

# Recent multiplicity-based measurements in jet physics

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Laboratory

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U.S. DEPARTMENT OF  
**ENERGY**

# Is there QGP creation in small systems?

## Experimental observation 1:

There is QGP formation in A-A collisions (large systems).

## Experimental observation 2:

Collective phenomena in high-multiplicity pp collisions:

- significant elliptic flow ( $v_n$ ),
- increased production of rare hadrons,
- ridge structure.

## Question: How do QGP signatures change with system size?

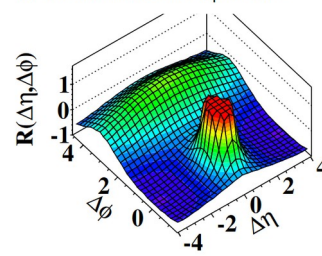
- Is there evidence for jet modification in small systems?

## QGP is not required to explain collective phenomena:

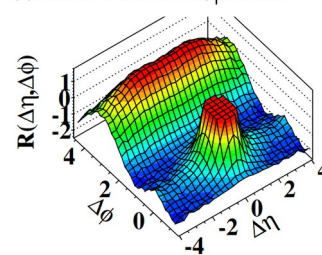
- These can be understood through **vacuum QCD processes** occurring at the boundary between the soft and hard regimes. e.g. **multiple-parton interactions** and **color reconnection**.

- We do not expect suppression, but the **shape** of the jets could be modified as a function of **multiplicity**.

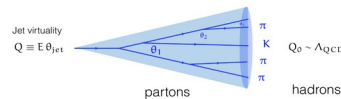
(b) CMS MinBias,  $1.0\text{GeV}/c < p_T < 3.0\text{GeV}/c$



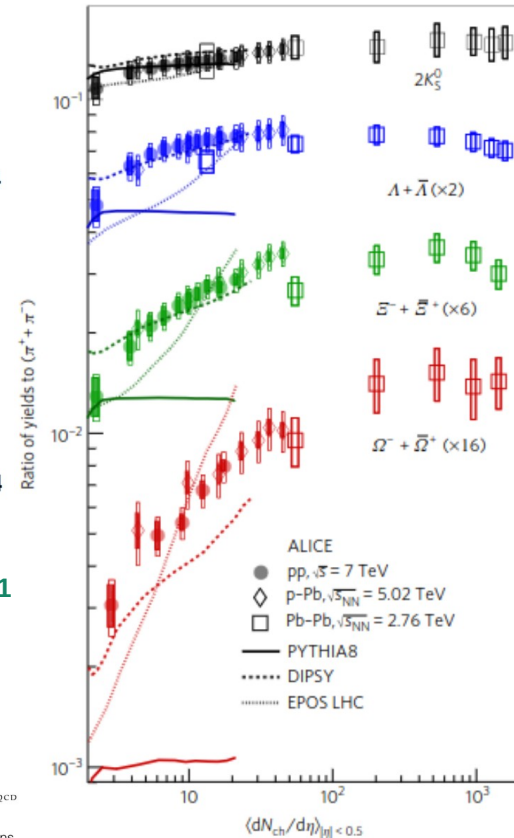
(d) CMS  $N \geq 110$ ,  $1.0\text{GeV}/c < p_T < 3.0\text{GeV}/c$



CMS, JHEP 09 (2010) 091



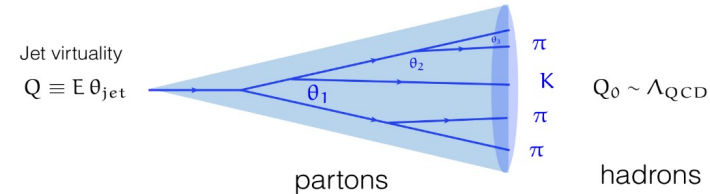
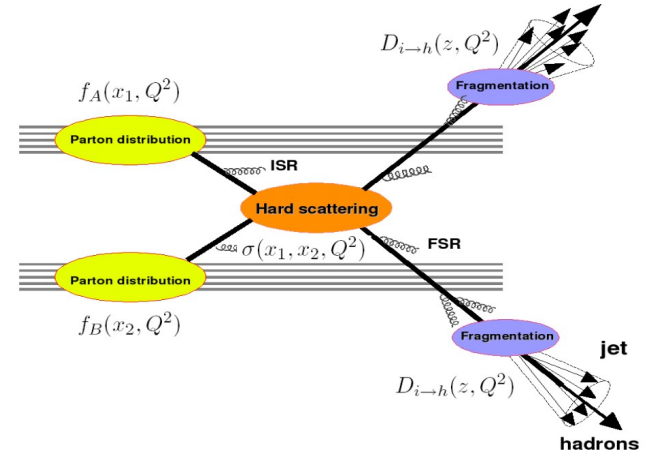
Nature Physics 13 (2017) 535-539



# Background

- **Jets:** collimated showers of particles produced by fragmentation and hadronization of hard-scattered partons.
- **Experimentally:** defined by a jet reconstruction algorithm and a jet resolution parameter  $R$ .
- An experimentally accessible **observable** to “capture” the directly unmeasurable **parton shower**.
- **Jet multiplicity** is the number of charged final-state particles within the jet cone.
- Important information to probe particle production models.

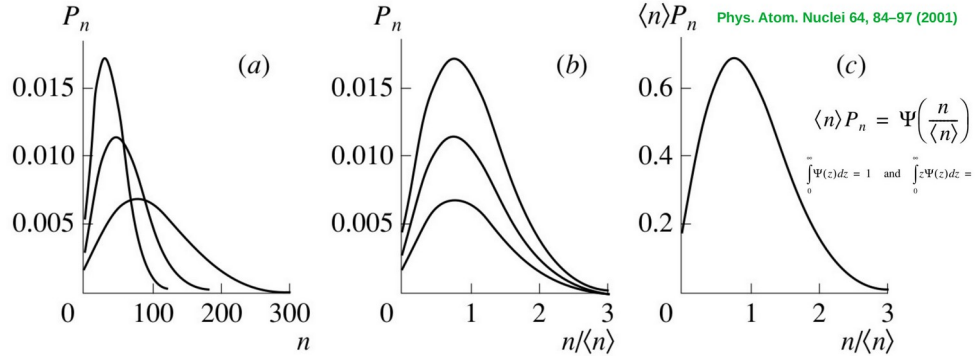
D. d’Enterria arXiv:0902.2011



# Koba-Nielsen-Olesen (KNO) scaling

## Koba-Nielsen-Olesen (KNO) scaling hypothesis:

- Influential contribution to the analysis of event multiplicities in high-energy particle collisions.



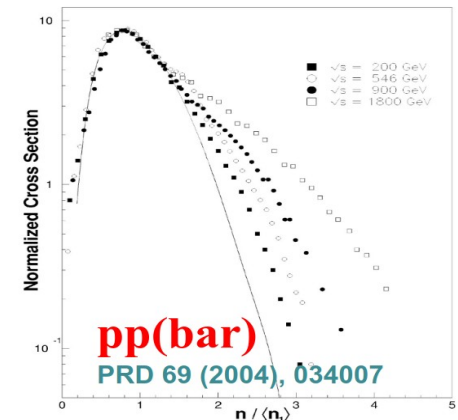
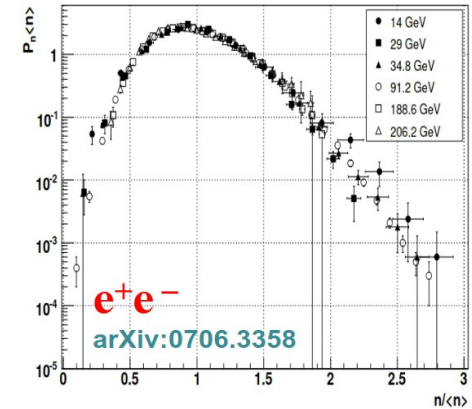
**(a):** many possible function shapes for given energy.

