

Quantifying Jet Quenching And Medium Response With Two Particle Correlations With PHENIX

Anthony Hodges for the PHENIX Collaboration

Hot Jets: Advancing the Understanding of High
Temperature QCD With Jets

January 8th, 2024



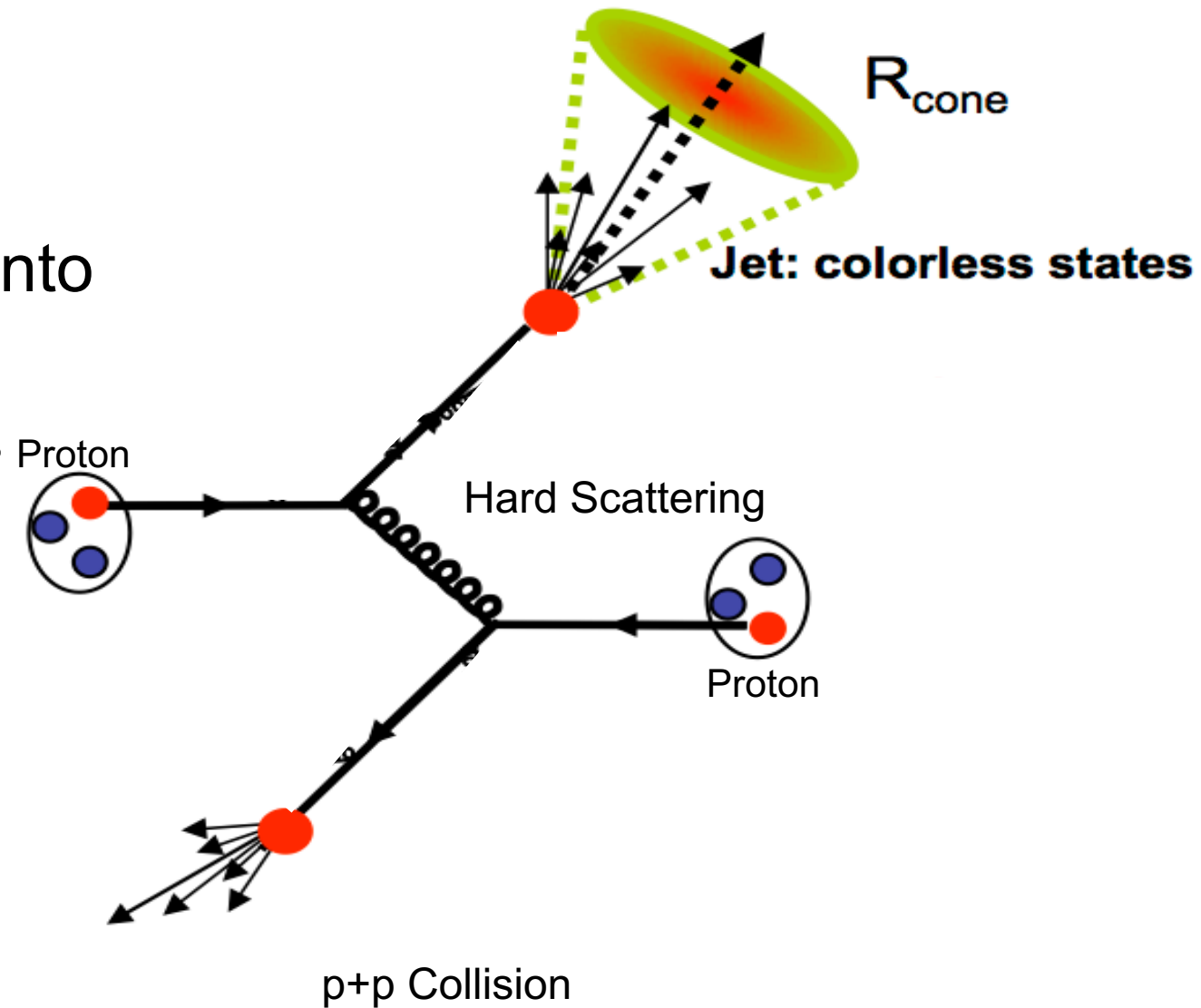
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NSF Ascend Fellow
#2316651
#1848162

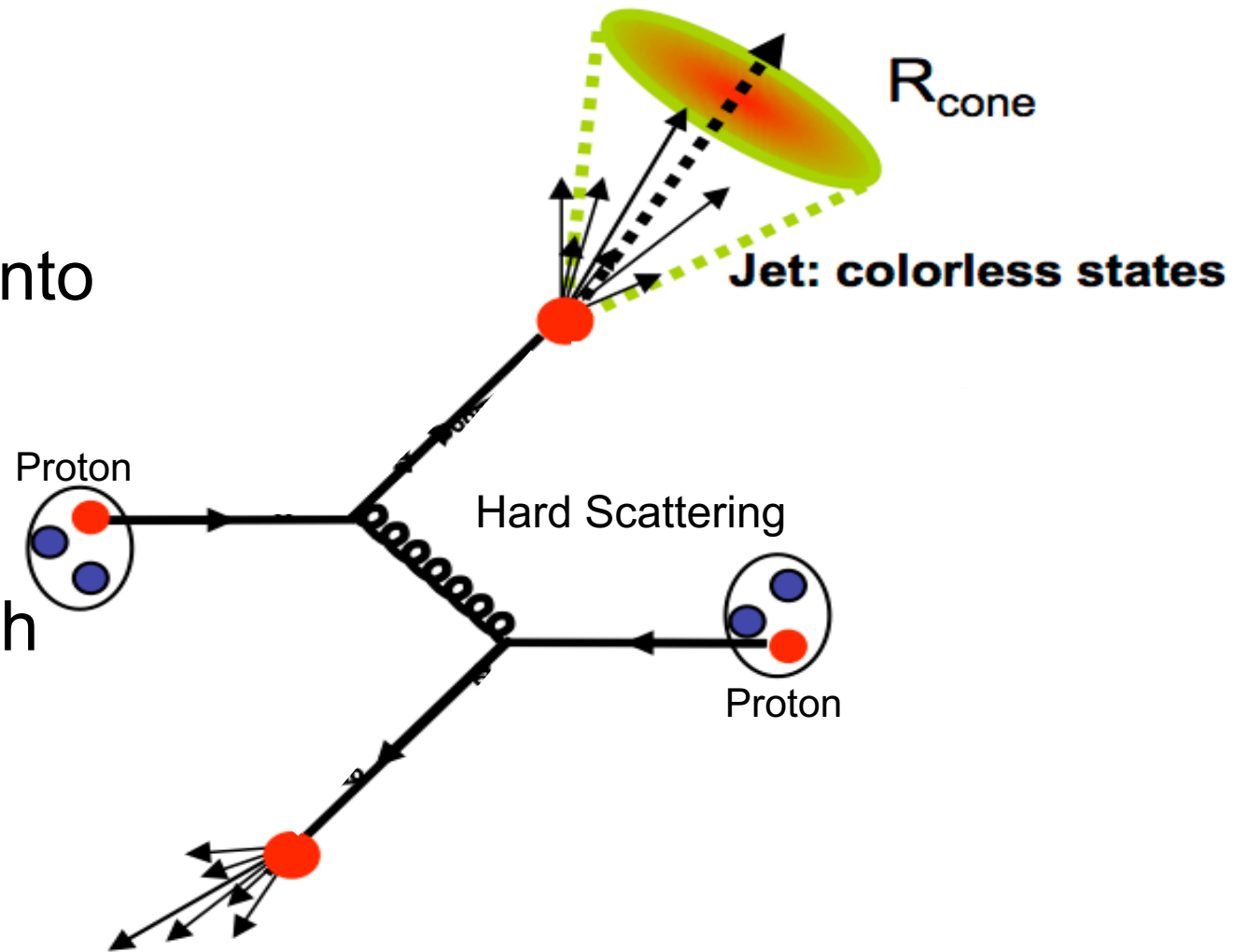
Jets as QCD Probes

- Partons undergo hard scattering
 - Large momentum transfer (q^2)
- Hard-scattered partons hadronize into jets
- Jet kinematics \approx parton kinematics good proxy



Jets as QCD Probes

- Partons undergo hard scattering
 - Large momentum transfer (q^2)
- Hard-scattered partons hadronize into jets
- Jet kinematics \approx parton kinematics good proxy
- Jets are accessible by pQCD, which makes them very theory-friendly

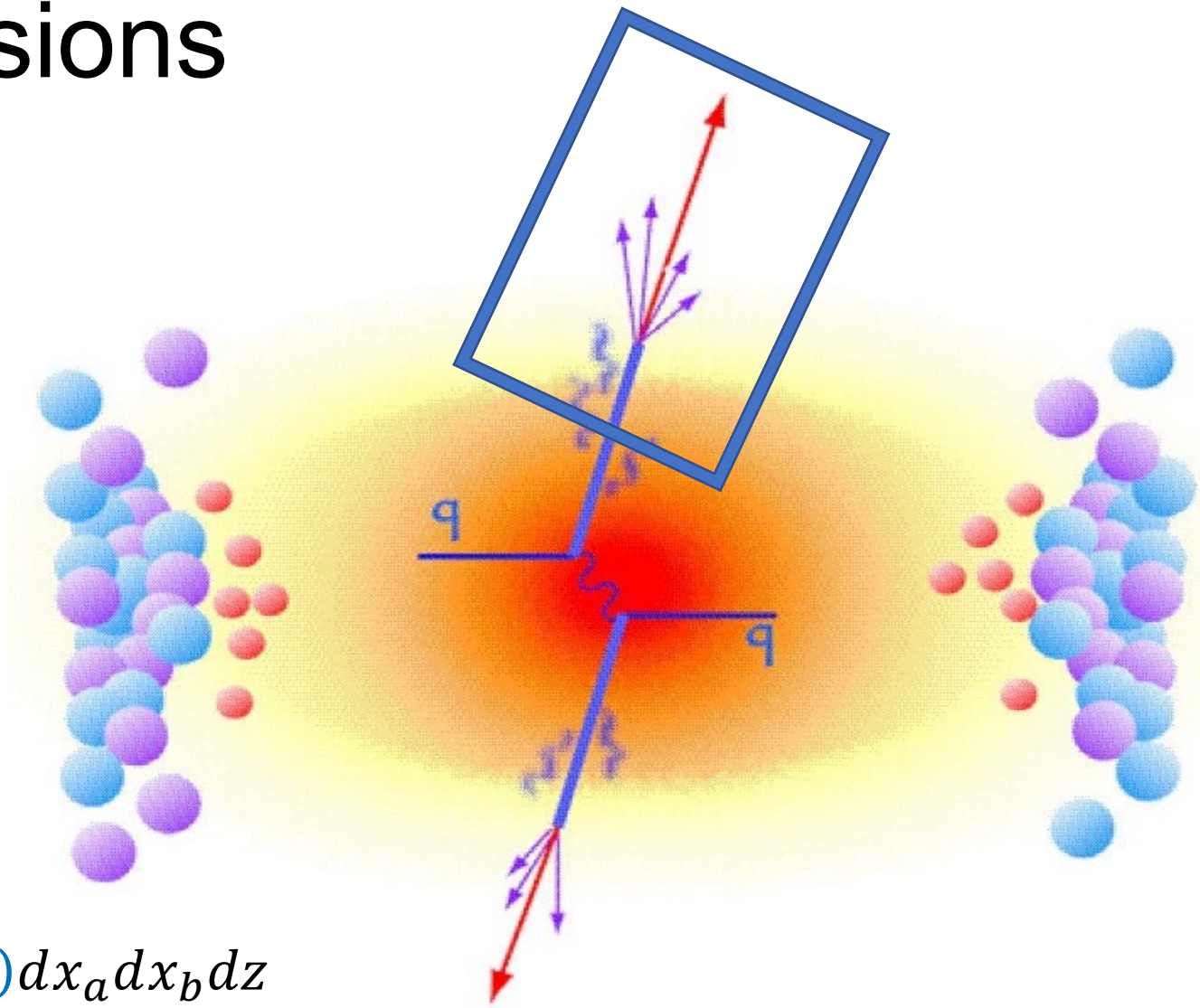


Differential cross section

$$d\sigma = \int \int \int f_a^A(x_a) f_b^B(x_b) \cdot d\sigma_{ab \rightarrow cX} \cdot D_c^h(z) dx_a dx_b dz \quad \text{p+p Collision}$$

Jets in Heavy Ion Collisions

- Hard-scattered partons now traverse dense QGP medium
- Energy lost to medium yields modified jets
- Measurement of jet modification allows us to study parton-medium interactions



Differential cross section

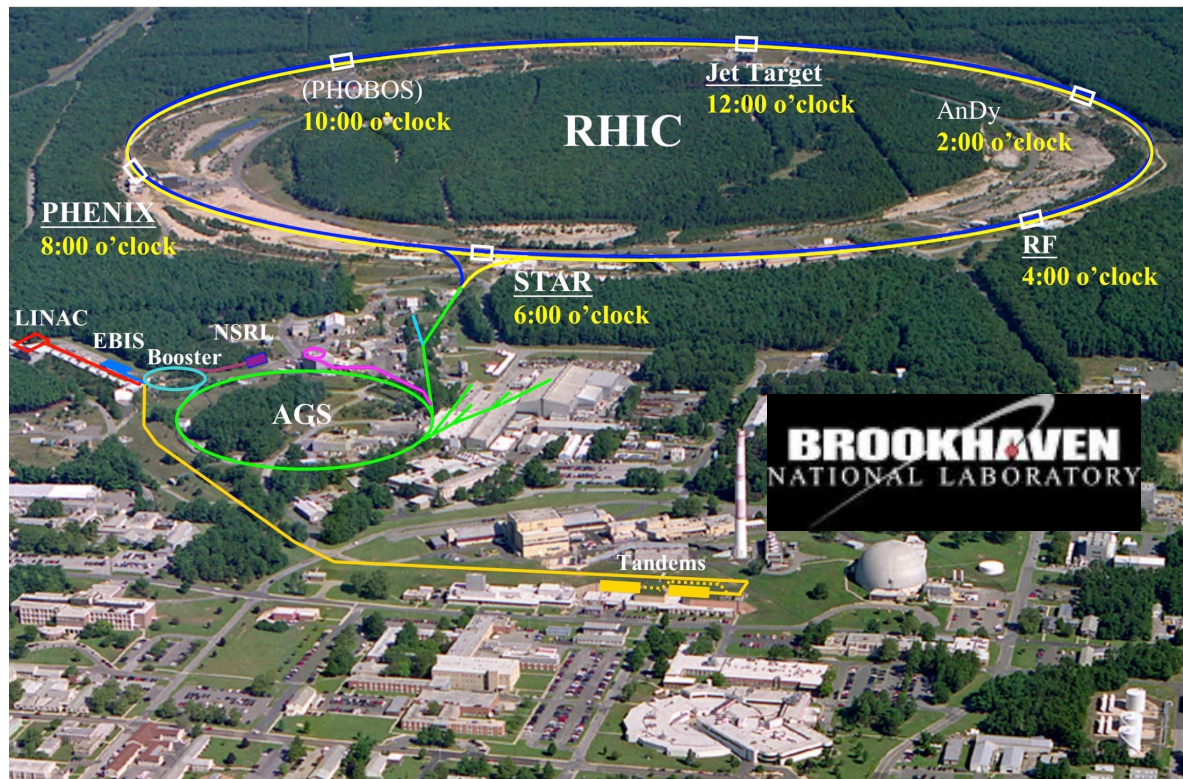
$$d\sigma = \int \int \int f_a^A(x_a) f_b^B(x_b) \cdot d\sigma_{ab \rightarrow cX} \cdot D_c^h(z) dx_a dx_b dz$$

Modification

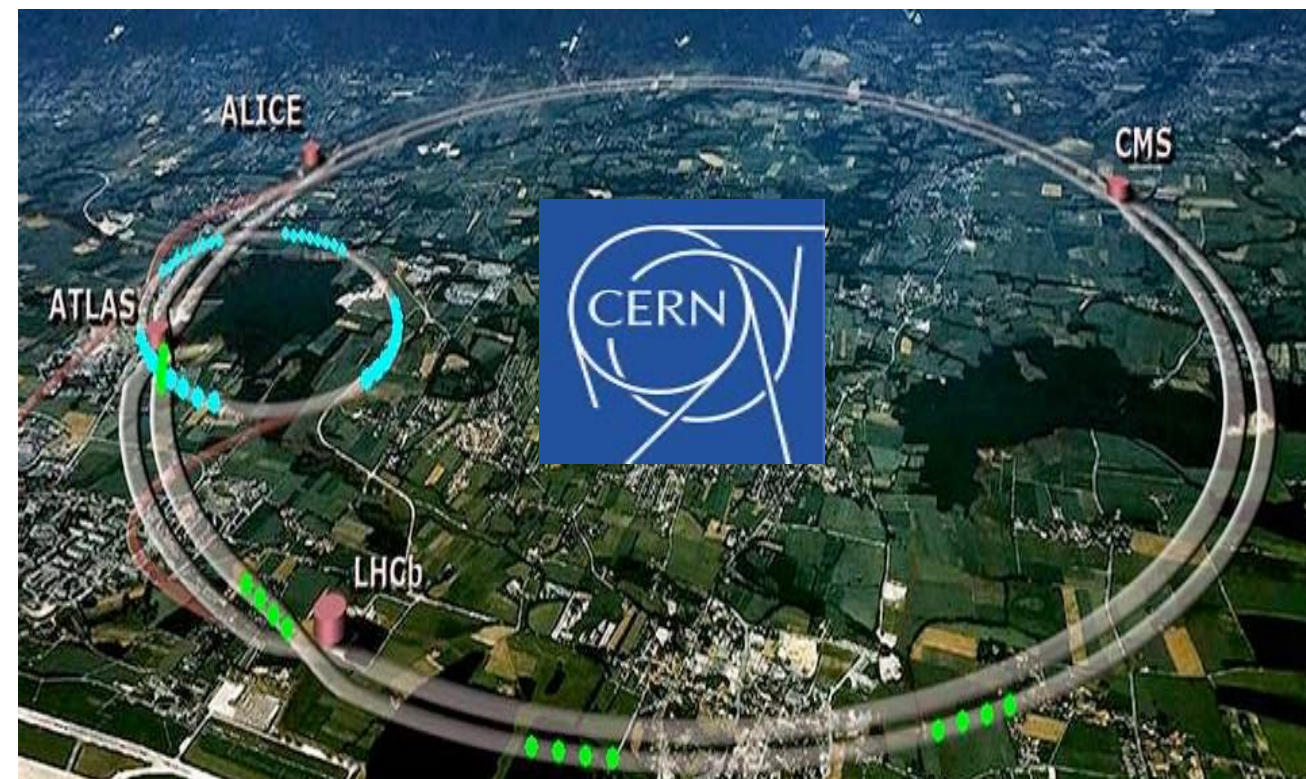
$$D_c^h(z) \rightarrow D_c^h(z') \text{ as } p^{\text{Parton}} \rightarrow p'^{\text{Parton}}$$

