

# How strange: Phase diagrams with 3 critical points

Jacquelyn Noronha-Hostler

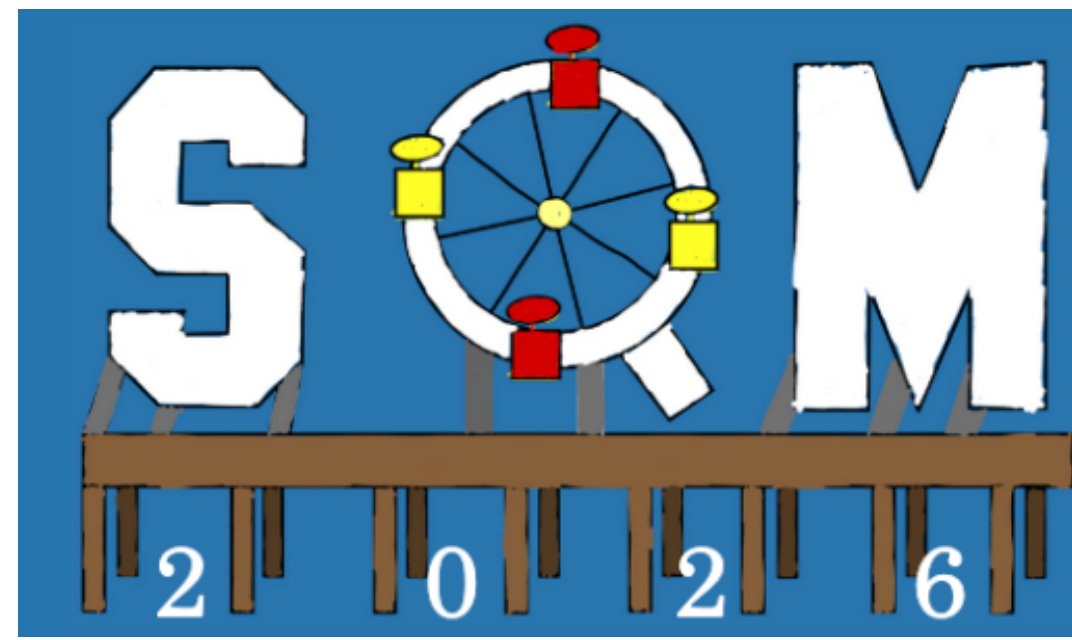
 | Illinois Center for Advanced Studies of the Universe

Based on:

$T = 0$  Cruz Camacho et al,  
*Phys.Rev.D* 111 (2025) 9, 094030;

$T > 0$  Reinke Pelicer, Cruz-Camacho,  
Kumar, Dexheimer, JNH, to appear soon





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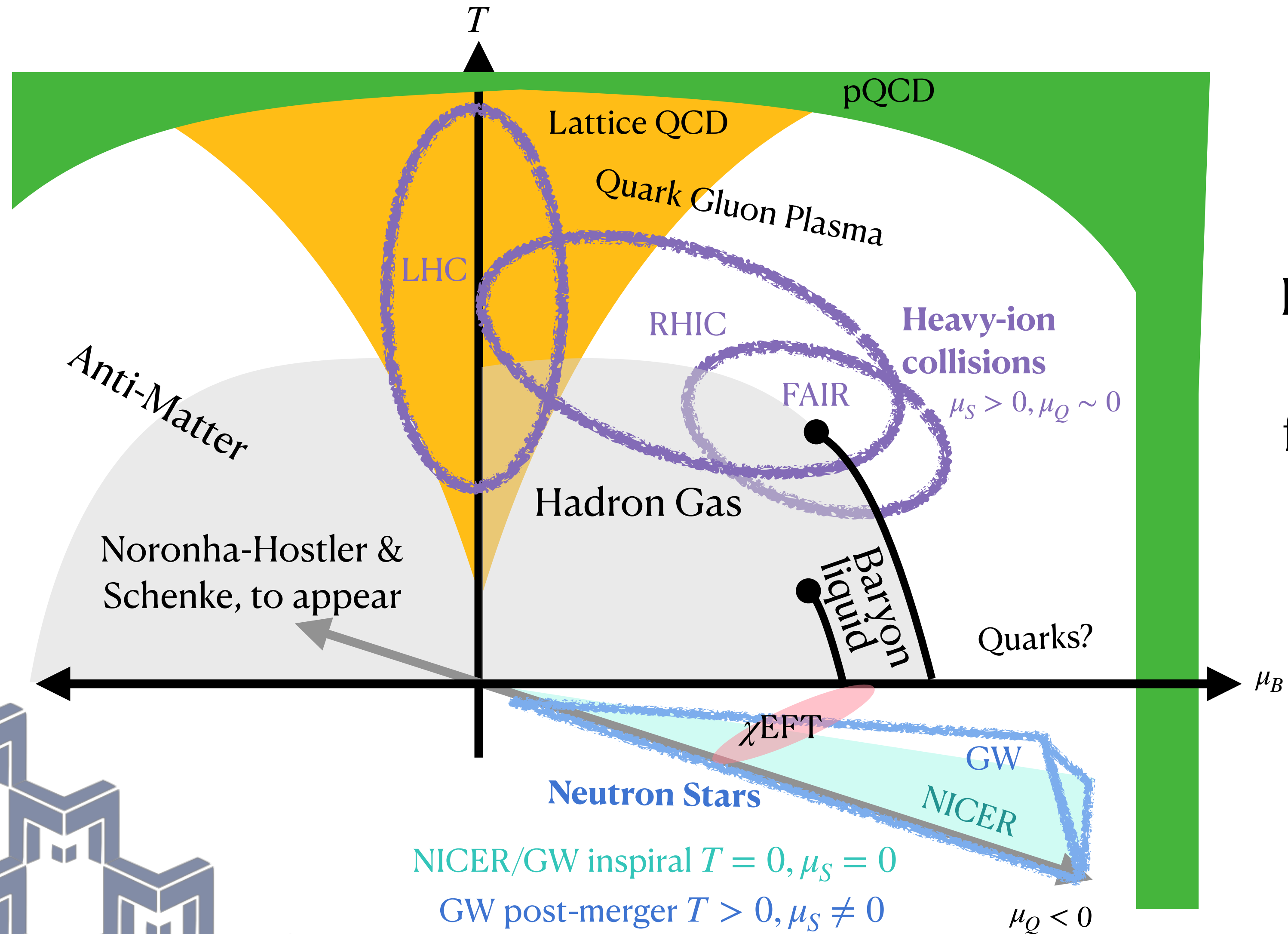
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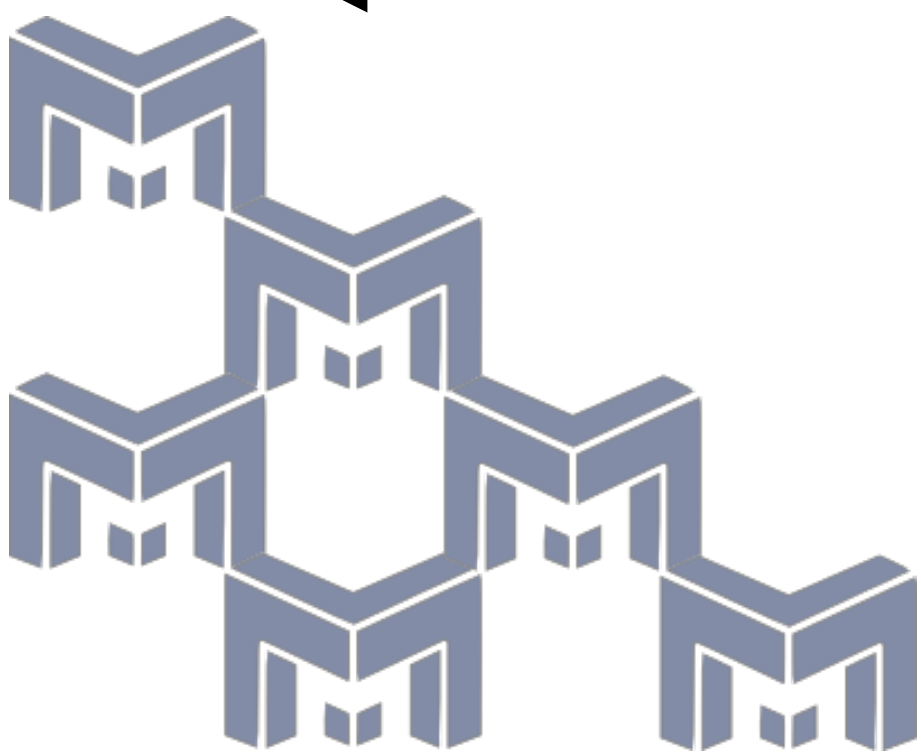
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# Mapping out the QCD phase diagram



How do we describe the equation of state (EOS) out the regime of first principle calculations (lattice QCD, pQCD,  $\chi$ EFT)?

