

Sphaleron rate for first-order electroweak phase transitions

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JA, Kari Rummukainen, Tuomas V.I. Tenkanen [arXiv:2506.04939]



Motivation: Electroweak Baryogenesis

- ◆ EWBG is a testable scenario with present and near future experiments: EDM, particle colliders, gravitational waves
- ◆ BSM physics is needed:
 - ✚ Needs new sources of CP violation
 - ✚ Needs a strong first order phase transition
- ◆ Standard Model has a source for baryon number violation: **Sphaleron**
 - ✚ This talk: **Sphaleron washout** lattice results
 - ✚ When is the phase transition strong enough?

Sphaleron rate Γ

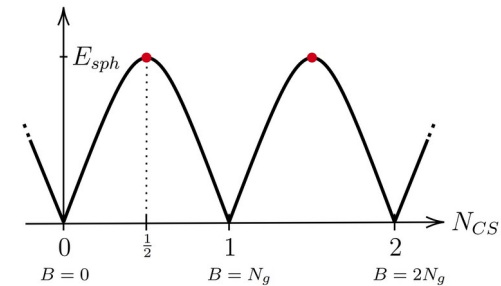
$$\Delta N_{\text{CS}} = \frac{g^2}{32\pi^2} \int_0^t dt \int d^3x \epsilon_{\alpha\beta\gamma\delta} F^{\alpha\beta} F^{\gamma\delta} = \frac{1}{3} \Delta B = \frac{1}{3} \Delta L$$

- ◆ The amount of baryon number violation can be quantified via the **sphaleron rate**: how rapidly sphaleron processes happen

$$\Gamma = \lim_{V, t \rightarrow \infty} \frac{\langle N_{\text{CS}}(t)^2 \rangle}{Vt}$$

- ◆ Low temperature Higgs phase: (sphaleron rate)

$$\Gamma_{\text{brk}} \sim \alpha_W^4 T^4 e^{-E_{\text{sph}}/T}$$



- ◆ High temperature Confinement like phase: ("hot sphaleron" rate)

$$\Gamma_{\text{sym}} \sim \alpha_W^5 \log(\alpha_W^{-1}) T^4$$

Sphaleron washout

- Given enough CP violation around the expanding bubble wall sphaleron processes generate baryon number in the high T confinement like phase.
- Baryon excess swallowed by the expanding bubble: **if Higgs phase sphalerons are suppressed fast enough** the baryon number freezes out, otherwise the asymmetry will be washed out.
- When sphaleron rate is comparable to Hubble rate: baryon number frozen

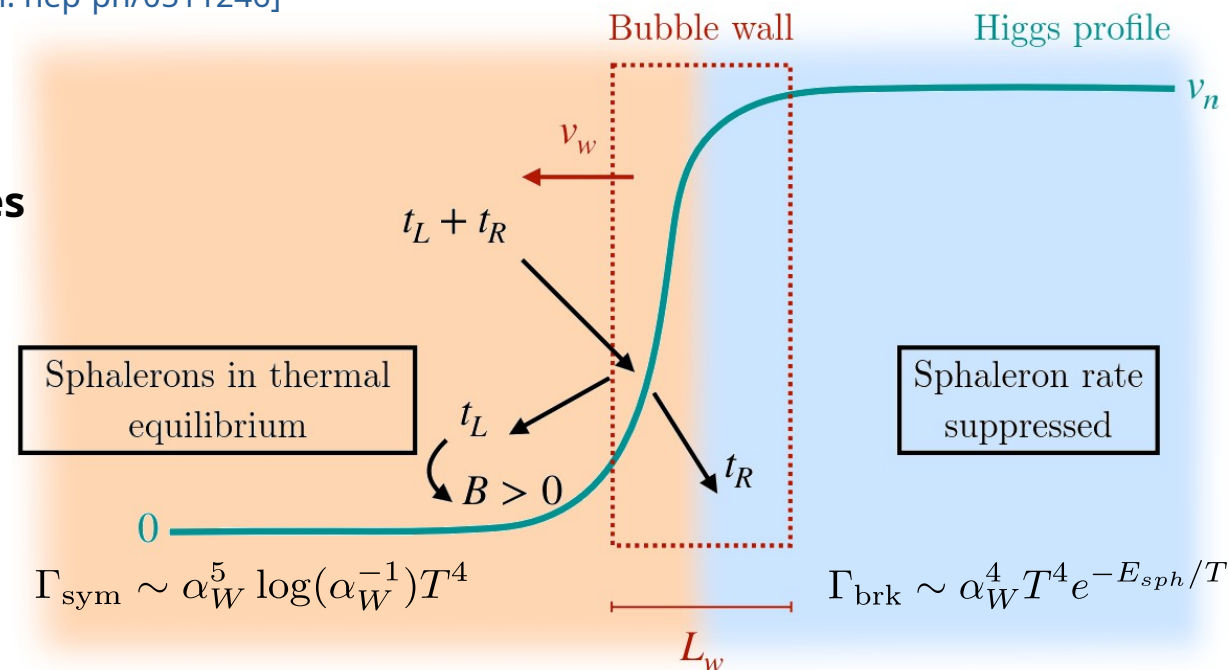
$$N_\rho \frac{\Gamma(T_f)}{T_f^3} = H(T_f) \quad [\text{Burnier et al. hep-ph/0511246}]$$

