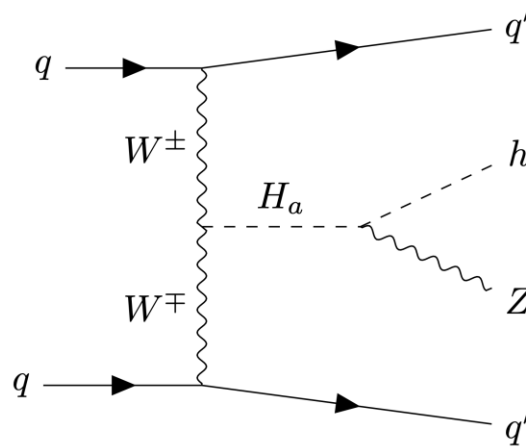
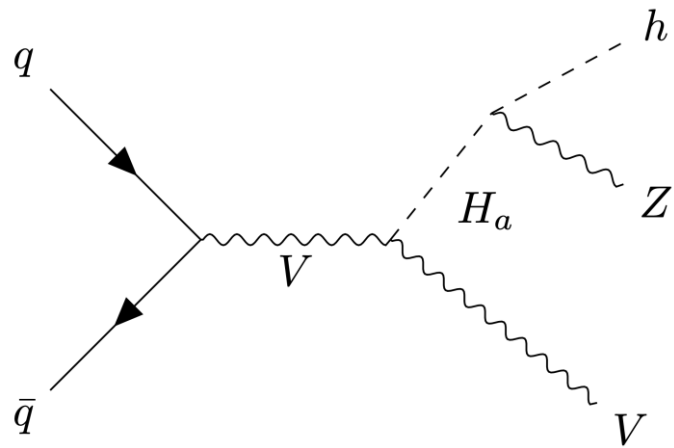




V-associated production & vector boson fusion as an LHC signature of CP violation

Álvaro Lozano Onrubia

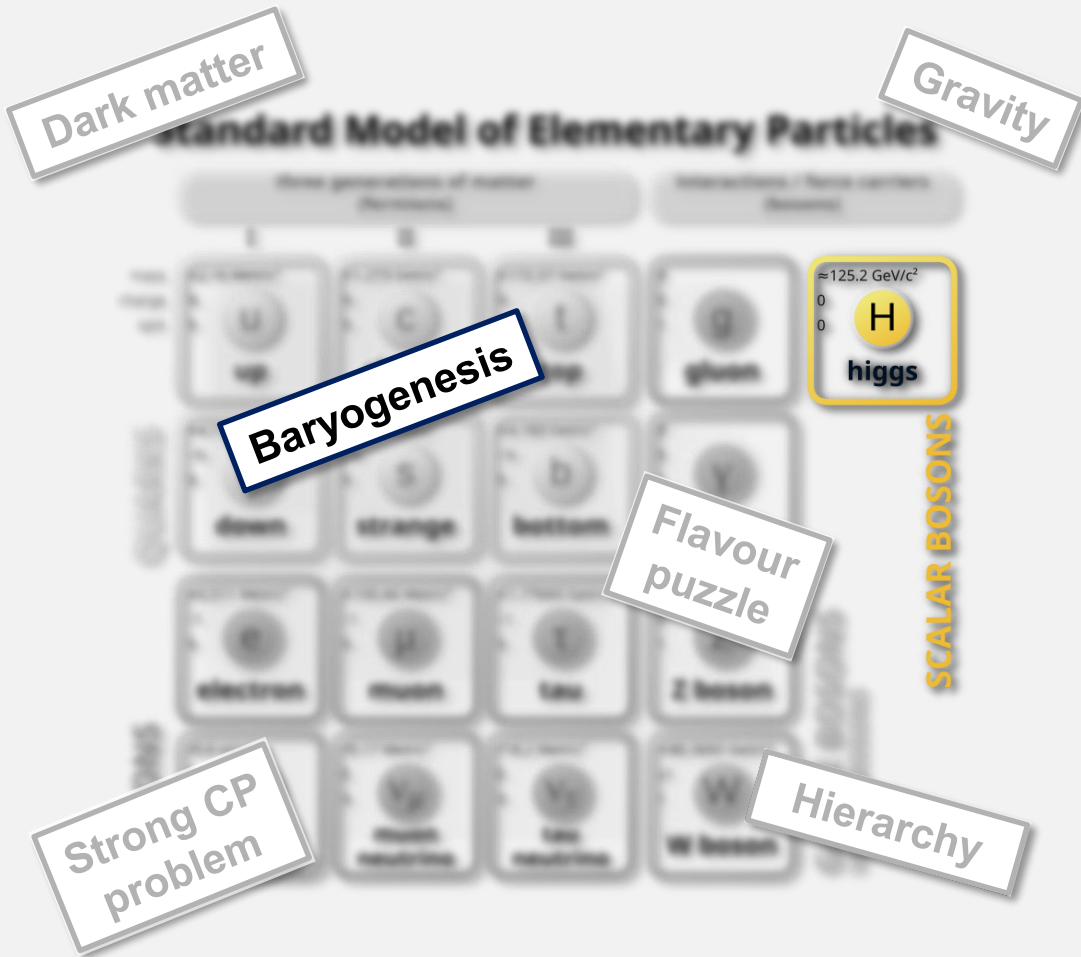
with Rodrigo Capucha, Luca Merlo, José Miguel No, Rui Santos

Based on 2507.05942

PRD 112 (2025) 9, 095005

CP violation... in the Higgs sector?

A personal and non-exhaustive take



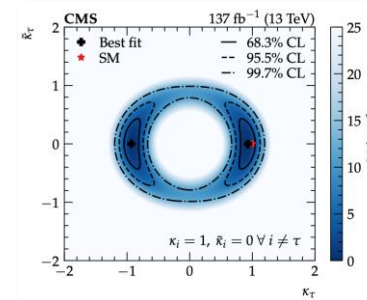
Matter-antimatter asymmetry

CP violation (CPV) in SM insufficient!

Gavela, M. B. et al. (1994). Nucl. Phys. B, 430(2), 345-381.

Gavela, M. B. et al. (1994). Nucl. Phys. B, 430(2), 382-426.

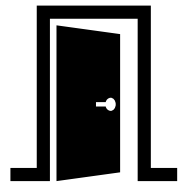
CP properties of SM Higgs



Yukawa & gauge couplings...

... SM-like!

CMS – 2110.04836

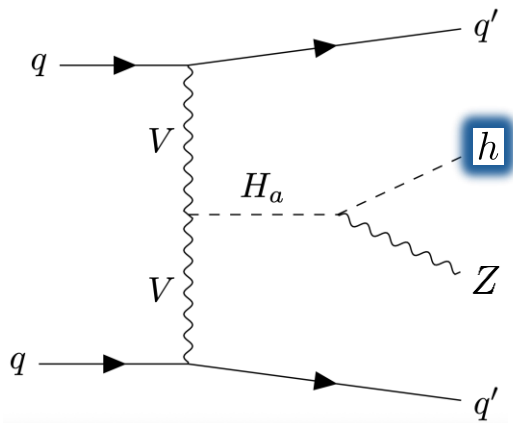


CPV in BSM Higgs sector?

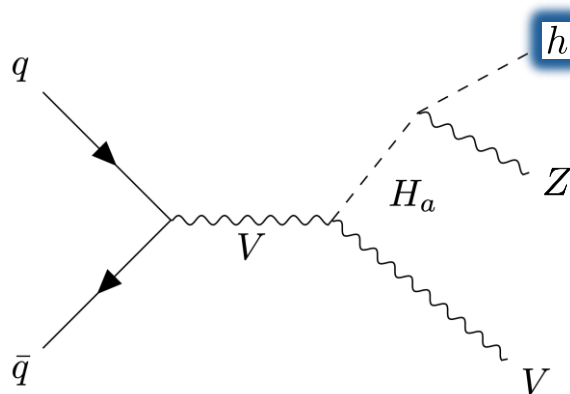
Unlocking the power of the LHC

Searches of CPV in extended scalar sectors

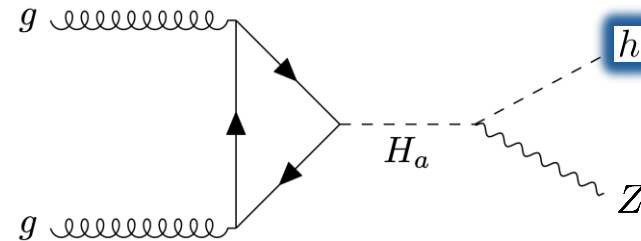
Vector boson fusion (VBF)



V-associated production (V-AP)



Gluon fusion assisted (ggF)

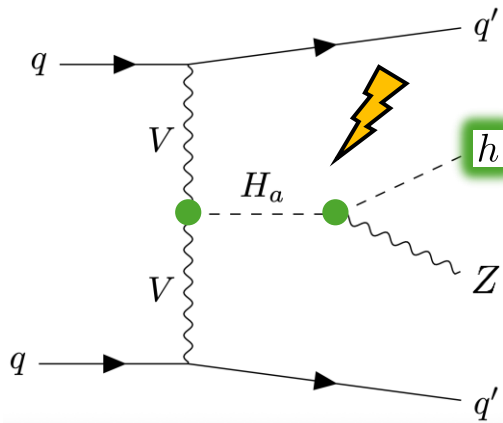


Throughout the talk:
 h is the SM-like, a priori CP -even Higgs with $m_h = 125$ GeV

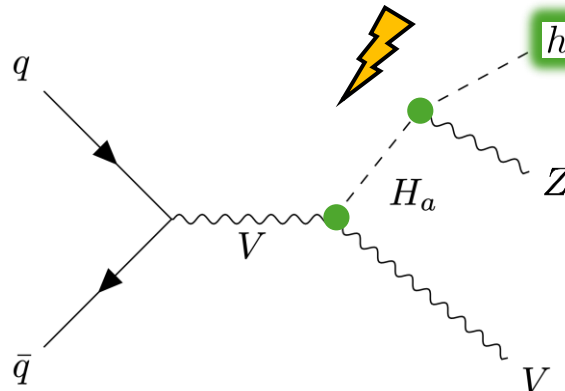
Unlocking the power of the LHC

Searches of *CPV* in extended scalar sectors

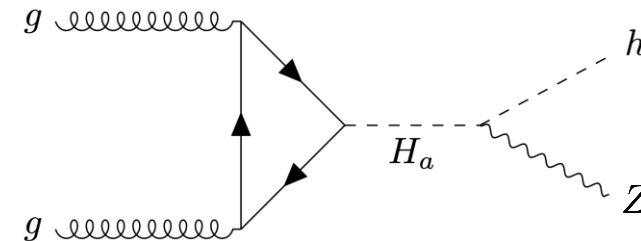
Vector boson fusion (VBF)



V-associated production (V-AP)



Gluon fusion assisted (ggF)



1) $H_a \rightarrow VV$: H_a is **CP-even**! (provided loop-effects are small)

2) $H_a \rightarrow Zh$: **opposite CP assignments** in renormalizable theory

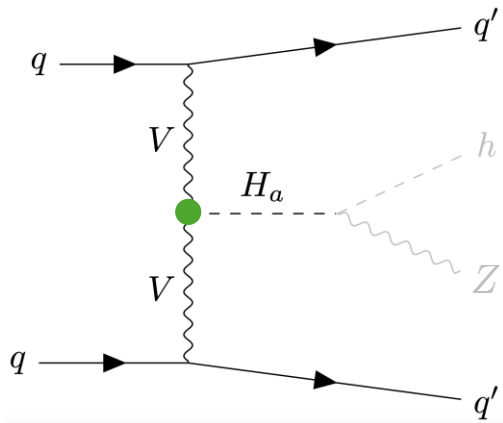
3) h is a priori **CP-even**

➡ **Signals of CPV!**

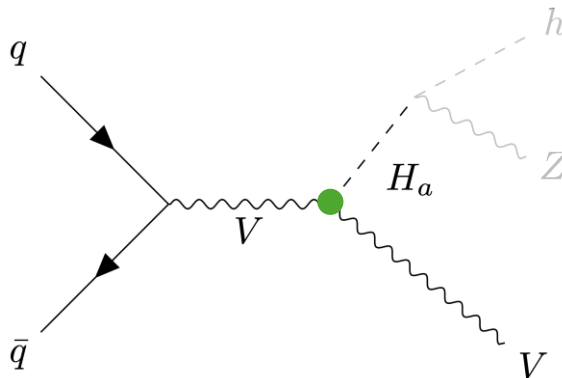
Unlocking the power of the LHC

Searches of CPV in extended scalar sectors

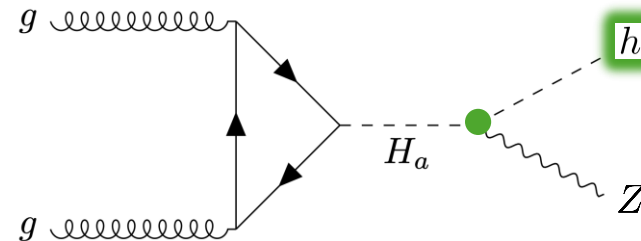
Vector boson fusion (VBF)



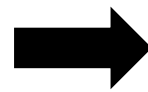
V-associated production (V-AP)



Gluon fusion **assisted** (ggF)



ggF-assisted CPV signal



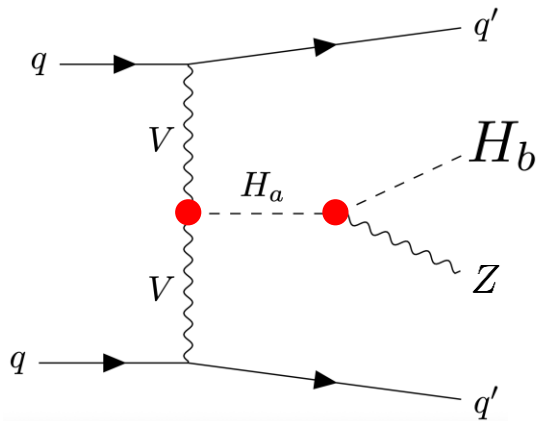
Benefit from **larger cross sections...**

... but need **concurrent pure VBF or V-AP of H_a !**

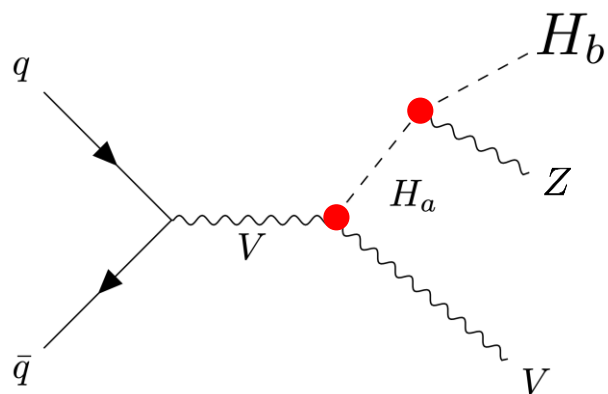
Unlocking the power of the LHC

Venturing beyond the 125 GeV Higgs

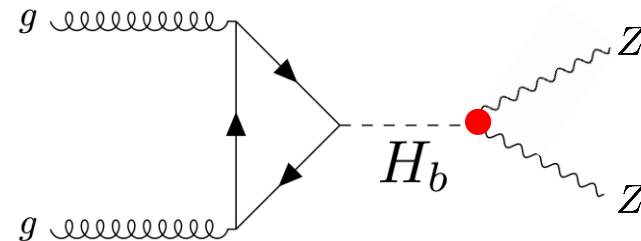
Vector boson fusion (VBF)



V-associated production (V-AP)



Gluon fusion assisted (ggF)



Same spiel with **two BSM scalars** **BUT:**
trade CP knowledge of h **for another process**

Unlocking the power of the LHC

LHC joint signatures of CPV featuring V-AP and/or VBF

(i) $pp \rightarrow VH_a, H_a \rightarrow Zh$

(ii) VBF production of $H_a, H_a \rightarrow Zh$

(iii) $pp \rightarrow VH_a, H_a \rightarrow ZH_b$ as well as $pp \rightarrow VH_b$

(iv) VBF production of $H_a, H_a \rightarrow ZH_b$ as well as VBF of H_b

(v) ggF production of H_a decaying via $H_a \rightarrow ZH_b$, as well as $pp \rightarrow VH_{a,b}$ or VBF of both $H_{a,b}$

(vi) ggF production of H_a decaying both via $H_a \rightarrow ZH_b$ and $H_a \rightarrow VV$ (or $H_a \rightarrow H_bH_b$), as well as $pp \rightarrow VH_b$ (or VBF of H_b)

(vii) ggF production of H_a decaying via $H_a \rightarrow Zh$, as well as $pp \rightarrow VH_a$ (or VBF production of H_a)

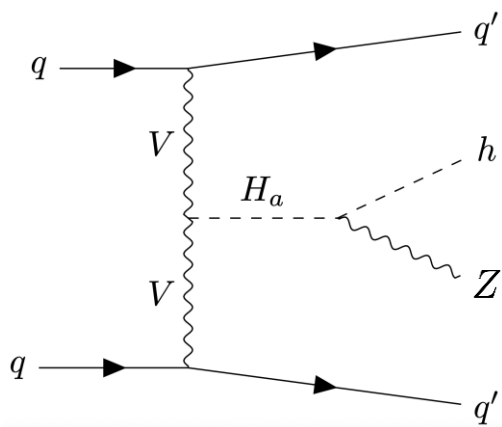
Each set (i) – (vii) is a **self-sufficient joint signature** of CPV.

