

Indirect Characterization of Photoinjector Parameters & their Correlations

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Advanced Accelerator Concepts 2026

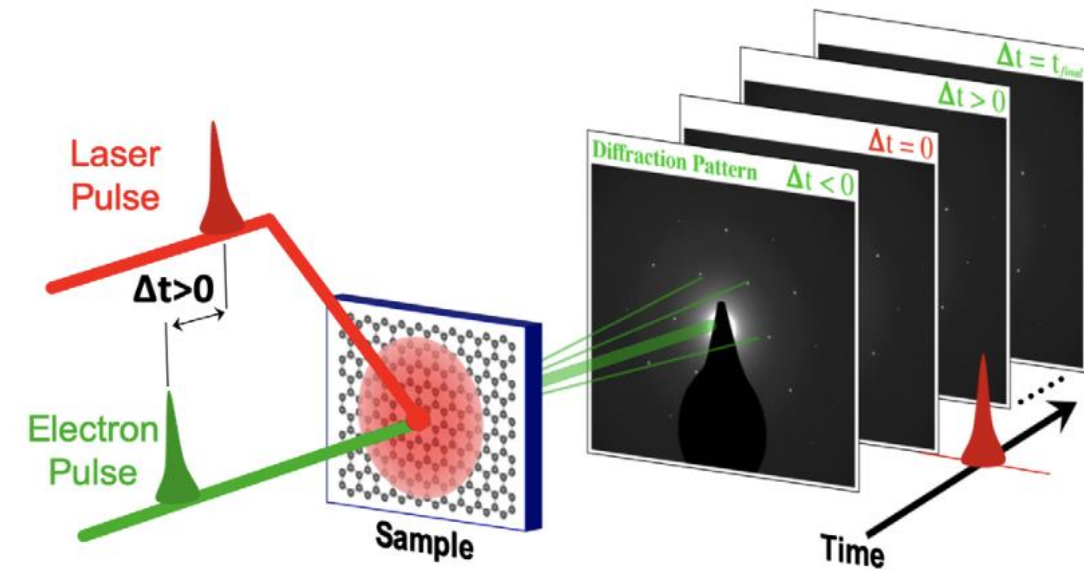
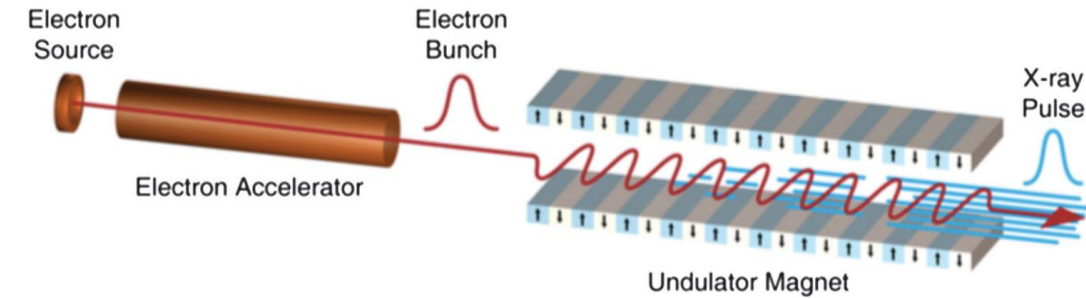
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The Center for
**BRIGHT
BEAMS**
A National Science Foundation
Science & Technology Center

- ❑ Motivation
- ❑ Measurement of the cathode re-entrance and intrinsic emittance
 - Degeneracies
- ❑ Surrogate model + MCMC technique
- ❑ Extracting the photocathode intrinsic emittance from a beam with:
 - Projected emittance growth
 - Space charge emittance growth
- ❑ Towards a measurement of the cathode time response
- ❑ Conclusion & Outlook

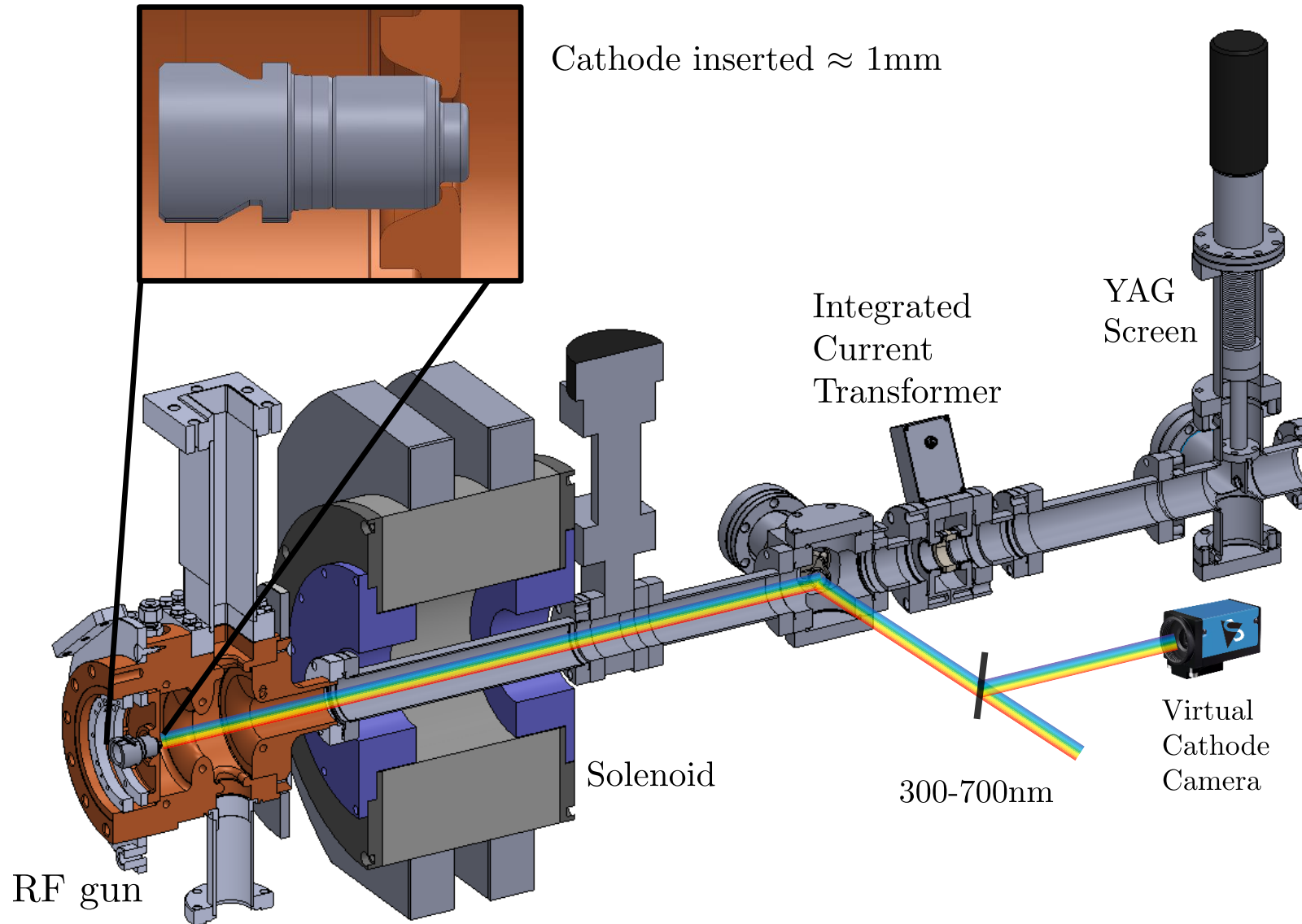
- ❑ Electron sources for next generation FELs & UED machines must be well tuned in order to provide high brightness beams
- ❑ Characterization of hard to measure parameters often requires additional diagnostics
 - Intrinsic emittance: apertures/slits
 - Pulse length: transverse deflecting cavity



Pegasus S-band Gun



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- ❑ OPA produces a short pulse tunable through the visible and UV spectrum
- ❑ Photoelectrons accelerated to 3MeV & focused using solenoid
- ❑ Charge, laser energy & transverse distribution at both the cathode & YAG screen are recorded