- Need to compute:

1. T_p when the transition completes

2. T_f when baryon number freezes

If $T_p \leq T_f \Rightarrow$ no washout



[Figure from: van de Vis et al. 2508.09989]

EWBG needs BSM physics

- ♦ Lots of BSM models, computing the sphaleron rate (T_f) and bubble nucleation rate (T_p) is a lot of work.
- ♦ Many BSM models can be mapped to a dimensionally reduced SU(2)+Higgs 3D EFT.
 - ✚ Get sphaleron and nucleation rates model independently in the EFT using lattice methods

BSM models

$\{\lambda_i\}$

Maps to the EFT

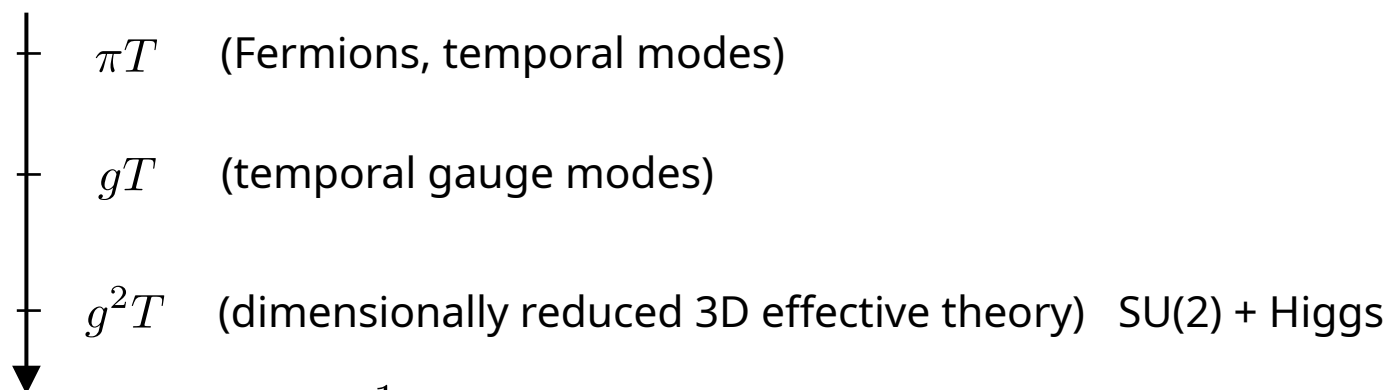
Simpler 3D **EFT** with
only 2 parameters

$\{x, y\}$

$\Gamma = \Gamma(x, y)$

SU(2)+Higgs Dimensionally reduced 3D EFT

1. Start with the Standard Model or some BSM model
2. High T hierarchy of scales: integrate out heavy modes



$$L_{3D} = \frac{1}{4} F_{ij} F_{ij} + (D_i \phi)^\dagger (D_i \phi) + m_3^2 \phi^\dagger \phi + \lambda_3 (\phi^\dagger \phi)^2,$$

