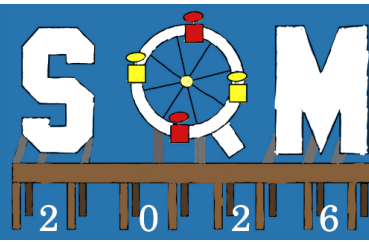


The 22<sup>nd</sup> International Conference on  
**Strangeness in Quark Matter**  
22-27 March, 2026, Los Angeles, CA



# Hyperon Spin Observables in Au+Au Collisions at RHIC BES-II: Global and Local Polarization, Spin Correlations

Tong Fu (付瞳) for the STAR Collaboration

Shandong University (山东大学)

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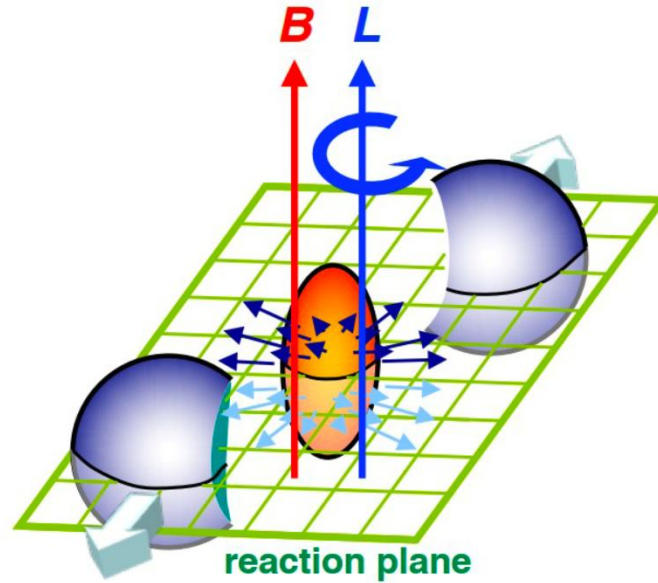
山东大学

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Tong Fu



- Global polarization
  - Global polarization of  $\Lambda$
  - Global polarization of Multi-strange hyperons( $\Xi, \Omega$ )
  
- Local polarization
  - Local polarization of  $\Lambda$
  
- Measurement status of  $\Lambda$ –(anti)– $\Lambda$  hyperon spin correlation in Au+Au Collisions
  
- Summary



Z.-T. Liang and X.-N. Wang, Phys. Rev. Lett. 94, 102301 (2005)

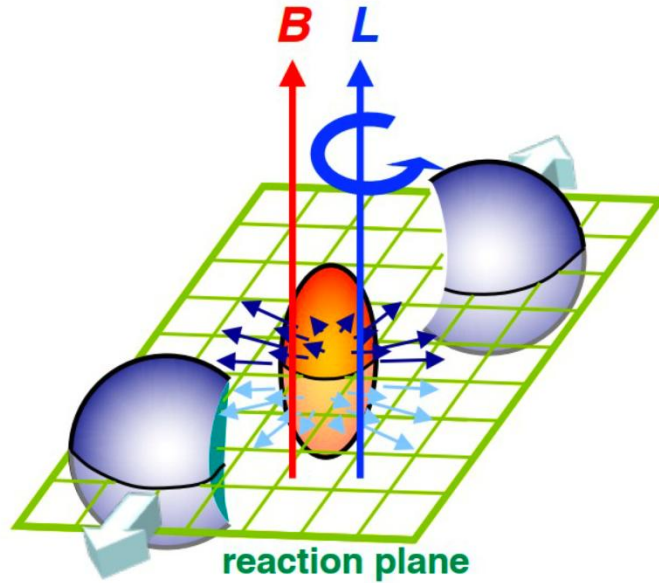
- Non-central HICs have large initial angular momentum and magnetic field



- Polarize quarks due to “spin-orbit” interaction



- Polarization of the final-state hadrons



Z.-T. Liang and X.-N. Wang, Phys. Rev. Lett. 94, 102301 (2005)

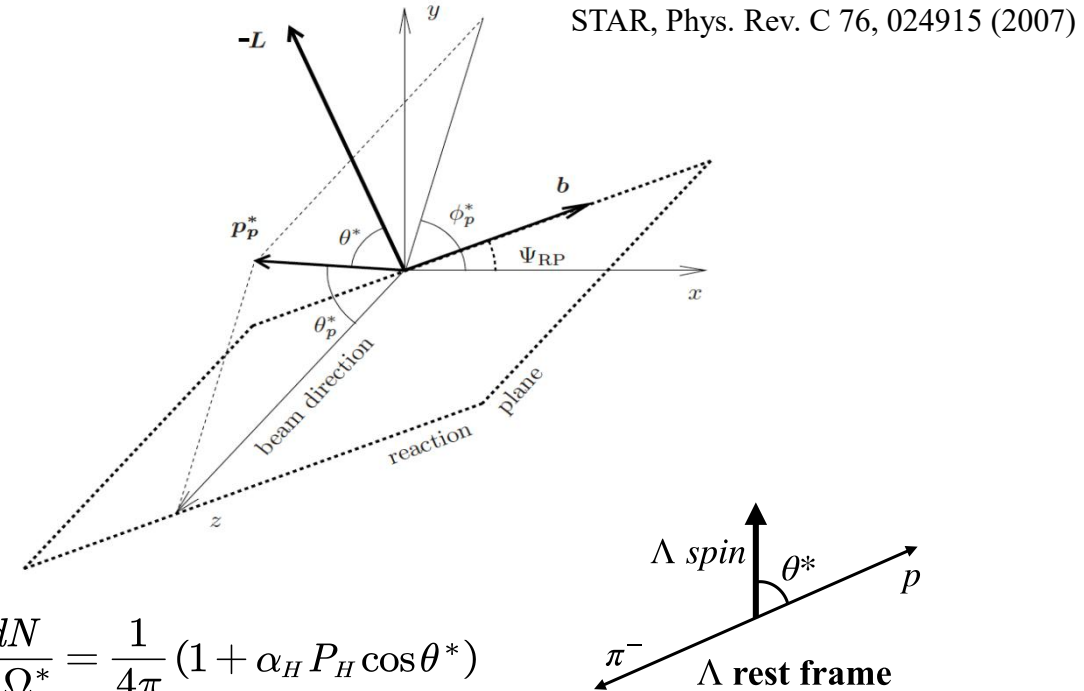
- Non-central HICs have large initial angular momentum and magnetic field



- Polarize quarks due to “spin-orbit” interaction



- Polarization of the final-state hadrons



STAR, Phys. Rev. C 76, 024915 (2007)

$$\frac{dN}{d\Omega^*} = \frac{1}{4\pi} (1 + \alpha_H P_H \cos\theta^*)$$



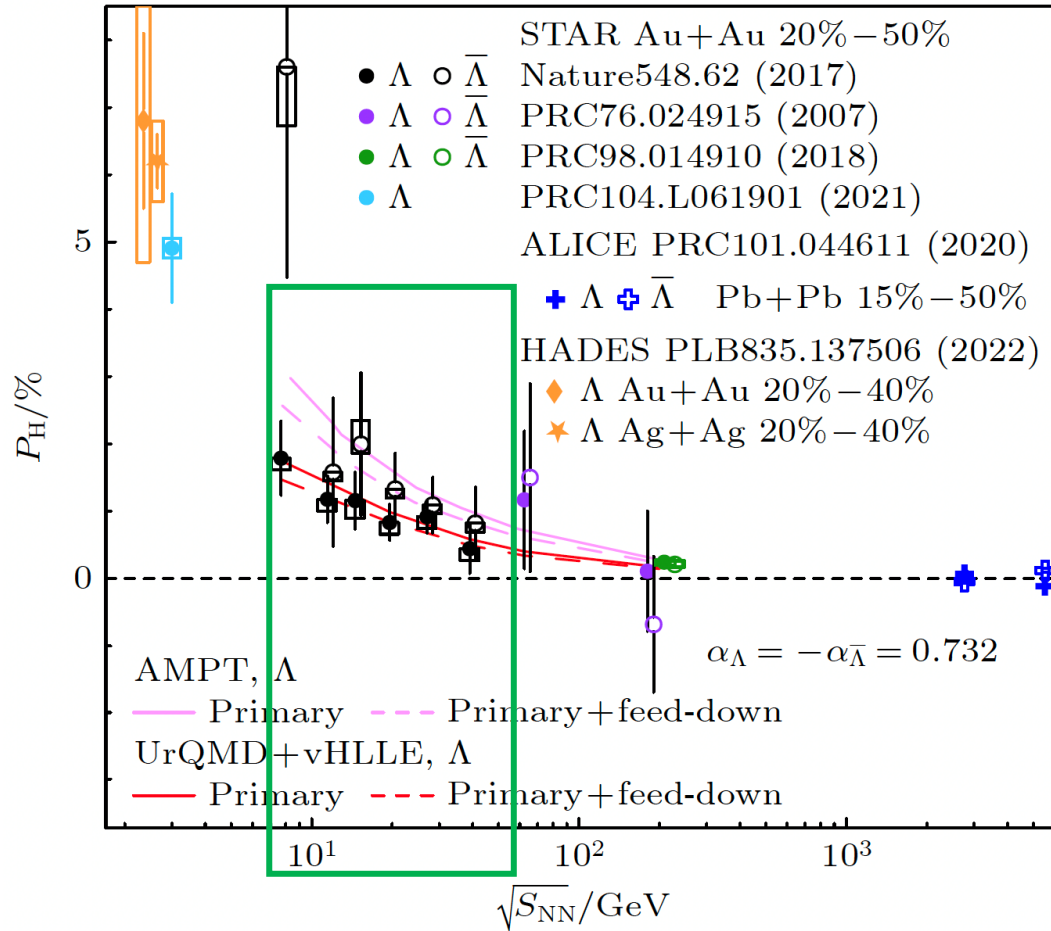
$$P_H = \frac{8}{\alpha\pi} \frac{1}{A_0} \frac{\langle \sin(\Psi_1 - \phi_p^*) \rangle}{Res(\Psi_1)}$$