RF gun

Solenoid

Integrated
Current
Transformer

YAG
Screen

Virtual
Cathode
Camera

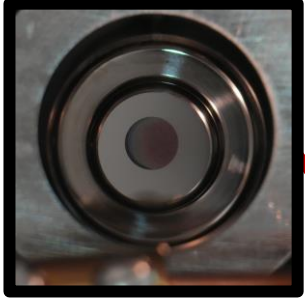
300-700nm

Semiconductor Photocathodes

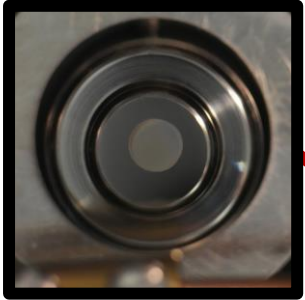


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UC04: Cs₃Sb



UC01: Cs-Te

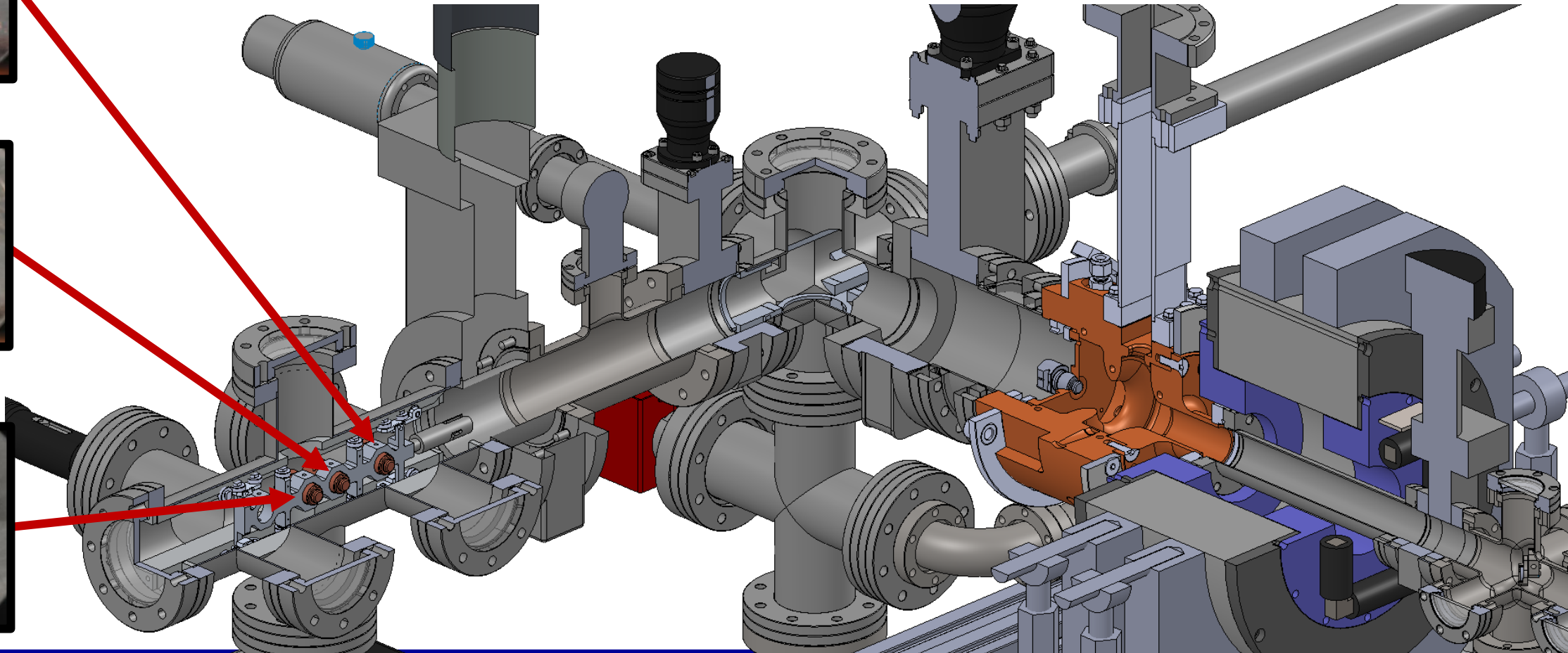


UC03: Cs₃Sb



□ Load-lock chamber behind RF permits exchange of air sensitive cathode under UHV conditions (<1e-9 Torr)

- 1 Cs-Te & 2 Cs₃Sb cathodes grown at SLAC & transported to UCLA

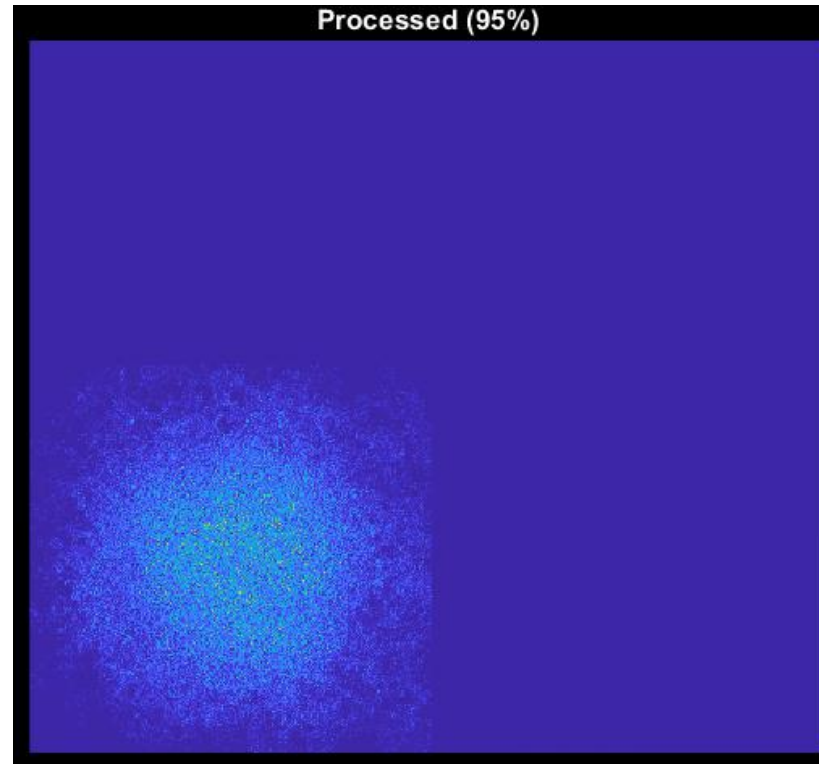
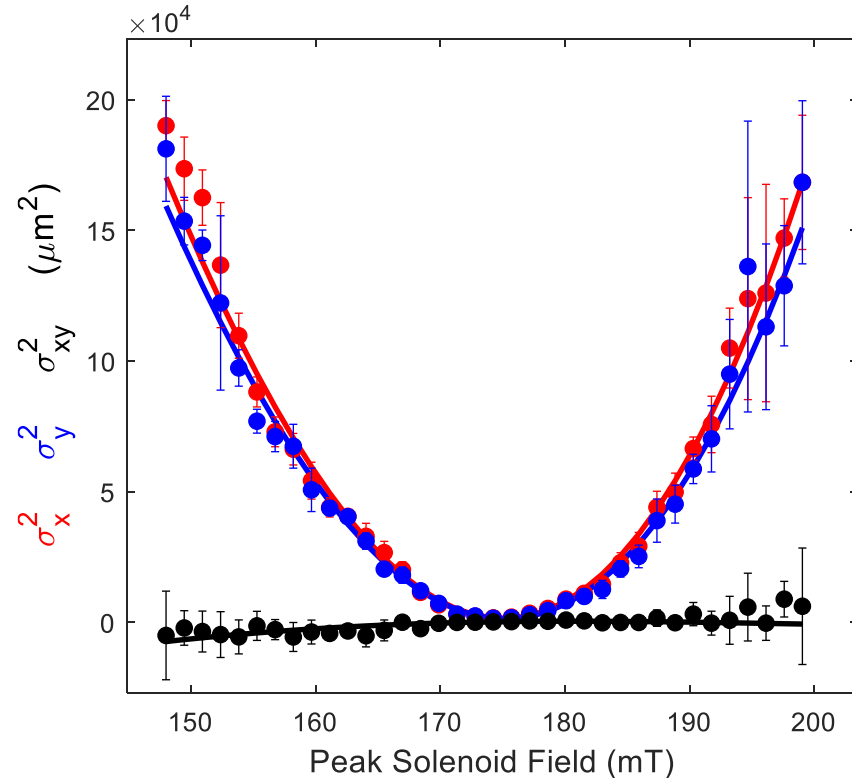