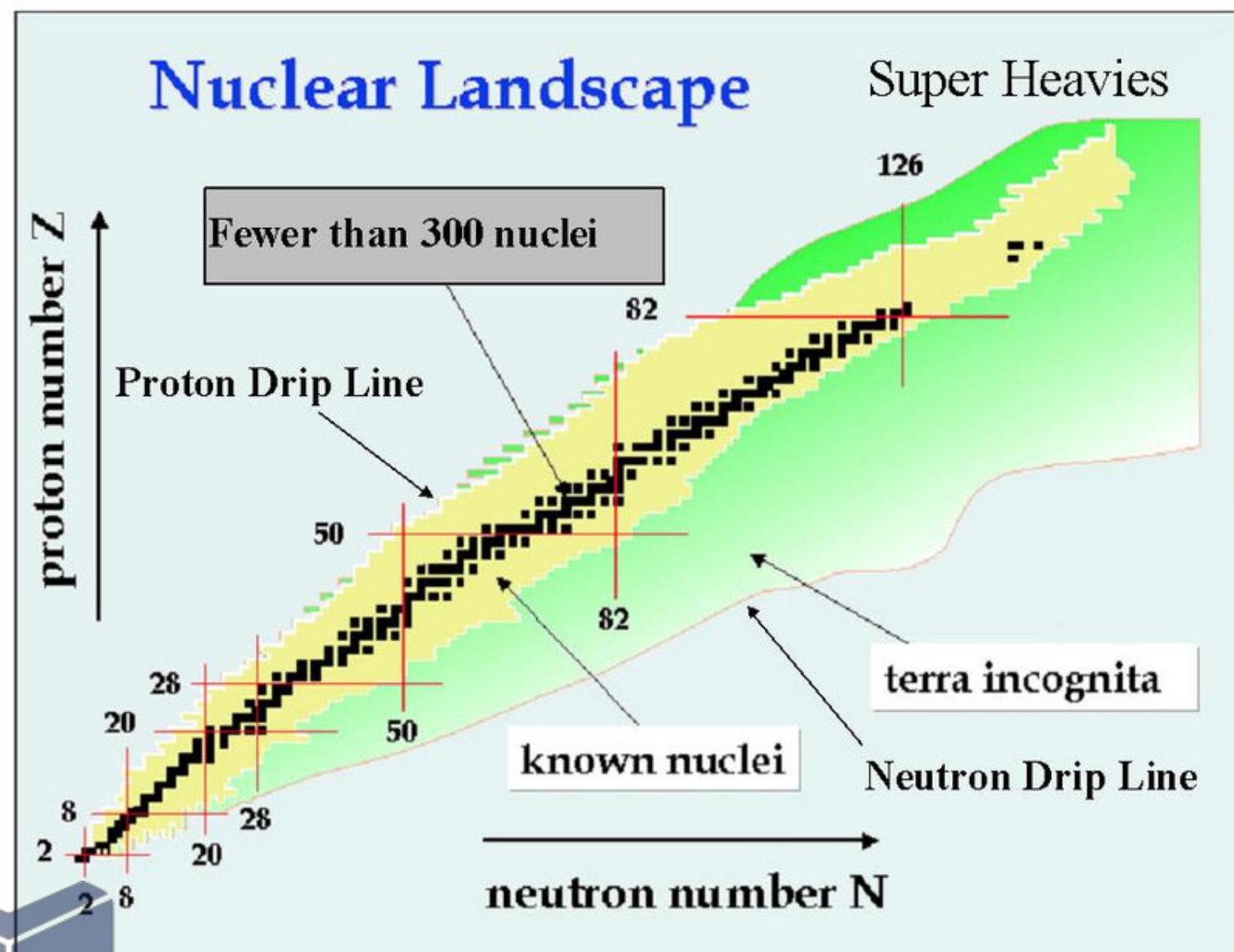


# Instead of waiting on direct QCD calculations

Microscopic models build in symmetries/properties of QCD that are solvable

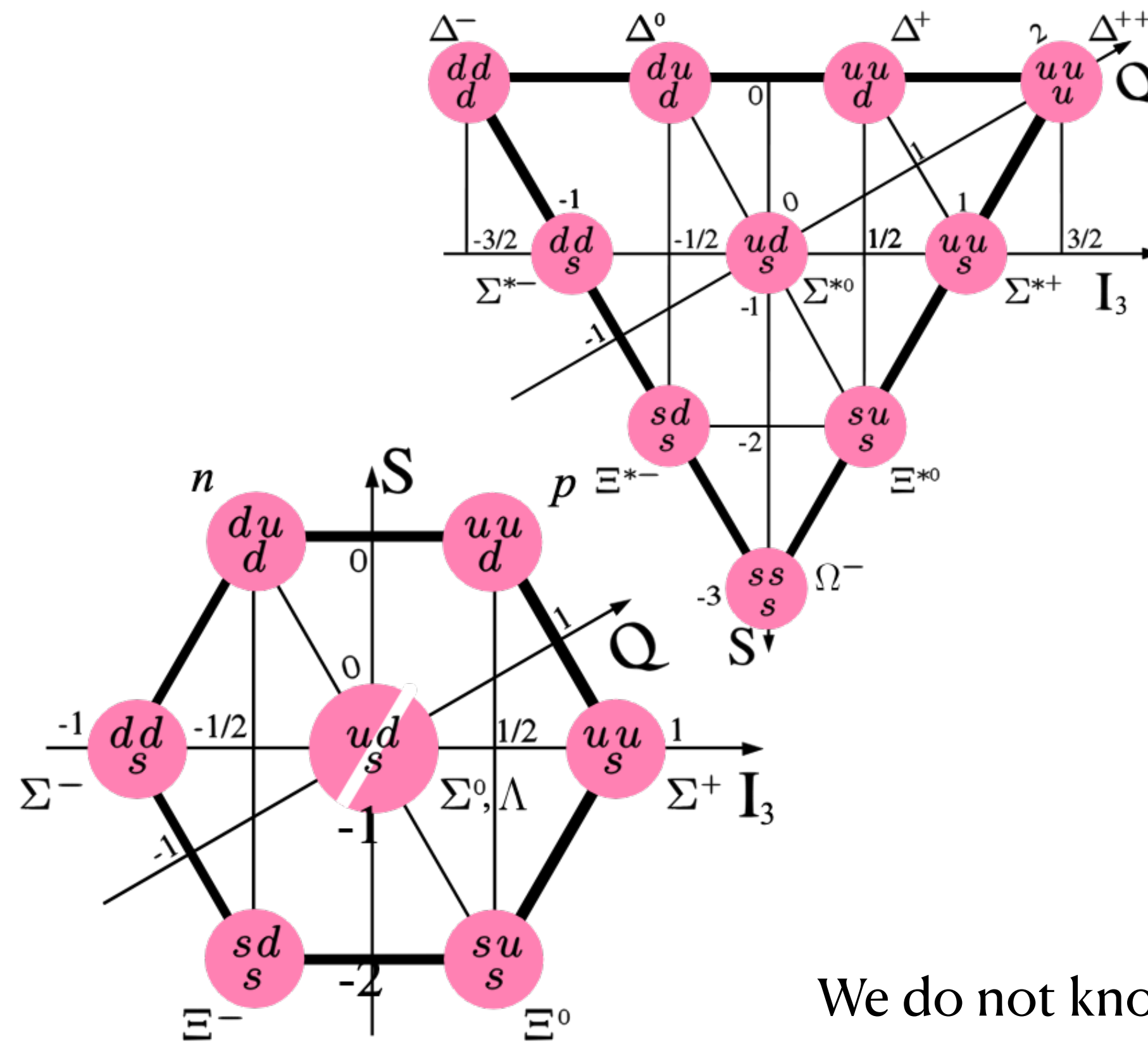
LOW  $n_B$

Nuclei before the liquid-gas phase transition



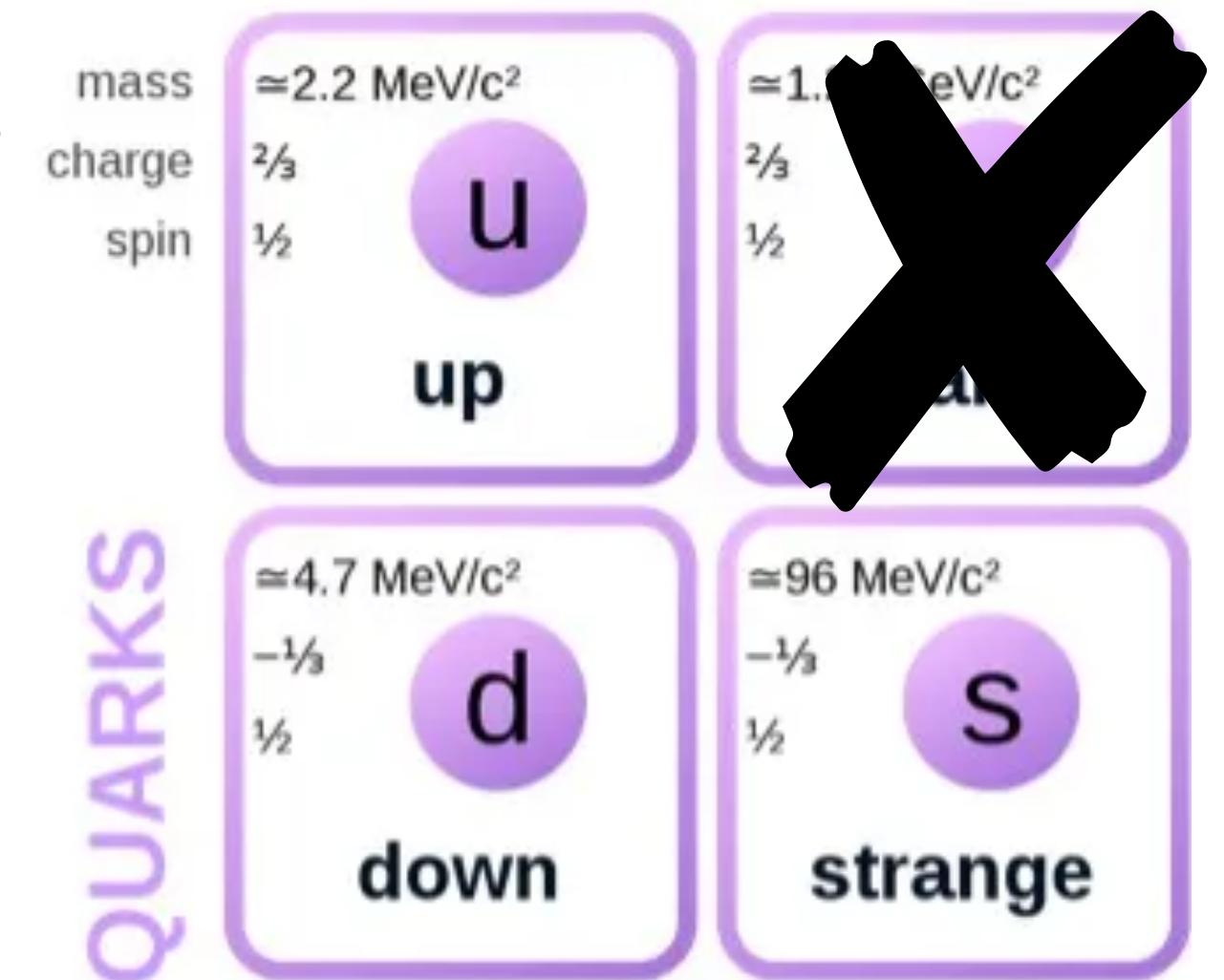
Intermediate  $n_B$

Hadrons (Baryons) before the deconfinement transition

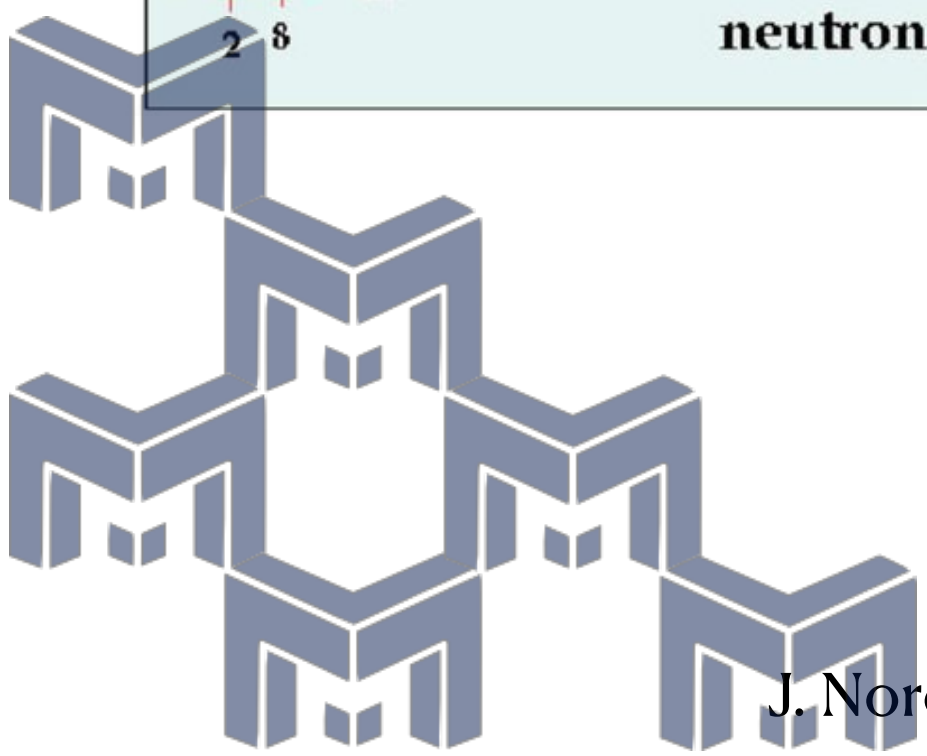


High  $n_B$

Quarks at the highest  $n_B$



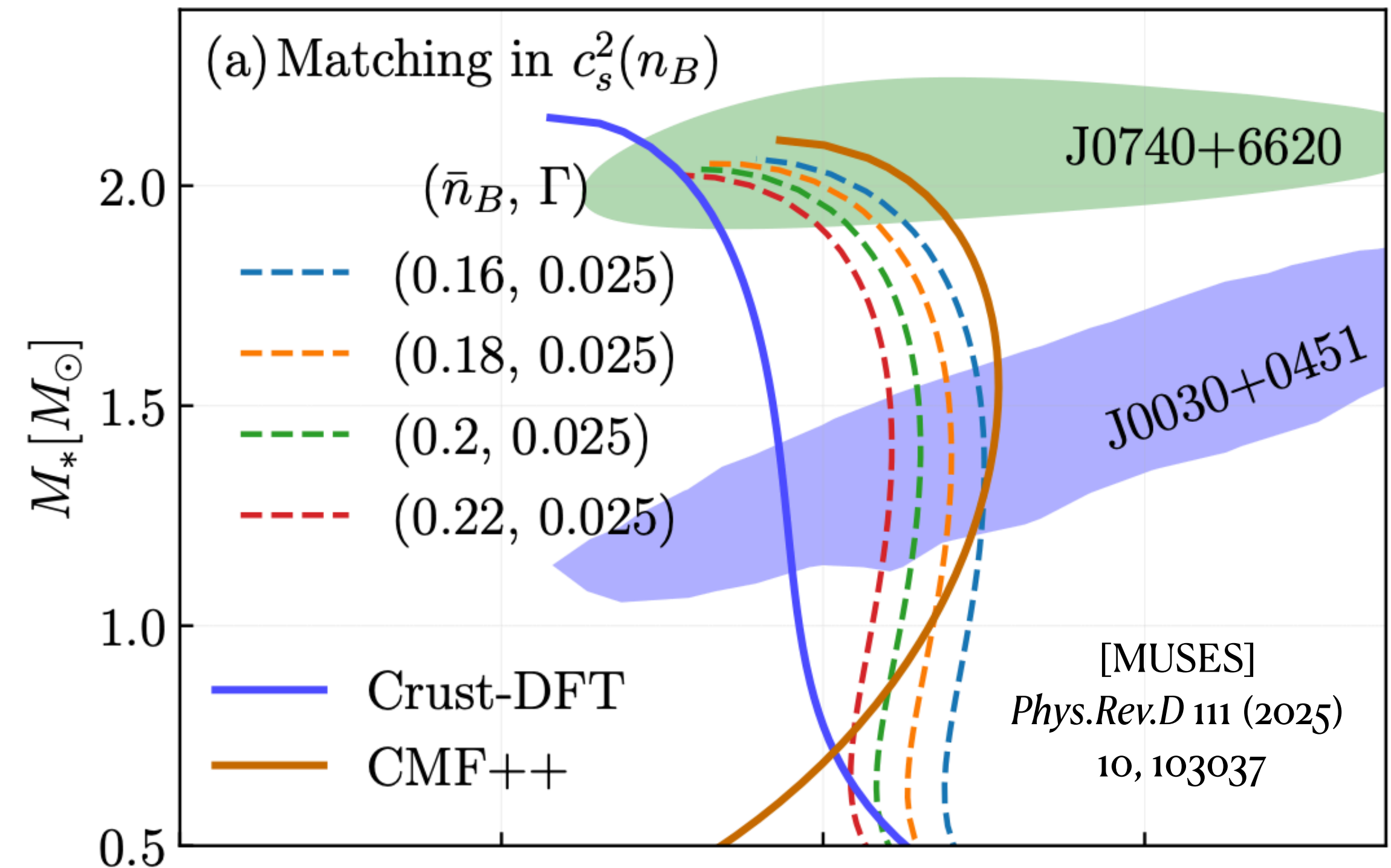
We do not know if quarks will appear before the maximum  $n_B^{max}$  possible for a neutron star, otherwise they might just be inside black holes!



# Chiral Mean Field (CMF) model

Dexheimer & Schramm *Astrophys.J.* 683 (2008) 943-948

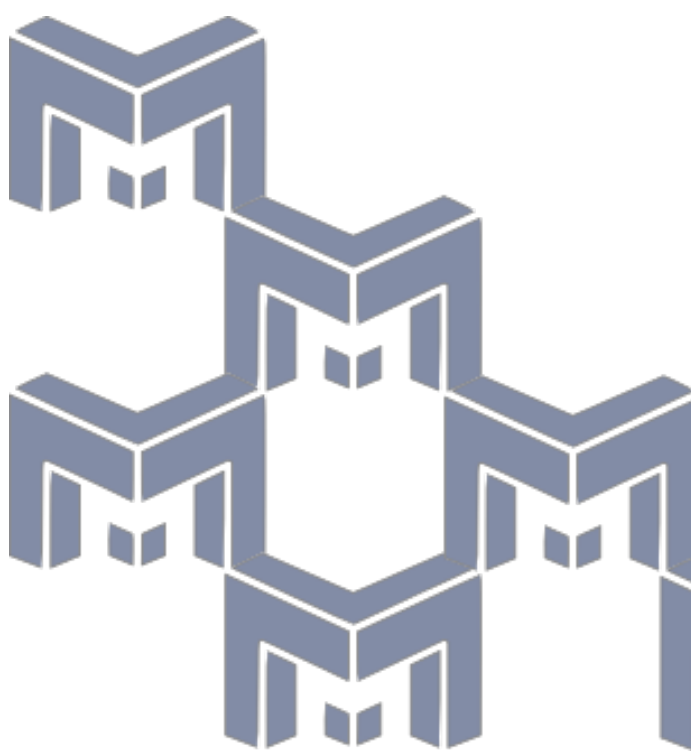
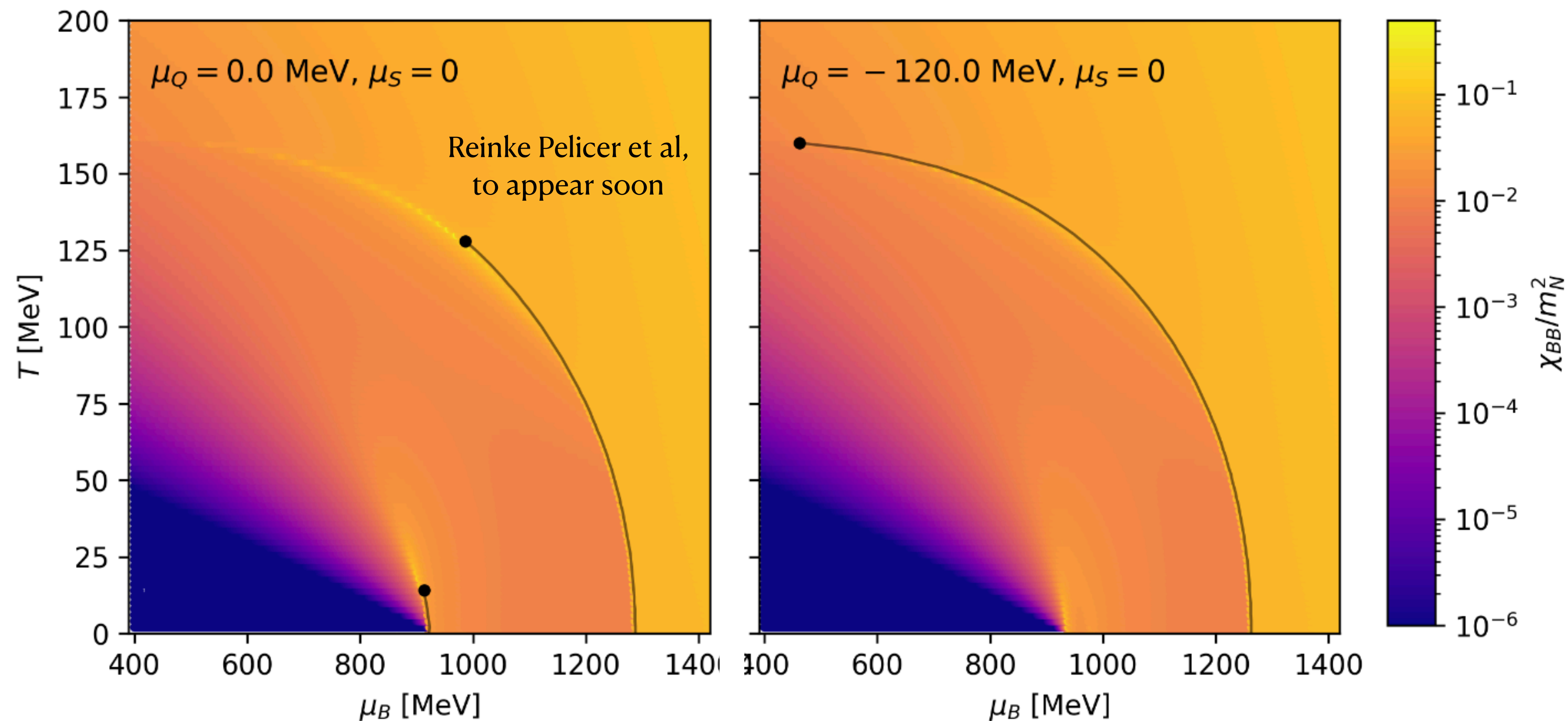
- Includes full SU(3) baryon octet and decuplet, quarks, Polyakov loop to describe deconfinement phase transition
- Mesons mediate interactions (includes self-interactions)
- Calculable in the full  $\{T, \mu_B, \mu_S, \mu_Q\}$  plane
- Can reproduce quenched QCD results  
Kumar et al, *Phys.Rev.D* 109 (2024) 7, 074008
- Reproduces vacuum masses, saturation properties, astrophysical constraints
- 20-50 free parameters (see **Cruz Camacho Tues**)  
Cruz Camacho et al, *2603.16019* [nucl-th]



Open-source, fast, modular c++  
version (CMF++) found here! [https://  
ce.musesframework.io/](https://ce.musesframework.io/)

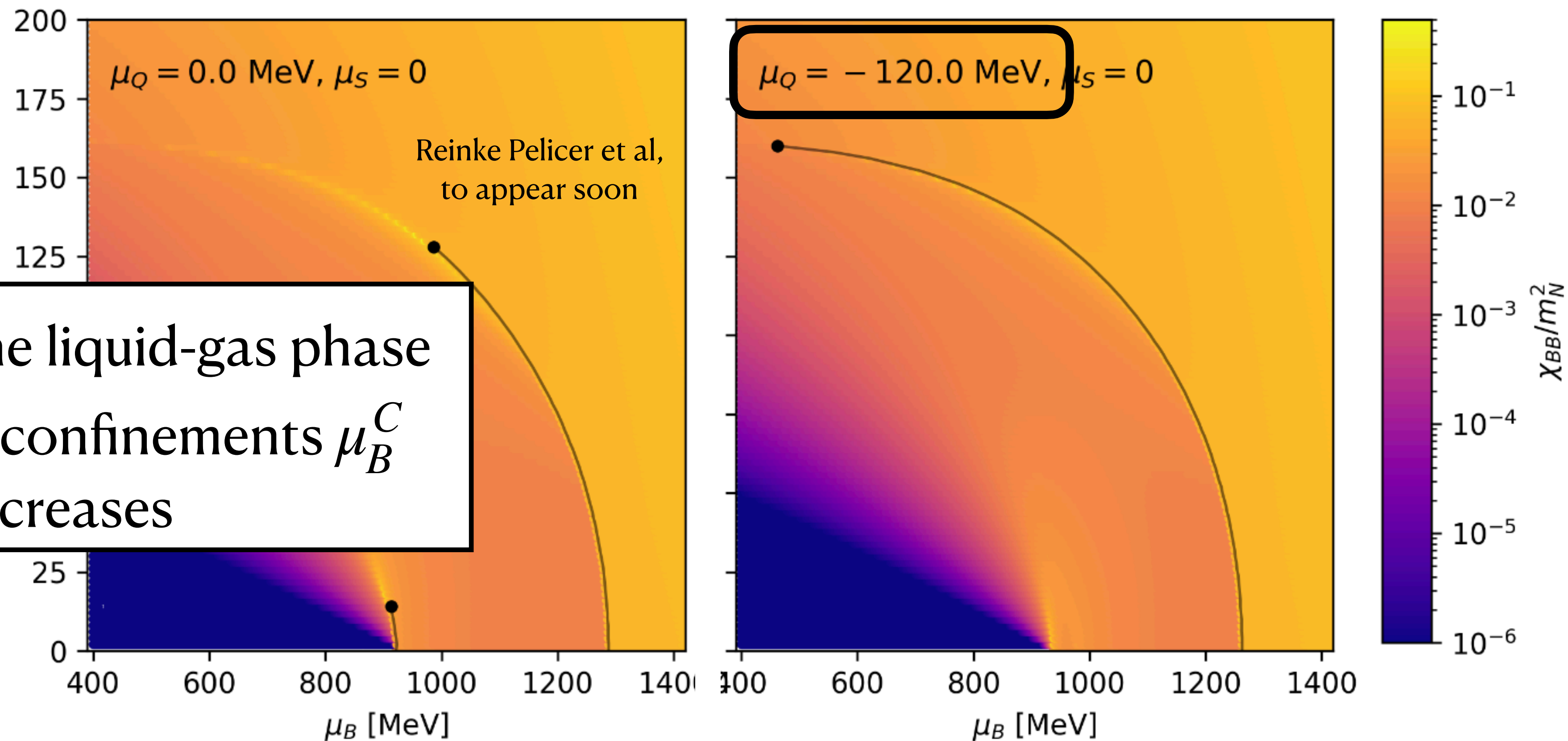
# Caveat: interpretation of these results

- CMF with octet+decuplet+quarks has ~50 parameters+ Ansätze for the couplings
- C4 coupling produces significantly fewer hyperons and has the “standard” QCD phase diagram with just the liquid gas and deconfinement phase transitions.