Heavy Ion Collision

Jet Factories: RHIC and the LHC

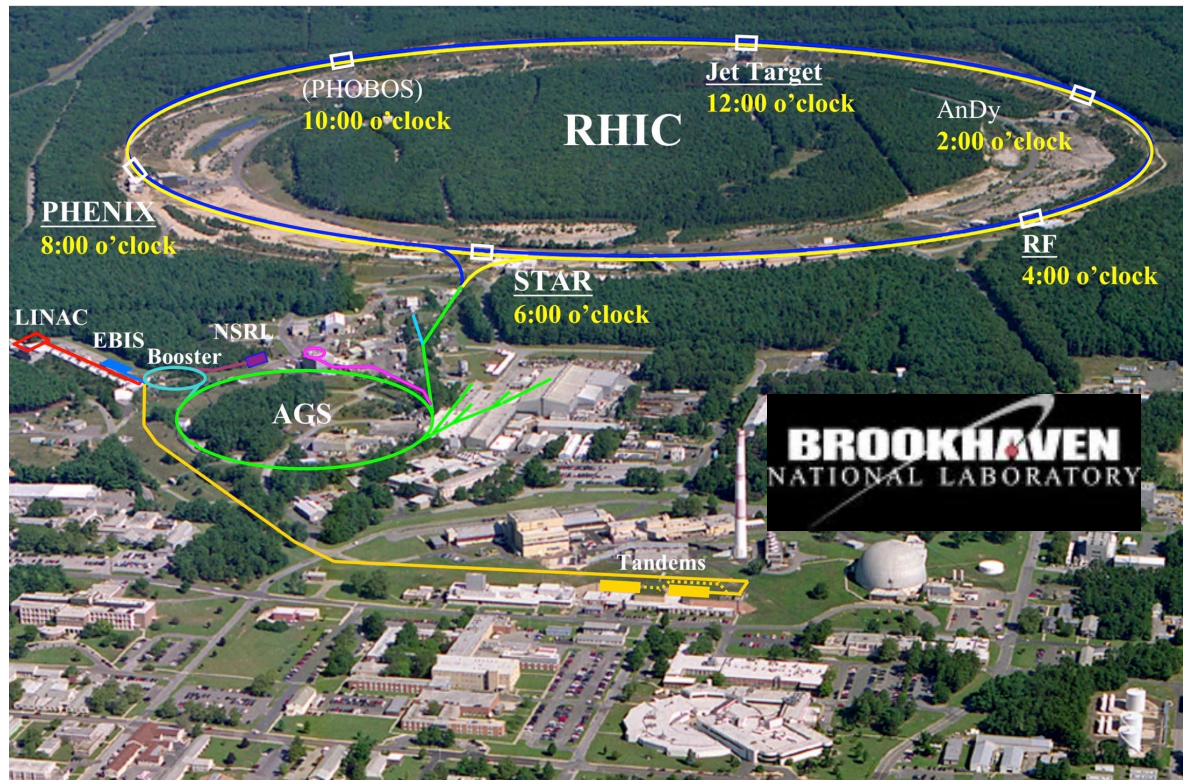


Relativistic Heavy Ion Collider



Large Hadron Collider

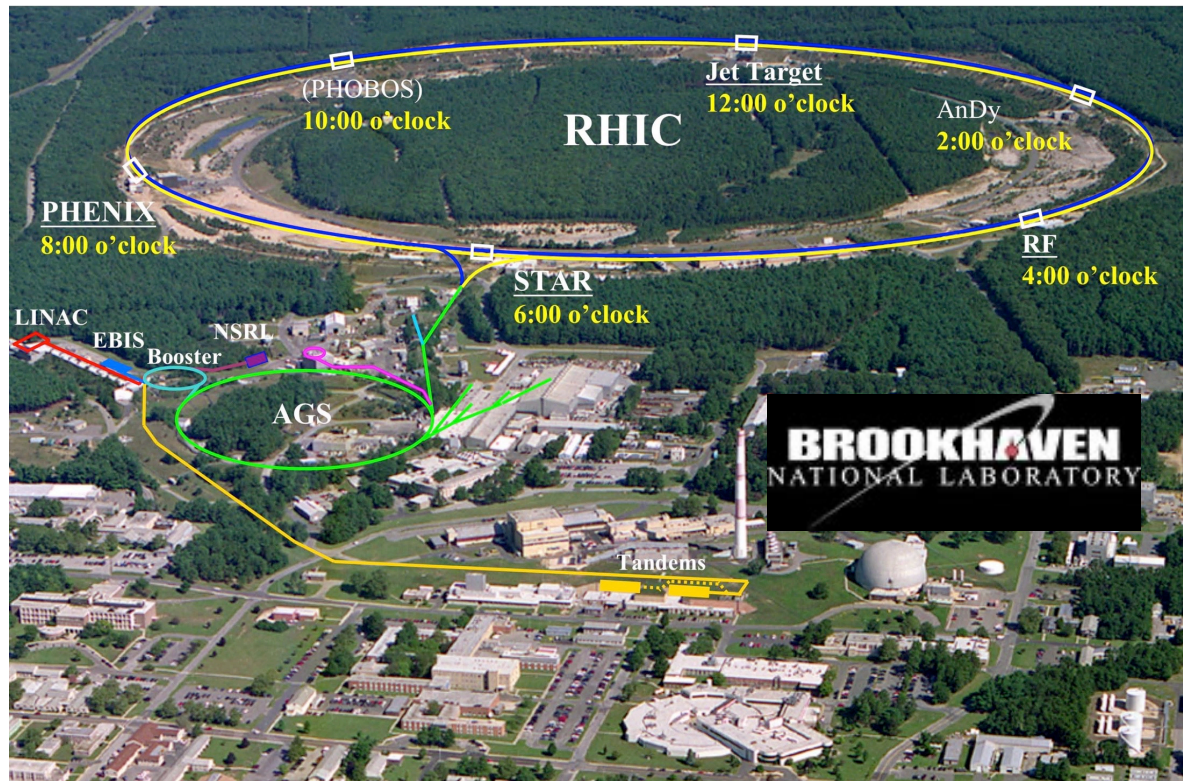
The Relativistic Heavy-Ion Collider



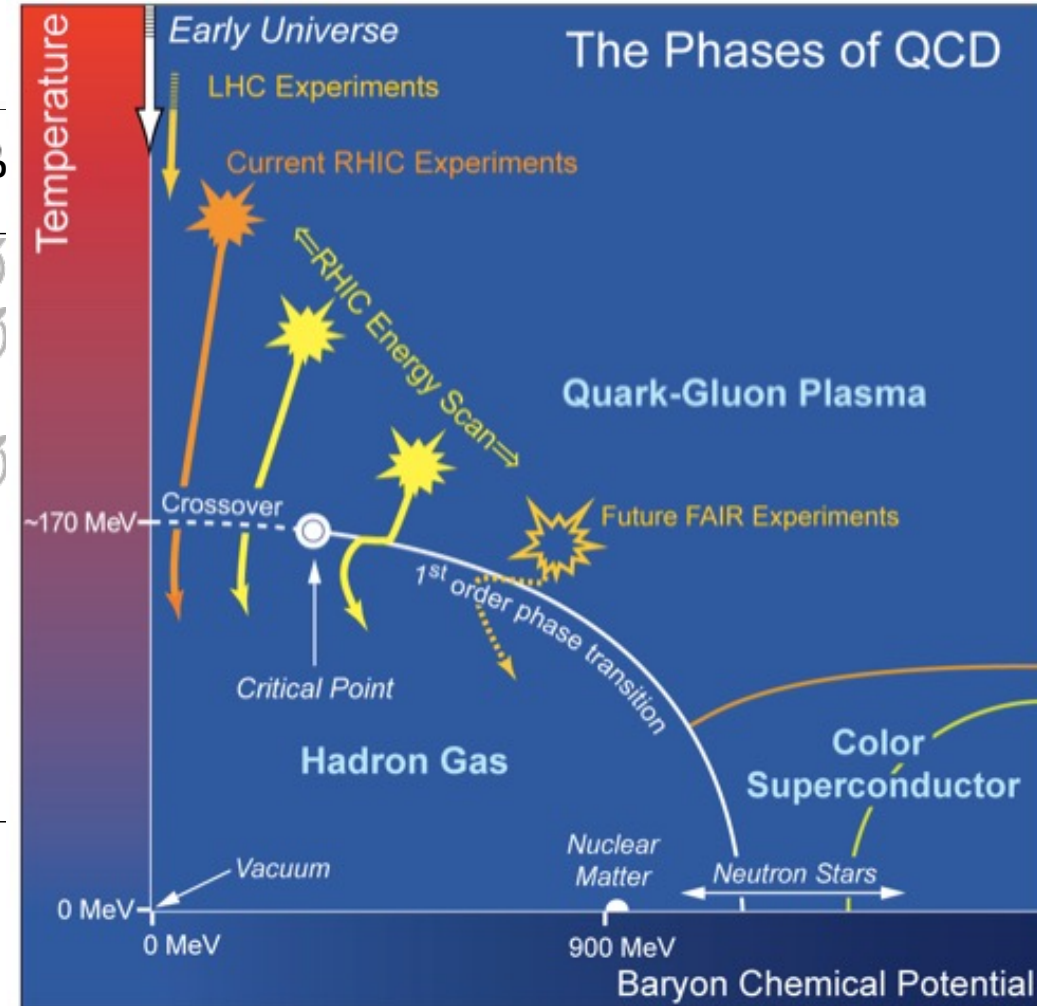
\sqrt{s} [GeV]	p+p	p+Al	p+Au	d+Au	³ He+Au	Cu+Cu	Cu+Au	Au+Au	U+U
510	✓								
200	✓	✓	✓	✓	✓	✓	✓	✓	✓
130								✓	
62.4	✓			✓		✓		✓	
39				✓				✓	
27				✓				✓	
20				✓		✓		✓	
14.5								✓	
7.7								✓	

- Highly versatile collider
- Decades of expansive data collected

The Relativistic Heavy-Ion Collider



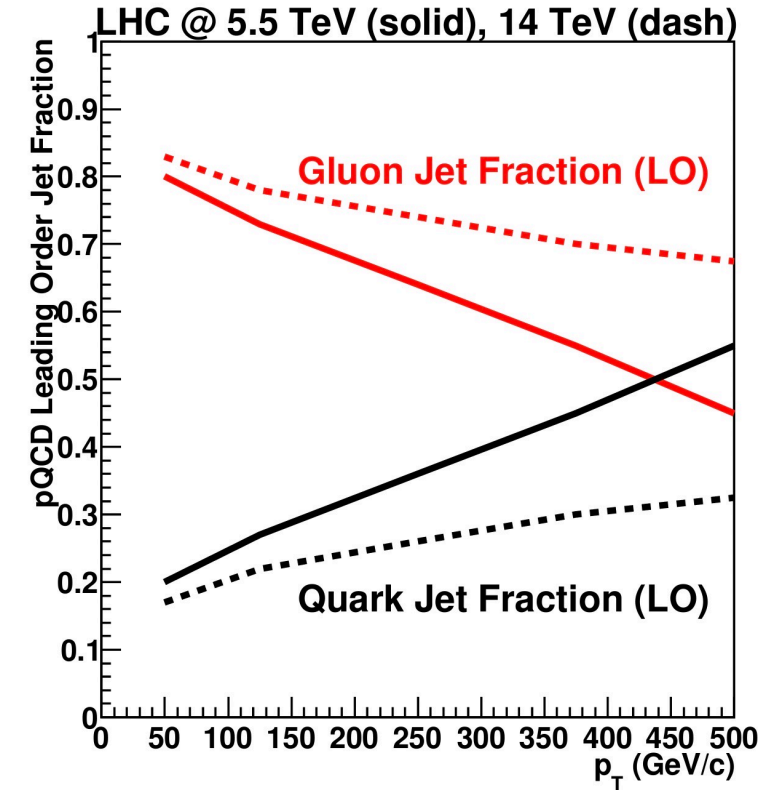
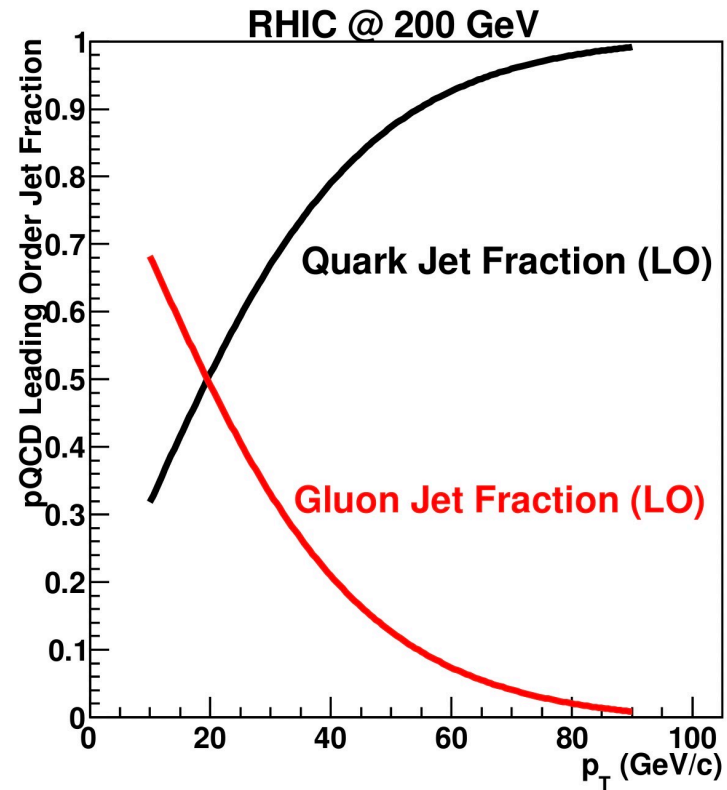
\sqrt{s} [GeV]	$p+p$
510	<input checked="" type="checkbox"/>
200	<input checked="" type="checkbox"/>
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62.4	<input checked="" type="checkbox"/>
39	<input type="checkbox"/>
27	<input type="checkbox"/>
20	<input type="checkbox"/>
14.5	<input type="checkbox"/>
7.7	<input type="checkbox"/>



- Capable of probing large swathe of QCD phase space!

