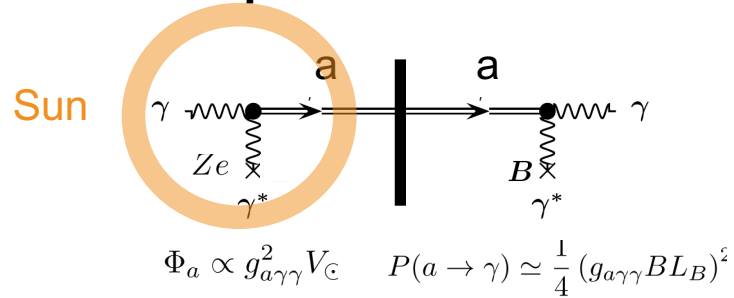


# Axion Experiments Not Relying on Axion Dark Matter

## Searching for solar or home-made axions

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- **Helioscopes:** [Sikivie 83]



- **CAST @ CERN (2003 - 2015)**



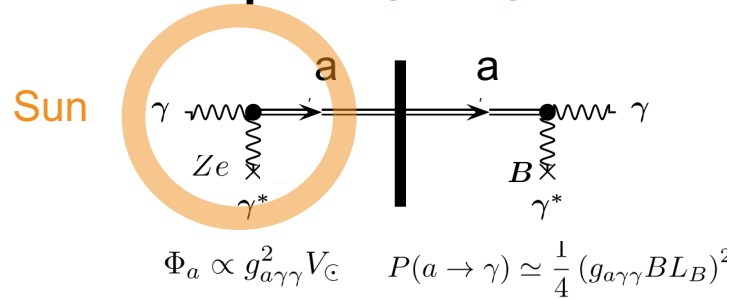
- LHC test dipole magnet (B ~ 9 T, L ~ 10 m)

# Axion Experiments Not Relying on Axion Dark Matter

## Searching for solar or home-made axions

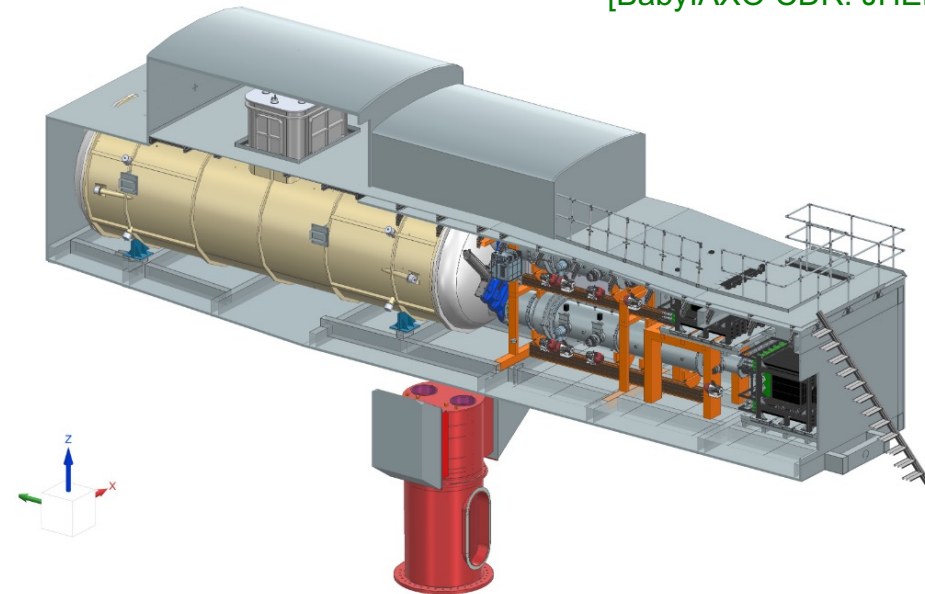
$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- **Helioscopes:** [Sikivie 83]



- **BabyIAXO** in construction @ DESY: (1st data 2031?)

[BabyIAXO CDR: JHEP 05 (2021) 137]



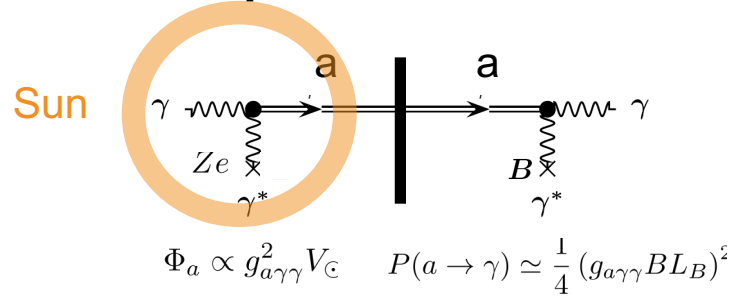
- Superconducting magnet featuring two flat racetrack coils made from Niobium-Titanium wire  
 $B \sim 2 \text{ T}$ ,  $L \sim 10 \text{ m}$ , 2 bores with 70 cm diameter
- 2 X-ray telescopes + 2 detection systems
- Movable platform with service

# Axion Experiments Not Relying on Axion Dark Matter

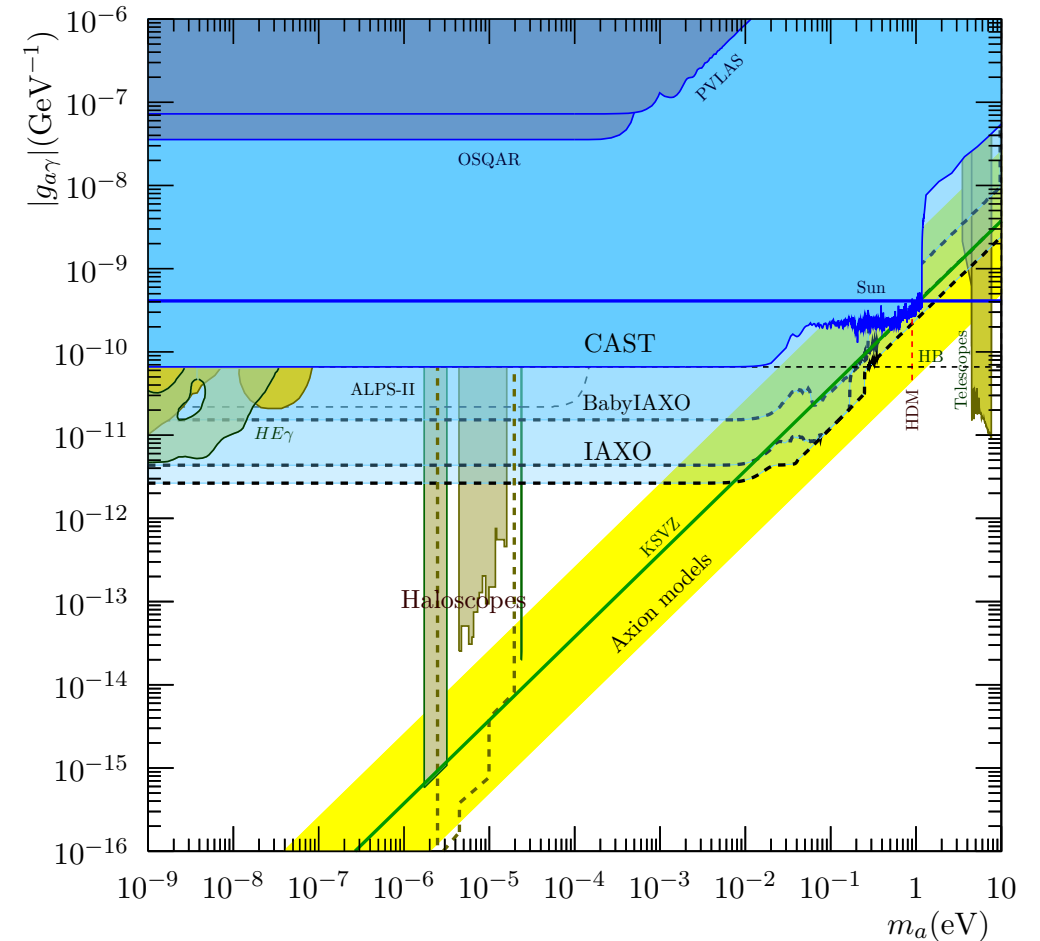
Searching for solar or home-made axions

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- Helioscopes: [Sikivie 83]



- BabyIAXO sensitivity forecast:



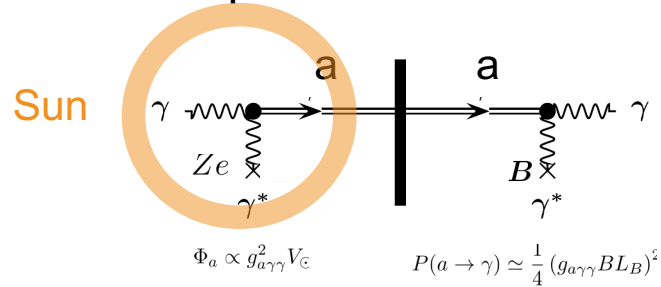
[IAXO Collaboration, JCAP 06 (2019) 047]

# Axion Experiments Not Relying on Axion Dark Matter

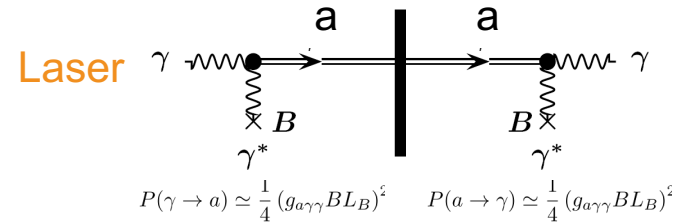
## Searching for solar or home-made axions

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- Helioscopes: [Sikivie 83]



- Light-shining through walls: [Anselm 85; van Bibber 87]

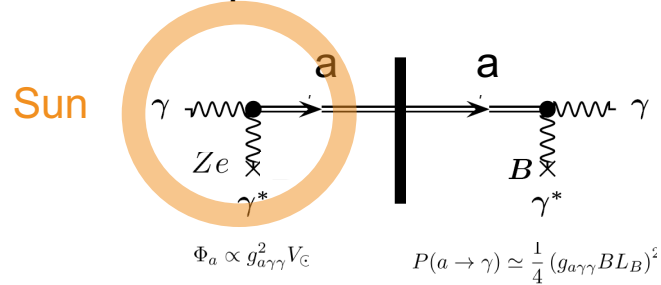


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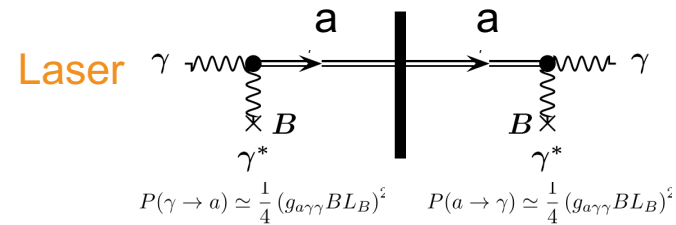
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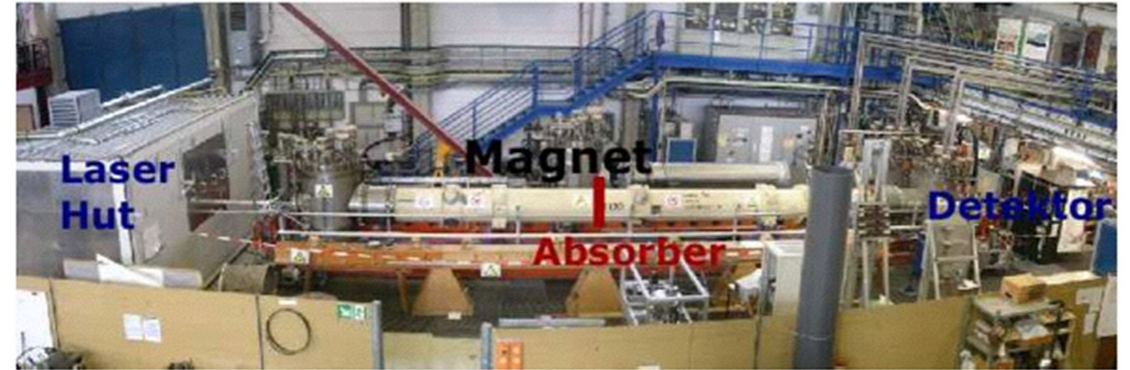
- Helioscopes: [Sikivie 83]



- Light-shining through walls: [Anselm 85; van Bibber 87]



- ALPS @ DESY (2007 - 2009)



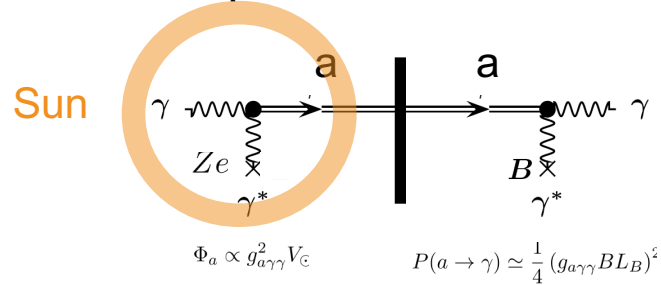
- HERA dipole magnet (B ~ 5 T, L ~ 9 m)

# Axion Experiments Not Relying on Axion Dark Matter

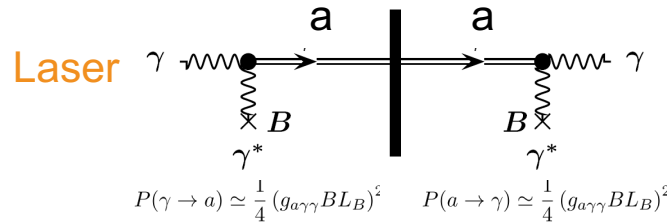
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$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- Helioscopes: [Sikivie 83]



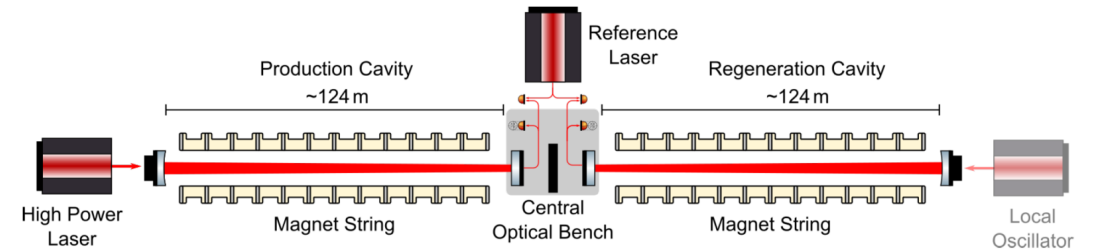
- Light-shining through walls: [Anselm 85; van Bibber 87]



- ALPS II @ DESY [Bähre et al (ALPS II TDR) 13]

12 + 12 straightened HERA magnets

[AR, Phys. Lett. B 569 (2003) 51]



Optical cavities both at production and regeneration sites

[Hogeveen, Ziegenhagen 91; Sikivie, Tanner, van Bibber 07]

Two detection techniques:

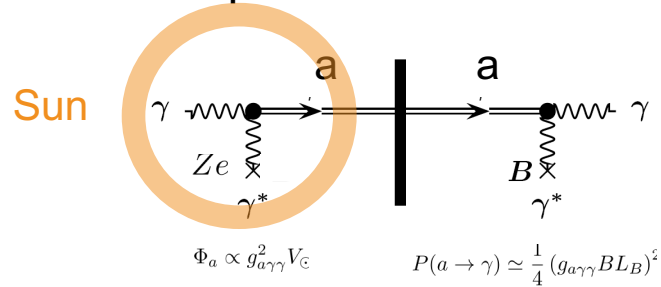
- Heterodyne
- Transition Edge Sensor (TES)

# Axion Experiments Not Relying on Axion Dark Matter

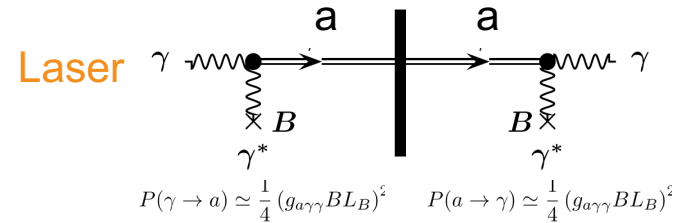
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- ALPS II @ DESY [Bähre et al (ALPS II TDR) 13]

Installation in HERA tunnel 2019-2022



First runs in 2023-2024 with cavity only at regeneration side

First science paper 2025

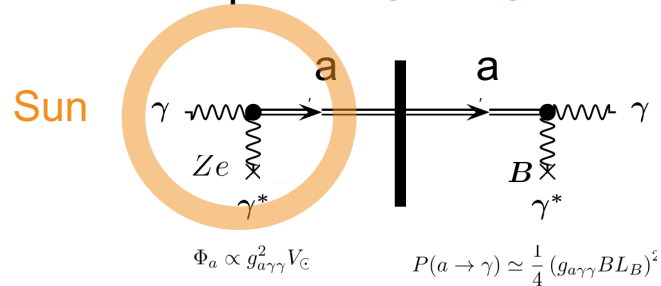
[ALPS II Collaboration, 2512.14110 [hep-ex]]

# Axion Experiments Not Relying on Axion Dark Matter

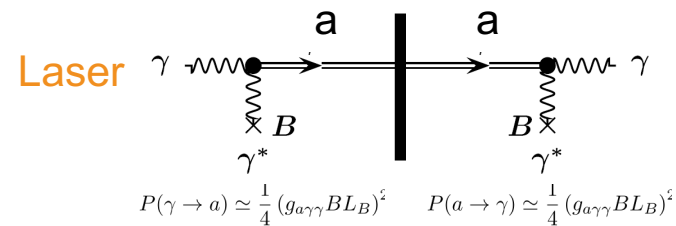
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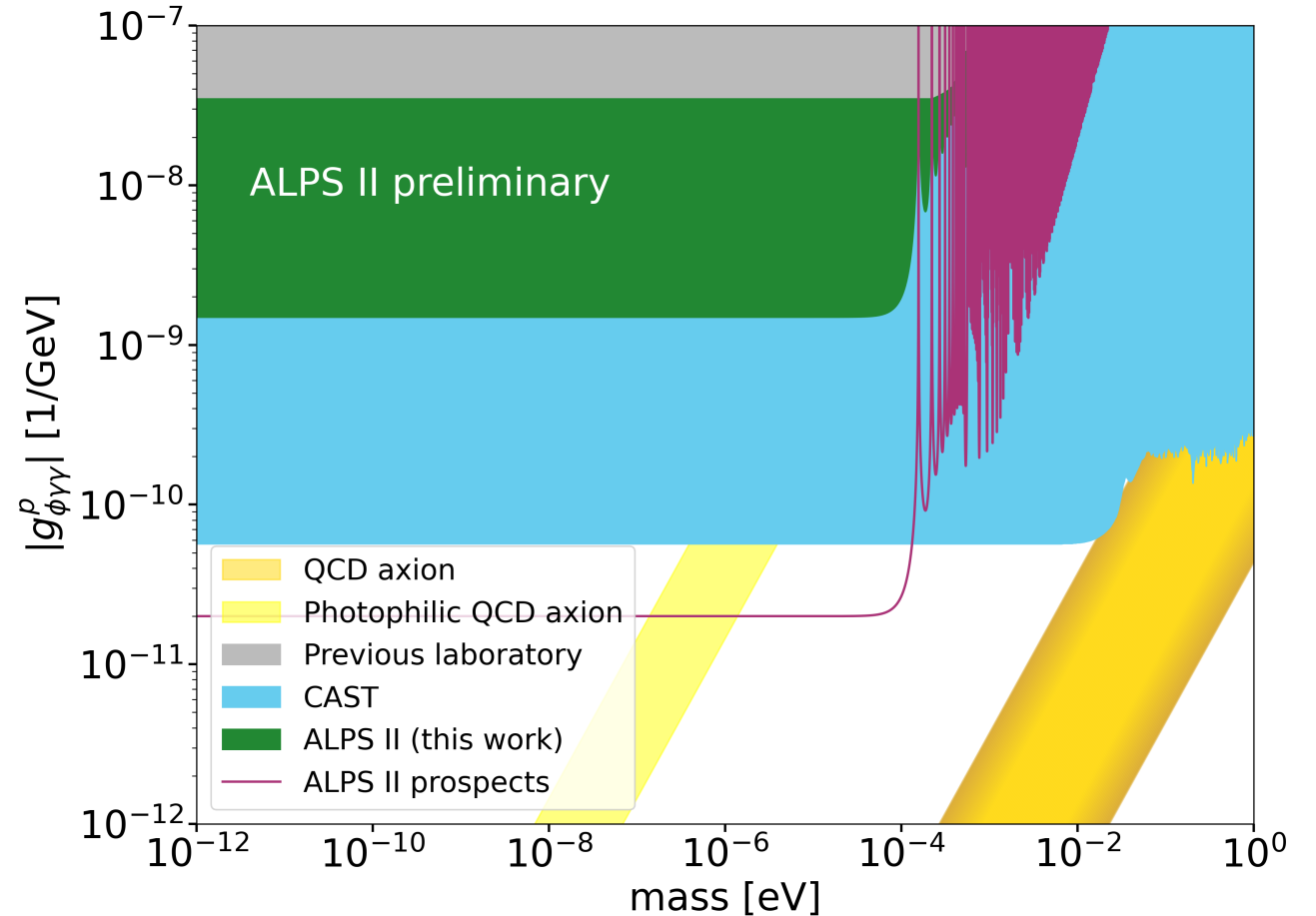
- Helioscopes: [Sikivie 83]



- Light-shining through walls: [Anselm 85; van Bibber 87]



- ALPS II current bound:



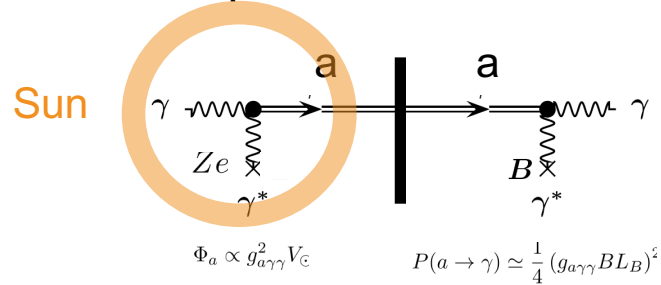
[ALPS II Collaboration, 2512.14110 [hep-ex]]

# Axion Experiments Not Relying on Axion Dark Matter

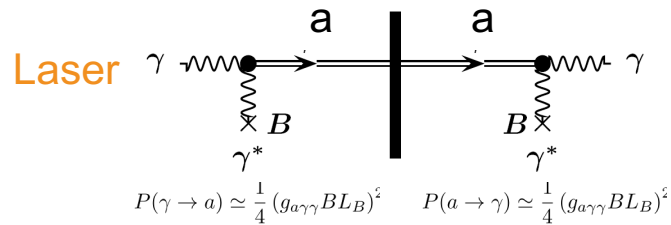
## Searching for solar or home-made axions

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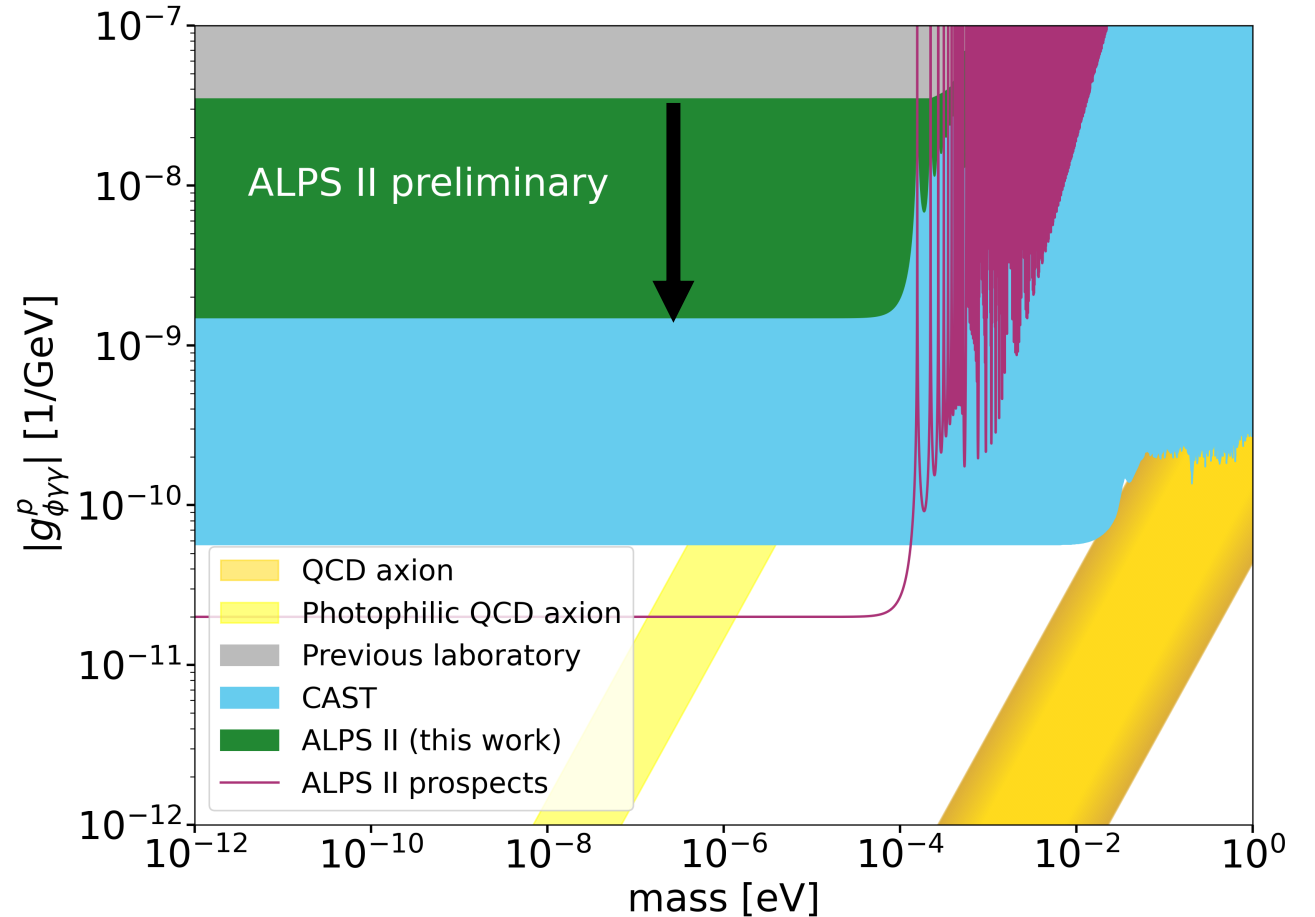
- Helioscopes: [Sikivie 83]



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- ALPS II current bound:



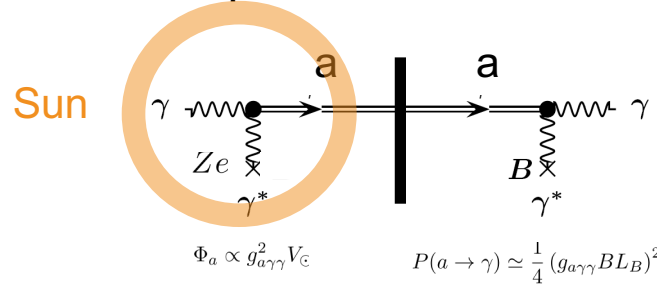
[ALPS II Collaboration, 2512.14110 [hep-ex]]

# Axion Experiments Not Relying on Axion Dark Matter

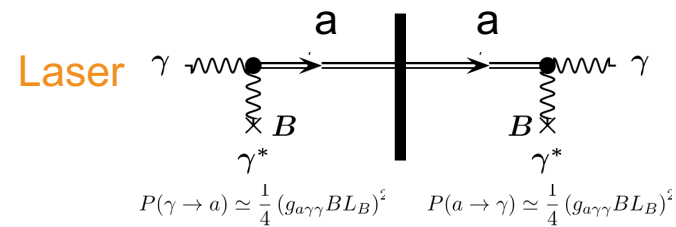
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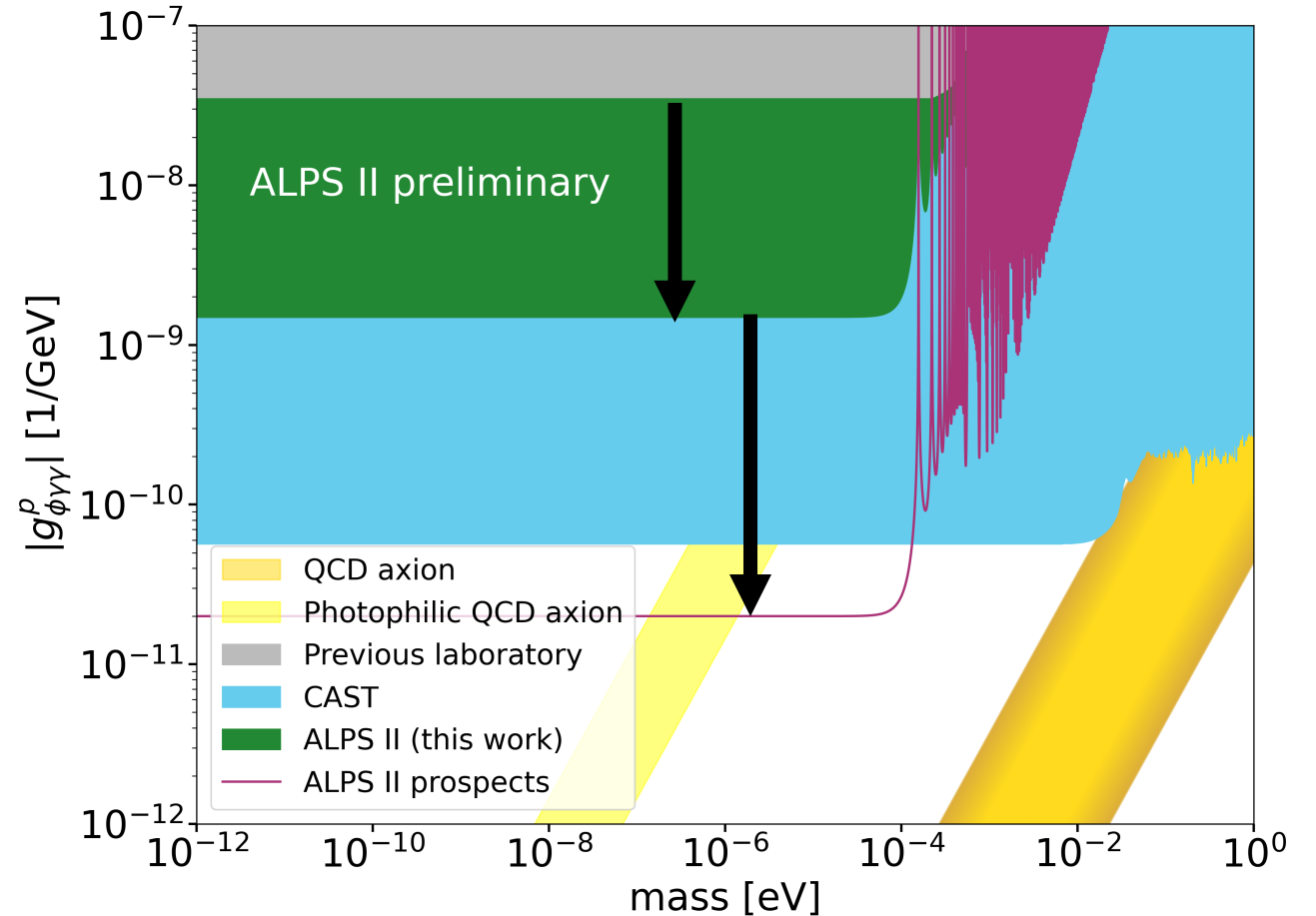
- Helioscopes: [Sikivie 83]



- Light-shining through walls: [Anselm 85; van Bibber 87]



- ALPS II current bound and 2028(?) sensitivity:



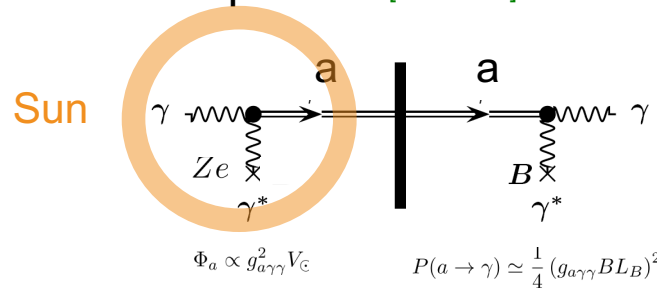
[ALPS II Collaboration, 2512.14110 [hep-ex]]

# Axion Experiments Not Relying on Axion Dark Matter

## Searching for solar or home-made axions

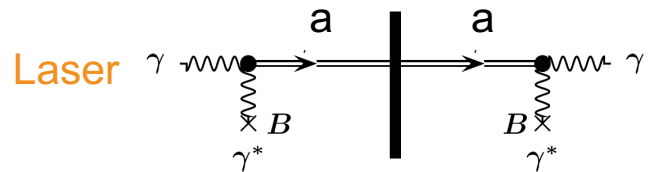
$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- Helioscopes: [Sikivie 83]



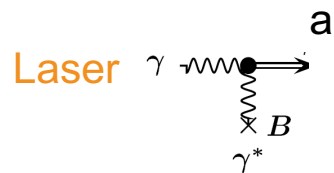
$$P(a \rightarrow \gamma) \simeq \frac{1}{4} (g_{a\gamma\gamma} B L_B)^2$$

- Light-shining through walls: [Anselm 85; van Bibber 87]



$$P(\gamma \rightarrow a) \simeq \frac{1}{4} (g_{a\gamma\gamma} B L_B)^2 \quad P(a \rightarrow \gamma) \simeq \frac{1}{4} (g_{a\gamma\gamma} B L_B)^2$$

- Photon-disappearance:



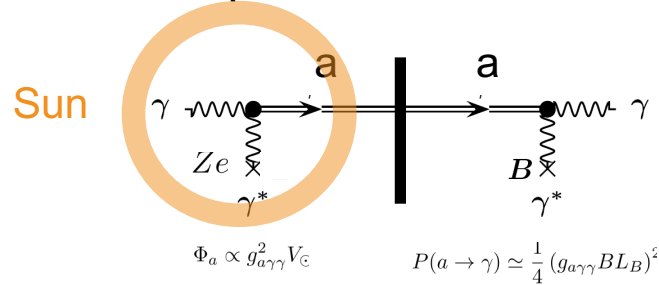
$$P(\gamma \rightarrow a) \simeq \frac{1}{4} (g_{a\gamma\gamma} B L_B)^2$$

# Axion Experiments Not Relying on Axion Dark Matter

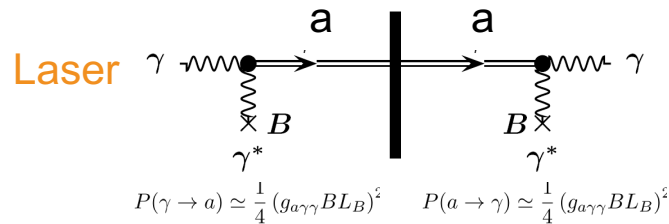
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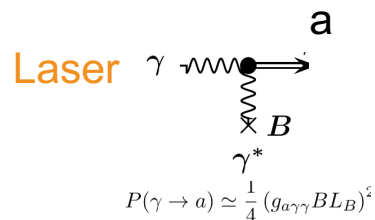
- Helioscopes: [Sikivie 83]



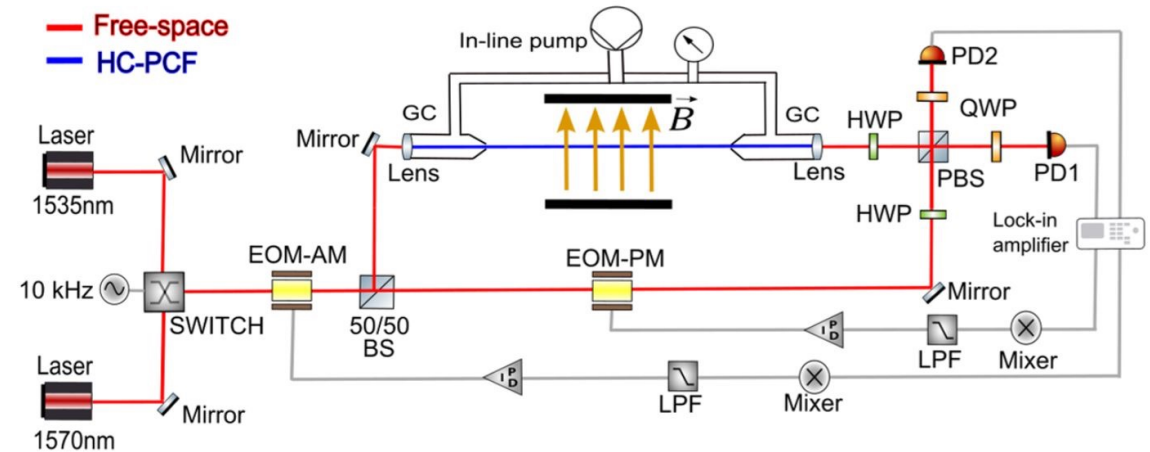
- Light-shining through walls: [Anselm 85; van Bibber 87]



- Photon-disappearance:



- WISPF1 @ University of Hamburg [Batllori et al., 2305.12969]



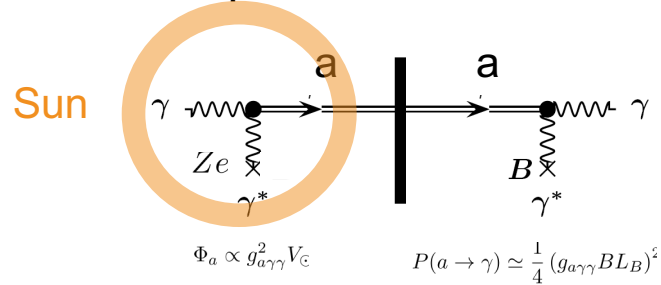
- Mach-Zehnder interferometer with one arm in a strong transverse magnetic field

# Axion Experiments Not Relying on Axion Dark Matter

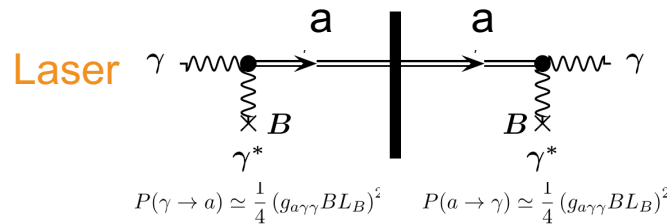
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$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

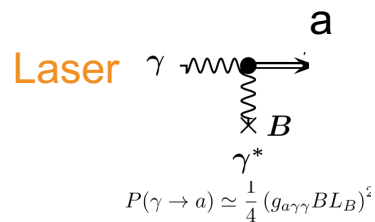
- Helioscopes: [Sikivie 83]



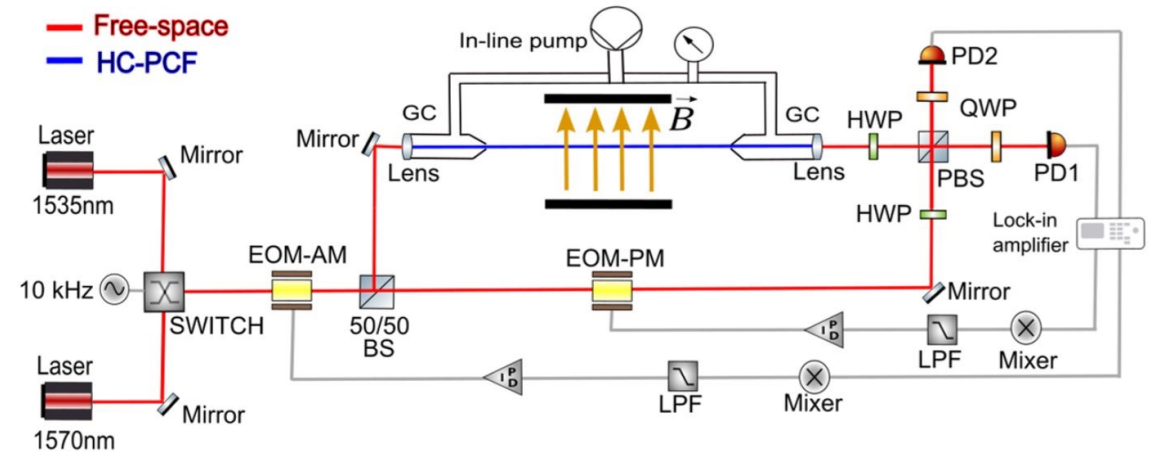
- Light-shining through walls: [Anselm 85; van Bibber 87]



- Photon-disappearance:



- WISPMI @ University of Hamburg [Batllori et al., 2305.12969]



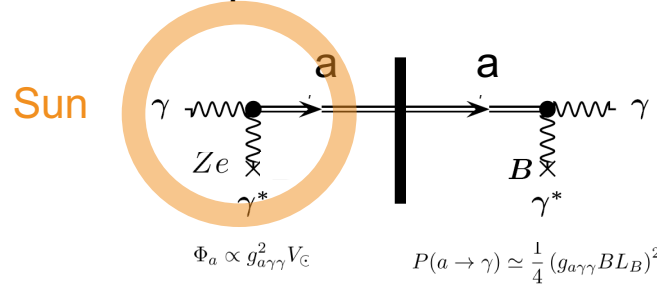
- Mach-Zehnder interferometer with one arm in a strong transverse magnetic field
- Photon-to-axion conversion causes photon disappearance in this arm

