

# High-precision measurement of the $B_d^0$ meson lifetime at the ATLAS experiment

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WIN 2025, University of Sussex, Brighton, 9-13 June 2025

# A new record for precision on B-meson lifetimes

As direct searches for physics beyond the Standard Model continue to push frontiers at the LHC, the b-hadron physics sector remains a crucial source of insight for testing established theoretical models.

The ATLAS collaboration recently published a new measurement of the  $B^0$  lifetime using  $B^0 \rightarrow J/\psi K^0$  decays from the entire Run-2 dataset it has recorded at 13 TeV. The result improves the precision of previous world-leading measurements by the CMS and LHCb collaborations by a factor of two.

Studies of b-hadron lifetimes probe our understanding of the weak interaction. The lifetimes of b-hadrons can be systematically computed within the heavy-quark expansion (HQE) framework, where b-hadron observables are expressed as a perturbative expansion in inverse powers of the b-quark mass.

ATLAS measures the "effective"  $B^0$  lifetime, which represents the average decay time incorporating effects from mixing and CP contributions, as  $\tau(B^0) = 1.5053 \pm 0.0012$  (stat.)  $\pm 0.0035$  (syst.) ps. The result is consistent with previous measurements published by ATLAS and other experiments, as summarised in figure 1. It also aligns with theoretical predictions from HQE and lattice QCD, as well as with the experimental world average.

The analysis benefitted from the large Run-2 dataset and a refined trigger selection, enabling the collection of an extensive sample of 2.5 million  $B^0 \rightarrow J/\psi K^0$  decays. Events with a  $J/\psi$  meson decaying into two muons with sufficient transverse momentum are cleanly identified in the ATLAS Muon Spectrometer by the first-level hardware trigger. In the next-level software trigger, exploiting the full detector information, these muons are

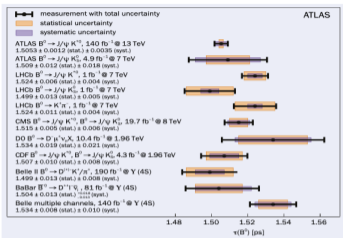


Fig. 1. A comparison of the current ATLAS result for the  $B^0$  lifetime with the previous ATLAS result in the  $B^0 \rightarrow J/\psi K^0$  channel, and with those from other experiments.

then combined with two tracks measured by the Inner Detector, ensuring they originate from the same vertex.

The  $B^0$ -meson lifetime is determined through a two-dimensional unbinned maximum-likelihood fit, utilising the measured  $B^0$ -candidate mass and decay time, and accounting for both signal and background components. The limited hadronic particle-identification capability of ATLAS requires careful modelling of the significant backgrounds from other processes that produce  $J/\psi$  mesons. The sensitivity of the fit is increased by estimating the uncertainty of the decay-time measurement provided by the ATLAS tracking and vertexing algorithms on a per-candidate basis. The resulting lifetime measurement is limited by sys-

tematic uncertainties, with the largest contributions arising from the correlation between  $B^0$  mass and lifetime, and ambiguities in modelling the mass distribution.

ATLAS combined its measurement with the average decay width ( $\Gamma$ ) of the light and heavy  $B_s$ -meson mass eigenstates, also measured by ATLAS, to determine the ratio of decay widths as  $\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022$  (stat.)  $\pm 0.0036$  (syst.)  $\pm 0.0057$  (ext.). The result is consistent with unity and provides a stringent test of QCD predictions, which also support a value near unity.

#### Further reading

ATLAS Collab. 2024, arXiv:2411.09962, ATLAS Collab. 2021, Eur. Phys. J. C **81** 342.

- "Precision measurement of the  $B_d^0$  meson lifetime performed at the ATLAS using Run2 data" [arXiv:2411.09962v1](https://arxiv.org/abs/2411.09962v1)-hep-ex-15 Nov 2024 ; accepted by JHEP, April 2025.
- ATLAS  $B^0$  lifetime is the most precise measurement to date
- Also measured the decay rate  $\Gamma_d$  of  $B^0$
- $\Gamma_d/\Gamma_s$  determined using ATLAS  $\Gamma_s$  measured in  $B_s^0 \rightarrow J/\psi \phi$  [Eur. Phys. J. C 81 \(2021\) 342](#)
- Recognised by ATLAS [ATLAS Briefing](#) ATLAS Briefing 25 November 2024
- Recognised by CERN [LHC CERN Seminar](#) and Article in [CERN Courier](#) p. 15).