3. Obtain relations of starting model parameters to EFT parameters

$g_3, \lambda_3, m_3^2 = \text{Functions of } T \text{ and the starting 4D model's parameters}$

**EFT with two
dimensionless
couplings:**



$$x(T) = \frac{\lambda_3}{g_3^2}, \quad y(T) = \frac{m_3^2}{g_3^4}$$

[Kajantie et al. hep-ph/9508379]

Largely automated by DRalgo code

[Ekstedt et al. 2205.08815]

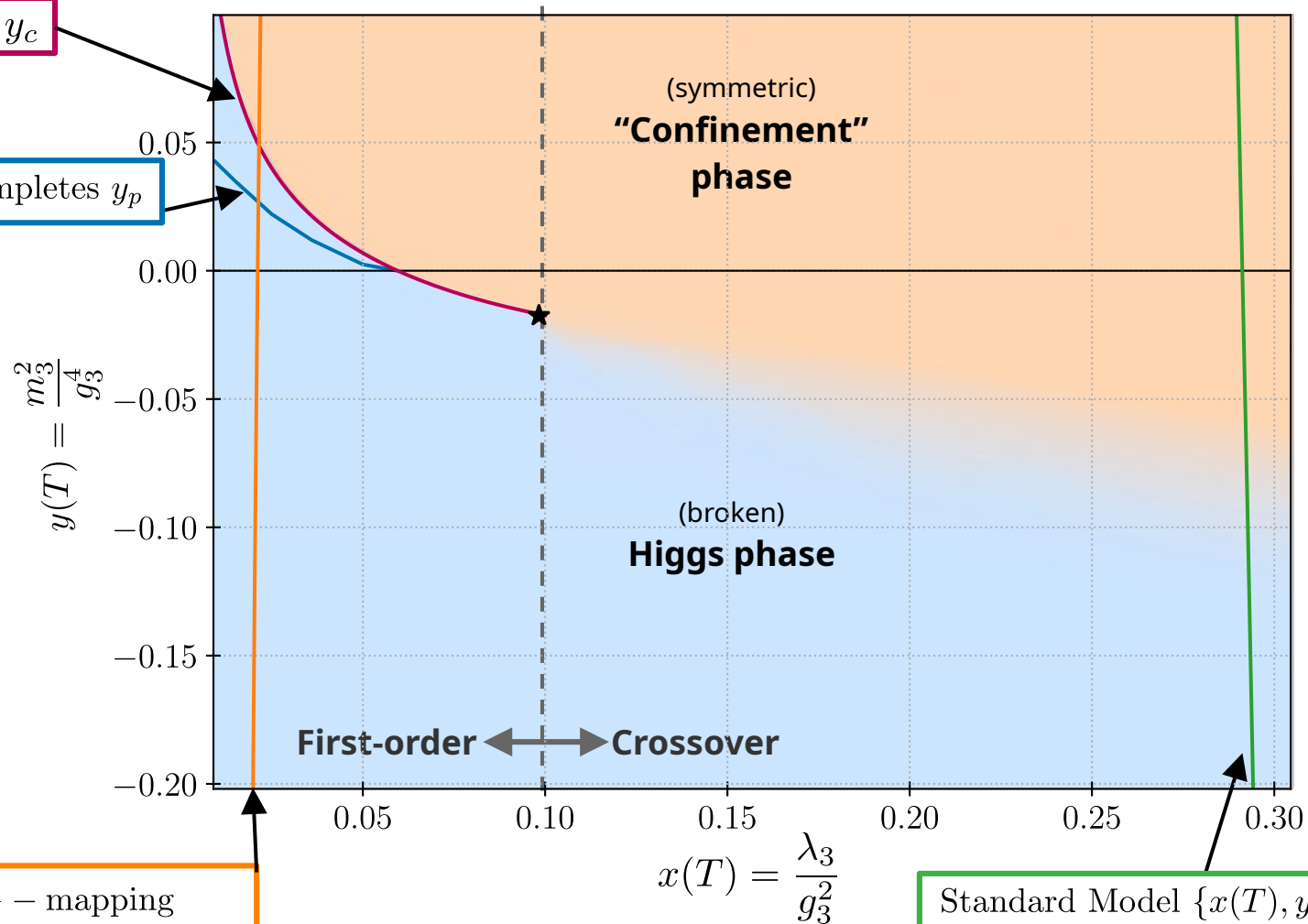
3D EFT in the xy-plane

- Study the EFT in its 2 parameter space on the lattice:

[Kajantie et al. hep-lat/9510020]
[Moore, Rummukainen hep-ph/0009132]

Critical y_c

Phase transition completes y_p



y_c and y_p curves computed on the lattice in [Gould et al., 2205.07238 refs therein]

Sphaleron rate on the lattice

- ♦ Computation of the (hot)sphaleron rate in the 3D EFT well established. confinement like phase and Higgs phase computations:
 - first-order transition region [[Moore et al. Hep-ph/0009132](#), [hep-ph/9805264](#), ...]
 - crossover region, SM parameters [[Moore hep-ph/9805264](#), [D'Onofrio et al. 1207.0685](#), [1404.3565](#), [JA et al. 2301.08626](#), ...]

N_{CS} defined on the lattice with gradient flow cooling

$$\tilde{\Gamma}(x, y) = \underbrace{\frac{P(|N_{CS} - \frac{1}{2}| < \frac{\epsilon}{2})}{\epsilon V}}_{\text{statistical}} \underbrace{\left\langle \left| \frac{\Delta N_{CS}}{\Delta t} \right| \right\rangle d}_{\text{dynamical}},$$

Dynamical evolution described by an effective Langevin evolution [[Bödeker, arXiv:hep-ph/9801430](#)]

$$\Gamma(x, y) = \frac{(g_3^2)^5}{\sigma_{\text{el}}} \tilde{\Gamma}(x, y),$$

Sphaleron freeze-out curve in xy-plane $y_f(x)$

- ♦ Comparing the sphaleron rate to Hubble rate gives the sphaleron freeze out curve $y_f(x)$ in the xy-plane:

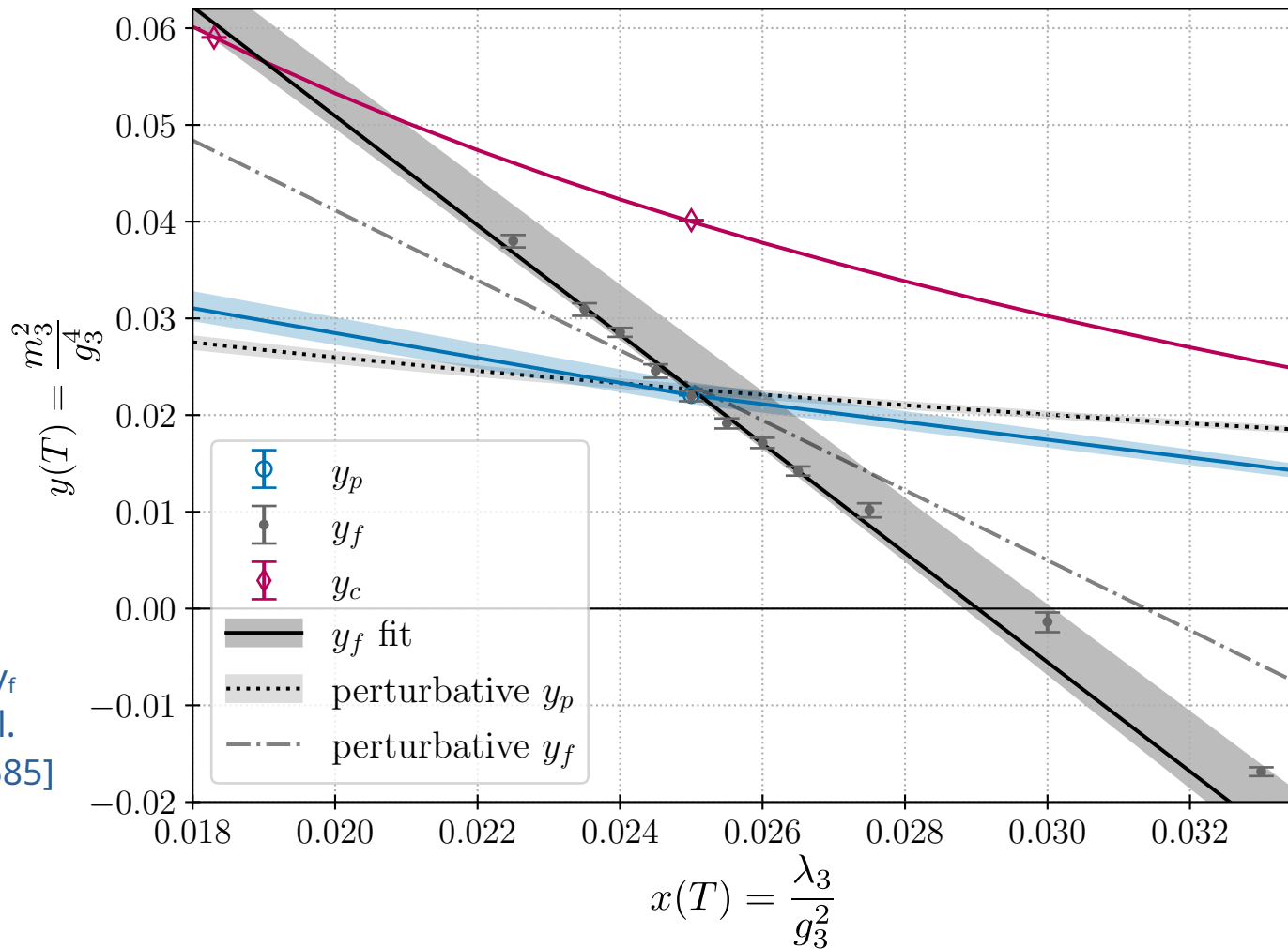
$$N_\rho \frac{\Gamma(T_f)}{T_f^3} = H(T_f)$$

$$\rightarrow \ln \Gamma(x, y_f) = \ln \tilde{\Gamma}(x, y_f) + \underbrace{\ln \left(\frac{(g_3^2)^5}{\sigma_{\text{el}} T_f^4} \right)}_{\text{Approx. constants}} = \underbrace{\ln \frac{H(T_f)}{N_\rho T_f}}_{\text{Approx. constants}} \approx -38$$

$$\rightarrow y_f(x) \quad (\text{Sphaleron freeze out curve})$$

$$\rightarrow \text{Baryon number preserved when:} \\ y_p(x_c) \leq y_f(x_c) \quad (\rightarrow T_p \leq T_f)$$

