

Chiral imbalances in neutron star mergers

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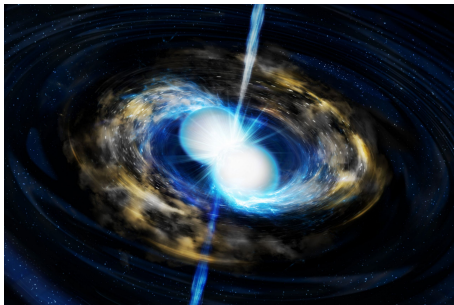
SPH & S. Sen arXiv:2602.21294

Strong & Electroweak Matter
August 21, 2026



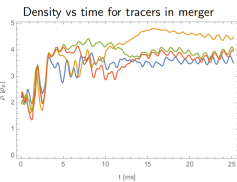
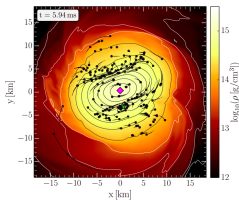
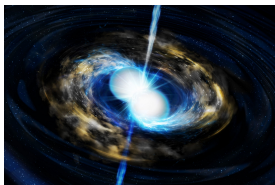
Outline

1. Shock/compression pushes matter out of beta equilibrium
2. Beta equilibration ensues, developing a chiral imbalance among the electrons
3. If the imbalance is large enough, and lasts long enough, it could give rise to novel effects (chiral plasma instability, heating, ...)

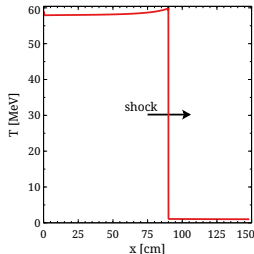
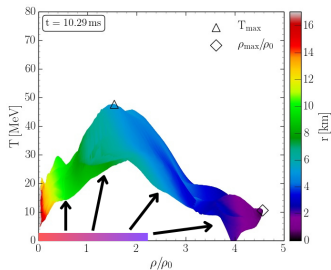


Compressions/shocks in NS mergers

Nuclear matter in neutron star mergers experiences both secular compression and superimposed density oscillations.



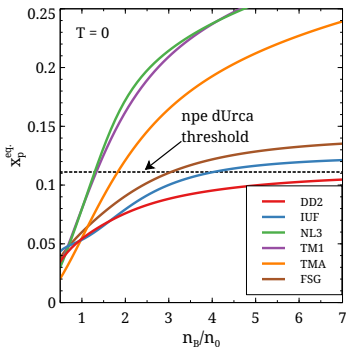
Simulation results from M. Hanauske and the Rezzolla group (Frankfurt)



Microscopic considerations: x_p and x_5

Consider dense matter as charge-neutral Fermi liquid of n , p , e^- .

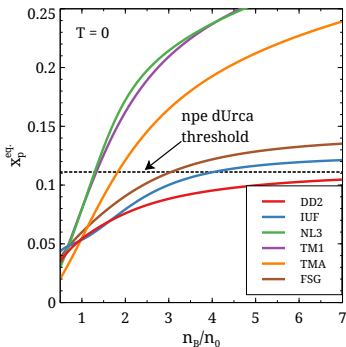
- At a given $\{n_B, T\}$, the proton fraction $x_p = n_p/n_B$ is set by the history of the system. In beta-equilibrium, $x_p = x_p^{\beta\text{-eq.}}(n_B, T)$.



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Chirality in dense matter

- ▶ Left- & right-handed fermions are present. Thermodynamically, they are identical ($\varepsilon, P, \dots$).
- ▶ $x_5 \equiv x_R - x_L$. In equilibrium, $x_5 = 0$.
- ▶ If this is so, can a **chiral imbalance** ever **develop**?
 - ▶ Yes, through the **weak interaction**.
 - ▶ Imbalance suppressed $\sim m_{\text{fermion}}$.
Focus on the lightest fermion (e^-).

