

TBA: Throwing Bubbles Around,
and finding out how they produce GWs

Strong and Electro-Weak Matter 2026

Helsinki, Finland (Aug. 20, 2026)



Alberto Roper Pol
University of Geneva



SNSF Ambizione grant (2023–2027): “*Exploring the early universe with gravitational waves and primordial magnetic fields.*”

Gravitational waves from cosmological phase transitions: *sound waves and MHD turbulence*

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Cosmological Gravitational Waves

- Gravity is the *weakest fundamental force*. Hence, GWs are difficult to detect but they propagate freely carrying *clean information of the source*.
- Cosmological GWs from the early Universe have the potential to provide us with *direct information on early universe physics* that is *not accessible via electromagnetic observations, complementary to collider experiments*.

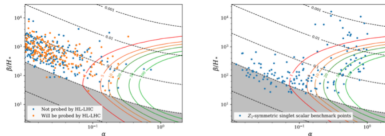
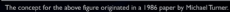


Figure 4: LEFT: Predicted values of α and β/H for $m_2 = 170$ GeV and 240 GeV (combined in one plot) in the general singlet model obtained by linearly varying the free parameters of the model and imposing requirements as described in the text. The mixing angles considered were $\sin \theta = 0.01$ (blue points) and 0.1 (orange points). The models in blue (orange) are unlikely (likely) to be probed by the high-luminosity LHC. The expected LISA sensitivities correspond to $T_* = 50$ GeV. RIGHT: Sensitivity curve for the Z_2 symmetric singlet extension. In both the left and right panels we have taken $v_w = 1$.

From [LISA CosWG], Caprini et al., *Detecting gravitational waves from cosmological phase transitions with LISA: an update (2019)*

pre-BBN magnetic fields and gravitational wave background (GWB)



Supported by DOE

Probing the early Universe with GWs

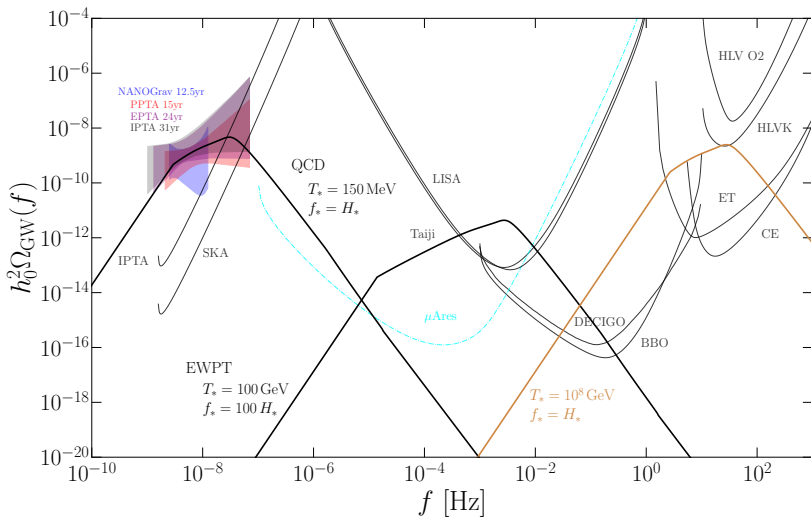
Cosmological (pre-recombination) GW background

- Why background? Individual sources are not resolvable, superposition of single events occurring in the whole Universe.

$$f_* \simeq 1.64 \times 10^{-3} \frac{0.01}{R_* \mathcal{H}_*} \frac{T_*}{100 \text{ GeV}} \text{ Hz}$$

- Probing different energy scales
 - Ground-based detectors (LVK, ET, CE) sensitive to 10–1000 Hz
Peccei-Quinn, B-L, left-right symmetries $\sim 10^7, 10^8$ GeV.
 - Space-based detectors (**LISA**) sensitive to 10^{-5} – 10^{-2} Hz
Electroweak phase transition ~ 100 GeV
 - Pulsar Timing Array (**PTA**) sensitive to 10^{-9} – 10^{-7} Hz
Quark confinement (QCD) phase transition ~ 100 MeV

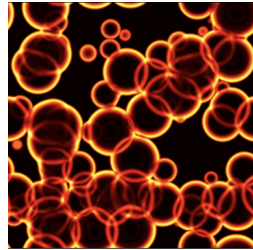
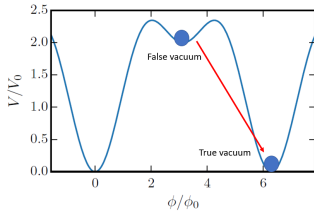
Gravitational spectrum from decaying turbulence¹



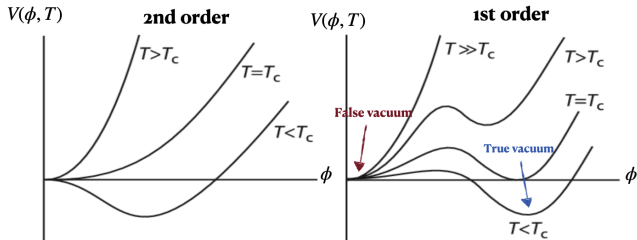
¹ ARP, C. Caprini, A. Neronov, D. Semikoz, *PRD* **105**, 123502 (2022), arXiv:2201.05630.
 A. Neronov, ARP, C. Caprini, D. Semikoz, *PRD* **103**, L041302 (2021), arXiv:2009.14174.
 ARP et al., *A&A* **708** (2026) A337, arXiv:2307.10744.