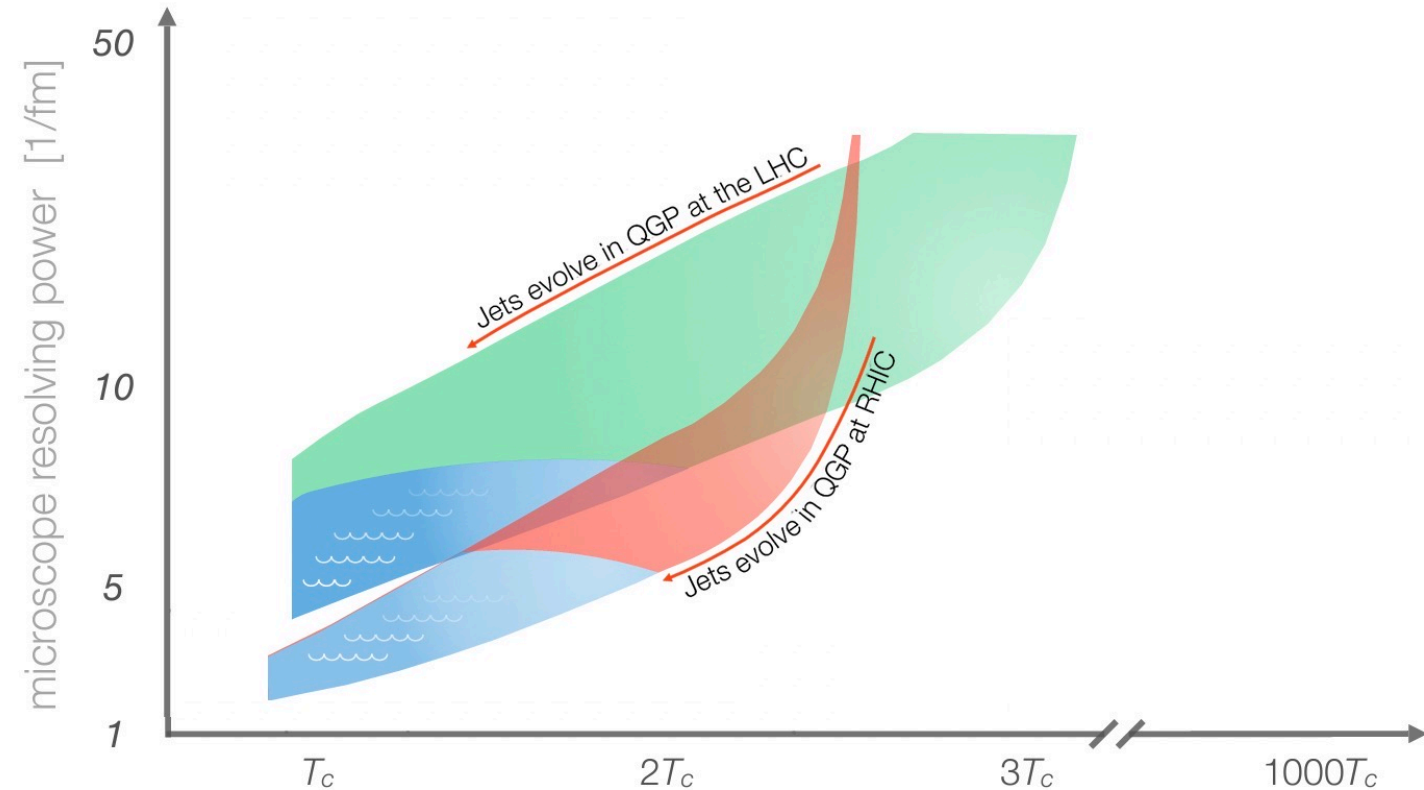
Why Jets at RHIC?

- Different jet populations!
- LHC jets largely initiated by gluons
- RHIC jets dominated by quark-initiated jets

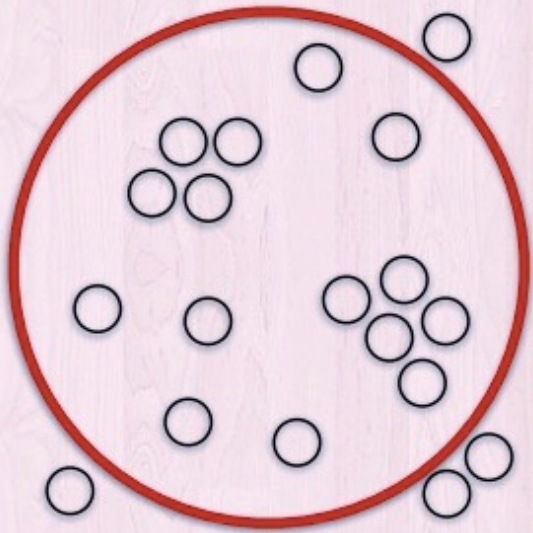


Why Jets in Heavy-Ion Collisions at RHIC?

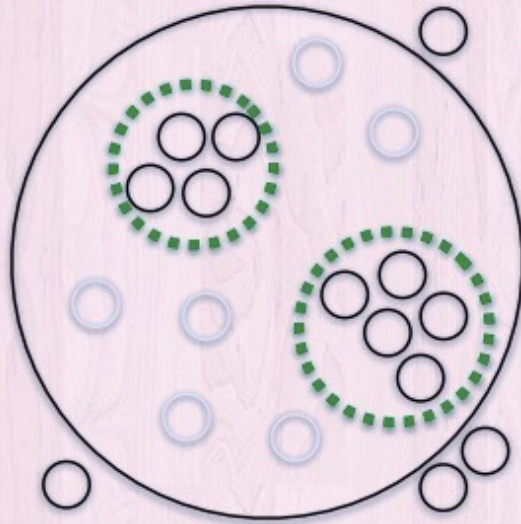
- Jet-medium interactions at RHIC and the LHC have subtle differences
- Jet energies at RHIC typically closer to medium energy scale
- RHIC jets also spend larger part of their evolution in the medium than LHC jets



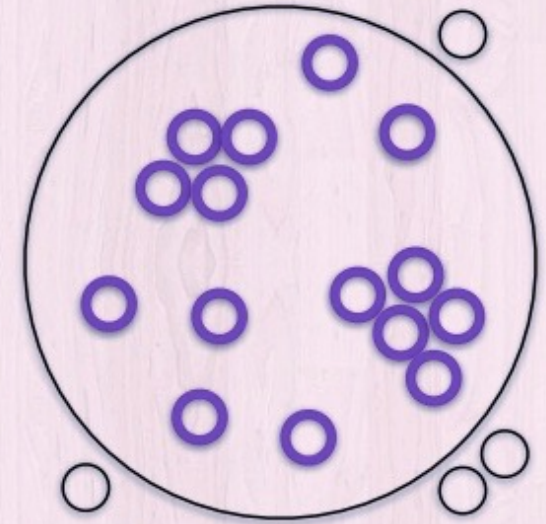
How We Study Jets



Full jet



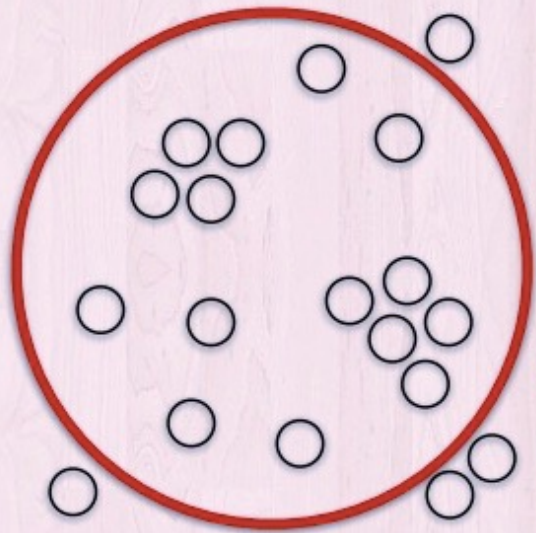
Intermediate structure



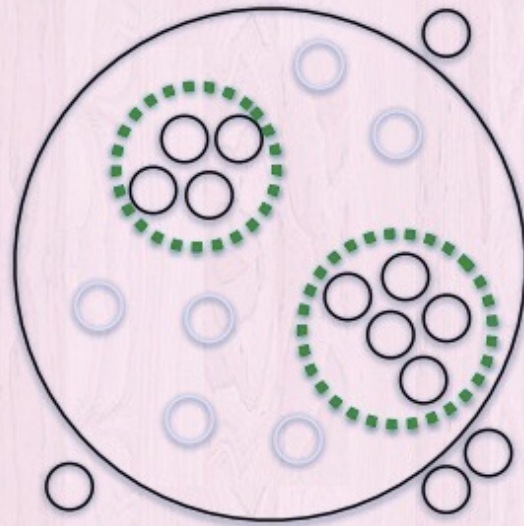
Constituent

PC: [Yi Chen, QM2019](#)

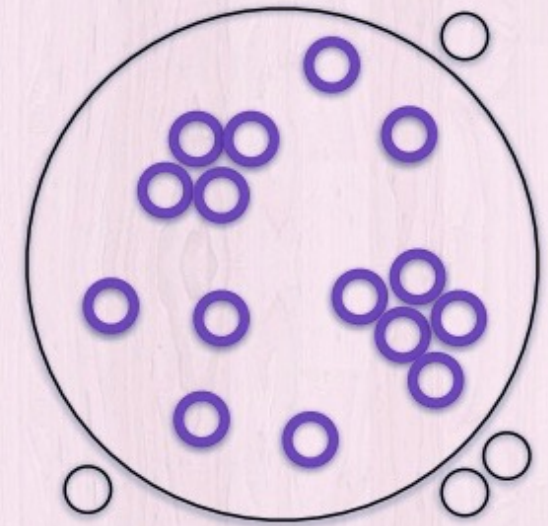
How We Study Jets



Full jet



Intermediate structure

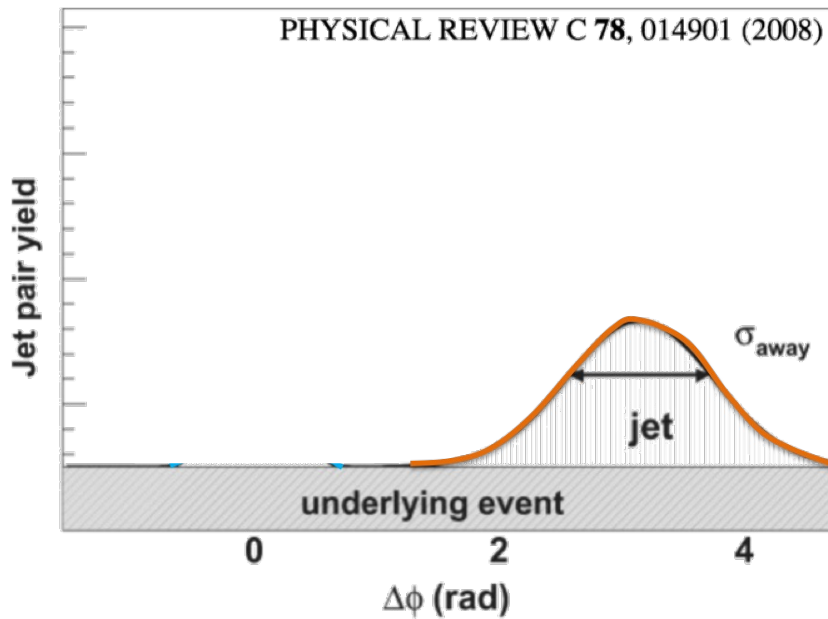


Constituent

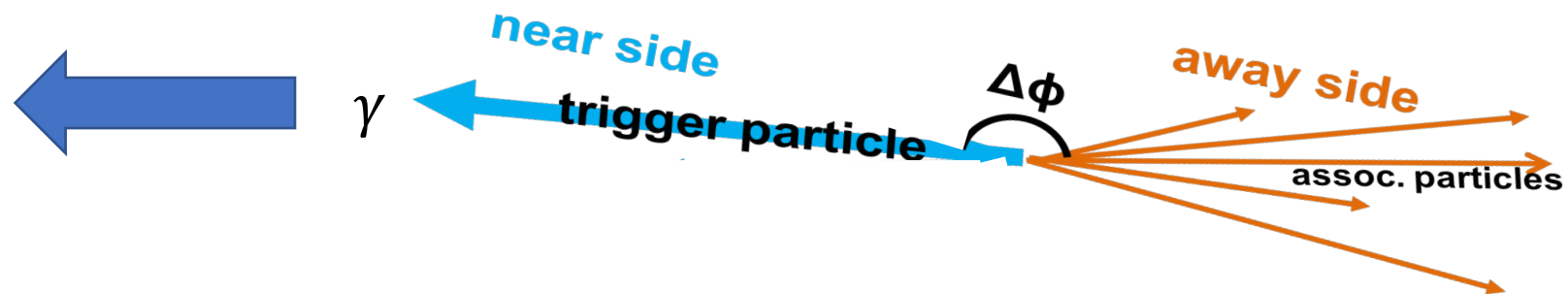
PC: [Yi Chen, QM2019](#)

Measuring Jets – Two Particle Correlations

- Direct photons: colorless, well-calibrated probe, but statistically limited



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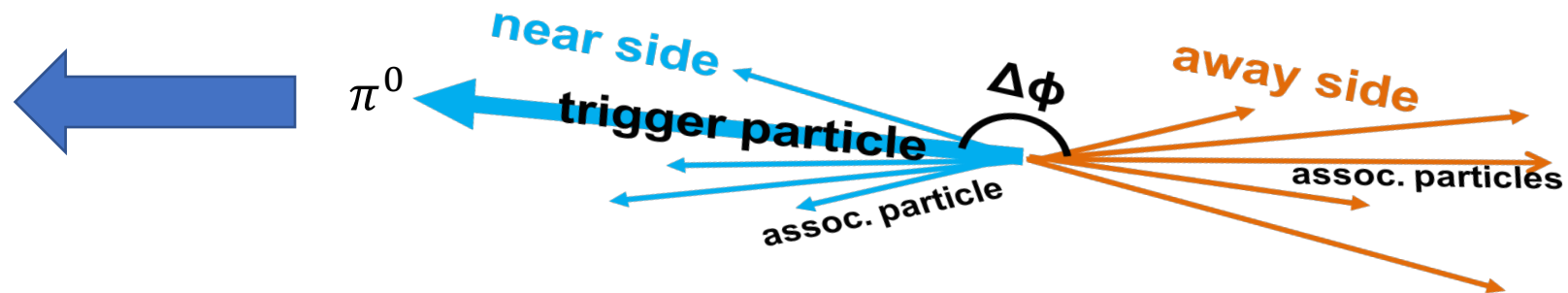
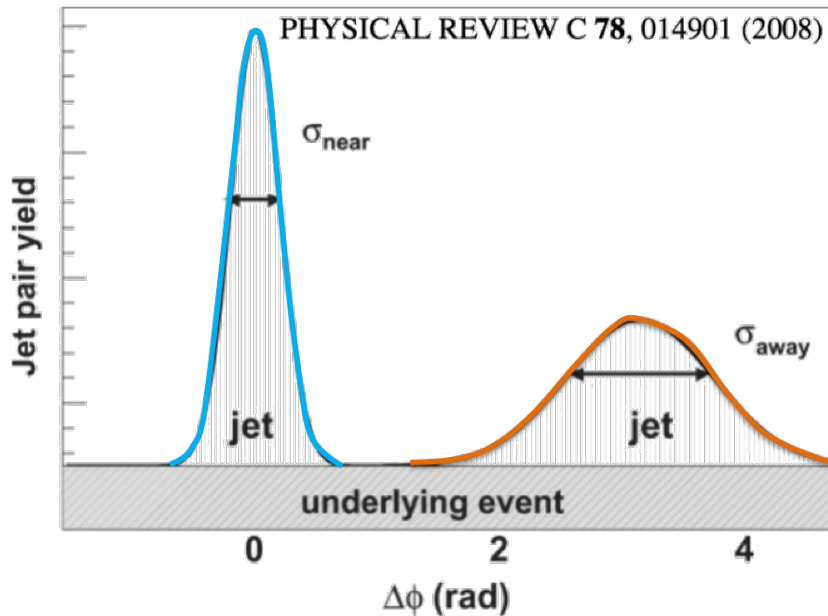


Anthony Hodges, NSF Ascend Fellow, UIUC

12

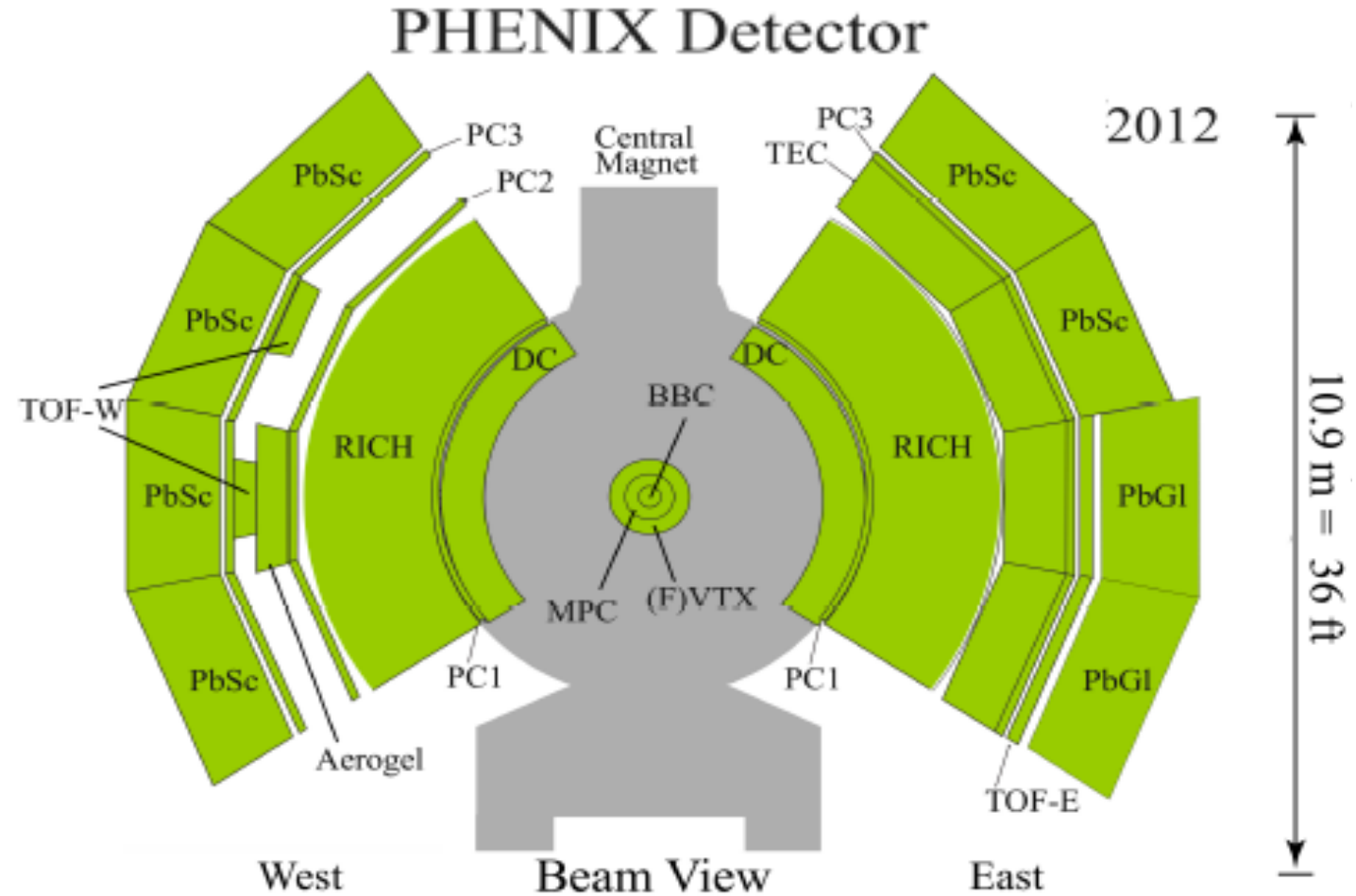
Measuring Jets – Two Particle Correlations