- $\alpha_H$  is the hyperon decay parameter,  $\alpha_\Lambda = 0.732 \pm 0.014$
- $\phi^*$  is the azimuthal angle of the daughter proton in  $\Lambda$  rest frame
- $A_0$  is an acceptance correction factor,  $A_0 = \langle \sin\theta_p^* \rangle$

# Observation of $\Lambda$ global polarization



X. Sun et al., Acta Phys. Sin. Vol. 72, No. 7(2023)

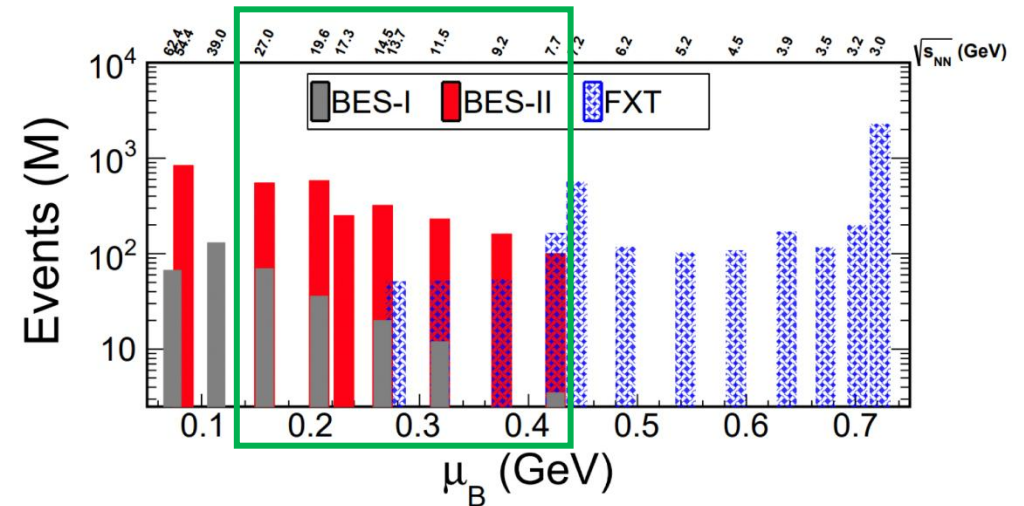


Strongest vorticity observed in nature

$$\omega \sim 10^{21} \text{ s}^{-1}$$

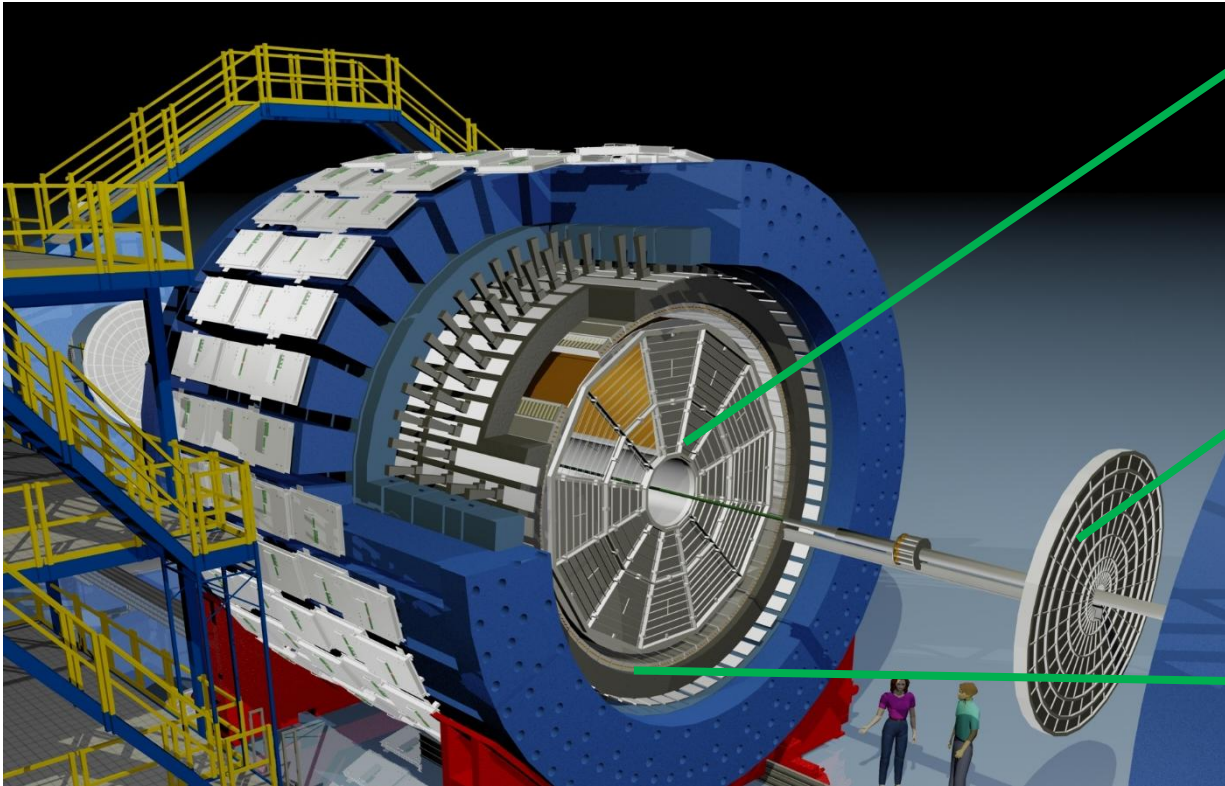
How does the late-stage magnetic field affect global polarization?

$$|B| \approx \frac{T_s |P_\Lambda - P_{\bar{\Lambda}}|}{2 |\mu_\Lambda|} \quad \text{F. Becattini et al., Phys. Rev. C 95.054902(2017)}$$



The BES-II by STAR collected an order of magnitude more data compared to BES-I, and collected two additional energy points ( 9.2, 17.3 GeV )

## Sub-system relevant to this analysis:



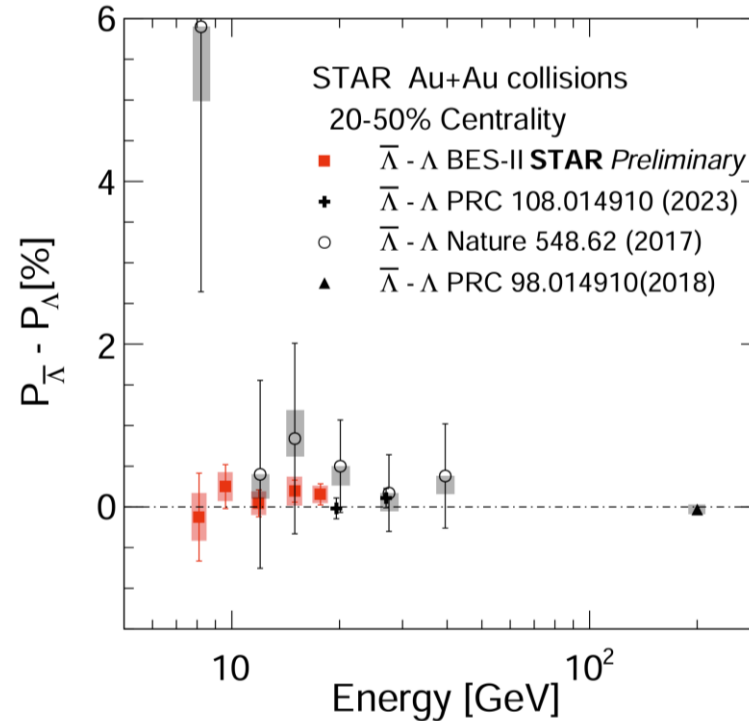
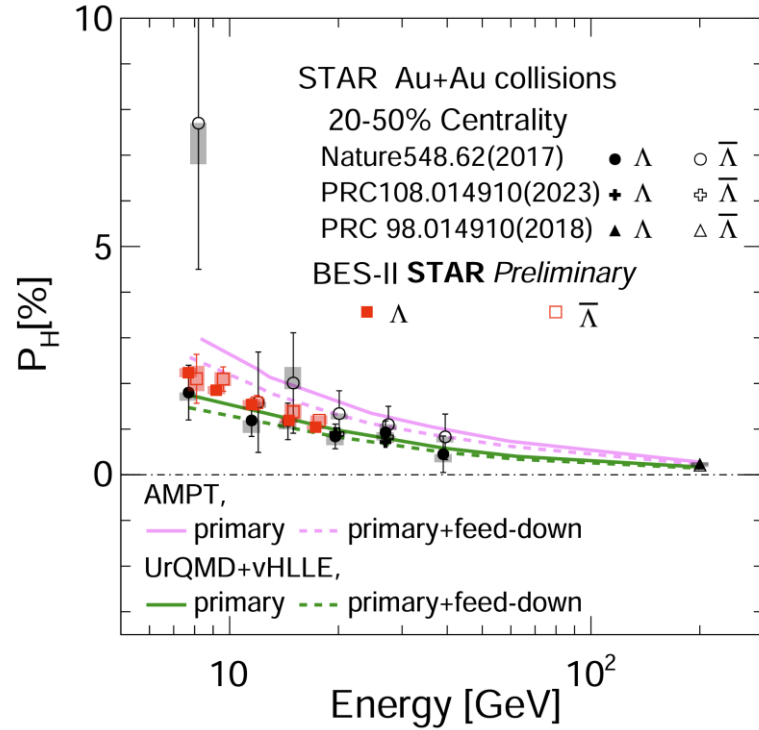
- **Time Projection Chamber**
  - Particle reconstruction
  - The iTPC upgrade extended the pseudo-rapidity coverage from  $|\eta| < 1$  to  $|\eta| < 1.5$
- **Event Plane Detector**
  - Event plane reconstruction
  - $2.1 < |\eta| < 5.1$
  - Improved the event plane resolution
- **Time Of Flight**
  - PID via particle velocity
  - $|\eta| < 0.9$