**(b):** Contract linearly along horizontal axis in proportion to  $\langle n \rangle$ .

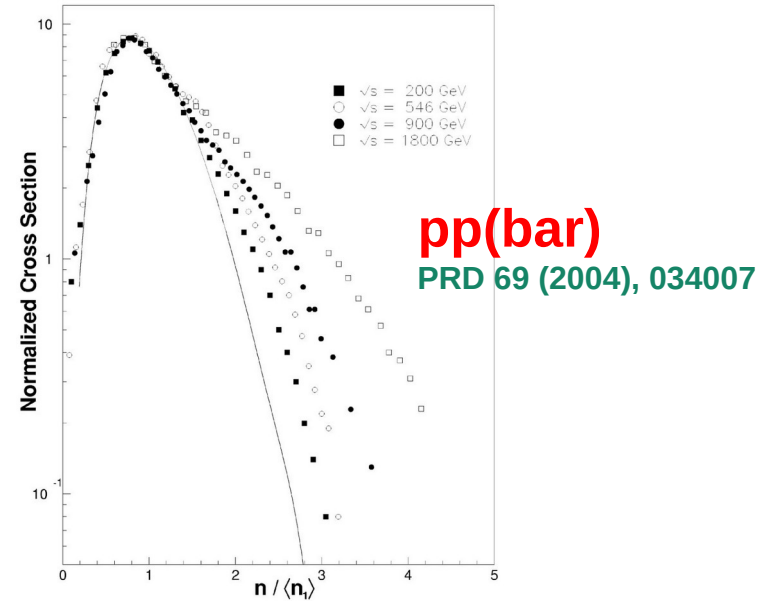
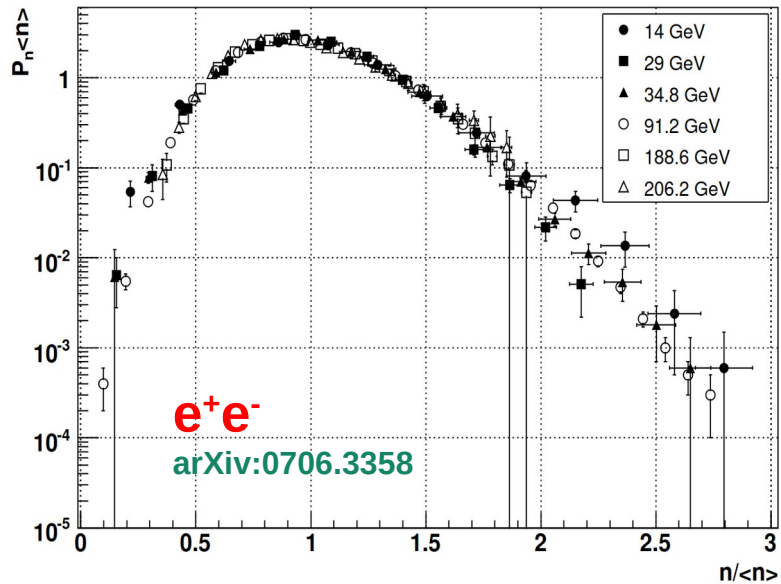
**(c):** Extend along vertical axis in proportion to  $\langle n \rangle$ .

**KNO-scaling:** these curves (multiplicity distributions) will coincide at each point.

**Recent phenomenological studies:** a similar scaling may hold within single jets for **jet multiplicity as a function of jet  $p_T$** .



# Koba-Nielsen-Olesen (KNO) scaling

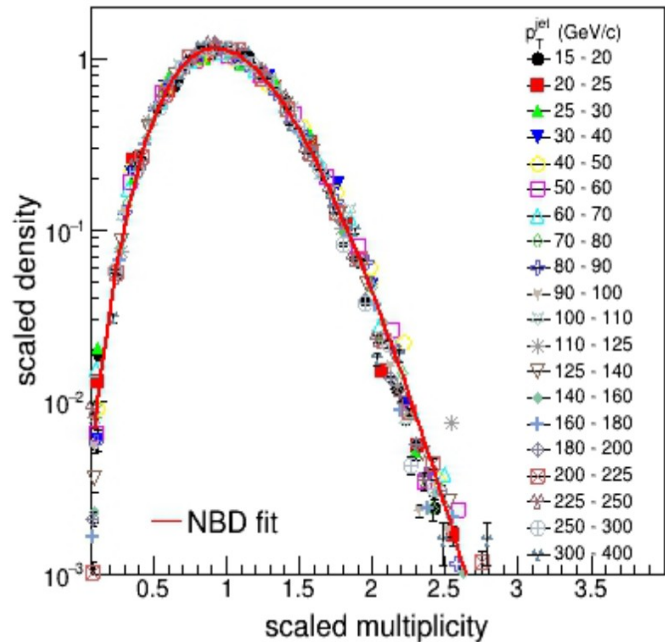


- **KNO scaling**: observed in  **$e^+e^-$  collisions**, and in  **$p$ - $\bar{p}$**  collisions.
- The scaling is violated toward higher energies and larger rapidity windows.
- Reason not fully understood, it might be violated by **MPI** and **overlapping color strings**.

Walker PRD 69, 034007 (2004); Abramovsky et al., arXiv:0706.3358

- **Is KNO scaling valid in jets? Origin of the scaling? How is it affected by MPI and CR?**
- **Flavor dependence: initial pQCD process or parton shower?**

# Simulated pp collisions – a KNO-like scaling

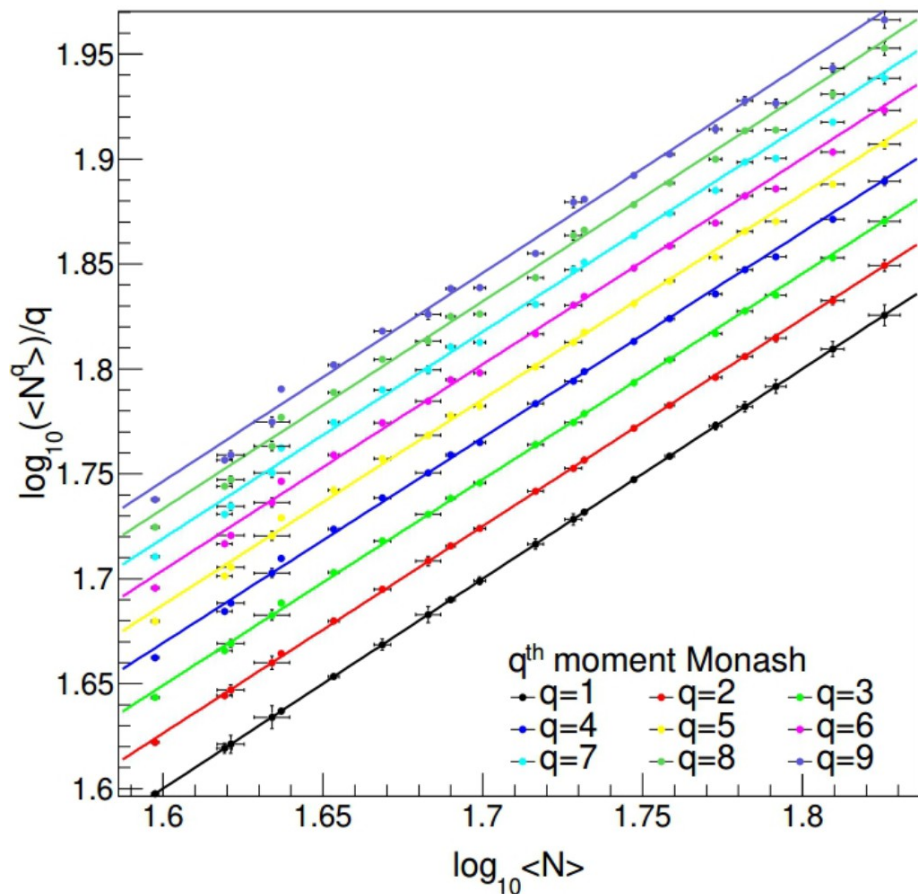


**Pythia 8 standalone simulations:**  
**A KNO-like scaling is observed within single jets for a wide range of jet  $p_T$  values!**