- In the Heavy Quark Expansion (HQE) theory the decay rate  $\Gamma = 1/\tau$  is calculated:

$$\Gamma(\mathcal{B}_q) = \Gamma_3 + \delta\Gamma(\mathcal{B}_q)$$

↙ leading
↘ subleading

**Free  $b$ -quark decay:**

- + free of non-perturbative uncertainties
- 0 Looks like the muon decay

$$\Gamma_3 \propto \frac{G_F^2 m_b^5}{192\pi^3} V_{cb}^2$$

- Quark masses are difficult to define, huge dependence on definition can be reduced by higher order **perturbative corrections**

**Power-suppressed terms on the HQE:**

- + suppressed with at least 2 powers of  $1/m_b \Rightarrow$  small
- 0 Individual contributions are products of **perturbative** Wilson coefficients and **non-perturbative matrix elements** (determined with lattice-QCD, sum rules and/or from fits of experimental data of inclusive semi-leptonic decays -  $V_{cb}$ )

- $\Gamma$  suffers from huge uncertainties of the leading term  $\Gamma_3$ ;  $\Gamma_d = 0.63_{-0.07}^{+0.11} \text{ ps}^{-1}$  [Lenz, et al. 2023](#) ↗
- In rates ratios the leading term  $\Gamma_3$  exactly cancels leaving just small uncertainties from sub-leading terms;  $\Gamma_d/\Gamma_s = 1.003 \pm 0.006$  [Lenz, Piscopo, Rusov 2023](#) ↗
- Lifetimes measurements also test models of New Physics (BSM) contributing to both leading and sub-leading terms. Excluding BSM effects larger than 5%, will already constrain many BSM scenarios [Lenz2021](#) ↗
- Since the theory predictions are more precise for ratios - we also extracted  $\Gamma_d/\Gamma_s$ .

## Method of derivation of decay rate $\Gamma_d$ and ratio $\Gamma_d/\Gamma_s$ from measured $B_d^0$ lifetime.

The lifetime we measured in  $B^0 \rightarrow J/\psi K^{*0}$  is the effective lifetime  $\tau_{B_d^0}$  related to decay widths:  $\Gamma_L$ ,  $\Gamma_H$  of the light and heavy mass eigenstates of  $B_d^0-\bar{B}_d^0$  system via: [Fleischer et al, 2011](#)

$$\tau_{B^0} = \frac{1}{\Gamma_d} \frac{1}{1-y^2} \left( \frac{1+2Ay+y^2}{1+Ay} \right), \quad (1)$$

$\Gamma_d = (\Gamma_L + \Gamma_H)/2$ ;  $y = \Delta\Gamma_d/(2\Gamma_d) = (\Gamma_L - \Gamma_H)/(2\Gamma_d)$ . The asymmetry  $A$ :

$$A = \frac{R_H^f - R_L^f}{R_H^f + R_L^f}.$$

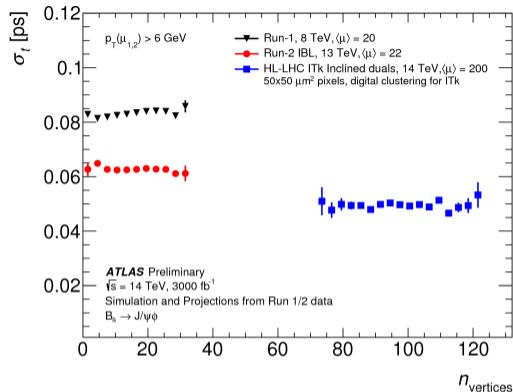
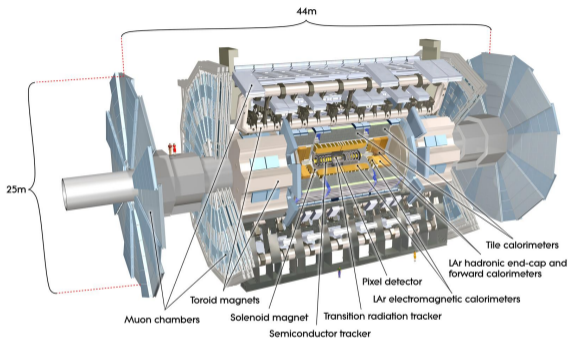
amplitudes  $R_L^f$  and  $R_H^f$  are defined via the summed decay rate of the members of the  $B_d^0-\bar{B}_d^0$  system: [Fleischer et al, 2011](#)

$$\langle \Gamma(B^0(t)) \rangle = \Gamma(B^0(t)) + \Gamma(\bar{B}^0(t)) = R_H^f \exp(-\Gamma_H t) + R_L^f \exp(-\Gamma_L t).$$

Using the values of  $y$  and  $A$  from Heavy Flavour Averaging group (HFLAV) Ref. [HFLAV:2023](#) , our  $\tau_{B_d^0}$  value and Eq(1) allows  $\Gamma_d$  to be extracted.

ATLAS measured  $\Gamma_s = 0.6703 \pm 0.0014(\text{stat.}) \pm 0.0018(\text{syst.}) \text{ ps}^{-1}$  from  $B_s^0 \rightarrow J/\psi \phi$  [Eur. Phys. J. C 81 \(2021\) 342](#) . This result combined with  $B_d^0 \rightarrow J/\psi K^*$  allowed us to determine the ratio  $\Gamma_d/\Gamma_s$ .