- Direct photons: colorless, well-calibrated probe, but statistically limited
- Neutral pion (π^0): not colorless, but abundant
 - Good for high-precision, differential measurements



The PHENIX Detector

- Central arms
 - π coverage in azimuth
 - Pseudorapidity coverage of $|\eta| < 0.35$
- Electromagnetic calorimeter
 - Photon and electron energy
- Drift/Pad chambers
 - Charged hadron momentum
- Beam-beam counters (BBC)
 - Event characterization



Two-Particle Correlation Anatomy

- $$\frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

Two-Particle Correlation Anatomy

- $$\frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

Per-trigger (jet) yield of hadrons (what we want to measure)

Two-Particle Correlation Anatomy

$$\bullet \frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

What we actually measure, the correlation function

Two-Particle Correlation Anatomy

$$\bullet \frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

Correction for detector inefficiencies

Two-Particle Correlation Anatomy

$$\bullet \frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

Correction for detector effects

Two-Particle Correlation Anatomy

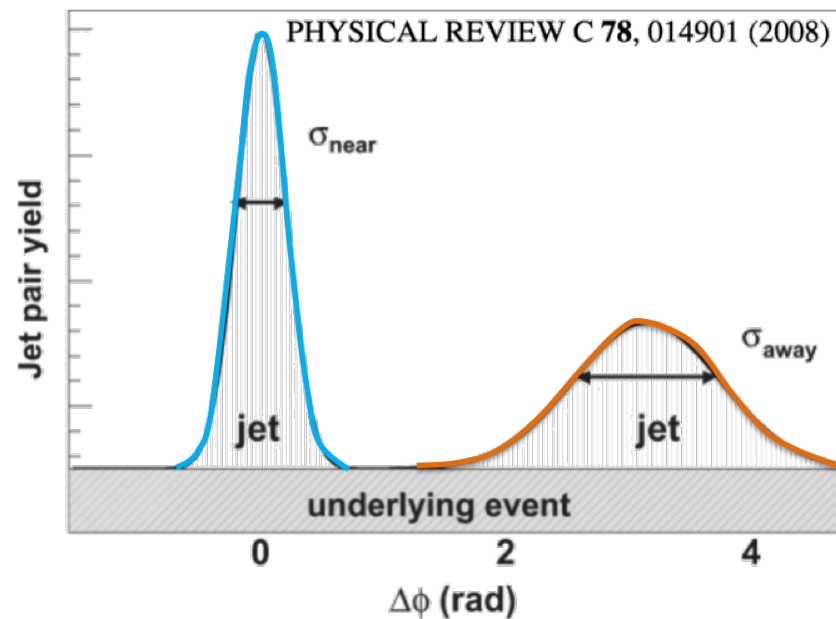
$$\bullet \frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

Underlying event subtraction

Two-Particle Correlation Anatomy

$$\bullet \frac{1}{N_{Trig}} \frac{dN^{Pair}}{d\Delta\phi} = \frac{1}{N_{Trig}} \frac{N^{Pair}}{\epsilon^{Hadron} \int \Delta\phi} \left\{ \frac{dN_{Real}^{Pair}/d\Delta\phi}{dN_{Mix}^{Pair}/d\Delta\phi} - b(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi)) \right\}$$

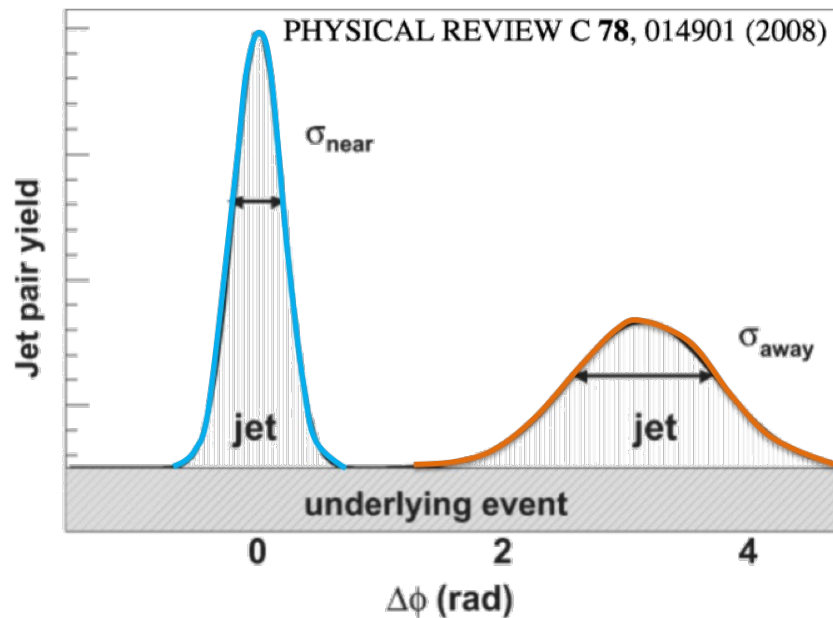
In p+p, nice and flat



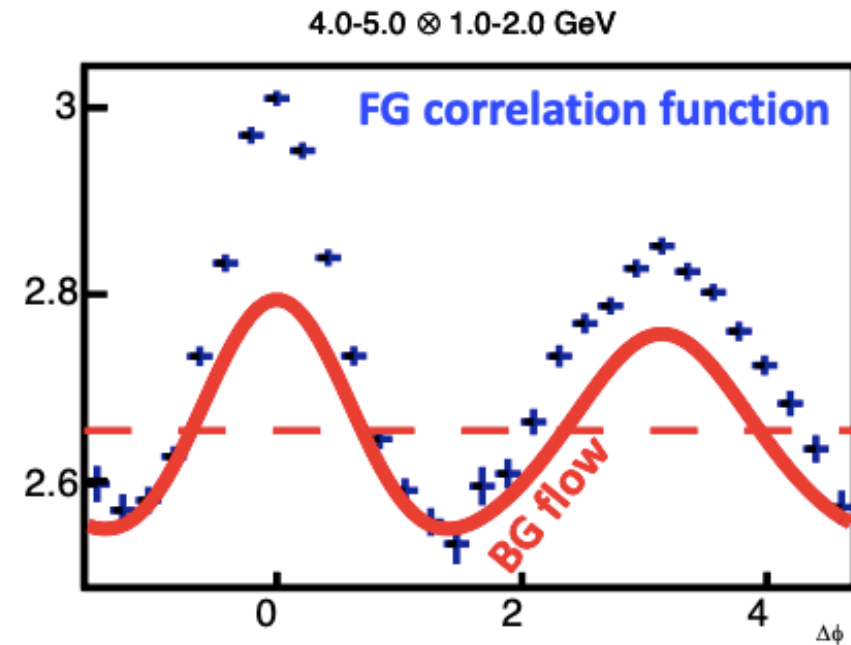
Two-Particle Correlation Anatomy

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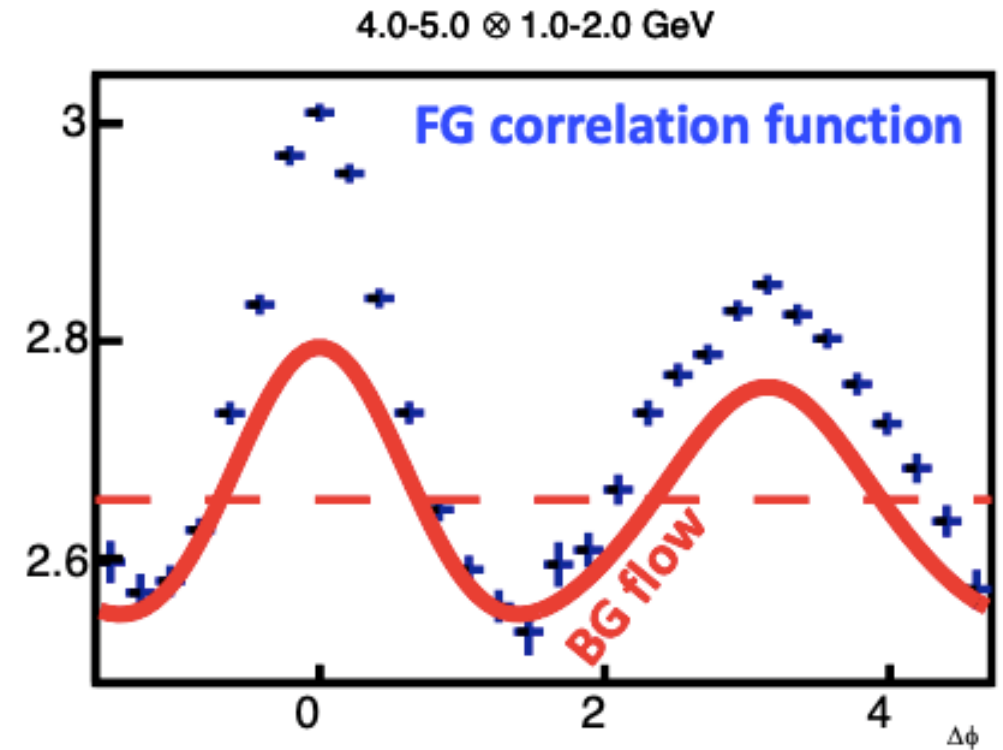
in A+A...



Underlying Event Subtraction

$$\frac{dN}{d\Delta\phi} = \mathbf{b} \left(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi) \right)$$

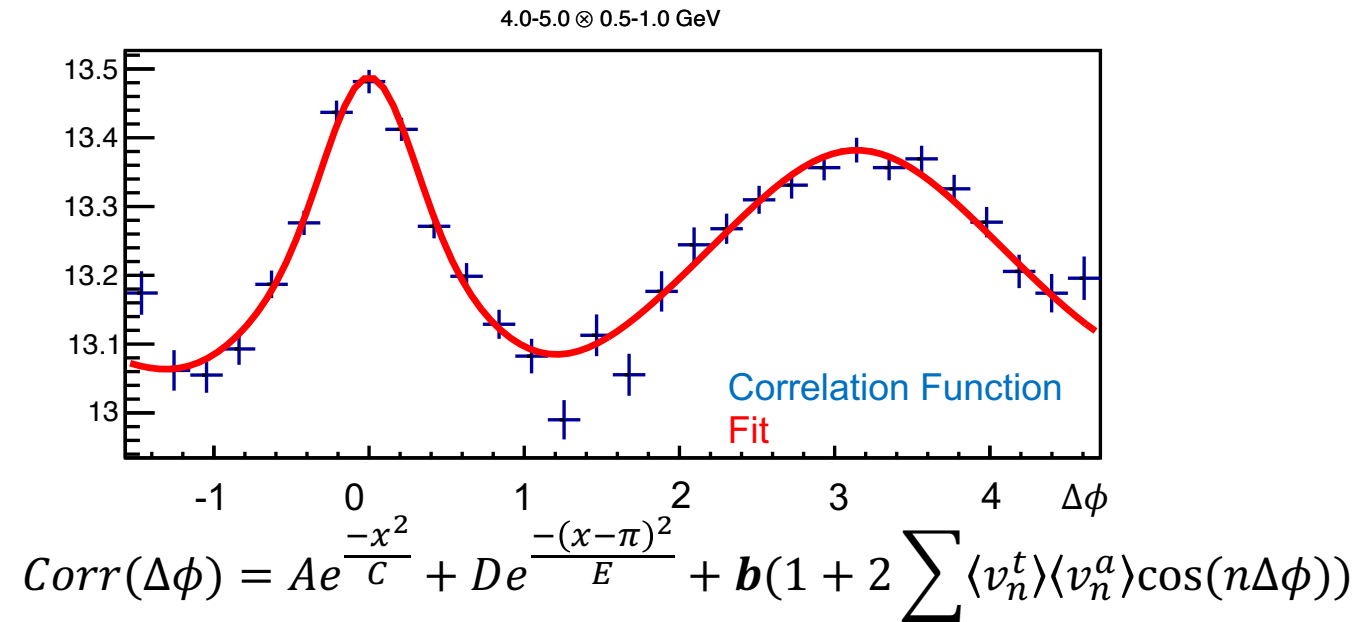
$$b = \frac{\xi \langle N_{Trig} \rangle \langle N_{h^\pm} \rangle}{\langle N_{Pairs} \rangle}$$



- Amplitude given by Absolute Background Subtraction method for $p_T^{Hadron} > 1\text{GeV}/c$
 - Phys. Rev. C **81**, 014908

Underlying Event Subtraction

$$\frac{dN}{d\Delta\phi} = \mathbf{b} \left(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi) \right)$$

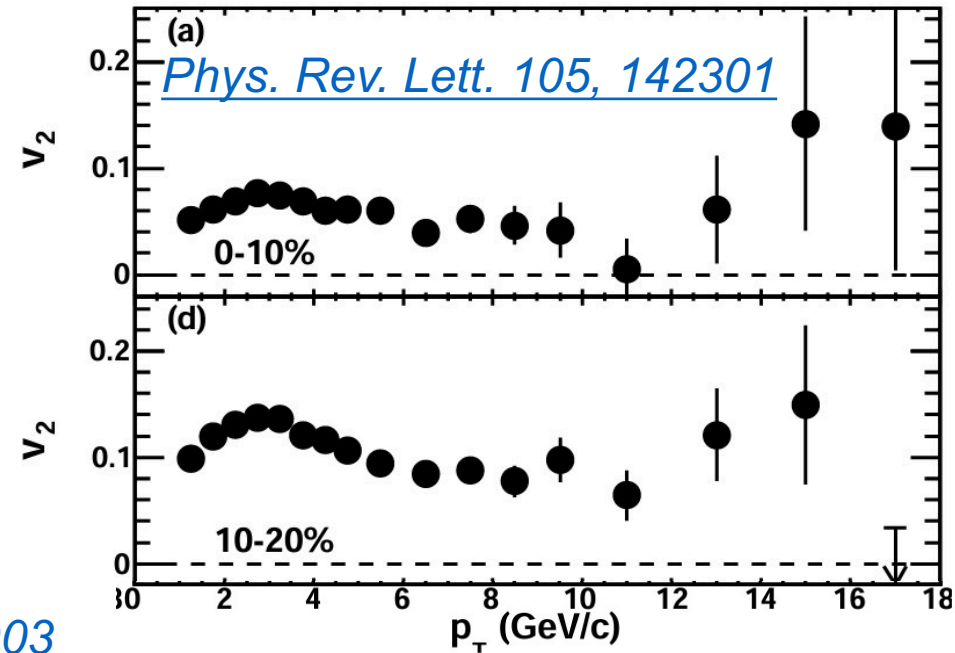


- b from ZYAM (Zero Yield At Minimum) for $p_T^{Hadron} < 1\text{GeV}/c$ to account for over-subtraction