- Multiplicity (dominated by the jet multiplicity) vs. jet momentum  $p_{T}^{\text{jet}}$
- Parametrized with a NBD
- Distributions at all  $p_{T}^{\text{jet}}$  fit well on a single NBD curve

$$P_N = \frac{\Gamma(Nk + a)}{\Gamma(a)\Gamma(Nk + 1)} p^{Nk} (1 - p)^a$$

# Simulated pp collisions – statistical moments



Quantifying how well the scaling is fulfilled:

- q<sup>th</sup> statistical moment

$$\langle N^q \rangle = \sum_{N=1}^{\infty} P_N N^q$$

- sensitive to goodness of scaling
- insensitive to fluctuations
- no need to parametrize and fit

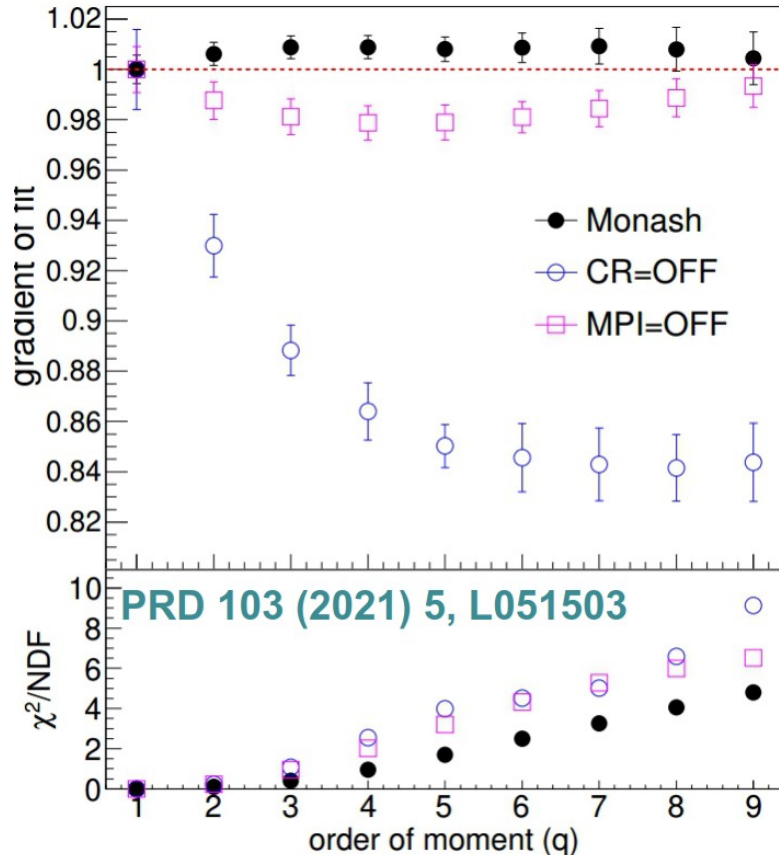
- Scaling:

$$\langle N^q(p_T^{\text{jet}}) \rangle = \lambda^q(p_T^{\text{jet}}) \langle N^q(p_0) \rangle \quad \lambda(p_0) = 1$$

$$\log \langle N^q \rangle / q \approx \log \langle N \rangle$$

**Scaling is fulfilled in the whole jet  $p_T$  range!**

# Multiplicity vs $p_T^{\text{jet}}$ : moments



## Quantifying how well the scaling is fulfilled:

- $q^{\text{th}}$  statistical moment

$$\langle N^q \rangle = \sum_{N=1}^{\infty} P_N N^q$$

- sensitive to goodness of scaling
- insensitive to fluctuations
- no need to parametrize and fit

- **Scaling:**

$$\langle N^q(p_T^{\text{jet}}) \rangle = \lambda^q(p_T^{\text{jet}}) \langle N^q(p_0) \rangle \quad \lambda(p_0) = 1$$

- **Origin of scaling (PYTHIA)**

- Physical case (Monash): All 9 moments are consistent with unity, slope within ~1%
- No CR: Scaling is broken by ~15%
- No MPI (also no CR): Scaling fulfilled to ~2%



# Heavy-flavor jets

**Heavy-flavor (HF) jets:** created via hard pQCD processes

- LO flavor creation
  - NLO gluon splitting + flavor excitation
- } contributions are of similar magnitude

S. Cao, G.Y. Qin, X.N. Wang,  
Phys.Rev.C 93 (2016) 2, 024912

**Jet production depends on quark flavor:**

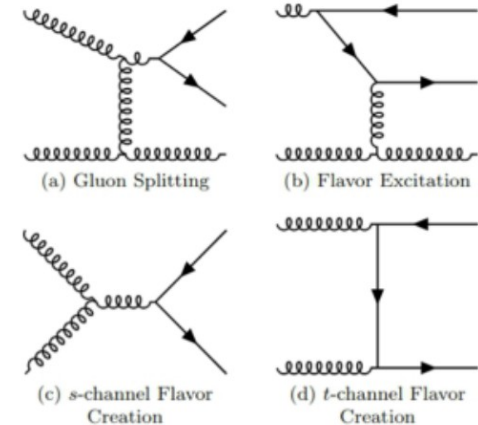
- mass dependence: harder fragmentation (dead-cone)
- color-dependence: HF initiated by quark jets only

**Comparing the scaling for LO vs. NLO:**

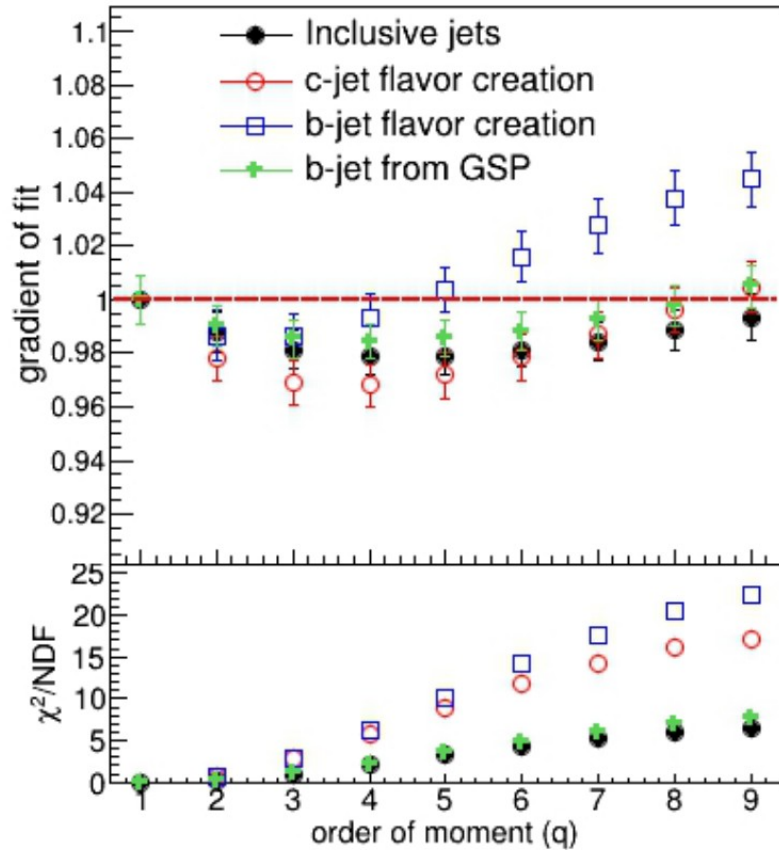
- **sensitivity to its origin** (i. e. hard QCD vs. jet development)