# Caveat: interpretation of these results

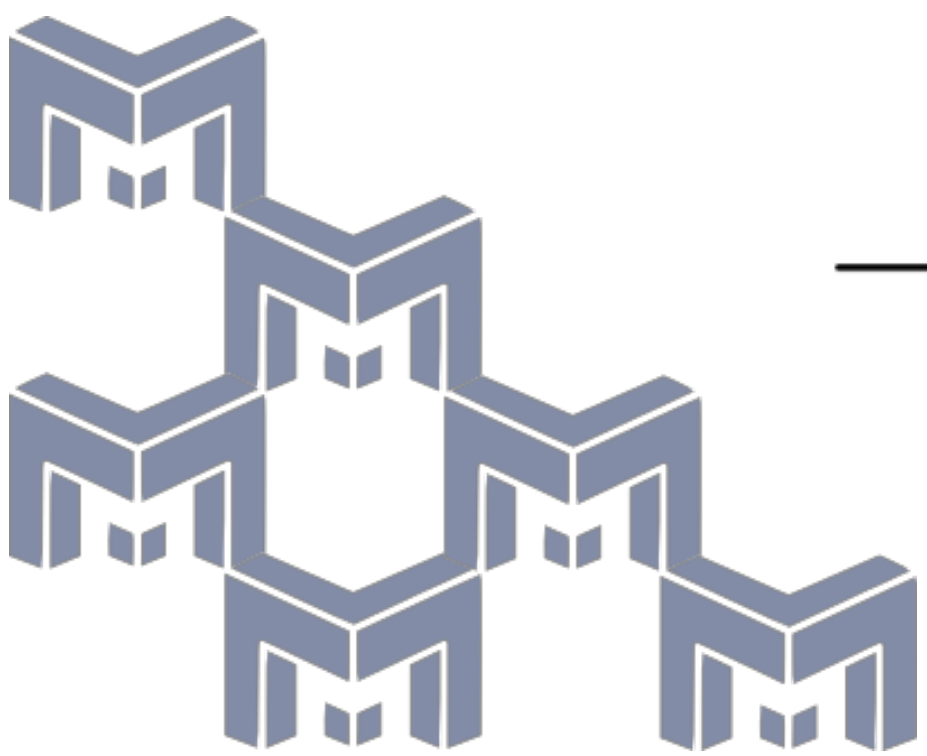
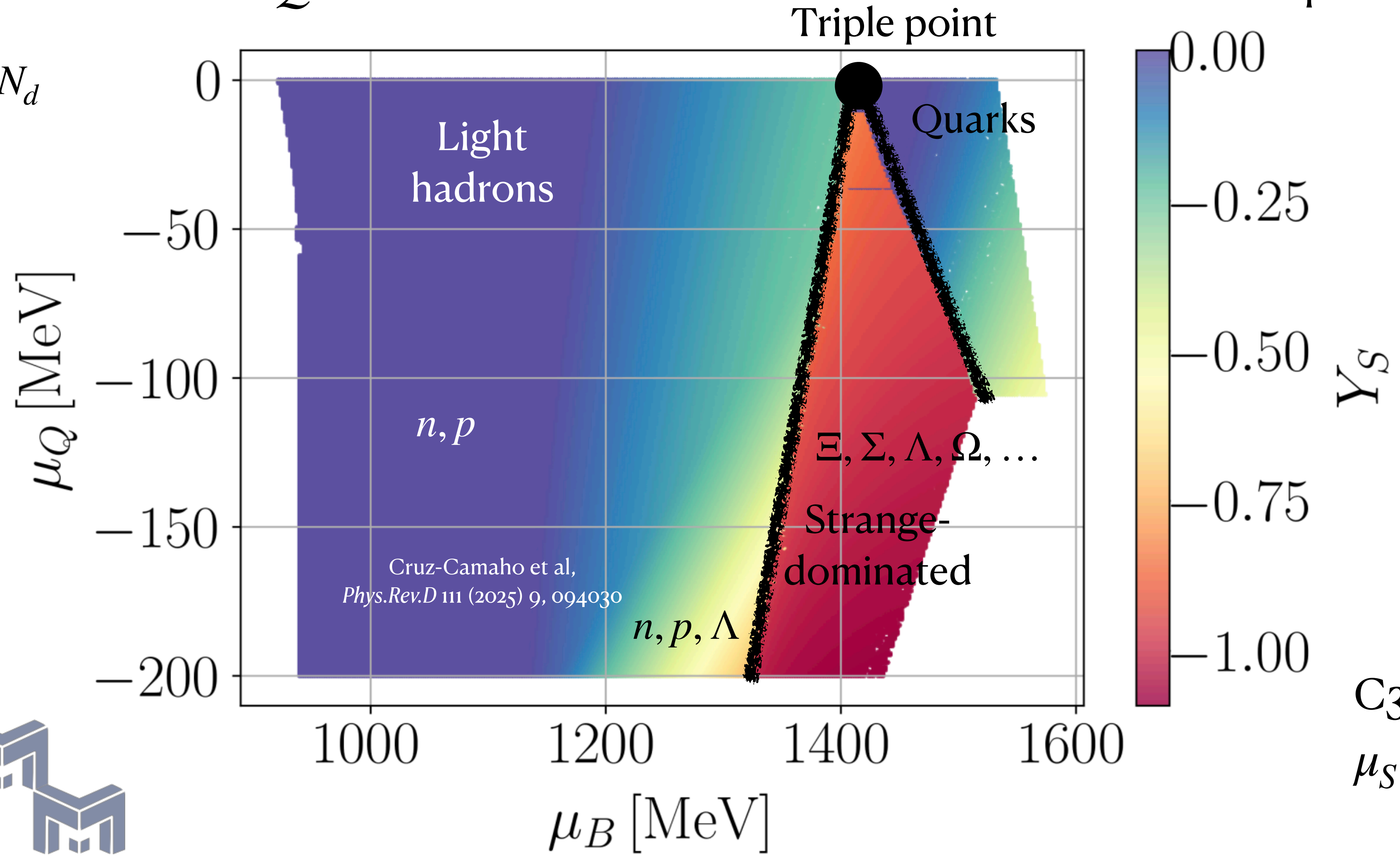
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# At finite $\mu_Q < 0$ , similar to heavy-ions!

