

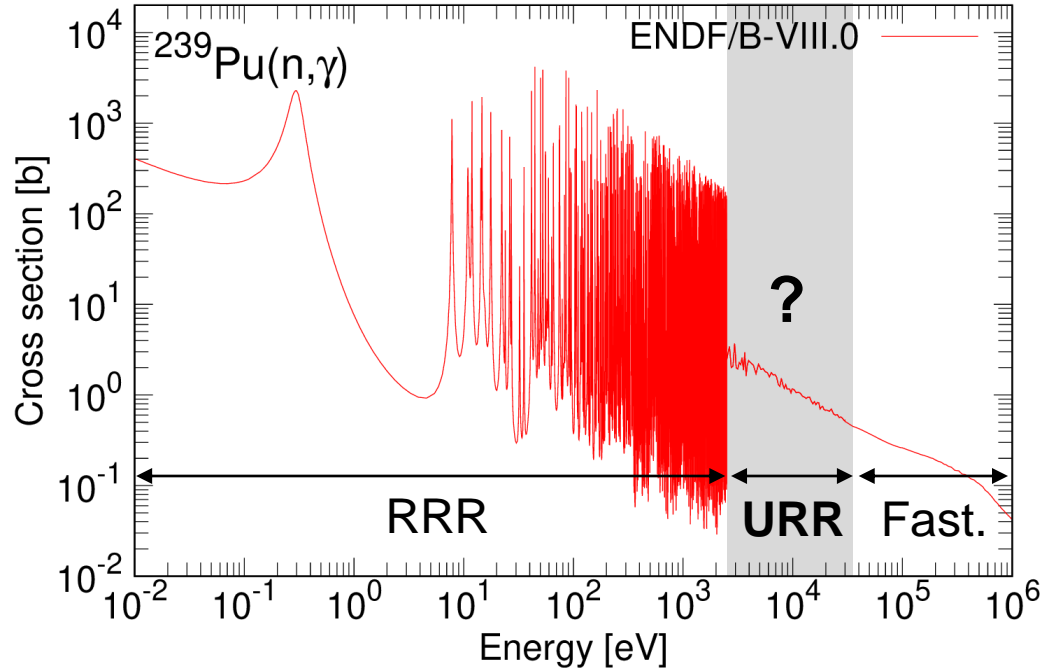
# Random-matrix approach for generating cross sections and probability table in the unresolved resonance region

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# Cross section in theoretical approaches



## Resolved resonance region (RRR)

- Individual resonances can be seen
- Expressed in R-matrix formalism by using resonance parameters

## Fast energy region

- Smooth cross section because resonance width is bigger than the spacing
- Calculated by Hauser-Feshbach theory

## Unresolved resonance region (URR)

- Resonances overlap and cannot be separated from one another
- **No commonly used theoretical model**

# Importance of URR and probability table

Self-shielding: Neutrons are easily captured at resonance peaks, and the neutron flux decreases at resonance energies. → Impacts on reaction rate

**We need to take into account resonances in URR as well !**

- A probability table which represents the probabilistic distribution of cross sections is employed in Monte-Carlo approaches
- Assume Wigner and chi-square distributions for the resonance spacing and width, respectively