**KNO-like scaling in HF jets may provide information on the origin of KNO scaling!**



# Heavy-flavor jet scaling



Comparison of scaling LO and NLO:

- Sensitivity to its origin (hard QCD vs. jet development)

All slopes are around unity within 5%

- Flavor creation (LO): mass-dependent deviation from inclusive jets

- Gluon splitting (NLO): Follows inclusive (mostly gluon) jets

Scaling is driven by initial hard processes!

What do we see in DATA?

- Analysis on published ATLAS data has similar conclusions.

G. R. Germano, F. S. Navarra, G. Wilk, Z. Włodarczyk Phys.Rev.D 110 (2024) 3, 034026

- Analysis in ALICE is underway!

# Charm baryon enhancement

Perturbative quantum chromodynamics (pQCD) calculations have been successful in describing the production of heavy-flavor mesons for several collision energies at the LHC.

$$d\sigma_{AB \rightarrow C}^{\text{hard}} = \sum_{a,b} f_{a/A}(x_a, Q^2) \otimes f_{b/B}(x_b, Q^2) \otimes d\sigma_{ab \rightarrow c}^{\text{hard}}(x_a, x_b, q^2) \otimes D_{c \rightarrow C}(z, Q^2)$$

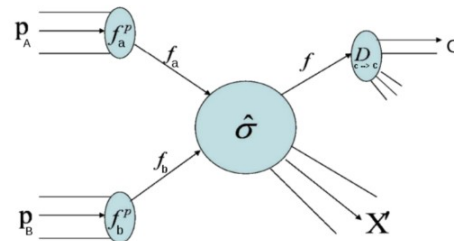
Parton Distribution Function (PDF)

Partonic hard scattering  
cross-section

Fragmentation  
function

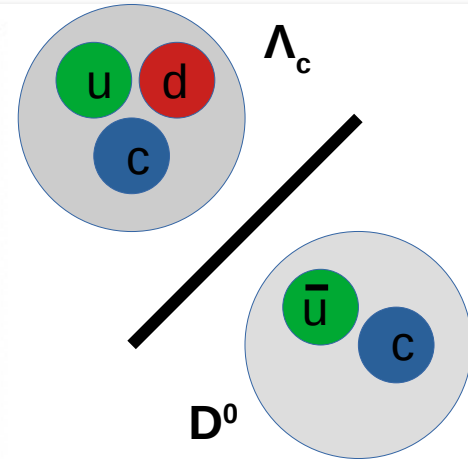
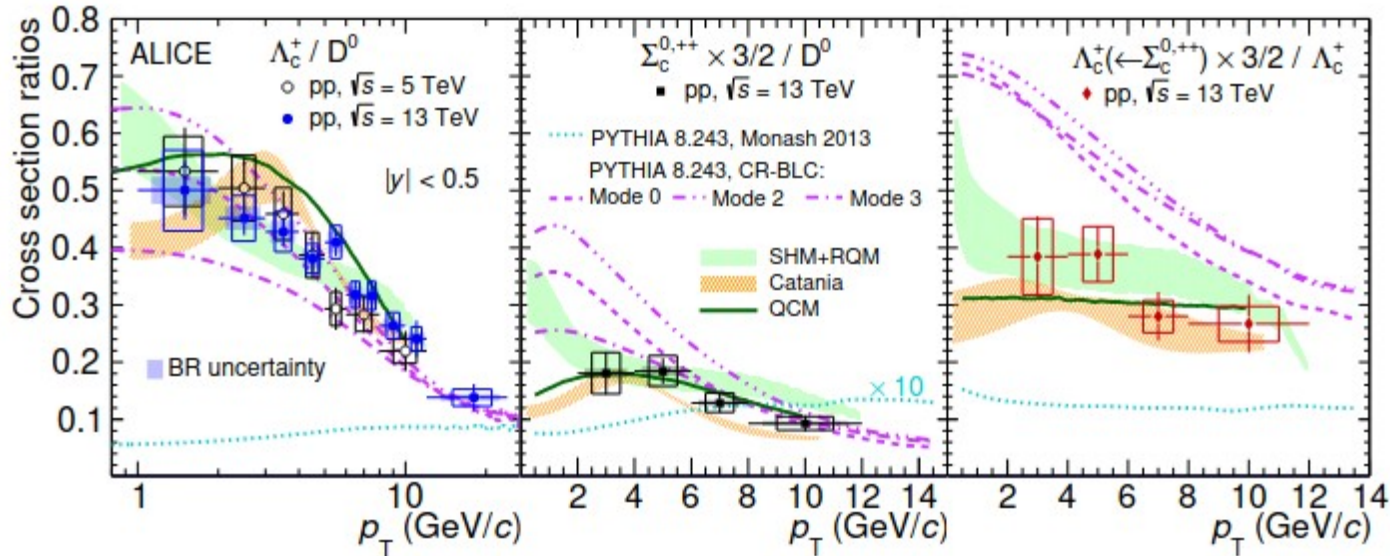
The usual description relies on the **factorization approach**: production cross section of heavy-flavor hadrons is calculated as a convolution of

- parton density functions (PDFs) of the colliding hadrons,
- **cross section of the hard-scattering process** and
- heavy-quark fragmentation function.



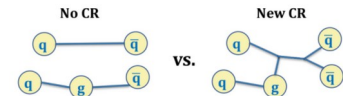
**Fragmentation function**: usually taken from  $e^-e^+$  collisions on the assumption of universality, e.g. Monash tune in PYTHIA 8. → **Experimental evidence suggest this is not true!**

