

Towards a minimal gCP modular GUT

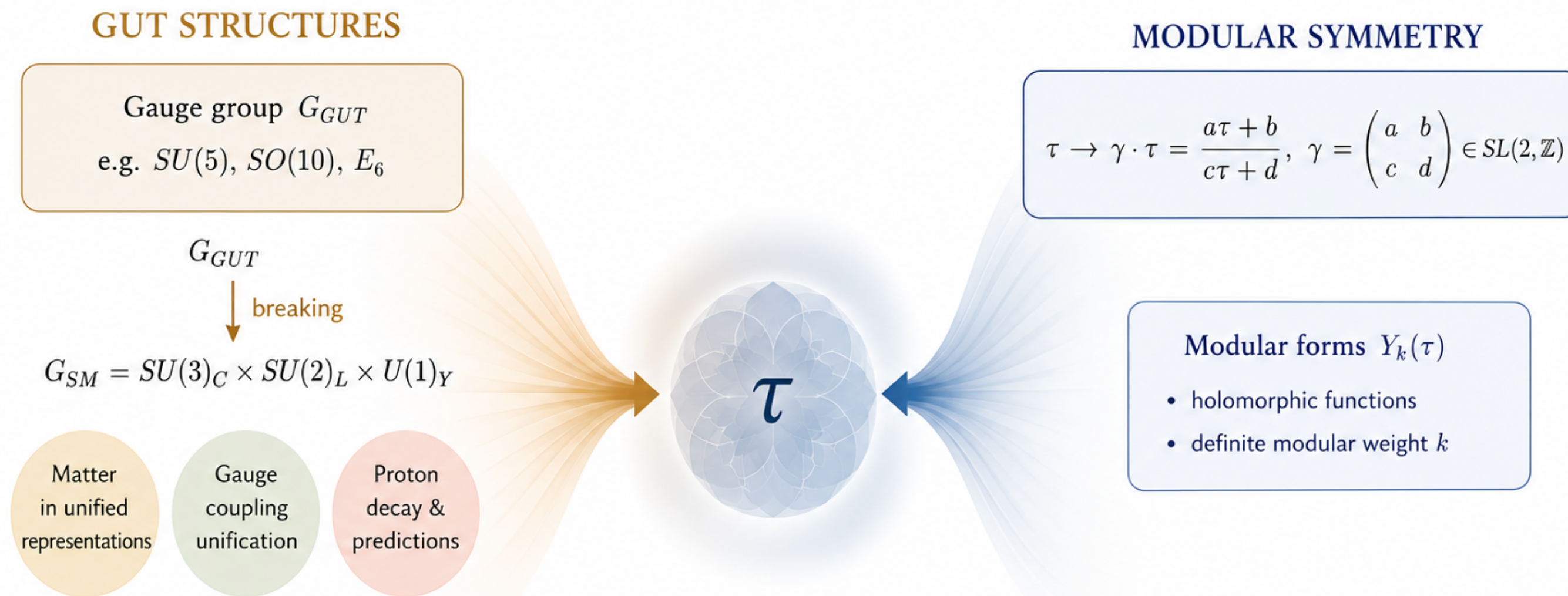
Marco Carducci



GUT \otimes MODULAR SYMMETRY

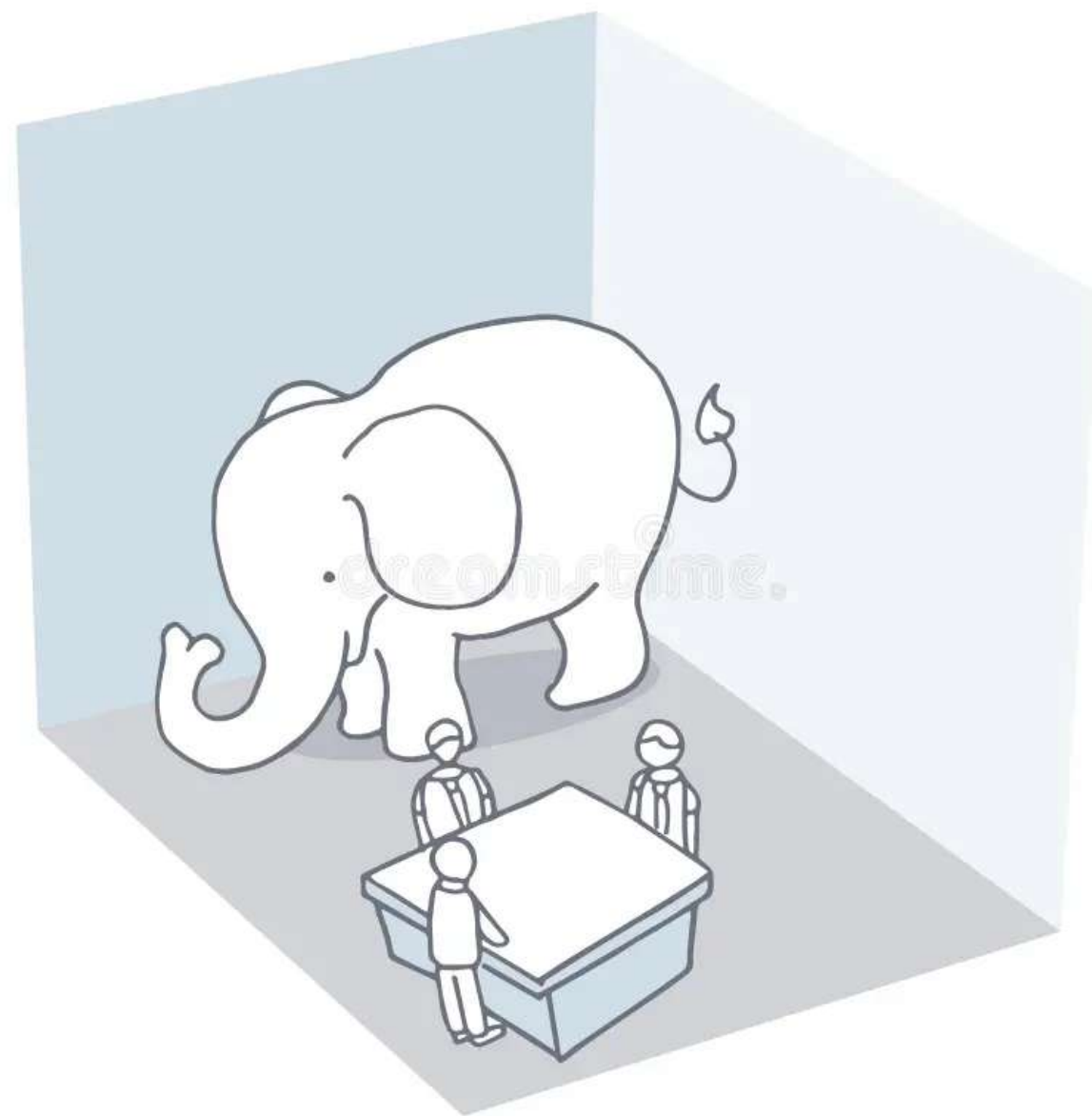
flavour from the interplay between horizontal and vertical symmetries