# Axion Experiments Not Relying on Axion Dark Matter

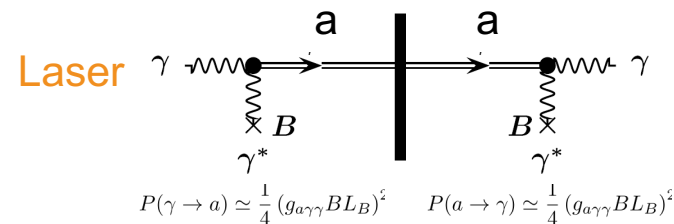
## Searching for solar or home-made axions

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

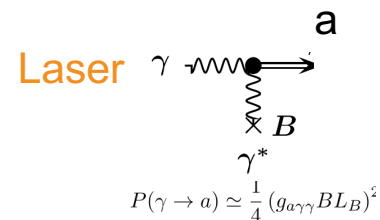
- Helioscopes: [Sikivie 83]



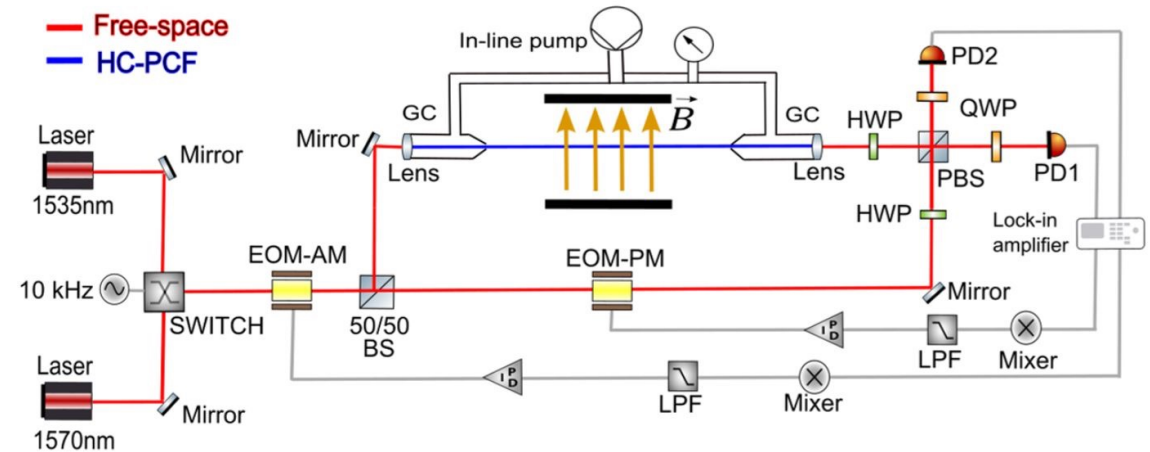
- Light-shining through walls: [Anselm 85; van Bibber 87]



- Photon-disappearance:



- WISPMI @ University of Hamburg [Batllori et al., 2305.12969]



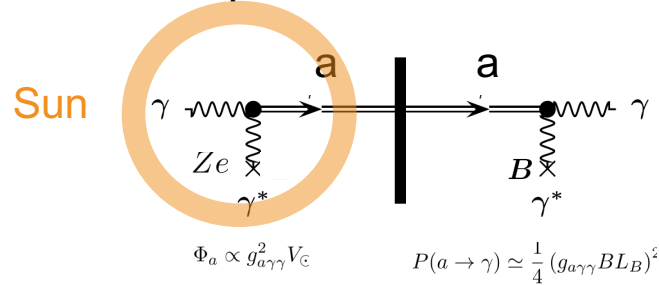
- Mach-Zehnder interferometer with one arm in a strong transverse magnetic field
- Photon-to-axion conversion causes photon disappearance in this arm
- Switching between two laser wavelengths modulates any axion-induced photon disappearance at a known frequency

# Axion Experiments Not Relying on Axion Dark Matter

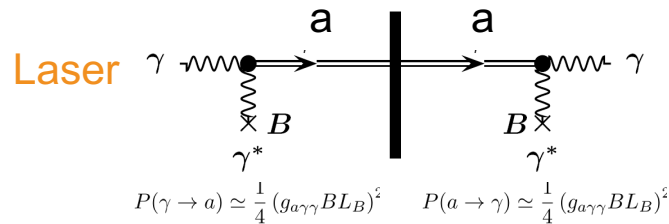
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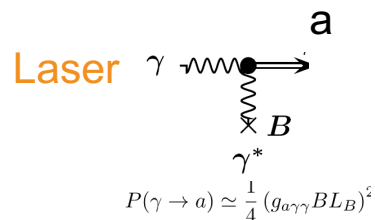
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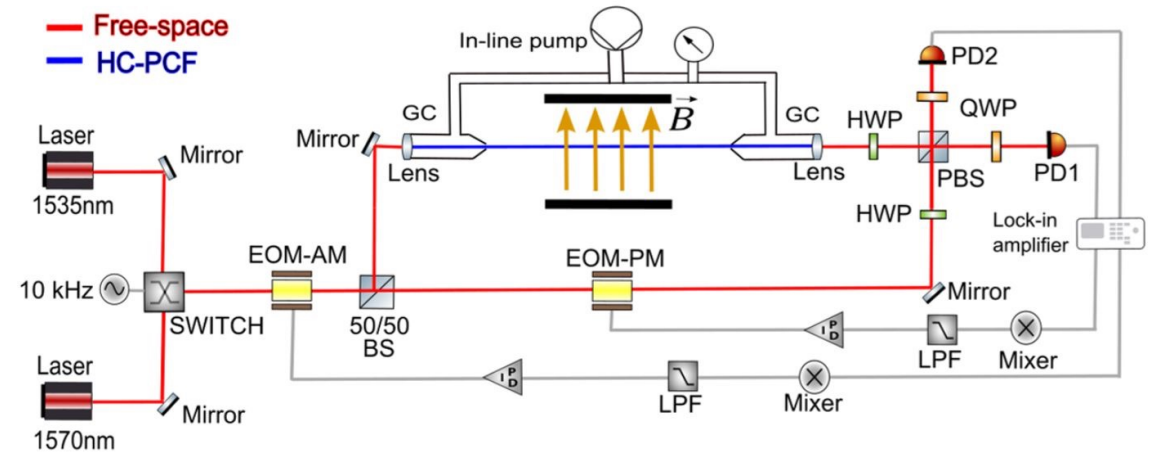
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- Photon-disappearance:



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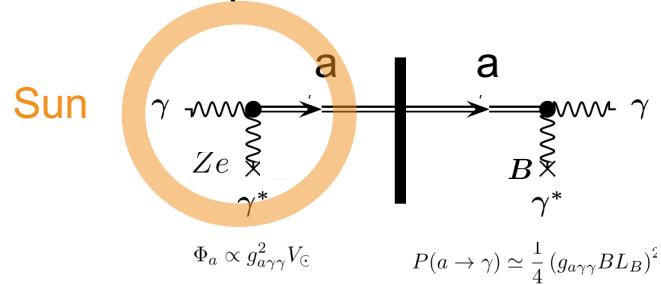
- Mach-Zehnder interferometer with one arm in a strong transverse magnetic field
- Photon-to-axion conversion causes photon disappearance in this arm
- Switching between two laser wavelengths modulates any axion-induced photon disappearance at a known frequency
- Hollow-core fiber (with effective mode index < 1) enables resonant photon-axion conversion

# Axion Experiments Not Relying on Axion Dark Matter

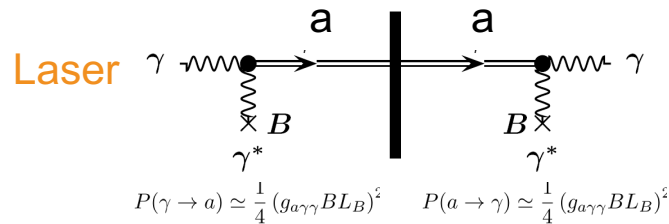
## Searching for solar or home-made axions

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

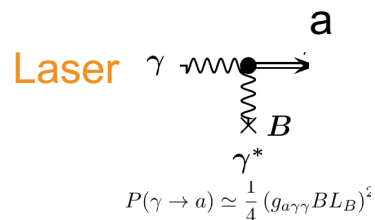
- Helioscopes: [Sikivie 83]



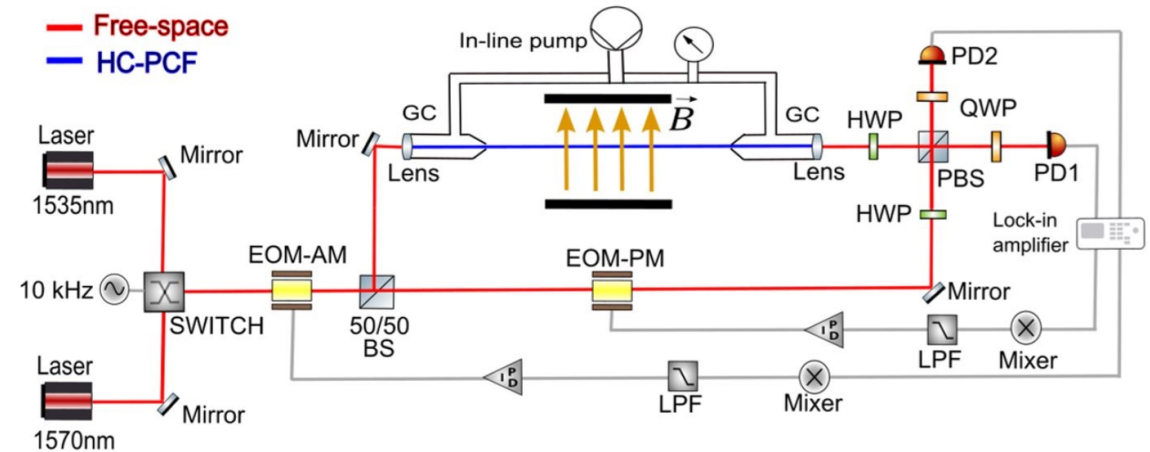
- Light-shining through walls: [Anselm 85; van Bibber 87]



- Photon-disappearance:



- WISPMI @ University of Hamburg [Batllori et al., 2305.12969]



- Mach-Zehnder interferometer with one arm in a strong transverse magnetic field
- Photon-to-axion conversion causes photon disappearance in this arm
- Switching between two laser wavelengths modulates any axion-induced photon disappearance at a known frequency
- Hollow-core fiber (with effective mode index < 1) enables resonant photon-axion conversion
- Gas pressure in the fiber tunes the resonance over 28 to 100 meV

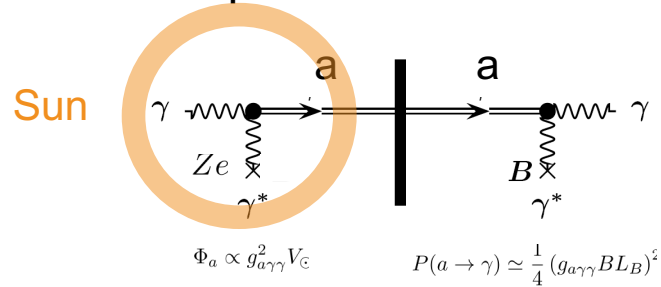
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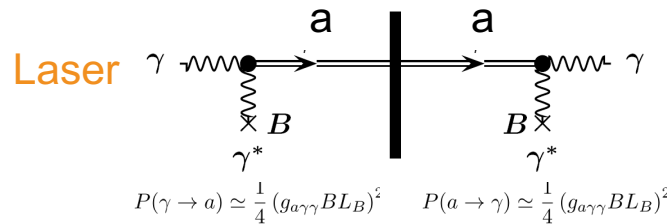
- WISPF1 prototype sensitivity forecast: [Batllori et al., 2510.01221]

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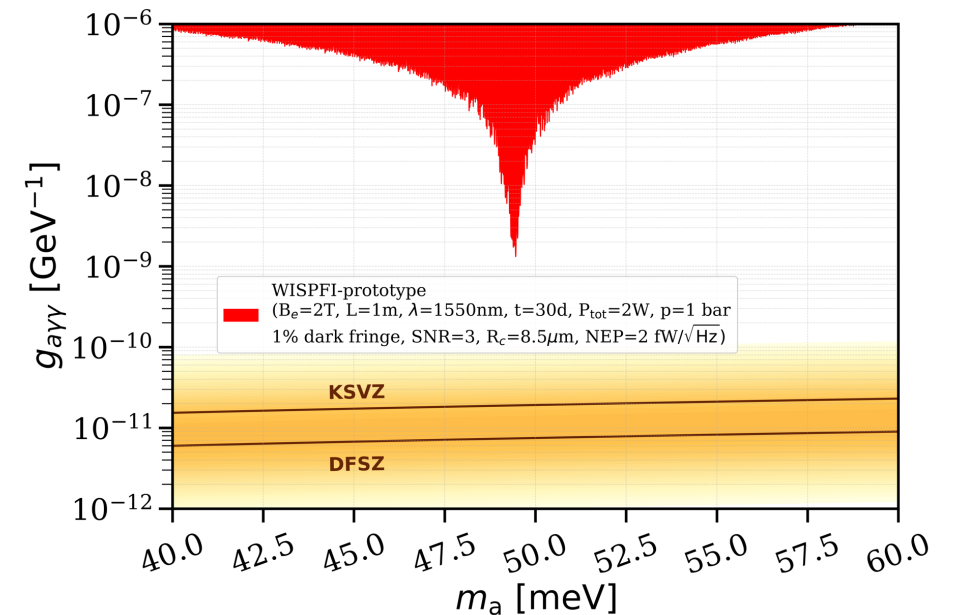
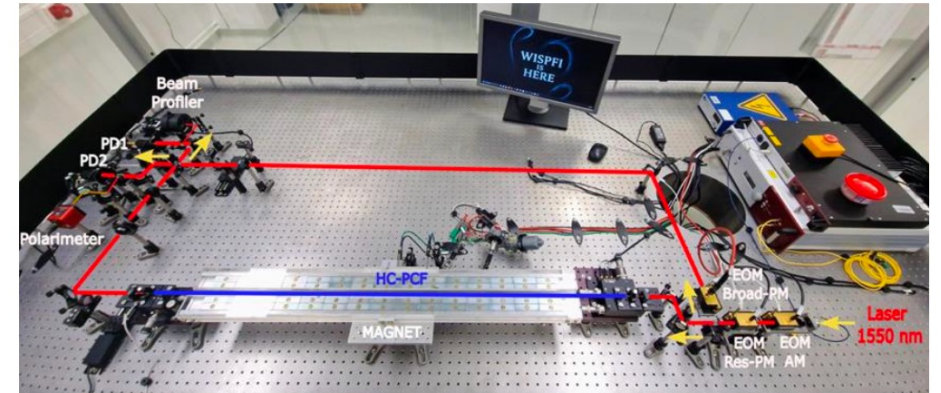
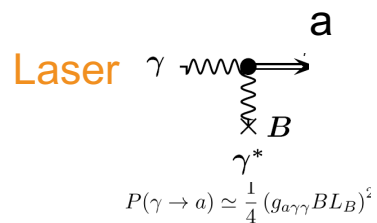
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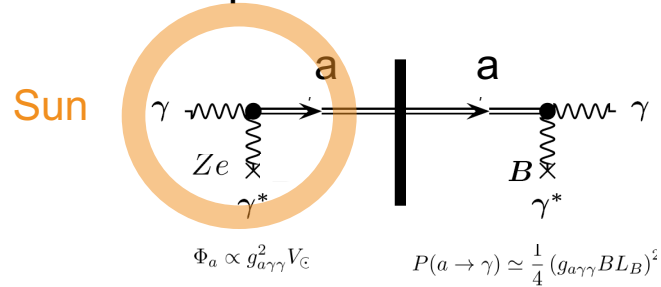
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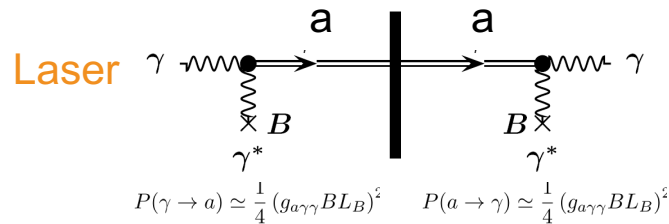
- WISPMI prosp. with 9 T along 100 m: [Batllori et al., 2510.01221]

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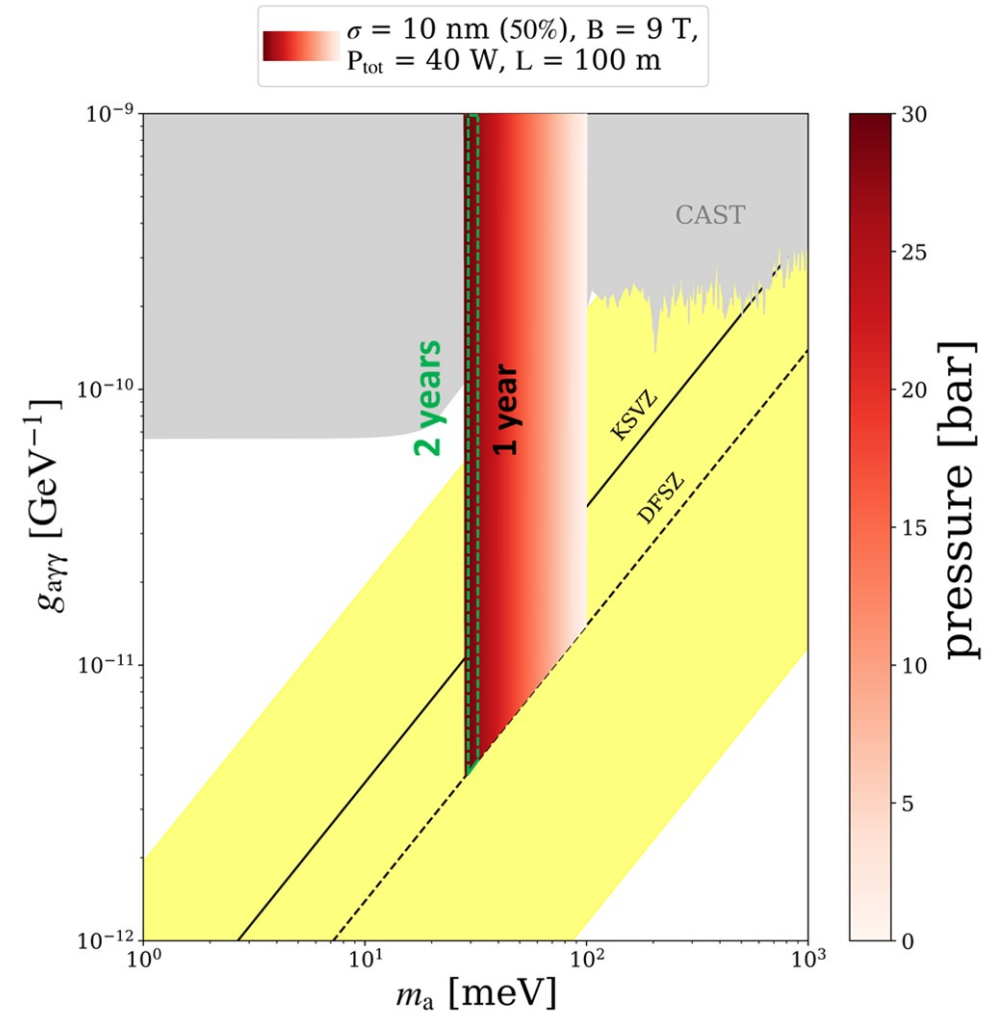
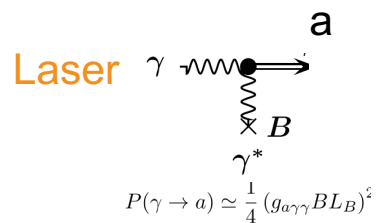
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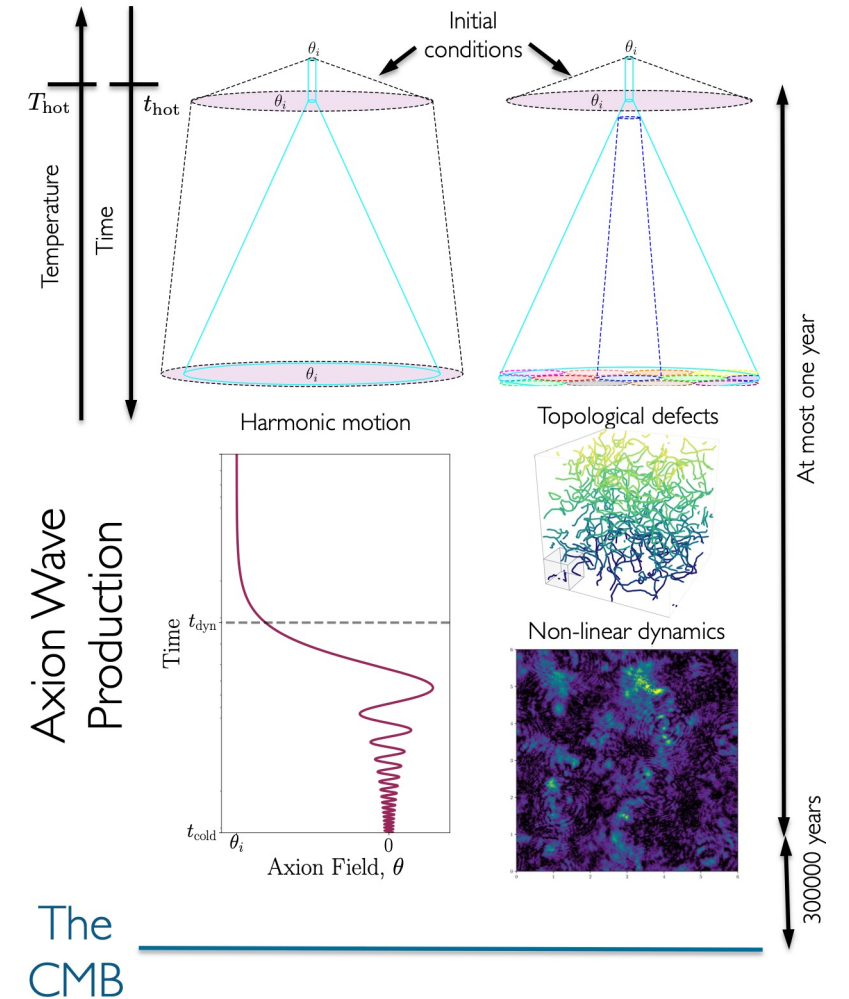
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# Axion Dark Matter

# Axion Dark Matter

DM production mechanisms support entire possible mass range

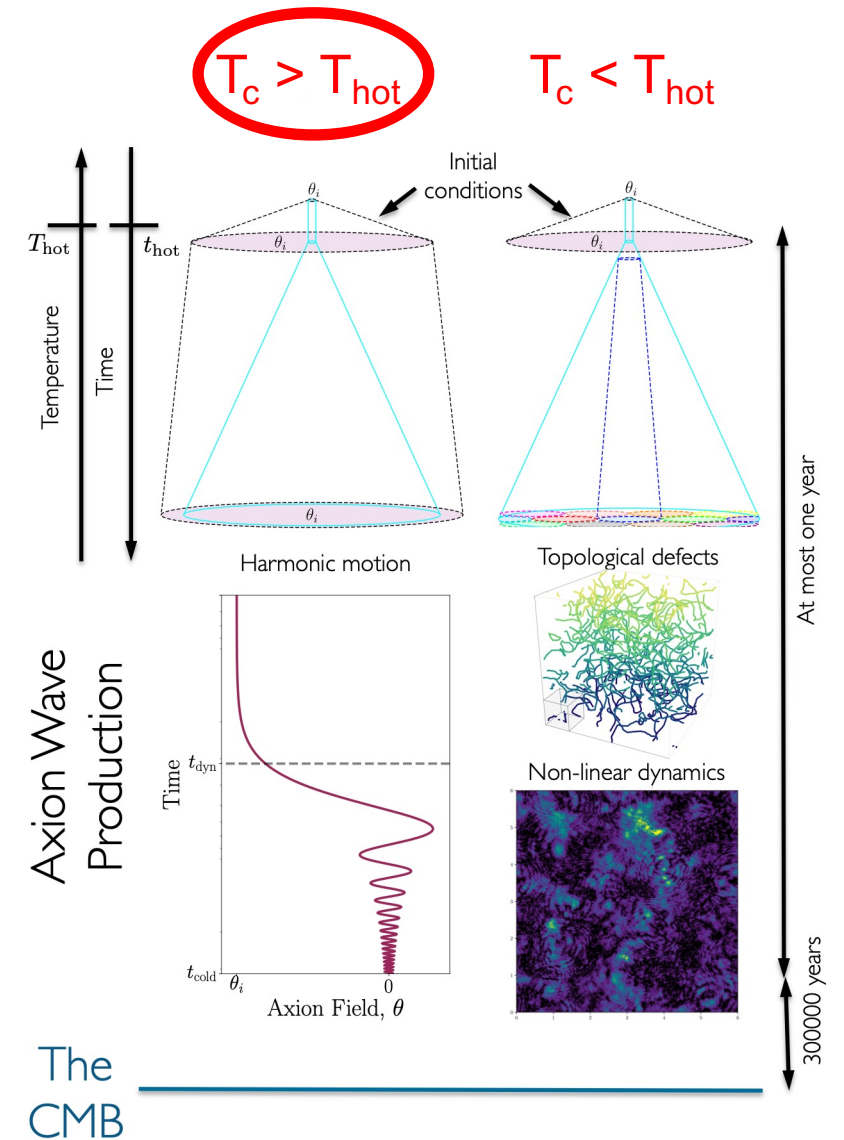


# Axion Dark Matter

DM production mechanisms support entire possible mass range

- **Preinflationary symmetry breaking:** misalignment mechanism (conventional, kinetic, trapped, ...)
- Postinflationary symmetry breaking: decay of topological defects

[Preskill,Wise,Wilczek `83; Abbott,Sikivie `83; Dine,Fischler `83; ... Co,Hall,Harigaya `20; ... Di Luzio et al. `21; ... Davis `86; ...]



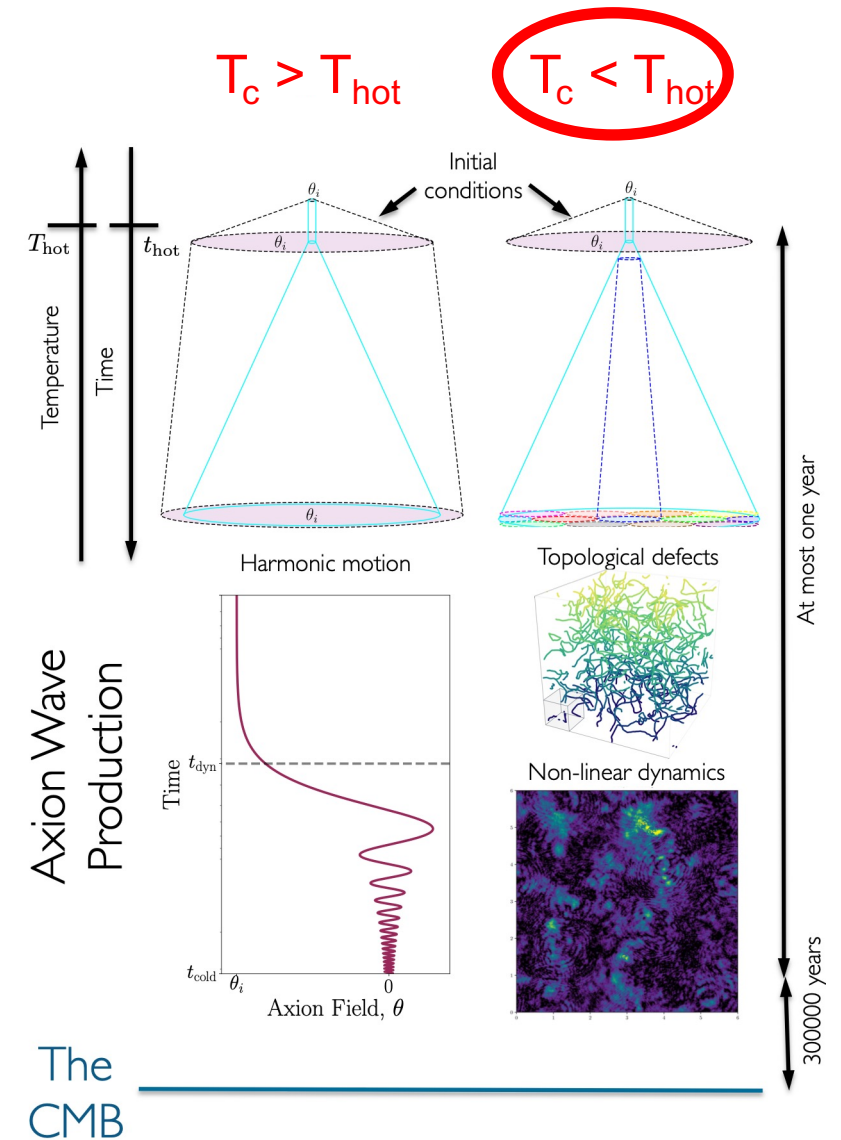
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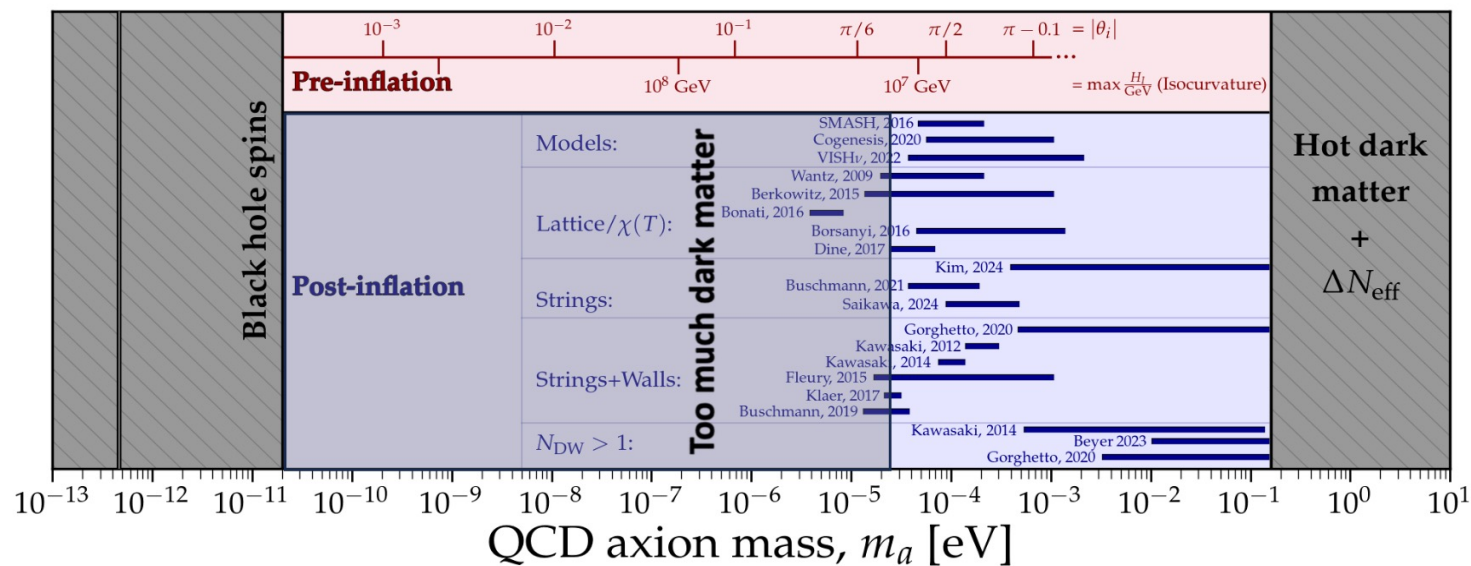
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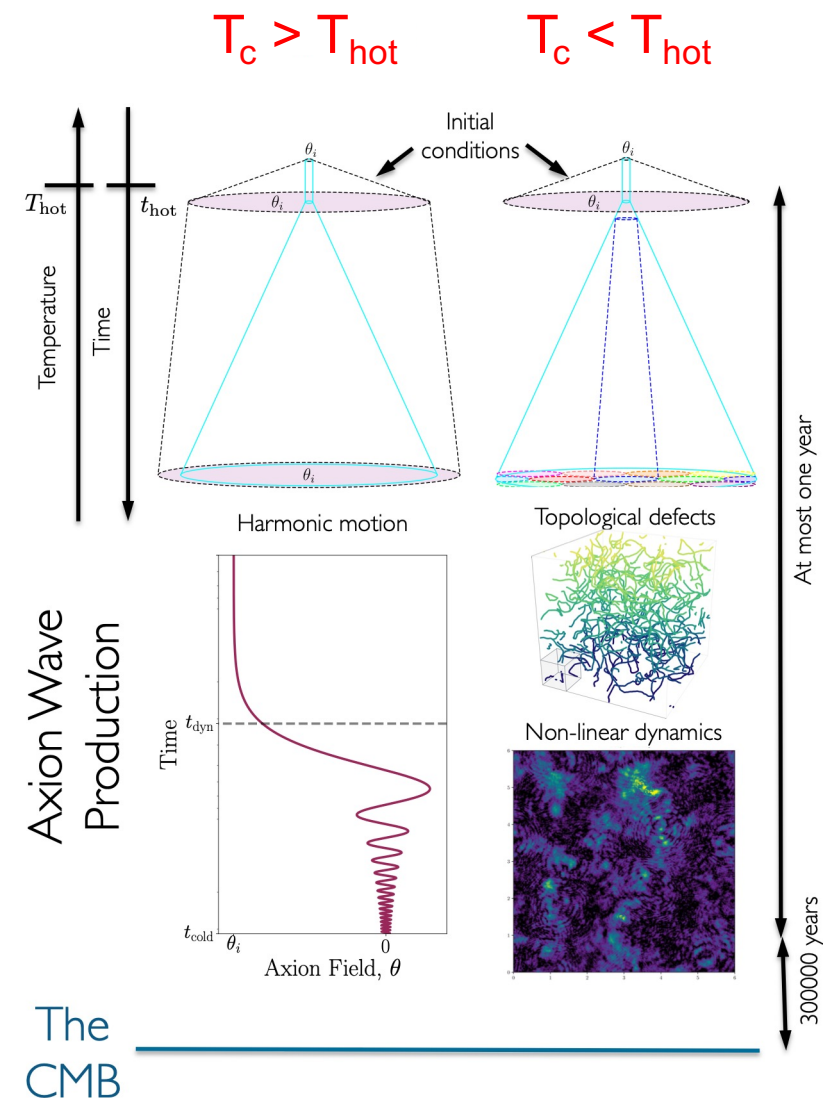
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[adapted from [https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots\\_png/AxionMass.png](https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots_png/AxionMass.png)]



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$$N_a |_{\text{dB}} = n_a \lambda_{\text{dB}}^3 \sim 10^{30} \left( \frac{\rho_a}{0.3 \text{ GeV/cm}^3} \right) \left( \frac{\mu\text{eV}}{m_a} \right)^4$$

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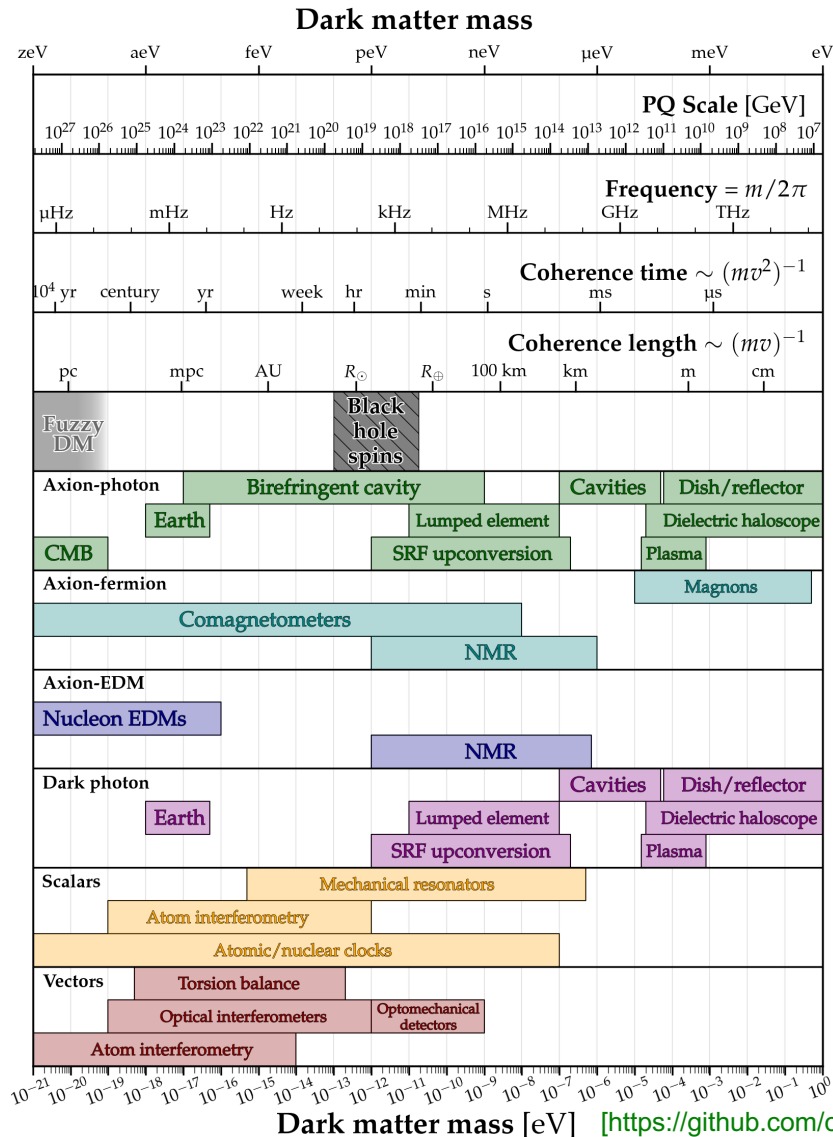
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- Correspondingly, halo dark-matter axions behave as an approximately spatially homogeneous and monochromatic classical oscillating field,

$$a(t) \simeq \sqrt{2\rho_a} \cos(m_a t)/m_a$$

# Axion Dark Matter Experiments

## Variety of experimental techniques



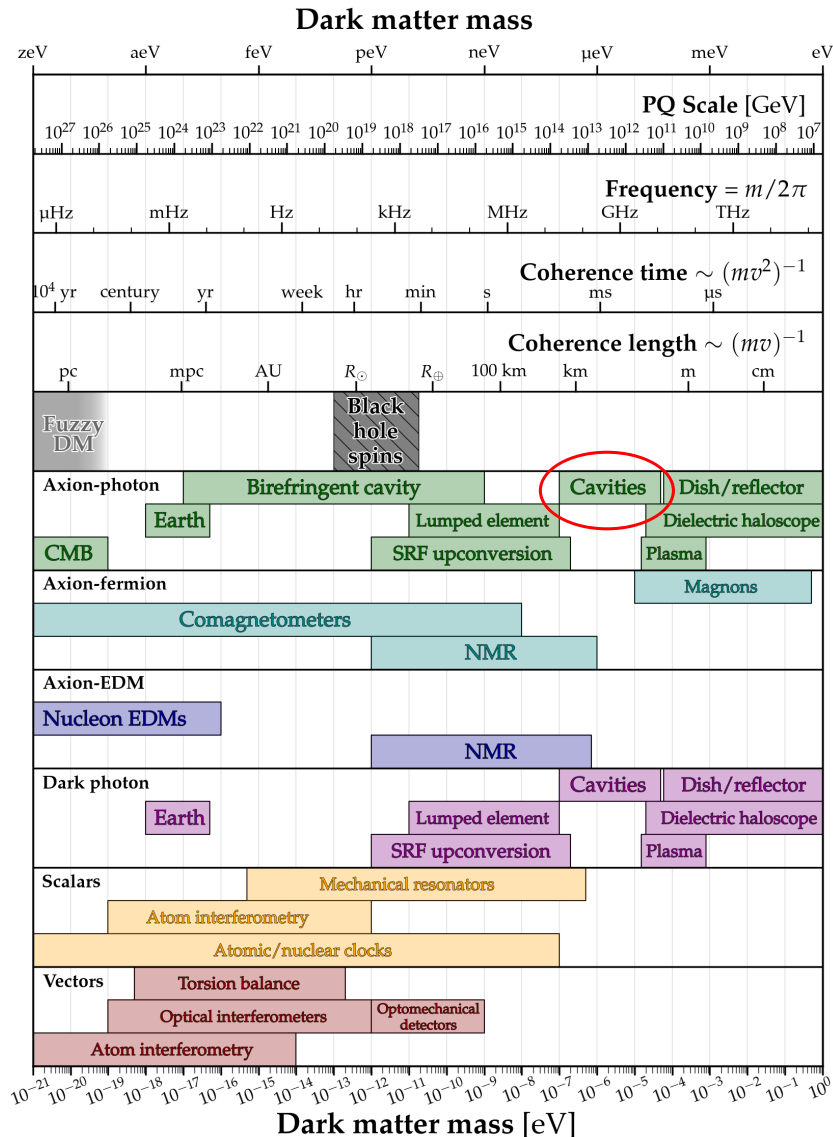
$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu}$$

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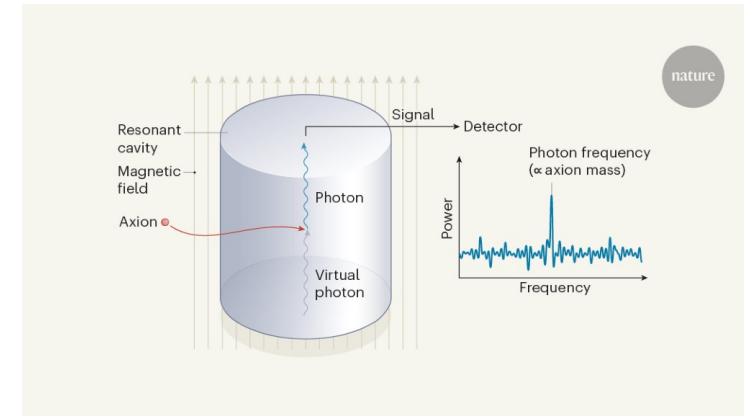
$$\mathcal{L} \supset -\frac{i}{2} g_{aN\gamma} a \bar{\psi}_N \sigma_{\mu\nu} \gamma_5 \psi_N F^{\mu\nu}$$