# $\Lambda$ global polarization



UrQMD+vHLLLE : I. Karpenko, F. Becattini Eur. Phys. J. C 77, 213 (2017)  
 AMPT: H. Li, L. Pang, Q. Wang, X. Xia Phys. Rev. C 96, 054908(2017)

F. Becattini et al., PRC.95.054902(2017)



$$P_{\Lambda} \approx \frac{1}{2} \frac{\omega}{T} + \frac{\mu_{\Lambda} B}{T}$$

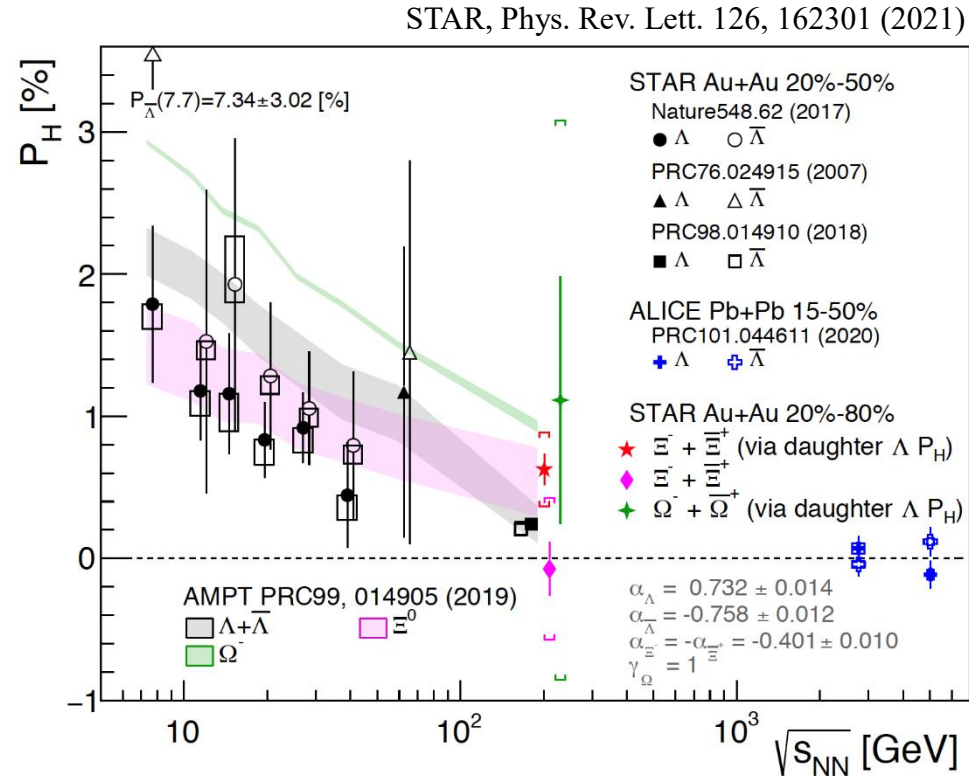
$$P_{\bar{\Lambda}} \approx \frac{1}{2} \frac{\omega}{T} - \frac{\mu_{\Lambda} B}{T}$$

$$\Delta P_H = |P_{\bar{\Lambda}} - P_{\Lambda}| \approx \frac{2|\mu_{\Lambda}|B}{T}$$

$$T = 150 \text{ MeV}, \mu_{\Lambda} = -1.93 \times 10^{-1} \text{ MeV}/T$$

- ❑ Significant improvement in precision was achieved, collision energy dependence consistent with BES-I
  - $P_H = 1.17 \pm 0.40(stat) \pm 0.27(syst) \%$  [BES-I]  $\longrightarrow$   $1.19 \pm 0.04(stat) \pm 0.05(syst)\%$  [BES-II] at 14.6 GeV
- ❑ No obvious splitting between  $\Lambda$  and  $\bar{\Lambda}$  global polarization with high precision. Upper limit on late-stage magnetic field
  - $B \lesssim 10^{13} \text{ T}$  (95% confidence level) STAR, Phys. Rev. C 108, 014910 (2023)

Hyperon	Spin
$\Lambda(uds)$	1/2
$\Xi^-(dss)$	1/2
$\Omega^-(sss)$	3/2



- ❑ Multi-strange hyperons are expected to be polarized in similar way
- ❑ The  $\Xi$  polarization is measured to be slightly larger than  $\Lambda$ . The  $\Omega$  global polarization shows a hint of being even larger
- ❑ Is there a dependence of  $\Xi$  and  $\Omega$  global polarization on collision energy?

# Extraction of $\Xi$ and $\Omega$ global polarization



PDG2021

Hyperon	Decay mode	$\alpha_H$	Spin
$\Lambda(uds)$	$\Lambda \rightarrow p + \pi^-$	0.732	1/2
$\Xi^-(dss)$	$\Xi^- \rightarrow \Lambda + \pi^-$	-0.401	1/2
$\Omega^-(sss)$	$\Omega^- \rightarrow \Lambda + K^-$	0.0157	3/2

## Two methods for extracting the $\Xi$ and $\Omega$ polarization

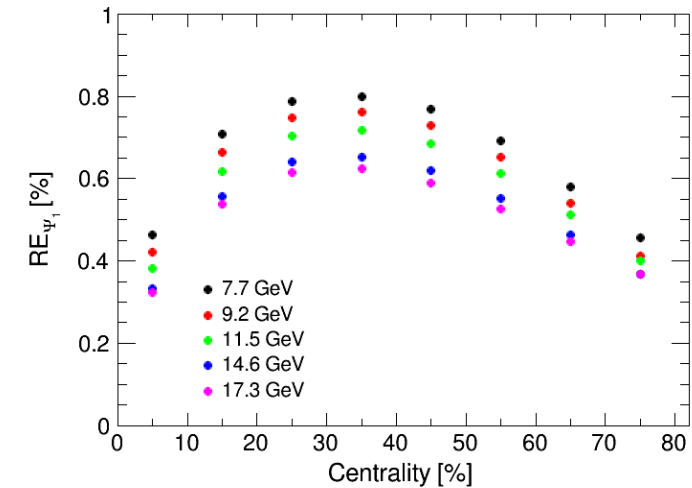
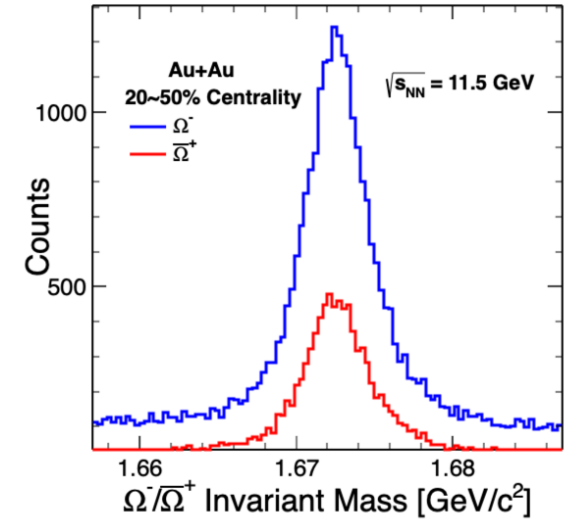
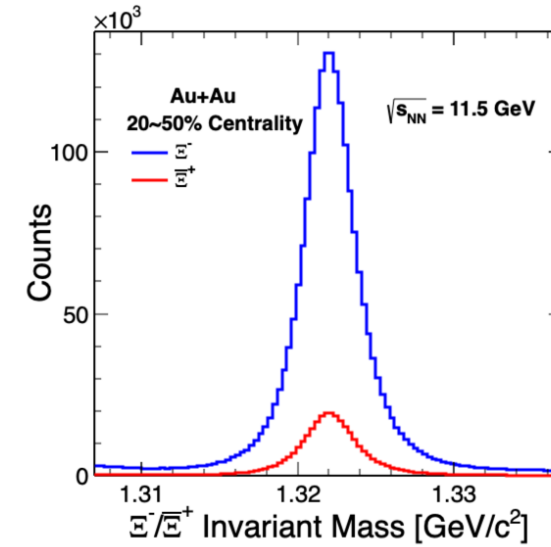
- Measured via daughter  $\Lambda$  angle distribution in  $\Xi$  rest frame

