# RELATIVISTIC (A)CAUSALITY IN HYDRODYNAMICS AND ITS EFFECT ON BAYESIAN ANALYSES

#### Matthew Luzum

References: T.S.Domingues, R.Krupczak, J.Noronha, T.N.da Silva, J-F.Paquet, ML; Phys.Rev.C 110 (2024) 6, 064904; arXiv:2409.17127 Arthur Lopez, ML; work in progress

University of São Paulo

Hot Jets: Advancing the Understanding of High Temperature QCD with Jets January 9, 2025

#### HYDRODYNAMIC VALIDITY

- Hydrodynamics central to simulations
- Validity of fluid description not always clear
- Typically derived as expansion around equilibrium
- Sometimes system is far from equilibrium (early times, near jets)
- Not definitive: hydro can be valid far from equilibrium
- Relativistic causality: definitive test
- How important is this issue? Quantify with Bayesian analysis

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- Bayesian inference: used to extract physical properties from data with systematic treatment of uncertainty
- JETSCAPE performed large-scale analysis of soft sector *Phys.Rev.C* 103 (2021) 5, 054904
- How are these results affected if we don't allow hydro to be used in acausal regime?

Norm. Pb-Pb 2.76 TeV Norm. Au-VA 200 GeV generalized mean nucleon width min. dist. btw. nucleons multiplicity fluctuation free-streaming time scale free-streaming time scale free-streaming energy dep. particitization temperature	N[2.76  TeV] N[0.2  TeV] p w $d_{min}^{3}$ $\sigma_{k}$ $\tau_{R}$ $\alpha$ $T_{SW}$	[10, 20] [3, 10] [-0.7, 0.7] [0.5, 1.5] fm [0, 1, 7 <sup>3</sup> ] fm <sup>3</sup> [0,3, 2.0] [0,3, 2.0] fm/c [-0.3, 0.3] [0,135, 0.165] GeV	temperatur $(\eta / s)$ at k low temp. s high temp. shear relax maximum of temperatur width of ( $\zeta$ asymmetry
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e of (n/s) kink ink slope of (n/s)slope of  $(\eta / s)$ ation time factor of ( ć / s) e of (c /s) peak /s) peak of  $(\zeta / s)$  peak

 $T_n$ [0.13, 0.3] GeV [0.01, 0.2]  $(n/s)_{ink}$ [-2, 1] GeV-1 alow [-1, 2] GeV-1 ahigh b<sub>π</sub> [2.8] (C/S)max [0.01. 0.25] TC [0.12, 0.3] GeV [0.025, 0.15] GeV wζ [-0.8. 0.8] DQC

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Norm. Au-Au 200 GeV
generalized mean
nucleon width
min. dist. btw. nucleons
multiplicity fluctuation
free-streaming time scale
free-streaming energy dep.
particlization temperature

N[2.76 TeV]
N[0.2 TeV]
р
w
a <sup>3</sup> min
$\sigma_k$
<sup>T</sup> R
α
T <sub>SW</sub>

[10, 20]
[3, 10]
[-0.7, 0.7]
[0.5, 1.5] fm
[0, 1.7 <sup>3</sup> ] fm <sup>3</sup>
[0.3, 2.0]
[0.3, 2.0] fm/c
[-0.3, 0.3]
[0.135, 0.165] Ge

temperature of $(\eta/s)$ kink $(\eta/s)$ at kink
low temp. slope of $(\eta/s)$
high temp. slope of $(\eta/s)$
shear relaxation time factor maximum of $(\zeta/s)$ temperature of $(\zeta/s)$ peak width of $(\zeta/s)$ peak asymmetry of $(\zeta/s)$ peak

[0.13, 0.3] GeV  $T_n$  $(n/s)_{kink}$ [0.01. 0.2] [-2, 1] GeV - 1 alow [-1, 2] GeV-1 ahigh b<sub>π</sub> [2.8] (C/S)max [0.01. 0.25] [0.12, 0.3] GeV [0.025, 0.15] GeV [-0.8. 0.8] nar

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eused	0.25 0.20 0.15 0.10 0.05 0.00 0.10	0.15 0.20 0.25 T [GeV]	0.30 0.35 0.4
e of $(\eta/s)$ kink hk lope of $(\eta/s)$ slope of $(\eta/s)$ ation time facto f $(\zeta/s)$	or	$T_{\eta} (\eta/s)_{kink}$ $a_{low}$ $a_{high}$ $b_{\pi} (\zeta/s)_{max}$	[0.13, 0.3] GeV [0.01, 0.2] [-2, 1] GeV <sup>-1</sup> [-1, 2] GeV <sup>-1</sup> [2, 8] [0.01, 0.25]

TC

w

0.35 -- Prior 60% C.I.

0.30

Specific shear viscosity posterior

Without causality analysis 60% C.L.