Is there any set where *all* processes are observable?

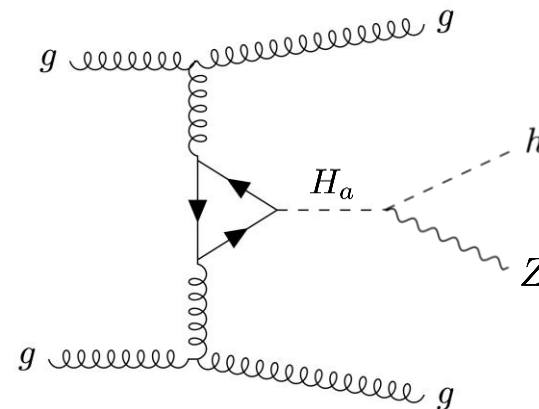
Unlocking the power of the LHC

A few words of caution for *VBF*

Vector boson fusion (*VBF*)



Gluon fusion + 2 gluon jets

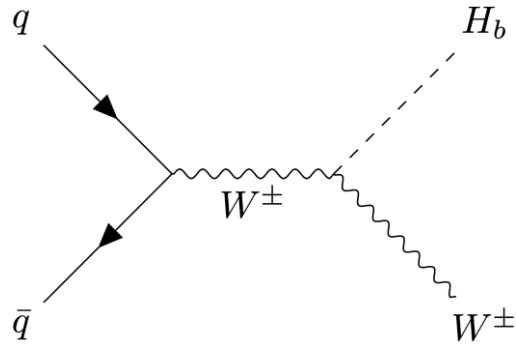


- ⚠ Contamination from ***ggF+jj*** process which **does not fix would-be CP properties of H_a**
- ❓ Possible improvement through **cuts** as in SM... but ***ggF-to-VBF ratio*** unknown *a priori* (BSM)
- ✓ Optimized ***quark-gluon jet tagging*** might solve this

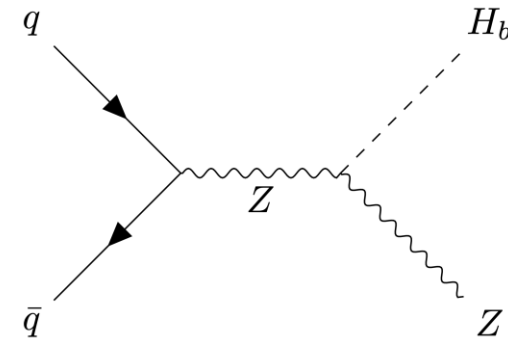
Unlocking the power of the LHC

A few words of caution for V-AP

V-associated production (V-AP)



V-associated production (V-AP)



Avoids background from $ggF \rightarrow$ intermediate **neutral scalars**

Only small background from $qq \rightarrow$ intermediate **charged scalars** (e.g. 2HDMs)

Benchmarking the LHC

Complex Two-Higgs-Doublet Model (C2HDM)

$$V = m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - \left(m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.} \right) + \frac{\lambda_1}{2} |\Phi_1|^4 + \frac{\lambda_2}{2} |\Phi_2|^4 + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 (\Phi_1^\dagger \Phi_2)(\Phi_2^\dagger \Phi_1) + \left[\frac{\lambda_5}{2} (\Phi_1^\dagger \Phi_2)^2 + \text{h.c.} \right]$$

Two $SU(2)$ doublets

$$\Phi_j = \begin{pmatrix} \phi_j^+ \\ \frac{v_j + \rho_j + i\eta_j}{\sqrt{2}} \end{pmatrix} \quad \text{with} \quad \tan \beta \equiv \frac{v_2}{v_1}$$

CPV phases

1 independent phase

$$\delta \equiv \text{Arg}[(m_{12}^2)^2 \lambda_5^*]$$

Scalar spectrum

3 neutral + CP-indefinite

$$H_1, H_2, H_3$$

$$m_{H_1} \leq m_{H_2} \leq m_{H_3}$$

2 charged

$$H^\pm \longrightarrow m_{H^\pm}$$

+ 3 Goldstones

Mass basis

$$R = \begin{pmatrix} c_1 c_2 & s_1 c_2 & s_2 \\ -(c_1 s_2 s_3 + s_1 c_3) & c_1 c_3 - s_1 s_2 s_3 & c_2 s_3 \\ -c_1 s_2 c_3 + s_1 s_3 & -(c_1 s_3 + s_1 s_2 c_3) & c_2 c_3 \end{pmatrix}$$

$$s_i, c_i \equiv \sin \alpha_i, \cos \alpha_i$$

$$-\pi/2 < \alpha_{1,2,3} \leq \pi/2$$

hep-ph/0211371, 1408.2534, 1502.01720, 1711.09419, 2403.02425

Benchmarking the LHC

Complex Two-Higgs-Doublet Model (C2HDM)

$$V = m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - \left(m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.} \right) + \frac{\lambda_1}{2} |\Phi_1|^4 + \frac{\lambda_2}{2} |\Phi_2|^4 + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \left[\frac{\lambda_5}{2} (\Phi_1^\dagger \Phi_2)^2 + \text{h.c.} \right]$$

Angular parameters

$$\tan \beta \equiv \frac{v_2}{v_1}$$

$$s_i, c_i \equiv \sin \alpha_i, \cos \alpha_i$$

Yukawa couplings

	u-type	d-type	leptons
Type I	Φ_2	Φ_2	Φ_2
Type II	Φ_2	Φ_1	Φ_1
Lepton-Specific	Φ_2	Φ_2	Φ_1
Flipped	Φ_2	Φ_1	Φ_2

	u-type	d-type	leptons
Type I	$\frac{R_{12}}{s_\beta} - i \frac{R_{13}}{t_\beta} \gamma_5$	$\frac{R_{12}}{s_\beta} + i \frac{R_{13}}{t_\beta} \gamma_5$	$\frac{R_{12}}{s_\beta} + i \frac{R_{13}}{t_\beta} \gamma_5$
Type II	$\frac{R_{12}}{s_\beta} - i \frac{R_{13}}{t_\beta} \gamma_5$	$\frac{R_{11}}{c_\beta} - i t_\beta R_{13} \gamma_5$	$\frac{R_{11}}{c_\beta} - i t_\beta R_{13} \gamma_5$
Lepton-Specific	$\frac{R_{12}}{s_\beta} - i \frac{R_{13}}{t_\beta} \gamma_5$	$\frac{R_{12}}{s_\beta} + i \frac{R_{13}}{t_\beta} \gamma_5$	$\frac{R_{11}}{c_\beta} - i t_\beta R_{13} \gamma_5$
Flipped	$\frac{R_{12}}{s_\beta} - i \frac{R_{13}}{t_\beta} \gamma_5$	$\frac{R_{11}}{c_\beta} - i t_\beta R_{13} \gamma_5$	$\frac{R_{12}}{s_\beta} + i \frac{R_{13}}{t_\beta} \gamma_5$

1711.09419

Mass basis

$$R = \begin{pmatrix} c_1 c_2 & s_1 c_2 & s_2 \\ -(c_1 s_2 s_3 + s_1 c_3) & c_1 c_3 - s_1 s_2 s_3 & c_2 s_3 \\ -c_1 s_2 c_3 + s_1 s_3 & -(c_1 s_3 + s_1 s_2 c_3) & c_2 c_3 \end{pmatrix}$$

Gauge couplings

$$[H_i, W_\mu^+, W_\nu^-] = i g M_W g_{\mu\nu} [R_{i1} \cos(\beta) + R_{i2} \sin(\beta)] \equiv i g M_W g_{\mu\nu} C_i$$

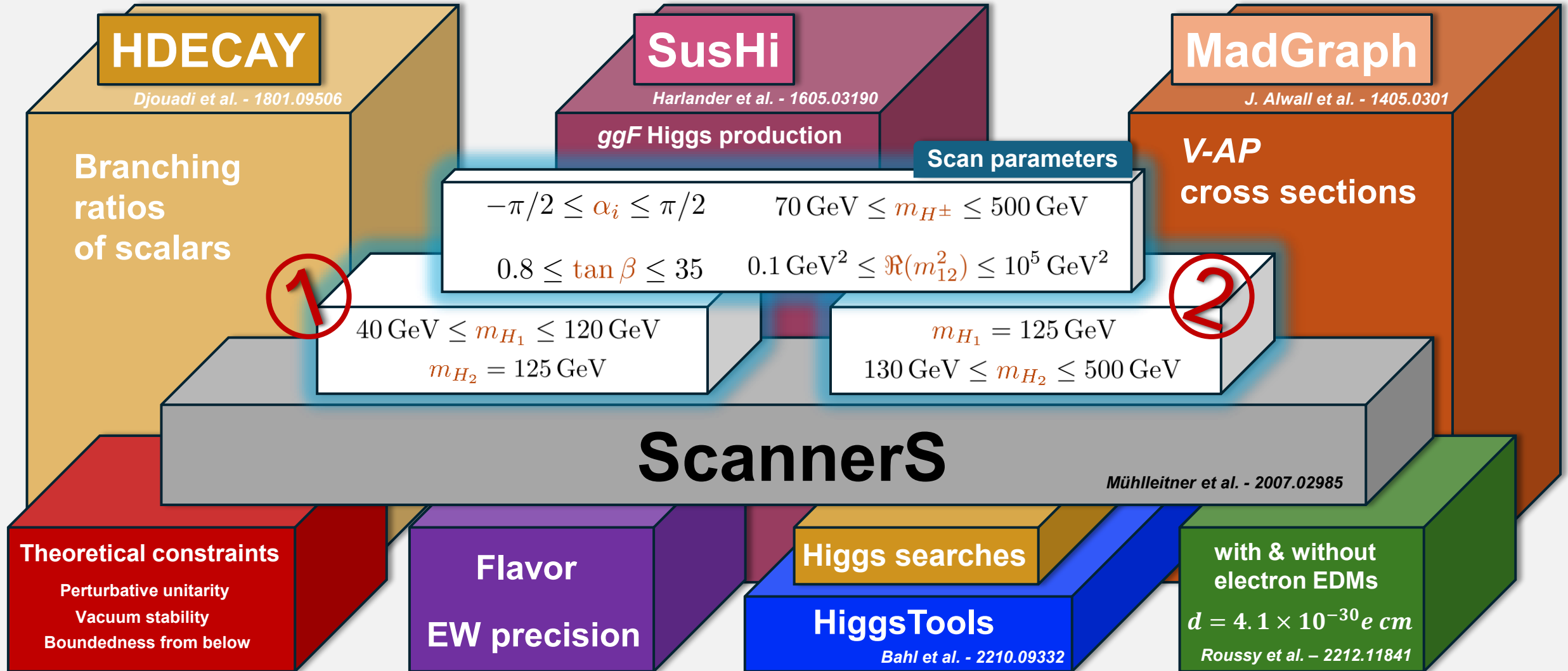
$$[H_i, Z_\mu, Z_\nu] = i \frac{g M_Z}{c_W} g_{\mu\nu} [R_{i1} \cos(\beta) + R_{i2} \sin(\beta)] \equiv i \frac{g M_Z}{c_W} g_{\mu\nu} C_i$$

1711.09419

hep-ph/0211371, 1408.2534, 1502.01720, 1711.09419, 2403.02425

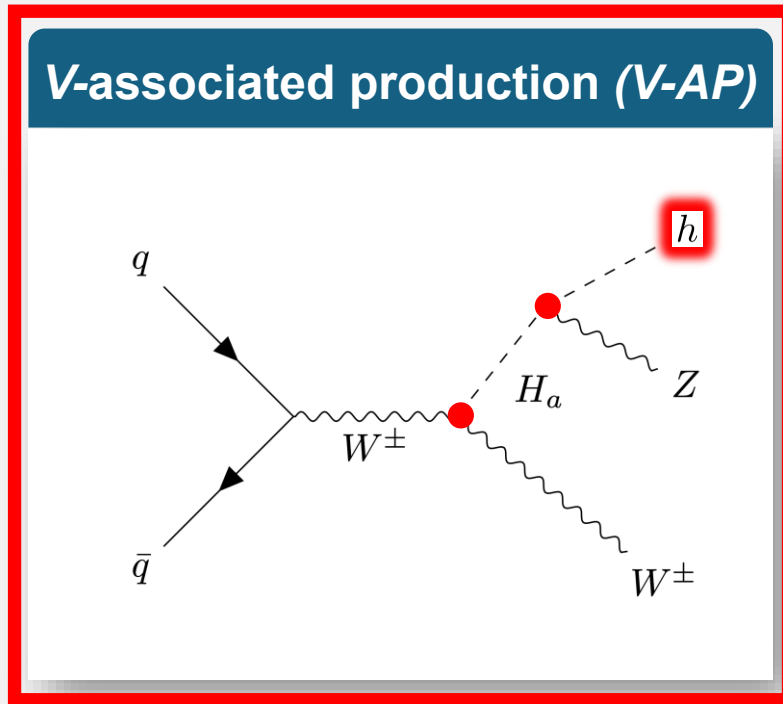
Benchmarking the LHC

Setting up the scan



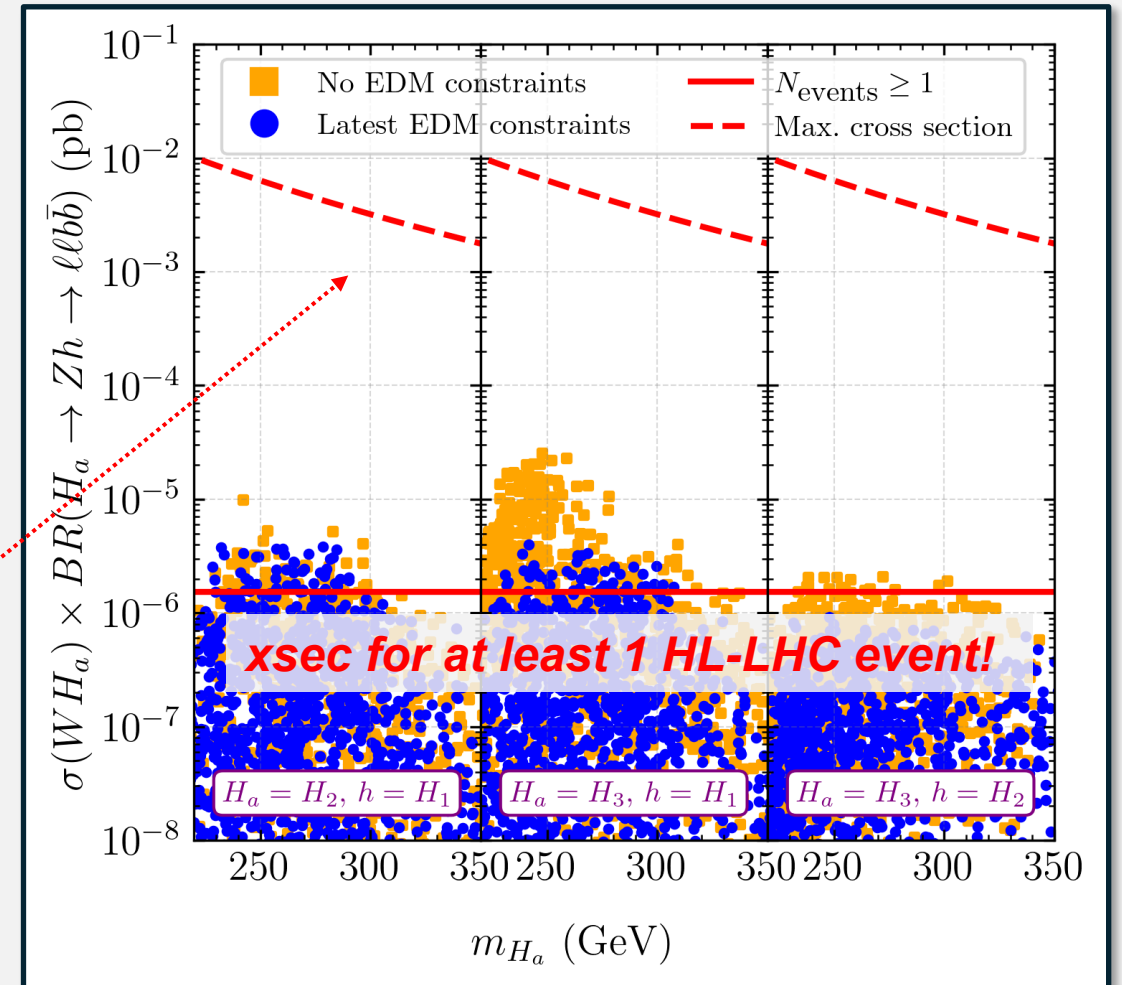
Results

Set (i): Probing CPV with poster child W-AP ?