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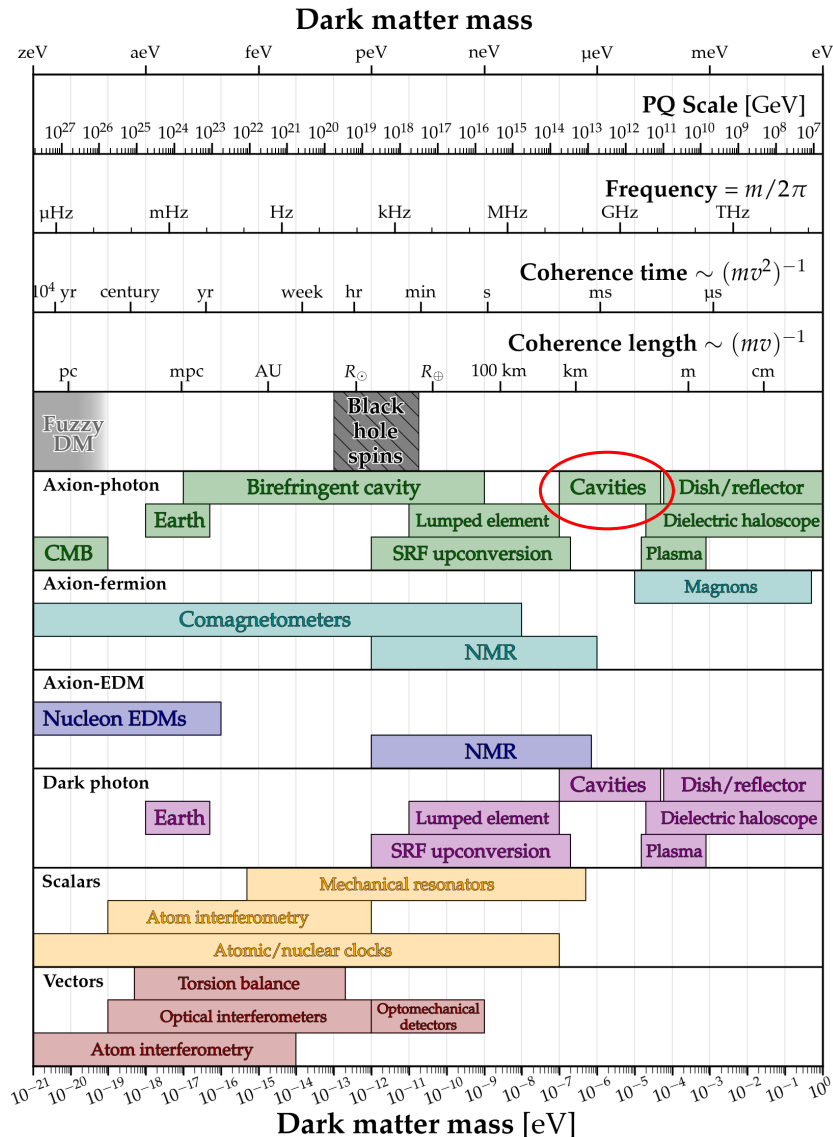


- **Cavity concept:** DM axion converts into photon in micro-wave cavity placed in magnetic field [Sikivie 83]

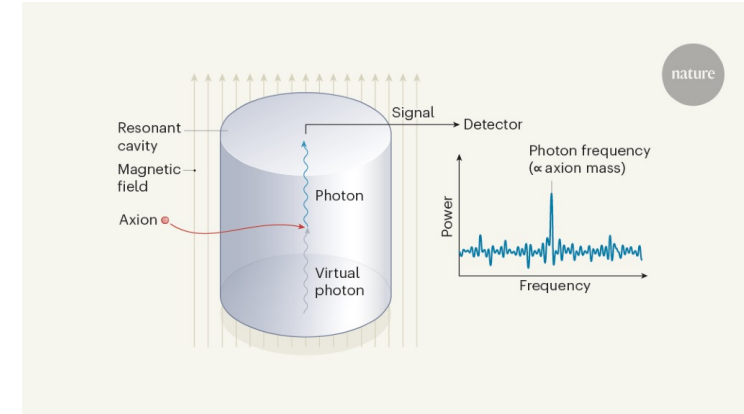


# Axion Dark Matter Experiments

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- **Cavity concept:** DM axion converts into photon in microwave cavity placed in magnetic field [Sikivie 83]



- If axion mass matches resonance frequency of cavity,

$$m_a = 2\pi\nu_{\text{res}} \sim 4 \mu\text{eV} \left( \frac{\nu_{\text{res}}}{\text{GHz}} \right)$$

power output

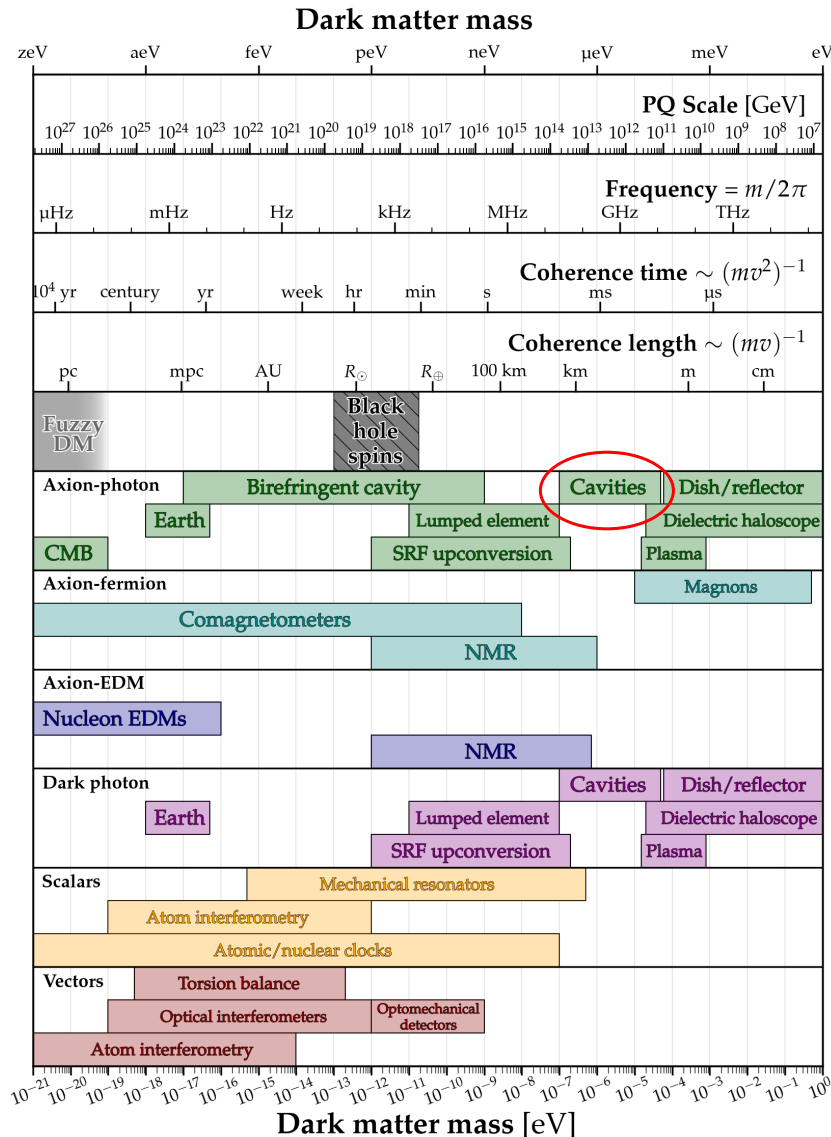
$$P_{\text{out}} \sim g_{a\gamma}^2 \rho_a B_0^2 V Q$$

enhanced by quality factor

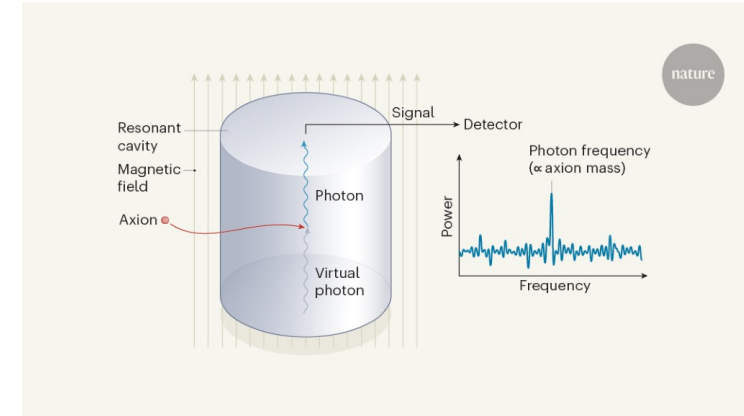
$$Q \sim 10^5$$

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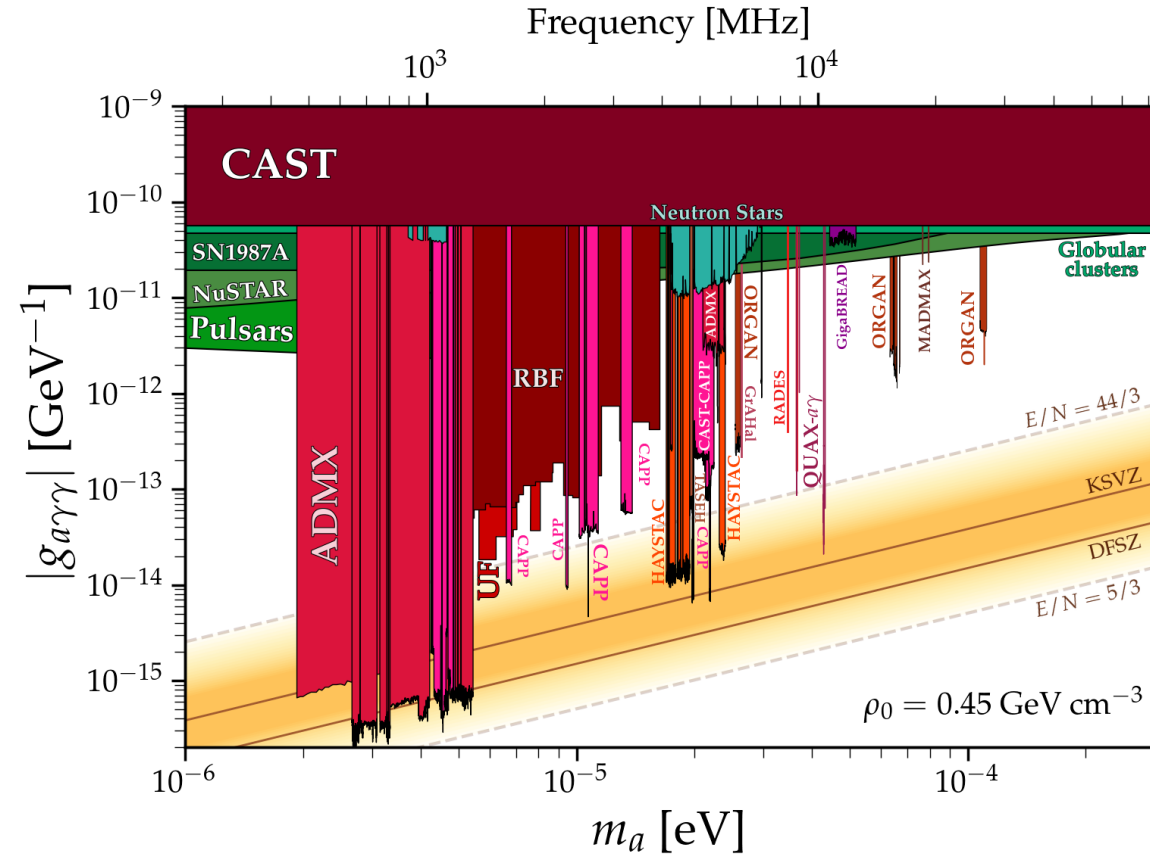
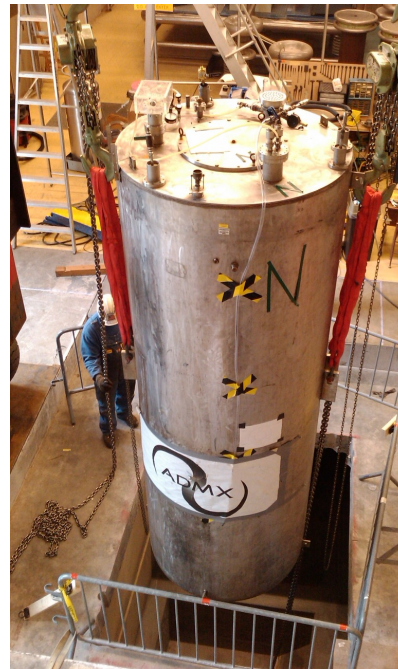
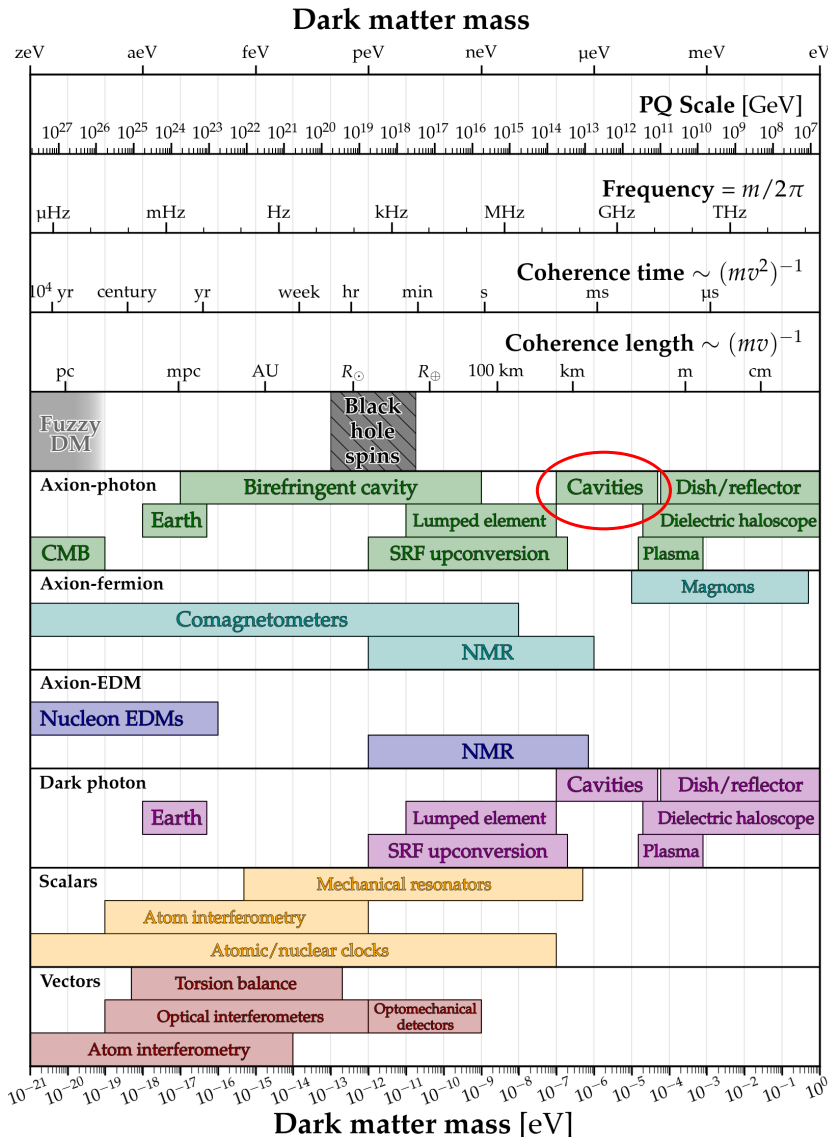
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- Need to scan by tuning resonance frequency

# Axion Dark Matter Experiments

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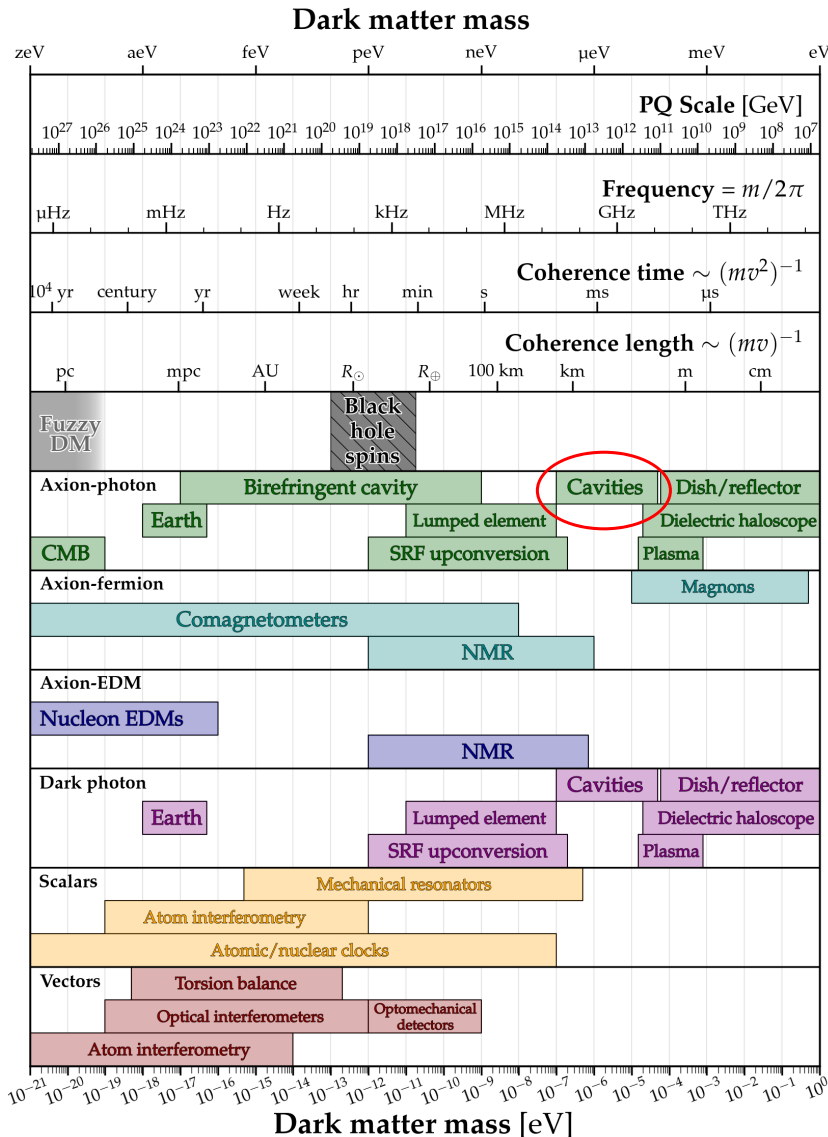
Axion DM cavity experiments have dipped already deep in axion band:



[https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots\\_png/AxionPhoton\\_RadioFreqCloseup.png](https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots_png/AxionPhoton_RadioFreqCloseup.png)

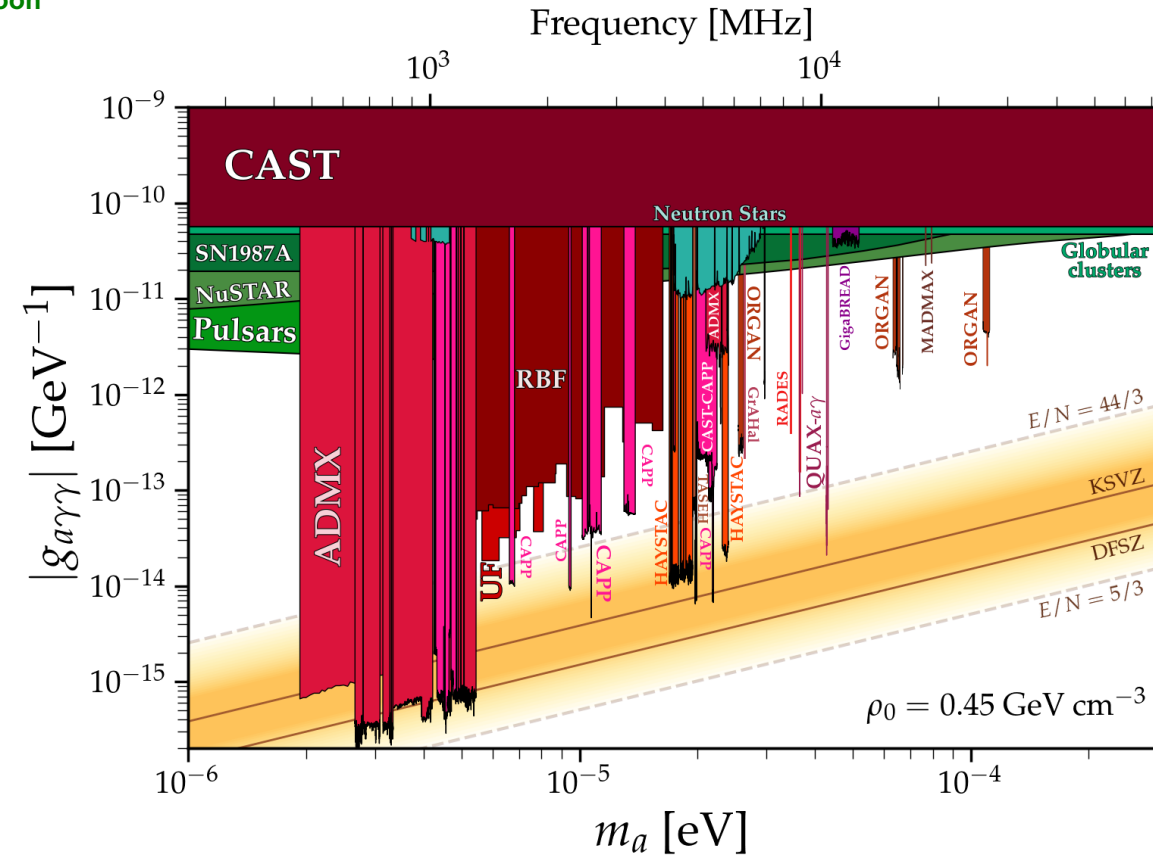
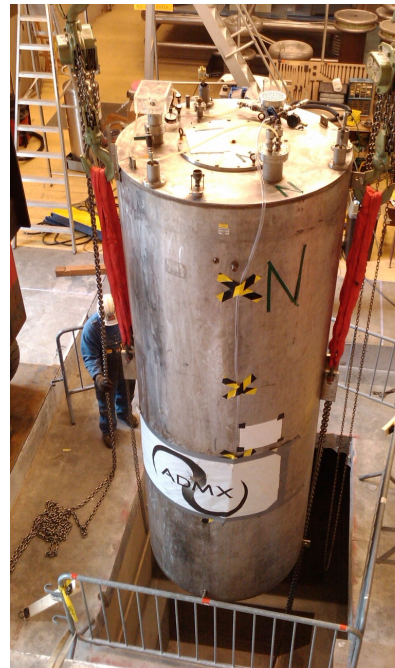
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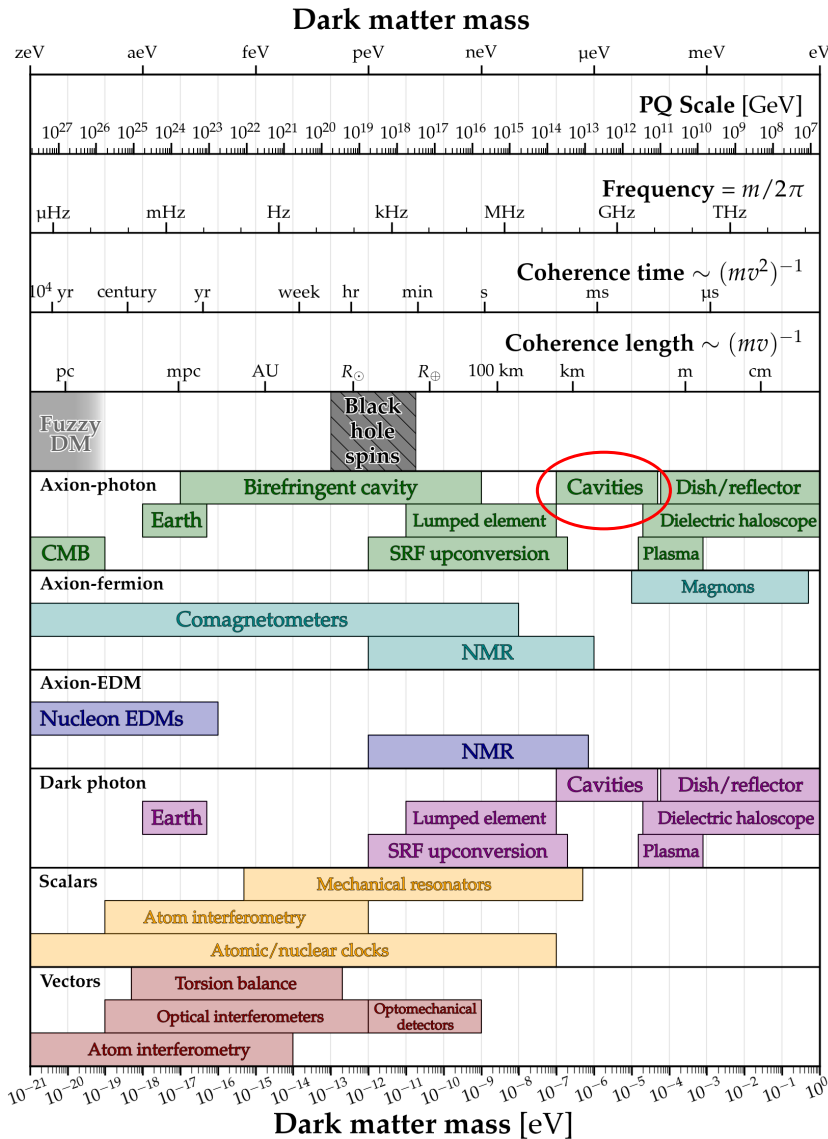
Small-group visits of Quantum Sensors for the Hidden Sector (QSHS) laboratory on Wed 24/06 afternoon (sign up at registration desk)



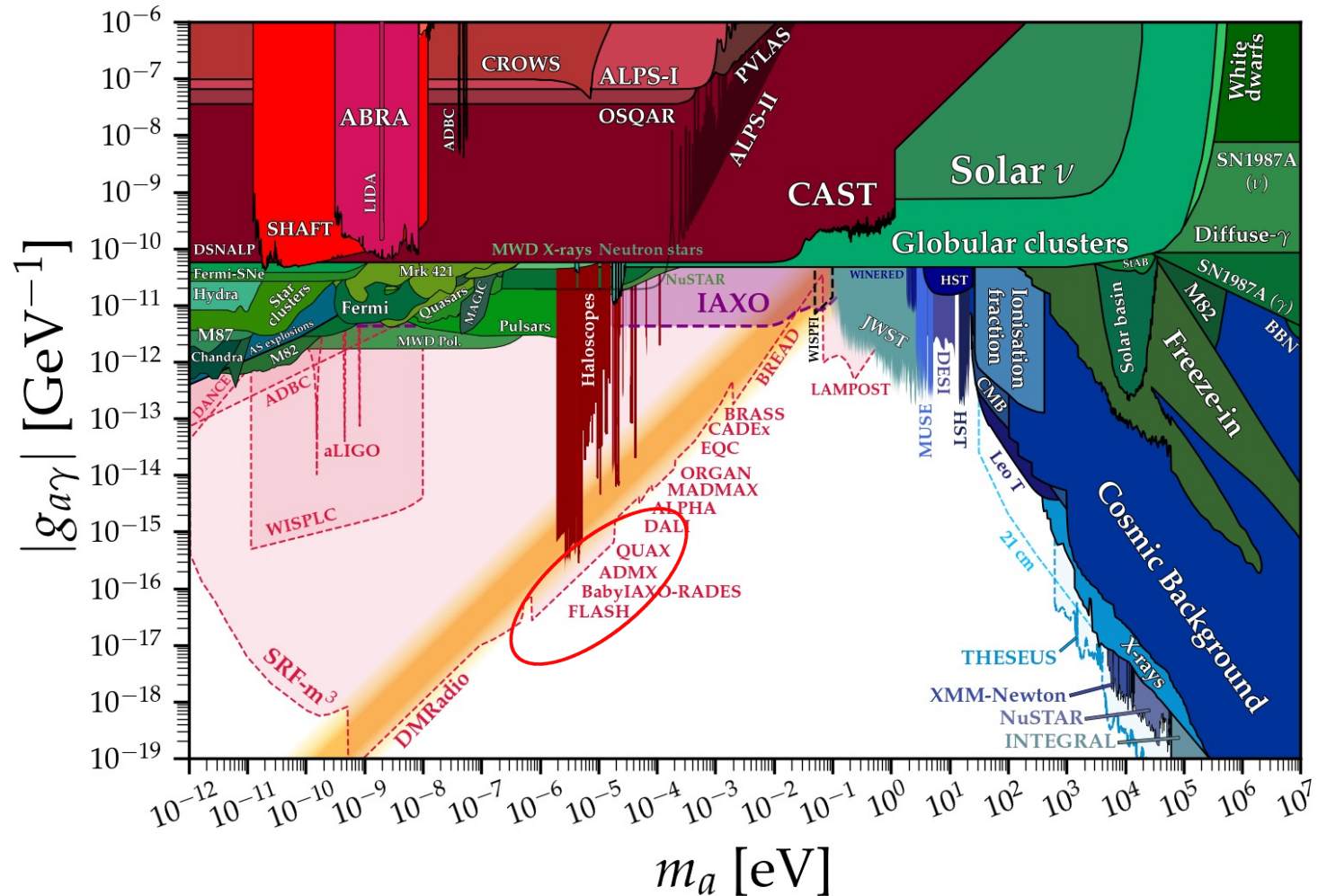
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# Axion Dark Matter Experiments

## Variety of experimental techniques



Axion DM cavity experiments have good prospects to probe entire band in mass range  $\mu\text{eV} \lesssim m_a \lesssim 100 \mu\text{eV}$ :

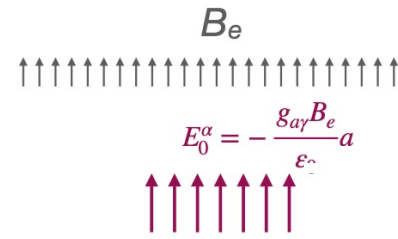
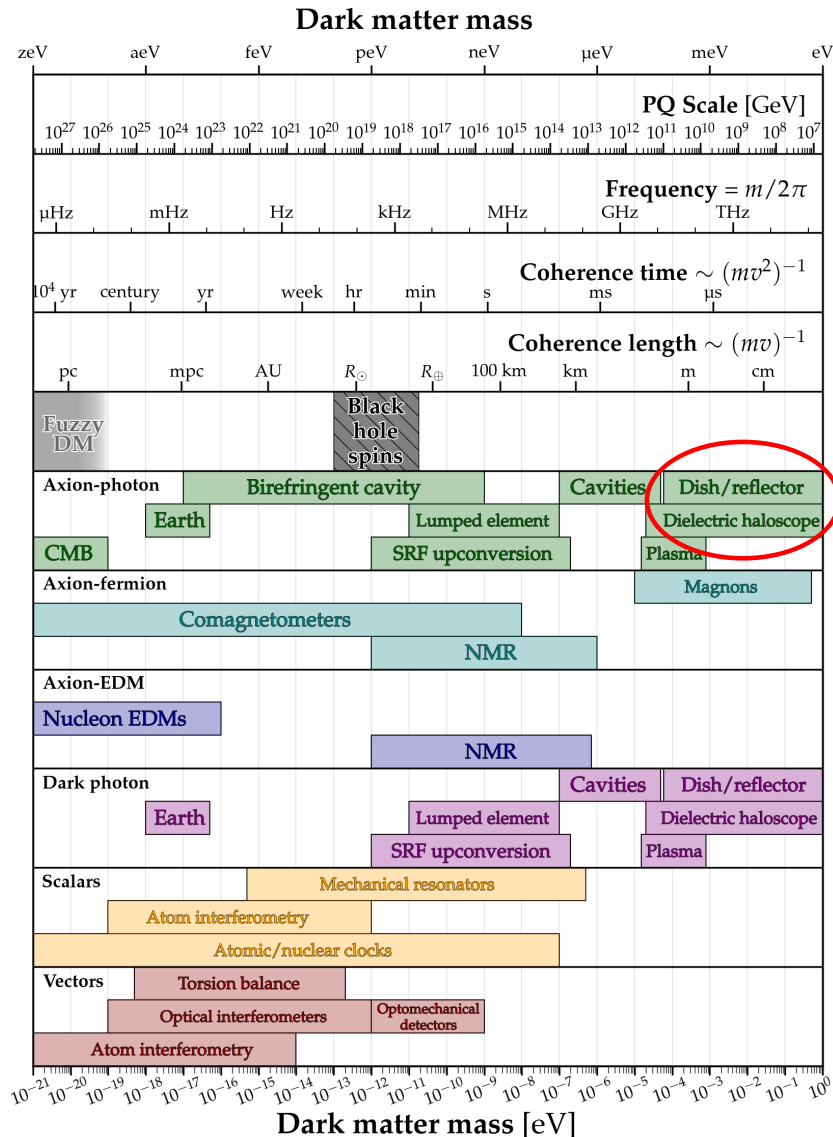