Example  $T = 0$   
phase diagram

$$N_u = N_d$$

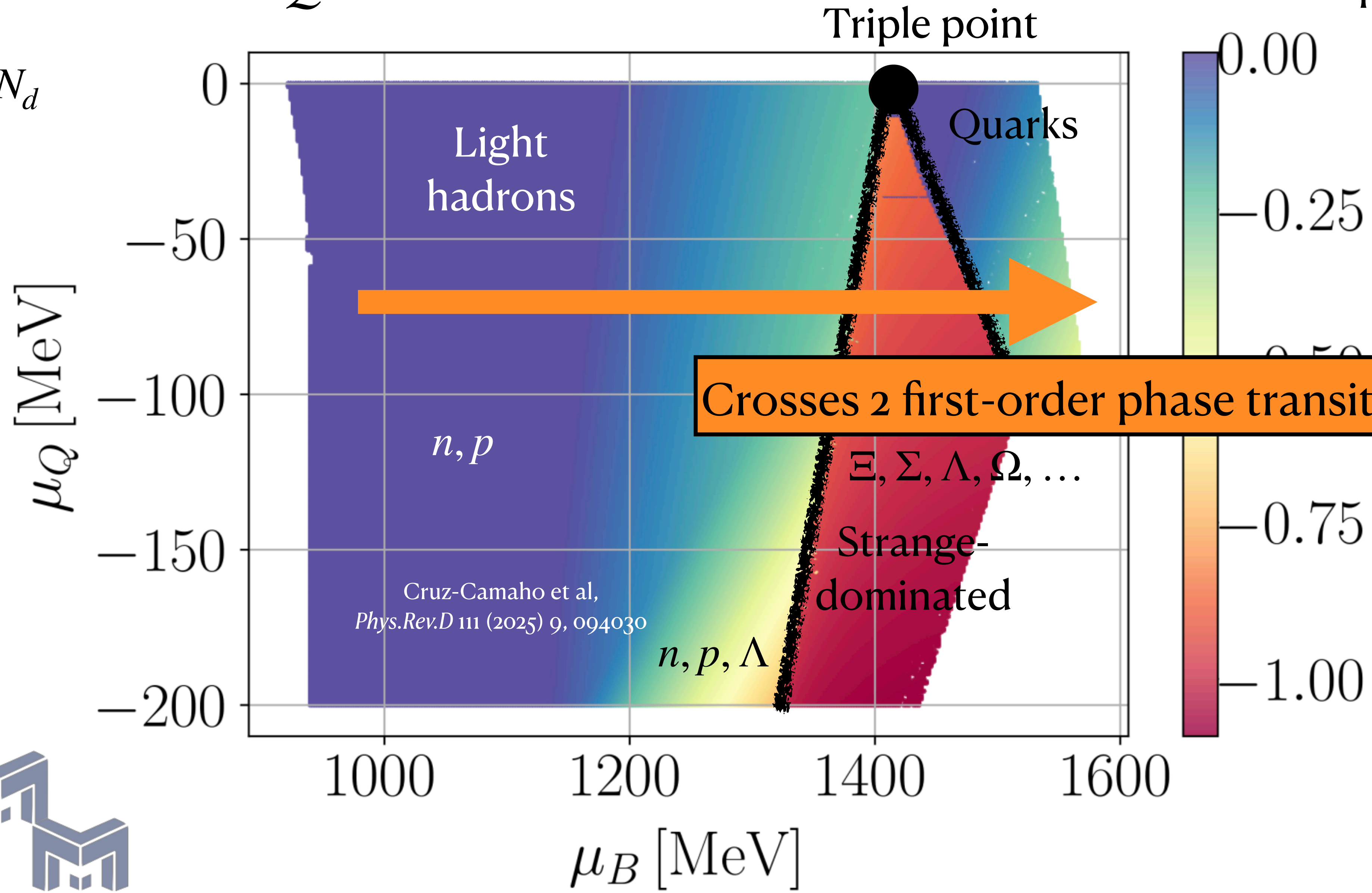


$C_3$  coupling  
 $\mu_S = 0$

# At finite $\mu_Q < 0$ , similar to heavy-ions!

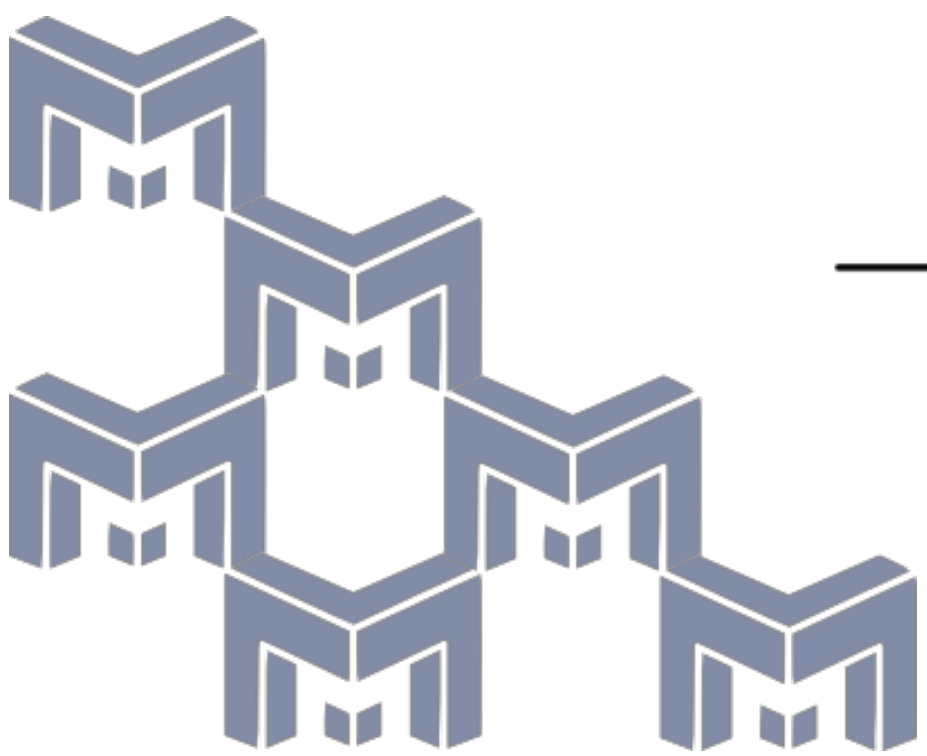
Example  $T = 0$   
phase diagram

$$N_u = N_d$$

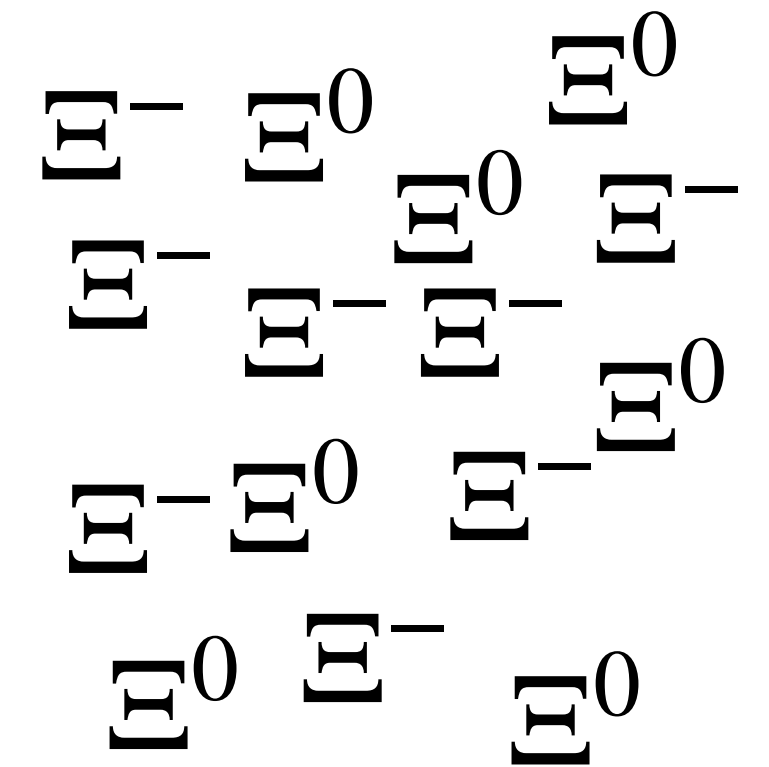
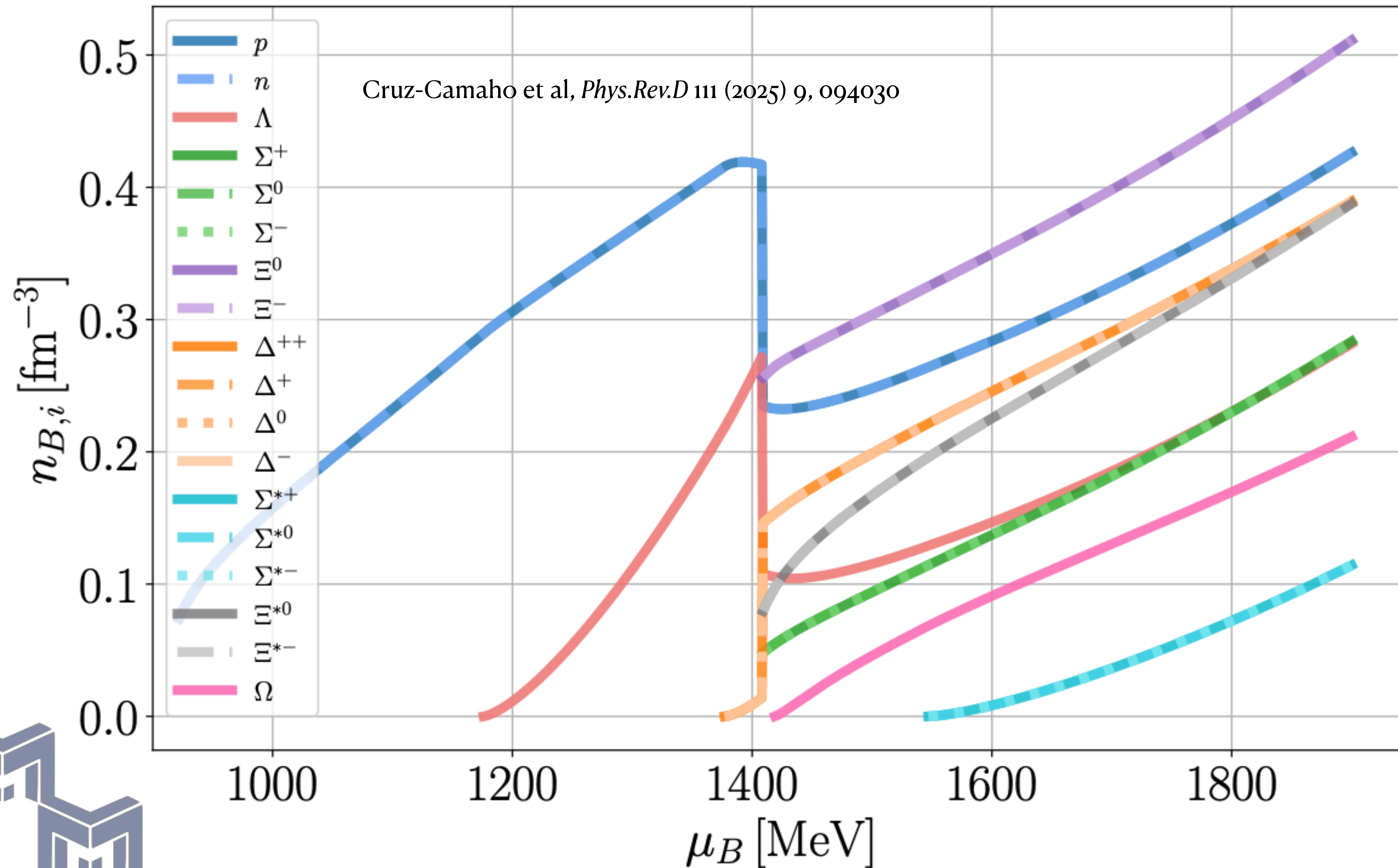


Crosses 2 first-order phase transitions!

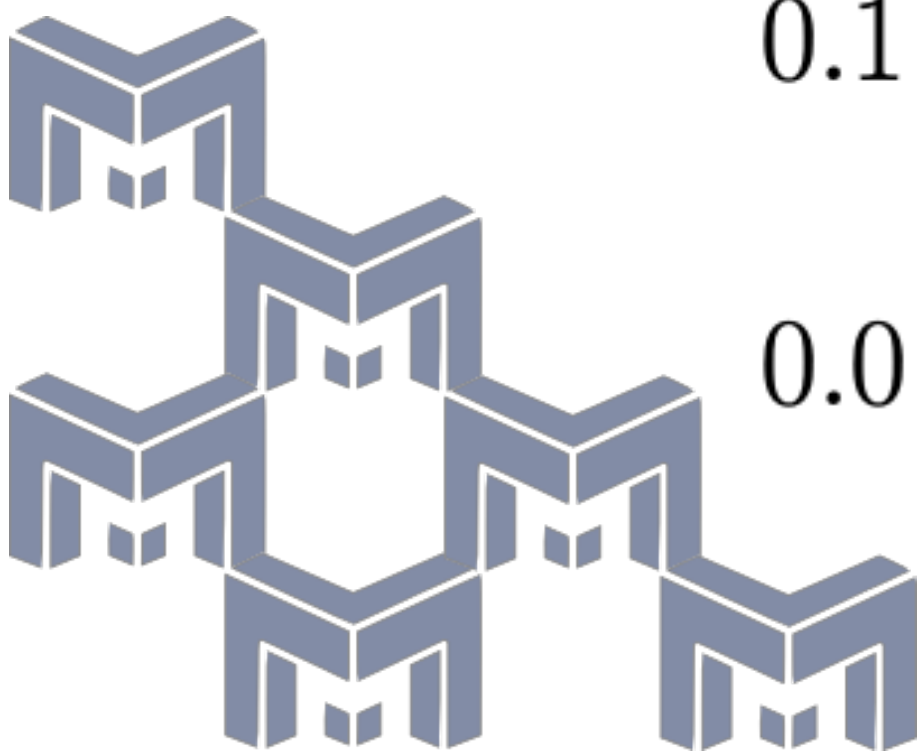
$C_3$  coupling  
 $\mu_S = 0$



# What's happens inside the strangeness dominated phase?

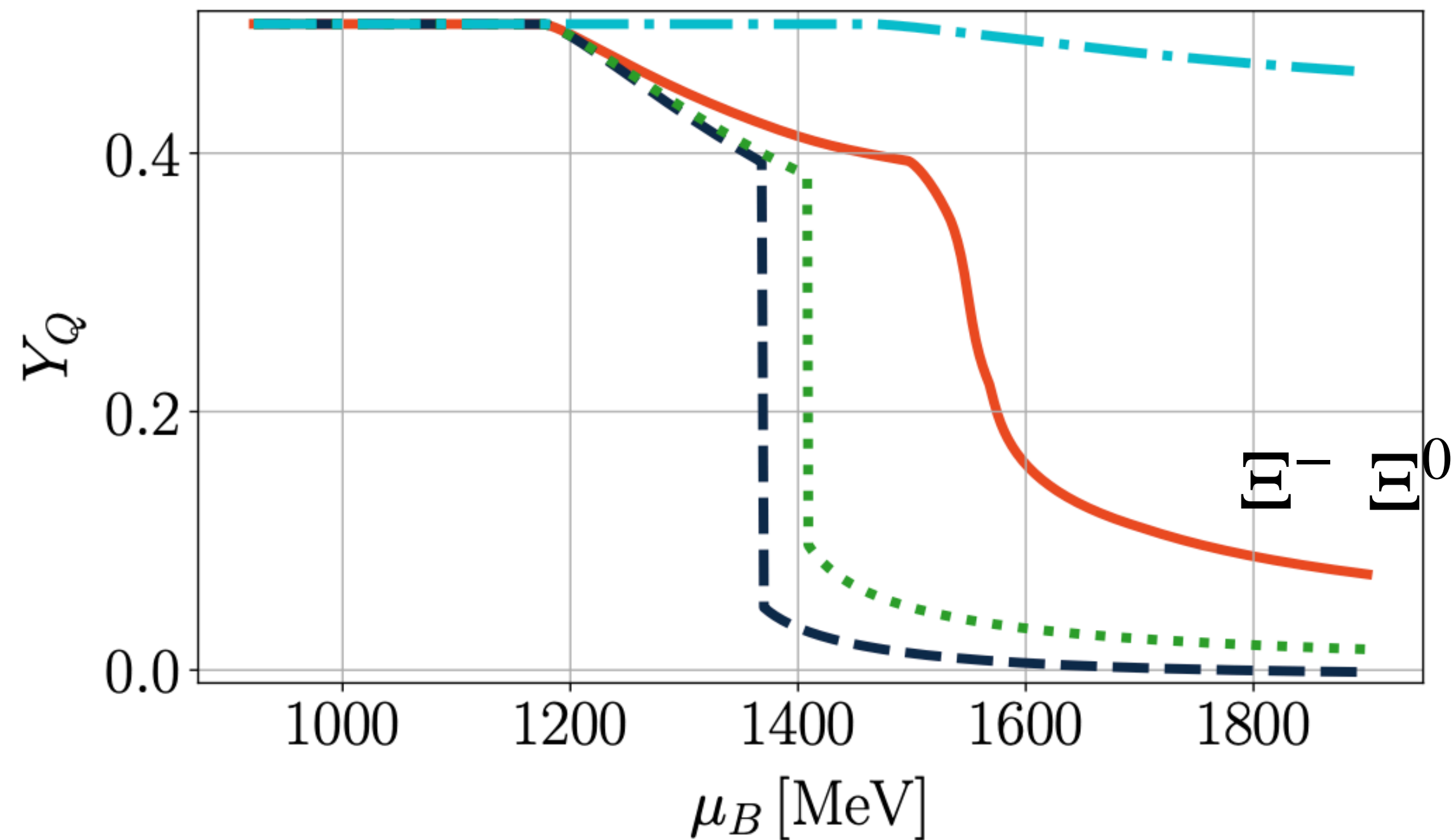
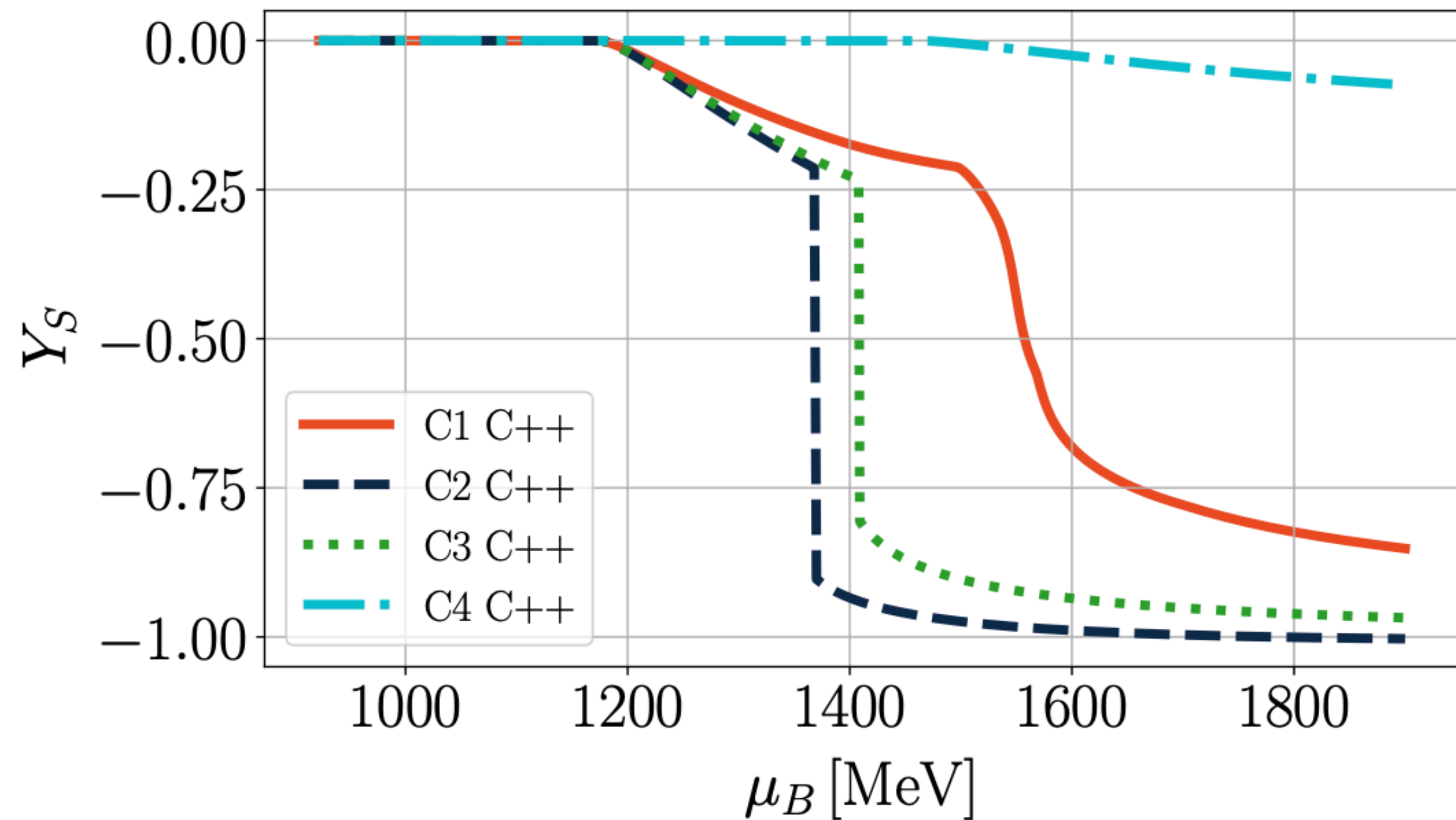


$$S = -2, Q = 0, -1$$



# What's the order parameter?

Isospin restoration, but not where you'd expect it!

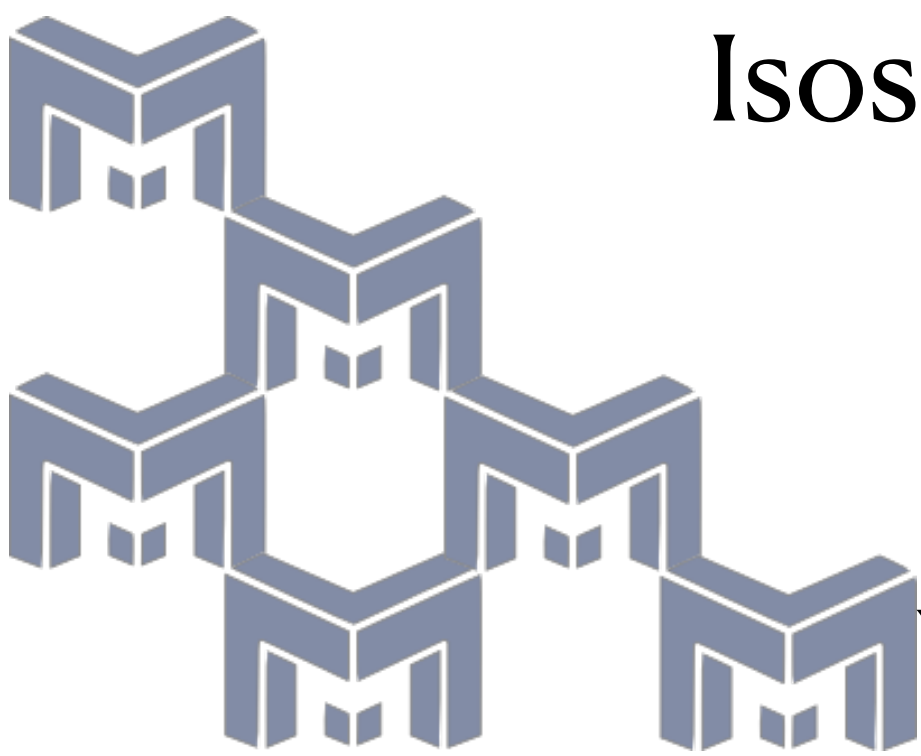


Isospin asymmetry:  $I_3 = Q - \frac{1}{2} (B + S)$  Gell-Mann–Nishijima formula

Divide by B,  $Y_{I_3} = Y_Q - \frac{1}{2} (1 + Y_S)$ .

See [Yumu Yang Weds](#)

When  $Y_S \rightarrow -2$ ,  $Y_Q \rightarrow -1/2$ , then  $Y_{I_3} = 0$  (isospin restored!)