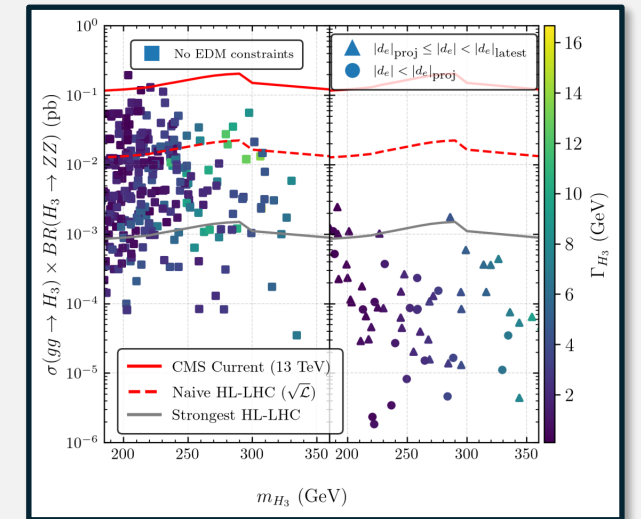
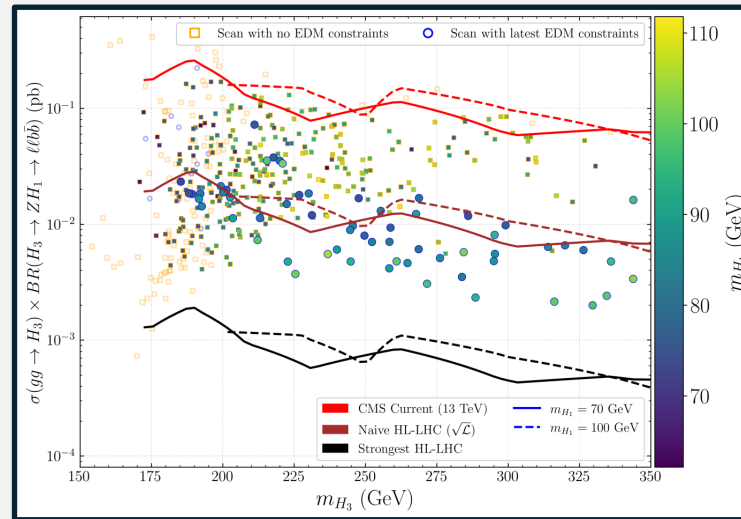
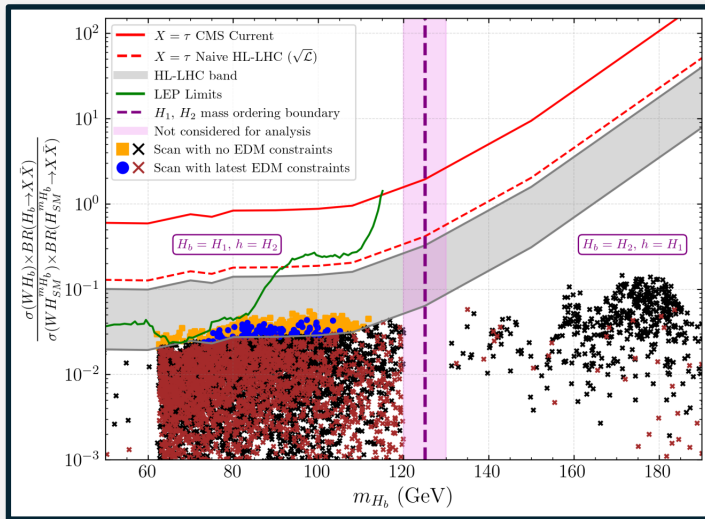
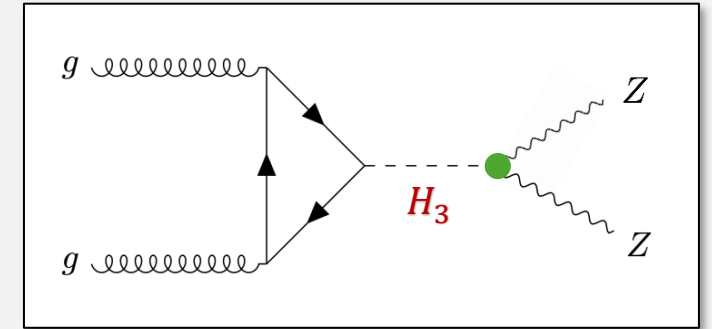
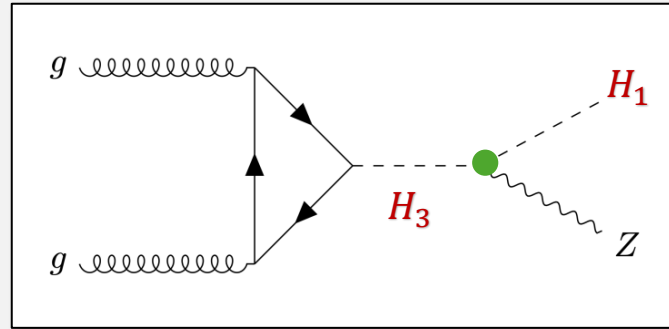
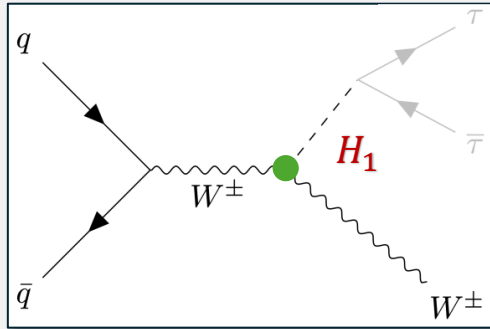
Theoretically maximal cross sections!

$$\left. \begin{array}{l} L_{int} = 3 \text{ ab}^{-1} \\ W \rightarrow \ell \nu \end{array} \right\} \mathcal{O}(1) \text{ events}$$



Results

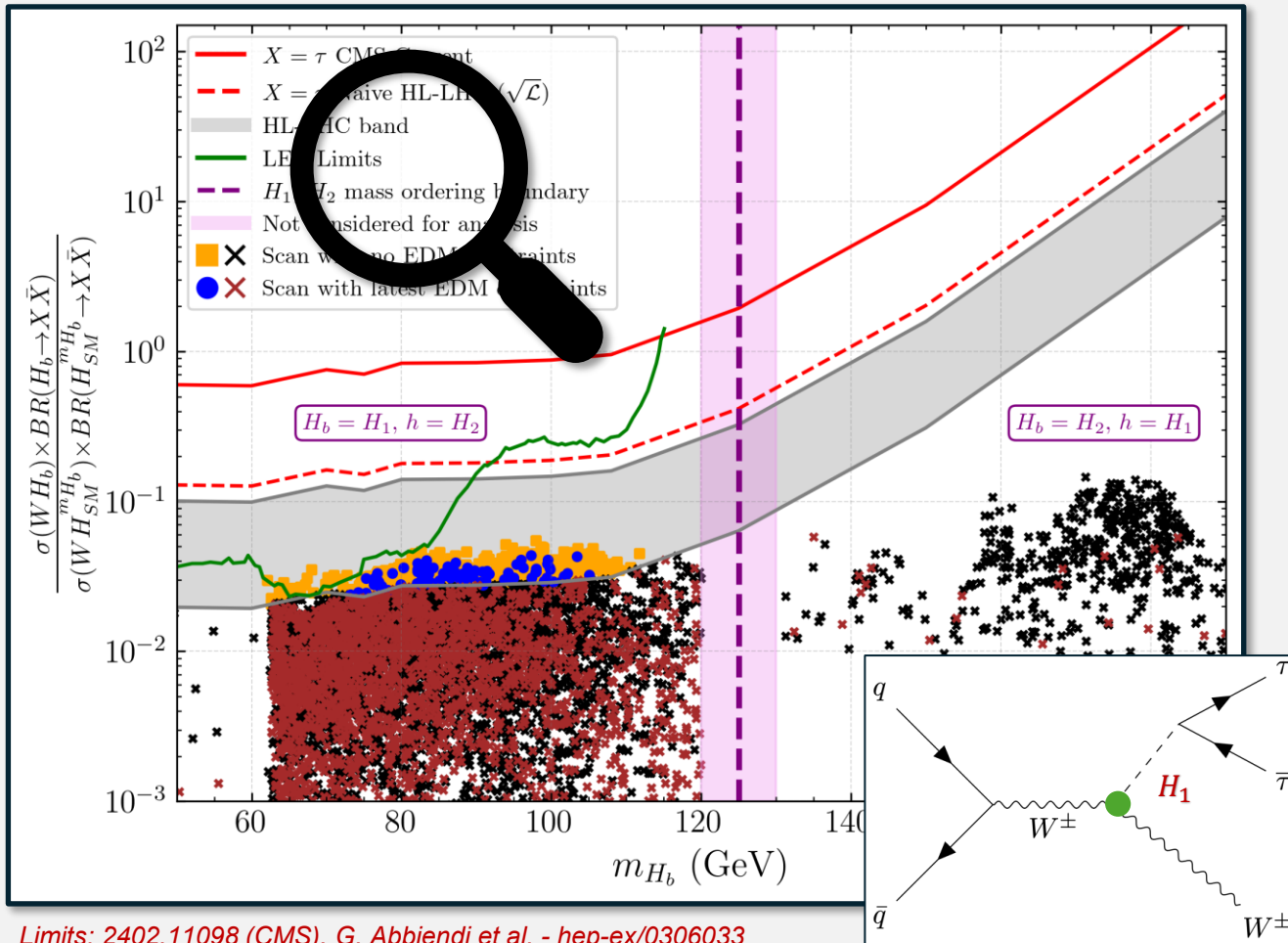
Set (vi): Example of a joint CPV signature



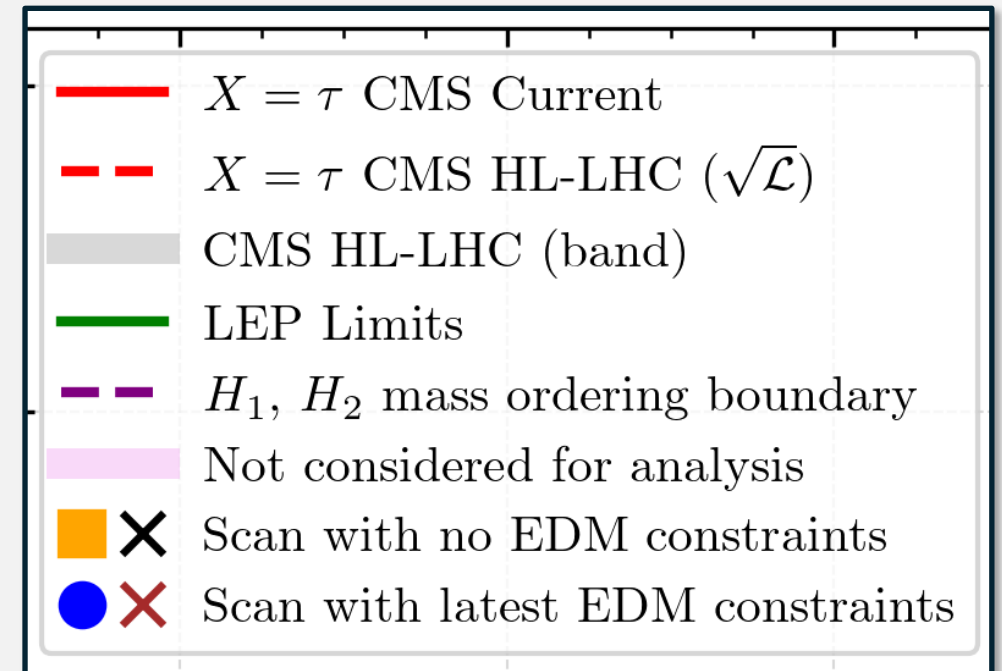
Results

Set (vi): W -AP of lightest non-SM Higgs in the C2HDM

W -AP in C2HDM normalized to W -AP of same-mass SM Higgs



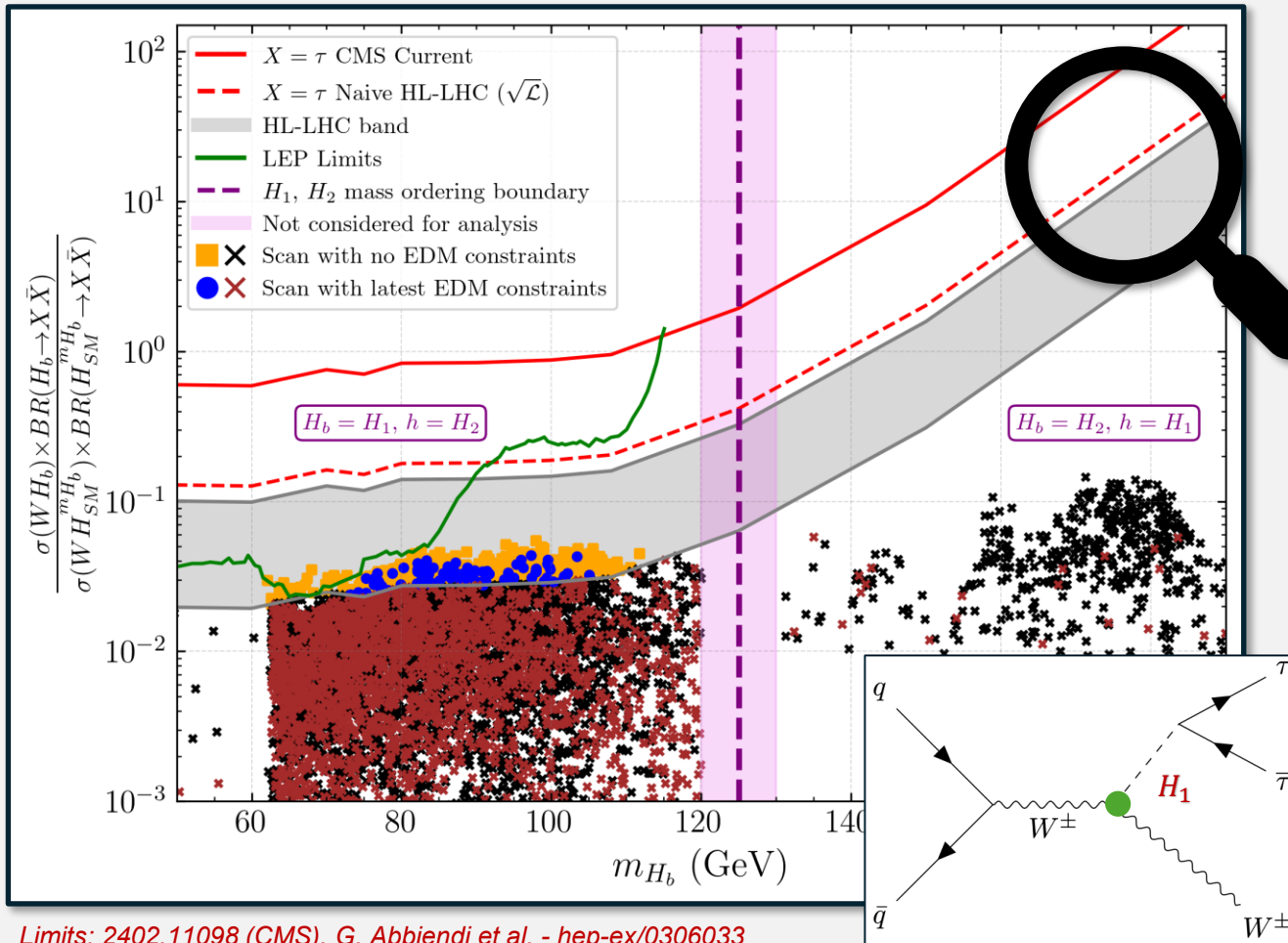
Limits: 2402.11098 (CMS), G. Abbiendi et al. - hep-ex/0306033



Results

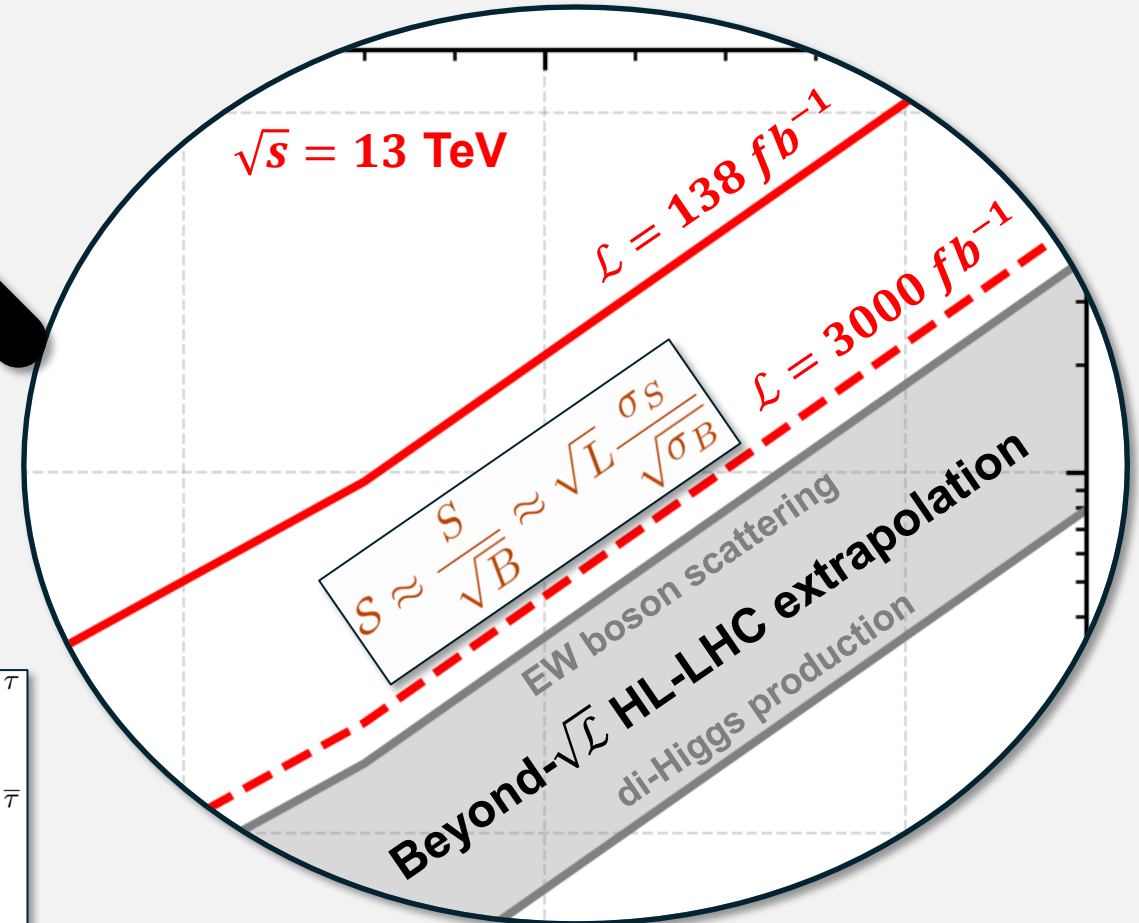
Set (vi): W-AP of lightest non-SM Higgs in the C2HDM