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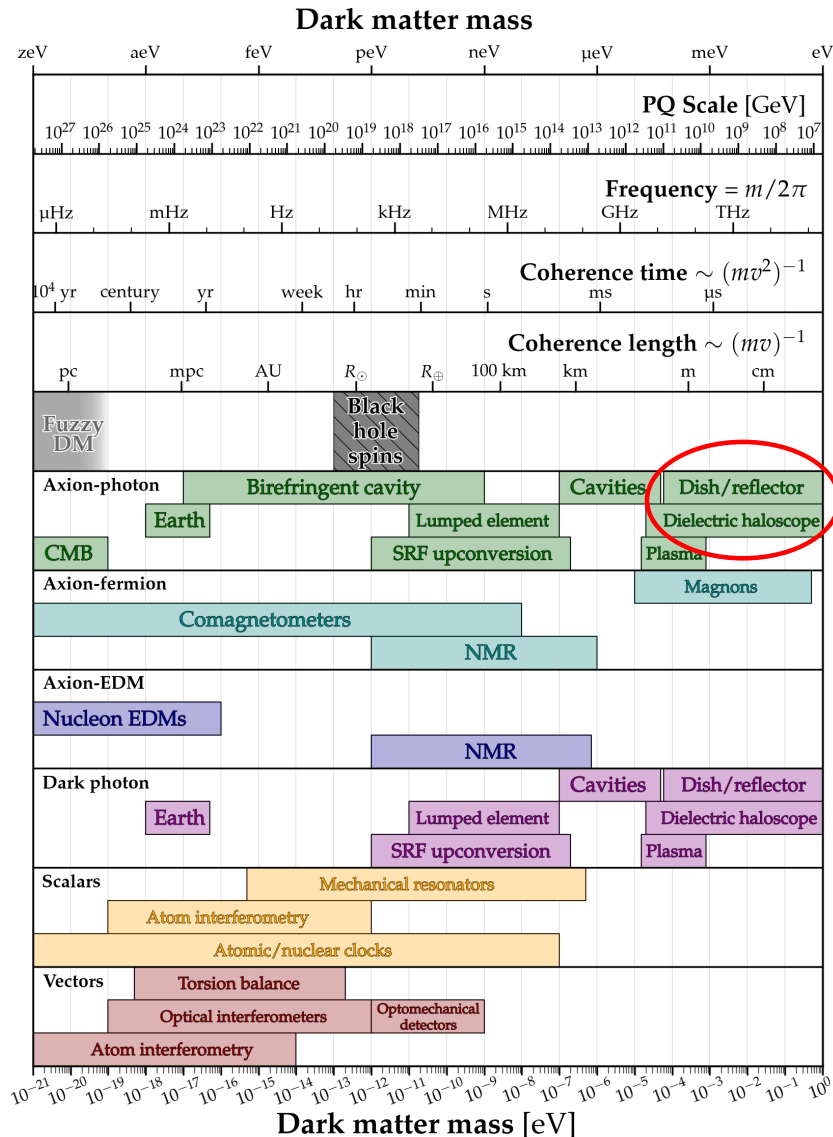
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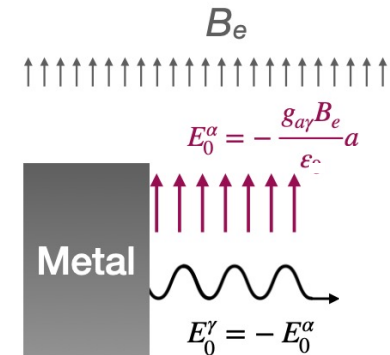
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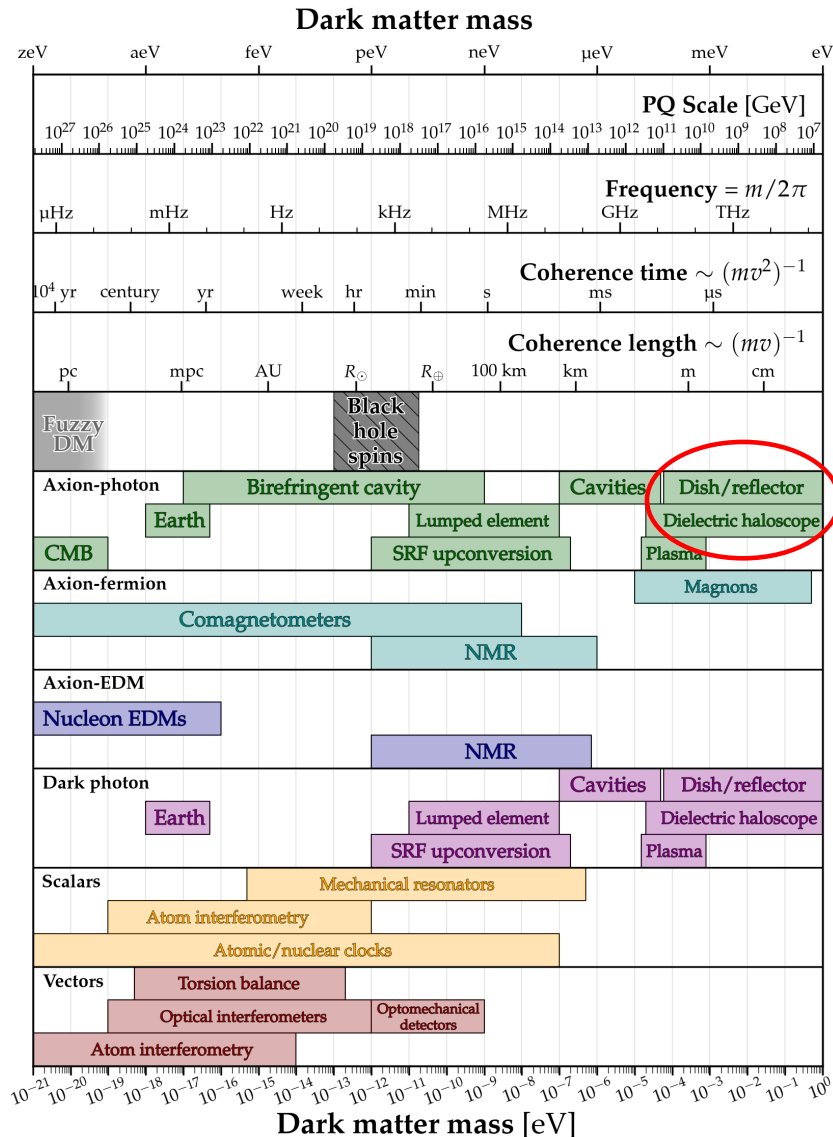
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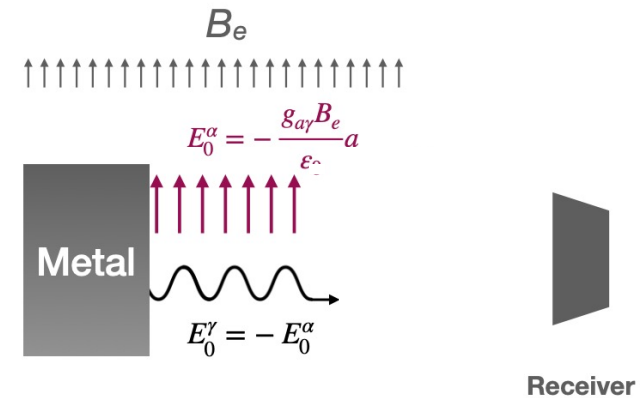
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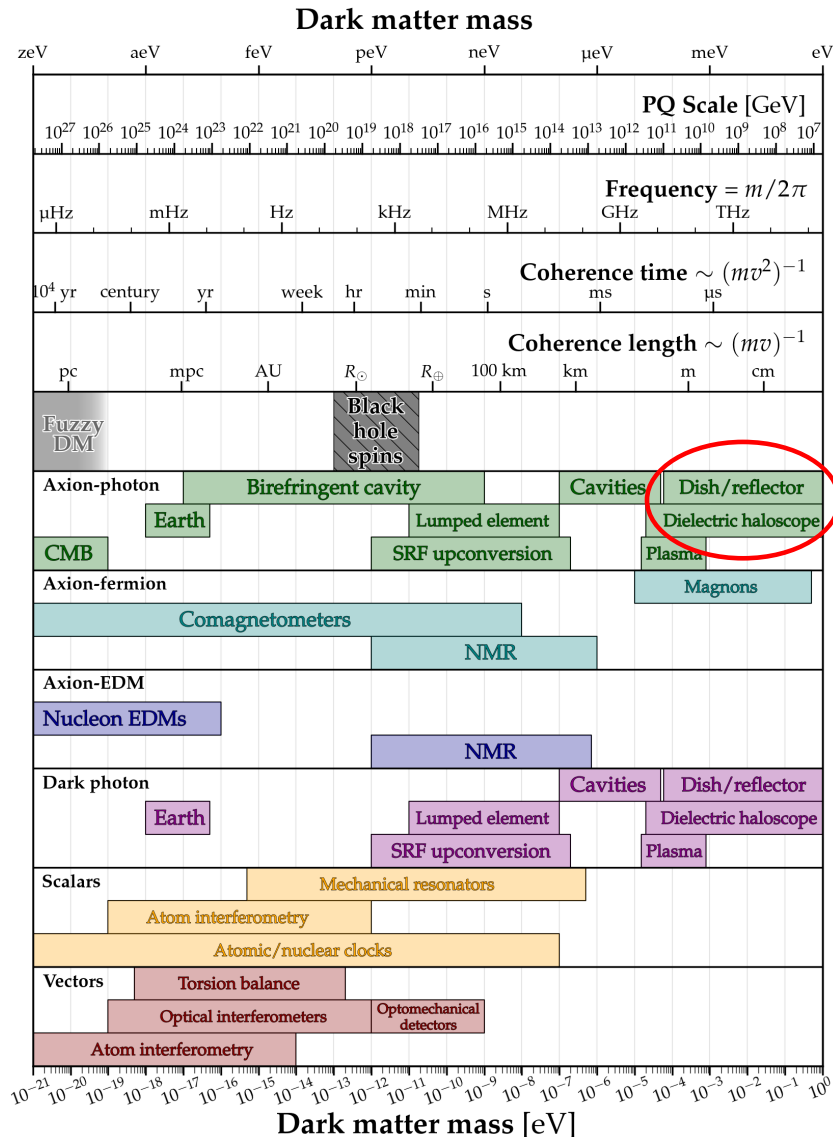
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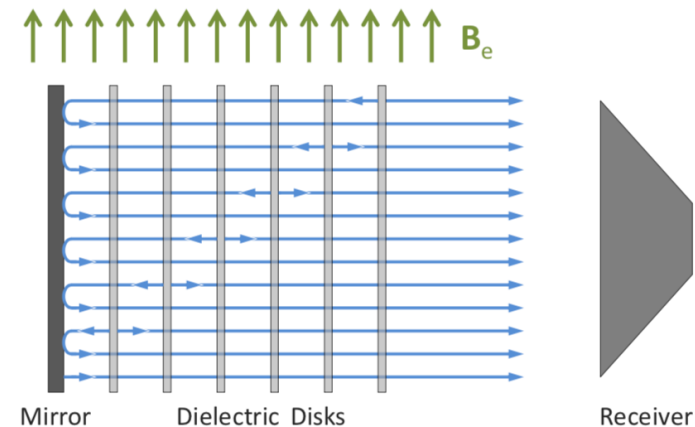
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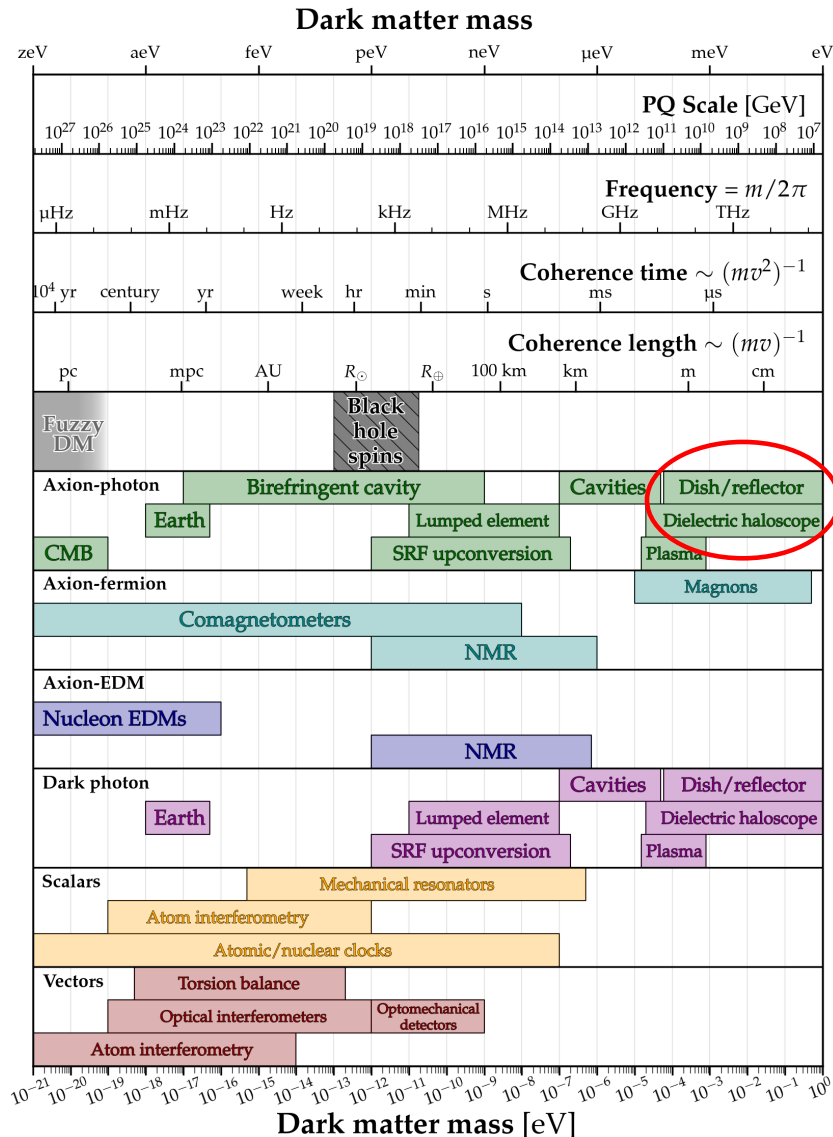
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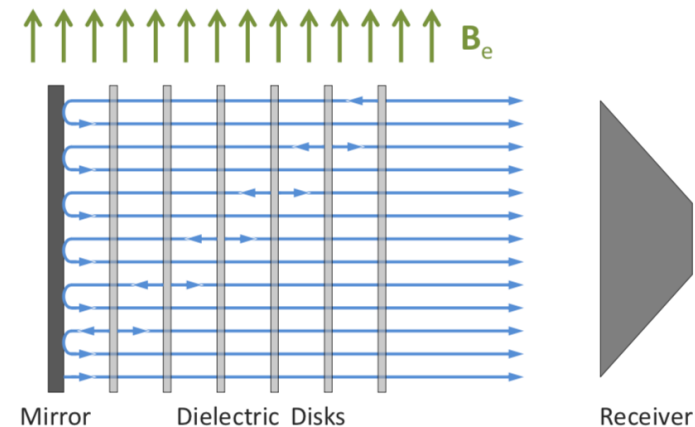
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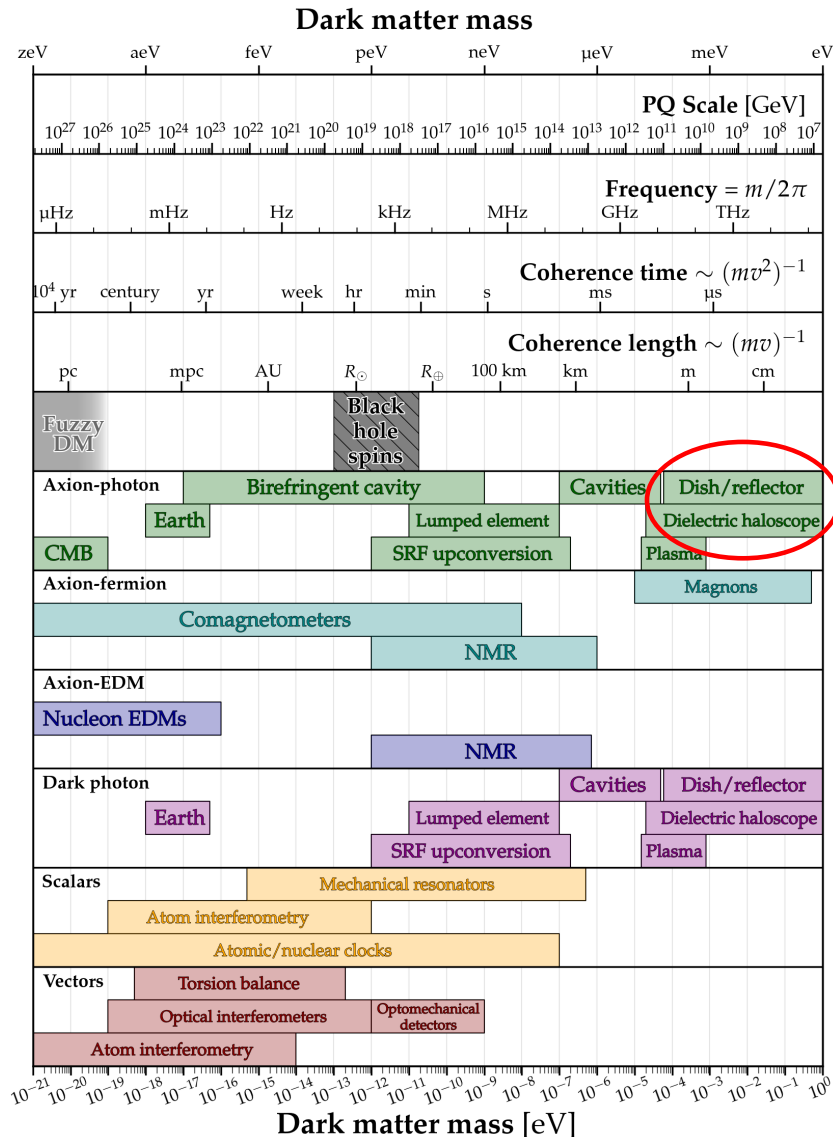
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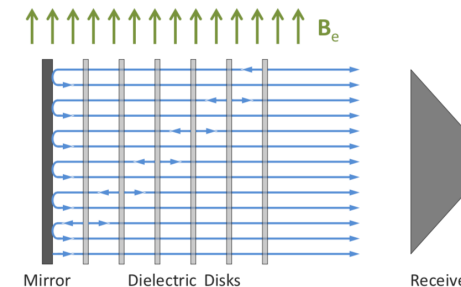
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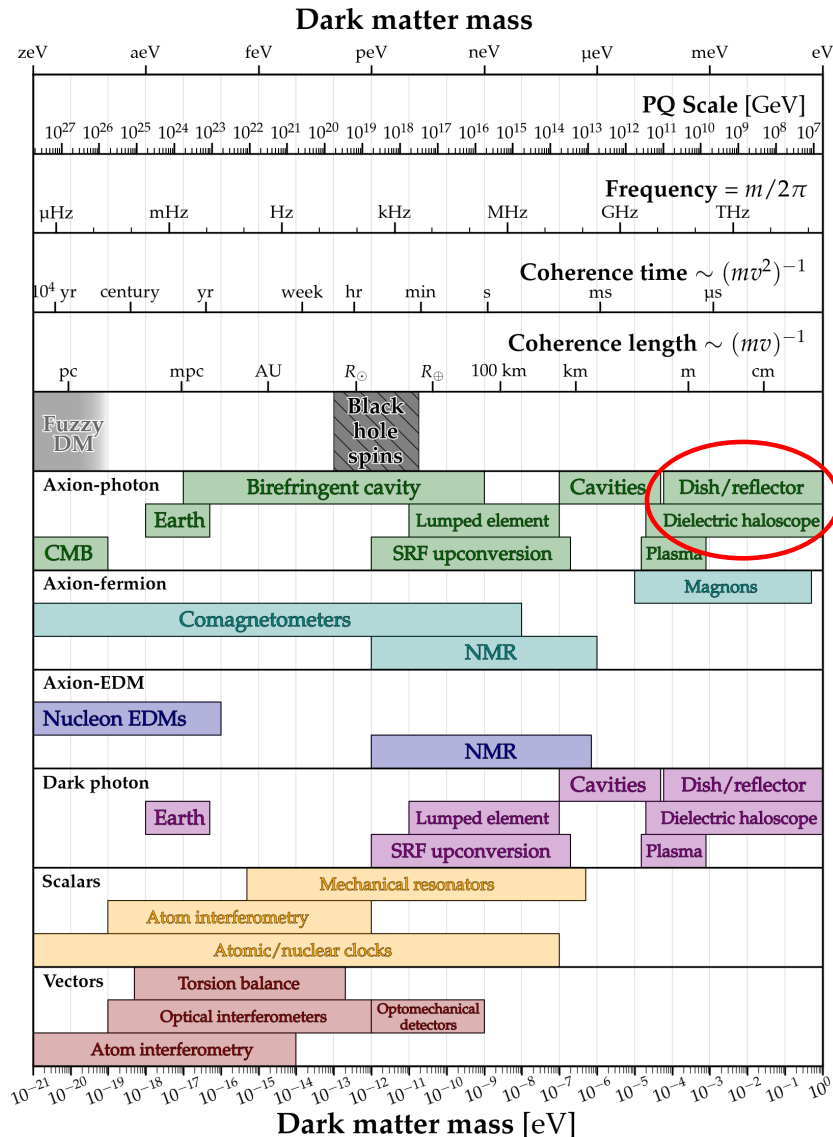
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## Variety of experimental techniques



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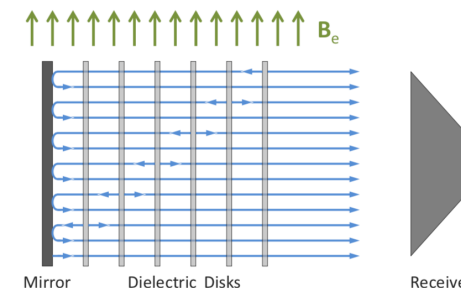
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### Dielectric haloscope concept:

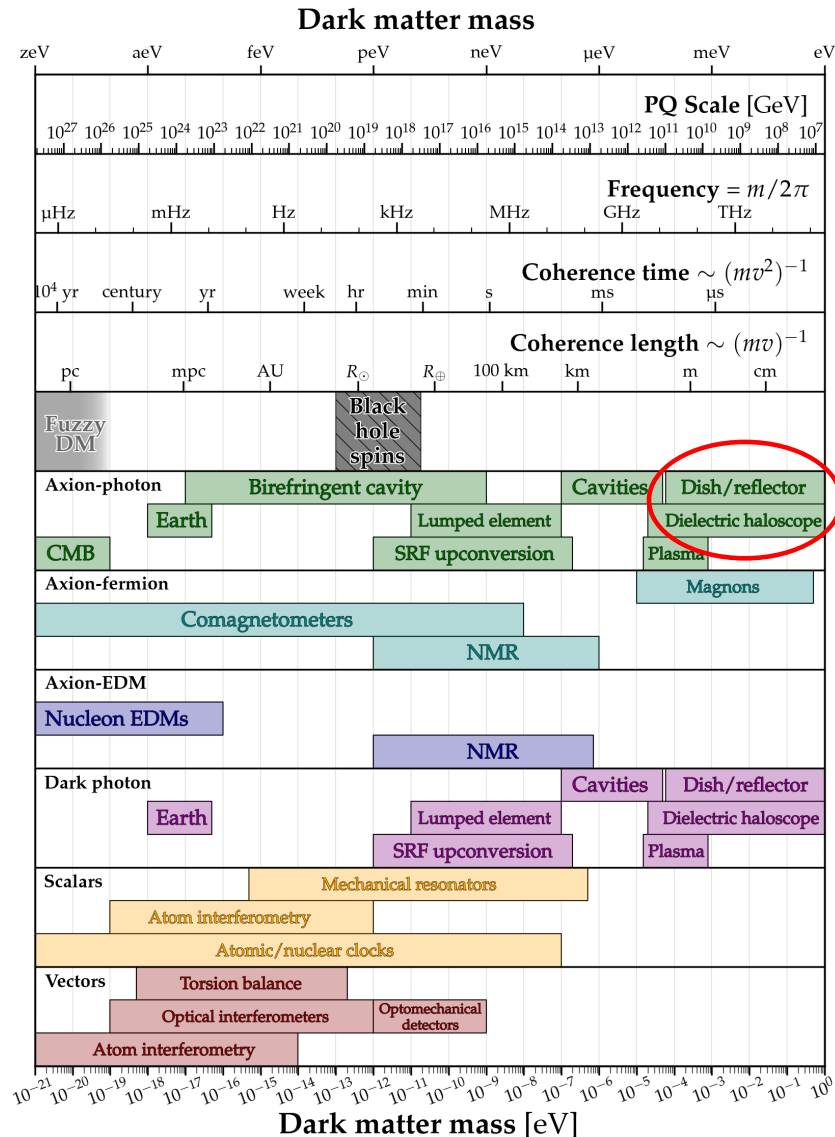
[Jaeckel, Redondo 13; Millar, Raffelt, Redondo, Steffen 16]

- Interfaces between media with different permittivities  $\epsilon$  produce discontinuities in the axion-induced field  $E_a$
- Each interface emits EM waves; the waves are transmitted and reflected by the disk stack
- By tuning the disk spacings, the emitted waves interfere constructively over a chosen frequency band, boosting the signal strength
- Moving the disks shifts the boosted band and scans the axion mass



# Axion Dark Matter Experiments

## Variety of experimental techniques

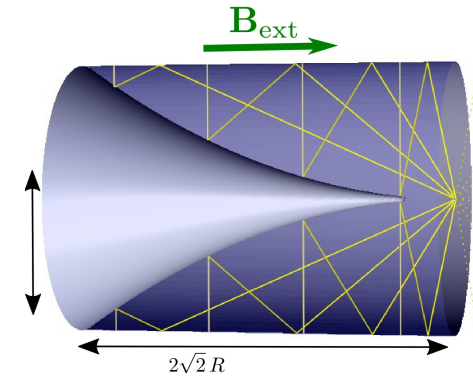


## Dish/reflector example:

### BREAD proposal

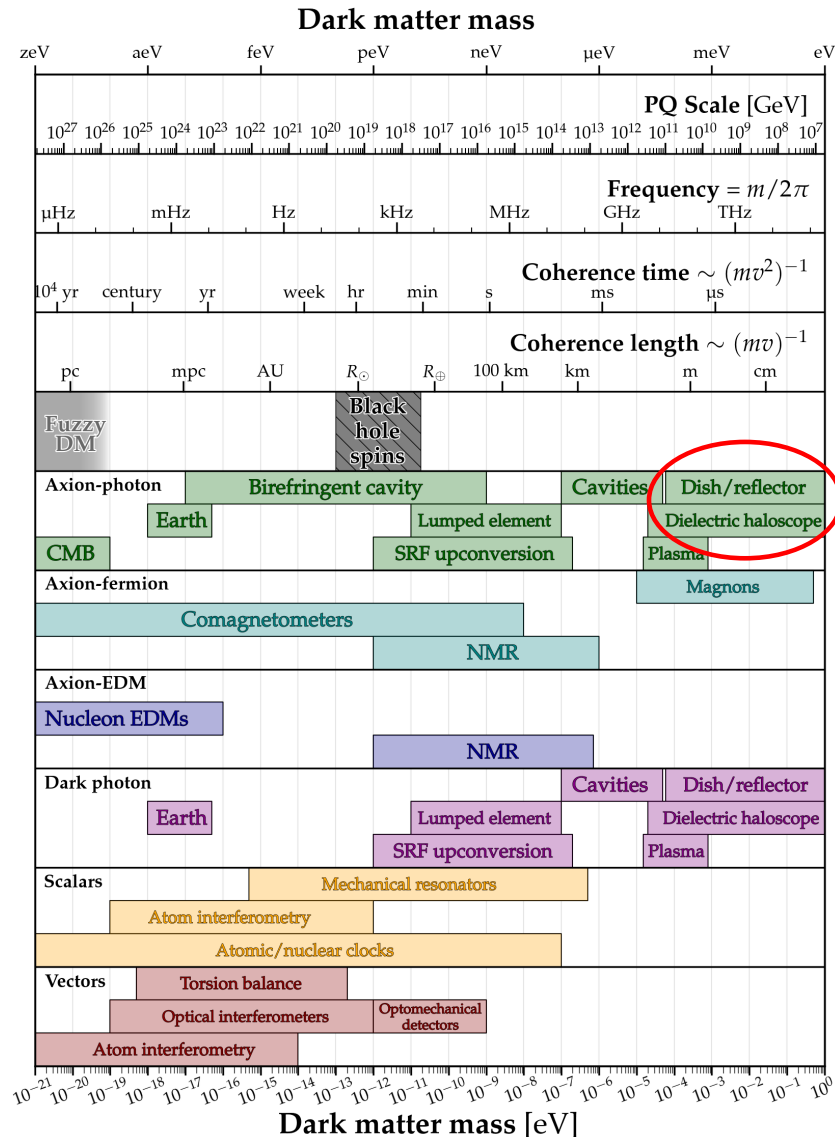
[Liu et al., 22]

- Coaxial dish antenna allows use of solenoidal magnetic field



# Axion Dark Matter Experiments

## Variety of experimental techniques

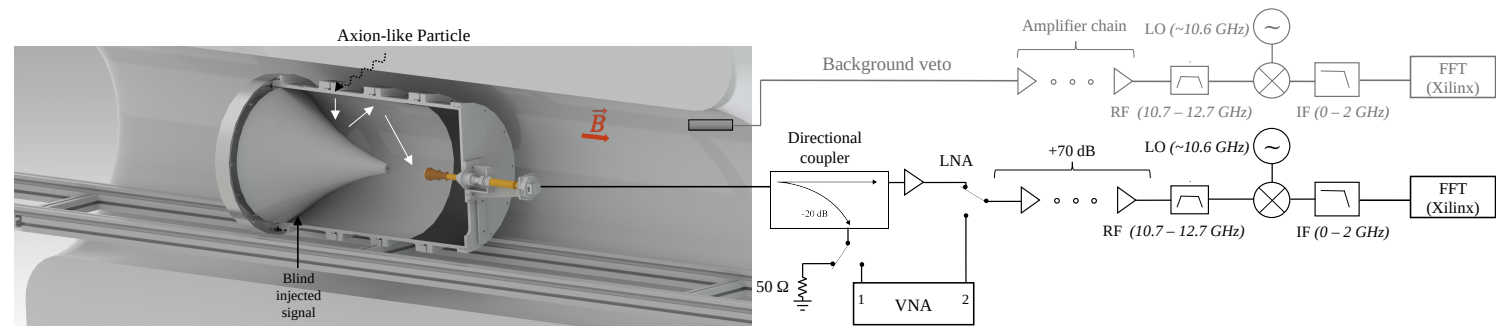


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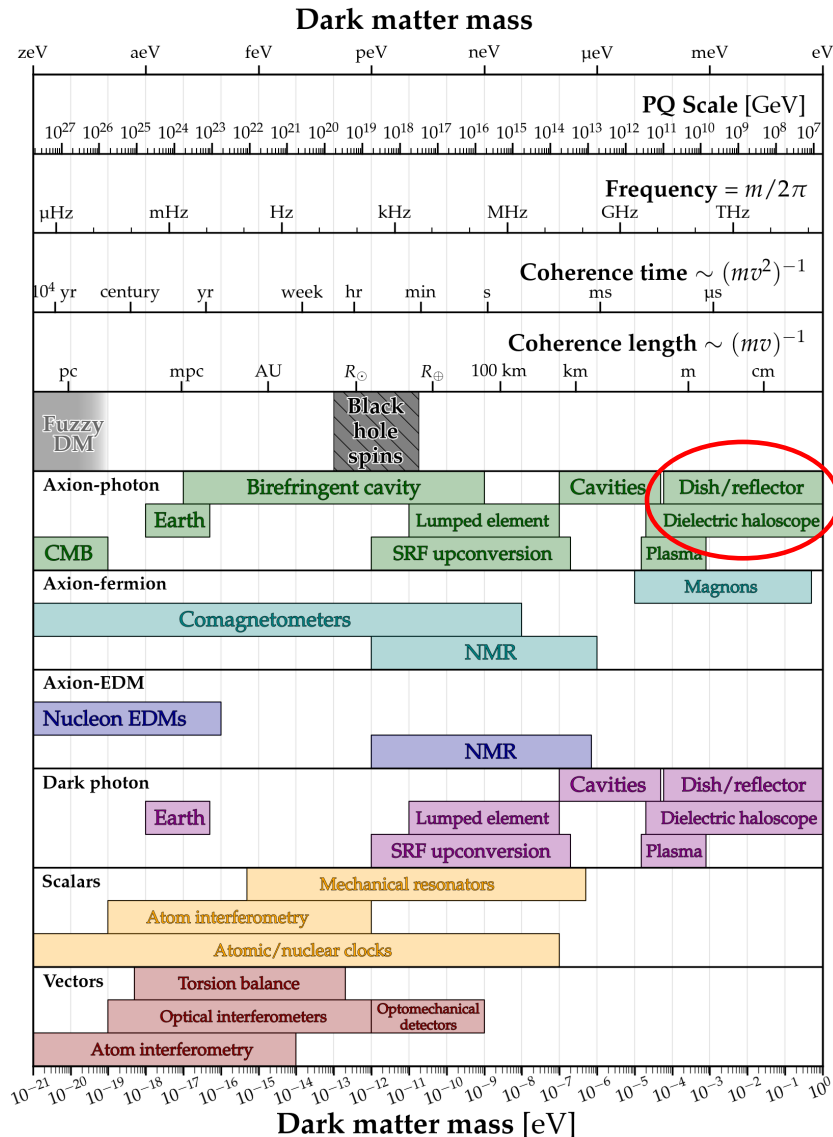
[Liu et al., 22]

- Coaxial dish antenna allows use of solenoidal magnetic field
- GigaBREAD @ Argonne:** [GigaBREAD Coll., PRL 134, 171002 (2025)]
- Room temperature pilot experiment within 3.9 T solenoid



# Axion Dark Matter Experiments

## Variety of experimental techniques



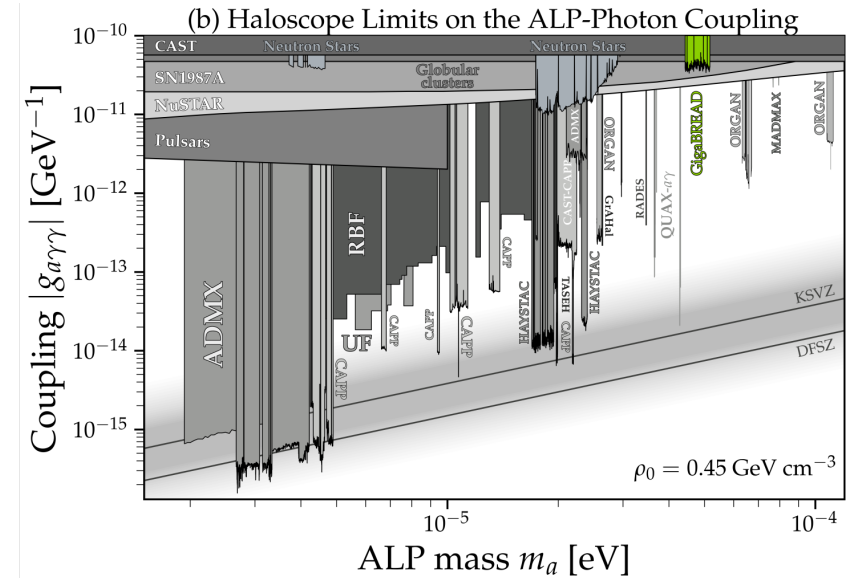
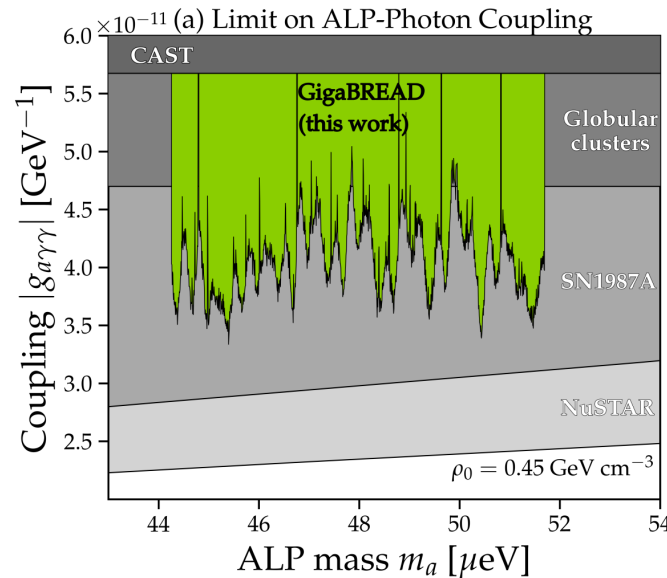
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[Liu et al., 22]

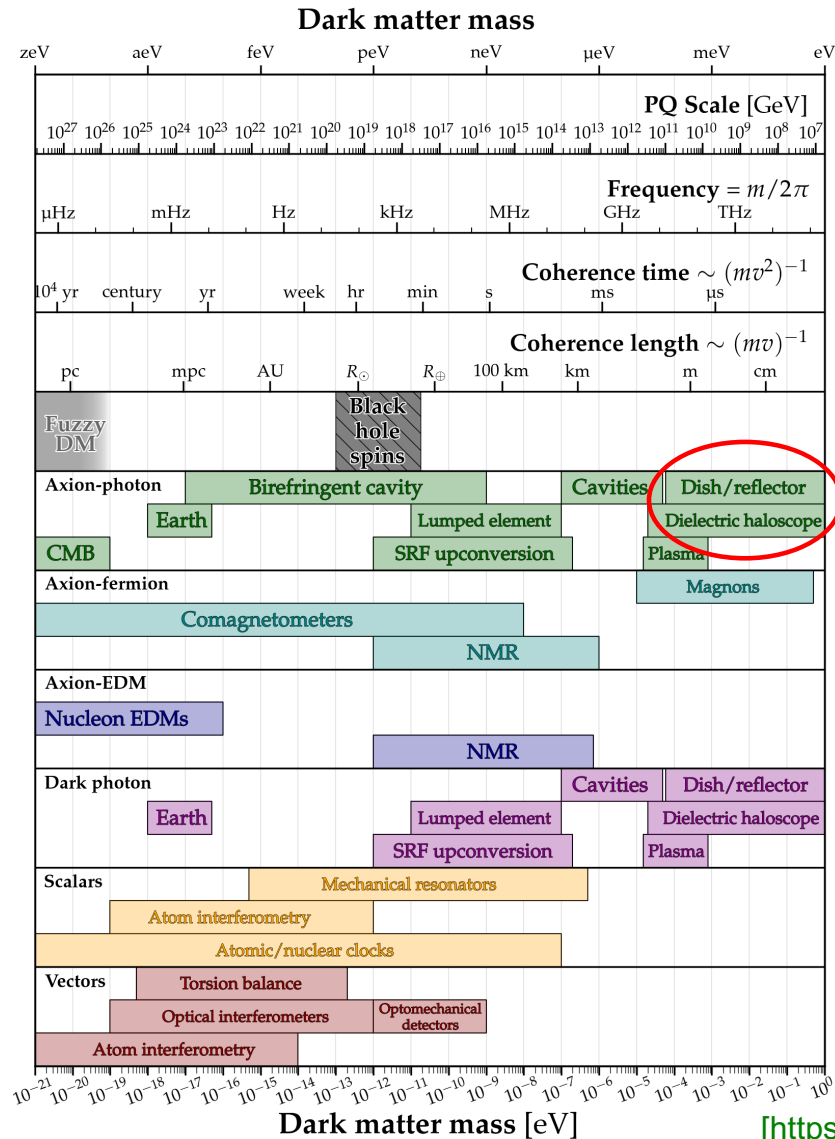
- Coaxial dish antenna allows use of solenoidal magnetic field
- GigaBREAD** @ Argonne demonstrated feasibility of proposal:

[GigaBREAD Coll., PRL 134, 171002 (2025)]

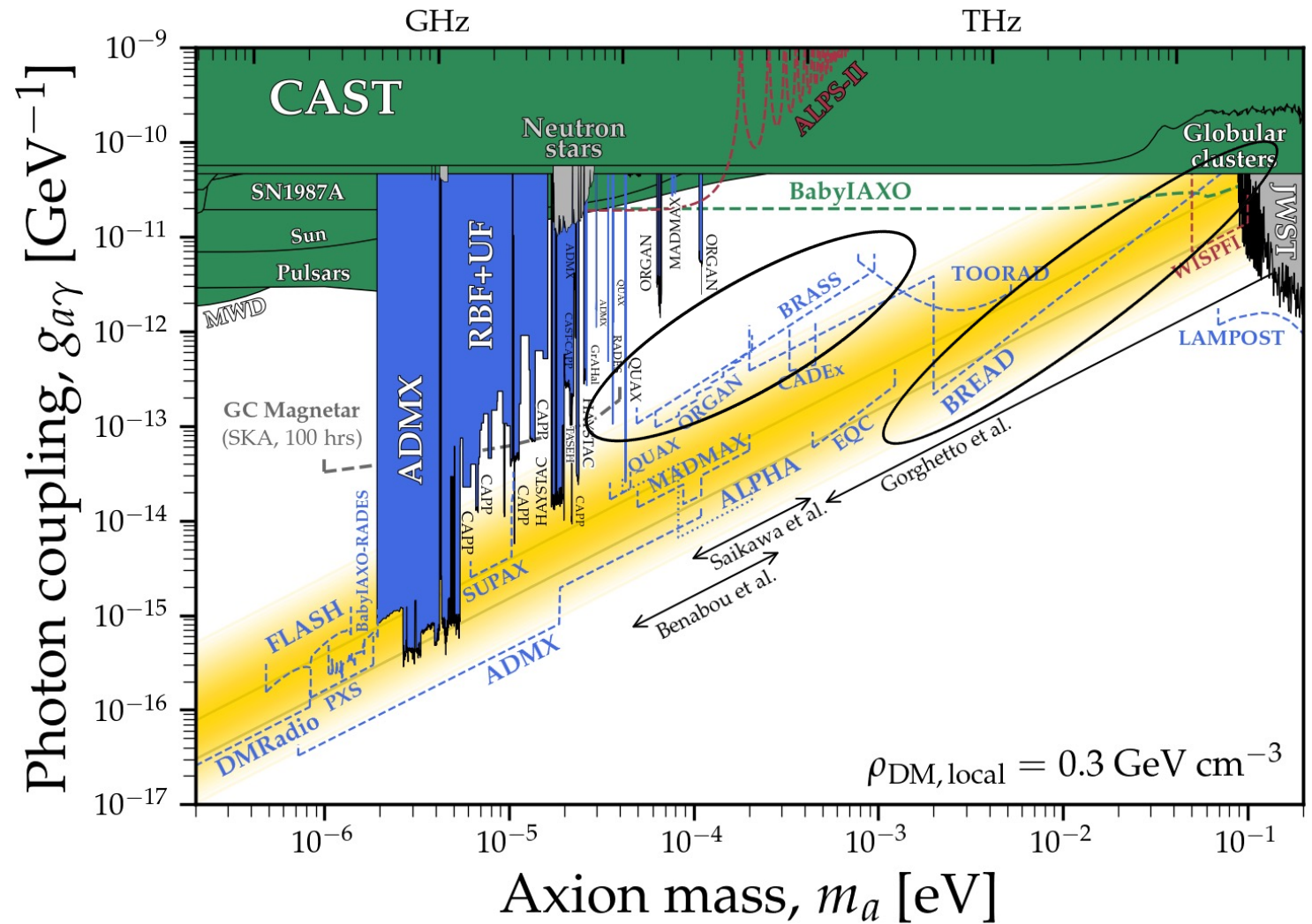


# Axion Dark Matter Experiments

## Variety of experimental techniques

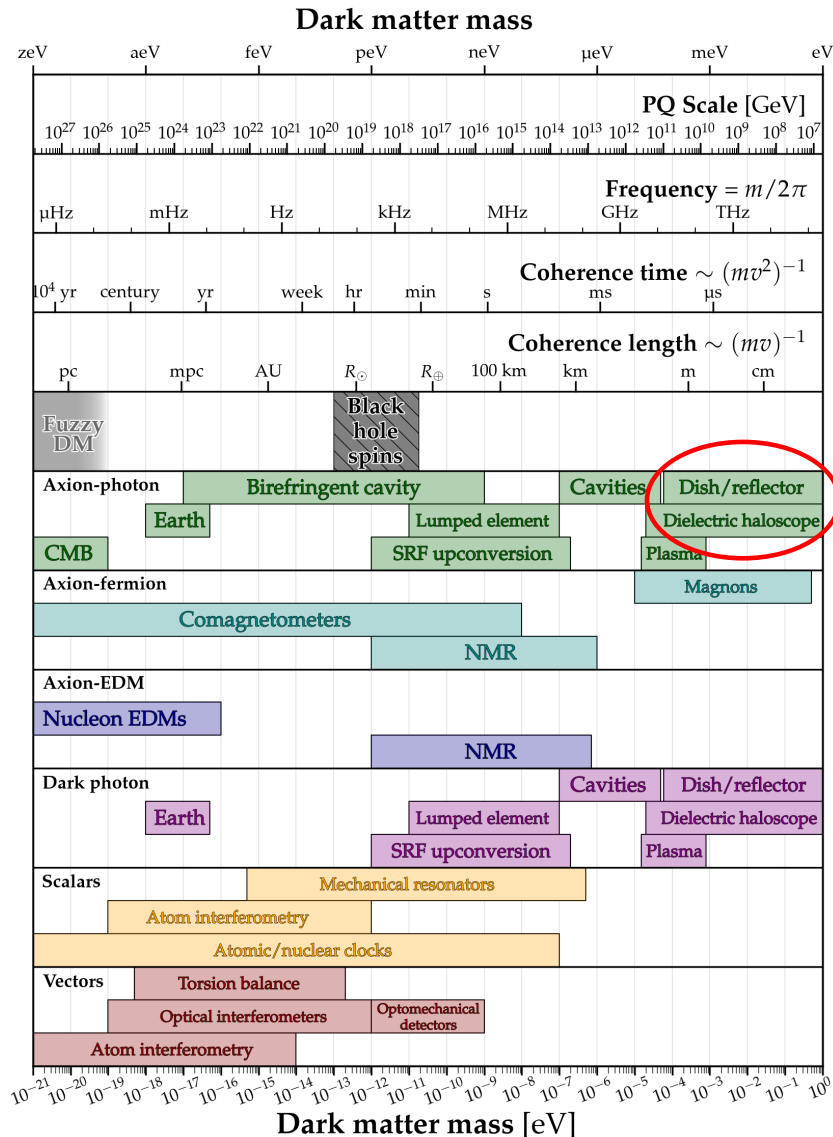


## Dish/reflector prospects:



# Axion Dark Matter Experiments

## Variety of experimental techniques

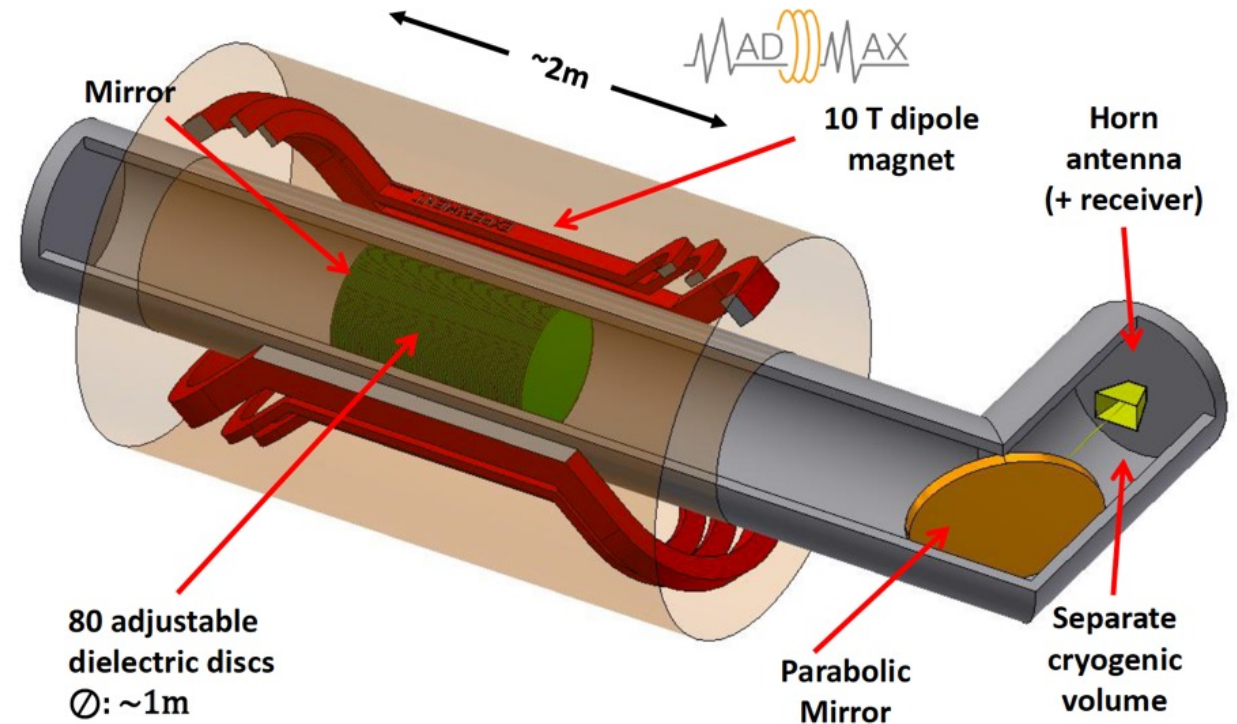


## Dielectric haloscope example:

### MADMAX

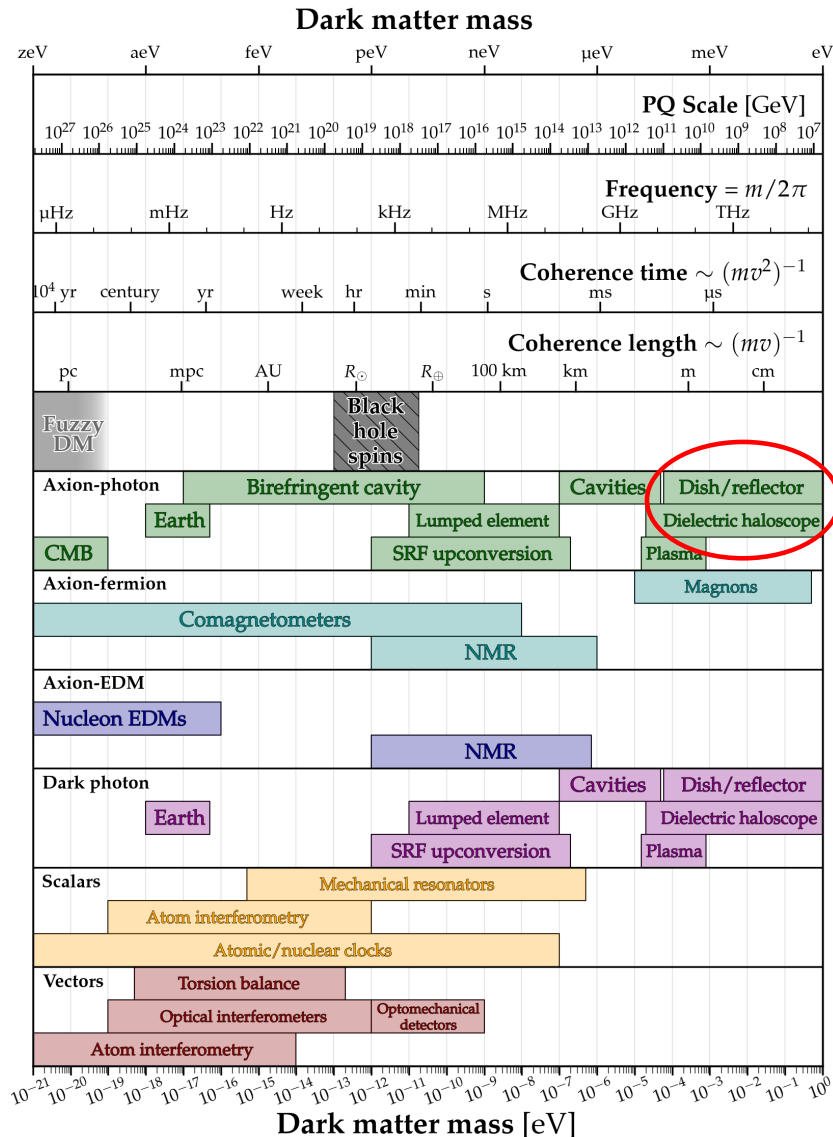
[MADMAX Coll., Eur.Phys.J.C 79 (2019) 3, 186]