First-order phase transition

$$V(\phi, T) = \frac{1}{2}M^2(T)\phi^2 - \frac{1}{3}\delta(T)\phi^3 + \frac{1}{4}\lambda\phi^4$$

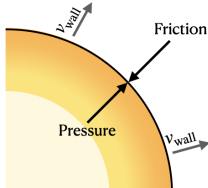


Credits: I. Stomberg



Hydrodynamics of first-order phase transitions²

- Broken-phase bubbles are nucleated and expand
- Friction from particles yield a terminal velocity ξ_w of the bubbles
- The bubble can run away when the friction is not enough to stop the bubble's acceleration



Boltzmann equation:

$$\left(p^\mu \nabla_\mu + \frac{1}{2} \partial_\mu m_X^2 \partial_{p_\mu} \right) f_X \delta^2(p^2 - m^2) = -C[f_X]$$

$$\Rightarrow \nabla_\mu T_{\text{fluid}}^{\mu\nu} = -\partial^\nu \phi \sum_X \frac{dm_X}{d\phi} \int \frac{d^3\vec{p}}{(2\pi)^3} \frac{1}{2E_X} f_X(\vec{x}, \vec{p})$$

$$\begin{array}{l} f_{\text{eq}} \rightarrow -\frac{\partial V_T}{\partial \phi} \partial^\nu \phi \\ \delta f \rightarrow -\eta u^\mu \partial_\mu \phi \partial^\nu \phi \end{array}$$

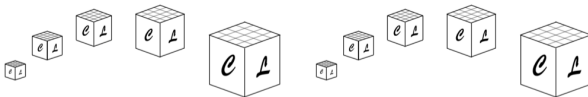
$$\nabla_\mu T_\phi^{\mu\nu} = \nabla_\mu (T_\phi^{\mu\nu} + T_{\text{fluid}}^{\mu\nu}) = 0 \quad \Rightarrow$$

$$\nabla_\mu T_\phi^{\mu\nu} = \frac{\partial V_T}{\partial \phi} \partial^\nu \phi + \eta u^\mu \partial_\mu \phi \partial^\nu \phi$$

$$\nabla_\mu T_{\text{fluid}}^{\mu\nu} = -\frac{\partial V_T}{\partial \phi} \partial^\nu \phi - \eta u^\mu \partial_\mu \phi \partial^\nu \phi$$

²Espinosa, Konstandin, No, Servant, *JCAP* **06** (2010) 028.

Scalar-Gauge-Fluid dynamics in cosmology



The art of simulating the early Universe

*(A dissertation series about lattice techniques for simulating
the dynamics of interacting fields in an expanding Universe)*

arXiv:2607.27141

Part III. Scalar-Gauge-Fluid Dynamics



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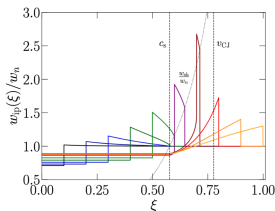
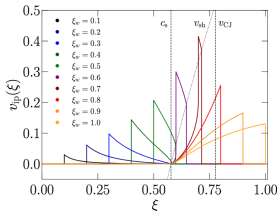
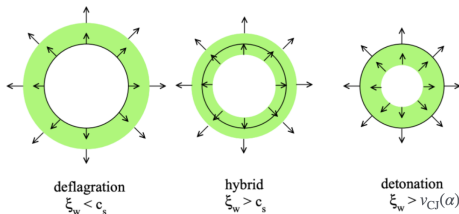
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Hydrodynamics of first-order phase transitions³

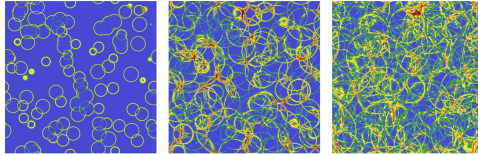
Depending on the wall velocity ξ_w and the strength of the phase transition $\alpha = \epsilon/\rho_{\text{rad}}$, we can have 3 types of solution:

subsonic deflagrations, supersonic deflagrations (hybrids) or detonations.