W-AP in C2HDM normalized to W-AP of same-mass SM Higgs



Limits: 2402.11098 (CMS), G. Abbiendi et al. - hep-ex/0306033

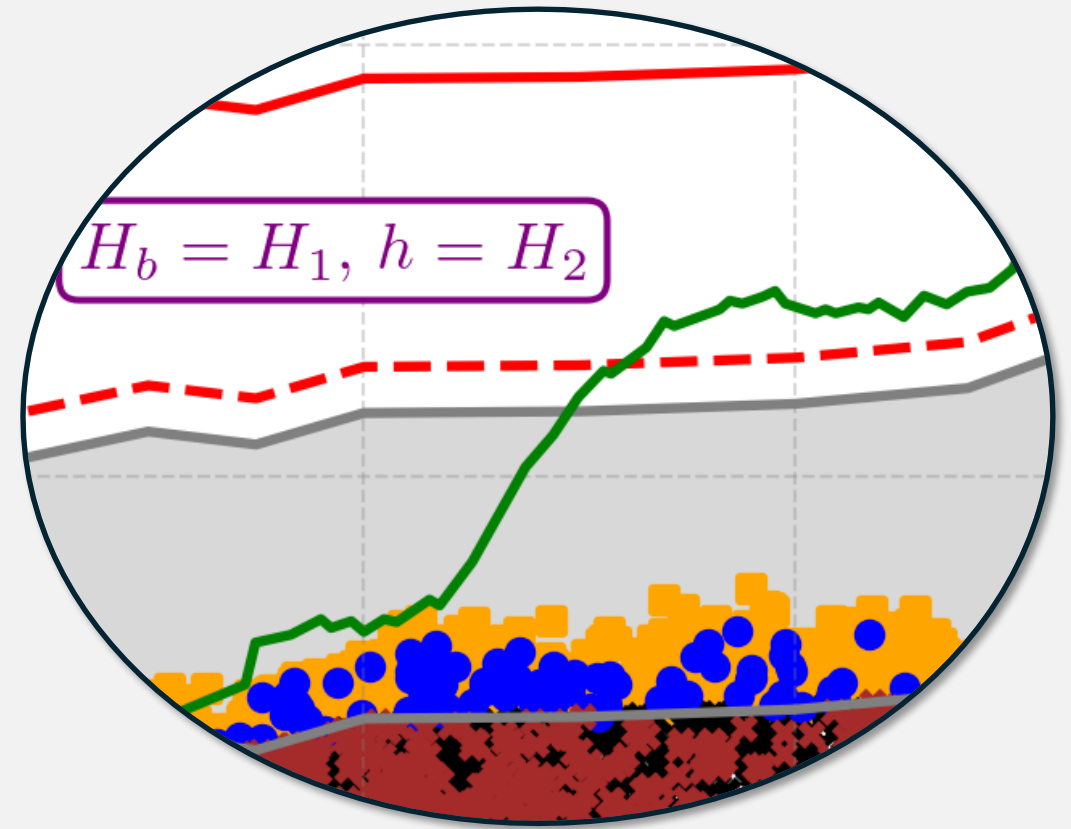
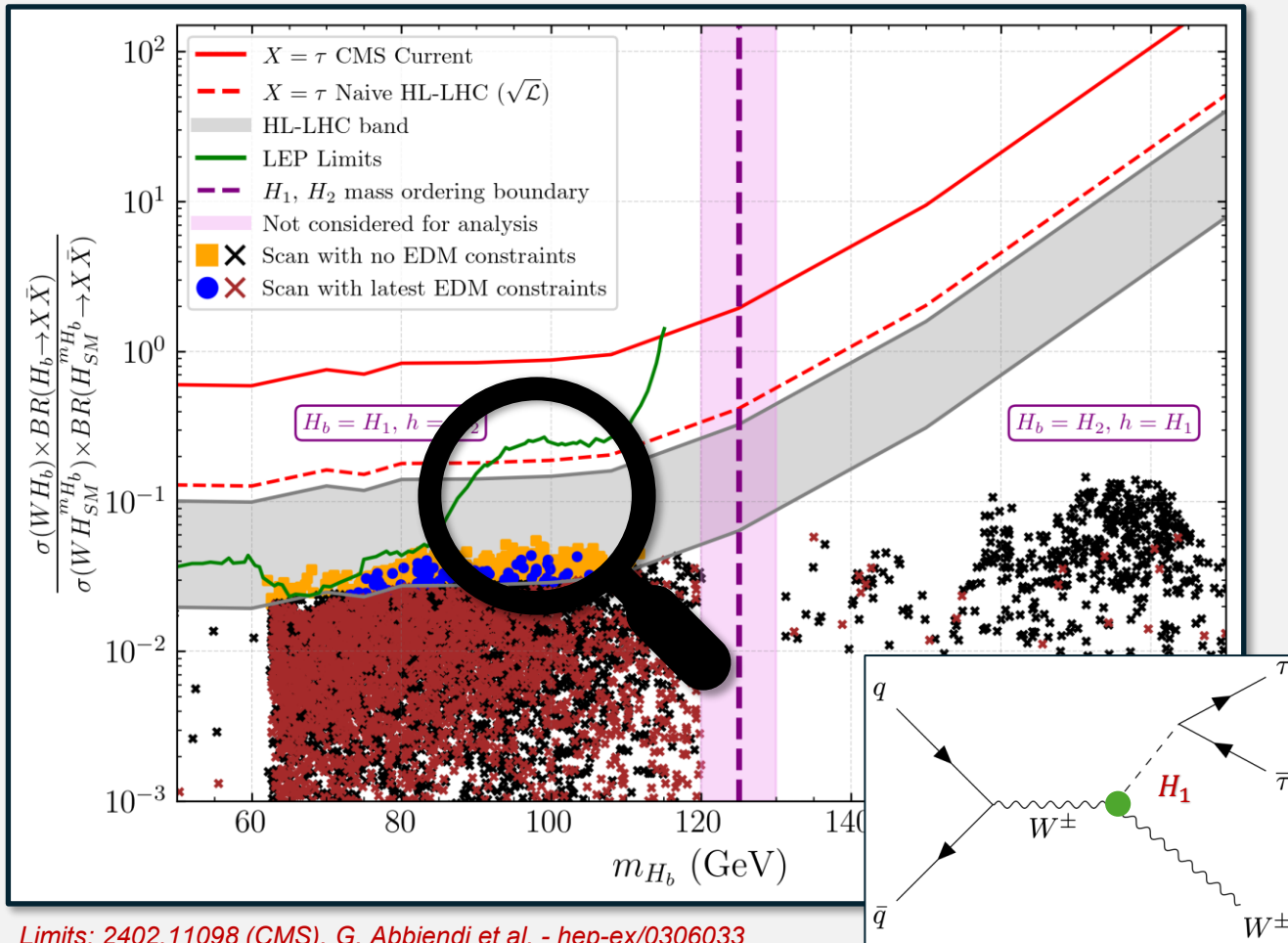
Statistics vs heuristics



Results

Set (vi): W -AP of lightest non-SM Higgs in the C2HDM

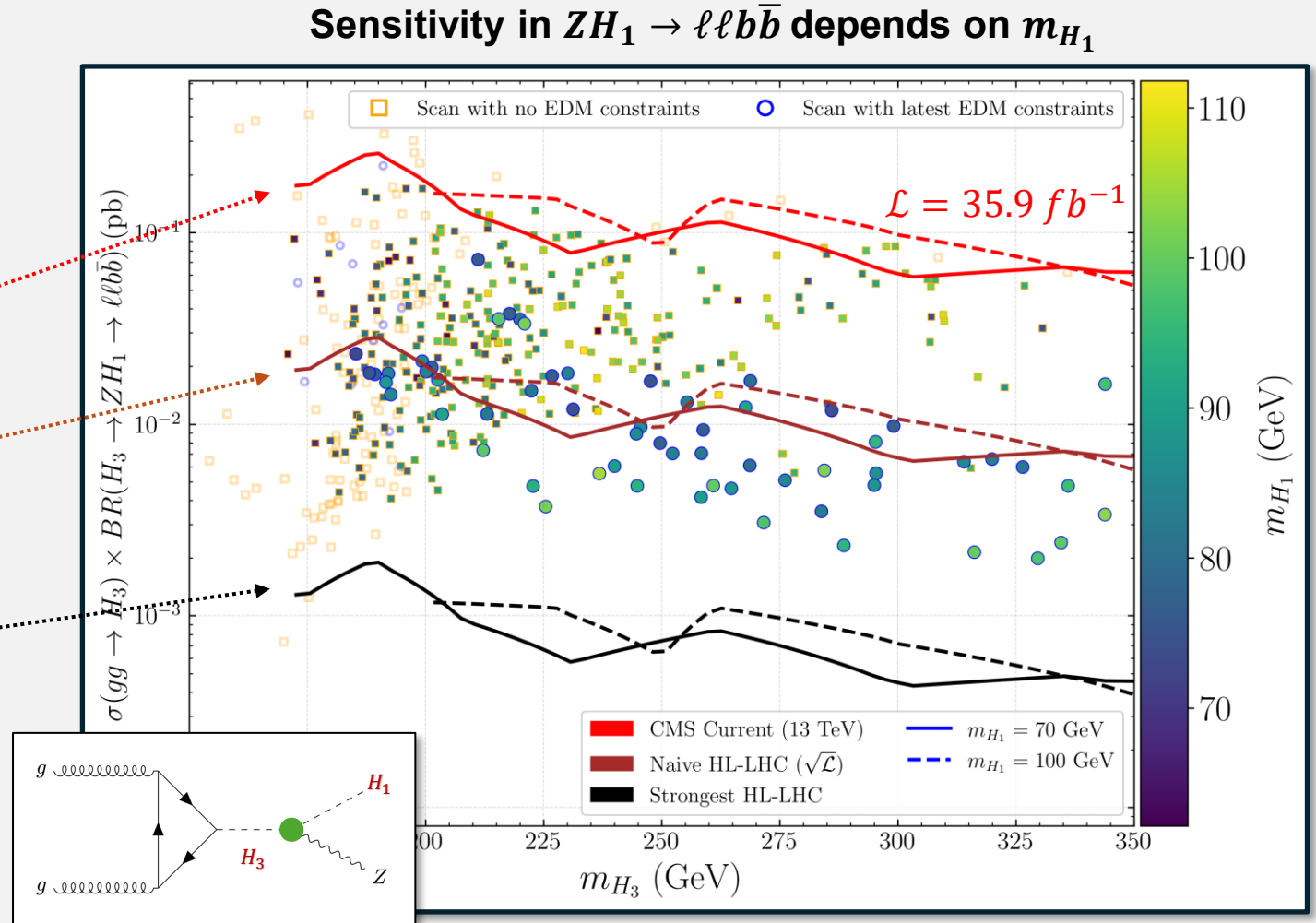
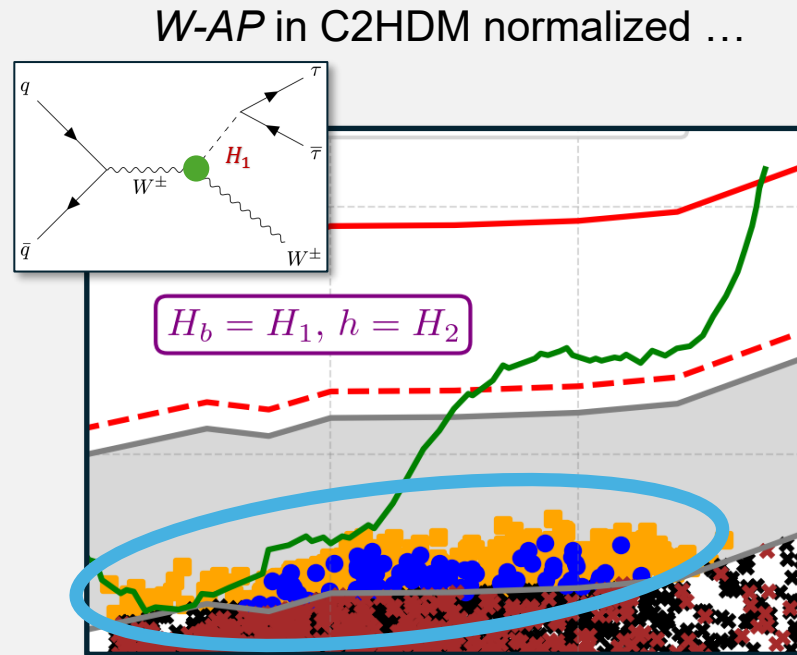
W -AP in C2HDM normalized to W -AP of same-mass SM Higgs



LEP limits: “Where does suppression come from?”

Results

Set (vi): reliance on $gg \rightarrow H_3 \rightarrow ZH_1$

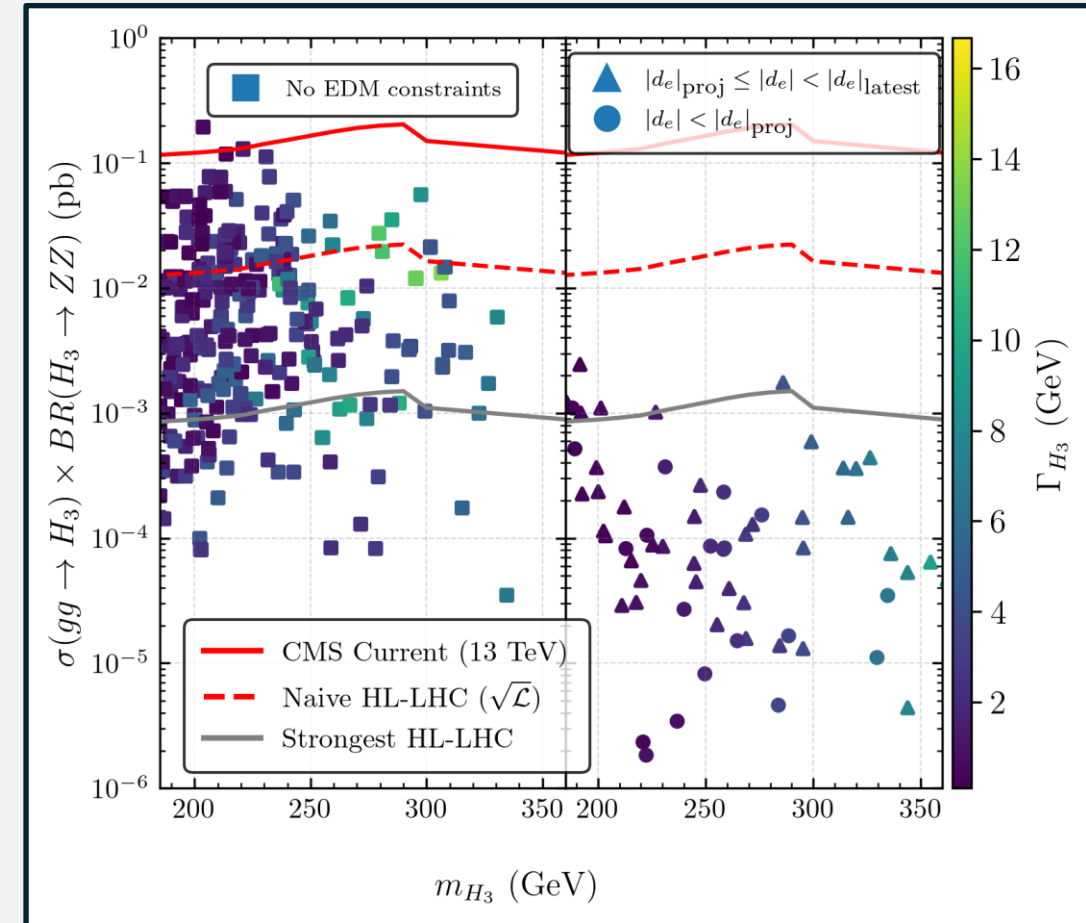
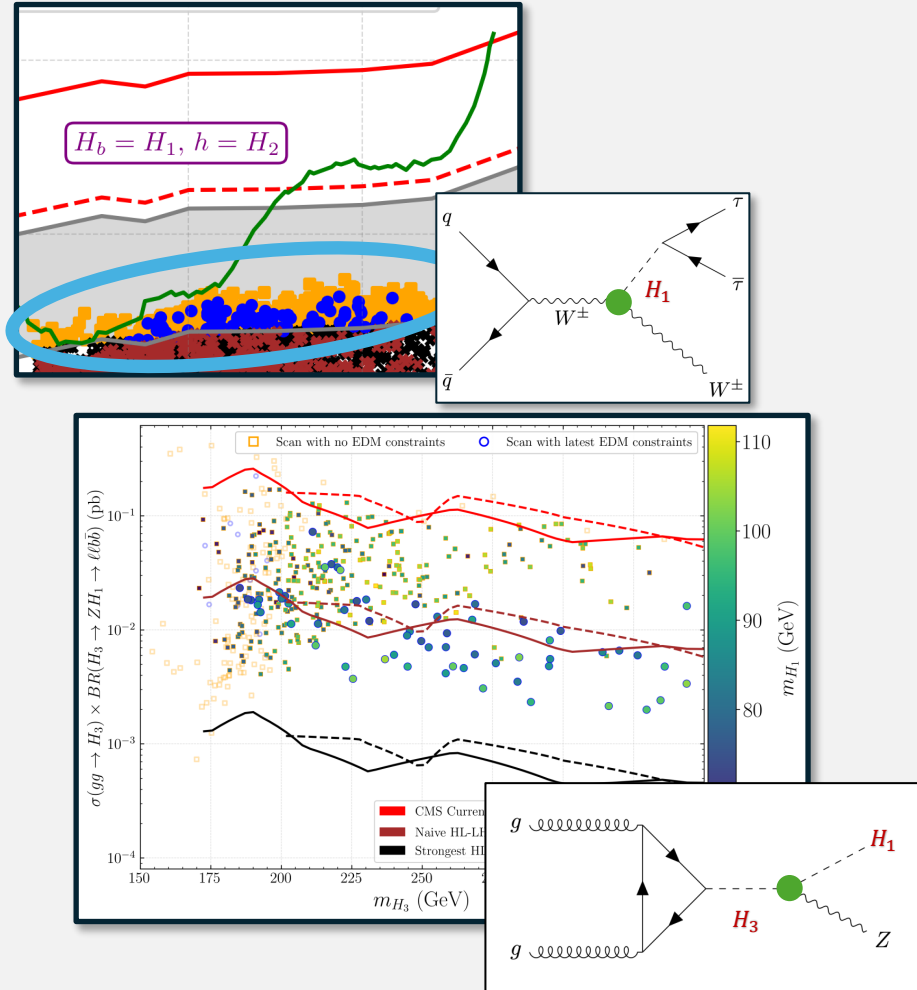


Limits: 1911.03781 (CMS)

Results

Set (vi): sensitivity to $pp \rightarrow H_3 \rightarrow ZZ$

Feasible in principle – up to Nature...