Analytic Analysis of Solenoid Scan

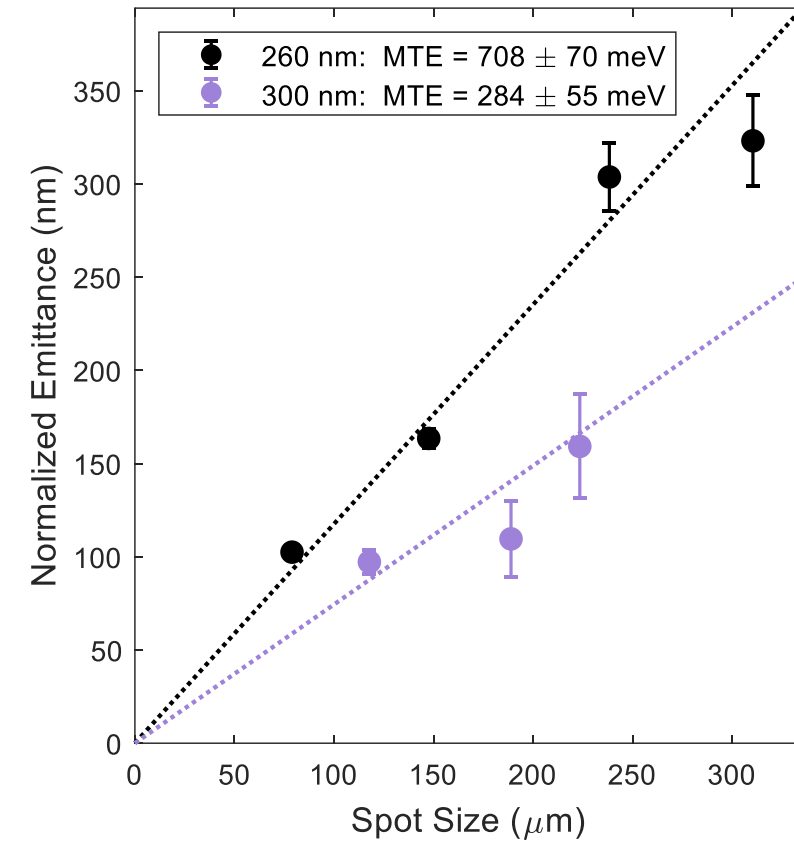
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95% emittance extracted from solenoid scan using linear transport formalism

- Low charge (<100fC) and short pulse (<1ps) minimize emittance growth from Coulomb force and RF field



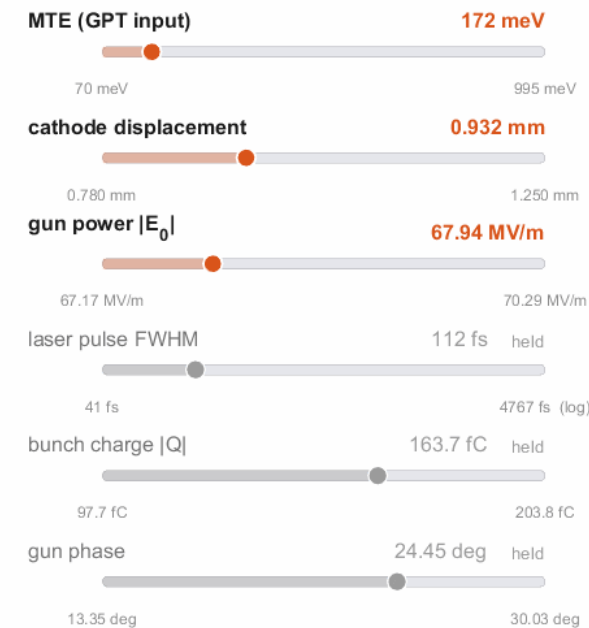
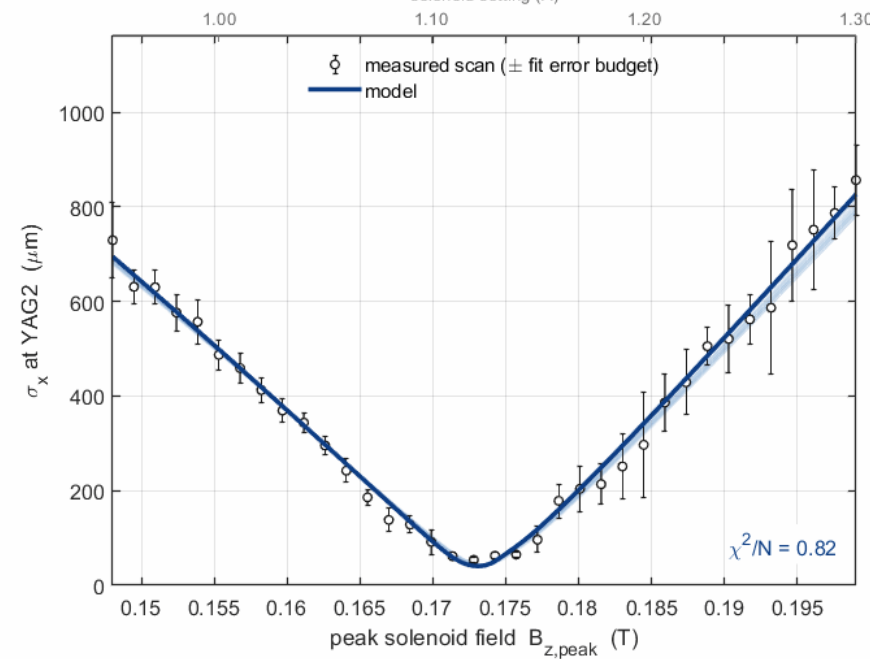
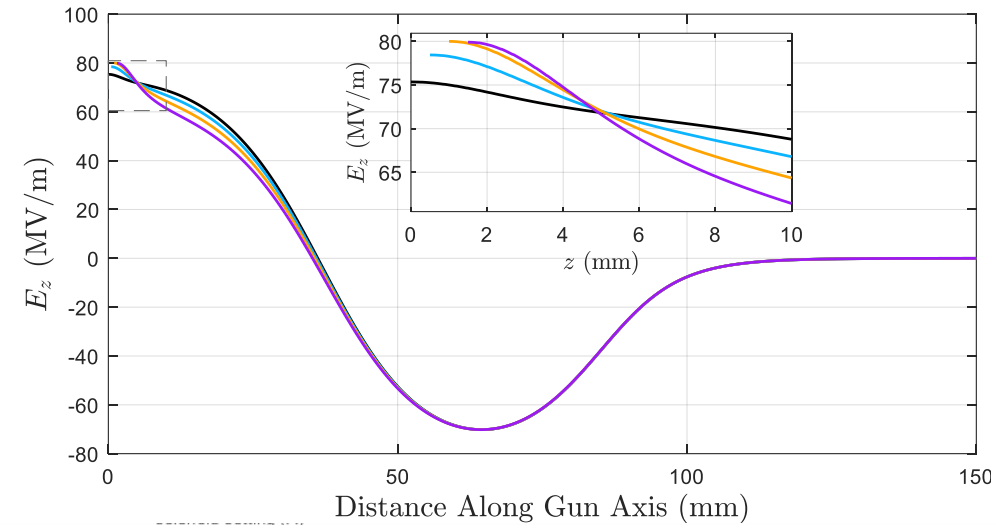
$$\varepsilon_n = \sigma_x \sqrt{\frac{\text{MTE}}{mc^2}}$$



Surrogate Model & Degeneracies

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- ❑ RF focusing from the re-entrance of the cathode plug is degenerate with the MTE
- ❑ Neural Net trained on a library of General Particle Tracer simulations that span the input space of the RF gun gives a fast surrogate model of the RF gun
 - Laser shape & duration (FWHM)
 - MTE & charge
 - Gun phase & power
 - Cathode plug position