**A theoretical approach is required to calculate the cross sections in URR without these kinds of assumptions**

	cumulative probability	total [b]	elastic [b]	fission [b]	capture [b]
1	0.00717	0.8067E+01	0.6145E+01	0.1220E+01	0.7018E+00
2	0.07133	0.9996E+01	0.8752E+01	0.9580E+00	0.2861E+00
3	0.34225	0.1166E+02	0.9970E+01	0.1405E+01	0.2863E+00
4	0.53333	0.1387E+02	0.1053E+02	0.2738E+01	0.6087E+00
5	0.64767	0.1664E+02	0.1076E+02	0.4663E+01	0.1218E+01
6	0.72067	0.2004E+02	0.1086E+02	0.7148E+01	0.2030E+01
7	0.77967	0.2408E+02	0.1123E+02	0.9497E+01	0.3352E+01
8	0.82000	0.2895E+02	0.1162E+02	0.1222E+02	0.5102E+01
9	0.85467	0.3481E+02	0.1218E+02	0.1630E+02	0.6329E+01
10	0.88833	0.4206E+02	0.1252E+02	0.2033E+02	0.9215E+01
11	0.91342	0.5036E+02	0.1386E+02	0.2378E+02	0.1272E+02
12	0.93617	0.6042E+02	0.1501E+02	0.2906E+02	0.1635E+02
13	0.95200	0.7228E+02	0.1652E+02	0.3567E+02	0.2008E+02
14	0.96642	0.8747E+02	0.2035E+02	0.4176E+02	0.2536E+02
15	0.97850	0.1054E+03	0.2452E+02	0.4657E+02	0.3435E+02
16	0.98667	0.1267E+03	0.2775E+02	0.6243E+02	0.3649E+02
17	0.99092	0.1526E+03	0.3403E+02	0.7385E+02	0.4471E+02
18	0.99467	0.1835E+03	0.4693E+02	0.8406E+02	0.5252E+02
19	0.99750	0.2200E+03	0.6376E+02	0.8705E+02	0.6918E+02
20	1.00000	0.2866E+03	0.1027E+03	0.7617E+02	0.1078E+03

Table 1. Probability table for <sup>239</sup>Pu at an incident neutron energy of 430 eV and a temperature of 300 degrees Kelvin [based on ENDF/B-V, Rev. 2]

[L. L. Carter et al., LA-UR-98-0026 \(1998\).](#)

# Research objective and approach

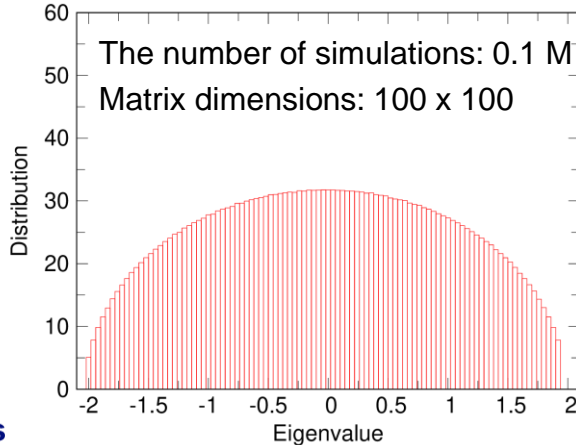
Develop a theoretical model to obtain cross sections in URR and employ them to create the probability table

## Random-matrix approach

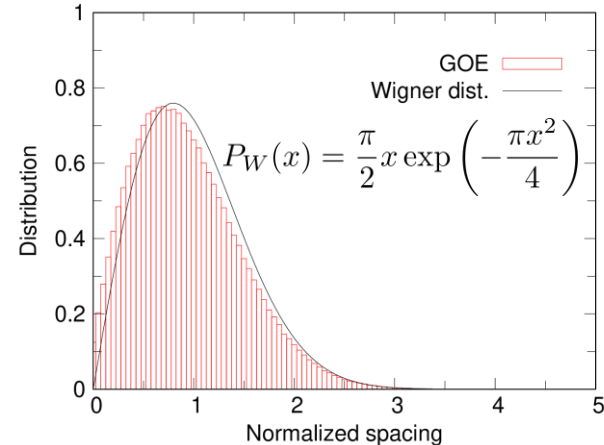
Gaussian Orthogonal Ensemble (GOE)

- Real symmetric matrix
- Matrix elements are Gaussian distributed random numbers

Eigenvalue distribution



Distribution of spacing of each eigenvalue



# S matrix calculated using GOE

$$S_{ab}^{(\text{GOE})} = \delta_{ab} - 2i\pi \sum_{\mu\nu} W_{a\mu} (D^{-1})_{\mu\nu} W_{\nu b}$$

$E$  : energy in units of  $\lambda$   
(Does not represent incident neutron energy!)