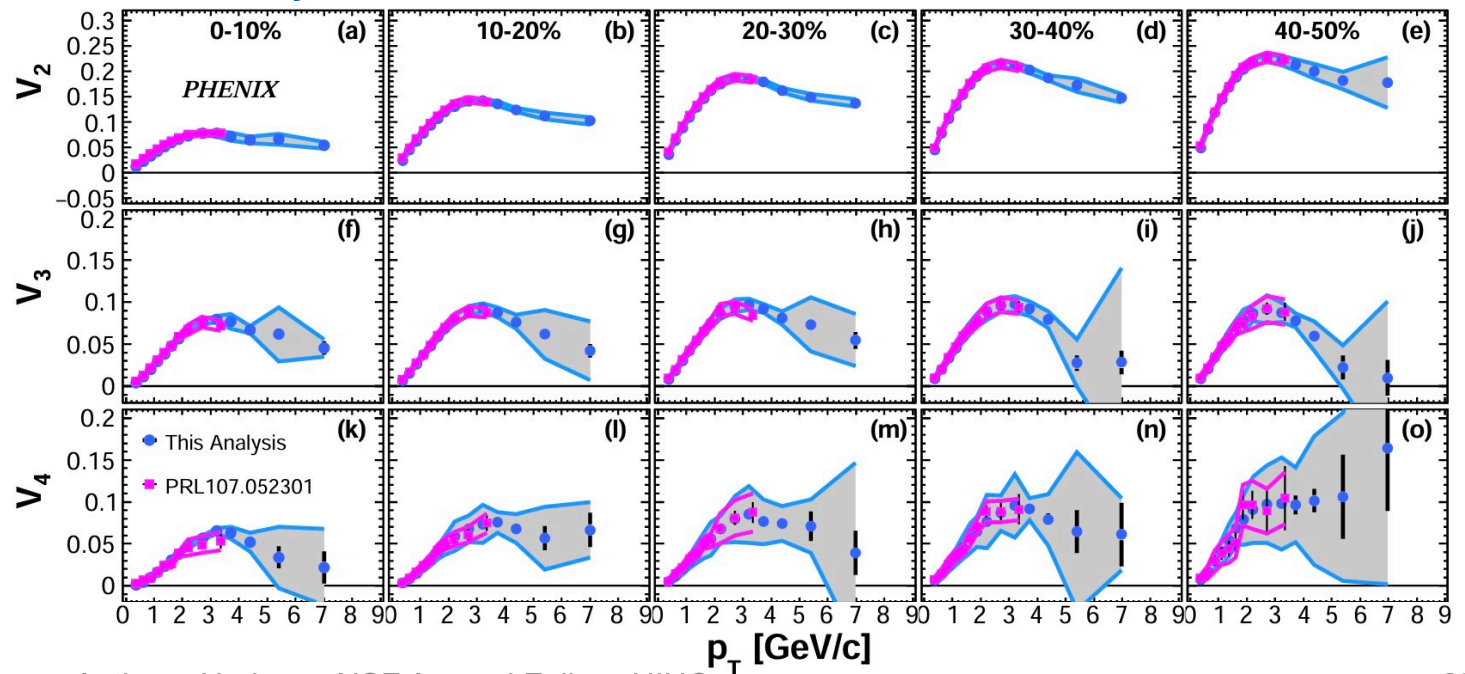
Underlying Event – Flow

$$\frac{dN}{d\Delta\phi} = b \left(1 + 2 \sum \langle v_n^t \rangle \langle v_n^a \rangle \cos(n\Delta\phi) \right)$$

- v_n terms quantify background shape, come from previous PHENIX analyses



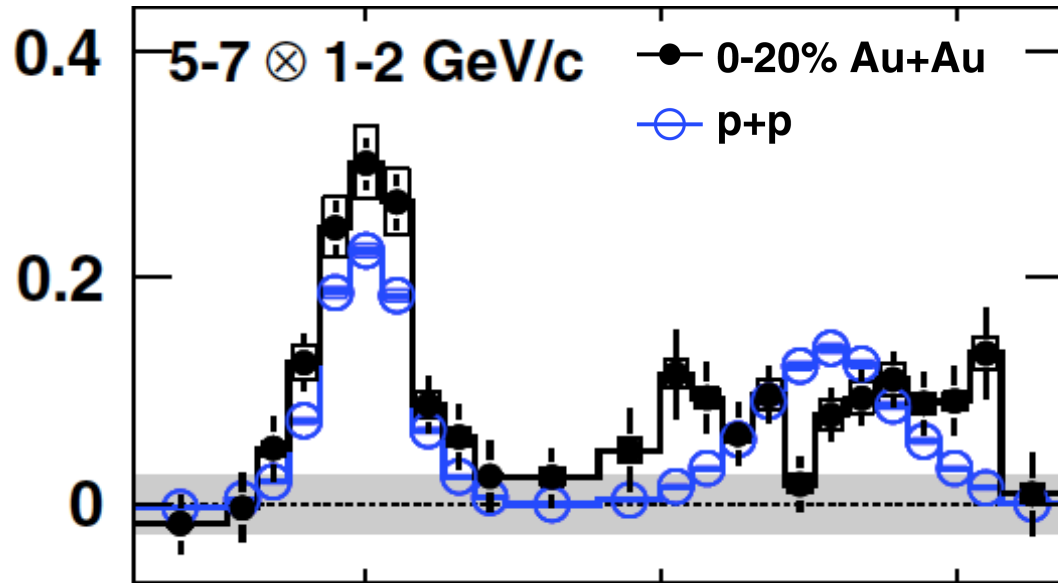
Phys. Rev. C 99, 054903



Improved Background Subtraction

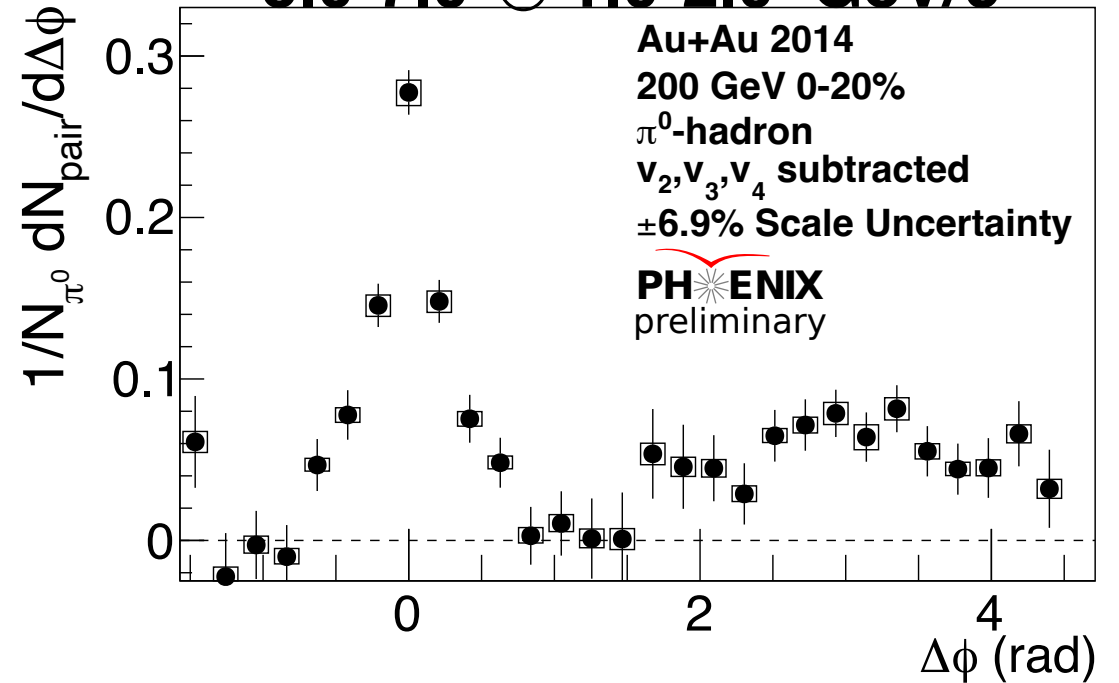
Previous

PRL 104 252301 (2010)



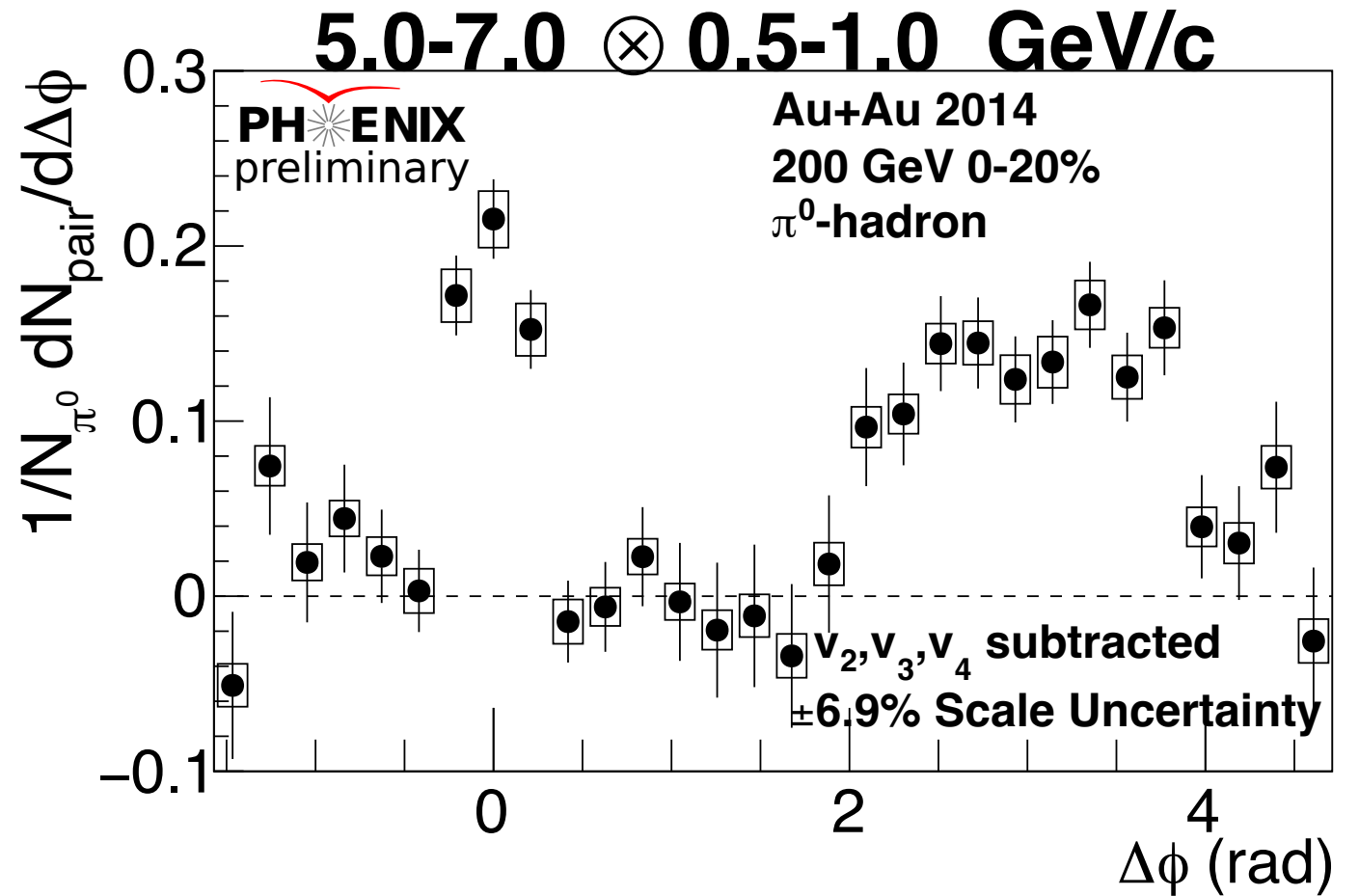
New

5.0-7.0 \otimes 1.0-2.0 GeV/c



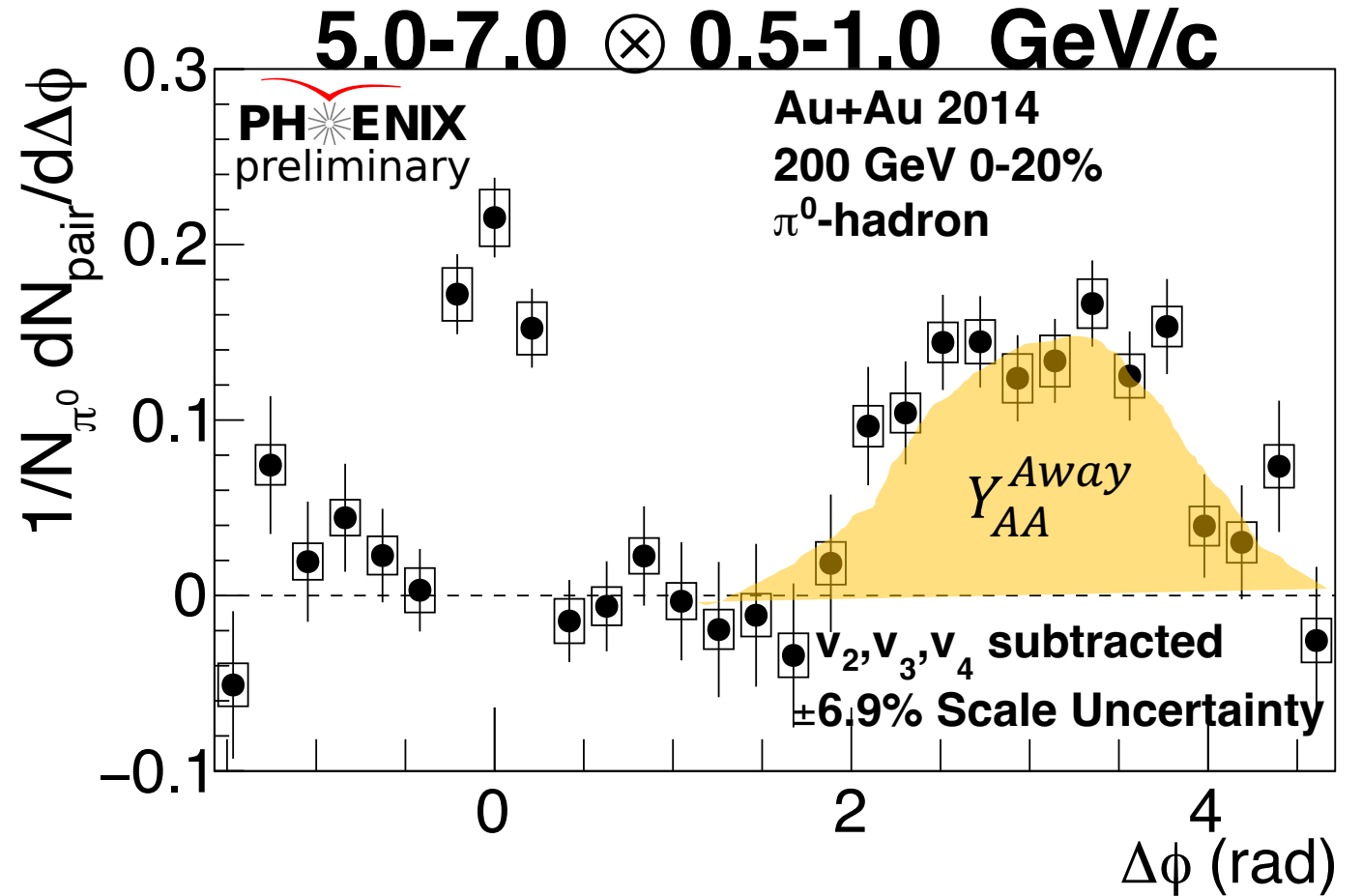
- New results (right) from data from 2014, 3.3x more trigger π^0 's
- v_2, v_3 , and v_4 subtracted in new results \rightarrow flow contamination removed

Jet Modification: $I_{AA}(p_T)$



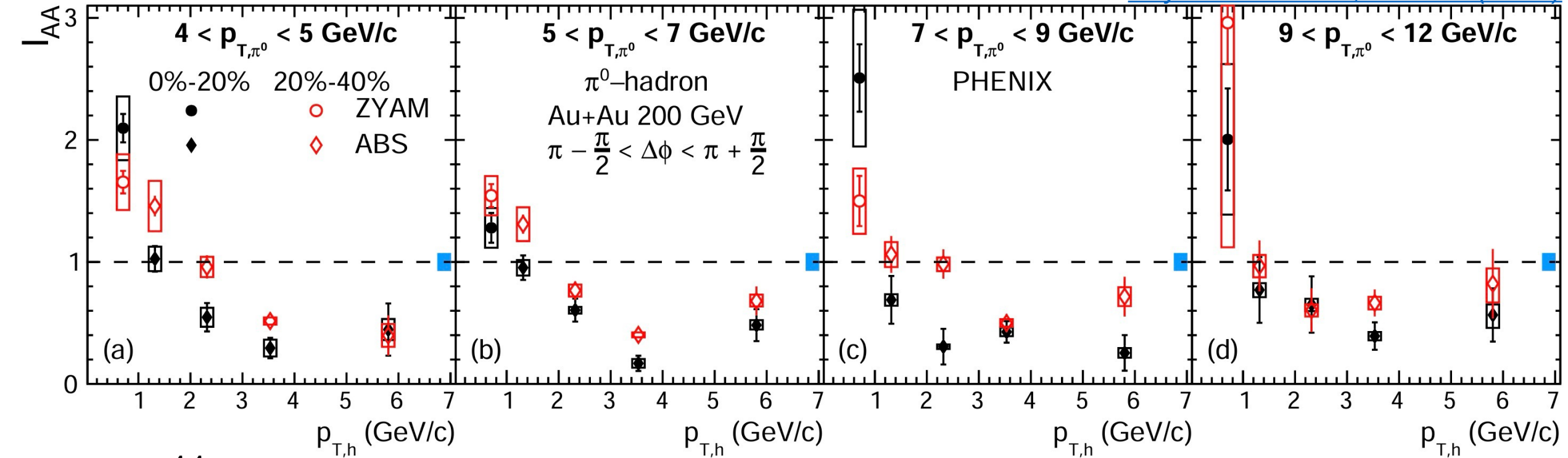
Jet Modification: $I_{AA}(p_T)$

- $I_{AA} = \frac{Y_{AA}^{Away}}{Y_{pp}^{Away}}$, bread-and-butter 2PC measurement
- $I_{AA} = 1 \rightarrow$ No modification
- $I_{AA} < 1 \rightarrow$ Suppression
- $I_{AA} > 1 \rightarrow$ Enhancement
- $I_{AA} = \frac{Y_{AA}}{Y_{pp}} \approx \frac{D_{AA}(z')}{D_{pp}(z)}$



π^0 -Hadron Correlations - $I_{AA}(p_T)$

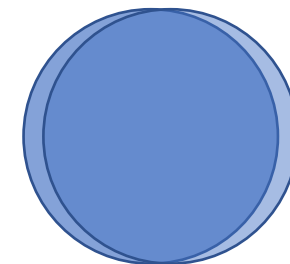
[Phys. Rev. C 110, 044901 \(2024\)](#)



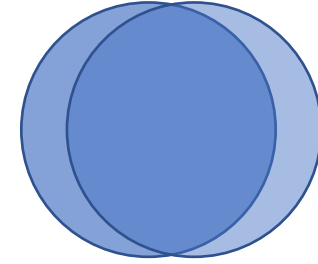
$$I_{AA} = \frac{Y_{Away}^{AA}}{Y_{Away}^{pp}}$$

Newly Published!