- Markov Chain Monte Carlo (MCMC) methods use Bayesian inference to retrieve posterior probability distributions of the model parameters θ_M given the data set D
 - Information on measured parameters (charge, gun phase, ...) enter through prior probabilities

$$\underbrace{P(\theta_M | D)}_{\text{Posterior probability}} \propto \underbrace{P(D | \theta_M)}_{\text{Likelihood}} \underbrace{P(\theta_M)}_{\text{Prior probability}}$$

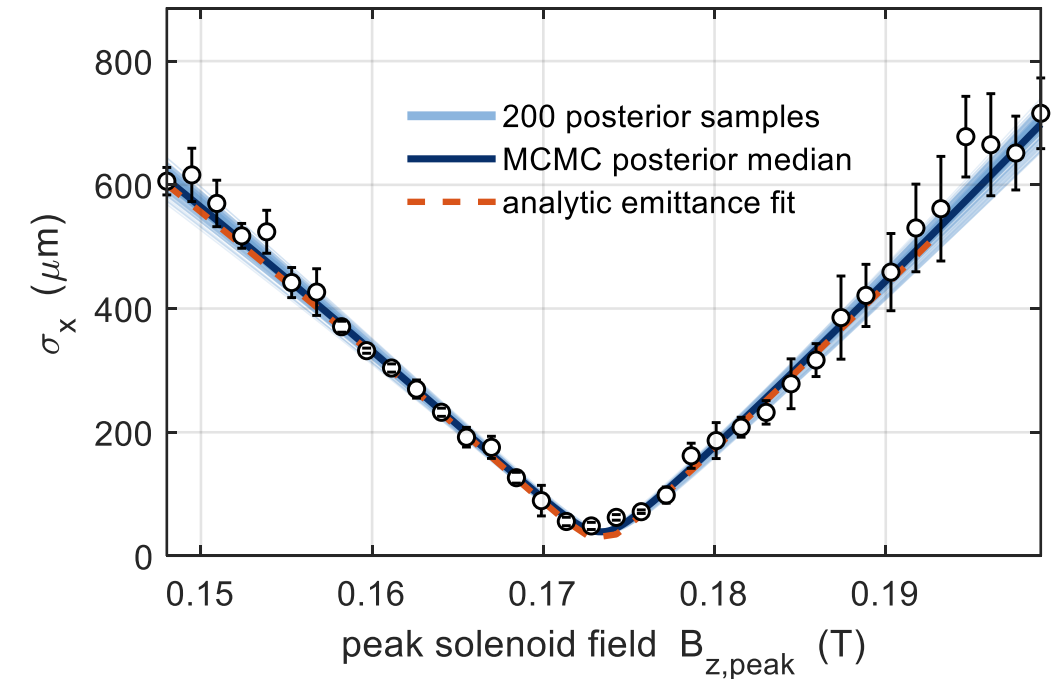
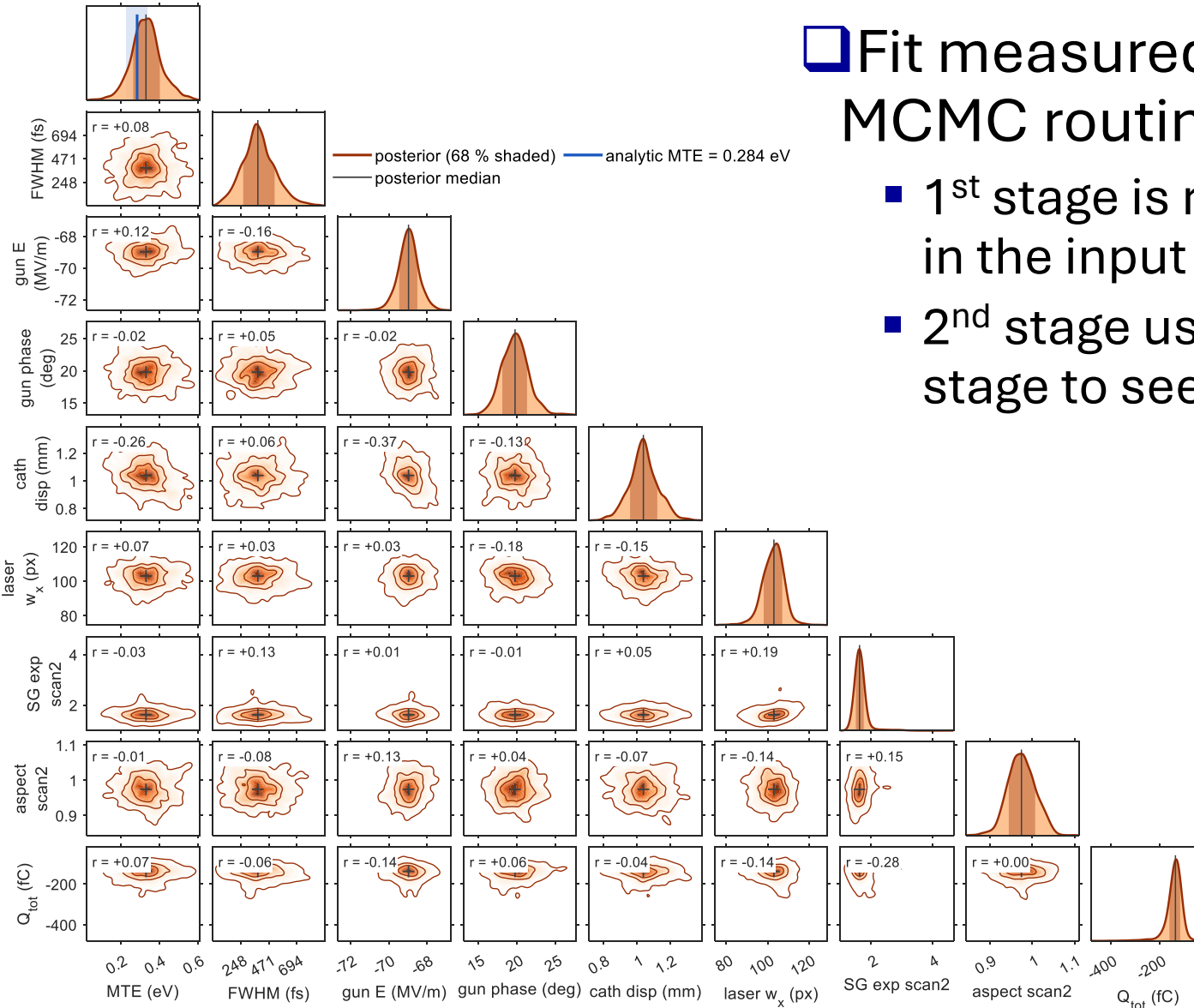
- GW-MCMC algorithm uses an ensemble of simultaneous walkers
 - Walkers talk to each other & give collective effects
 - Affine invariance means the algorithm works well on oddly shaped spaces