Norm. Pb-Pb 2.76 TeV	N[2.76 TeV]	[10, 20]	temperature of $(\eta/s)$ kink $(\eta/s)$ at kink
Norm. Au-Au 200 GeV	N[0.2 TeV]	[3, 10]	
generalized mean	p	[-0.7, 0.7]	low temp. slope of $(\eta/s)$
nucleon width	w	[0.5, 1.5] fm	high temp. slope of $(\eta/s)$
min. dist. btw. nucleons	d <sup>3</sup> min	[0, 1.7 <sup>3</sup> ] fm <sup>3</sup>	shear relaxation time factor
multiplicity fluctuation	<sup>σ</sup> k	[0.3, 2.0]	maximum of $(\zeta / s)$
free-streaming time scale	<sup>τ</sup> B	[0.3, 2.0] fm/ <i>c</i>	temperature of $(\zeta / s)$ peak
free-streaming energy dep.	α	[–0.3, 0.3]	width of $(\zeta/s)$ peak
particlization temperature	T <sub>SW</sub>	[0.135, 0.165] GeV	asymmetry of $(\zeta/s)$ peak

[0.12, 0.3] GeV

[0.025, 0.15] GeV [-0.8, 0.8]

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	e of (1
generalized mean p [-0.7, 0.7] low temp. slope	· · · ·
nucleon width w [0.5, 1.5] fm high temp. slope	be of (
min. dist. btw. nucleons d <sup>3</sup> <sub>min</sub> [0, 1.7 <sup>3</sup> ] fm <sup>3</sup> shear relaxation	n time
multiplicity fluctuation $\sigma_k$ [0.3, 2.0] maximum of $(\zeta)$	;/s)
free-streaming time scale $\tau_R^{-1}$ [0.3, 2.0] fm/c temperature of (	(ζ/s
free-streaming energy dep. $\alpha$ [-0.3, 0.3] width of ( $\zeta$ /s) p	peak
particlization temperature T <sub>SW</sub> [0.135, 0.165] GeV asymmetry of (4	$\zeta/s$



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## ISRAEL-STEWART HYDRODYNAMICS

• Modern hydrodynamic theory used in simulations:

$$T^{\mu\nu} = \epsilon u^{\mu} u^{\nu} - (P + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu}$$
  

$$\partial_{\mu} T^{\mu\nu} = \mathbf{0},$$
  

$$\tau_{\Pi} \dot{\Pi} + \Pi = -\zeta \theta - \delta_{\Pi\Pi} \Pi \theta + \lambda_{\Pi\pi} \pi^{\mu\nu} \sigma_{\mu\nu}$$
  

$$\tau_{\pi} \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = \mathbf{2}\eta \sigma^{\mu\nu} - \delta_{\pi\pi} \pi^{\mu\nu} \theta + \varphi_{7} \pi^{\langle\mu}_{\alpha} \pi^{\nu\rangle\alpha} - \tau_{\pi\pi} \pi^{\langle\mu}_{\alpha} \sigma^{\nu\rangle\alpha} + \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu}$$

• Common parameterizations:

$$\tau_{\Pi} = b_{\Pi} \frac{\zeta}{\left(\frac{1}{3} - c_s^2\right)^2 (\epsilon + p)}$$
$$\tau_{\pi} = b_{\pi} \frac{\eta}{sT}$$

$$\frac{\delta_{\Pi\Pi}}{\tau_{\Pi}} = \frac{2}{3}$$
$$\frac{\delta_{\pi\pi}}{\tau_{\pi}} = \frac{4}{3}$$
$$\frac{\tau_{\pi\pi}}{\tau_{\pi}} = \frac{10}{7}$$

$$egin{aligned} &rac{\lambda_{\pi\Pi}}{ au_{\pi}}=rac{6}{5}\ &rac{\lambda_{\Pi\pi}}{ au_{\Pi}}=rac{8}{5}\left(rac{1}{3}-c_{s}^{2}
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$$\tau_{\pi} \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = 2\eta \sigma^{\mu\nu} - \delta_{\pi\pi} \pi^{\mu\nu} \theta + \varphi_{7} \pi^{\langle\mu}_{\alpha} \pi^{\nu\rangle\alpha} - \tau_{\pi\pi} \pi^{\langle\mu}_{\alpha} \sigma^{\nu\rangle\alpha} + \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu}$$

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$$\begin{aligned} \frac{\lambda_{\pi\Pi}}{\tau_{\pi}} &= \frac{6}{5} \\ \frac{\lambda_{\Pi\pi}}{\tau_{\Pi}} &= \frac{8}{5} \left( \frac{1}{3} - c_s^2 \right) \end{aligned}$$

- Simplest situation: infinitessimal perturbation around static global equilibrium
- Linearize equations of motion. Demanding signal propagation *v* < *c* gives condition