Strong Baryon number preservation condition

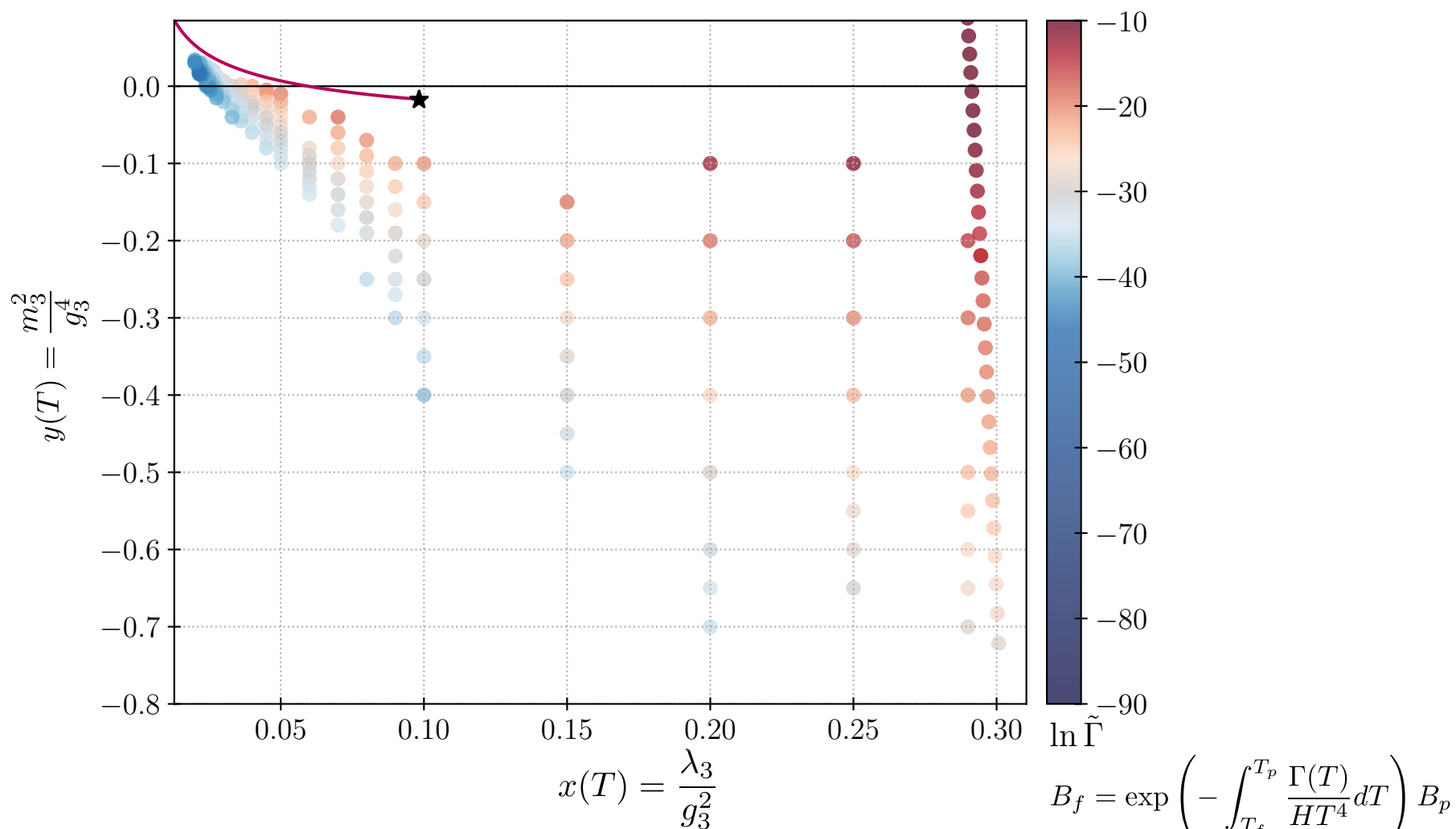


- ♦ Sphaleron rate is smaller than the Hubble rate when the transition completes if:

$$x_c \lesssim \bar{x} \approx 0.025, \quad v/T_c \equiv \sqrt{2\Delta\langle\phi^\dagger\phi\rangle}/T_c \gtrsim 1.33$$

none of the generated baryon number is washed out.

Less strict bounds can be made from the data



- ♦ If you allow for partial washout you can have slightly less strong transitions...
washout factor of bigger than $\exp(-100)$:

$$x_c \lesssim \hat{x} \approx 0.036,$$

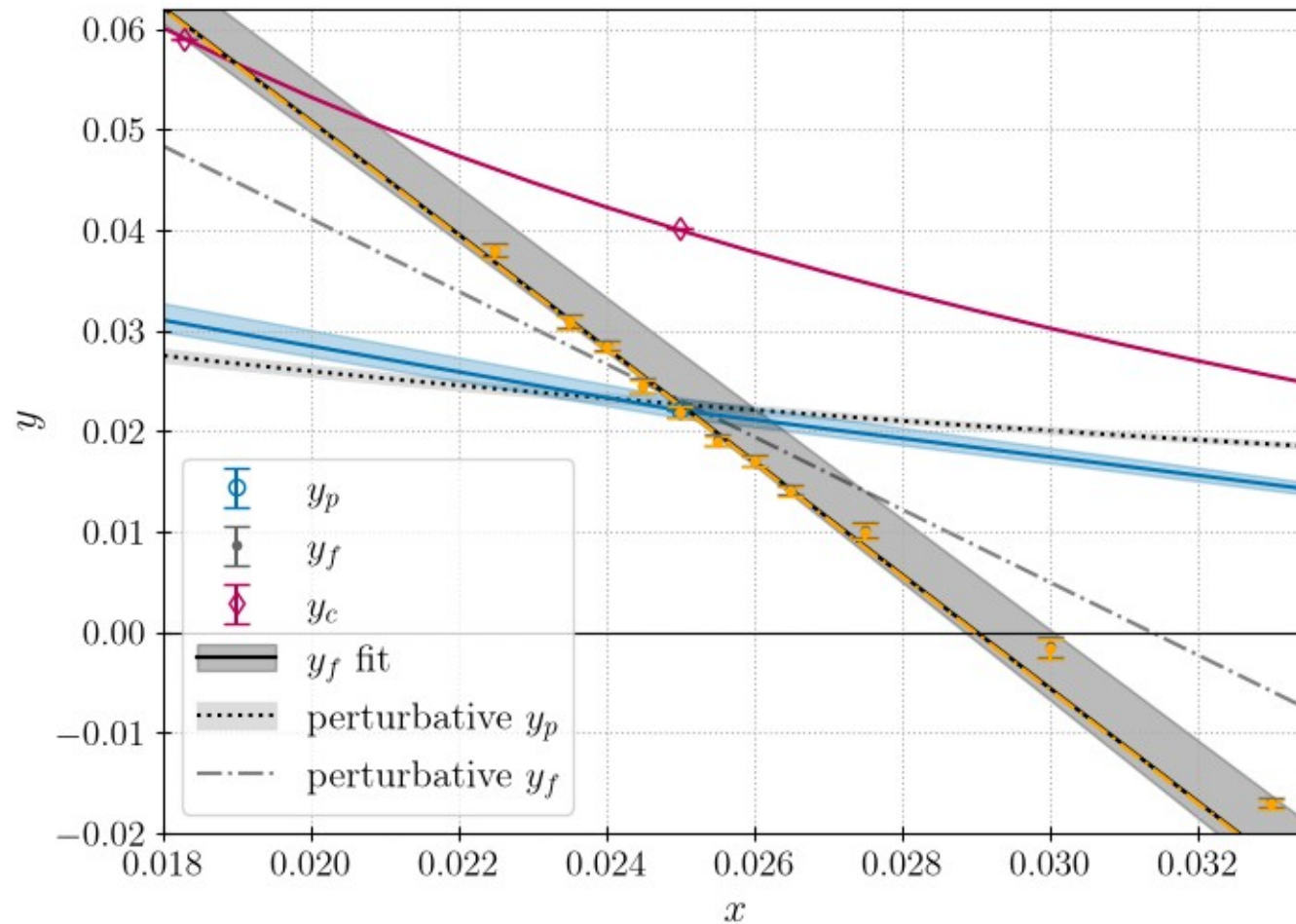
[more details: Wu et al. arxiv:2506.01585]

Summary

- ◆ BSM models that map to the SU(2)+Higgs EFT can be efficiently studied on the lattice without resulting into more complex simulations with the extra field content.
- ◆ We obtained a sphaleron freeze-out curve in the xy-plane
- ◆ Using previous determination of the critical and percolation curves we obtained a strong condition for baryon number preservation, which is slightly stricter than previous estimates $x_c \lesssim \bar{x} \approx 0.025$.