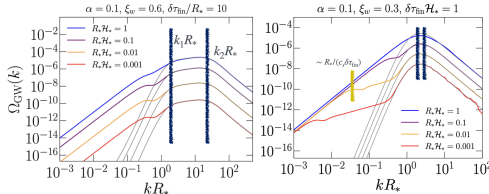


GWs from sound waves

- Numerical simulations of the scalar + fluid system performed by the Sussex/Helsinki group via an effective friction term indicate sound-wave regime to dominate for weak/intermediate phase transitions.⁴



- Sound-shell model.⁵ Three scales are found to determine the GW spectrum: $k_1 = R_*$, $k_2 = \Delta R_*$ (sound-shell thickness), and $\delta\tau_{\text{fin}}$ (sound-wave regime duration).



$$\Omega_{\text{GW}}(k) = 3 \tilde{\Omega}_{\text{GW}} K^2 \Upsilon(\tau_{\text{fin}}) (H_* R_*) S(k R_*)$$

$$\Upsilon(\tau_{\text{fin}}) = \frac{H_* \tau_{\text{fin}}}{1 + H_* \tau_{\text{fin}}}$$

$$S(k R_*) = \mathcal{N} \left(\frac{k}{k_1} \right)^{\alpha_1} \left[1 + \left(\frac{k}{k_1} \right)^{\alpha_1} \right]^{-\frac{-\alpha_1 + \alpha_2}{\alpha_1}} \left[1 + \left(\frac{k}{k_2} \right)^{\alpha_2} \right]^{-\frac{-\alpha_2 + \alpha_3}{\alpha_2}}$$

⁴ Hindmarsh *et al.*, 2013, 2015, 2017, Cutting *et al.*, 2019, Correia *et al.*, 2025.

⁵ Hindmarsh, 2016, Hindmarsh, Hijazi, 2019, ARP, Procacci, Caprini, 2023.

Scales and intermediate slope in the velocity spectrum⁶

- Identified two peaks directly from the velocity spectrum: $\tilde{k}_1$ and $\tilde{k}_2$, determined by positions of discontinuities in the 1D profiles, and by nucleation history.

$$\xi_1 \equiv \frac{2\pi}{z_1} \simeq \begin{cases} \frac{4}{3}(\xi_f + \xi_w) & \text{for deflagrations,} \\ \frac{4}{3}(\xi_w + c_s) & \text{for detonations.} \end{cases} \quad \xi_2 \equiv \frac{2\pi}{z_2} \simeq \begin{cases} 2(\xi_f - \xi_w) & \text{for deflagrations,} \\ \frac{4}{3}(\xi_w - c_s) & \text{for detonations.} \end{cases}$$

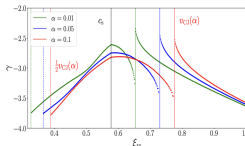
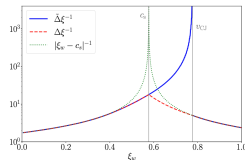
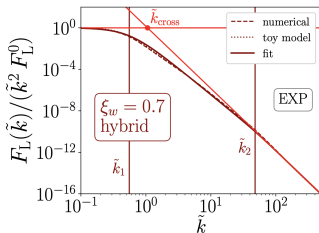
$$\frac{\tilde{k}_{1,2}^{\text{exp}}}{z_{1,2}} \simeq \{0.17, 0.21\}, \quad \frac{\tilde{k}_{1,2}^{\text{sim}}}{z_{1,2}} \simeq \{0.38, 0.25\}.$$



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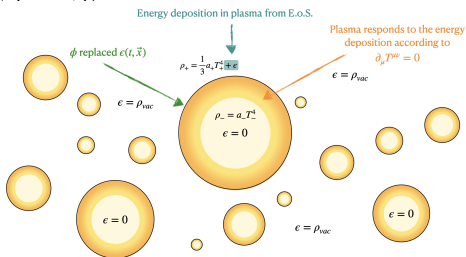
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GWs from sound waves: Higgsless simulations⁸

- Difficulty on simulations is due to the different scales of the scalar field ϕ and the fluid shell, so one can consider a nucleation history and set the pressure and energy density by knowing the value of ϵ and setting it during the simulation.
- Effect of bubble collisions on GWs is subdominant when sound waves are produced, so one can ignore the scalar field.
- Prescribes a terminal wall velocity, so slow down of the bubbles found, e.g., for strong deflagrations cannot be modeled.⁷
- Nucleation history is produced from an exponential probability distribution

$$P(t) \propto \exp[\beta(t - t_*)].$$

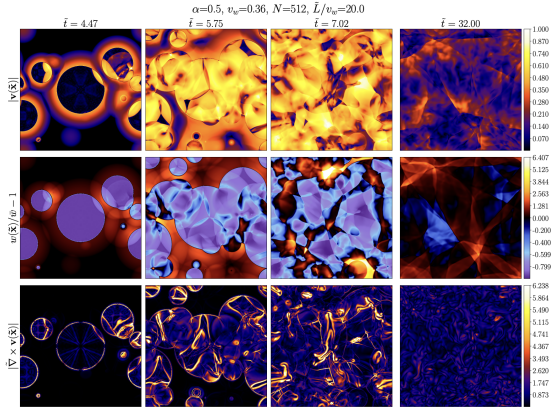


Credit: I. Stomberg

⁷ Cutting et al., 2019; Correia et al., 2025.