Creating a chiral imbalance

Urca process^a

- ▶ $n \rightarrow p + e_L^- + \bar{\nu}$ and $e_L^- + p \rightarrow n + \nu$
- ▶ Weak interaction - slow-ish
- ▶ Only involves left-handed e^-

Rutherford scattering

- ▶ $e_L^- + p \leftrightarrow e_R^- + p$
- ▶ Electromagnetic interaction - fast!
- ▶ Chirality flipping - goes as m_e^2

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Define net rates:

$$\Gamma_{\text{Urca}}^{\text{net}} \equiv \Gamma_{n \rightarrow p + e^- + \bar{\nu}} - \Gamma_{e^- + p \rightarrow n + \nu}$$

$$\Gamma_{L \leftrightarrow R}^{\text{net}} \equiv \Gamma_{e_R + p \rightarrow e_L + p} - \Gamma_{e_L + p \rightarrow e_R + p}$$

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$$\frac{dn_{e,L}}{dt} = \Gamma_{L \leftrightarrow R}^{\text{net}} + \Gamma_{\text{Urca}}^{\text{net}}$$

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$$\begin{aligned} \frac{dn_5}{dt} &= \frac{dn_{e,R}}{dt} - \frac{dn_{e,L}}{dt} \\ &= -2\Gamma_{L \leftrightarrow R}^{\text{net}} - \Gamma_{\text{Urca}}^{\text{net}} \end{aligned}$$

Creating a chiral imbalance (2)

For small deviations from chemical equilibrium:

$$\Gamma_{\text{Urca}}^{\text{net}} \propto \delta\mu \equiv \mu_n - \mu_p - \mu_{e,L} \quad \text{and} \quad \Gamma_{L \leftrightarrow R}^{\text{net}} \propto \mu_5 \equiv \mu_{e,R} - \mu_{e,L}.$$

Evolution equations

$$\frac{dn_p}{dt} = \underbrace{\Gamma_{\text{Urca}}^{\text{net}}}_{\sim \delta\mu}$$

$$\frac{dn_5}{dt} = -2 \underbrace{\Gamma_{L \leftrightarrow R}^{\text{net}}}_{\sim \mu_5} - \underbrace{\Gamma_{\text{Urca}}^{\text{net}}}_{\sim \delta\mu}$$

Possibility of a plateau^a in n_5 .

^aFirst raised by Grabowska, Kaplan, & Reddy arXiv:1409.3602. We improve upon their Urca treatment.

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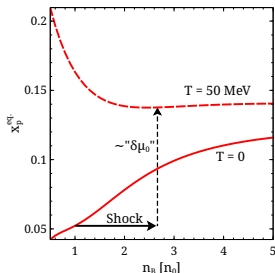
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Initial conditions

► Matter is shocked right before $t = 0$. Yields...

► $\delta\mu(t = 0) = \delta\mu_0$

► $\mu_5(t = 0) = 0$



Chiral imbalance plateau

$$\delta\mu(t=0) = \delta\mu_0.$$

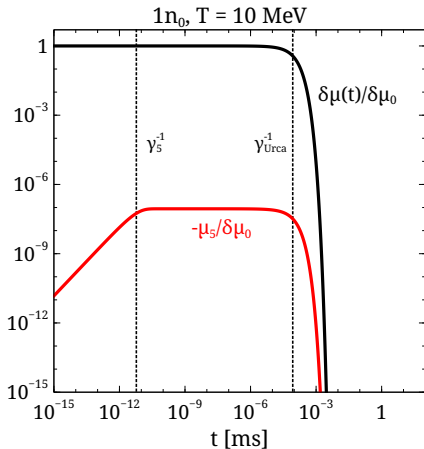
$$\mu_5(t=0) = 0.$$

$$\frac{dn_p}{dt} = \underbrace{\Gamma_{\text{Urca}}^{\text{net}}}_{\sim \delta\mu}$$