# Charm baryon enhancement



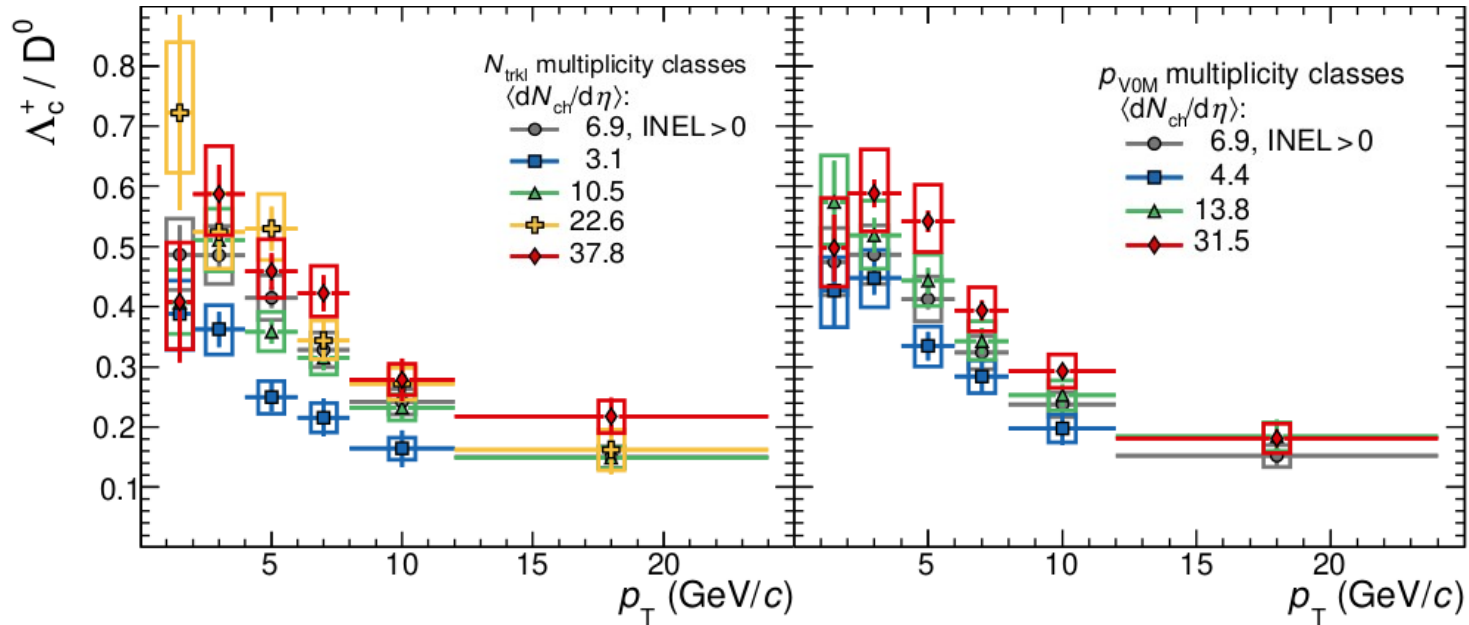
- Charm baryon-to-meson ratios: **sensitive probes of fragmentation!**
- $\Lambda_c/D^0$  and  $\Sigma_c/D^0$  underestimated by models: **HF fragmentation universality broken!**

- **PYTHIA 8 CR-BLC**: string formation beyond leading color approximation,
- **SH model + RQM**: feed-down from augmented set of charm-baryon states,
- **Catania**: fragmentation + coalescence of charm and light quarks,
- **QCM**: coalescence model based on statistical weights + equal quark-velocity.



# Multiplicity dependence of $\Lambda_c/D^0$

ALICE Coll. Phys.Lett.B 829 (2022) 137065



- The enhancement in  $\Lambda_c/D^0$  depends on the final state multiplicity at mid- and forward rapidity.
- **Goal: Understand the origin of the enhancement with detailed event activity studies.**
- **Does it originate in jet processes or the underlying event? How to measure it?**
- Using standalone PYTHIA 8 to test the observable effects of the **CR-BLC model**.

# Event classifiers

The collisions can be categorized with different **event characteristics**, which help distinguish how much the **jets** or the **background** event dominate in a given event.

- $N_{CH}$  : **charged-hadron multiplicity** at central rapidity ( $|\eta| < 1$ ):
  - number of final-state charged particles,
  - global parameter that does not take leading process into account.
- $N_{fw}$  : **forward multiplicity** ( $2 < \eta < 5$ ).

## - $R_T$ : relative transverse multiplicity (UE activity)

- exclude jets from leading process,
- high- $p_T$  trigger hadron required.

$$R_T = \frac{N_{CH}^{transverse}}{\langle N_{CH}^{transverse} \rangle} \quad \frac{\pi}{3} < |\Delta\phi| < \frac{2\pi}{3}$$

## - $R_{NC}$ : relative near-cone multiplicity (jet region activity)

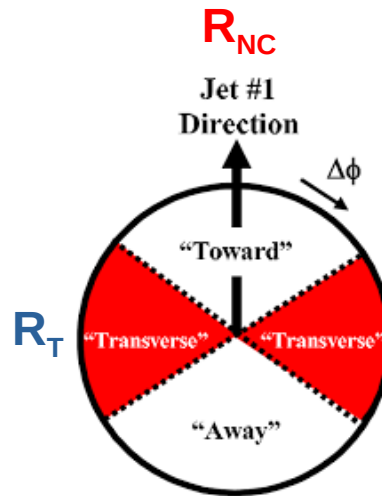
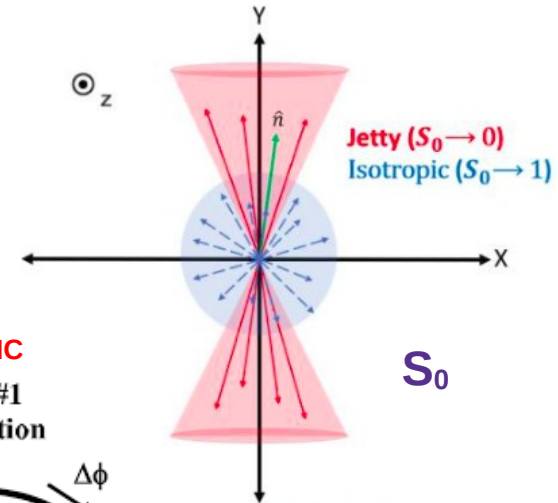
- includes leading process,
- high- $p_T$  trigger required.

$$R_{NC} = \frac{N_{CH}^{near-side cone}}{\langle N_{CH}^{near-side cone} \rangle} \quad \sqrt{(\Delta\phi)^2 + (\Delta\eta)^2} < 0.5$$

## - $S_0$ : sphericity

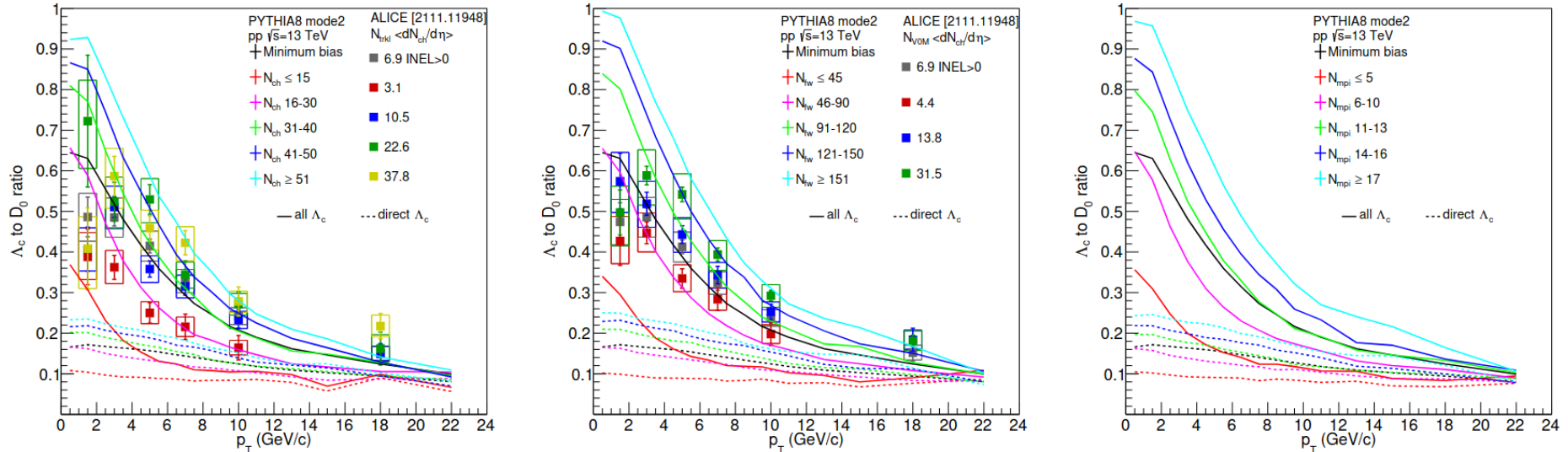
- jetty or isotropic event,
- no need for trigger hadron.
- only mid-rapidity.