MCMC results



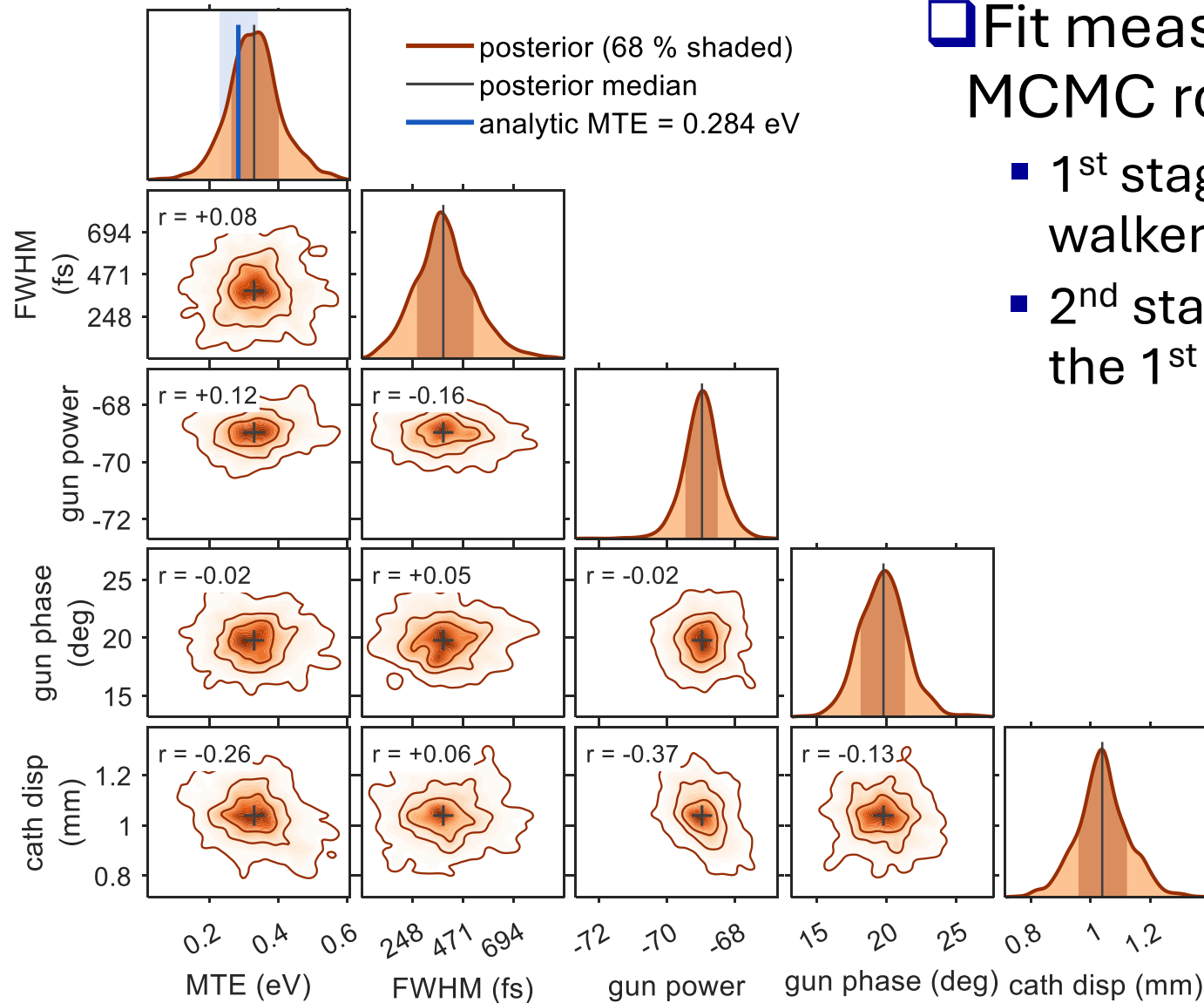
Fit measured data from Cs-Te with a 2 stage MCMC routine:

- 1st stage is run with a uniform distribution of walkers in the input space
- 2nd stage used regions of high likelihood from the 1st stage to seed 1/2 of the walkers



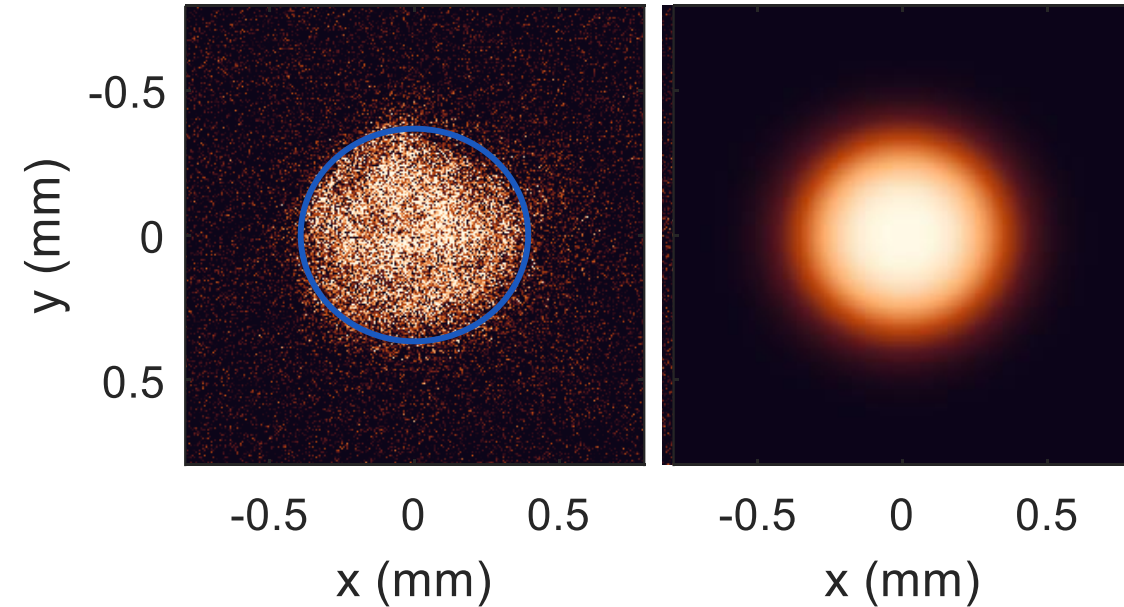
MCMC results

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Fit measured data from Cs-Te with a 2 stage MCMC routine:

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Space Charge Emittance Growth

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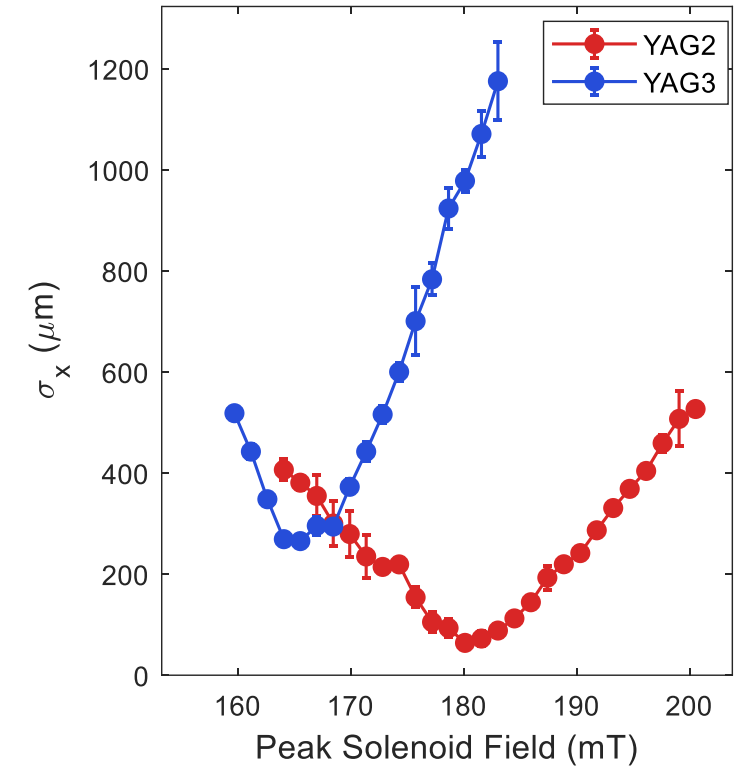
❑ Space charge emittance growth is only negligible for a small set of circumstances

- Our far YAG screen cannot reconstruct the MTE because the beam is too dense for too long
- Beam out of VHF guns is too soft (<1MeV)