In collaboration with D. Meloni and J. T. Penedo

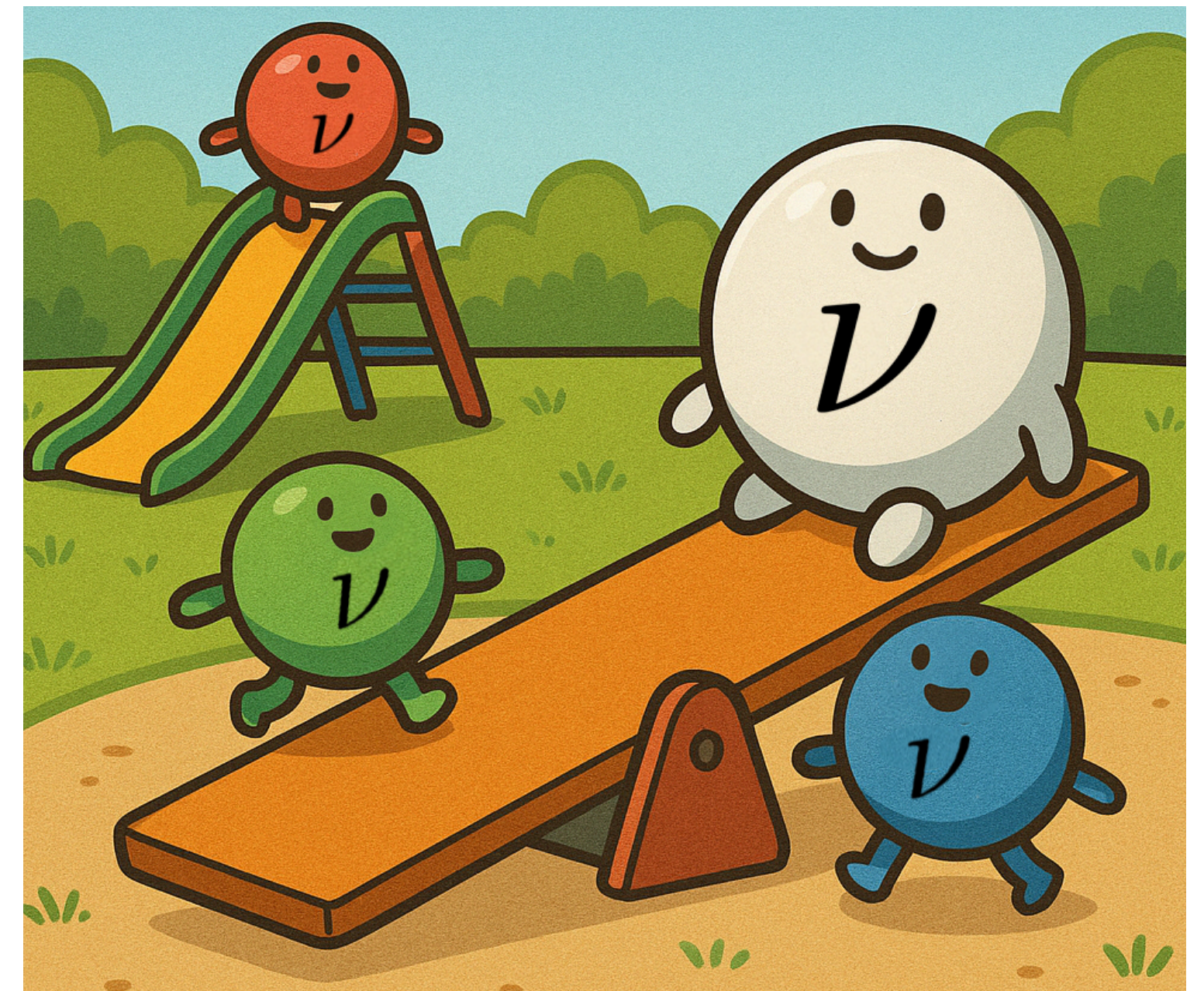


PASCOS 2026 (23 June)

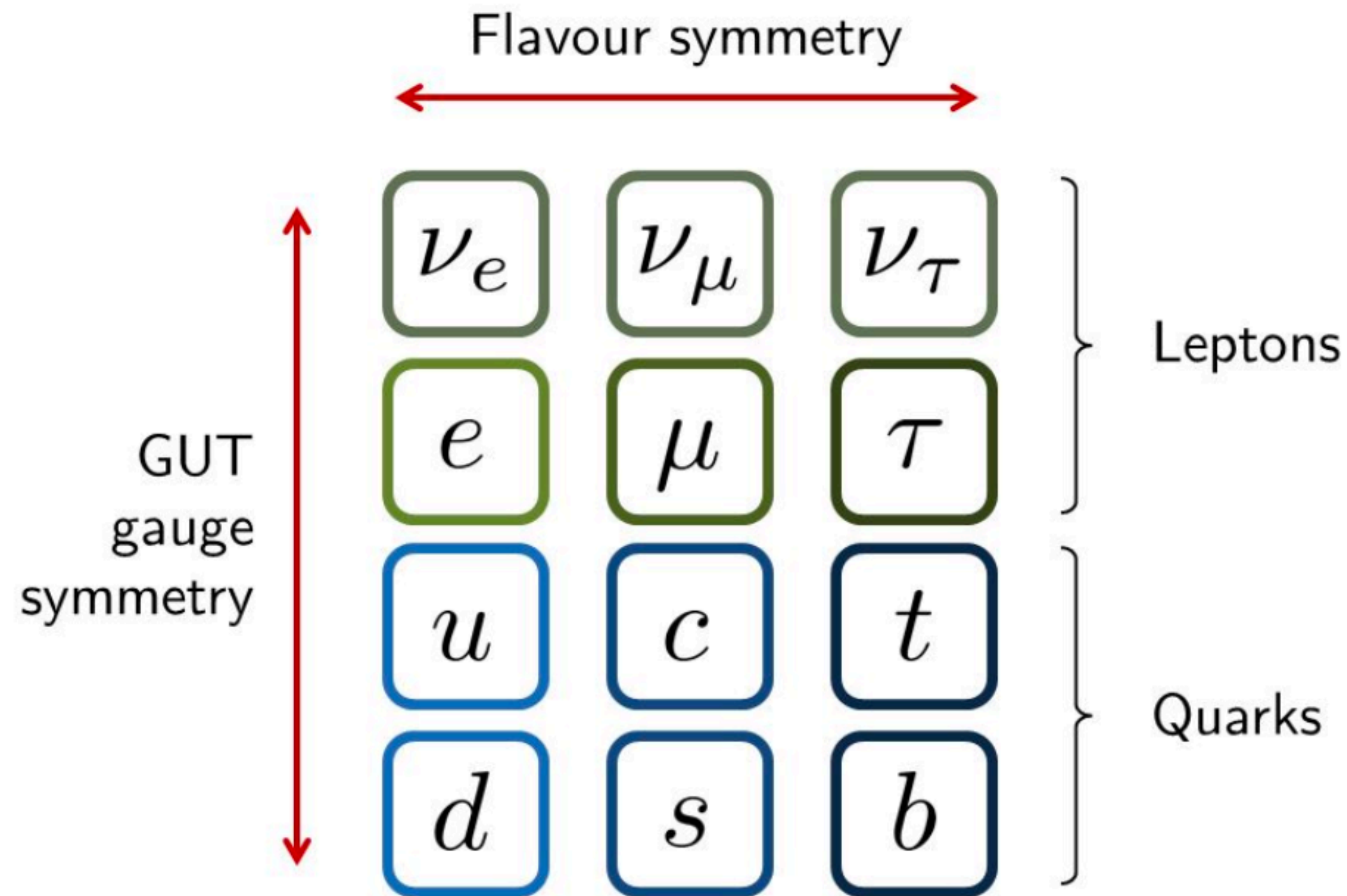
Flavour puzzle



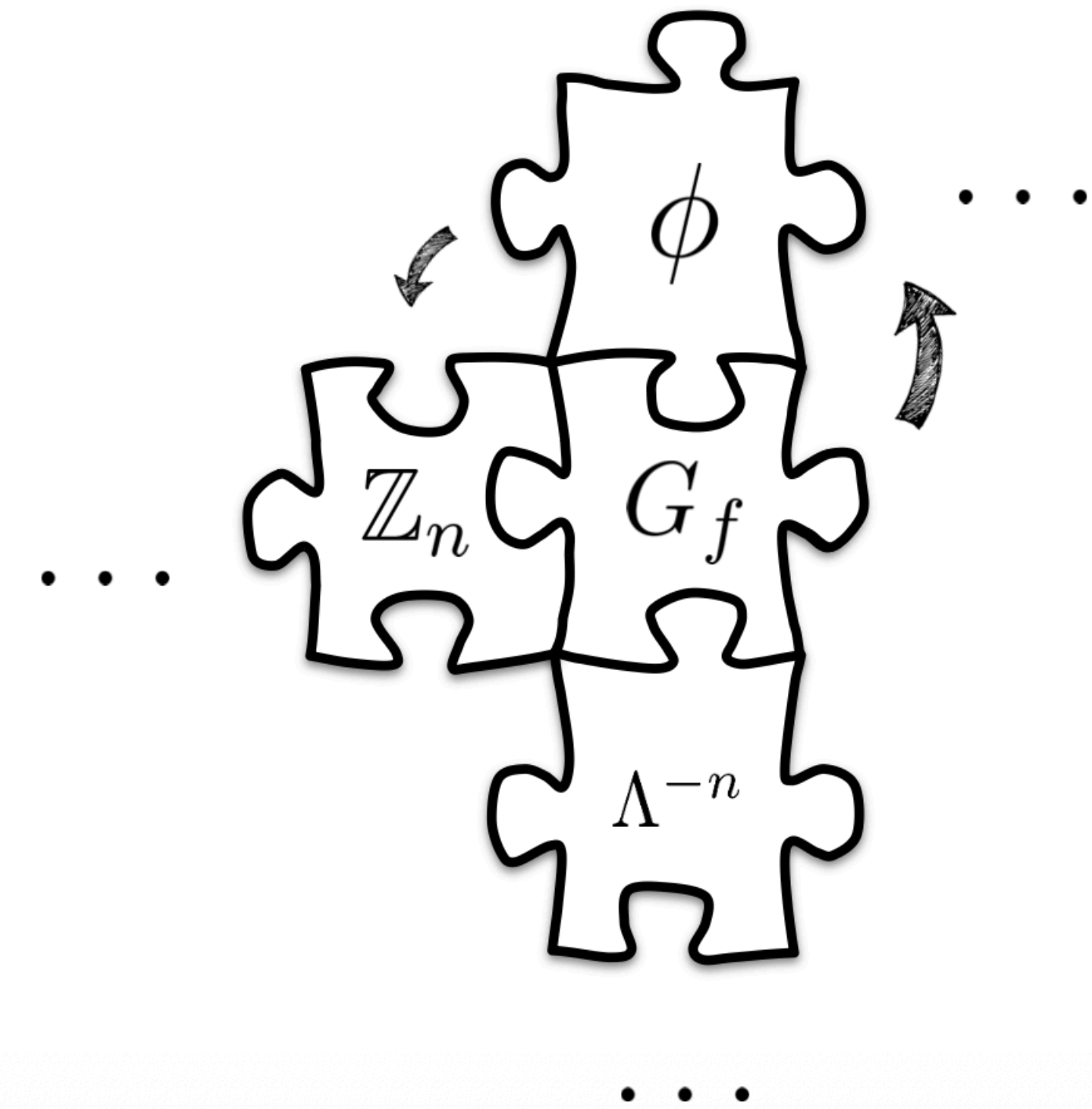
$$\|V_{CKM}\| = \begin{pmatrix} \text{black} & \text{light gray} & \text{white} \\ \text{light gray} & \text{black} & \text{white} \\ \text{white} & \text{white} & \text{black} \end{pmatrix} \quad \|U_{PMNS}\| = \begin{pmatrix} \text{black} & \text{medium gray} & \text{light gray} \\ \text{light gray} & \text{medium gray} & \text{dark gray} \\ \text{light gray} & \text{medium gray} & \text{dark gray} \end{pmatrix}$$