- Enhancement of yield ($I_{AA} > 1$) at low associate particle momentum
- Depletion ($I_{AA} < 1$) at high associate particle momentum



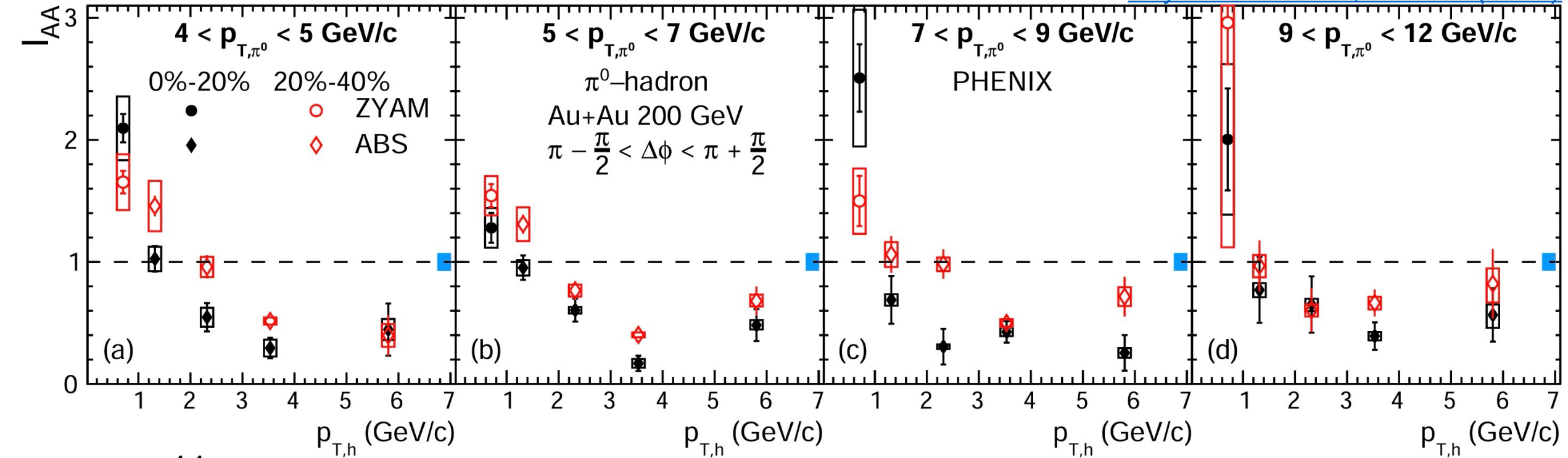
0-20%



20-40%

π^0 -Hadron Correlations - $I_{AA}(p_T)$

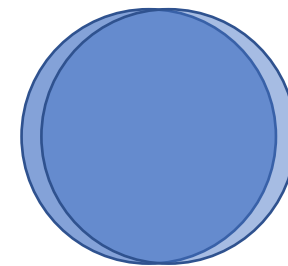
[Phys. Rev. C 110, 044901 \(2024\)](#)



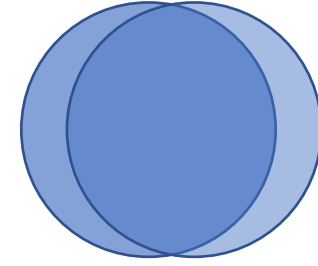
$$I_{AA} = \frac{Y_{Away}^{AA}}{Y_{Away}^{pp}}$$

Newly Published!

- Suppression phenomenon is consistently more severe in most central collisions (black points)



0-20%

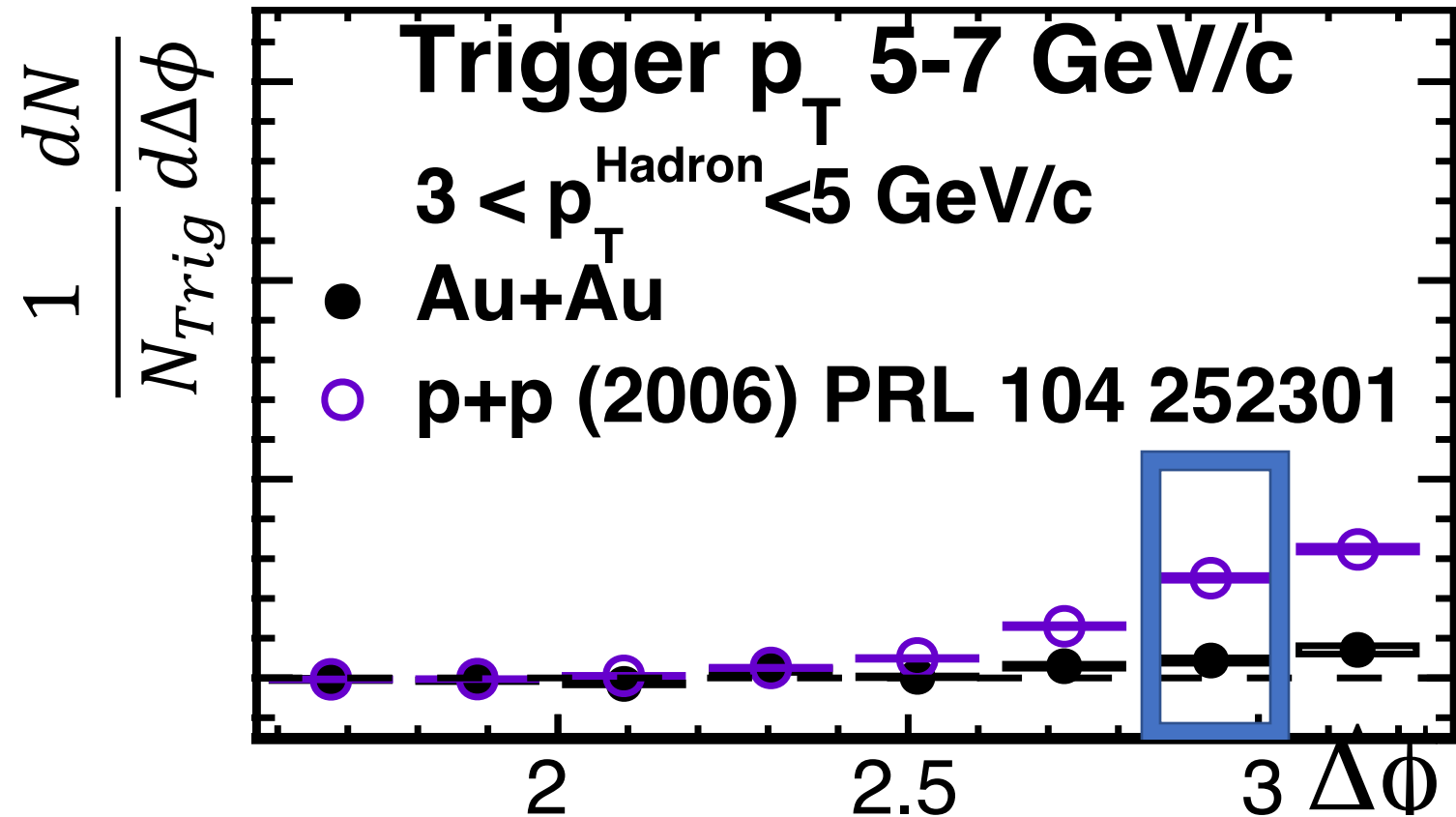


20-40%

Jet Structure Modification: Δ_{AA} vs. $\Delta\phi$

[Phys. Rev. C 110, 044901 \(2024\)](#)

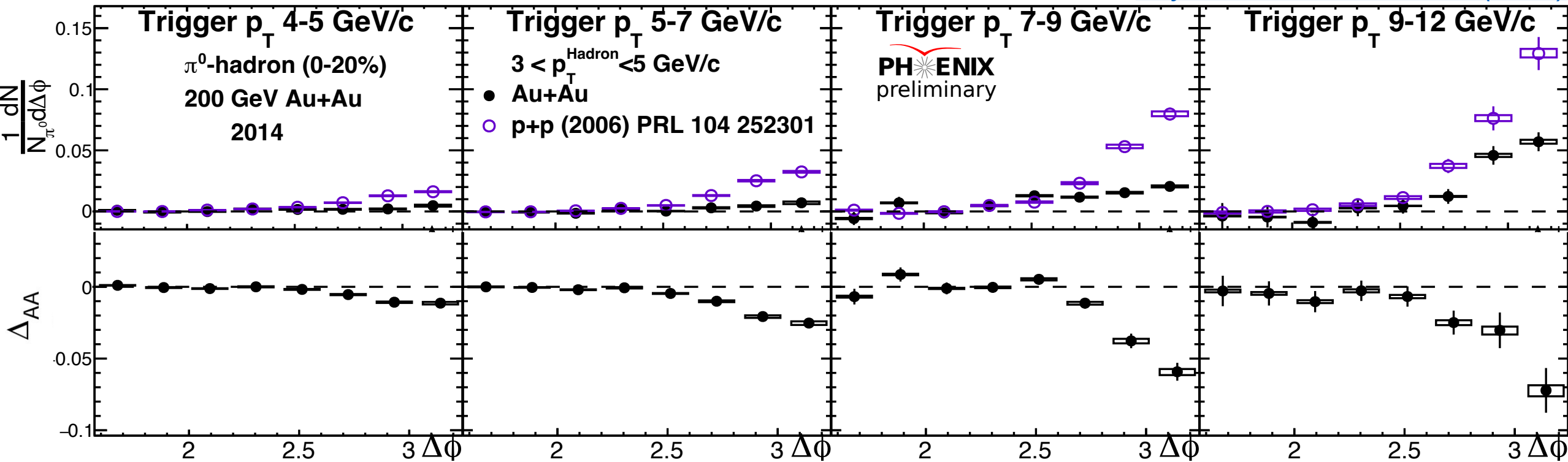
- Can we look at modification “inside” the jet?
- Can take difference between *differential* yields rather than integrated
- How is the actual distribution of particles changed within a jet?



Jet Structure Modification: Δ_{AA} vs. $\Delta\phi$

- $\Delta_{AA} = Y_{AA} - Y_{pp}$
- Captures jet modification across wide range of $\Delta\phi$ values
- Difference, rather than ratio, better behaved for yields ~ 0

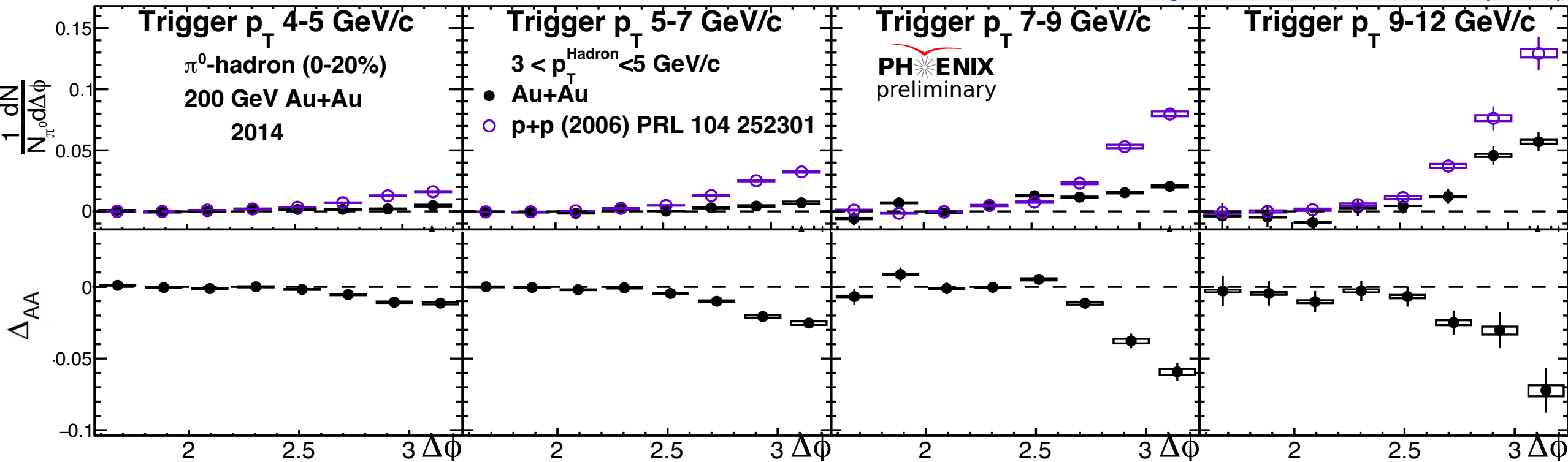
[Phys. Rev. C 110, 044901 \(2024\)](#)



Jet Structure Modification: Δ_{AA} vs. $\Delta\phi$

- For high p_T constituents, can see suppression ($\Delta_{AA} < 0$)
- Suppression is most severe at jet core ($\Delta\phi \sim \pi$)

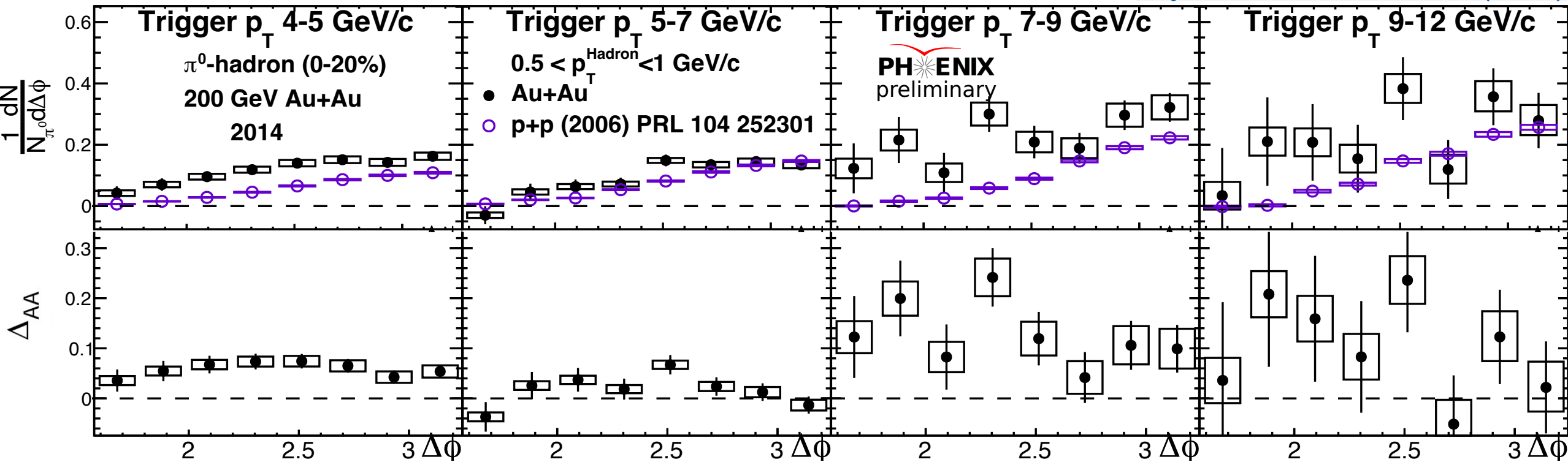
[Phys. Rev. C 110, 044901 \(2024\)](#)



Jet Structure Modification: Δ_{AA} vs. $\Delta\phi$

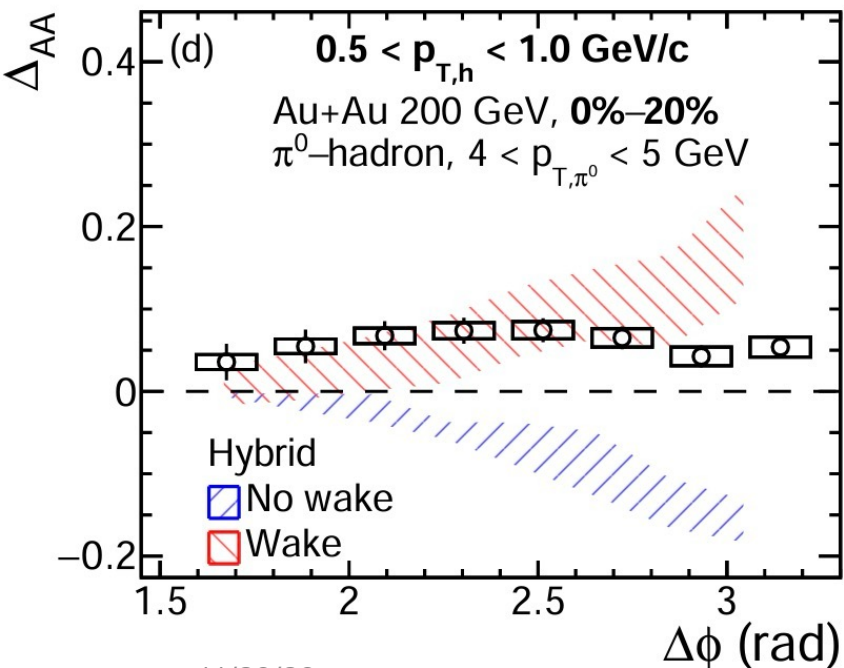
- Enhancement of soft particles seen at wide angles
- For softest jets, this enhancement even appears at the core