# ATLAS detector and feature important for this measurement

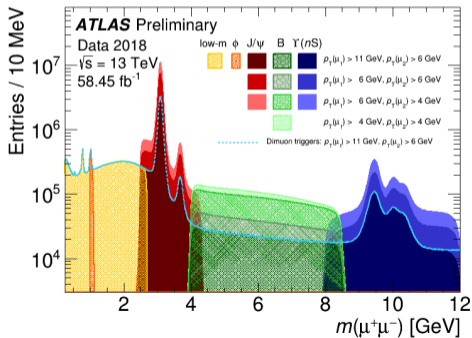


Time resolution of  $B_s^0 \rightarrow J/\psi\phi$  for different numbers of reconstructed PV in the same bunch crossing [ATL-PHYS-PUB-2013-010](https://arxiv.org/abs/2013.01010)

- Inner Detector: PIX, SCT and TRT,  $p_T > 0.5 \text{ GeV}$ ,  $|\eta| < 2.5$ 
  - Run2: new IBL 25% improvement of time resolution with respect to Run1.
  - Time, mass resolutions remain stable within increasing pileup in Run 2
- Muon Spectrometer: triggering ( $|\eta| < 2.4$ ), precision tracking ( $|\eta| < 2.7$ )

# Data-taking conditions for this analysis

- The measurement uses 2015-2018  $pp$  collision data at  $\sqrt{s} = 13$  TeV.
- Triggers:  $J/\psi \rightarrow \mu^+ \mu^-$ , the muon  $p_T$  thresholds varying:
  - 4 GeV, 6 GeV and 11 GeV.
  - Low  $p_T$  thresholds were activated in the end of fills when the instantaneous luminosity decreases.
- For events accepted for this analysis, an average number of  $pp$  interactions per bunch crossing (pile-up) was 31.
- No displaced  $J/\psi$  vertex cuts applied. Trigger tracking: transverse impact parameter  $d_0 < 10$ mm on all tracks.



2D unbinned maximum-likelihood fit applied simultaneously to mass and proper decay time of  $B_d^0$  candidates. Likelihood Function formula:

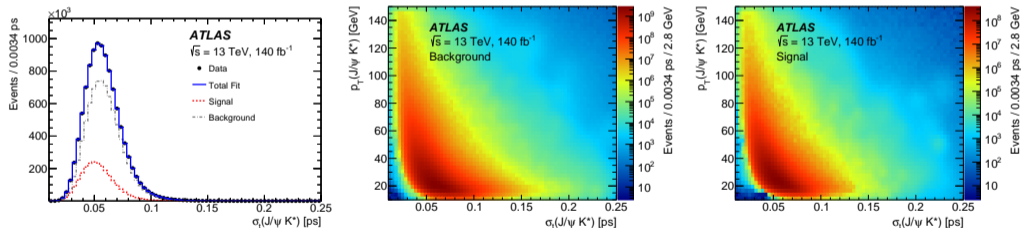
$$\ln L = \sum_{i=1}^N w(t_i) \ln [f_{\text{sig}} \mathcal{M}_{\text{sig}}(m_i) \mathcal{T}_{\text{sig}}(t_i, \sigma_{t_i}, p_{T_i}) + (1 - f_{\text{sig}}) \mathcal{M}_{\text{bkg}}(m_i) \mathcal{T}_{\text{bkg}}(t_i, \sigma_{t_i}, p_{T_i})].$$

- $f_{\text{sig}}$  - fraction of signal events in the total number of events,  $N$ .
- $\mathcal{M}_{\text{sig}}, \mathcal{T}_{\text{sig}}$  mass and time signal PDFs
- $\mathcal{M}_{\text{bkg}}, \mathcal{T}_{\text{bkg}}$  mass and time Background PDFs
- The mass  $m_i$ , the proper decay time  $t_i$ , its uncertainty  $\sigma_i$  and the  $B_d^0$  candidate transverse momentum  $p_{T_i}$  are the values measured from the data for each event  $i$ .
- Weight  $w_i$  accounts for event selection efficiency, based on MC.

# Conditional probability of the time uncertainty $\sigma_{t_i}$ and $p_{T_i}$ of $B_d^0$

- Per-candidate time errors  $\sigma_{t_i}$ , extracted from data in the vertex fit of each  $B_d^0$  are used for deconvolution of proper decay times  $t_i$ .
- $\sigma_{t_i}$  are different for signal and background, see fig Left. Same it true for  $p_{T_i}$  of  $B_d^0$ .
- To account for these differences, the 2D probability terms  $C_j(\sigma_{t_i}, p_{T_i})$  are introduced into Signal and Bg Time PDF terms  $\mathcal{T}$ . Method first used in [G.Punzi](#)

$$\mathcal{T}_j(t_i, \sigma_{t_i}, p_{T_i}) = P_j(t_i, \sigma_{t_i}, p_{T_i}) \cdot C_j(\sigma_{t_i}, p_{T_i}), \quad (2)$$



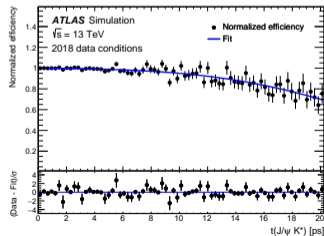
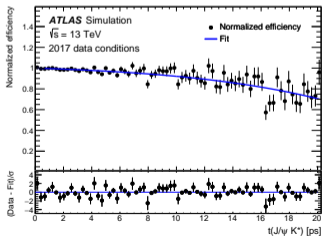
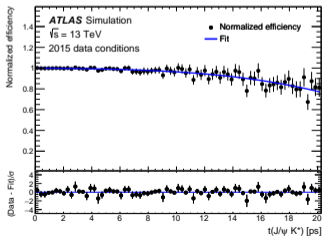
Left:  $\sigma_{t_i}$  for Signal and Bg. Middle and Right: 2D probability terms  $C_j(\sigma_{t_i}, p_{T_i})$  for Bg. and Signal. The sPlot technique [Pivk:2004ty](#) is used to separate Signal and Bg.