$$S_0 = \frac{\pi^2}{4} \times \min_{\hat{n} = (n_x, n_y, 0)} \left( \frac{\sum_i |\vec{p}_{T_i} \times \hat{n}|}{\sum_i \vec{p}_{T_i}} \right)^2$$



# Multiplicity dependence of $\Lambda_c/D^0$

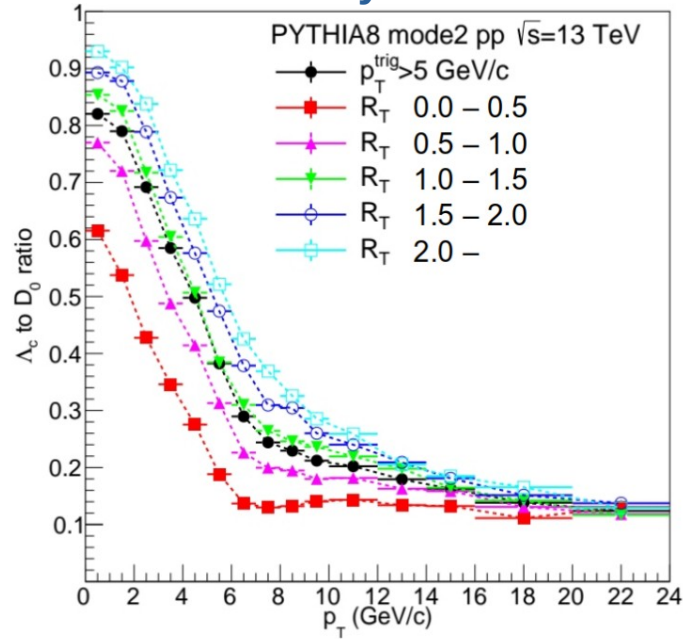
Z. Varga, R. Vértési, A. Misák, J. Phys. G: Nucl. Part. Phys. 50 (2023) 075002



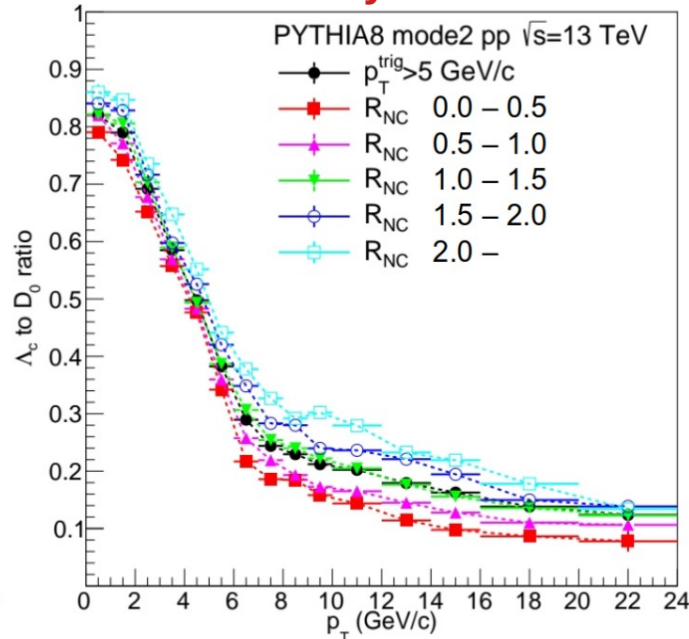
- **Simulations are in agreement with ALICE experiment** (minbias + multiplicity-differential).
- Rapidity gap: reduces correlation between leading hard processes and the multiplicity.
- **Multiplicity dependence not driven by charm production inside jets.**

# $\Lambda_c/D^0$ yield for triggered events

## UE activity classifier

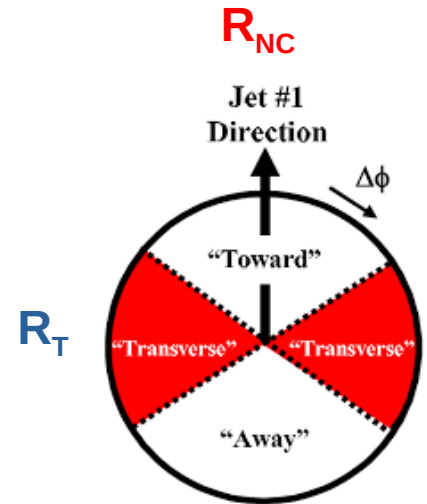


## Jet activity classifier



Z. Varga, R. Vártesi, J. Phys. G: Nucl. Part. Phys. 49 (2022) 075005

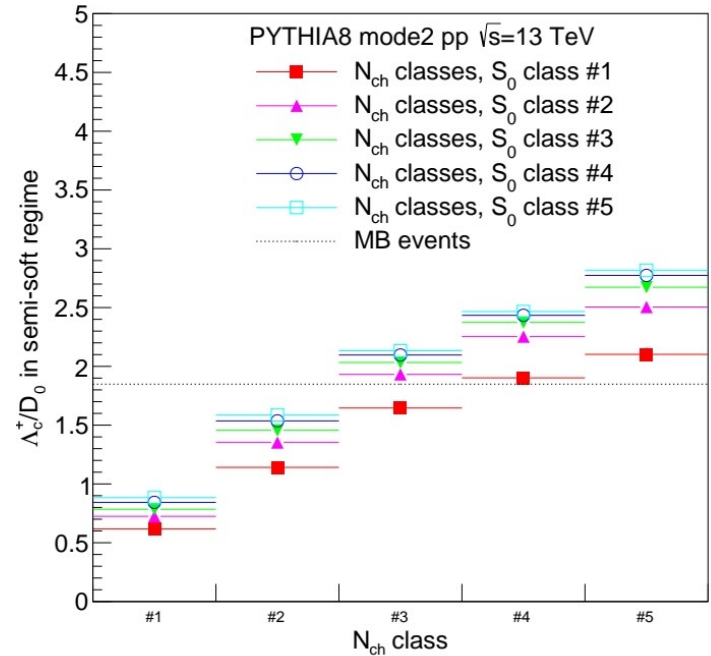
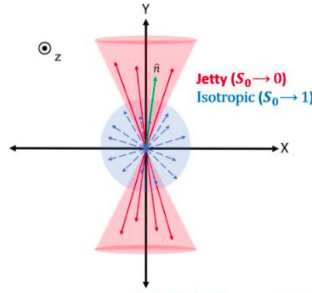
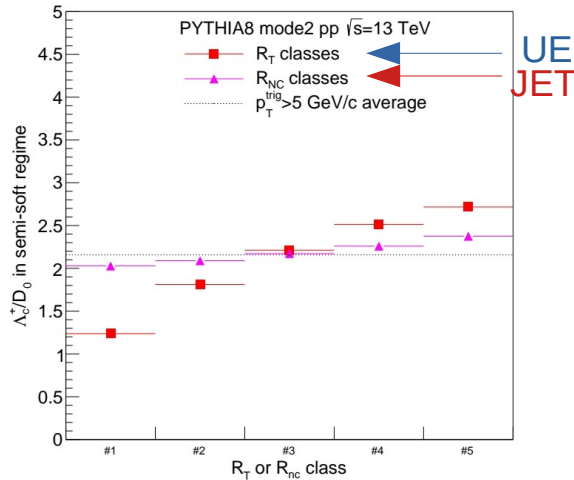
$$\Lambda_c(qqc), l = 0$$



- Events require  $p_T > 5$  GeV/c **hadron trigger**.
- Significant difference is observable in case of  $R_T$  (UE activity).
- No significant difference when classified by  $R_{NC}$  classes (jet activity).