- 10 T magnet
- Large number of adjustable dielectric disks



# Axion Dark Matter Experiments

## Variety of experimental techniques



Dielectric haloscope example:

**MADMAX**

[MADMAX Coll., Eur.Phys.J.C 79 (2019) 3, 186]

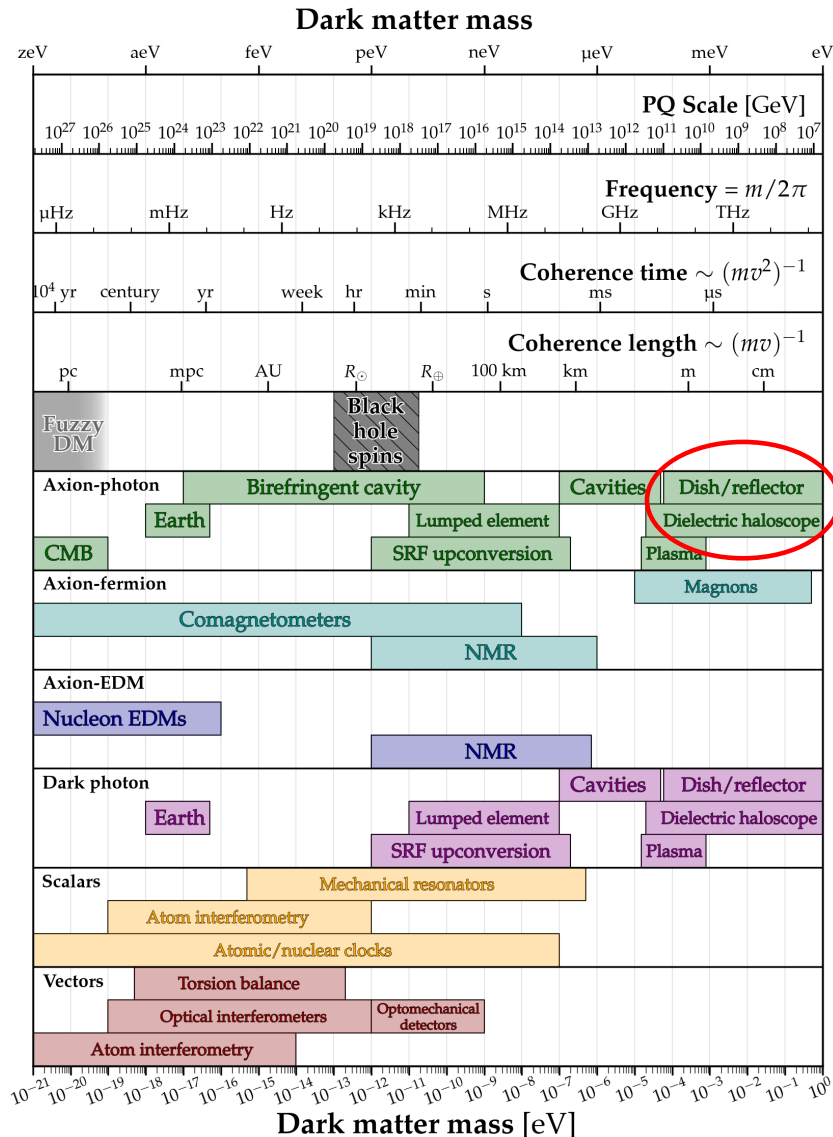
Prototype tests and science runs with MORPURGO magnet @ CERN



$\sim 1.5$  T

# Axion Dark Matter Experiments

## Variety of experimental techniques



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[MADMAX Coll., Eur.Phys.J.C 79 (2019) 3, 186]

Prototype tests and science runs with MORPURGO magnet @ CERN

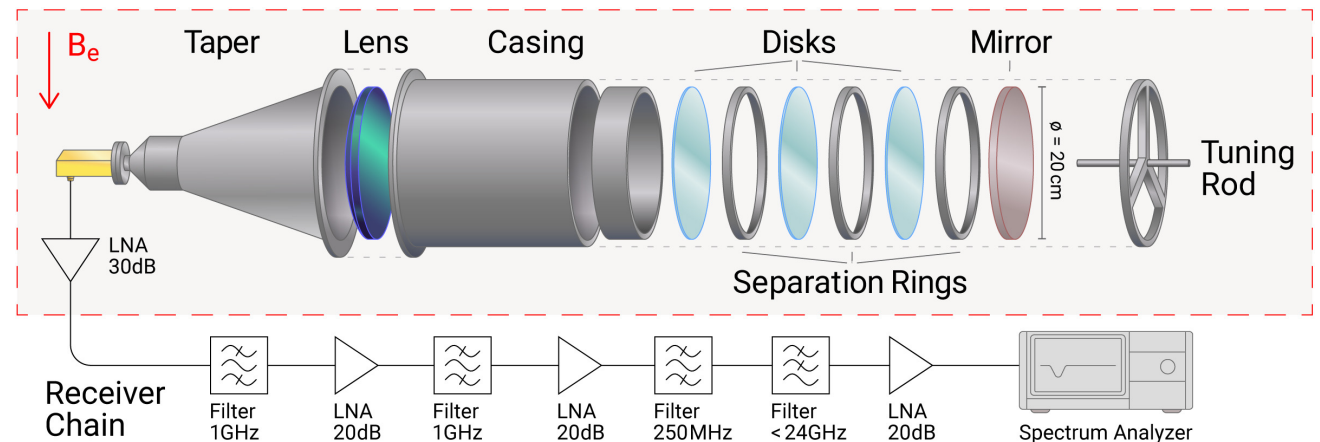
First search published:

PHYSICAL REVIEW LETTERS 135, 041001 (2025)

### First Search for Axion Dark Matter with a MADMAX Prototype

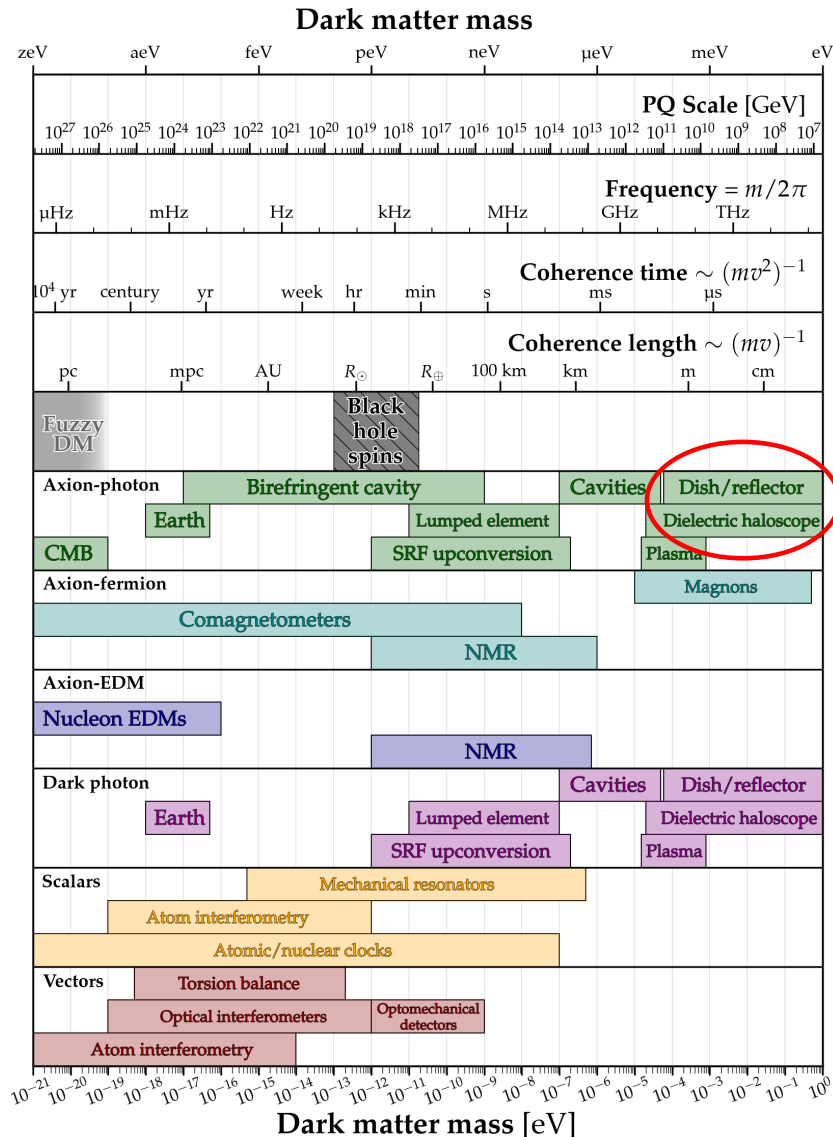
B. Ary dos Santos Garcia,<sup>1</sup> D. Bergermann,<sup>1</sup> A. Caldwell,<sup>2</sup> V. Dabhi,<sup>3</sup> C. Diaconu,<sup>3</sup> J. Diehl,<sup>2</sup> G. Dvali,<sup>2</sup> J. Egge,<sup>4</sup> E. Garutti,<sup>4</sup> S. Heyminck,<sup>5</sup> F. Hubaut,<sup>3</sup> A. Ivanov,<sup>2</sup> J. Jochum,<sup>6</sup> S. Knirck,<sup>7</sup> M. Kramer,<sup>5</sup> D. Kreikemeyer-Lorenzo,<sup>2</sup> C. Krieger,<sup>4</sup> C. Lee,<sup>2,\*</sup> D. Leppla-Weber,<sup>8</sup> X. Li,<sup>2,†</sup> A. Lindner,<sup>8</sup> B. Majorovits,<sup>2</sup> J. P. A. Maldonado,<sup>2</sup> A. Martini,<sup>8</sup> A. Miyazaki,<sup>9</sup> E. Öz,<sup>1</sup> P. Pralavorio,<sup>3</sup> G. Raffelt,<sup>2</sup> J. Redondo,<sup>10</sup> A. Ringwald,<sup>8</sup> J. Schaffran,<sup>8</sup> A. Schmidt,<sup>1</sup> F. Steffen,<sup>2</sup> C. Strandhagen,<sup>6</sup> I. Usherov,<sup>6</sup> H. Wang,<sup>1</sup> and G. Wieching<sup>5</sup>

(MADMAX Collaboration)



# Axion Dark Matter Experiments

## Variety of experimental techniques



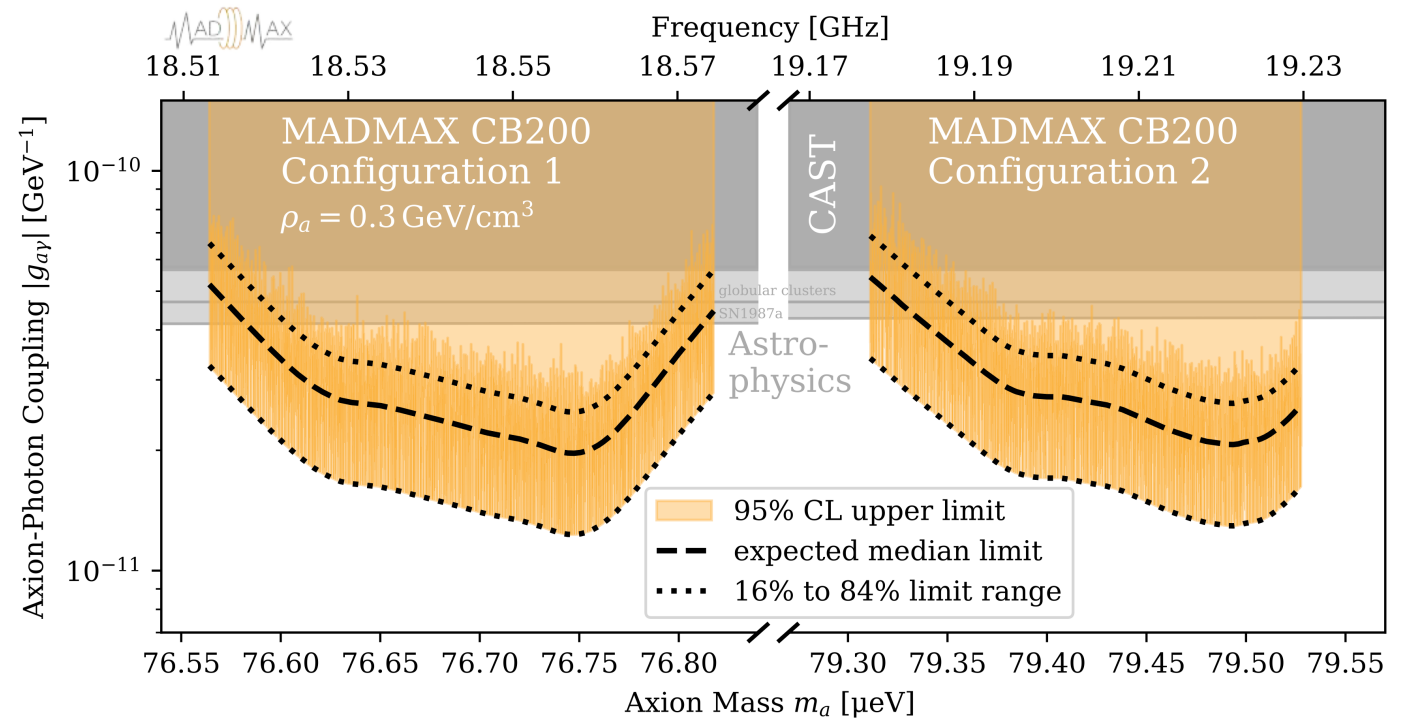
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Prototype tests and science runs with MORPURGO magnet @ CERN

First search published:

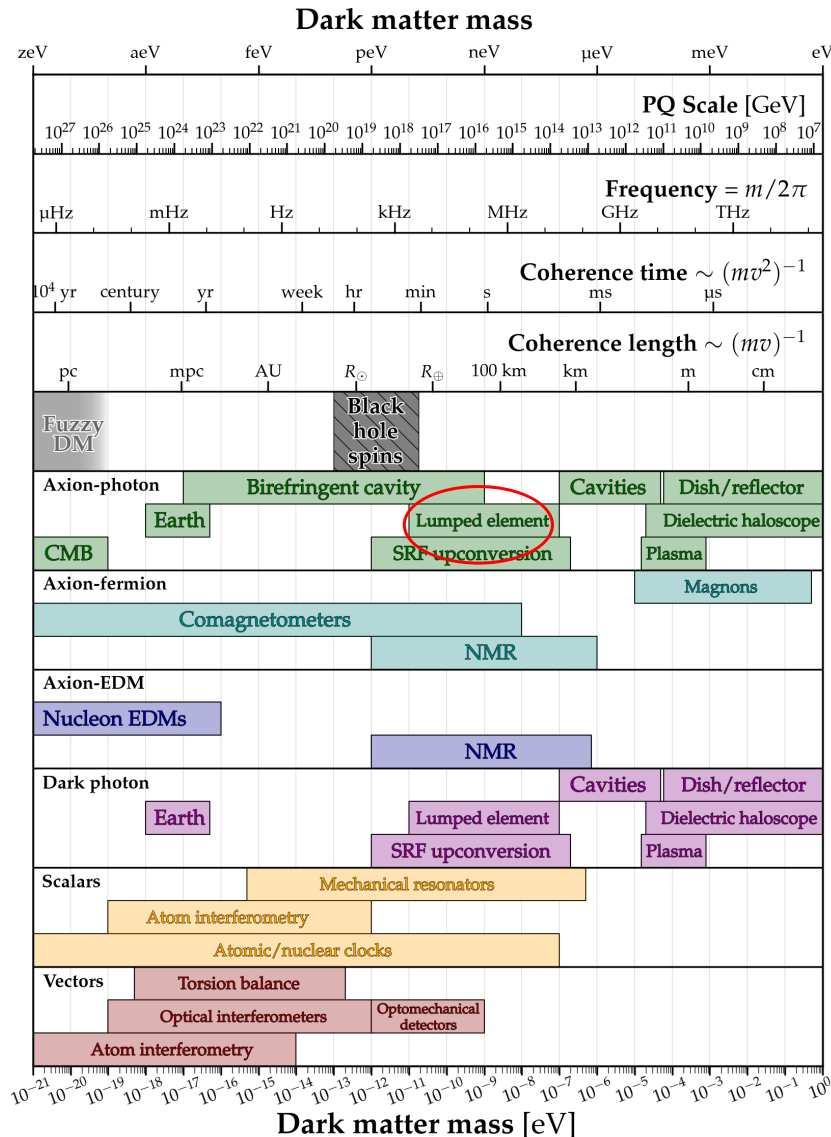


[MADMAX Collaboration, PRL 135, 041001 (2025)]



# Axion Dark Matter Experiments

## Variety of experimental techniques



### Lumped element concept:

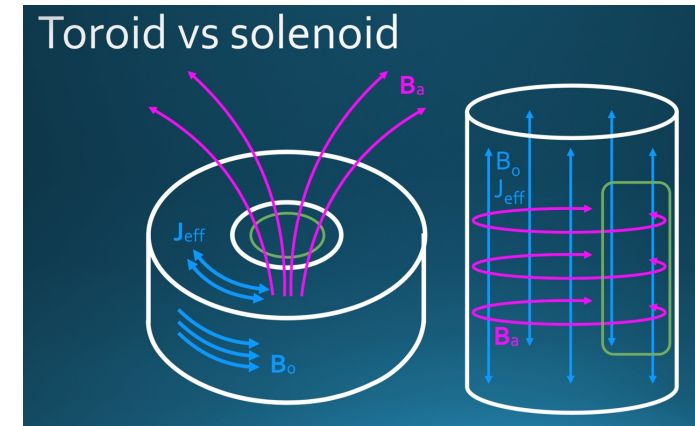
[Sikivie, Sullivan, Tanner 14; Kahn, Safdi, Thaler '16]

- In a static magnetic field  $B_0$ , the oscillating axion DM field acts as an effective current density

$$\mathbf{j}_a = -g_{a\gamma} \mathbf{B}_0 \dot{a}$$

which, in the magneto-quasistatic limit, sources a tiny oscillating magnetic field  $B_a$ , with approximately

$$\nabla \times \mathbf{B}_a = \mathbf{j}_a$$

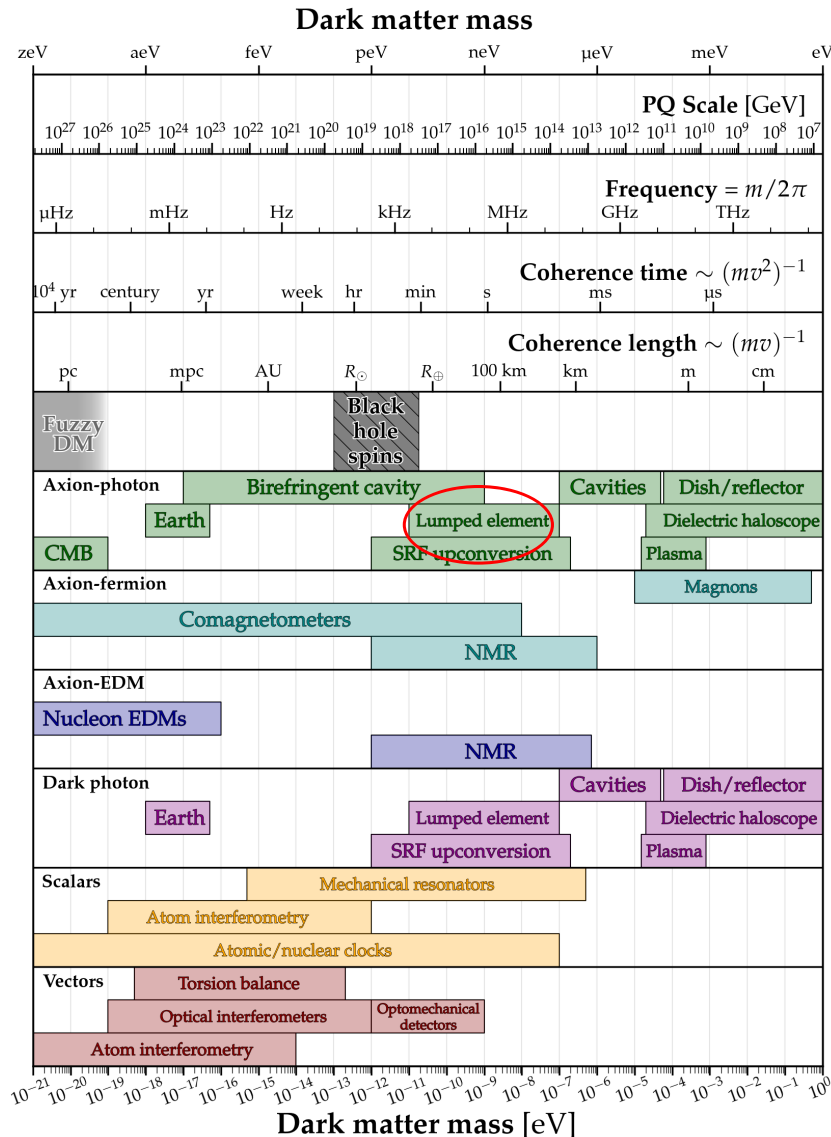


[Salemi '21]

$$\mathbf{J}_{eff} = g_{a\gamma\gamma} \sqrt{2\rho_{DM}} \cos(m_a t) \mathbf{B}$$

# Axion Dark Matter Experiments

## Variety of experimental techniques



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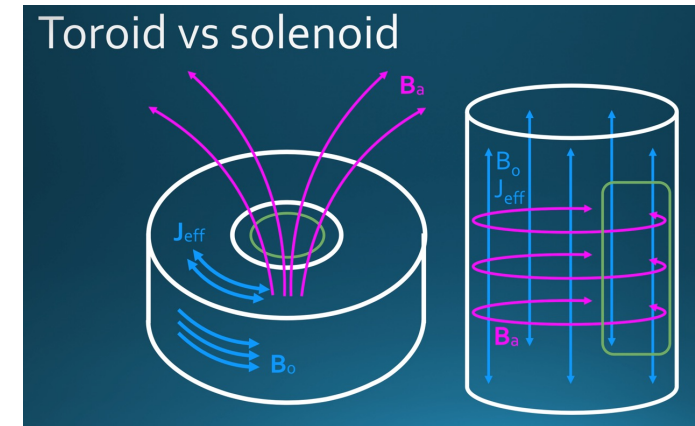
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$$\nabla \times \mathbf{B}_a = \mathbf{j}_a$$

- The induced oscillating flux is coupled into a superconducting pickup loop, amplified/read out by a lumped-element circuit, and detected with a SQUID



[Salemi '21]

$$\mathbf{J}_{eff} = g_{a\gamma} \sqrt{2\rho_{DM} \cos(m_a t)} \mathbf{B}$$

# Axion Dark Matter Experiments

## Variety of experimental techniques

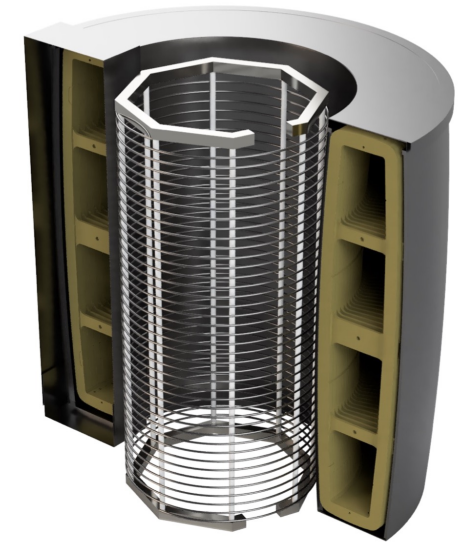
## Lumped element examples:

### DMRadio

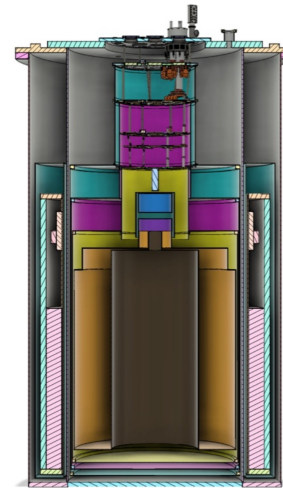
Talk by Andrew Yi on Tue 23/06 at 14:00 in Particles: Dark Matter parallel session

- Ambitious suite of experiments in the US

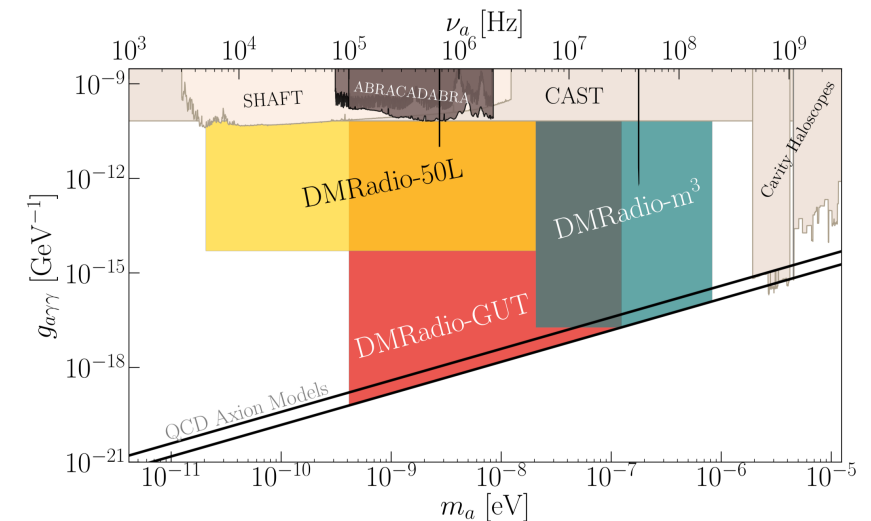
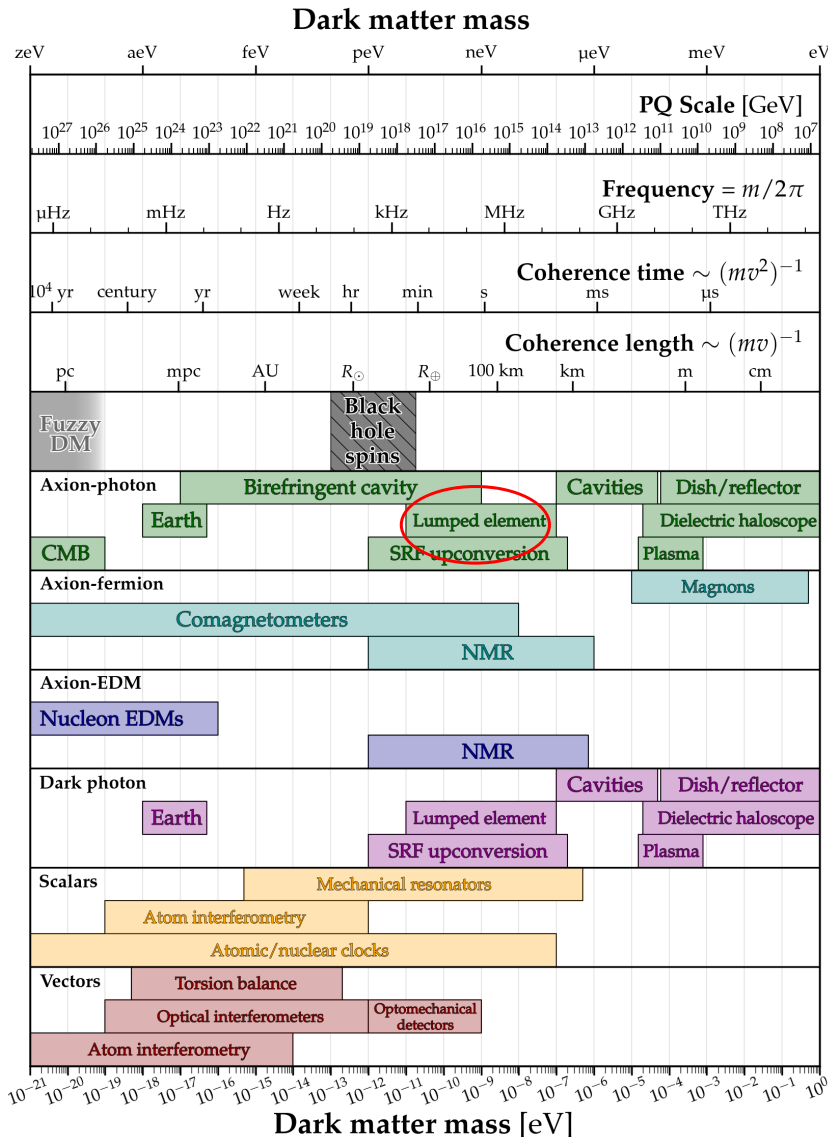
- DMRadio-50L
- DMRadio-m<sup>3</sup>
- DMRadio-GUT



50L

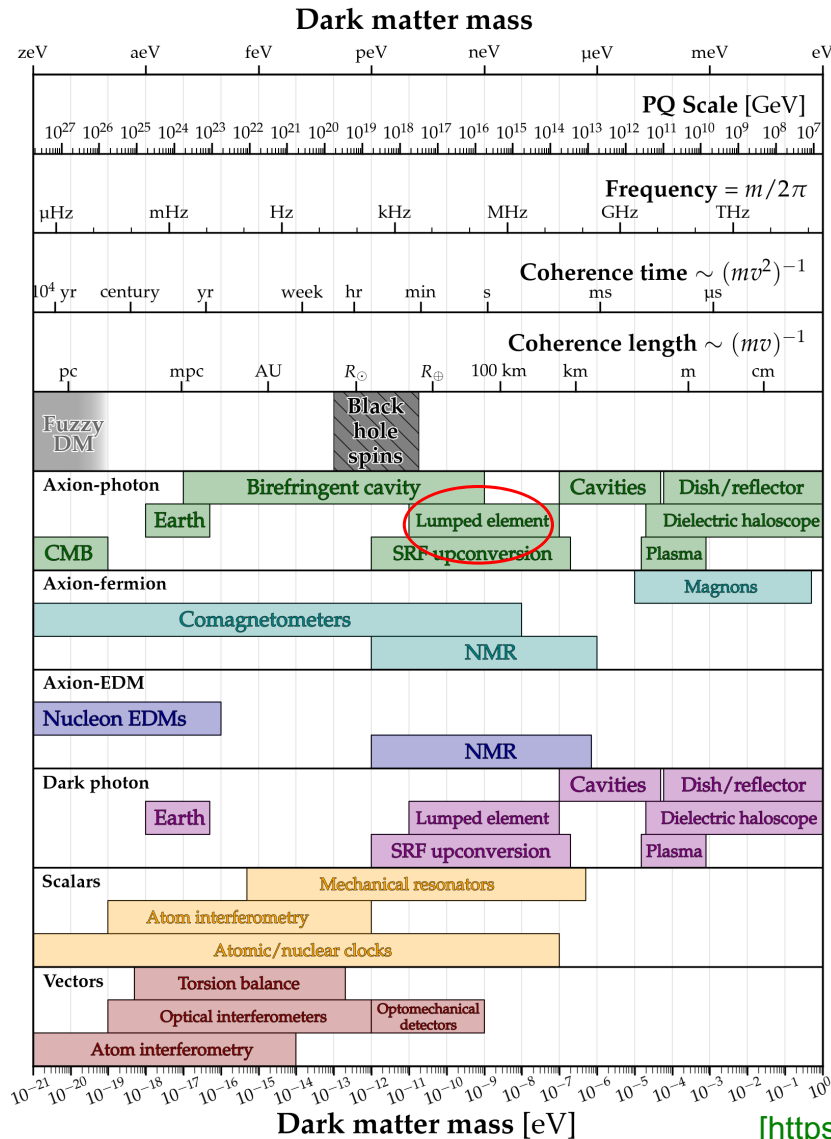


m<sup>3</sup>

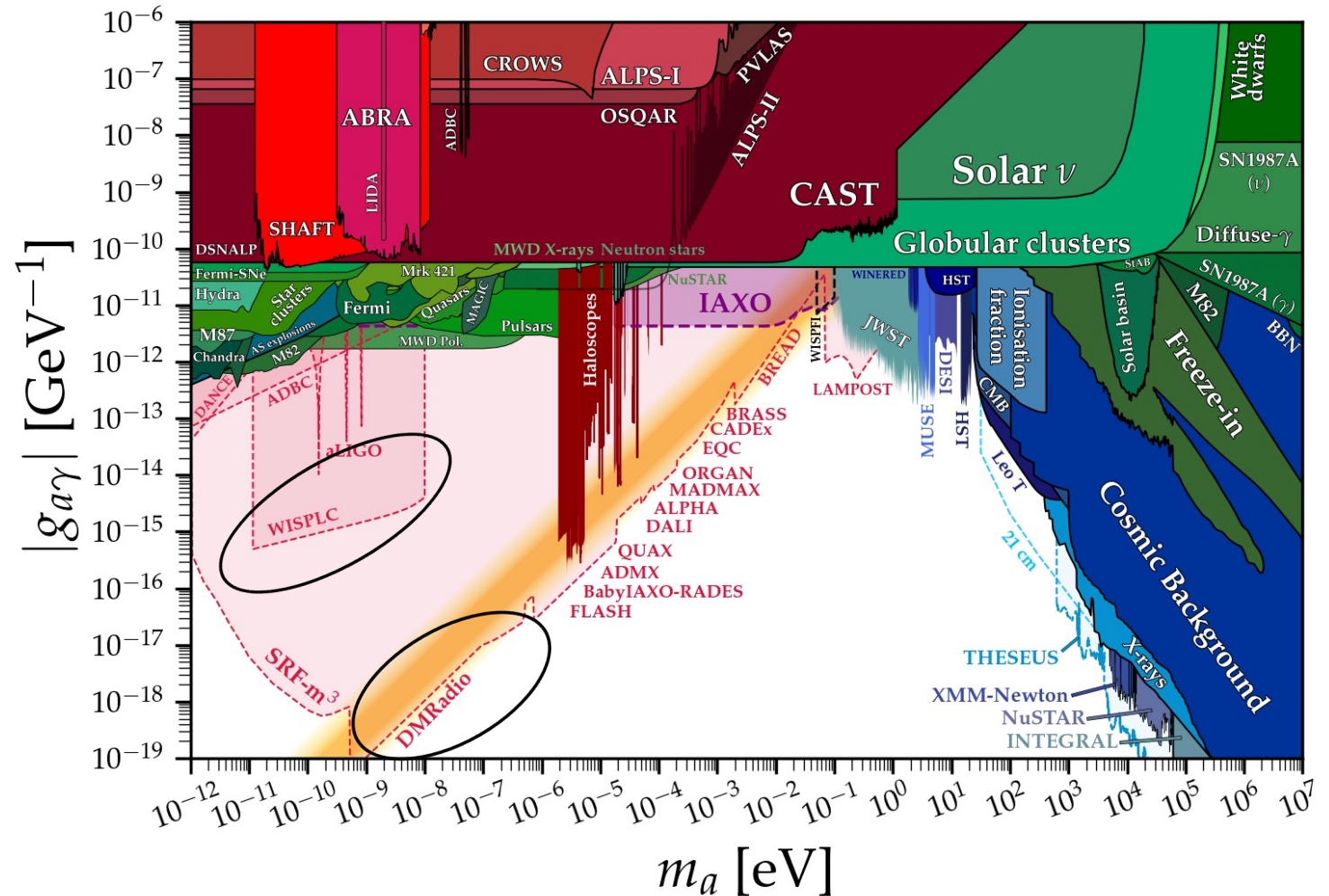


# Axion Dark Matter Experiments

## Variety of experimental techniques



## Lumped element prospects:



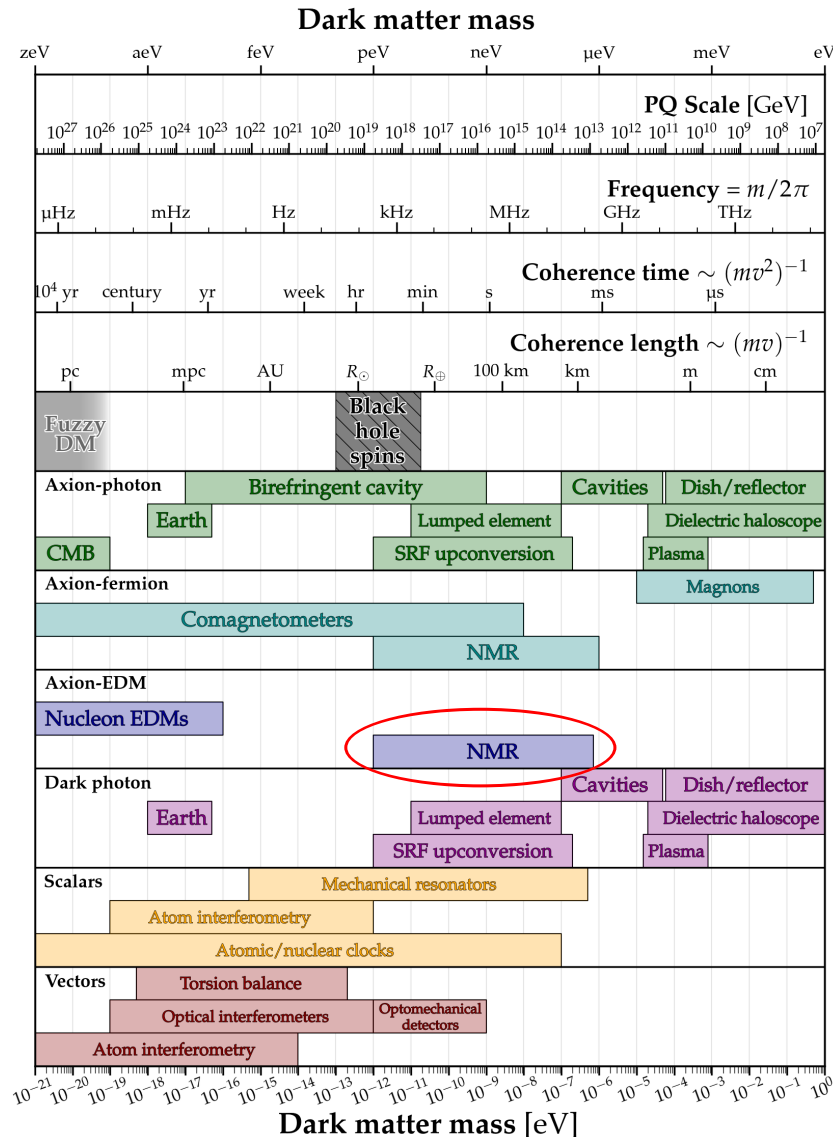
[[https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots\\_png/AxionPhoton\\_Closeup\\_AltColours.png](https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots_png/AxionPhoton_Closeup_AltColours.png)]