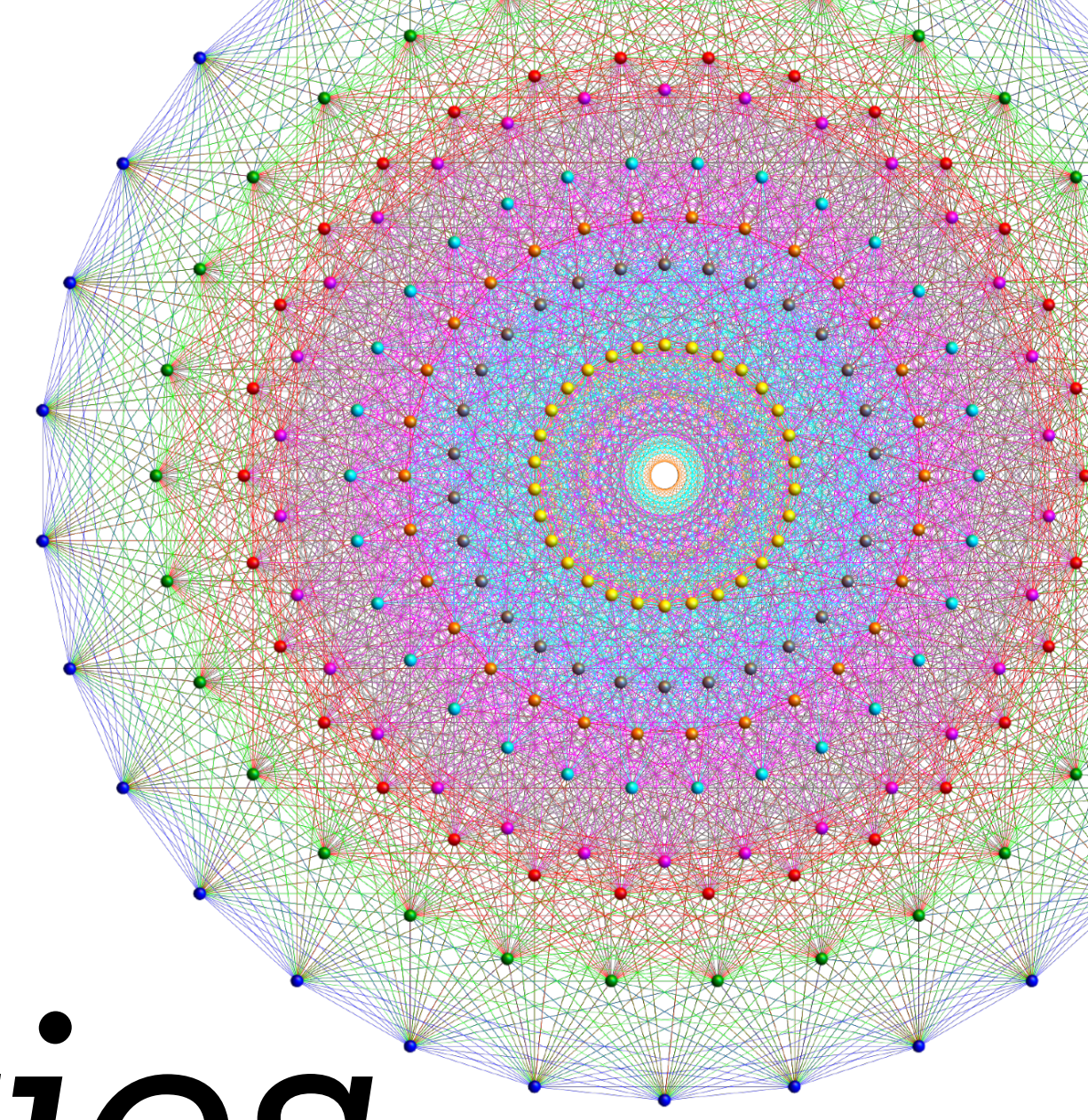


Flavour puzzle



Attempts...





Grand Unified Theories

STANDARD MODEL

u	c	t	q	H
d	s	b	g	
ν_e	ν_μ	ν_τ	Z	
e	μ	τ	W	

STANDARD MODEL interaction

GRAND UNIFIED THEORY

ALIAS

GUTs

GUT SCALE

DOCUMENT FILE

b

NAME	Bottom Quark
ALIAS	beauty
DOB	1977
MASS	4.18 GeV
CHARGE	1/3
BIRTHPLACE	Fermilab
SPIN	1/2

Predicted before Discovered

Z
W

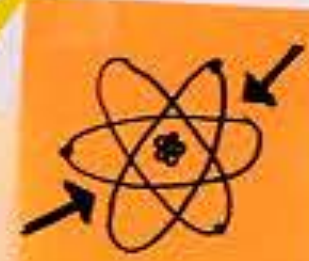
WEAK NUCLEAR FORCE

- particles decay
- short ranged

ELECTROMAGNETIC FORCE

dictates the structure of atoms and the behavior of light

e⁻
electron



electroweak interaction

HYPER K



3 Neutrino TYPES

- ν_e electron
- ν_τ tau
- ν_μ muon



STRONG NUCLEAR FORCE

Binds quarks together into things like the protons and neutrons in our atoms

u, c, t, d, s, b

Additional Higgses?

Complicated interactions and more PARTICLES



PROTON DECAY ?

SUSY

NAME	SUPERSYMMETRY
NICKNAME	SUSY
SUSPECT ACCOMPLICE	SPARTICLE

CONFIDENTIAL

Messy

DUNE



Supersymmetric $SU(5)$: Particle Content

H. Georgi and S. L. Glashow

Phys. Rev. Lett.



SM group is of **rank 4**: Embed into simple groups of rank 4 or more.



SU(5) Each generation a 10 and $\bar{5}$ (without RH neutrino)

Minimal non-SUSY model ruled out by $\sin^2 \theta_W$

Supersymmetric $SU(5)$: Particle Content

B. Bajc and G. Senjanovic
[hep-ph/0612029]-JHEP



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$$F \equiv \bar{5}_F = \begin{pmatrix} d^c \\ \epsilon_2 \ell \end{pmatrix} \quad 10_F = \begin{pmatrix} \epsilon_3 u^c & q \\ -q^T & \epsilon_2 e^c \end{pmatrix} \equiv T$$

$$\Sigma(24) = \begin{pmatrix} \Sigma_8 & \Sigma_{(3,2)} \\ \Sigma_{(\bar{3},2)} & \Sigma_3 \end{pmatrix} + \frac{1}{2\sqrt{15}} \begin{pmatrix} 2 \mathbf{1}_3 & 0 \\ 0 & -3 \mathbf{1}_2 \end{pmatrix} \Sigma_{24}$$

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Label \curvearrowright Irrep \curvearrowright

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Adjoint \curvearrowright

$$\Sigma(24) = \begin{pmatrix} \Sigma_8 & \Sigma_{(3,2)} \\ \Sigma_{(\bar{3},2)} & \Sigma_3 \end{pmatrix} + \frac{1}{2\sqrt{15}} \begin{pmatrix} 2 \mathbf{1}_3 & 0 \\ 0 & -3 \mathbf{1}_2 \end{pmatrix} \Sigma_{24}$$

Supersymmetric $SU(5)$: Superpotential

D. Emmanuel-Costa and S. Wiesenfeldt

[hep-ph/0302272]-Nucl. Phys. B

$$W_Y = \sqrt{2} \left(Y_5^{ij} \bar{H}_a T_i^{ab} F_{jb} + h_1^{ij} \bar{H}_a \frac{\Sigma_b^a}{\Lambda} T_i^{bc} F_{jc} + h_2^{ij} \bar{H}_a T_i^{ab} \frac{\Sigma_b^c}{\Lambda} F_{jc} \right) + \\ + \frac{1}{4} \epsilon_{abcde} \left(Y_{10}^{ij} T_i^{ab} T_j^{cd} H^e + f_1^{ij} T_i^{ab} T_j^{cd} \frac{\Sigma_f^e}{\Lambda} H^f + f_2^{ij} T_i^{ab} T_j^{cd} H^d \frac{\Sigma_f^e}{\Lambda} \right)$$

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$$+ \frac{1}{4} \epsilon_{abcde} \left(Y_{10}^{ij} T_i^{ab} T_j^{cd} H^e + f_1^{ij} T_i^{ab} T_j^{cd} \frac{\sum_f^e}{\Lambda} H^f + f_2^{ij} T_i^{ab} T_j^{cd} H^d \frac{\sum_f^e}{\Lambda} \right)$$