- Trigger, offline reconstruction, event selections bias the reconstructed proper-decay time distribution.
- Trigger and offline tracking impose  $|d_0| < 10$  mm, for all four final-state tracks of  $B^0 \rightarrow J/\psi K^{*0}$ , resulting in inefficiency at large times.
- Inefficiencies determined by signal MC, passed through simulation of detector response & triggers, offline tracking and event selections - as data.

Ratio of proper decay time distributions before and after the whole chain, were fitted by:

$$1/w(t_i) = p_0 \cdot [1 - p_1 \cdot (\text{Erf}((t_i - p_3)/p_2) + 1)] \quad (3)$$

Erf is error function,  $p_0, p_1, p_2, p_3$  determined in the fit. Weights  $w(t_i)$  are used to re-weight each event in Likelihood fit.



Source of uncertainty	Systematic uncertainty [ps]
ID alignment	0.00108
Choice of mass window	0.00104
Time efficiency	0.00130
Best-candidate selection	0.00041
Mass fit model	0.00152
Mass-time correlation	0.00229
Proper decay time fit model	0.00010
Conditional probability model	0.00070
Fit model test with pseudo-experiments	0.00002
Total	0.0035

- The statistical error of this  $B_d^0$  lifetime measurement is 0.0012 (stat.) ps.

- ID misalignment effects: dominated by global length scale biases originating from ID geometry radial and longitudinal distortions along track trajectory.
- These manifest themselves as a shift in the reconstructed masses of known resonances, e.g.  $J/\psi \rightarrow \mu^+ \mu^-$  [Eur. Phys. J. C 80 \(2020\) 1194](#) ↗.

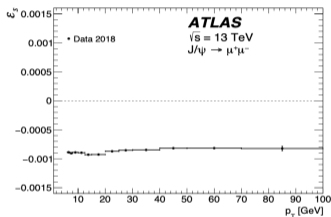


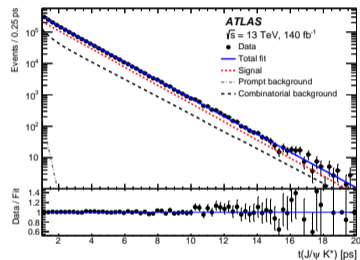
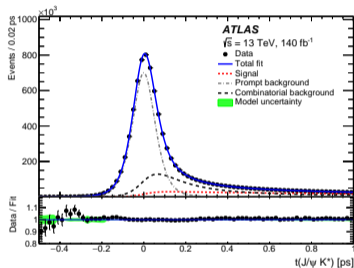
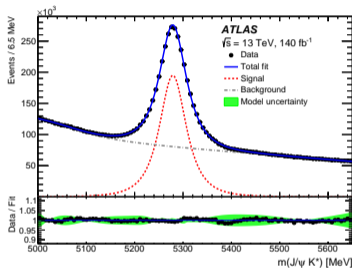
Fig showing a bias of the  $J/\psi \rightarrow \mu^+ \mu^-$  mass, measured as a relative difference from PDG value, as a function of a muon track  $p_T$ .

- In our analysis  $B_d^0$  tracks are re-fitted with the  $J/\psi$  mass constrained to PDG, effectively removing the misalignment effect. The impact of misalignment - estimated by alternative fit:  $B_d^0$  vertex re-fitted without PDG constrain.
- Additionally, to account for the momentum bias affecting hadrons from  $K^{*0}$  - their  $p_T$  are altered by  $-0.085\%$  [Eur. Phys. J. C 80 \(2020\) 1194](#) ↗.
- The two effects summed in quadrature, give the systematics:  $0.9 \sigma_{stat}$

# Results: The $B^0$ effective lifetime and the mass and time projections of the likelihood fit

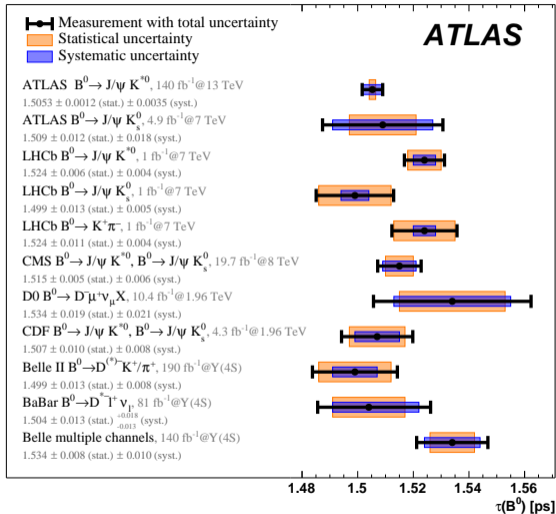
The  $B^0$  effective lifetime value measured with a total of  $2\,450\,500 \pm 2400$   $B^0 \rightarrow J/\psi K^{*0}$  signal events. The measured effective lifetime is

$$\tau = 1.5053 \pm 0.0012 \text{ (stat.)} \pm 0.0035 \text{ (syst.) ps.}$$




- Mass fit projection (left). Proper decay time fit projections in two different ranges:  $(-0.5; 1.0)$  ps (Middle) and  $(1; 20)$  ps (Right).
- Solid blue line - total fit, dashed red line - signal.
- The lower panels: ratio of data point to the fit value. The green band - the envelope of model variations included in the systematic uncertainty, the bars on the data points indicate statistical uncertainties. Plot -right - the model variation band too small to be visible.