- Prior (and posterior) allows violation of linear causality!
- Small dependence of observables on τ<sub>π</sub> and τ<sub>Π</sub> gives flat posterior and no strong effect on conclusions about other parameters



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• Recently-derived (necessary) conditions for general, nonlinear case: *Phys. Rev. Lett. 126, 222301 (2021)* 

$$n_1 \equiv rac{2}{b_\pi} + rac{\lambda_\pi \Pi}{ au_\pi} rac{\Pi}{arepsilon + P} - rac{ au_\pi \pi}{2 au_\pi} rac{|\Lambda_1|}{arepsilon + P} \ge 0,$$

$$n_2\equiv 1-rac{1}{b_\pi}+\left(1-rac{\lambda_{\pi\Pi}}{2 au_\pi}
ight)rac{\Pi}{arepsilon+eta}-rac{ au_{\pi\pi}}{4 au_\pi}rac{\Lambda_3}{arepsilon+eta}\geq 0, \ \geq 0,$$

$$n_3\equiv rac{1}{b_\pi}+rac{\lambda_{\pi\Pi}}{2 au_\pi}rac{\Pi}{arepsilon+P}-rac{ au_{\pi\pi}}{4 au_\pi}rac{\Lambda_3}{arepsilon+P} \ge 0,$$

$$n_4 \equiv 1 - rac{1}{b_\pi} + \left(1 - rac{\lambda_{\pi\Pi}}{2 au_\pi}
ight) rac{\Pi}{arepsilon + P} + \left(1 - rac{ au_{\pi\pi}}{4 au_\pi}
ight) rac{\Lambda_a}{arepsilon + P} - rac{ au_{\pi\pi}}{4 au_\pi} rac{\Lambda_d}{arepsilon + P} \ge 0,$$

$$n_5 \equiv c_s^2 + \frac{4}{3} \frac{1}{b_\pi} + b_\Pi \left(\frac{1}{3} - c_s^2\right)^2 + \left(\frac{2}{3} \frac{\lambda_{\pi\Pi}}{\tau_\pi} + \frac{\delta_{\Pi\Pi}}{\tau_\Pi} + c_s^2\right) \frac{\Pi}{\varepsilon + P} \left(\frac{3\delta_{\pi\pi} + \tau_{\pi\pi}}{3\tau_\pi} + \frac{\lambda_{\Pi\pi}}{\tau_\Pi} + c_s^2\right) \frac{\Lambda_1}{\varepsilon + P} \ge 0.$$

$$n_6 \equiv 1 - \left(c_s^2 + \frac{4}{3}\frac{1}{b_\pi} + b_\Pi \left(\frac{1}{3} - c_s^2\right)^2\right) + \left(1 - \frac{2}{3}\frac{\lambda_{\pi\Pi}}{\tau_\pi} - \frac{\delta_{\Pi\Pi}}{\tau_\Pi} - c_s^2\right)\frac{\Pi}{\varepsilon + P} + \left(1 - \frac{3\delta_{\pi\pi} + \tau_{\pi\pi}}{3\tau_\pi} - \frac{\lambda_{\Pi\pi}}{\tau_\Pi} - c_s^2\right)\frac{\Lambda_3}{\varepsilon + P} \ge 0.$$

• In practice, *n*<sub>6</sub> is the most stringent condition

## QUANTIFYING ACAUSALITY



- We perform a *b* = 0 simulation and quantify the fraction of system (defined by totel energy) that is in an acausal regime at onset of hydrodynamics
- What happens if we make cuts on the posterior?



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# OBSERVABLES (MAP WITH AND WITHOUT ACAUSALITY CUTS)



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- Ability to fit data not destroyed by stringent causality demands
- $\bullet\,$  Maximum probability (Maximum a Posteriori, MAP) of original posterior is  $\sim$  3 times as likely as best-fit after the strongest cut







• Demanding causality alters posteriors





 Demanding causality alters posteriors



### **1D** PARAMETER POSTERIORS

- 1D marginalized posterior distributions
- Demanding causality alters posteriors



### VISCOSITY POSTERIORS



- Shear viscosity not signifantly affected
- Large bulk viscosity disfavored by causality cuts

### VISCOSITY POSTERIORS



- Hydodynamic simulations typically enter acausal regimes, at least sometimes
- Demanding limits on acauasality has nonnegligible effects on existing Bayesian analyses
- In the era of precision heavy-ion physics, it is an issue that should be addressed
  - Improve pre-hydrodynamic description
  - Further developments in hydrodynamic theory

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# BONUS: JET/MEDIUM INTERACTION

- Early times: system is far from equilibrium and must thermalize/hydrodynamize
- Same considerations near jets: energy lost by the jet must thermalize/hydrodynamize
- May have observable affects?



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## EXTRA SLIDES

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