[Phys. Rev. C 110, 044901 \(2024\)](#)

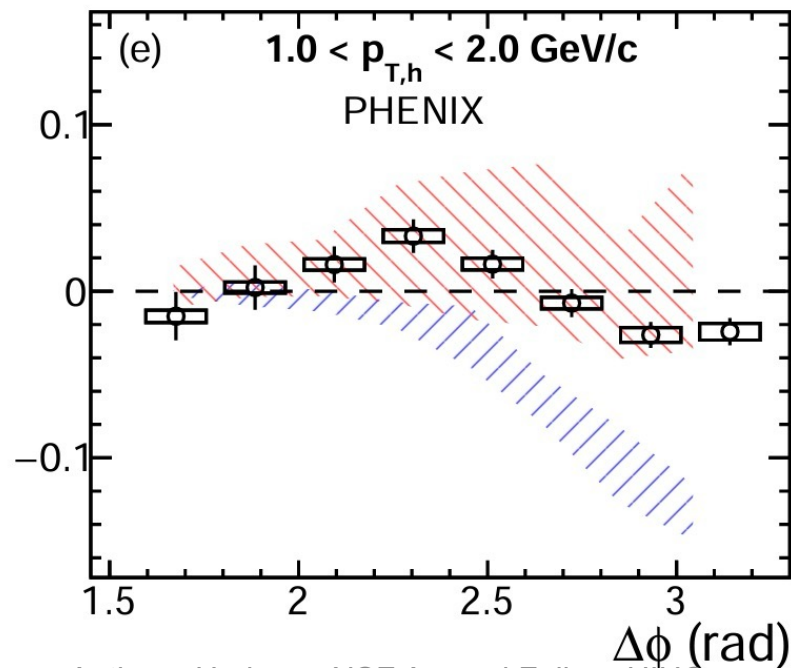


Model Comparisons

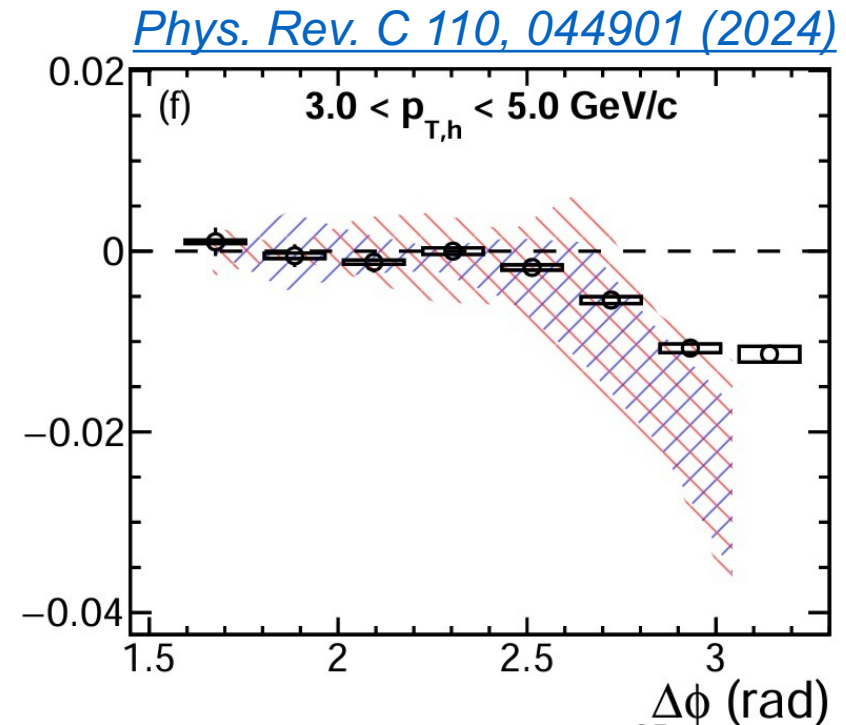
- Here, Hybrid model is shown with two configurations
- **Wake:** includes a medium response in the form a hydrodynamic wake of low p_T particles
- **No wake:** energy loss via pQCD only



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35

The Importance of Medium Response

- Jet modification is an interplay between energy loss *and* the medium's response to that energy

[Phys. Rev. C 110, 044901 \(2024\)](#)

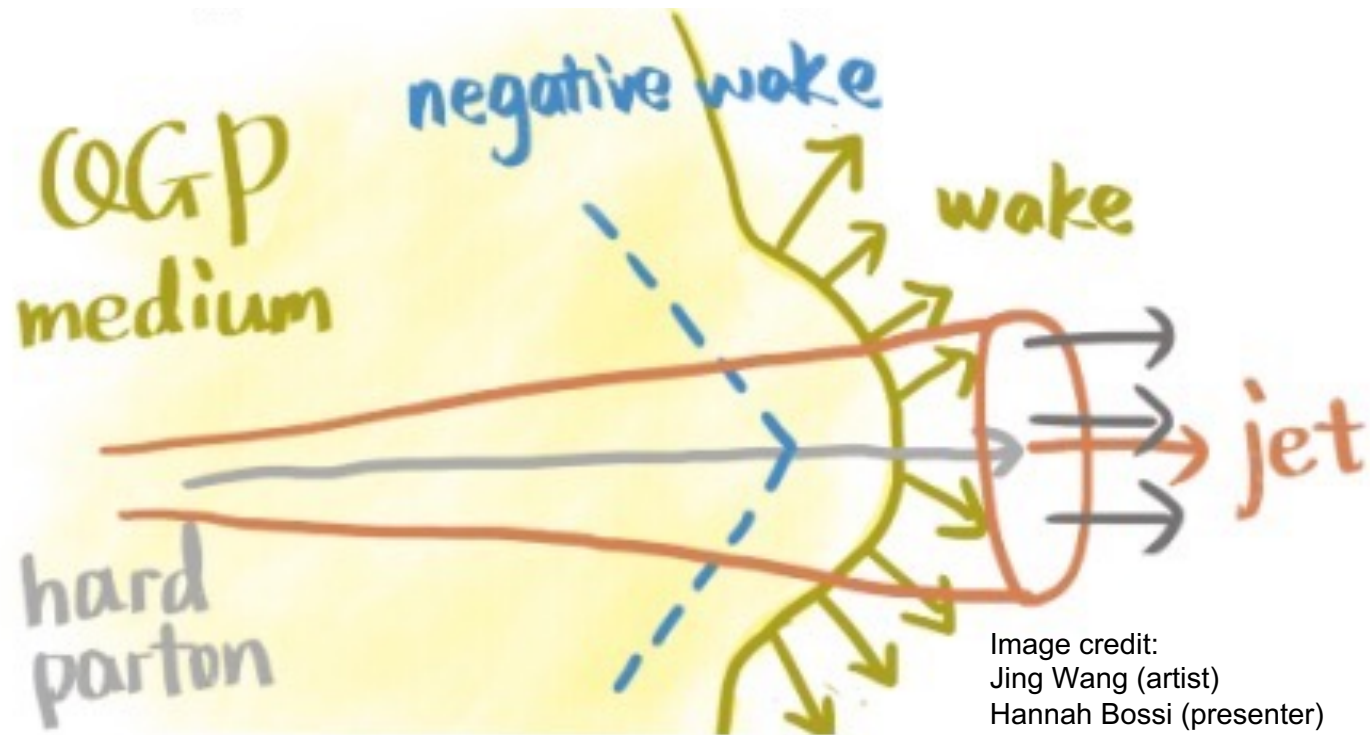
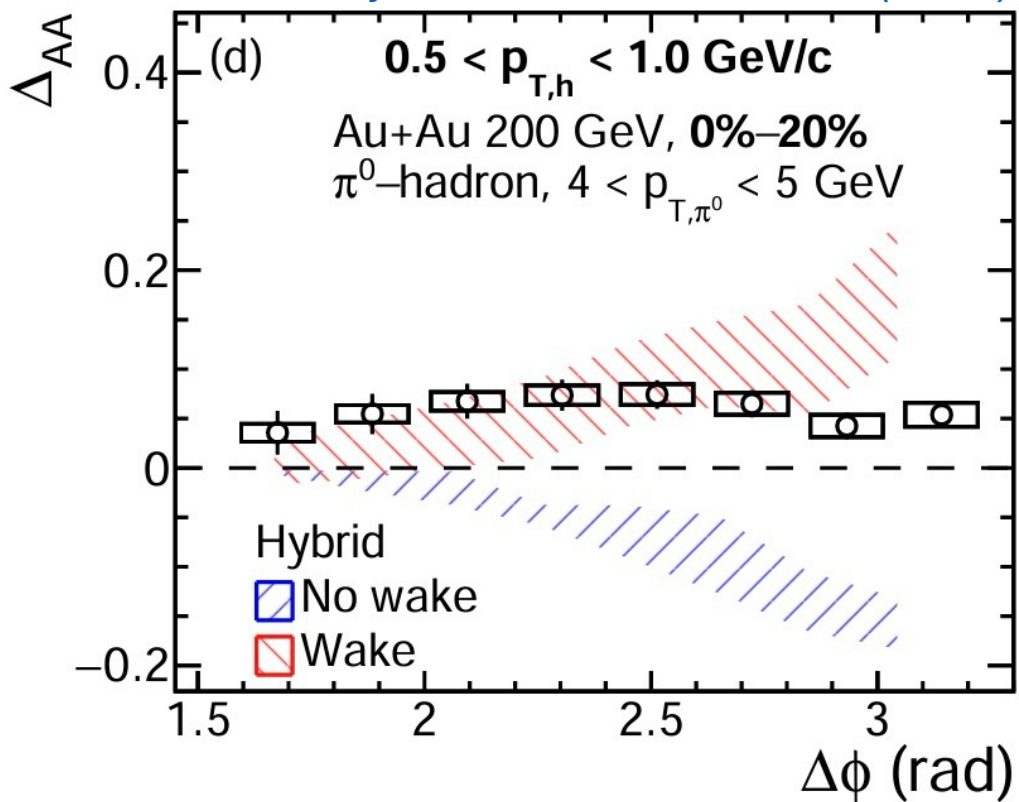
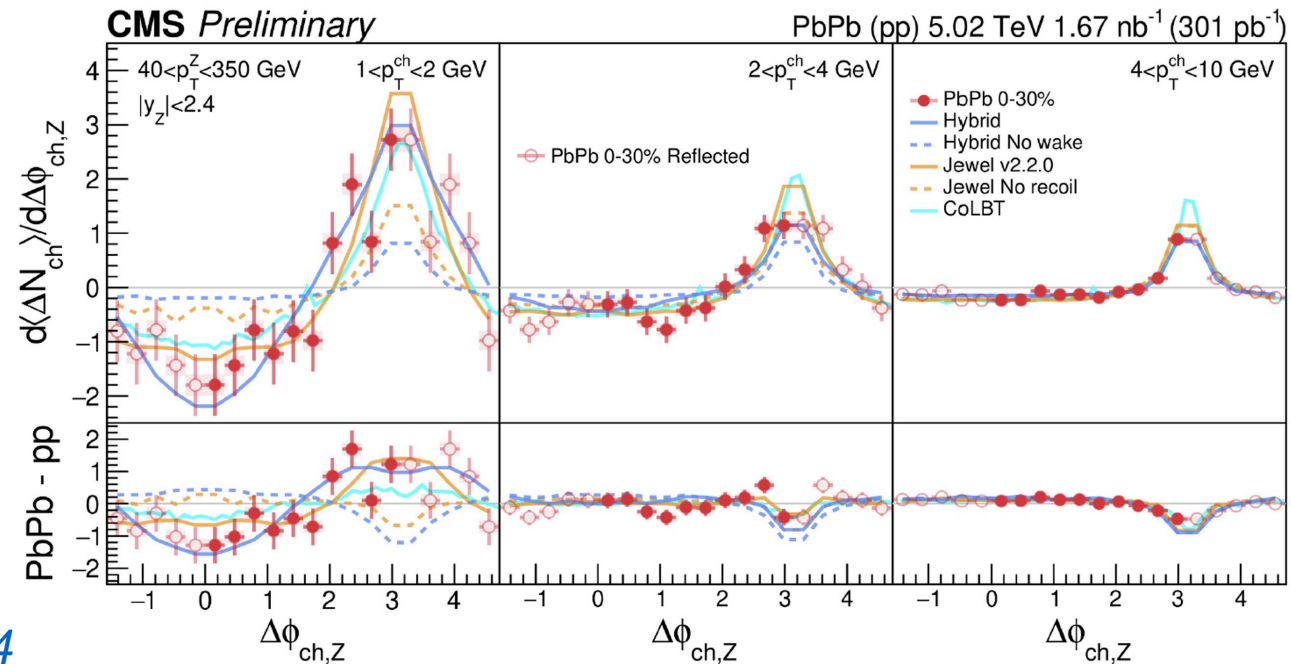


Image credit:
Jing Wang (artist)
Hannah Bossi (presenter)

Turning 180°: The Diffusion Wake

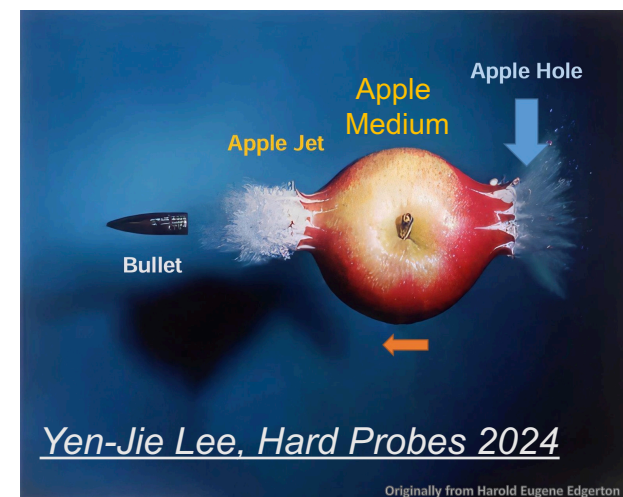
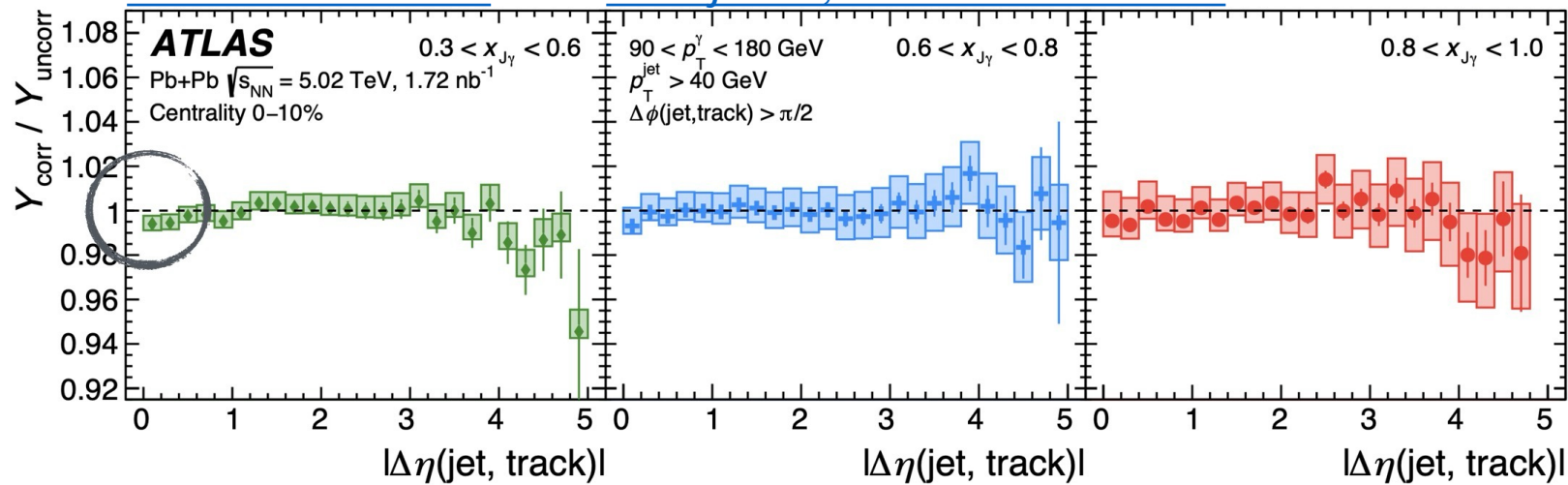
[CMS-PAS-HIN-23-006](#)

- In addition to enhancement wake, recent studies have also sought a corresponding depletion wake near $\Delta\phi \sim 0$



[arXiv:2408.08599](#)