# A comparison ATLAS $B^0$ lifetime with the latest results of other experiments.




- The current ATLAS result in  $B^0 \rightarrow J/\psi K^{*0}$  channel.
- The previous ATLAS result [Phys. Rev. D 87 \(2013\) 032002](#) in the  $B^0 \rightarrow J/\psi K_S^0$  channel.
- Latest LHCb results [JHEP 04 \(2014\) 114](#) in  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow J/\psi K_S^0$  decays and [Phys. Lett. B 736 \(2014\) 446](#) in the  $B^0 \rightarrow K^+ \pi^-$ .
- Latest CMS [Eur. Phys. J. C 78 \(2018\) 457](#) combined result for  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow J/\psi K_S^0$  decays.
- Tevatron experiments: D0 [Phys. Rev. Lett. 114 \(2015\) 062001](#) in the  $B^0 \rightarrow D^- \mu^+ \nu_\mu$  channel, and CDF [Phys. Rev. Lett. 106 \(2011\) 121804](#) with a combined result for  $B^0 \rightarrow J/\psi K^{*0}$  and  $B^0 \rightarrow J/\psi K_S^0$ .
- $e^+ e^-$  colliders: Belle II [Phys. Rev. D 107 \(2023\) L091102](#) in the  $B^0 \rightarrow D^{(*)-} K^+ / \pi^+$  channel and the last result from BaBar [Phys. Rev. D 73 \(2006\) 012004](#) in the  $B^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell$ . Belle [PhysRevD.71.07990, 2005](#) this combination includes  $B_d^0$  decays to  $D^{*-} \ell^+ \nu$ ,  $D^{*-} \pi^+$ ,  $D^- \pi^+$ ,  $D^{*-} \rho^+$ ,  $J/\psi K^{*0}$ ,  $J/\psi K_S^0$ .

## $\Gamma_d$


- We determine  $\Gamma_d$  from our measured effective lifetime  $\tau_{B^0}$ , using Eq (1) and input values  $2y = \Delta\Gamma_d/\Gamma_d = 0.001 \pm 0.010$  and asymmetry  $A = -0.578 \pm 0.136$  from [HFLAV:2023](#) :

$$\Gamma_d = 0.6639 \pm 0.0005 \text{ (stat.)} \pm 0.0016 \text{ (syst.)} \pm 0.0038 \text{ (ext.) ps}^{-1}$$

The uncertainty denoted 'ext.' originates from the HFLAV.

- The value  $\Gamma_d$  is in agreement with HQE theory of  $0.63_{-0.07}^{+0.11}$  ps<sup>-1</sup> [Lenz et al. 2023](#) 

## $\Gamma_d/\Gamma_s$

- Using  $\Gamma_s = 0.6703 \pm 0.0014$  (stat.)  $\pm 0.0018$  (syst.) ps<sup>-1</sup> measured by the ATLAS [Eur. Phys. J. C 81 \(2021\) 342](#)  the resulting  $\Gamma_d/\Gamma_s$  ratio is

$$\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022 \text{ (stat.)} \pm 0.0036 \text{ (syst.)} \pm 0.0057 \text{ (ext.)}$$

- The statistical, systematic and external uncertainties are propagated from the quantities above. In  $\Gamma_d/\Gamma_s$  - systematic uncertainties of the ATLAS measurements of  $\tau_{B^0}$  and  $\Gamma_s$  primarily come from different sources. They are therefore treated as uncorrelated.
- $\Gamma_d/\Gamma_s$  agrees with theory HQE and lattice QCD models prediction.

- ATLAS performed a measurement of  $B_d^0$  effective lifetime and the average decay width  $\Gamma_d$  using  $B^0 \rightarrow J/\psi K^{*0}$  events reconstructed from a  $140 \text{ fb}^{-1}$  data sample of  $pp$  collisions collected with the ATLAS detector during the  $\sqrt{s}=13 \text{ TeV}$  LHC run.
- The  $B^0$  effective lifetime is measured to be  $1.5053 \pm 0.0012(\text{stat}) \pm 0.0035(\text{syst}) \text{ ps}$ . This result is compatible with other experimental measurements and is the most precise measurement to date.
- The measured average decay width of the heavy and light  $B^0$  mass eigenstates is  $\Gamma_d = 0.6639 \pm 0.0005 \text{ (stat.)} \pm 0.0016 \text{ (syst.)} \pm 0.0038 \text{ (ext.) ps}^{-1}$ . This value is in good agreement with the theory prediction.
- The measured average decay width  $\Gamma_d$  is combined with the average decay width  $\Gamma_s$  measured previously by ATLAS to obtain the ratio  $\Gamma_d/\Gamma_s = 0.9905 \pm 0.0022 \text{ (stat.)} \pm 0.0036 \text{ (syst.)} \pm 0.0057 \text{ (ext.)}$ . This result is compatible with the theory predictions from HQE and lattice QCD calculations as well as with the experimental average.