NR operators

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$$+ \frac{1}{4} \epsilon_{abcde} \left(Y_{10}^{ij} T_i^{ab} T_j^{cd} H^e + f_1^{ij} T_i^{ab} T_j^{cd} \frac{\Sigma_f^e}{\Lambda} H^f + f_2^{ij} T_i^{ab} T_j^{cd} H^d \frac{\Sigma_f^e}{\Lambda} \right)$$

$$\langle \Sigma \rangle = v_\Sigma \text{diag} (2, 2, 2, -3, -3)$$

Supersymmetric $SU(5)$: Superpotential

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$$+ \frac{1}{4} \epsilon_{abcde} \left(Y_{10}^{ij} T_i^{ab} T_j^{cd} H^e + f_1^{ij} T_i^{ab} T_j^{cd} \frac{\Sigma_f^e}{\Lambda} H^f + f_2^{ij} T_i^{ab} T_j^{cf} H^d \frac{\Sigma_f^e}{\Lambda} \right)$$

$$\langle \Sigma \rangle = v_\Sigma \text{diag}(2, 2, 2, -3, -3)$$

$$Y_u = Y_{10} + 3 \frac{v_\Sigma}{\Lambda} f_1^S + \frac{1}{4} \frac{v_\Sigma}{\Lambda} (3f_2^S + 5f_2^A)$$

$$Y_d = Y_5 - 3 \frac{v_\Sigma}{\Lambda} h_1 + 2 \frac{v_\Sigma}{\Lambda} h_2$$

$$Y_e = Y_5 - 3 \frac{v_\Sigma}{\Lambda} h_1 - 3 \frac{v_\Sigma}{\Lambda} h_2$$

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$$+ \frac{1}{4} \epsilon_{abcde} \left(Y_{10}^{ij} T_i^{ab} T_j^{cd} H^e + f_1^{ij} T_i^{ab} T_j^{cd} \frac{\Sigma_f^e}{\Lambda} H^f + f_2^{ij} T_i^{ab} T_j^{cd} H^d \frac{\Sigma_f^e}{\Lambda} \right)$$

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NR corrections

Supersymmetric $SU(5)$: Superpotential

$$\frac{1}{\Lambda^2}$$

$$\frac{1}{\Lambda^2}$$

$$\frac{1}{\Lambda^2}$$

$$\frac{1}{\Lambda^2}$$

As a further natural step, one can also consider additional NR operators

$$\frac{1}{\Lambda^2}$$

$$\frac{1}{\Lambda^2}$$

$$\frac{1}{\Lambda^2}$$

$$\frac{1}{\Lambda^2}$$

Supersymmetric $SU(5)$: Superpotential

$$W_{Y_u}^{(2)} = \frac{1}{4} \epsilon_{abcde} \left(k_1 T^{ab} T^{cd} H^e \frac{\text{Tr} \Sigma^2}{\Lambda^2} + k_2 T^{ap} T^{bf} H^c \frac{\sum_p^d \sum_f^e}{\Lambda^2} + \right. \\ \left. + k_3 T^{ab} T^{cd} H^p \frac{\sum_f^d \sum_p^e}{\Lambda^2} + k_4 T^{ab} T^{nm} H^c \frac{\sum_n^d \sum_m^e}{\Lambda^2} + k_5 T^{ab} T^{cg} H^d \frac{\sum_g^f \sum_f^e}{\Lambda^2} \right)$$

$$W_{Y_{de}}^{(2)} = \sqrt{2} \left(g_1 \bar{H}_a T^{ab} F_b \frac{\text{Tr} \Sigma^2}{\Lambda^2} + g_2 \bar{H}_a \frac{\sum_b^a \sum_c^b}{\Lambda^2} T^{cd} F_d + \right. \\ \left. + g_3 \bar{H}_a T^{ba} \frac{\sum_b^c \sum_c^d}{\Lambda^2} F_d + g_4 \bar{H}_a \frac{\sum_b^a \sum_c^d}{\Lambda^2} T^{bc} F_d \right)$$

LiE code

Supersymmetric $SU(5)$: Neutrinos

C. Biggio and L. Calibbi

[1007.3750]-JHEP

$$W_{Y_\nu} \propto F_a \Sigma_{Fb}^a H^b \rightarrow \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & m_1 \end{pmatrix}$$

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$$W_{Y_\nu} \propto F_a \Sigma_{Fb}^a H^b \rightarrow \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & m_1 \end{pmatrix} \quad \times$$

Second viable mass

$$\begin{aligned} W_{Y_\nu}^{NR} = & n_1 F_a \Sigma_{Fb}^a \frac{\Sigma_c^b}{\Lambda} H^c + n_2 F_a \frac{\Sigma_b^a}{\Lambda} \Sigma_{Fc}^b H^c + n_3 F_a H^a \Sigma_{Fd}^c \frac{\Sigma_c^d}{\Lambda} + \\ & + n_4 F_a \Sigma_{Fb}^a H^b \frac{Tr \Sigma^2}{\Lambda^2} + n_5 F_a \Sigma_{Fb}^a \frac{\Sigma_c^b \Sigma_d^c}{\Lambda^2} H^d + n_6 F_a \frac{\Sigma_b^a \Sigma_d^c}{\Lambda^2} \Sigma_{Fc}^b H^d + \\ & + n_7 F_a \frac{\Sigma_b^a \Sigma_c^b}{\Lambda^2} \Sigma_{Fd}^c H^d + n_8 F_a \frac{\Sigma_b^a \Sigma_c^d}{\Lambda^2} H^b \Sigma_{Fd}^c + n_9 F_a H^a \frac{\Sigma_c^b \Sigma_d^c}{\Lambda^2} \Sigma_{Fb}^d \end{aligned}$$

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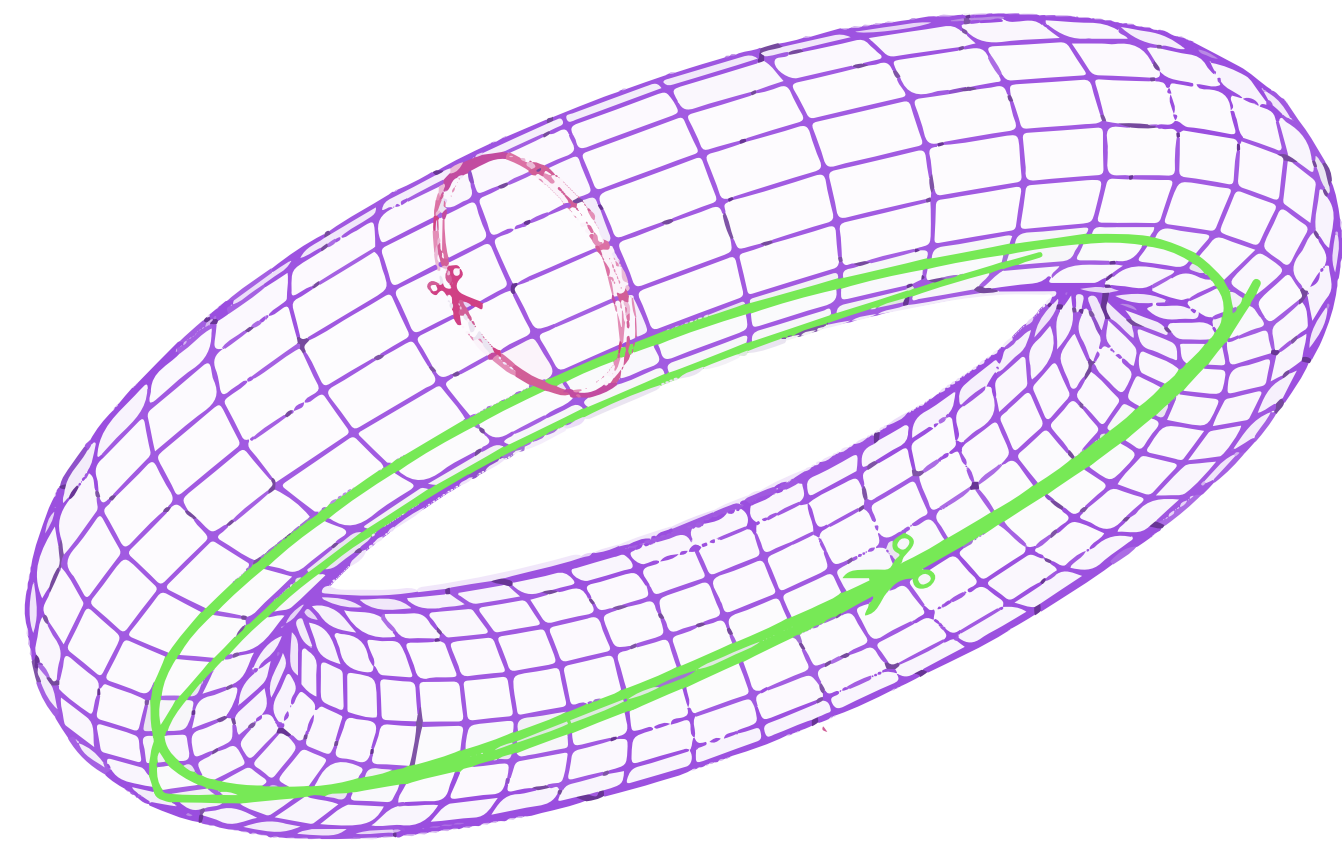
Second viable mass

$$W_{Y_\nu}^{NR} = n_1 F_a \Sigma_{Fb}^a \frac{\Sigma_c^b}{\Lambda} H^c + n_2 F_a \frac{\Sigma_b^a}{\Lambda} \Sigma_{Fc}^b H^c + n_3 F_a H^a \Sigma_{Fd}^c \frac{\Sigma_c^d}{\Lambda} +$$