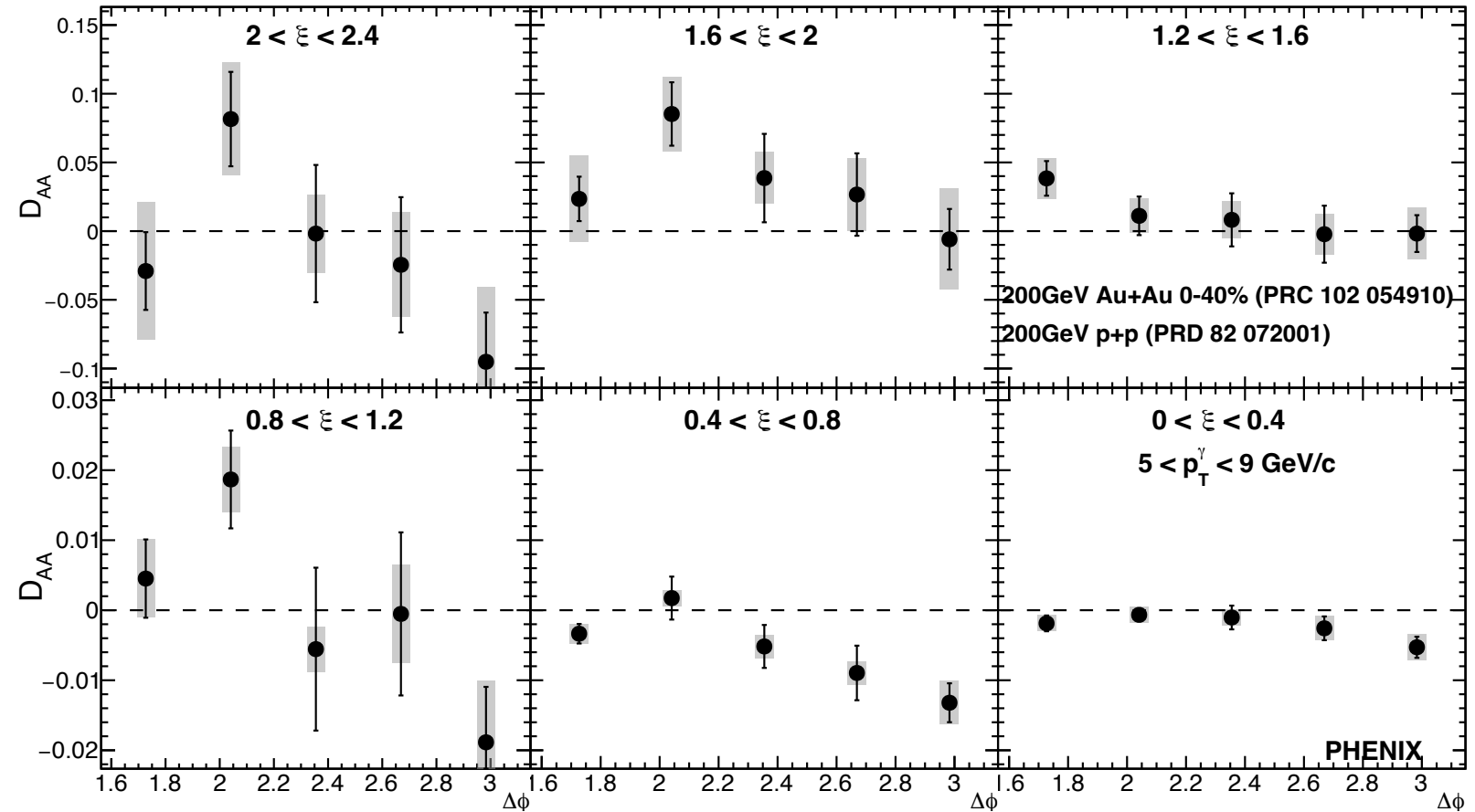
[Yeonju Go, Hard Probes 2024](#)



Future PHENIX Results

[Phys. Rev. C 102, 054910 \(2020\)](#)

- Δ_{AA} extracted from previously published results
- Familiar depletion signal for low ξ (high p_T) bins
- High ξ (low p_T) bins will benefit from larger Run 14 dataset in ongoing analysis



High $\xi \rightarrow$ low p_T

$$\xi = \log(1/z_t)$$

Low $\xi \rightarrow$ high p_T

Summary and Conclusion

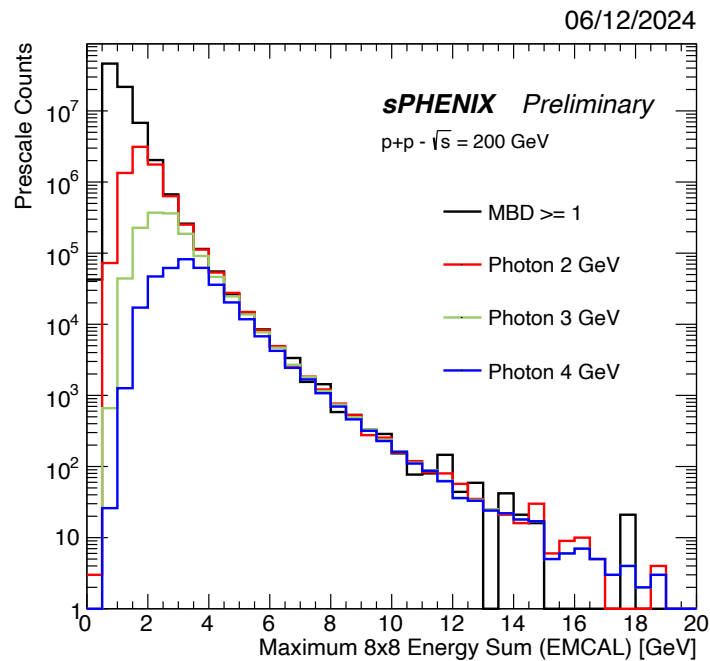
- New two-particle correlation results from PHENIX probe modification to the structure of recoil jets
- The yield of high p_T hadrons associated with jets is found to be suppressed, especially at the jet core
- The yield of low p_T hadrons is enhanced, with this enhancement appearing as far as $\frac{\pi}{2}$ [rad] away from the jet peak
- Enhancement phenomenon is captured well by HYBRID model, where it is owed to a hydrodynamic wake of low p_T particles
- Future PHENIX measurements will employ the large-statistics Run 14 dataset to probe jet modification via prompt photon-hadron correlations

Thank You! Questions?

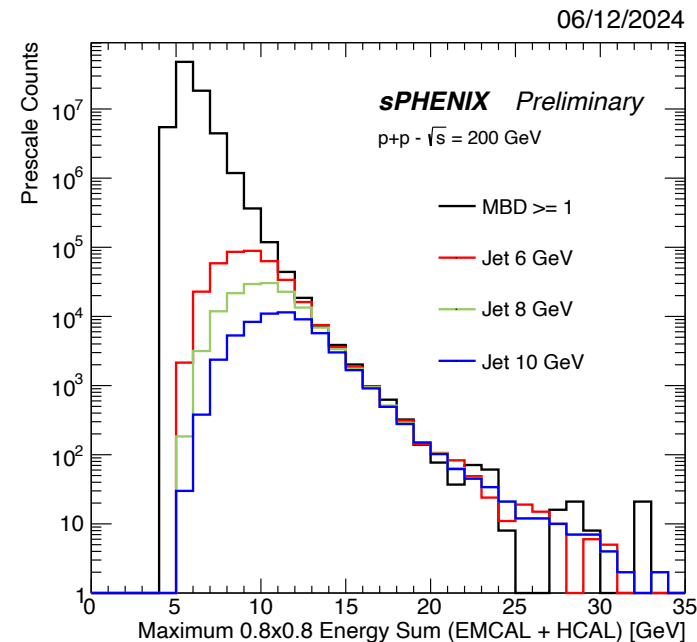
Back-Up

Future Jet-Medium Studies at RHIC

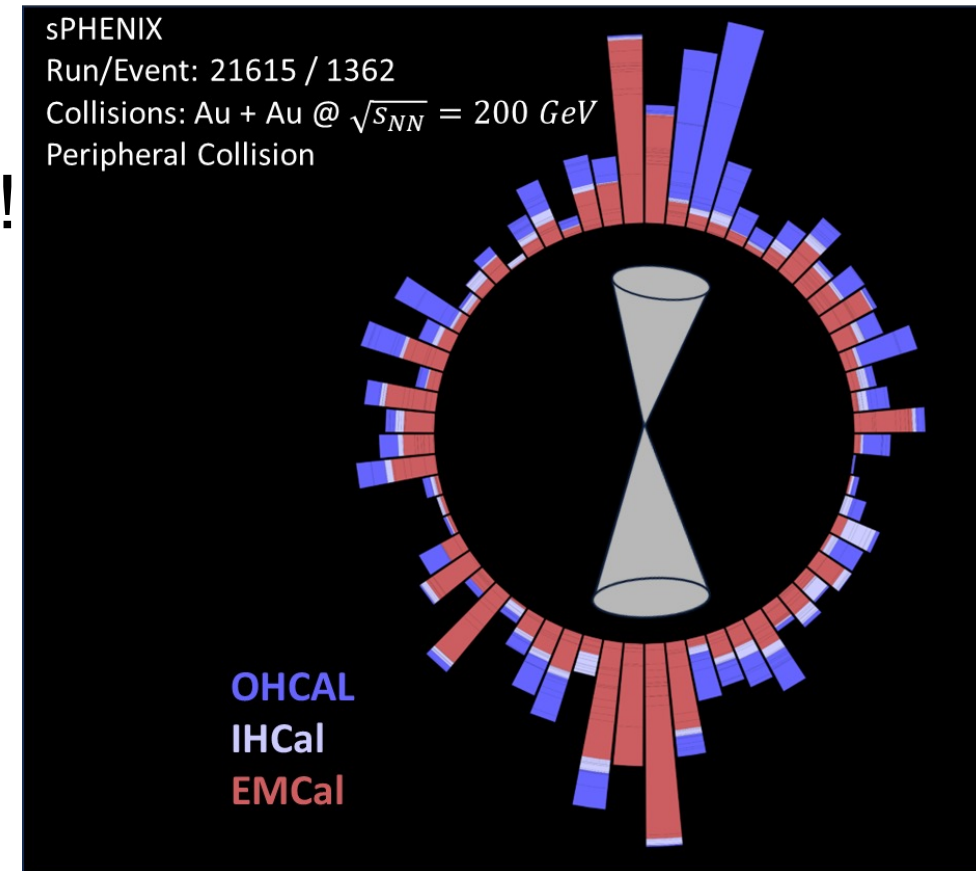
- sPHENIX completed reference $p + p$ run in 2024 exceeding luminosity projections by 200%
- $Au + Au$ run in a matter of months!
- Exciting era of jet-medium interactions awaits!



11/29/23

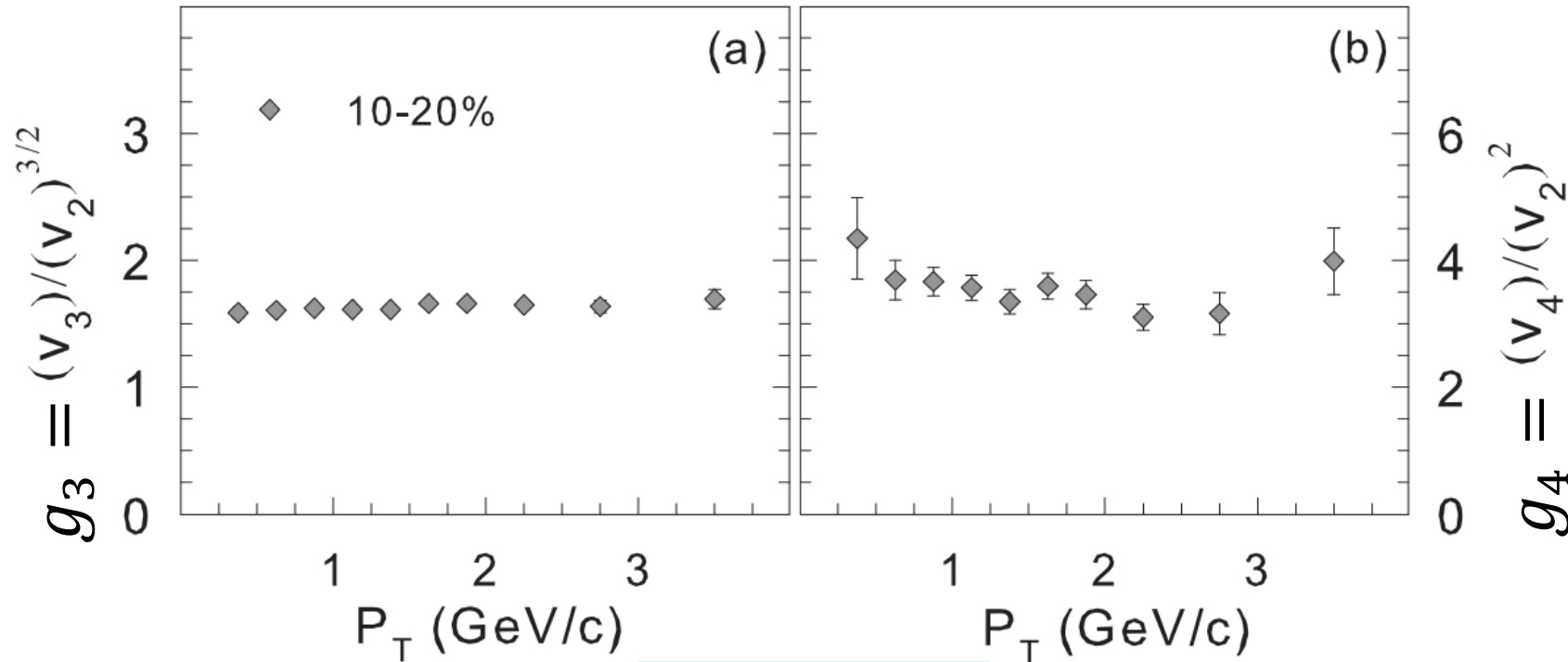


Anthony Hodges, NSF Ascend Fellow, UIUC



42

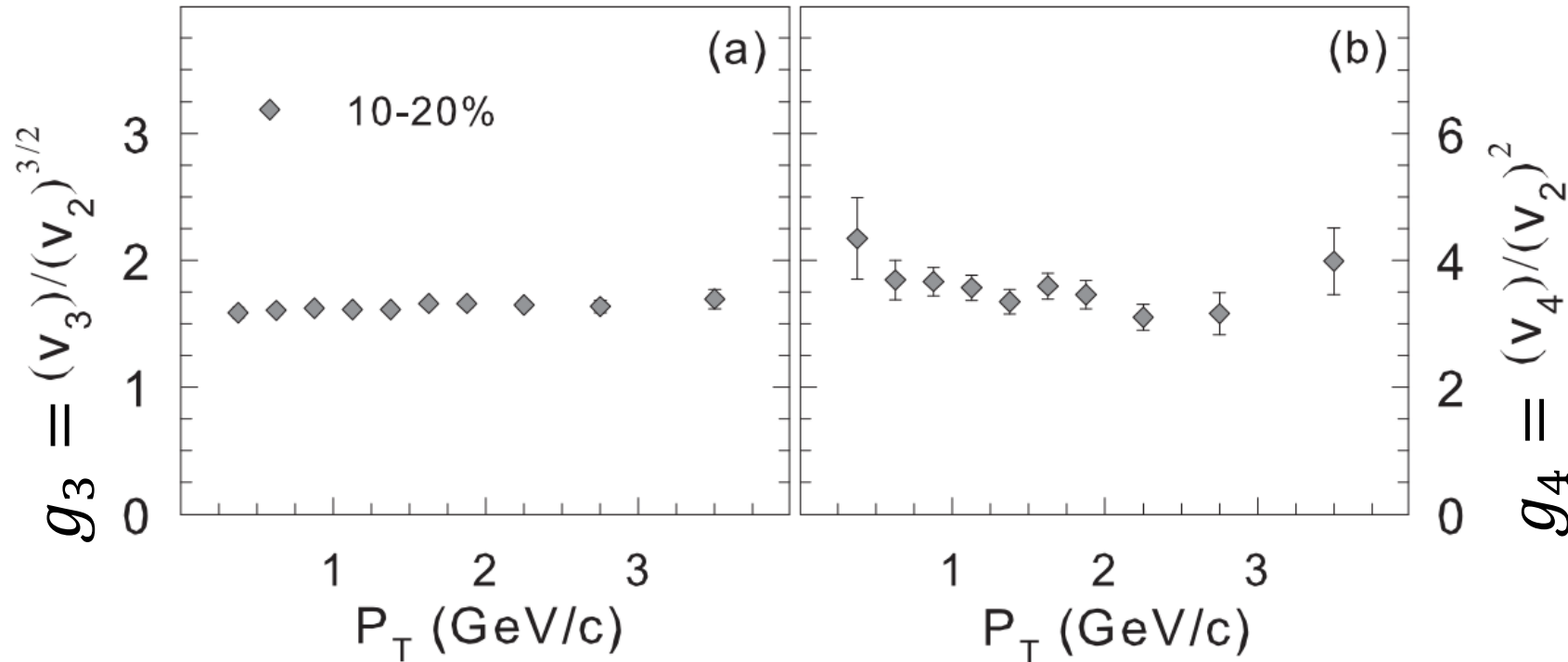
Flow Subtraction – Acoustic Scaling



$$g_n = \frac{v_n}{(v_2)^{n/2}} \quad \text{arXiv:1105.3782v2}$$

- Have charged hadron v_n^a for ($n = 2,3,4$) from PHENIX results
- No π^0 v_3 or v_4 measured at RHIC energies
- v_n harmonics can be scaled to one another via value g_n

Flow Subtraction – Acoustic Scaling



arXiv:1105.3782v2

$$v_n^{\pi^0} = g_n^h (v_2^{\pi^0})^{n/2}$$

- Can calculate $\pi^0 v_3, v_4$ by scaling $\pi^0 v_2$ with charged hadron g_n

Jets in Heavy Ion Collisions

Inside the medium...