⁸ Jinno et al. JCAP 02 (2023) 011, 2209.04369,

Higgsless simulations of strong PTs⁹



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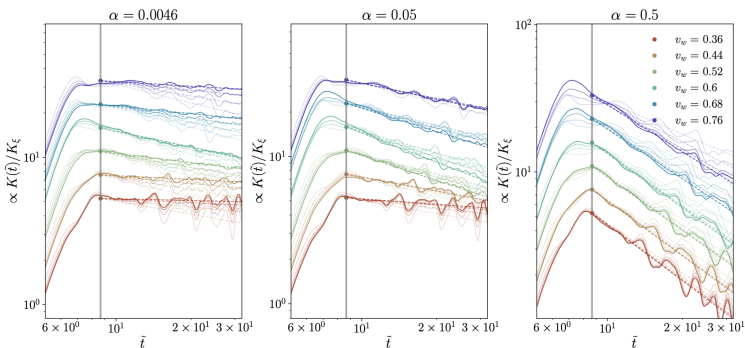
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Higgsless simulations (results)¹⁰

- Kinetic energy decay is observed in the simulations.
- For weak and strong PTs, increasing discretization enhances the decay.
- Potential indication of the development of non-linearities (turbulence).



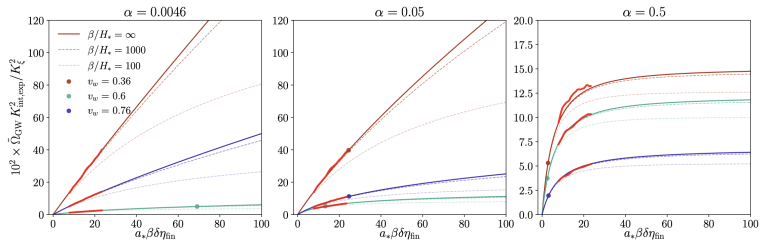
Higgsless simulations (results)¹¹

- In the literature, the GW spectrum from sound waves is usually assumed to be

$$\Omega_{\text{GW}}(f) = 3 \tilde{\Omega}_{\text{GW}} K^2 \mathfrak{T}(\tau_{\text{sw}}) (H_* R_*) S(f R_*)$$

- The linear growth, which only appears when expansion is neglected, is modified when the decay of the source is significant (e.g., due to the development of non-linearities).
- Extended model to proposed locally stationary UETC

$$\Omega_{\text{GW}}(f) = 3 \tilde{\Omega}_{\text{GW}} \mathfrak{K}_{\text{int,exp}}^2 (H_* R_*) S(f R_*)$$



Higgsless simulations (results)¹⁰

- In the literature, the GW spectrum from sound waves is usually assumed to be

$$\Omega_{\text{GW}}(f) = 3 \tilde{\Omega}_{\text{GW}} K^2 \mathcal{T}(\tau_{\text{sw}}) (H_* R_*) S(\epsilon, \beta)$$

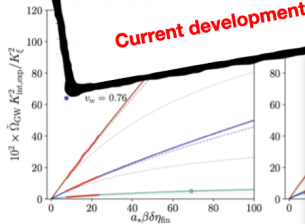
- The linear growth, which ends when the decay of the sound waves is modified

- sources based on Higgsless simulations available using public code CosmoGW

Template for sound-shell model and template for non-linear
sources based on Higgsless simulations available using public
code CosmoGW
(<https://github.com/cosmoGW/cosmoGW>)

proceedings on arXiv with I. Stomberg (2508.04263) and tutorial
available on CosmoGW documentation

For use: pip install cosmoGW
Current developments undergoing to extend CosmoGW to
other sources



Gravitational wave spectra for cosmological phase transitions with non-linear decay of the fluid motion

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COSMOGW review paper and official release (2609.xxyy)

Classical and Quantum Gravity

CONTRIBUTION TO "FOCUS ON THE LISA MISSION AT THE TIME OF ADOPTION"

COSMOGW: A public library to study cosmological backgrounds of gravitational waves. Part I — Phase transitions: Sound waves and MHD turbulence

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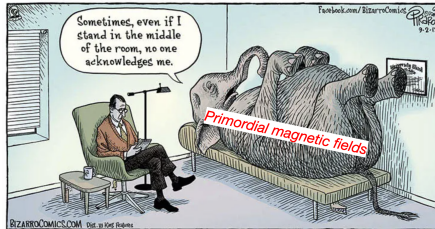
[Website](#)

Documentation available on Read the Docs, source code on GitHub.