Backup Slides - 1 More details on  $B^0 \rightarrow J/\psi K^{*0}$  ATLAS analysis

## $B^0 \rightarrow J/\psi K^{*0}$ candidates

- At least one  $J/\psi \rightarrow \mu^+ \mu^-$  with  $\chi^2/\text{ndof} < 10$ ; within mass window retaining 99.7%  $J/\psi$  candidates.
- $K^*$ : Out of two hypothesis  $K^+ \pi^- / K^- \pi^+$  the one closer to  $K^*$  PDG mass selected.
- $J/\psi$  and  $K^*$  - fitted to a common vertex, constrained by fixing di-muon mass to  $J/\psi$  PDG. Only  $\chi^2/\text{ndof} < 3$  retained.
- 10% events have multiple  $J/\psi K^*$  candidates (in average 2.1); the one with smallest  $\chi^2/\text{ndof}$  selected.

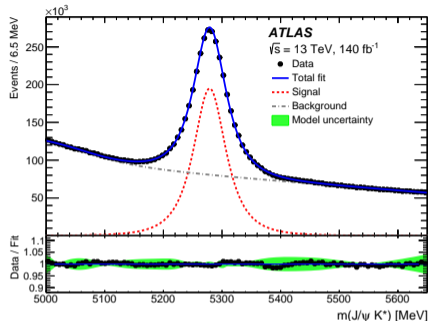
## Primary vertex (PV) selection

- For selected events the average pileup is 31: need to choose the best PV candidate where  $B_d^0$  is produced.
- PV positions are recalculated after removing any tracks used in the  $B_d^0$ .
- The PV candidate with the smallest 3D impact parameter,  $a_0$  (min. distance between PV and the line extrapolated from the  $B_d^0$  vertex in  $B_d^0$  momentum direction), is used.

For each  $B_d^0$  candidate, the proper decay time  $t$  is determined:

$$t = \frac{L_{xy} m_B}{p_{\Gamma B}}, \quad (4)$$

$p_{\Gamma B}$  is transverse momentum of  $m_B = \text{PDG } B_d^0$  mass. The transverse decay length,  $L_{xy}$ , is the distance in the transverse plane from PV to the  $B_d^0$  decay vertex, projected onto  $p_{\Gamma B}$ .



Signal mass is modelled with a Johnson  $S_U$ -distribution [Johnson](#)

$$\mathcal{M}_{\text{sig}}(m_i) = \frac{\delta}{\lambda \sqrt{2\pi} \sqrt{1 + \left(\frac{m_i - \mu}{\lambda}\right)^2}} \exp \left[ -\frac{1}{2} \left( \gamma + \delta \sinh^{-1} \left( \frac{m_i - \mu}{\lambda} \right) \right)^2 \right]$$

where  $\mu$ ,  $\gamma$ ,  $\delta$  and  $\lambda$  are free parameters of fit.

**Background** has two components: 1. **prompt** -  $J/\psi$  produced in  $pp \rightarrow J/\psi X$ , combined with random  $K^*$ . 2. **combinatorial** -  $J/\psi$  from any  $b$ -hadron decay combined with random  $K^*$ . They are modelled by sum of linear and sigmoid functions:

$$\mathcal{M}_{\text{bkg}}(m_i) = f_{\text{poly}}(1 + p_0 \cdot m_i) + (1 - f_{\text{poly}}) \left( 1 - \frac{s(m_i - m_0)}{\sqrt{1 + (s(m_i - m_0))^2}} \right)$$

where  $f_{\text{poly}}$  relative size of the two components and  $m_0$ ,  $s$  and  $p_0$  are parameters of the fit.

**Signal PDF:** exponential function convolved by Resolution function  $R$ .

$$P_{\text{sig}}(t_i | \sigma_{t_i}, p_{T_i}) = E(t', \tau_{B^0}) \otimes R(t' - t_i, \sigma_{t_i})$$

$E(t', \tau_{B^0}) = (1/\tau_{B^0}) \exp(-t'/\tau_{B^0})$  for  $t' \geq 0$ , with  $\tau_{B^0}$  - the fitted  $B_d^0$  lifetime.

**Background PDF:**

$$P_{\text{bkg}}(t_i | \sigma_{t_i}, p_{T_i}) = \left( f_{\text{prompt}} \cdot \delta_{\text{Dirac}}(t') + (1 - f_{\text{prompt}}) \sum_{k=1}^3 b_k \prod_{l=1}^{k-1} (1 - b_l) E(t', \tau_{\text{bkg}_k}) \right) \otimes R(t' - t_i, \sigma_{t_i})$$

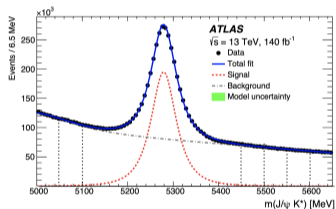
Dirac function  $\delta_{\text{Dirac}}$ : direct background; 3 exponentials  $E(t', \tau_{\text{bkg}_k})$ : components of combinatorial background;  $\tau_{\text{bkg}_k}$ , and fractions:  $f_{\text{prompt}}$  and  $b_k$  are free parameters of fit.