$$D_{\mu\nu} = E\delta_{\mu\nu} - H_{\mu\nu}^{(\text{GOE})} + i\pi \sum_c W_{\mu c} W_{c\nu}$$

$H_{\mu\nu}^{(\text{GOE})}$  : Random matrix

[J. J. M. Verbaarschot et al., Phys. Rep. 129, 367 \(1985\).](#)

[T. Kawano et al., Phys. Rev. C92, 044617 \(2015\).](#)

## Input parameters

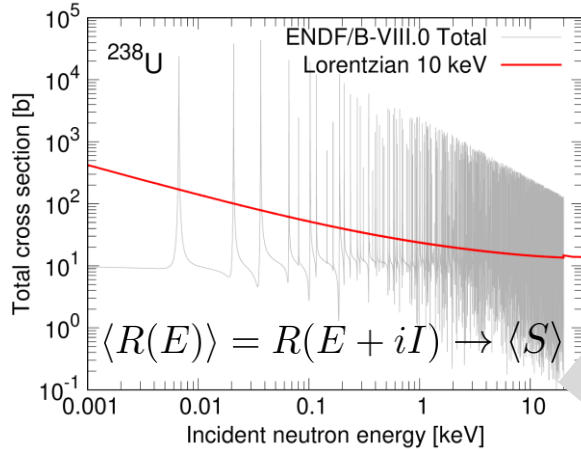
- The number of levels  $N$
- The number of channels  $\Lambda$ 
  - The number of inelastic scattering channels is determined by possible number of neutron partial waves and the corresponding excited levels
- Neutron, photon, and fission transmission coefficients  $T_n, T_\gamma, T_f$

# How to smoothly connect 2 regions

## The **S** matrix smoothly connects between RRR and fast energy regions

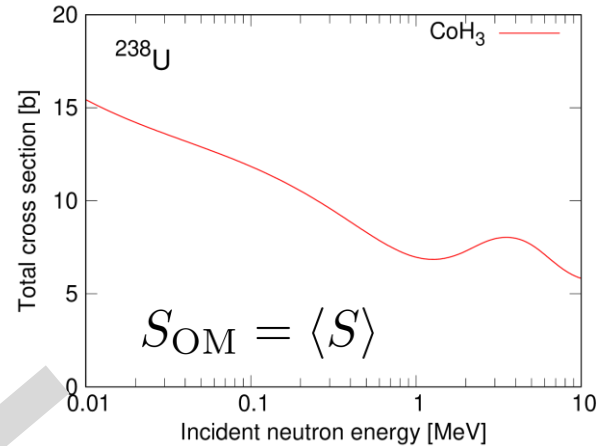
Between RRR and URR

1. Obtain the energy-averaged R matrix in RRR (ENDF/B-VIII.0)
2. Determine the optical model parameters to reproduce  $\langle R(E) \rangle$



Between URR and fast energy regions

- ✓ We can consistently use the same treatments as those used in the Hauser-Feshbach model and (coupled-channel) optical model in  $\text{CoH}_3$



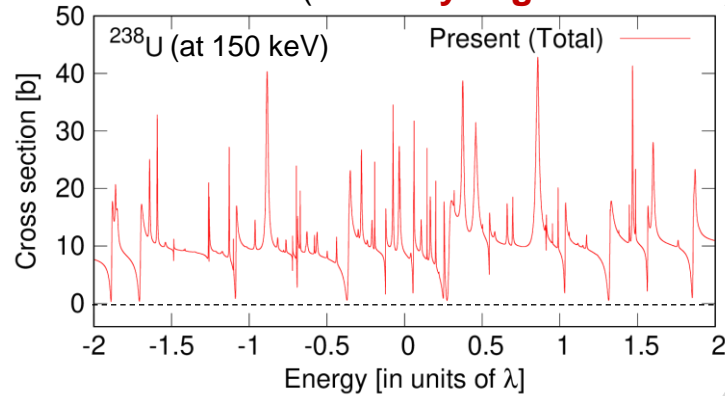
Key input parameters:

**Transmission coefficients**  $T = 1 - |\langle S \rangle|^2$

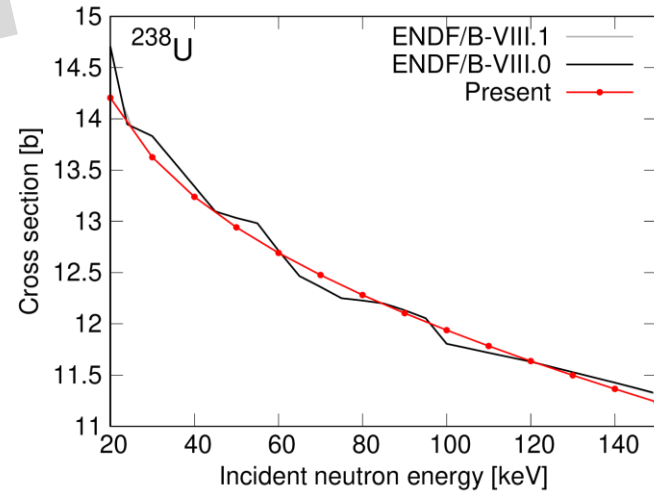
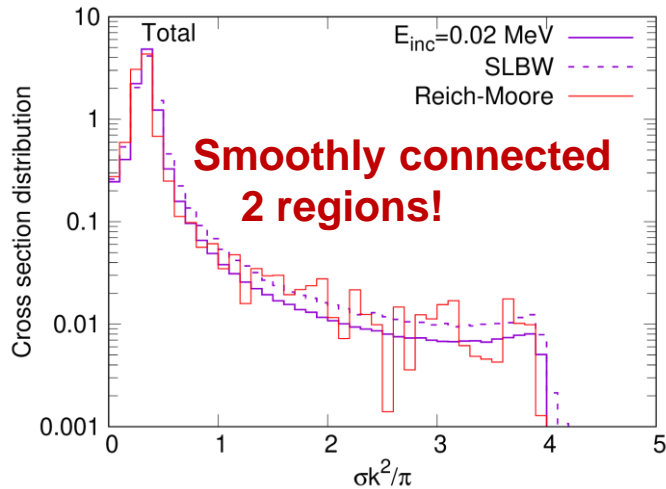
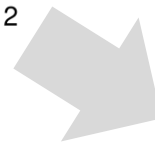
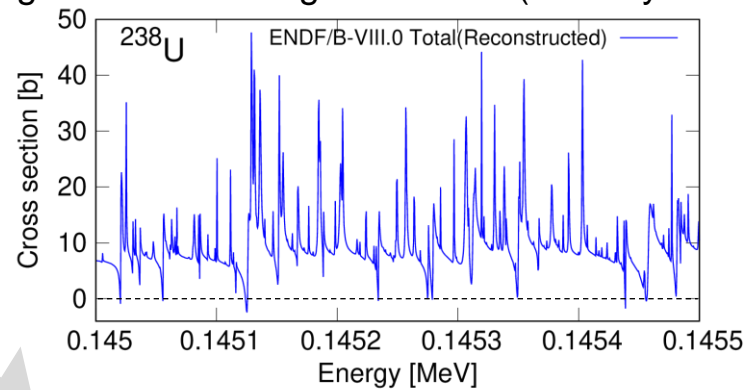
# Calculated cross sections

[Phys. Rev. C112, 064608 \(2025\).](#)

GOE-S-matrix model (**unitarity is guaranteed!**)