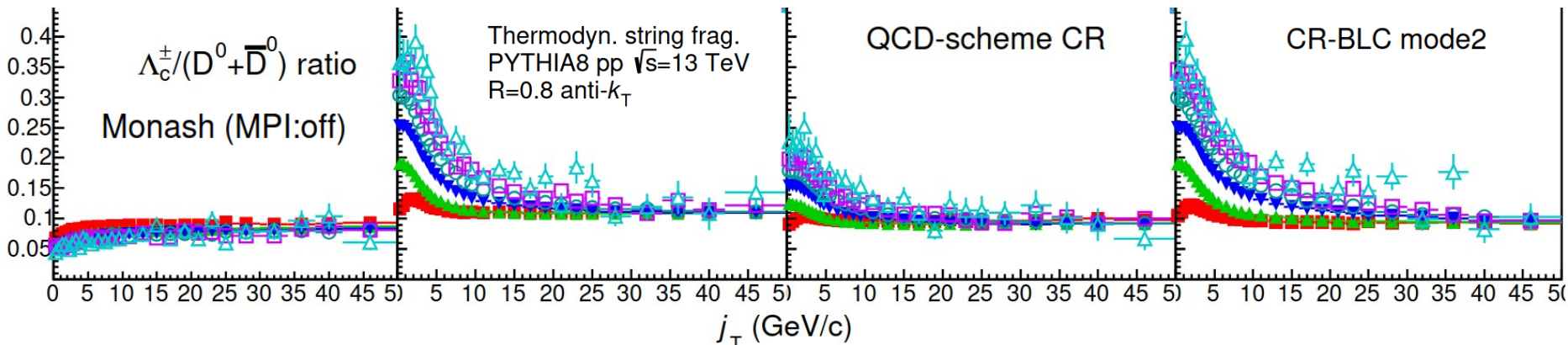


# $\Lambda_c^+/D^0$ yield ratios – trigger vs. minbias



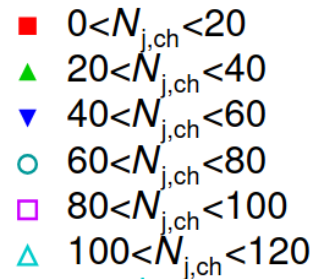
- In case we require a hard process ( $P_T^{\text{trigger}} > 5$  GeV/c):
  - **Strong dependence** of ratios on the **UE activity**,
  - **No pronounced dependence** on the **jet multiplicity**.
- In minimum-bias events
  - In case of high final-state multiplicity, ratio depends on jettiness,
  - Dependence is minute for low final-state multiplicity.
- For sphericity  $S_0$ , dependence on jettiness observable in minimum-bias events. No need to use a trigger that biases the sample and decreases available statistics.

# String fragmentation in jets

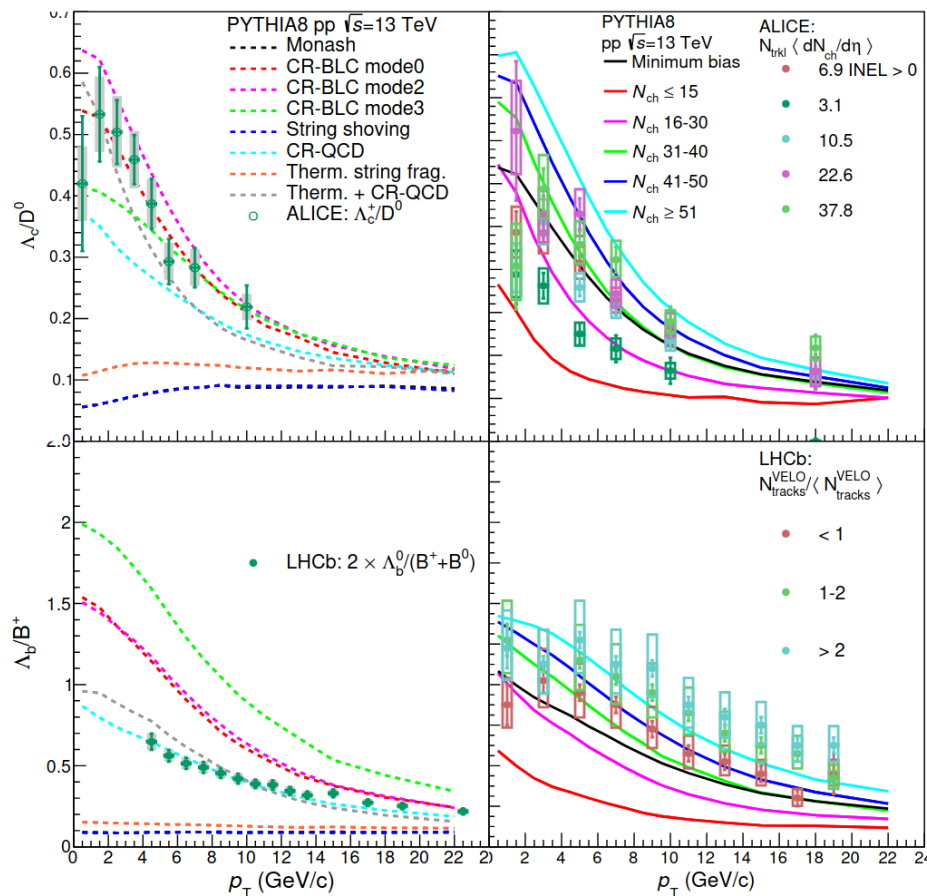


R. Vártesi, A. Ortiz, arXiv:2408.06340

- $j_T$ : momentum component perpendicular to the jet axis.
- $N_{j,ch}$ : jet charged-constituent multiplicity.
- $\Lambda_c/D^0$  shows an increase at low  $j_T$  with increasing  $N_{j,ch}$ . This effect is similar to the multiplicity dependence of the ratio as a function of  $p_T$  reported by the ALICE Collaboration.
- The ratio is nearly flat for the lowest  $N_{j,ch}$  class, and consistent with the corresponding ratios measured in  $e^-e^+$  collisions at LEP.
- Utilizing jet substructure measurements provide more information on the fragmentation of heavy flavor.