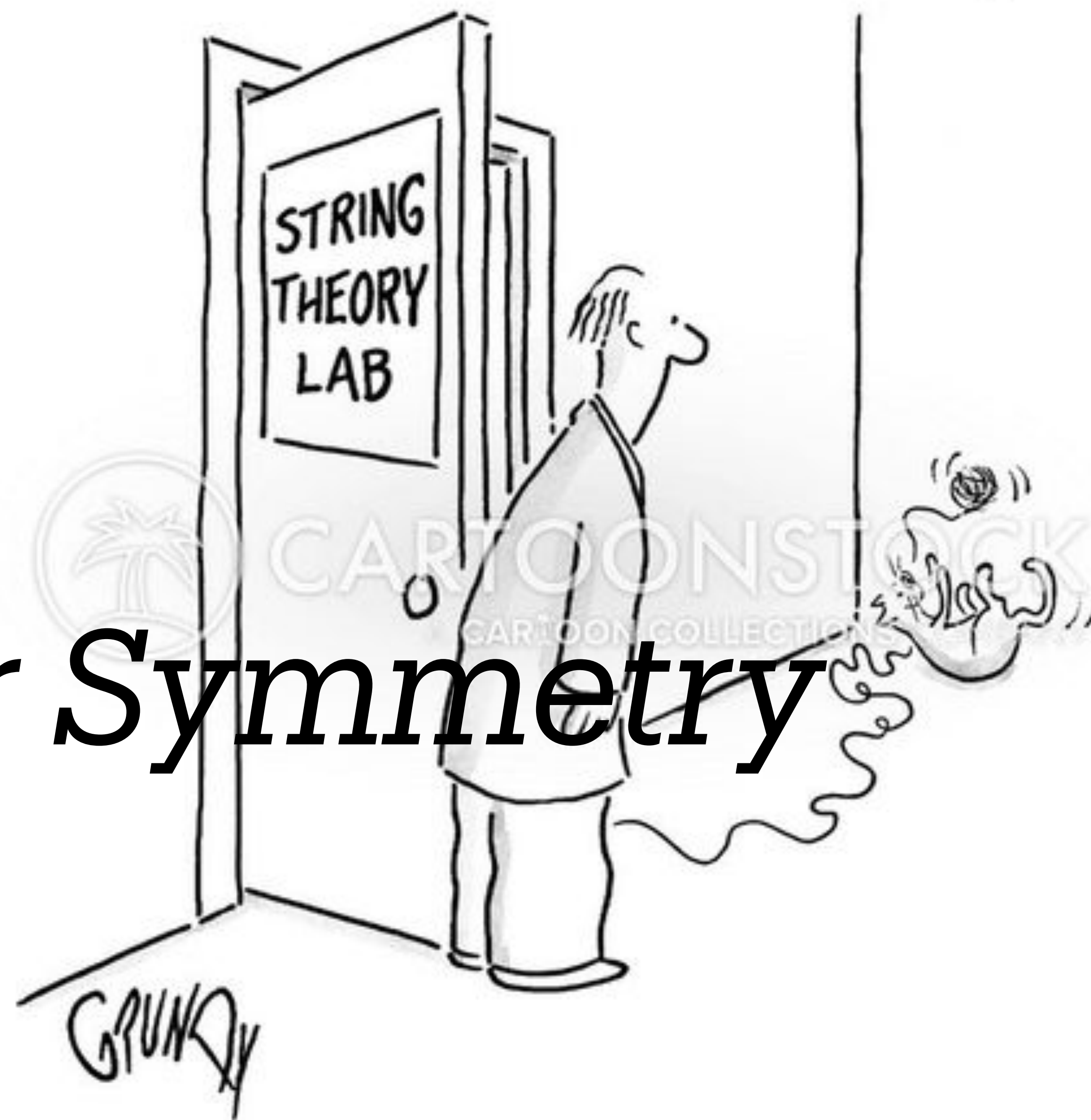
$$+ n_4 F_a \Sigma_{Fb}^a H^b \frac{Tr \Sigma^2}{\Lambda^2} + n_5 F_a \Sigma_{Fb}^a \frac{\Sigma_c^b \Sigma_d^c}{\Lambda^2} H^d + n_6 F_a \frac{\Sigma_b^a \Sigma_d^c}{\Lambda^2} \Sigma_{Fc}^b H^d +$$

$$+ n_7 F_a \frac{\Sigma_b^a \Sigma_c^b}{\Lambda^2} \Sigma_{Fd}^c H^d + n_8 F_a \frac{\Sigma_b^a \Sigma_c^d}{\Lambda^2} H^b \Sigma_{Fd}^c + n_9 F_a H^a \frac{\Sigma_c^b \Sigma_d^c}{\Lambda^2} \Sigma_{Fb}^d$$