$$P_H = \frac{8}{\alpha_H \pi} \frac{1}{A_0} \frac{\langle \sin(\Psi_1 - \phi^*) \rangle}{Res(\Psi_1)}$$

$\Xi^-$  rest frame

- via daughter  $\Lambda$  polarization with spin transfer factor ( $C_{\Xi^- \rightarrow \Lambda} = 0.944, C_{\Omega^- \rightarrow \Lambda} = 1.0$  is assumed)

## First-order event plane reconstructed by EPD

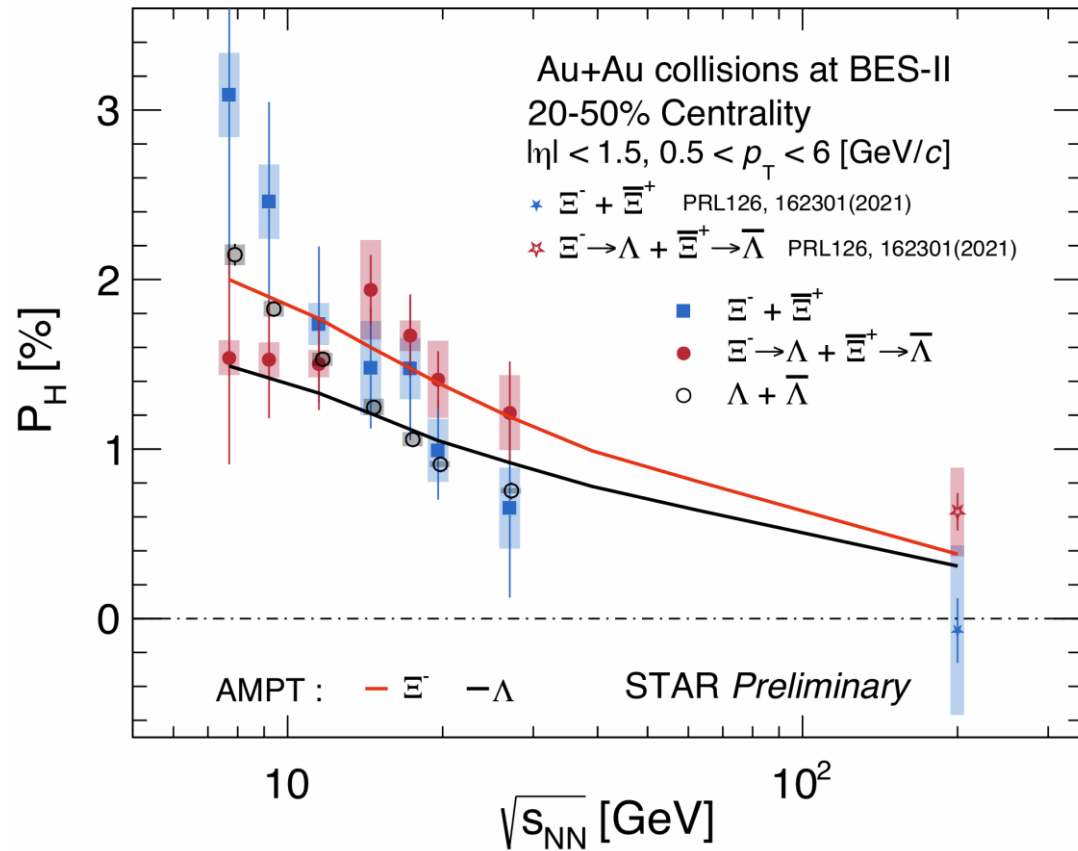


# $\Xi^- + \bar{\Xi}^+$ global polarization



Model calculation:

H. Li, X. Xia, X. G. Huang, H. Z. Huang Phys. Lett. B 827, 136971 (2022)



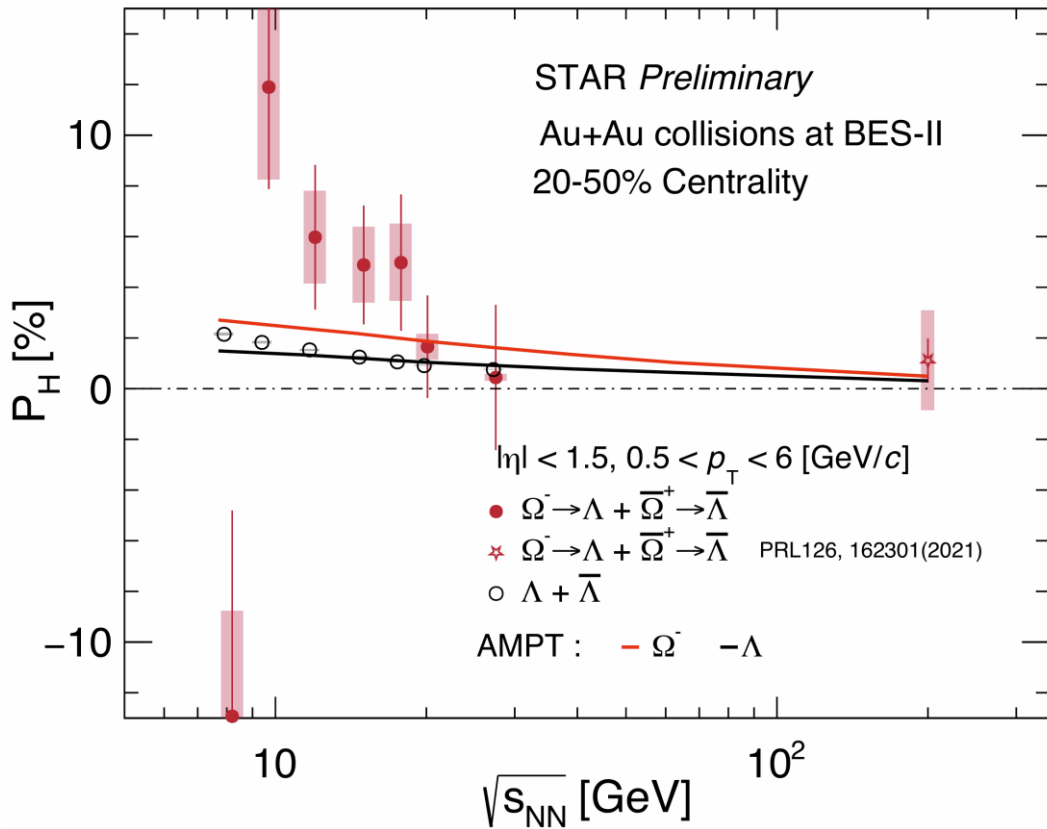
- For the first time, significant global polarization of  $\Xi^- + \bar{\Xi}^+$  has been observed ( $\sim 5 \sigma$ )
  - $P_H = 1.940 \pm 0.205(stat.) \pm 0.293(syst.)$  at 14.6 GeV
- Global polarization of  $\Xi^- + \bar{\Xi}^+$  decrease with collision energy
- $\Xi^- + \bar{\Xi}^+$  global polarization are consistent between direct and indirect measurement methods
- No obvious difference between  $\Lambda + \bar{\Lambda}$  and  $\Xi^- + \bar{\Xi}^+$  global polarization within uncertainties

# $\Omega^- + \bar{\Omega}^+$ global polarization



Model calculation:

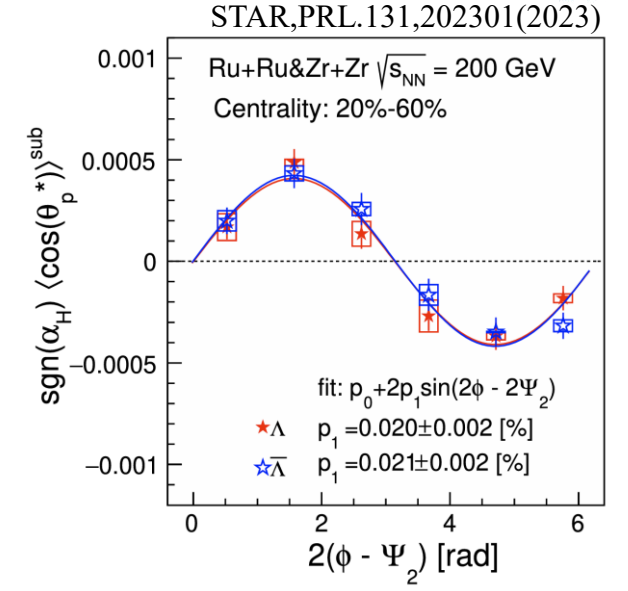
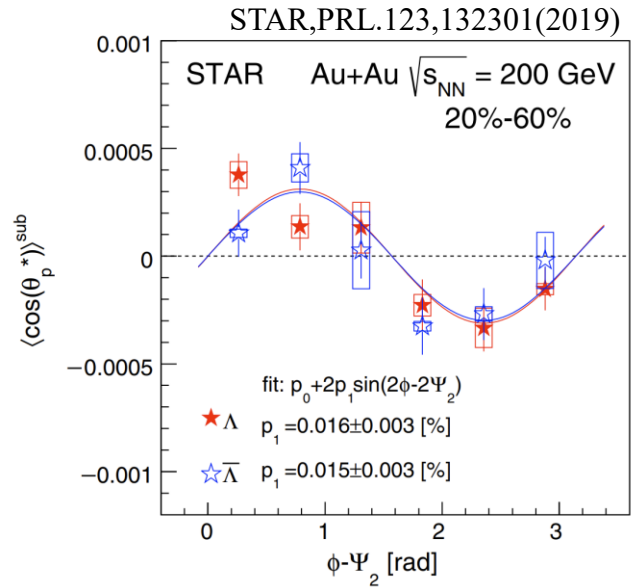
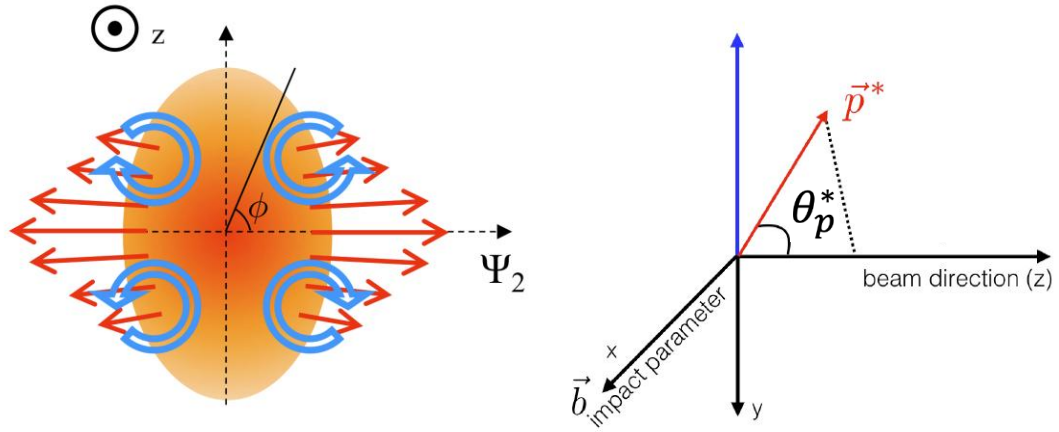
H. Li, X. Xia, X. G. Huang, H. Z. Huang Phys. Lett. B 827, 136971 (2022)



- Global polarization of  $\Omega^- + \bar{\Omega}^+$  tend to decrease with collision energy
- A hint of larger  $\Omega^- + \bar{\Omega}^+$  polarization than  $\Lambda + \bar{\Lambda}$  in lower energies
- Possible  $\Lambda, \Xi, \Omega$  global polarization difference?

$$P_\Lambda \cong P_S, \text{ assuming that } P_{u,d} \sim P_S \longrightarrow P_\Xi \sim P_\Lambda, P_\Omega \sim \frac{5}{3} P_\Lambda$$

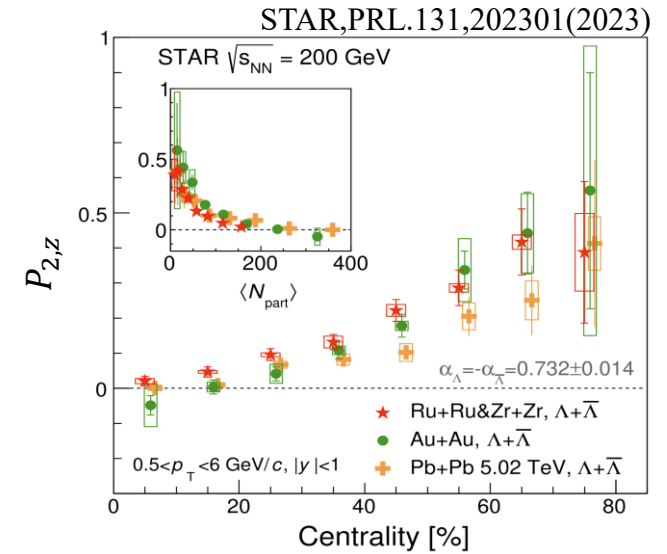
F. Becattini et al., Phys. Rev. C 95.054902(2017)

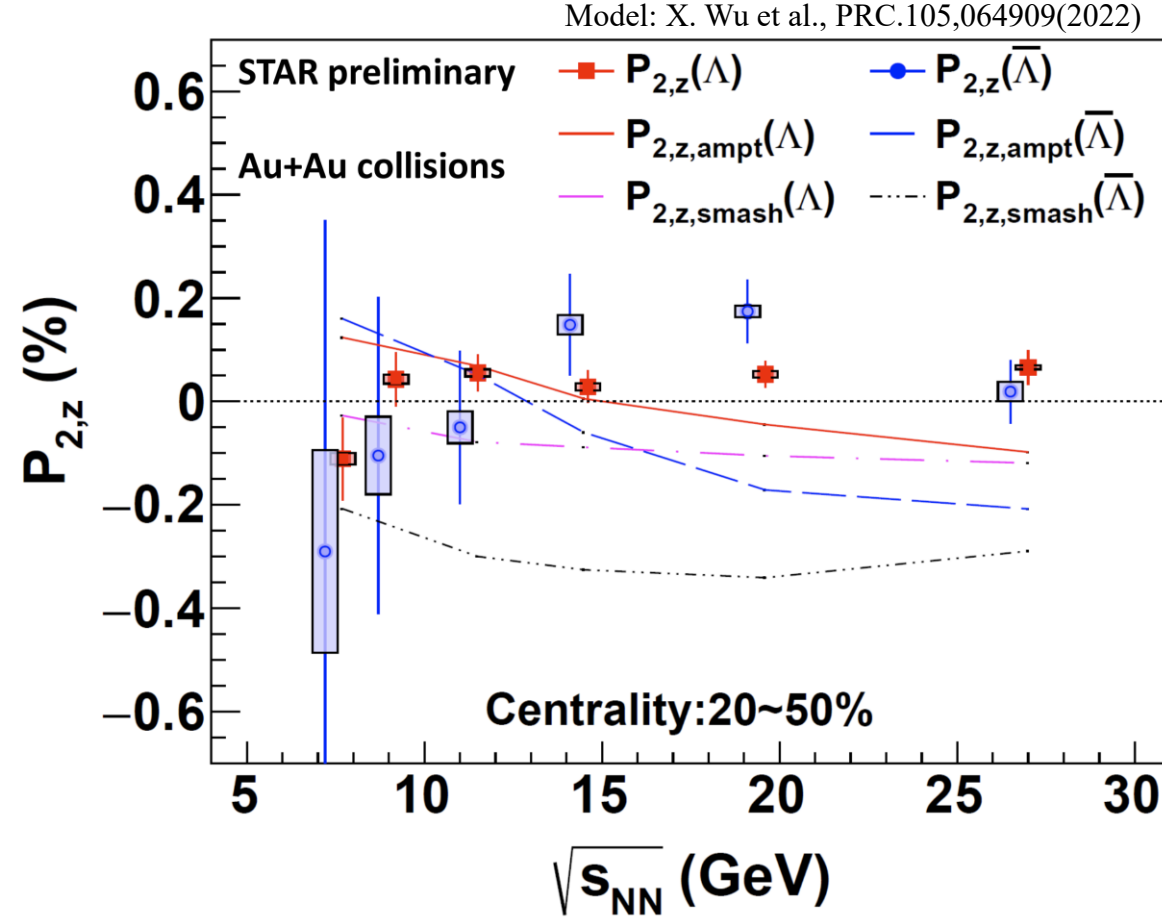


□ Anisotropic expansion of QGP leads to vorticity and particle polarization

$$P_{2,z} = \frac{1}{\alpha_H \langle \cos^2 \theta_p^* \rangle} \frac{\langle \cos \theta_p^* \sin [2(\phi_\Lambda - \Psi_2)] \rangle}{\text{Res} \{ \Psi_2 \}}$$

□  $\Lambda$  local polarization has been observed by STAR in Au+Au and isobar collisions at 200 GeV.