Limits: 1804.01939 (CMS)

Results

The (not so) dark clouds of eEDM

Present constraints: $|d_e| \leq 1.1 \times 10^{-29} e \cdot \text{cm}$

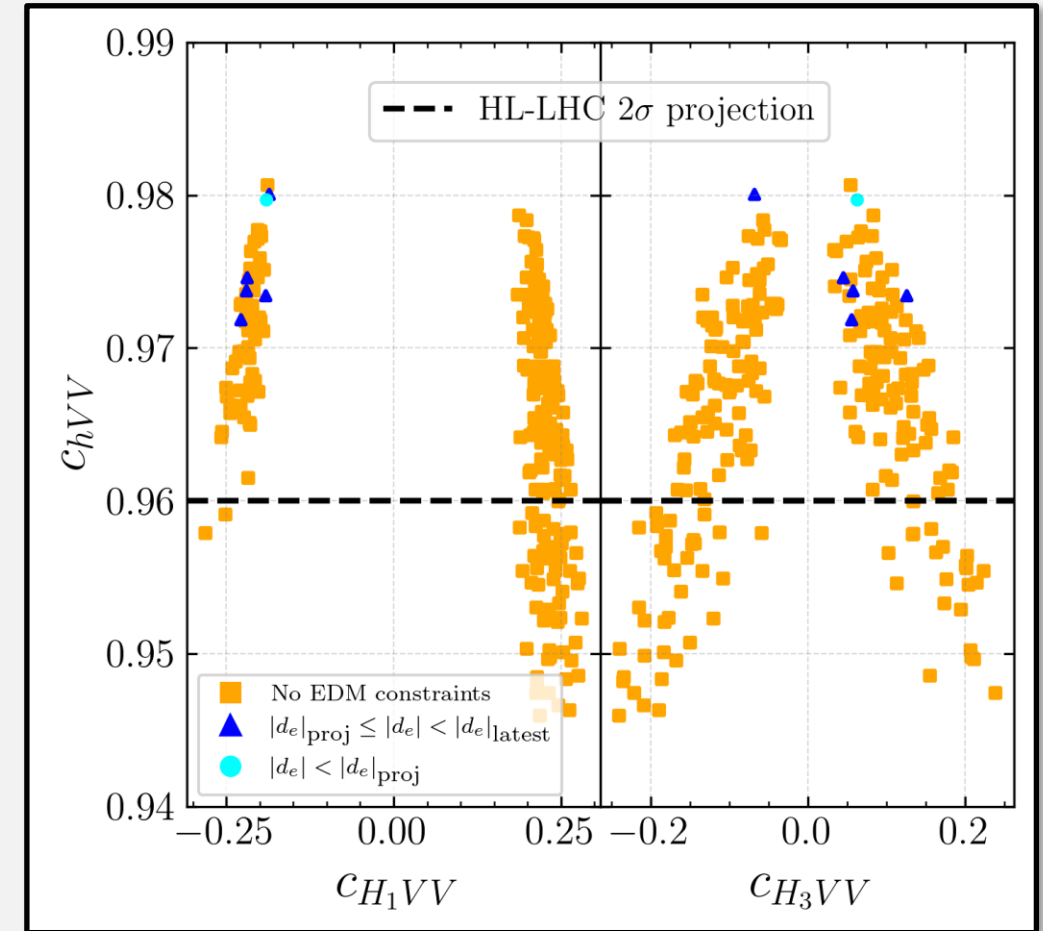
Future sensitivity projections: $|d_e| \leq 1.0 \times 10^{-30} e \cdot \text{cm}$

Kill off significant amount of testable points at LHC...!

Nonetheless, points still allowed under projections:

very SM-like Higgses h even by HL-LHC prospects!

Results of this work and upcoming spin-off: **remain relevant!**



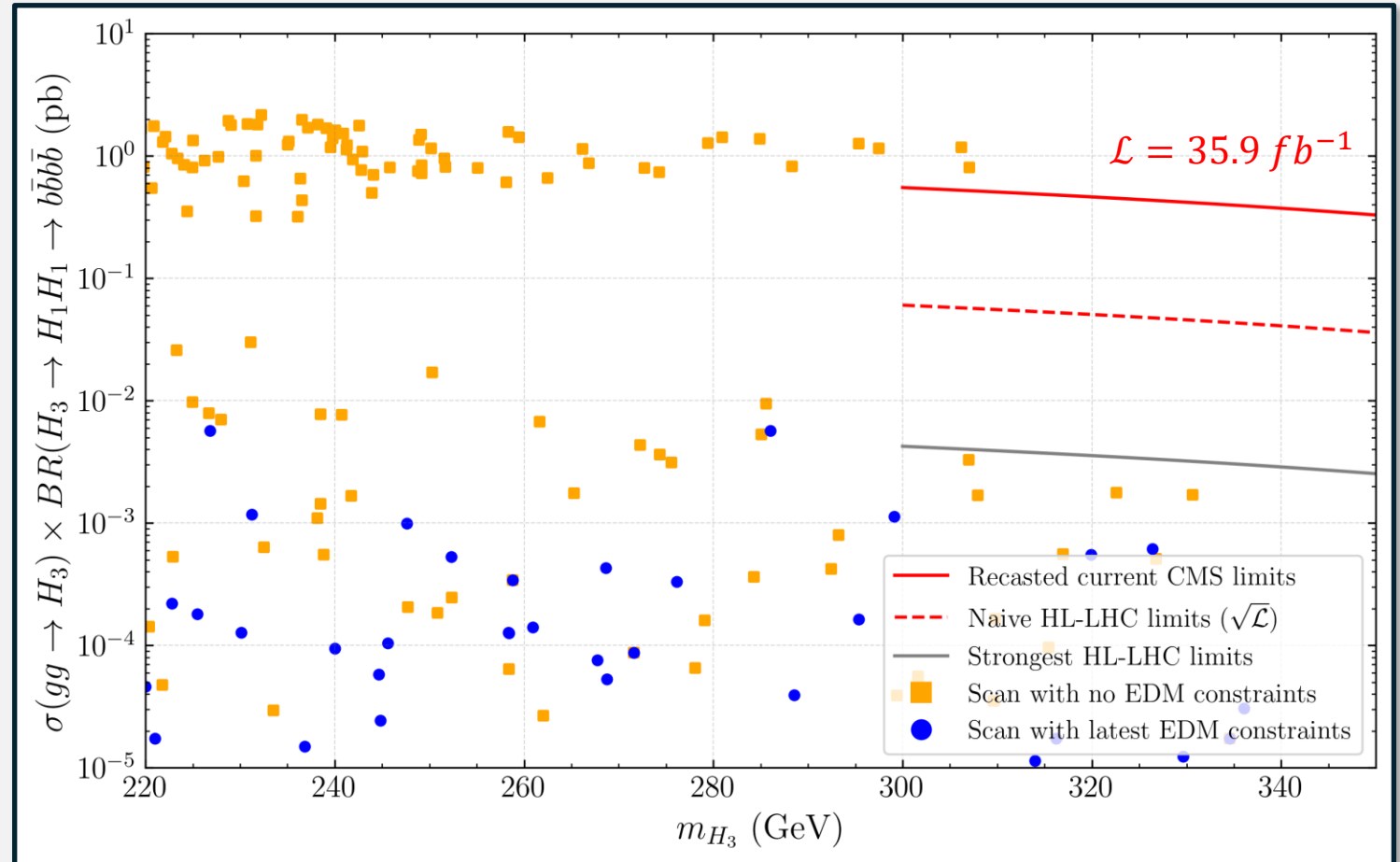
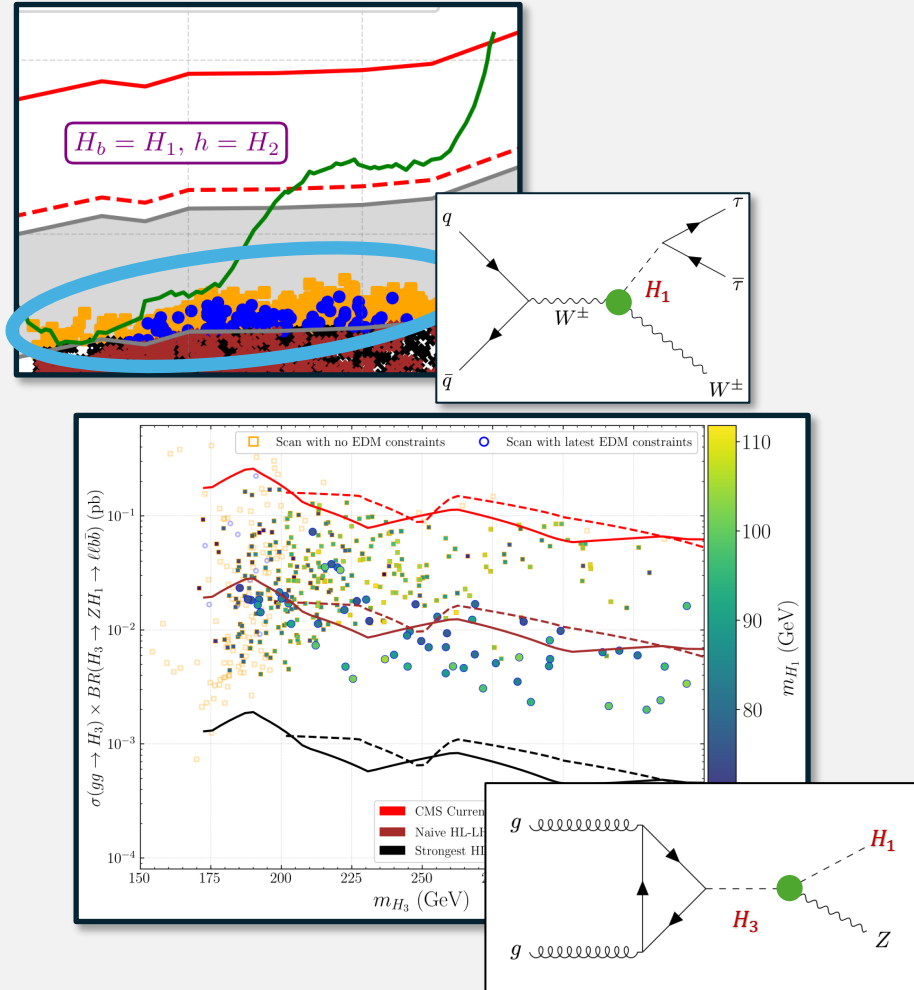
Gauge-Higgs couplings for points passing set (vi)

Limits: 2207.00043

Results

Set (vi): sensitivity to $pp \rightarrow H_3 \rightarrow H_1 H_1$

Limits need to be extended... but likely feasible



Limits: 1910.08574 (Barducci et al.)

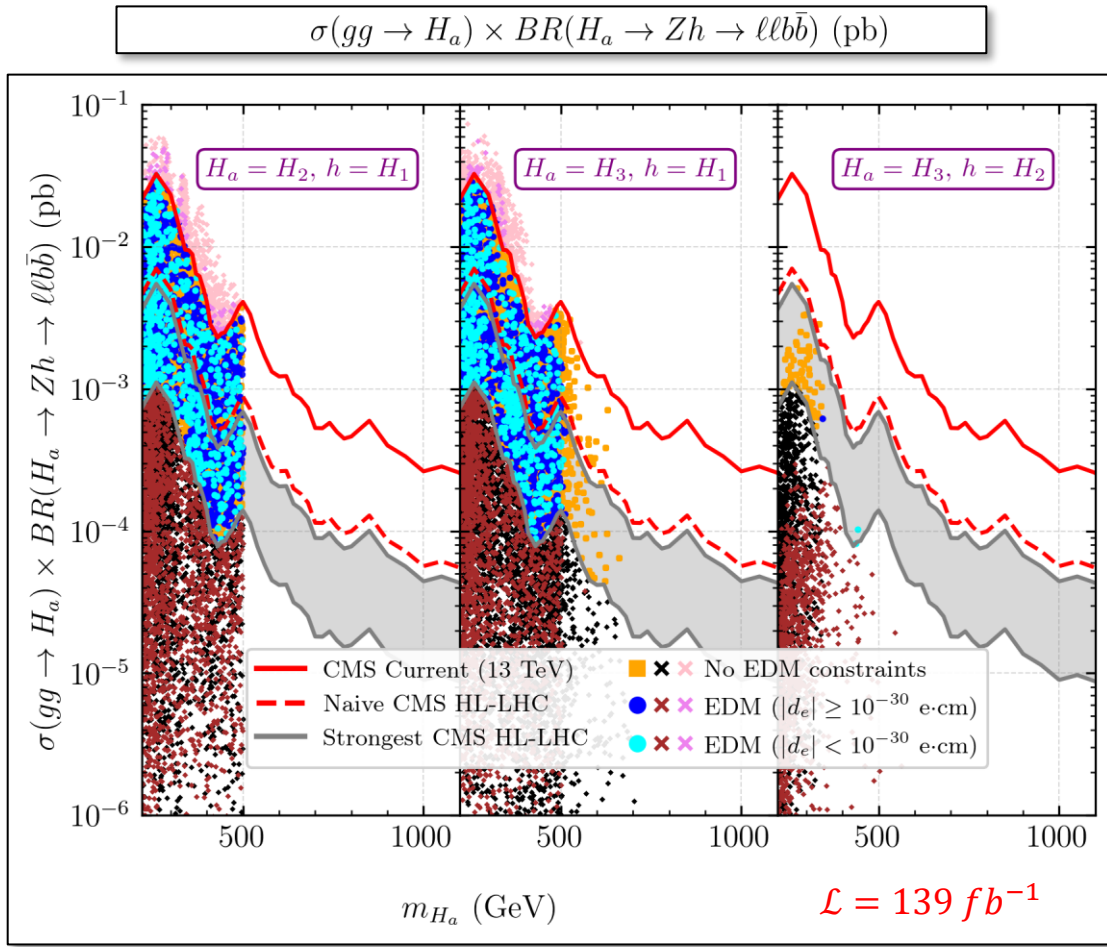
Unlocking the power of the LHC

Outlook and open issues: what can the community do?