# Axion Dark Matter Experiments

## Variety of experimental techniques

### NMR for EDM coupling concept:

[Graham, Rajendran 13; Budker et al. 14]



- Axion DM field acts as an oscillating effective NEDM:

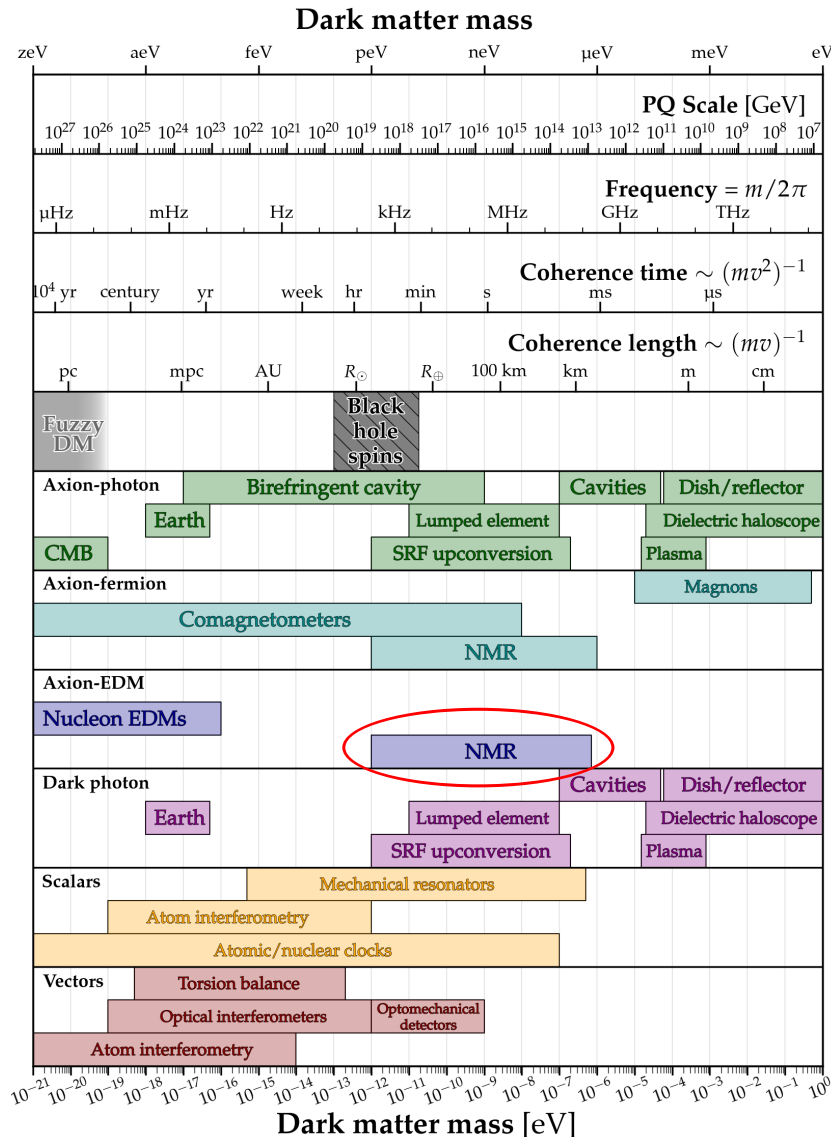
$$d_N(t) = g_{aN\gamma} \sqrt{2\rho_a} \cos(m_a t) / m_a$$

# Axion Dark Matter Experiments

## Variety of experimental techniques

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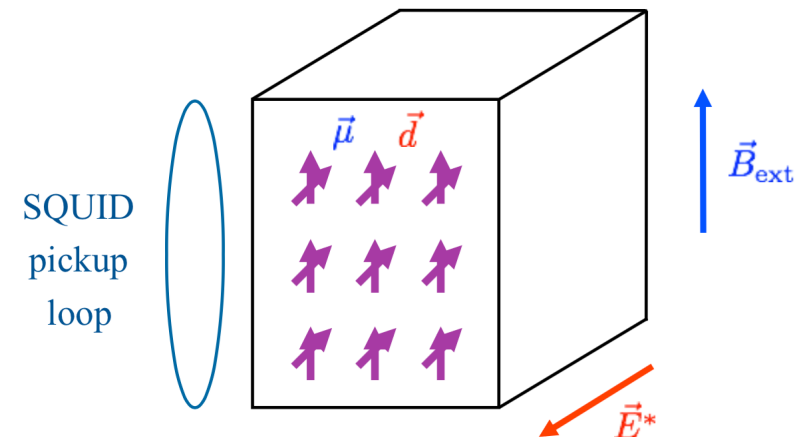
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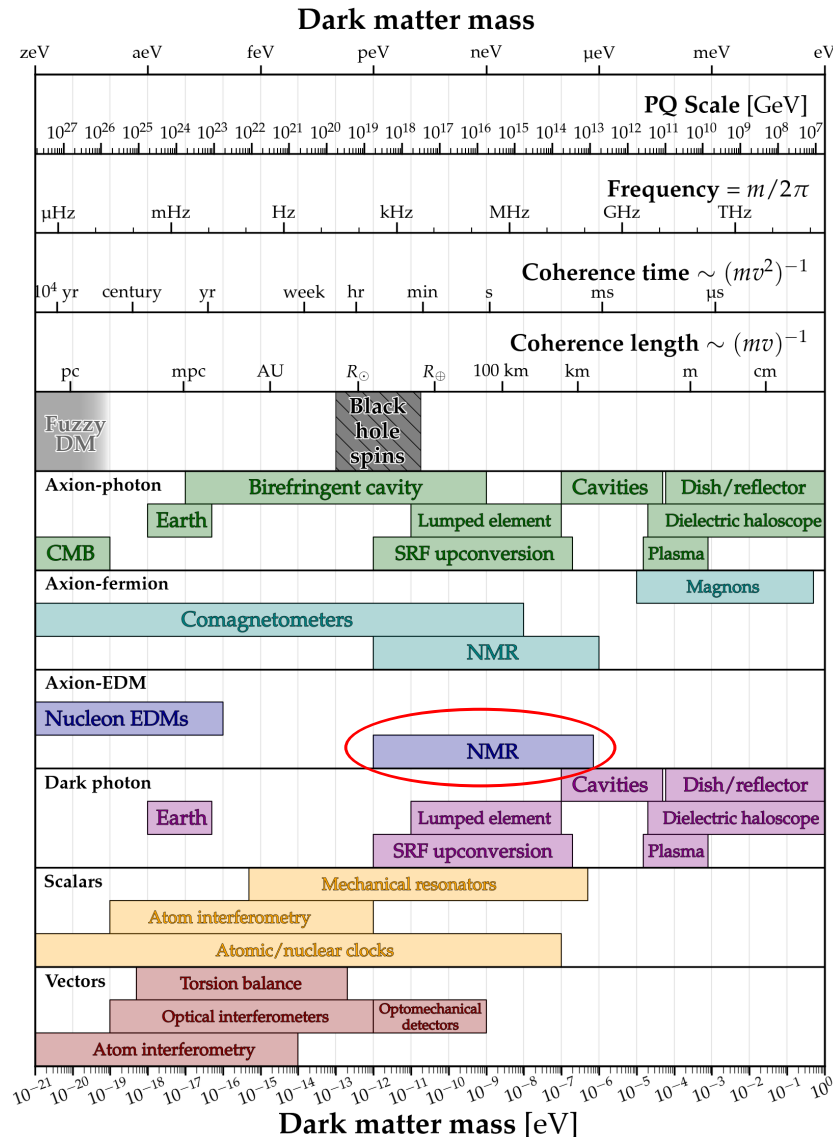
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# Axion Dark Matter Experiments

## Variety of experimental techniques



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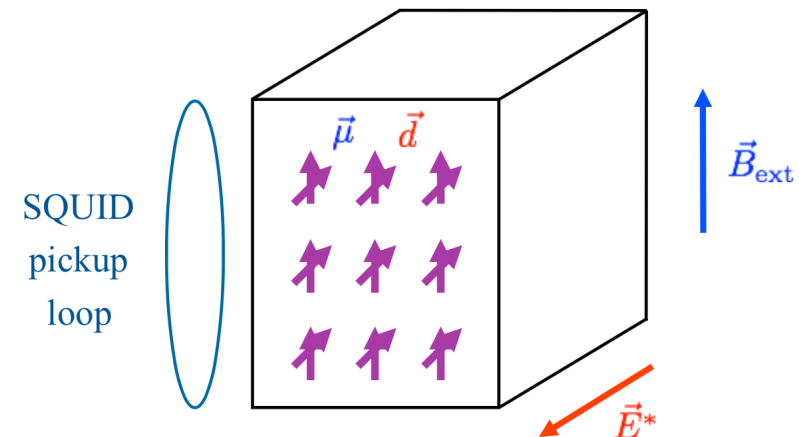
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- Nuclear spins are polarised along  $\vec{B}_{\text{ext}}$ , and precess at Larmor frequency

$$\omega_L = \gamma_N B_{\text{ext}}$$

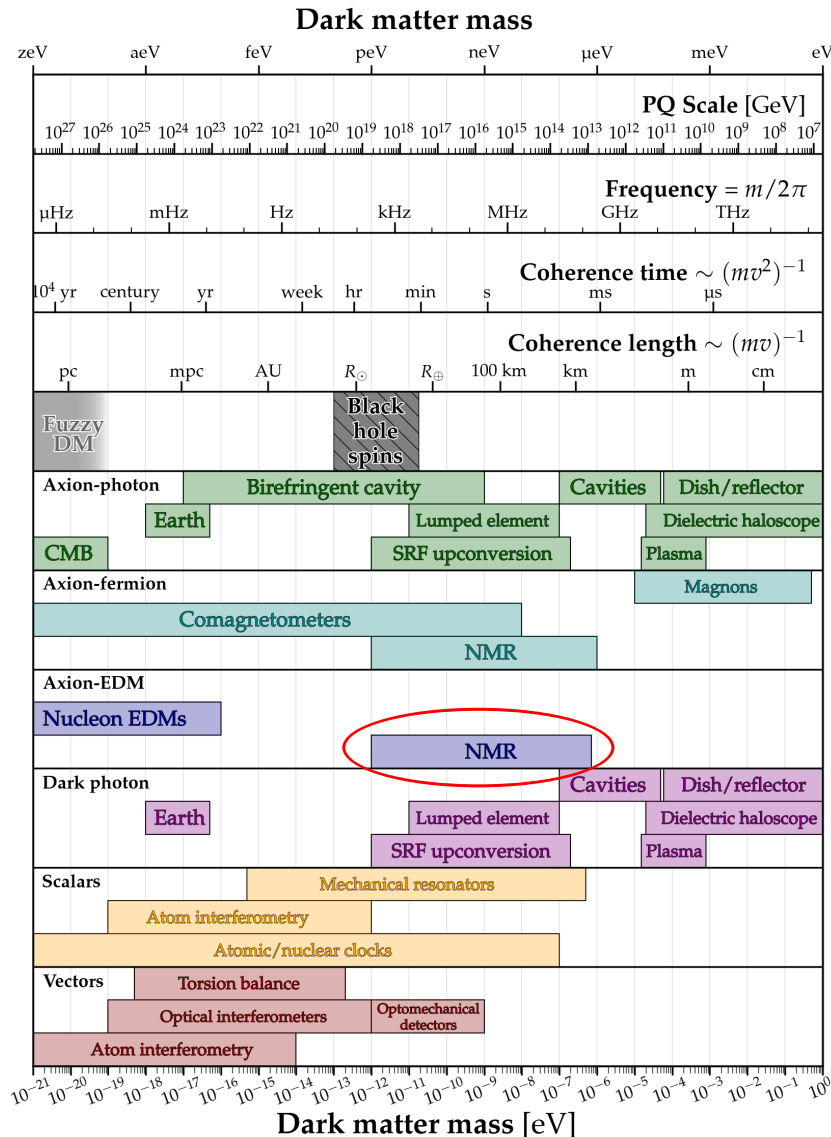


# Axion Dark Matter Experiments

## Variety of experimental techniques

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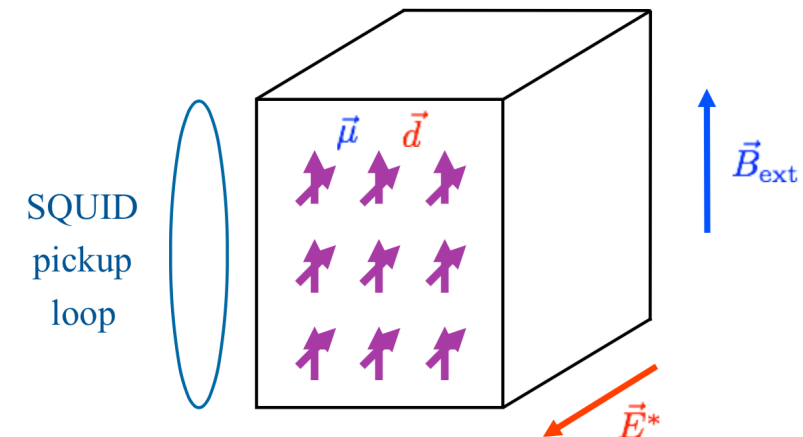
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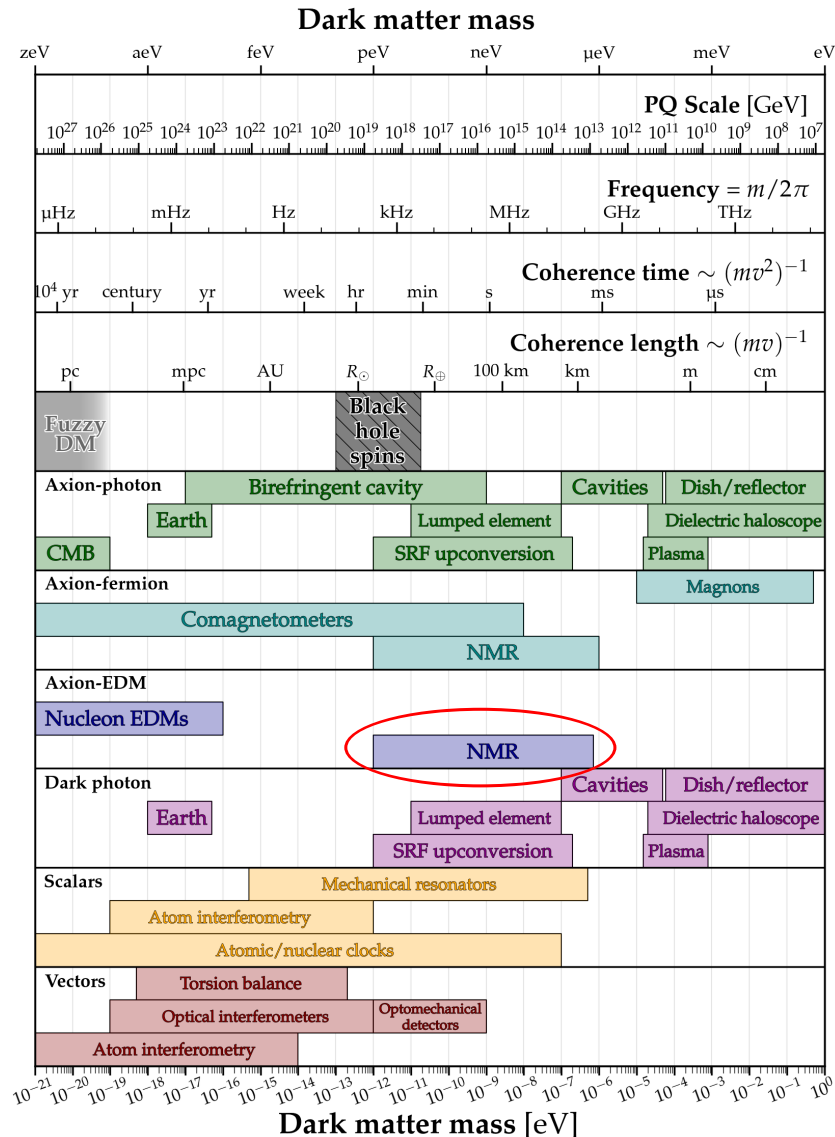
$$\omega_L = \gamma_N B_{\text{ext}}$$

- EDM interaction,  $\epsilon_S \vec{d}_N(t) \cdot \vec{E}^*$ , drives the spins and resonantly builds up transverse magnetization when  $\omega_L \simeq m_a$



# Axion Dark Matter Experiments

## Variety of experimental techniques



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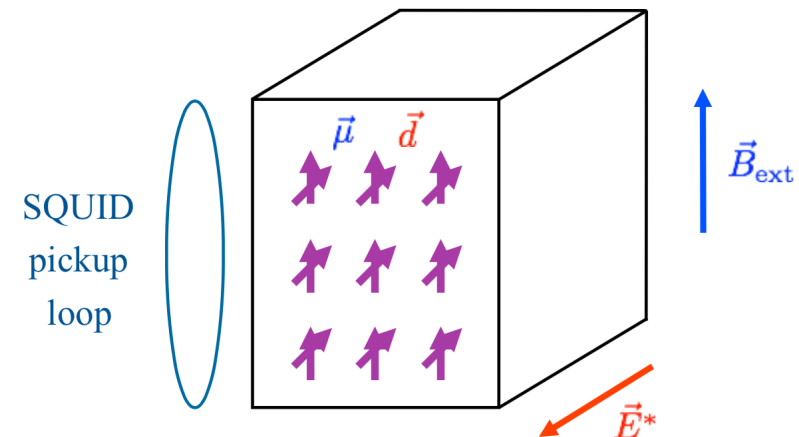
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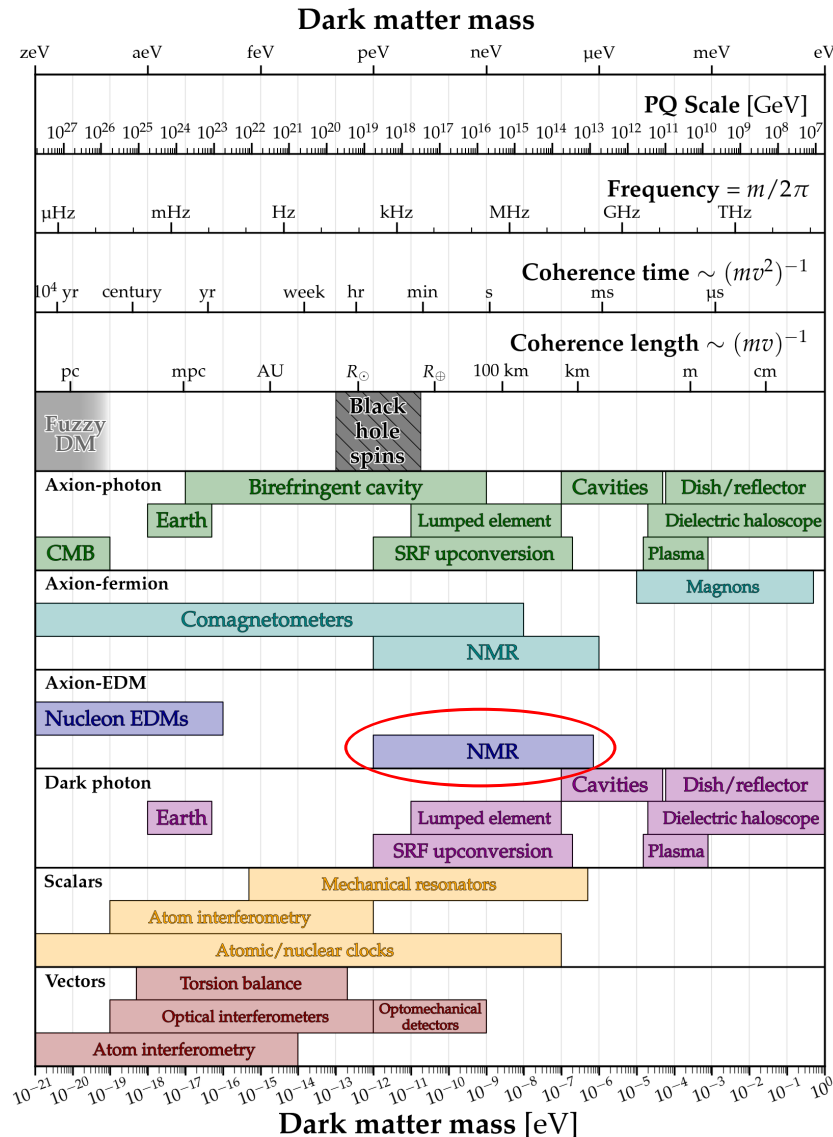
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- Sweeping  $\vec{B}_{\text{ext}}$  scans the axion mass; the transverse magnetization is read out with a pickup loop/SQUID



# Axion Dark Matter Experiments

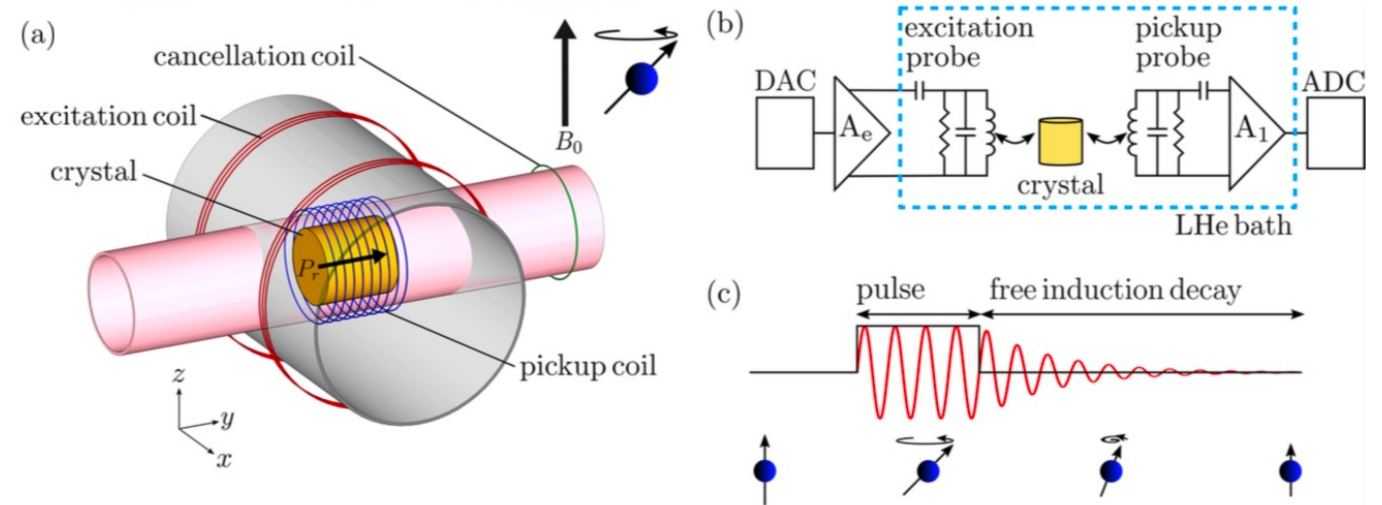
## Variety of experimental techniques



## NMR for EDM coupling example:

### CASPER-Electric in Boston

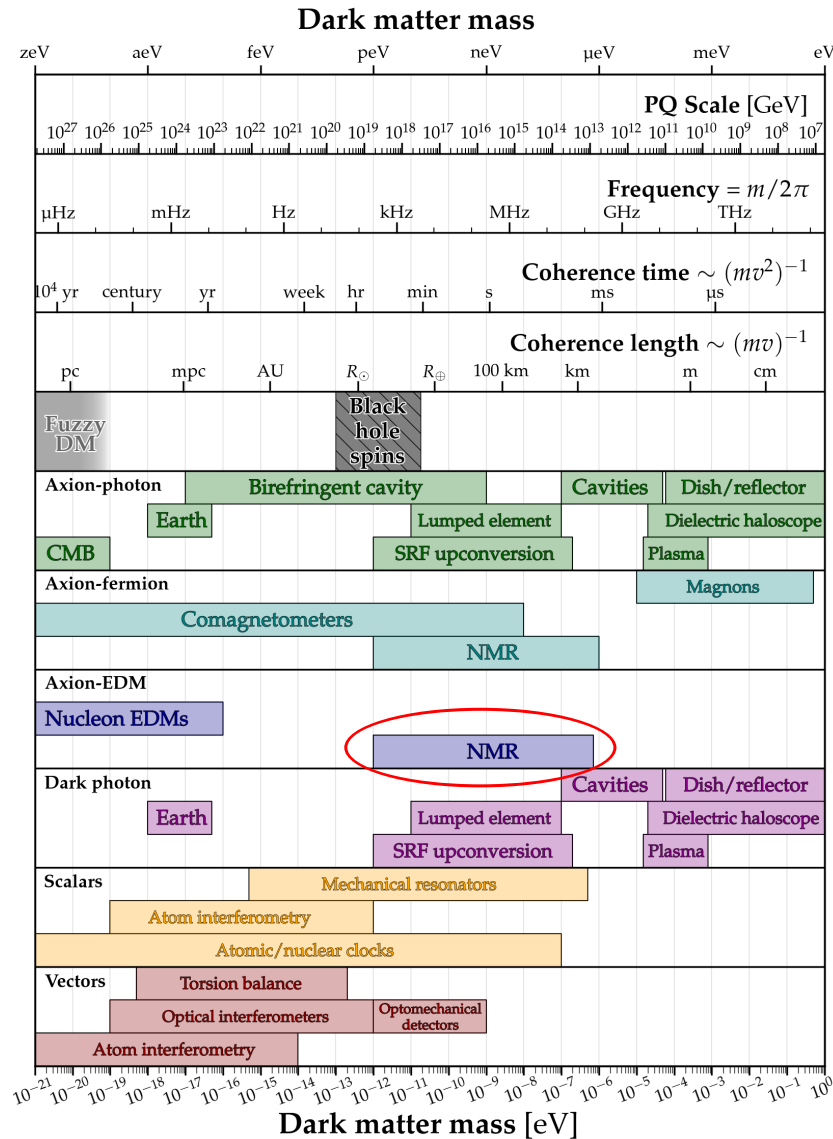
- Initial pathfinder experiment in 2021



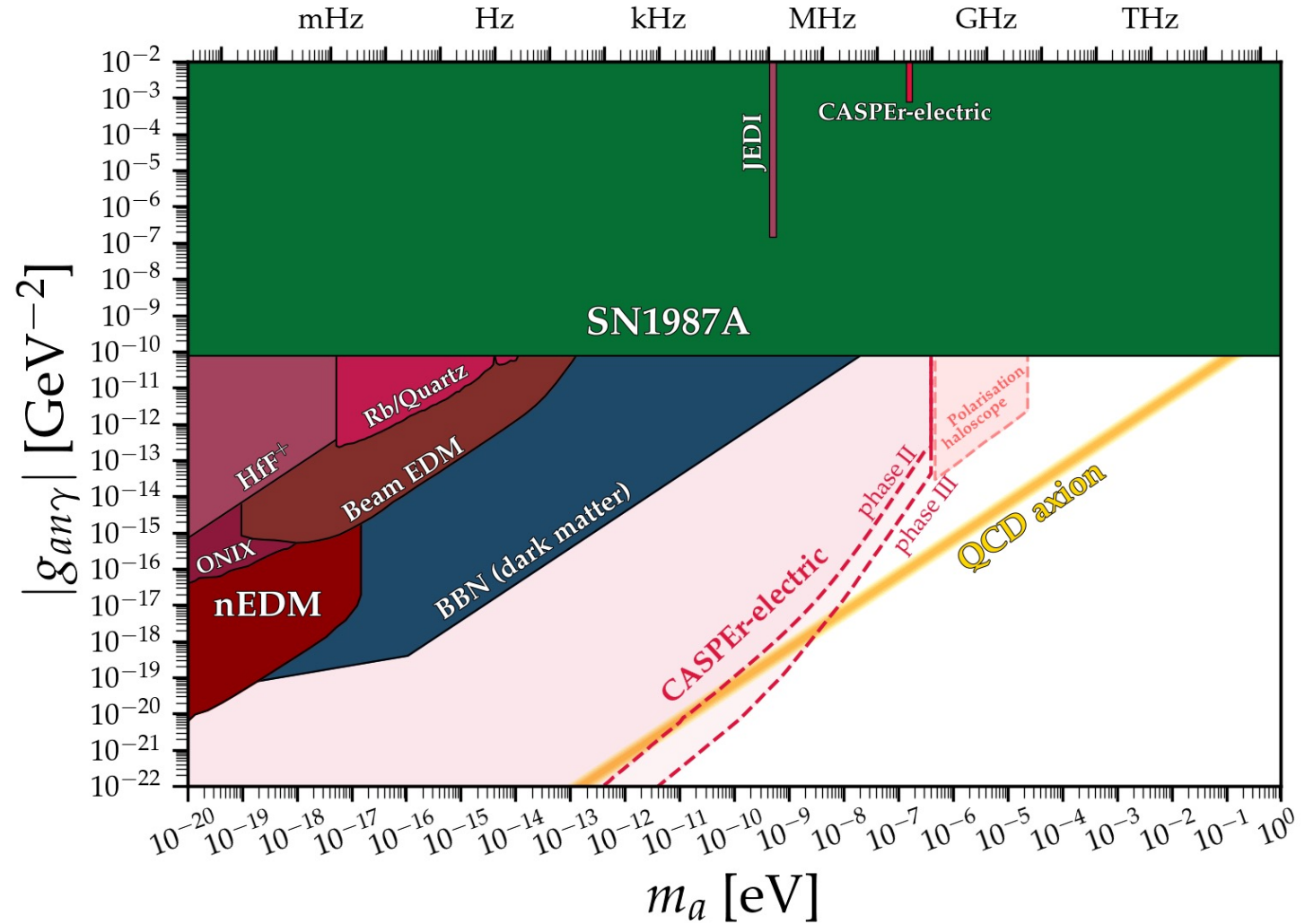
[Aybas et al., 2101.01241]

# Axion Dark Matter Experiments

## Variety of experimental techniques



## NMR for EDM coupling prospects:



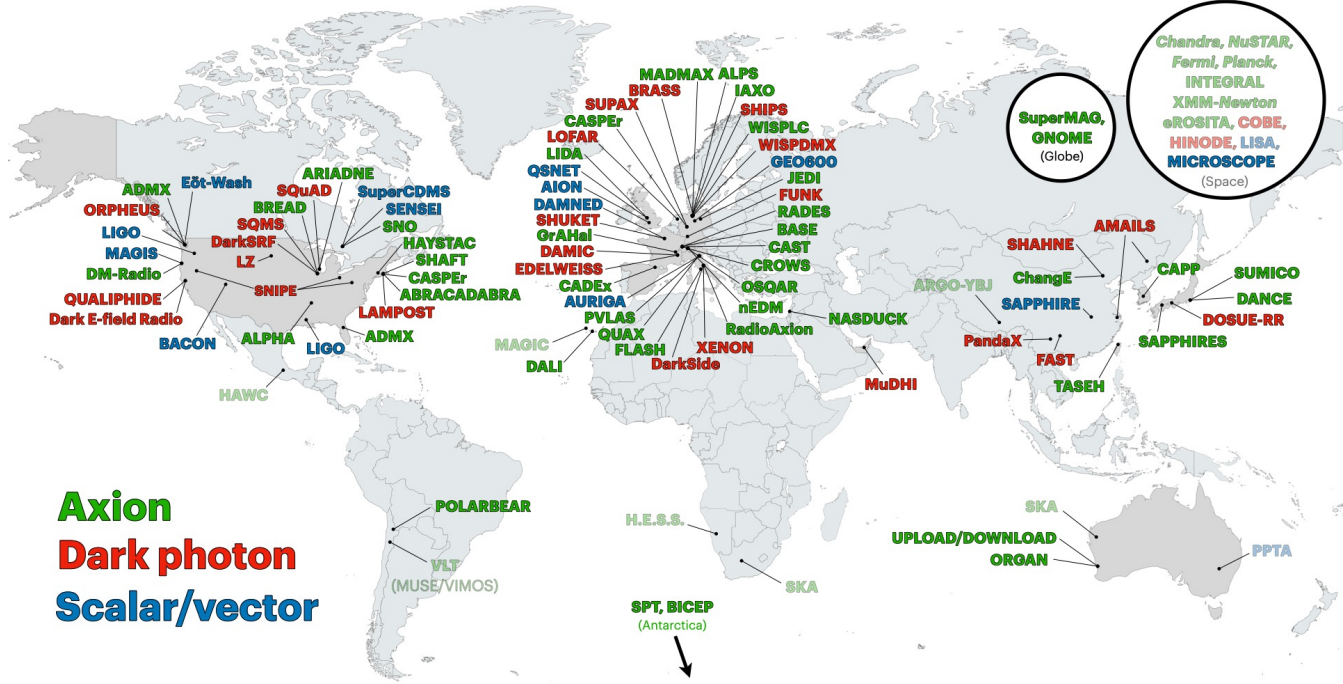
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# Conclusions

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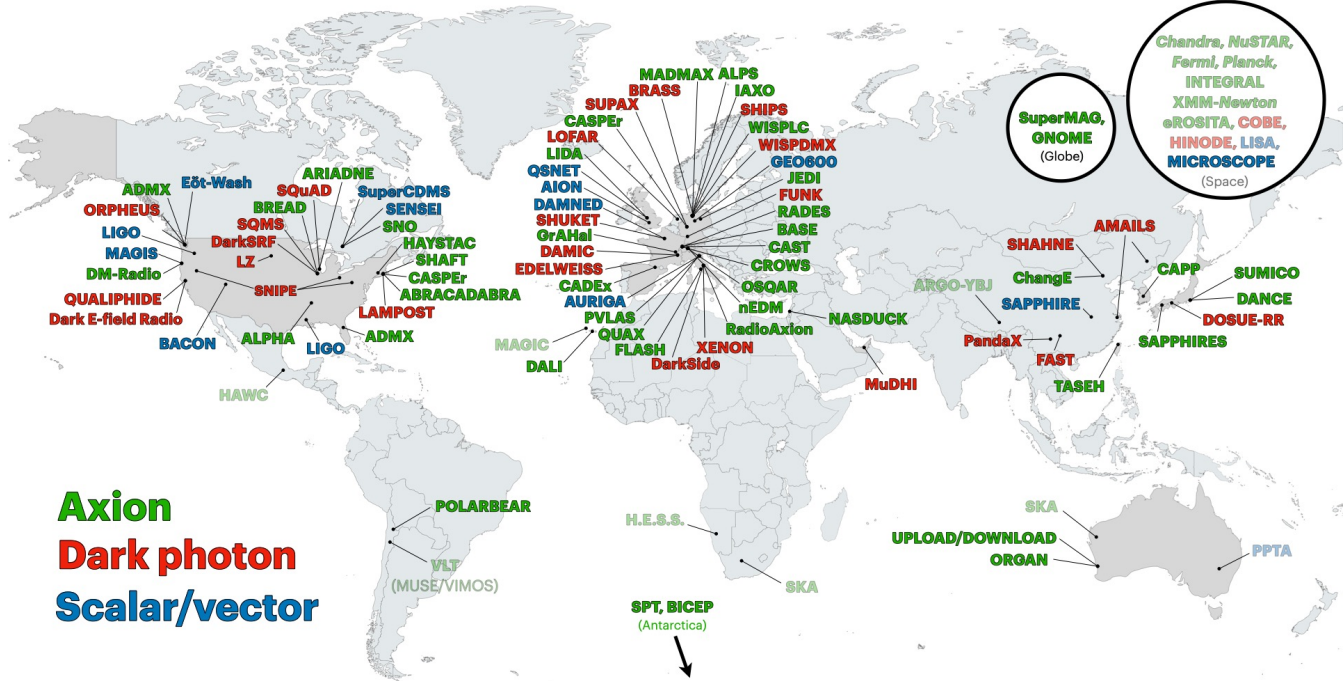
- Axion remains one of the theoretically best motivated dark matter candidates
- World-wide vigorous experimental activity on axion searches exploiting different techniques and couplings:



[[https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/Wavy\\_Map.png](https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/Wavy_Map.png)]

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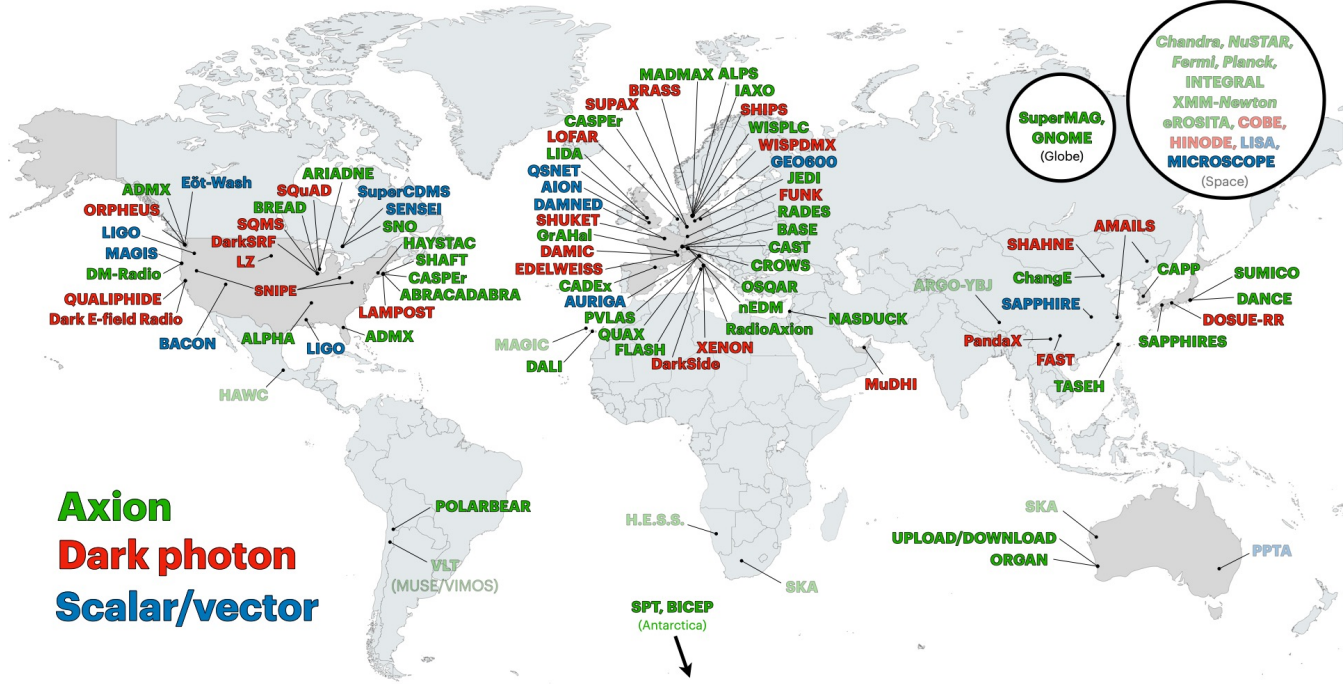


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# Conclusions

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**Stay tuned!**

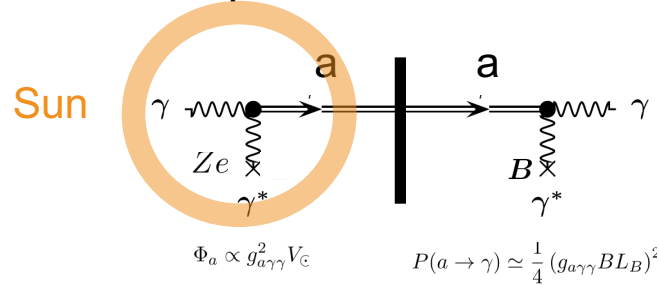
# Backup Slides

# Axion Experiments Not Relying on Axion Dark Matter

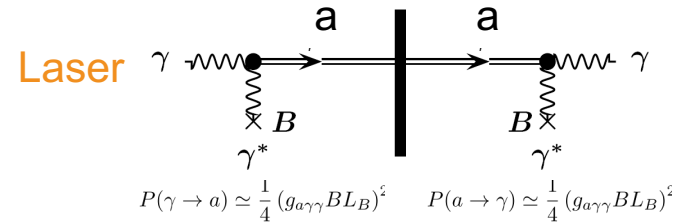
## Searching for solar or home-made axions

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- Helioscopes: [Sikivie 83]



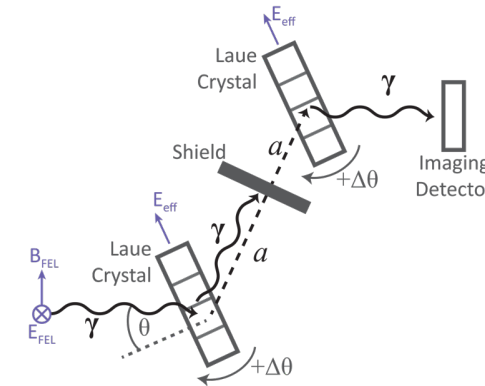
- Light-shining through walls: [Anselm 85; van Bibber 87]



## Crystal-assisted X-ray LSW concept:

[Buchmüller & Hoogeveen, Phys. Lett. B 237 (1990) 278]

- Use intense X-ray beam instead of an optical laser
- In a crystal, strong microscopic electric field of the lattice acts as external field for Primakoff conversion:



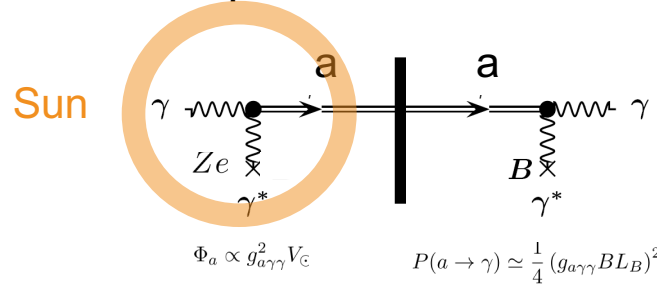
- Coherent conversion occurs when the crystal satisfies a Bragg/Laue momentum-matching condition.
- An absorber blocks ordinary X-rays
- Second aligned crystal converts axions back into X-rays:
- Look for regenerated X-rays downstream of the wall.
- Rotating/detuning the crystals scans the ALP mass.