- Routines to generate GW backgrounds from sound waves and turbulence associated to phase transitions
- Reading routines from numerical simulations (e.g., PENCIL CODE)
- Efficient numerical integration techniques
- Cosmology calculations
- Interferometry, detector sensitivities

Primordial magnetic fields

- Magnetic fields can either be produced at or present during cosmological phase transitions.
- The magnetic fields are strongly coupled to the primordial plasma and effectively produce vortical motion, inevitably leading to the development of MHD turbulence.¹²
- Present magnetic fields can be amplified by primordial turbulence via dynamo.¹³



¹²J. Ahonen and K. Enqvist, *Phys. Lett. B* **382**, 40 (1996).

¹³A. Brandenburg *et al.* (incl. ARP), *Phys. Rev. Fluids* **4**, 024608 (2019).

Generation of primordial magnetic fields during phase transitions

- Parity-violating processes during the electroweak phase transition ($T \sim 100$ GeV) are predicted by SM extensions that account for baryogenesis and can produce helical magnetic fields through sphaleron decay or B+L anomalies.¹⁴

$$\mathbf{B} = \nabla \times \mathbf{A} - i \frac{2 \sin \theta_w}{g v^2} \nabla \Phi^\dagger \times \nabla \Phi$$

- Also after inflation, axion fields can amplify and produce magnetic fields.¹⁵ For example, the QCD axion could oscillate and produce magnetic fields around the QCD scale ($T \sim 100$ MeV).¹⁶

$$\mathcal{L} \supset \frac{\phi}{f} F_{\mu\nu} \tilde{F}^{\mu\nu}$$

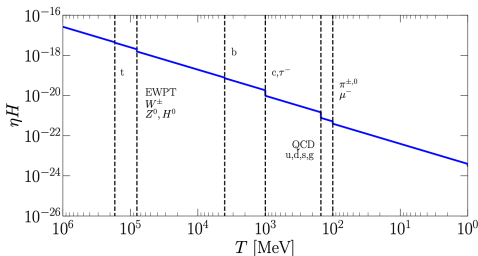
¹⁴T. Vachaspati, *Phys. Rev. B* **265**, 258 (1991), T. Vachaspati, *Phys. Rev. Lett.* **87**, 251302 (2001), J. M. Cornwall, *Phys. Rev. D* **56**, 6146 (1997).

¹⁵M. M. Forbes and A. R. Zhitnitsky, *Phys. Rev. Lett.* **85**, 5268 (2000).

¹⁶Miniati, Gregori, Reville & Sarkar (2018).

MHD turbulence (large conductivity)

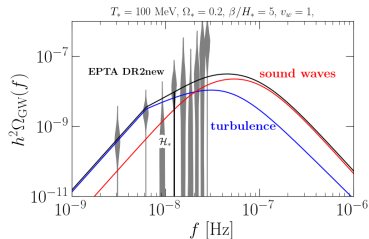
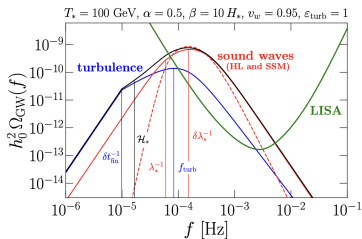
- High-conductivity (low magnetic diffusivity $\eta = 1/\sigma$) of the early universe leads to a high-coupling between magnetic and velocity fields.¹⁷



GW sources in the early Universe

- Magnetohydrodynamic (MHD) sources of GWs:
 - Sound waves generated from first-order phase transitions.
 - (M)HD turbulence from first-order phase transitions.
 - Primordial magnetic fields.

ARP *et al.*, 2307.10744, 2308.12943; [EPTA/InPTA] (incl. ARP), 2306.16227





Primordial magnetic fields and GWs from MHD turbulence

- During the radiation-dominated era, after they are produced, magnetic fields will decay following MHD turbulence.
- Direct numerical simulations using the `PENCIL CODE`¹⁸ to solve:
 - 1 Relativistic MHD equations adapted for radiation-dominated era (after electroweak symmetry is broken).
 - 2 Gravitational waves equation.
- In general, large-scale simulations are necessary to solve the MHD nonlinearities (e.g., unequal-time correlators UETC and non-Gaussianities, which require simplifying assumptions in analytical studies).

¹⁸Pencil Code Collaboration, JOSS **6**, 2807 (2020),
<https://github.com/pencil-code/>

MHD description

Right after the electroweak phase transition we can model the plasma using continuum MHD.

- Charge-neutral, electrically conducting fluid.
- Relativistic magnetohydrodynamic (MHD) equations.
- Radiation-dominated fluid

$$p = \rho/3,$$

i.e., $c_s^2 = 1/3$ (ultrarelativistic EoS).