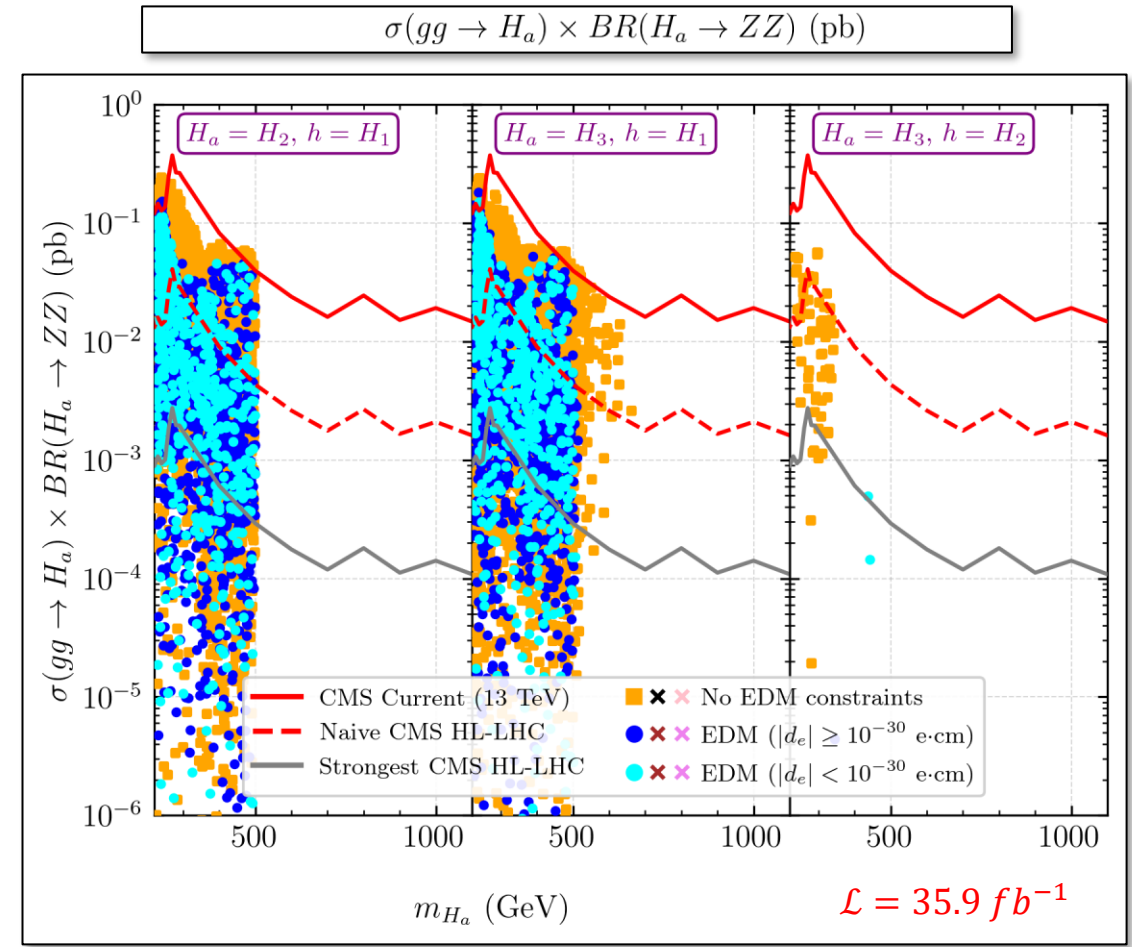
- (i) $pp \rightarrow VH_a, H_a \rightarrow Zh$
 - (ii) VBF production of $H_a, H_a \rightarrow Zh$
 - (iii) $pp \rightarrow VH_a, H_a \rightarrow ZH_b$ as well as $pp \rightarrow VH_b$
 - (iv) VBF production of $H_a, H_a \rightarrow ZH_b$ as well as VBF of H_b
 - (v) ggF production of H_a decaying via $H_a \rightarrow ZH_b$, as well as $pp \rightarrow VH_{a,b}$ or VBF of both $H_{a,b}$
 - (vi) ggF production of H_a decaying both via $H_a \rightarrow ZH_b$ and $H_a \rightarrow VV$ (or $H_a \rightarrow H_bH_b$), as well as $pp \rightarrow VH_b$ (or VBF of H_b)
 - (vii) ggF production of H_a decaying via $H_a \rightarrow Zh$, as well as $pp \rightarrow VH_a$ (or VBF production of H_a)
- No experimental limits on some processes...
 - ... or available in inconvenient formats!
 - Address quark-gluon jet tagging for VBF signals!

For the skeptics out there...

Preliminary results of the sequel, ggF-only work



Limits from 2207.00230



Limits from 1804.01939

Concluding remarks

CPV may be hiding in a scalar sector beyond the SM...

... the LHC is a promising place to find out...

... but hard work lies ahead.



Thank you!

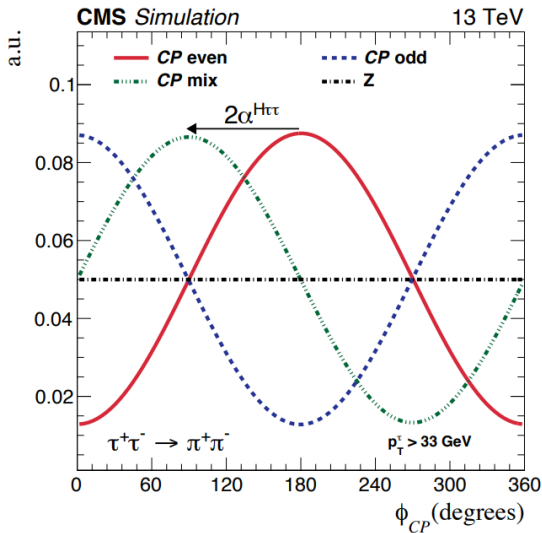
CP violation... in the Higgs sector?

A personal and non-exhaustive take

Fermion sector

$$-\frac{m_\tau}{v}H(\kappa_\tau\bar{\tau}\tau + \tilde{\kappa}_\tau\bar{\tau}i\gamma_5\tau)$$

CP-even CP-odd



CMS – 2110.04836

Scalar sector

Bosonic field	J^{PC}	J^P
γ	1^{--}	
Z	1^{--}	
$H_{\text{CP-even}}$	0^{++}	
$H_{\text{CP-odd}}, G^0$	0^{+-}	
$W^\pm$		1^-
$H^\pm, G^\pm$		0^+

$H_3 \rightarrow ZH_2$	$H_2 \rightarrow Zh_1$	
$H_2 \rightarrow ZH_1$	$H_1 \rightarrow ZZ$	...
$H_3 \rightarrow ZH_1$	$H_2 \rightarrow ZZ$	

1506.06755, 2206.09643, ...

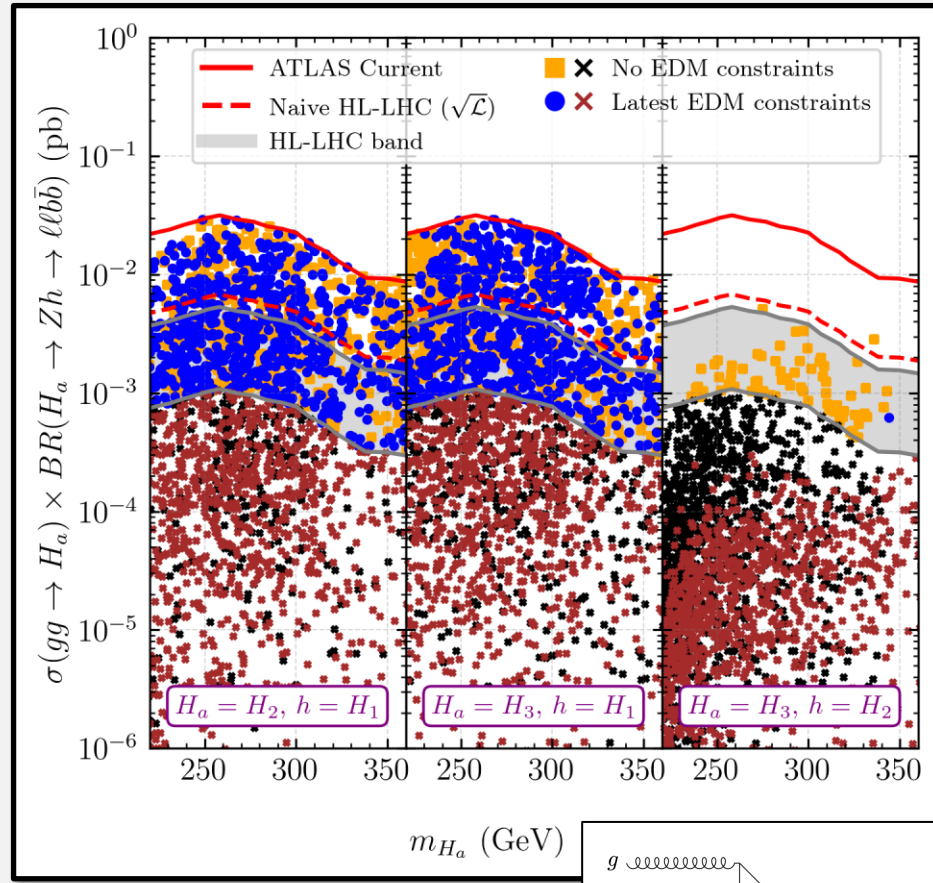
CP states *ad absurdum*?

➡ CPV

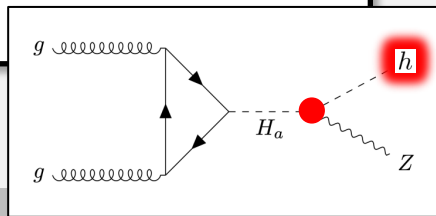
Results

Set (vii): Example of a joint signature

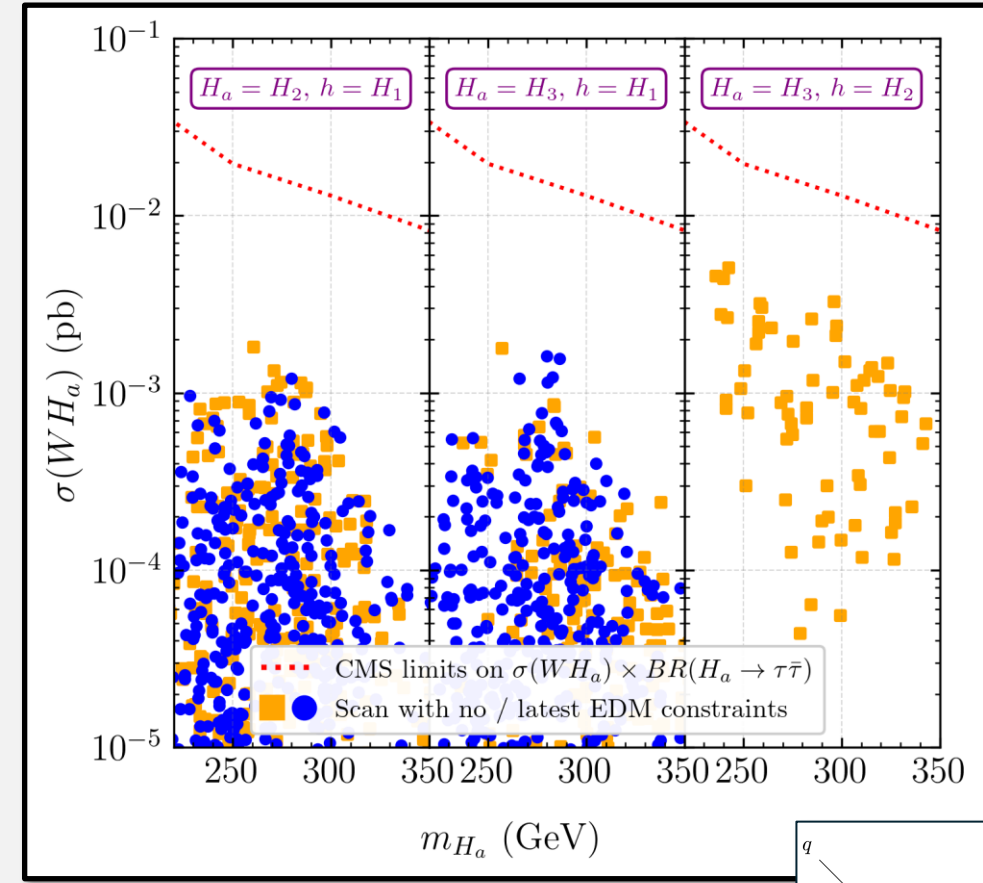
Sensitivity to $gg \rightarrow H_a, H_a \rightarrow Zh \rightarrow \ell\ell b\bar{b}$



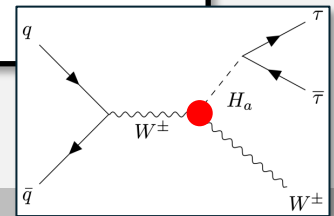
Limits: 2207.00230 (ATLAS)



Sensitivity to $pp \rightarrow WH_a$



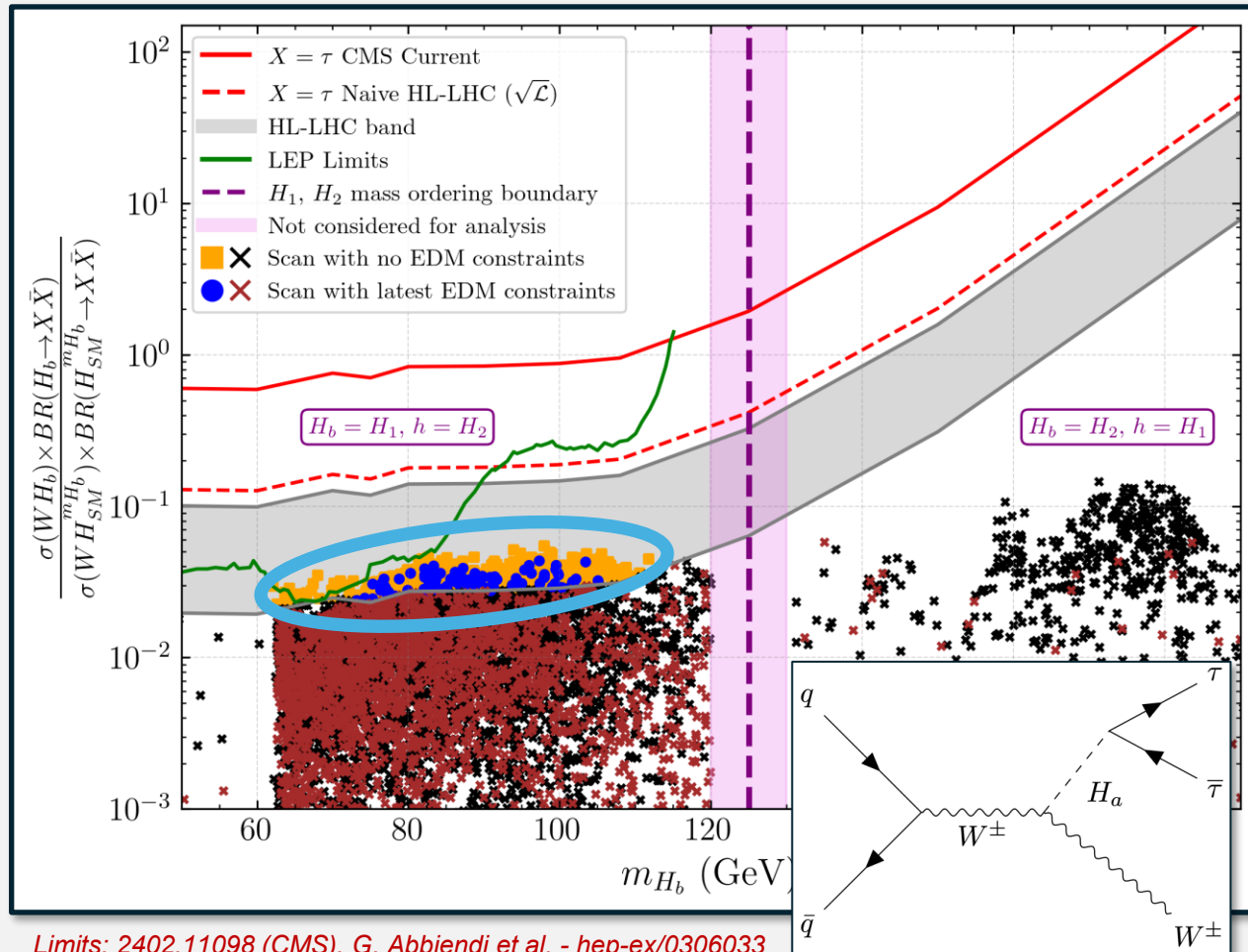
Limits: 2402.11098 (CMS), G. Abbiendi et al. - hep-ex/0306033