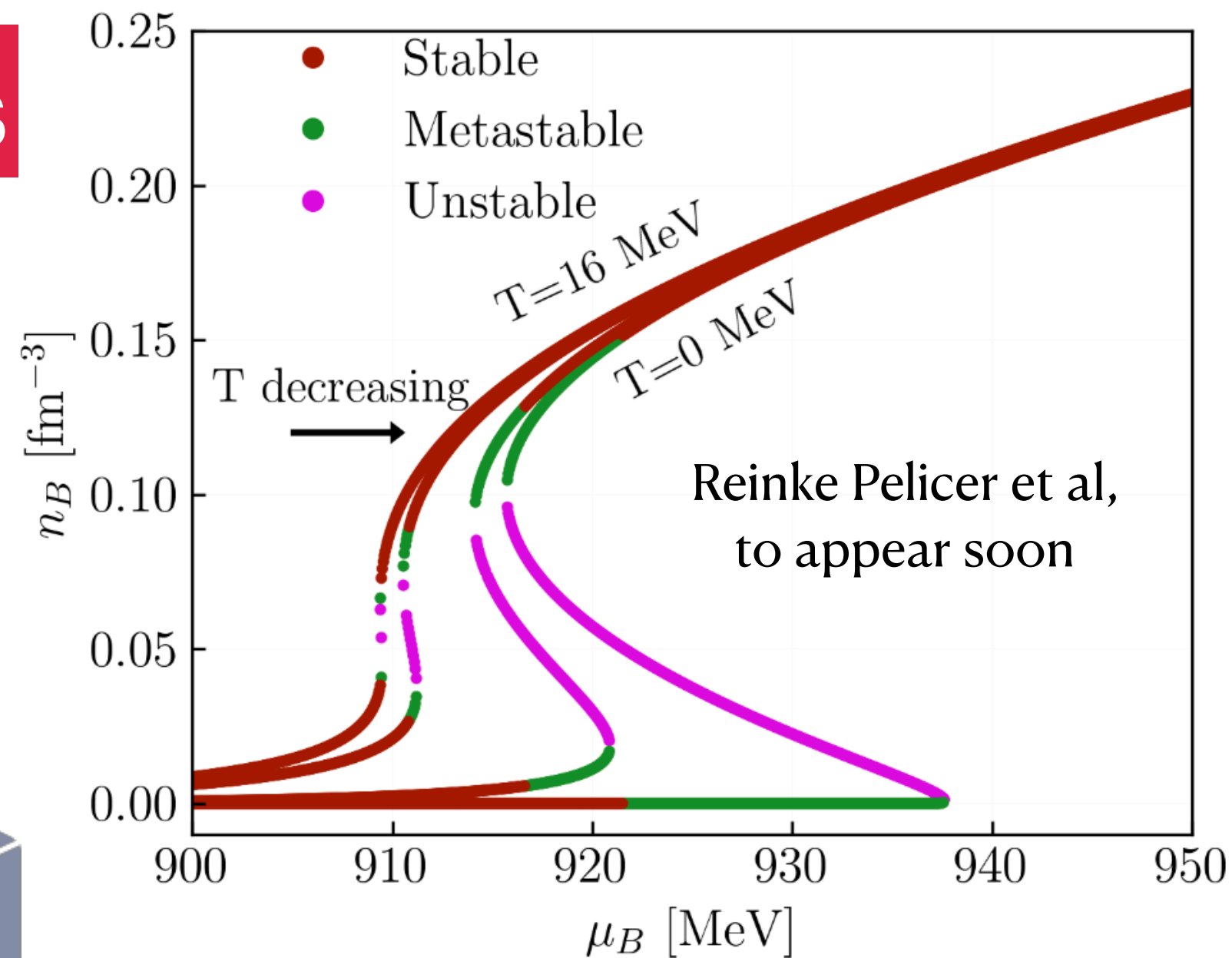
# What about finite T?

- New challenge: determining stable, metastable, unstable phases. Must be a positive matrix:

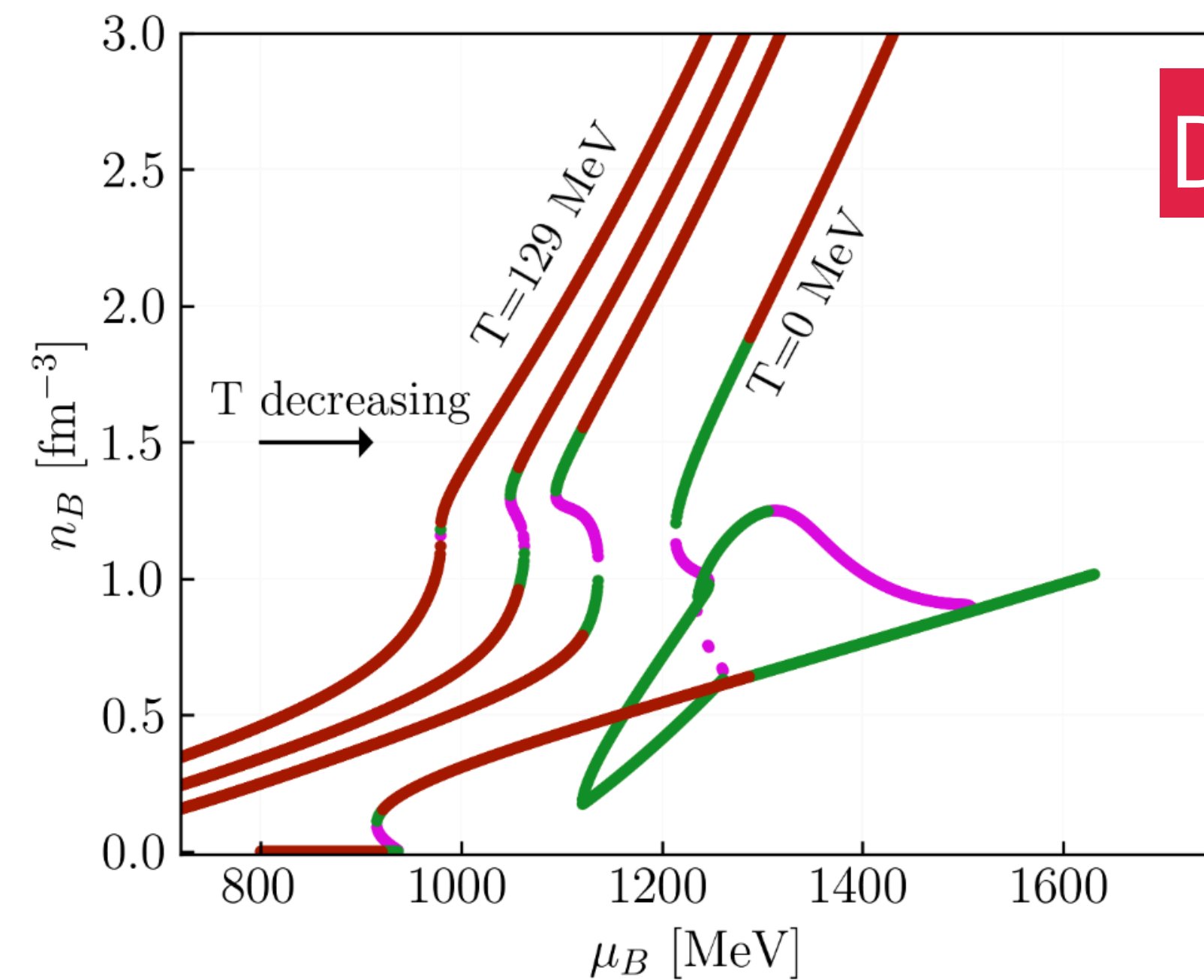
$$M = \begin{bmatrix} \left. \frac{\partial s}{\partial T} \right|_{\vec{\mu}} & \left. \frac{\partial s}{\partial \mu_B} \right|_{T, \mu_S, \mu_Q} & \left. \frac{\partial s}{\partial \mu_S} \right|_{T, \mu_B, \mu_Q} & \left. \frac{\partial s}{\partial \mu_Q} \right|_{T, \mu_B, \mu_S} \\ \left. \frac{\partial n_B}{\partial T} \right|_{\vec{\mu}} & \chi_2^B & \chi_{11}^{BS} & \chi_{11}^{BQ} \\ \left. \frac{\partial n_S}{\partial T} \right|_{\vec{\mu}} & \chi_{11}^{SB} & \chi_2^S & \chi_{11}^{SQ} \\ \left. \frac{\partial n_Q}{\partial T} \right|_{\vec{\mu}} & \chi_{11}^{QB} & \chi_{11}^{QS} & \chi_2^Q \end{bmatrix}$$

- Calculate susceptibilities in 4D phase space

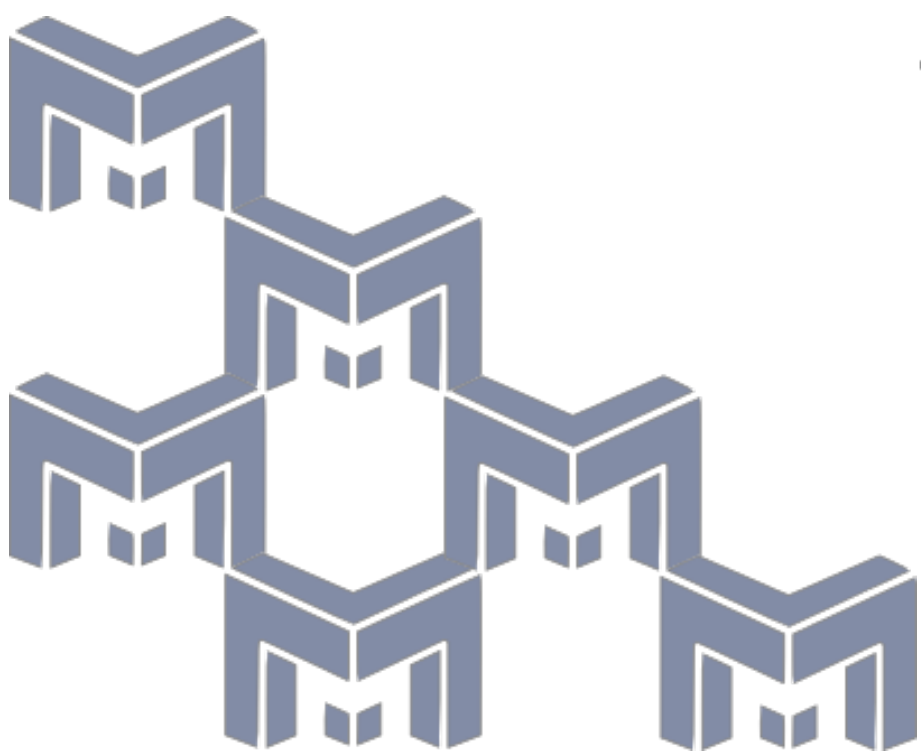
Liquid-gas



Deconfinement



C4 coupling



# What about finite T?

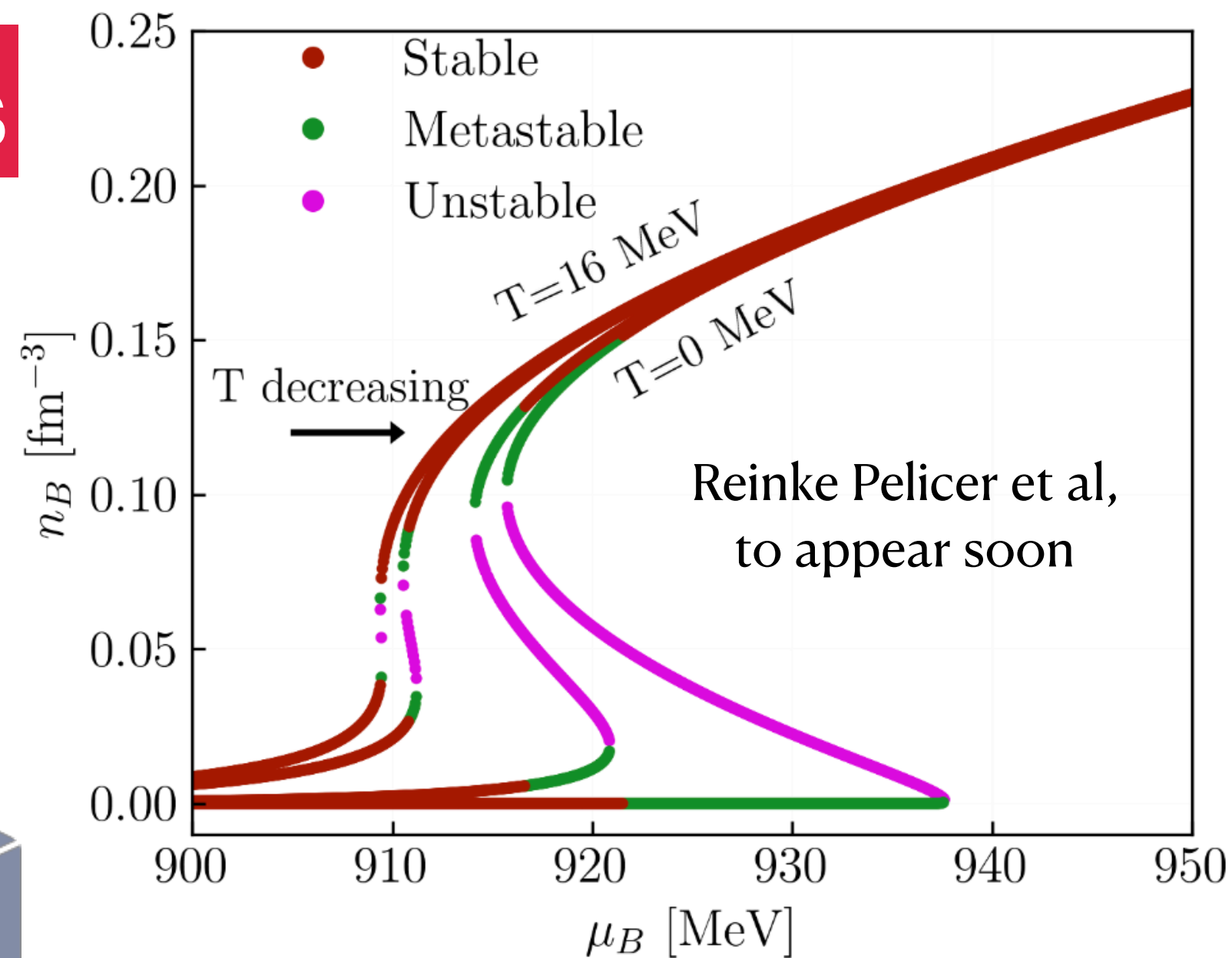
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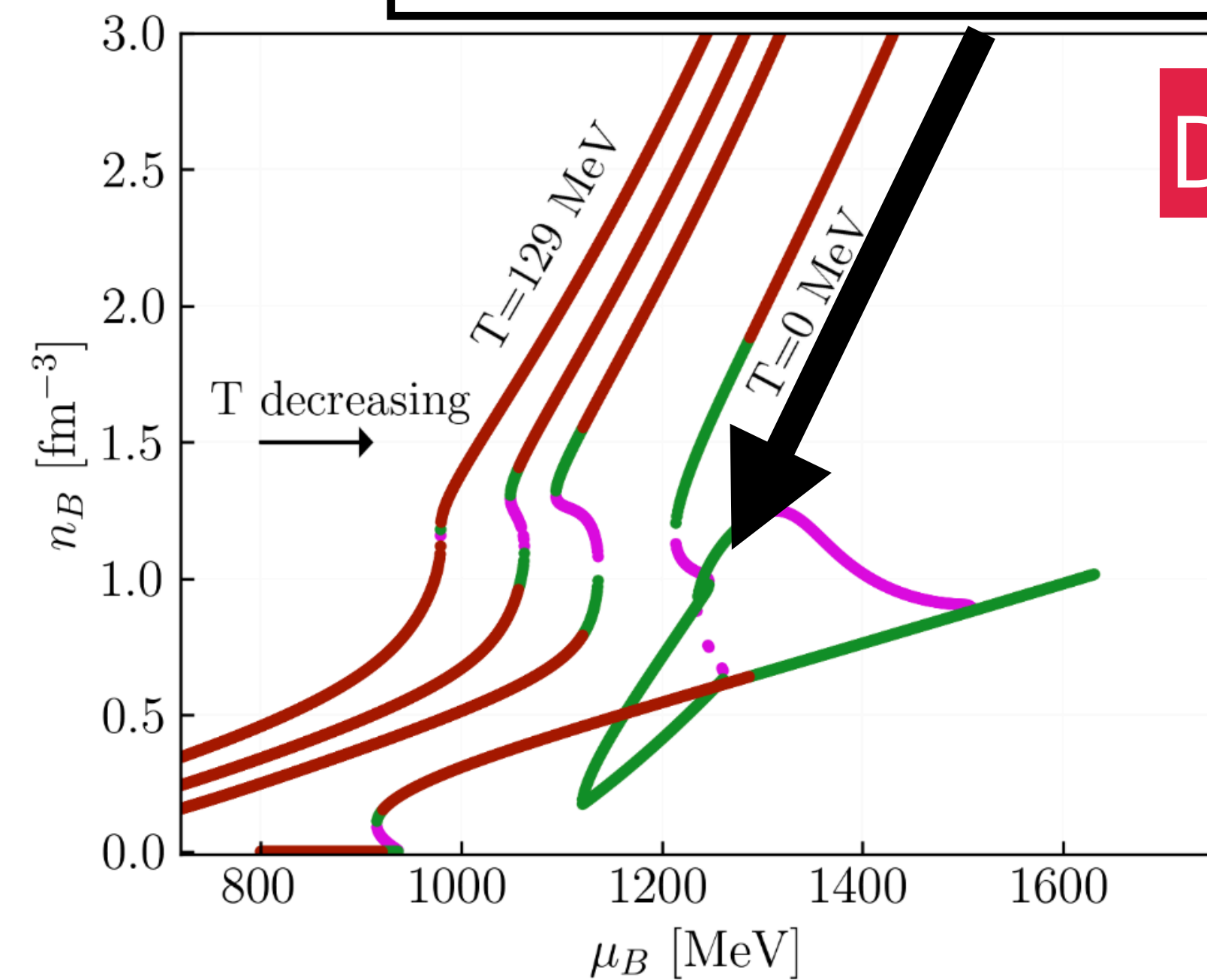
Strangeness dominated metastable phase in between quarks and hadrons