- Friedmann–Lemaître–Robertson–Walker metric

$$g_{\mu\nu} = a^2 \eta_{\mu\nu}$$

Conservation laws for MHD turbulence

$$\nabla_\nu T^{\mu\nu} = 0, \quad \nabla_\nu F^{\mu\nu} = -J^\mu, \quad \nabla_\nu \mathcal{F}^{\mu\nu} = 0$$

In the limit of subrelativistic bulk flow:

$$\gamma^2 \sim 1 + (v/c)^2 + \mathcal{O}(v/c)^4$$

Relativistic MHD equations are reduced to¹⁹

$$\begin{aligned} \frac{\partial \ln \rho}{\partial t} &= -\frac{4}{3} \left(\nabla \cdot \mathbf{u} + \frac{1}{2} \mathbf{u} \cdot \nabla \ln \rho \right) + \frac{1}{\rho} [\eta \mathbf{J}^2 - \mathbf{u} \cdot (\mathbf{J} \times \mathbf{B})], \\ \frac{D\mathbf{u}}{Dt} &= \frac{\mathbf{u}}{3} \left(\nabla \cdot \mathbf{u} + \frac{1}{2} \mathbf{u} \cdot \nabla \ln \rho \right) - \frac{\mathbf{u}}{\rho} \left[\eta \mathbf{J}^2 + \frac{1}{2} \mathbf{u} \cdot (\mathbf{J} \times \mathbf{B}) \right] \\ &\quad - \frac{1}{4} \nabla \ln \rho + \frac{3}{4\rho} \mathbf{J} \times \mathbf{B} + \frac{2}{\rho} \nabla \cdot (\rho \nu \mathbf{S}), \\ \frac{\partial \mathbf{B}}{\partial t} &= \nabla \times (\mathbf{u} \times \mathbf{B} - \eta \mathbf{J}), \quad \mathbf{J} = \nabla \times \mathbf{B}, \end{aligned}$$

for a flat expanding universe with comoving and normalized

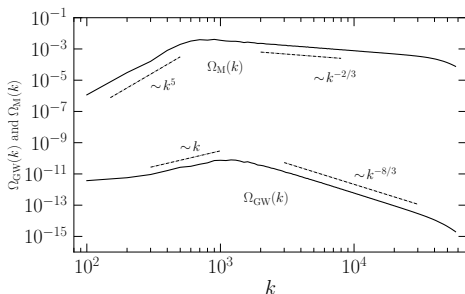
$\rho = a^4 \rho_{\text{phys}}, \rho = a^4 \rho_{\text{phys}}, B_i = a^2 B_{i,\text{phys}}, u_i$, and conformal time t ($dt = a dt_c$).

¹⁹ A. Brandenburg, *et al.*, *Phys. Rev. D* **54**, 1291 (1996).

ARP, Midiri, *Relativistic Magnetohydrodynamics in the early Universe*, *JCAP* **06** (2026) 070, arXiv:2501.05732 (2025).

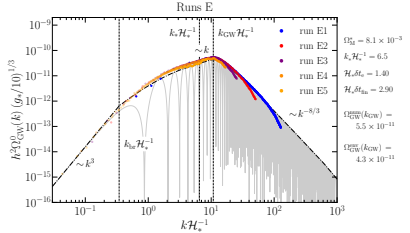
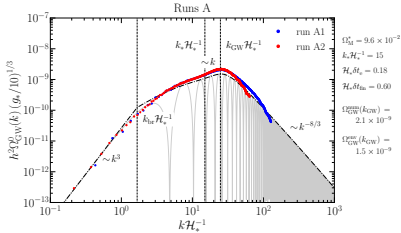
Numerical results for decaying MHD turbulence²⁰

$$1152^3, k_* = 2\pi \times 100, \Omega_M \sim 10^{-2}, \sigma_M = 1$$



- **Characteristic k scaling in the subinertial range for the GW spectrum.**
- k^2 expected at scales $k < k_*$ and k^3 at $k < H_*$ according to the “top-hat” model (Caprini *et al.*, 2020).

Numerical results for decaying MHD turbulence²¹



run	Ω_M^*	$k_* \mathcal{H}_*^{-1}$	$\mathcal{H}_* \delta t_e$	$\mathcal{H}_* \delta t_{\text{fin}}$	$\Omega_{\text{GW}}^{\text{num}}(k_{\text{GW}})$	$[\Omega_{\text{GW}}^{\text{num}}/\Omega_{\text{GW}}^{\text{num}}](k_{\text{GW}})$	n	$\mathcal{H}_* L$	$\mathcal{H}_* t_{\text{end}}$	$\mathcal{H}_* \eta$
A1	9.6×10^{-2}	15	0.176	0.60	2.1×10^{-9}	1.357	768	6π	9	10^{-7}
A2	—	—	—	—	—	—	768	12π	9	10^{-6}
E1	8.1×10^{-3}	6.5	1.398	2.90	5.5×10^{-11}	1.184	512	4π	8	10^{-7}
E2	—	—	—	—	—	—	512	10π	18	10^{-7}
E3	—	—	—	—	—	—	512	20π	61	10^{-7}
E4	—	—	—	—	—	—	512	30π	114	10^{-7}
E5	—	—	—	—	—	—	512	60π	234	10^{-7}