**Resolution function  $R$ :** modelled as a sum of three Gaussian distributions with widths:  $S^{(k)} \sigma_{\tau_i}$

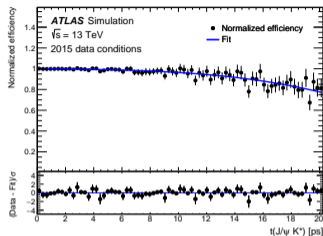
$$R(t' - t_i, \sigma_{t_i}) = \sum_{k=1}^3 f_{\text{res}}^{(k)} \frac{1}{\sqrt{2\pi} S^{(k)} \sigma_{t_i}} \exp\left(\frac{-(t' - t_i)^2}{2(S^{(k)} \sigma_{t_i})^2}\right)$$

three  $f_{\text{res}}^{(k)}$  fractions and scale factors  $S^{(k)}$  are free parameters of fit,  $\sigma_{\tau_i}$  is the per-candidate time error, extracted from data in the vertex fit of each  $B^0 \rightarrow J/\psi K^{*0}$ .

- Correlations between the invariant mass and the pseudo-proper decay time, and their potential impact on the fit results, are studied.
- No mass-time correlation in signal PDF proven by MC.
- For background the correlations are studied in data. First, applying the default time Background PDF  $P_{\text{bkg}}(t_i, \sigma_i, \rho_{T_i})$  fit in each of the 6 mass sideband bins in data. Fractions  $f_{\text{prompt}}$  and  $b_1, b_2$  determined in each bin were then fitted to extract their dependence on mass. The best description - achieved by linear functions.



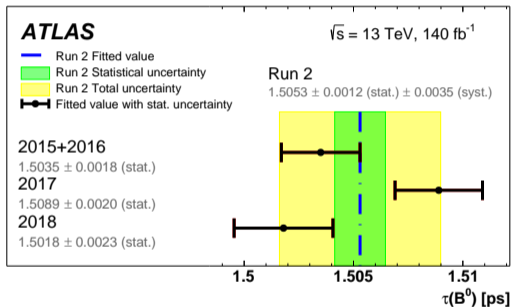
- Based on this information, an alternative fit model, in which the background PDF term  $P_{\text{bkg}}(t_i, \sigma_i, \rho_{T_i})$  is constructed with the parameters accounting for mass-dependence:  
 $f_{\text{prompt}}(m_i) = a + b(m_i - 5.279 \text{ GeV})$ ,  $b_1(m_i) = c_1 + d_1(m_i - 5.279 \text{ GeV})$  and  
 $b_2(m_i) = c_2 + d_2(m_i - 5.279 \text{ GeV})$ .  $a, b, c_1, d_1, c_2$  and  $d_2$  are free parameters of the fit. A difference of lifetime w.r.t default fit is  $0.00228 \text{ ps}$  ( $1.9 \sigma$ )



- Two alternative functions fitting the time efficiency histogram, replacing the default error function, were: hyperbolic tangent function or  $(x^2 + 1)^{-1/2}$ . The larger of the changes in the fitted lifetime, of size  $0.5\sigma_{\text{stat}}$  - taken as a systematic uncertainty.
- Systematic effects due to limited MC statistic, used to build the time efficiency histogram, was estimated by repeating the fit with a large number of alternative time efficiency functions, obtained by smearing the number of MC events in the time bins, leading to  $0.5\sigma_{\text{stat}}$ .
- An alternative fit using data events with  $t < 8$  ps, is performed to validate the modelling of the efficiency for high lifetimes, where the efficiency decrease is large, lead to small systematics  $0.5\sigma_{\text{stat}}$ .





## Results: A consistency and stability test over data taking periods

As a consistency and stability test, the  $B^0$  lifetime value was fitted separately for each data-taking period (2015+2016, 2017 and 2018). Figure shows the degree of stability over time. The  $p$ -value for consistency of the three individual results, accounting for just statistical uncertainties is 0.038.



The fitted values of the  $B^0$  lifetime, measured with  $B^0 \rightarrow J/\psi K^{*0}$  decays, for the 2015+2016, 2017 and 2018 subsamples compared to the value for the whole sample. The  $B^0$  lifetime value for each subsample is shown by a black point, with the error bar indicating the statistical uncertainty.