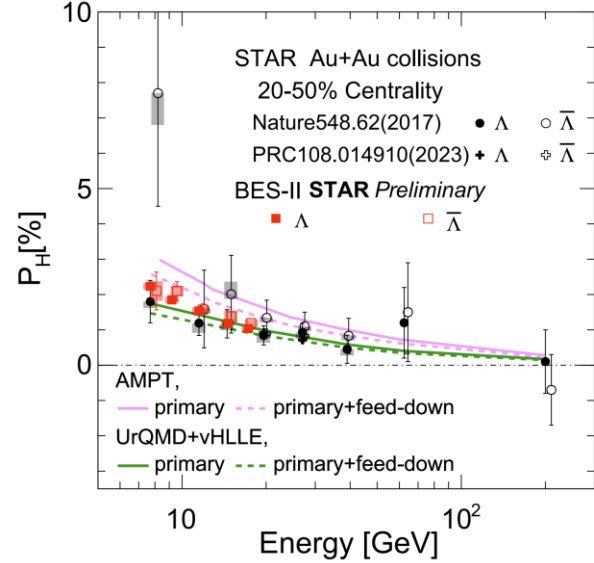




- $\square$   $\Lambda$  local polarization measured in Au+Au collisions at BES-II energies
- $\square$  The magnitude of  $P_{2,z}$  is small, with no significant collision energy dependence observed.

## global polarization

## spin alignment of $\phi$ -meson



$\phi$  - meson ( $|y| < 1.0$  &  $1.2 < p_T < 5.4$  GeV/c)

$\sqrt{s_{NN}}$	$\rho_{00}$
11.5	$0.365 \pm 0.014$ stat. $\pm 0.008$ syst.
19.6	$0.359 \pm 0.008$ stat. $\pm 0.007$ syst.
27	$0.3608 \pm 0.0025$ stat. $\pm 0.004$ syst.

STAR Collaboration. Nature 614, 244–248 (2023).

$$P_q = P_{\bar{q}} \sim 10^{-2}$$

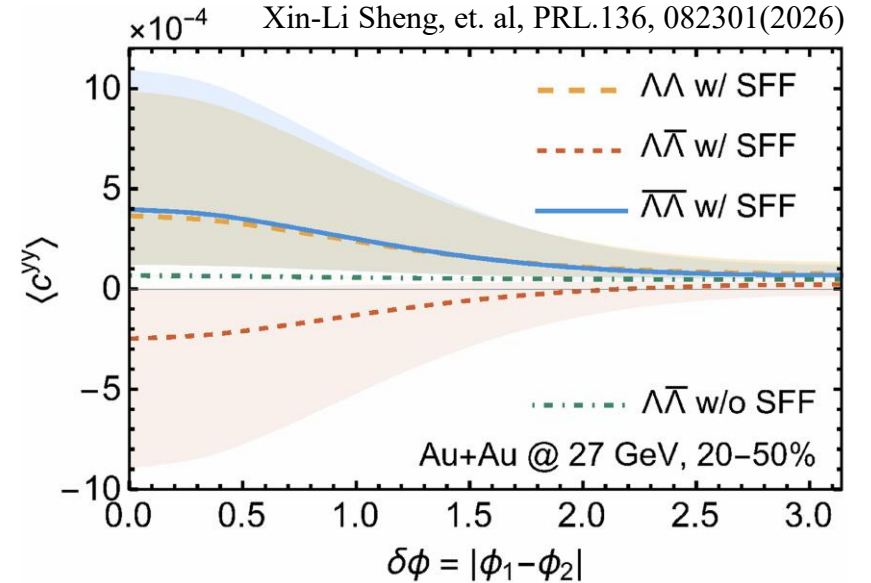
$$\rho_{00} - \frac{1}{3} \sim 10^{-2} \gg P_q^2 (10^{-4})$$



$$\langle P_q P_{\bar{q}} \rangle \neq \langle P_q \rangle \langle P_{\bar{q}} \rangle$$

This indicates a correlation between  $P_q$  and  $P_{\bar{q}}$

J.-p. Lv, Z.-h. Yu, Z.-t. Liang, Q. Wang, and X.-N. Wang, Phys. Rev. D 109, 114003 (2024)

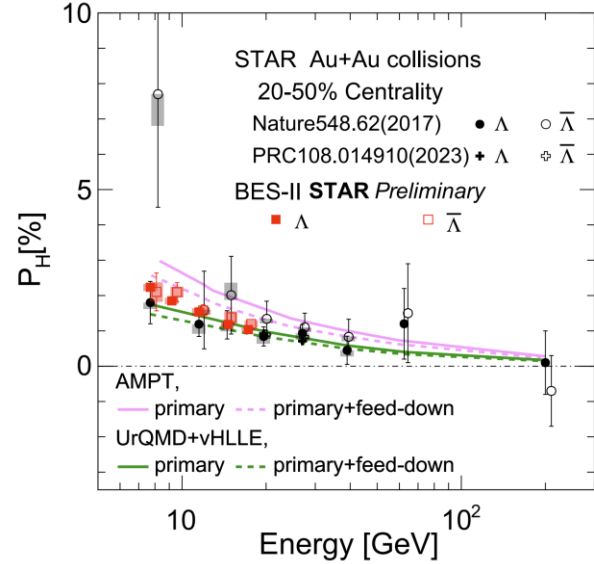


Models have predicted the spin correlation of  $\Lambda$  in heavy-ion collisions.

## global polarization

## spin alignment of $\phi$ -meson

## Spin correlations in pp measurement.



$\phi$  - meson ( $|y| < 1.0$  &  $1.2 < p_T < 5.4$  GeV/c)

$\sqrt{s_{NN}}$	$\rho_{00}$
11.5	$0.365 \pm 0.014$ stat. $\pm 0.008$ syst.
19.6	$0.359 \pm 0.008$ stat. $\pm 0.007$ syst.
27	$0.3608 \pm 0.0025$ stat. $\pm 0.004$ syst.

STAR Collaboration. Nature 614, 244–248 (2023).

$$\frac{dN}{d\cos(\theta_{12}^*)} = 1 + \alpha_1 \alpha_2 P_{\Lambda_1 \Lambda_2} \cos(\theta_{12}^*)$$

Gong, W. et al., PhysRevD.106.L031501(2022)

$\theta_{12}^*$  is the angle between the momentum of the protons, each boosted to the rest frame of their parent particle

A nonzero correlation is observed in short-range  $\Lambda\bar{\Lambda}$  pairs.

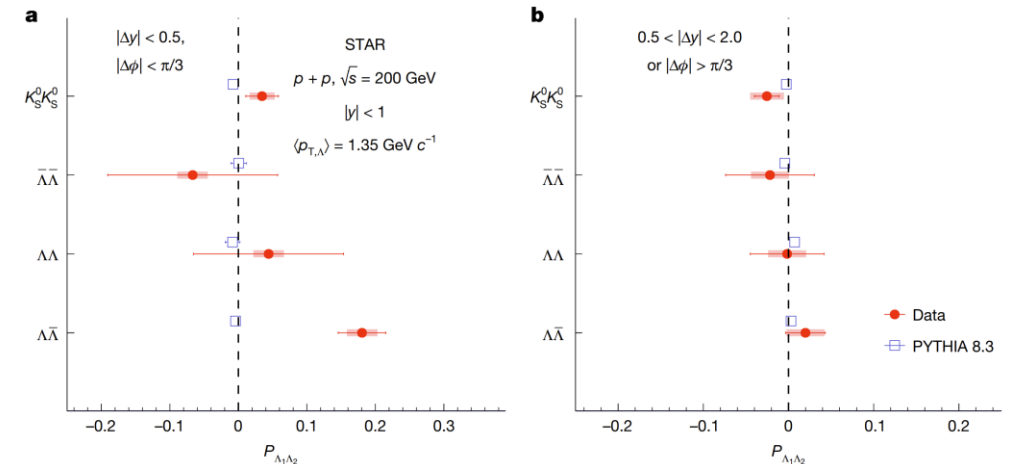
$$P_q = P_{\bar{q}} \sim 10^{-2}$$

$$\rho_{00} - \frac{1}{3} \sim 10^{-2} \gg P_q^2 (10^{-4})$$

$$\langle P_q P_{\bar{q}} \rangle \neq \langle P_q \rangle \langle P_{\bar{q}} \rangle$$