❑ Rho parameter quantifies when a beam is emittance dominated ($\rho < 1$) or space charge dominated ($\rho > 1$)

$$\rho = \frac{I \sigma_r^2}{2 I_o \beta \gamma \epsilon_n^2}$$

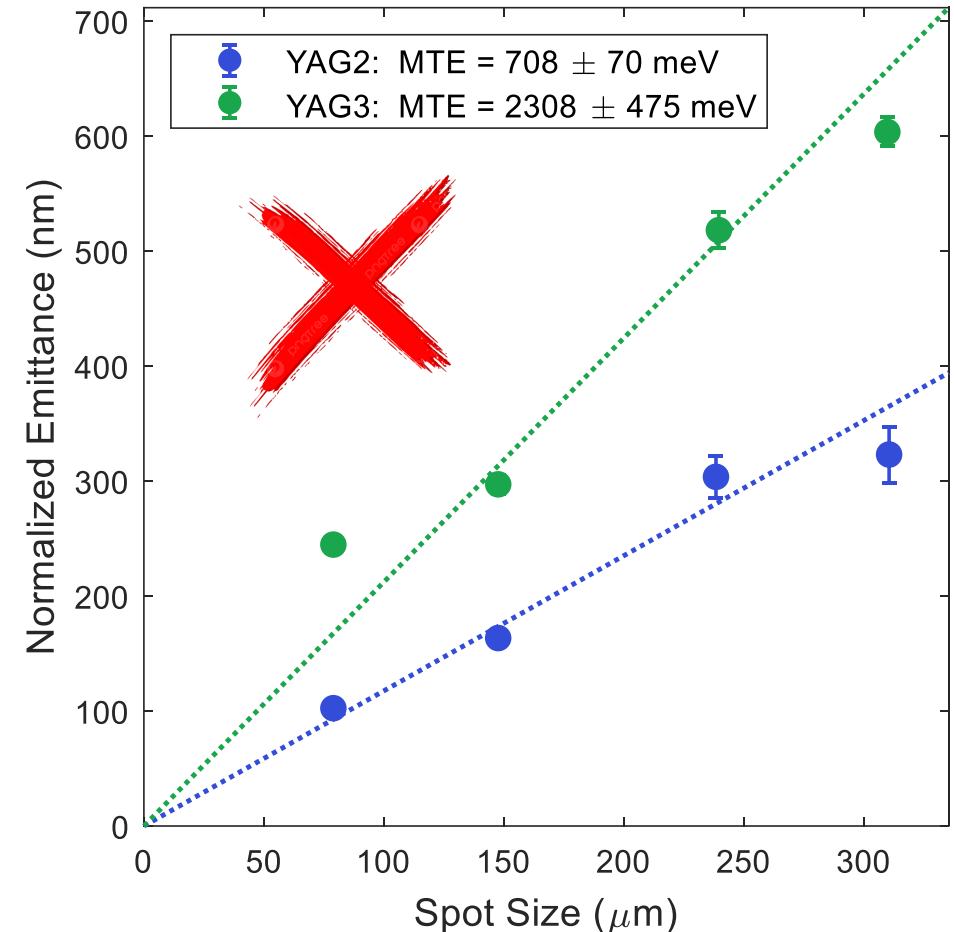
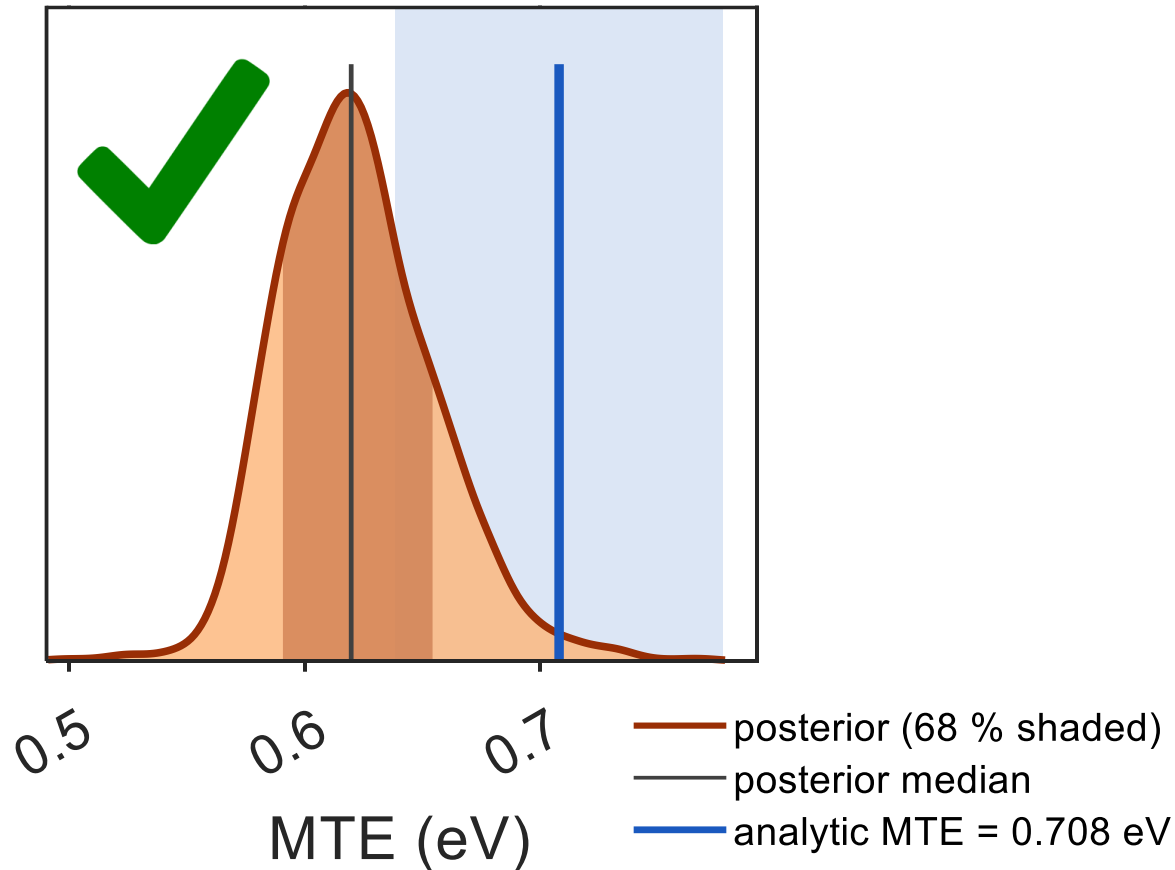
Screen	Position (m)	ρ at waist	Regime
YAG 2	1.3	0.42	Emittance
YAG 3	2.8	3.41	Space charge