Analytical CiT model for GWs from decaying turbulence

- Assumption: magnetic or velocity field evolution $\delta t_e \sim 1/(u_* k_*)$ is slow compared to the GW dynamics ($\delta t_{\text{GW}} \sim 1/k$) at all $k \gtrsim u_* k_*$.
- We can derive an analytical expression for nonhelical fields of the envelope of the oscillations²² of $\Omega_{\text{GW}}(k)$.

$$\Omega_{\text{GW}}(k, t_{\text{fin}}) \approx 3 \left(\frac{k}{k_*} \right)^3 \Omega_{\text{M}}^*{}^2 \frac{\mathcal{C}(\alpha)}{\mathcal{A}^2(\alpha)} p_{\Pi} \left(\frac{k}{k_*} \right) \\ \times \begin{cases} \ln^2[1 + \mathcal{H}_* \delta t_{\text{fin}}] & \text{if } k \delta t_{\text{fin}} < 1, \\ \ln^2[1 + (k/\mathcal{H}_*)^{-1}] & \text{if } k \delta t_{\text{fin}} \geq 1. \end{cases}$$

- p_{Π} is the anisotropic stress spectrum and depends on spectral shape, can be approximated for a von Kármán spectrum as²³

$$p_{\Pi}(k/k_*) \simeq \left[1 + \left(\frac{k}{2.2 k_*} \right)^{2.15} \right]^{-11/(3 \times 2.15)}$$

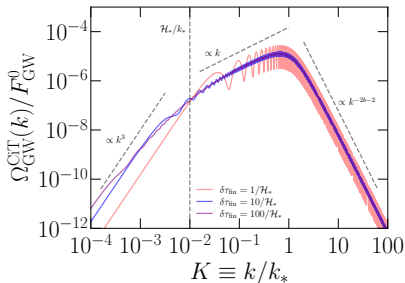
²² ARP *et al.*, *Phys. Rev. D* **105**, 123502 (2022); ARP, Midiri, Salomé, Caprini, in preparation.

²³ ARP *et al.*, arXiv:2307.10744 (2023).

Constant-in-Time (CiT) model (arXiv:2609.xxzz)

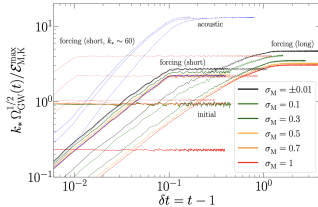
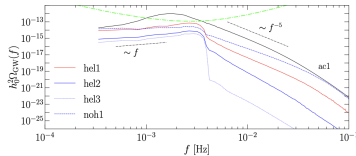
[PRELIMINARY work with A. Midiri, M. Salomé, C. Caprini]

- Applicable to generic slowly-decaying sources
- Extension to include decay of the source impacts IR part of the spectrum
- Used to compute circular polarization, compatible with numerical results.
- Applicable to generic vector and scalar fields sourcing GWs.
- Application to sound waves when $c_s \ll 1$ and to turbulence when $u_{dc} \ll 1$.



Acoustic vs vortical sourcing

- Numerical simulations show that acoustic forcing is more efficient than vortical in producing GWs²⁴



- Counterintuitive since vortical motion contains more energy in the tensor-mode of stresses²⁵.

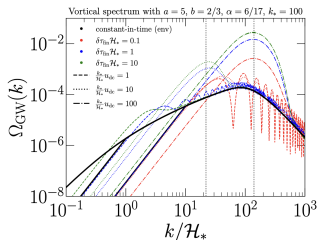
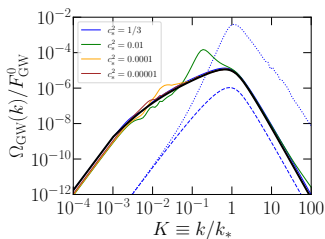
²⁴ ARP *et al.*, *Phys.Rev.D* **102** (2020) 8, 083512
 ARP *et al.*, *JCAP* **04** (2022) 04, 019.