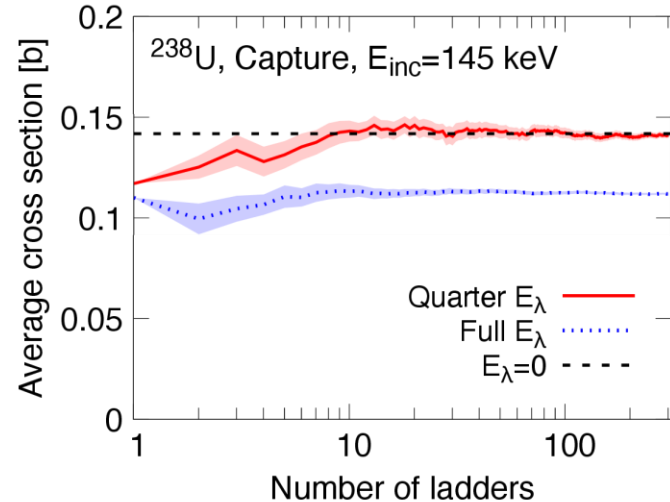
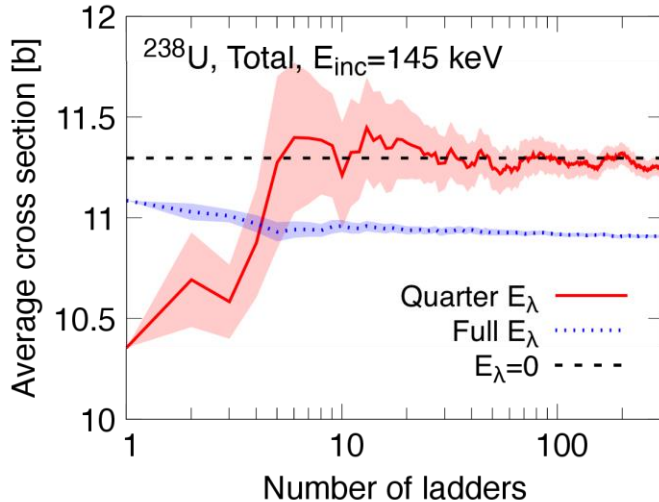
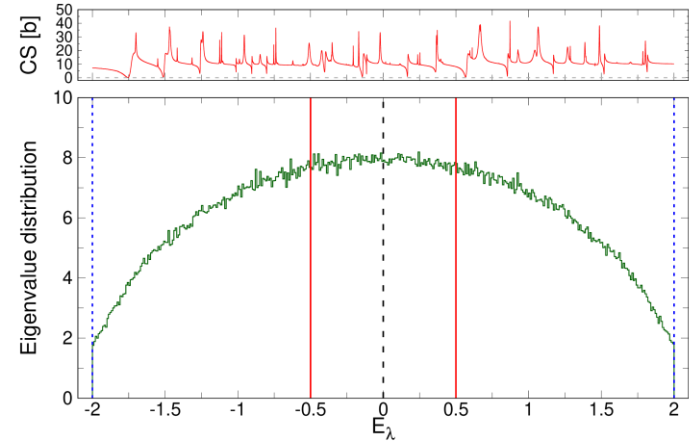


Single-level Breit-Wigner formula (unitarity is not guaranteed)