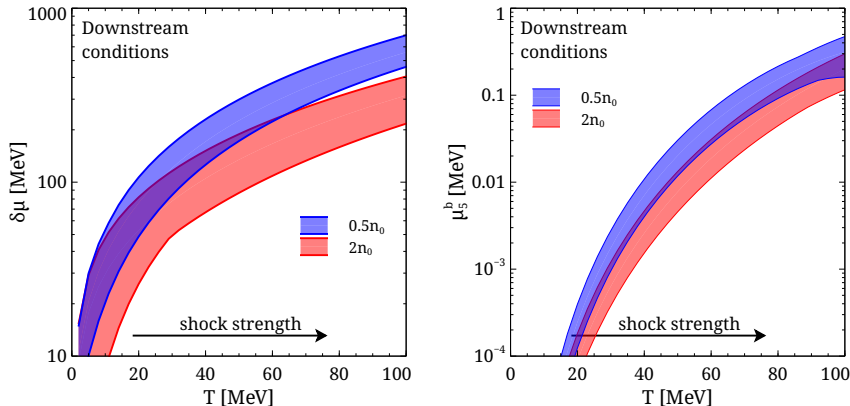
$$\frac{dn_5}{dt} = -2 \underbrace{\Gamma_{L \leftrightarrow R}^{\text{net}}}_{\sim \mu_5} - \underbrace{\Gamma_{\text{Urca}}^{\text{net}}}_{\sim \delta\mu}$$

Assuming $\gamma_5 \gg \gamma_{\text{Urca}}$, **plateau**:

► Duration: $\gamma_{\text{Urca}}^{-1}$ Height: $-\delta\mu_0 \frac{\gamma_{\text{Urca}}}{\gamma_5}$



Chiral imbalance plateau (2)



Chiral imbalances of $\mu_5 \gtrsim 10 - 100$ keV could potentially be generated by the strongest of shocks, lasting briefly.

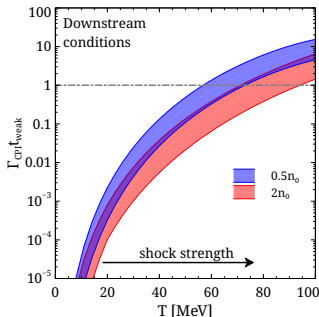
- Look for interesting effects that happen in regions of finite μ_5 .

Chiral plasma instability

$$\nabla \times \mathbf{B} - \frac{\partial \mathbf{E}}{\partial t} = \sigma \mathbf{E} + \underbrace{\frac{\alpha}{\pi} \mu_5 \mathbf{B}}_{\text{CME current}}$$

Supports exponentially growing mode if $\mu_5 \neq 0$.

Can mode grow sufficiently during the time that the μ_5 plateau lasts?



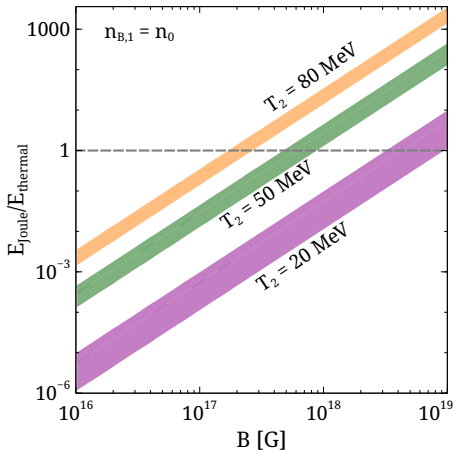
Yes, for very strong shocks. Potentially leads to very strong B-fields!

Joule Heating

The CME current presents an additional Joule heating source.

$$\frac{dW}{dt} \sim \mathbf{E} \cdot \mathbf{J}_{CME} \sim \frac{\alpha^2 \mu_5^2}{\sigma} \mathbf{B}^2$$

Joule heating could exceed shock heating if $B \gtrsim 10^{17}$ G and if the shock is strong enough to generate a sizable μ_5 .



Conclusions

- ▶ A chiral imbalance is generated wherever matter is out of beta equilibrium
 - ▶ Shocks in NS mergers push matter out of beta equilibrium.
 - ▶ Urca & chirality flipping balance, leading to plateau in μ_5 .
Plateau grows with
 - ▶ Shock strength
 - ▶ Ratio of Urca rate / chirality-flip rate
- ▶ Strong shocks could produce large enough μ_5 to cause chiral plasma instability and/or sizable Joule heating.
- ▶ Future:
 - ▶ Consider neutrino-trapping
 - ▶ Put reactions into neutron star merger simulation to see how μ_5 develops.
 - ▶ Study the electromagnetic consequences of μ_5