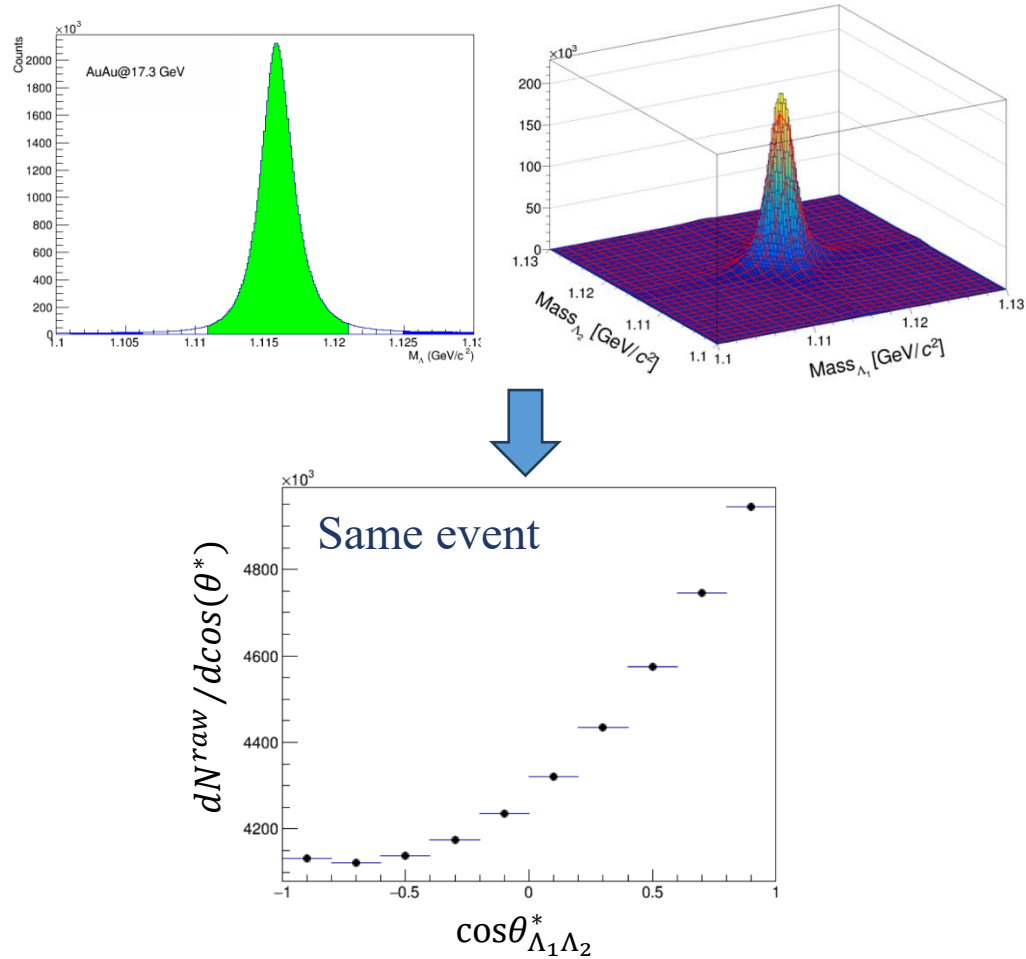
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J.-p. Lv, Z.-h. Yu, Z.-t. Liang, Q. Wang, and X.-N. Wang, Phys. Rev. D 109, 114003 (2024)

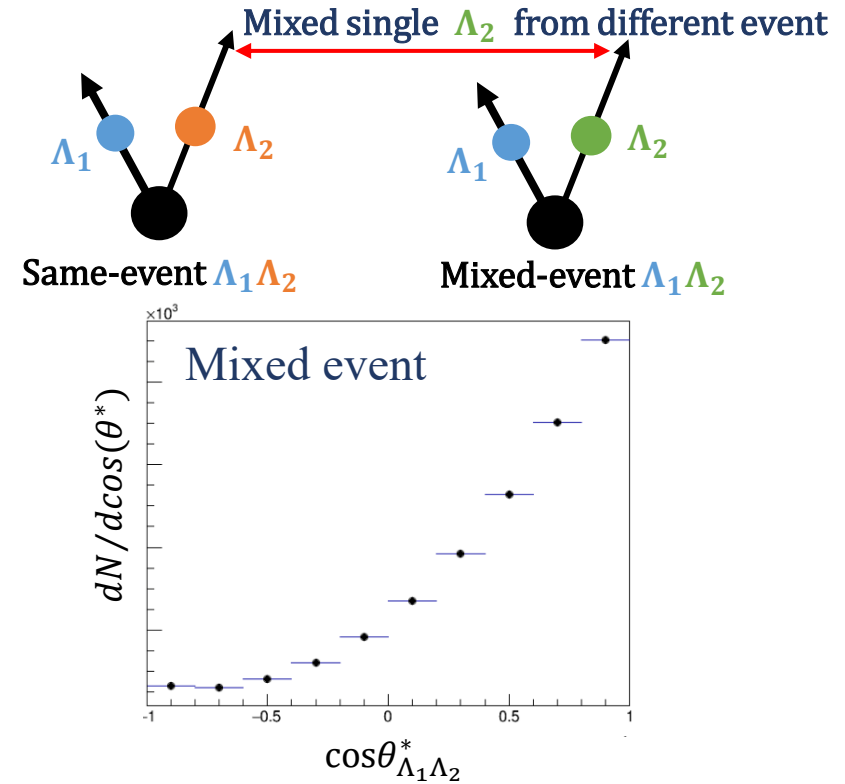


STAR Collaboration. Nature 650, 65–71 (2026).

- Measured  $dN^{raw}/d\cos(\theta^*)$  distributions.



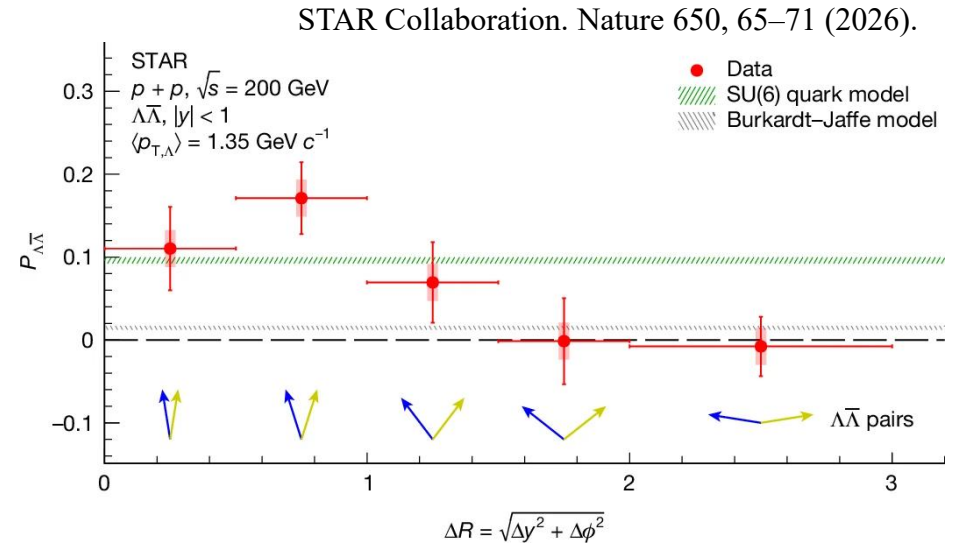
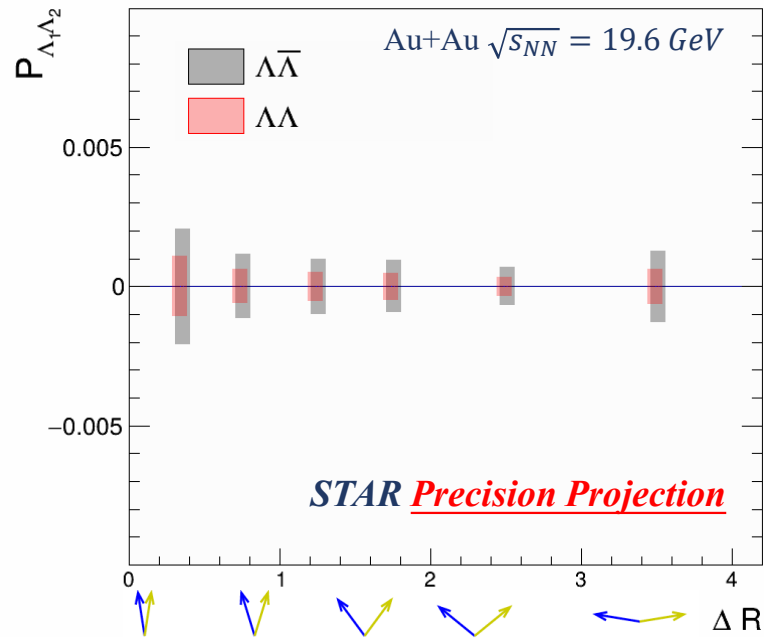
- Using the mixed-event method to correct for detector effects.



The spin correlation signal can be extracted by fitting the  $\cos\theta^*$  distribution after mixed-event correction.

$$\frac{dN}{d\cos(\theta_{12}^*)} = 1 + \alpha_1 \alpha_2 P_{\Lambda_1\Lambda_2} \cos(\theta_{12}^*)$$

# $\Lambda\Lambda$ and $\Lambda\bar{\Lambda}$ spin correlation statistical projection



- Measurement precision in Au+Au collisions is better than in pp collisions, detailed analysis is on-going.