Intrinsic Emittance with Space Charge



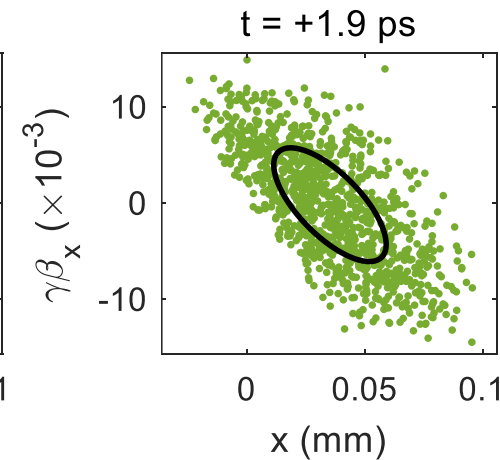
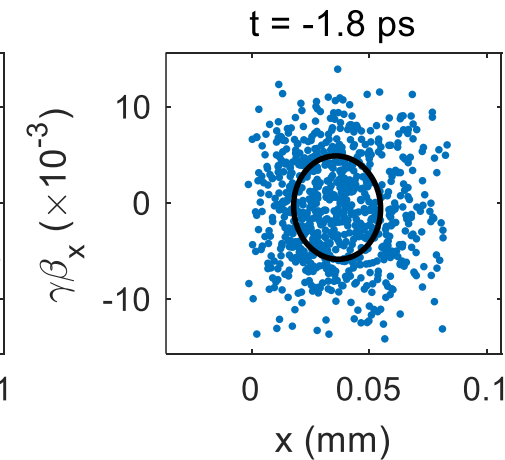
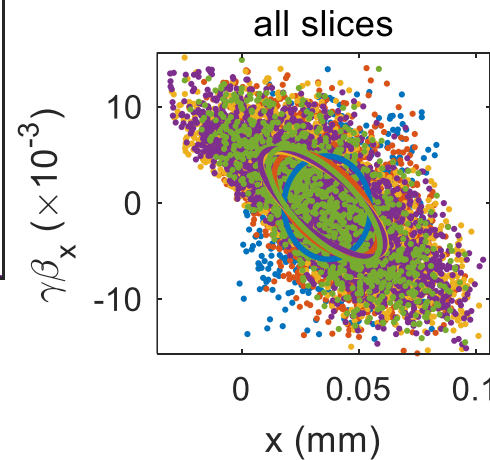
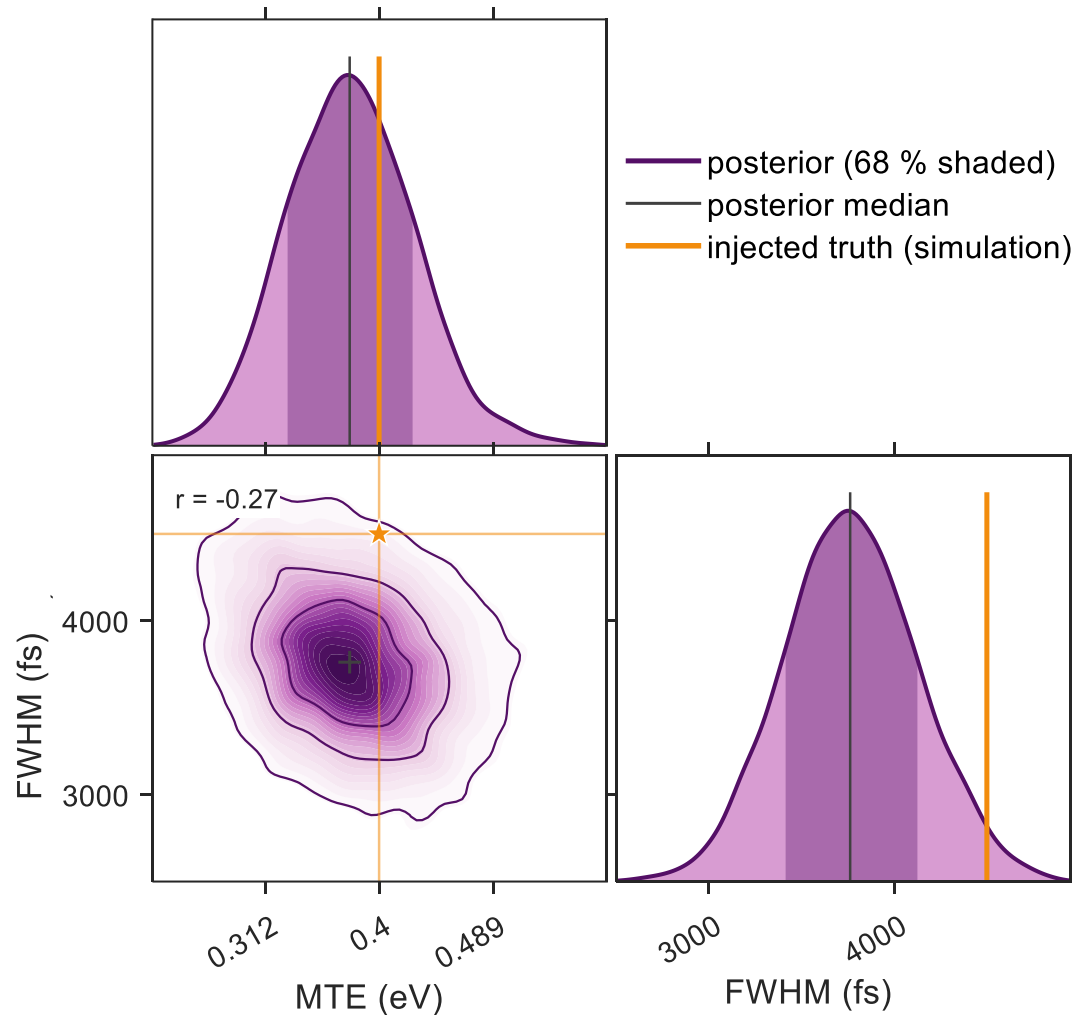
- Analytic reconstruction of the MTE fails for YAG3, but the MCMC routine can retrieve the MTE for scans taken on YAG3



Intrinsic Emittance with a Long Pulse

□ Simulate stretching the laser from 100 fs to 4.5 ps

- Suppress nonlinear photoemission near threshold
- The RF gun will give a time varying transverse kick, resulting in a projected emittance growth for long beams

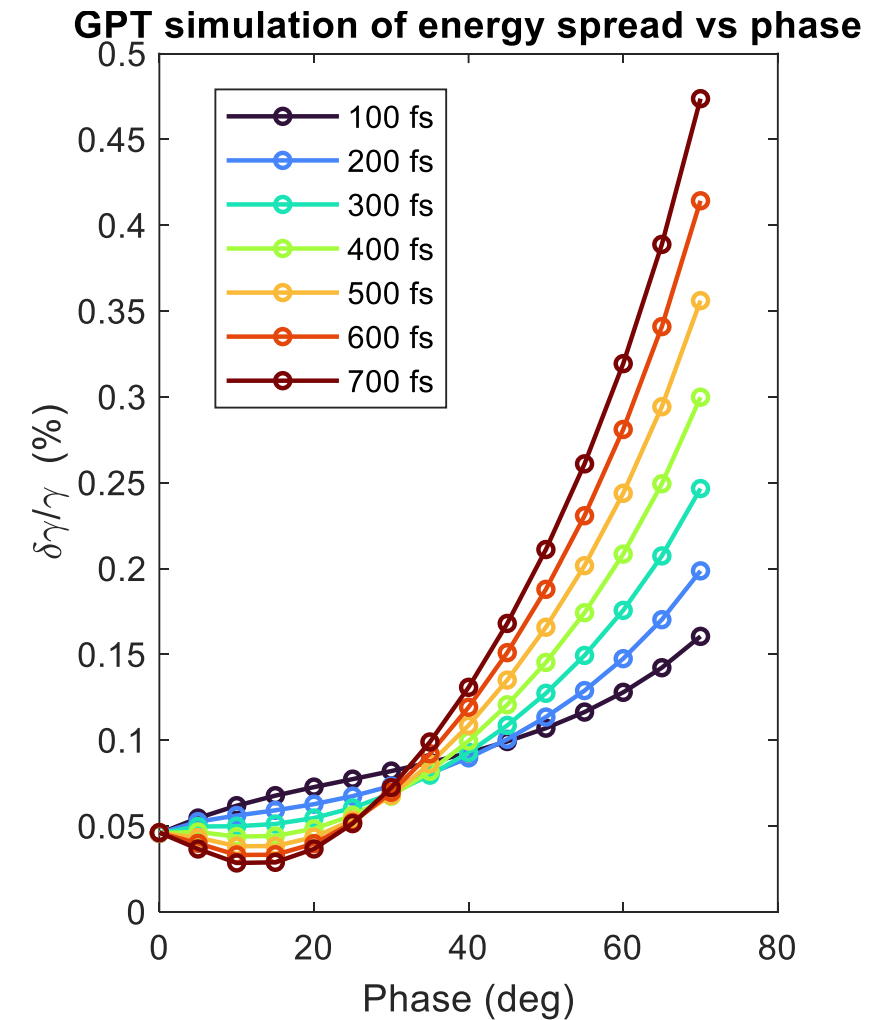
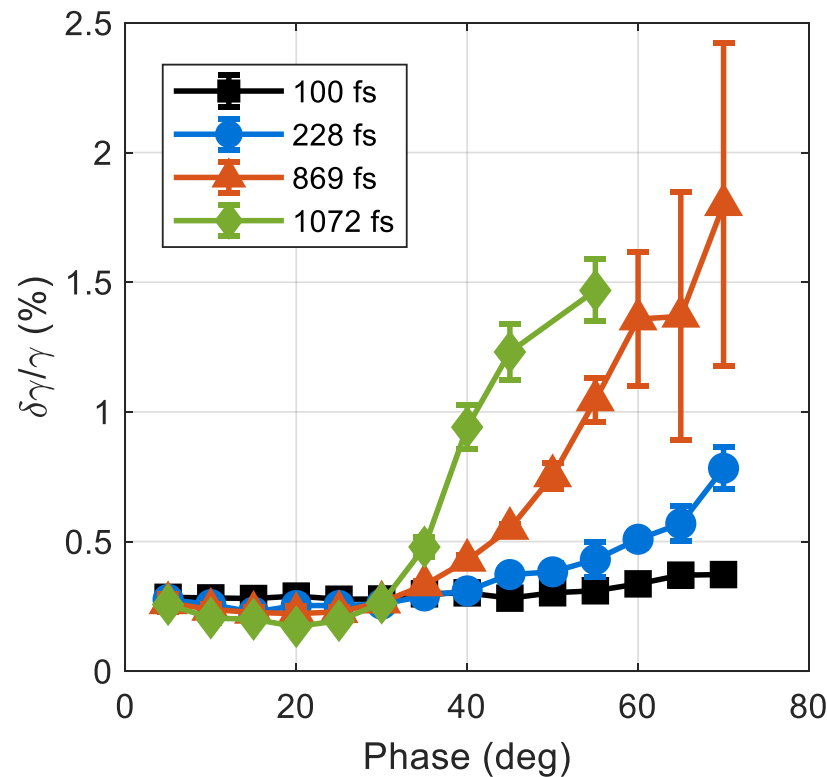
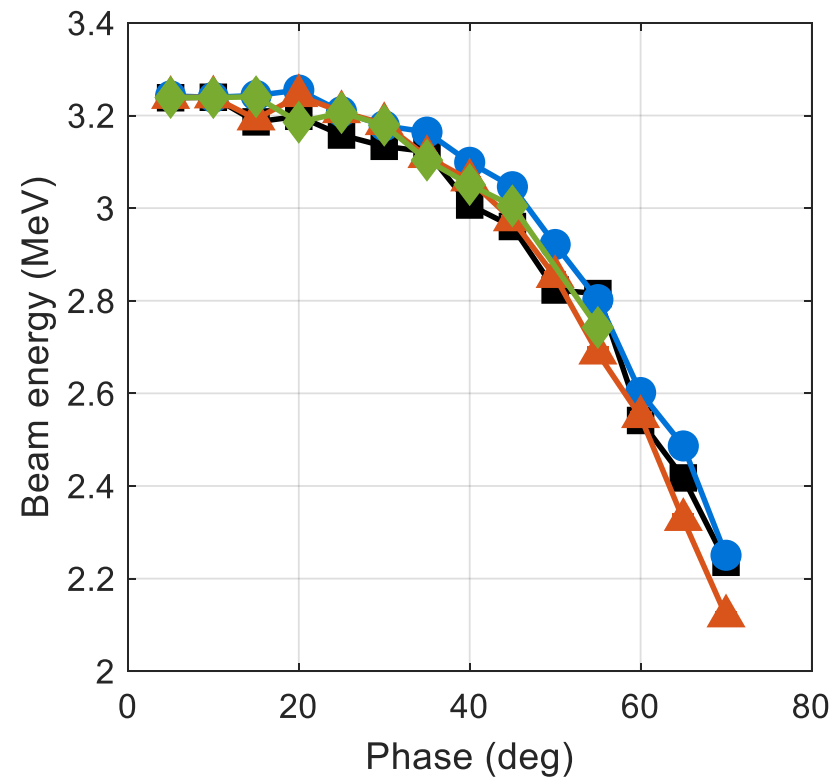