# Axion Experiments Not Relying on Axion Dark Matter

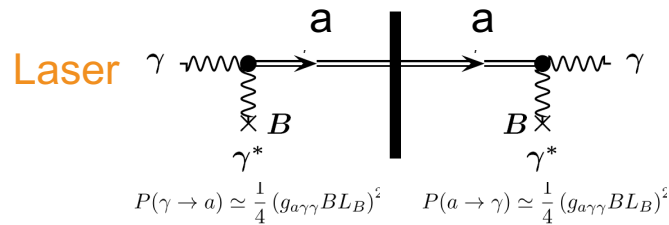
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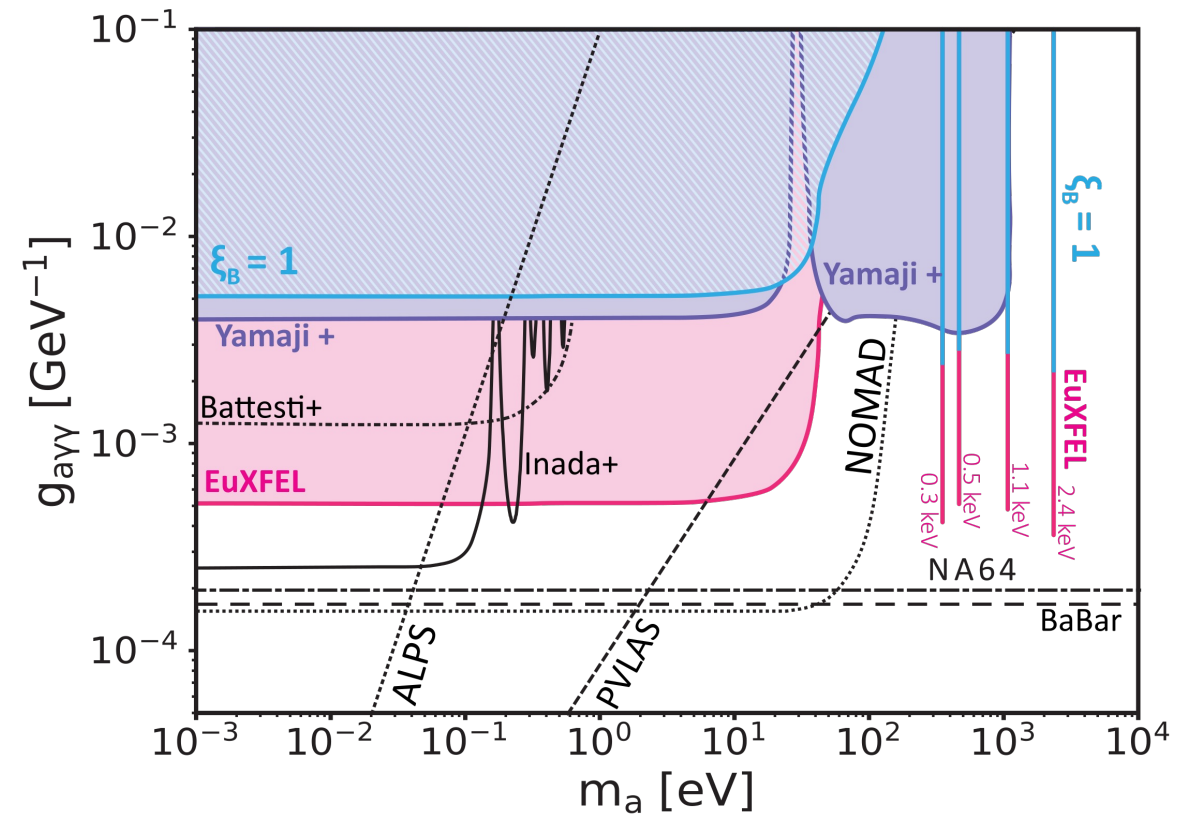


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## Experimental status of crystal-assisted X-ray LSW:

[Yamaji et al., Phys. Lett. B 782 (2018) 523;  
Halliday et al., PRL 134 (2025) 055001;  
Heaton et al., 2603.15808]



[Halliday et al., PRL 134 (2025) 055001]

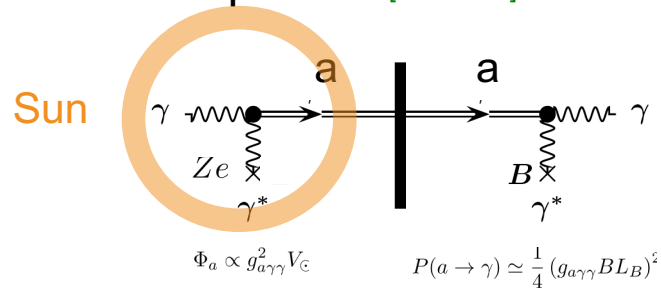
# Axion Experiments Not Relying on Axion Dark Matter

Searching for solar or home-made axions

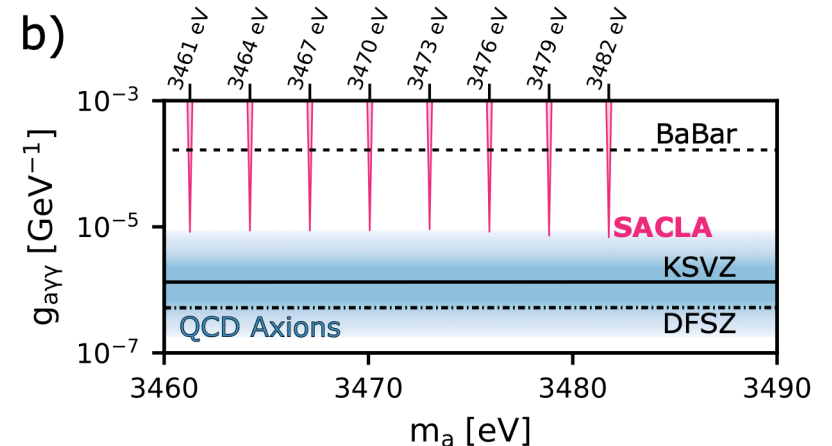
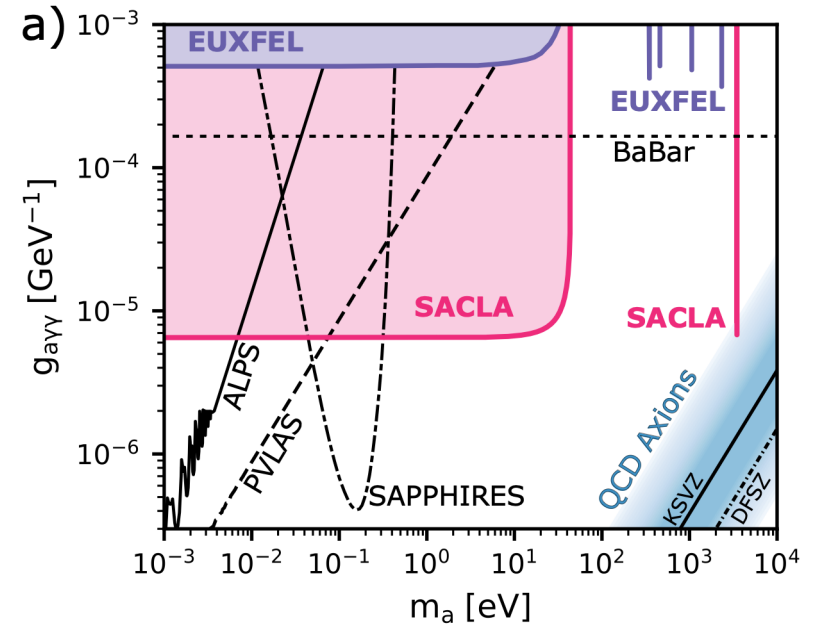
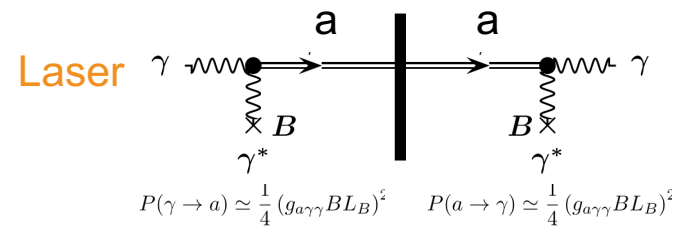
Experimental status of crystal-assisted X-ray LSW:

$$\mathcal{L} \supset \frac{g_{a\gamma\gamma}}{4} a F_{\mu\nu} \tilde{F}^{\mu\nu} \equiv g_{a\gamma\gamma} a \mathbf{E} \cdot \mathbf{B}$$

- Helioscopes: [Sikivie 83]



- Light-shining through walls: [Anselm 85; van Bibber 87]



# Axion Experiments Not Relying on Axion Dark Matter

## Searches for axion mediated forces

- Experiments searching for axion mediated forces particularly effective in meV mass range
- Monopole-dipole interaction between nucleon and fermion:

$$U_{\text{mon-dip}}(r) = \frac{g_{aNN} g_{af\bar{f}}}{8\pi m_f} \left( \frac{m_a}{r} + \frac{1}{r^2} \right) e^{-m_a r} (\hat{\sigma} \cdot \hat{r})$$

$$\mathcal{L}_{\text{int}} = g_{aNN} a \bar{N} N - i g_{af\bar{f}} a \bar{f} \gamma_5 f$$

- Proposed ARIADNE experiment searches for forces between a rotating cylinder, made of unpolarized material, and a vessel containing hyperpolarized  $^3\text{He}$  gas
  - Since  $^3\text{He}$  magnetic moment dominated by neutron contribution: sensitive to monopole-dipole interaction between nucleus and neutrons,  $|g_{aNN} g_{an\bar{n}}|$

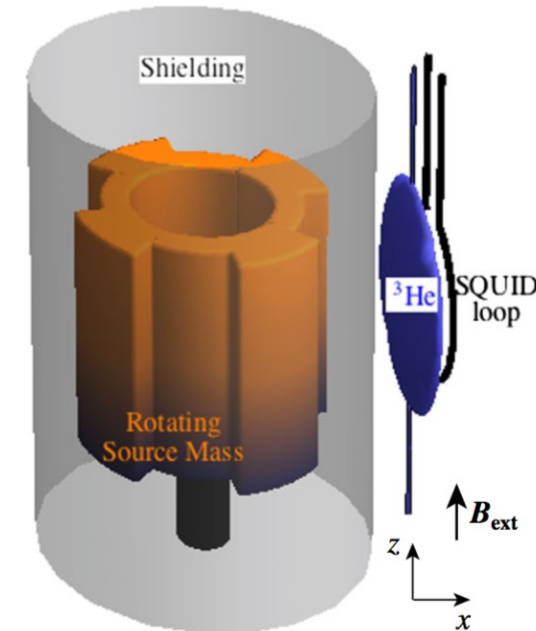


FIG. 1 (color online). A source mass consisting of a segmented cylinder with  $n$  sections is rotated around its axis of symmetry at frequency  $\omega_{\text{rot}}$ , which results in a resonance between the frequency  $\omega = n\omega_{\text{rot}}$  at which the segments pass near the sample and the resonant frequency  $2\vec{\mu}_N \cdot \vec{B}_{\text{ext}}/\hbar$  of the NMR sample. Superconducting cylinders screen the NMR sample from the source mass and (not shown) the setup from the environment.

[Arvanitaki, Geraci 14]

# Axion Experiments Not Relying on Axion Dark Matter

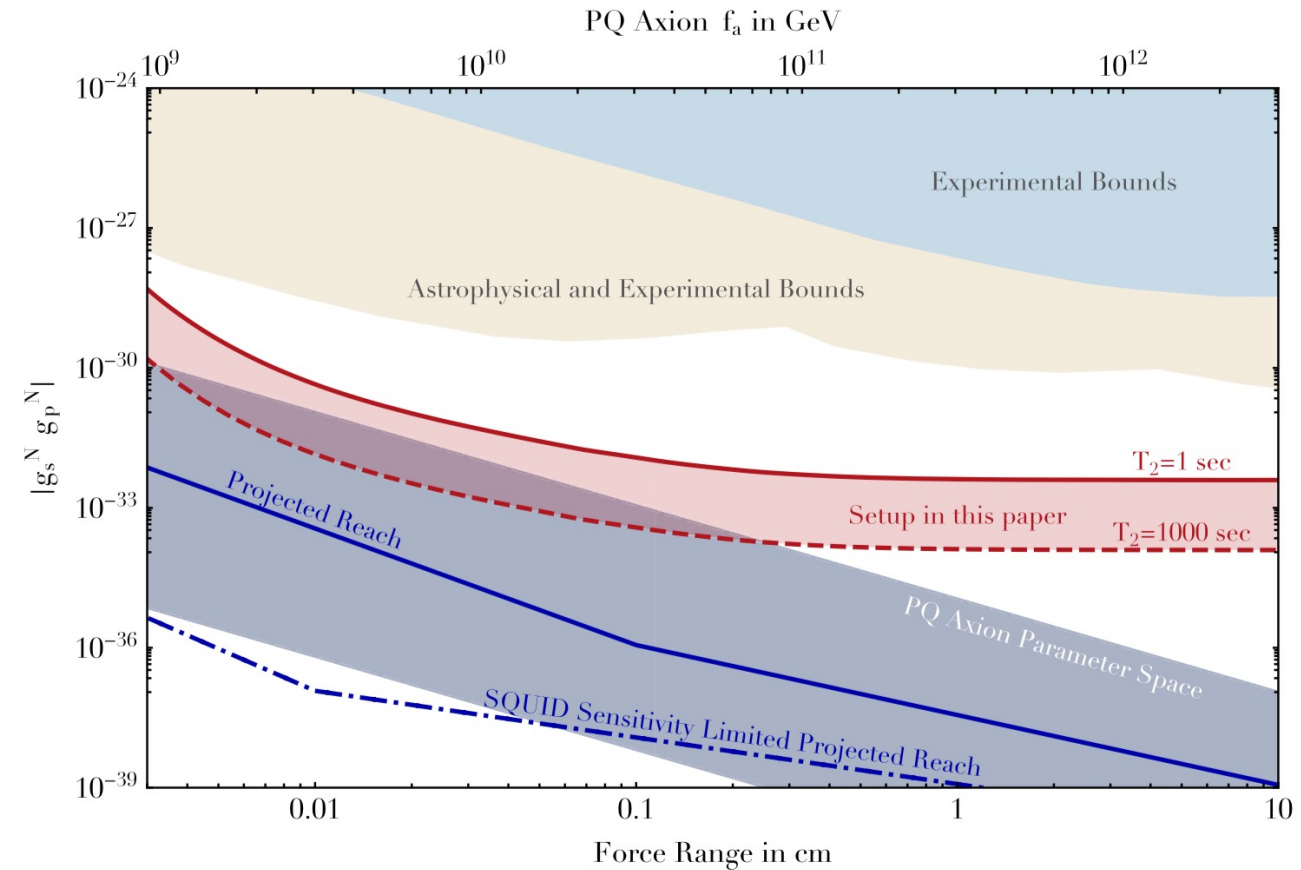
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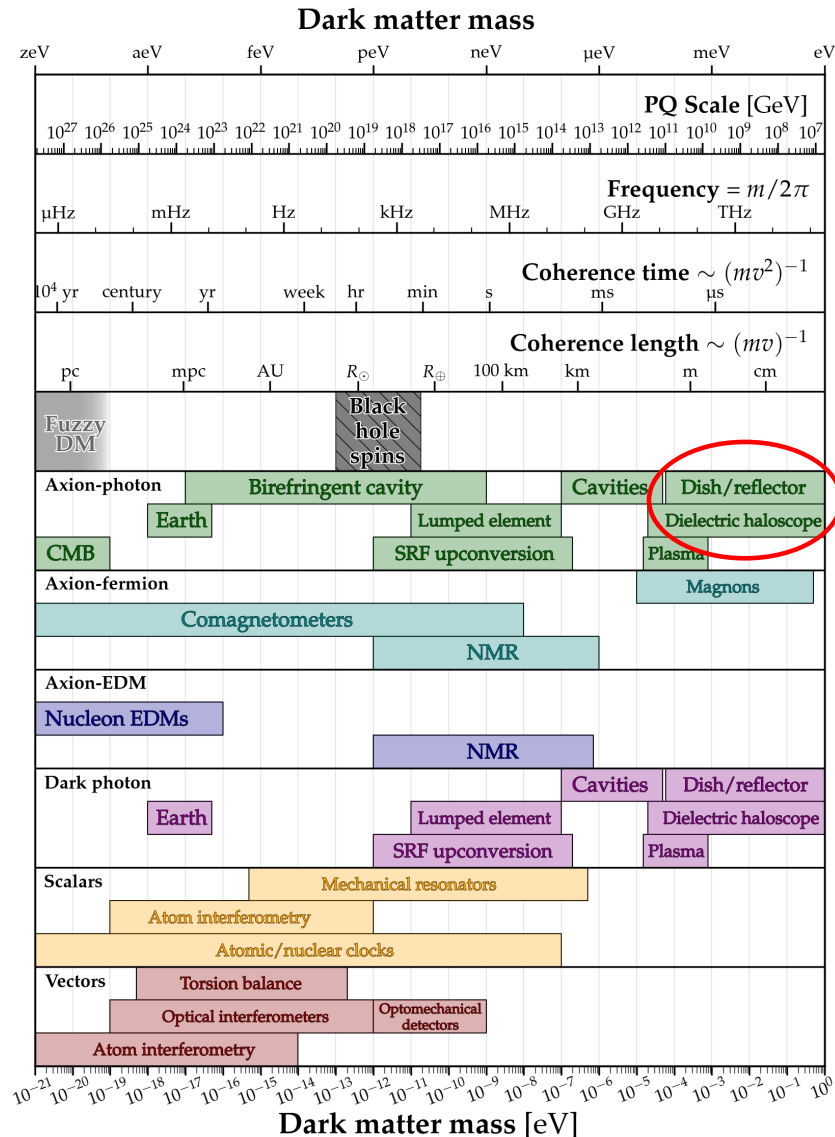
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# Axion Dark Matter Experiments

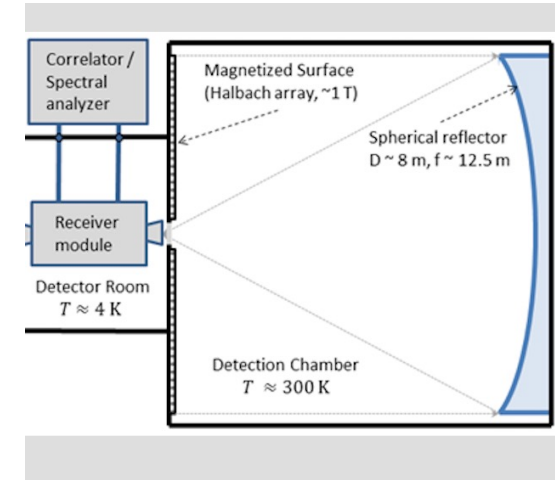
## Variety of experimental techniques



## Dish/reflector example:

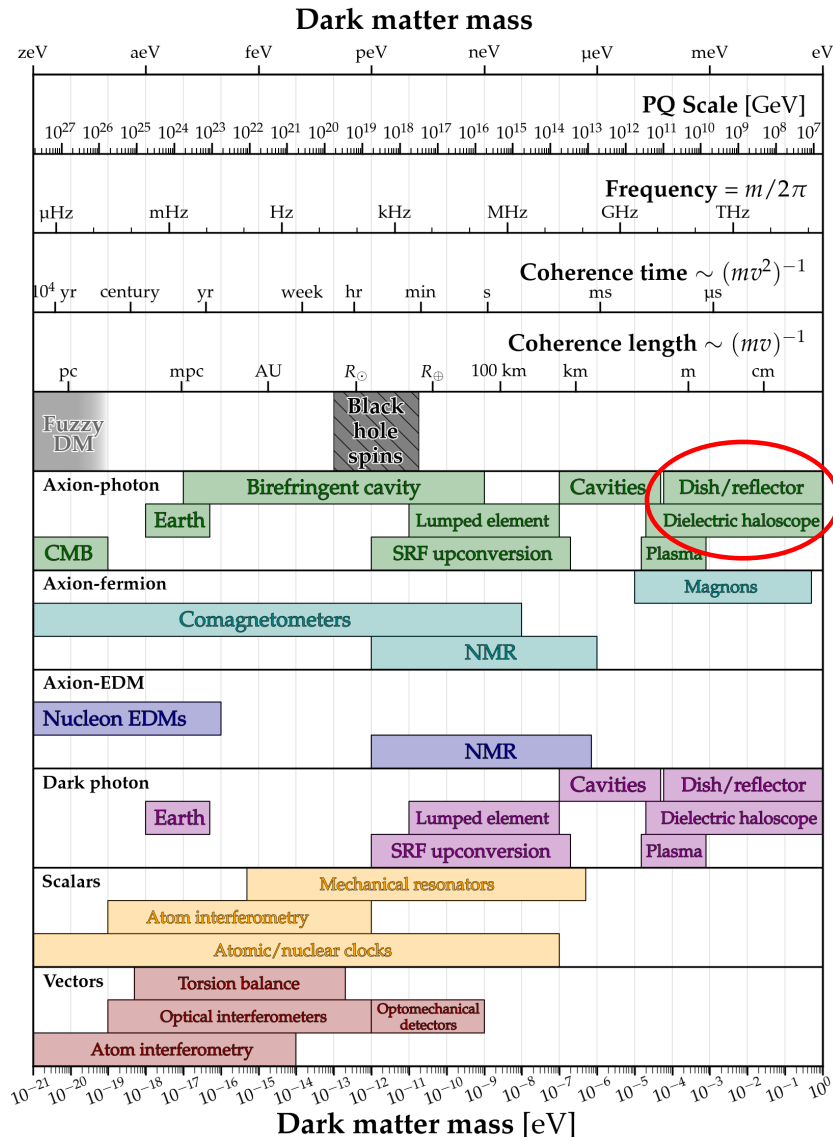
### BRASS @ UHH

- Plane permanently magnetized conversion panel
- Spherical reflector



# Axion Dark Matter Experiments

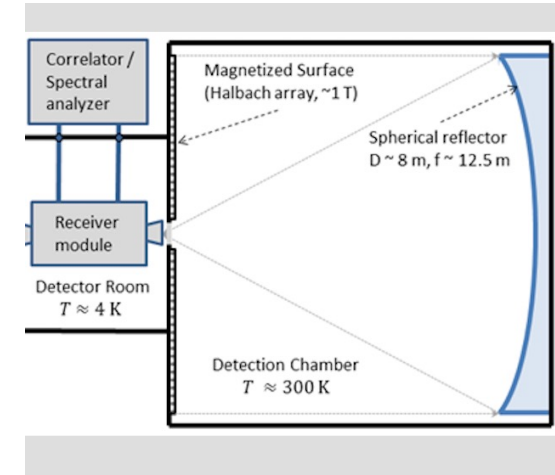
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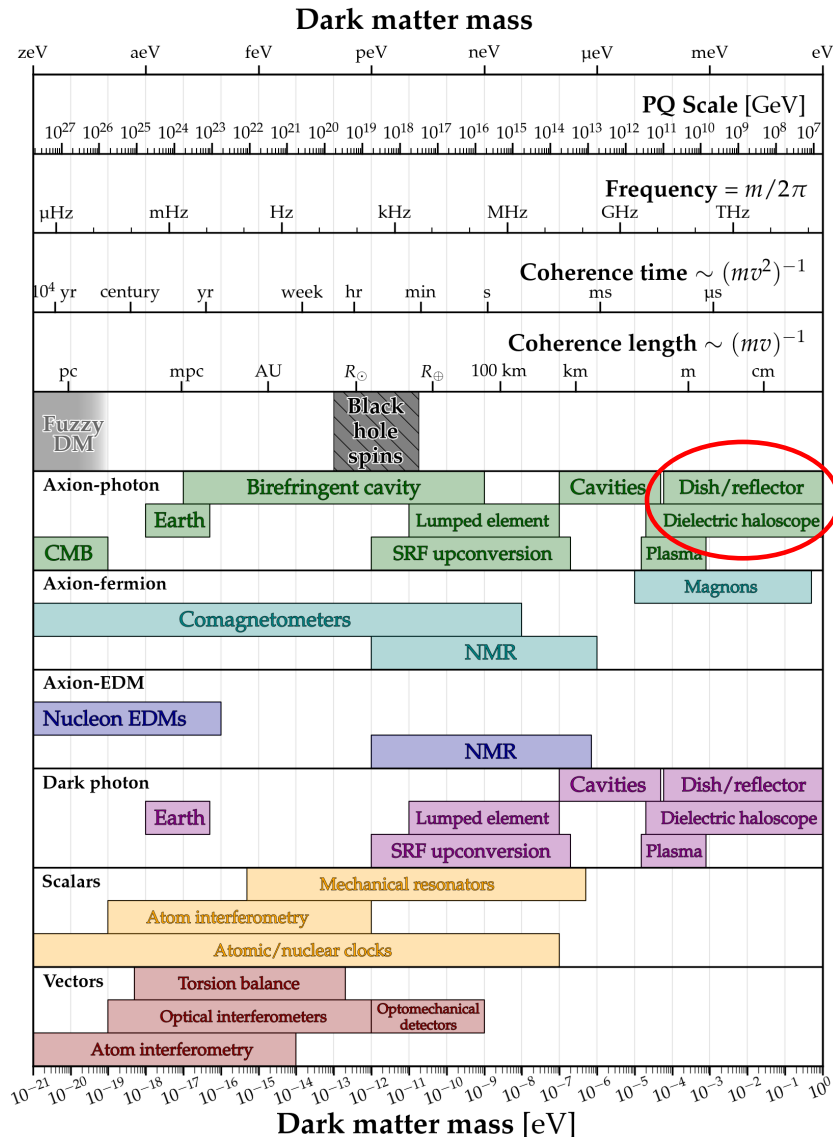


### BRASS-p pioneering run

- Conversion panel not magnetized
- Dark photon search [Bajjali et al., JCAP 08 (2023) 077]

# Axion Dark Matter Experiments

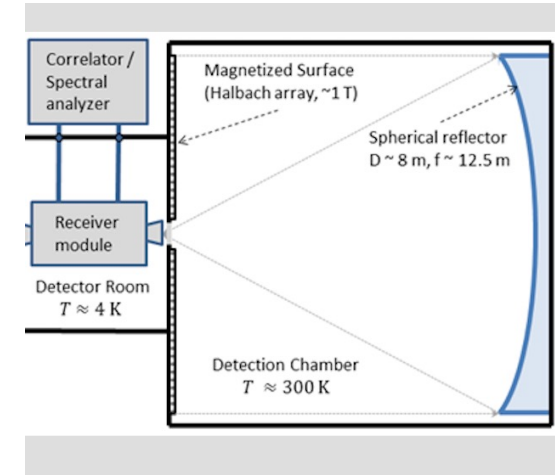
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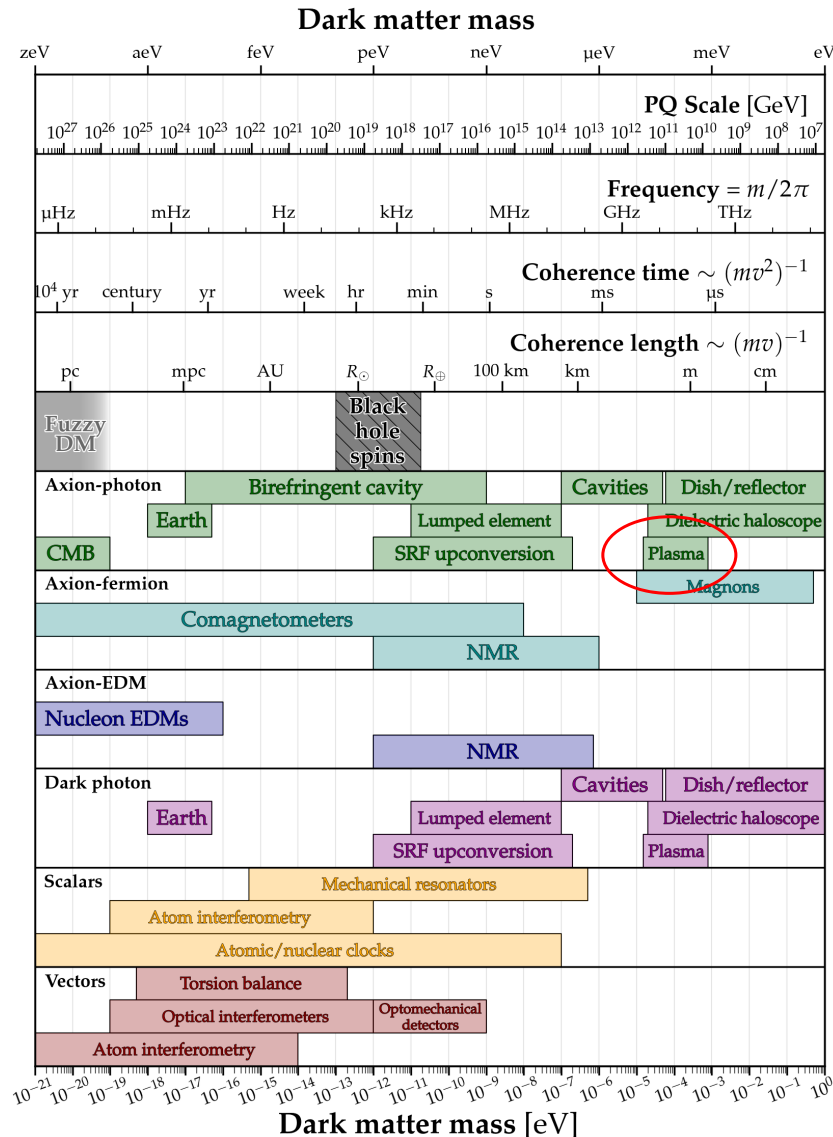
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BRASS magnetization on the way ...

# Axion Dark Matter Experiments

## Variety of experimental techniques



## Plasma haloscope concept:

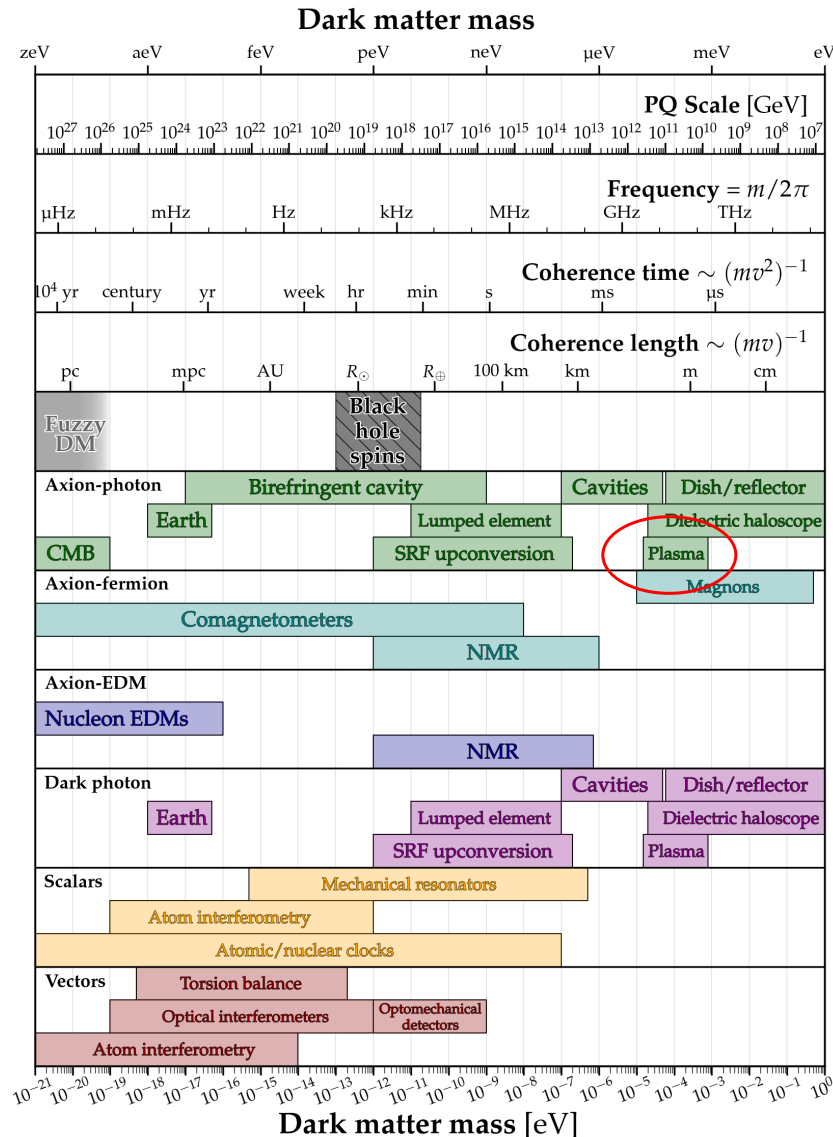
[Lawson, Millar, Pancaldi, Vitagliano, Wilczek, 19]

- In a magnetized plasma, oscillating axion DM induces plasmon excitations,

$$\mathbf{E} = -g_{a\gamma} \mathbf{B}_e a \left( 1 - \frac{\omega_p^2}{\omega_a^2 - i\omega_a \Gamma} \right)^{-1}$$

# Axion Dark Matter Experiments

## Variety of experimental techniques



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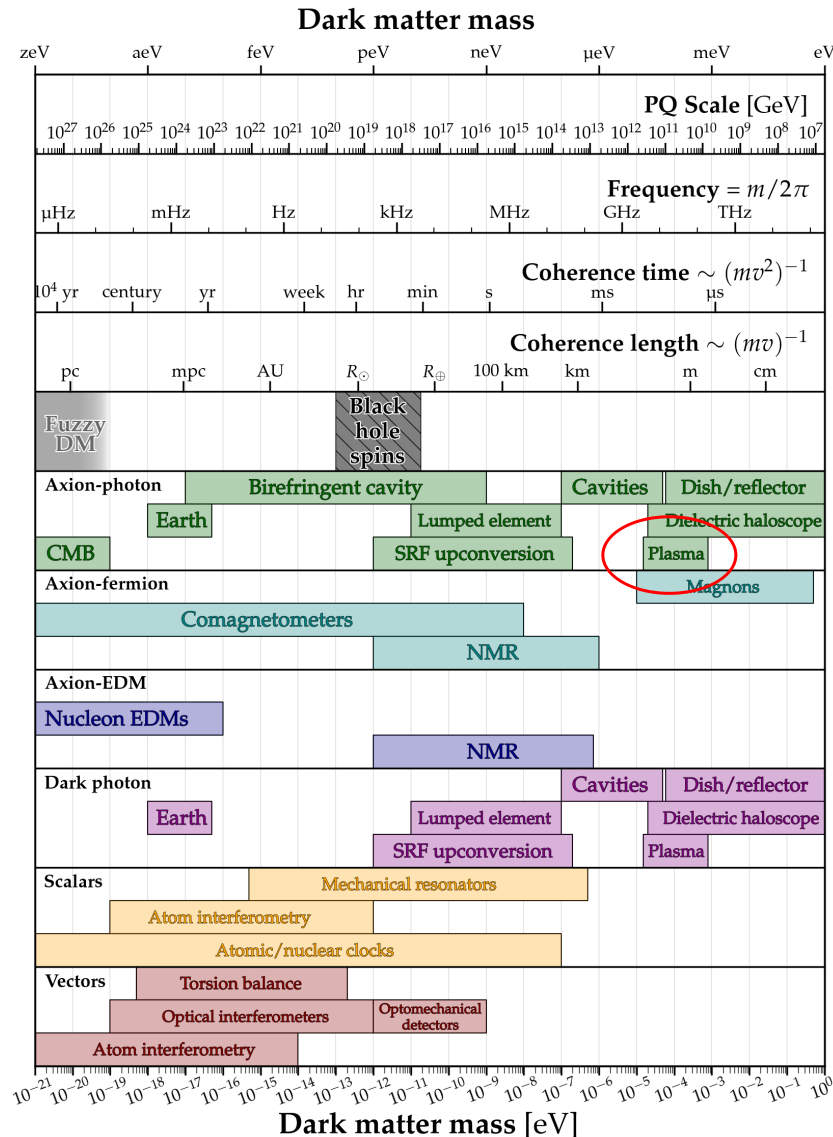
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$$\omega_p = \omega_a \approx m_a$$

# Axion Dark Matter Experiments

## Variety of experimental techniques



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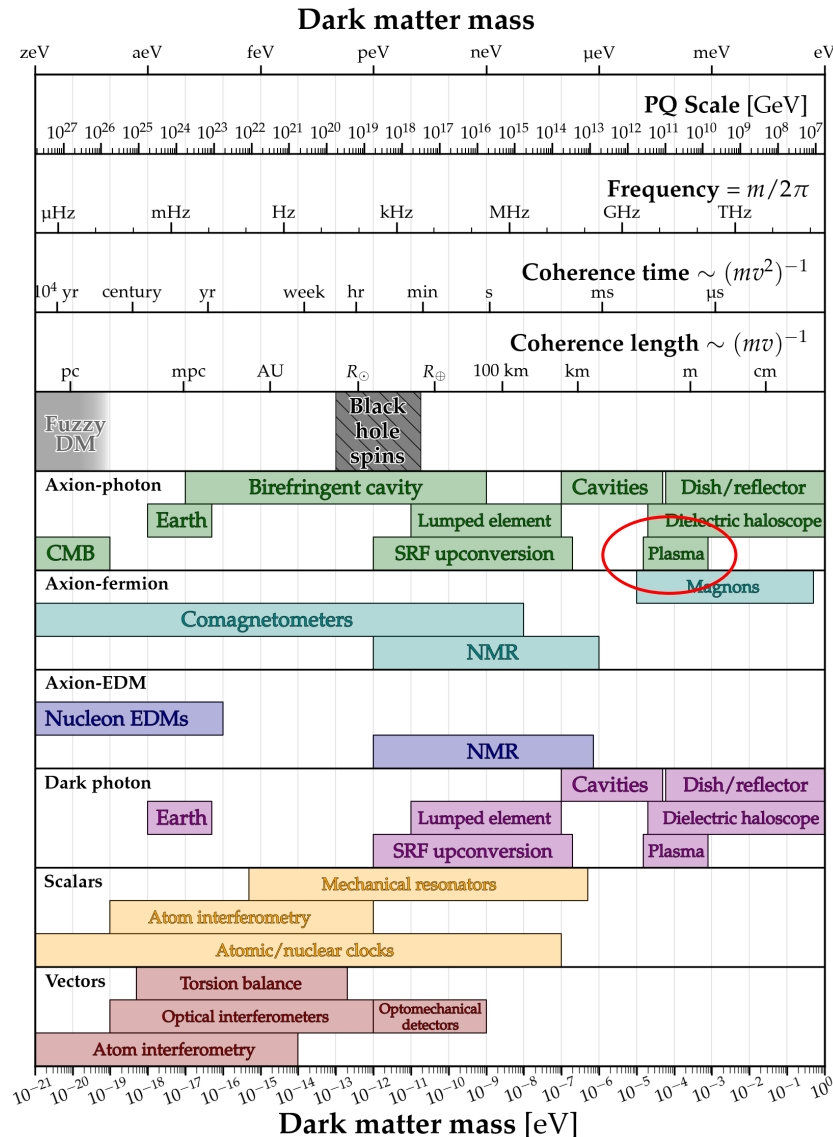
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# Axion Dark Matter Experiments