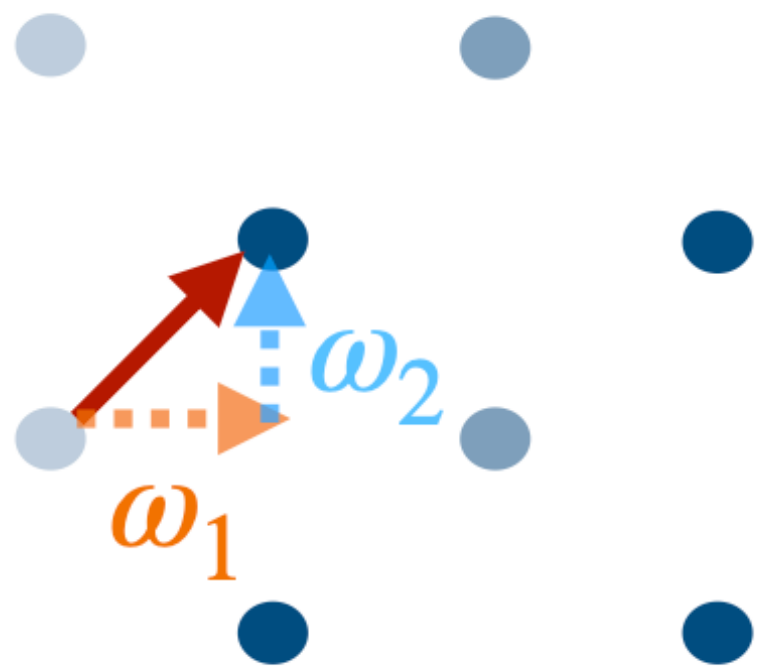
NR



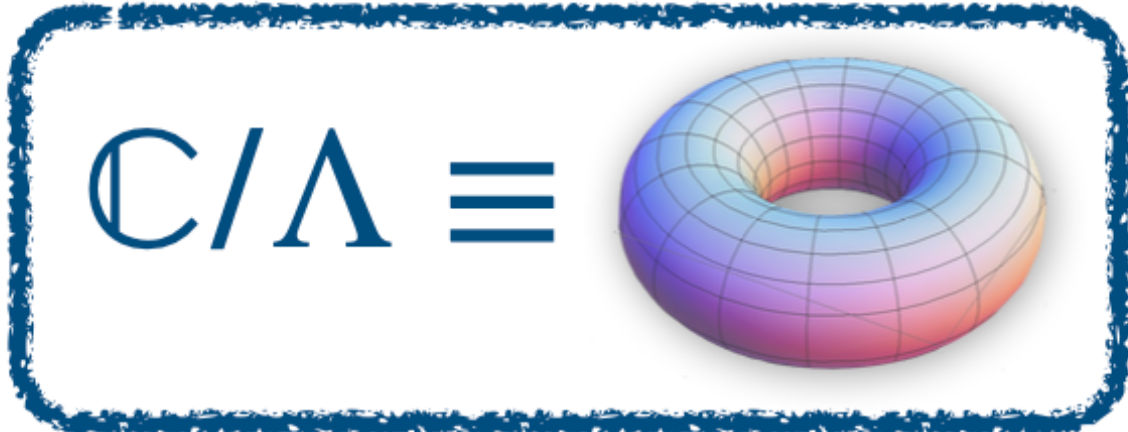
Modular Symmetry



Modular symmetry: Lattice

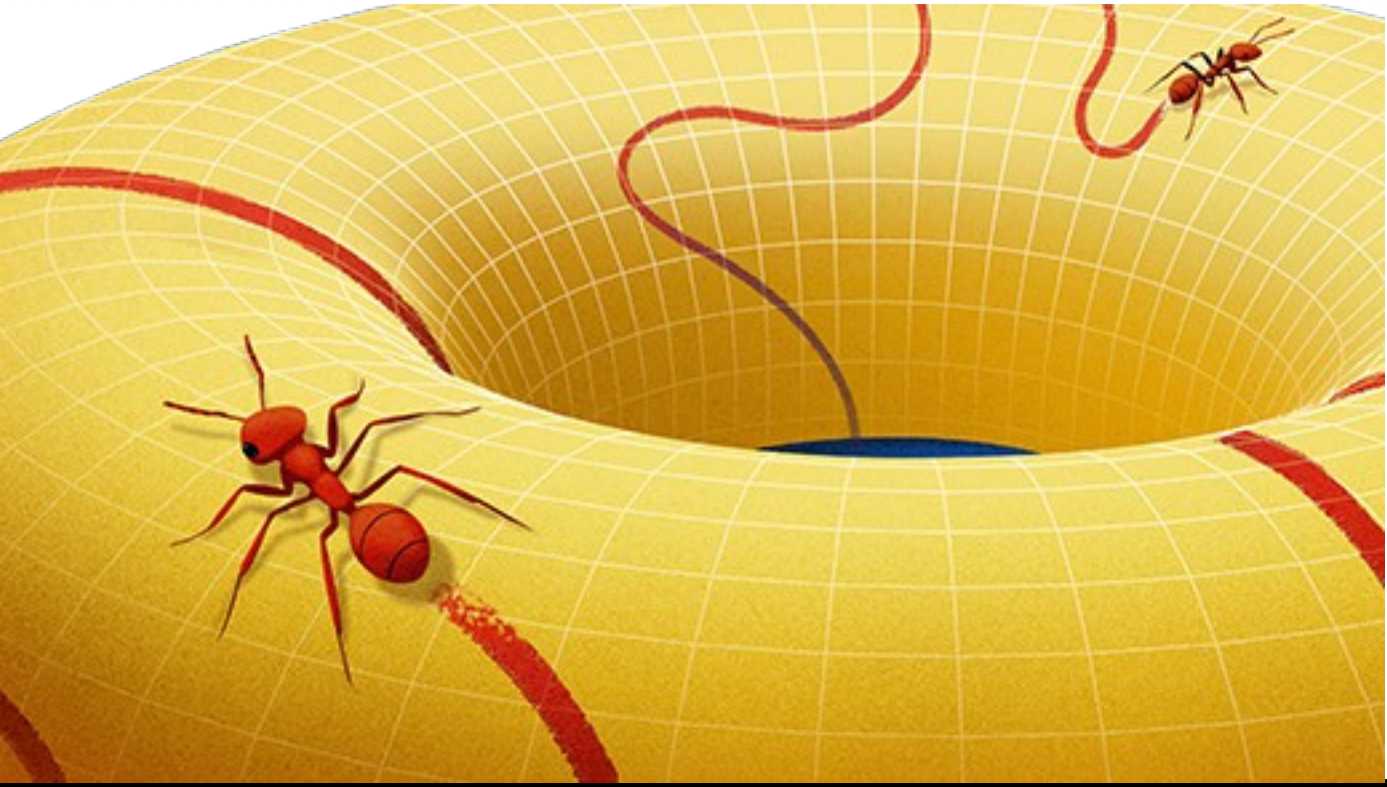
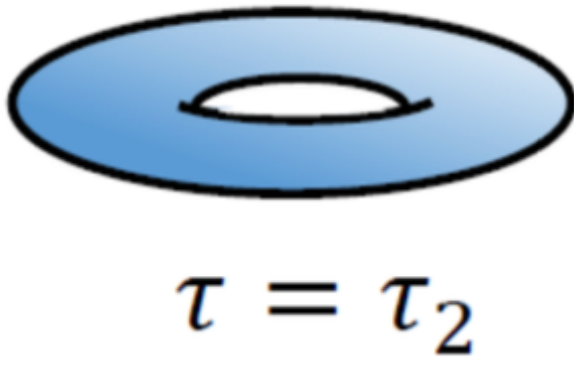
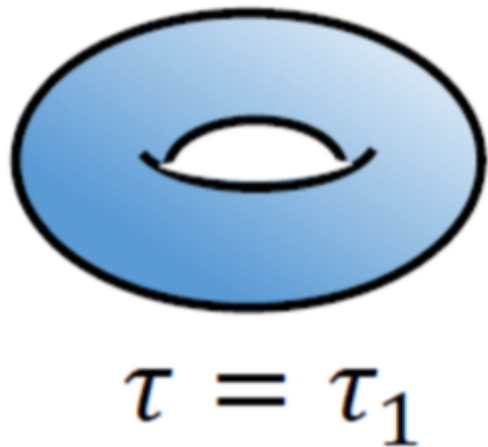


Discrete lattice Λ
 $\omega_1, \omega_2 \in \mathbb{C}$



$$\tau \equiv \frac{\omega_2}{\omega_1}$$

Modulus



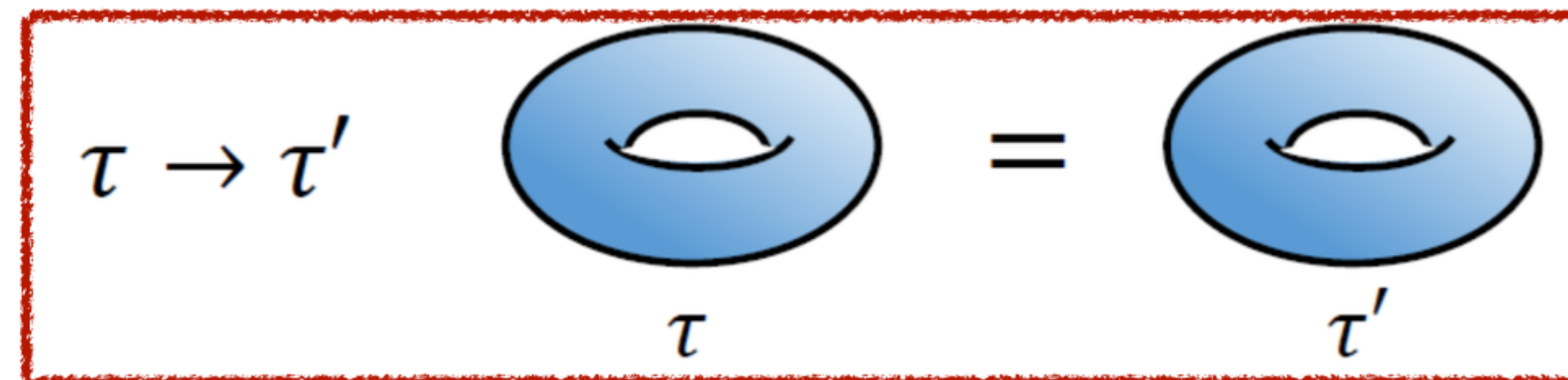
Modular symmetry: Modular Transf.

Change basis?

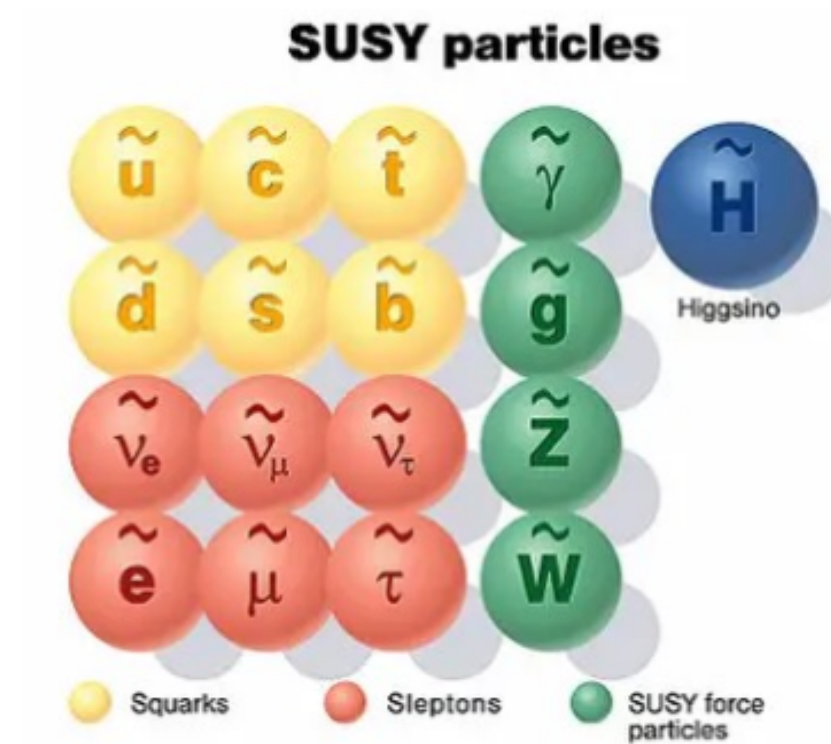
$$\begin{pmatrix} \omega'_2 \\ \omega'_1 \end{pmatrix} = \gamma \begin{pmatrix} \omega_2 \\ \omega_1 \end{pmatrix} \equiv \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \omega_2 \\ \omega_1 \end{pmatrix} \quad a, b, c, d \in \mathbb{Z}$$

$\gamma \in \text{SL}(2, \mathbb{Z}) \equiv \text{Modular group} \equiv \Gamma$

$$\tau \equiv \frac{\omega_2}{\omega_1} \xrightarrow{\text{SL}(2, \mathbb{Z})} \tau' = \frac{a \omega_2 + b \omega_1}{c \omega_2 + d \omega_1} = \frac{a\tau + b}{c\tau + d}$$



Superfields transformations



$$\left\{ \begin{array}{l} \tau \rightarrow \gamma(\tau) = \frac{a\tau + b}{c\tau + d} \quad \gamma \in \Gamma \\ \varphi^{(I)} \rightarrow (c\tau + d)^{-k_I} \rho^{(I)}(\gamma) \varphi^{(I)} \end{array} \right.$$

Usual matter fields

Unitary irrep. of $\Gamma_N \subset \Gamma$

$N = 1, 2, 3, \dots$ "Level"