Results

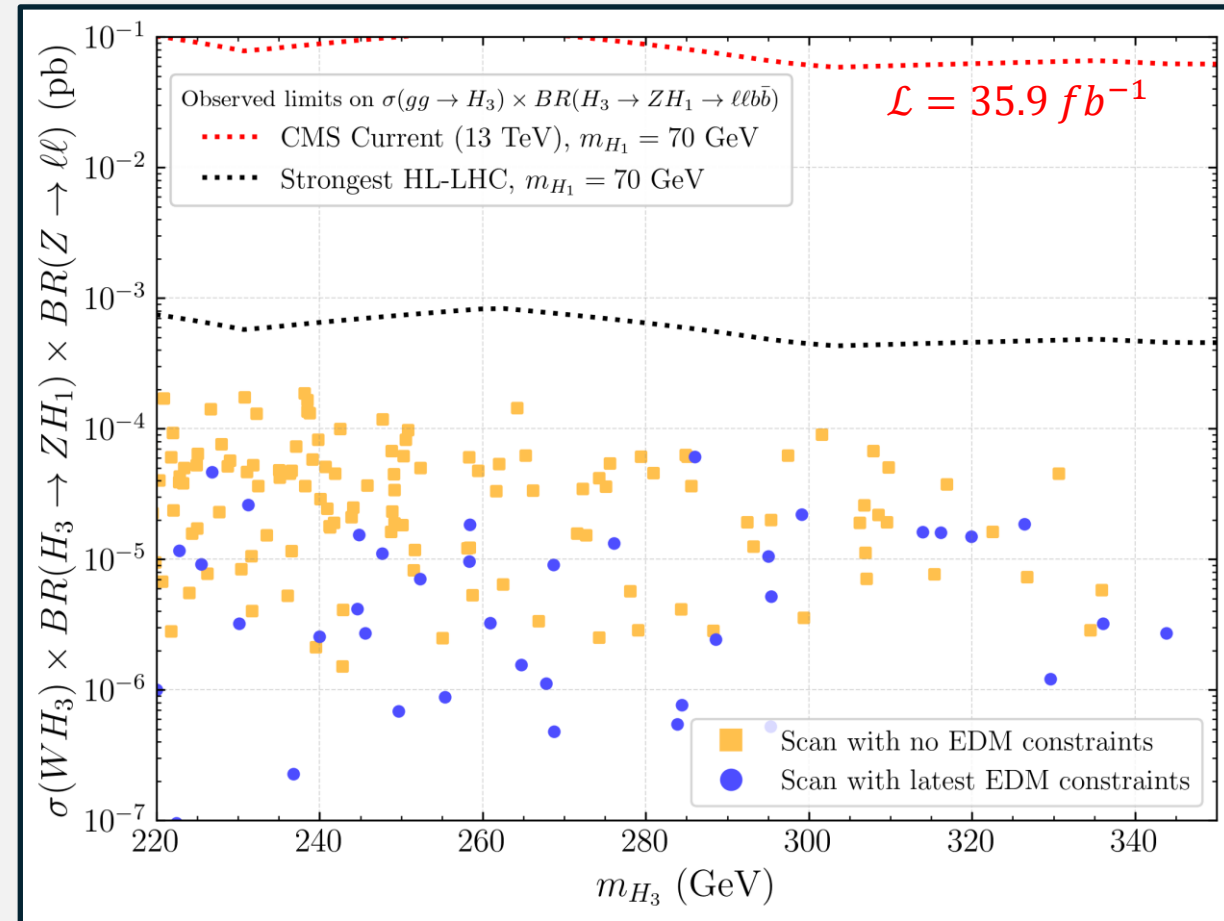
Set (iii): $pp \rightarrow VH_a$, $H_a \rightarrow ZH_b$ and $pp \rightarrow VH_b$

W -AP in C2HDM normalized to W -AP of same-mass SM Higgs



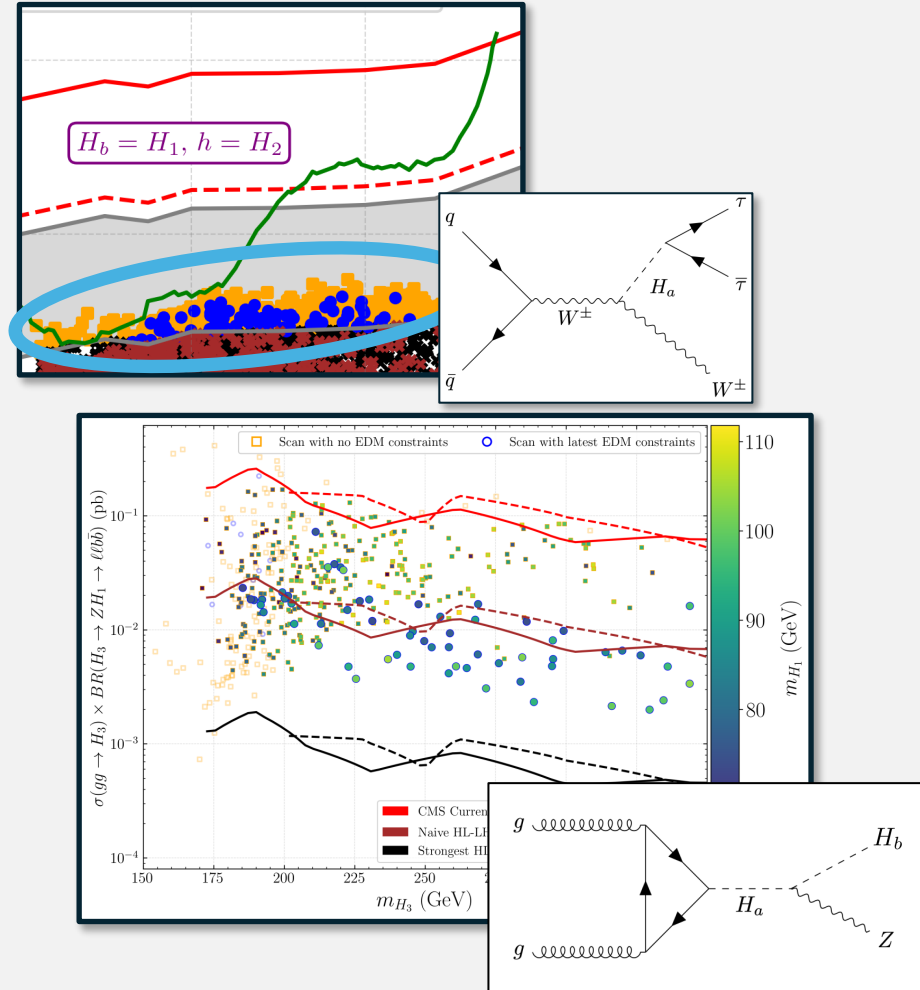
Limits: 2402.11098 (CMS), G. Abbiendi et al. - hep-ex/0306033

No existing limits on $pp \rightarrow VH_3$, $H_3 \rightarrow H_1(Z \rightarrow \ell\ell)$!

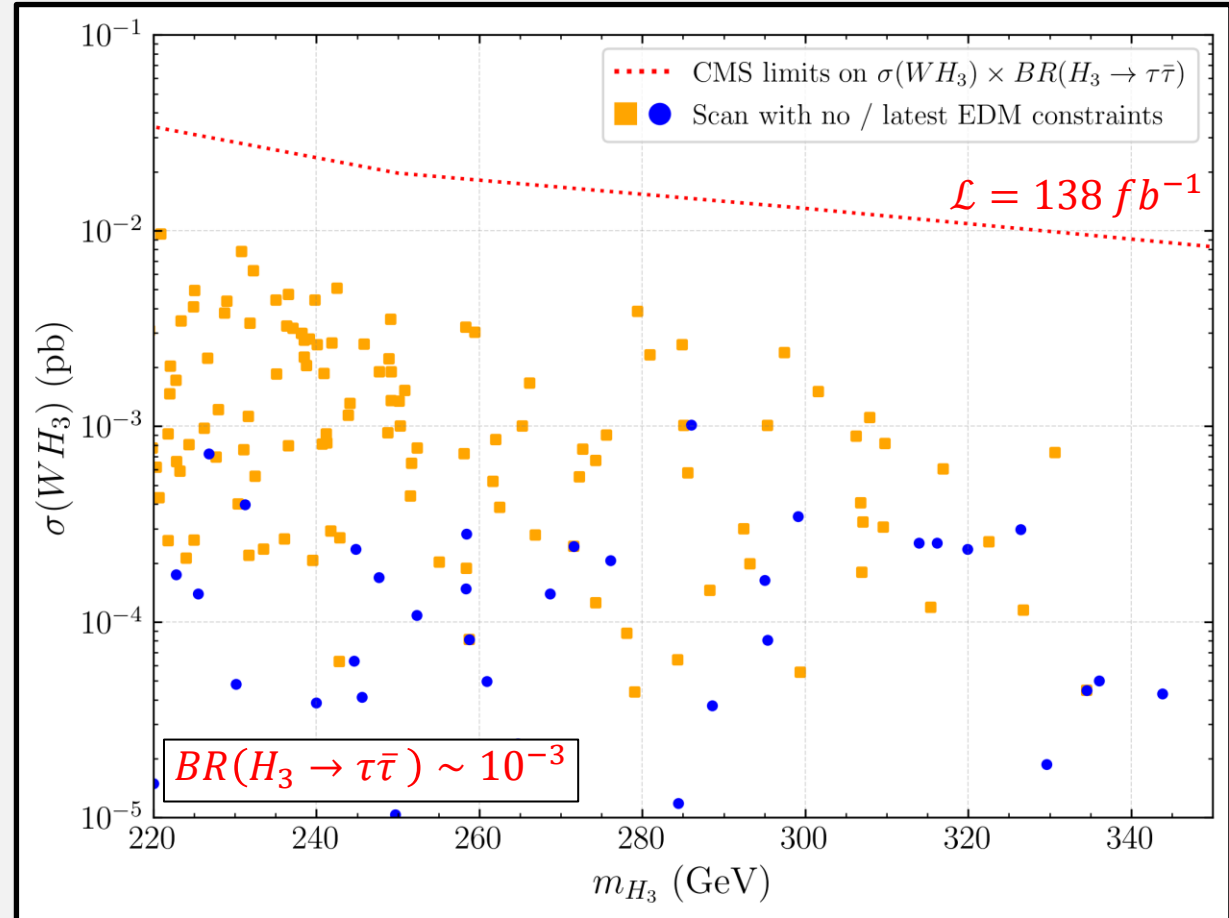


Results

Set (v): sensitivity to $pp \rightarrow W(H_3 \rightarrow \tau\bar{\tau})$



Full cross section well below sensitivity to $\tau\bar{\tau}$ final state

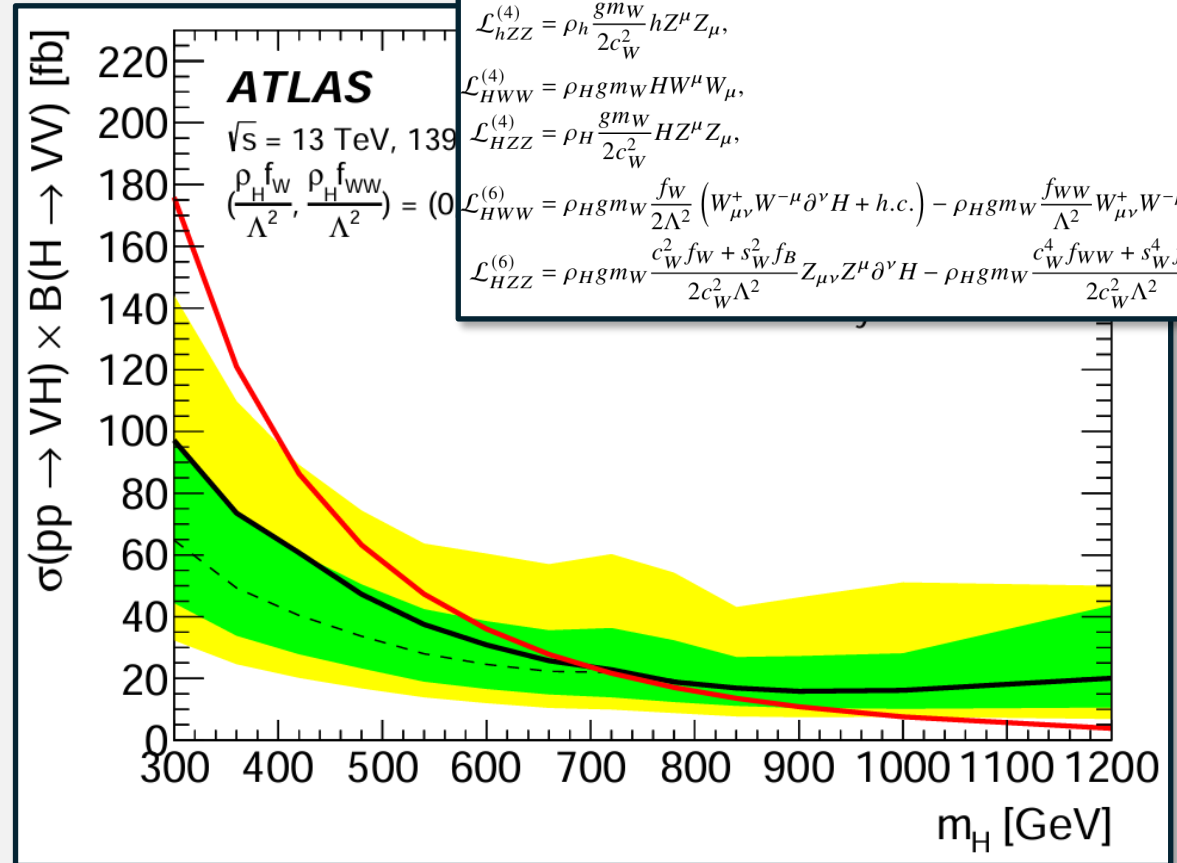
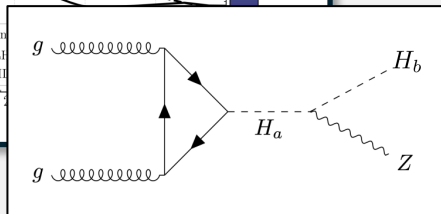
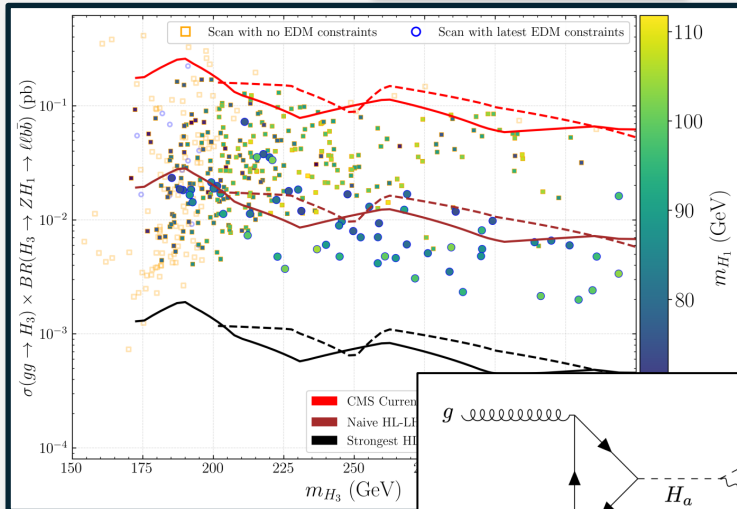
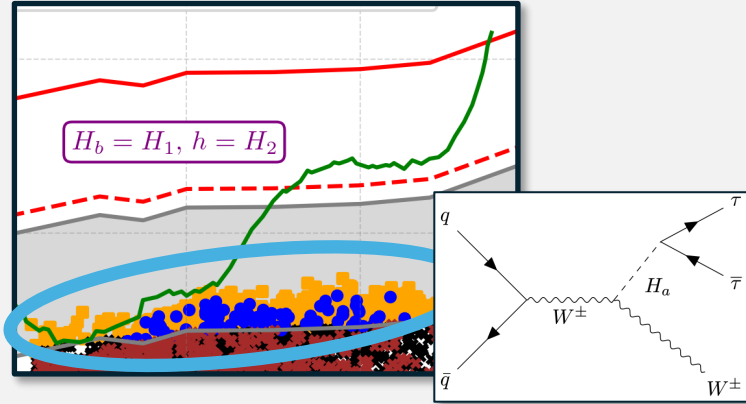


Limits: 2402.11098 (CMS), G. Abbiendi et al. - hep-ex/0306033

Results

Set (v): sensitivity to $pp \rightarrow W(H_3 \rightarrow WW \rightarrow 2\ell^{+/-} + 2\nu)$

Cannot use existing limits!