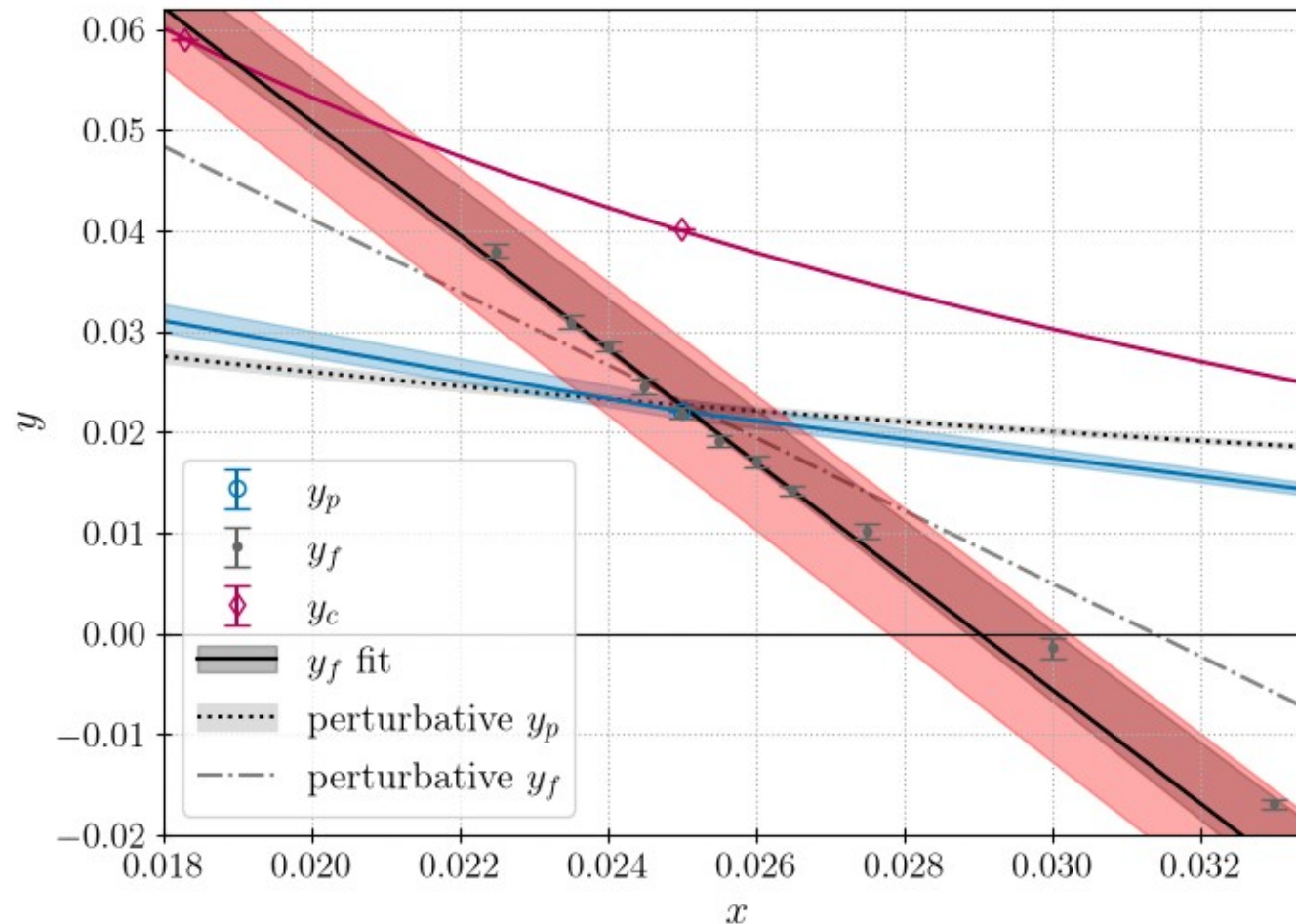
Thank you for your attention!

Uncertainties



- ◆ Assume g_3 , σ_{el} , $H(T_f)$ do not change much by BSM physics.
(Orange: decrease σ_{el} by 6%) [Bödecker, Klose arxiv:2510.20594]

Uncertainties



- ♦ Using the Langevin dynamics deep in the Higgs phase is not well justified. However, the exponential suppression from the statistical part remains the dominant contribution. (Red band: 1000% error to dynamical part.)