- Calculate susceptibilities in 4D phase space

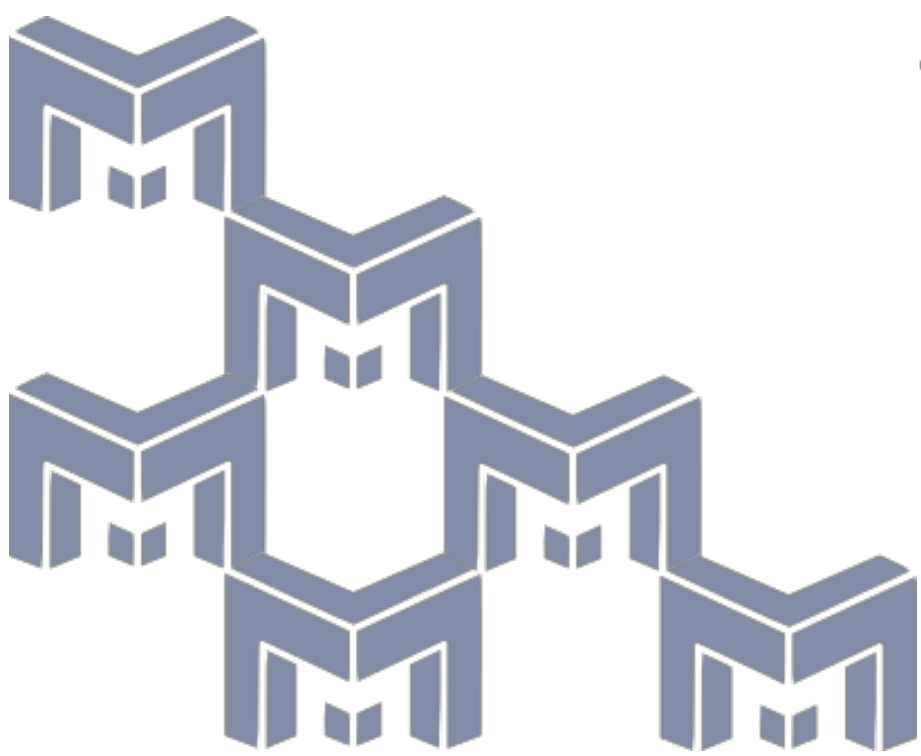
Liquid-gas



Deconfinement



C4 coupling

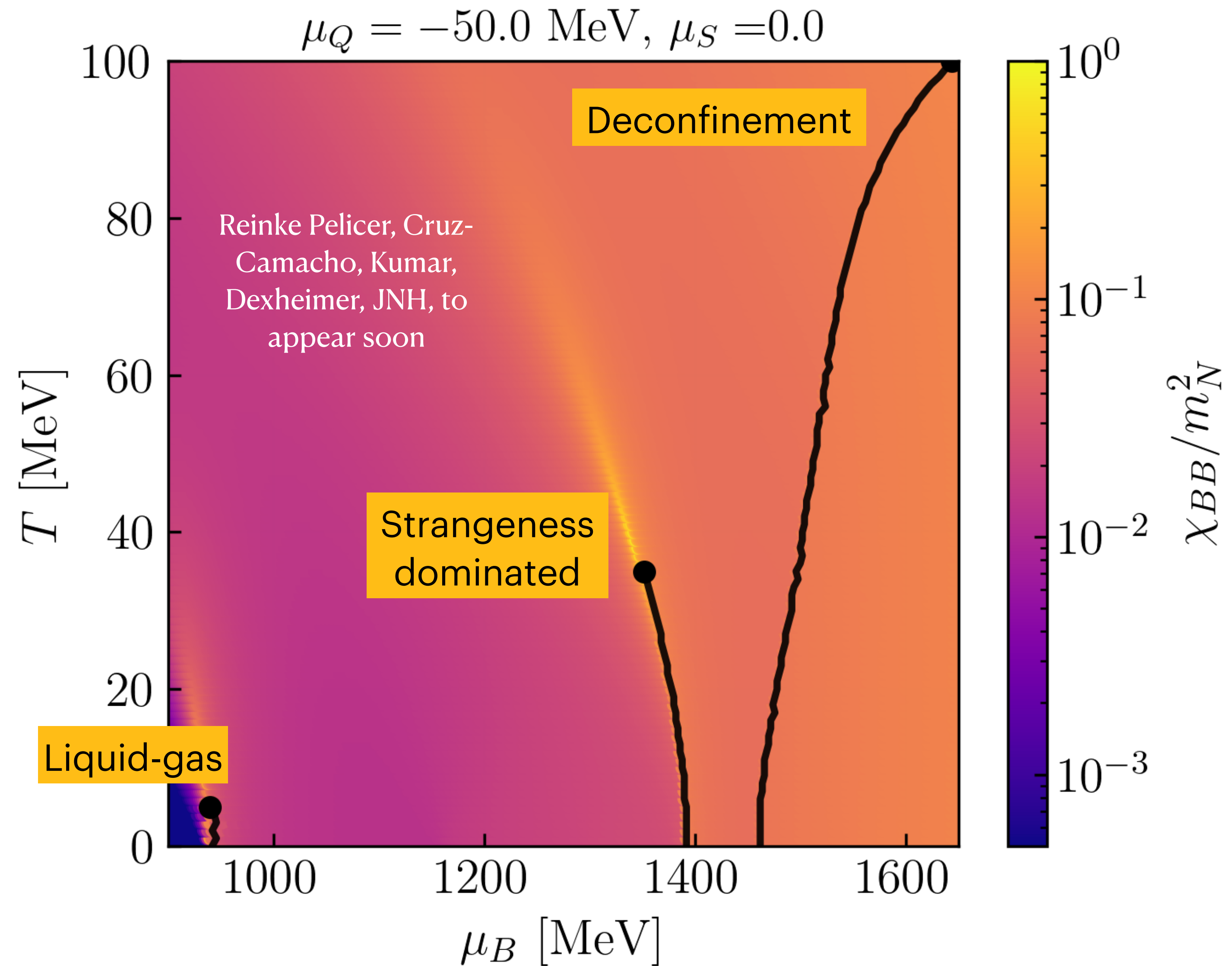


# Turn up the heat!

We have 3 critical points!

We could see this for isospin asymmetric systems at very low  $\sqrt{s_{NN}}$  or binary neutron star mergers

Note this is mean-field theory, so we need to think about the universality class some more.



# Can you give me a better intuition for $\mu_S$ ?

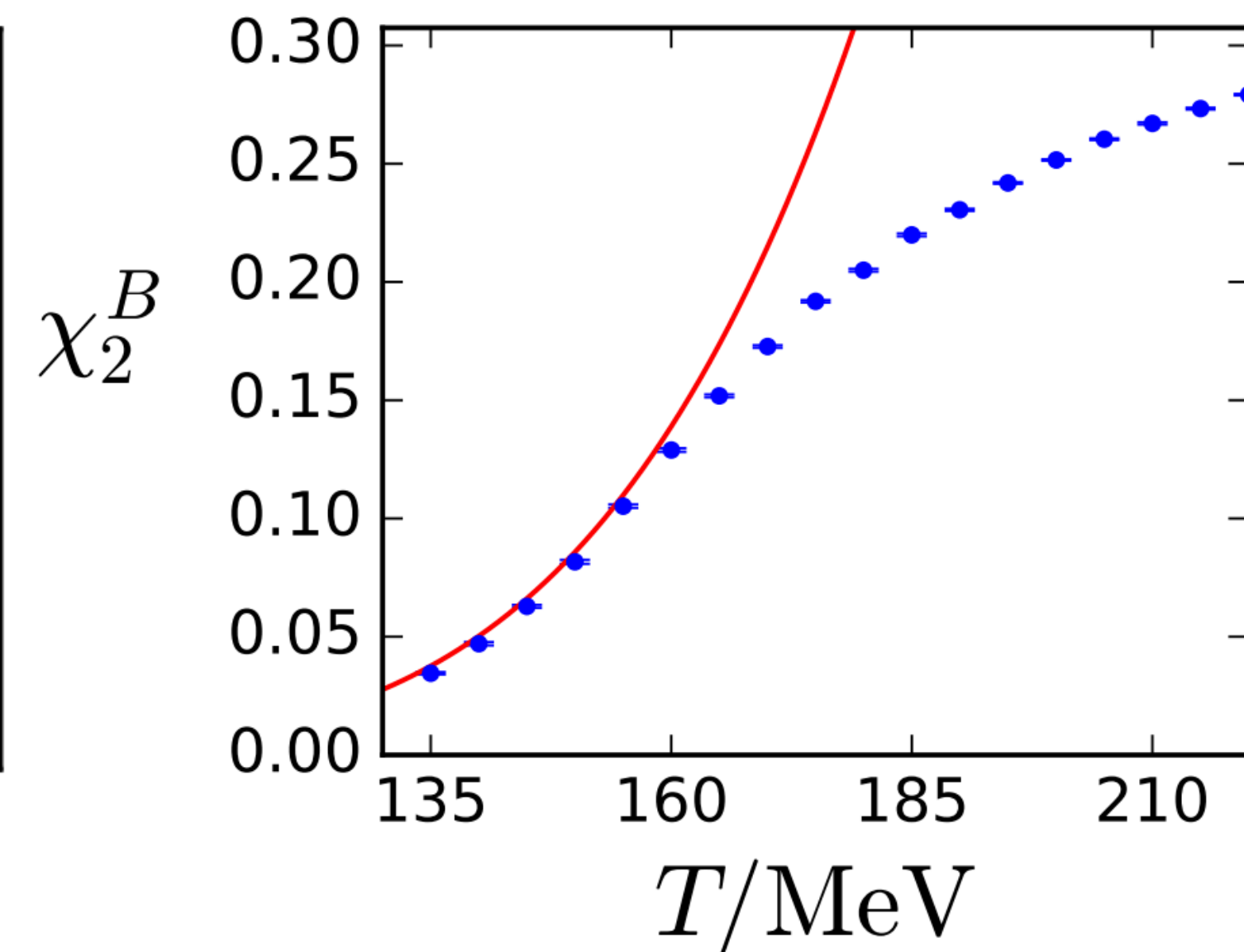
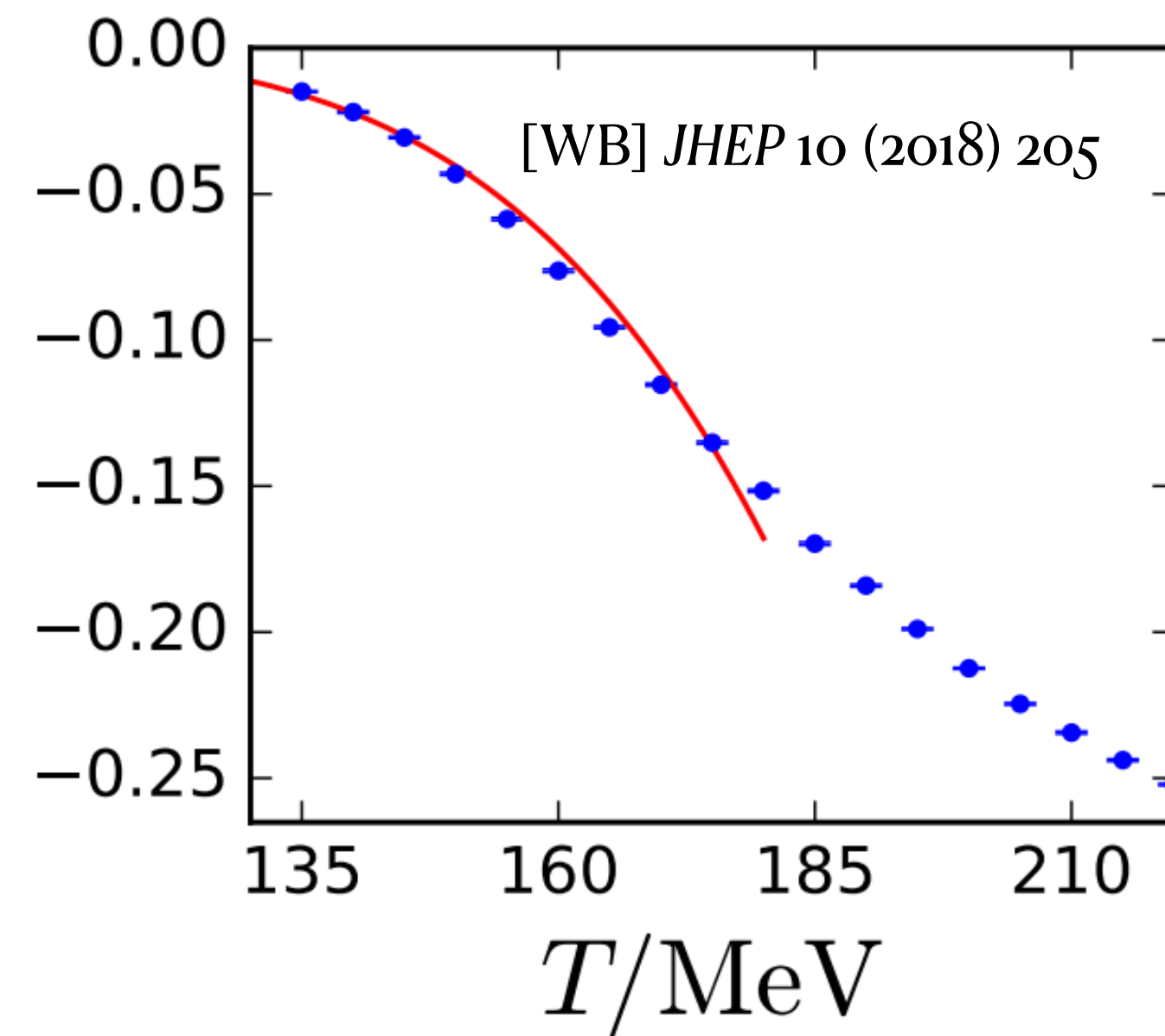
Sure!