## Backup slides - 2 Other B lifetimes measured in ATLAS

- Measurement of effective lifetime in  $B_{(s)}^0 \rightarrow \mu^+ \mu^-$   
[JHEP 09 \(2023\) 199 W](#) 
- Measurement of  $\Lambda_b^0$  lifetime in  $\Lambda_b^0 \rightarrow J/\psi \Lambda^0$  and  $B_d^0$  lifetime in  $B_d^0 \rightarrow J/\psi K_S^0$  decays  
[PhysRevD.87.032002, 2013](#) 
- Measurement of the relative width difference  $\Delta\Gamma_d/\Gamma_d$  of the  $B_d^0$ - $\overline{B}_d^0$  system with the ATLAS detector  
[JHEP06\(2016\)081](#) 
- Measurement of the CP-violating phase, the width difference  $\Delta\Gamma_s$  between the meson mass eigenstates and the average decay width  $\Gamma_s$  in the  $B_s \rightarrow J/\psi \phi$  decay.  
[Eur. Phys. J. C 81 \(2021\) 342](#) 

- In Standard Model (SM) only the  $CP$ -odd heavy-mass eigenstate in the  $B_s^0 - \bar{B}_s^0$  pair decays into  $\mu^+ \mu^-$   
M. Beneke JHEP 10 (2019) 232 and errata JHEP 11 (2022) 099.
- Beyond the Standard Model (BSM) such as minimal supersymmetric Standard Model extensions D. M. Straub II Nuovo Cimento C 35 (2012) 249 can potentially perturb the effective lifetime in  $B_{(s)}^0 \rightarrow \mu^+ \mu^-$  decays. These perturbations can be significant also in absence of measurable BSM effects on the  $B_{(s)}^0 \rightarrow \mu^+ \mu^-$  branching fraction (BR).

- The effective  $B_{(s)}^0 \rightarrow \mu^+ \mu^-$  lifetime is defined as

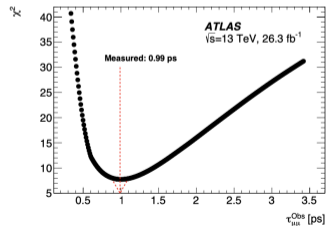
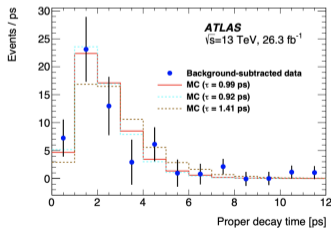
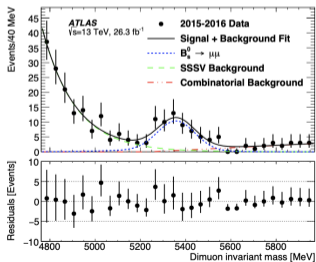
$$\tau_{\mu\mu} = \frac{\int_0^\infty t \Gamma(B_s^0(t) \rightarrow \mu\mu) dt}{\int_0^\infty \Gamma(B_s^0(t) \rightarrow \mu\mu) dt}, \text{ where: } \Gamma(B_s(t) \rightarrow \mu\mu) = \Gamma(B_s^0(t) \rightarrow \mu\mu) + \Gamma(\bar{B}_s^0(t) \rightarrow \mu\mu)$$


and  $t$  is the proper decay time of the  $B_s^0$  and  $\bar{B}_s^0$  mesons.

- In the SM hypothesis  $\tau_{\mu\mu}$  coincides with the lifetime of the heavy  $B_s^0$  eigenstate  $\tau_{B_s^H}$ .
- The experimental average of the  $B_s^0 - \bar{B}_s^0$  lifetimes and their difference Phys. Rev. D 107 (2023) 052008 yields the prediction  $\tau_{\mu\mu}^{SM} = (1.624 \pm 0.009)$  ps, with new physics effects perturbing it at most by the difference between the heavy and light eigenstate lifetimes (0.193 ps).

# Measurement of effective lifetime in $B_{(s)}^0 \rightarrow \mu^+ \mu^-$ , cont 1

- Data from 2015-2016 are used in this measurement.
- Un-binned maximum likelihood fit to candidates in the [4766 – 5966] MeV mass region (Left Fig), yielding  $58 \pm 13$  (stat. only)  $B_{(s)}^0 \rightarrow \mu^+ \mu^-$  signal events.
- Signal and backgrounds weights calculated from the result of the mass fit are used to construct the proper decay time histogram - background-subtracted employing per-event weights calculated according to the *sPlot* technique (Middle Fig).
- The lifetime measurement is obtained by minimising the binned  $\chi^2$  between the data histogram and lifetime-dependent pure signal MC templates extracted from MC simulated samples, as illustrated in the Middle and Right Fig.



- The statistical uncertainty is derived from the Neyman CL band construction Figure Left. The  $\chi^2$  minimum and the Neyman belt construction yield  $\tau_{\mu\mu} = 0.99_{-0.07}^{+0.42}$  (stat) ps. The imbalance between positive and negative statistical uncertainties is already suggested by the asymmetry in the  $\chi^2$  scan.
- The systematic errors are dominated by data-MC discrepancies, followed by uncertainties in backgrounds lifetime modelling.
- The result using 2015-2016 data is  $\tau_{\mu\mu} = 0.99_{-0.07}^{+0.42}$  (stat.)  $\pm 0.17$  (syst.) ps. It is consistent with the SM prediction  $\tau_{\mu\mu}^{\text{SM}} = (1.624 \pm 0.009)$  ps [Phys. Rev. D 107 \(2023\) 052008](#)  as well as with the other available experimental results.