Indirect Measure of Pulse Length

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□ The energy of the electron beam leaving the RF gun is correlated with the launch phase

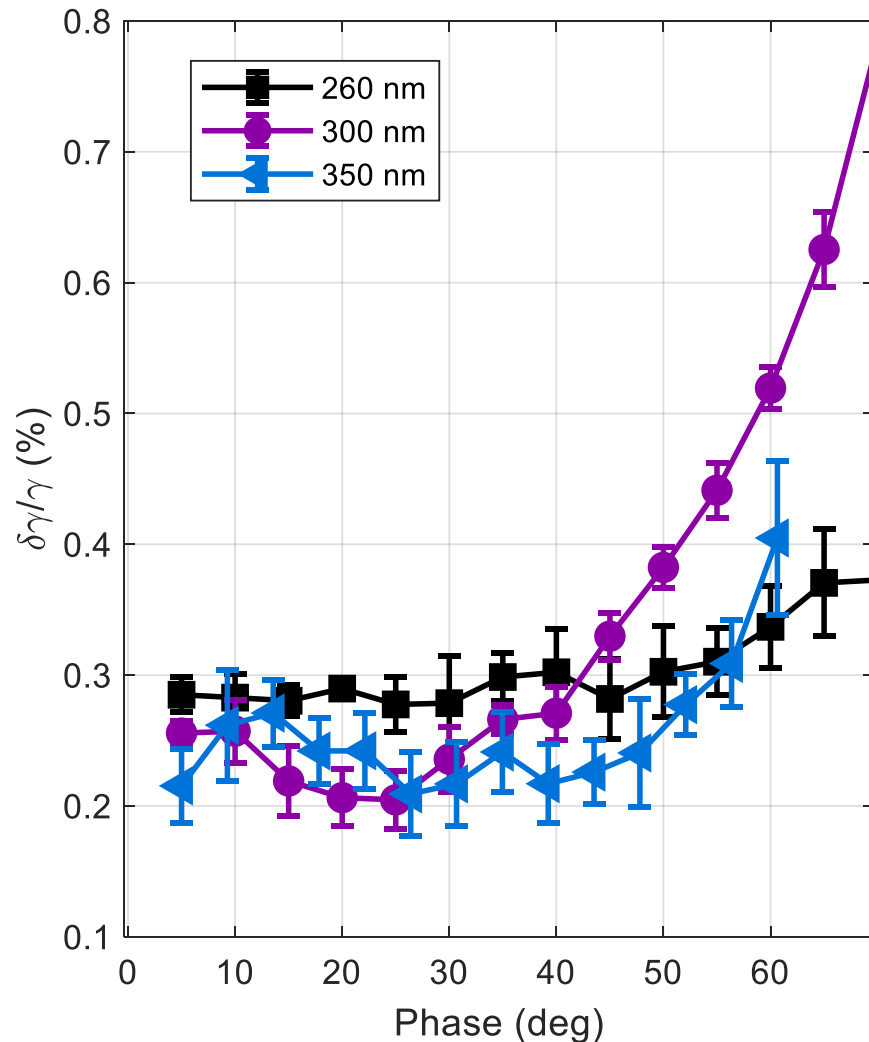
- A longer pulse will sample many different energies, especially at high phase, resulting in a larger energy spread



Time Response vs. Wavelength

- The FWHM (τ_{e-}) of the electron beam emitted at each wavelength is fit using GPT simulations to calculate the time response (τ_{TR})
- Temporal broadening near threshold transitions to prompt emission when nonlinear

$$\tau_{e-} = \sqrt{\tau_{\text{laser}}^2 + \tau_{TR}^2}$$



Wavelength (nm)	Laser Pulse Length (fs)	Electron Pulse Length (fs)	Time Response (fs)
260	100 ± 10	450 ± 20	460 ± 40
300	50 ± 20	1000 ± 200	985 ± 120
350	50 ± 20	$300 - 700$	$350 - 700$

- ❑ We plan to use these techniques to characterize advanced semiconductor cathodes in the Pegasus RF gun to push towards a 6D characterization of photocathode performance in an RF gun
 - Most useful for long pulse analysis close to threshold
 - Next batch of cathodes coming to UCLA early fall from both SLAC & ASU
- ❑ Future expansions
 - More accurate treatment of the initial charge distribution
 - Train NN that spans a larger input space

Acknowledgments



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□ PBPL

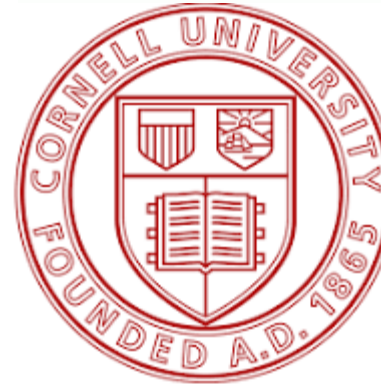
- Advisor: Pietro Musumeci
- Support from post docs & students
- Past PBPL students who have worked on this project: Eric Cropp, Gilles Silvi

□ Collaborators who have or will grow cathodes for testing at UCLA

- SLAC: T. Vecchione & G. Adhikari
- Cornell: Maxson Lab
- ASU: Karkare Lab

□ Funding

- CBB/NSF
- DOE



Boersch Effect



□ Coulomb repulsion of the electron beam will increase the energy spread if too dense

■ Correlation plots show this effect becomes negligible as the beam is stretched