Using a Taylor's series (a la lattice QCD) up to  $\mathcal{O}(\mu/T)^2$ :

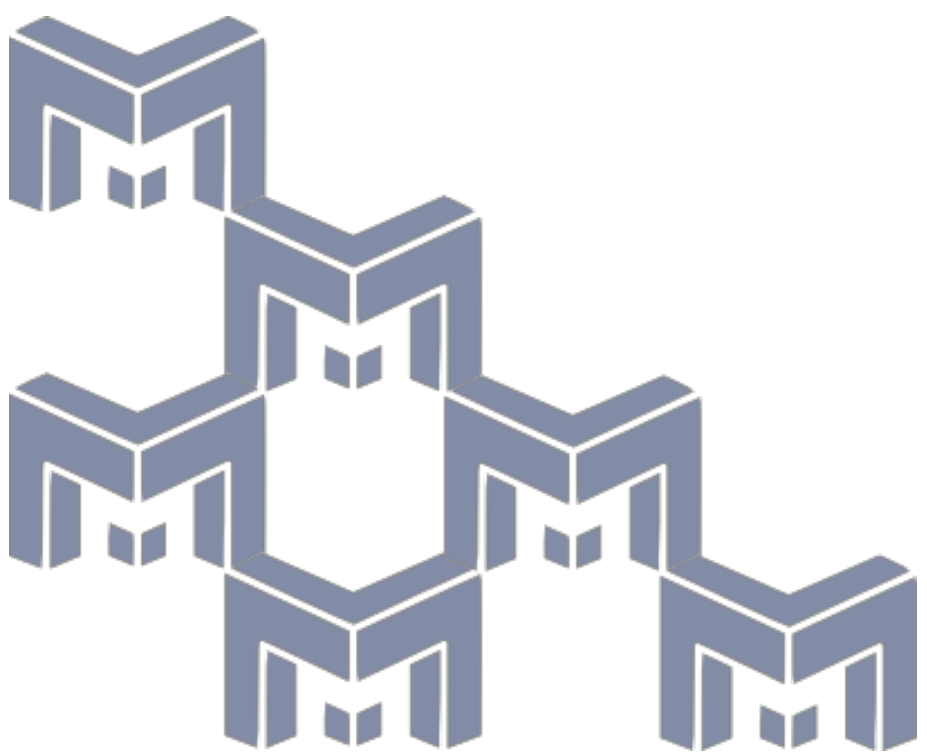
$$n_S(T, \mu_B, \mu_S) = \chi_S^2 \left( \frac{\mu_S}{T} \right) + \chi_{BS}^{11} \left( \frac{\mu_B}{T} \right)$$

For strangeness neutrality  $n_S = 0$ ,

$$\text{so } \mu_S = - \frac{\chi_S^2}{\chi_{BS}^{11}} \mu_B$$



Since  $\chi_S^2 > 0$  (for stability) and  $\chi_{BS}^{11} < 0$  (since  $S = -1$ ) then  $\mu_S$  has the same sign as  $\mu_B$  for strangeness neutrality



# At finite $\mu_S \neq 0$ , tricritical point

Example  $T = 0$   
phase diagram

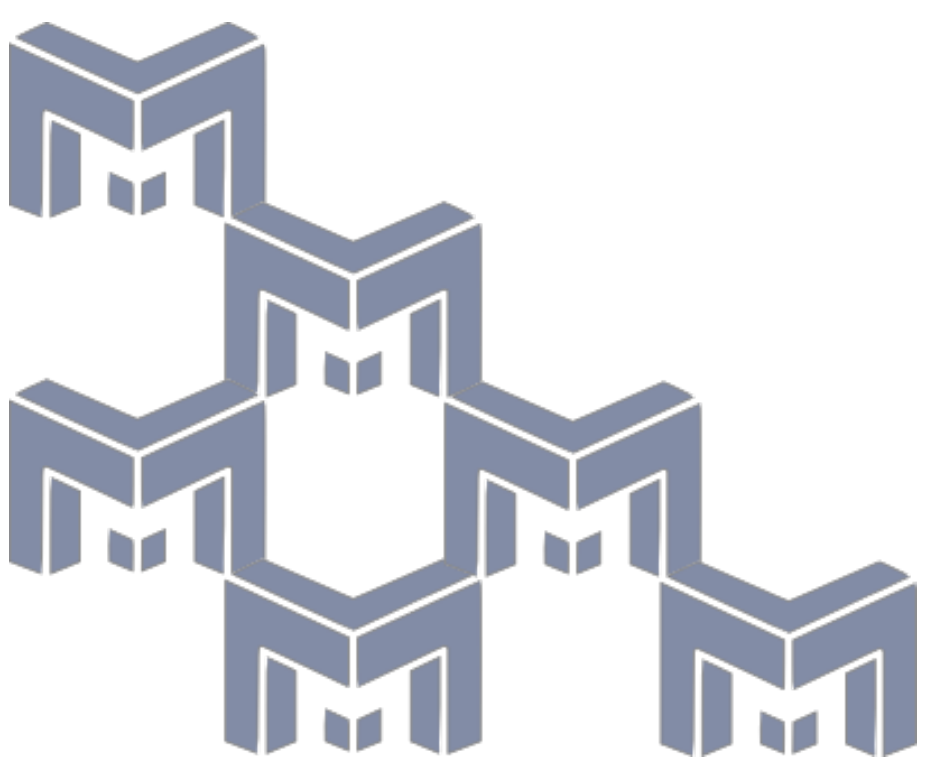
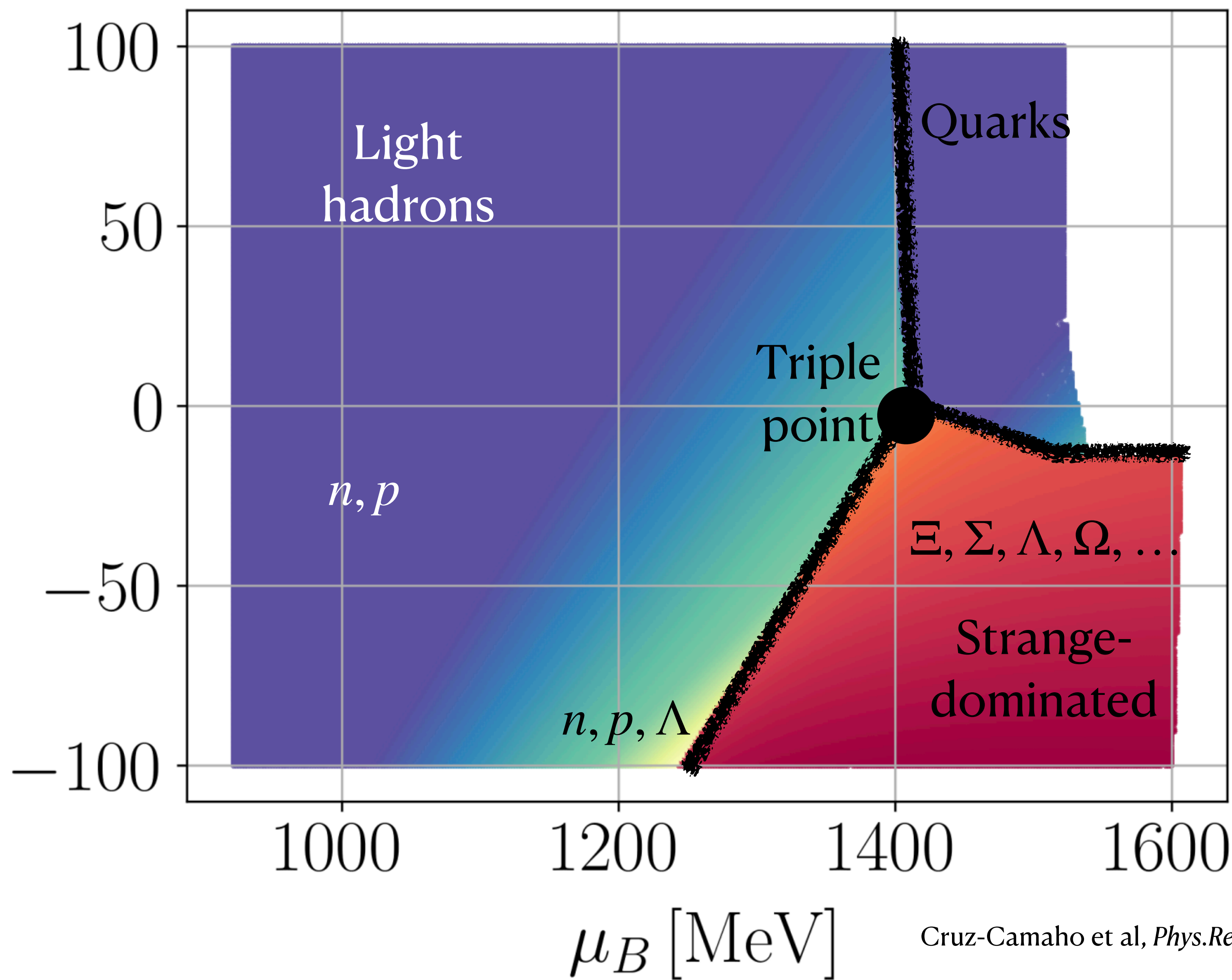
Heavy-ions

$$\mu_S > 0$$

Electroweak  
equilibrium

$$\mu_S = 0$$

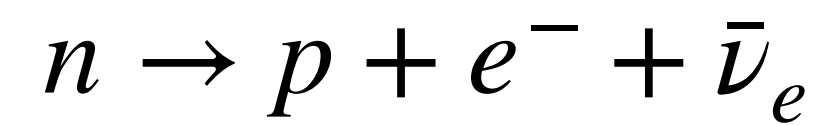
$\mu_S$  [MeV]



# Strangeness-dominated phase in neutron stars?

Need fluctuations in  $\mu_S$

Due to long time scales in neutron stars, they are in electroweak equilibrium i.e.



$$\mu_n = \mu_p + \mu_e$$

that leads to

$$\mu_S = 0$$

However, binary neutron stars mergers can *perturb* the system out-of-electroweak-equilibrium

Measurements: Hegade *Nature Astron.* 8 (2024) 10, 1277-1283; *Phys.Rev.D* 108 (2023) 10, 10 ; *Phys.Rev.D* 110 (2024) 4, 044041;  
*Phys.Rev.D* 109 (2024) 10, 104064

Calculations: Alford et al, *Phys.Rev.D* 110 (2024) 6, L061303;

*Phys.Rev.D* 108 (2023) 8, 083019;

*Phys.Rev.C* 109 (2024) 5, 055803;

*Phys.Rev.D* 104 (2021) 10, 103027;

*Phys.Rev.C* 103 (2021) 4, 045810;

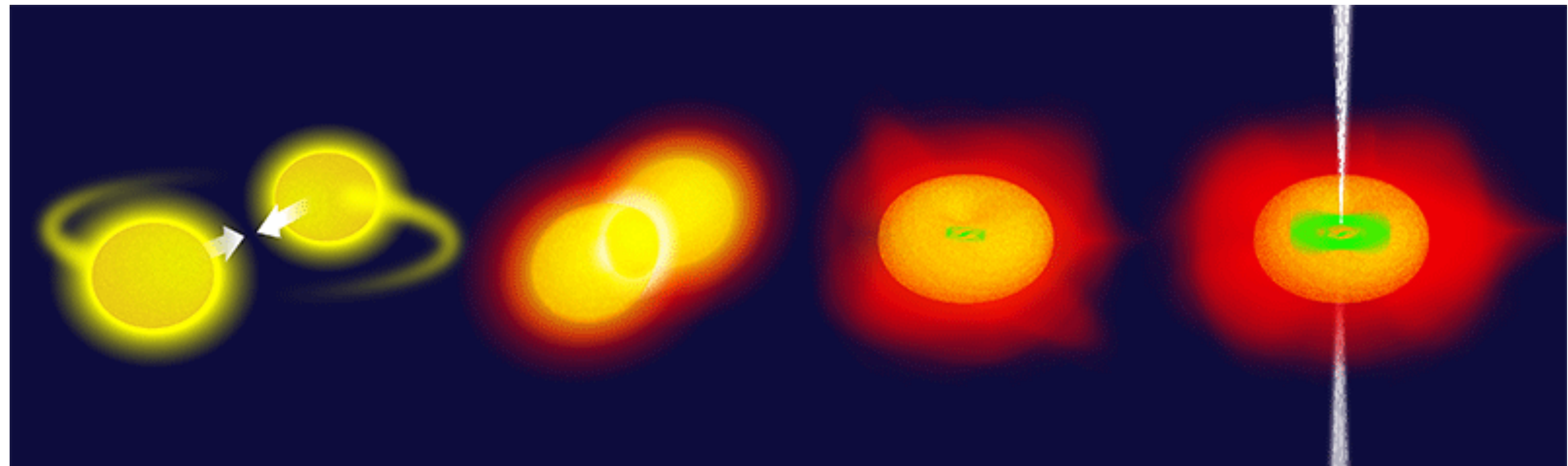
*Phys.Rev.D* 100 (2019) 10, 103021

Most et al, *Astrophys.J.Lett.* 967 (2024) 1, L14;

*Mon.Not.Roy.Astron.Soc.* 509 (2021) 1, 1096-1108

Yang et al, *Phys.Rev.Lett.* 135 (2025) 22, 222702;

*Phys.Rev.C* 109 (2024) 1, 015805



# Coupled electroweak rate equations

Storbacka, Wu, Dong, Cruz Camacho, Haber, Reinke Pelicer, Dexheimer, Most, Noronha-Hostler, to appear

Rate equations:

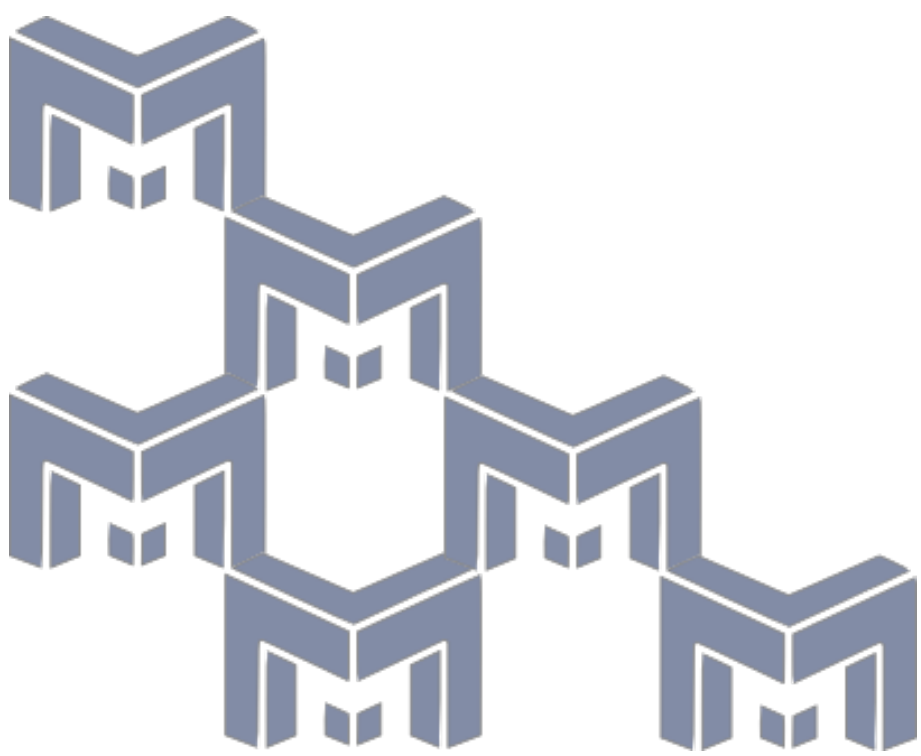
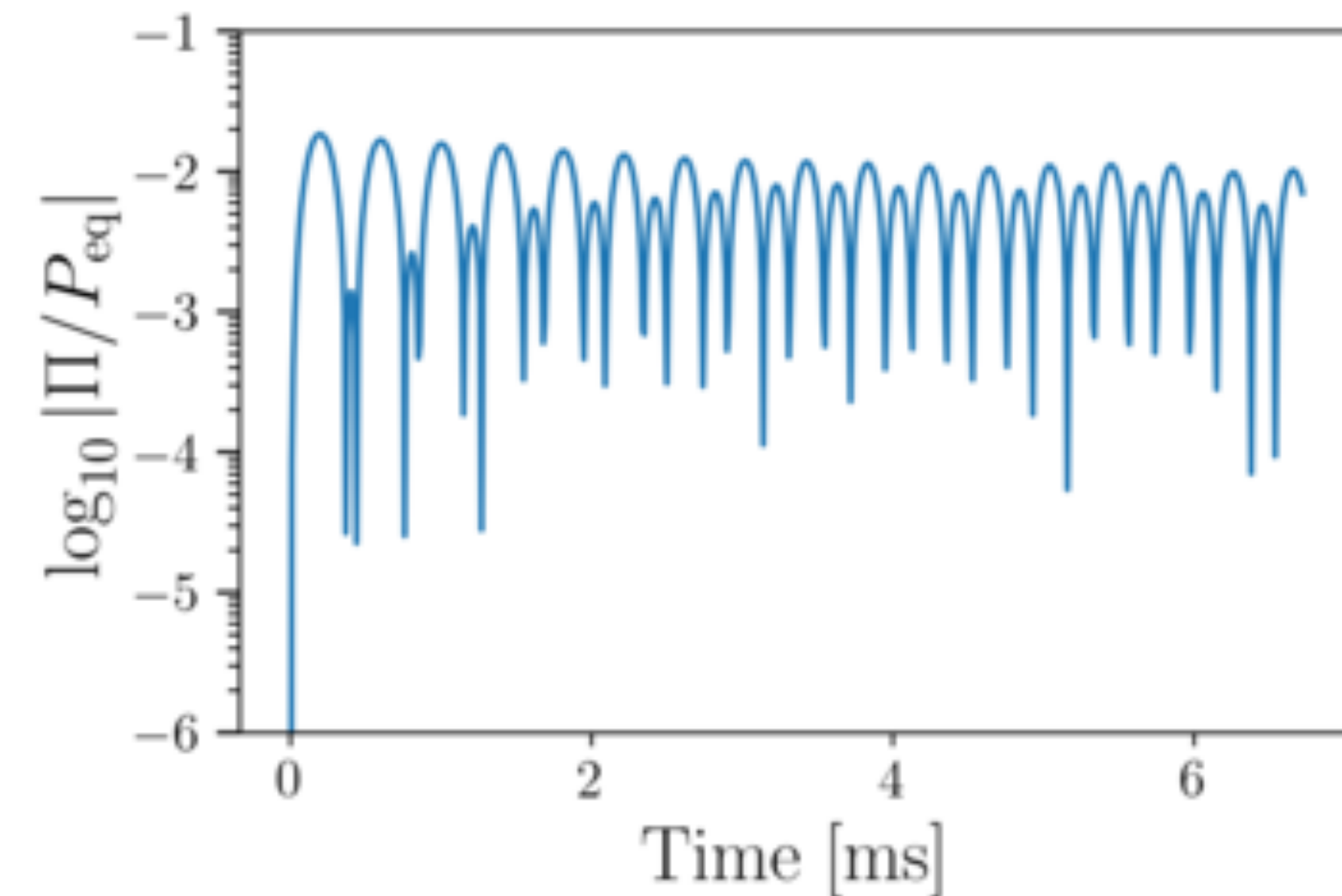
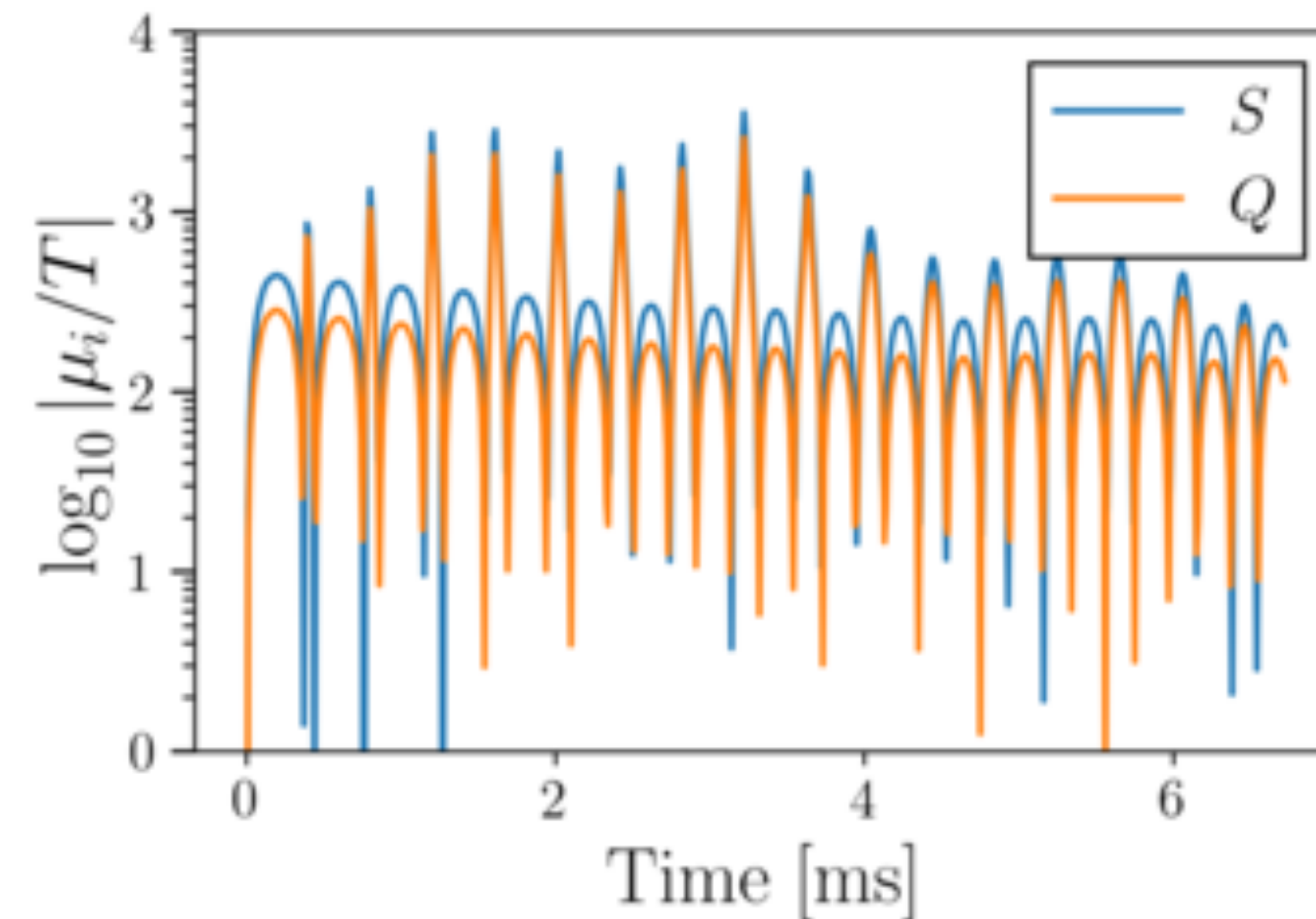
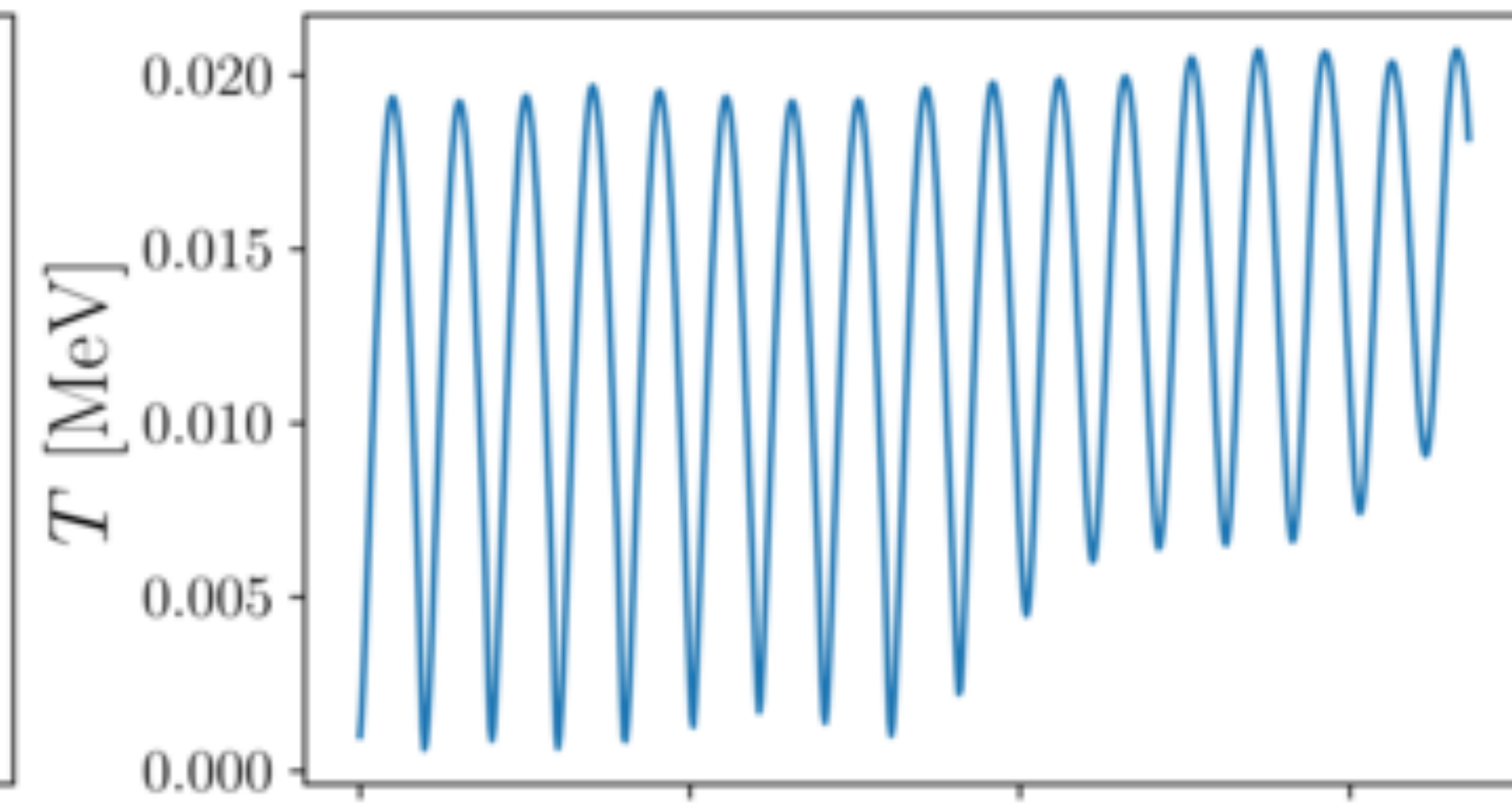
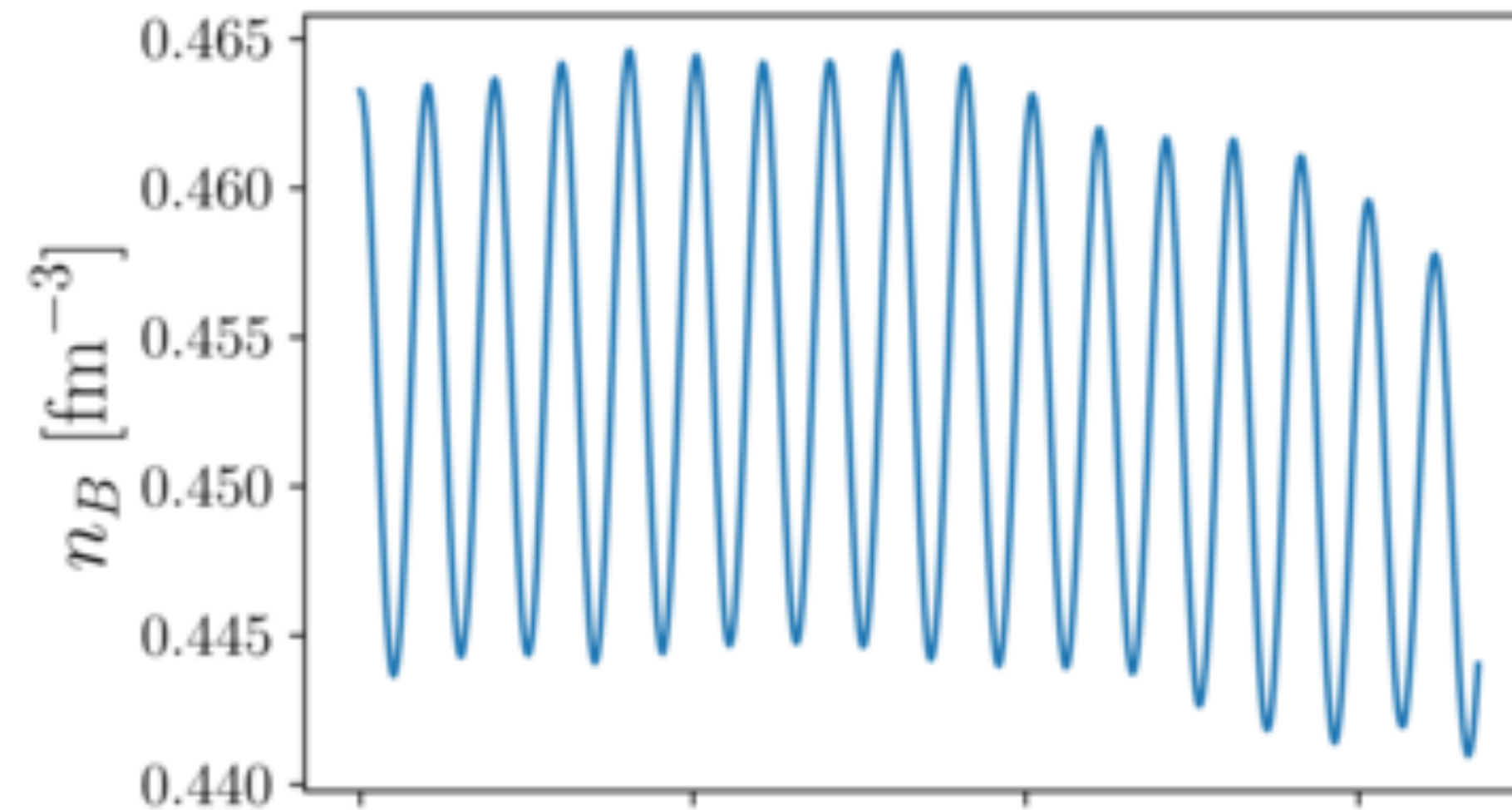
$$\dot{n} = \Gamma [\text{gain} - \text{loss}] + \dots$$

for each particle

In medium reaction rates

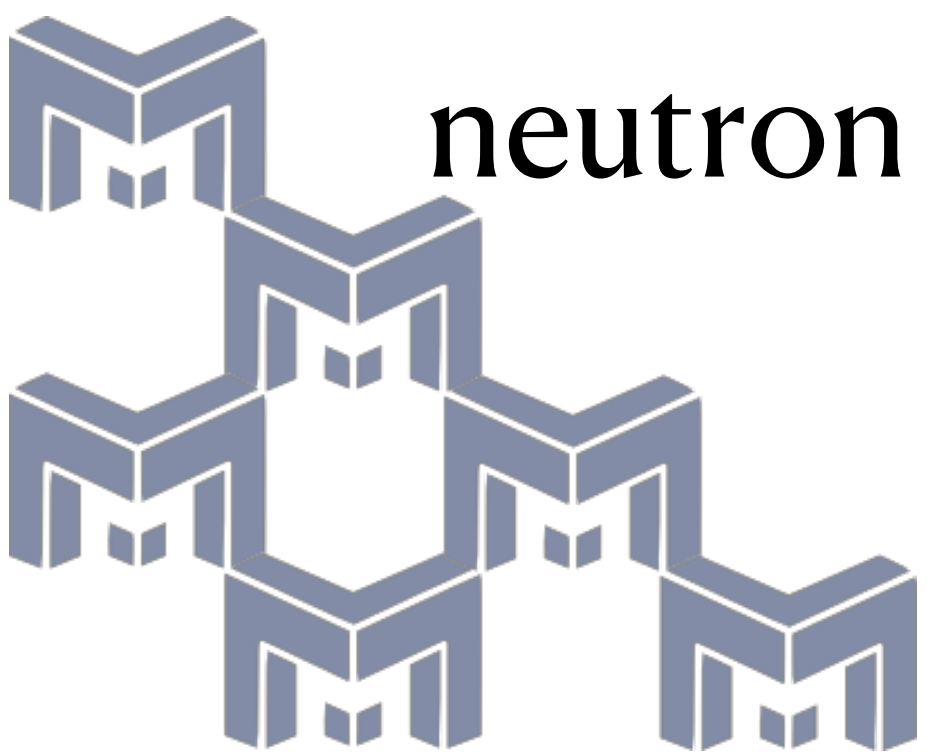
Equation of state

$$\left\{ T, \mu_B, \mu_S, \mu_Q \right\}$$



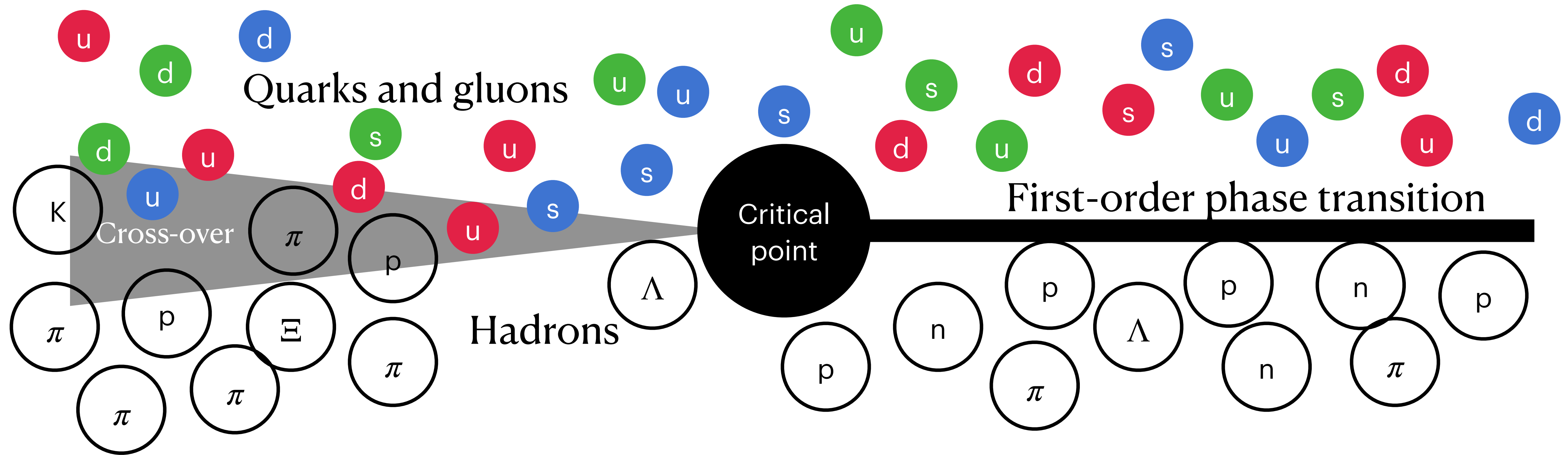
# Conclusions and Outlook

- Using a chiral mean field model, we find a new possible phase of matter at large densities: strangeness dominated
- This phase of matter is associated with isospin symmetry restoration
- Depending on the parameters of the equation of state, this phase may either be a metastable phase between light hadrons and deconfinement or a first-order phase transition with a triple point and critical point
- Now possible to explore coupled out-of-electroweak equilibrium in binary neutron star mergers, find large fluctuations in  $\mu_Q, \mu_S$ !



# Searching for the QCD critical point

Boundary between a first-order phase transition and cross-over



Finding the QCD critical point would be the first...

1. in a fundamental theory of nature,
2. in a truly ultra-relativistic many-body system,
3. accessed through relativistic hydrodynamic evolution