## □ Global polarization of $\Lambda$ , $\Xi$ and $\Omega$ at BES-II

- High precision measurements of  $\Lambda$  polarization with BES-II, no splitting between  $P_\Lambda$  and  $P_{\bar{\Lambda}}$  within uncertainties
- A significant  $\Xi^- + \bar{\Xi}^+$  global polarization is observed ( $\sim 5\sigma$ ), and also decrease with collision energy. Hint of larger  $\Omega^- + \bar{\Omega}^+$  polarization at lower energies observed.

## □ Local polarization of $\Lambda$ at BES-II

- The magnitude of  $P_{2,z}$  is small, with no significant collision energy dependence observed.

## □ Spin correlation of $\Lambda$ - (anti)- $\Lambda$ in Au+Au collisions

- Detector acceptance corrected using mixed-event method. Further checks are on-going.

# *Back Up*

# Why $C_{\Omega^- \rightarrow \Lambda} = 1.0$



K. B. Luk et al. Phys. Rev. D 38, 19 (1988)

The weak decay  $\Omega^- \rightarrow \Lambda + K^-$  proceeds to the final states of orbital angular momentum  $L=1,2$  through the amplitudes  $A_L$ . The asymmetry parameters in the decay process can be written as:

$$\alpha = \frac{2 \operatorname{Re}(A_1^* A_2)}{|A_1|^2 + |A_2|^2}, \quad \beta = \frac{2 \operatorname{Im}(A_1^* A_2)}{|A_1|^2 + |A_2|^2}, \quad \gamma = \frac{|A_1|^2 - |A_2|^2}{|A_1|^2 + |A_2|^2}, \quad \alpha^2 + \beta^2 + \gamma^2 = 1.$$

When the joint probability distribution for the decay chain is integrated over the  $\Lambda$  angular distribution, the angular distribution of the daughter proton is given by:

$$\frac{dn}{d\Omega_p} = \frac{1}{4\pi} (1 + \alpha_\Lambda \mathbf{P}_\Lambda \cdot \hat{\mathbf{p}}), \quad \frac{dN}{d\Omega_p} = \frac{1}{4\pi} \left[ 1 + \frac{\alpha_\Lambda}{2(J+1)} [1 + (2J+1)\gamma_\Omega] \mathbf{P}_\Omega \cdot \hat{\mathbf{p}} \right]. \quad \Rightarrow \quad \mathbf{P}_\Lambda = \frac{1}{2(J+1)} [1 + (2J+1)\gamma_\Omega] \mathbf{P}_\Omega.$$

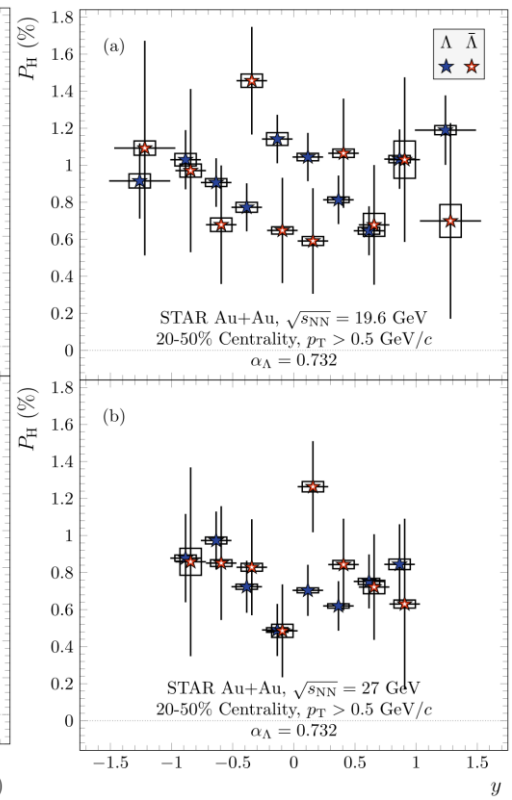
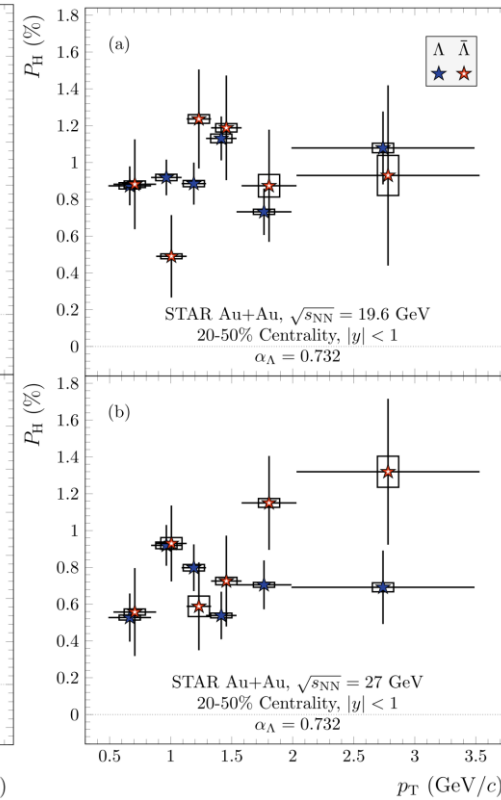
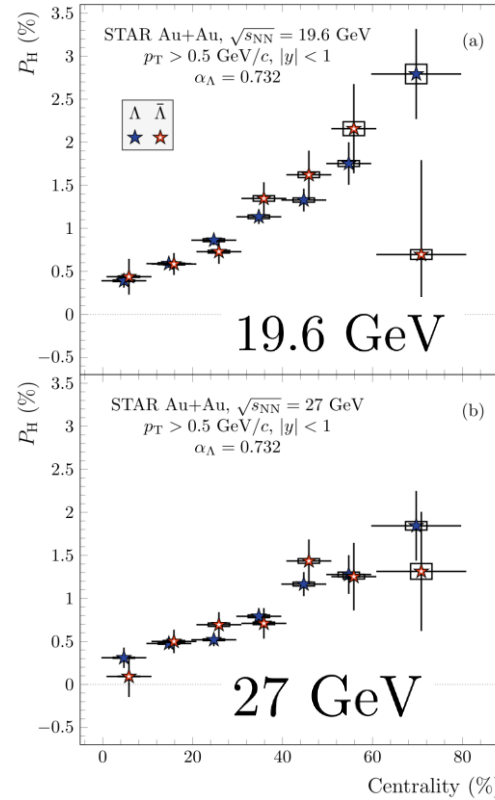
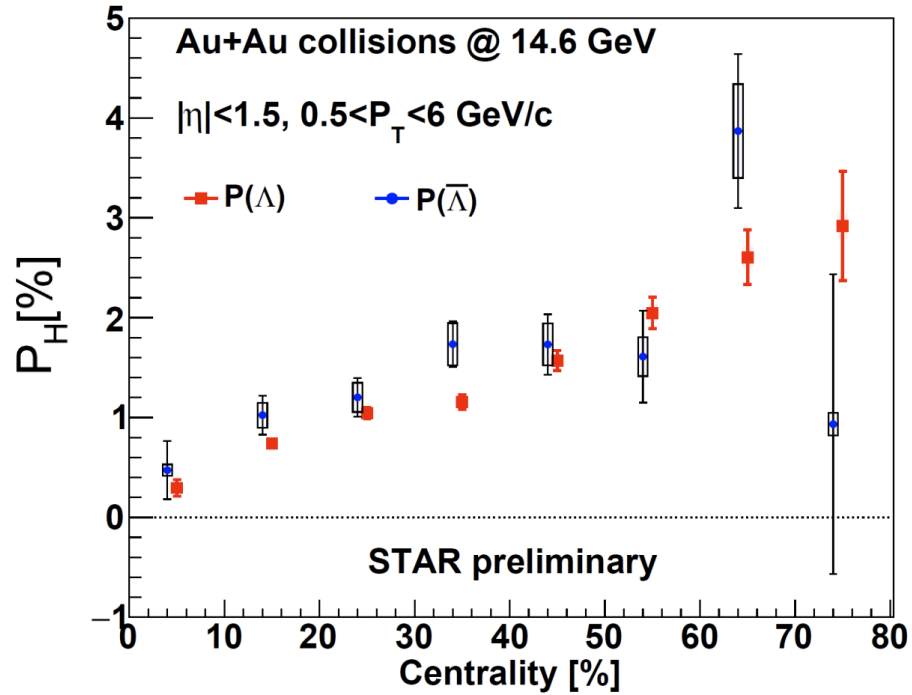
Using  $J = \frac{3}{2}$ :

$$\mathbf{P}_\Lambda = \begin{cases} \mathbf{P}_\Omega & \text{if } \gamma_\Omega = 1, \\ -0.6\mathbf{P}_\Omega & \text{if } \gamma_\Omega = -1 \end{cases}$$

# Result of Global Polarization from BES-II



STAR, Phys. Rev. C 108, 014910 (2023)



- Clear centrality dependence of  $\Lambda$  and  $\bar{\Lambda}$
- Trend consistent with expectation from vorticity