Finite modular group

$$\Gamma_N$$

for $N \leq 5$ isomorphic to

$$N = 2$$

$$S_3$$

$$N = 3$$

$$A_4$$

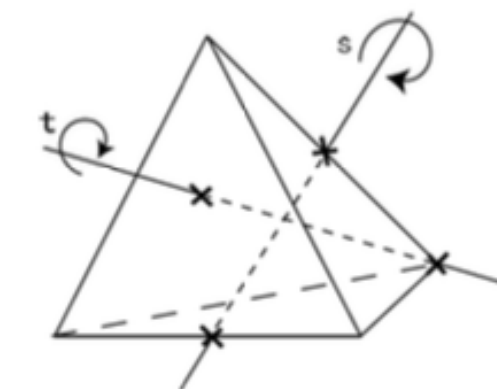
$$N = 4$$

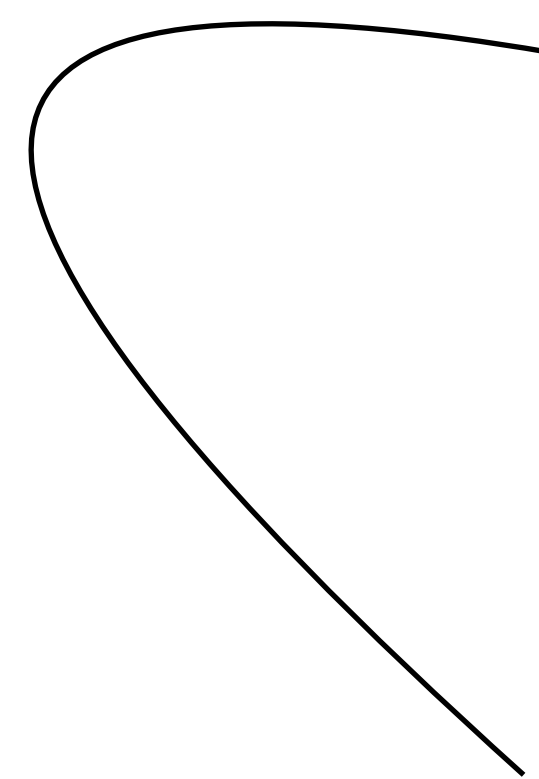
$$S_4$$

$$N = 5$$

$$A_5$$

non-abelian discrete groups




$$Y(\tau) \rightarrow (c\tau + d)^{k_Y} \rho_Y(\gamma) Y(\tau)$$

Yukawa couplings are now promoted to modular forms

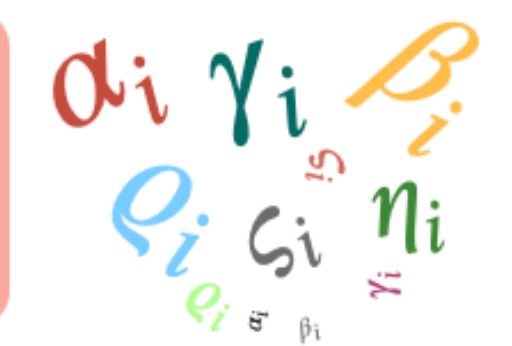
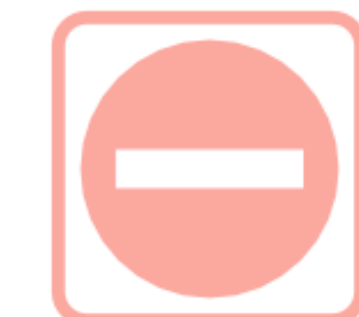
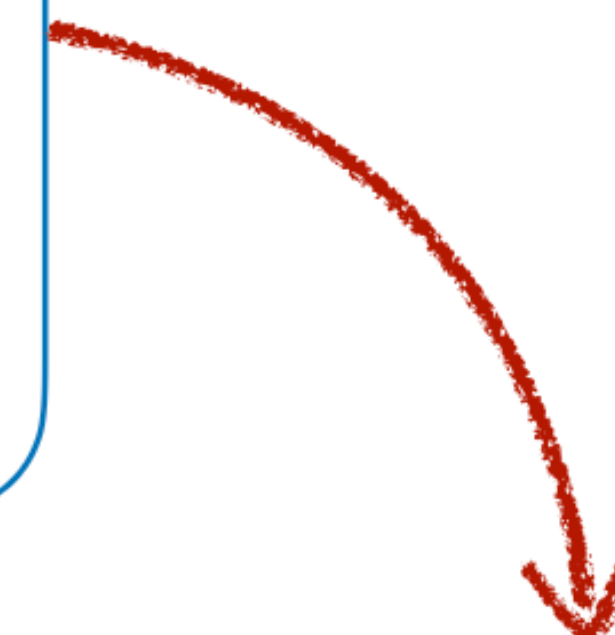
Modular symmetry: gCP

Impose CP symmetry on the model

► $gCP \Rightarrow \alpha_i \in \mathbb{R}$

P. Novichkov, J. Penedo, S. Petcov, A. Titov

Journal of High Energy Physics **2019** no. 7, (Jul, 2019)



Only source of CPV in the model is the VEV of τ

$$\langle \tau \rangle = \text{Re } \tau + i \text{Im } \tau$$

The Model

	T	T_3	F_1	F_2	F_3	H_5	\bar{H}_5	Σ
$SU(5)$	10	10	$\bar{\mathbf{5}}$	$\bar{\mathbf{5}}$	$\bar{\mathbf{5}}$	5	$\bar{\mathbf{5}}$	24
$\Gamma_2 \simeq S_3$	2	$\mathbf{1}'$	$\mathbf{1}'$	1	1	1	1	$\mathbf{1}'$
Weight	-1	6	-1	-3	0	0	0	3

The model: Superpotential

$$\begin{aligned} W_{TT} = & a_1(TT)_2 \frac{\Sigma^2}{\Lambda^2} H_5 Y_2^{(4)} + a_2(TT) \frac{\Sigma^2}{\Lambda^2} H_5 Y_1^{(4)} + a_3(TT_3) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_4(TT_3) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} \\ & + a_5(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_6(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} + a_7(T_3T_3) H_5 Y_1^{(12)} \end{aligned}$$

The model: Superpotential

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 W_{TT} = & a_1(TT)_2 \frac{\Sigma^2}{\Lambda^2} H_5 Y_2^{(4)} + a_2(TT) \frac{\Sigma^2}{\Lambda^2} H_5 Y_1^{(4)} + a_3(TT_3) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_4(TT_3) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} \\
 & + a_5(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_6(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} + a_7(T_3T_3) H_5 Y_1^{(12)}
 \end{aligned}$$

$$\downarrow \begin{pmatrix} \epsilon_{2 \times 2}^2 & \epsilon \\ \epsilon & 1 \end{pmatrix}, \quad \epsilon \equiv \frac{v_\Sigma}{\Lambda}$$

$$Y_u \propto \begin{pmatrix} -\frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_1 + \frac{v_\Sigma^2}{\Lambda^2} a_2 Y_{\mathbf{1}}^{(4)} & \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_2 & \frac{v_\Sigma}{\Lambda} a_3 (Y_{\mathbf{2},1}^{(8)})_1 + \frac{v_\Sigma}{\Lambda} a_4 (Y_{\mathbf{2},2}^{(8)})_1 \\ \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_2 & \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_1 + \frac{v_\Sigma^2}{\Lambda^2} a_2 Y_{\mathbf{1}}^{(4)} & \frac{v_\Sigma}{\Lambda} a_3 (Y_{\mathbf{2},1}^{(8)})_2 + \frac{v_\Sigma}{\Lambda} a_4 (Y_{\mathbf{2},2}^{(8)})_2 \\ \frac{v_\Sigma}{\Lambda} a_5 (Y_{\mathbf{2},1}^{(8)})_1 + \frac{v_\Sigma}{\Lambda} a_6 (Y_{\mathbf{2},2}^{(8)})_1 & \frac{v_\Sigma}{\Lambda} a_5 (Y_{\mathbf{2},1}^{(8)})_2 + \frac{v_\Sigma}{\Lambda} a_6 (Y_{\mathbf{2},2}^{(8)})_2 & a_7 Y_{\mathbf{1}}^{(12)} \end{pmatrix}^T$$

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 & + a_5(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_6(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} + a_7(T_3T_3) H_5 Y_1^{(12)}
 \end{aligned}$$

$$\downarrow \left(\begin{array}{c} \epsilon_{2 \times 2}^2 \\ \epsilon \end{array} \begin{array}{c} \epsilon \\ \diamond \\ 1 \end{array} \right), \quad \epsilon \equiv \frac{v_\Sigma}{\Lambda}$$

$$Y_u \propto \left(\begin{array}{ccc} -\frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_1 + \frac{v_\Sigma^2}{\Lambda^2} a_2 Y_{\mathbf{1}}^{(4)} & \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_2 & \frac{v_\Sigma}{\Lambda} a_3 (Y_{\mathbf{2},1}^{(8)})_1 + \frac{v_\Sigma}{\Lambda} a_4 (Y_{\mathbf{2},2}^{(8)})_1 \\ \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_2 & \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_1 + \frac{v_\Sigma^2}{\Lambda^2} a_2 Y_{\mathbf{1}}^{(4)} & \frac{v_\Sigma}{\Lambda} a_3 (Y_{\mathbf{2},1}^{(8)})_2 + \frac{v_\Sigma}{\Lambda} a_4 (Y_{\mathbf{2},2}^{(8)})_2 \\ \frac{v_\Sigma}{\Lambda} a_5 (Y_{\mathbf{2},1}^{(8)})_1 + \frac{v_\Sigma}{\Lambda} a_6 (Y_{\mathbf{2},2}^{(8)})_1 & \frac{v_\Sigma}{\Lambda} a_5 (Y_{\mathbf{2},1}^{(8)})_2 + \frac{v_\Sigma}{\Lambda} a_6 (Y_{\mathbf{2},2}^{(8)})_2 & \diamond a_7 Y_{\mathbf{1}}^{(12)} \end{array} \right)^T$$

O(1)

The model: Superpotential

$$\begin{aligned}
 W_{TT} = & a_1(TT)_2 \frac{\Sigma^2}{\Lambda^2} H_5 Y_2^{(4)} + a_2(TT) \frac{\Sigma^2}{\Lambda^2} H_5 Y_1^{(4)} + a_3(TT_3) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_4(TT_3) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} \\
 & + a_5(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,1}^{(8)} + a_6(T_3T) \frac{\Sigma}{\Lambda} H_5 Y_{2,2}^{(8)} + a_7(T_3T_3) H_5 Y_1^{(12)}
 \end{aligned}$$