# Selected energy range

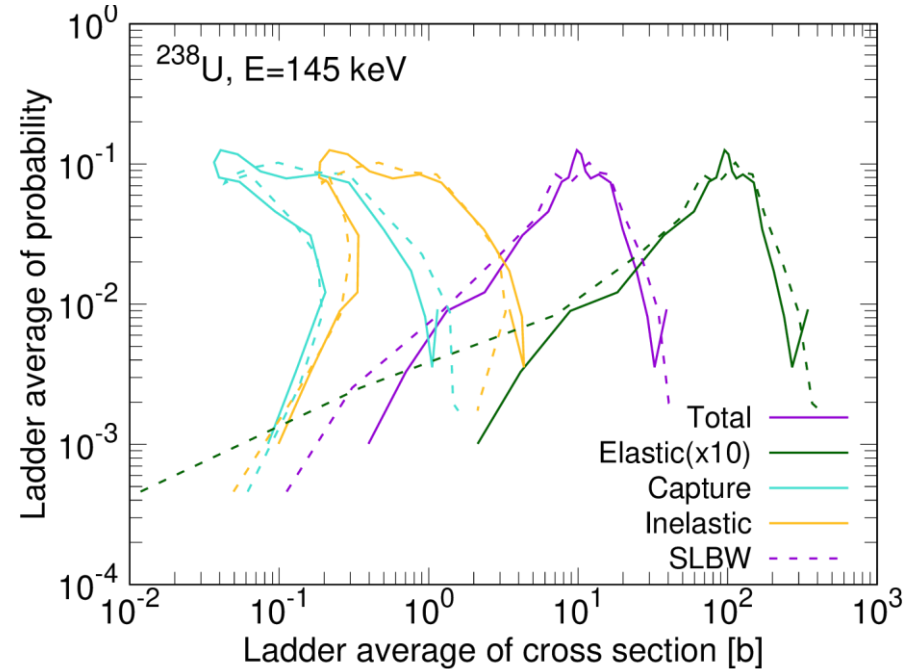
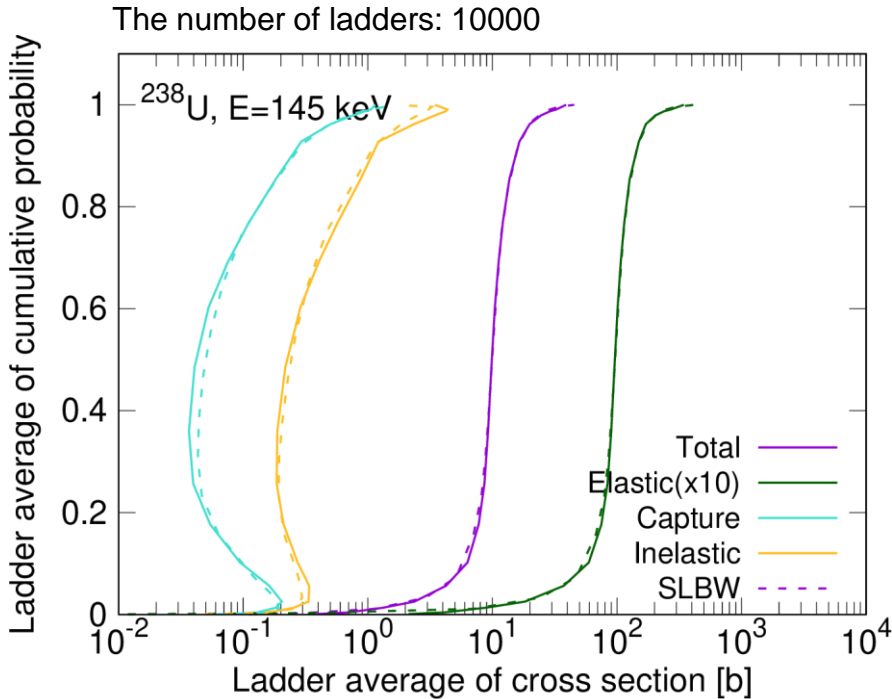
- The eigenvalue distribution corresponds to the level density.
- The convergence behavior depends on the level density.
- Compared the convergence of the cross sections using the full  $E_\lambda$  range and the quarter  $E_\lambda$  range.



**We restrict the energy range to  $E_\lambda = -N/2$  to  $N/2$ .**

# Calculated probability tables

[J. Nucl. Sci. Technol. \(published\)](#)



- No significant differences are observed between the GOE calculation and the single-level Breit-Wigner result.
- Small ladder average of cross sections are seen in the SLBW results due to the negative cross sections.

# Conclusions

- **We developed a new theoretical model to obtain cross sections in URR and employ them to create the probability table.**
- The input transmission coefficients are determined that the S matrix smoothly connects between RRR and higher energy regions.
- The average cross sections converge to evaluate data by restricting the range to  $E_\lambda = -\lambda/2$  to  $\lambda/2$ .
- The main difference in probability table between our new approach and the SLBW comes from the unitarity condition.

# Future work

- Introduce the temperature effect (Doppler broadening)

# Acknowledgements

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