# $\Lambda_b/B^+$ yield for triggered events



L. V. Földvári, Z. Varga, R. Vértesi, arXiv:2408.16447

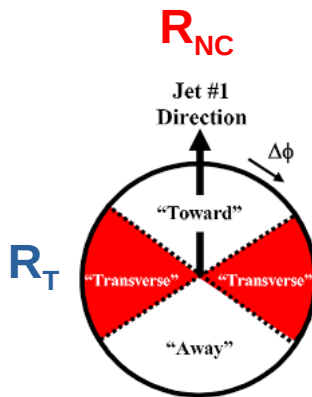
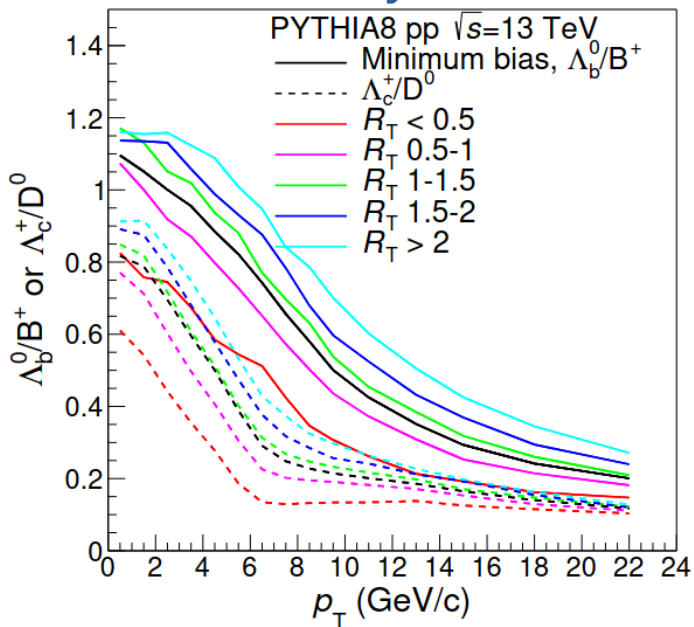
- Performance of multiple models compared for  $\Lambda_c/D^0$ ,
- and also to recent  $\Lambda_b/B^+$  measurements by LHCb.

LHCb Collab. Phys. Rev. Lett. 132, 081901 (2024)

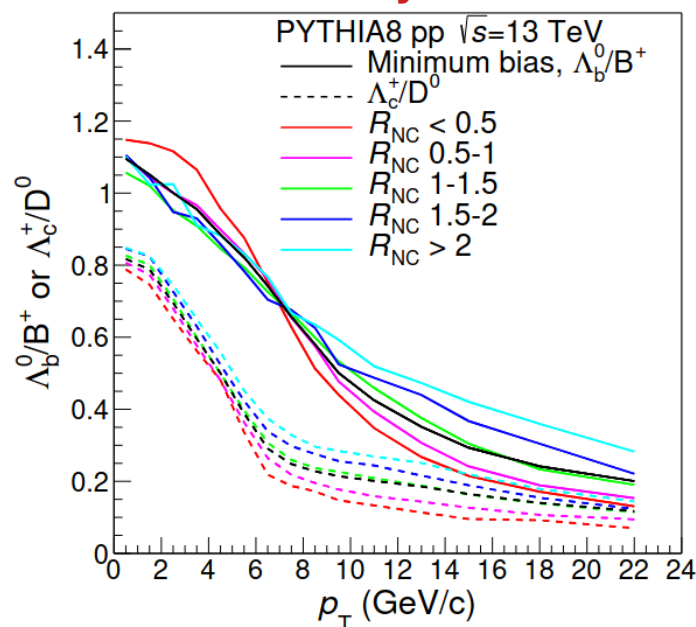
- **CR-BLC overall outperform other models for charm, however overestimates beauty production.**
- New high-luminosity data from LHC Run 3 data can further constrain heavy-flavor fragmentation mechanisms.

# $\Lambda_b^0/B^+$ yield for triggered events

## UE activity classifier



## Jet activity classifier



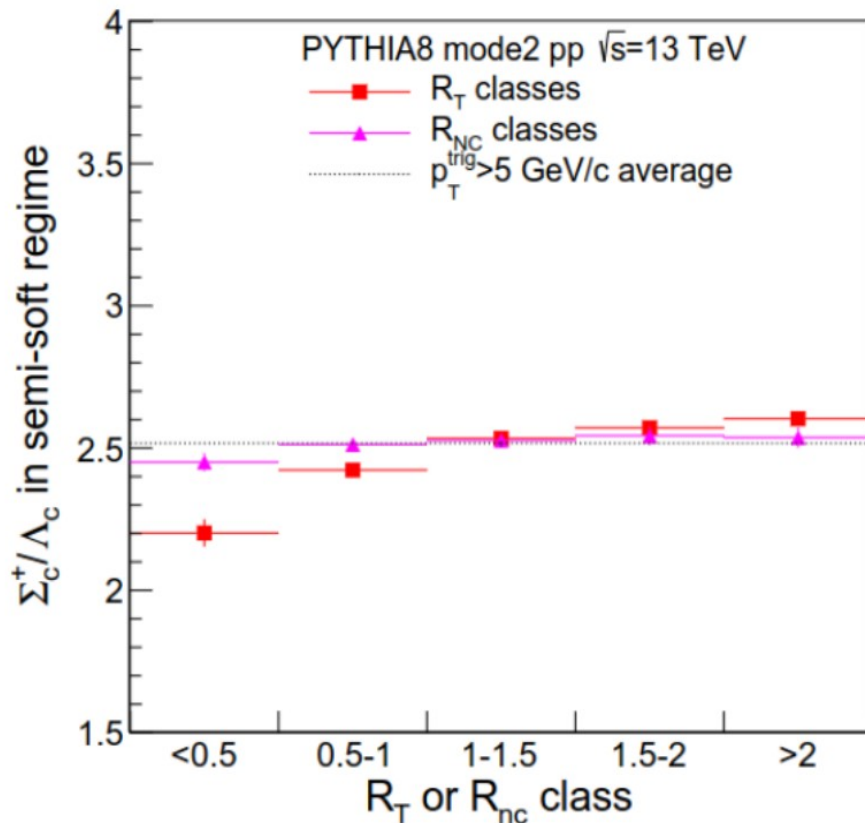
- Similar conclusions for beauty baryon enhancement as for charm.
- Significant difference is observable in case of  $R_T$  (UE activity).
- No significant difference when classified by  $R_{NC}$  classes (jet activity).

# Production of Excited Charm and Charm-strange Baryon States

## Previous studies extended for several excited charm baryon states:

- investigating the production of charmed baryons with different **isospin** and **strangeness** content,
- comparison to both charmed  $D^0$  mesons and  $\Lambda_c^+$  baryons in pp collisions at LHC energies.
- Conclusion 1: **Strangeness content has no further sensitivity to event-property descriptors.**
- Conclusion 2: **Charm enhancement driven by different mechanism than strange enhancement.**
- Conclusion 3: **The isospin of the charmed-baryon state has a strong impact on the enhancement.**

Z. Varga, R. Vértési, A. Misák, J. Phys. G: Nucl. Part. Phys. 50 (2023) 075002



# Summary

- **KNO scaling:** Event multiplicity distributions scale with center-of-mass energy.
    - Violated at higher energies, and not fully understood.
  - Recent phenomenological studies: **a KNO-like scaling may be present within jets.**
    - Jet multiplicity distributions are NBD and can be collapsed onto a universal distribution.
    - KNO scaling is likely violated by complex QCD processes outside the jet development, such as single and double-parton scatterings or softer MPI.
  - **Testing for this scaling behavior can be an important element in model development.**
  - **In CR-BLC model the charm baryon enhancement originates in the underlying event** → the processes inside the jets do not play a significant role.
    - Good quantitative description of both charm and beauty baryon enhancement with the same model settings is still missing, but **demonstrated a good sensitivity to certain event-activity observables.**
    - Utilizing jet substructure measurements may provide more information on the hadronization.
- LHC Run 3 data with increased statistics will be essential.**

**Thank you!**