Rank one \downarrow $\begin{pmatrix} \epsilon_{2 \times 2}^2 & \epsilon \\ \epsilon & 1 \end{pmatrix}, \quad \epsilon \equiv \frac{v_\Sigma}{\Lambda}$

$$Y_u \propto \begin{pmatrix} -\frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_1 + \frac{v_\Sigma^2}{\Lambda^2} a_2 Y_{\mathbf{1}}^{(4)} & \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_2 & \frac{v_\Sigma}{\Lambda} a_3 (Y_{\mathbf{2},1}^{(8)})_1 + \frac{v_\Sigma}{\Lambda} a_4 (Y_{\mathbf{2},2}^{(8)})_1 \\ \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_2 & \frac{v_\Sigma^2}{\Lambda^2} a_1 (Y_{\mathbf{2}}^{(4)})_1 + \frac{v_\Sigma^2}{\Lambda^2} a_2 Y_{\mathbf{1}}^{(4)} & \frac{v_\Sigma}{\Lambda} a_3 (Y_{\mathbf{2},1}^{(8)})_2 + \frac{v_\Sigma}{\Lambda} a_4 (Y_{\mathbf{2},2}^{(8)})_2 \\ \frac{v_\Sigma}{\Lambda} a_5 (Y_{\mathbf{2},1}^{(8)})_1 + \frac{v_\Sigma}{\Lambda} a_6 (Y_{\mathbf{2},2}^{(8)})_1 & \frac{v_\Sigma}{\Lambda} a_5 (Y_{\mathbf{2},1}^{(8)})_2 + \frac{v_\Sigma}{\Lambda} a_6 (Y_{\mathbf{2},2}^{(8)})_2 & a_7 Y_{\mathbf{1}}^{(12)} \end{pmatrix}^T$$

The model: Superpotential

$$W_{TF} = B_1(TF_1) \frac{\Sigma^2}{\Lambda^2} \bar{H}_5 Y_2^{(4)} + B_2(TF_2) \frac{\Sigma^2}{\Lambda^2} \bar{H}_5 Y_2^{(2)} + B_3(TF_3) \frac{\Sigma}{\Lambda} \bar{H}_5 Y_2^{(2)} + \\ + B_4(T_3 F_2) \frac{\Sigma}{\Lambda} \bar{H}_5 Y_1^{(6)} + B_5(T_3 F_3) \bar{H}_5 Y_1^{(6)}$$

$$W_\nu = D_1(F_1 \Sigma) \frac{\Sigma}{\Lambda} H_5 Y_1^{(8)} + D_2(F_2 \Sigma) \frac{\Sigma^2}{\Lambda^2} H_5 Y_1^{(6)} + D_3(F_3 \Sigma) \frac{\Sigma}{\Lambda} H_5 Y_1^{(6)}$$

The model: Superpotential

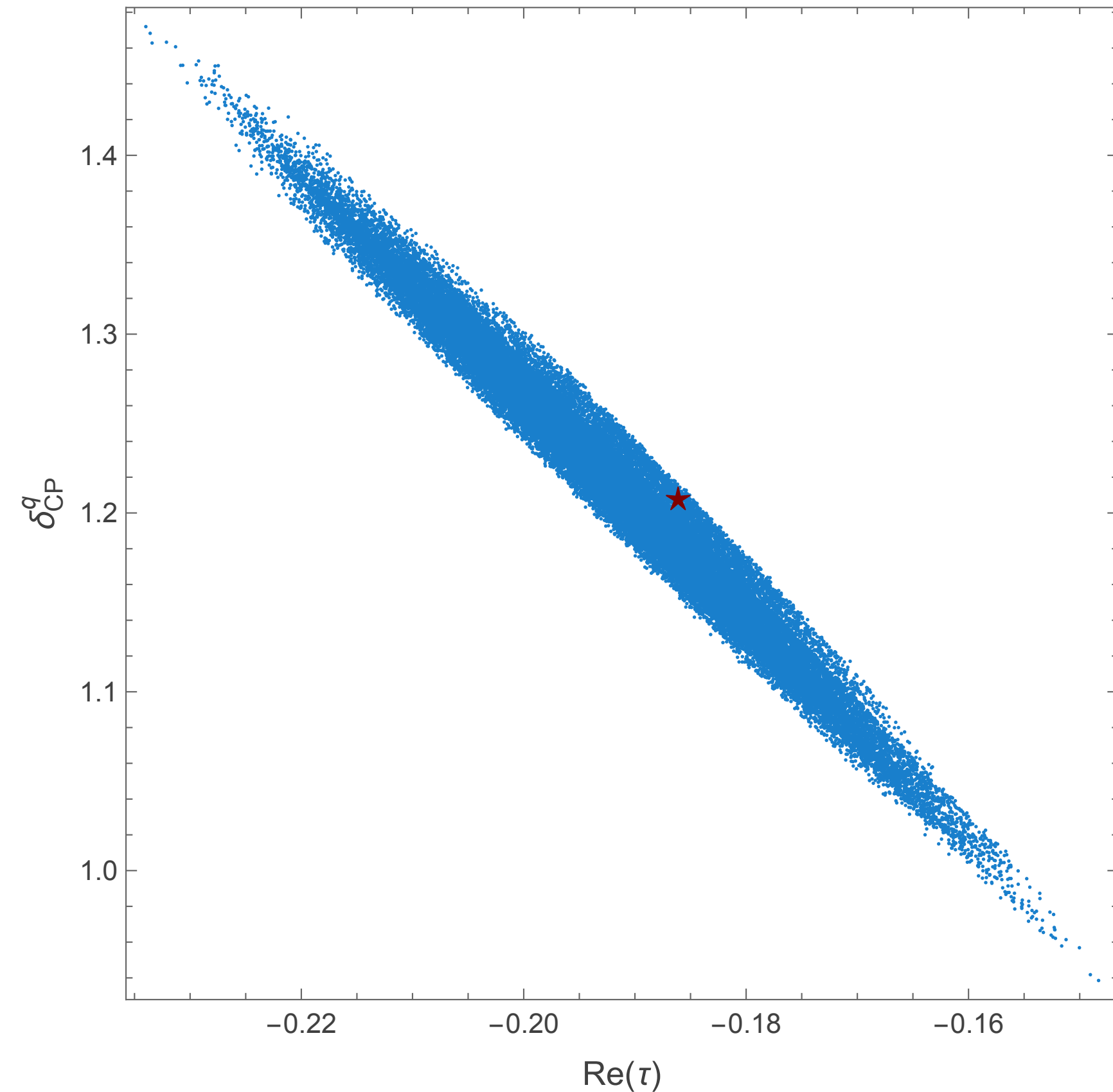
$$\begin{aligned}
 W_{TF} = & B_1(TF_1) \frac{\Sigma^2}{\Lambda^2} \bar{H}_5 Y_2^{(4)} + B_2(TF_2) \frac{\Sigma^2}{\Lambda^2} \bar{H}_5 Y_2^{(2)} + B_3(TF_3) \frac{\Sigma}{\Lambda} \bar{H}_5 Y_2^{(2)} + \\
 & + B_4(T_3 F_2) \frac{\Sigma}{\Lambda} \bar{H}_5 Y_1^{(6)} + B_5(T_3 F_3) \bar{H}_5 Y_1^{(6)}
 \end{aligned}
 \longrightarrow
 Y_d \propto \begin{pmatrix} \frac{v_\Sigma^2}{\Lambda^2} b_1 (Y_2^{(4)})_2 & \frac{v_\Sigma^2}{\Lambda^2} b_2 (Y_2^{(2)})_1 & \frac{v_\Sigma}{\Lambda} b_3 (Y_2^{(2)})_2 \\ -\frac{v_\Sigma^2}{\Lambda^2} b_1 (Y_2^{(4)})_1 & \frac{v_\Sigma^2}{\Lambda^2} b_2 (Y_2^{(2)})_2 & -\frac{v_\Sigma}{\Lambda} b_3 (Y_2^{(2)})_1 \\ \diamond 0 & \frac{v_\Sigma}{\Lambda} b_4 Y_1^{(6)} & b_5 Y_1^{(6)} \end{pmatrix}^T$$

Modular Zero

$$W_\nu = D_1(F_1 \Sigma) \frac{\Sigma}{\Lambda} H_5 Y_1^{(8)} + D_2(F_2 \Sigma) \frac{\Sigma^2}{\Lambda^2} H_5 Y_1^{(6)} + D_3(F_3 \Sigma) \frac{\Sigma}{\Lambda} H_5 Y_1^{(6)} \longrightarrow m_\nu \propto Y_\nu^* M^{-1} Y_\nu^\dagger$$

$M = \text{diag}(G_{N^c}, G_{\Sigma_F^c})$
Type I+III seesaw

The model: Numerical Procedure



$$\chi^2 = 0.363$$

Deformed Bottom-Tau unification

$$y_b/y_\tau \sim 0.789$$

Correlation region in the quark sector between the CP-violating phase δ_{CP}^q and the real part of τ , which constitutes the sole source of CP violation

The model: Outlook

First successful gCP implementation in a Modular-GUT scenario

The model: Outlook

First successful gCP implementation in a Modular-GUT scenario

MORE TO BE DONE...

- Top-Down \leftrightarrow Bottom-Up Unification
- Moduli Stabilization
- Choice of Modular Weights

Thank You



Roma tre Neutrino Theory Group