²⁵ Mandal *et al.* (incl. ARP), *Class. Quant. Grav.* **38** (2021) 14, 145002

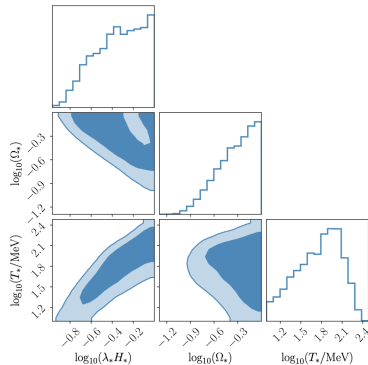
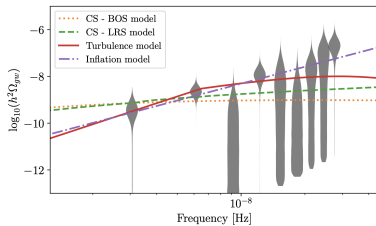
Acoustic vs vortical sourcing

[PRELIMINARY work with A. Midiri, M. Salomé, C. Caprini]

- Key point to solve this question is the applicability (or non-applicability) of the CiT model to sound waves and turbulence.

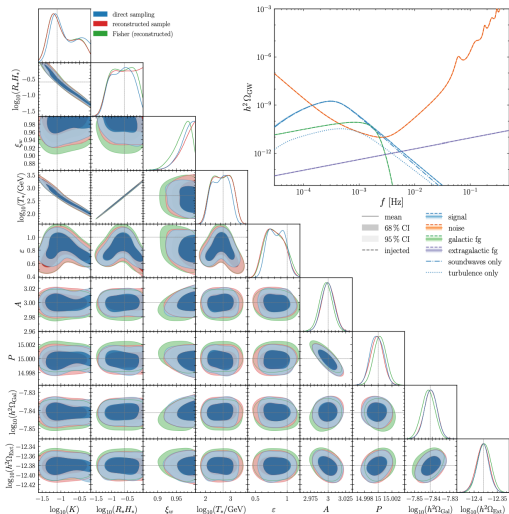


Primordial magnetic fields constraints with PTA²⁶



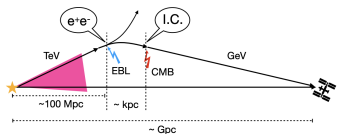
²⁶[EPTA/InPTA collab.] (incl. ARP), arXiv:2306.16227 (2023).

Parameter reconstruction of GW signal for sound waves and turbulence from a PT with LISA²⁷

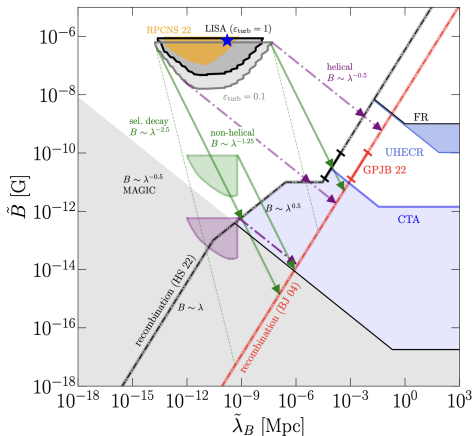


Multi-messenger studies of primordial magnetic fields³⁰

- Primordial magnetic fields would evolve through the history of the universe up to the present time and could explain the lower bounds in cosmic voids derived by the Fermi collaboration.³¹



- Maximum amplitude of primordial magnetic fields is constrained by the big bang nucleosynthesis.³²
- Additional constraints from CMB, Faraday Rotation, ultra-high energy cosmic rays (UHECR).



³⁰ ARP et al., arXiv:2307.10744 (2023).

³¹ A. Neronov and I. Vovk, *Science* **328**, 73 (2010).

³² V. F. Shvartsman, *Pisma Zh. Eksp. Teor. Fiz.* **9**, 315 (1969).

Conclusions

- Velocity and magnetic fields in the early universe can significantly contribute to the cosmological GW background (GWB) via sound waves and (M)HD turbulence.
- A clean observation of a GWB from the early Universe would probe BSM physics that are not accessible to electromagnetic observations and that can be complementary to collider experiments in some BSM theories.
- The GWB produced by non-linear fluid (and magnetic field) dynamics requires, in general, performing high-resolution numerical simulations, which can be done using open-source codes like `PENCIL CODE` or `CosmoLattice`.
- We expect the dynamics of first-order phase transitions to be affected by gauge fields, with its exact impact on the resulting GWB and on the fluid dynamics sourcing GWs being at the moment uncertain.
- Primordial magnetic fields produced in the early Universe can be probed by other cosmological (CMB) and astrophysical (γ -ray, radio observations at voids, filaments of the LSS), providing us a multi-messenger approach to study early Universe physics.
- To extract clean information from a GWB observation in LISA, PTA, or ET, we will need to have accurate modeling of the resulting GWB, since the reconstruction of the signal is very challenging (multiple cosmo sources, astrophysical background, detectors noise). Current state-of-the-art models provide good estimates but still require improvement for extracting clean information of the underlying source. Python packages like `COSMOGW` available.



Thank You!

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github.com/AlbertoRoper/cosmoGW
albertoroper.github.io



Exploring Cosmological Phase Transitions with Gravitational Waves

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Co-organized with
M. Nemevšek, D. Racco, P. Schicho, B. Świeżewska.