## Variety of experimental techniques



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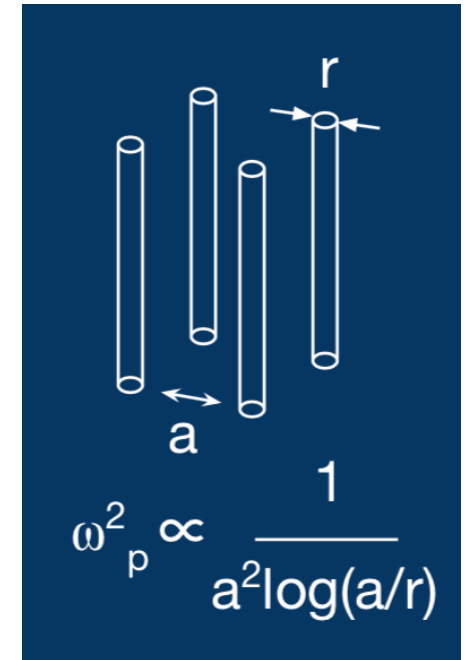
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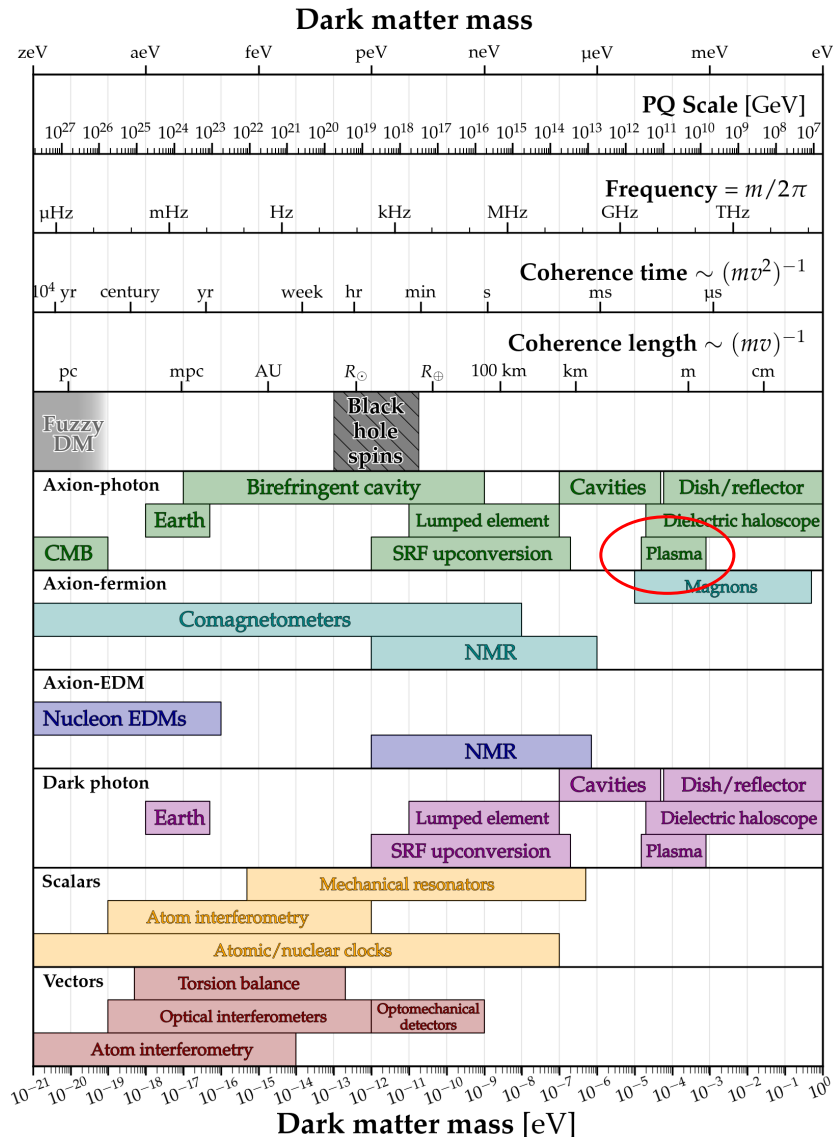
$$\omega_p = \omega_a \approx m_a$$

- limited by losses ( $\Gamma$ )
- A plasma with tunable plasma frequency in the GHz range can be realised by a wire array with variable interwire spacing (“wire metamaterial”)



# Axion Dark Matter Experiments

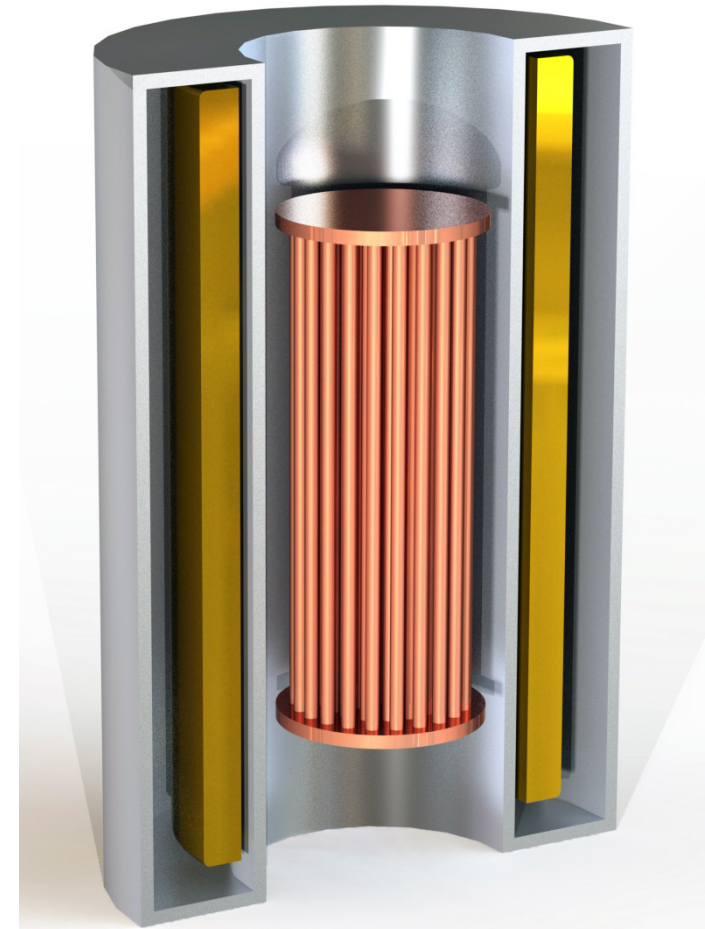
## Variety of experimental techniques



## Plasma haloscope example:

### ALPHA @ Yale

- Goal: building tunable, cryogenic plasma haloscope
- ALPHA Phase I
  - 17 cm bore 9T solenoid
  - dark photon search planned in this year (2026)



[<https://axion-dm.yale.edu/alpha-experiment>]

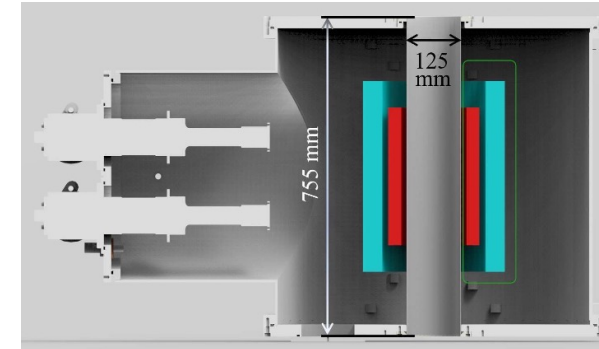
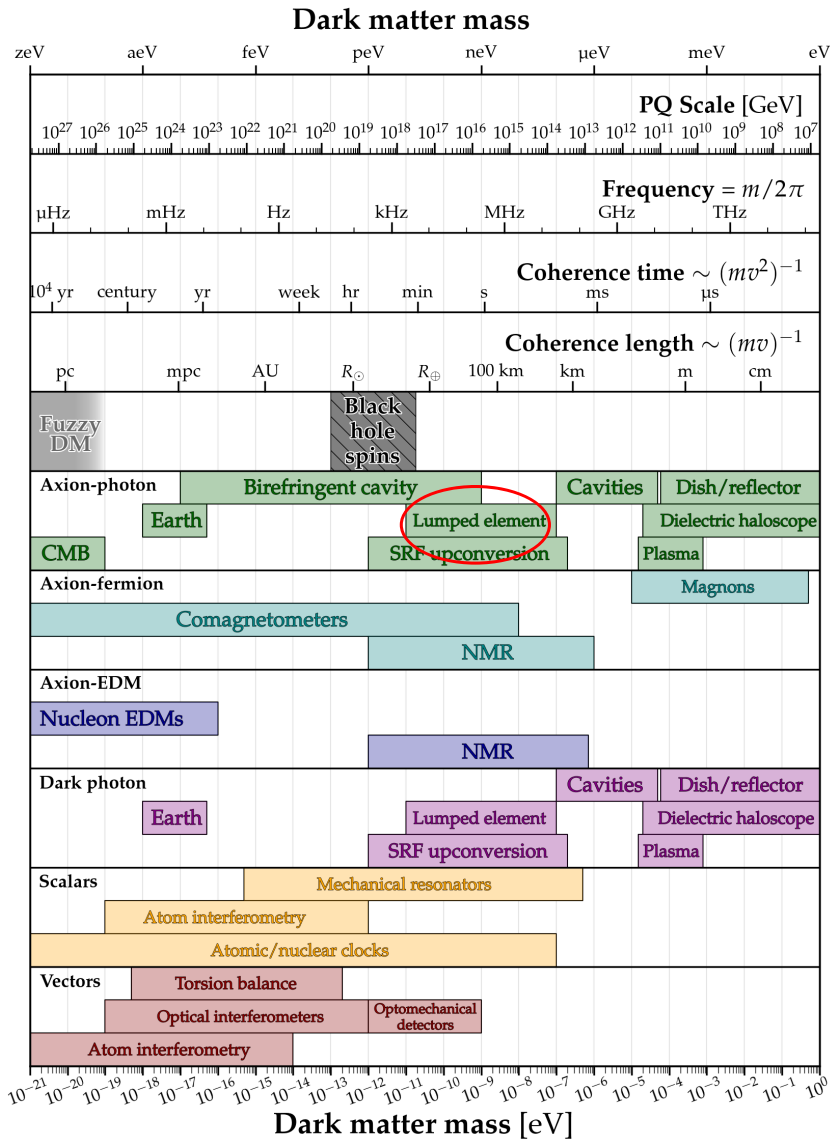
# Axion Dark Matter Experiments

## Variety of experimental techniques

## Lumped element examples:

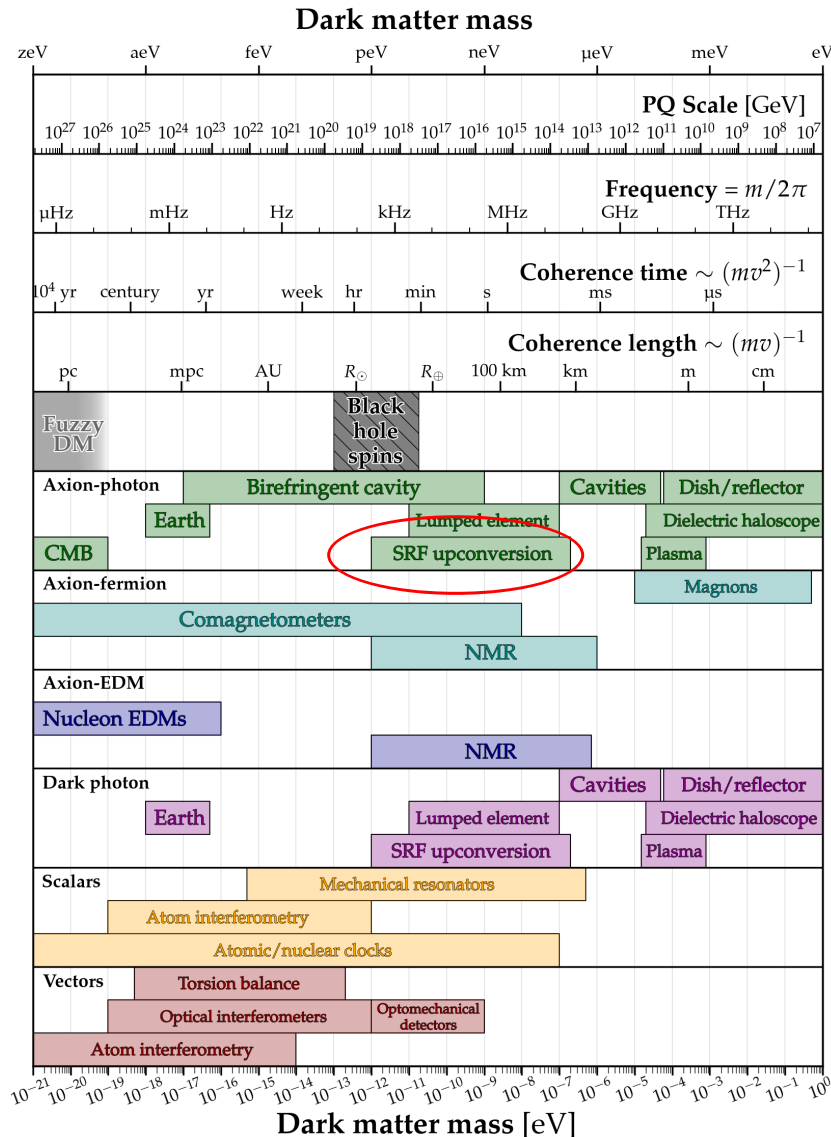
### WISPLC @ UHH

- Proposal  
[Zhang,Horns,Ghosh, Phys.Rev.D 106 (2022) 2, 023003]
- First pioneering run with WISPLC prototype in 6 T magnet in 2024
- Second run ongoing



# Axion Dark Matter Experiments

## Variety of experimental techniques



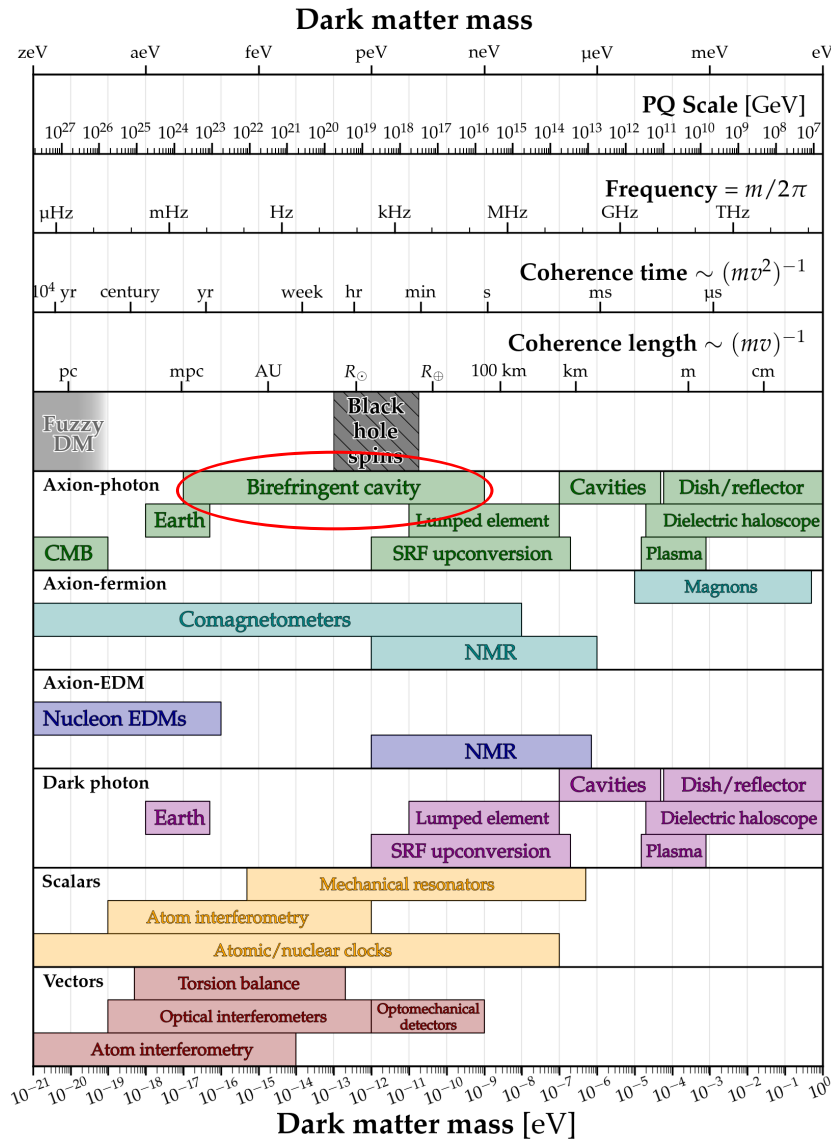
## SRF up-conversion /heterodyne concept:

[Berlin et al., JHEP 07 (2020) 088;  
 Berlin et al., PRD 104 (2021) L111701;  
 Giaccone et al., 2207.11346]

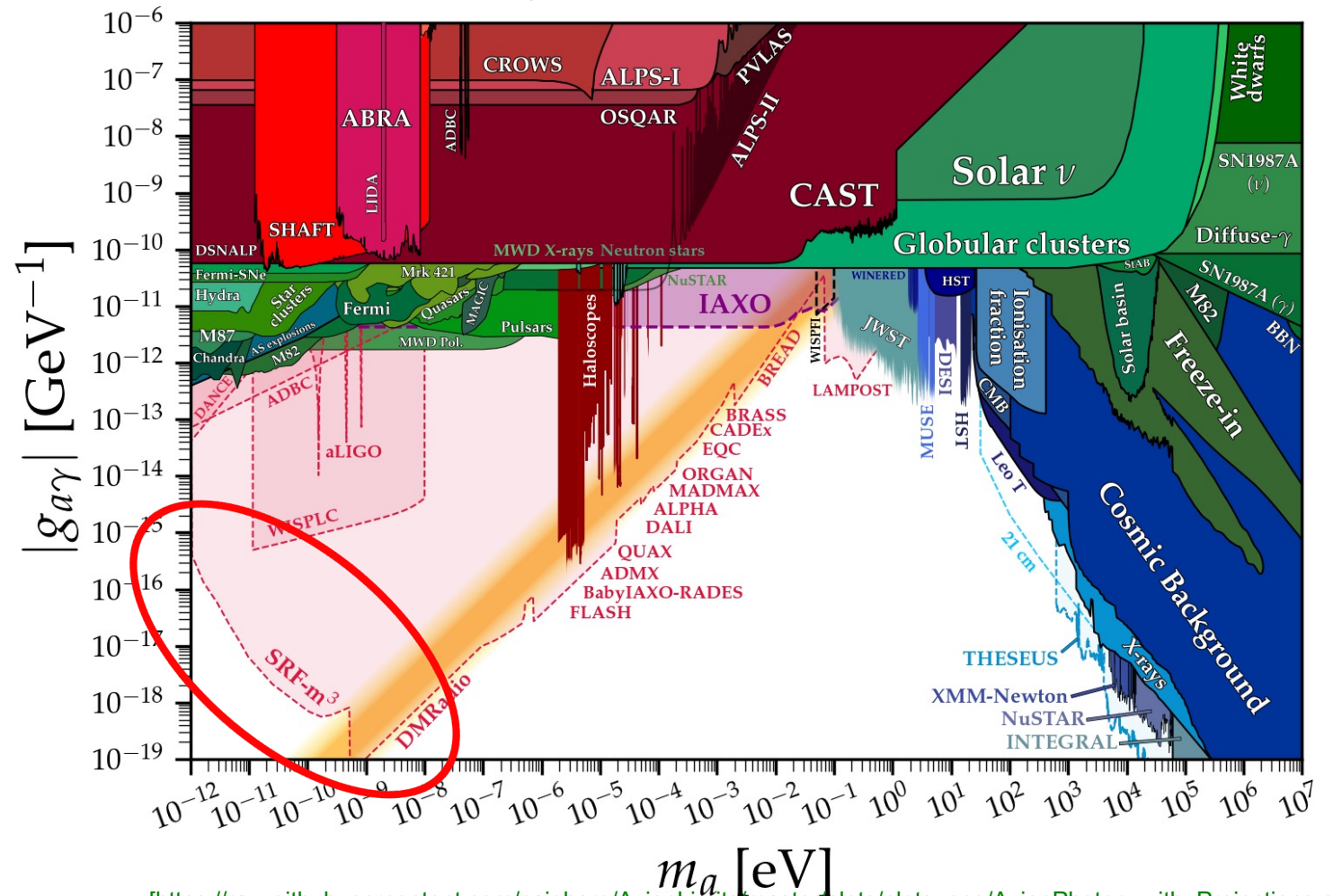
- Axion DM coupled to photons induces transitions between two electromagnetic cavity modes.
- Drive a high-Q pump mode at frequency  $\omega_0$
- Axion DM converts a tiny fraction of the stored pump power into a signal mode at  $\omega_1 \simeq \omega_0 \pm m_a$
- The axion mass is encoded in the mode splitting:  $m_a \simeq |\omega_1 - \omega_0|$
- Using superconducting RF cavities gives ultra-high quality factors, potentially avoiding the large-volume problem of low-mass conventional haloscopes.
- Main challenges: pump leakage, mode mixing, mechanical noise, microphonics, tuning and stability of nearly degenerate modes.
- Experimental R&D is underway in SHADE/SQMS; no published axion-DM limit yet

# Axion Dark Matter Experiments

## Variety of experimental techniques



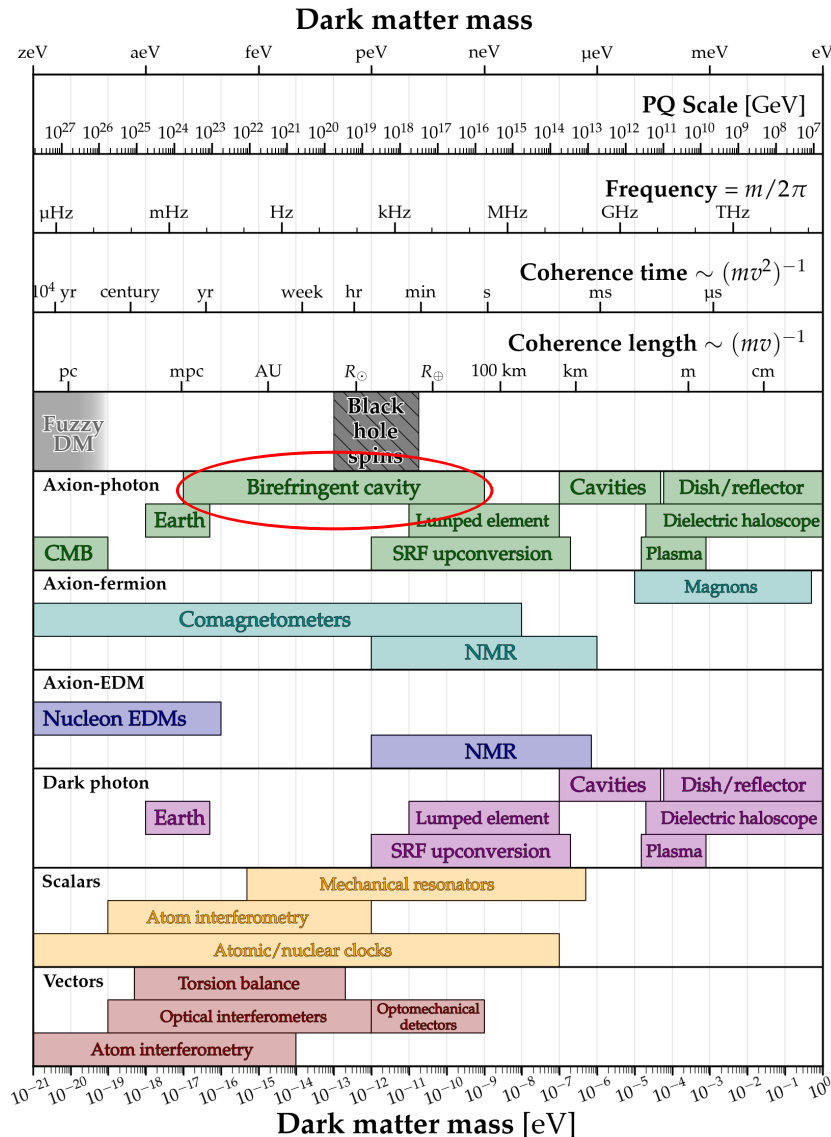
## SRF up-conversion /heterodyne prospects:



[[https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots\\_png/AxionPhoton\\_with\\_Projections.png](https://raw.githubusercontent.com/cajohare/AxionLimits/master/plots/plots_png/AxionPhoton_with_Projections.png)]

# Axion Dark Matter Experiments

## Variety of experimental techniques



## Birefringent optical-cavity concept:

[Obata et al.,1805.11753; Liu et al.,1809.01656]

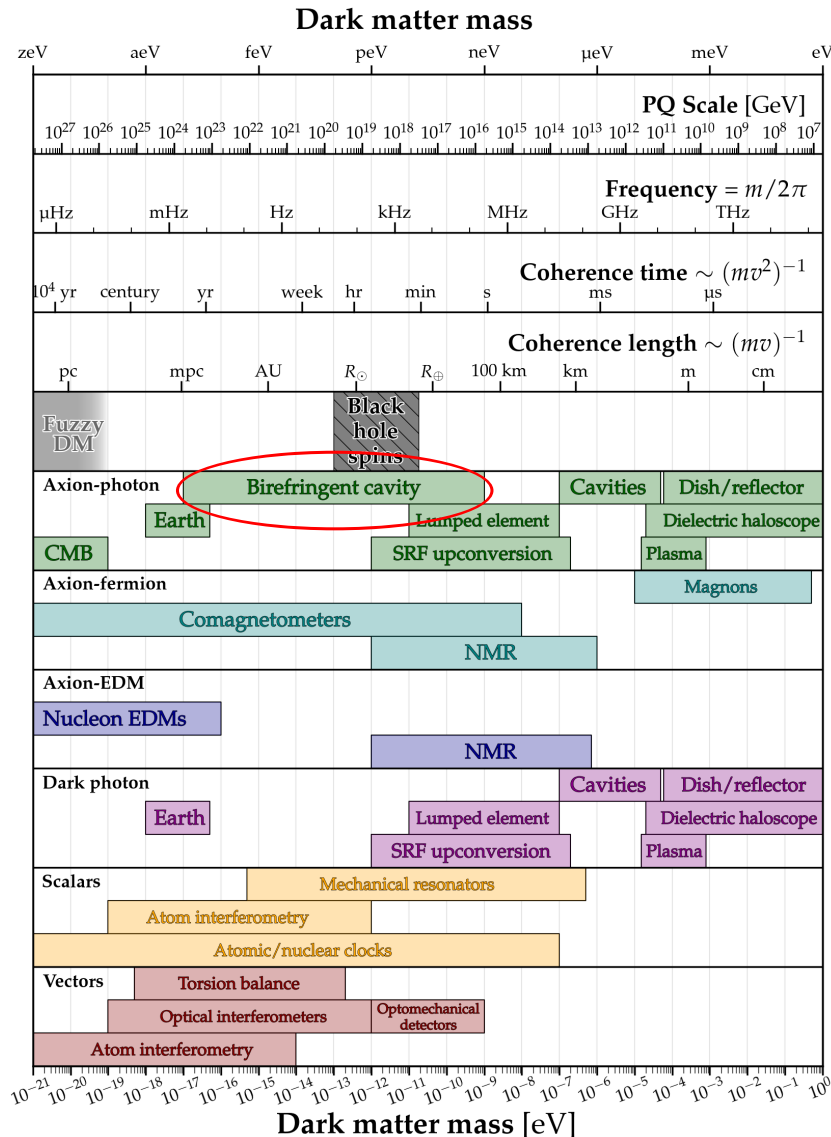
- The axion background acts as a tiny circularly birefringent medium:

$$c_{L/R} = 1 \pm \delta c(t), \quad \delta c(t) \simeq \frac{g_{a\gamma\gamma} \dot{a}(t)}{2\omega_L}$$

- A linearly polarized laser is a superposition of left- and right-circular polarizations; the axion-induced phase shift rotates its polarization at frequency  $\omega_a \simeq m_a$
- Equivalently, the axion converts a strong carrier field at ( $\omega_0$ ) into weak sidebands in the orthogonal polarization at  $\omega_0 \pm m_a$
- A high-finesse optical cavity enhances the interaction time / stored optical power; engineered birefringence, ring geometry, or wave plates can make the axion-induced sidebands build up coherently
- Search for a narrow oscillating polarimetric signal — no external magnetic field required.

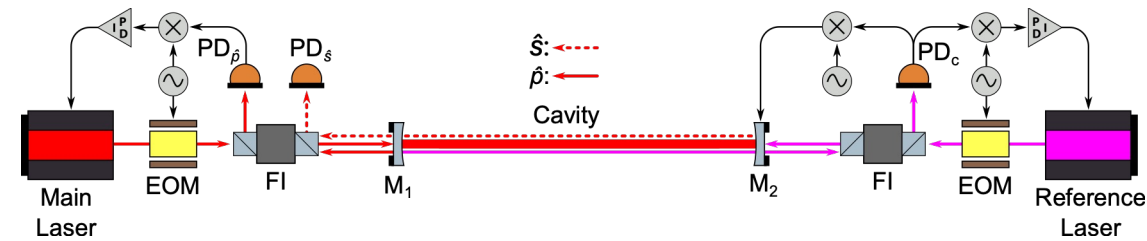
# Axion Dark Matter Experiments

## Variety of experimental techniques



## Example: ALPS II optical cavity as axion-DM polarimeter

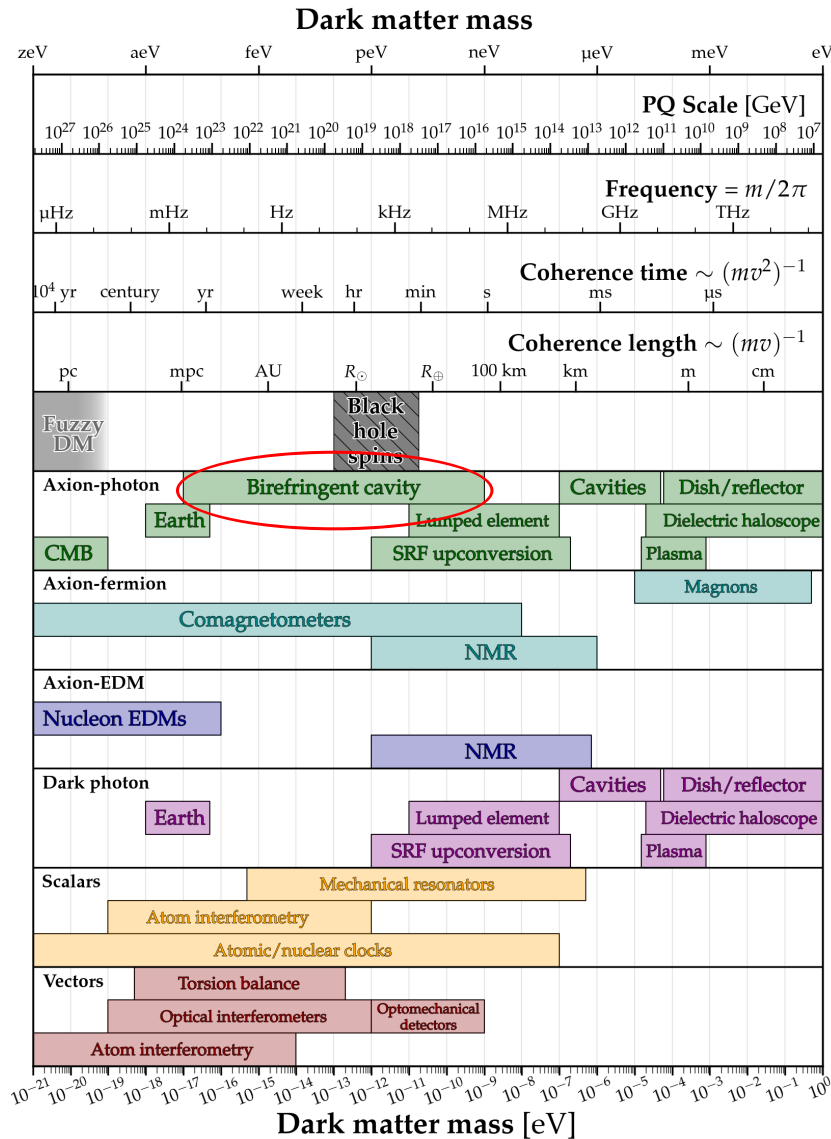
[Garcia-Cely, Marsili, AR, Spector, 2501.08382]



- ALPS II already provides long, high-finesse optical cavities with large stored laser power
- Axion DM would induce a tiny oscillating polarization component orthogonal to the injected laser polarization
- Polarimetric readout of the cavity output searches for sidebands at  $\omega_0 \pm m_a$
- Using different ALPS II cavity configurations, including quarter-wave plates, gives resonant or broader sensitivity over  $m_a \sim (10^{-9} - 10^{-6})$  eV
- Same formalism also applies to high-frequency gravitational-wave backgrounds.

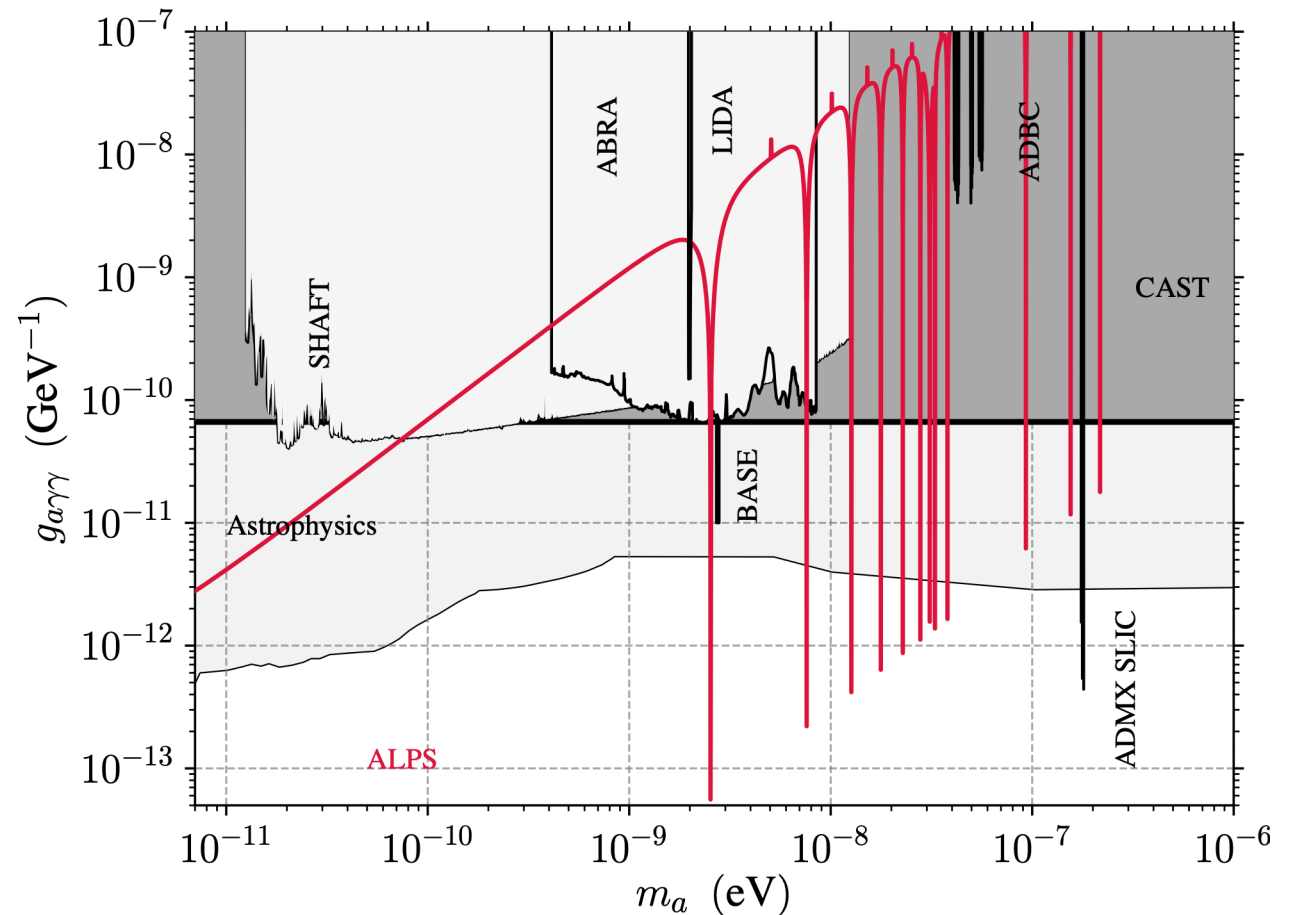
# Axion Dark Matter Experiments

## Variety of experimental techniques



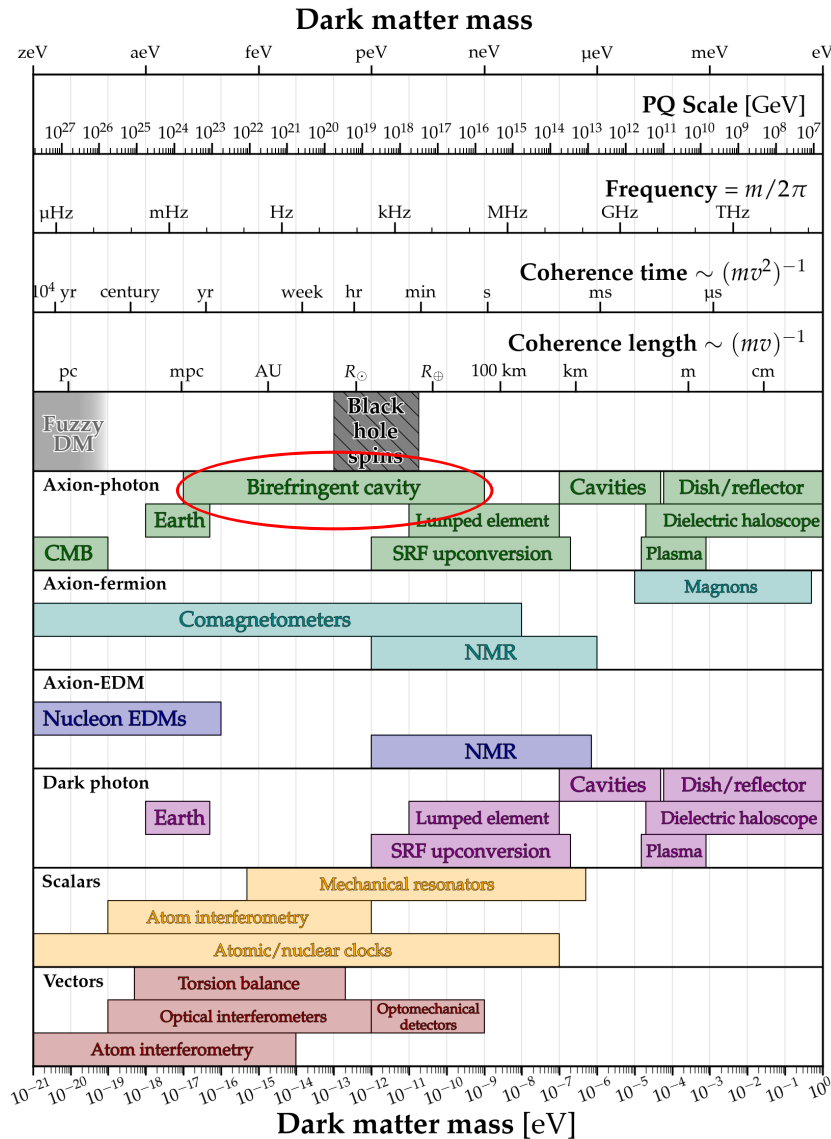
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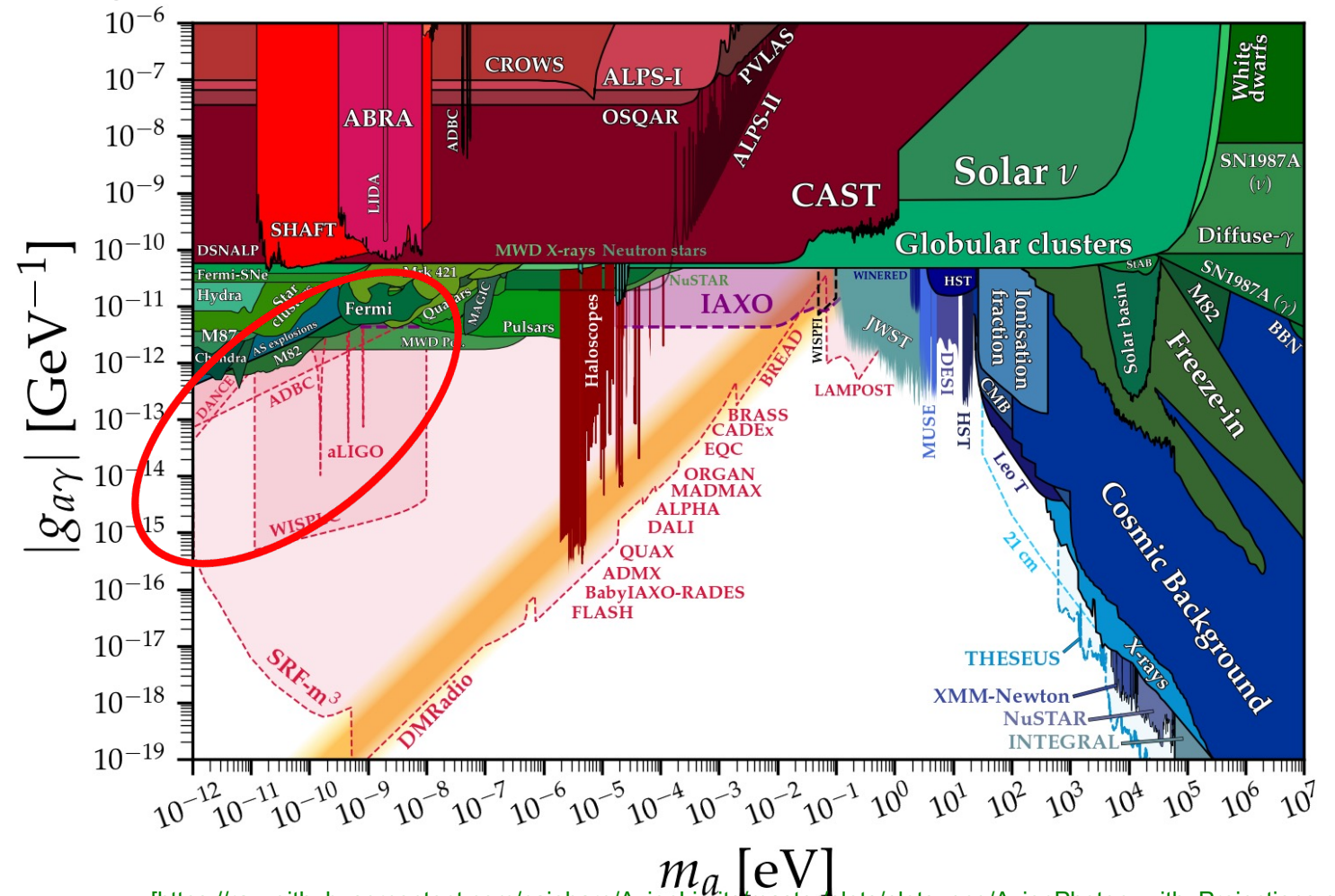


# Axion Dark Matter Experiments

## Variety of experimental techniques



## Birefringent cavities prospects:



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