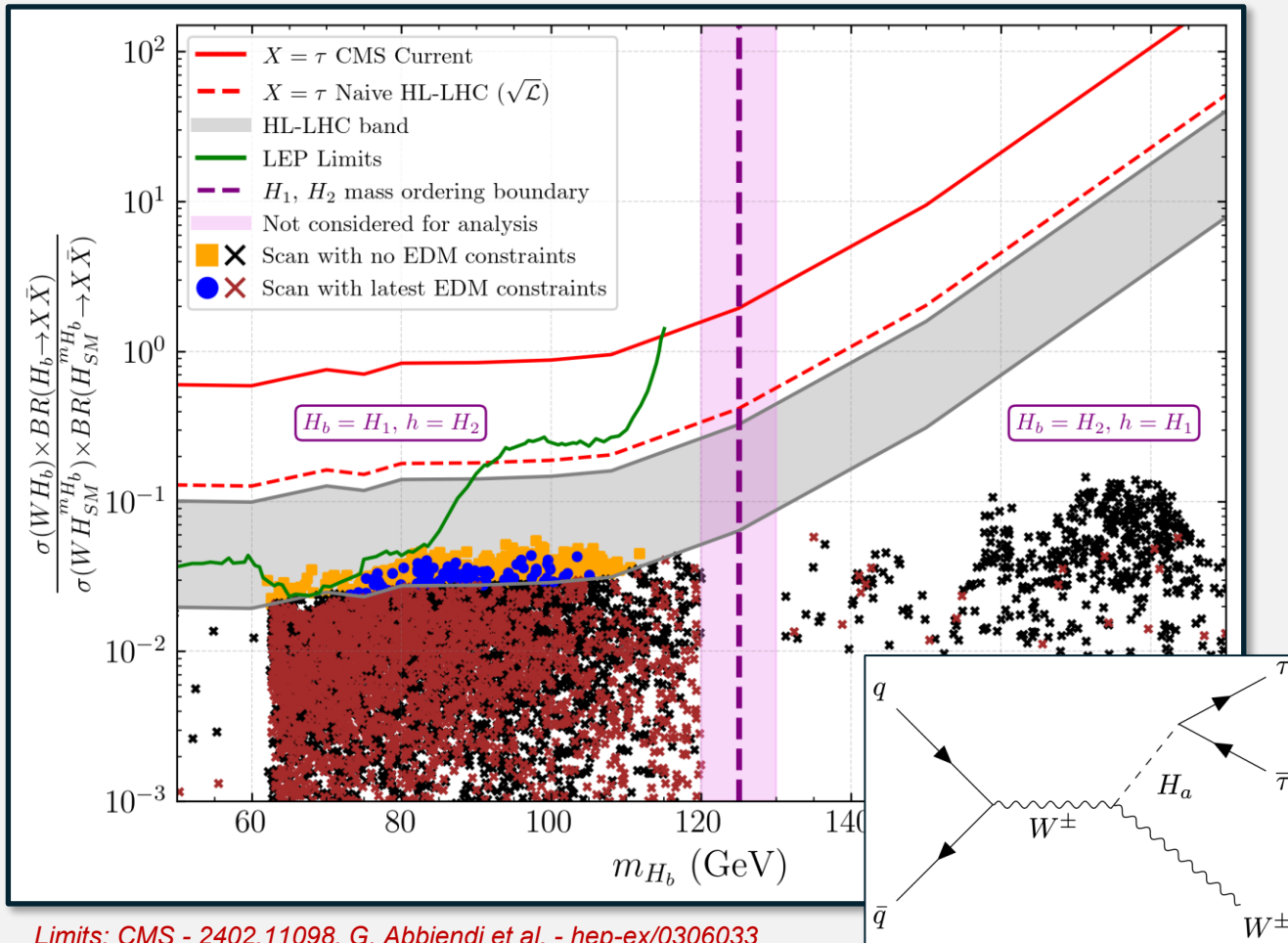
$$\begin{aligned}\mathcal{L}_{hWW}^{(4)} &= \rho_h g m_W h W^\mu W_\mu, \\ \mathcal{L}_{hZZ}^{(4)} &= \rho_h \frac{g m_W}{2c_W^2} h Z^\mu Z_\mu, \\ \mathcal{L}_{HWW}^{(4)} &= \rho_H g m_W H W^\mu W_\mu, \\ \mathcal{L}_{HZZ}^{(4)} &= \rho_H \frac{g m_W}{2c_W^2} H Z^\mu Z_\mu, \\ \mathcal{L}_{HWW}^{(6)} &= \rho_H g m_W \frac{f_W}{2\Lambda^2} (W_{\mu\nu}^+ W^{-\mu} \partial^\nu H + h.c.) - \rho_H g m_W \frac{f_{WW}}{\Lambda^2} W_{\mu\nu}^+ W^{-\mu\nu} H, \\ \mathcal{L}_{HZZ}^{(6)} &= \rho_H g m_W \frac{c_W^2 f_W + s_W^2 f_B}{2c_W^2 \Lambda^2} Z_{\mu\nu} Z^\mu \partial^\nu H - \rho_H g m_W \frac{c_W^4 f_{WW} + s_W^4 f_{BB}}{2c_W^2 \Lambda^2} Z_{\mu\nu} Z^{\mu\nu} H\end{aligned}$$

Plot and formulae taken from 2211.02617 (ATLAS)

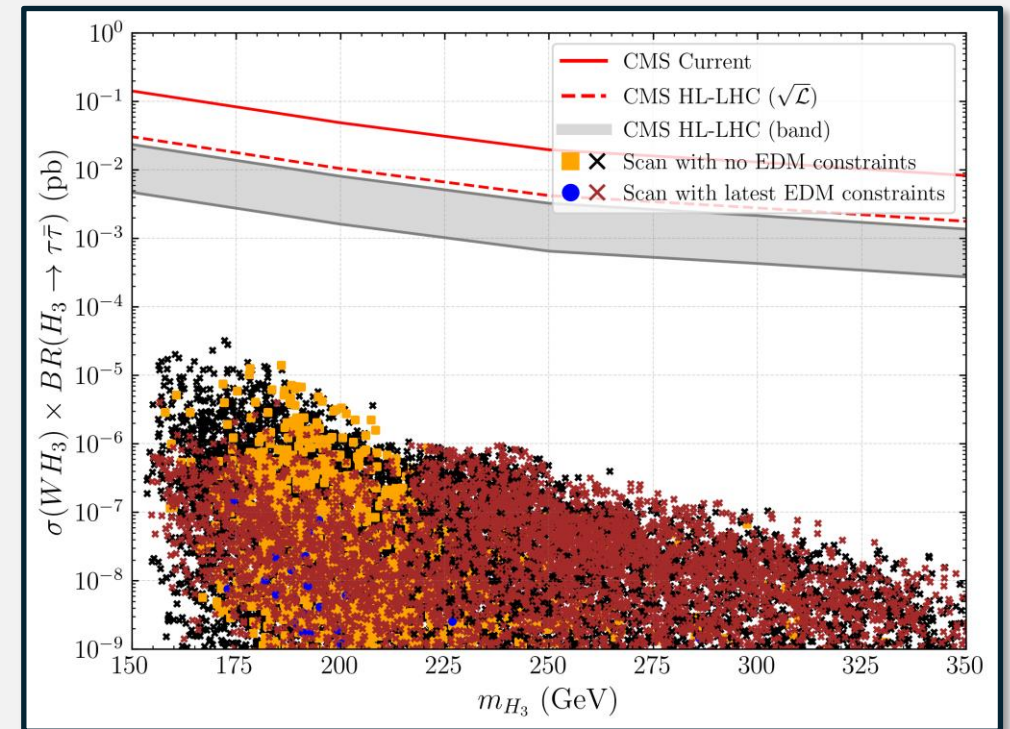
Exemplary results

Sets (ii) – (vi): *W-AP* of lightest non-SM Higgs in the C2HDM

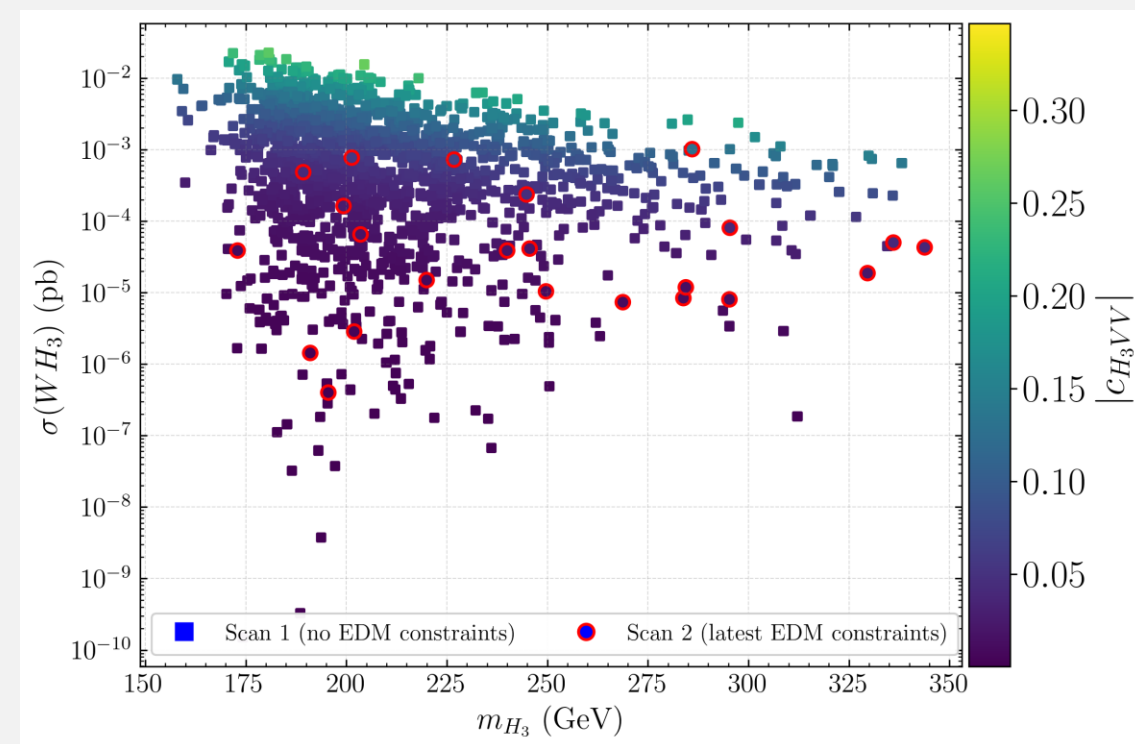
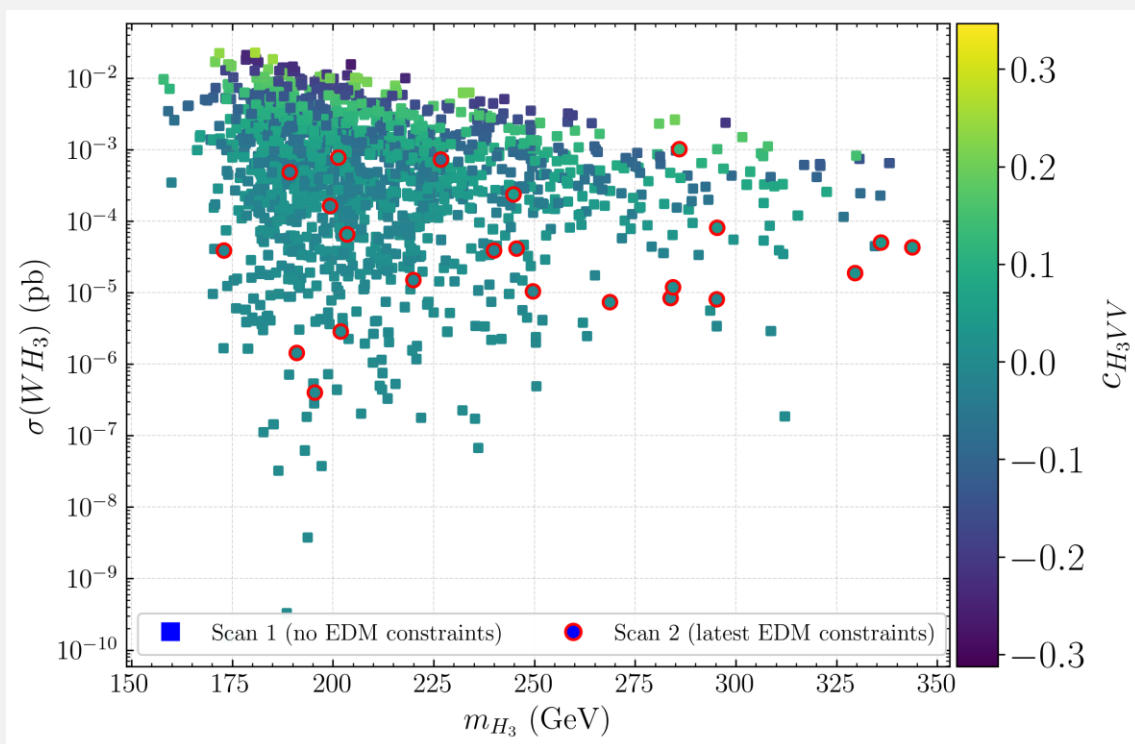
W-AP in C2HDM normalized to *W-AP* of same-mass SM Higgs



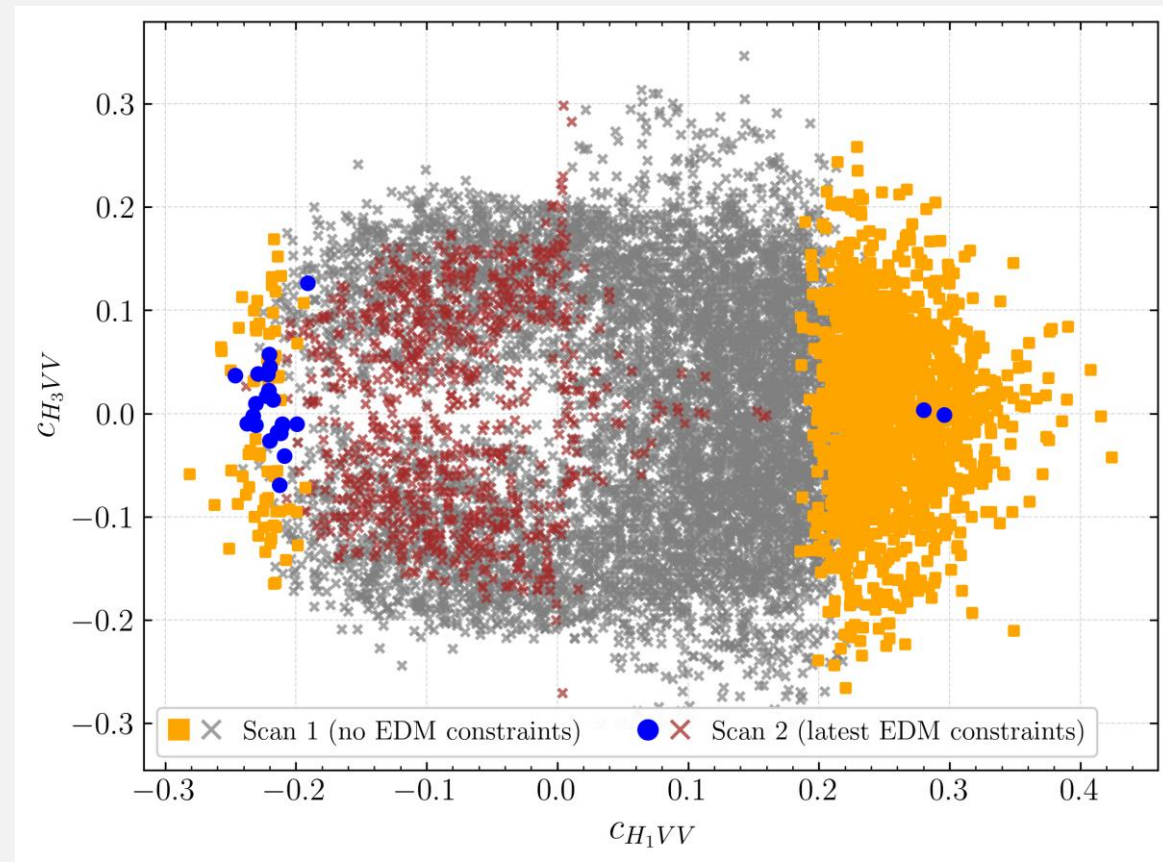
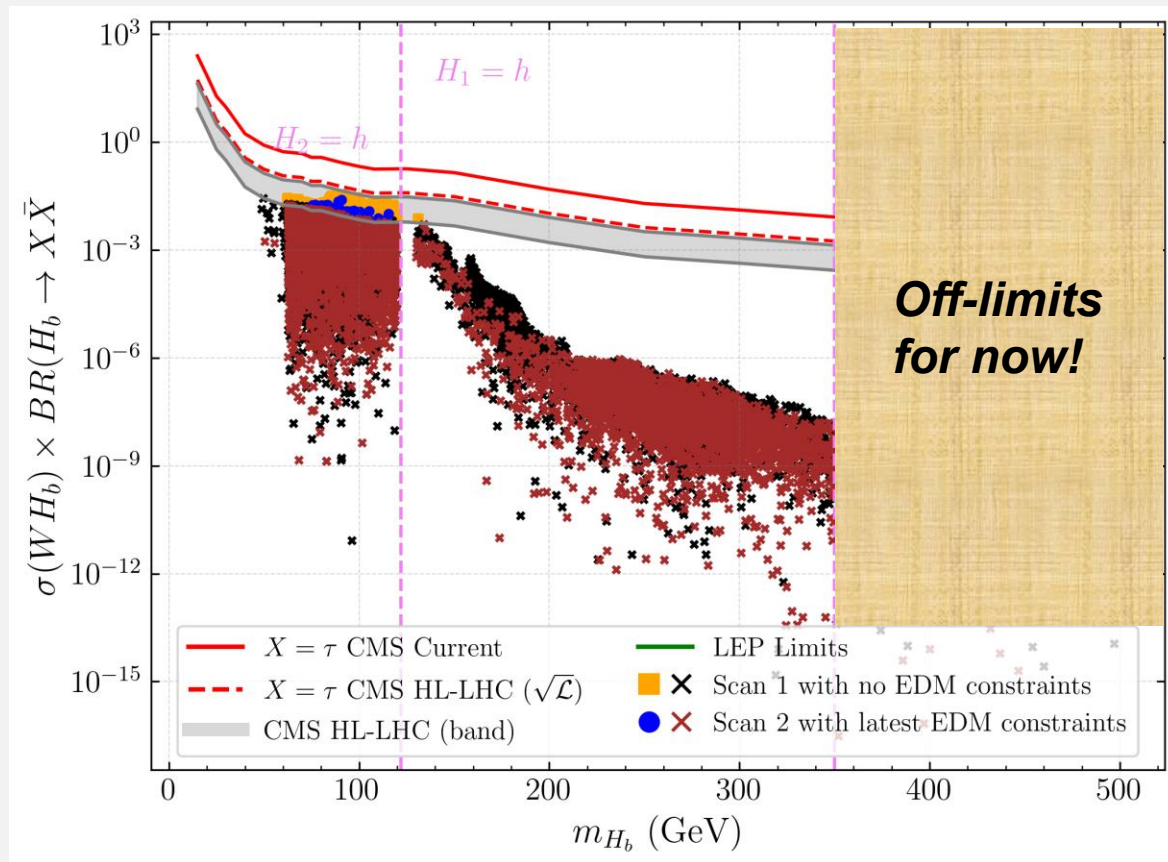
Compare to *W-AP* of H_3 ...



Backup



Backup



Backup

